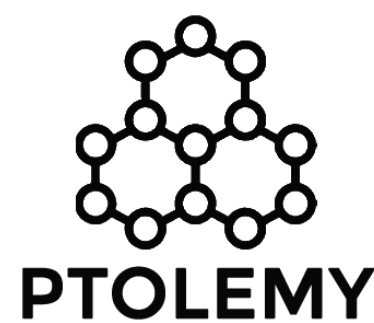
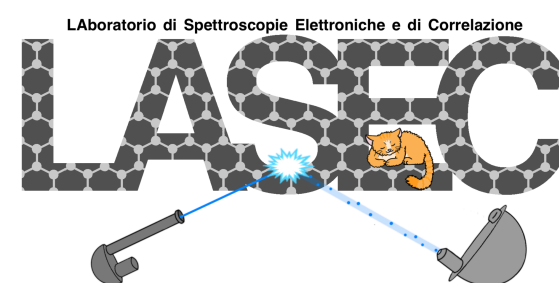


Depth Profile of the Hydrogenation of Multi Walls Carbon Nanotubes with Soft-to-Hard X-ray Photoemission Spectroscopy

FisMat2025, Jul 9th 2025

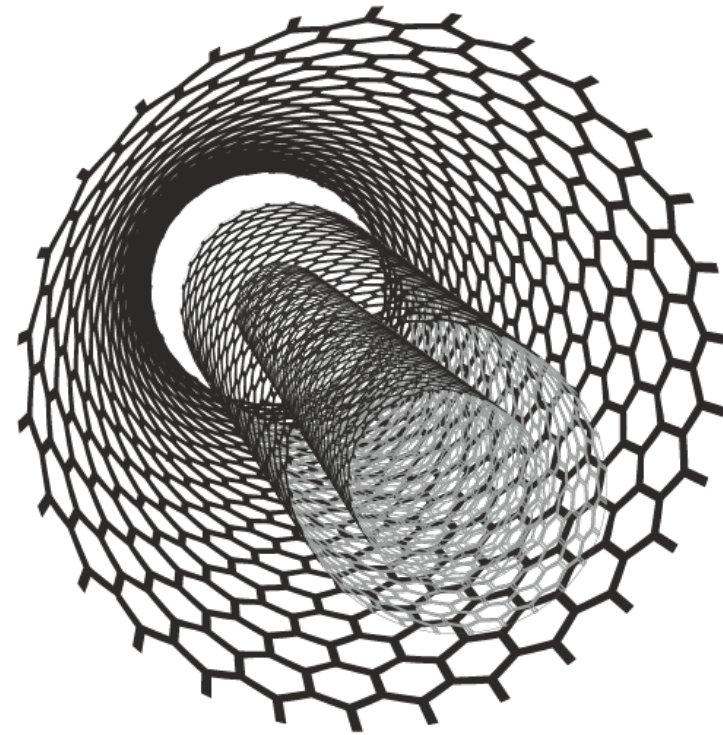
Orlando Castellano, Alice Apponi, Daniele Paoloni, Luca Cecchini, Francesco Pandolfi, Ilaria Rago, Carlo Mariani, Gianluca Cavoto, Tien-Lin Lee, Pardeep Kumar Thakur, Francesco Offi, Alessandro Ruocco



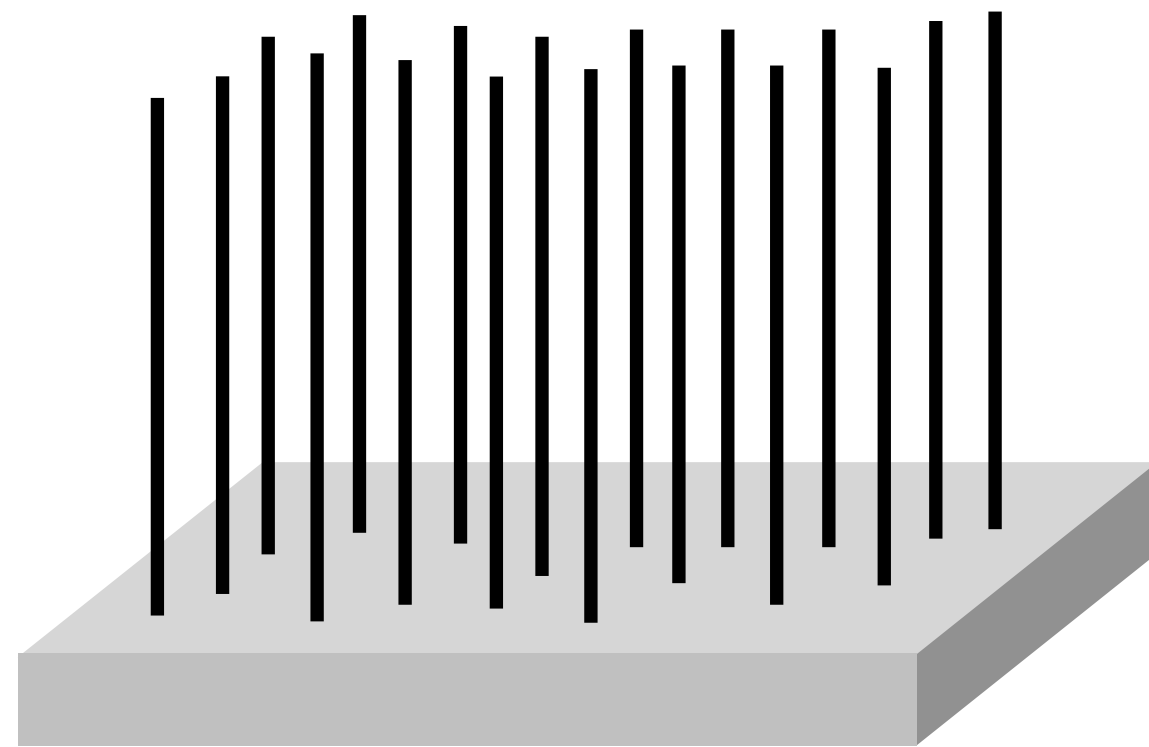
**FisMat
2025**



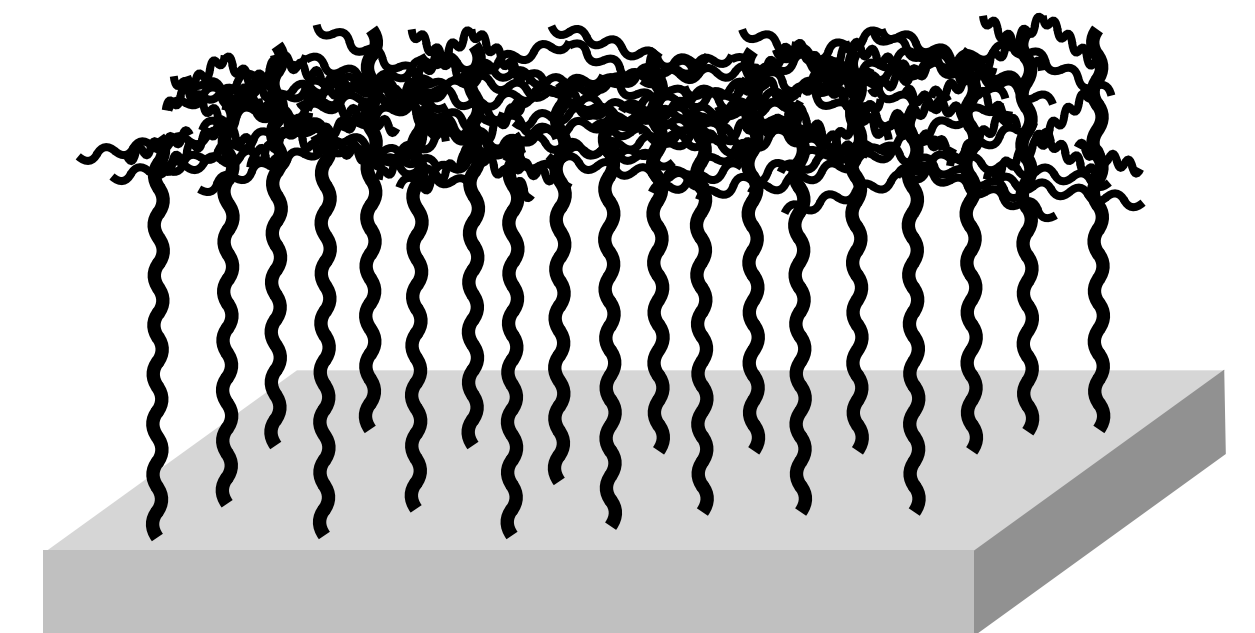
Multi wall CNT



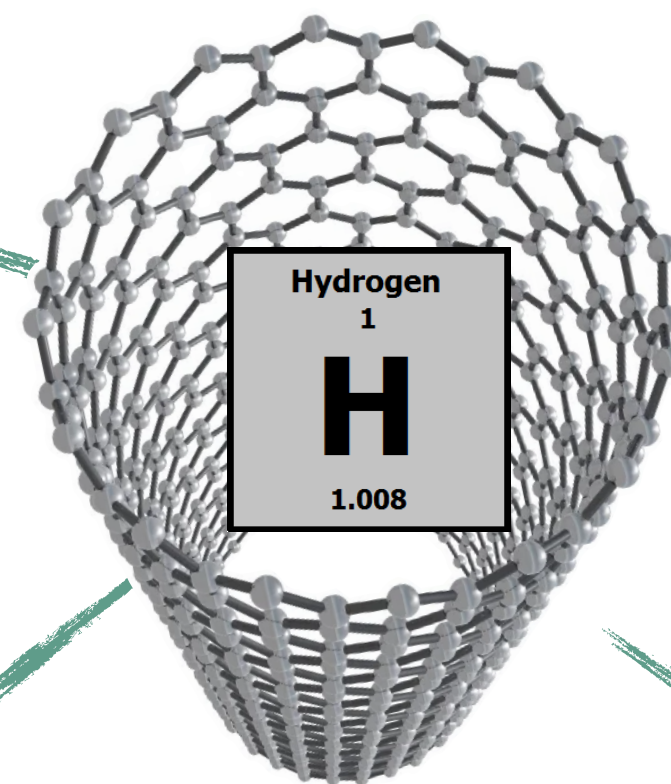
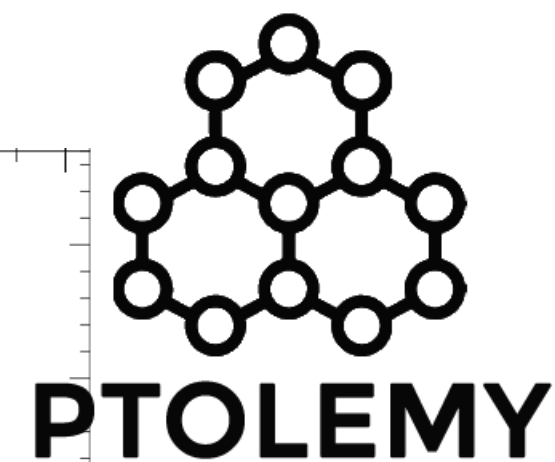
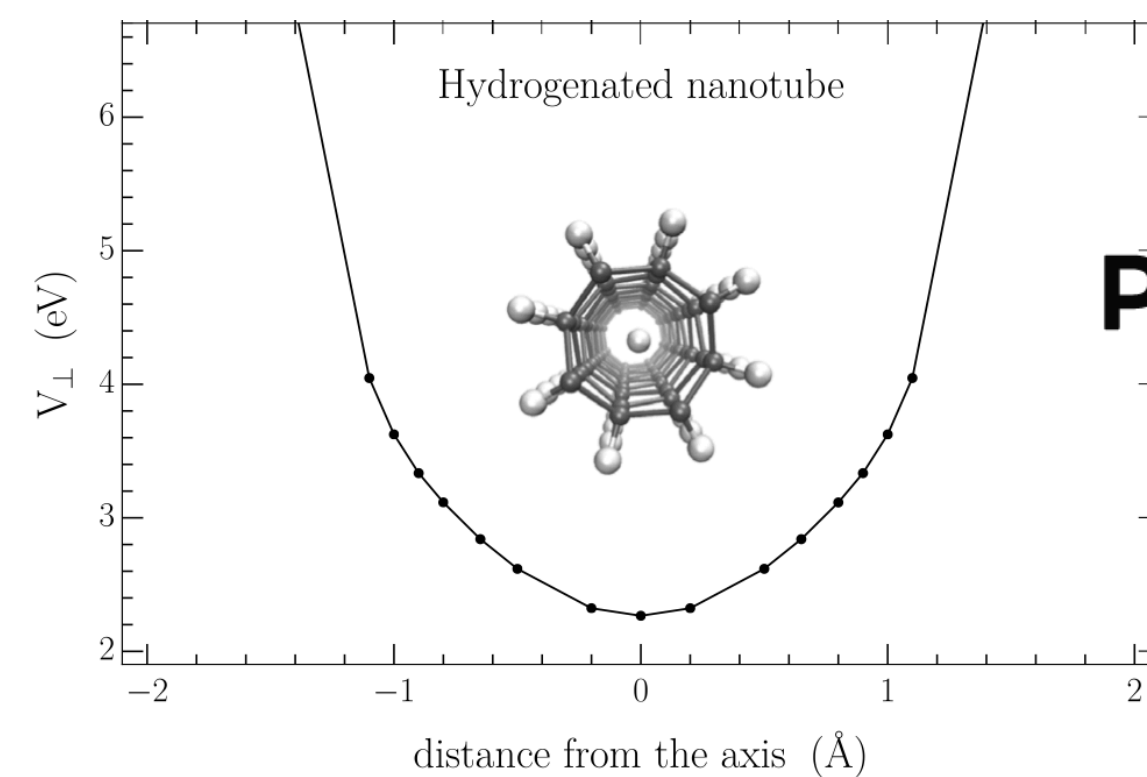
Carbon Nanotubes (CNTs) are rolled-up sheets of graphene forming cylinders with nanometer-scale diameter



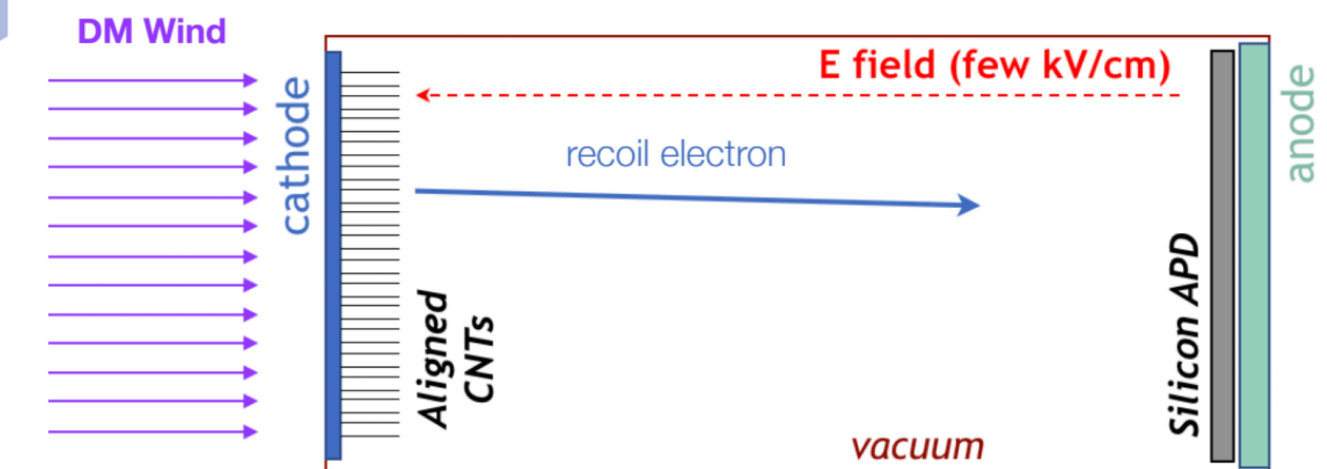
Ideal vs realistic Vertically Aligned CNT forest



Solid Tritium target for Neutrino Mass Measurement



Directional Dark Matter Detector



Tunability of Electronic Properties

Hydrogen Storage

Two Samples and a Plasma Recipe for Hydrogenation

4



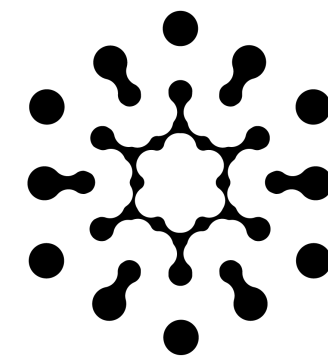
CNT pristine

Annealing

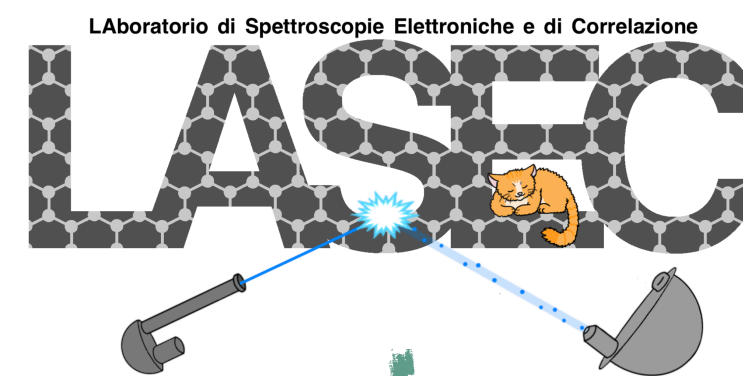
Used as a reference



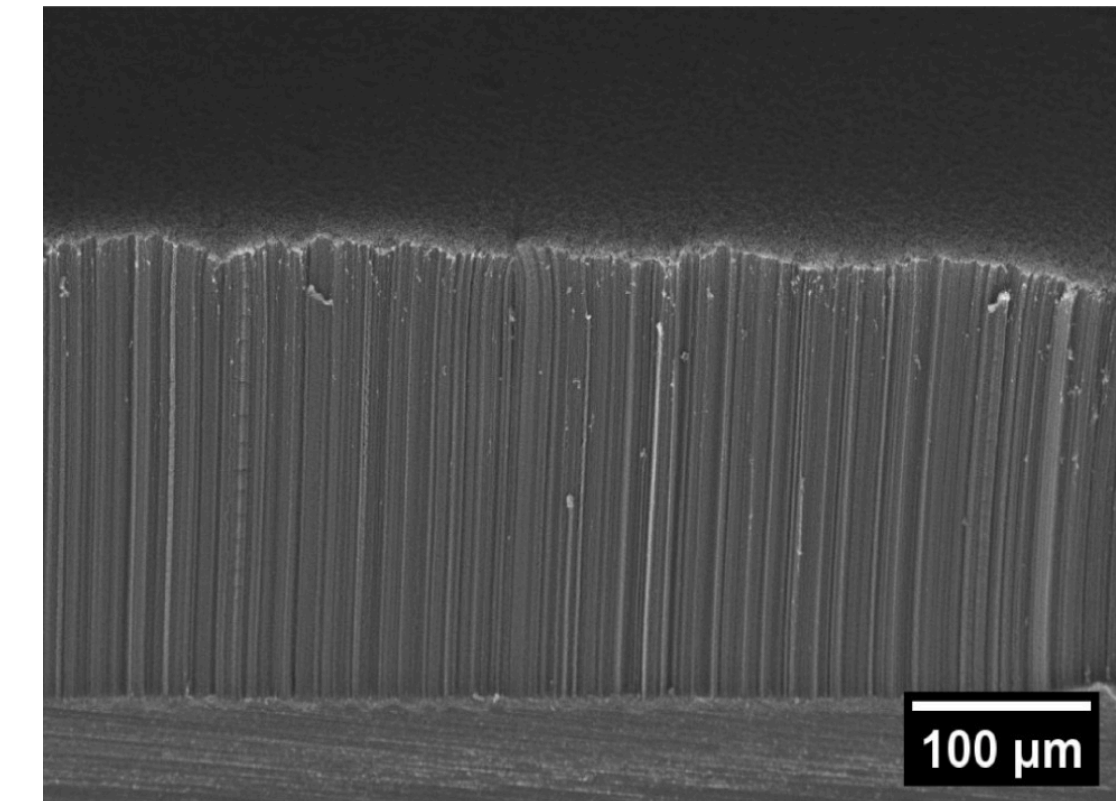
Two samples from the same batch
Annealing 650 °C



TITAN LAB
TECHNOLOGY INNOVATION THROUGH
ADVANCED NANOSTRUCTURES



 **diamond**
I09 beamline



CNT plasma H

Exposed to hydrogen plasma:

- 100 W
- 0.7 mbar H₂, 300 sccm
- 1 h

Kept in low vacuum during transfer

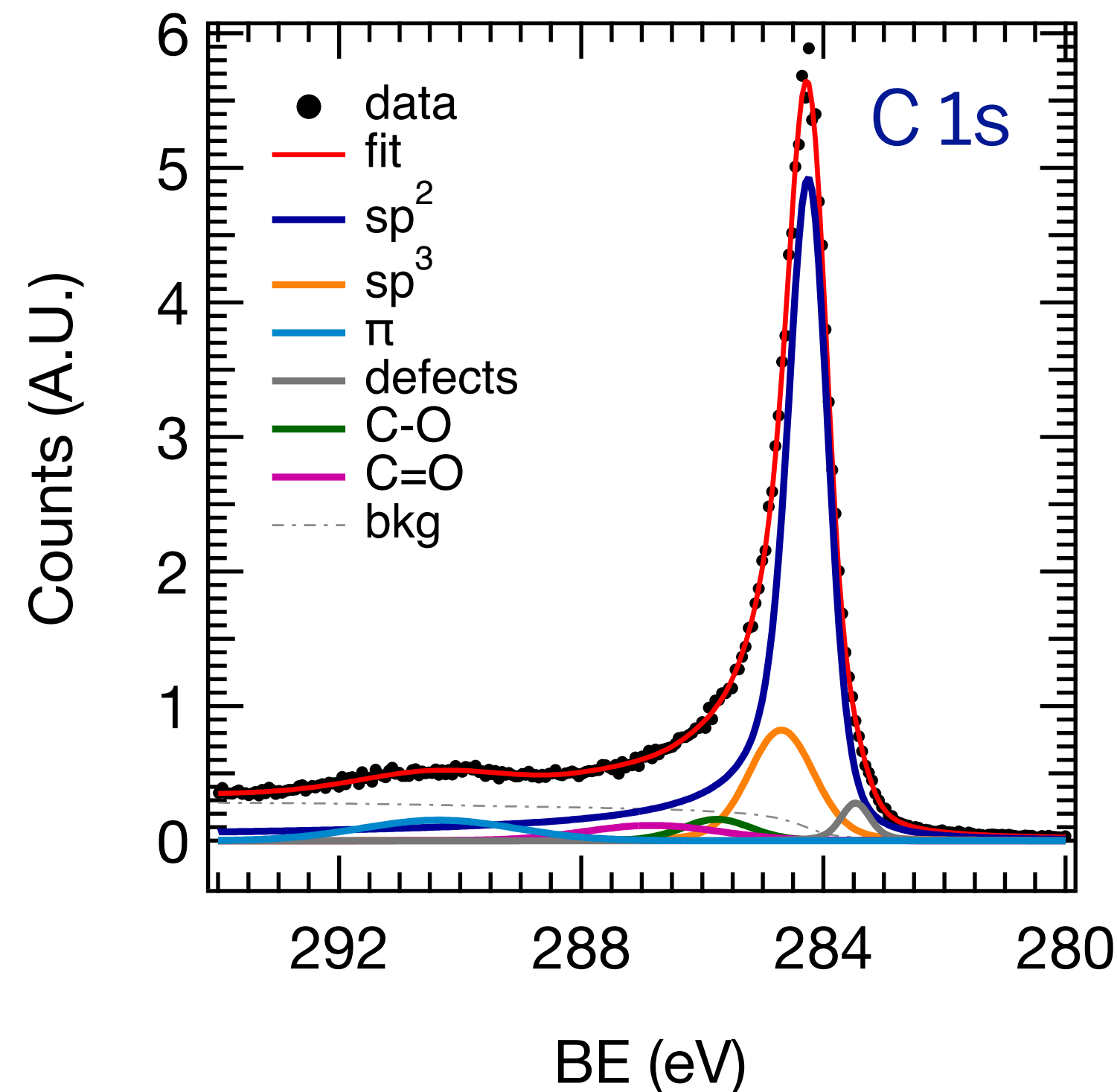
Annealing at $T < 300$ °C



XPS

Monochromatized Al K α = 1487.6 eV

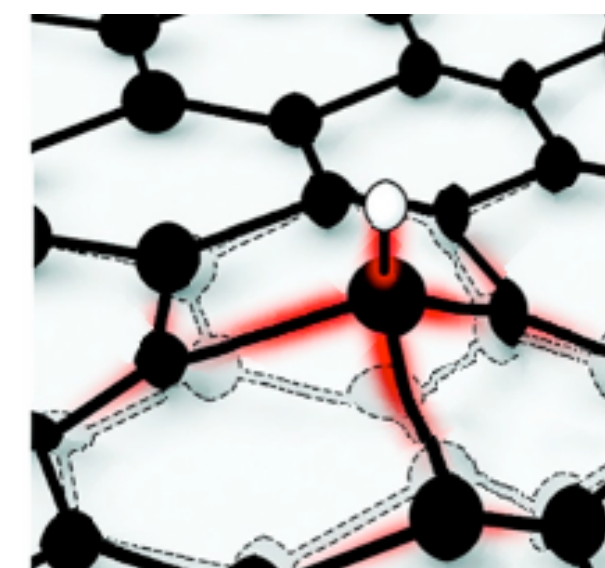
CNT pristine



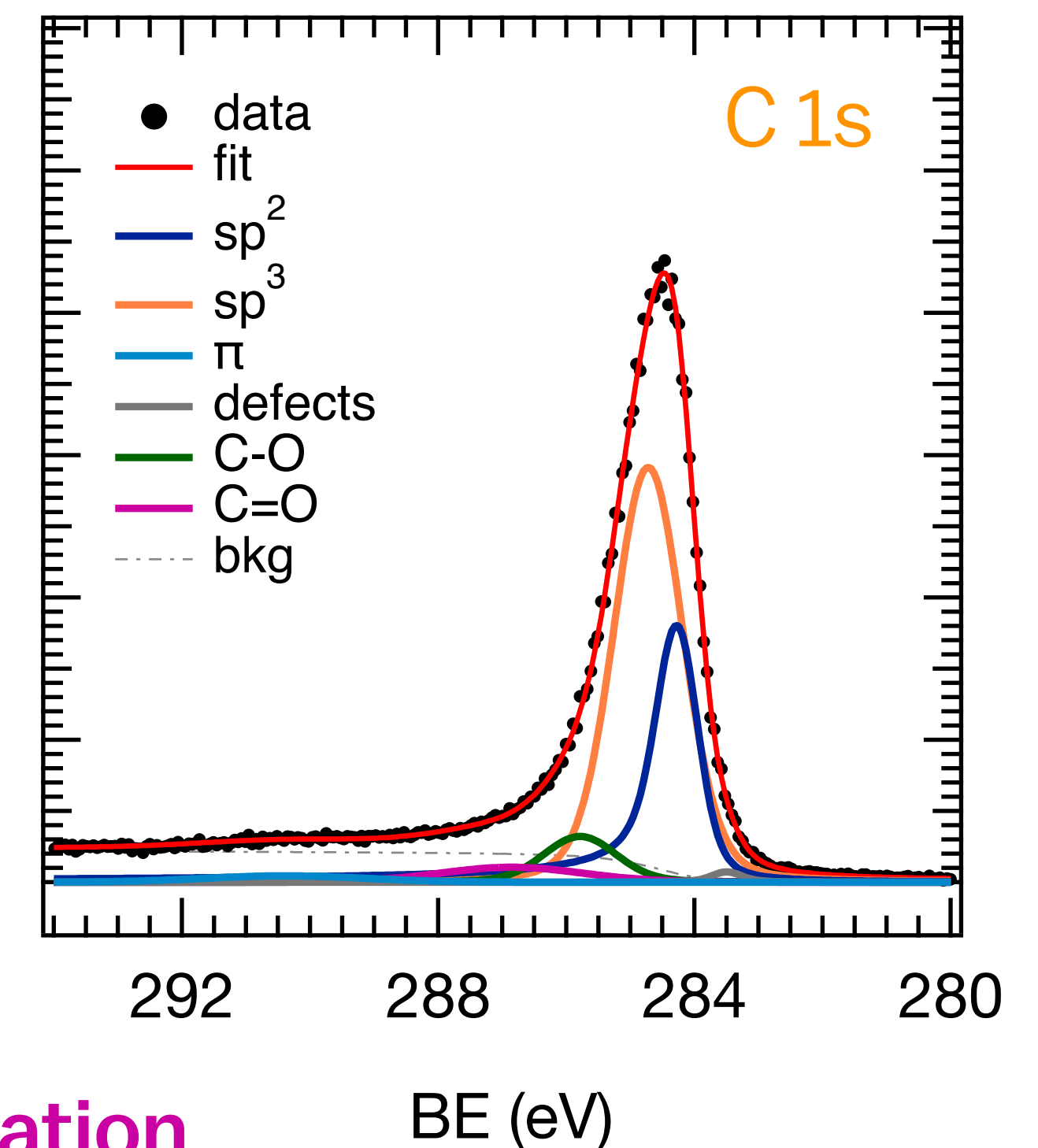
○ sp^2 -to- sp^3 carbon hybridization

$$\frac{I_{sp^3}}{I_{sp^3} + I_{sp^2}} = 0.67$$

○ Hydrogenation causes sp^2 to sp^3 transitions



CNT plasma H

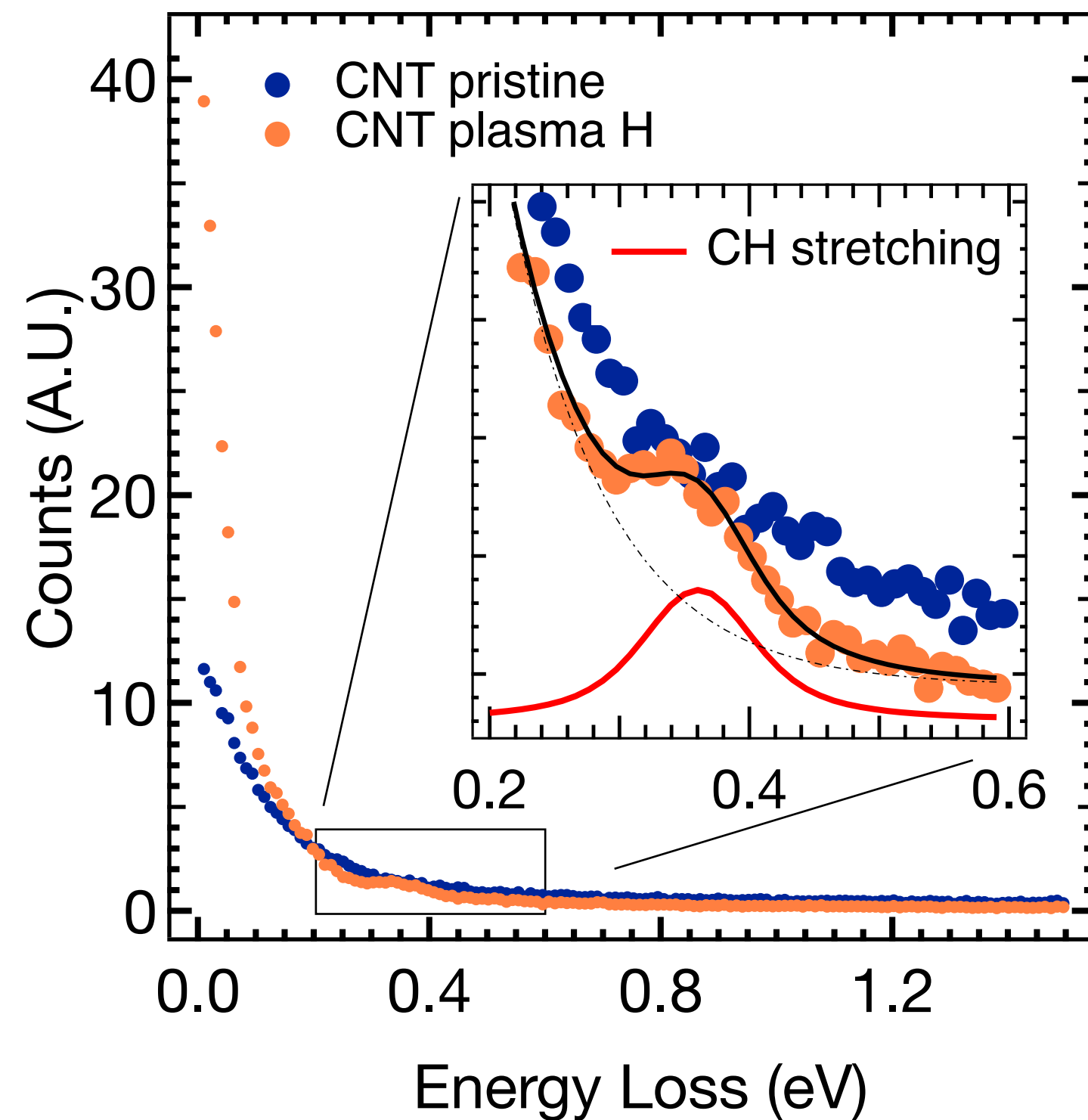


○ Negligible increase of oxides contamination

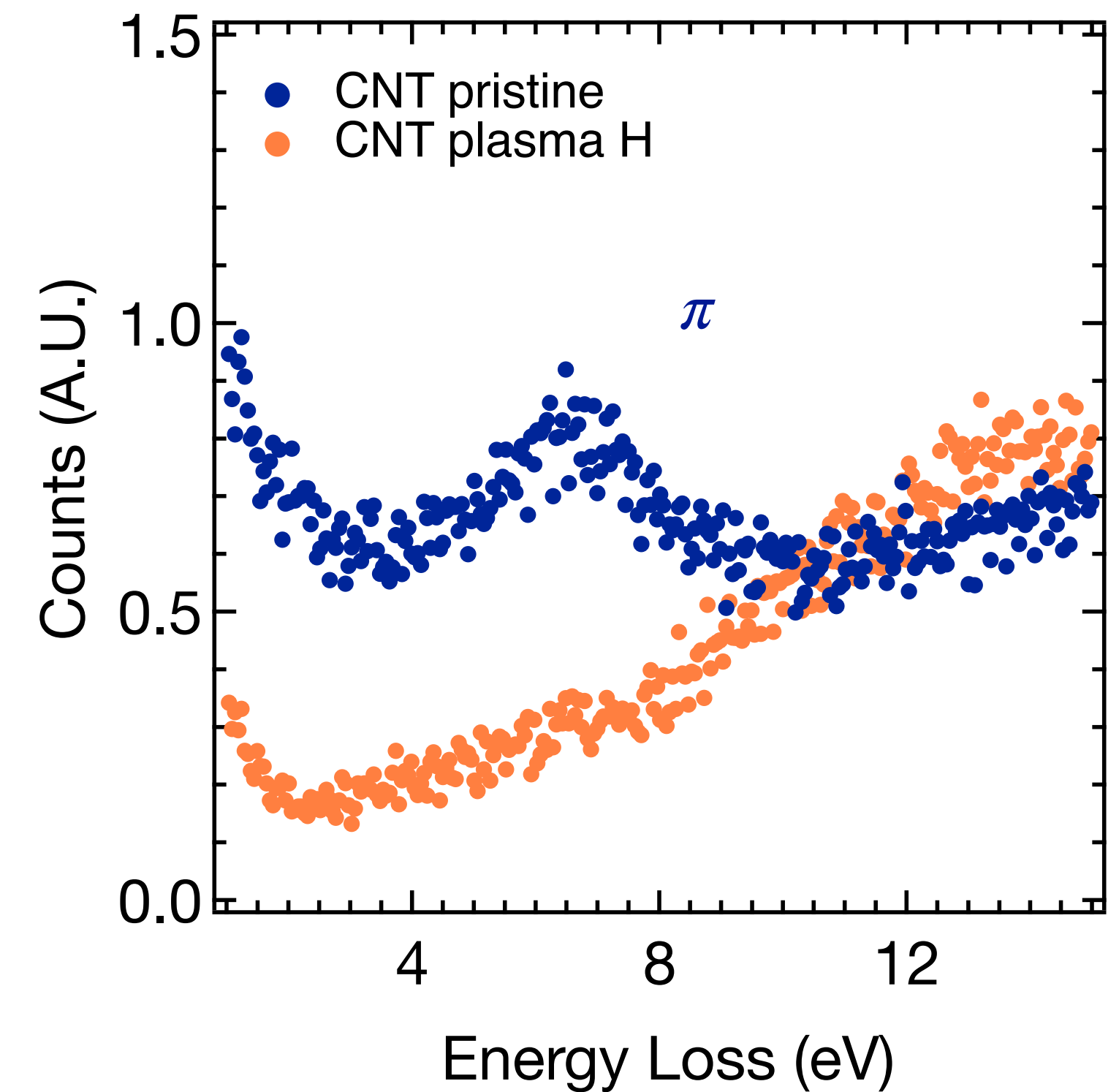
○ 0.5 eV lower Work function in H-CNT

EELS

$$E_{\text{primary}} = 90 \text{ eV}$$

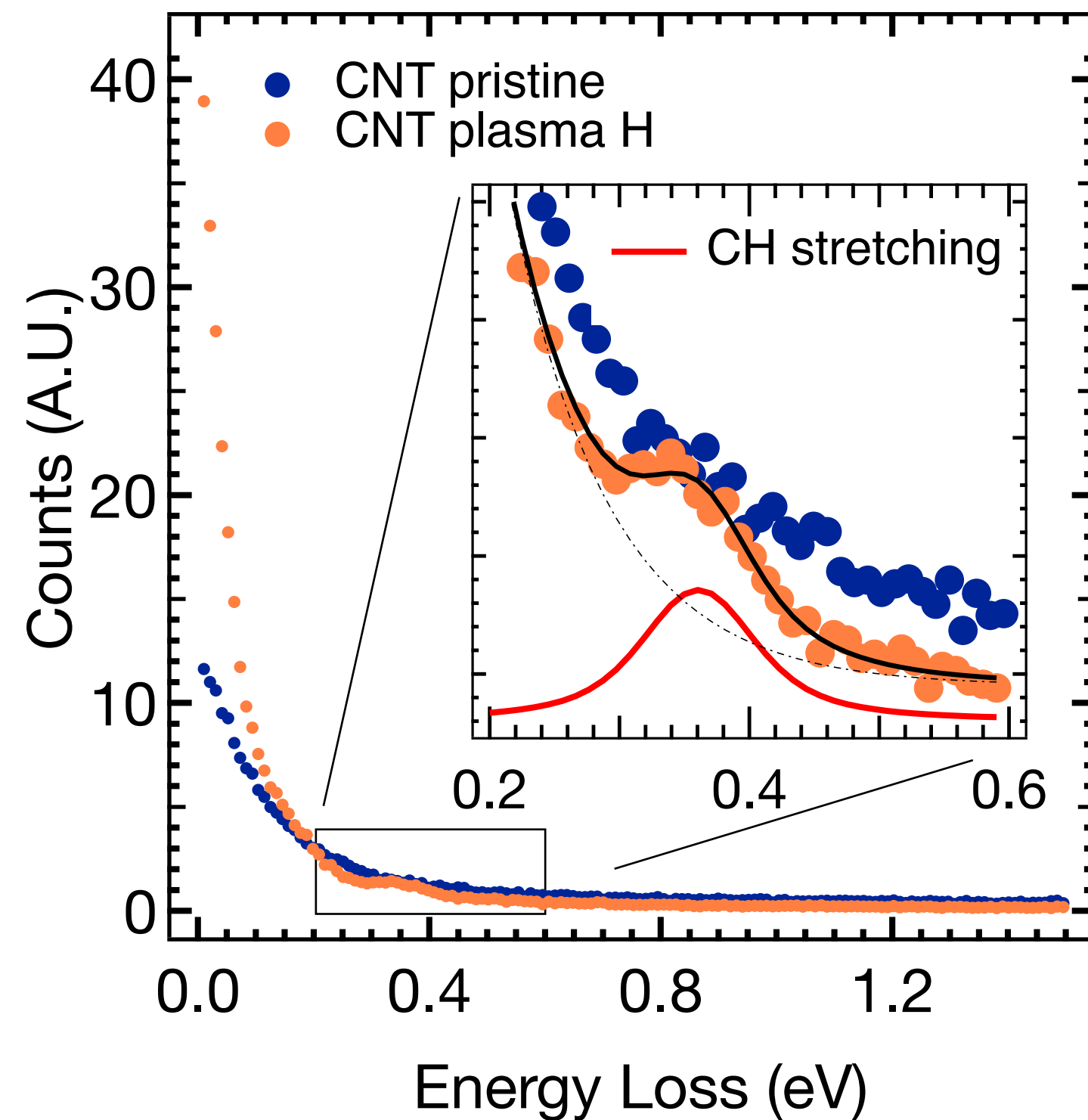


- C-H stretching
- π plasmon quenching
- Band Gap opening

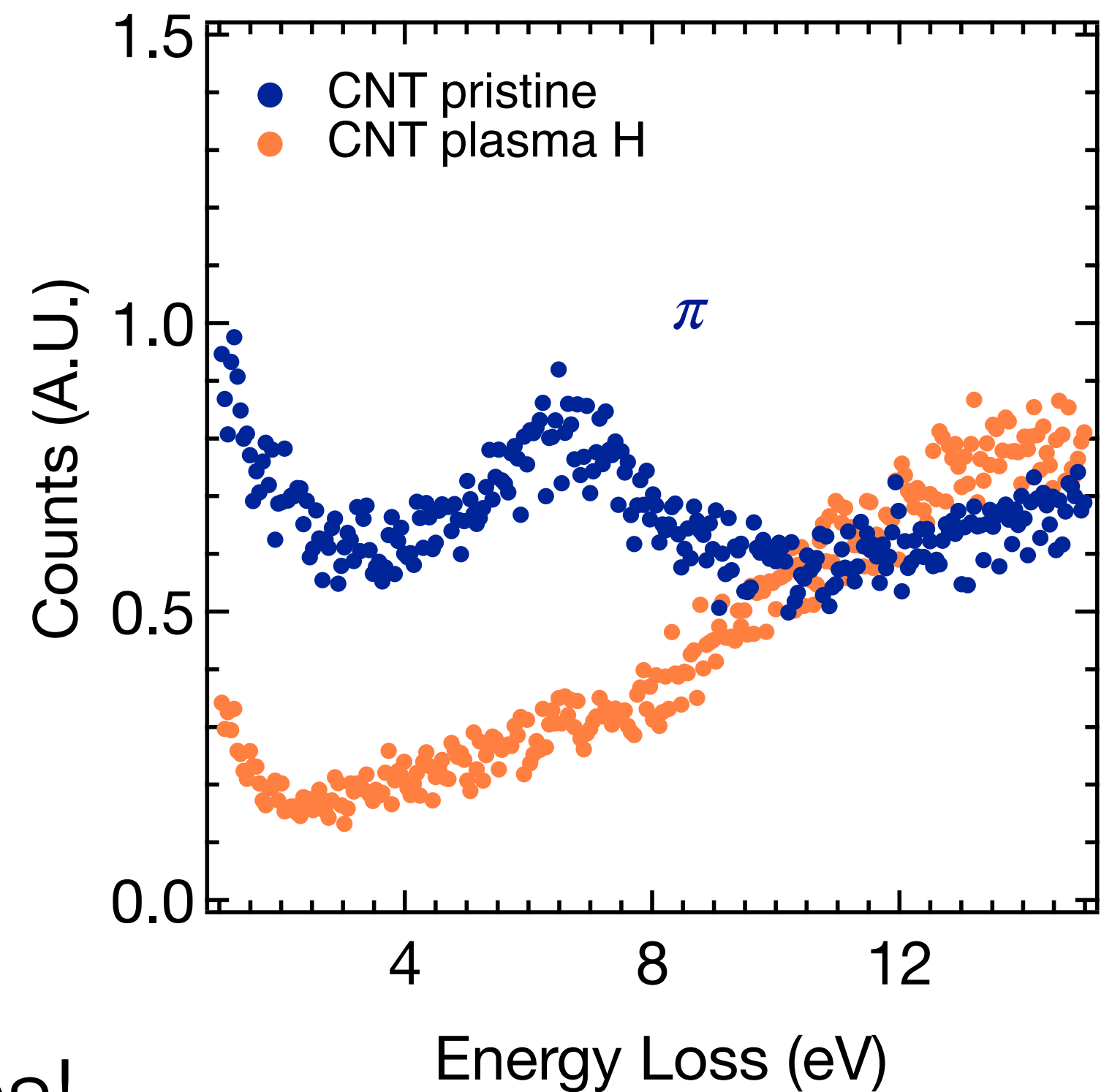


EELS

$$E_{\text{primary}} = 90 \text{ eV}$$

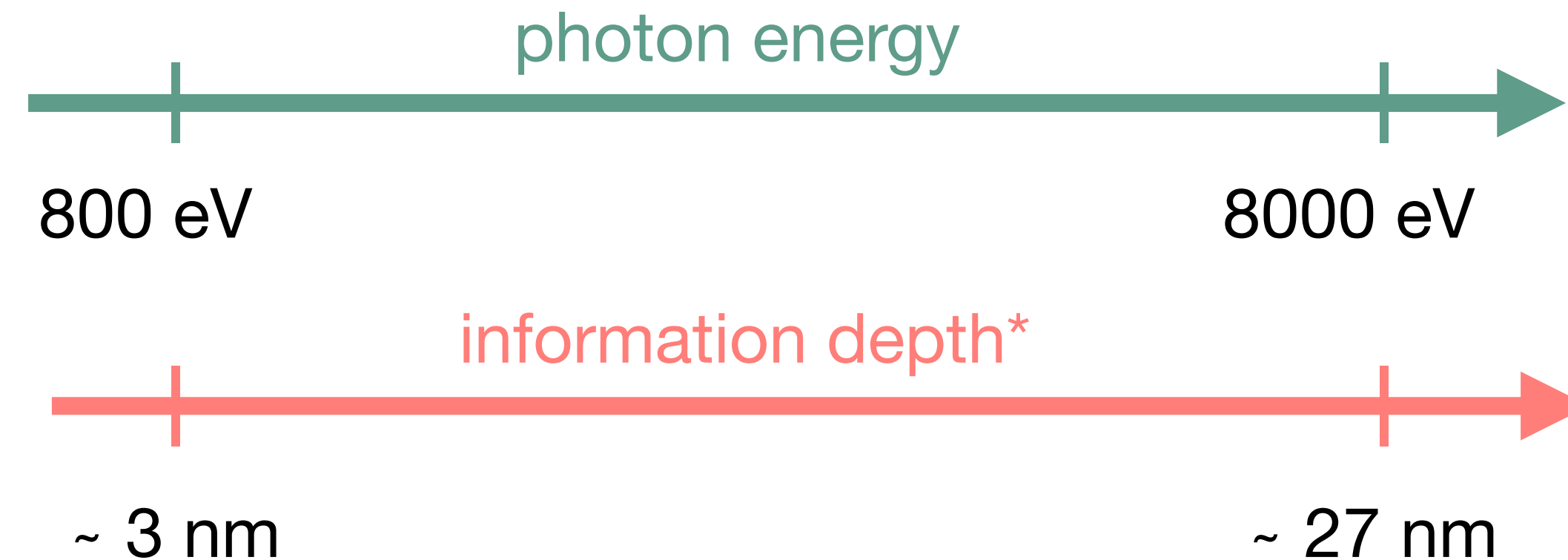


- C-H stretching
- π plasmon quenching
- Band Gap opening



Surface sensitive techniques!

- **Variable photon energy** allows probing from surface to inner layers
- Lineshape changes provide insight into **sp²** vs **sp³** carbon hybridization



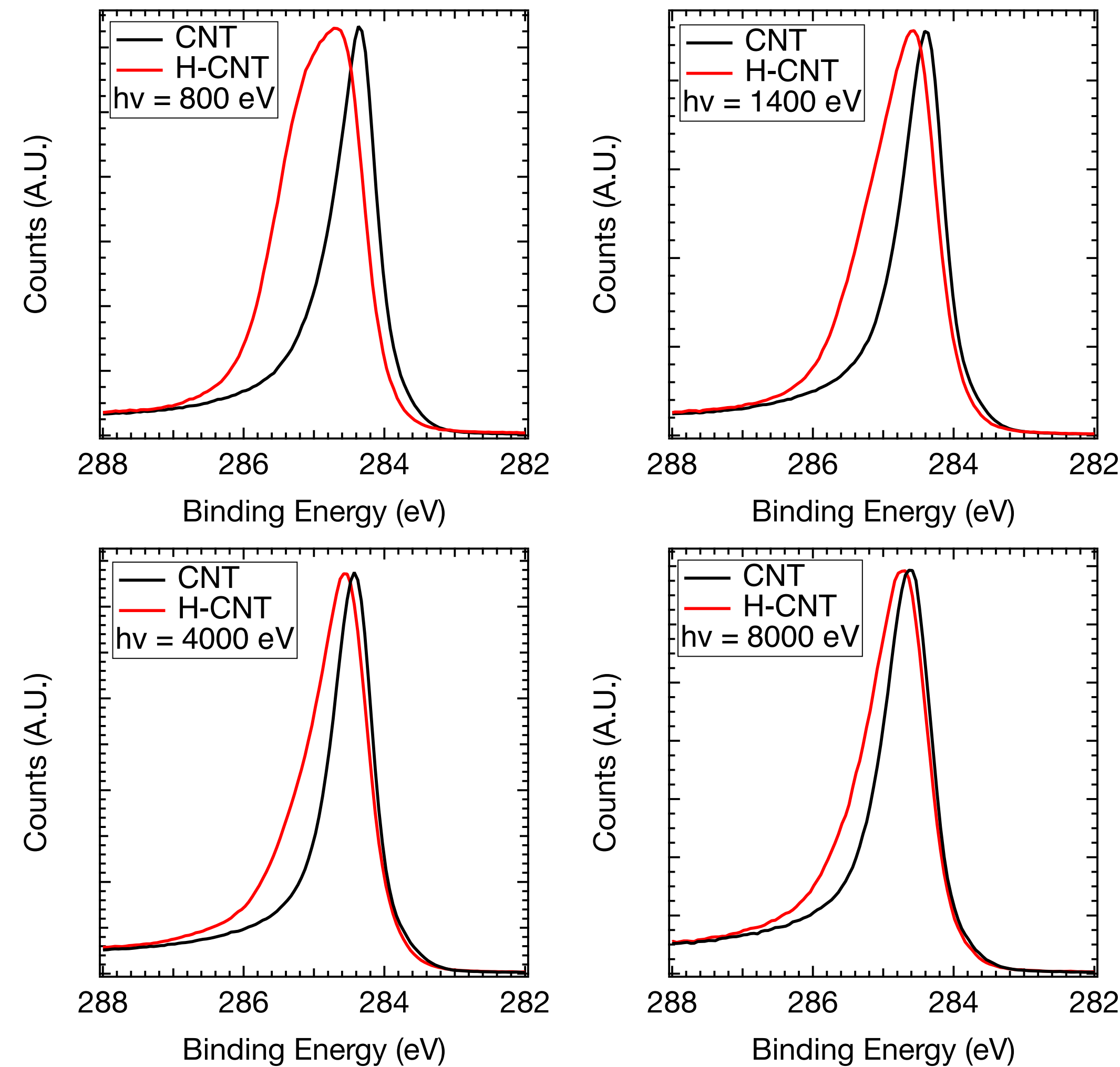
*3 x electron Inelastic Mean Free Path in graphite

C 1s

Inelastic Mean Free Path

h ν (eV)	$\lambda_{\text{graphite}}$ (nm)
800	1
1400	2
4000	4.7
8000	9.1

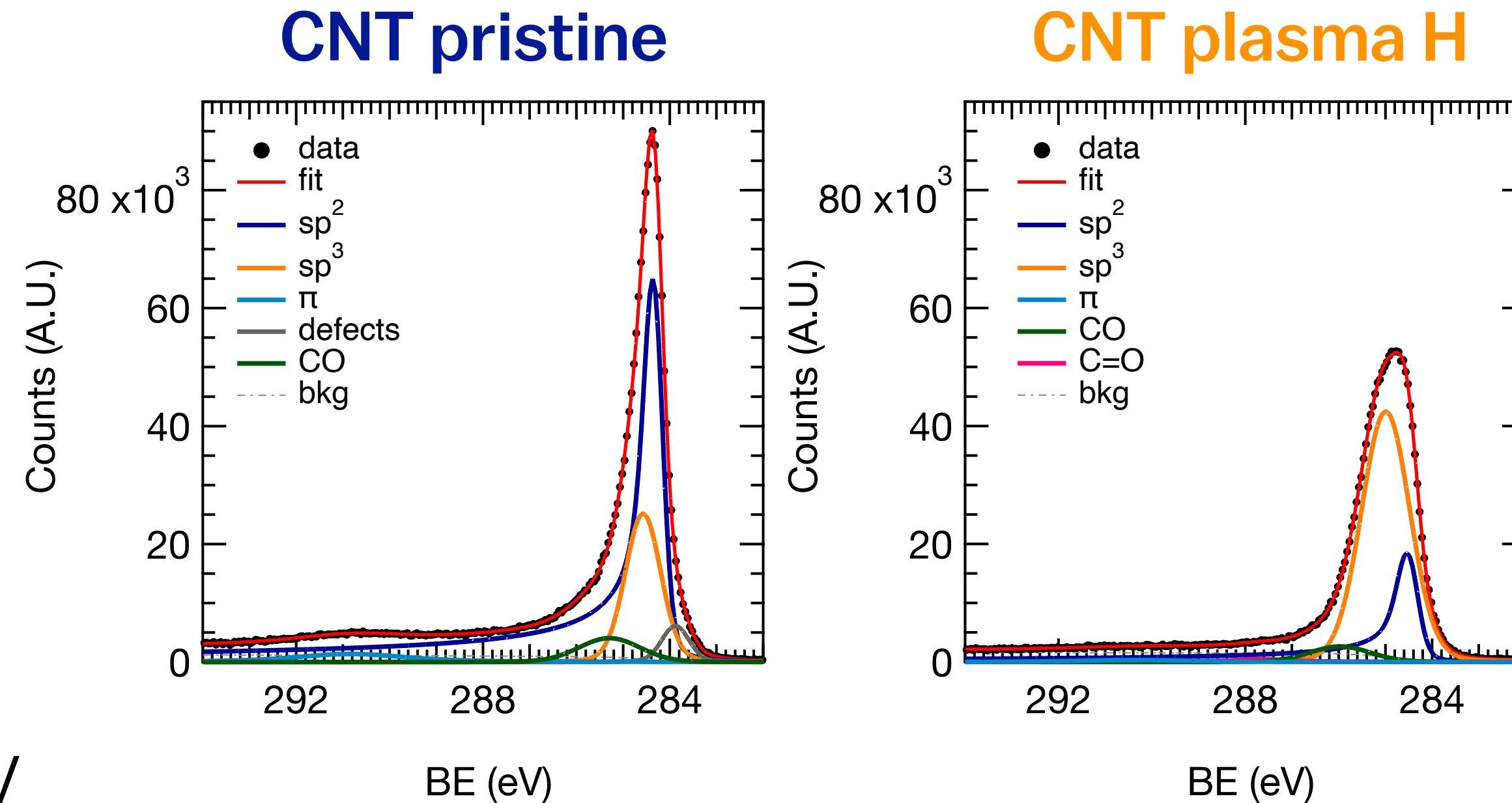
- Line shape differences between CNT and H-CNT **decrease with depth sensitivity**
- With **8 keV** photon **there is still a sign of hydrogenation**



Quantifying the Sp³ Content with Soft Photons

$h\nu = 800 \text{ eV}$

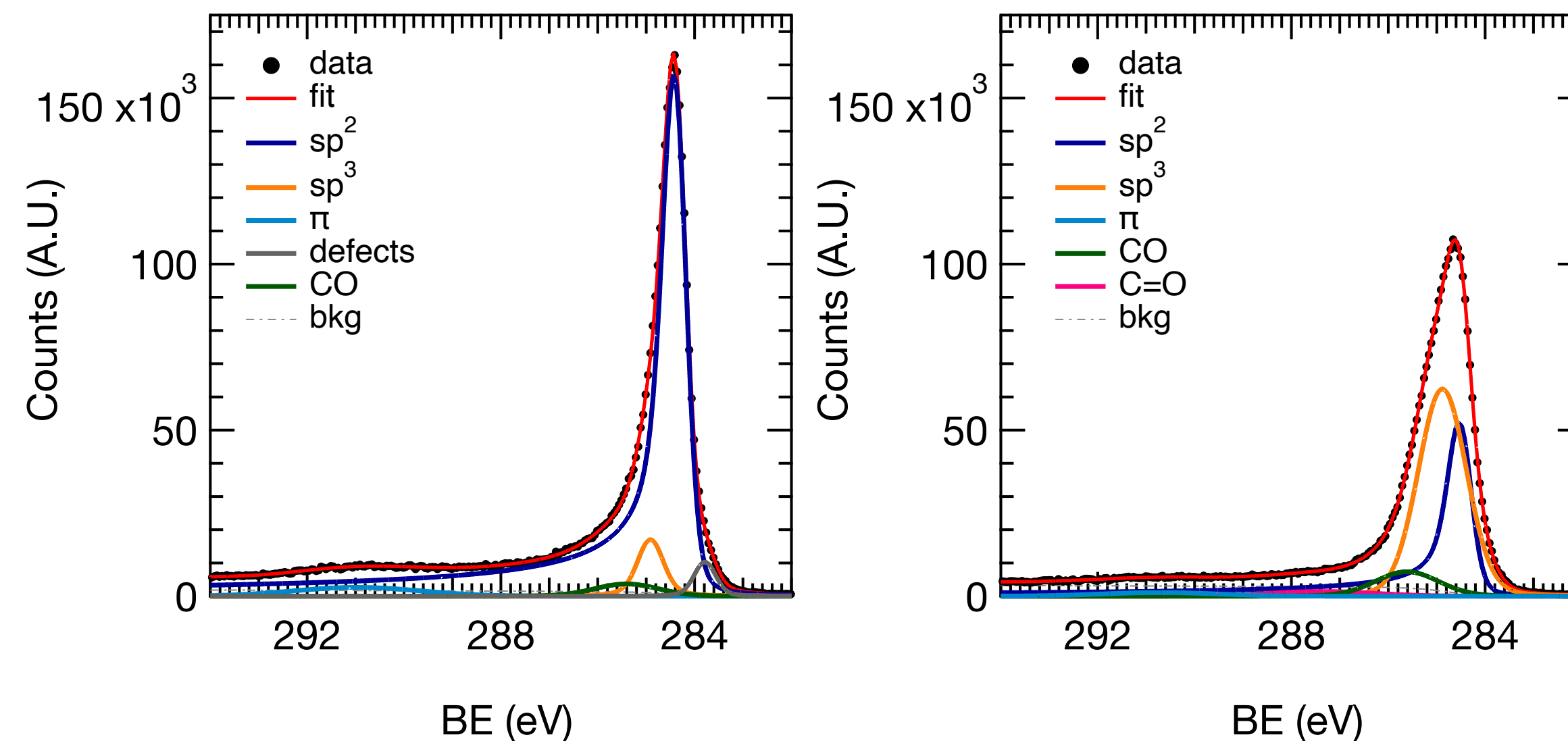
$\lambda = 1 \text{ nm}$



$$\frac{I_{sp^3}}{I_{sp^3} + I_{sp^2}} = 0.74$$

$h\nu = 1400 \text{ eV}$

$\lambda = 2 \text{ nm}$



$$\frac{I_{sp^3}}{I_{sp^3} + I_{sp^2}} = 0.61$$

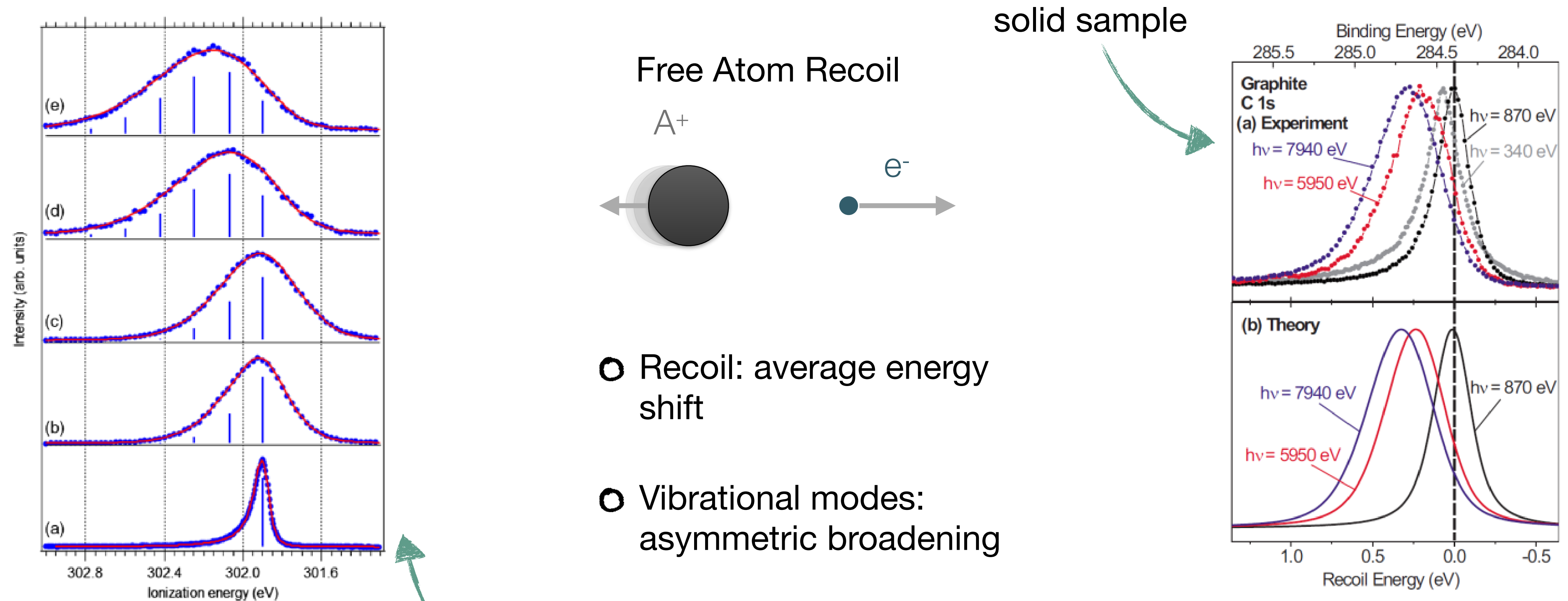
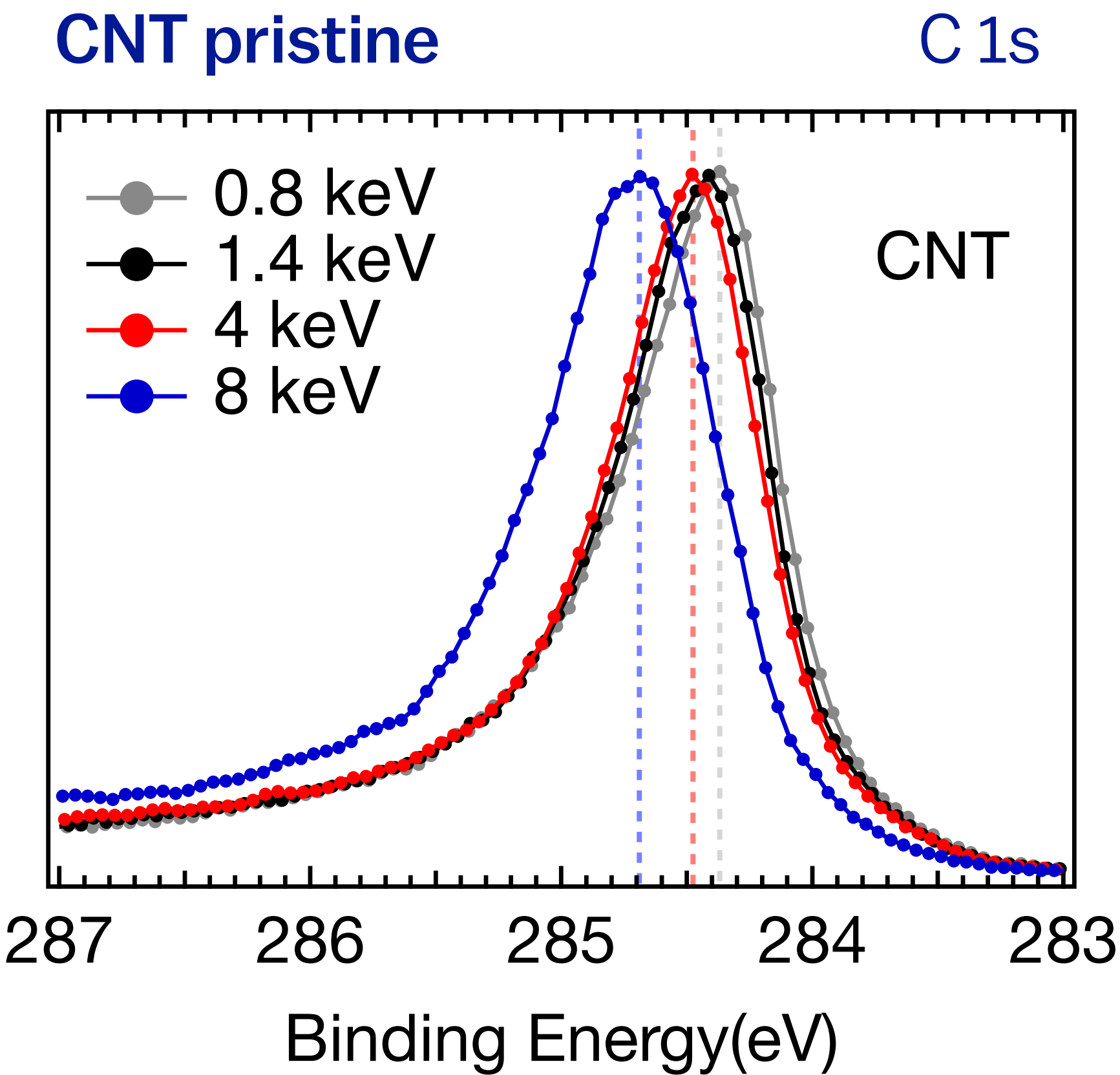


FIG. 1. Carbon 1s photoelectron spectra of CF_4 , measured at different photon energies: (a) 330 eV, (b) 2.3 keV, (c) 3.0 keV, (d) 6.9 keV, and (e) 8.5 keV. Dots, experimental data points; continuous red line, least-squares curve-fitting result; vertical sticks, positions and relative intensities of the vibrational peaks.

Kukk *et al.*, Phys. Rev. Lett. 121, 073002 (2018)

FIG. 1. (Color online) (a) Photon energy dependence of C 1s core-level spectra of graphite. The soft x ray ($h\nu=340$ and 870 eV) and hard x ray ($h\nu=5950$ and 7940 eV) are measured at the emission angles of 90° and 85° relative to the sample surface. (b) Theoretically obtained spectra taking into account the recoil effect in a Debye model with $\hbar\omega_{b,D}=75$ meV.

Takata *et al.*, Phys. Rev. B 75, 233404 (2007)



Measured energy shift wrt C1s at $h\nu = 800$ eV:

- $h\nu$ 4000 eV: ~ 100 meV
- $h\nu$ 8000 eV: ~ 300 meV

Free Atom Recoil:

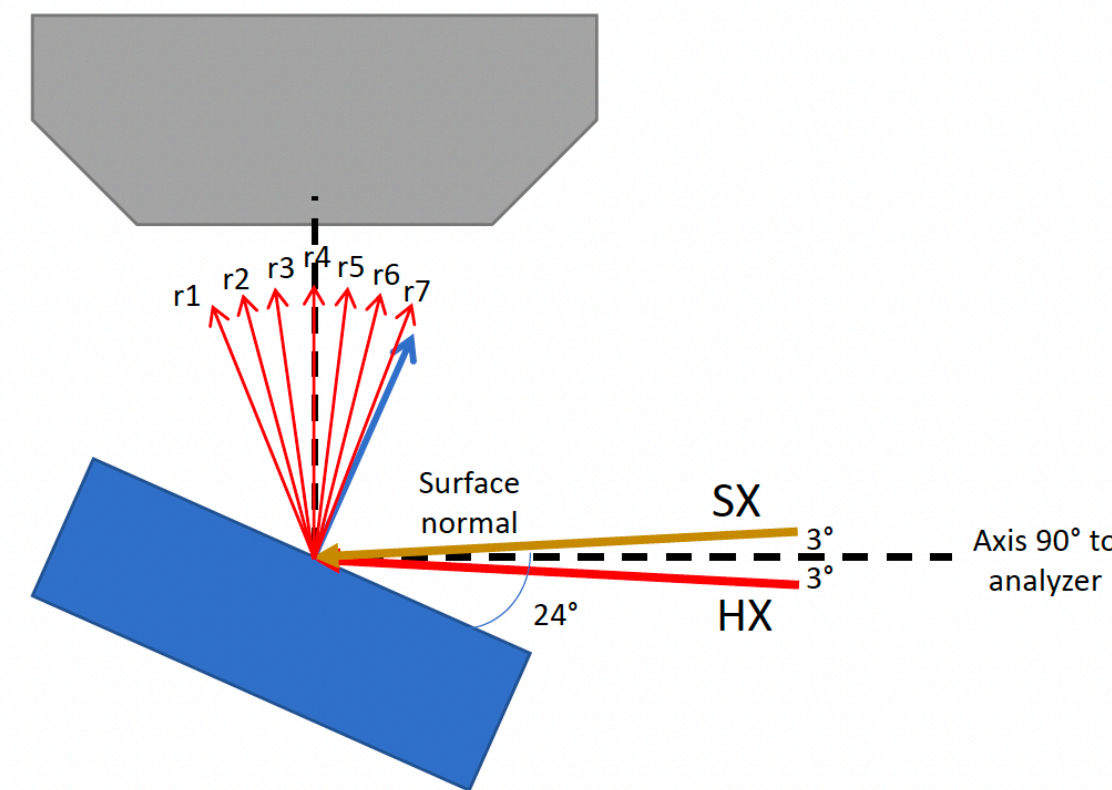
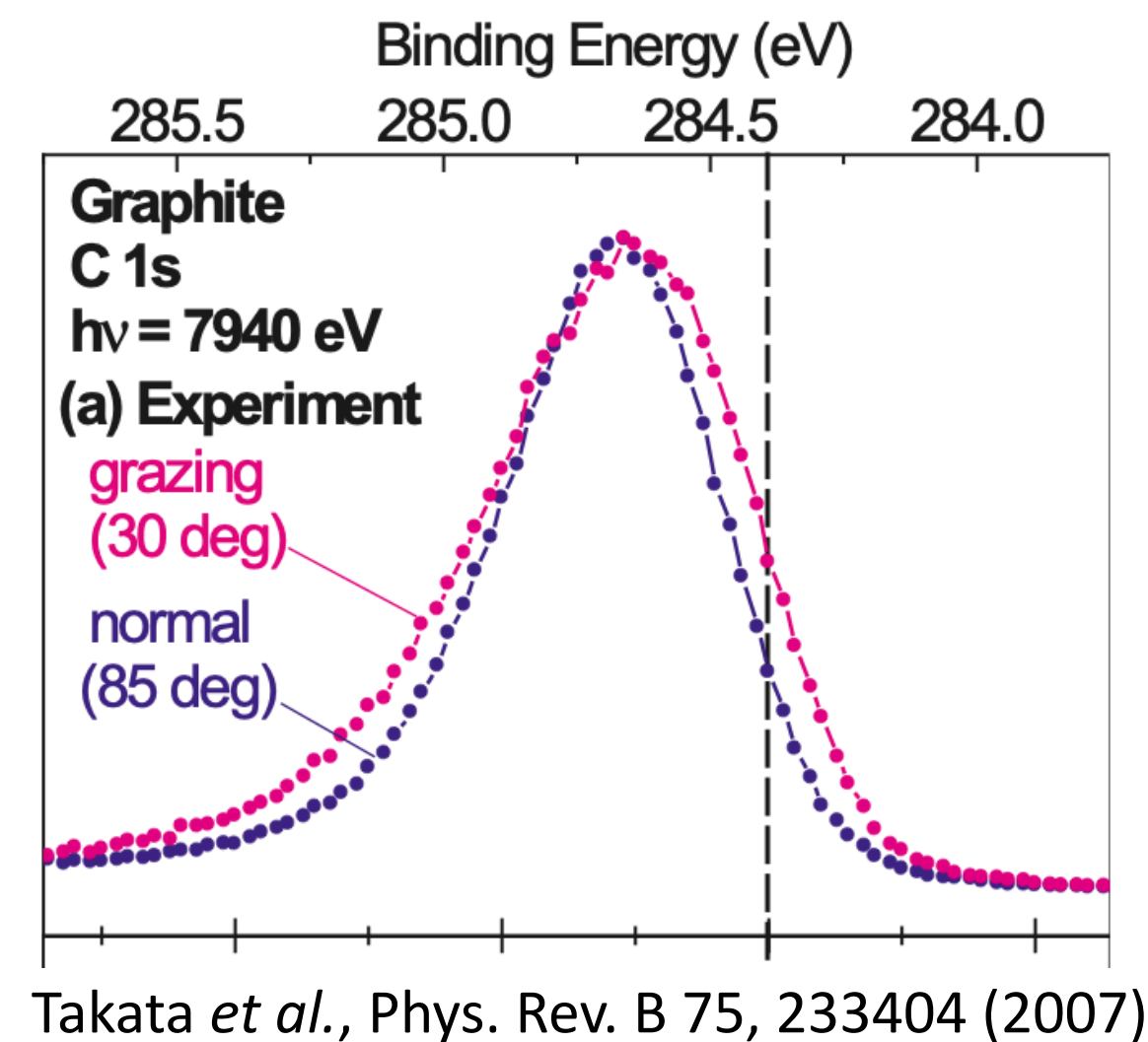
$$E_r = \frac{m_e}{M} E_k \quad \frac{m_e}{M_C} \sim 4.6 \times 10^{-5}$$

$h\nu$ 800 eV	$E_r \sim 0.02$ eV
$h\nu$ 4 KeV	$E_r \sim 0.17$ eV
$h\nu$ 8 KeV	$E_r \sim 0.37$ eV

Graphite

Grazing emission: in plane stretching mode, broader peak

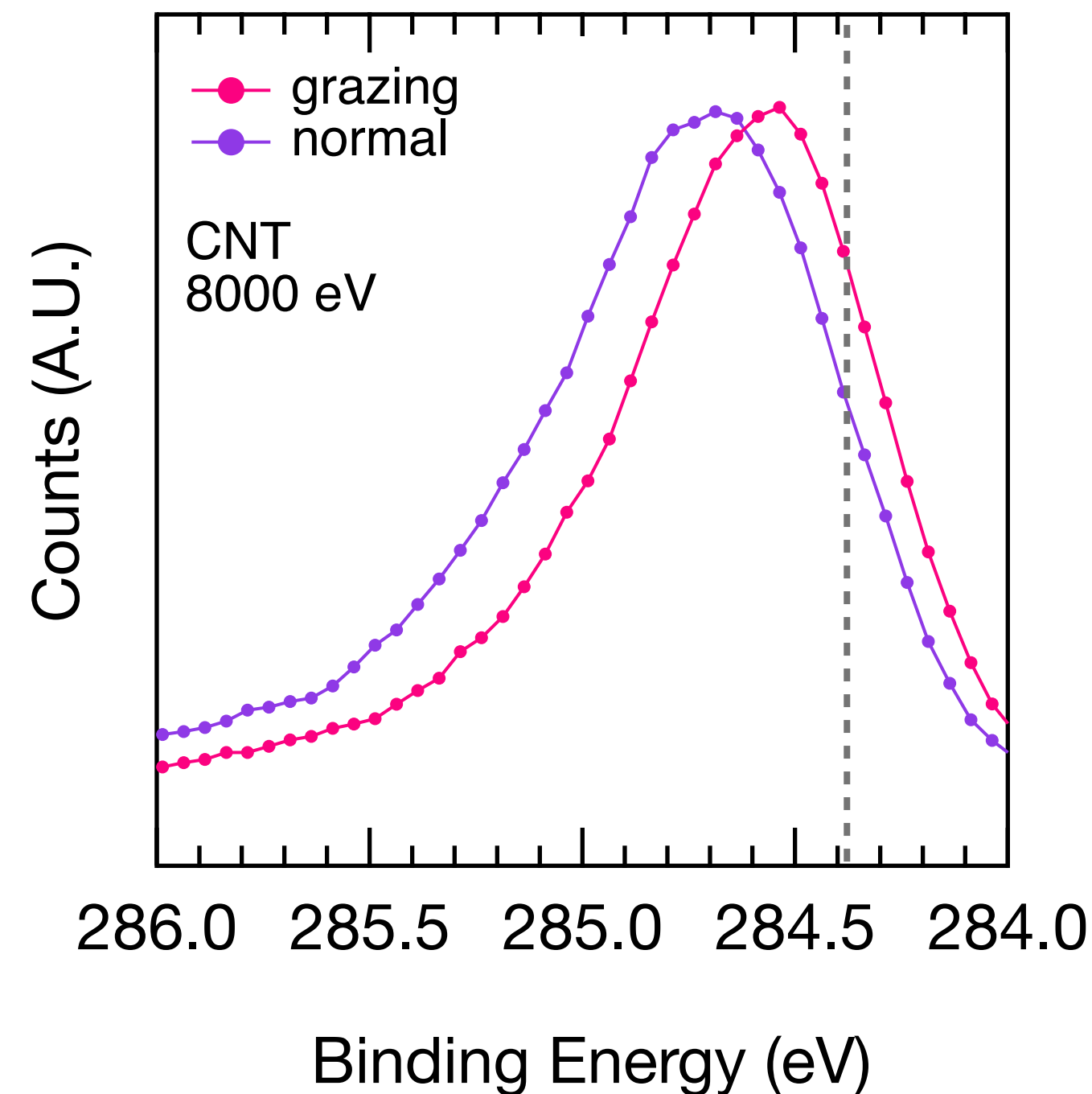
Normal emission: out-of-plane bending mode



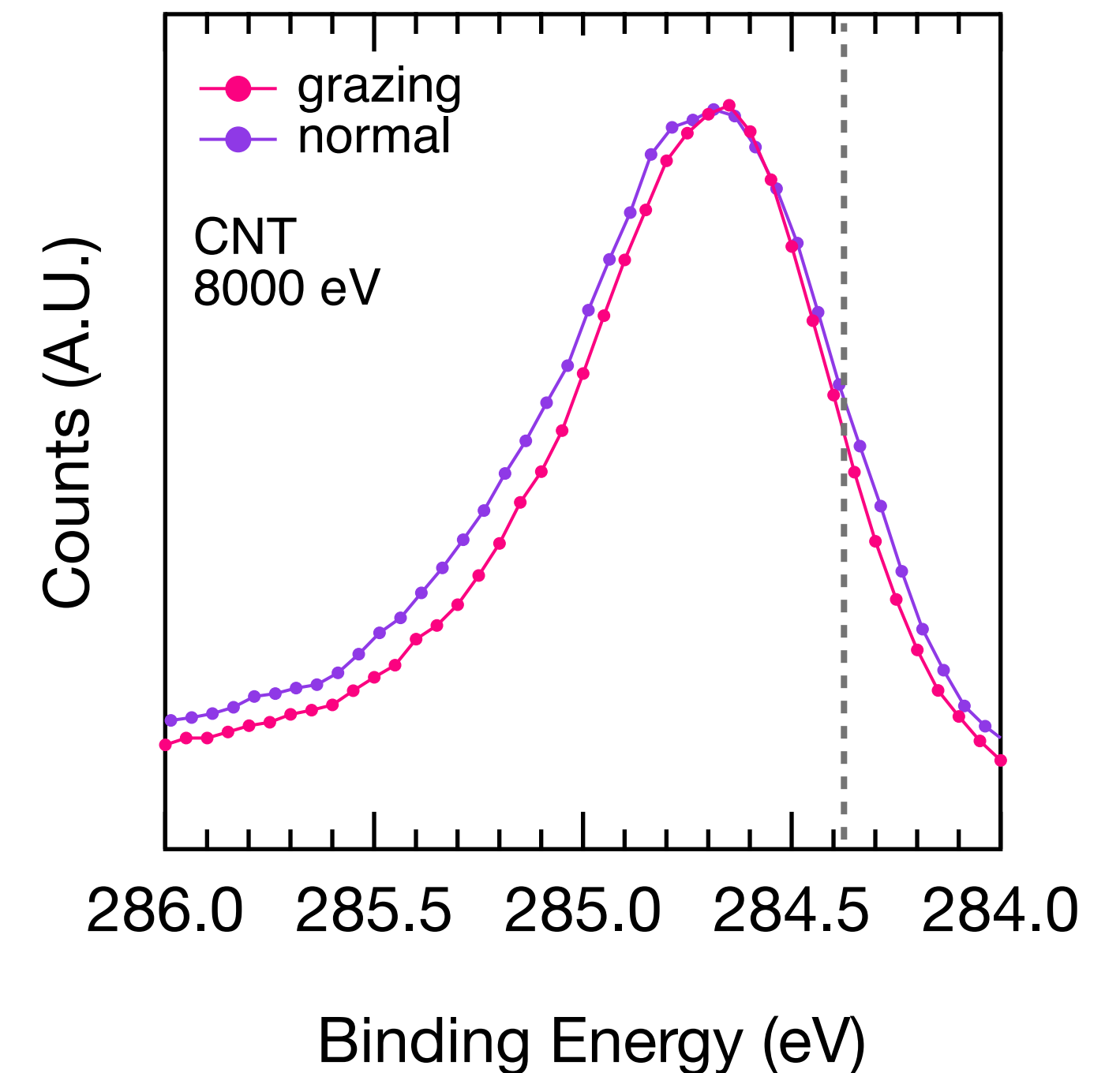
VA-CNT

○ **Normal emission:** broader peak

○ Average shift changes with emission angle



manually aligned



✓ **Highly hydrogenated CNTs** with minimal oxygen contamination

✓ **Hydrogenation** involves mainly the **outer layers**, but is **not limited to the very surface**

Can we quantify hydrogen distribution in depth?

✓ **Recoil effects observed in pristine CNTs**, average energy shift consistent with graphite

How can we explain the angular dependence?

What about H-CNTs? Can we study C-H vibrations?

✓ **Highly hydrogenated CNTs** with minimal oxygen contamination

✓ **Hydrogenation** involves mainly the **outer layers**, but is **not limited to the very surface**

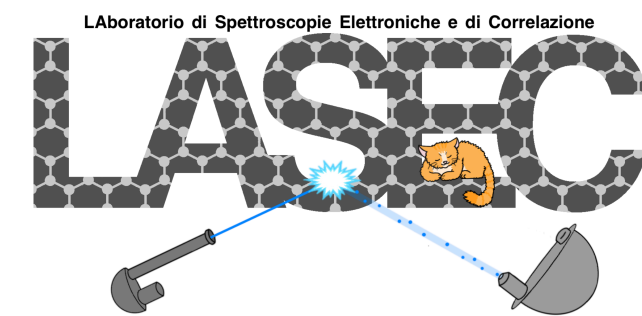
Can we quantify hydrogen distribution in depth?

✓ **Recoil effects observed in pristine CNTs**, average energy shift consistent with graphite

How can we explain the angular dependence?

What about H-CNTs? Can we study C-H vibrations?

Work in progress!



Orlando Castellano, Daniele Paoloni,
Francesco Offi, Alessandro Ruocco



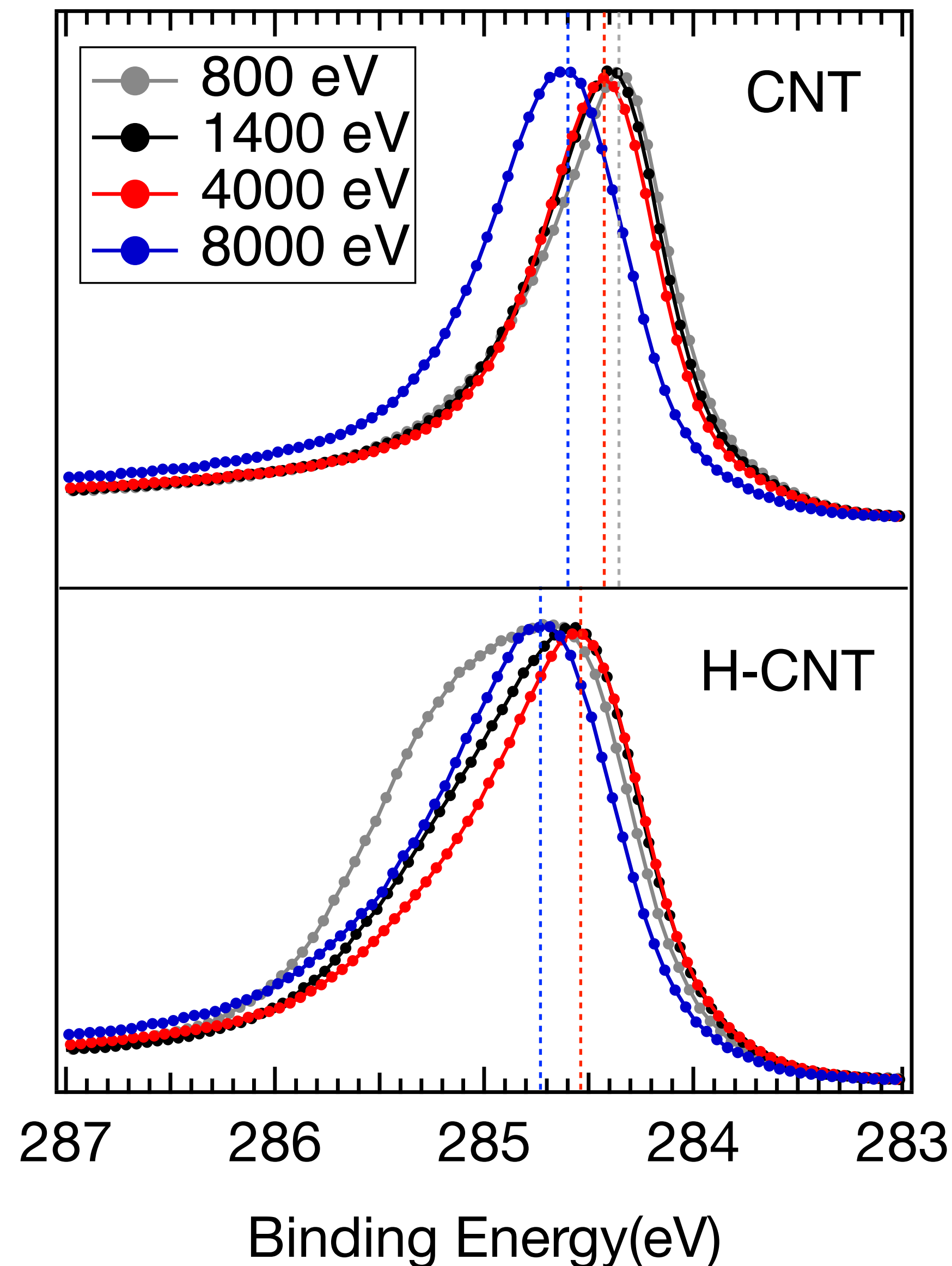
Alice Apponi, Luca Cecchini,
Francesco Pandolfi, Ilaria Rago



Tien-Lin Lee, Pardeep Kumar Thakur



Carlo Mariani, Gianluca Cavoto



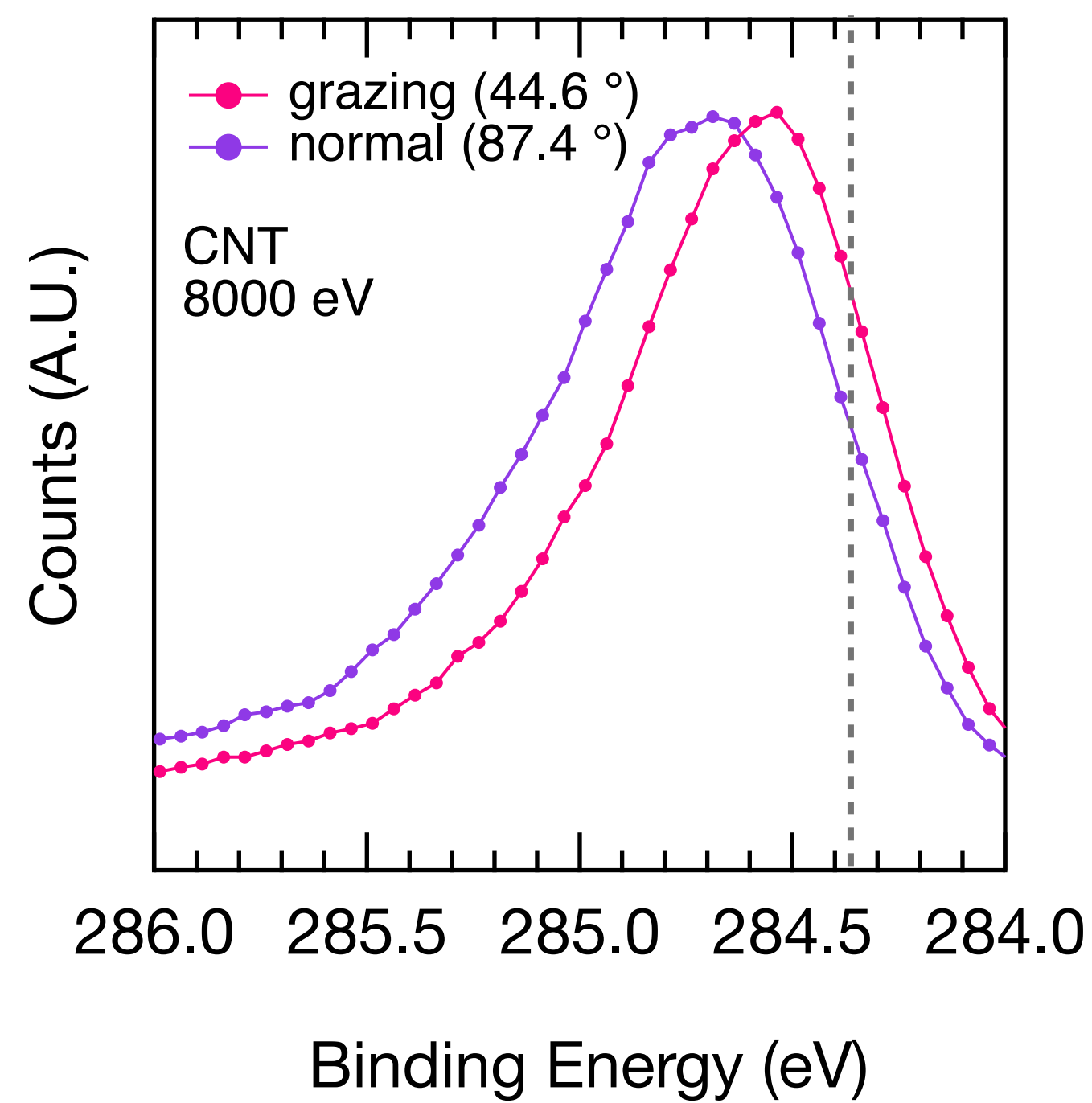
- 800 eV ~ 1400 eV no major C-C recoil effect
- 4000 eV energy shift ~ 100 meV
- 8000 eV energy shift ~ 250 meV

- Mix of depth inhomogeneous and recoil effect hard to disentangle

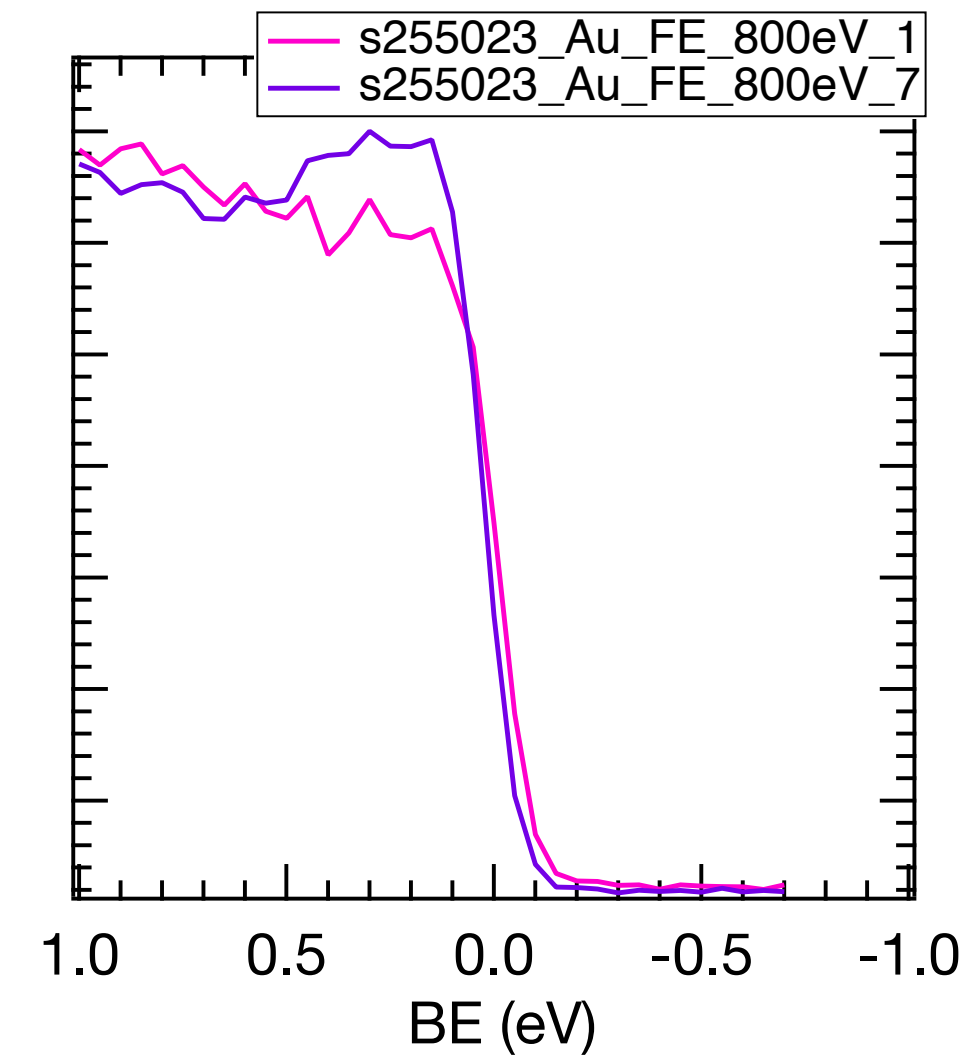
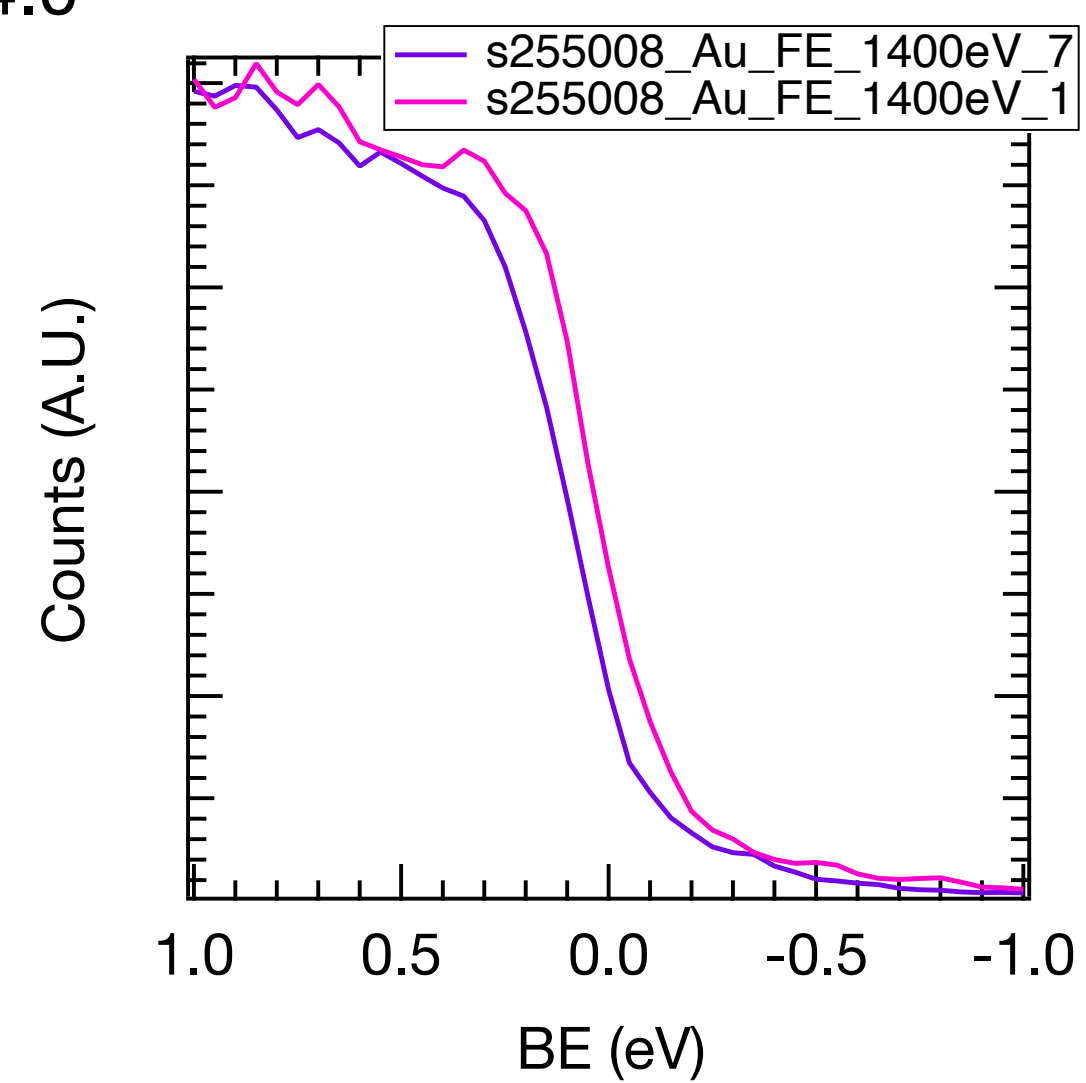
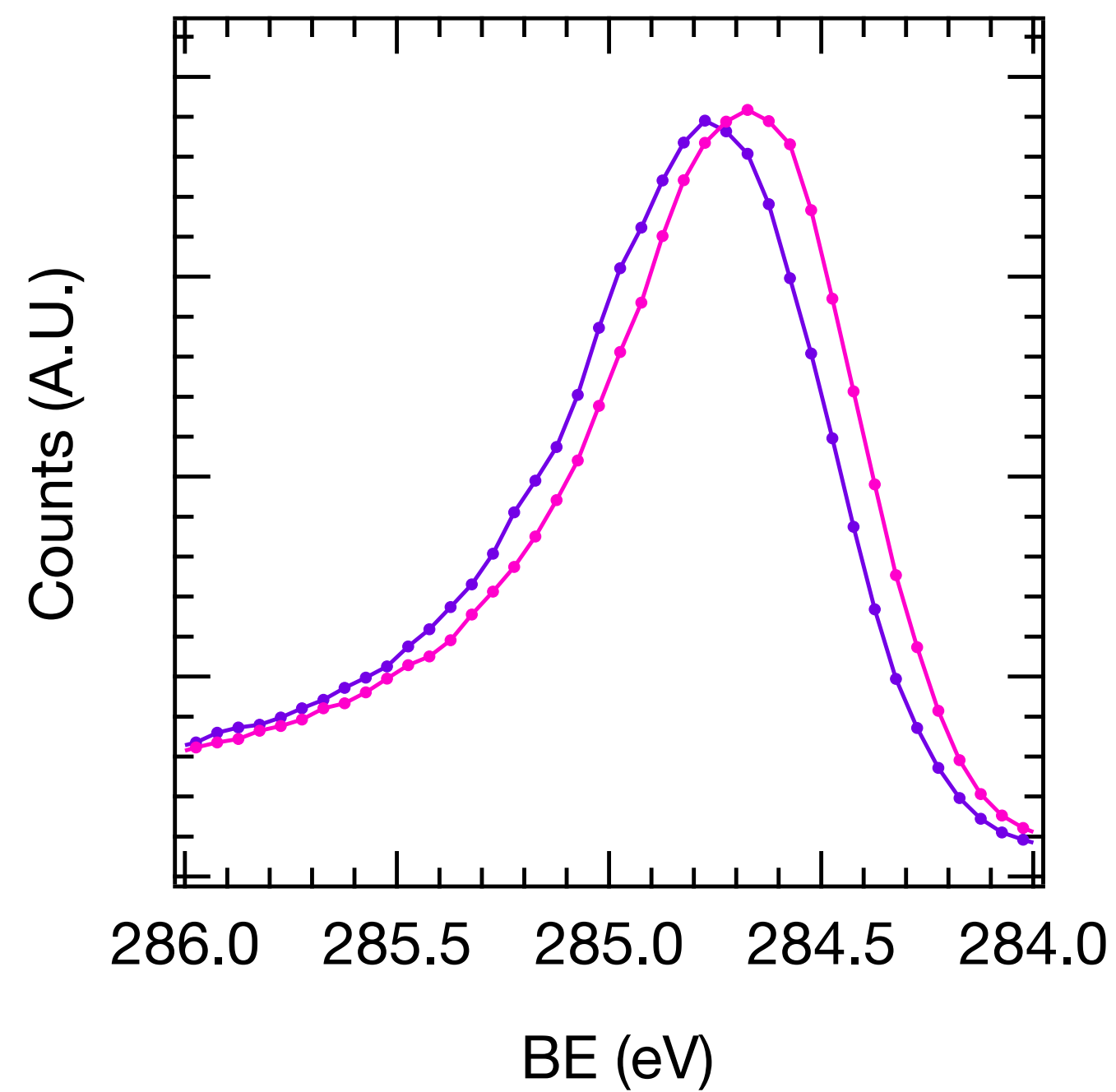
Sp3 lower at higher information depth, effective sharper line shape and shift to lower BE

Recoil increase with photon energy, peak shift to lower KE, higher BE

VA-CNT



NPG



(c) Semi-infinite substrate with uniform overlayer of thickness t —
Peak k from substrate with $E_{kln} \equiv E_k$:

$$N_k(\theta) = I_0 \Omega_0(E_k) A_0(E_k) D_0(E_k) \rho \, d\sigma_k/d\Omega \, \Lambda_e(E_k) \times \exp(-t/\Lambda_e'(E_k) \sin \theta) \quad (117)$$

Peak l from overlayer with $E_{kln} \equiv E_l$:

$$N_l(\theta) = I_0 \Omega_0(E_l) A_0(E_l) D_0(E_l) \rho' \, d\sigma_l/d\Omega \, \Lambda_e'(E_l) \times [1 - \exp(-t/\Lambda_e'(E_l) \sin \theta)] \quad (118)$$

where

$\Lambda_e(E_k)$ = an attenuation length in the substrate

$\Lambda_e'(E_k)$ = an attenuation length in the overlayer

ρ = an atomic density in the substrate

ρ' = an atomic density in the overlayer.

$$\frac{I_{sp3}}{I_{sp3} + I_{sp2}} = \frac{I_{\text{overlayer}}}{I_{\text{overlayer}} + I_{\text{substrate}}} = \frac{1 - e^{-d/\lambda \cos \theta}}{1 - e^{-d/\lambda \cos \theta} + e^{-d/\lambda \cos \theta}} = 1 - e^{-d/\lambda}$$

$$h\nu = 800 \text{ eV}$$

$$\lambda = 1\text{nm}, \frac{sp3}{sp3 + sp2} = 0.74$$

$$1 - e^{-d/\lambda} = 0.74 \Rightarrow d = -\ln(0.26) = 1.35 \text{ nm}$$

$$h\nu = 1400 \text{ eV}$$

$$\lambda = 2\text{nm}, \frac{sp3}{sp3 + sp2} = 0.61$$

$$1 - e^{-d/\lambda} = 0.61 \Rightarrow d = -2\ln(0.39) = 1.88 \text{ nm}$$

$$((1.88+1.35)/2=1.615$$

$$2*(1.88-1.35)/(1.88+1.35)=0.328$$

$$f_{sp3}(z) = f_0 \cdot e^{-z/d_0}$$

sp³ is 100% at the surface

d₀ characteristic sp³ length

$$f_{sp2}(z) = 1 - f_{sp3}(z)$$

$$I_{sp3} \propto \int_0^\infty f_{sp3}(z) \cdot e^{-z/\lambda} dz = \int_0^\infty e^{-z/d_0} \cdot e^{-z/\lambda} dz = \int_0^\infty e^{-z(1/d_0 + 1/\lambda)} dz I_{sp3} \propto \frac{1}{(1/d_0 + 1/\lambda)} = \frac{d_0 \lambda}{d_0 + \lambda}$$

$$I_{sp2} \propto \int_0^\infty [1 - f_{sp3}(z)] \cdot e^{-z/\lambda} dz = \int_0^\infty (1 - e^{-z/d_0}) \cdot e^{-z/\lambda} dz = \lambda \left(1 - \frac{d_0}{d_0 + \lambda} \right) = \lambda \left(\frac{\lambda}{d_0 + \lambda} \right) = \frac{\lambda^2}{d_0 + \lambda}$$

$$f_{sp^3}(z) = f_0 \cdot e^{-z/d_0}$$

sp³ is 100% at the surface

d₀ characteristic sp³ length

$$f_{sp^2}(z) = 1 - f_{sp^3}(z)$$

$$\frac{I_{sp^3}}{I_{sp^3} + I_{sp^2}} = \frac{d_0\lambda/(d_0 + \lambda)}{d_0\lambda/(d_0 + \lambda) + \lambda^2/(d_0 + \lambda)} = \frac{d_0}{d_0 + \lambda}$$

hν = 800 eV

hν = 1400 eV

$$\lambda = 1\text{ nm}, \frac{sp^3}{sp^3 + sp^2} = 0.74$$

$$\lambda = 2\text{ nm}, \frac{sp^3}{sp^3 + sp^2} = 0.61$$

$$d_0 = \frac{0.74}{1 - 0.74} = 2.85\text{ nm}$$

$$d_0 = \frac{0.61}{1 - 0.61} = 3.13\text{ nm}$$

$$(3.13 + 2.85)/2 = 2.99$$

$$2 \cdot (3.13 - 2.85)/(3.13 + 2.85) = 0.0936$$

Layer separation in graphite = 0.35 nm

$$d = 1.61 \text{ (uniform model)}$$

$$1.61/0.35 = 4.6$$

$$d_0 = 2.99 \text{ (exponential model)}$$

$$2.99/0.35 = 8.5$$

$$n_C \approx 3.82 \times 10^{13} \text{ atomi/mm}^2$$

$$N_H = 4.6 \cdot n_C = 4.6 \cdot 3.82 \times 10^{13} = 1.7572 \times 10^{14} \text{ atomi H/mm}^2$$

$$m_H = 1.0078 \text{ u} = 1.0078 \cdot 1.6605 \times 10^{-24} \text{ g} \approx 1.6736 \times 10^{-24} \text{ g}$$

$$m_H \approx 29 \text{ ng/cm}^2$$

$$m_H = 1 \text{ ug} \approx 60$$



$$n_0 = \frac{3.82 \times 10^{13}}{0.335} \approx 1.14 \times 10^{14} \text{ atomi/mm}^2/\text{nm}$$

$$N_H = \int_0^\infty n_H(z) dz = \int_0^\infty n_0 \cdot e^{-z/d_0} dz = n_0 \cdot d_0 = 1.14 \times 10^{14} \cdot 3 = 3.42 \times 10^{14} \text{ atomi/mm}^2$$

$$m_H \approx 57 \text{ ng/cm}^2$$

1 ug Hydrogen (exponential model)

27

Layer separation as in graphite = 0.35 nm

Electron IMFP as in graphite

Characteristic penetration depth: 8.5 layers

$$m_H \approx 57 \text{ ng/cm}^2$$

$$\frac{sp3}{sp3 + sp2} = \text{H:C stoichiometry}$$

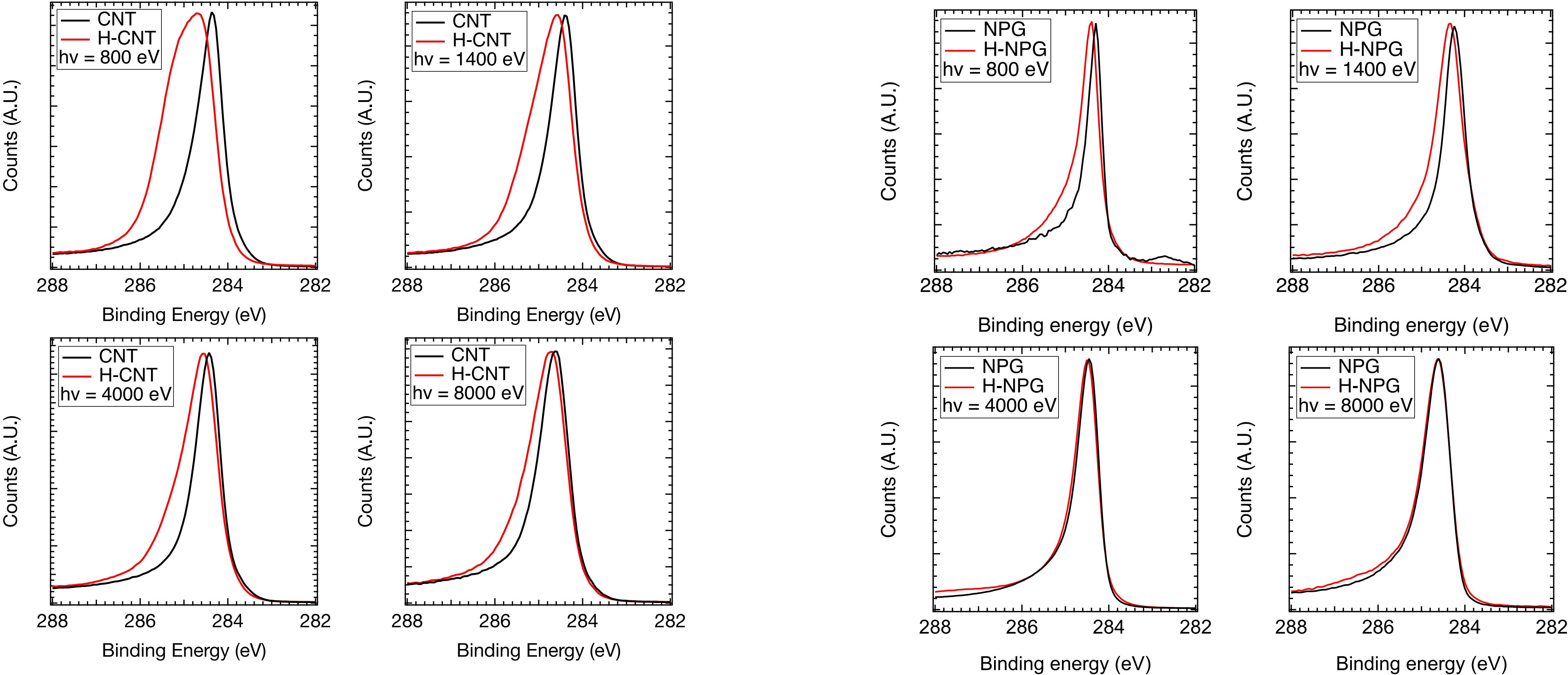
$$m_H = 1 \text{ ug} \approx 30$$

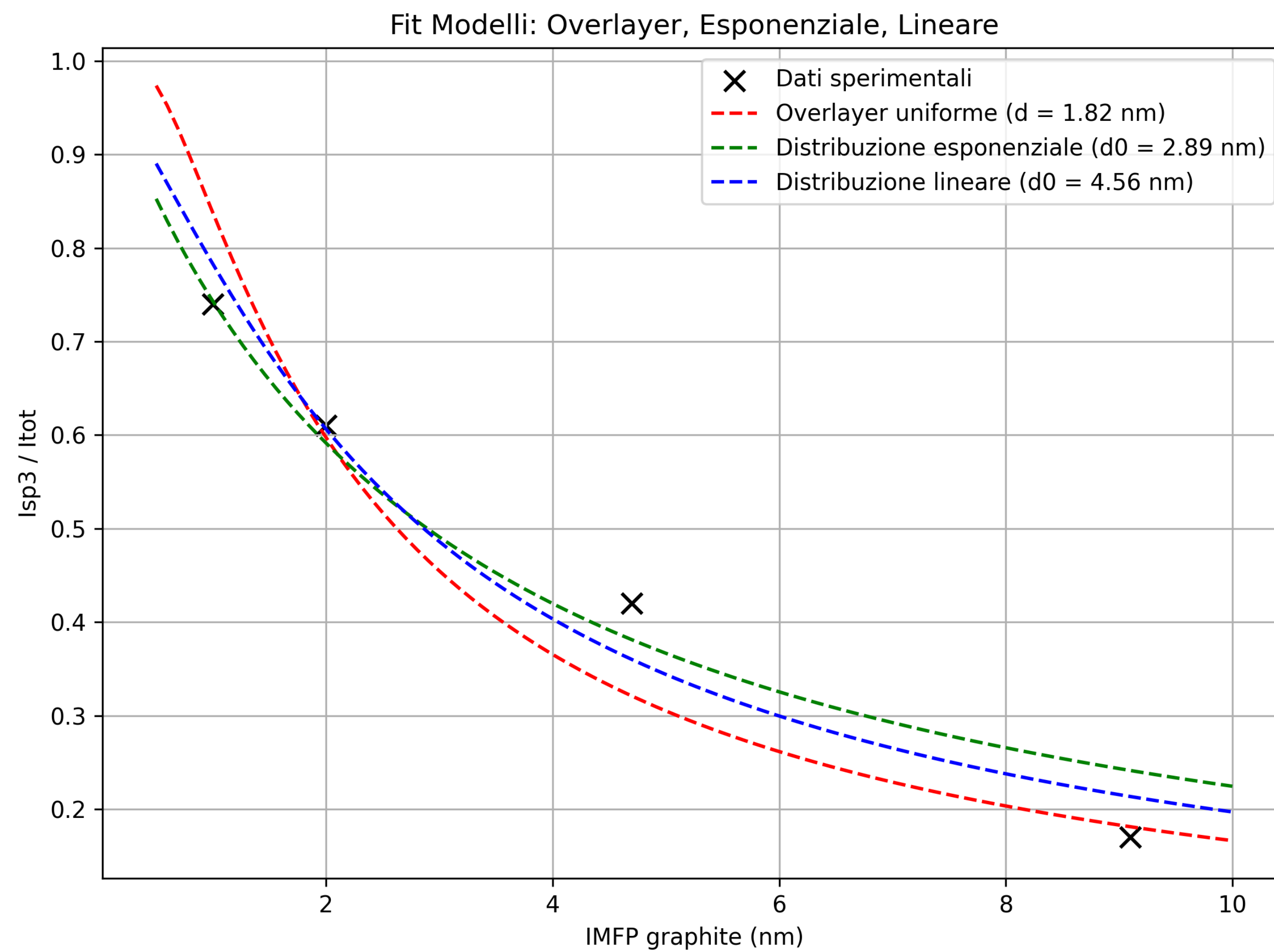


*

* 15 cm² total

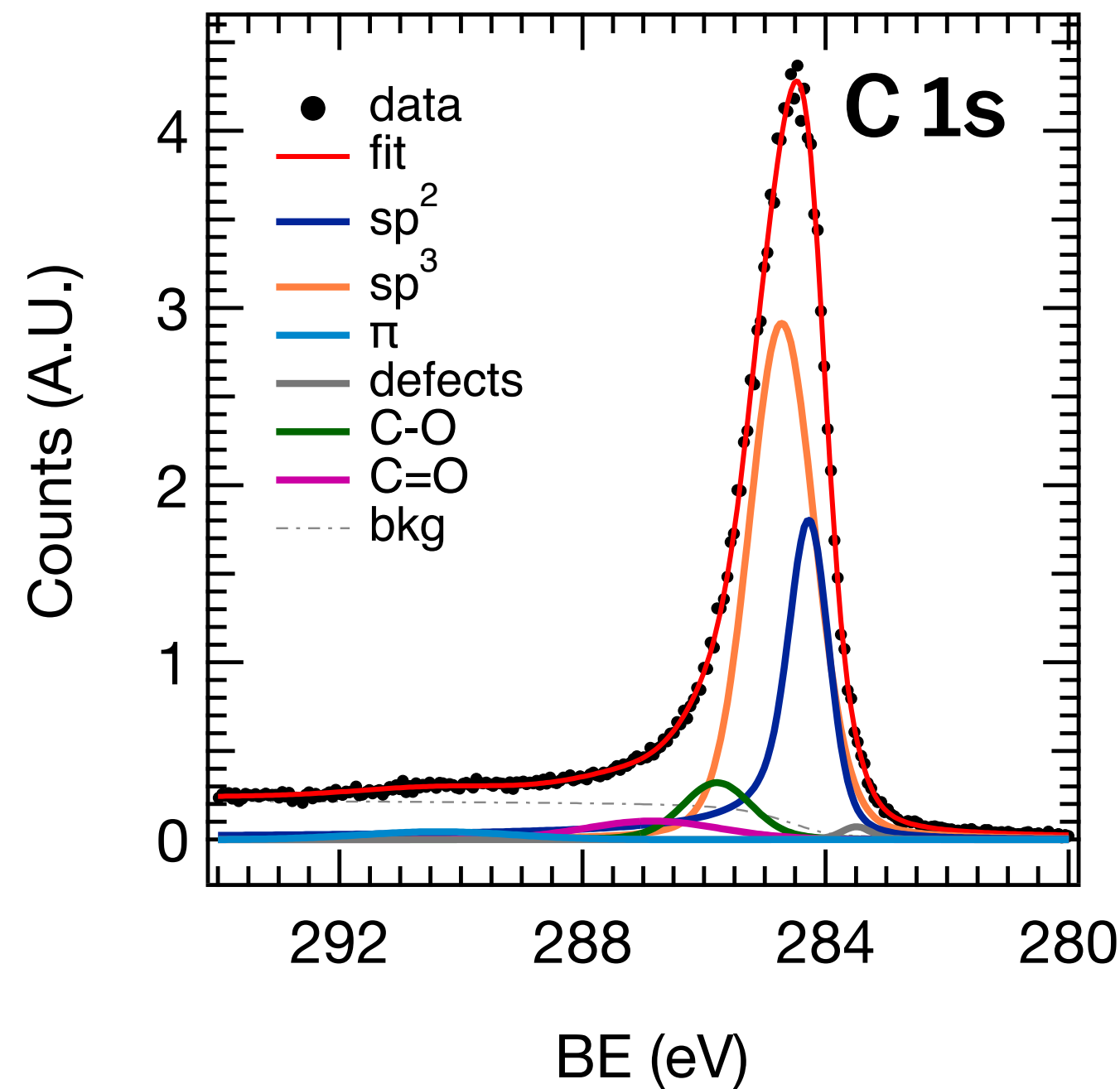
Not surface-limited, not fully bulk... but more H than in nano-porous graphene ²⁸





Hydrogen Plasma VS Thermal Cracking Hydrogenation

30



$$I_{sp3} / (I_{sp2} + I_{sp3}) \approx 67\%$$

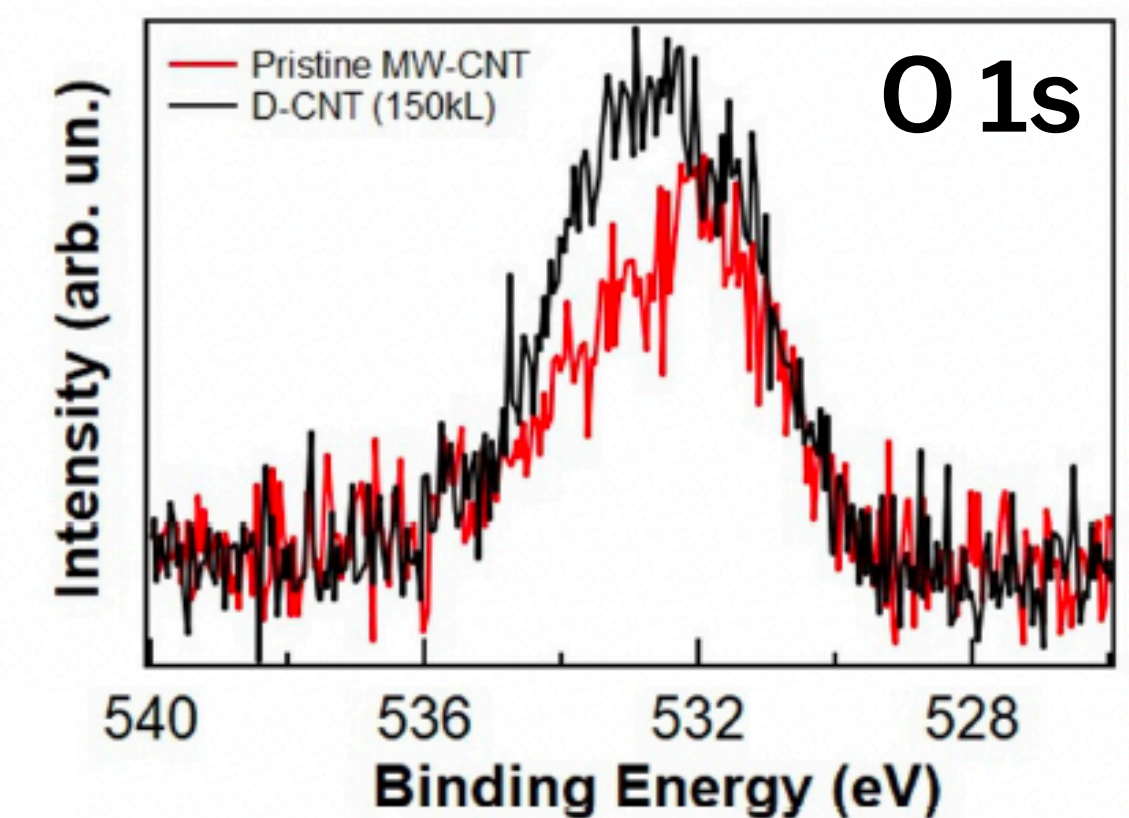
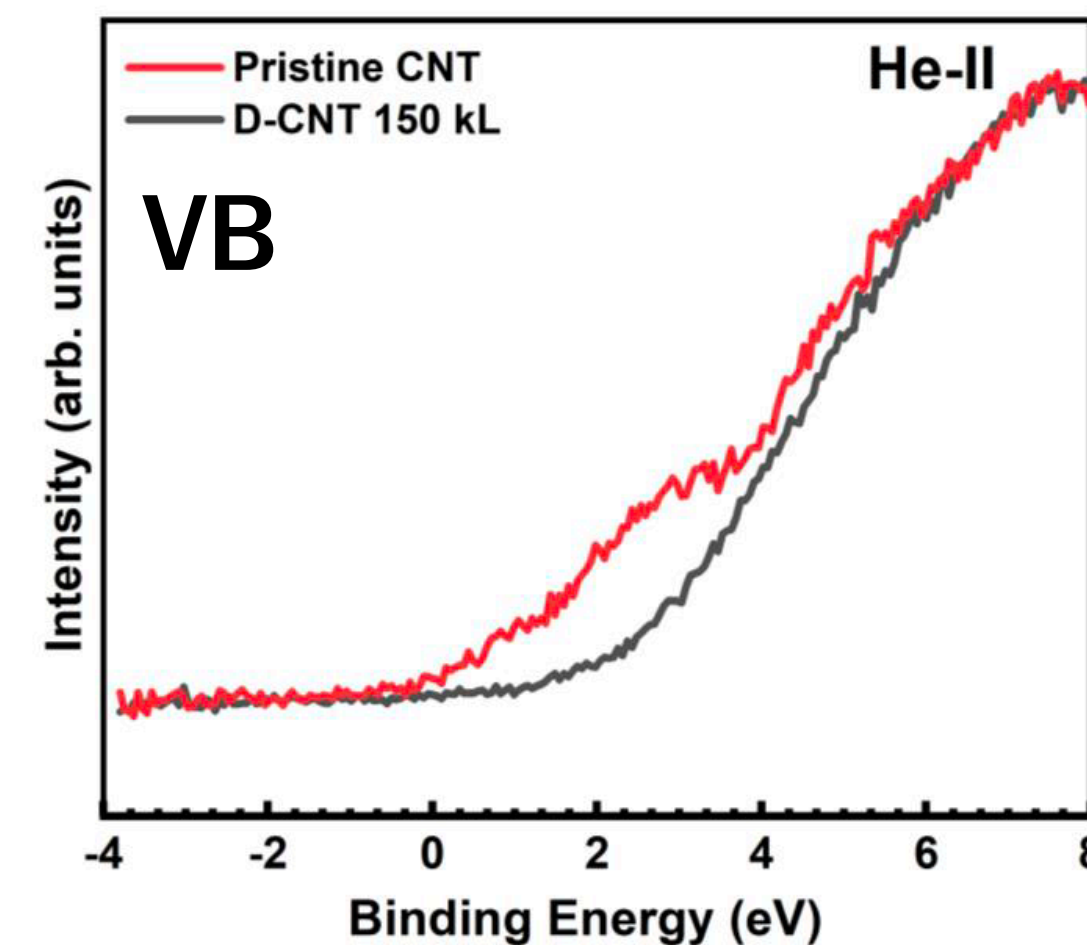
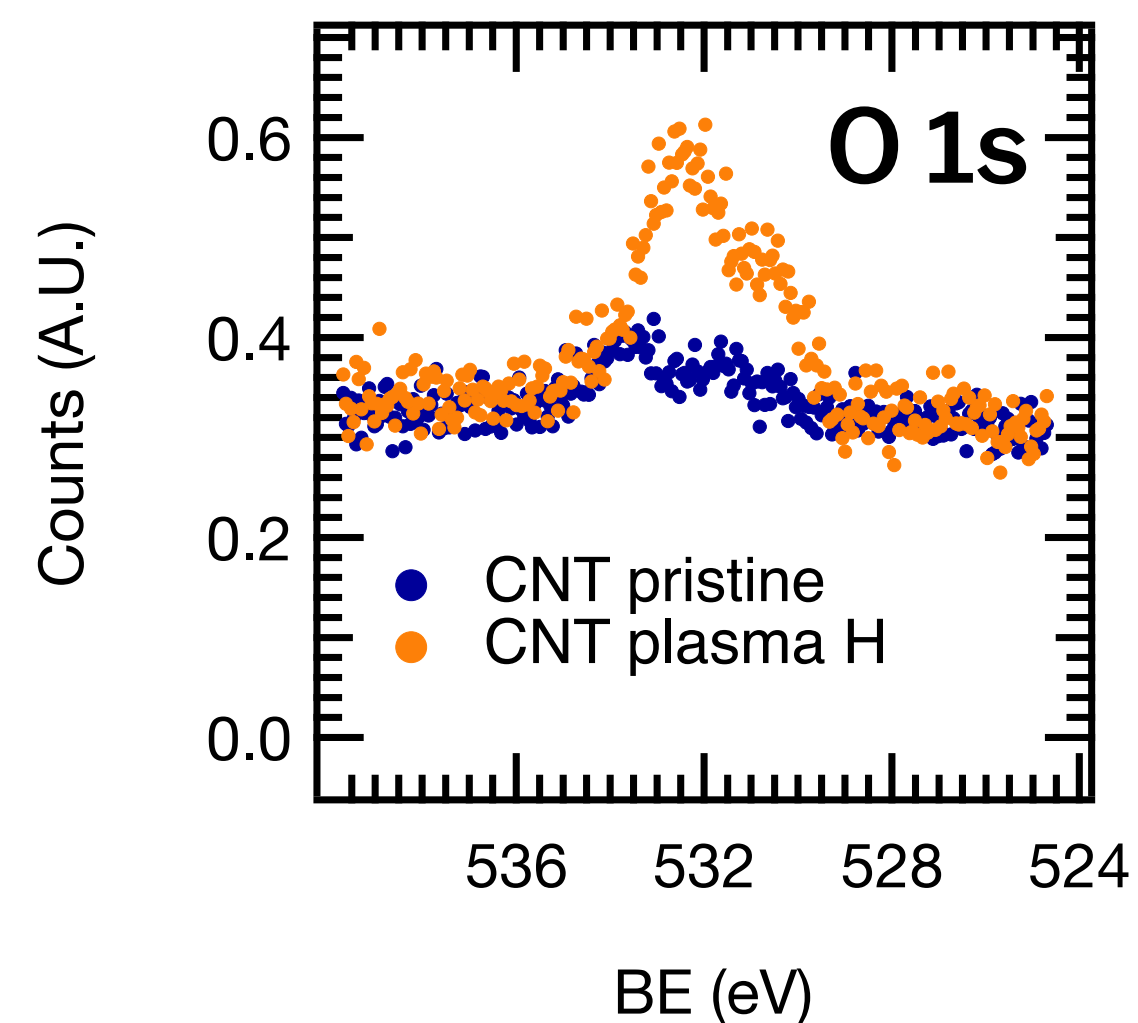
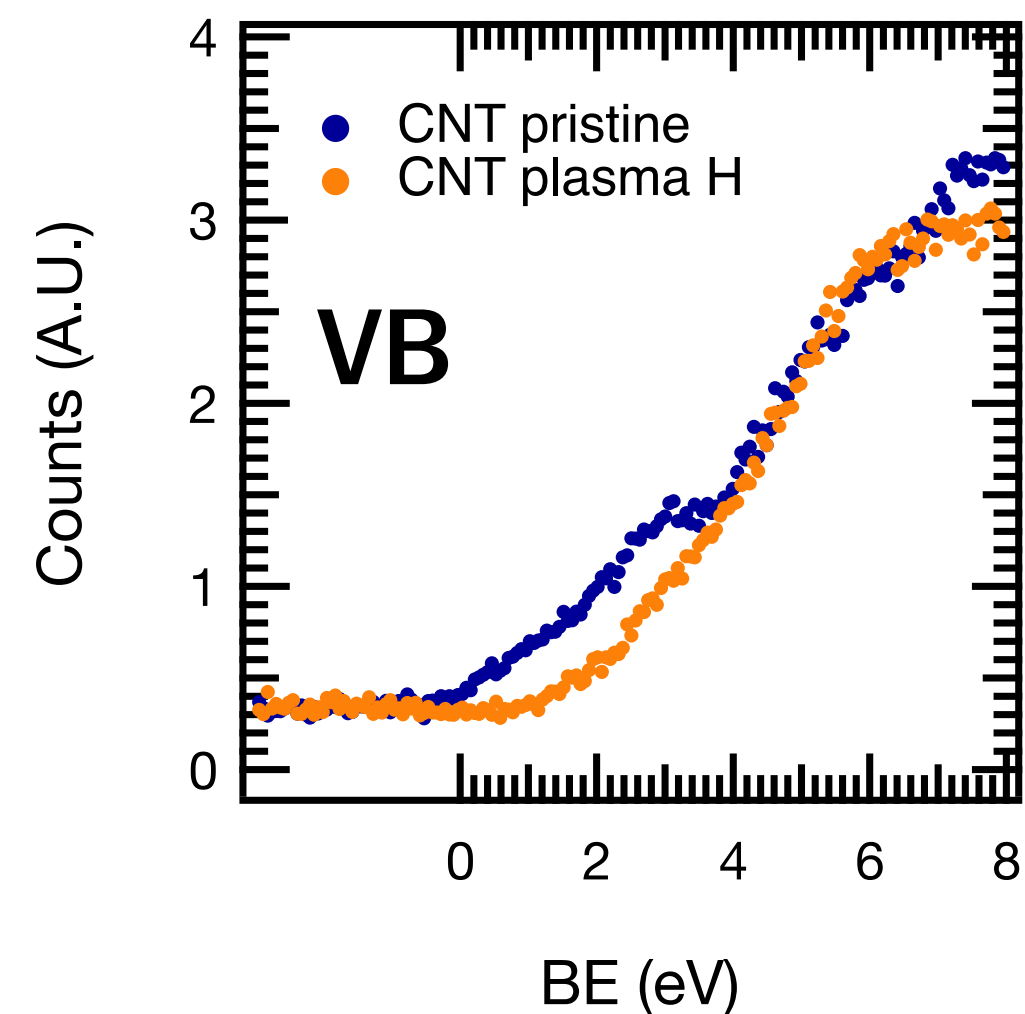
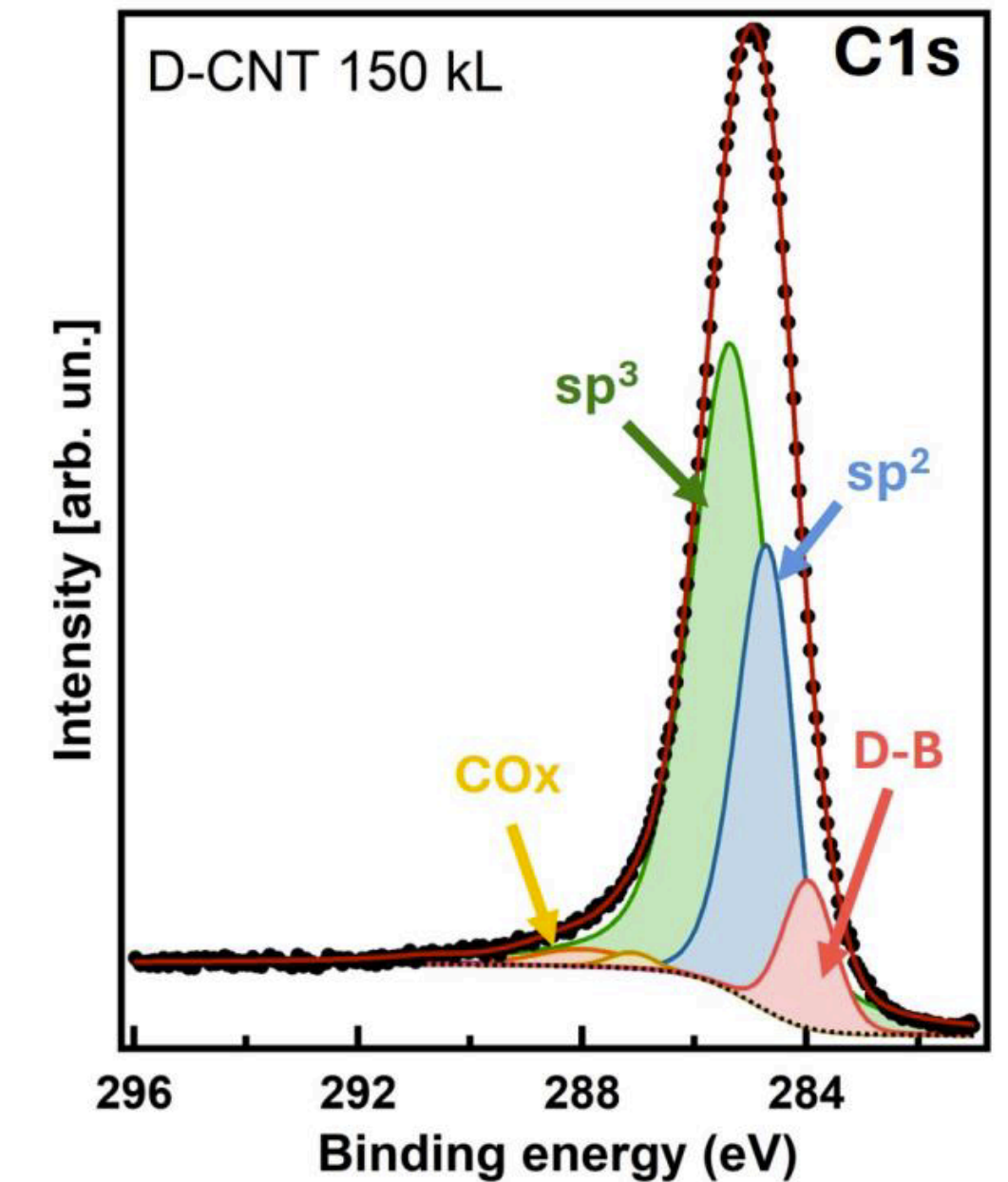
$$WF = 3.81 \pm 0.05 \text{ eV}$$

Atomic Deuterium bonding to Multi-Walled Carbon Nano Tubes

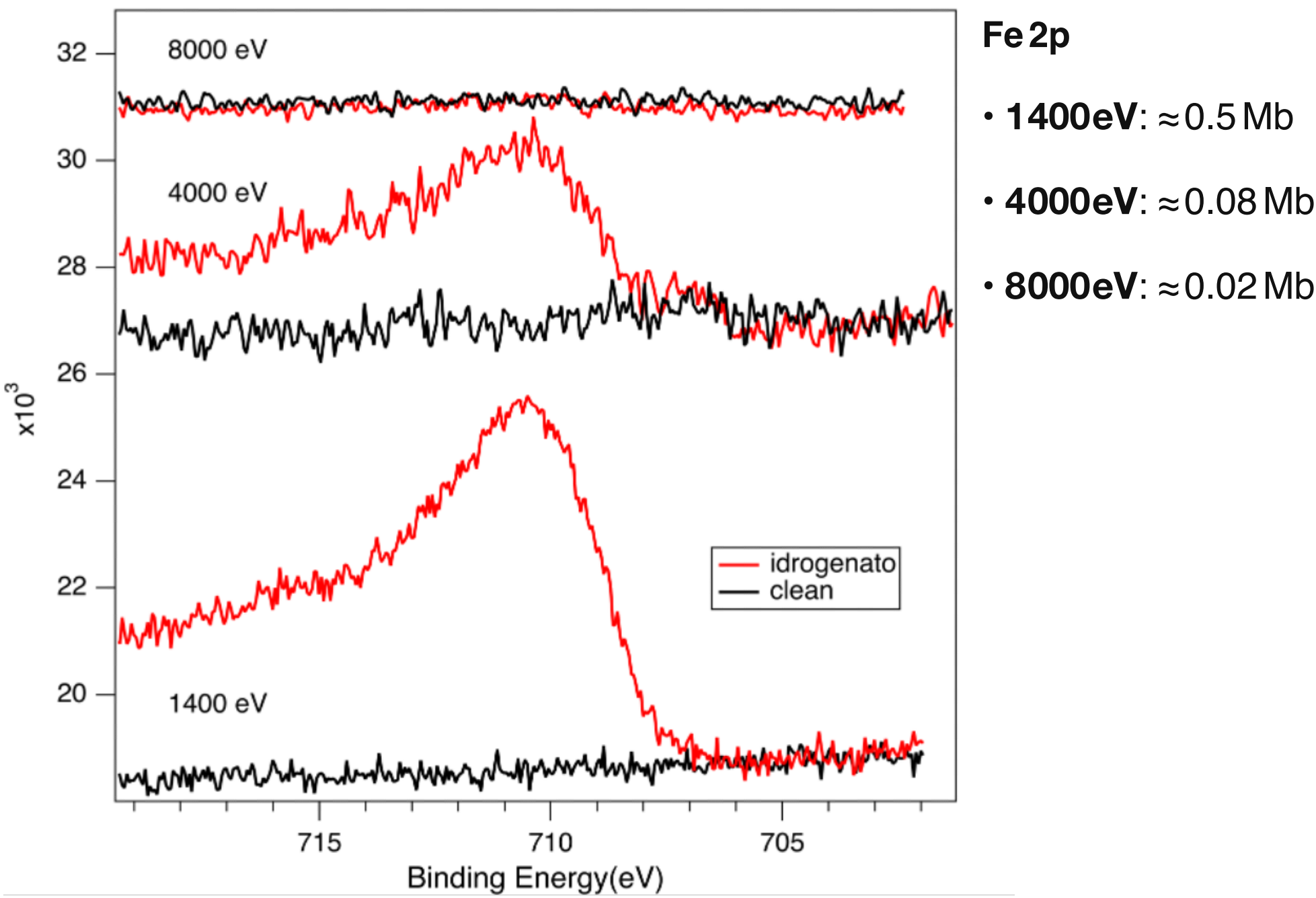
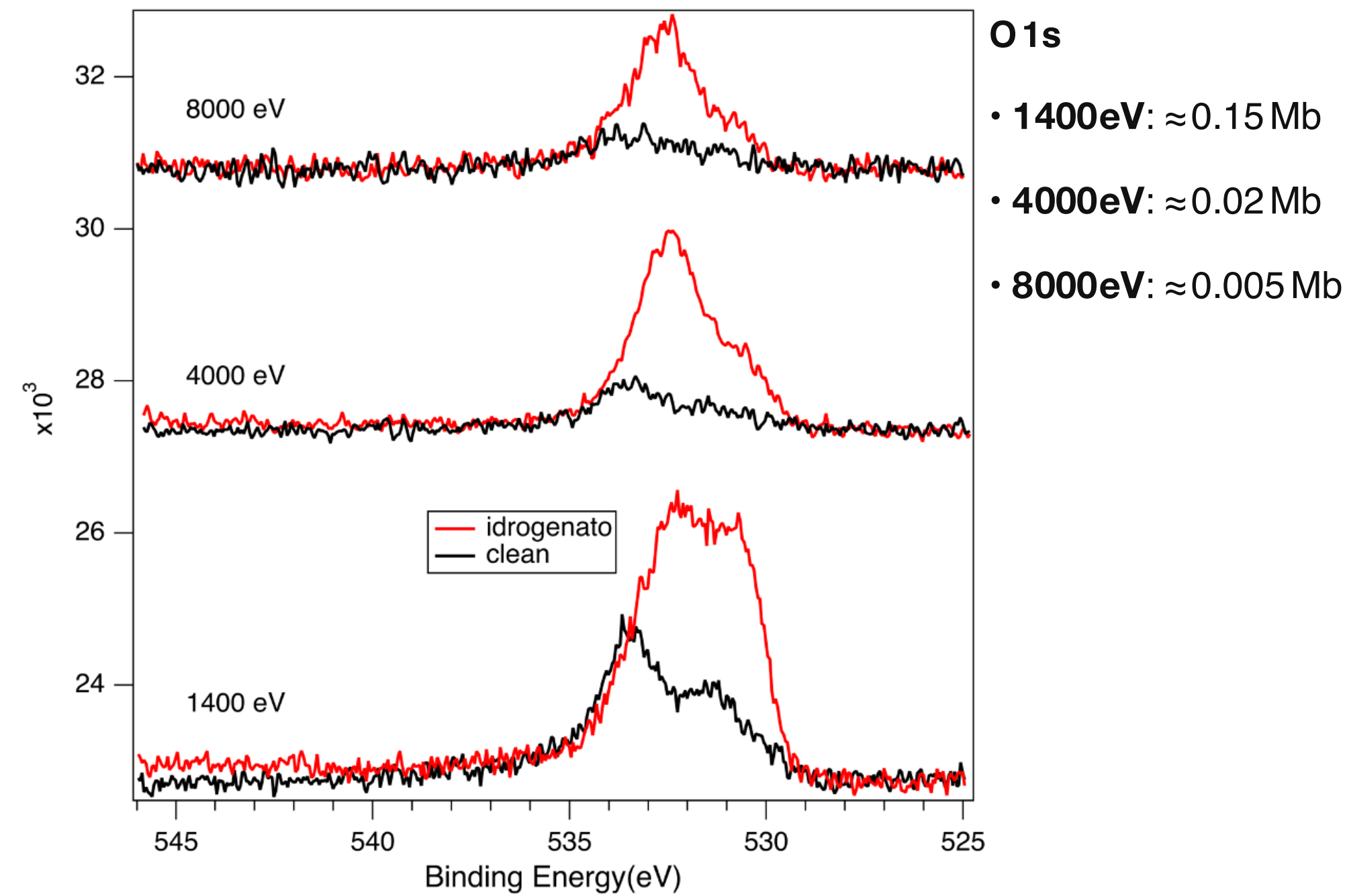
Sammar Tayyab et al., to be published

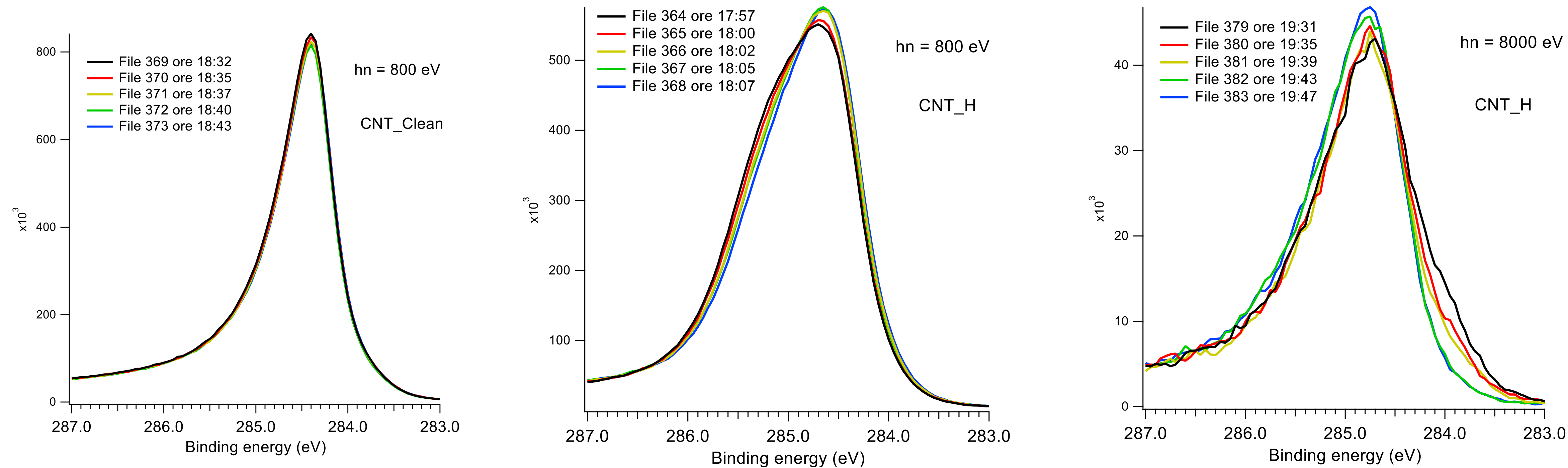
$$I_{sp3} / (I_{sp2} + I_{sp3}) \approx 70\%$$

$$WF = 3.84 \pm 0.05 \text{ eV}$$

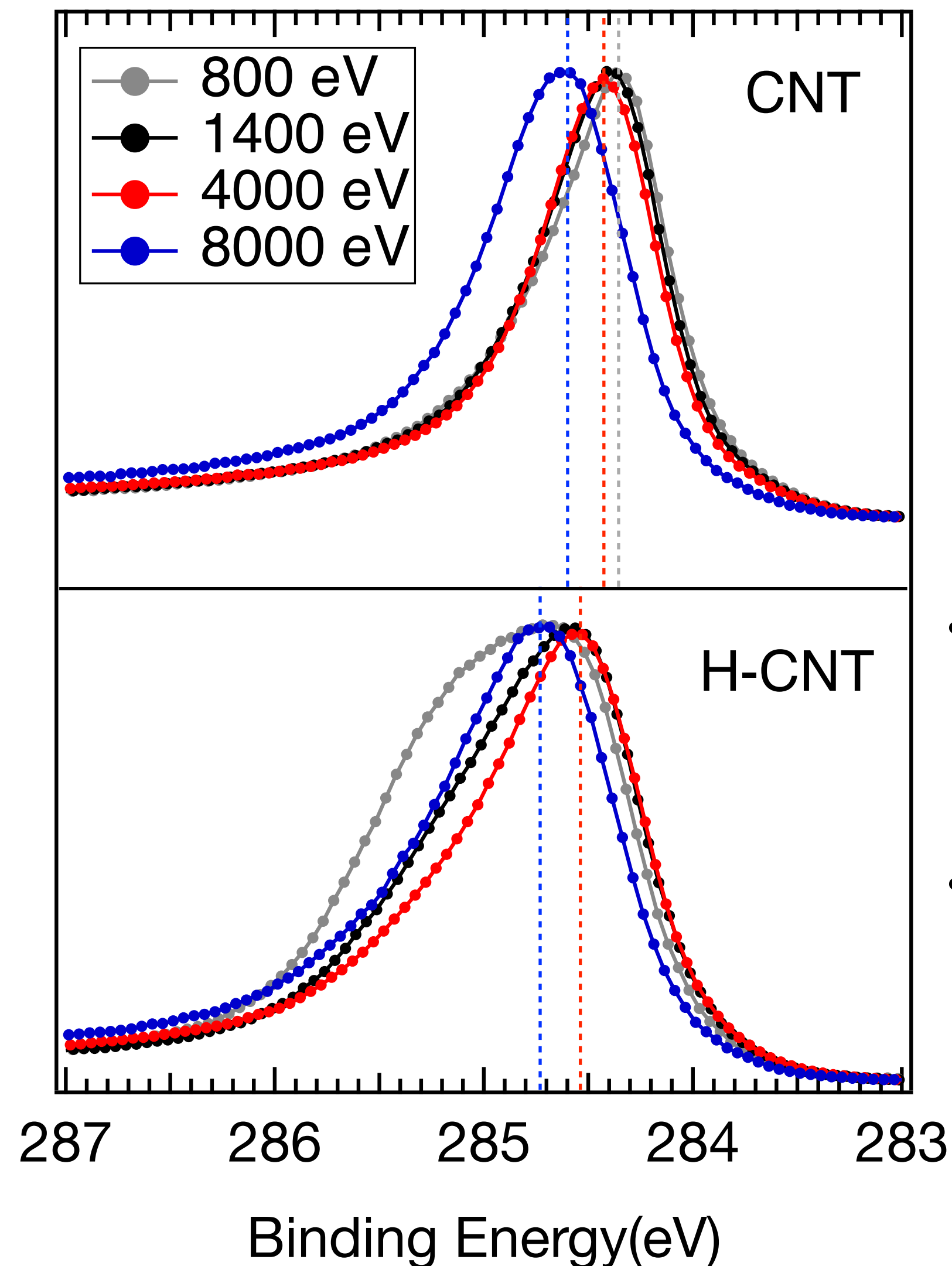


Cross section values
Yeh & Lindau (1985)
Scofield (1973)





- Recoil: shift
- Kind/number of phonon excitations: asymmetric broadening

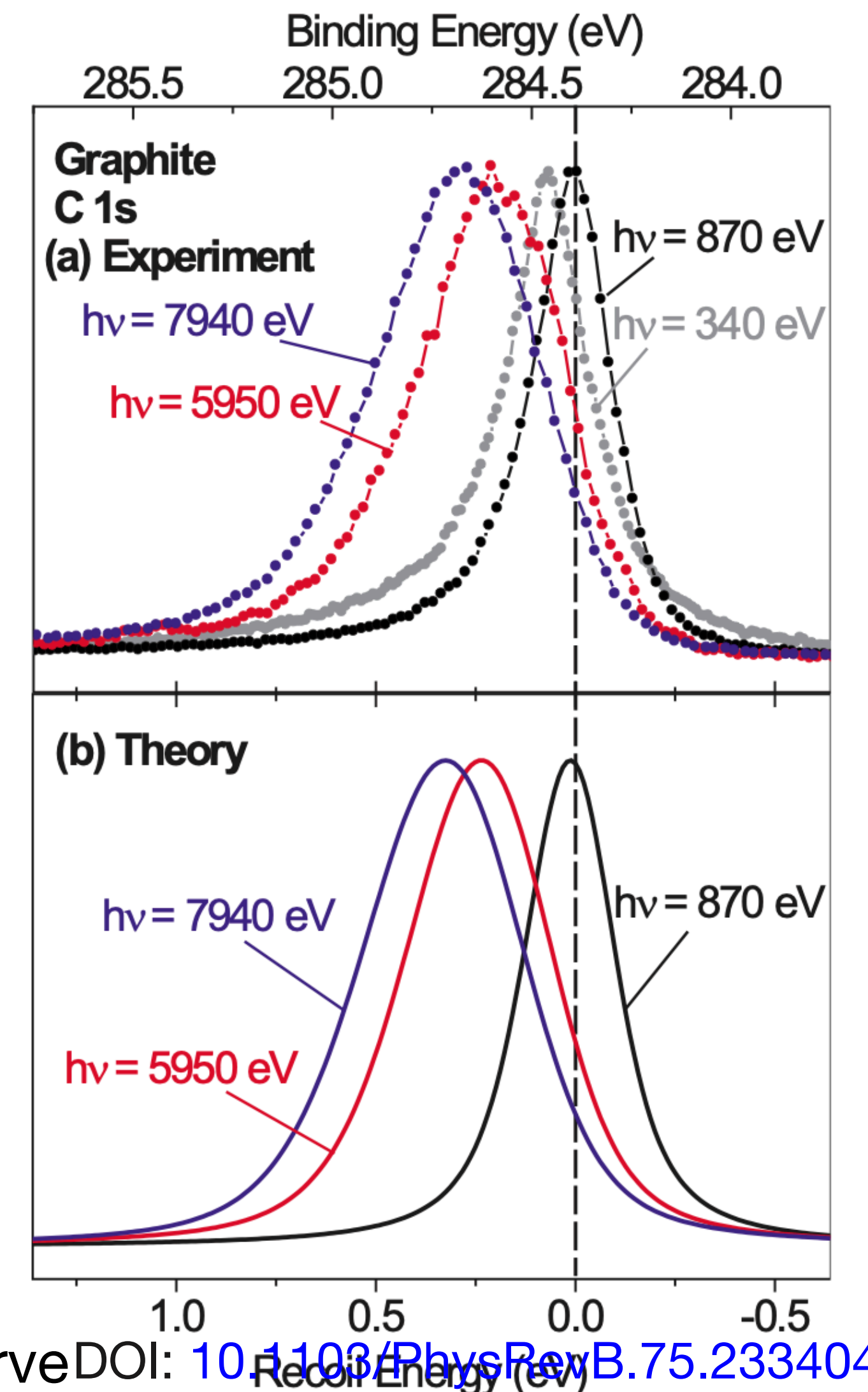


- 800 eV ~ 1400 eV no major C-C recoil effect
- 4000 eV energy shift ~ 100 meV
- 8000 eV energy shift ~ 250 meV

- prominent sp³ component gradually decreases with photon energy

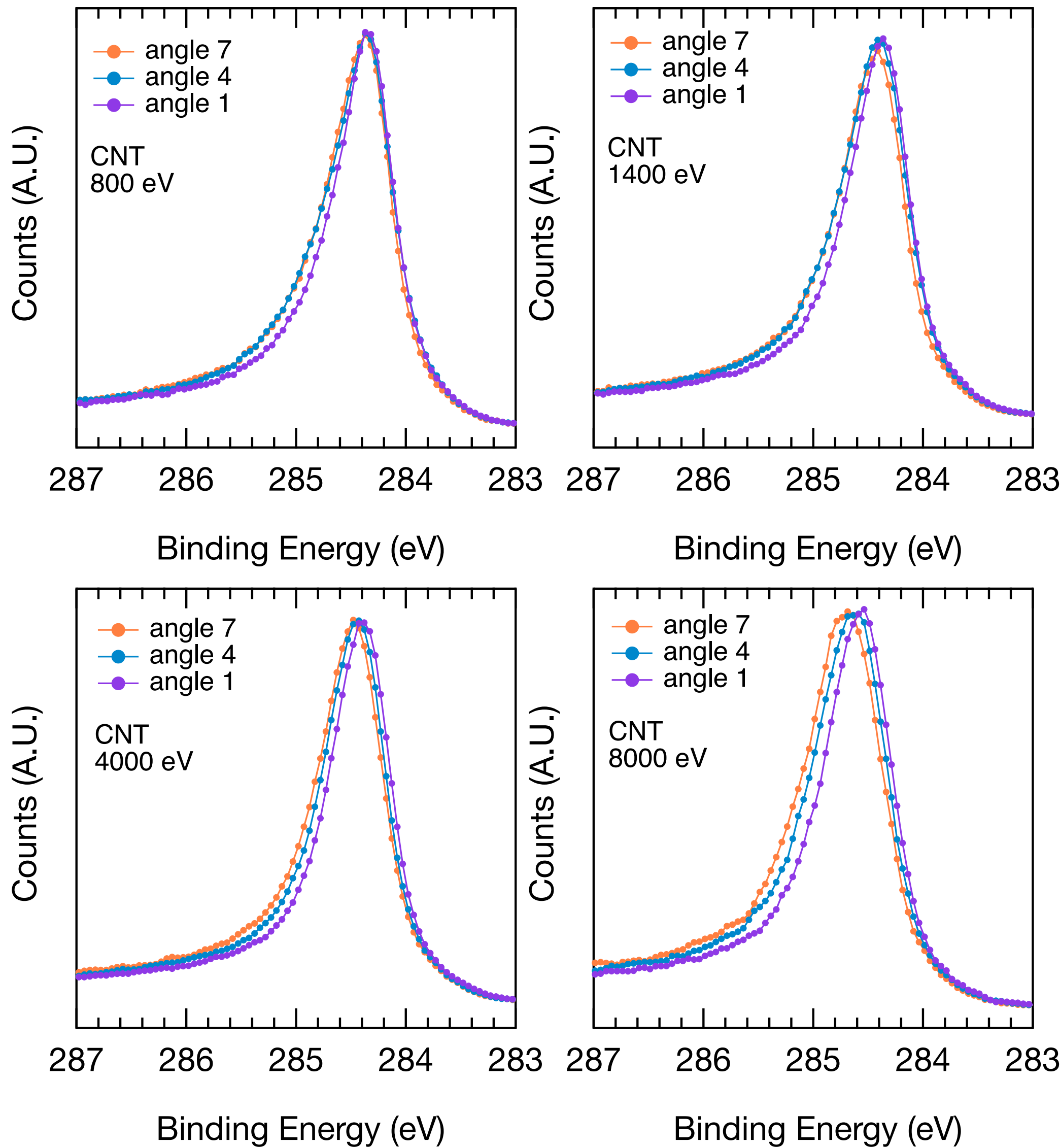
- From 4000 eV to 8000 eV increasing recoil effect

- Modello semiclassico
- Armonico (no dissipaz)
- Single phonon
- No coupling elettronico (curve simmetriche no coda di drude)

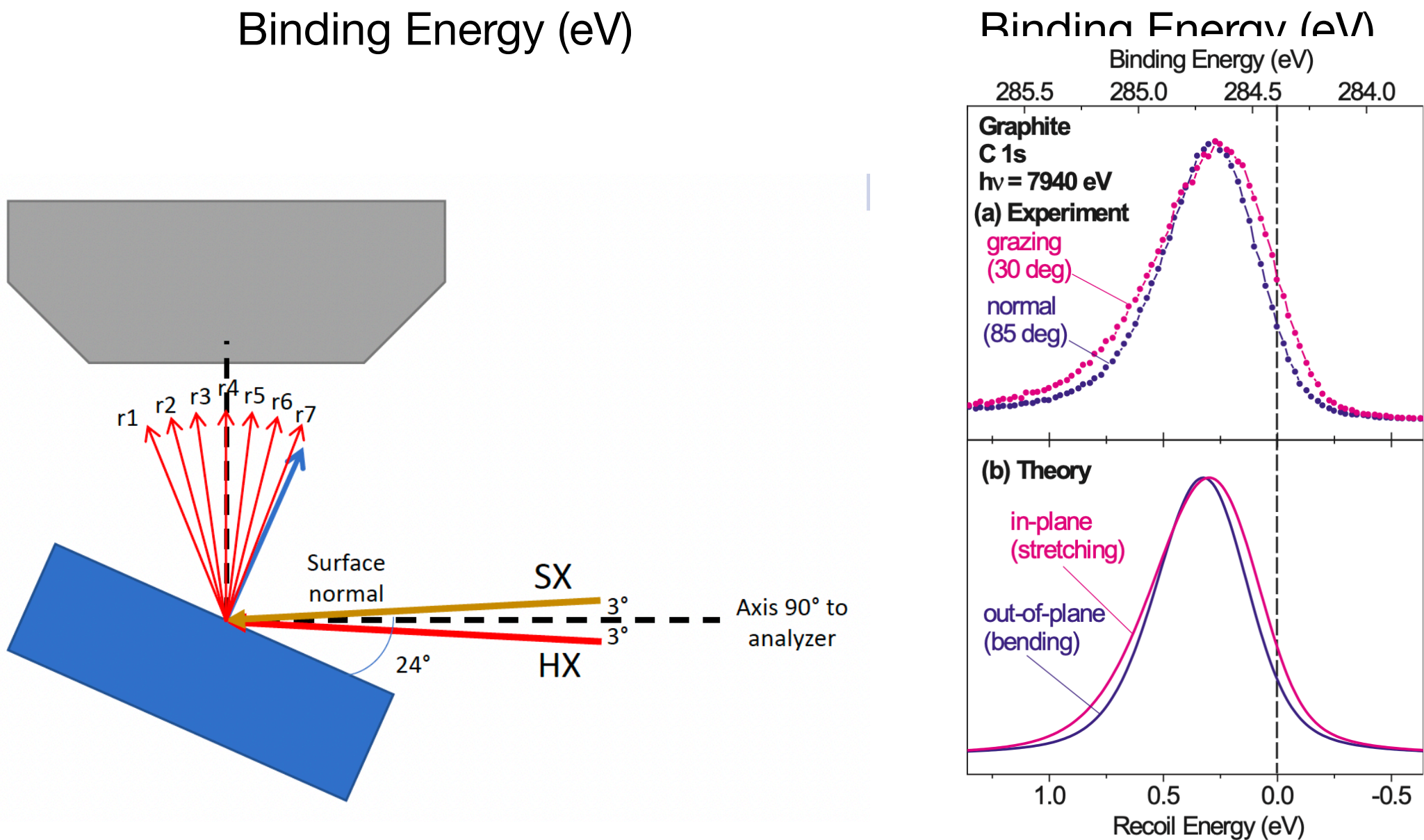
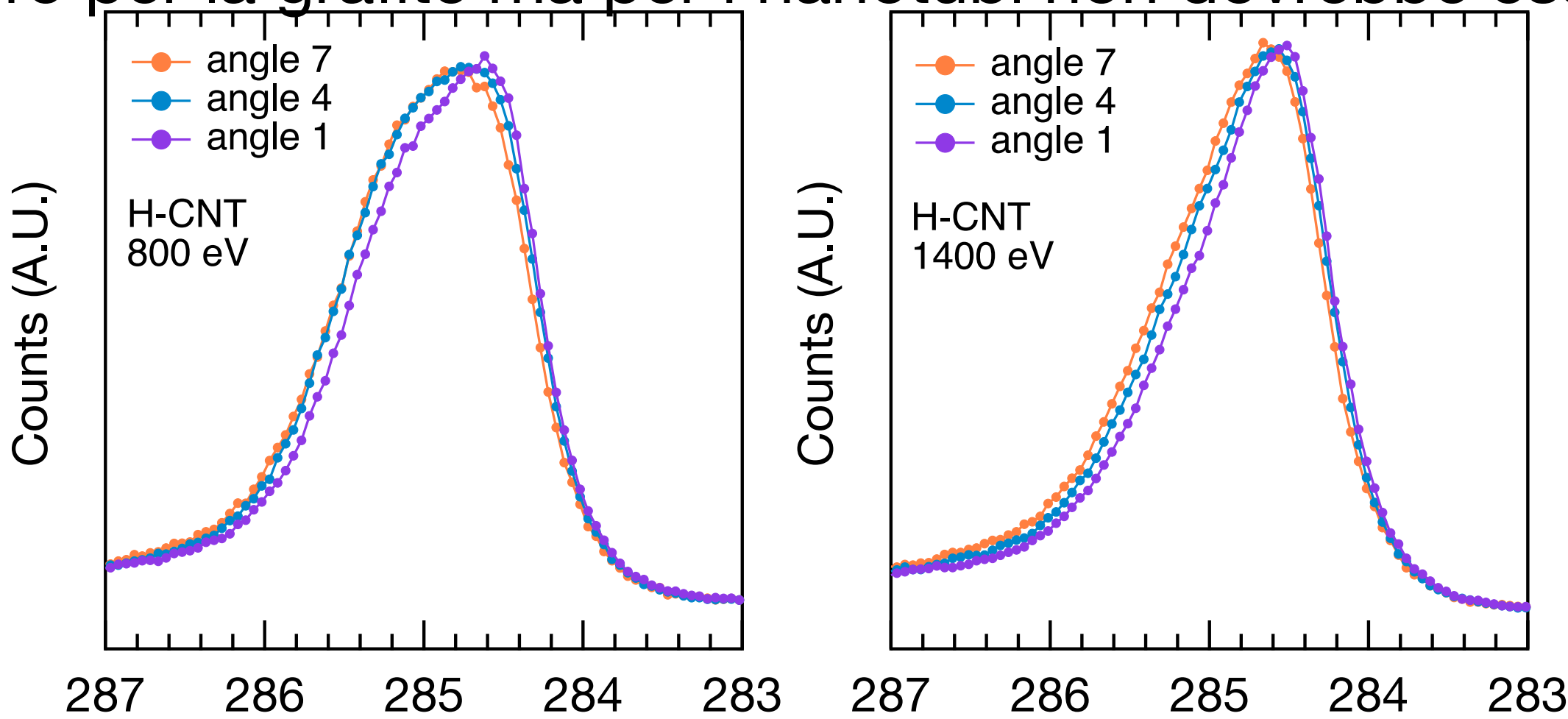


Bending vs Stretching modes with High Energy ARPES

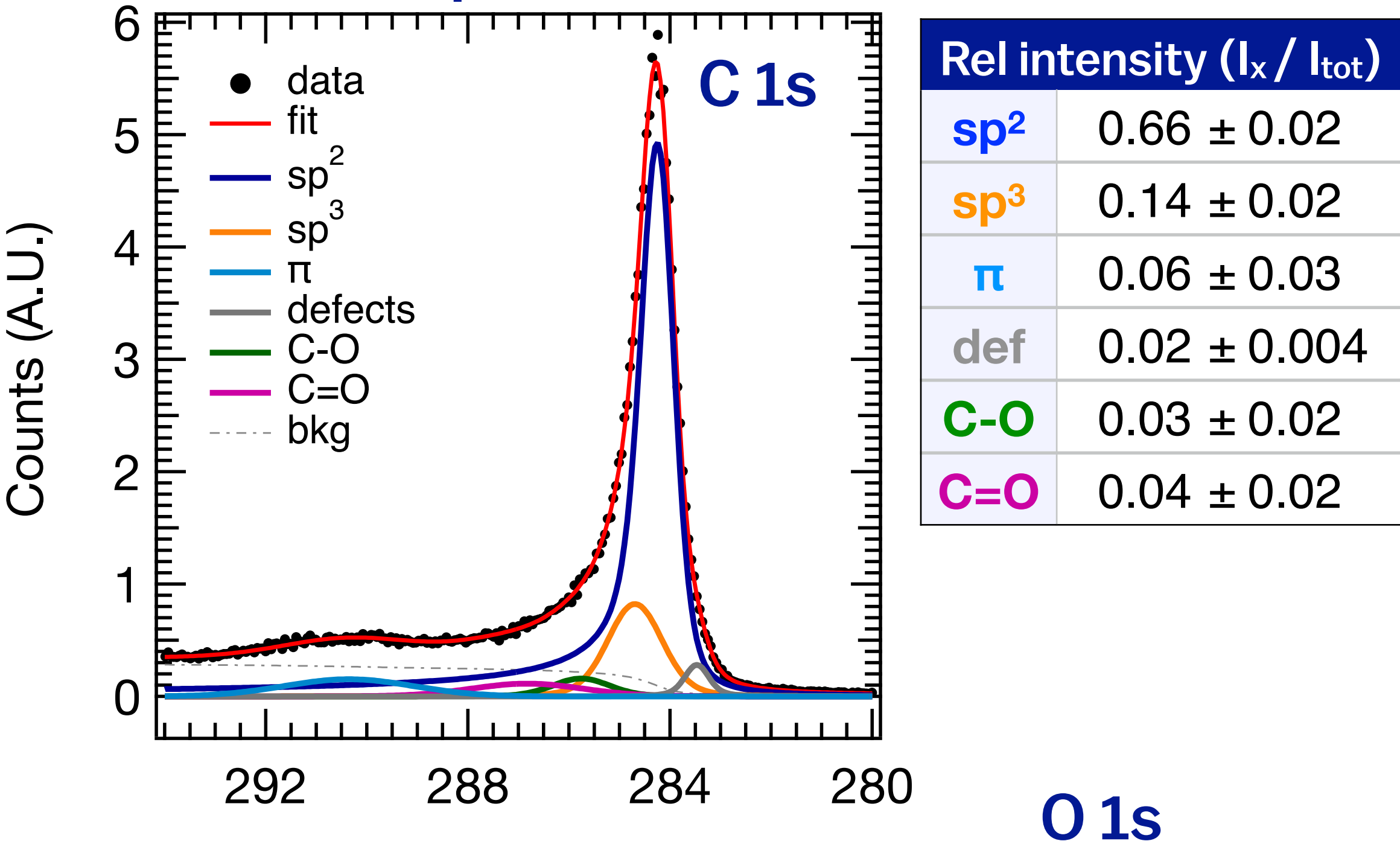
- Angle 7: normal emission



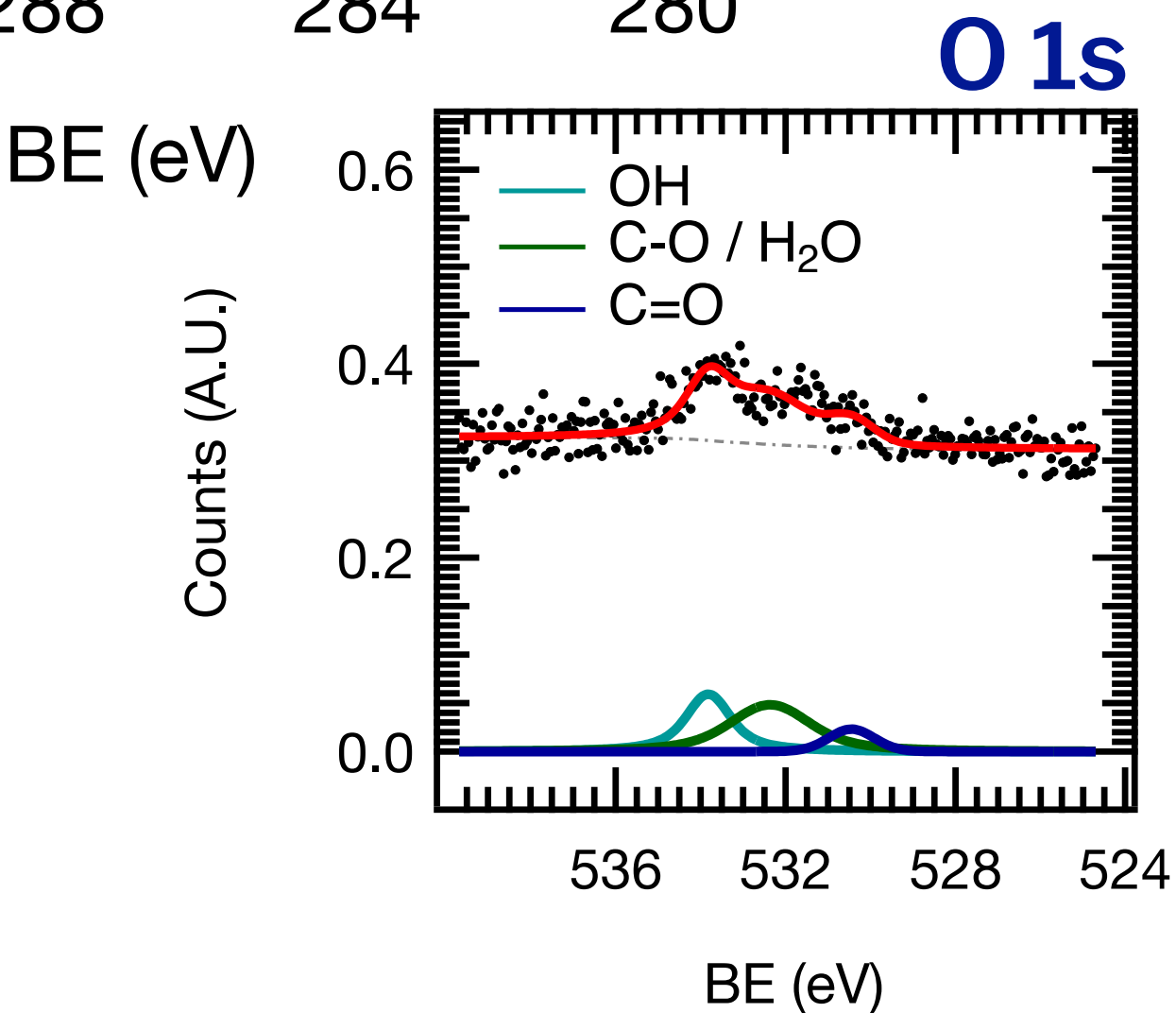
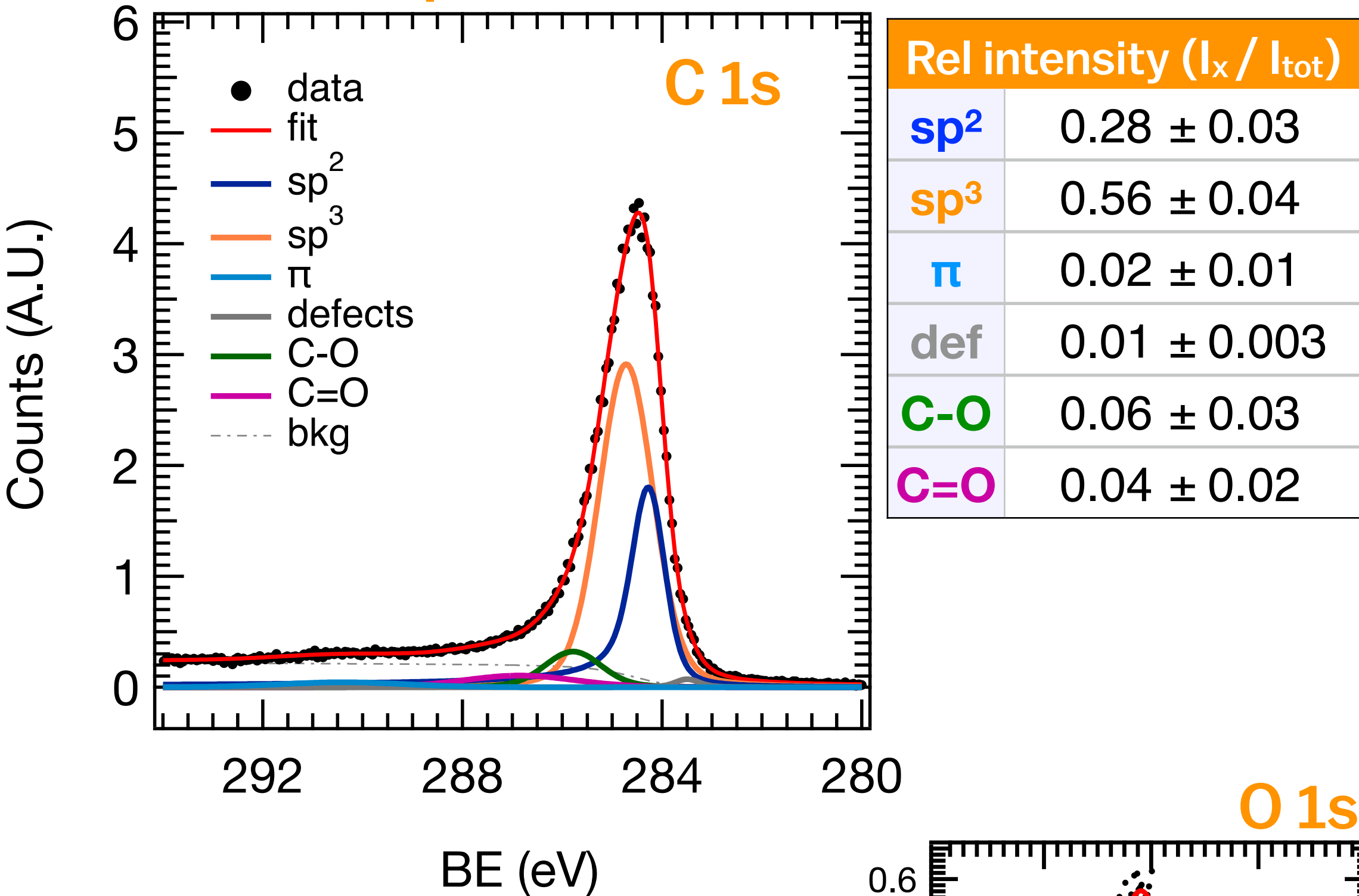
- Out of plane (normal emission): higher recoil energy
È vero per la grafite ma per i nanotubi non dovrebbe essere l'op



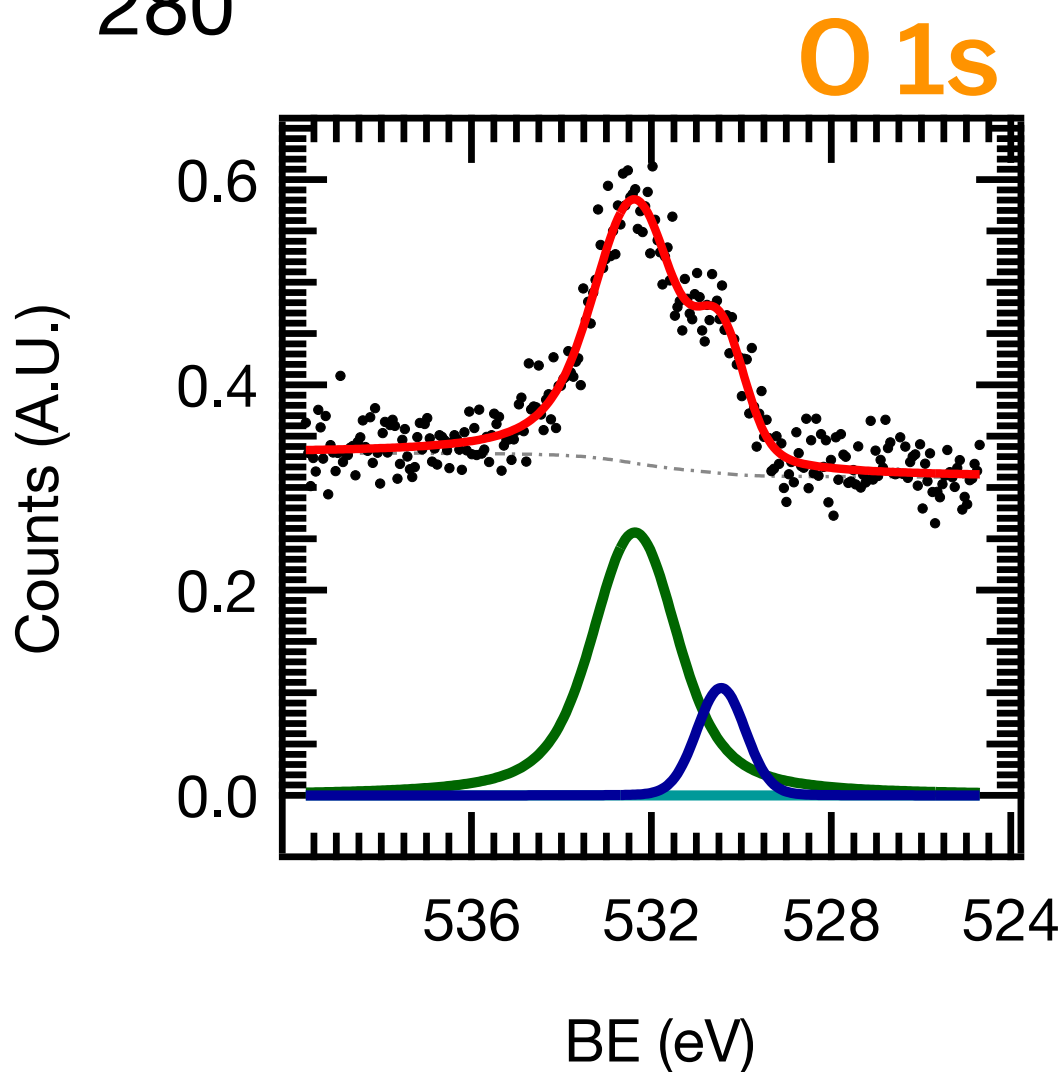
CNT pristine



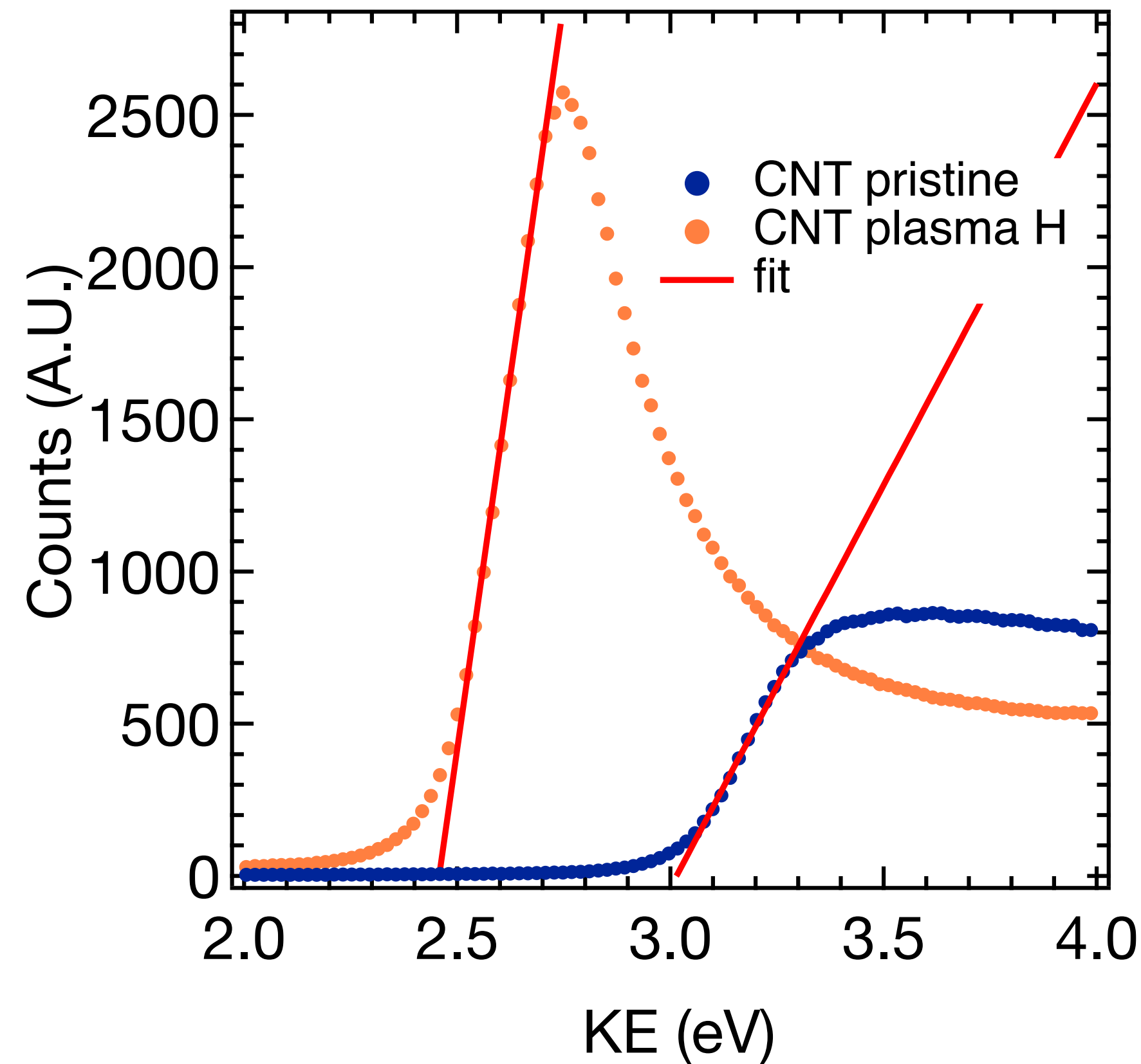
CNT plasma H



○ sp^2 decreases, sp^3 increases (x4)
○ C-O and C=O increases



Secondary Electron onset (SE)



$V_{\text{BIAS}} = -3 \text{ V}$

Work Function measure:

$$\phi_{\text{sample}} = h\nu - \Delta E_k = h\nu - (\text{FL} - \text{SE})$$

Surface dipole changes:

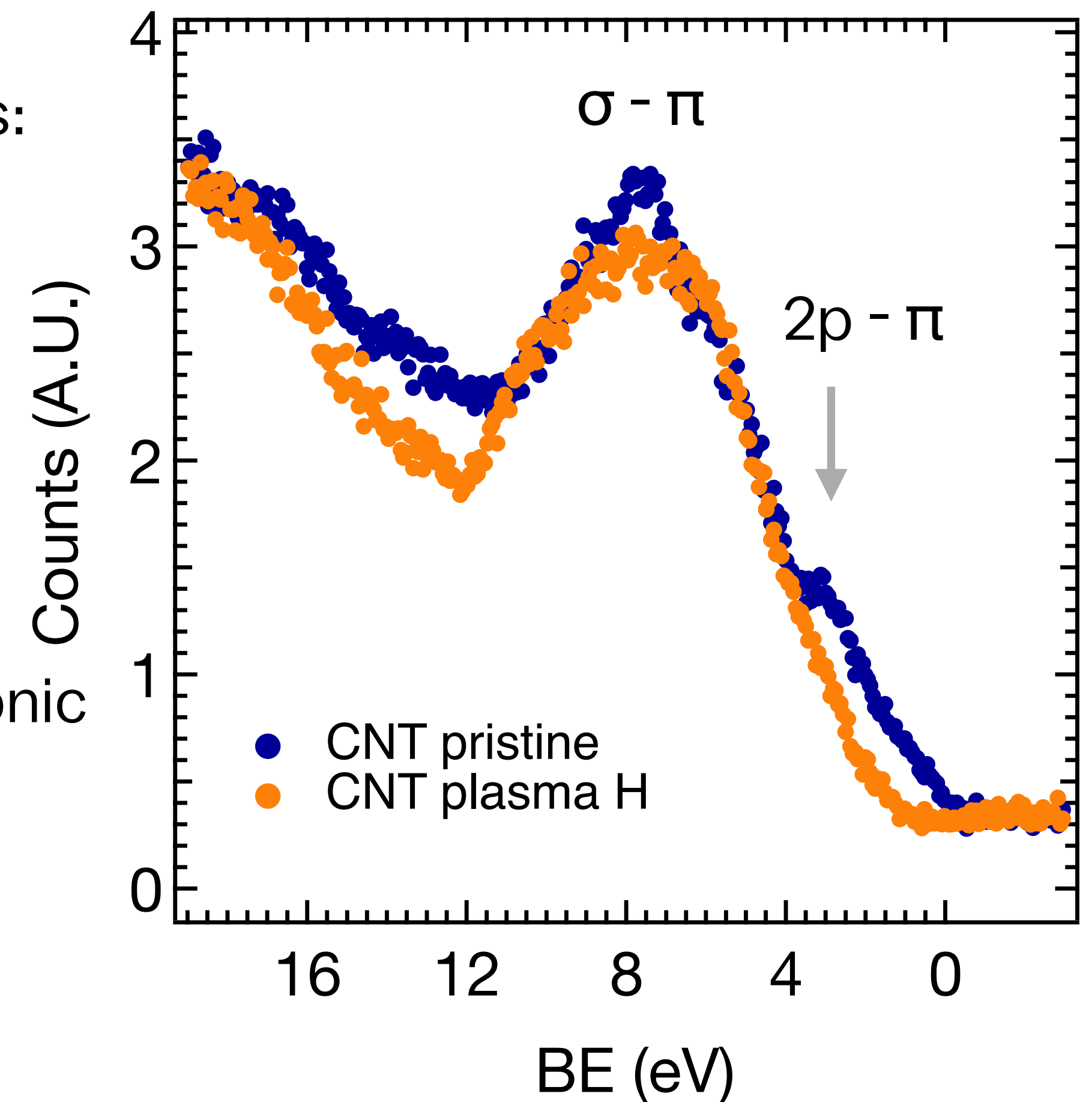
- **WF lowers** from $4.33 \pm 0.05 \text{ eV}$ to $3.81 \pm 0.05 \text{ eV}$

Depletion of the electronic levels near Fermi level

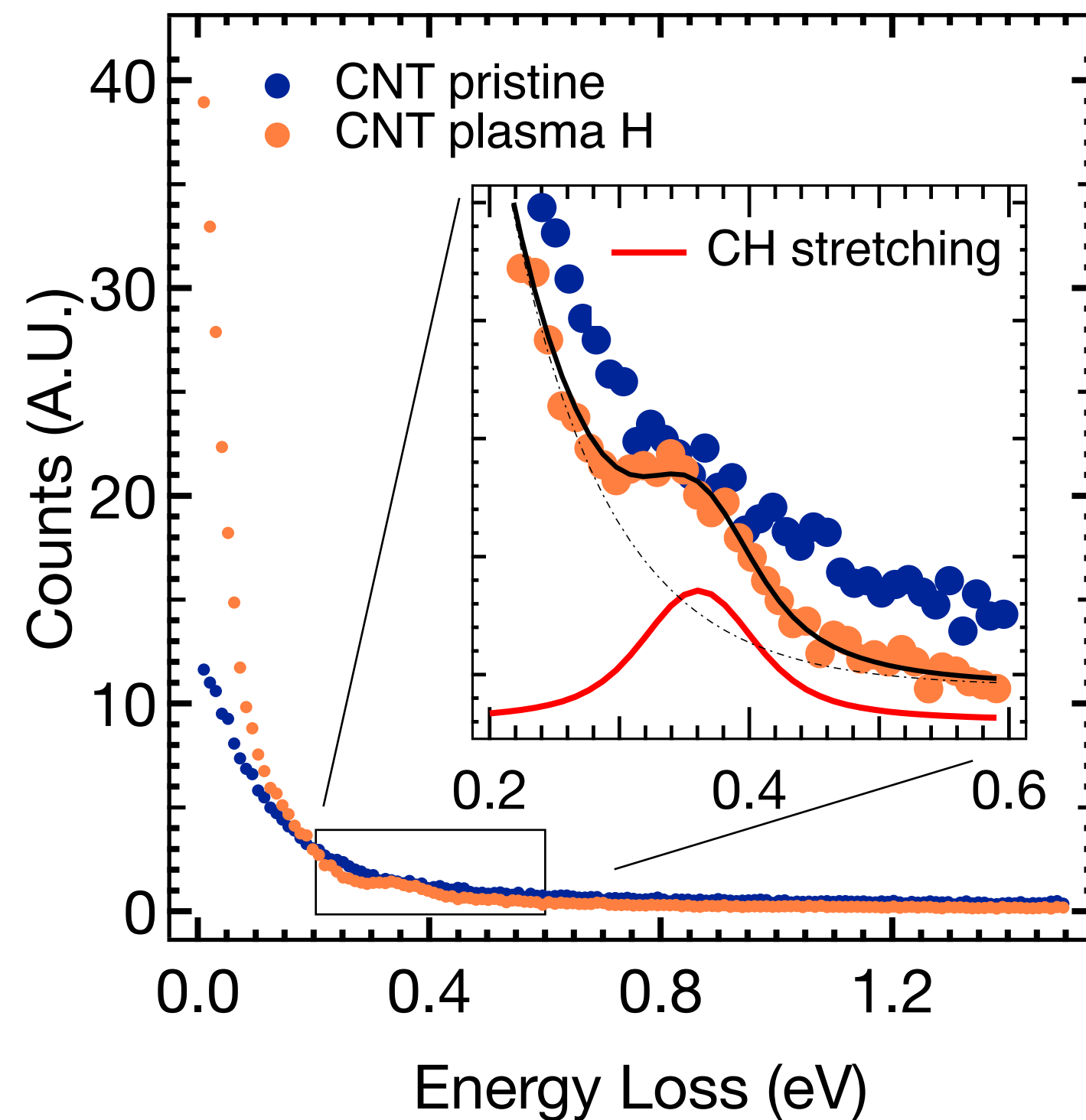
- **Gap opening**

Valence Band

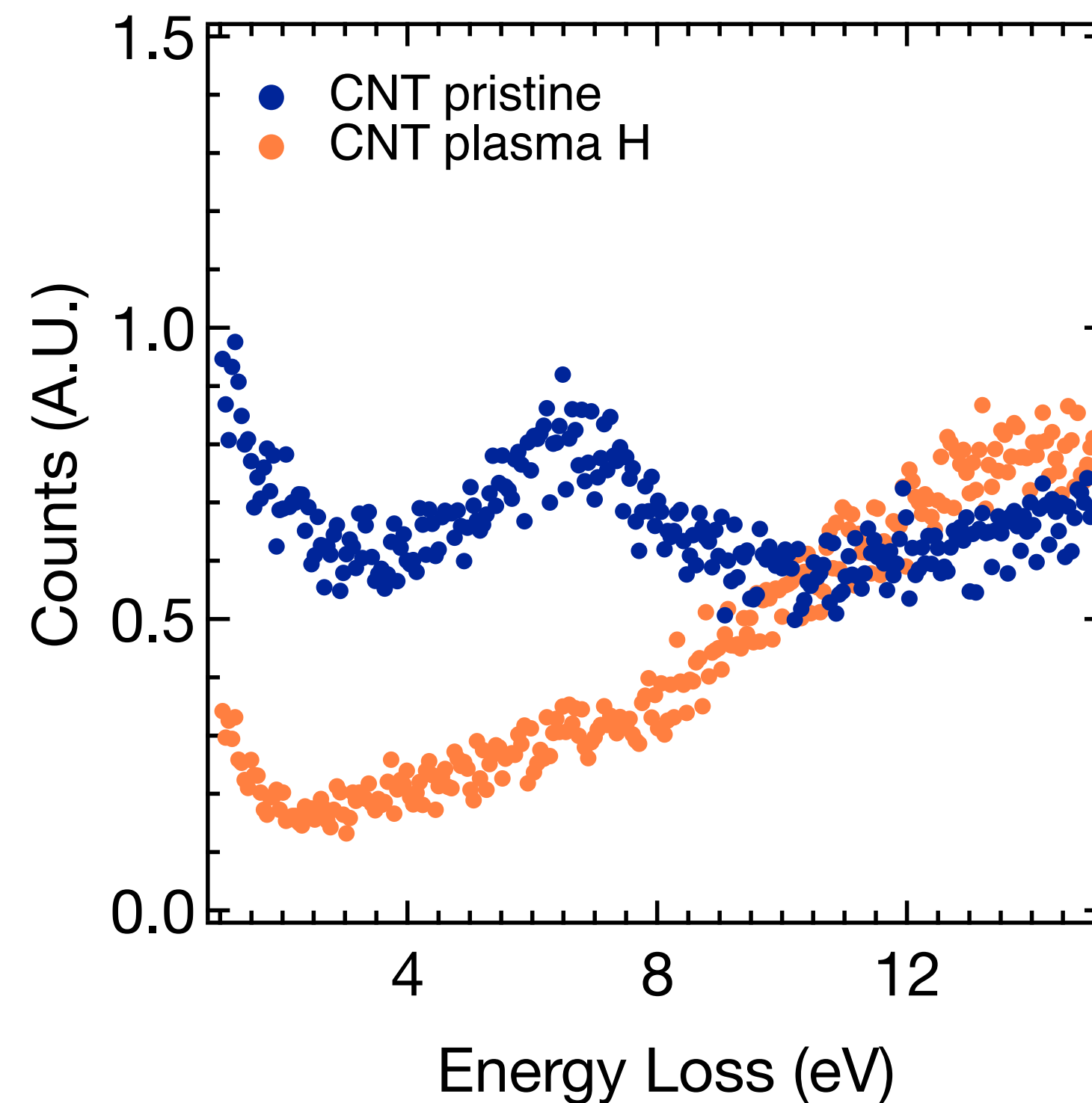
He II



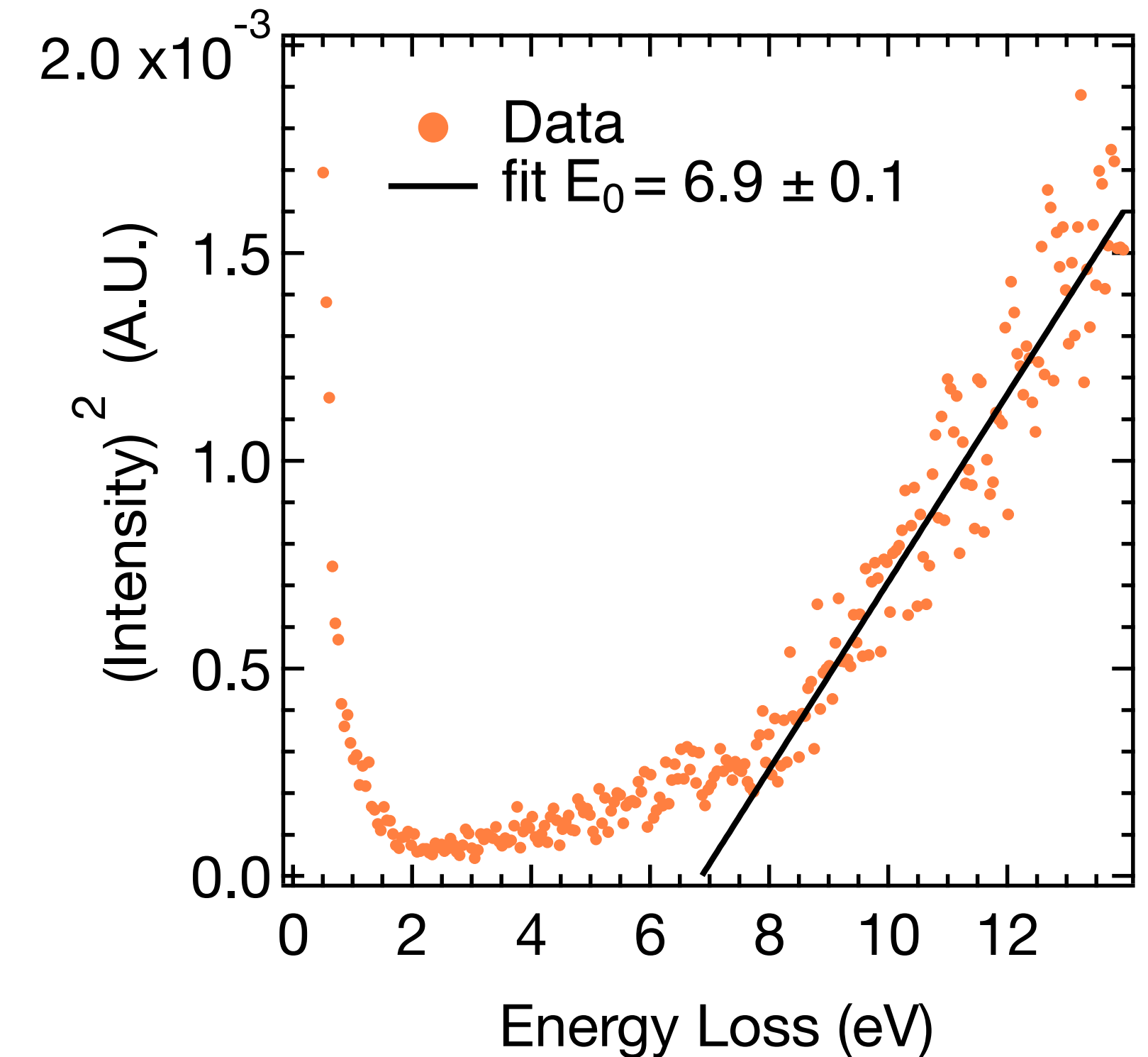
Electron Energy Loss Spectroscopy



○ **CH stretching** appears at ~ 0.36 eV



○ Quenching of the π plasmon and low energy losses



○ Wide **band gap** opening ~ 6.9 eV