

Spin precession across timescales

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Black-hole binary spin precession

1. A sweet relativistic effect full of surprises

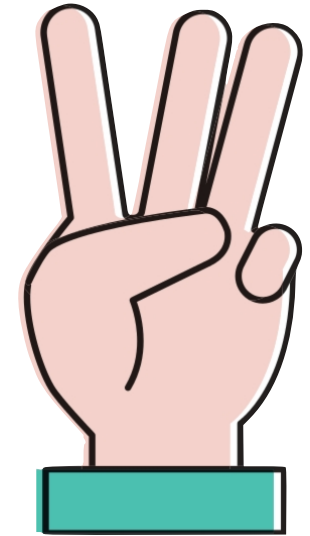
Waveform modulations, frame dragging, resonances, instabilities, superkicks

2. A treasure trove for black-hole astrophysics

Formation channels, supernova kicks, mass transfer, tides, core-envelope interactions, accretion disk physics



Three things today...



1. Precession-averaged PN dynamics

DG+ 2015, arXiv:1506.03492

DG+ 2023, arXiv:2304.04801

2. Up-down instability

DG+ 2015, arXiv:1506.09116

Mould, **DG** 2020, arXiv:2003.02281

3. What precession parameter

DG+ 2021, arXiv:2011.11948

(Eccentricity in the next talk by Giulia)

Fumagalli, **DG** 2023, arXiv:2310.16893

A tale of three timescales

1. **Orbital** motion $t_{\text{orb}} \propto (r/r_g)^{3/2}$
Kepler's third law
2. Spin & orbital-plane **precession** $t_{\text{pre}} \propto (r/r_g)^{5/2}$
Apostolatos et al 1994
3. GW emission and **inspiral** $t_{\text{RR}} \propto (r/r_g)^4$
Quadrupole formula
Peters & Matthews 1963

if (Post-)Newtonian $r \gg r_g = GM/c^2$: timescale hierarchy



BH binary multi-timescale analysis: **DG+** 2015, 2023

1. Solve the dynamics (hopefully analytically) on the shorter time
2. Quasi-adiabatic evolution ("average") on the longer time

Two-spin PN dynamics



- Three momenta, **9** components $\mathbf{L}, \mathbf{S}_1, \mathbf{S}_2$
- Spin magnitudes are const. Frame (3 constraints)

Spin precession is a **4D** problem: $r, \theta_1, \theta_2, \Delta\Phi$

- r and also $J = |\mathbf{L} + \mathbf{S}_1 + \mathbf{S}_2|$ vary on t_{RR}
- Effective spin is constant (at least) at 2PN Racine 2008

Spin precession is (actually) a **1D** problem!

DG+ 2015

...and it's analytically tractable!

DG+ 2023

$$\delta\chi = \frac{\chi_1 \hat{\mathbf{S}}_1 - q\chi_2 \hat{\mathbf{S}}_2}{1+q} \cdot \hat{\mathbf{L}}$$

These 1D solutions are now at the heart of the “twisting up” procedure in some of the modern precessing GW waveforms

Averaging the average



Now GW emission

- Quasi-adiabatic approach
- Only r and J vary on t_{RR}



Usual **orbit** average

$$\langle X \rangle_{\text{orb}} = \frac{\int d\psi \, X \, dt/d\psi}{\int d\psi \, dt/d\psi}$$

Some parameters for the dynamics (say Kepler's true anomaly) →

→ Orbital period

Now a **precession** average

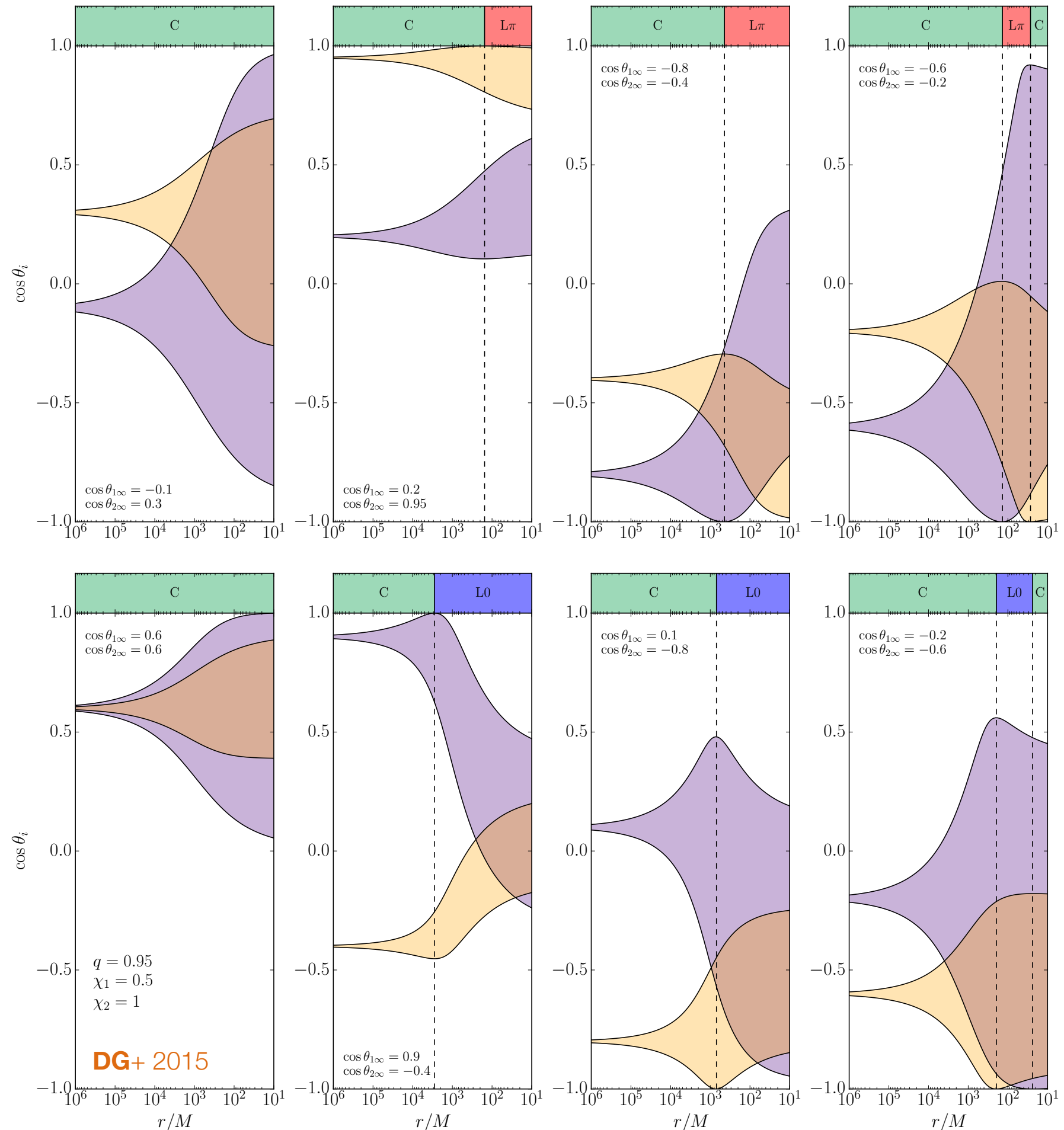
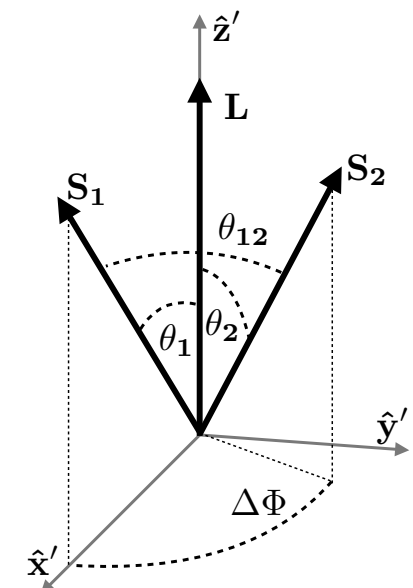
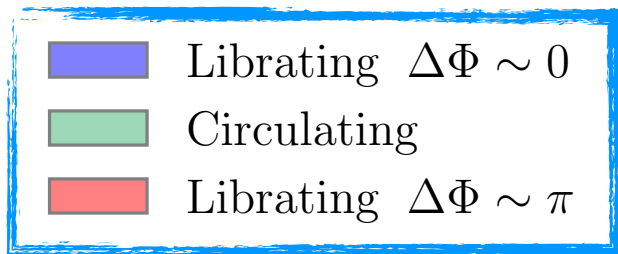
$$\langle X \rangle_{\text{pre}} = \frac{\int d\delta\chi \, \langle X \rangle_{\text{orb}} \, dt/d\delta\chi}{\int d\delta\chi \, dt/d\delta\chi}$$

Dynamics is now parametrized by $\delta\chi$ →

→ Precession period

Precession-averaged evolutions

- Only way evolve from infinitely large separations
- Takes seconds!
- Now LVK reports spins at infinity with this formalism



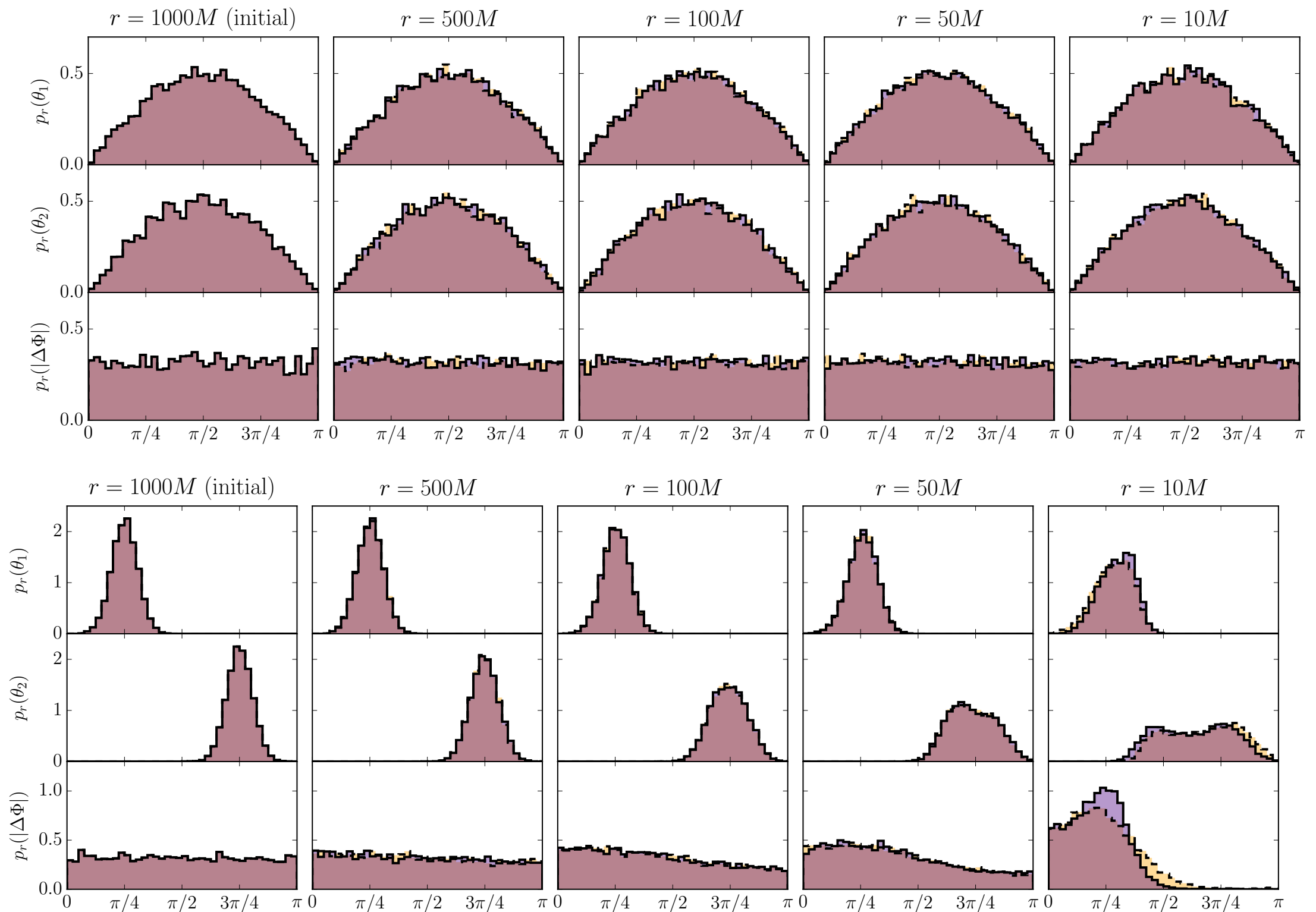
(Updates on this business in the very recent: **DG, Foroni+ 2025**)

It's accurate!

$$q = 0.7$$

$$\chi_1 = 0.8$$

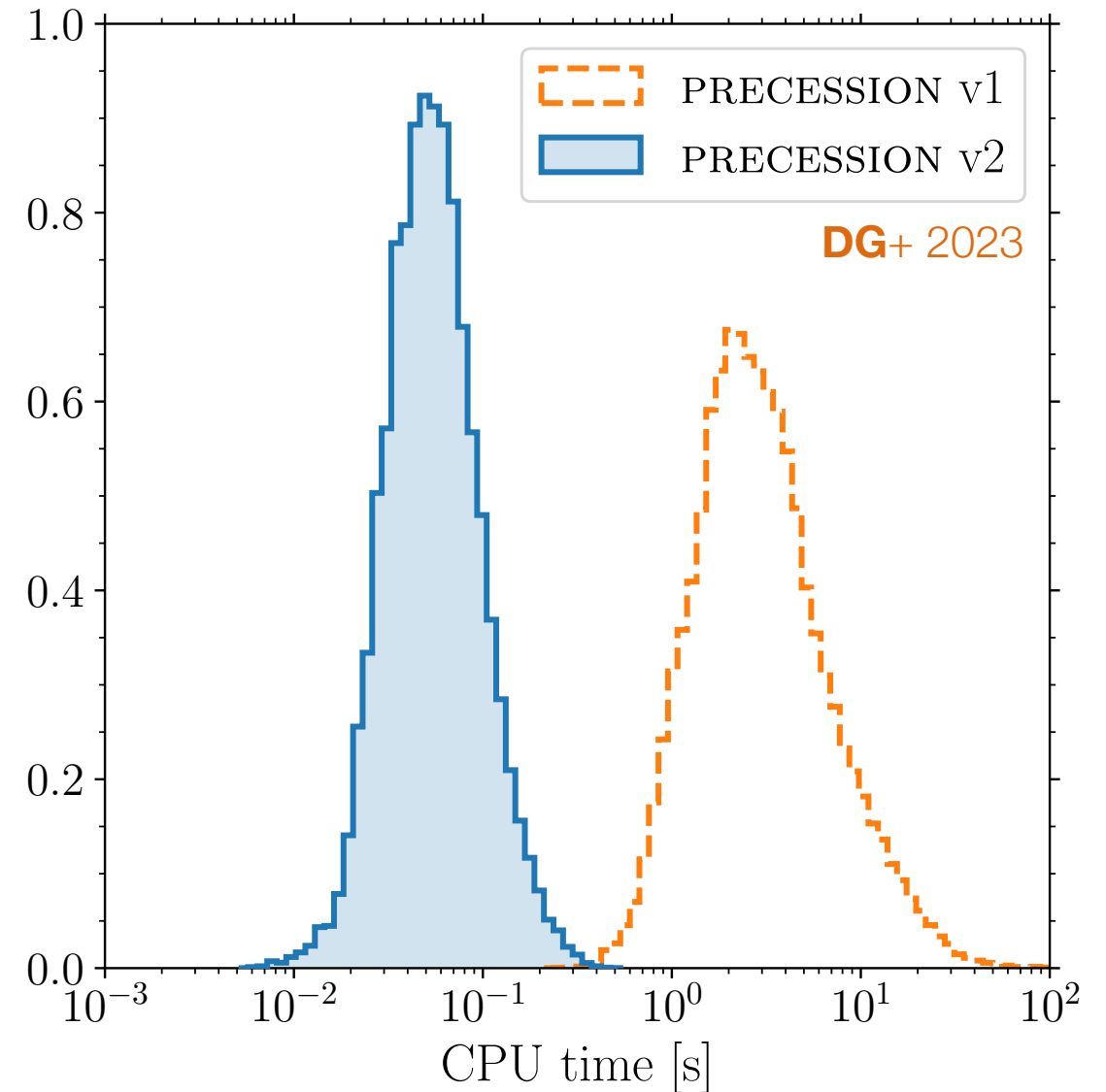
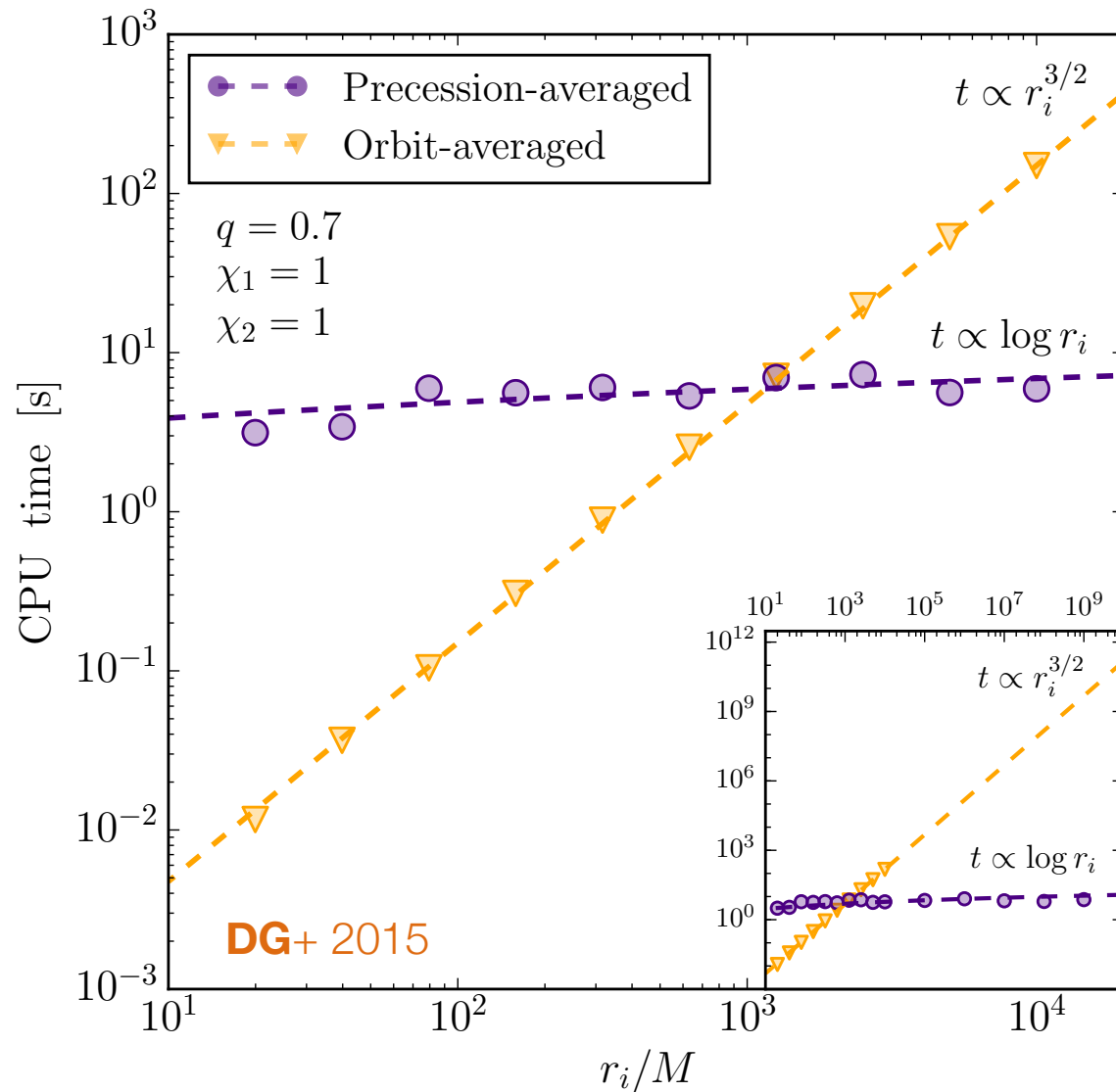
$$\chi_2 = 0.4$$



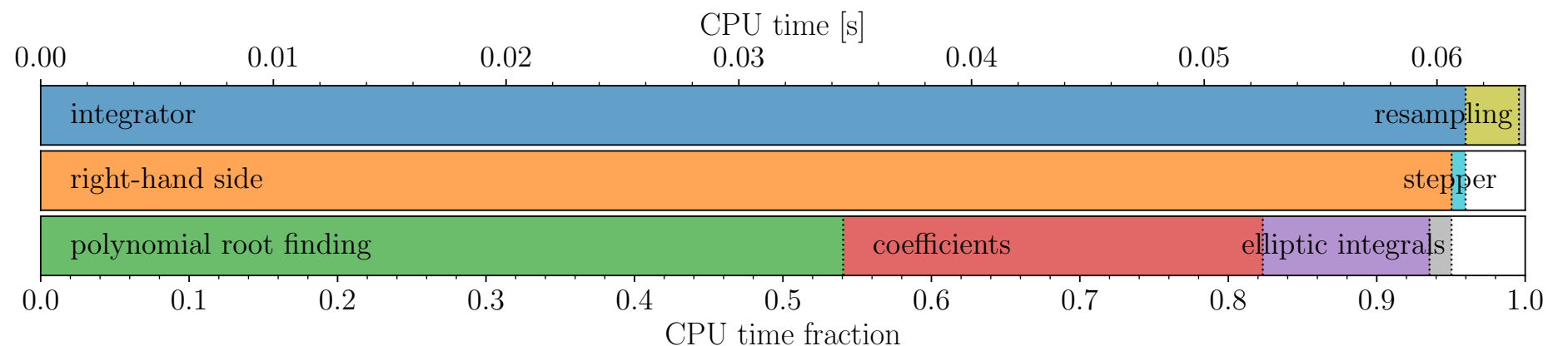
* = We might have a different definition of “accuracy”...

It's fast!

Incredibly favourable scaling with the separation
(can actually integrate all the way to/from infinity!)



Current performance is $O(0.1\text{ s})$ on a laptop,
but one can do better



Try this at home

precession: open-source python module

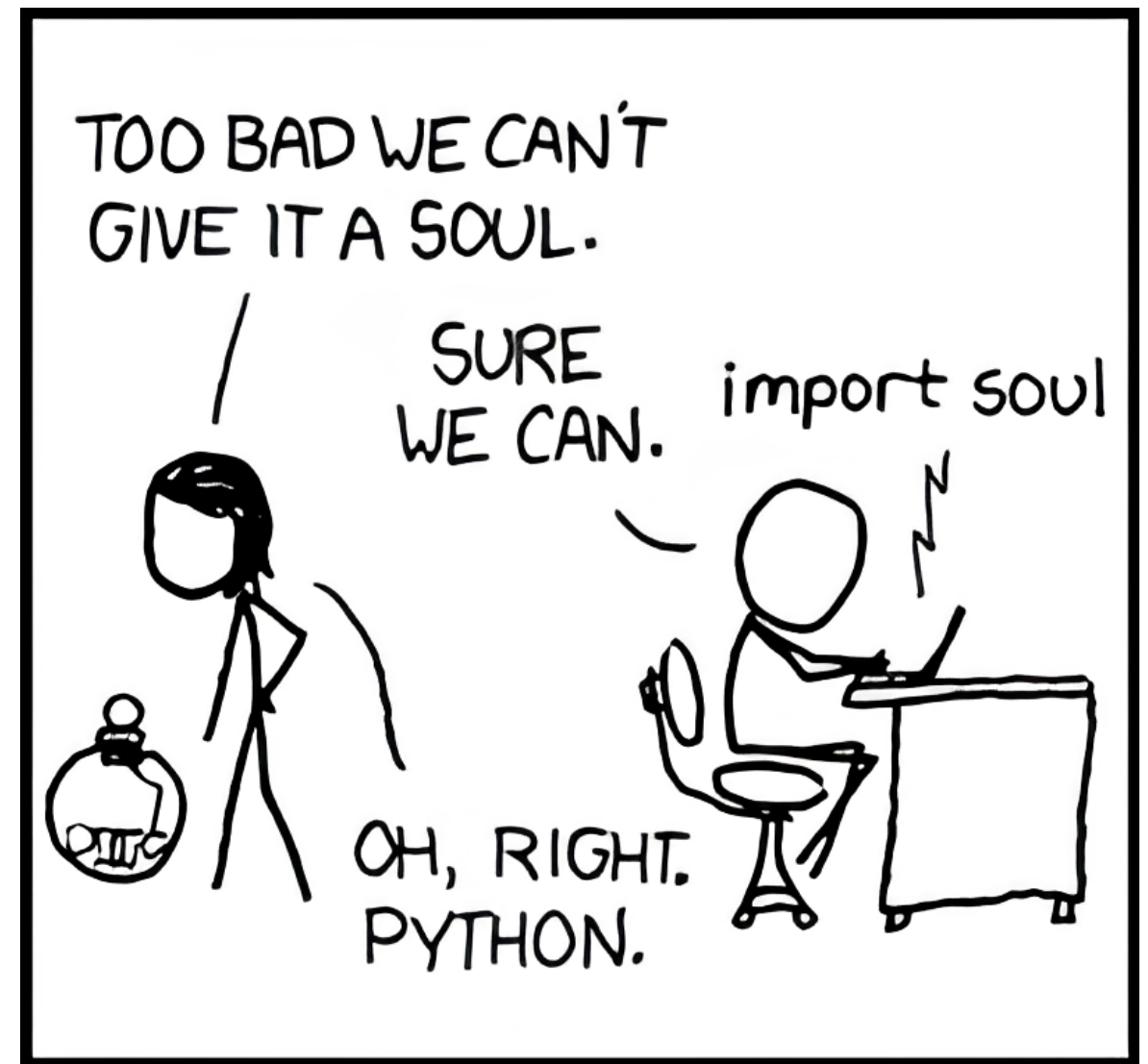
DG+ 2015, DG+ 2023, Fumagalli DG+ 2023

```
pip install precession  
>>> import precession
```

Features

1. Precessional dynamics
2. Orbit-averaged inspirals
3. Precession-averaged inspirals
4. Remnant predictions
5. API documentation
6. Tests and tutorial

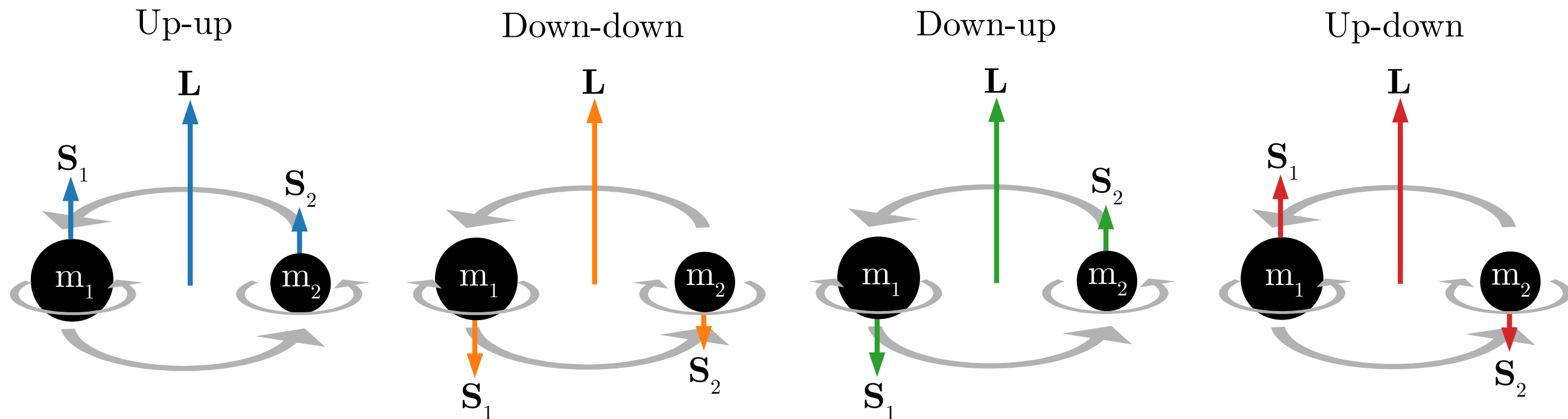
Used in >100 papers to date,
by many people



github.com/dgerosa/precession

Aligned configurations

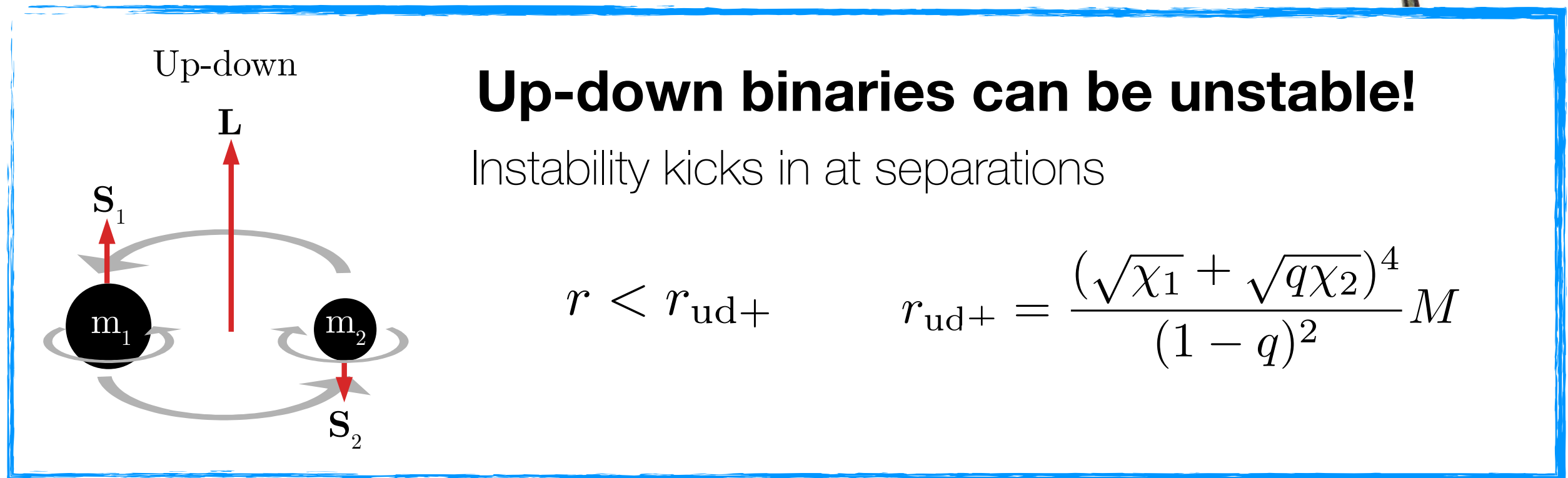
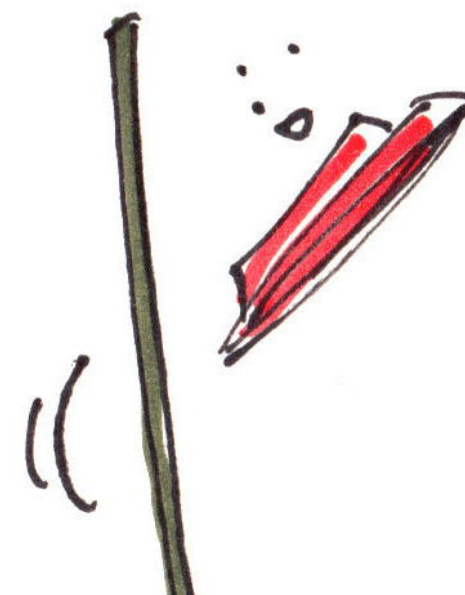
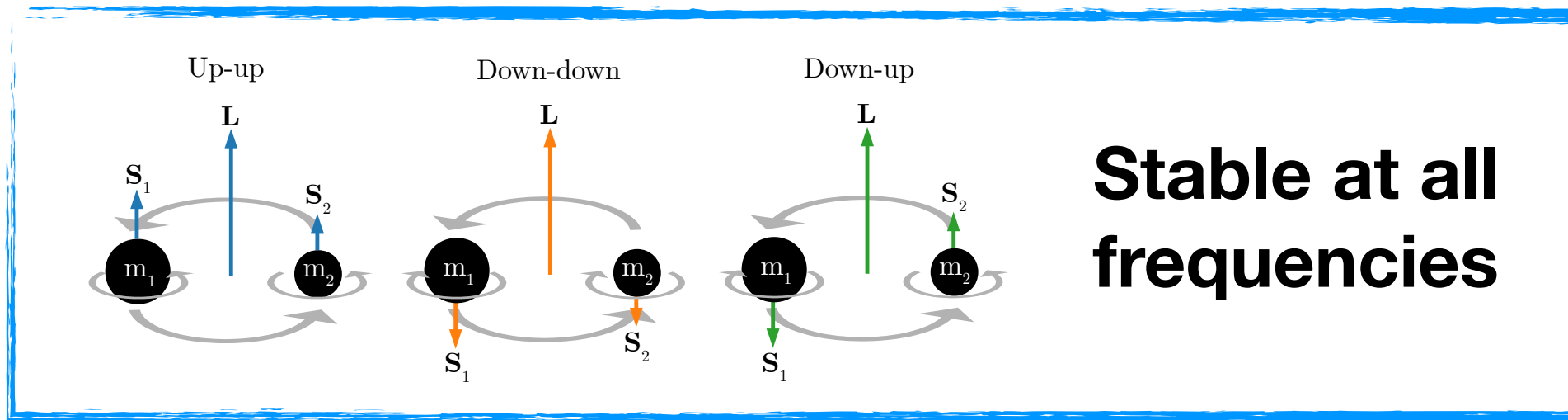
- Astrophysical processes drive binary here (maybe)
- Used as testbed for waveforms developing



All these configurations are solutions of the PN equations at any separation...

are they stable?

New precessional instability!



- First found using the effective potentials **DG+ 2015b**
- Then with explicit perturbation theory (still PN) **Mould DG 2020, Lousto Healy 2016**
- And then even with full numerical relativity! **Varma Mould DG 2021**
- Injections into LVK noise **De Renzi DG 2022**

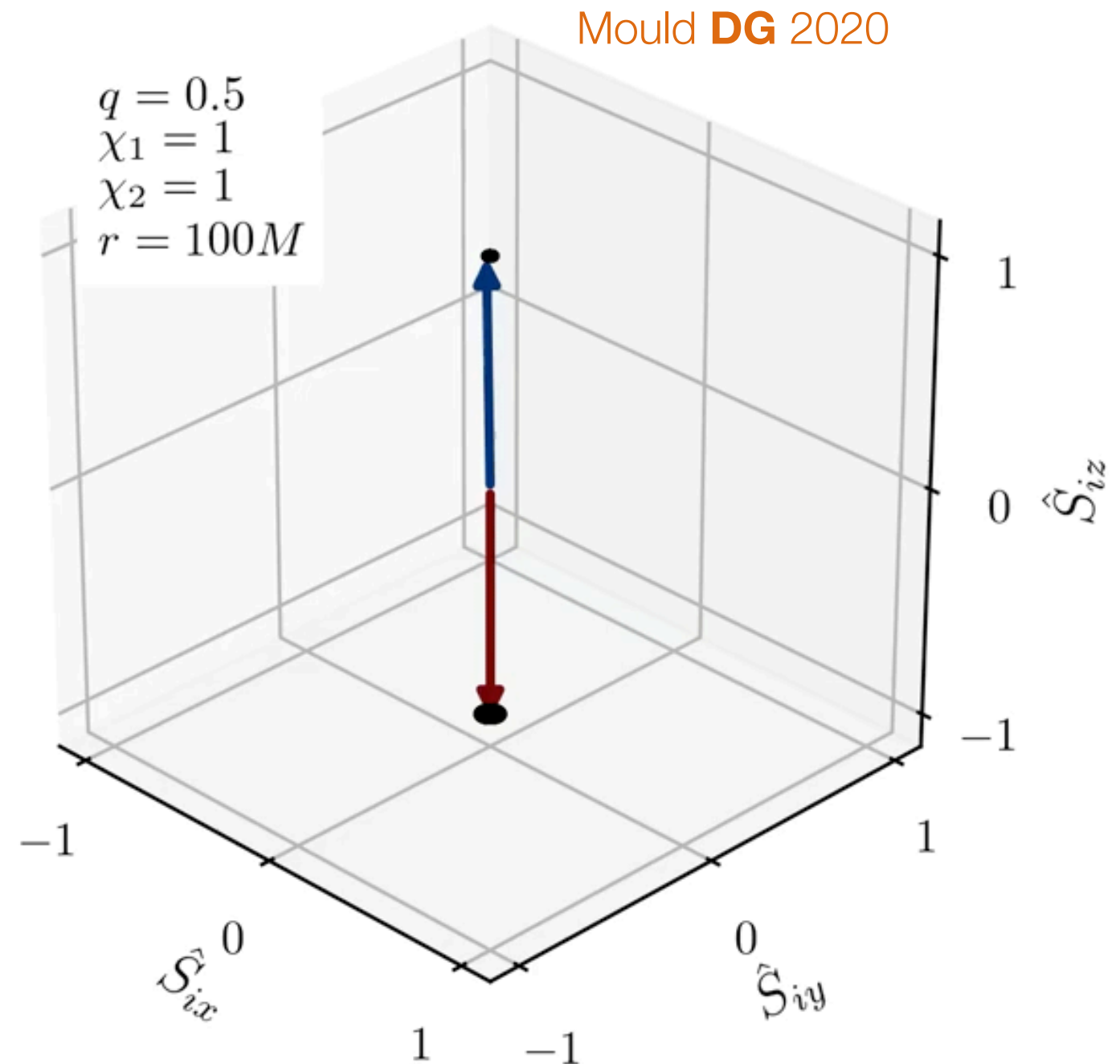
Observational predictions

- Binary likely underwent the instability before entering the LIGO band
- Mechanisms to deplete the aligned region
- Well defined endpoint

$$\hat{\mathbf{S}}_1 = \hat{\mathbf{S}}_2$$

$$\hat{\mathbf{S}}_1 \cdot \hat{\mathbf{L}} = \hat{\mathbf{S}}_2 \cdot \hat{\mathbf{L}} = \frac{\chi_1 - q\chi_2}{\chi_1 + q\chi_2}$$

- But how about binaries that start precessing while being observed?!



What's up with χ_p ?

precession

radiation reaction

$$\frac{d\mathbf{L}}{dt} = \frac{d\hat{\mathbf{L}}}{dt}L + \frac{dL}{dt}\hat{\mathbf{L}} = (\boldsymbol{\Omega}_L \times \hat{\mathbf{L}})L + \frac{dL}{dt}\hat{\mathbf{L}}$$

2PN spin precession equations:

$$\boldsymbol{\Omega}_L = \Omega_1 \chi_1 \hat{\mathbf{S}}_1 + \Omega_2 \chi_2 \hat{\mathbf{S}}_2$$

$$\Omega_1 = \frac{M^2}{2r^3(1+q)^2} \left[4 + 3q - \frac{3q\chi_{\text{eff}}}{(1+q)} \frac{M^2}{L} \right]$$

$$\Omega_2 = \frac{qM^2}{2r^3(1+q)^2} \left[4q + 3 - \frac{3q\chi_{\text{eff}}}{(1+q)} \frac{M^2}{L} \right]$$

$$\left| \frac{d\hat{\mathbf{L}}}{dt} \right|^2 = (\Omega_1 \chi_1 \sin \theta_1)^2 + (\Omega_2 \chi_2 \sin \theta_2)^2 + 2\Omega_1 \Omega_2 \chi_1 \chi_2 \sin \theta_1 \sin \theta_2 \cos \Delta\Phi$$

Schmidt+ 2015 Extremize $\cos \Delta\Phi = \pm 1$, take an arithmetic mean, and normalize by Ω_1

$$\text{"Heuristic"} \quad \chi_p \equiv \frac{1}{2\Omega_1} \left(\left| \frac{d\hat{\mathbf{L}}}{dt} \right|_+ + \left| \frac{d\hat{\mathbf{L}}}{dt} \right|_- \right) = \max \left(\chi_1 \sin \theta_1, q \frac{4q+3}{4+3q} \chi_2 \sin \theta_2 \right)$$

But:

- Often $\cos \Delta\Phi = \pm 1$ is not allowed geometrically
- The three angles $\theta_1, \theta_2, \Delta\Phi$ vary on the same timescale!

Two strategies to fix this

The common χ_p expression is a partial average over the precession timescale

Two things one can do:

1. Retain all variations on the precessional timescale. “Generalized” χ_p

$$\chi_p \equiv \left| \frac{d\hat{\mathbf{L}}}{dt} \right| \frac{1}{\Omega_1} = \left[(\chi_1 \sin \theta_1)^2 + \left(q \frac{4q+3}{4+3q} \chi_2 \sin \theta_2 \right)^2 + 2q \frac{4q+3}{4+3q} \chi_1 \chi_2 \sin \theta_1 \sin \theta_2 \cos \Delta\Phi \right]^{1/2}$$

same terms as before

this is new!

2. Precession-average them all “Averaged” $\langle \chi_p \rangle = \frac{\int \chi_p(\psi) \left(\frac{d\psi}{dt} \right)^{-1} d\psi}{\int \left(\frac{d\psi}{dt} \right)^{-1} d\psi}$

A better precession parameter

DG+ 2021

Replace $\chi_p = \max \left(\chi_1 \sin \theta_1, q \frac{4q+3}{4+3q} \chi_2 \sin \theta_2 \right) \in [0, 1]$

with $\chi_p = \left[(\chi_1 \sin \theta_1)^2 + \left(q \frac{4q+3}{4+3q} \chi_2 \sin \theta_2 \right)^2 + 2q \frac{4q+3}{4+3q} \chi_1 \chi_2 \sin \theta_1 \sin \theta_2 \cos \Delta \Phi \right]^{1/2} \in [0, 2]$

and average consistently $\langle \chi_p \rangle$

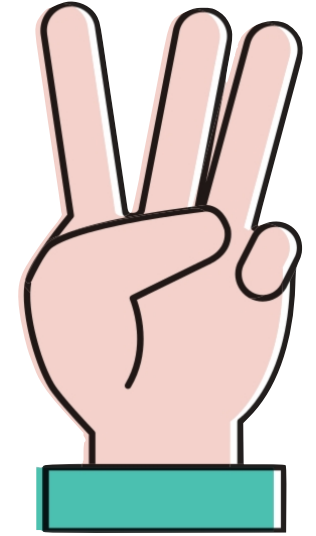
The $\chi_p \in [1, 2]$ region is exclusive to binaries with two precessing spins

$\chi_p > 0$ Evidence for **spin precession**

$\chi_p > 1$ Evidence for **two-spin effects**



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More spin things!

Spins and astrophysics Tidal interactions, supernova kicks, mass transfer accretion **DG+** 2013, 2018 and infinitely many more

Spins and waveforms Our solutions now somewhere in PhenomX
Chatziioannou+ 2017, Khan+ 2019, Pratten + 2023

Spins and population stats Need to go back to past time infinity!
Now LVK GWTC4 catalog quoting spins back in time Mould, **DG+** 2022

Spins and eccentricity Extremely interesting interplay
(notion of “effective circular binary”) Fumagalli, **DG+** 2023, 2024

See Giulia's talk...

Spins and three body Clustering in the orbital plane
e.g. Antonini+ 2017

Spins and tides What happens with neutron stars?
LaHaye+ 2022

Spin and kicks Big kicks only come from spins!
Rodrigues+ 2007, Campanelli+ 2007, **DG+** 2018

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