
GAMMA-RAY BURSTS: AFTERGLOW EMISSION

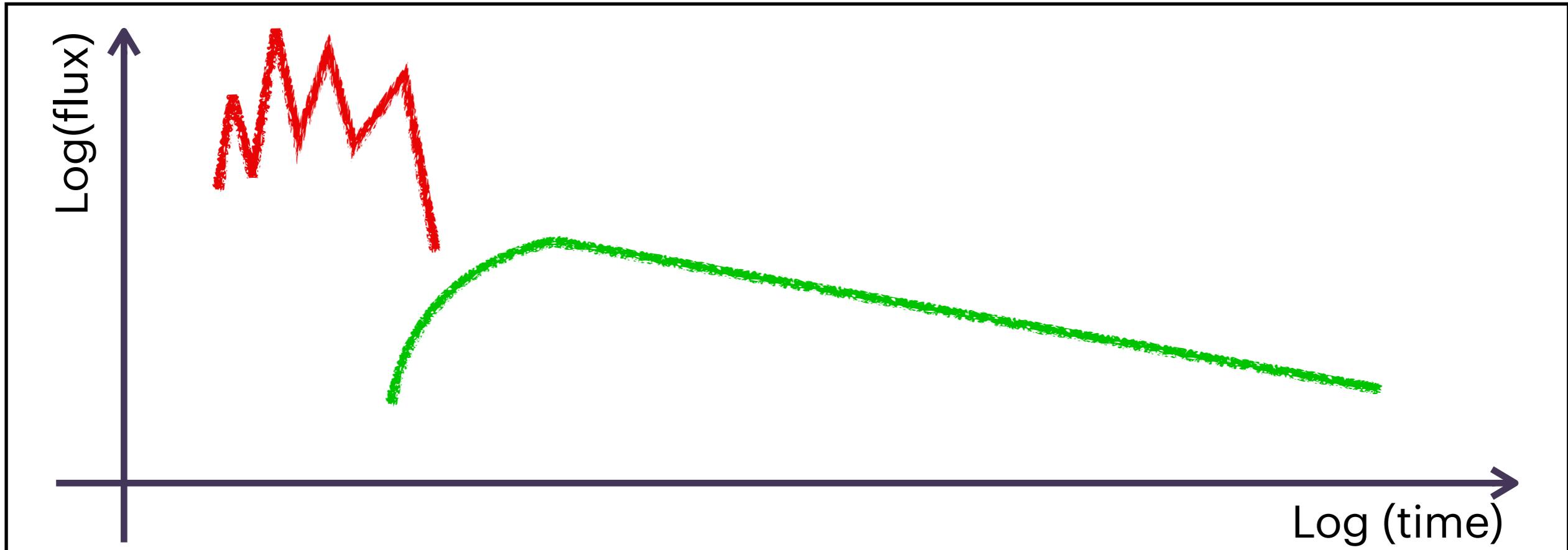


Lara Nava
INAF
Osservatorio Astronomico di Brera



EMISSION FROM GRBS: AN OVERVIEW

TWO DISTINCT EMISSION PHASES



PROMPT

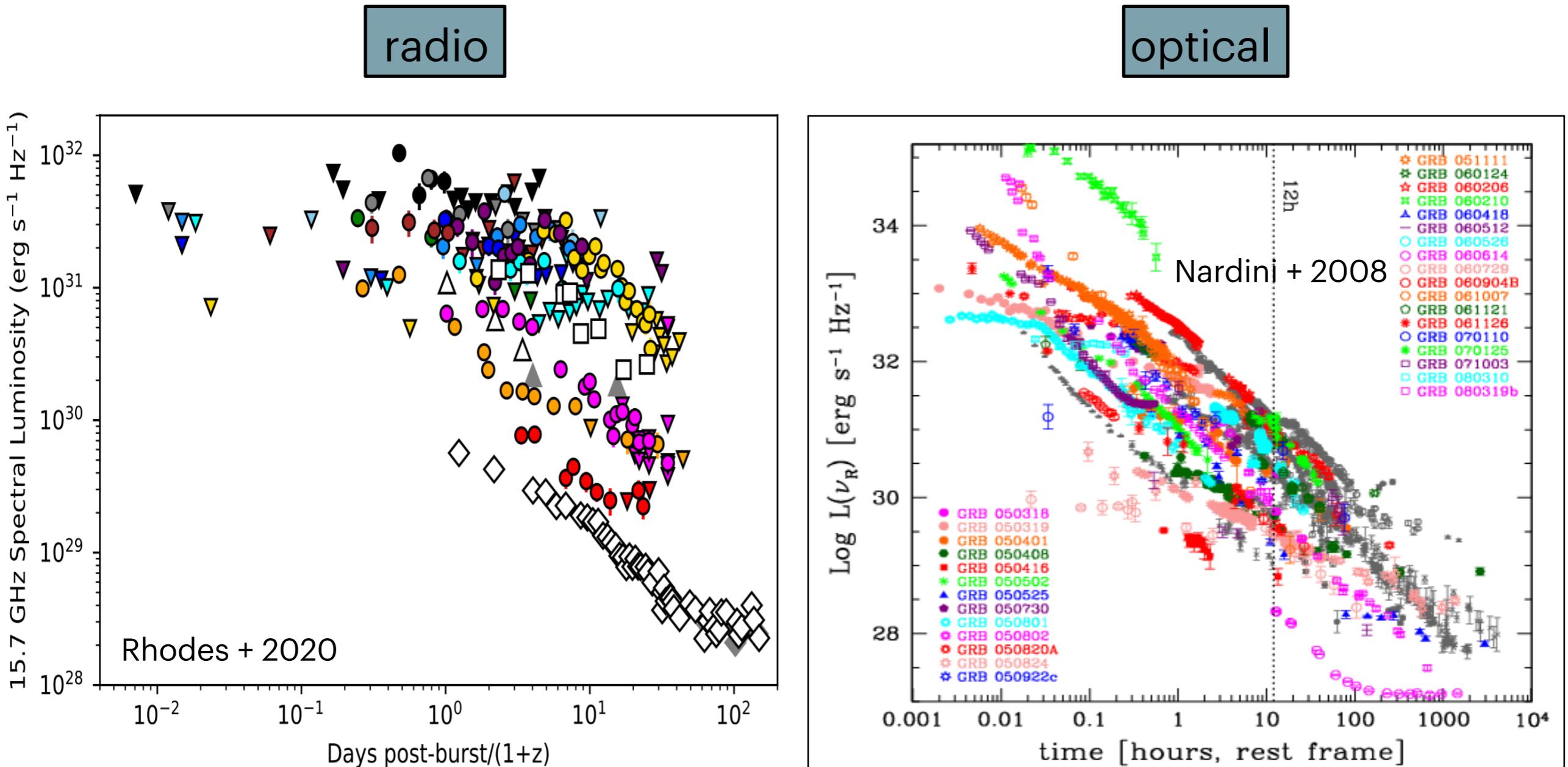
- from 10 keV to 10 MeV
- non-thermal spectra
- 0.1 seconds to 10^3 seconds
- highly variable flux

AFTERGLOW

- from radio to TeV
- non-thermal spectra
- days - weeks
- smooth (PL) lightcurve

AFTERGLOW EMISSION LIGHTCURVES

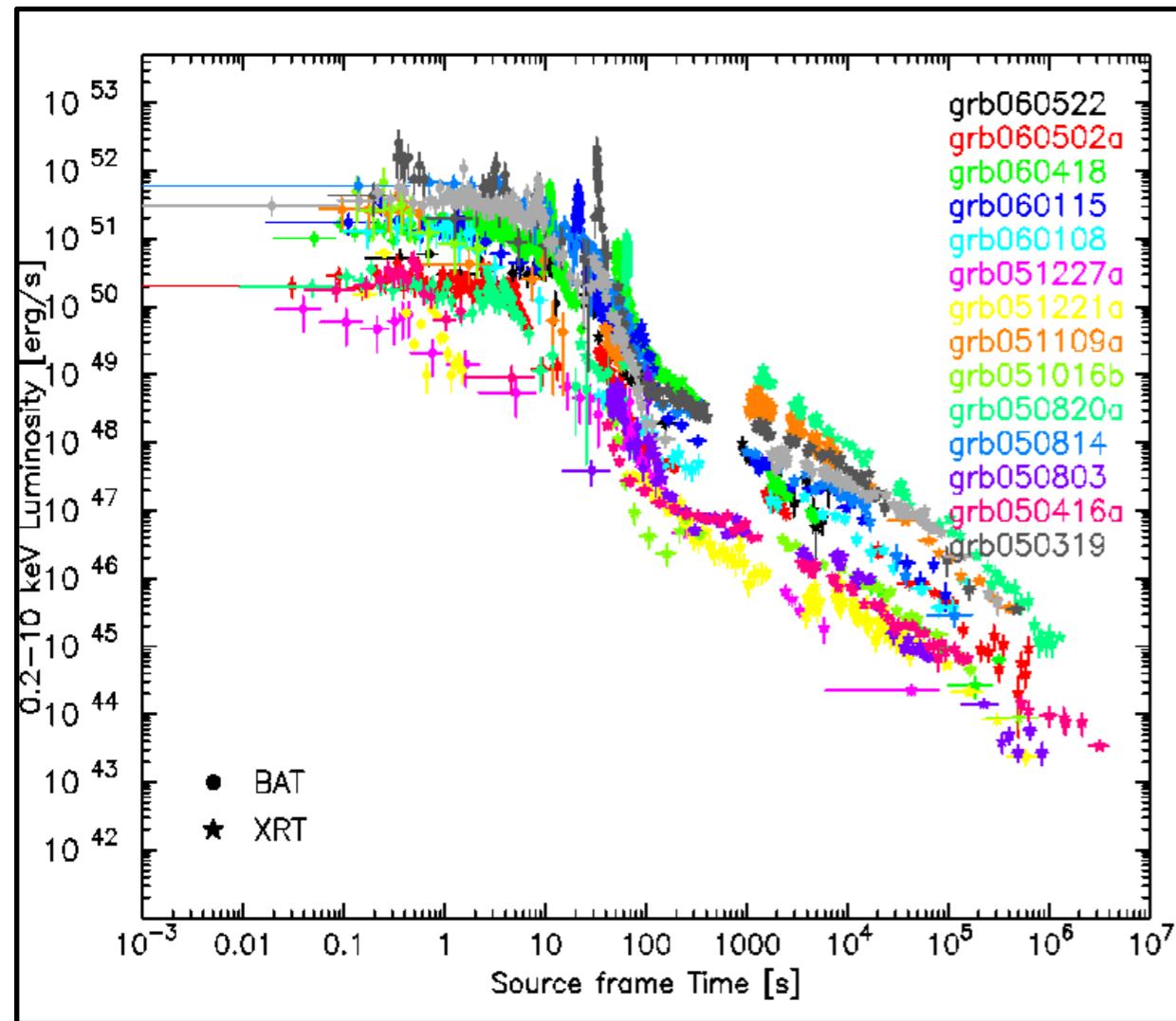
OBSERVATIONS AT DIFFERENT FREQUENCIES



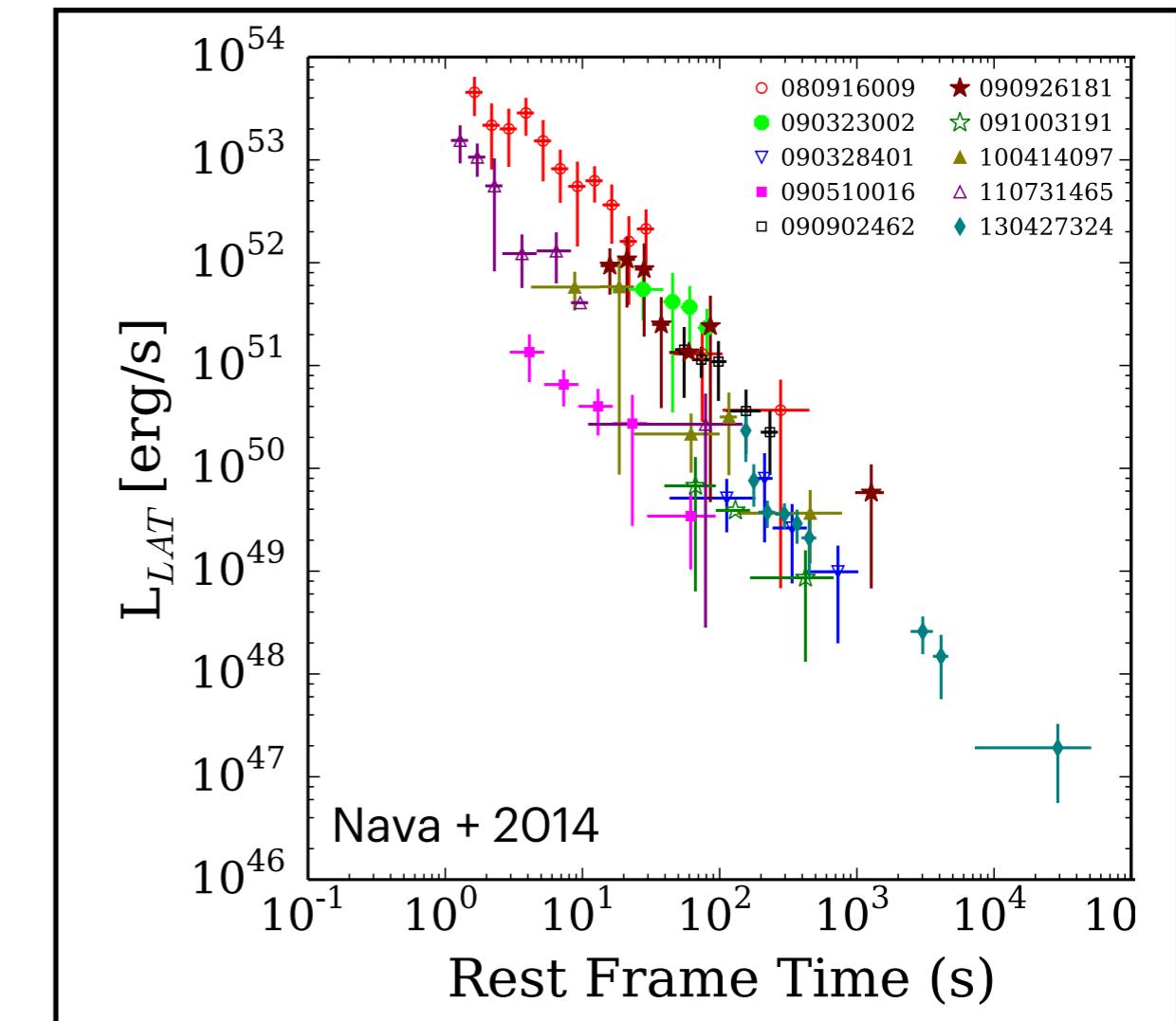
AFTERGLOW EMISSION LIGHTCURVES

OBSERVATIONS AT DIFFERENT FREQUENCIES

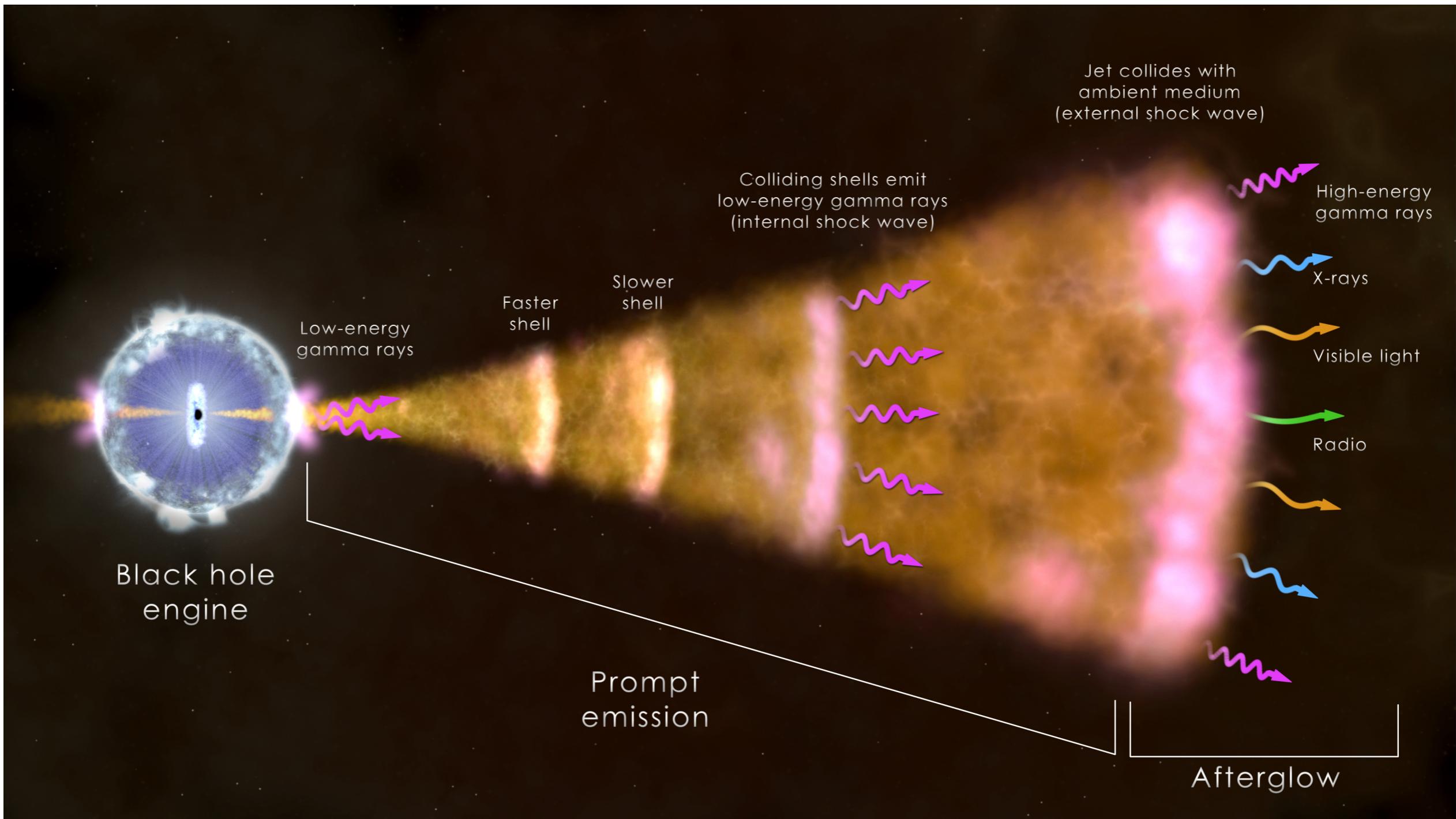
X-ray



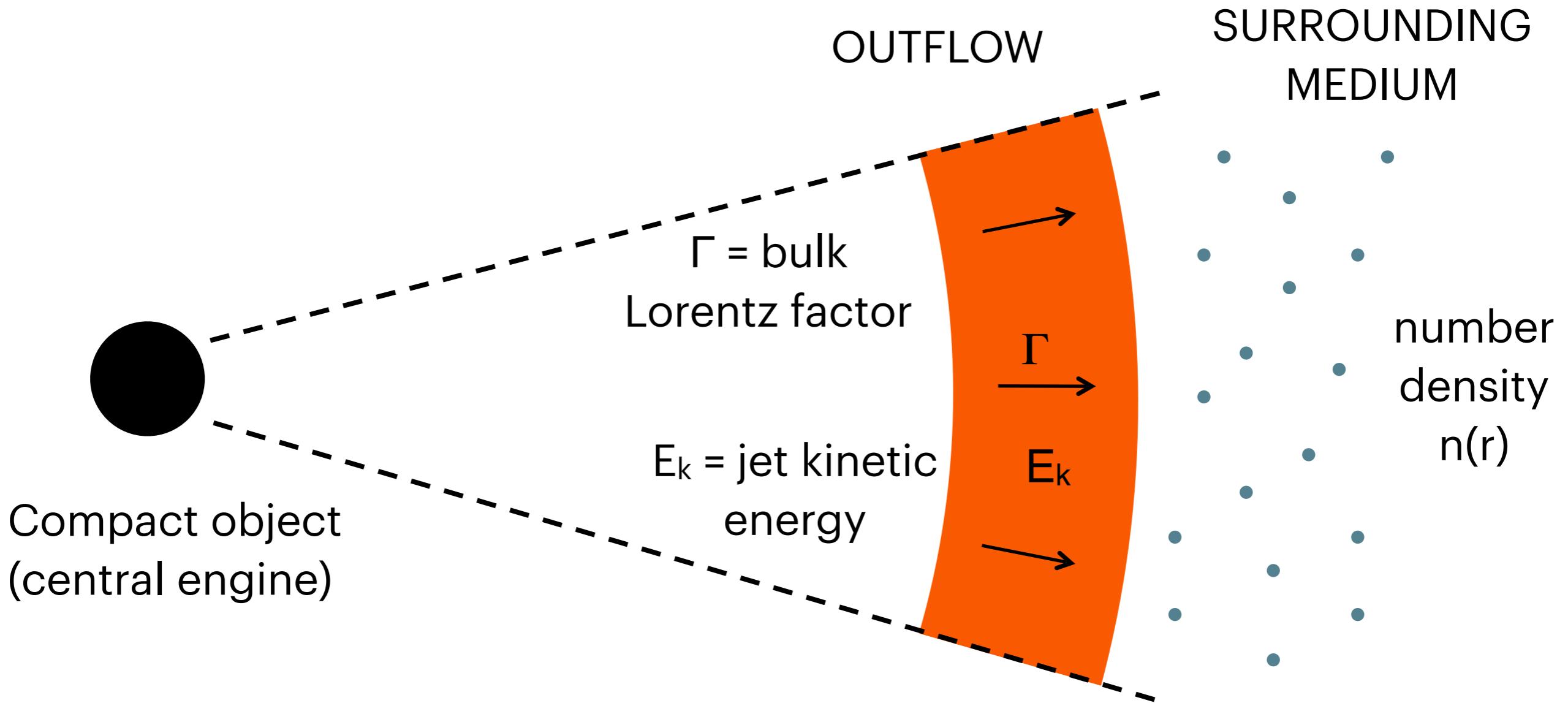
GeV (0.1-10 GeV)



GAMMA-RAY BURSTS - THE STANDARD MODEL



AFTERGLOW EMISSION: EXTERNAL SHOCK MODEL



AFTERGLOW EMISSION

To predict and model the radiative output, it is necessary to model the following three processes:

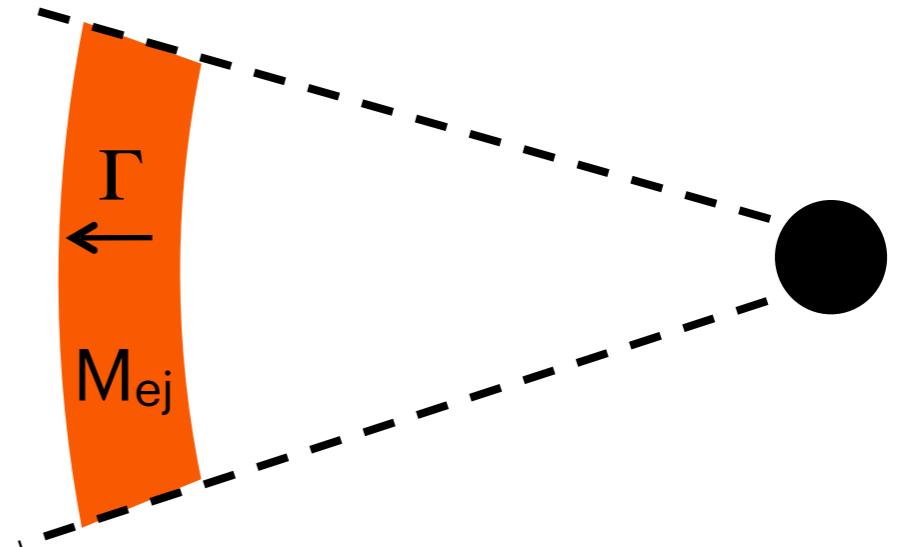
- Dynamics
how Γ evolves as a function of distance r and time
- Relativistic collisionless shocks
particle acceleration and B amplification
- Radiative process(es)
how the radiative output is produced

DYNAMICS

Collision between jet of kinetic energy $E_k = \Gamma_0 M_{ej} c^2$ and the external medium with number density $n(r) = A r^{-s}$

At distance r , the collected mass is

$$m(r) = \int 4\pi r^2 n(r) m_p dr \propto r^{3-s}$$



Conservation of energy and momentum

BEFORE COLLISION

$$\Gamma_0 M_{ej} c^2 + m(r)c^2$$

$$\Gamma_0 \beta_0 M_{ej} c$$

AFTER THE COLLISION

$$\Gamma(r) [M_{ej} + m(r) + \epsilon'(r)/c^2] c^2$$

$$\Gamma(r) \beta(r) [M_{ej} + m(r) + \epsilon'(r)/c^2] c$$

DYNAMICS

$$\Gamma(r) = \frac{\Gamma_0 M_{ej} + m(r)}{\sqrt{m^2(r) + 2\Gamma_0 M_{ej} m(r) + M_{ej}^2}}$$

- $m(r) < \frac{M_{ej}}{2\Gamma_0}$ $\Gamma(r) = \Gamma_0$ coasting phase
- $\frac{M_{ej}}{\Gamma_0} < m(r) < M_{ej}\Gamma_0$ $\Gamma(r) \simeq \sqrt{\frac{E_k}{m(r)}}$ deceleration phase
see Blandford & McKee 1976
- $m(r) > M_{ej}\Gamma_0$ $\Gamma(r) \rightarrow 1$ non-relativistic phase

DYNAMICS

A closer look to the deceleration phase
see Blandford & McKee 1976

$$\frac{M_{ej}}{\Gamma_0} < m(r) < M_{ej}\Gamma_0$$

→

$$\Gamma(r) \simeq \sqrt{\frac{E_k}{m(r)}} \propto r^{-\frac{(3-s)}{2}}$$

Given the relation
between radius
and time:

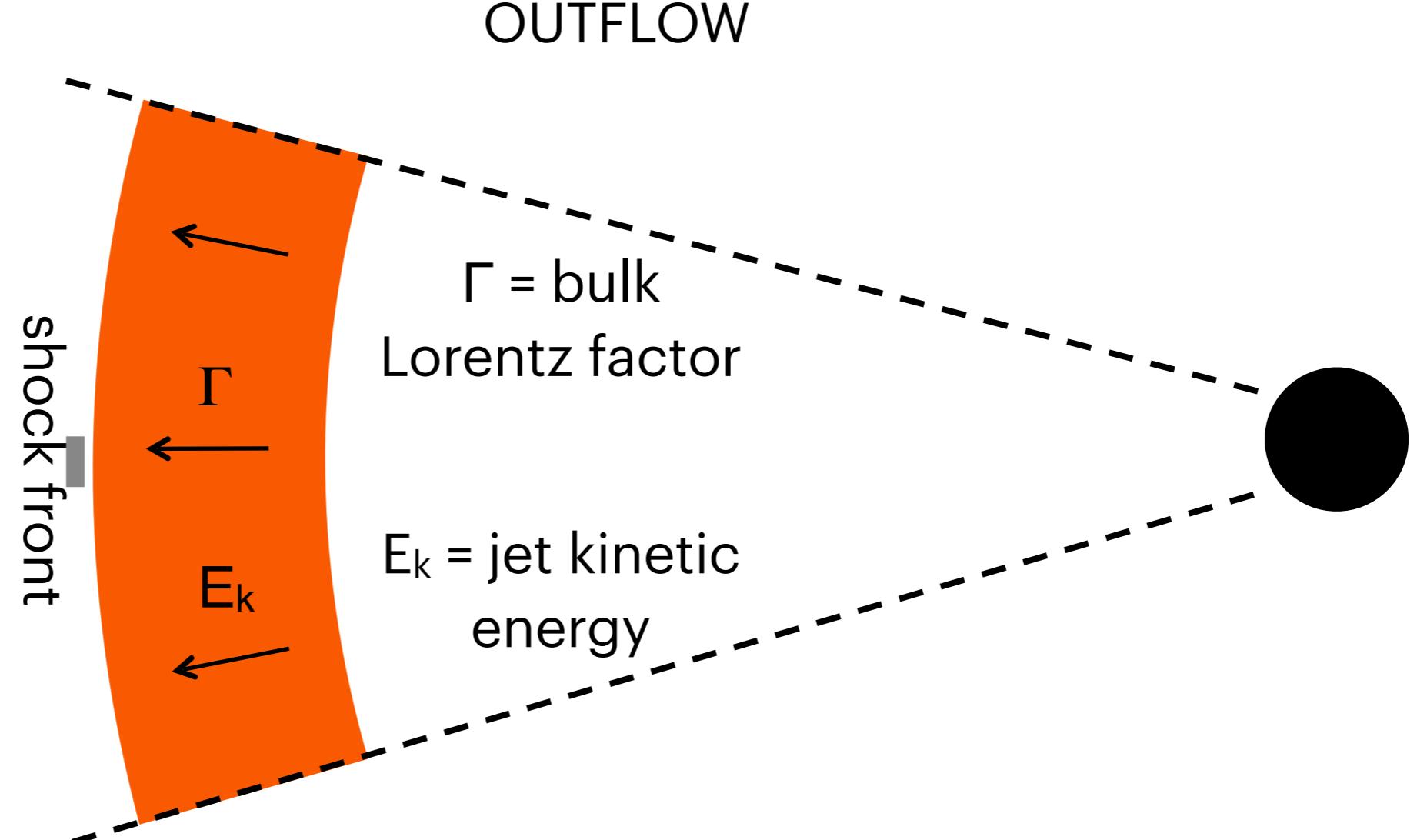
$$t_{obs} \simeq (1+z) \frac{r}{c\Gamma^2} \quad \rightarrow$$

$$\Gamma(t_{obs}) \propto r^{\frac{3-s}{8-2s}}$$

RELATIVISTIC SHOCKS

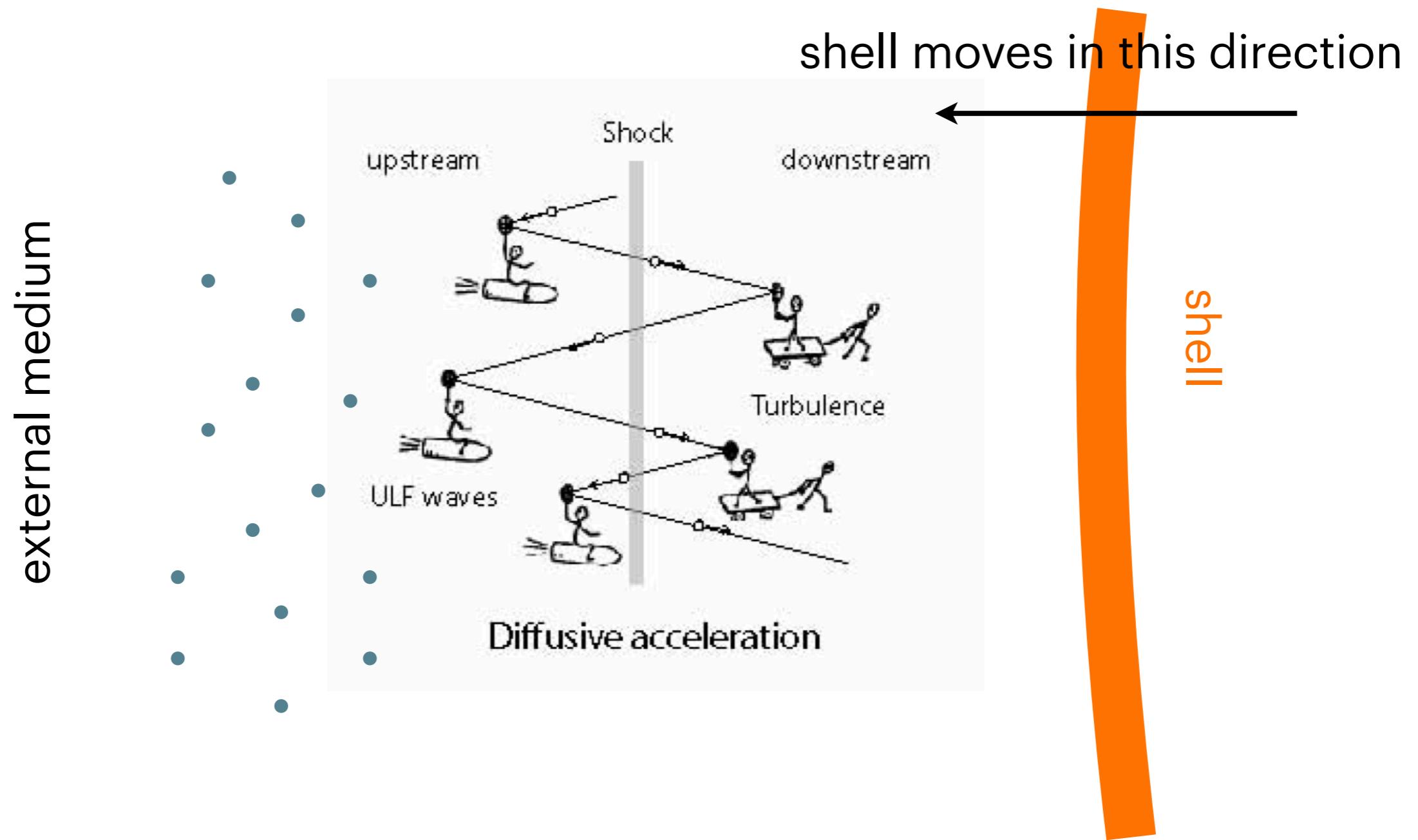
SURROUNDING
MEDIUM

number
density
 $n(r)$



RELATIVISTIC SHOCKS

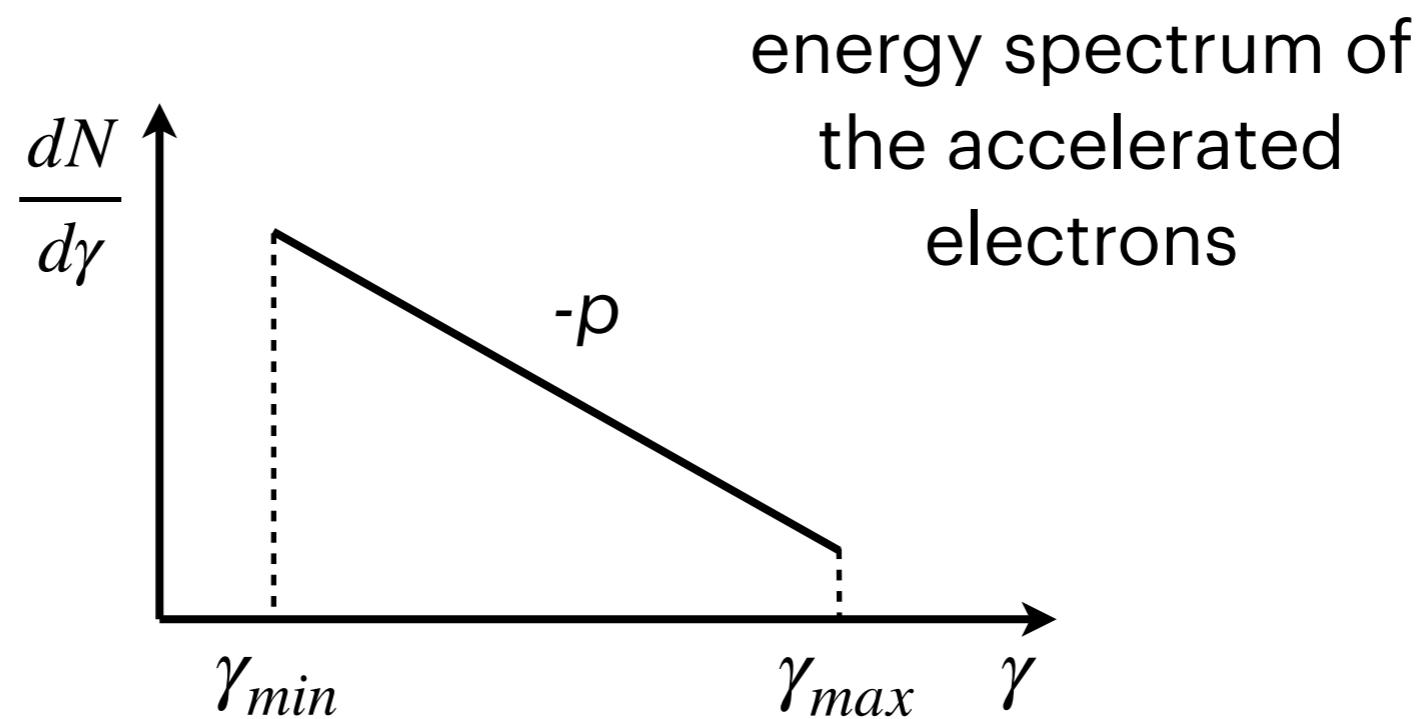
DIFFUSIVE SHOCK ACCELERATION (FERMI MECHANISM)



RELATIVISTIC SHOCKS

PARTICLE ACCELERATION

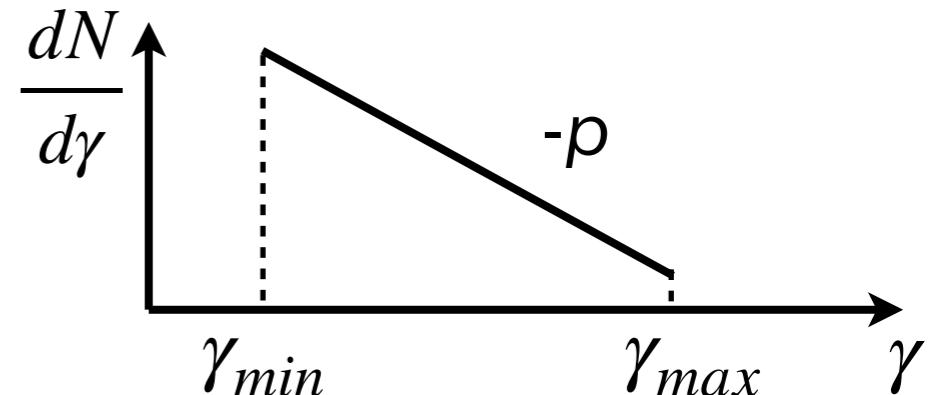
A fraction of the dissipated energy is used to accelerate particles through **collisionless** (no Coulomb collision) **shocks**. The output of this acceleration process is a power-law particle spectrum with index $p = 2.2 - 2.4$



RELATIVISTIC SHOCKS

PARTICLE ACCELERATION

Derivation of the average and minimum Lorentz factor ($\langle \gamma \rangle$ and γ_{\min}) of the electrons



$$\bullet N_e \langle \gamma \rangle m_e c^2 = \epsilon_e N_p m_p c^2 (\Gamma - 1)$$

$$\Rightarrow \langle \gamma \rangle = \epsilon_e \frac{m_p}{m_e} (\Gamma - 1) \simeq 2 \times 10^4 \epsilon_{e,-1} \Gamma_2$$

$$\bullet \langle \gamma \rangle = \frac{\int_{\gamma_{\min}}^{\gamma_{\max}} \frac{dN}{d\gamma} \gamma d\gamma}{\int_{\gamma_{\min}}^{\gamma_{\max}} \frac{dN}{d\gamma} d\gamma} = \frac{\int_{\gamma_{\min}}^{\gamma_{\max}} \gamma^{-p+1} d\gamma}{\int_{\gamma_{\min}}^{\gamma_{\max}} \gamma^{-p} d\gamma}$$

for $\gamma_{\max} \gg \gamma_{\min}$ and $p > 2$

$$\Rightarrow \gamma_{\min} = \epsilon_e \frac{m_p}{m_e} \frac{p-2}{p-1} (\Gamma - 1)$$

RELATIVISTIC SHOCKS

MAGNETIC FIELD AMPLIFICATION

Derivation of the magnetic field strength

$$\bullet \quad \epsilon_B n' m_p c^2 (\Gamma - 1) = \frac{B'^2}{8\pi}$$

$$n'_p = \frac{N_p}{V'} = \frac{N_p}{V/\Gamma} = n_p \Gamma$$

factor 4 due to shock compression $\Rightarrow n' = 4 \Gamma n$

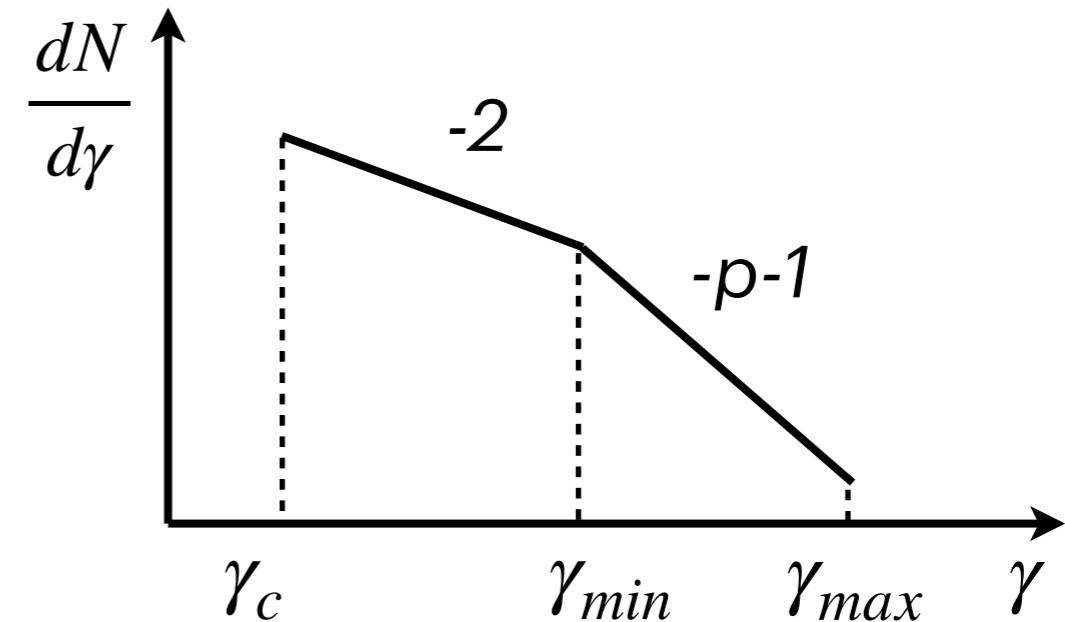
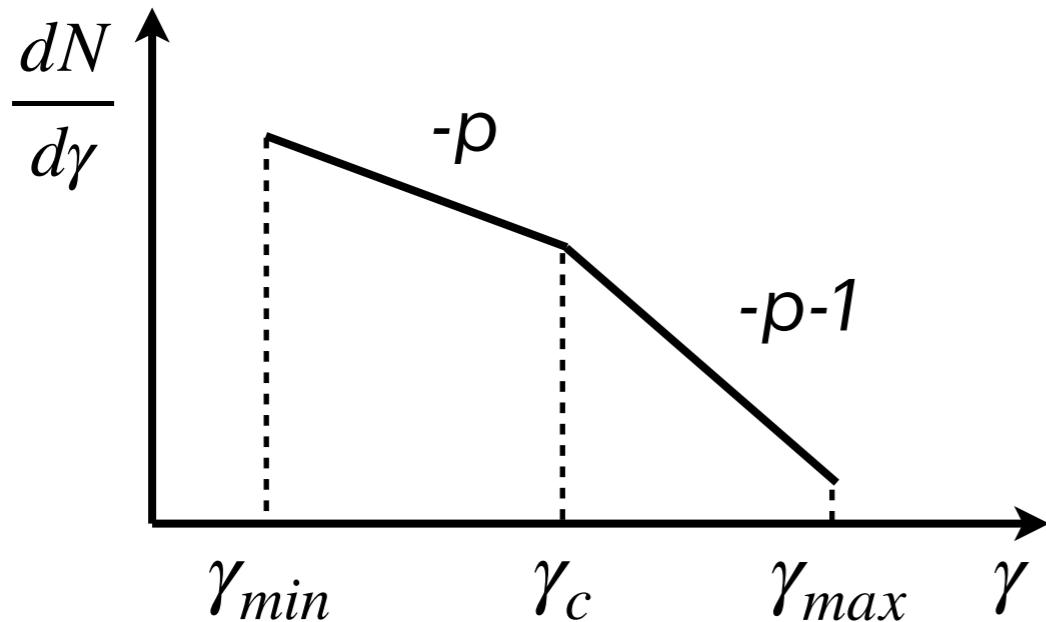
$$\Rightarrow B' = \sqrt{32\pi\epsilon_B n m_p c^2} \Gamma$$

SYNCHROTRON EMISSION

Synchrotron cooling time

$$t'_c(\gamma, B') = \frac{E(\gamma)}{P(\gamma, B')} = \frac{6 \pi m_e c}{\sigma_T \gamma B'^2}$$

$$\gamma_c = \frac{6 \pi m_e c}{\sigma_T t' B'^2}$$



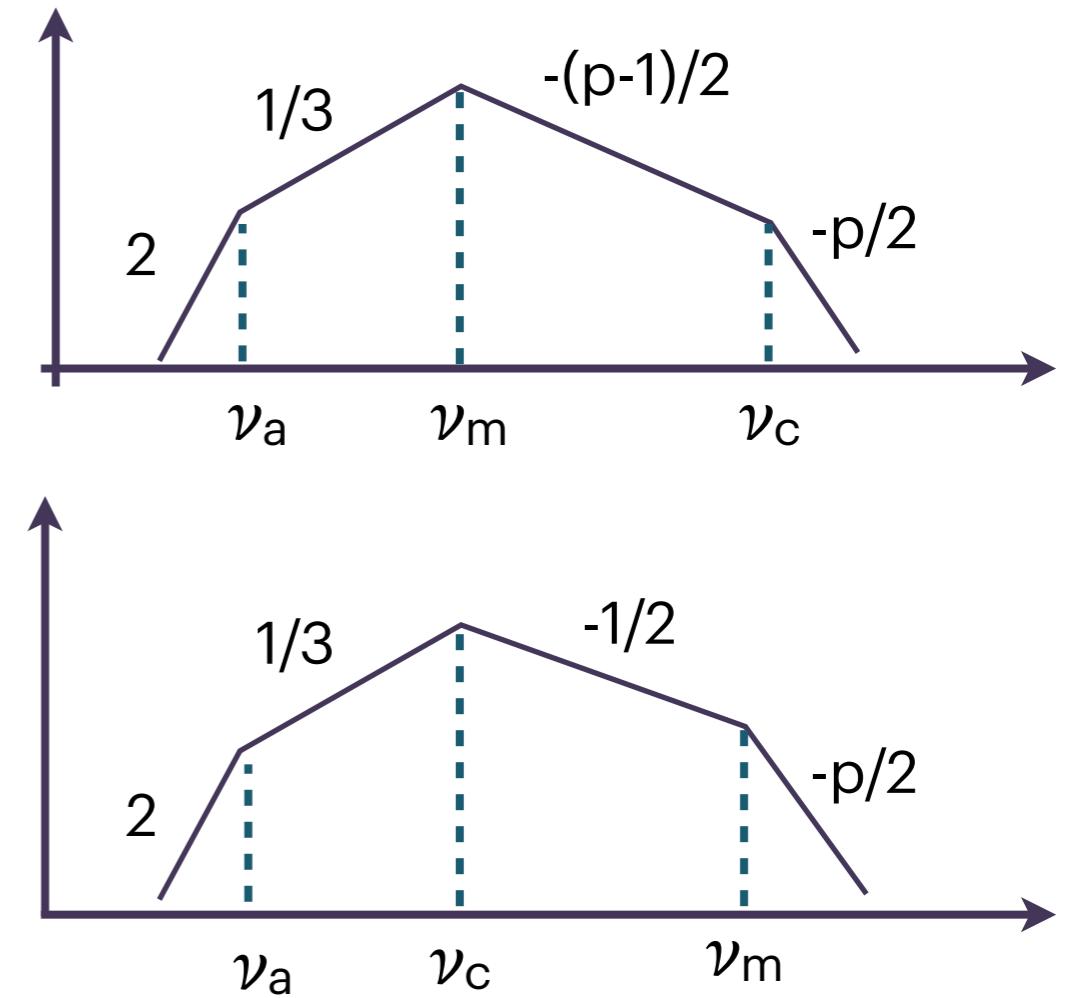
SYNCHROTRON SPECTRUM

SPECTRUM FROM A POPULATION OF ELECTRONS

$$\nu_{syn} = \gamma^2 \frac{q_e B}{2 \pi m_e c} \Gamma$$

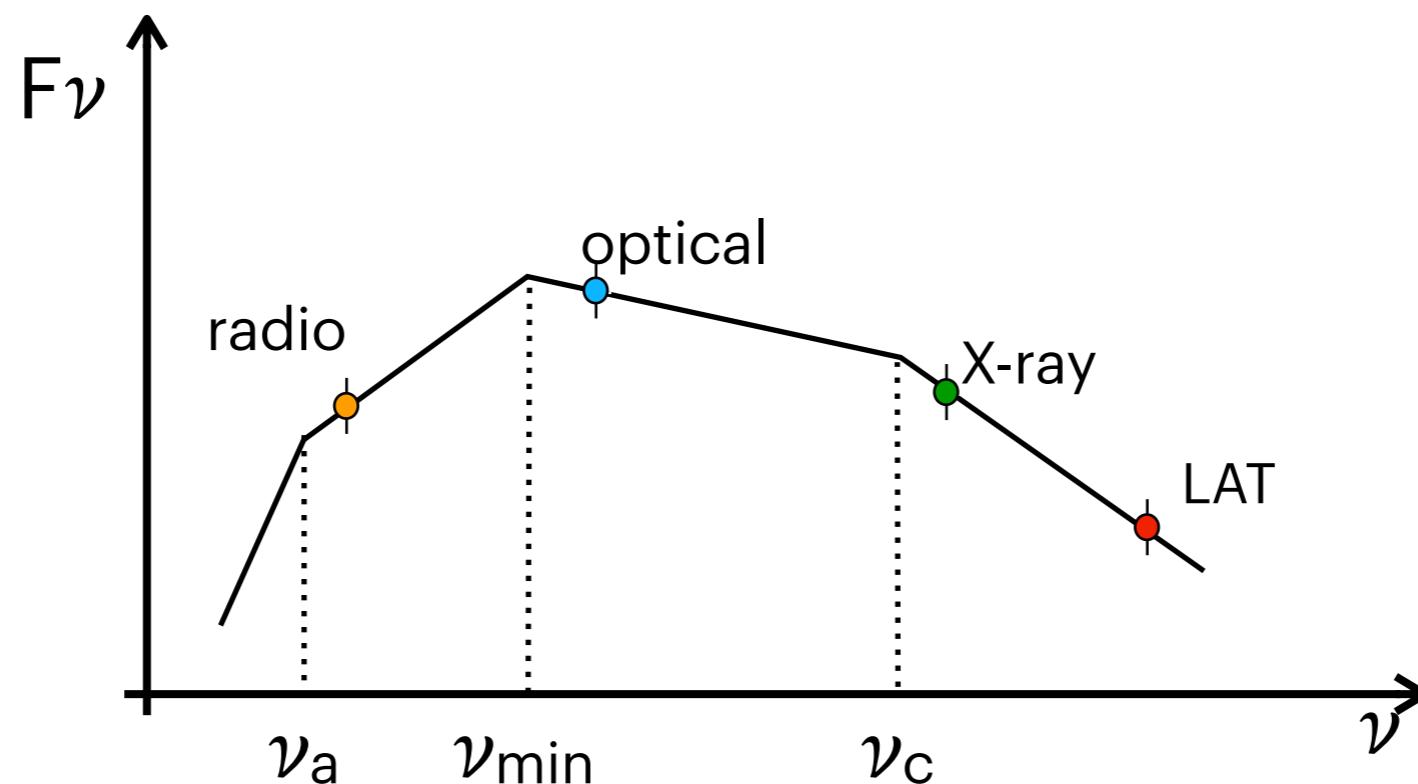
$$t < t_{cool} \Rightarrow F(\nu) \propto \begin{cases} \nu^2 & \nu < \nu_a \\ \nu^{1/3} & \nu_a \leq \nu < \nu_m \\ \nu^{-\frac{p-1}{2}} & \nu_m \leq \nu < \nu_{cool} \\ \nu^{-\frac{p}{2}} & \nu \geq \nu_{cool} \end{cases}$$

$$t > t_{cool} \Rightarrow F(\nu) \propto \begin{cases} \nu^2 & \nu < \nu_a \\ \nu^{1/3} & \nu_a \leq \nu < \nu_{cool} \\ \nu^{-\frac{1}{2}} & \nu_{cool} \leq \nu < \nu_m \\ \nu^{-\frac{p}{2}} & \nu \geq \nu_m \end{cases}$$



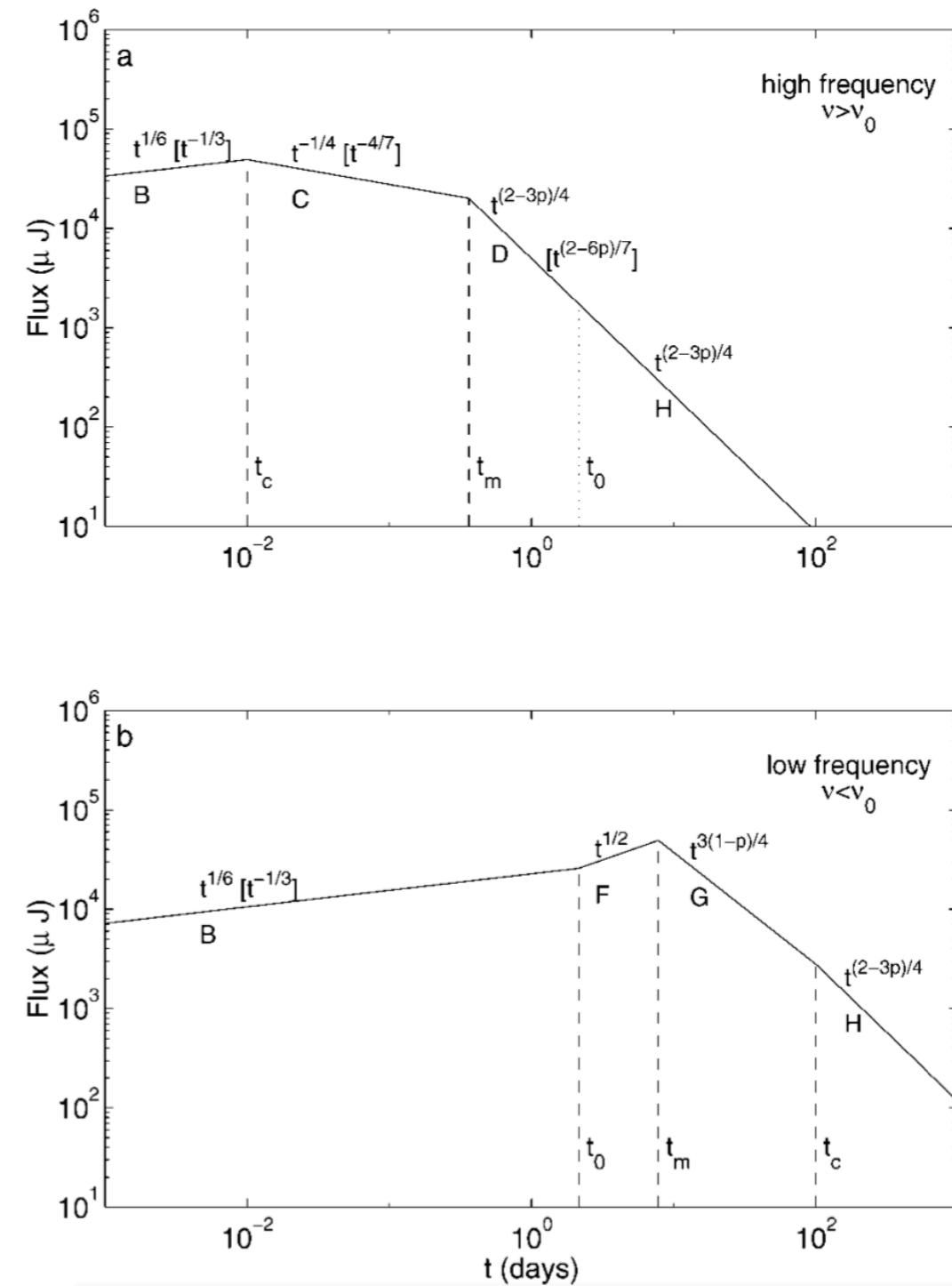
SYNCHROTRON SPECTRUM

synchrotron spectrum



SYNCHROTRON LIGHTCURVES

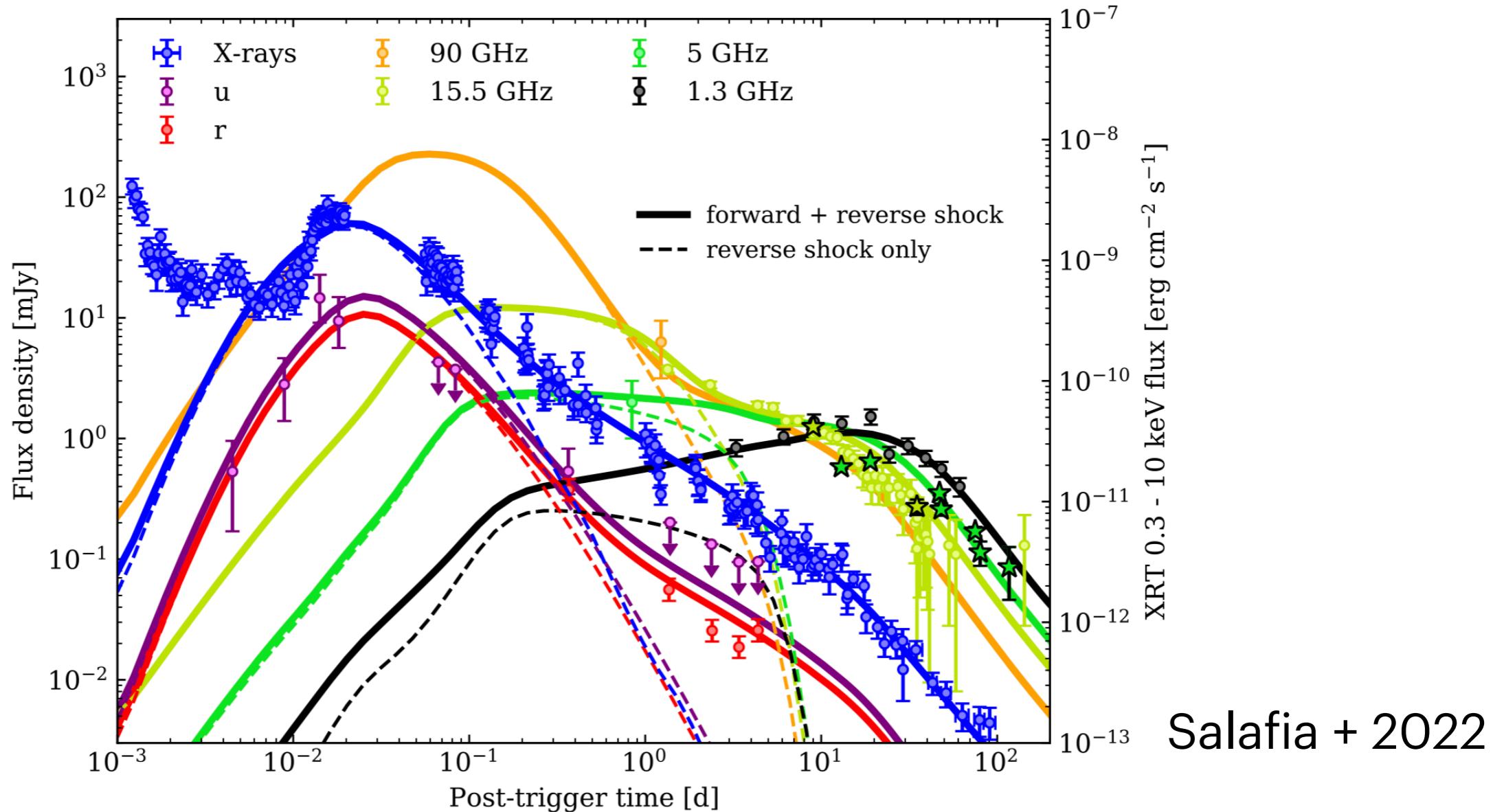
EXAMPLES OF EXPECTED AFTERGLOW LIGHTCURVES



Sari, Piran Narayan, 1998

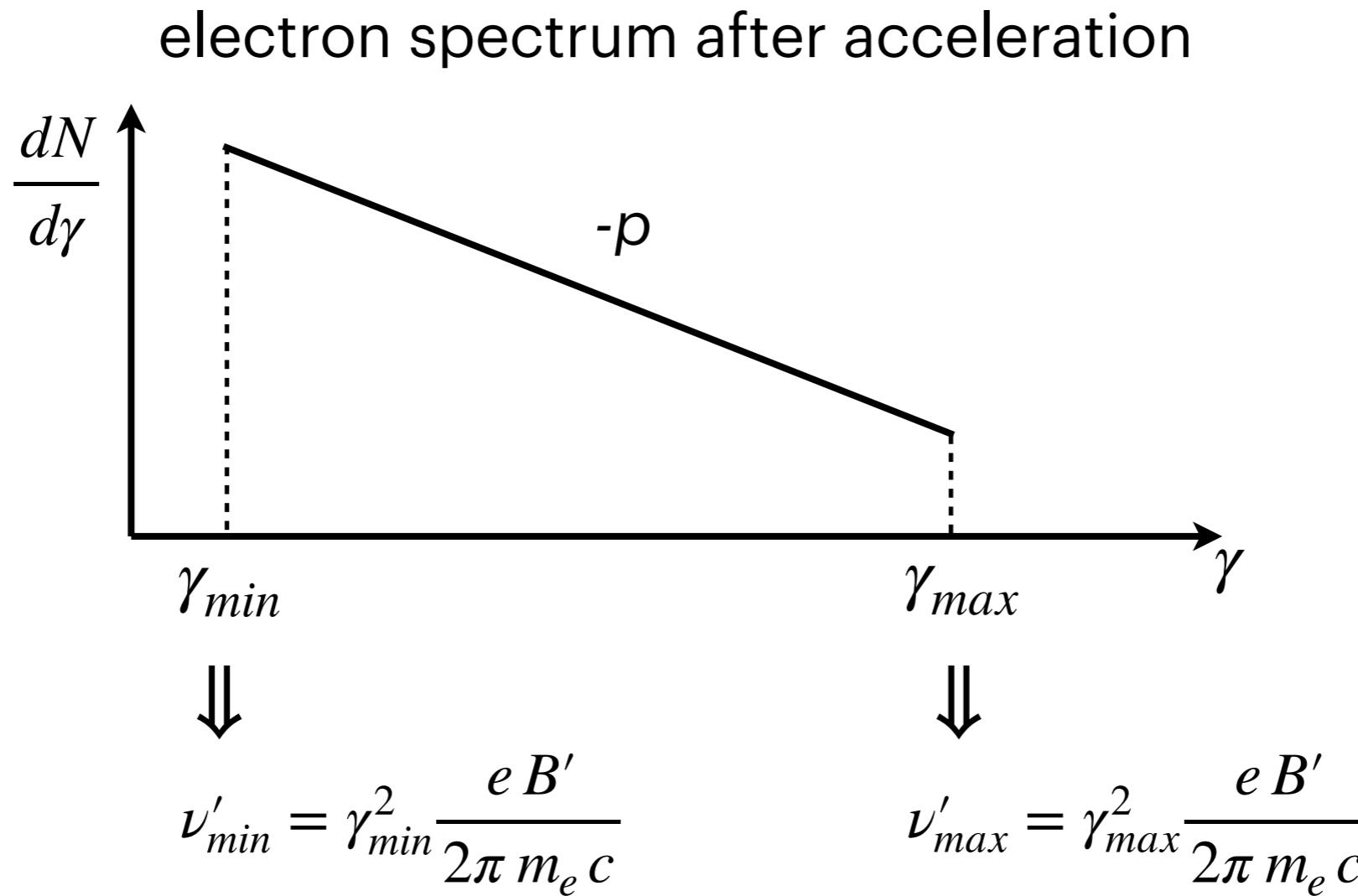
MODELING OF AFTERGLOW LIGHTCURVES

EXAMPLE OF MODELING OF MULTI-WAVELENGTH AFTERGLOW LIGHTCURVES



GRBS AT GEV ENERGIES

THE MAXIMAL SYNCHROTRON FREQUENCY



maximal energy of accelerated electrons implies
maximal energy of synchrotron photons

GRBS AT GEV ENERGIES

THE MAXIMAL SYNCHROTRON FREQUENCY

Maximal electron energy γ_{max} is reached when acceleration time is equal to cooling time: $t'_{acc}(\gamma) = t'_{syn}(\gamma)$

$$t'_{acc} \simeq \frac{r_L}{c} \simeq \frac{E'}{eB'c} = \frac{\gamma m_e c^2}{eB'c}$$
$$t'_{syn} = \frac{6\pi m_e c}{\sigma_T B'^2 \gamma}$$

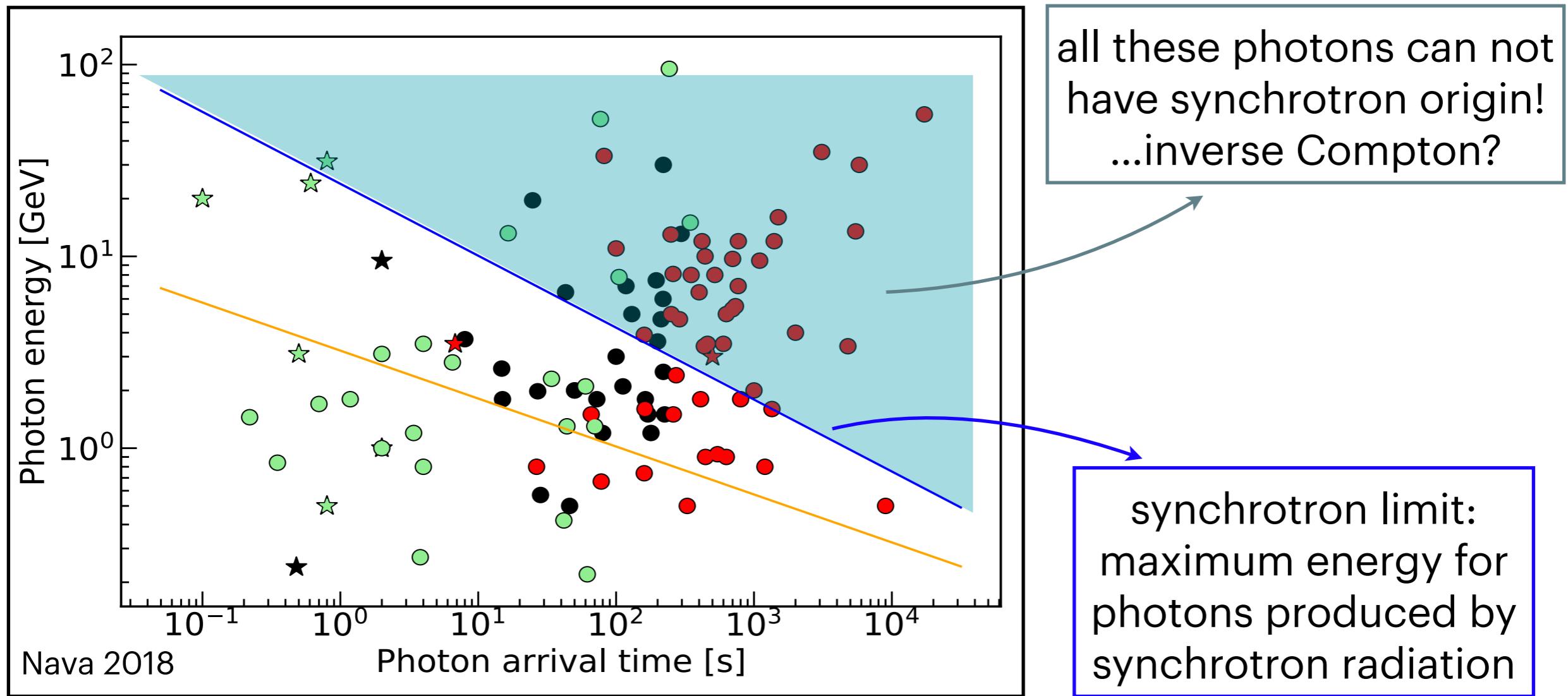
$$\Rightarrow \gamma_{max}^2 = \frac{6\pi e}{\sigma_T B'}$$

$$E_{max} = h\nu_{max} = h\nu'_{max} 2\Gamma = \gamma_{max}^2 \frac{eB'h}{2\pi m_e c} 2\Gamma = \frac{9m_e c^2 h}{4\pi e^2} \Gamma \simeq 150 MeV \times \Gamma$$

GRBS AT GEV ENERGIES

PROBLEM WITH THE INTERPRETATION

Many photons detected from GRBs have energies exceeding the maximal synchrotron energy



GRBS AT TEV ENERGIES???

PRESENCE OF SYNCHROTRON SELF COMPTON?

