VP: status and discussion



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- . Introduction
- .. HLMNT routine; status and comparison with other VP compilations; latest changes
- ... Narrow resonances: treatment and pitfalls
- Higgs fit

Introduction

• Why Vacuum Polarisation / running lpha corrections ?

Precise knowledge of VP / $\alpha(q^2)$ needed for:

- Corrections for data used as input for g-2: 'undressed' $\sigma_{\rm had}^0$ $a_{\mu}^{\rm had,LO} = \frac{1}{4\pi^3} \int_{m_{\pi}^2}^{\infty} \mathrm{d}s \, \sigma_{\rm had}^0(s) K(s) \,, \quad \text{with } K(s) = \frac{m_{\mu}^2}{3s} \cdot (0.63 \dots 1)$
- Determination of α_s and quark masses from total hadronic cross section $R_{\rm had}$ at low energies and of resonance parameters.
- Part of higher order corrections in Bhabha scattering important for precise Luminosity determination.
- $\alpha(M_Z^2)$ a fundamental parameter at the Z scale (the least well known of $\{G_\mu, M_Z, \alpha(M_Z^2)\}$), needed to test the SM via precision fits/constrain new physics.
- → Ingredient in MC generators for many processes.

• Dyson summation of Real part of one-particle irreducible blobs Π into the effective, real running coupling $\alpha_{\rm QED}$:

$$\Pi = \bigvee_{q}^{r} \bigvee_{q}$$

Full photon propagator $\sim 1 + \Pi + \Pi \cdot \Pi + \Pi \cdot \Pi \cdot \Pi + \dots$

$$\sim \alpha(q^2) = \frac{\alpha}{1 - \text{Re}\Pi(q^2)} = \alpha / (1 - \Delta\alpha_{\text{lep}}(q^2) - \Delta\alpha_{\text{had}}(q^2))$$

• The Real part of the VP, Re Π , is obtained from the Imaginary part, which via the *Optical Theorem* is directly related to the cross section, Im $\Pi \sim \sigma(e^+e^- \to hadrons)$:

$$\Delta \alpha_{\rm had}^{(5)}(q^2) = -\frac{q^2}{4\pi^2 \alpha} P \int_{m_{\pi}^2}^{\infty} \frac{\sigma_{\rm had}^0(s) \, \mathrm{d}s}{s - q^2} , \quad \sigma_{\rm had}(s) = \frac{\sigma_{\rm had}^0(s)}{|1 - \Pi|^2}$$

 $[\rightarrow \sigma^0 \text{ requires 'undressing', e.g. via } \cdot (\alpha/\alpha(s))^2 \rightsquigarrow \text{ iteration needed}]$

- Observable cross sections σ_{had} contain the |full photon propagator|², i.e. |infinite sum|².
 - \rightarrow To include the subleading Imaginary part, use dressing factor $\frac{1}{|1-\Pi|^2}$.

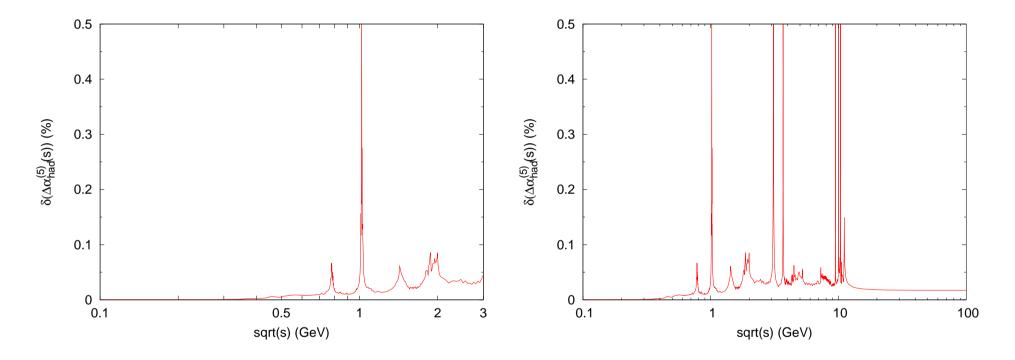
HLMNT routine; status and comparison

Features of the HLMNT VP code

- Latest routine VP_HLMNT_v2_1, version 2.1, available since 27 Jan. 2012
- Simple set of (standard) Fortran routines; completely standalone, no libs needed;
 all explanations in comment-headers
- Gives separately real and imaginary part ($\Delta \alpha(s)$ and R(s))
- Tabulation/interpolation of hadronic part, for both space- and time-like region, including errors; input 'data' included in routine (rhad not needed)
- Leptonic part coded analytically; all special function included (partly with custom made expansions)
- top contribution in the same way
- \rightarrow Flag to include or exclude narrow resonances J/ψ , ψ' , $\Upsilon(1-6S)$ [But ϕ always included via integral over final state data $(3\pi, KK)$. Is this o.k. for users?]

 $\bullet \quad \text{Typical accuracy } \delta \left(\Delta \alpha_{\mathrm{had}}^{(5)}(s) \right)$

Error of VP in the timelike regime at low and higher energies (HLMNT compilation):

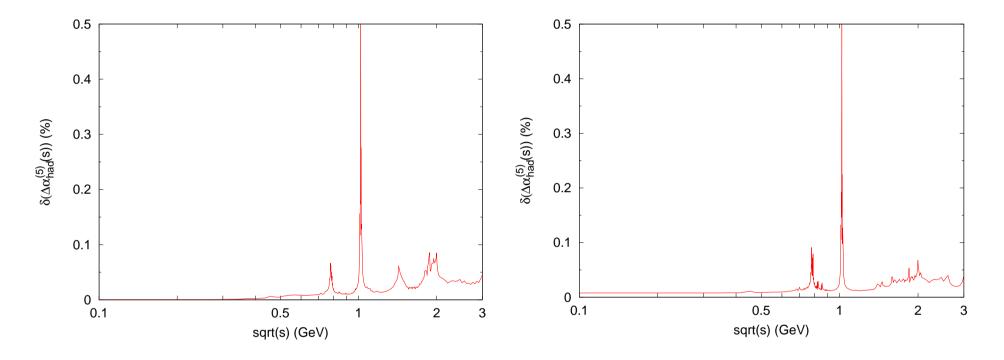


 \rightarrow Below one per-mille (and typically $\sim 5 \cdot 10^{-4}$), apart from Narrow Resonances where the bubble summation is not well justified.

Enough in the long term? Need for more work in resonance regions.

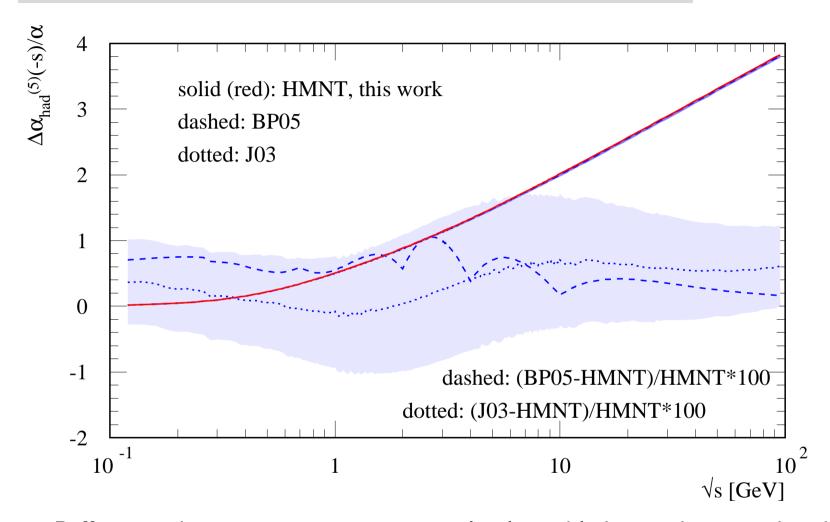
 $\bullet \quad \text{Typical accuracy } \delta \left(\Delta \alpha_{\mathrm{had}}^{(5)}(s) \right)$

Error of VP in the timelike regime: old vs. new HLMNT '11 compilation):



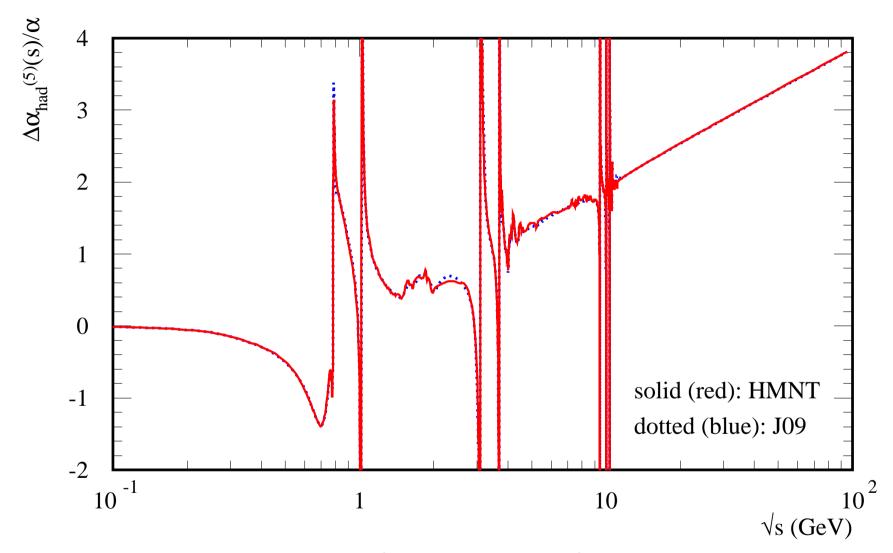
ightarrow region $1 < \sqrt{s} < 2$ GeV (and higher) improved, ho suffers from tension in 2π data (BaBar included).

• Comparison of Spacelike $\Delta \alpha_{\rm had}^{(5)}(-s)/\alpha$ (smooth $\alpha(q^2<0)$)

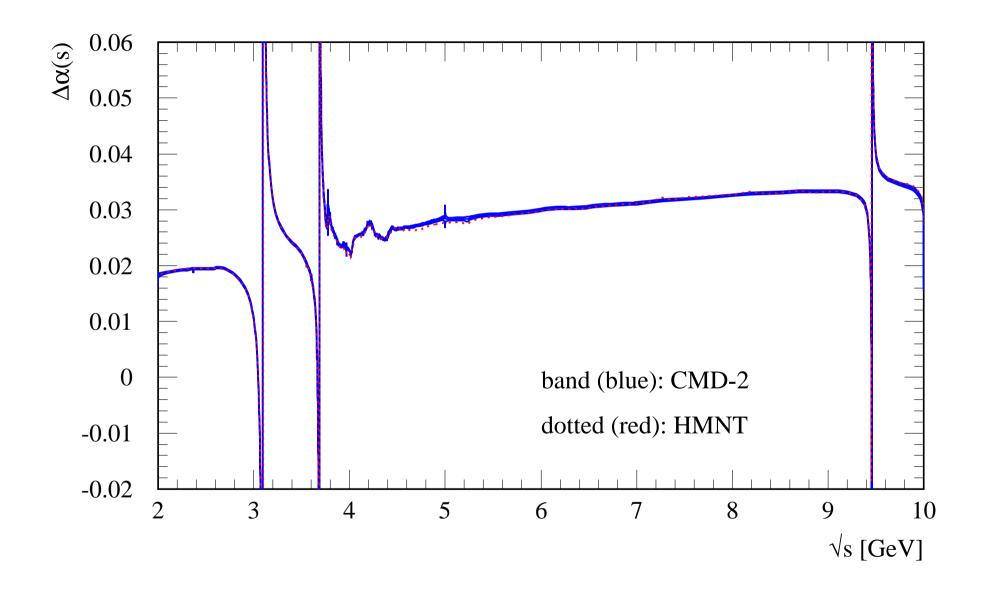


- Differences between parametrisations clearly visible but within error band (of HLMNT)
- Few-parameter formula from Burkhardt+Pietrzyk slightly 'bumpy' but still o.k.
- Encourage use of more accurate recent tabulations; $\Delta \alpha(M_Z^2)$

ullet $\Delta lpha(q^2)$ in the time-like: HLMNT compared to Fred Jegerlehner's new routines



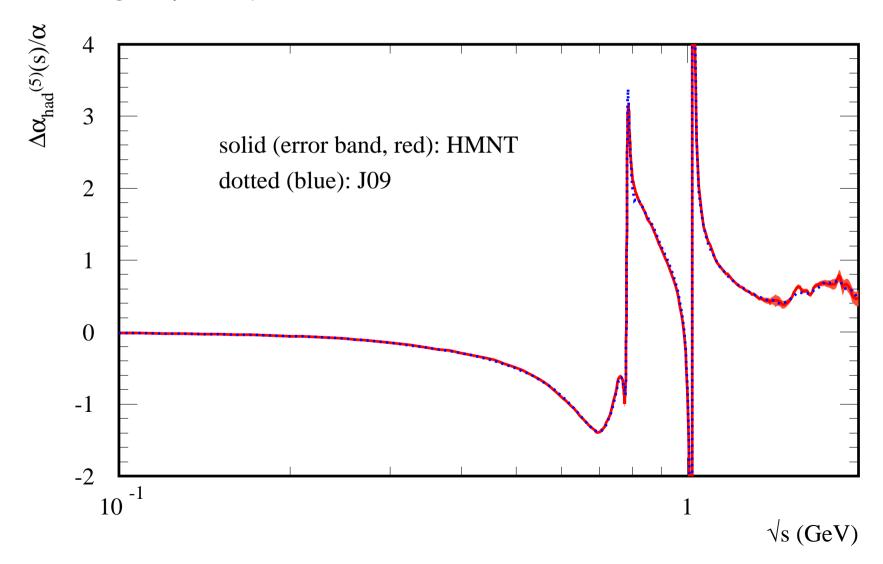
- → with new version big differences (with 2003 version) gone
 - smaller differences remain and reflect different choices, smoothing etc.



More comparison plots...

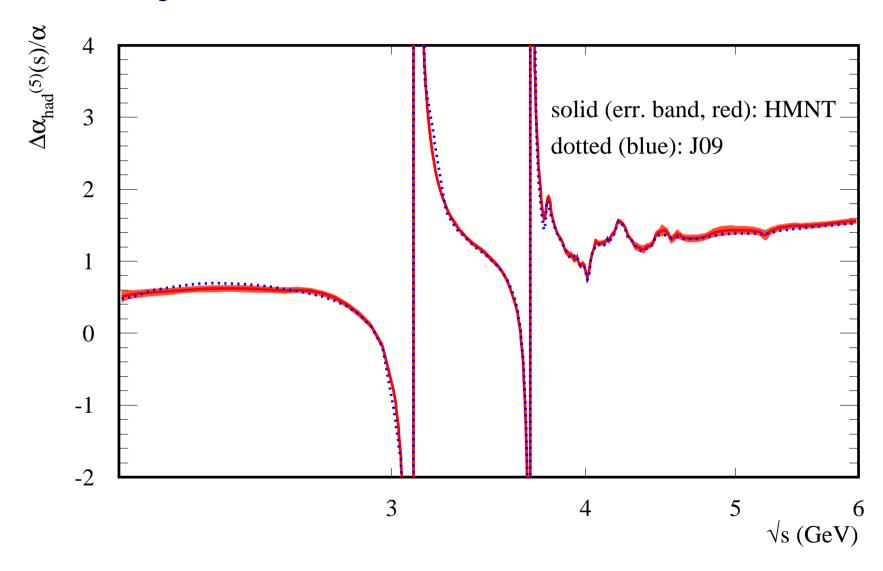
HLMNT compared to Fred Jegerlehner's new version: Detailed look

Low energies: ho and ϕ



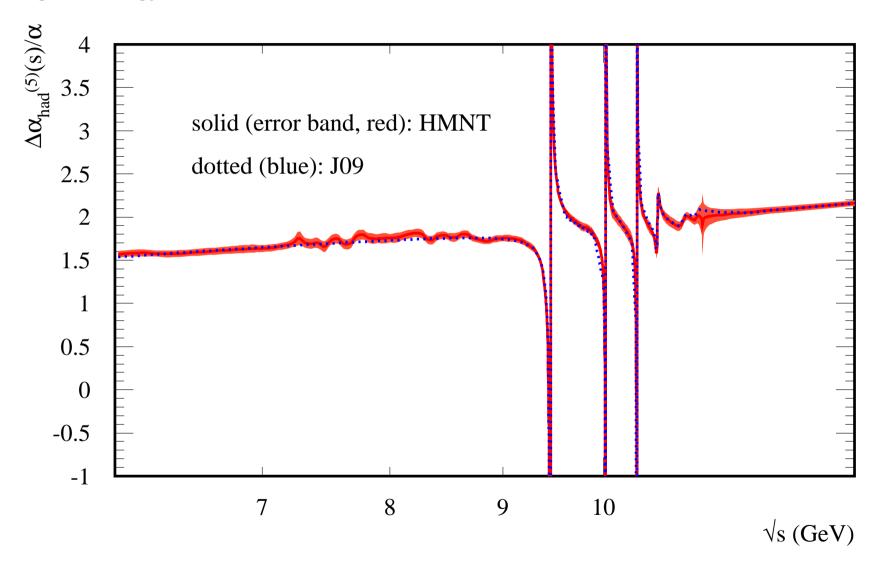
HLMNT compared to Fred Jegerlehner's new version: Detailed look

Medium energies: continuum and charm

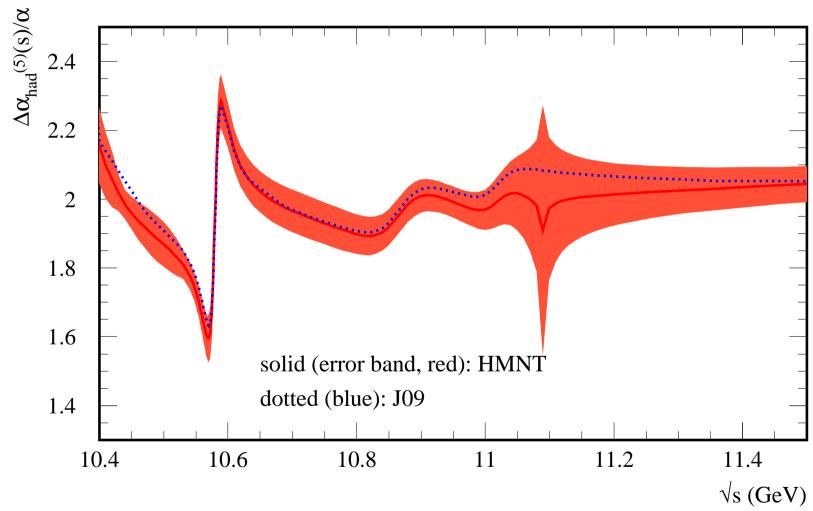


HLMNT compared to Fred Jegerlehner's new version: Detailed look

Higher energy continuum; bottom



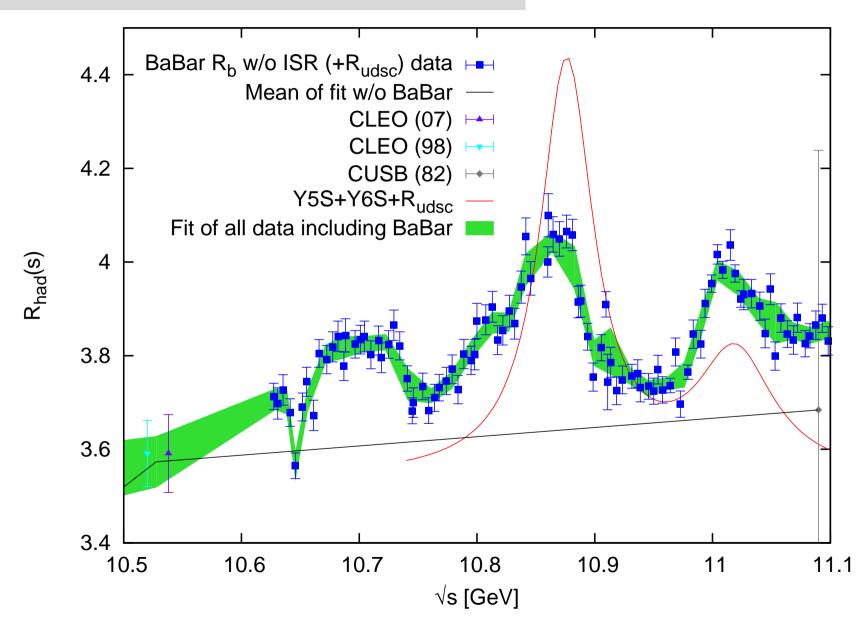
Details of higher $\Upsilon(4,5,6S)$ [10580, 10860, 11020] / open bottom region



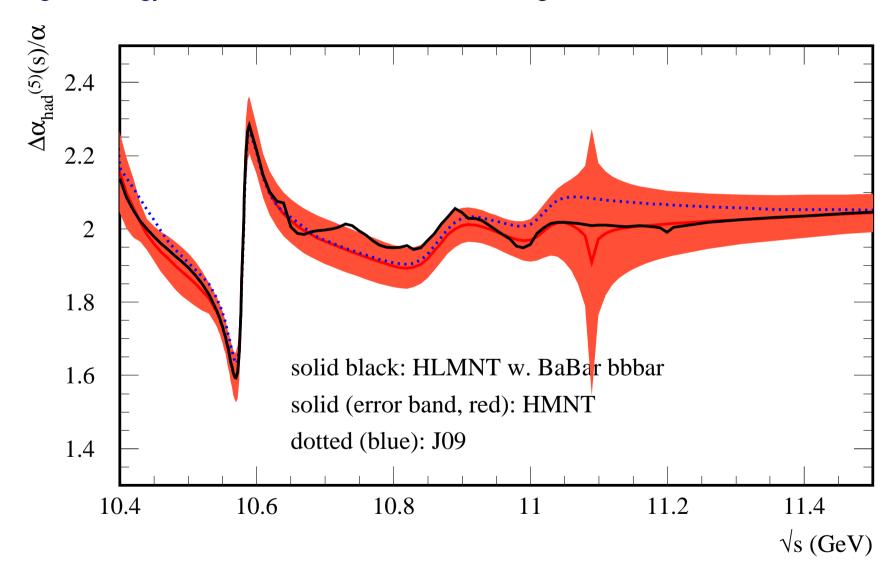
- ightarrow HLMNT still to include BaBar's $R_{b\bar{b}}$ data; ISR unfolding.. work in progress \checkmark
- expected to smooth and improve region above $11~\text{GeV}~\checkmark$

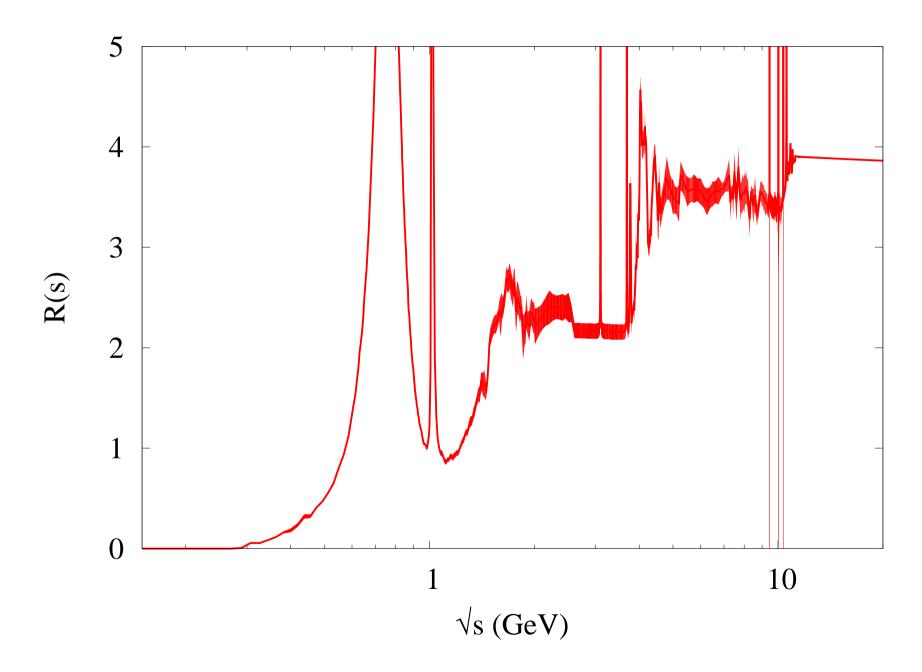
Latest changes, included in v2.1

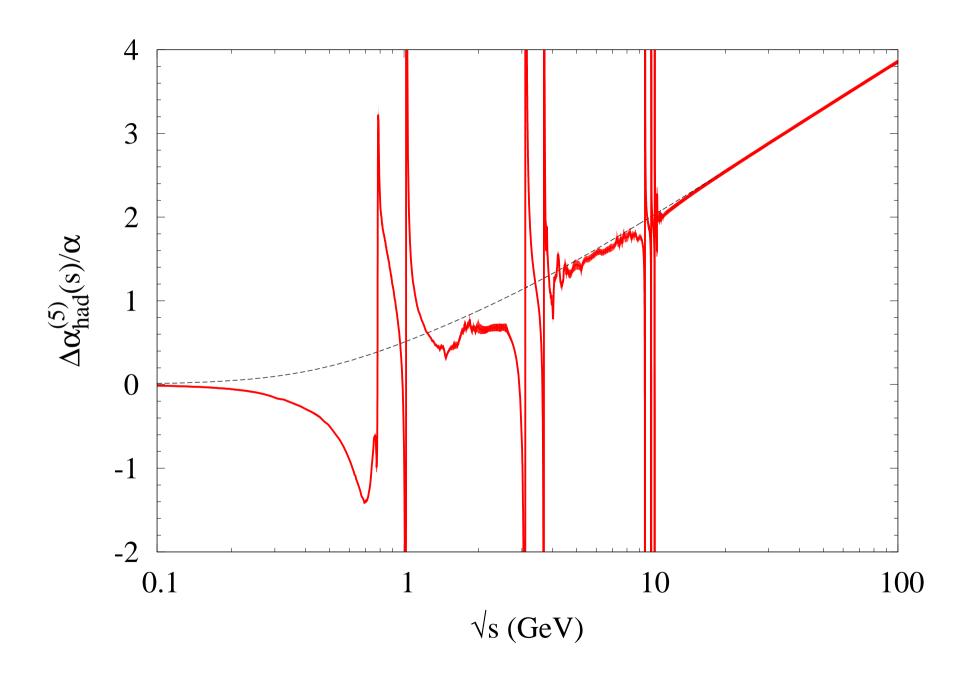
Inclusion of BaBar's $b\bar{b}$ after ISR deconvolution



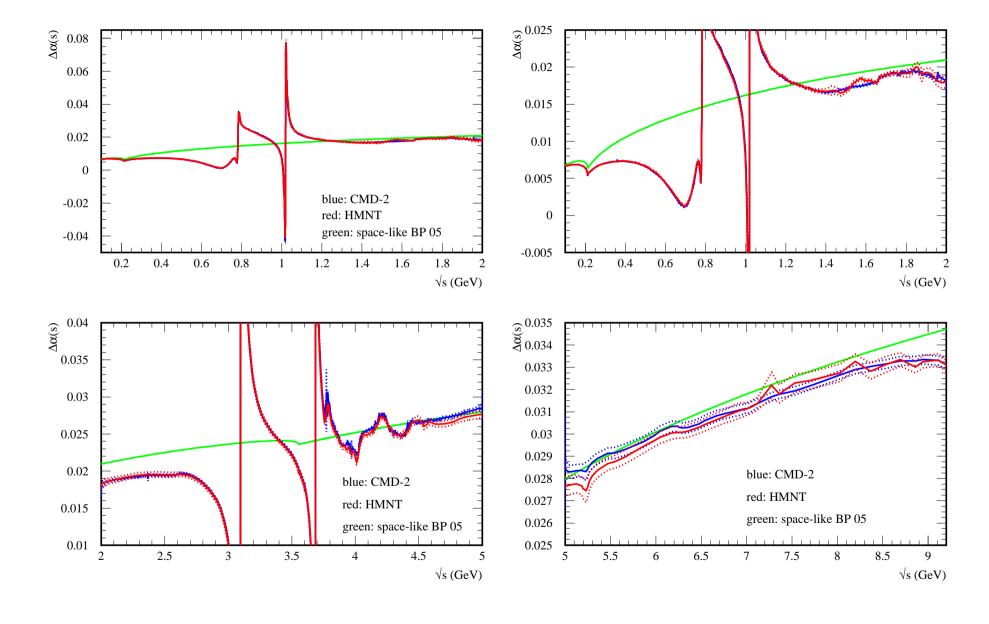
Higher energy continuum; bottom. No smoothing!



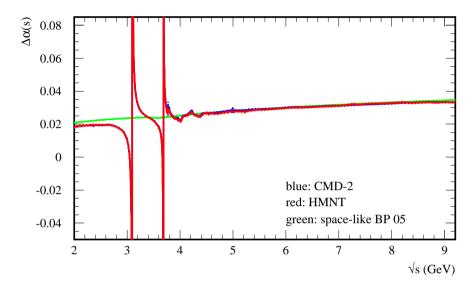


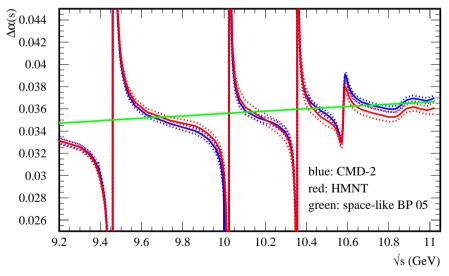


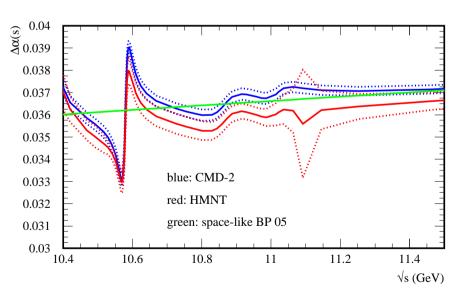
HLMNT compared to CMD-2's routine: Detailed looks



HLMNT compared to CMD-2's routine: three more zooms







Narrow Resonances: treatment and pitfalls

Note 1:

- For $\Delta \alpha$ or g-2, using NR or BW formulae with the dressed width Γ_{ee} for a resonance V is inconsistent and introduces sizeable effects (a few percent).
- Undressing via the smooth spacelike running $\alpha(-M_V)$ comes closer numerically but is not fully correct.
- Use undressing formula

$$\Gamma_{ee}^{0} = \frac{\left[\alpha/\alpha_{\text{no}V}(M_V^2)\right]^2}{1+3\alpha/(4\pi)} \Gamma_{ee},$$

where 'no V' means that the resonance V is excluded from the running α .

Note 2:

- Close to narrow resonance energies $|\Pi| \sim 1$ and the summation breaks down
- \hookrightarrow Need other formulation, e.g. Breit-Wigner resonance propagator interfering with γ :

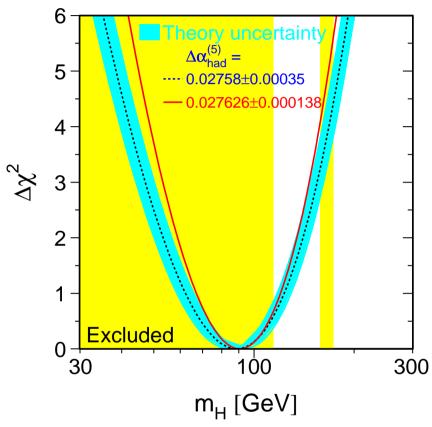
$$\left(\frac{\alpha(s)}{s}\right)^2 \rightarrow \frac{1}{s^2} \left| \alpha_{\text{no}\,V}(s) + \frac{3\Gamma_{ee}M_V}{s - M_V^2 + i\Gamma M_V} \right|^2.$$

The running QED coupling $\Delta lpha(M_Z^2)$... and the Higgs mass



- Vacuum polarisation leads to the 'running' of α from $\alpha(q^2=0)=1/137.035999084(51)$ to $\alpha(q^2=M_Z^2)\sim 1/129$
- $\alpha(q^2) = \alpha / (1 \Delta \alpha_{\text{lep}}(q^2) \Delta \alpha_{\text{had}}(q^2))$
- Hadronic uncertainties \leadsto α the least well known EW param. of $\{G_{\mu}, M_Z, \alpha(M_Z^2)\}$!
- ullet We find (HLMNT '11): $\Deltalpha_{
 m had}^{(5)}(M_Z^2)=0.027626\pm0.000138$ i.e. $lpha(M_Z^2)^{-1}=128.944\pm0.019$
- \bullet evaluation from Davier et~al.~ now very similar: $\Delta\alpha_{\rm had}^{(5)}(M_Z^2) \,=\, 0.02757 \pm 0.00010$ [now more use of pQCD can only slightly improve error]

Fit of the SM Higgs mass: LEP EWWG



Fit and Fig. thanks to M. Grünewald

$$\hookrightarrow M_H = 91^{+30}_{-23}\,\mathrm{GeV}$$
vs. 'default' $89^{+35}_{-26}\,\mathrm{GeV}$
[w. $m_t = (173.3 \pm 1.1)\,\mathrm{GeV}$]



Comparison of different compilations

ullet Timelike lpha(s) from Fred Jegerlehner's (2003 routine as available from his web-page)

$$\alpha(s = E^2) = \alpha / \left(1 - \Delta\alpha_{\text{lep}}(s) - \Delta\alpha_{\text{had}}^{(5)}(s) - \Delta\alpha^{\text{top}}(s)\right)$$

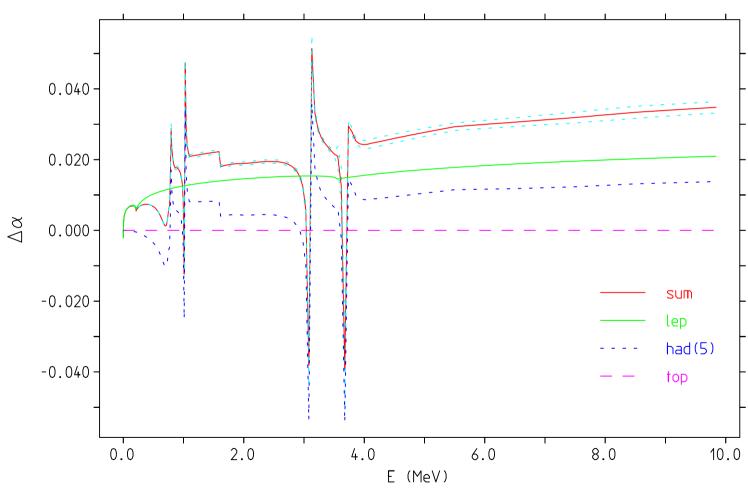
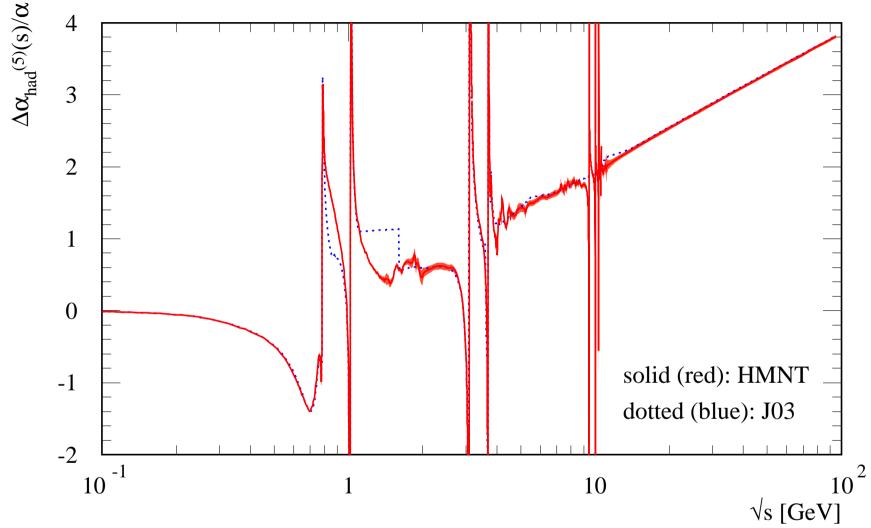


Figure from Fred Jegerlehner

Timelike $\alpha(s=q^2>0)$ follows resonance structure:

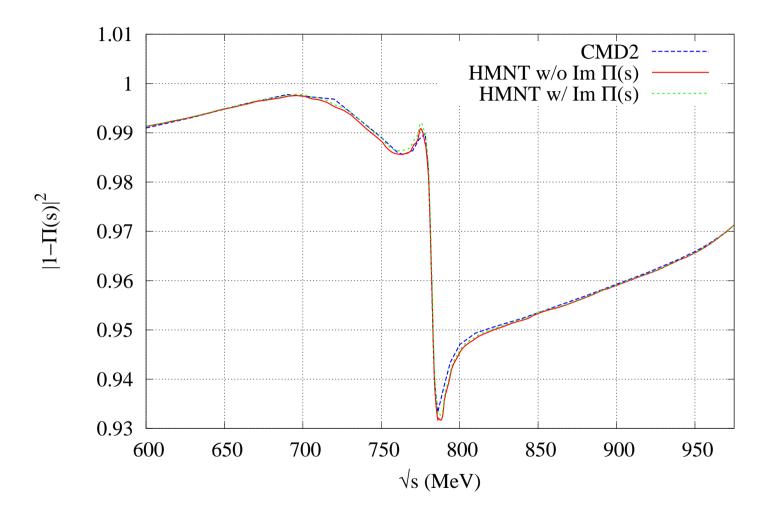


- Step below just a feature of unfortunate grid.
- Difference below 1 GeV not expected from data.

[Comparisons with other parametrisations confirm HMNT.]

• HMNT compared to Novosibirsk's parametrisation

Timelike $|1 - \Pi(s)|^2 \sim (\alpha(s)/\alpha)^2$ in ρ central energy region: A relevant correction!



→ Small but visible differences, as expected from independent compilations.