



5-th INTERNATIONAL CONFERENCE
“CHANNELING 2012” SEPTEMBER 23 - 28 2012

ALGHERO, ITALY

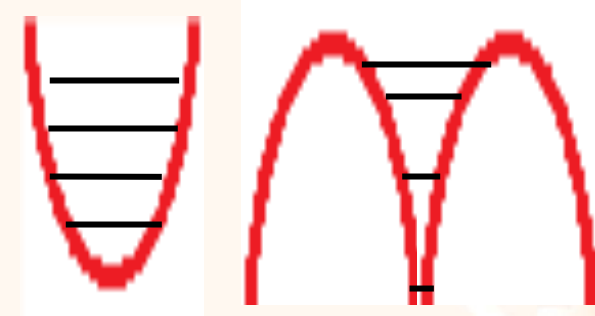
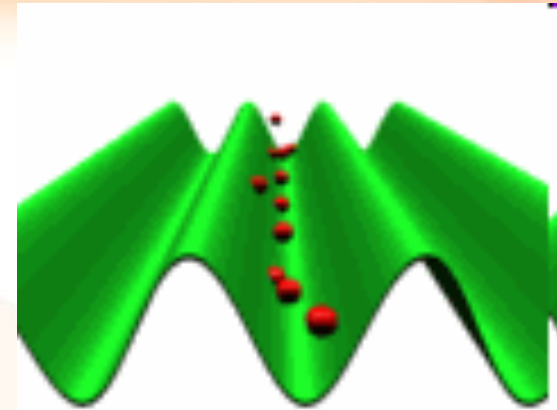
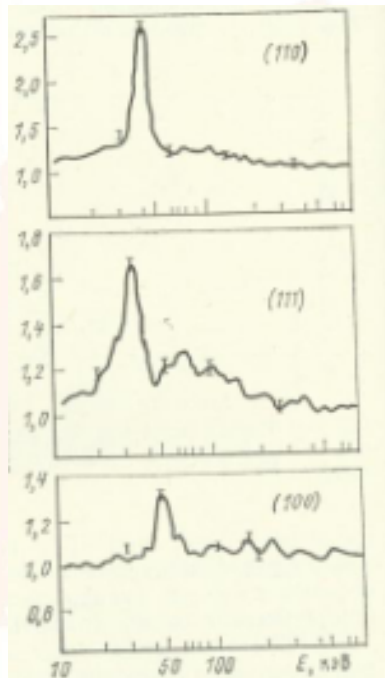
Radiation from Channeling Electrons, Stimulated by Laser Beam

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Standard approach



$$\Delta E \sim (U/E)^{1/2} (hc/d) \sim (10-100\text{eV})(mc^2/E)^{1/2};$$

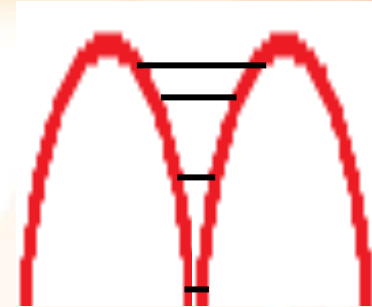
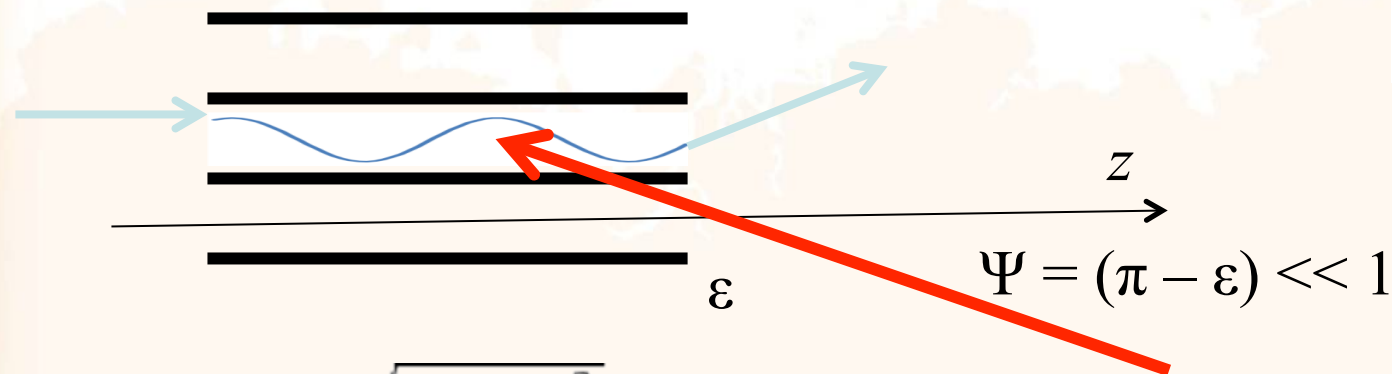
$$h\nu = \Delta E (E/mc^2)^2 \sim (10-100\text{eV})(E/mc^2)^{3/2}$$

$$U' = U_0 E/mc^2 \quad N \sim (E/mc^2)^{1/2} \quad \Delta E \sim U'/N \sim U_0 (E/mc^2)^{1/2}$$



Channeling radiation

Radiation, stimulated by laser beam



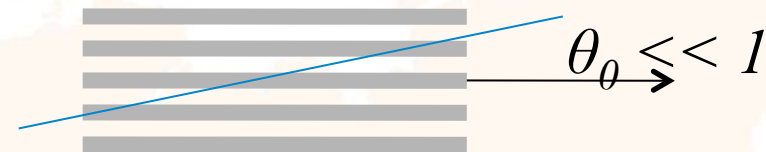
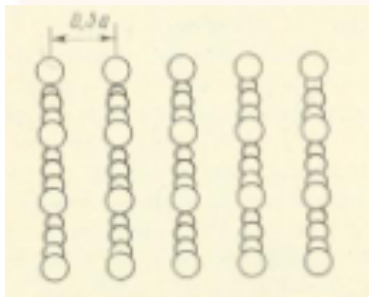
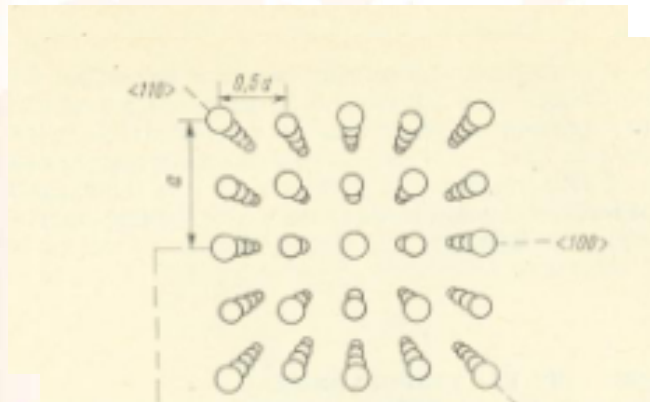
$$\omega = \omega_0 \cdot \frac{\sqrt{1 - \frac{v^2}{c^2}}}{1 + \frac{v}{c} \cdot \cos \theta} = \sim 2\omega_0 / (mc^2/E + E\Psi^2/mc^2) < 2\omega_0 E/mc^2$$

In the *accompanying system*, moving parallel with the electron beam

$$\omega = (2\pi\Delta E/h)(E/mc^2) \sim (2\pi c/d)(EU)^{1/2}/mc^2 < 2\omega_0 E/mc^2$$

$$\Delta E \sim (10-100eV)(mc^2/E)^{1/2} < h\omega_0/\pi \sim 5eV$$

Radiation, stimulated by periodically inhomogeneous planar channeling potential



*Electron in planar channeling.
Plane consisting of axes*

$$\Delta E \sim (U/E)^{1/2} (hc/d) \sim (hc/d) \theta_L (\text{planar})$$

$$\theta_0 > \theta_L (\text{axial}) \gg \theta_L (\text{planar})$$

$$h\nu = hc(\theta_0/d)(E/mc^2)^2 \gg \Delta E(E/mc^2)^2 \sim hc(\theta_L/d)(E/mc^2)^2$$



Kinematics

$$\vec{P}_i = \vec{P}_i(x) + \vec{P}_{i\perp}(\theta) = P_0 \left(1 - \frac{\theta_0^2}{2} \right) \vec{i} + P_0 \theta_0 \vec{j}$$



$$\frac{P_{i\perp}^2}{2m} \approx \langle U \rangle, \quad \frac{P_0^2 \theta_0^2}{2m} \approx \langle U \rangle, \quad \theta_0 \leq \frac{2m \langle U \rangle}{P_0^2} = \theta_L^2$$



“Accompanying” reference (coordinate) system

$$P'_x = \frac{P_x - \frac{V}{c^2} E}{\sqrt{1 - \frac{V^2}{c^2}}} = 0 \quad \Rightarrow V = \frac{P_x}{E} c^2$$

$$\begin{aligned} P'_y &= P_y \\ P'_z &= P_z \end{aligned} \quad \Rightarrow P'_\perp = P_\perp$$

$$E' = \frac{E - VP_x}{\sqrt{1 - \frac{V^2}{c^2}}} \Rightarrow m_0 c^2 + \frac{P_\perp^2}{2m_0}$$



In “accompanying” system

$$\omega = \omega_0 \frac{\sqrt{1 - \frac{v^2}{c^2}}}{1 + \frac{\vec{v} \cdot \vec{k}}{\omega_0}} = \omega_0 \frac{\sqrt{1 - \frac{v^2}{c^2}}}{1 + \frac{v}{c} \cos \psi}$$

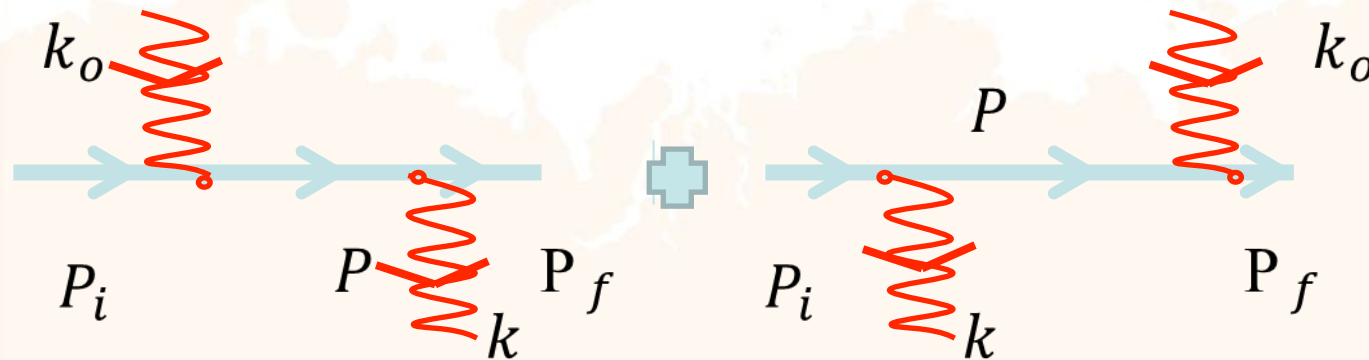
$$\psi = \pi - \varepsilon$$

$$\omega = \omega_0 \frac{\sqrt{1 - \frac{v^2}{c^2}}}{1 - \frac{v}{c} + \frac{v}{c} \frac{\varepsilon^2}{2}} \Rightarrow \omega = \omega_0 \frac{1}{\gamma^{-1} + \frac{\varepsilon^2}{2} \gamma}$$



Induce scattering laser photon by "bound" electron

Richard P. FEYMAN diagrams



$$P_i + k_0 = P_f + k$$

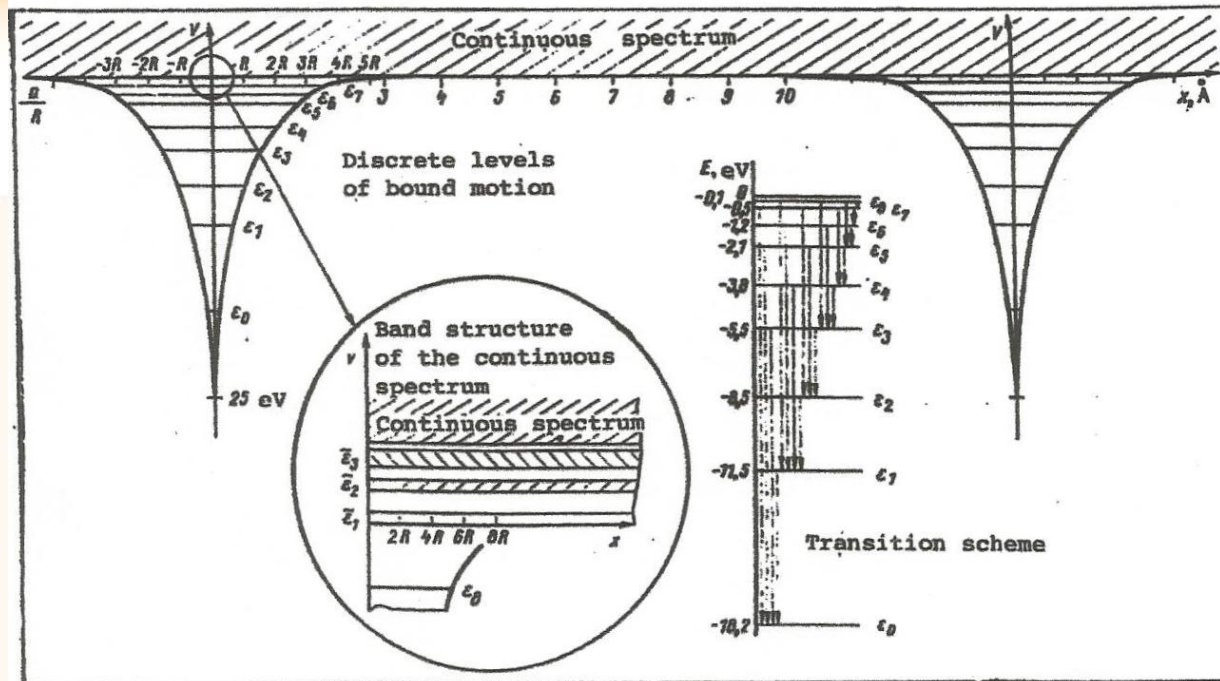
$$P_i \cos \theta_0 + \frac{\hbar \omega_0}{c} \cos \psi = P_f \cos \theta_f + \frac{\hbar \omega}{c} \cos \psi_f$$

$$E_i + \hbar \omega_0 = E_f + \hbar \omega$$

$$\hbar \omega = E_i - E_f + \hbar \omega_0$$



Structure of energy bands and radiative transitions of 56-MeV electrons channeled along the (110) plane in Si



$$\overline{\Delta E} \sim \frac{U_{acc}}{N} \sim \frac{U_0 \gamma}{\sqrt{\gamma}} \sim U_0 \sqrt{\gamma}$$

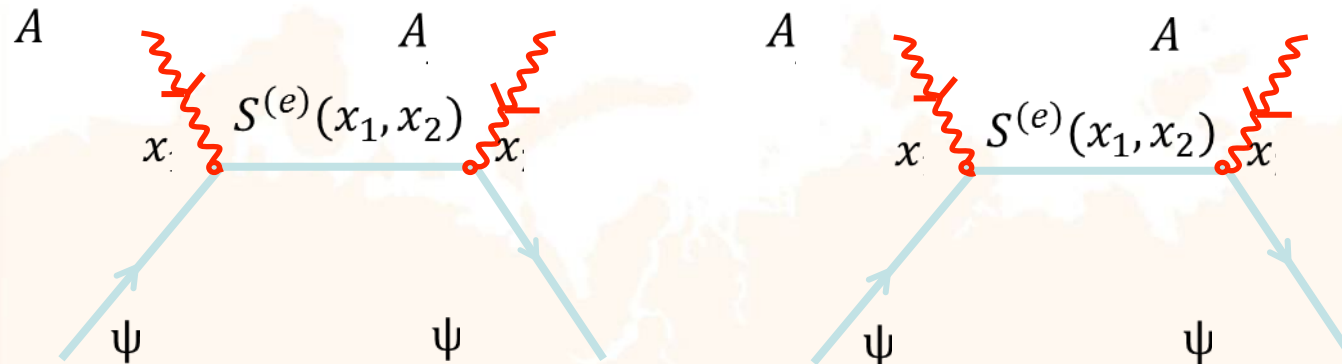
$$U_0(\text{Si}) \approx 25 \text{ eV} \quad R \approx 0,2 \text{ \AA}$$

$$\overline{\omega}_{12}(\text{lab}) \sim \gamma^{3/2}$$

$$\omega_{max}(\text{lab}) \sim U_0 \gamma^2$$



Scattering of a photon by a "bound" electron



$$\vec{A}_i(x) = \vec{A}_i(r) e^{-i\omega_i t} = \frac{\vec{e}_i}{\sqrt{2\omega_i}} e^{i(\vec{k}_i \vec{r} - \omega_i t)}$$

$$\vec{A}_f(x) = \vec{A}_f(r) e^{i\omega_f t} = \frac{\vec{e}_f}{\sqrt{2\omega_f}} e^{-i(\vec{k}_f \vec{r} - \omega_f t)}$$

$$\psi_i(x) = \psi_i(\vec{r}) e^{-iE_i t}$$

$$\psi_f(x) = \psi_f(\vec{r}) e^{-iE_f t}$$



Matrix element

$$U_{i \rightarrow f} = \frac{2\pi\alpha}{\sqrt{\omega_i\omega_f}} \sum_n \left\{ \frac{\left(f \left| \hat{e}_f e^{-i\vec{k}_f \vec{r}} \right| n \right) \left(n \left| \hat{e}_i e^{i\vec{k}_i \vec{r}} \right| i \right)}{E_n - E_i - \omega_i} \right. \\ \left. + \frac{\left(f \left| \hat{e}_i e^{i\vec{k}_i \vec{r}} \right| n \right) \left(n \left| \hat{e}_f e^{-i\vec{k}_f \vec{r}} \right| i \right)}{E_n - E_f - \omega_f} \right\}$$



Probability of photon scattering by a “bound” electron

$S_{i \rightarrow f} = -(2\pi)iU_{i \rightarrow f} \delta(E_i + \omega_i - E_f - \omega_f)$, where

$$U_{i \rightarrow f} = 2\pi\alpha e^{i(\vec{k}_i - \vec{k}_f)\vec{R}} \sqrt{\omega_i \omega_f} \sum_n \left\{ \frac{(\vec{r}\vec{e}_f)_{fn}(\vec{r}\vec{e}_i)_{ni}}{E_i - E_n + \omega_i} + \frac{(\vec{r}\vec{e}_i)_{fn}(\vec{r}\vec{e}_f)_{ni}}{E_i - E_n - \omega_f} \right\}$$

$$d\sigma = 2\pi |U_{if}|^2 \delta(E_i + \omega_i - E_f - \omega_f) \frac{d^3\vec{k}_f}{(2\pi)^3}$$

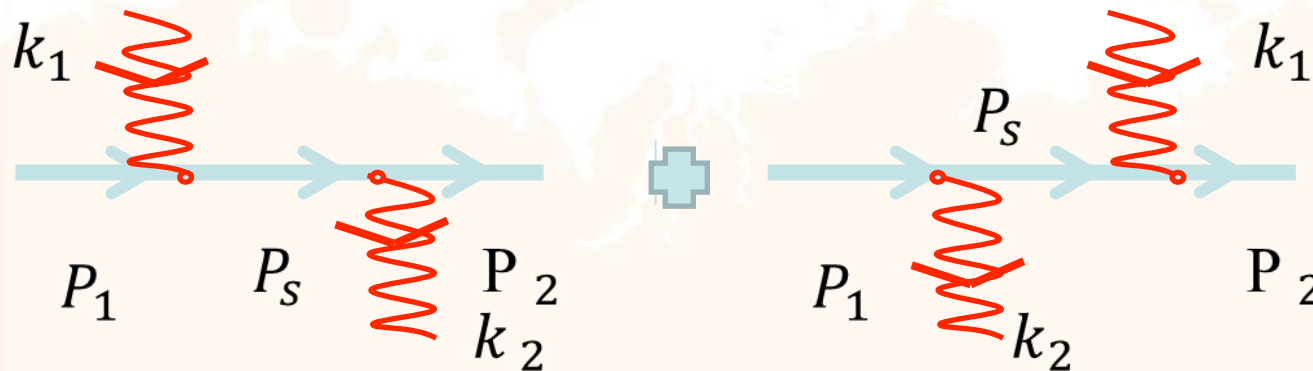
$$d\sigma = \left| \sum_n \left[\frac{(\vec{Q}\vec{e}_f)_{fn}(\vec{Q}\vec{e}_i)_{ni}}{E_i - E_n + \omega_i} + \frac{(\vec{Q}\vec{e}_i)_{fn}(\vec{Q}\vec{e}_f)_{ni}}{E_i - E_n - \omega_f} \right] \right|^2 \omega_i \omega_f^2 d\Omega_f$$



Resonance scattering laser photons

$$\omega_1 = E_s - E_1 \quad (\omega_1 = \omega_2)$$

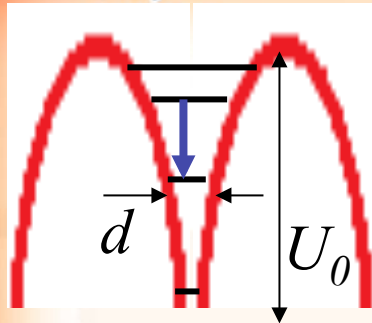
$$\psi_s \sim e^{-i(E_s - \frac{i}{2}r_s)t}$$



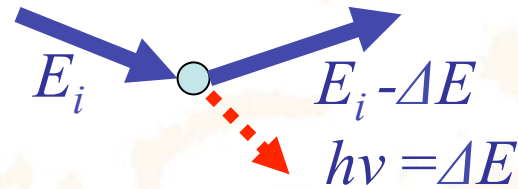
$$d\sigma = \omega_1 \omega_2^3 d\Omega_2 \sum \frac{\langle 2 | \vec{Q} \hat{e}_2 | s \rangle \langle s | \vec{Q} \hat{e}_1 | 1 \rangle}{(E_s - E_1 - \omega_1)^2 + \frac{1}{4} r_s^2}$$



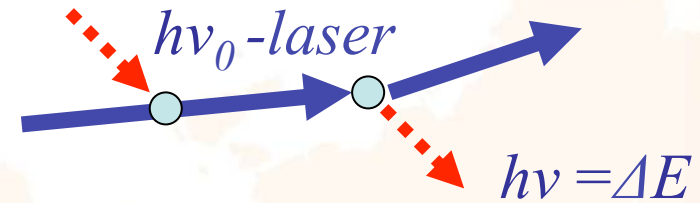
Channeling radiation



Spontaneous radiation



Stimulated radiation



1. Potential in lab. system: $U_0 \sim 20eV$; $d \sim 0,2-0,3 \text{ \AA}$ (Si);
in accompanying system ($v_x = 0$): $U = U_0 (E/mc^2)$

2. Number of levels: $N \sim P_{x \max} d / h \sim (EU_0)^{1/2} d / hc$

3. Distance between levels in acc. system: $\Delta E \sim U/N \sim (EU_0)^{1/2} (h/mcd)$

4. Laser photon energy in acc. system:

$$h\nu = 2h\nu_0 / (mc^2/E + E\psi^2/mc^2) < 2h\nu_0 E/mc^2$$

5. In resonance conditions stimulated radiation can be very effective:

$$\Delta E \sim (EU_0)^{1/2} (h/mcd) = h\nu < 2h\nu_0 E/mc^2 = 2hE/mc\lambda_0$$

6. Resonance can be reached by correctly orienting laser beam, if:

$$E/U_0 > (\lambda_0/2d)^2 \sim 10^{7-8}$$



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Thank You for Your Attention!

