

Crystal Undulators: from the Prediction to the Mature Simulations

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Research Institute for Nuclear Problems, Minsk, Belarus

Channeling 2012 Alghero September 27

Plan

CU radiation prediction and ultimate perspectives of its application

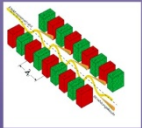
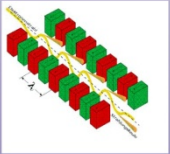
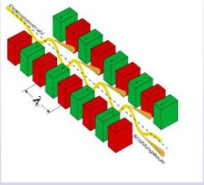
CU radiation simulation

experiments with electrons

predictions for positrons

CU radiation improvement

by crystal cut



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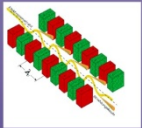
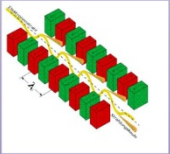
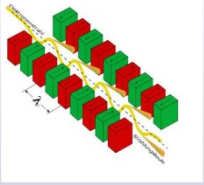
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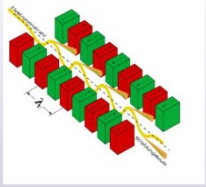
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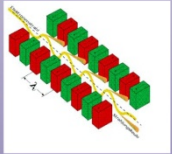




The prediction of Radiation in Crystal Undulators:

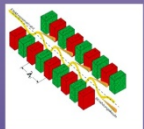
Baryshevsky V.G., Grubich A.O., Dubovskaya I. Ya.

Phys. Lett. A77(1980)61.



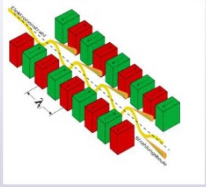
Voroviev S.A., Kaplin V.V., Plotnikov S.V.

Zh. Tekh. Fiz. 50(1980)1079.



Crystal plane undulations
can be exerted by either
ultrasonic or *laser wave*





X-ray laser on channeling needs $j \geq 10^{13} \text{ A/cm}^2$

Baryshevsky V.G., Feranchuk I.D.

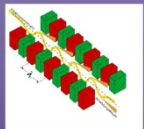
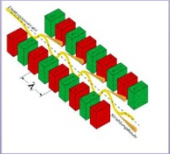
Phys. Lett. A102(1984)141.

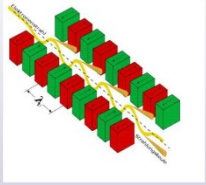
Baier V.N., Milstein A.I.

NIM B17(1986)25.

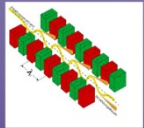
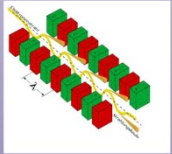
Bazylev V.A., Zhevago N.K.

*Fast Particle Radiation in Matter
and External Fields, 1987.*





γ -ray laser threshold current
can be reduced to $j \geq 10^8 \text{ A/cm}^2$
using distributed feedback
provided by *x-ray diffraction*



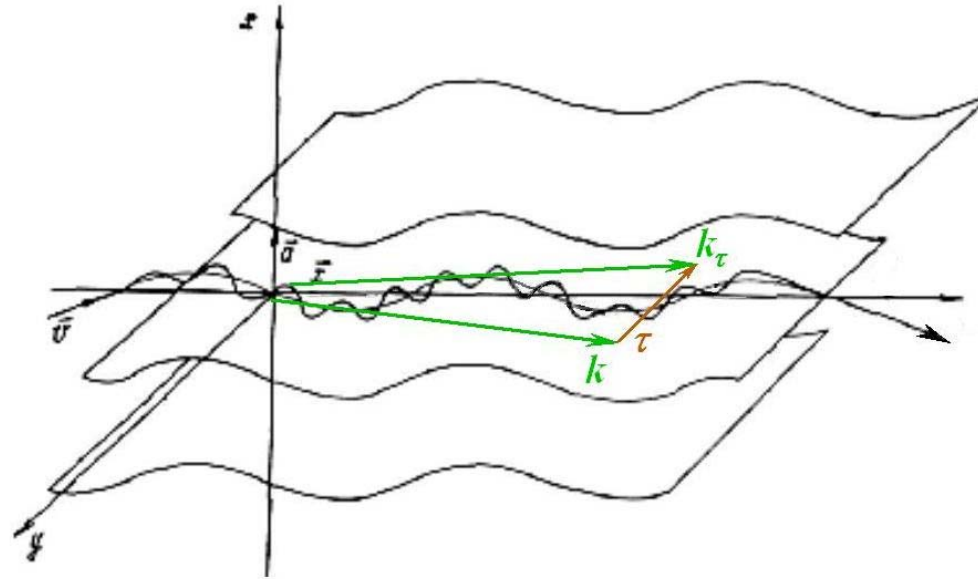
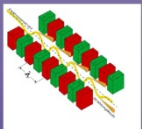
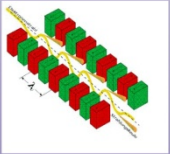
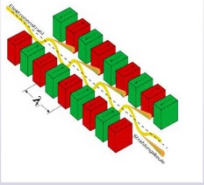
Baryshevsky V.G., Feranchuk I.D.

Phys. Lett. A102(1984)141.



Diffraction Radiation in Crystal Undulator

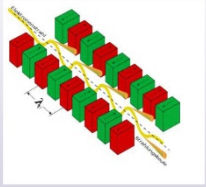
Baryshevsky V.G., Dubovskaya I. Ya. J. Phys. CM. 33(1991)2421.



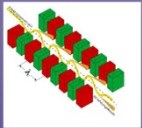
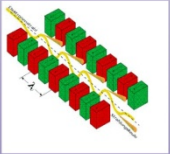
$$N_{\tau m}^s = [\pi Q^2 \sqrt{\beta_1} L |g_{\tau}^s(\omega_B)|^2 \omega_B^2 m^2 / 16 |\tau_z|] \quad (22)$$

$$\times a^2 \Omega'^2 [1 - 2(\omega_B / \omega_{\max}) + 4P_{\tau}^s (\omega_B / \omega_{\max})^2].$$

– is *more intensive* than the
Diffraction Radiation of Channeled Particles



Crystal Undulator radiation
used with distributed feedback
provided by *x-ray diffraction*
open up a real perspective
of *γ -laser* development



Plan

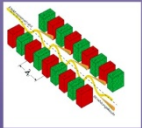
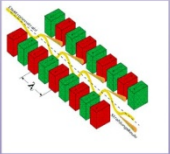
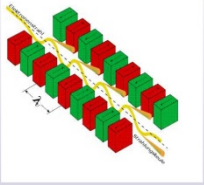
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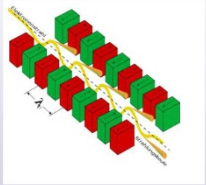
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We proceed from the genuine formula of *Baier and Katkov*

The general expression for radiation intensity

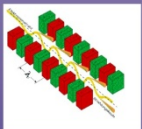
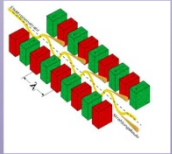
$$\frac{d^2 I}{d\omega d^2\theta} = \frac{\alpha\omega^2 d\omega}{8\pi^2 \varepsilon'^2} \times \int \int dt_1 dt_2 \left[(\varepsilon^2 + \varepsilon'^2) (\mathbf{v}_\perp(t_1) - \boldsymbol{\theta})(\mathbf{v}_\perp(t_2) - \boldsymbol{\theta}) + \omega^2/\gamma^2 \right] \\ \exp \left\{ i \frac{\omega\varepsilon}{2\varepsilon'} \left[\int_{-\infty}^{t_1} (\gamma^{-2} + (\mathbf{v}_\perp(t') - \boldsymbol{\theta})^2) dt' + \int_{-\infty}^{t_2} (\gamma^{-2} + (\mathbf{v}_\perp(t'') - \boldsymbol{\theta})^2) dt'' \right] \right\}$$

contains two integrals

$$A = \int \exp \left\{ i \frac{\omega\varepsilon}{2\varepsilon'} \int_{-\infty}^t [\gamma^{-2} + (\mathbf{v}_\perp(t') - \boldsymbol{\theta})^2] dt' \right\} dt,$$

$$\mathbf{B} = \int (\mathbf{v}_\perp(t) - \boldsymbol{\theta}) \exp \left\{ i \frac{\omega\varepsilon}{2\varepsilon'} \int_{-\infty}^t [\gamma^{-2} + (\mathbf{v}_\perp(t') - \boldsymbol{\theta})^2] dt' \right\} dt$$

and slowly decreases with radiation angle $\boldsymbol{\theta}$, complicating its numerical integration.



Key simulation points:

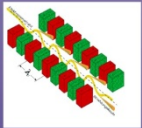
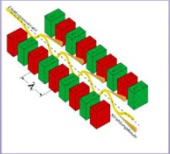
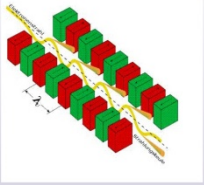
Trajectory simulations in most
realistic potentials

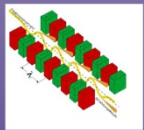
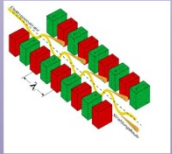
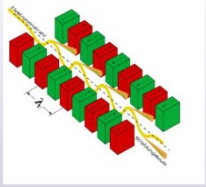
Simulation of **incoherent scattering** on
both nuclei and electrons

Separate simulation of **single**
and **multiple** scattering

Direct integration of
Baier-Katkov formula

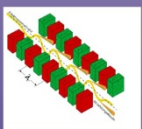
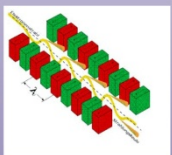
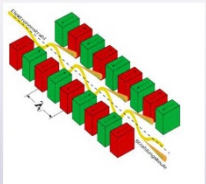
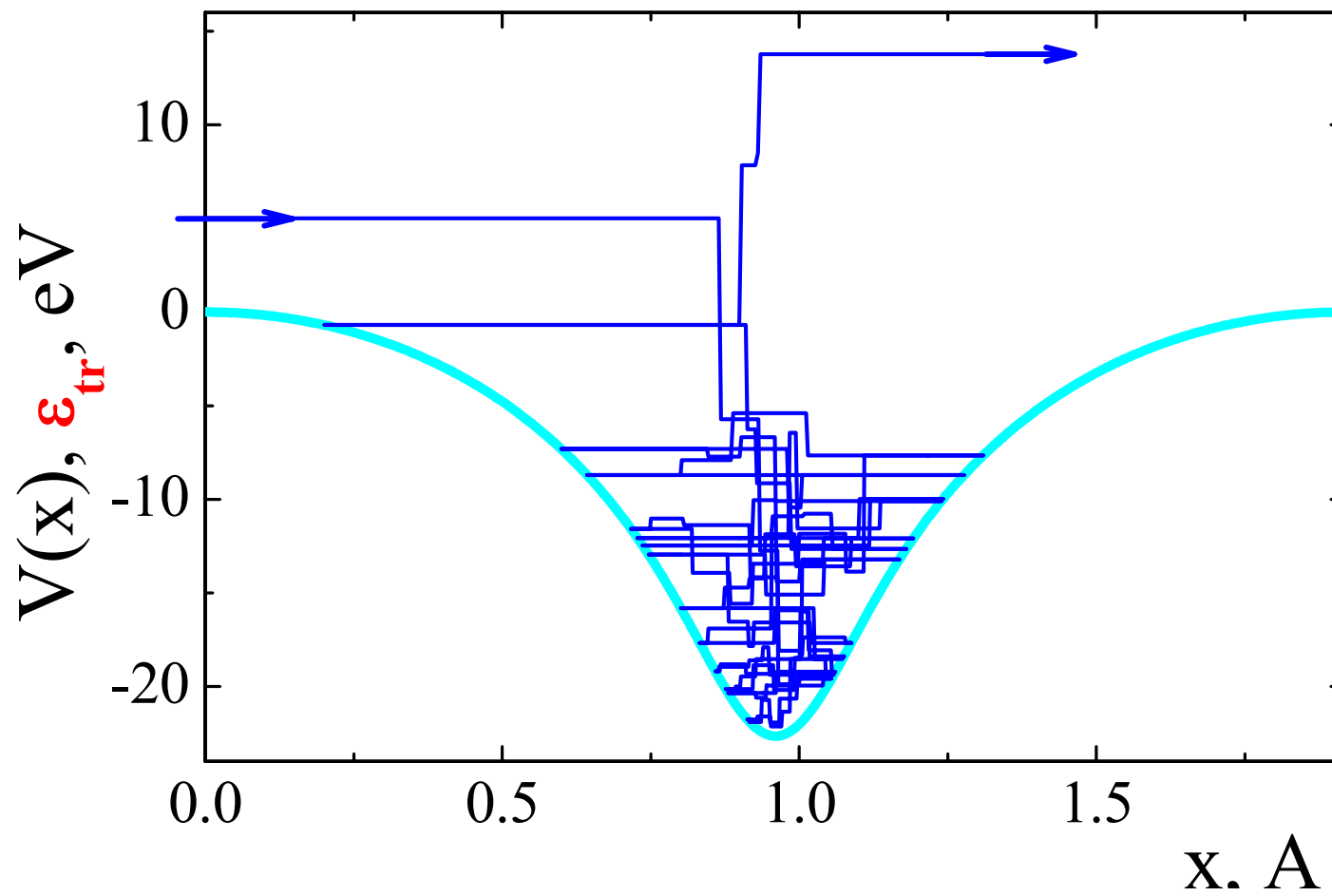
Infinite trajectories, **density** effect...

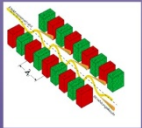
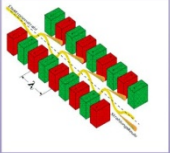
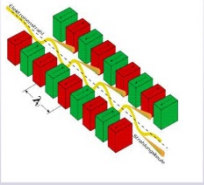




Electron radiation in CU
is essentially determined
by the influence
of *incoherent scattering*

Capture of a 287 MeV electron





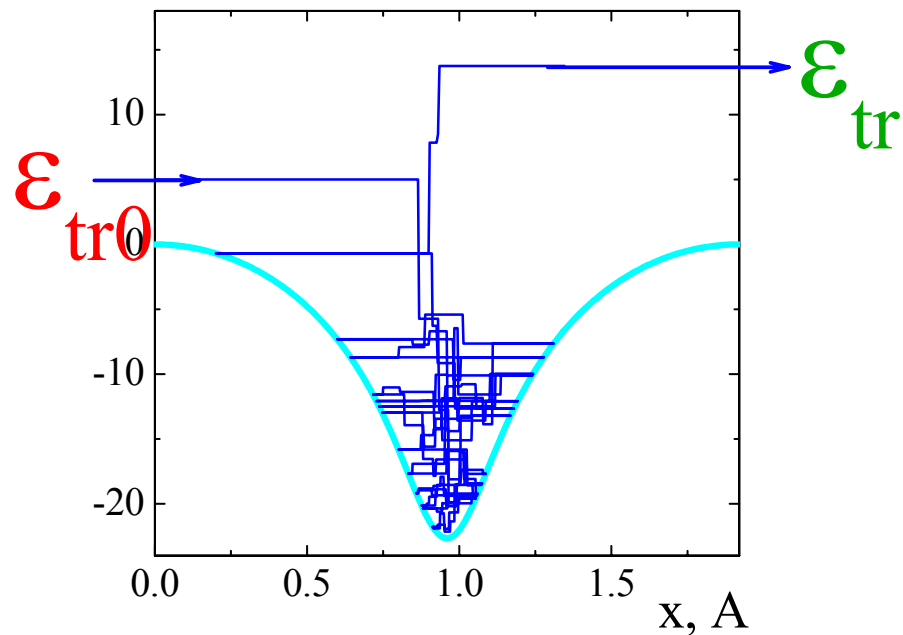
We use below

ε_{tr} – energy of *transverse particle motion*

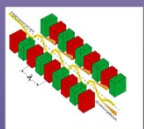
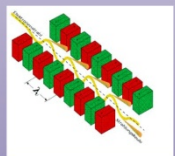
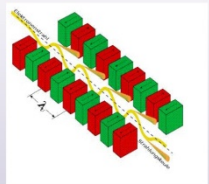
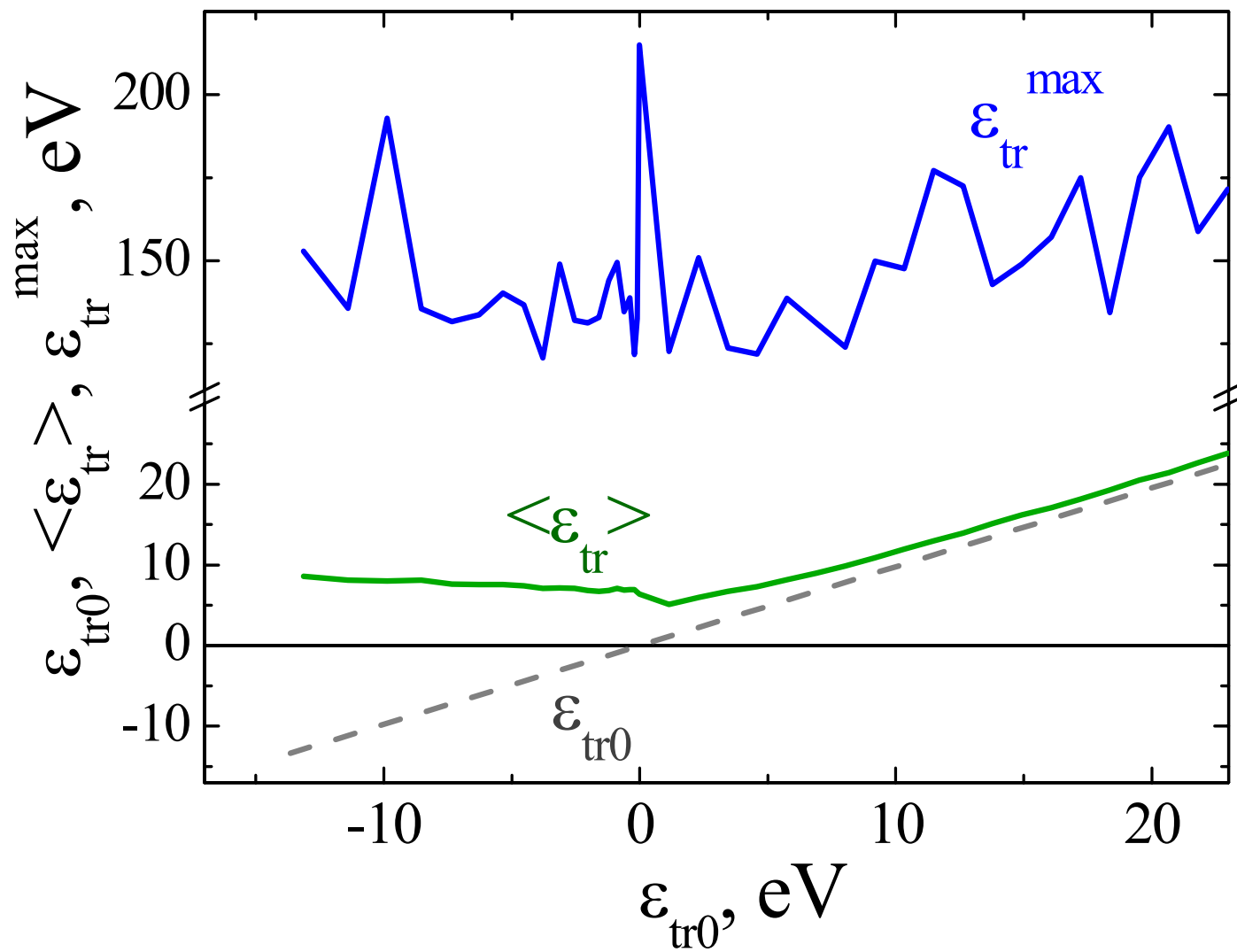
ε_{tr0} – initial _____//_____

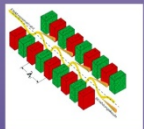
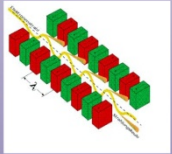
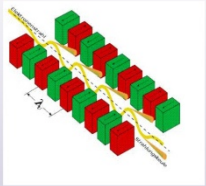
$\langle \varepsilon_{tr} \rangle$ – average _____//_____

ε_{tr}^{\max} – maximal _____//_____

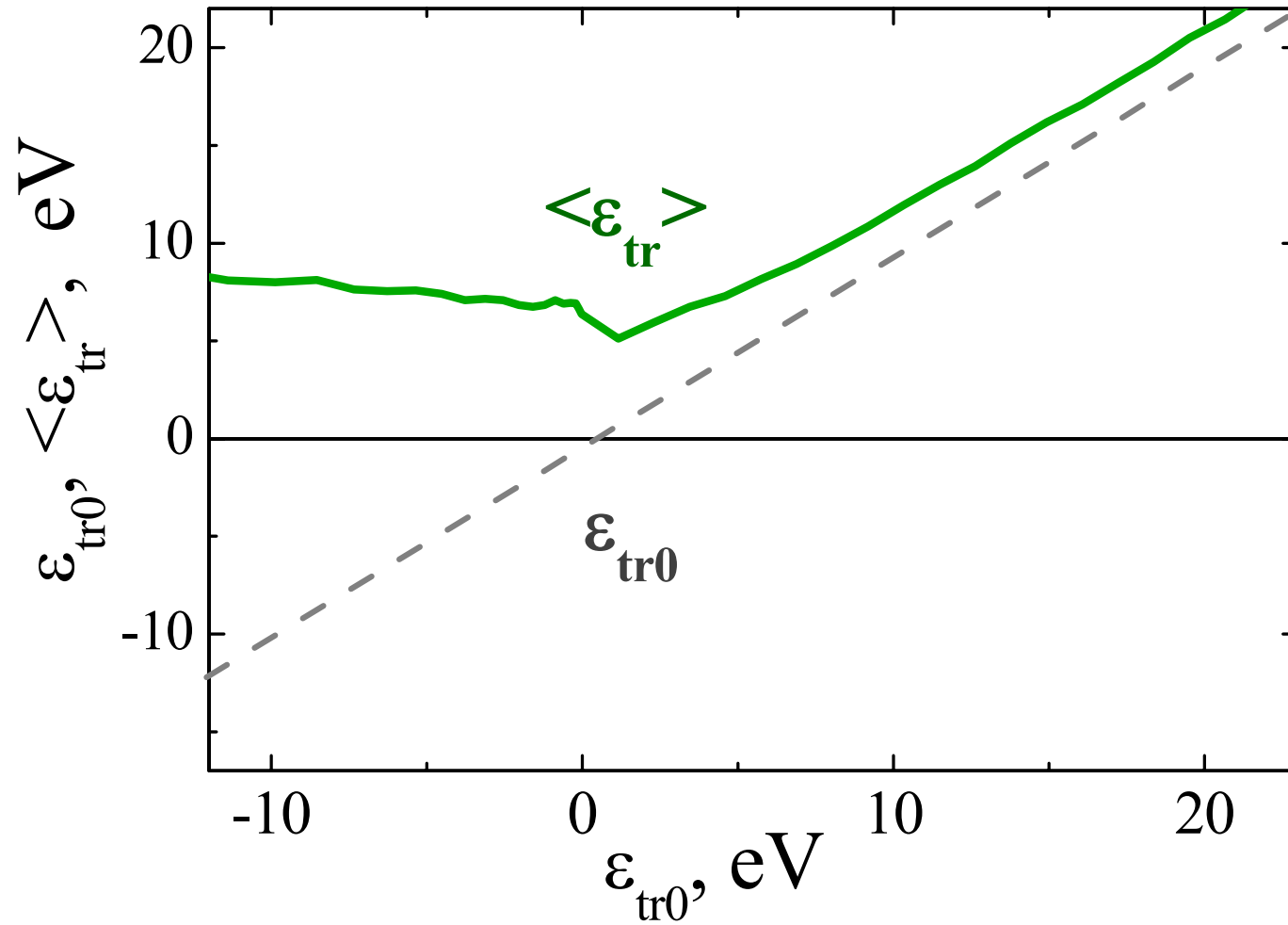


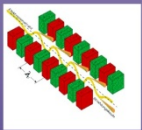
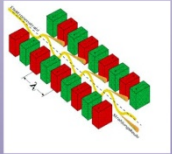
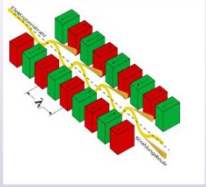
Average and maximum transverse energies



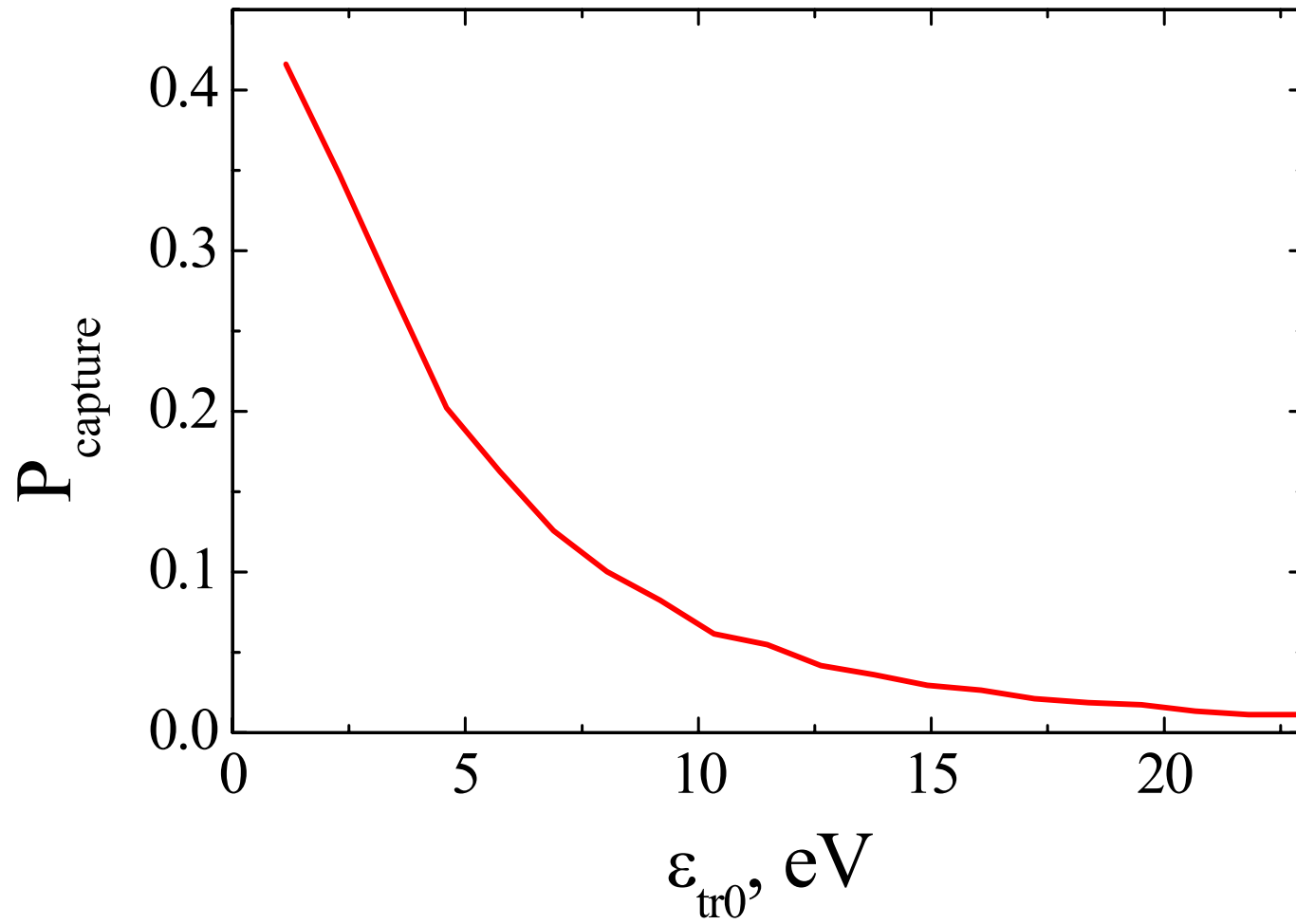


Average and maximum transverse energies

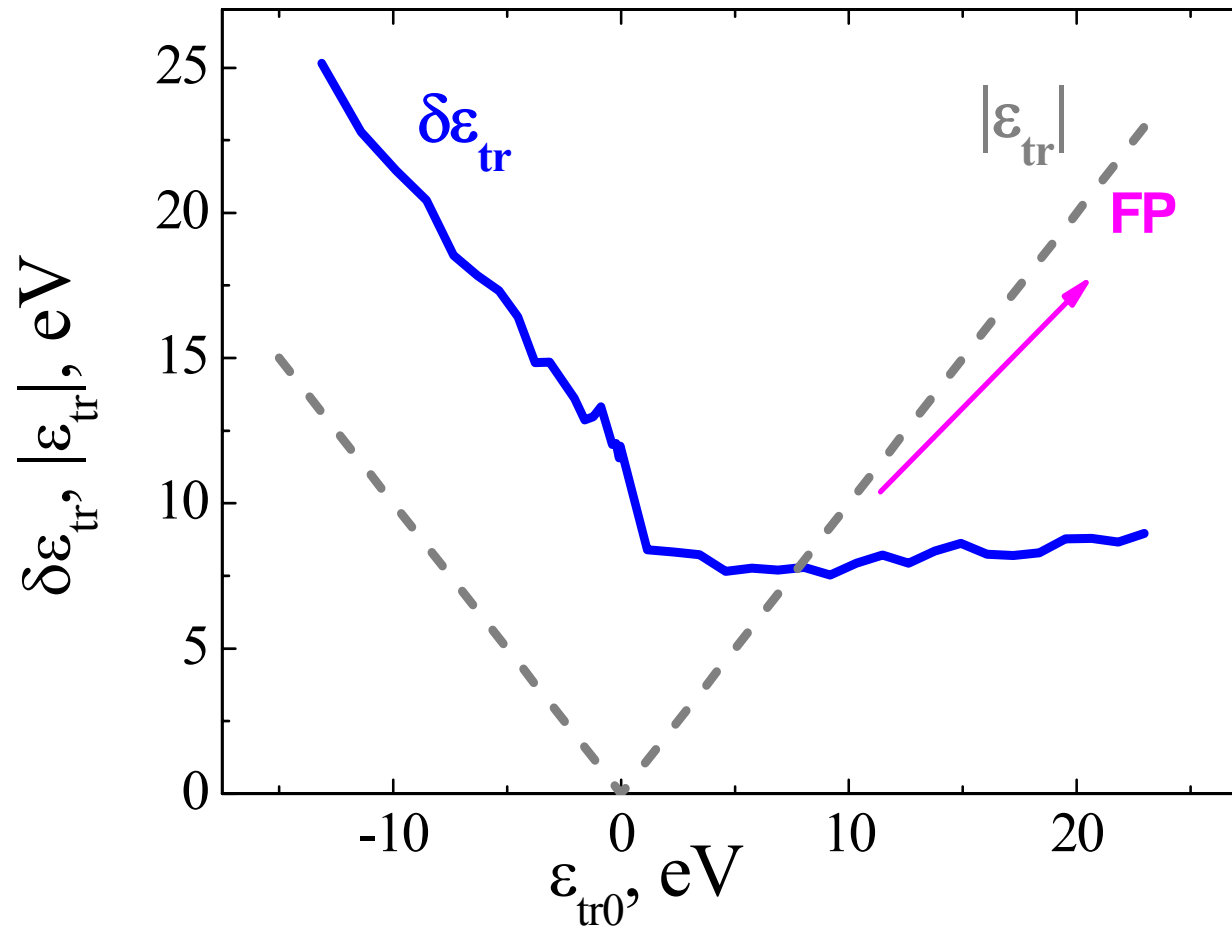




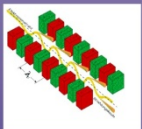
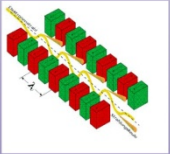
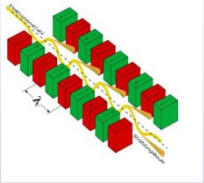
Probability of *immediate* capture of 287 MeV protons

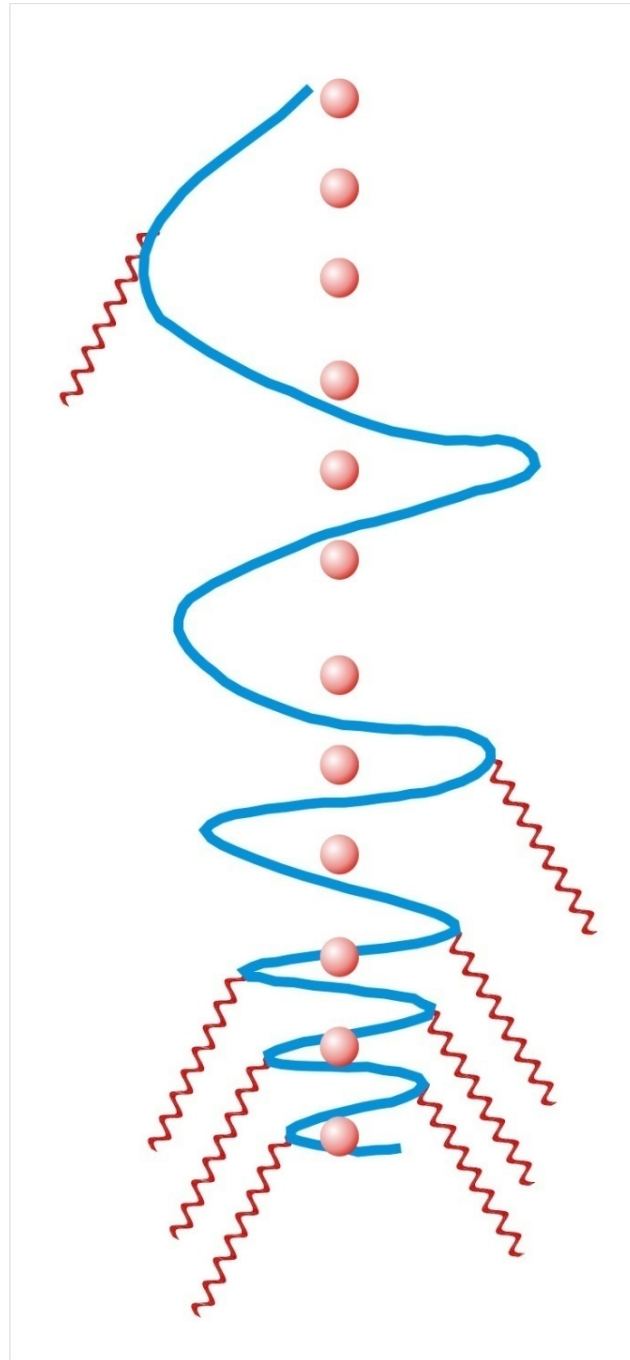
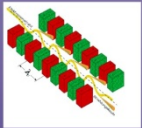
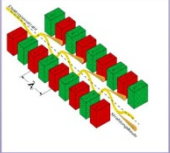
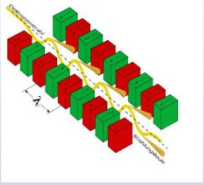


Transverse energy gains large dispersion



Fokker-Planck equation has very limited applicability





Lessons of
radiative
cooling
studies
work
after
25 years!

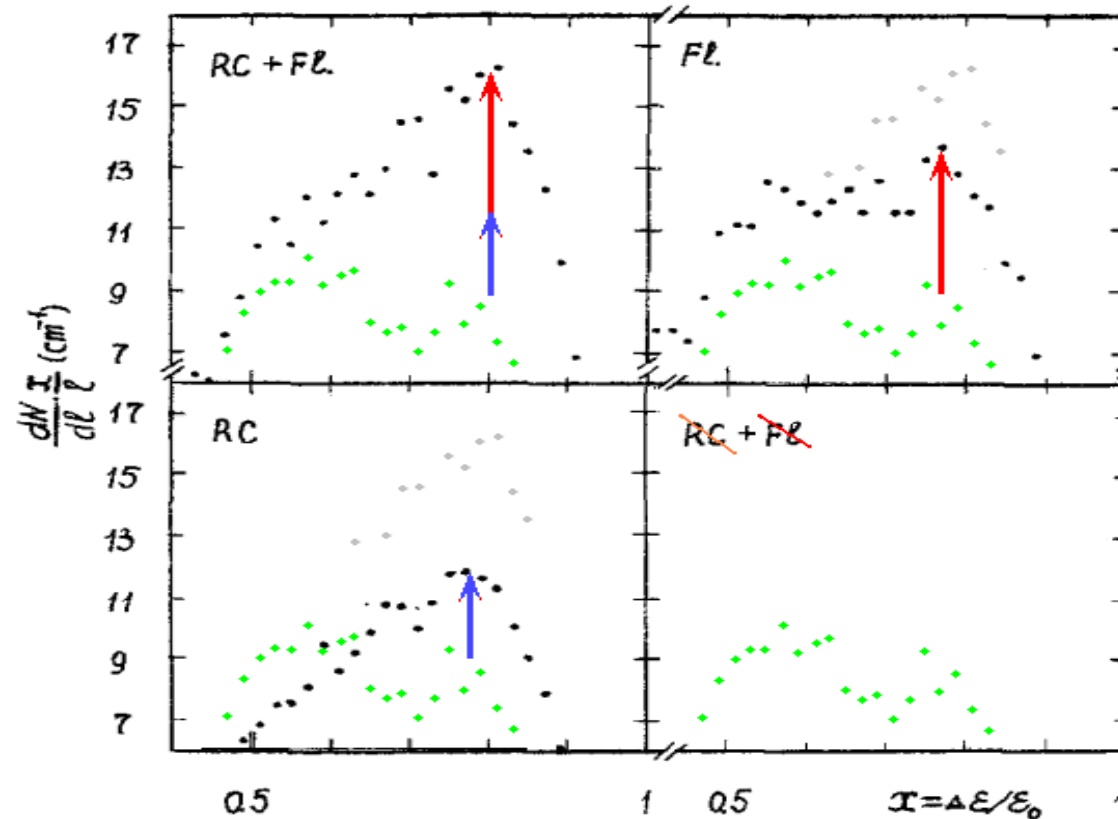
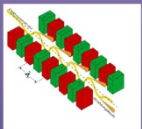
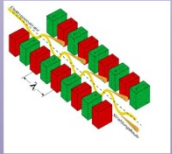
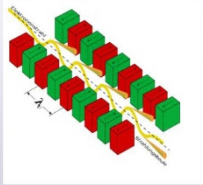
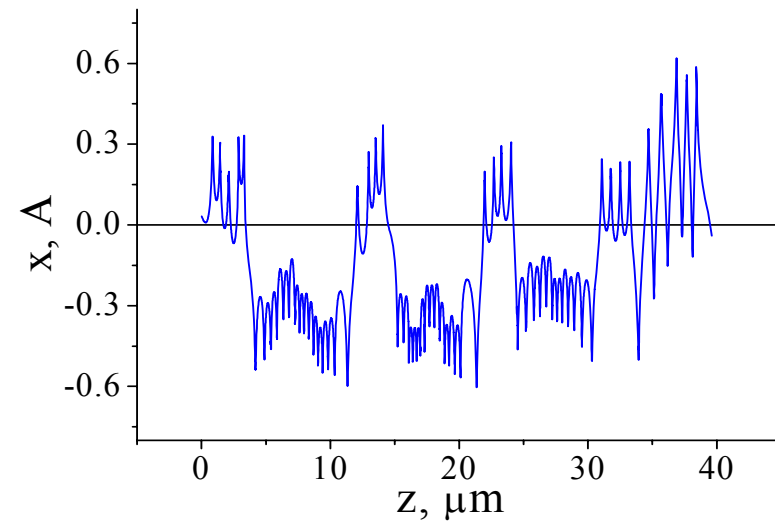
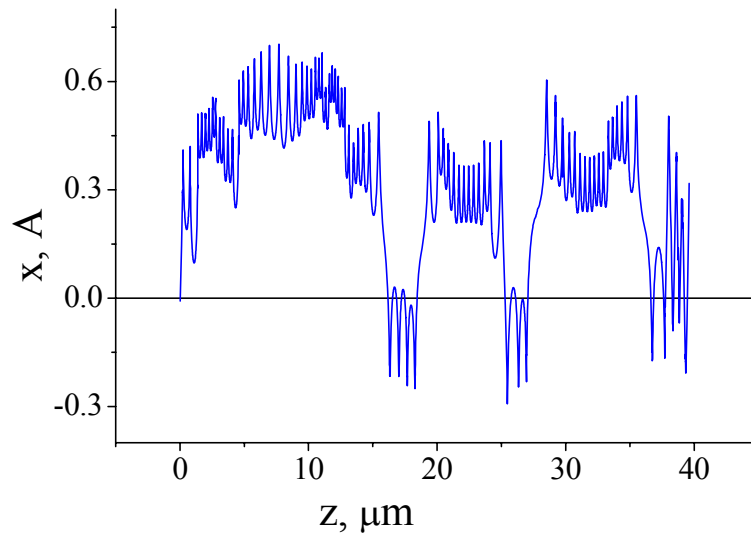
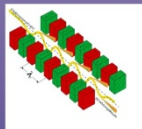
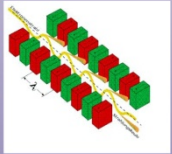
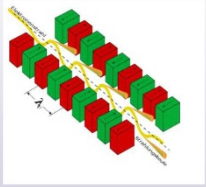


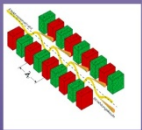
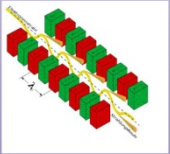
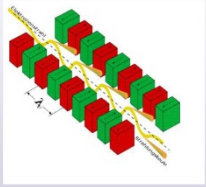
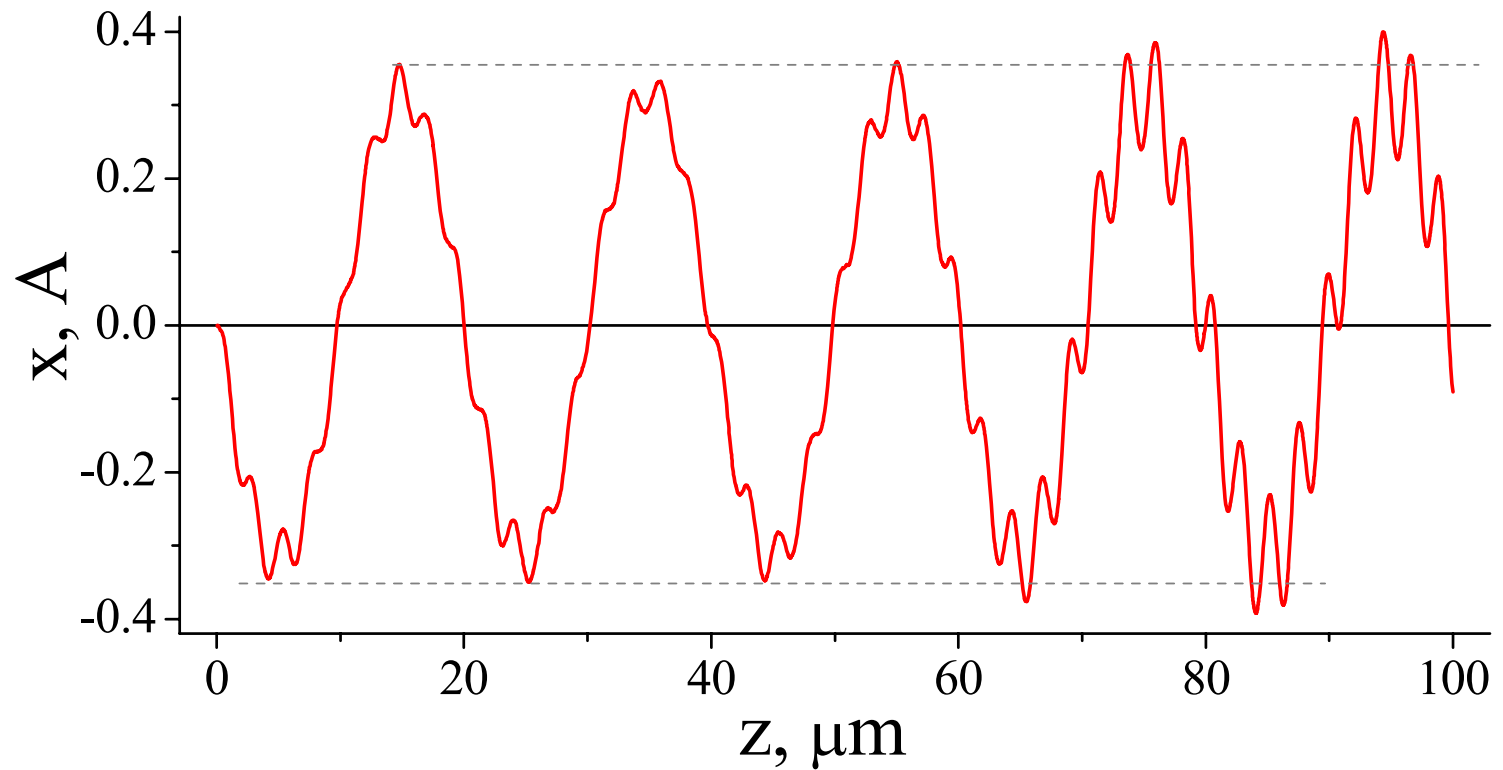
Fig. 2. Simulated energy loss spectra formed by e^- which emit only sufficiently hard γ -quanta with the energies satisfying the inequality (15) under the experimental conditions of ref. [1]. (a) Both the radiation cooling (RC) and ϵ_{\perp} fluctuations (FL) are allowed for; (b) the radiation cooling, respectively (c) the ϵ_{\perp} fluctuations are neglected; (d) both processes are neglected.

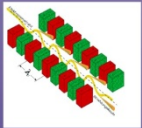
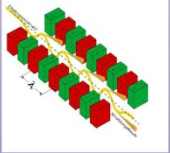
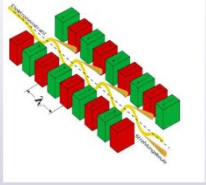
Trajectories of 270 MeV electrons in the $40=4 \times 10 \mu\text{m}$ CU



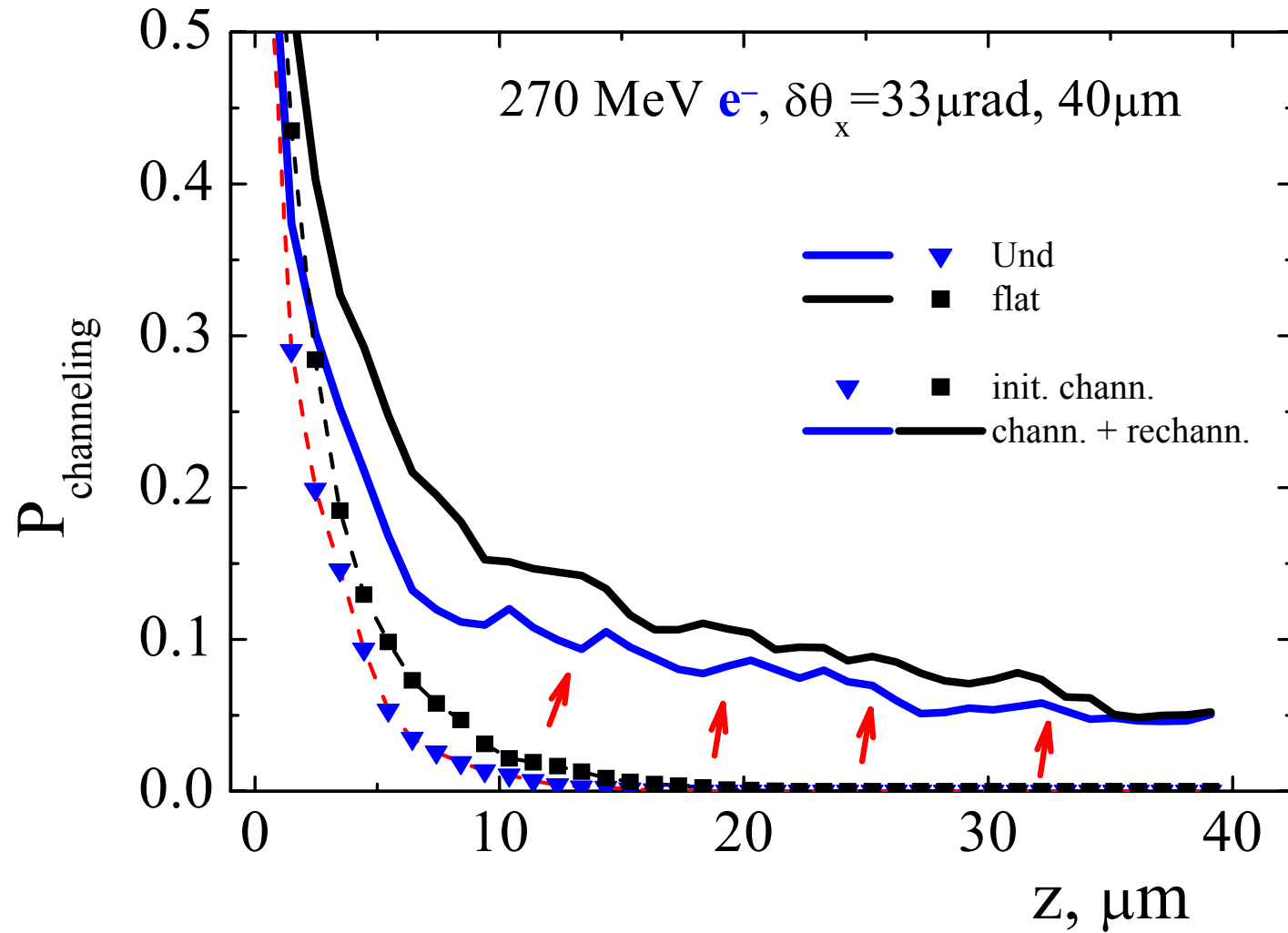
Electrons become channeled several
times for short parts of the period

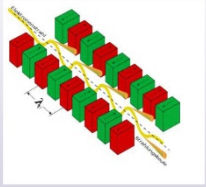
A “fortunate” 500 MeV positron trajectory





The role of RECHANNELING is enormous





Radiation process simulations from the “*First Principles*”

The general expression for radiation intensity

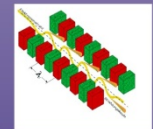
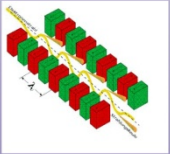
$$\frac{d^2 I}{d\omega d^2\theta} = \frac{\alpha\omega^2 d\omega}{8\pi^2 \epsilon'^2} \times \int \int dt_1 dt_2 \left[(\epsilon^2 + \epsilon'^2)(\mathbf{v}_\perp(t_1) - \boldsymbol{\theta})(\mathbf{v}_\perp(t_2) - \boldsymbol{\theta}) + \omega^2/\gamma^2 \right] \\ \exp \left\{ i \frac{\omega\epsilon}{2\epsilon'} \left[\int_{-\infty}^{t_1} (\gamma^{-2} + (\mathbf{v}_\perp(t') - \boldsymbol{\theta})^2) dt' + \int_{-\infty}^{t_2} (\gamma^{-2} + (\mathbf{v}_\perp(t'') - \boldsymbol{\theta})^2) dt'' \right] \right\}$$

contains two integrals

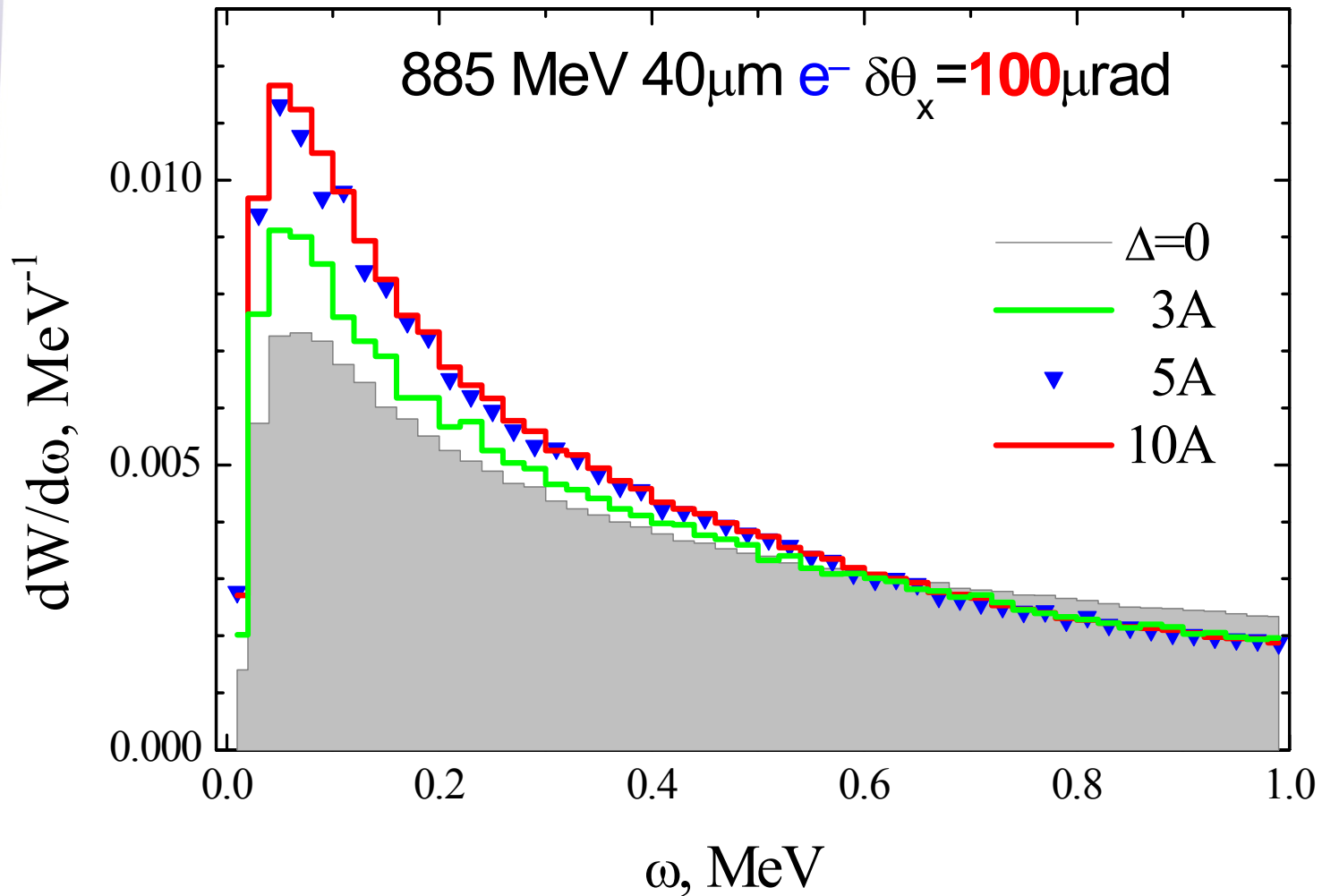
$$A = \int \exp \left\{ i \frac{\omega\epsilon}{2\epsilon'} \int_{-\infty}^t [\gamma^{-2} + (\mathbf{v}_\perp(t') - \boldsymbol{\theta})^2] dt' \right\} dt,$$

$$\mathbf{B} = \int (\mathbf{v}_\perp(t) - \boldsymbol{\theta}) \exp \left\{ i \frac{\omega\epsilon}{2\epsilon'} \int_{-\infty}^t [\gamma^{-2} + (\mathbf{v}_\perp(t') - \boldsymbol{\theta})^2] dt' \right\} dt$$

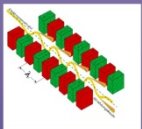
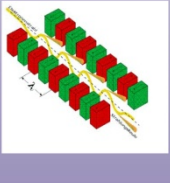
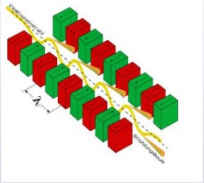
and slowly decreases with radiation angle $\boldsymbol{\theta}$, complicating its numerical integration.



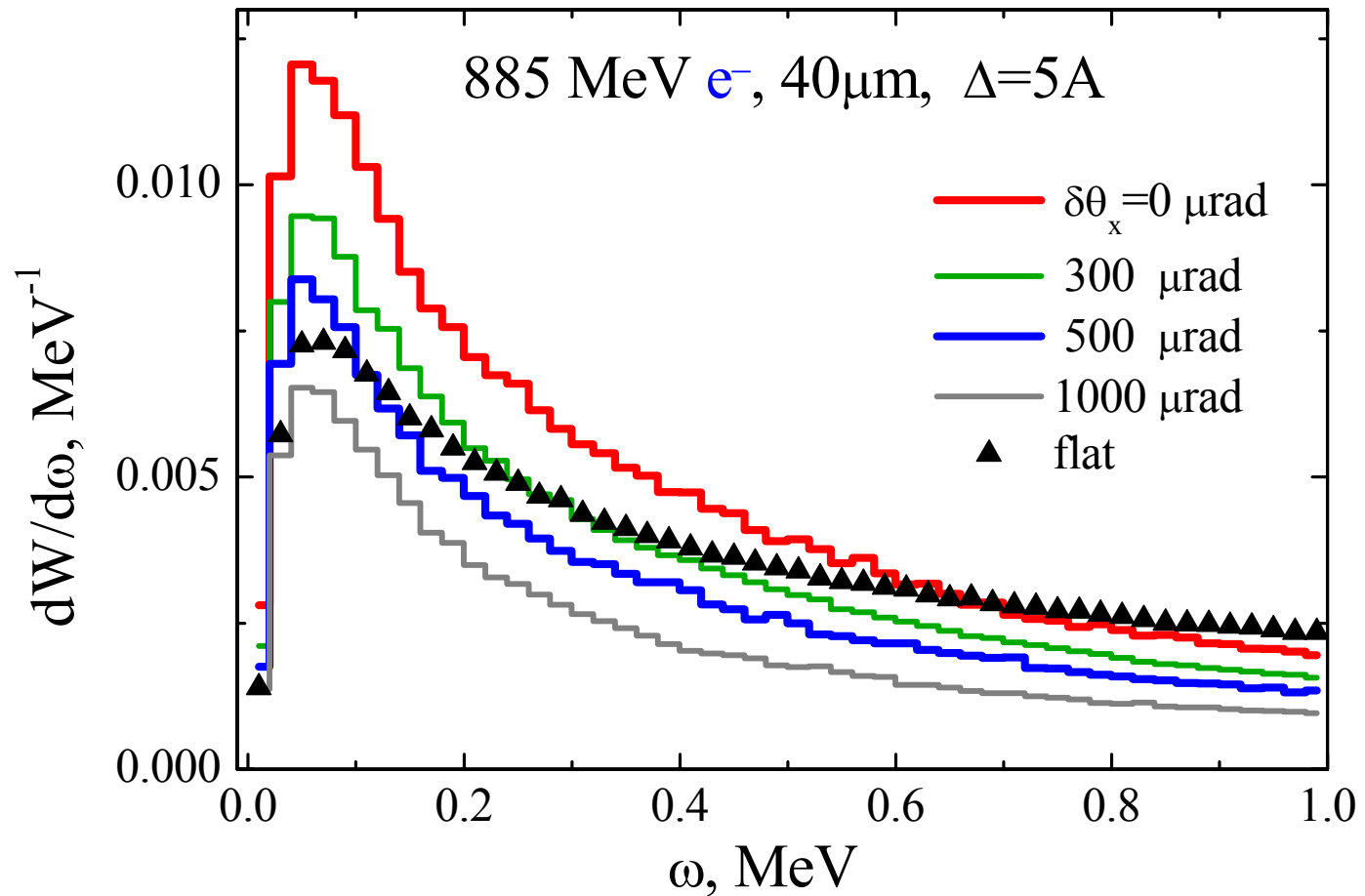
CU intensity dependence on CU modulation amplitude



Optimal modulation amplitude is about **5 Angstroms**



CU intensity dependence on beam divergence



Optimal beam divergence is about **100 μrad**

Plan

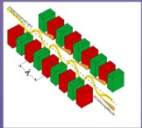
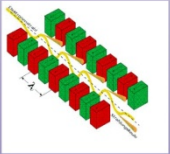
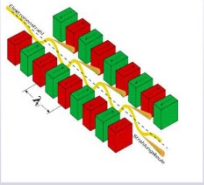
CU radiation prediction and ultimate
perspective of its application

CU radiation simulation

experiments with electrons

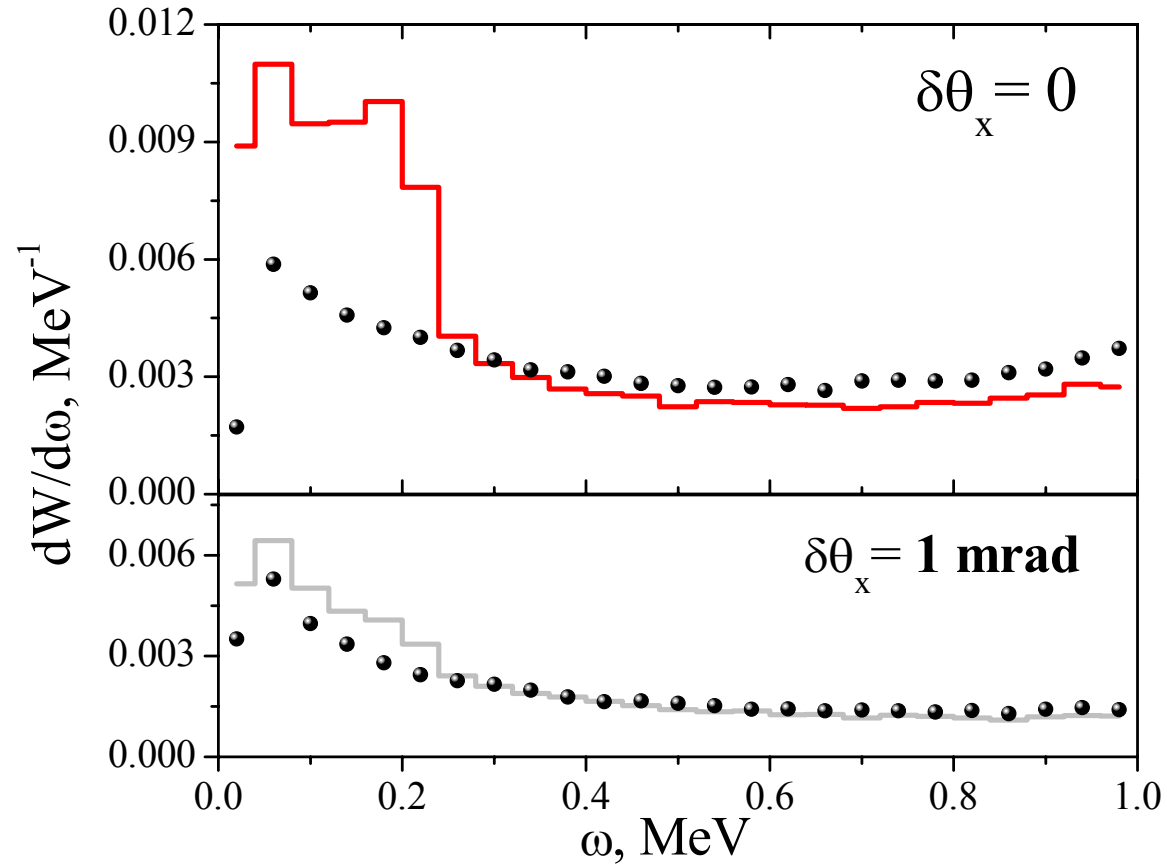
predictions for positrons

CU radiation improvement
by crystal cut



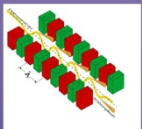
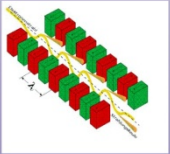
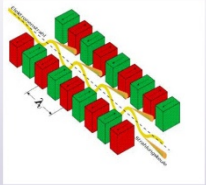
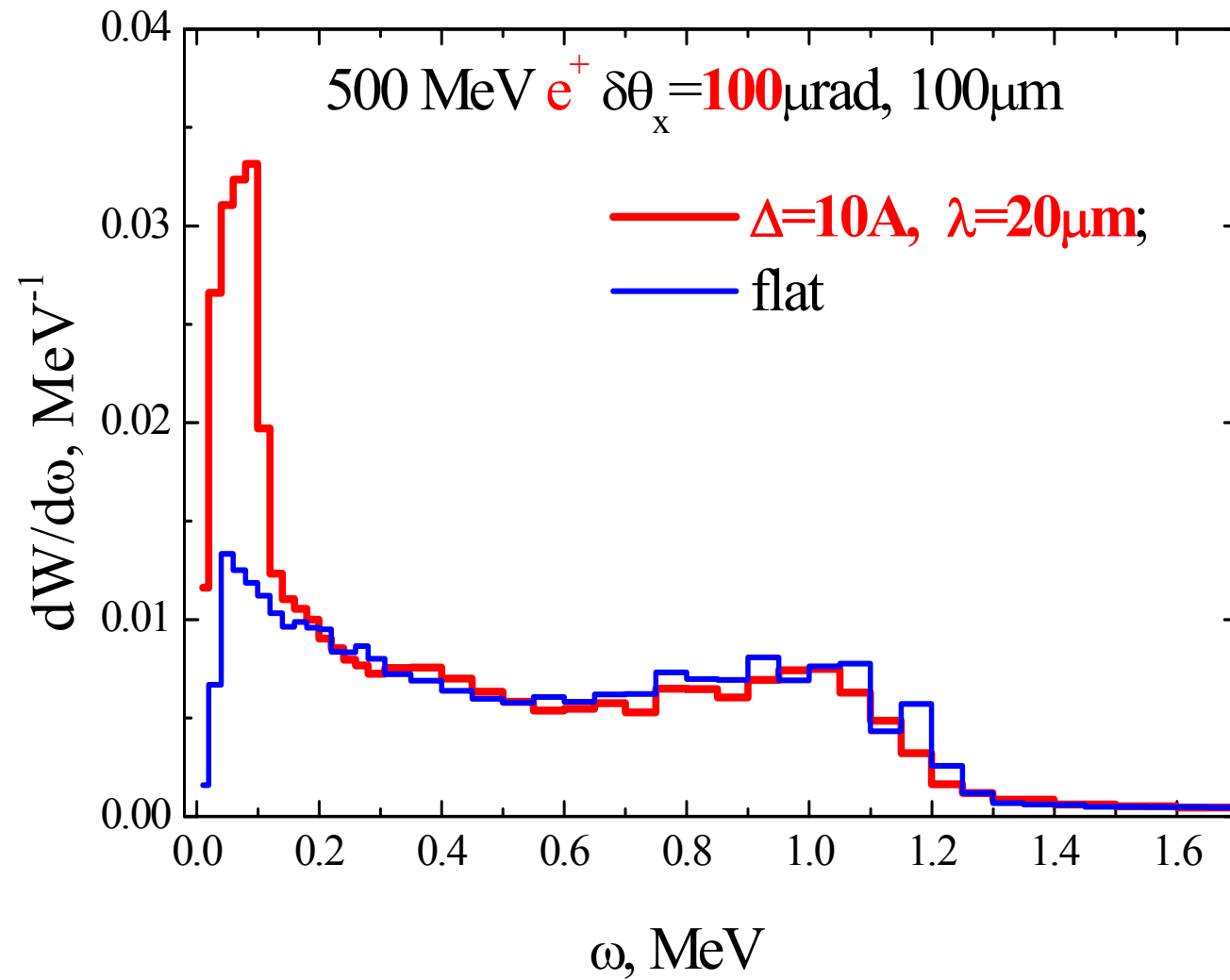
Positron CU radiation

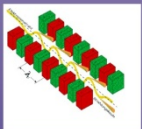
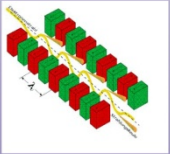
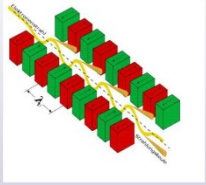
500 MeV Positrons in — 4x9.9 μ m Undulator, • flat crystal



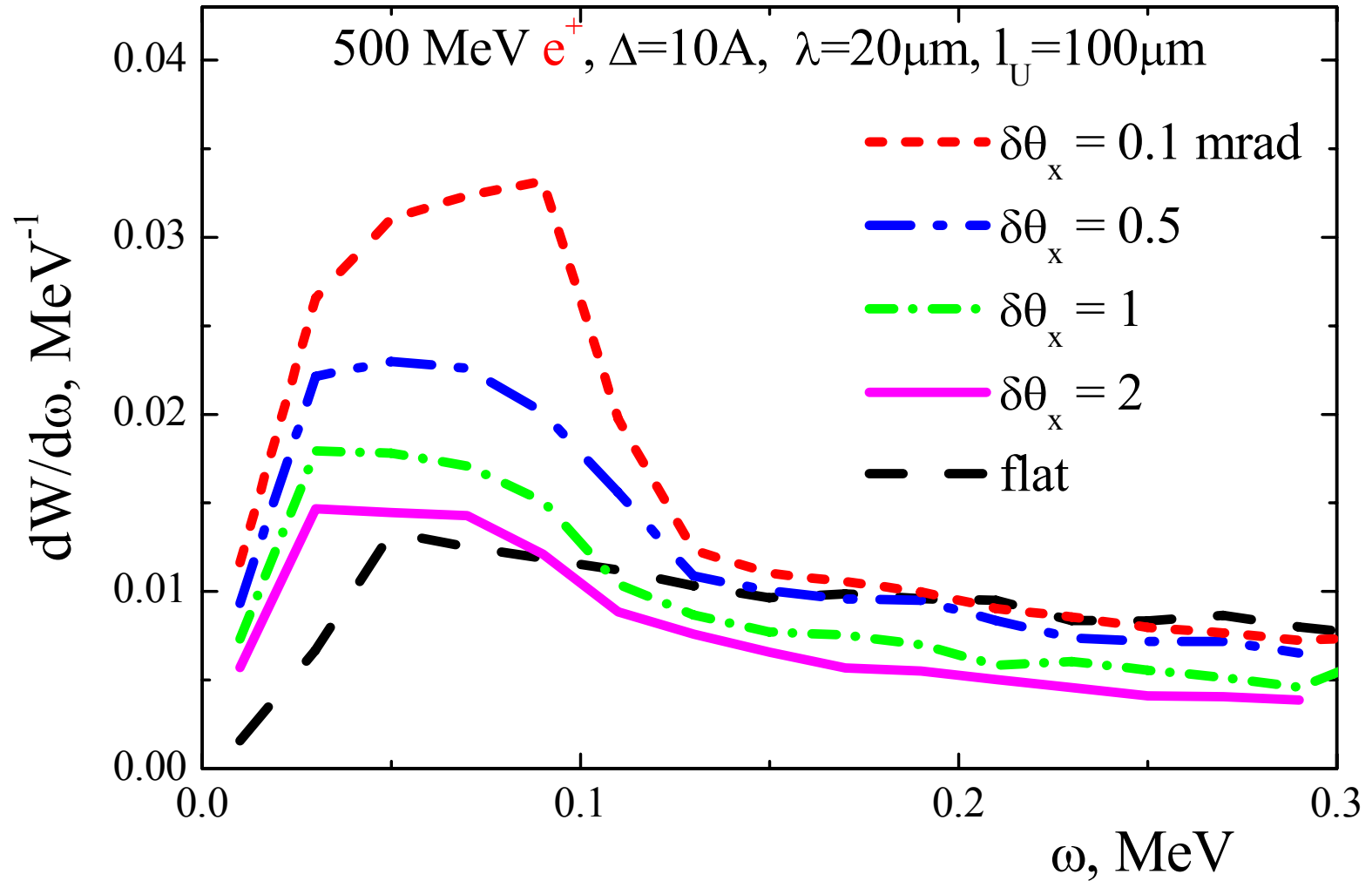
Positron CU radiation have good perspectives

Radiation in more effective CU





CU intensity dependence on beam divergence



Positron beam divergence can be in **milliradian range**

Plan

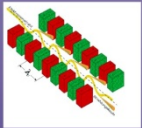
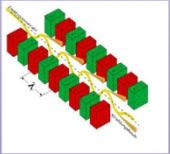
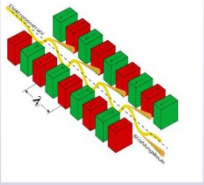
CU radiation prediction and ultimate perspective of its application

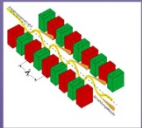
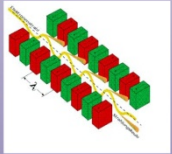
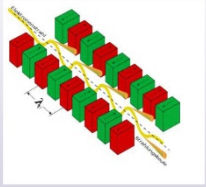
CU radiation simulation

experiments with electrons

predictions for positrons

**CU radiation improvement
by crystal cut**

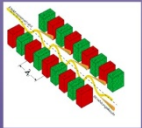
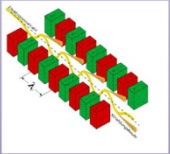
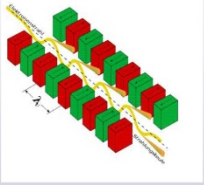
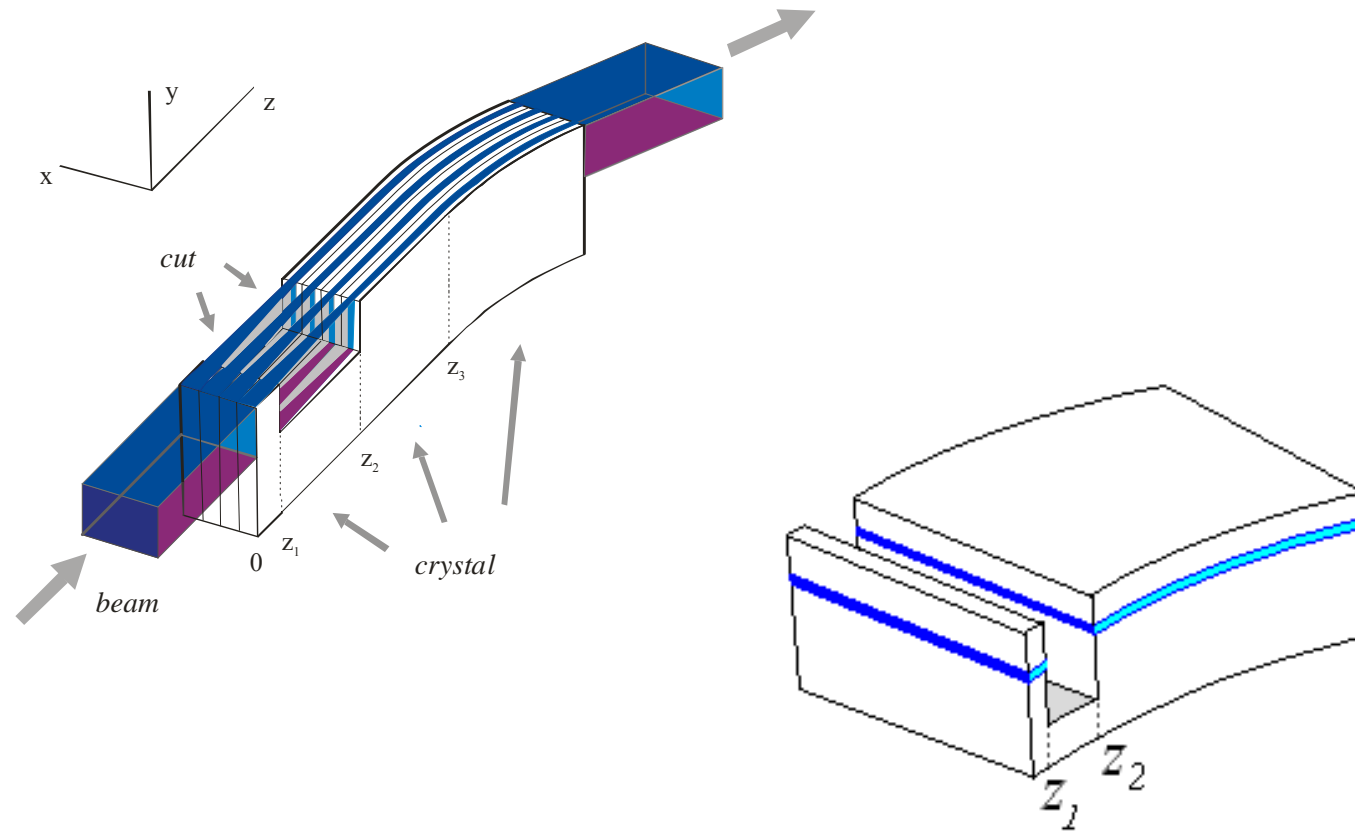




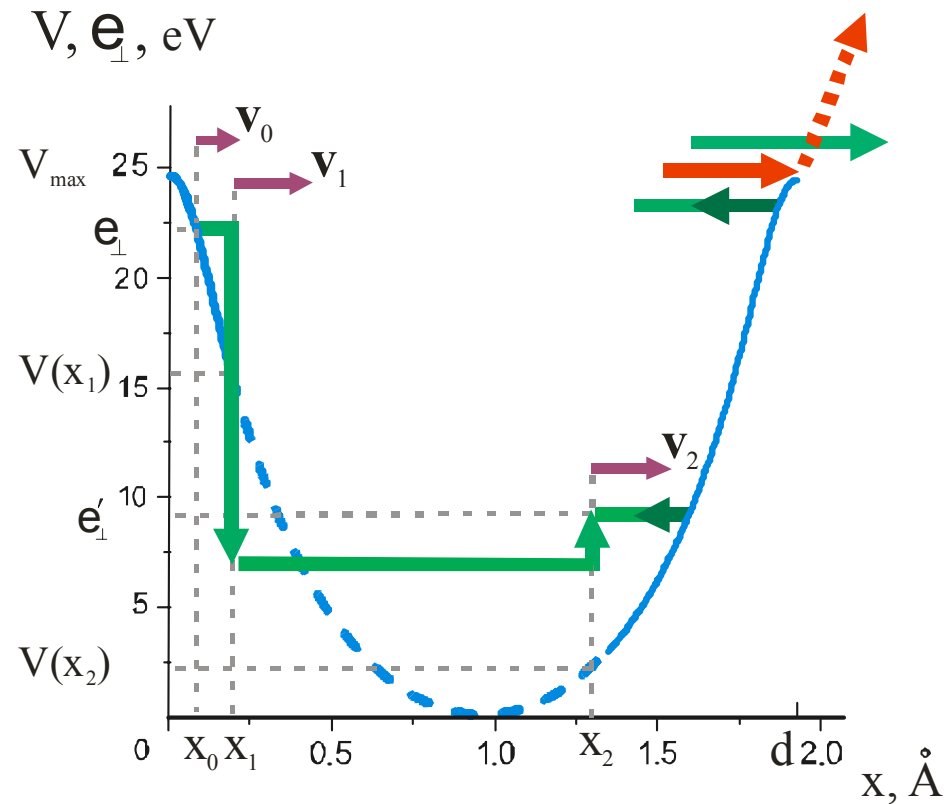
An idea for future development

The capture probability increase by crystal cut

V.V. Tikhomirov, *JINST*, 2(2007)P08006

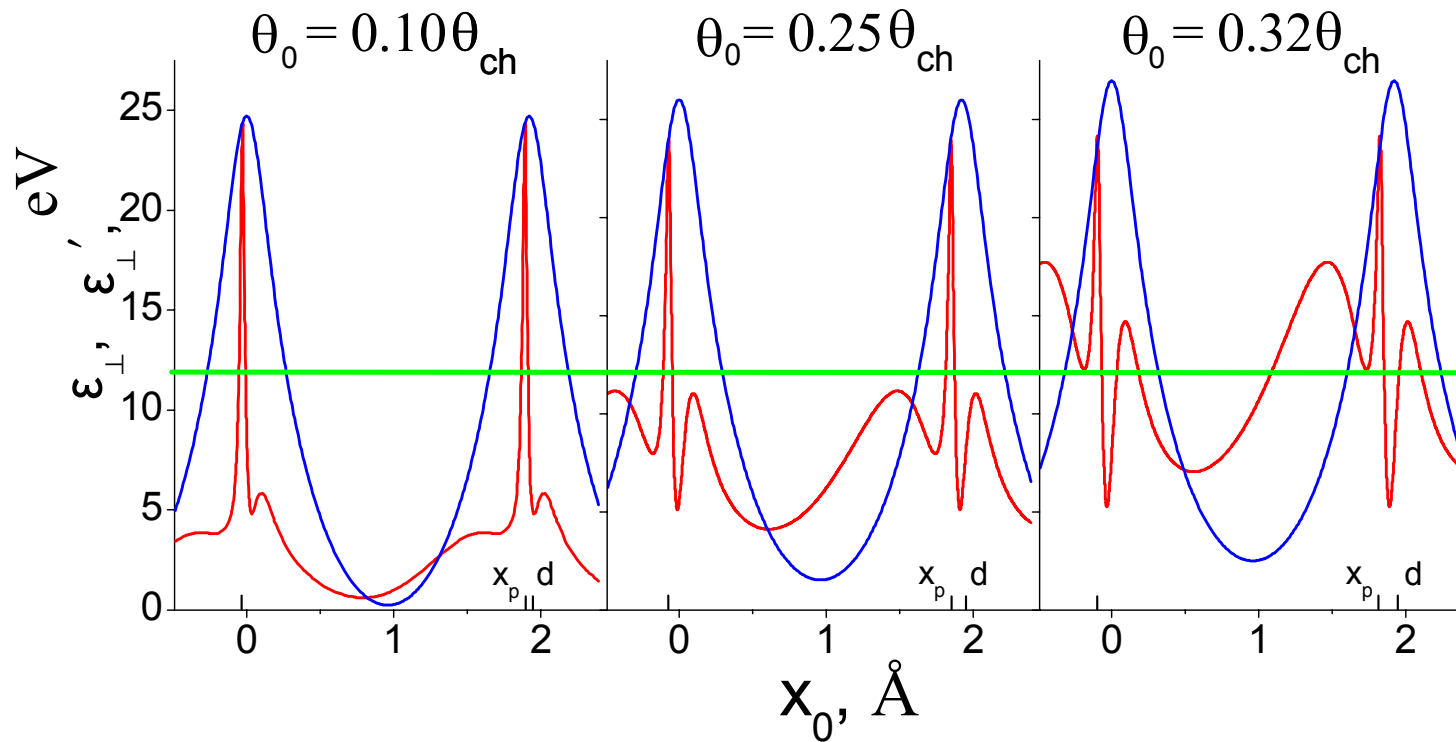


Transverse energy reduction *by the cut* - 1

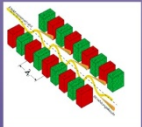
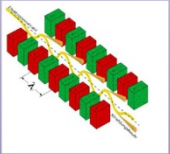
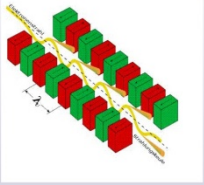


The cut diminishes the potential energy conserving the transverse kinetic one

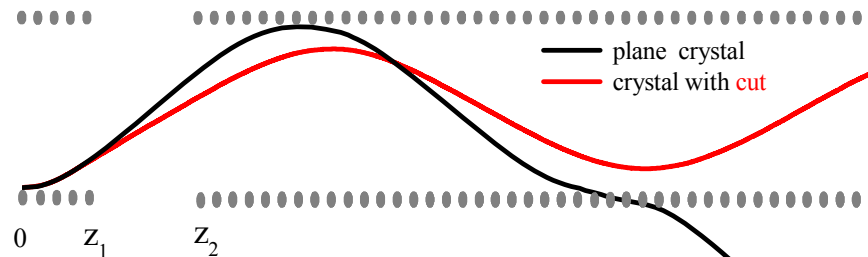
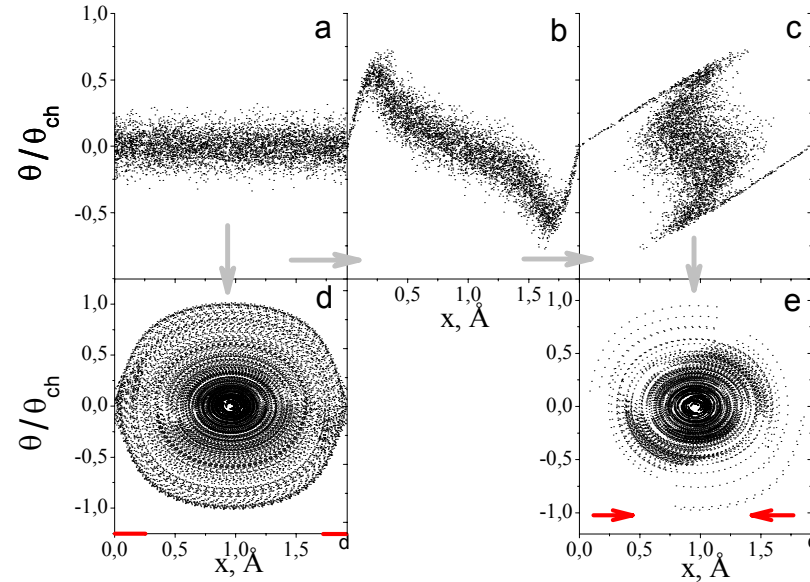
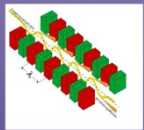
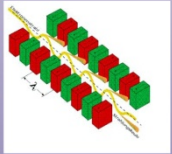
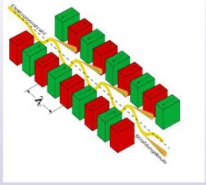
Transverse energy reduction *by the cut* - 2



Only **1-2%** of protons avoid drastic transverse energy reduction by the cut

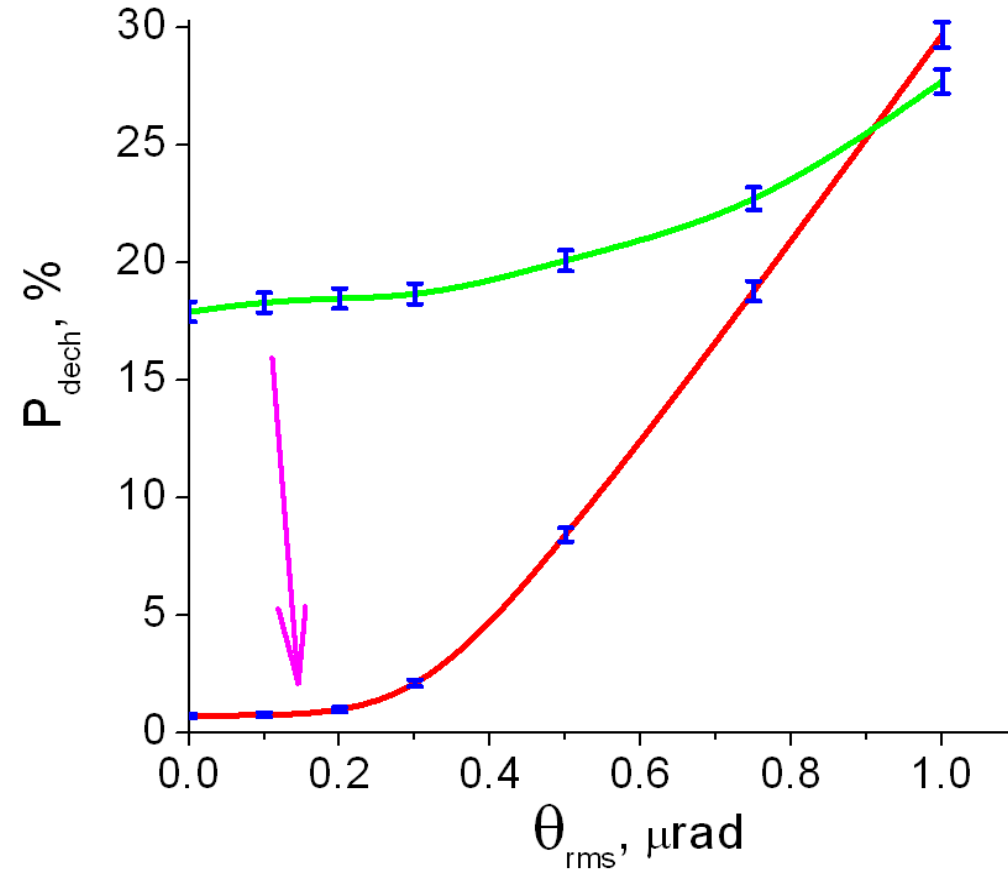
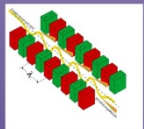
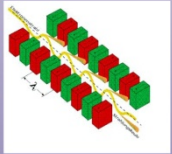
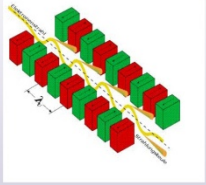


Phase space transformation by the cut

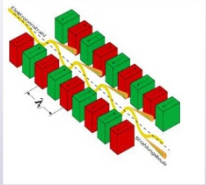


Protons cease to reach the high nuclear density **regions**

Channeling fraction increase by the cut



The cut increases channeling fraction **from 85 to 99%**



Cut formation method

(110) Silicon Etching for High Aspect Ratio Comb Structures

*Seong-Hyok Kim, Sang-Hun Lee, Hyung-Taek Lim, and Yong-Kweon Kim

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Seung-Ki Lee

Department of Electrical Engineering, Dankook University

8, Hannam-dong, Yongsan-gu, Seoul 140-714, Korea

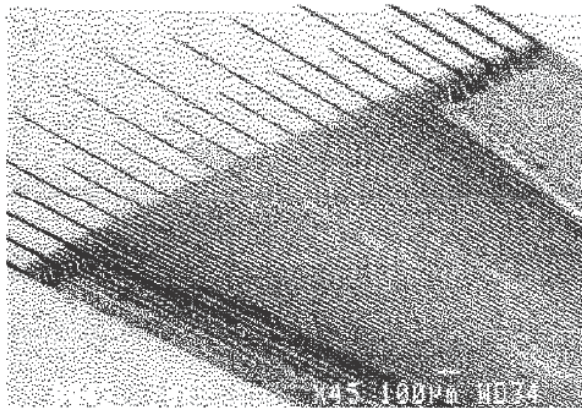
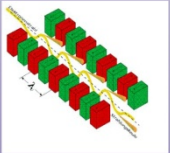


Fig 1. SEM photograph of alignment target after wet etching

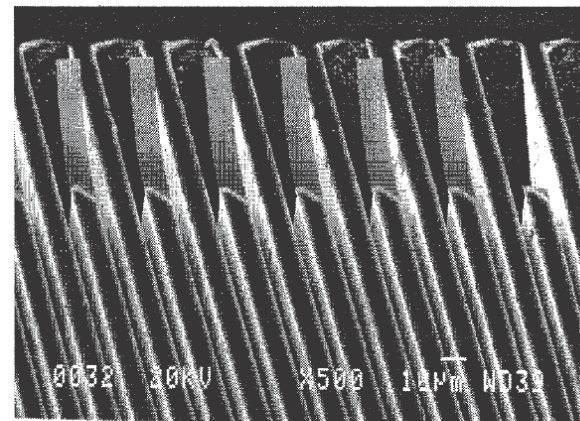
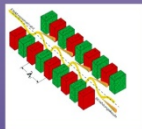
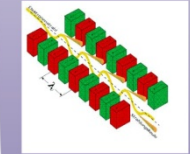
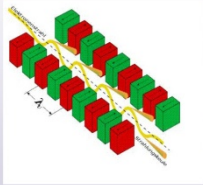


Fig. 12 Fabricated comb structures.
The width is $8\mu\text{m}$, gap is $7\mu\text{m}$ and height is about $150\mu\text{m}$

Crystal cut can be produced by
anisotropic etching





Ultimate Top-down Etching Processes for Future Nanoscale Devices: Advanced Neutral-Beam Etching

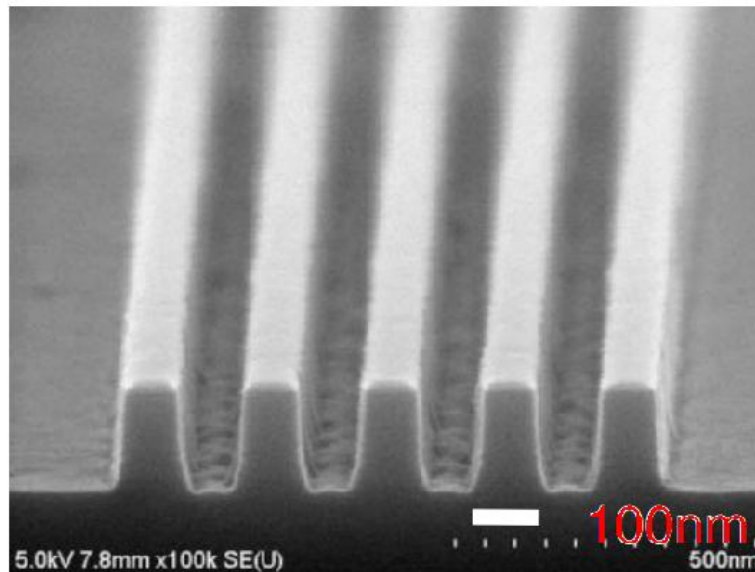
Seiji SAMUKAWA*

Institute of Fluid Science, Tohoku University, 2-1-1 Katahira, Aoba-ku, Sendai 980-8577, Japan

(Received August 4, 2005; accepted November 15, 2005; published online April 7, 2006)

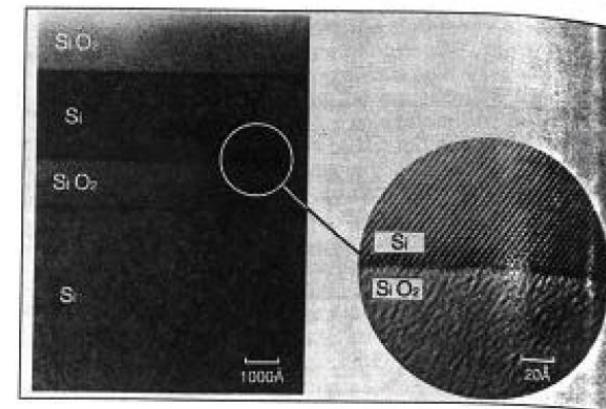
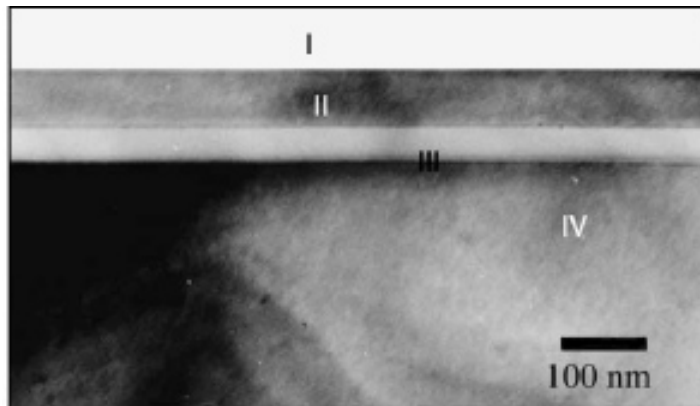
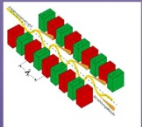
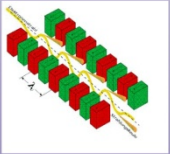
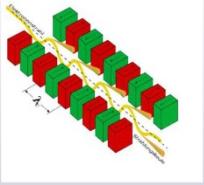
For the past 30 years, plasma etching technology has led in the efforts to shrink the pattern size of ultralarge-scale integrated (ULSI) devices. However, inherent problems in the plasma processes, such as charge buildup and UV photon radiation, limit the etching performance for nanoscale devices. To overcome these problems and fabricate sub-10-nm devices in practice, neutral-beam etching has been proposed. In this paper, I introduce the ultimate etching processes using neutral-beam sources and discuss the fusion of top-down and bottom-up processing for future nanoscale devices. Neutral beams can perform atomically damage-free etching and surface modification of inorganic and organic materials. This technique is a promising candidate for the practical fabrication technology for future nano-devices. [DOI: 10.1143/JJAP.45.2395]

KEYWORDS: top-down process, plasma etching, radiation damage, neutral beam, sub-10nm patterning



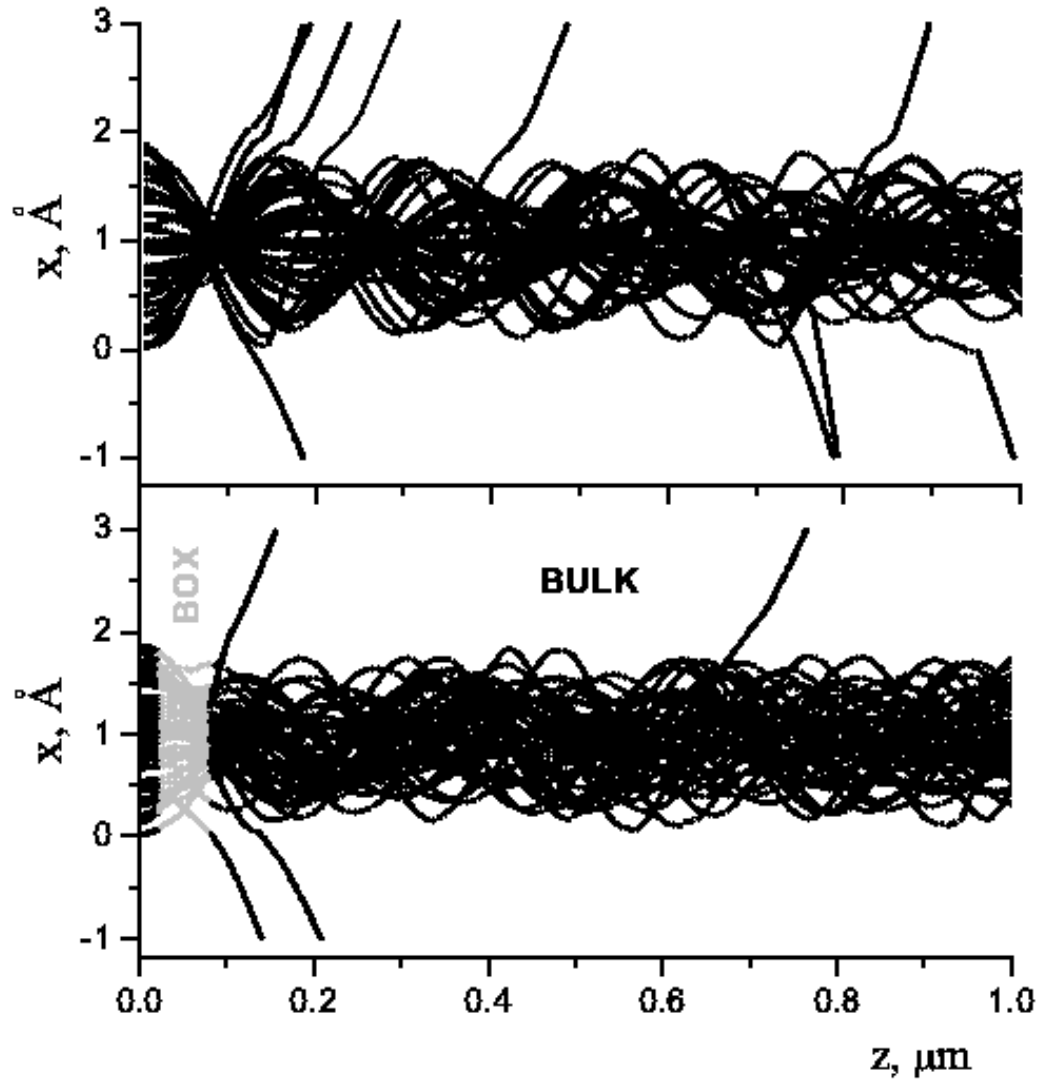
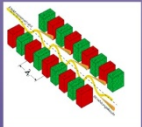
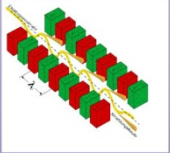
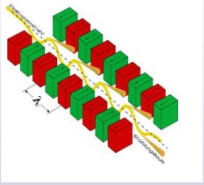
SIMOX Buried Oxide Layer can be used instead of crystal cut

V. Guidi, A. Mazzolari and V.V. Tikhomirov,
J. Phys. D: Appl. Phys. 42(2009) 165301



- Thermal annealing restores silicon crystalline quality and creates a buried SiO₂ layer,
- Interfaces between Si and SiO₂ are well terminated,
- Misalignment between silicon layers in available SIMOX structures: less than 0.7 Å/mm.





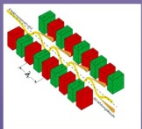
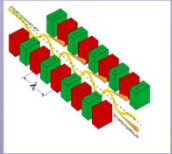
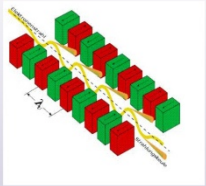
BOX layer
“focuses”
protons
like a **cut**
diminishing
their
transverse
energy

$$z_1 = 20 \text{ nm},$$

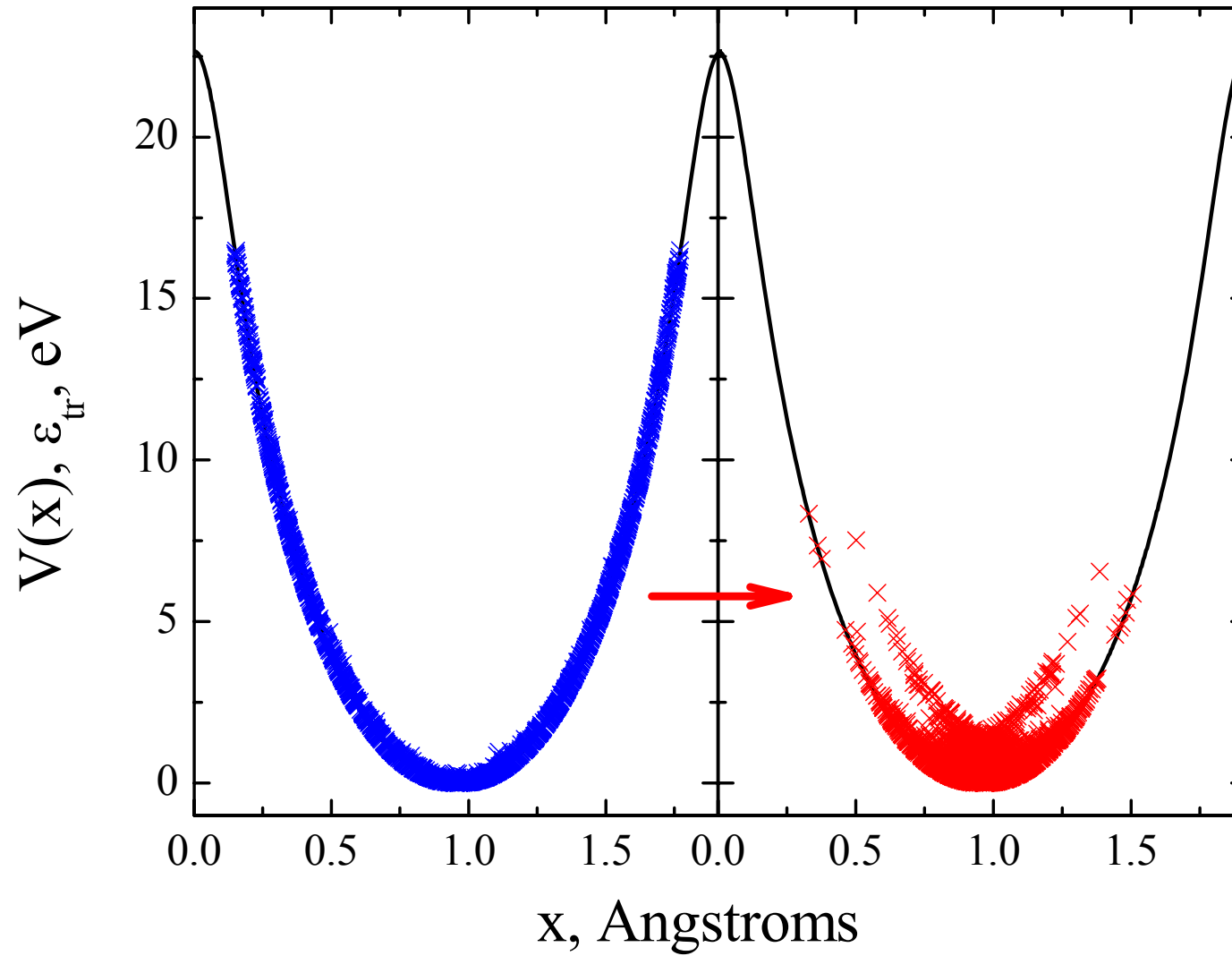
$$z_2 = 80 \text{ nm},$$

$$z_3 = 1 \text{ } \mu\text{m},$$

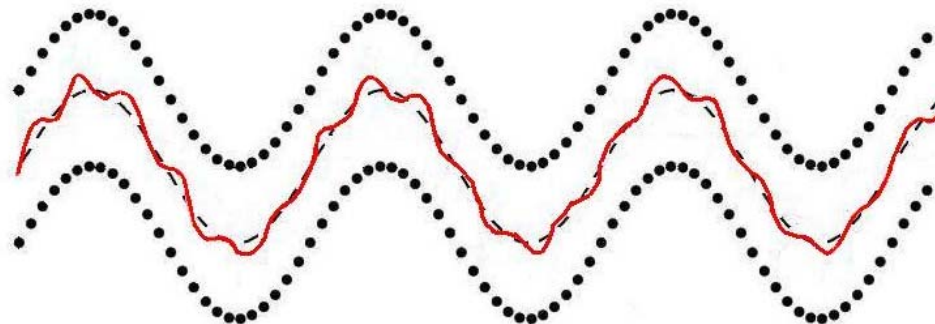
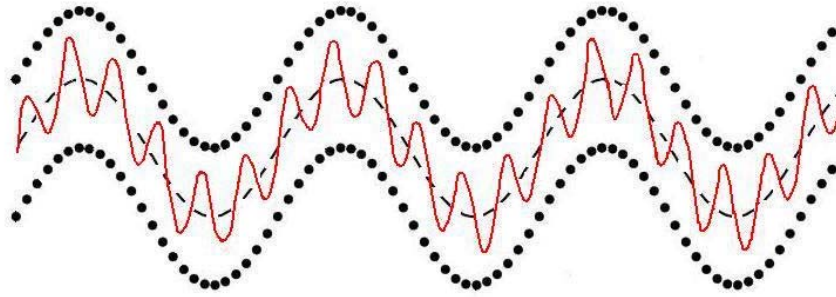
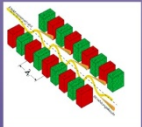
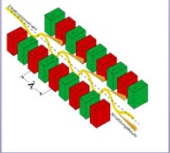
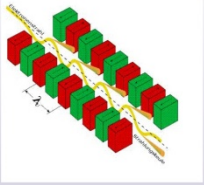
$$E_p = 7 \text{ MeV}$$

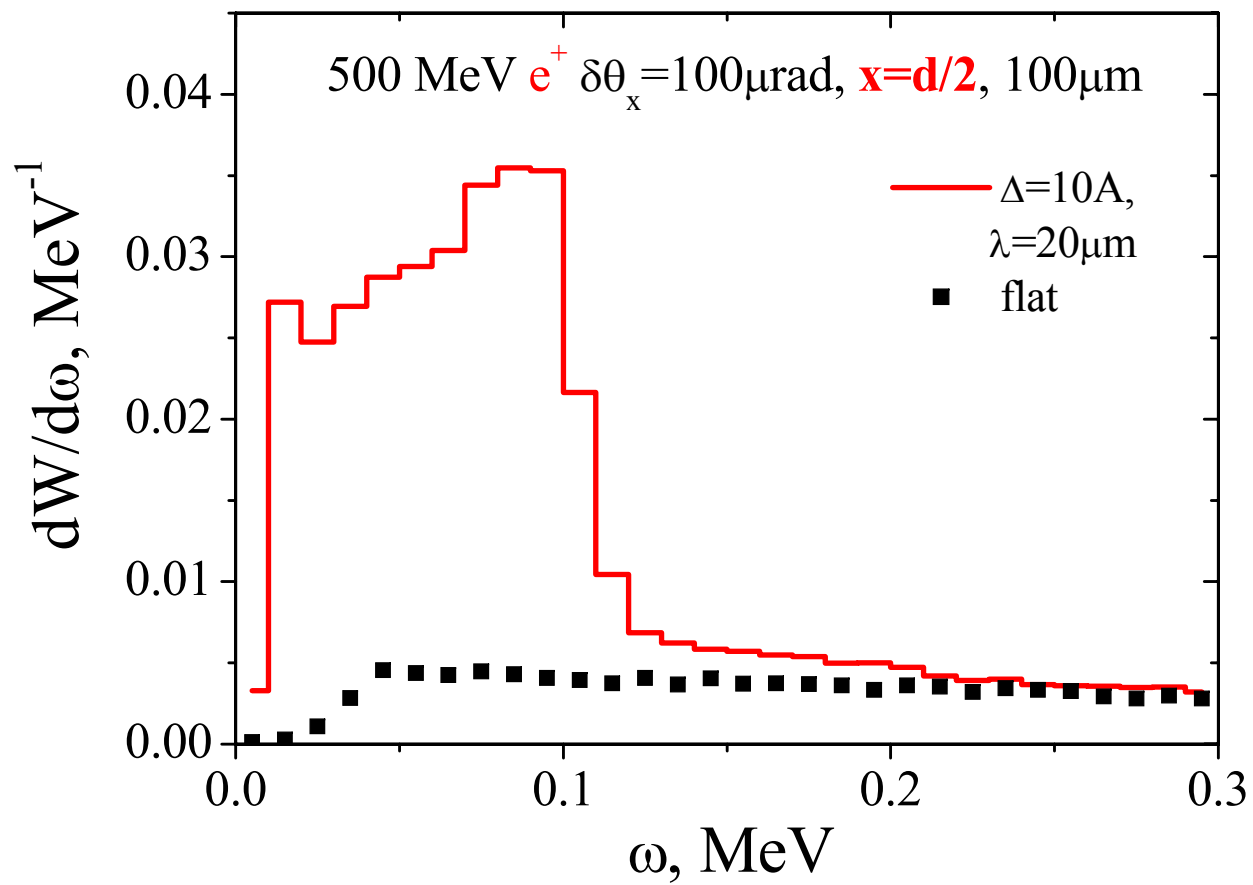
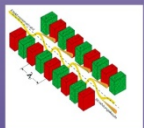
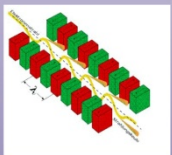
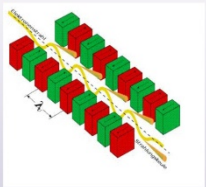


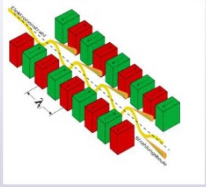
transverse energy reduction by **cut**



The **cut** allows to diminish channeling oscillations:



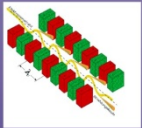
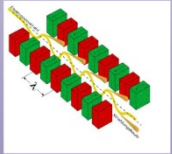




Essential CU improvement by diminishing of channeling oscillations by the crystal **cut**

- *increase of number of particles*
- *increase of radiation intensity*
- *decrease of undulator period*
- *narrowing of radiation spectrum*
- *elimination of energy losses
for channeling radiation*

**A real way to both
 γ -laser and high positron energies**



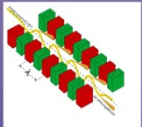
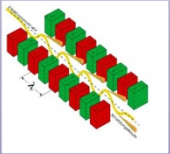
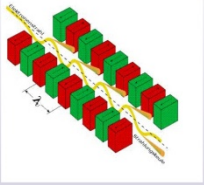
Conclusions

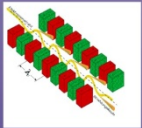
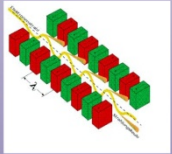
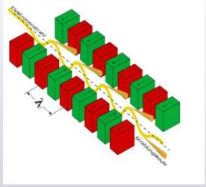
Electron radiation in CU represents a good testbed for simulation methods

Positron radiation in CU can be observed under reasonable beam quality

Crystal cut can improve the CU functioning in numerous directions

Positron CU radiation + photon diffraction open up a way to γ -laser





Thank you for attention!

