

Radiation and Particle Secondary Sources based on electron Laser Plasma Acceleration

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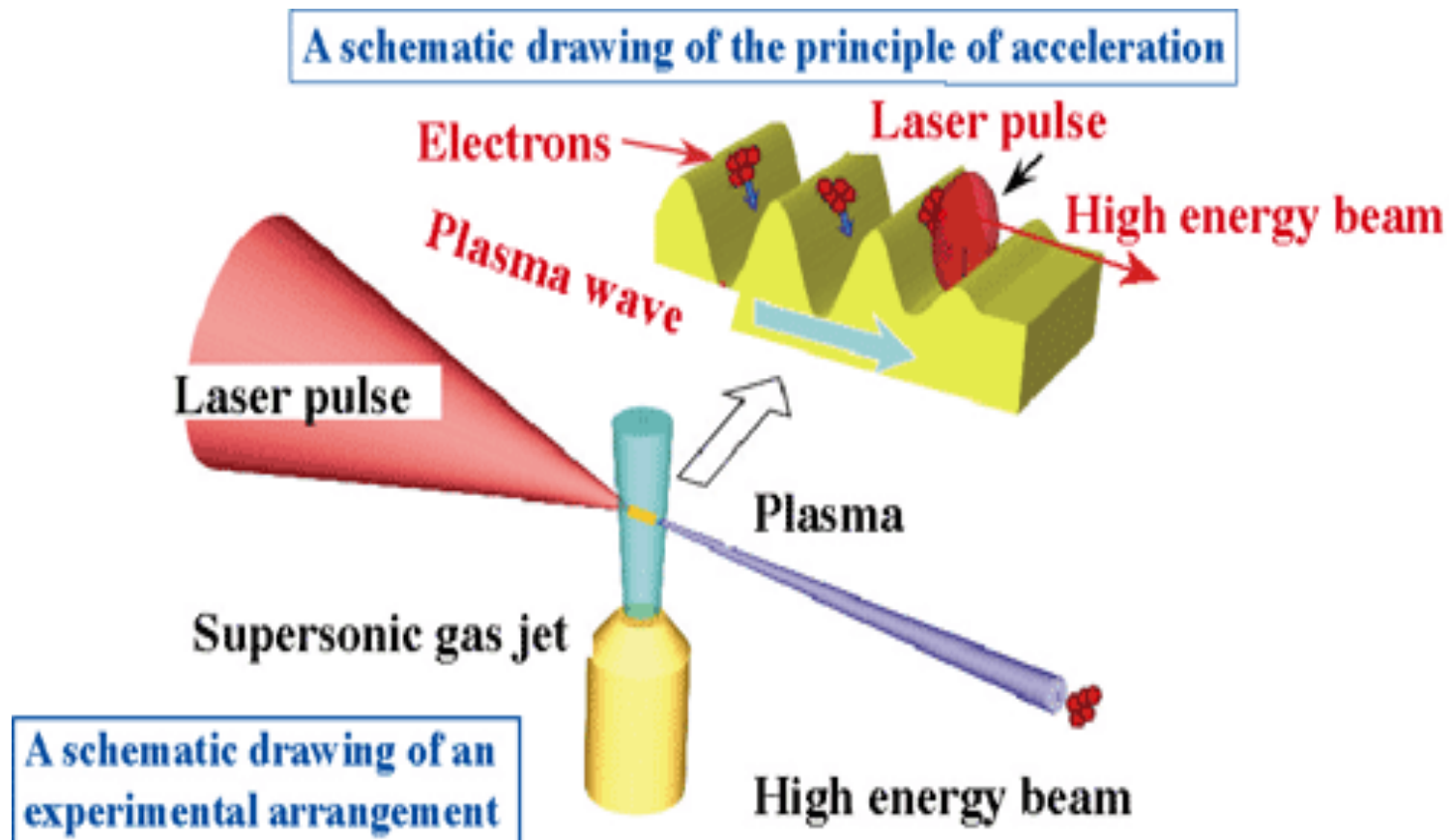
Alghero, Channeling 2012
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LASER-PLASMA ACCELERATION

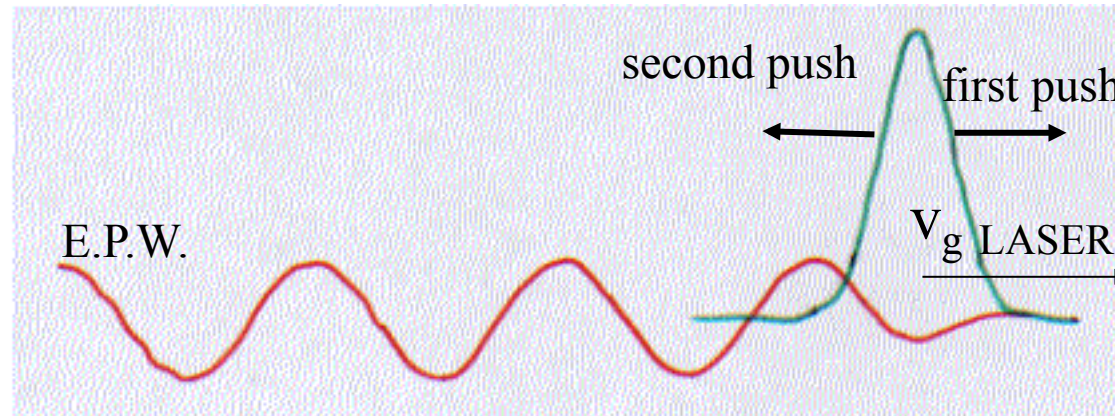
Ultra-short Ti:Sa laser pulses at relativistic intensities (greater than 10^{18} W/cm²), propagating in plasmas at densities of the order of 10^{18-19} el/cm³, can induce accelerating electric fields up to 10^4 times the maximum fields available in the conventional accelerators.

Refer to V. Malka and M. Borghesi contributions in this Conference

The **ponderomotive force** related to the **laser pulse** produces the **Wake Field** and the **Coulomb force** related to the **Electron Plasma Wave** accelerates the electrons



LASER WAKE FIELD



$$\tau \cdot c \approx \frac{\lambda_p}{2} \Leftrightarrow \tau \approx \frac{T_p}{2} \Rightarrow n_e (cm^{-3}) \approx \frac{3 \cdot 10^{-9}}{\tau_{(s)}^2}$$

$$example: \tau = 30 fs \Rightarrow n_e \approx 3.3 \cdot 10^{18} cm^{-3}$$

$$v_{\phi, epw} = v_{g, laser} = c \left(1 - \frac{\omega_{pe}^2}{\omega^2} \right)^{\frac{1}{2}}$$

ELECTRON ACCELERATION IN E.P.W.

1D MODEL

multi-GeV electrons in a few cm
@ $n_e \approx 10^{18} \text{cm}^{-3}$, with 100TW
class Ti:Sa lasers

for $\gamma_p \approx \frac{\omega}{\omega_{pe}} \gg 1 \Rightarrow \Delta W_{\text{max}} = 4\gamma_p^2 \frac{\delta n_e}{n_e} mc^2$ is the max. energy gain

along a distance $L_{\text{deph}} \approx \gamma_p^2 \lambda_p = \lambda_0 \left(\frac{n_c}{n_e} \right)^{\frac{3}{2}}$ $\lambda_p \approx \frac{2\pi c}{\omega_{pe}}$

$$\Delta W_{\text{max}} \approx eE_{\text{max}} \cdot L_{\text{deph}} \propto n_e^{\frac{1}{2}} \cdot \frac{1}{n_e} \cdot n_e^{-\frac{1}{2}} = \frac{1}{n_e}$$

LPA Secondary Sources

The drastic **reduction of the dimensions and costs** of the Laser Plasma Acceleration apparatus, open to **several applications**. In fact, once energetic electron bunches at rep rate of a few Hz are produced, X- γ radiation and particle secondary sources can be carried out.

Beside their compactness these sources are easily synchronized with other laser systems, fulfilling the best conditions for **femtosecond time resolved pump and probe experiments**.

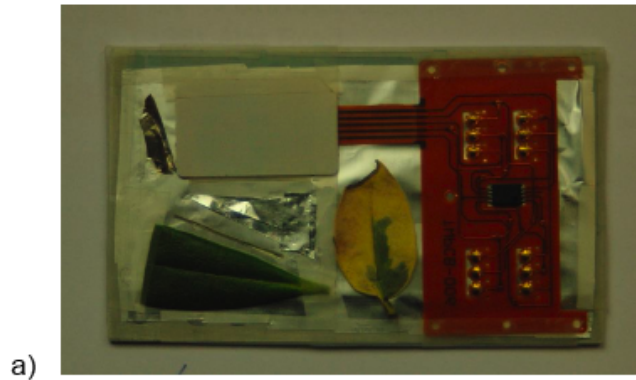
Radiation Secondary Sources

- Bremsstrahlung
- K- α radiation
- Betatron radiation
- Thomson Scattering (V. Malka contribution in this Conference)
- FEL
- Terahertz

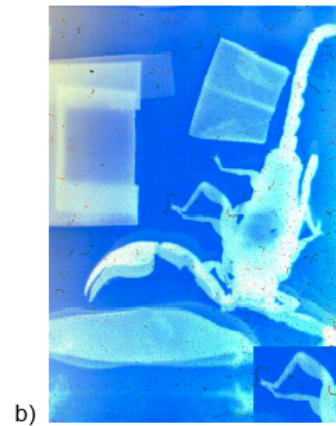
Particle Secondary Sources

- Ions (M. Borghesi talk; G.P. Cirrone et al poster PS2-02)
- $\gamma \rightarrow e^+ + e^-$
- Neutrons, μ ,

Electron Radiography



G.C. Bussolino et al,
submitted to Medical Physics



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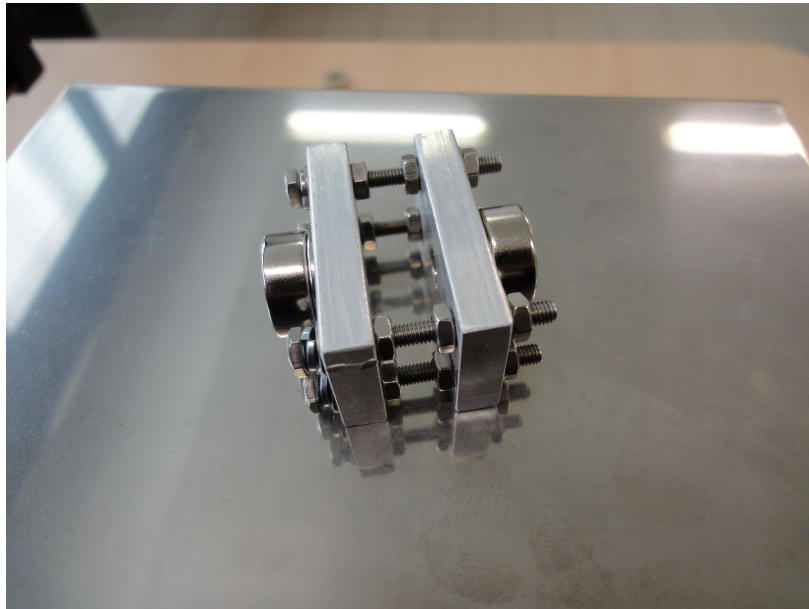
LPA: the Control of Electron Bunches

- Plasma shaping
- Laser guiding
- Ionization process
- Multi-stage LPA
- **Static Magnetic Field**
(T.Hosokai et al, PRL, 97, 075004, 2006)
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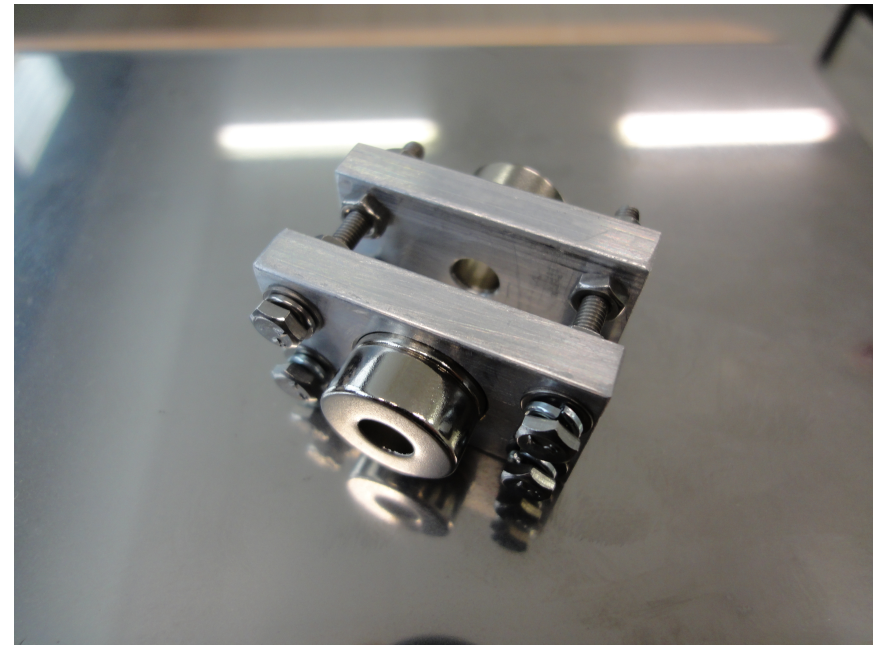


The magnetic device

$B \approx 0.2$ Tesla

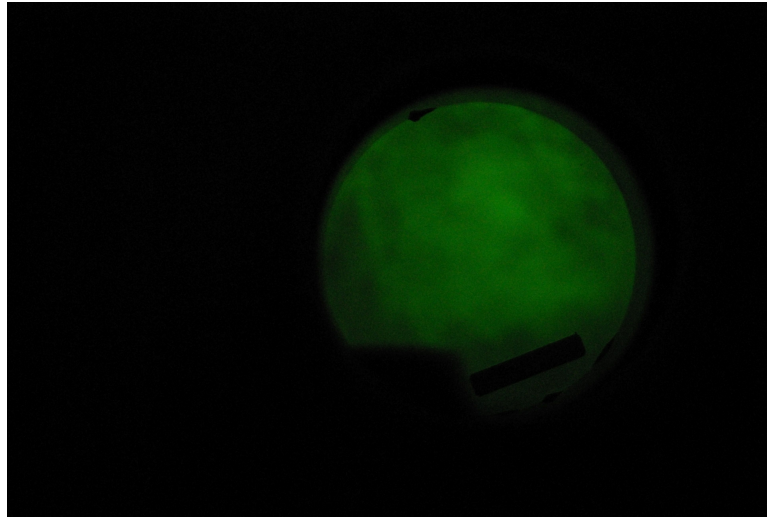


Equipment projected by Y. Oishi and
D. Giulietti at INFN-Pisa



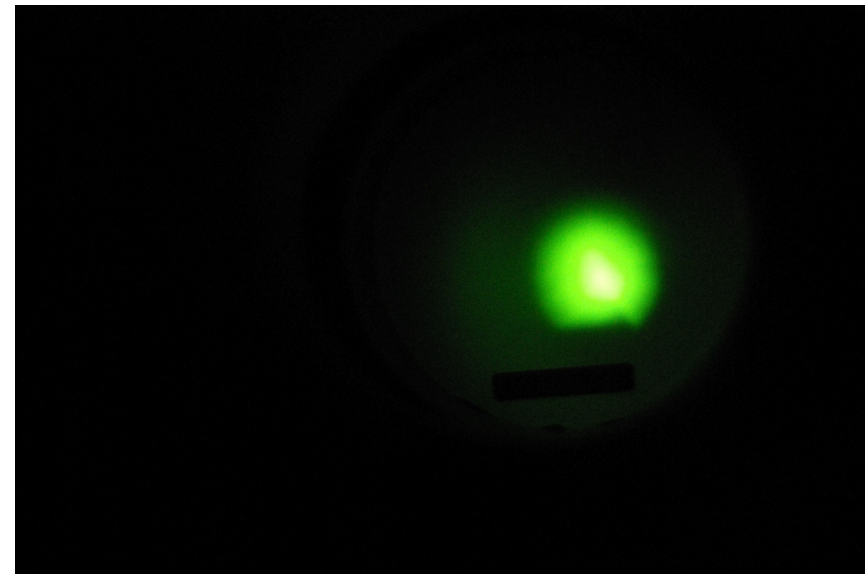
Static Magnetic Field in LPA Experiments

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$N_2=38$ bar

$B=0.2$ T in the interaction region, along the laser axis



Influence of static magnetic field on LPA mechanism ?

Comparing the **magnetic pressure** with the **plasma pressure**

$$\frac{B^2}{2\mu_0} = nKT \rightarrow B = \sqrt{2\mu_0 K T n} = 20 \div 50 \text{ Tesla}$$

We see that the magnetic field required to have confinement of typical **laser produced plasmas** ($n_e \approx 10^{19} \text{cm}^{-3}$, $T \approx 100 \text{eV}$) is of a few 10's Tesla: much more high the one we use in the magnetic device (0.1-0.2 Tesla).

Magnetized Plasmas ?

The condition for magnetized plasma holds in the regions where the temperature is high and the density low:

$$\omega_L \tau_{ei} \geq 1 \text{ or } \frac{l_{mfp}}{r_L} \geq 1$$
$$\omega_L = \frac{eB}{m}, \quad \tau_{ei} = \frac{3.4 \times 10^5 T_e^{\frac{3}{2}}}{Z^2 n_i \Lambda}$$

T_e (eV), n_i (cm^{-3}), Z = ioniz. number, Λ = Coulomb log

l_{mfp} = mean free path; r_L = Larmor radius

For $B \approx 0.2$ Tesla such conditions could be satisfied in low density plasma channels as those we evidenced in recent experiments performed in Pisa and CEA-Saclay, or in the pre-plasma produced by the pre-pulse accompanying the ultra-intense laser pulse as suggested by T. Hosokai in PRL, **97**, 075004, 2006.

LPA electrons affected by electromagnetic lens effects ?

When an electron passes through an electromagnetic lens it is subjected to two forces at any particular moment: a force (F_z) parallel to the core (Z axis) of the lens, and a force (F_r) parallel to the radius of the lens. These two forces are responsible for two different actions on the electrons, **spiraling** and **focusing**, as they pass through the lens. An electron passing through the lens parallel to the Z axis will experience the force F_z causing it to spiral through the lens. This spiraling causes the electron to experience the force F_r which causes the beam to be compressed toward the Z axis. The **magnetic field is inhomogeneous** in such a way that it is weak in the center of the gap and becomes stronger close to the bore. **Electrons close to the center are less strongly deflected than those passing the lens far from the axis.**

CONCLUSIONS

- New Acceleration Techniques
- High brightness sources of electrons, protons, ions, neutrons, positrons, X & γ -rays, ...
- Applications in HEP, medicine, ICF, material science, astrophysics, femtochemistry, attosecond science, ...
- **The control of the LPA process and the manipulation of produced electron bunches required**

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