

Plasma Discharge Undulator: a novel concept for plasma-based radiation sources

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INTRODUCTION

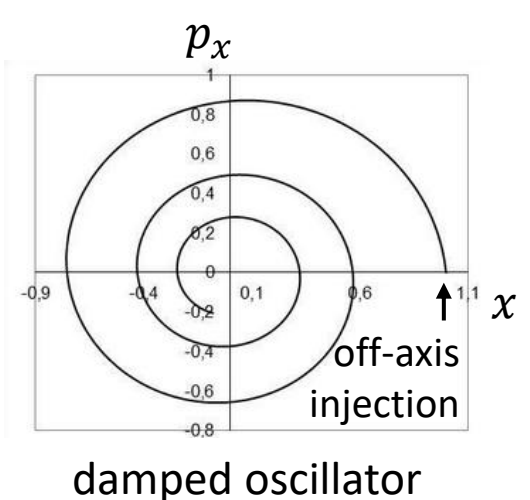
Plasma discharge devices have recently demonstrated their potential for compact particle beam manipulation. Building upon the Active Plasma Lens [1] and its extension to curved geometries (Active Plasma Bending [2]), new studies have revealed that chicane-like configurations can support sub-betatron oscillations of the beam. Motivated by this observation, the novel concept of the Plasma Discharge Undulator (PDU) is introduced. The PDU consists of a sequence of transversely displaced jointed plasma capillaries, carrying a high-current discharge. The resulting azimuthal magnetic field focuses the beam, while the periodic transverse displacement of the equilibrium axis acts as a geometric forcing term. This yields a well-defined oscillation at a wavelength λ_{PDU} distinct from the natural betatron wavelength λ_β . By proper injection, collective betatron oscillations can be suppressed, leaving only the forced λ_{PDU} beam centroid oscillation. This approach eliminates the intrinsic K-spread limitation of plasma undulators, achieving a narrow undulator strength distribution while preserving strong plasma focusing for beam matching. The theoretical framework, single-particle and beam dynamics analyses, first estimates of radiation emission from the PDU and preliminary particle-in-cell studies of segmented capillaries are presented. The PDU thus offers a pathway toward miniaturized, tunable, full-plasma-based radiation beamlines, with enhanced control over beam quality and radiation properties.

MOTIVATION

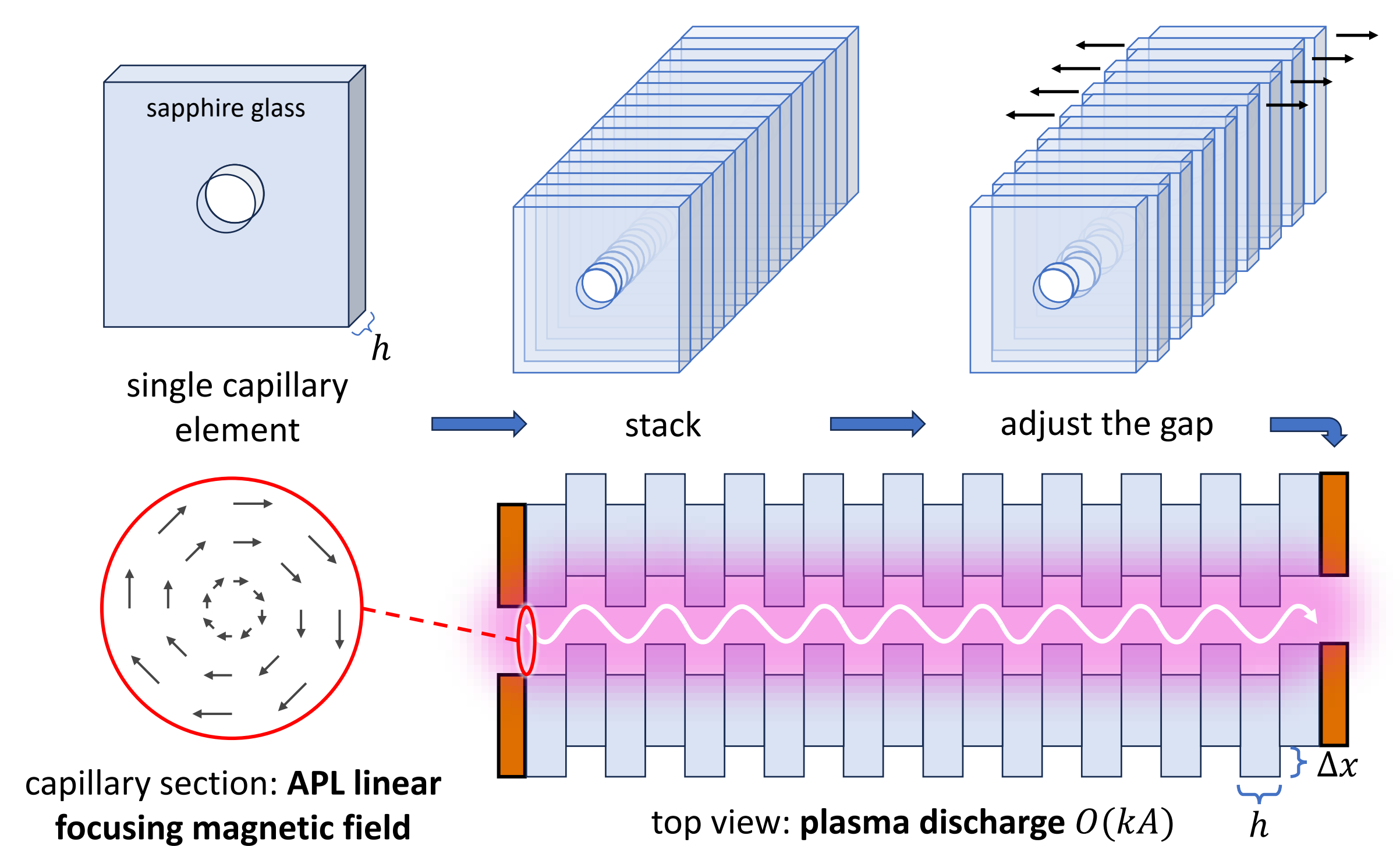
Betatron plasma undulators limitations:

- On axis injection: **K-spread constraint**
- Off axis injection: less K-spread but possibly **unstable + oscillation damping and/or emittance growth**
- Possible solution? **Add driving term**
How? **Geometrically**

$$\frac{\sigma_K}{\langle K \rangle} = \sqrt{\frac{4}{\pi} - 1} \approx 0.523 \quad (\text{Gaussian beam})$$



CONCEPT



ANALYTICAL MODEL : DRIVEN OSCILLATOR

$$x_0 = \frac{\Delta x}{2} \frac{k_\beta^2}{k_\beta^2 - k_{PDU}^2}$$

sub-oscillation amplitude

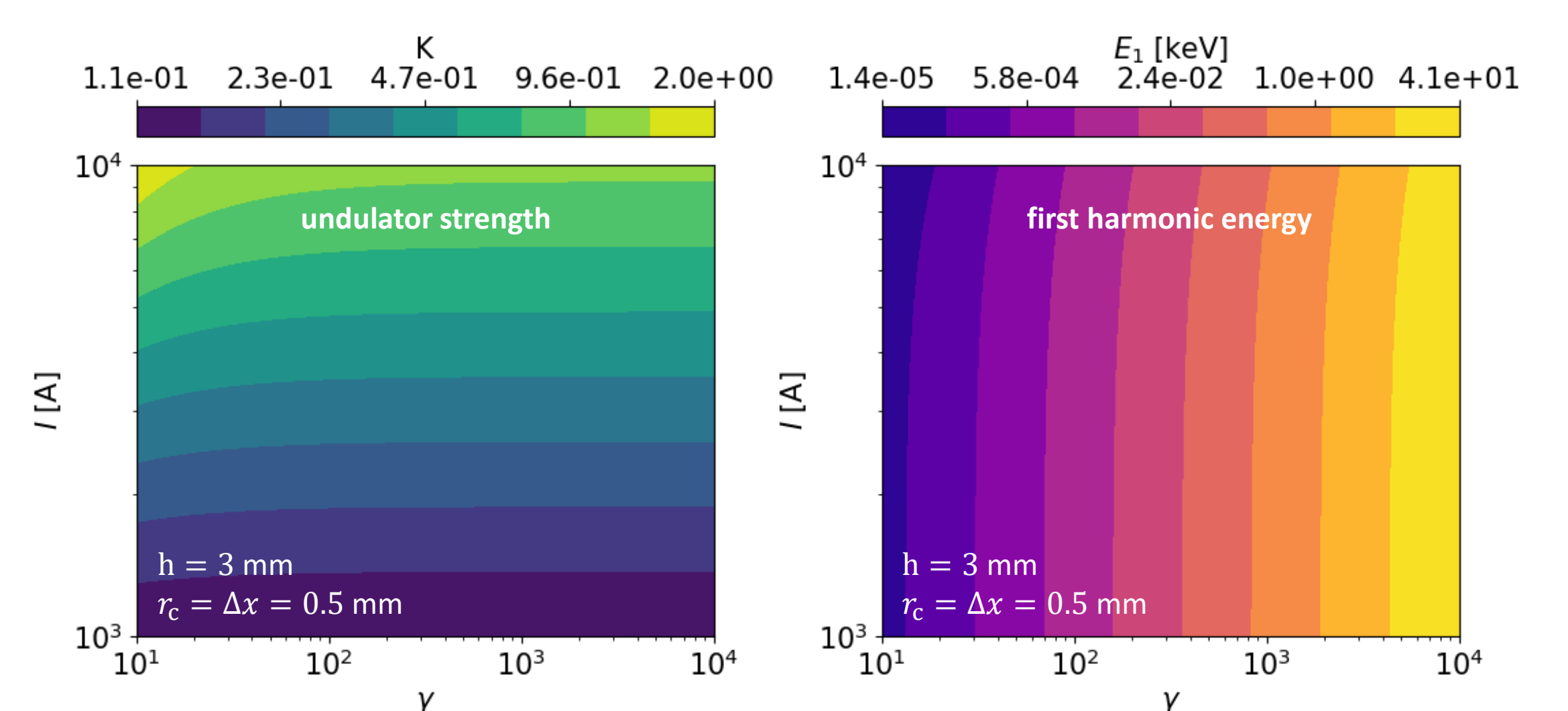
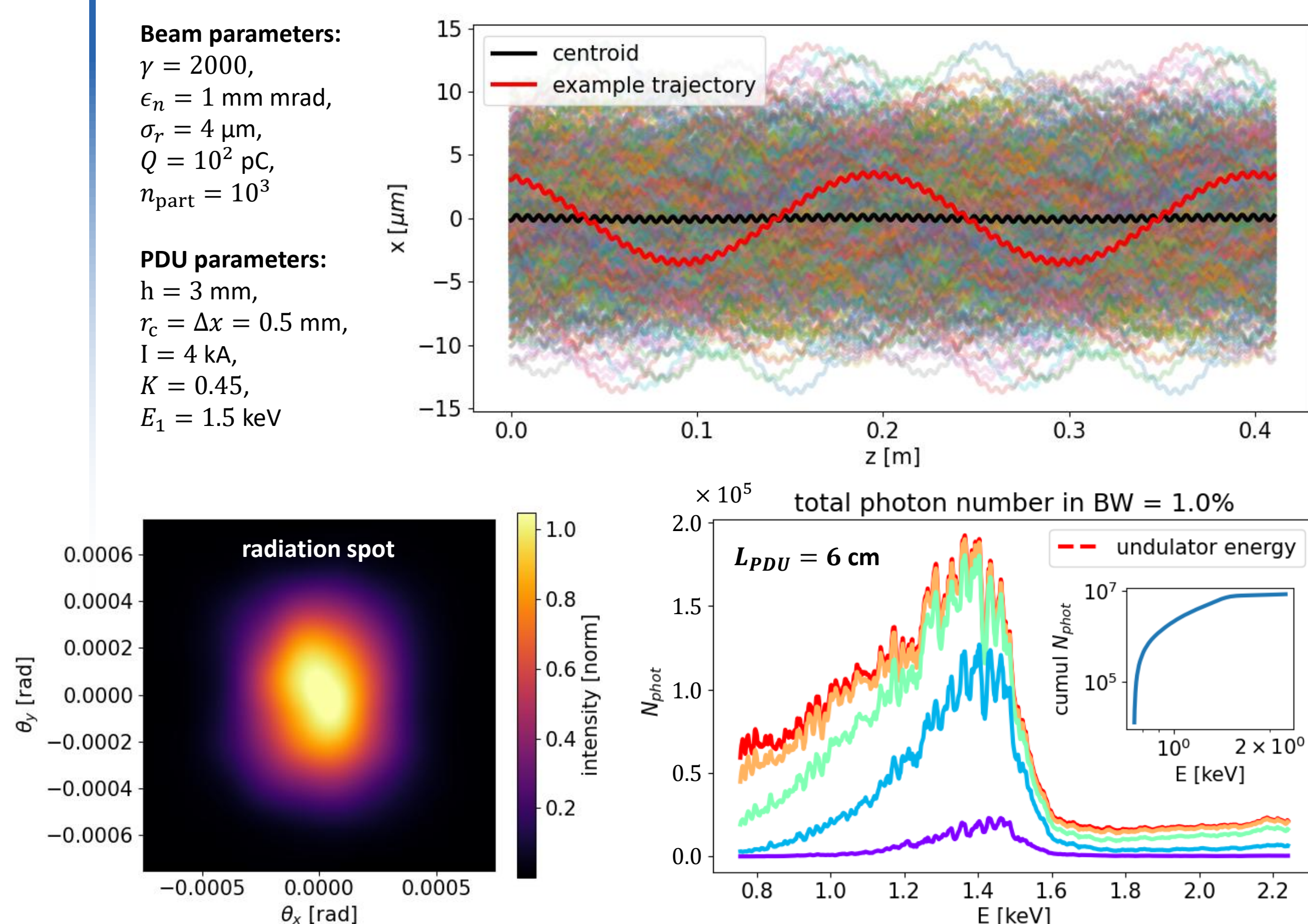
$$K_{PDU} = \frac{\gamma \Delta x}{2 \frac{h}{\pi} - \frac{4m_e c \gamma \pi}{e \mu_0 J \frac{h}{h}}}$$

undulator strength

$$\sigma_M = \left(\frac{2m_e c \gamma \epsilon_{rms}^2}{e \mu_0 J} \right)^{\frac{1}{4}}$$

APL matching

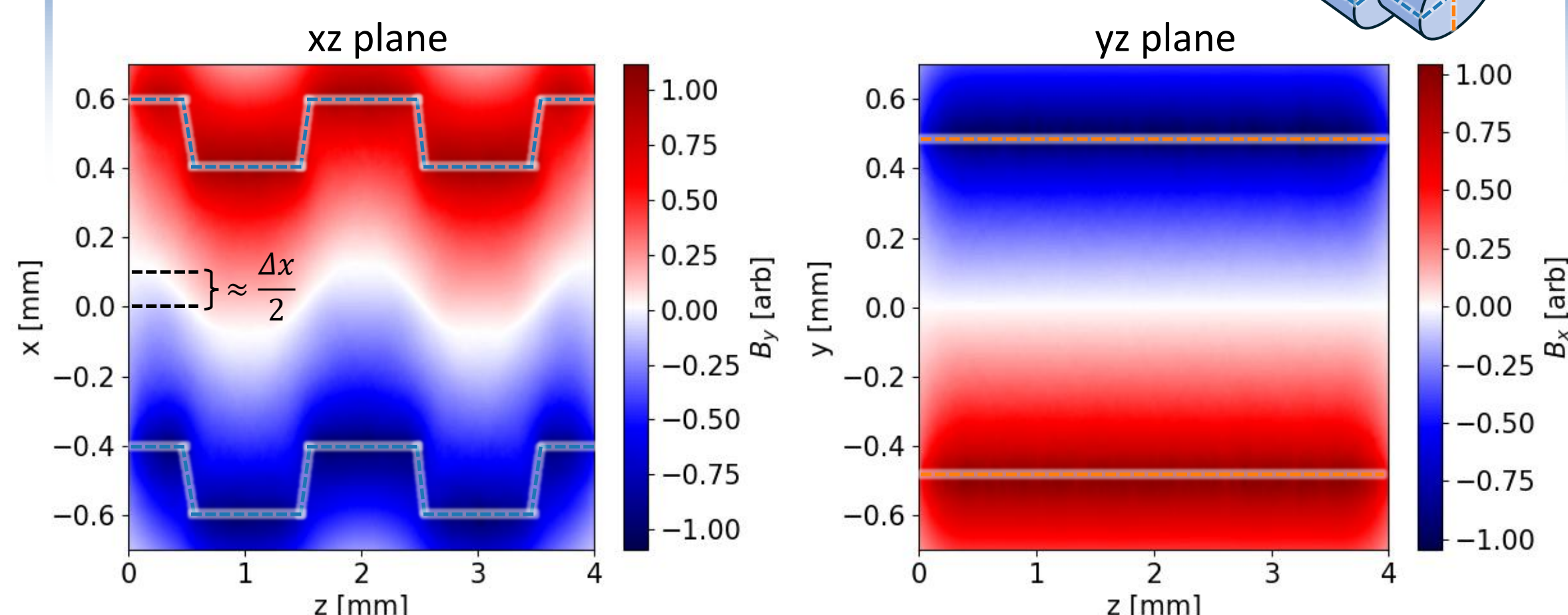
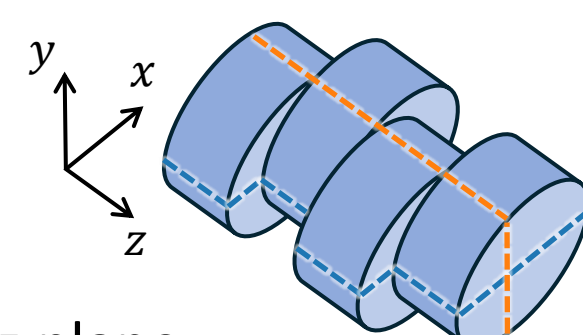
DYNAMICS AND RADIATION SIMULATIONS



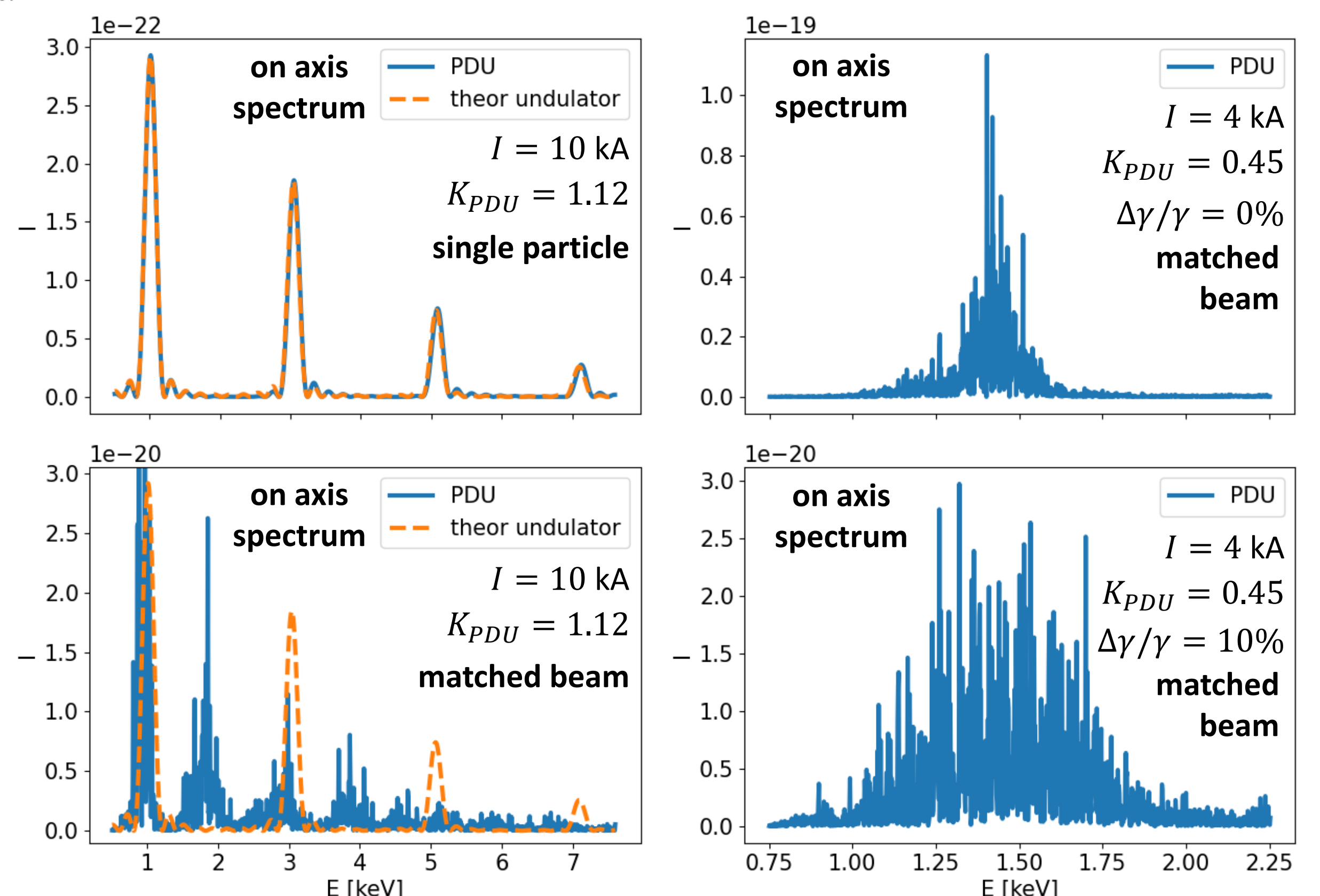
- Focusing magnetic field with oscillating zero: alternating dipolar error acts as effective undulator
- Magnetic field driven bending: the beam should feature longitudinal microbunching (transverse in ICL)
- Beam centroid injection at $x = x_0$ avoids λ_β collective oscillations, leaving just λ_{PDU} collective oscillations
- The wavelength λ_{PDU} appears as a modulation of amplitude x_0 over the main λ_β trajectory oscillation
- Intrinsically 2-color: λ_{PDU} acts as a conventional magnetic undulator while λ_β as a conventional plasma undulator
- Due to short λ_β , beam trajectories feature high average slope, so even harmonics contribution appear on axis too
- K gets greater for larger λ_{PDU} up to the threshold $\lambda_{PDU} = \lambda_\beta$, after which the behavior is opposite

PLASMA SIMULATIONS

- Just a preliminary pre-ionized, cold gas, short FBPIC simulation to evaluate current density evolution and magnetic field zero
- Thermal effects could help achieving larger field oscillation amplitude \Rightarrow larger K



$h = 1 \text{ mm}$, $r_c = 0.5 \text{ mm}$, $\Delta x = 0.2 \text{ mm}$, $n_p = 10^{15} \text{ cm}^{-3}$, $E_z = 2.5 \times 10^5 \text{ V/m}$, $n_{modes} = 6$



[1] "Direct measurement of focusing fields in active plasma lenses", DOI: <https://doi.org/10.1103/PhysRevAccelBeams.21.122801>
[2] "Guiding of Charged Particle Beams in Curved Plasma-Discharge Capillaries," DOI: <https://doi.org/10.1103/PhysRevLett.132.215001>