Simulation Studies of Temperature Effects for High-Repetition-Rate Plasma Wakefield Accelerators



I. Najmudin, J. Beinortaite, J. Cowley, M. Thévenet, A. Thomas, J. Wood, R. D'Arcy





High-Repetition-Rate Plasma Acceleration

> Goal: Implement plasma accelerators into operational facilities - FEL's and Linear Colliders

Collider Luminosity:
$$L \propto \frac{Q_{bunch}f_{rep}}{\sigma_{x}\sigma_{y}}$$
 FEL Instantaneous Brightness: $dB/dt \propto \frac{Q_{bunch}f_{rep}}{\sigma_{x}\sigma_{y}\sigma_{E}}$

> Requirement: Demonstrate stable operation of plasma accelerator stages at competitively high repetition rates.

High-Repetition-Rate Plasma Acceleration

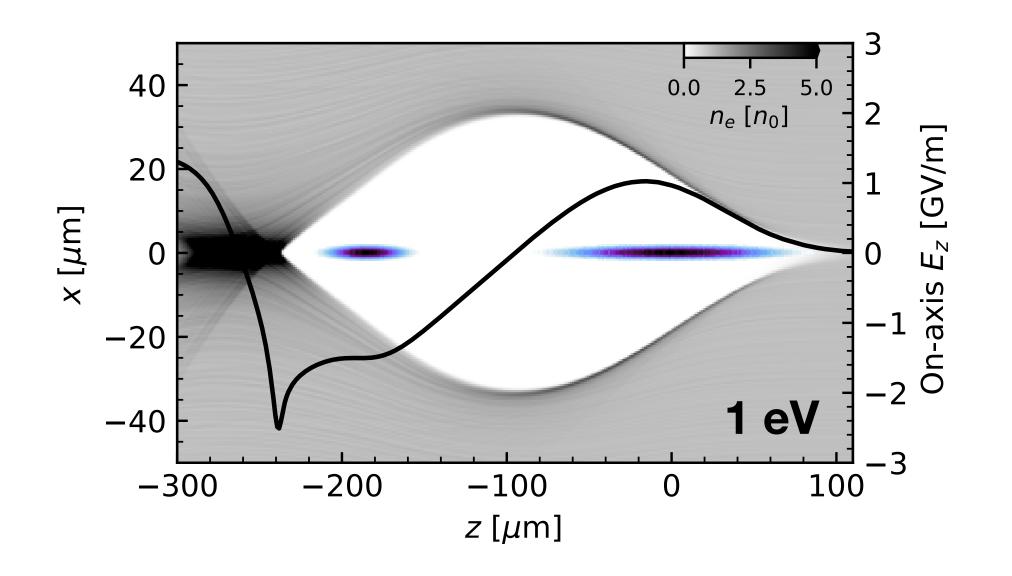
> Goal: Implement plasma accelerators into operational facilities - FEL's and Linear Colliders

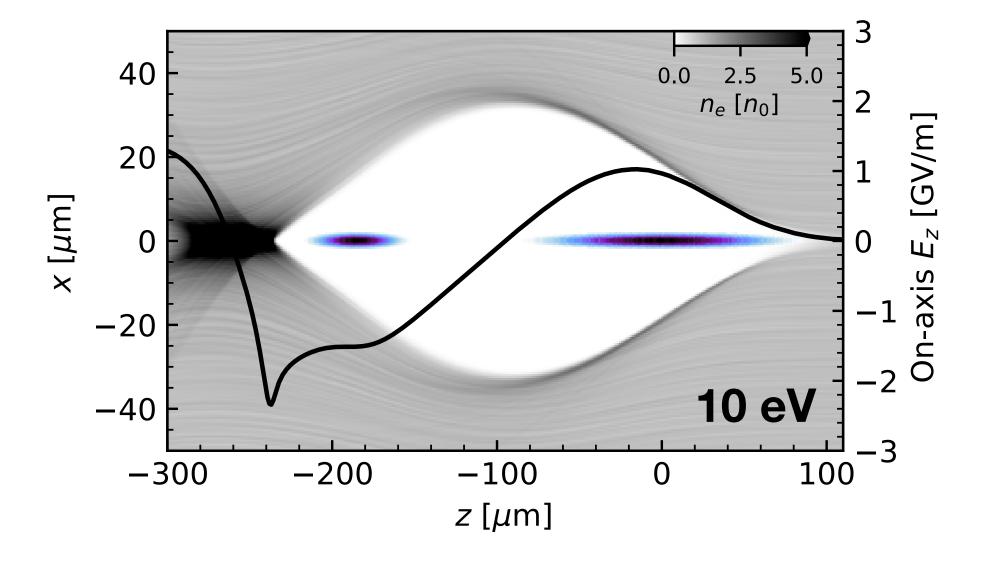
Collider Luminosity:
$$L \propto \frac{Q_{bunch}f_{rep}}{\sigma_x\sigma_y}$$
 FEL Instantaneous Brightness: $dB/dt \propto \frac{Q_{bunch}f_{rep}}{\sigma_x\sigma_y\sigma_E}$

- > Requirement: Demonstrate stable operation of plasma accelerator stages at competitively high repetition rates.
- > **Challenges:** For high repetition rates, we need to reuse the plasma, so to maintain consistent quality of acceleration for many bunches, we need,
 - > Consistent plasma densities / plasma profiles
 - > Durable plasma sources / containers
 - > To keep the plasma temperature low (or consistent?)
- > Why care about plasma temperature? Plasma acceleration is not totally efficient, so any energy from the driver that is not extracted by the witness, makes it's way into and through the plasma. An increased plasma temperature then effectively changes the plasma frequency e.g via the Bohm-Gross dispersion relation:

$$\omega^2 = \omega_p^2 + \frac{3k_B T}{m_e} k^2$$

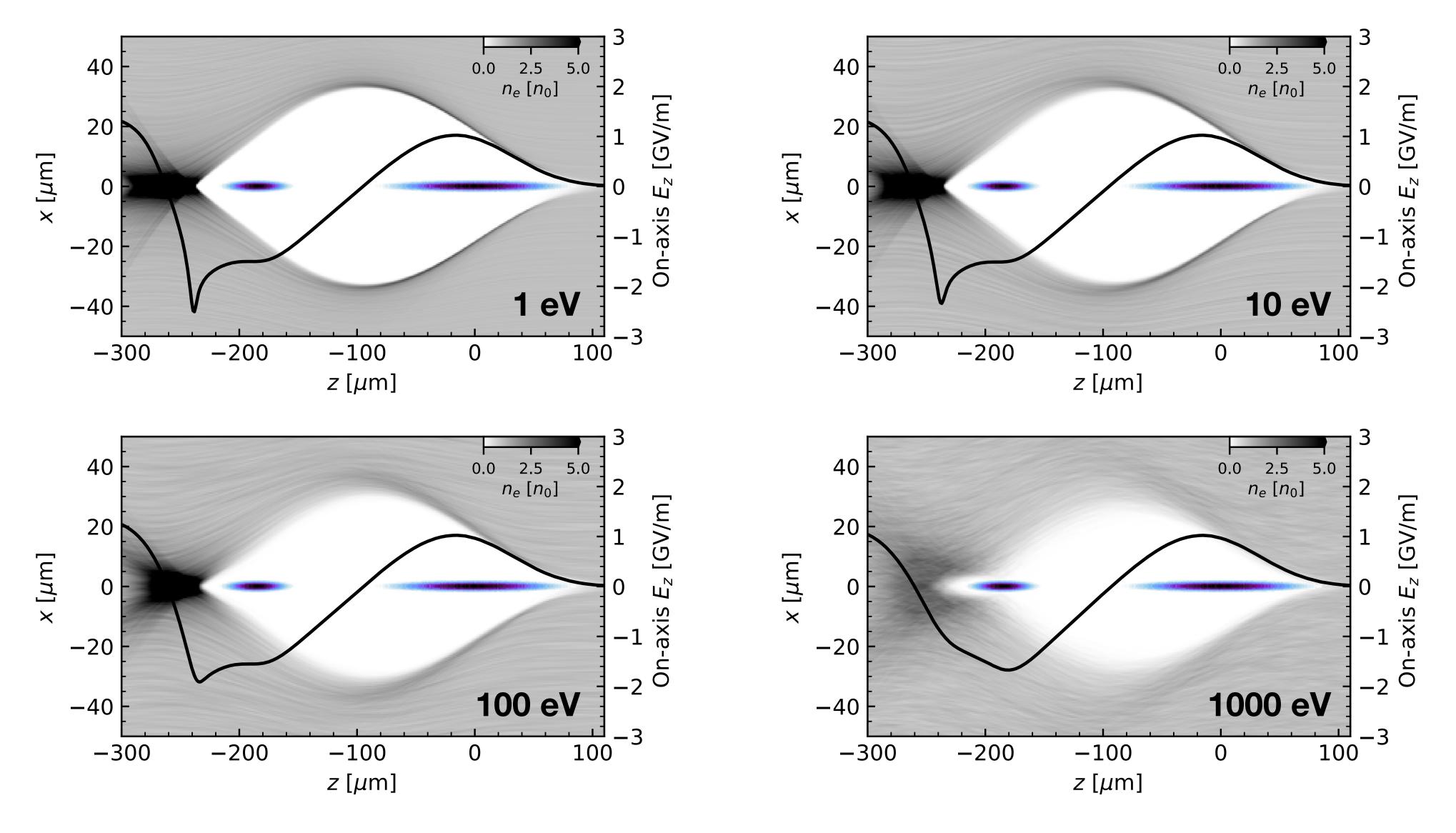
Effect of plasma temperature on non-linear wakefields







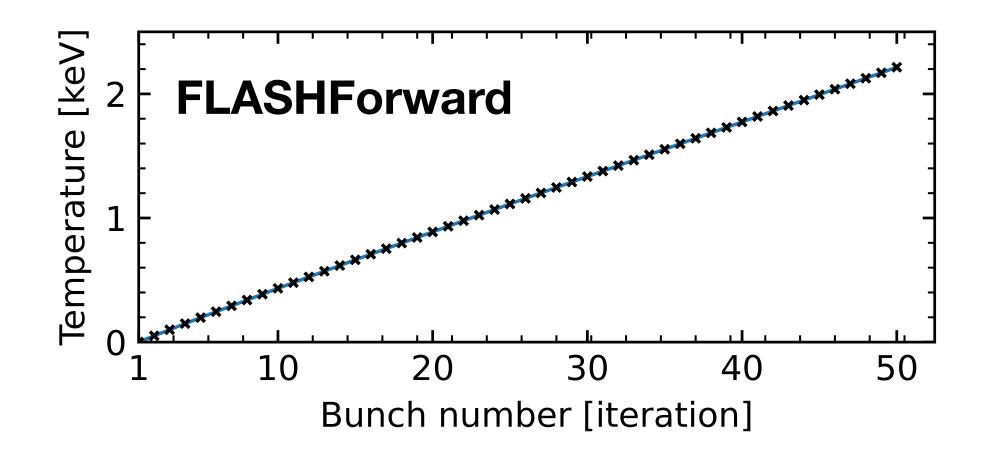
Effect of plasma temperature on non-linear wakefields

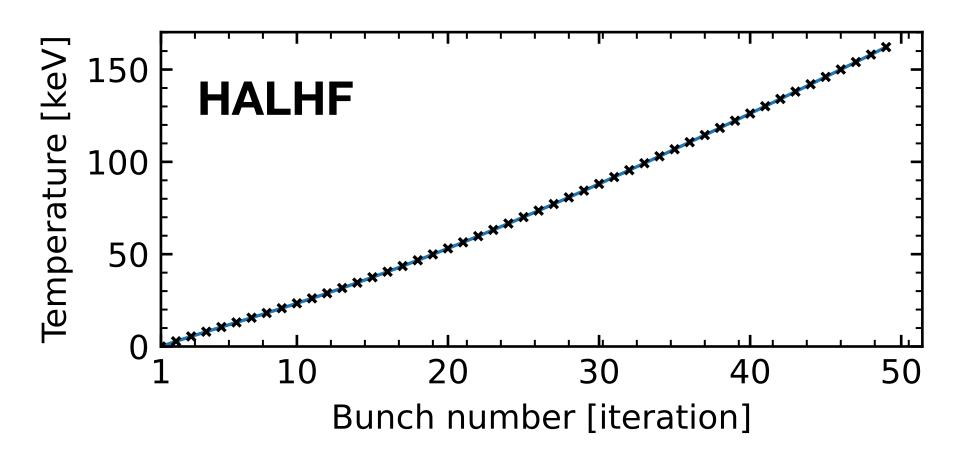




Naive assumptions to form an iterative PIC loop

- > What kind of temperatures are expected?
- > Worst case scenario: Plasma takes all deposited energy from the bunches and splits it evenly among the electrons and ions, leading to a corresponding rise in plasma temperature

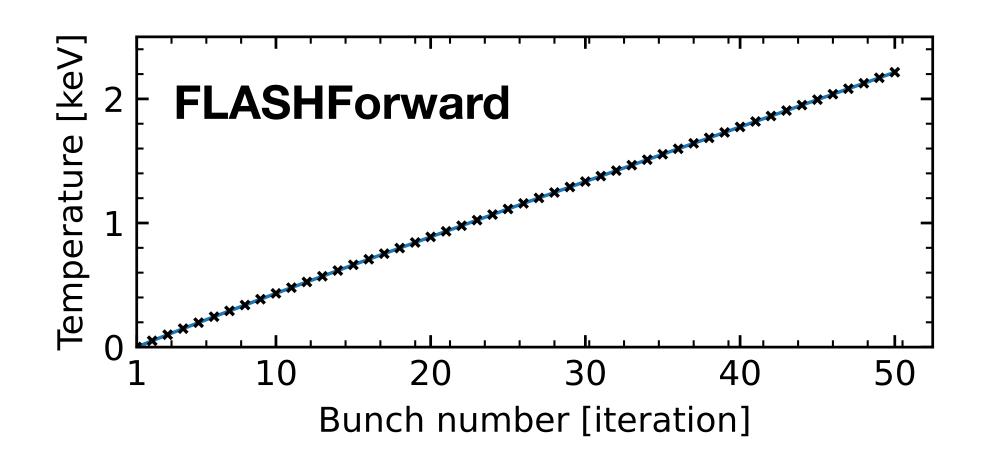


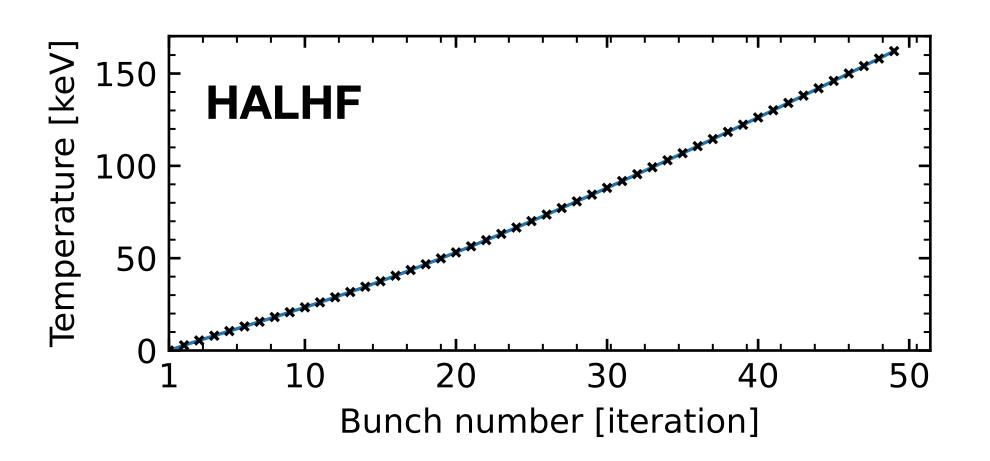




Naive assumptions to form an iterative PIC loop

- > What kind of temperatures are expected?
- > Worst case scenario: Plasma takes all deposited energy from the bunches and splits it evenly among the electrons and ions, leading to a corresponding rise in plasma temperature





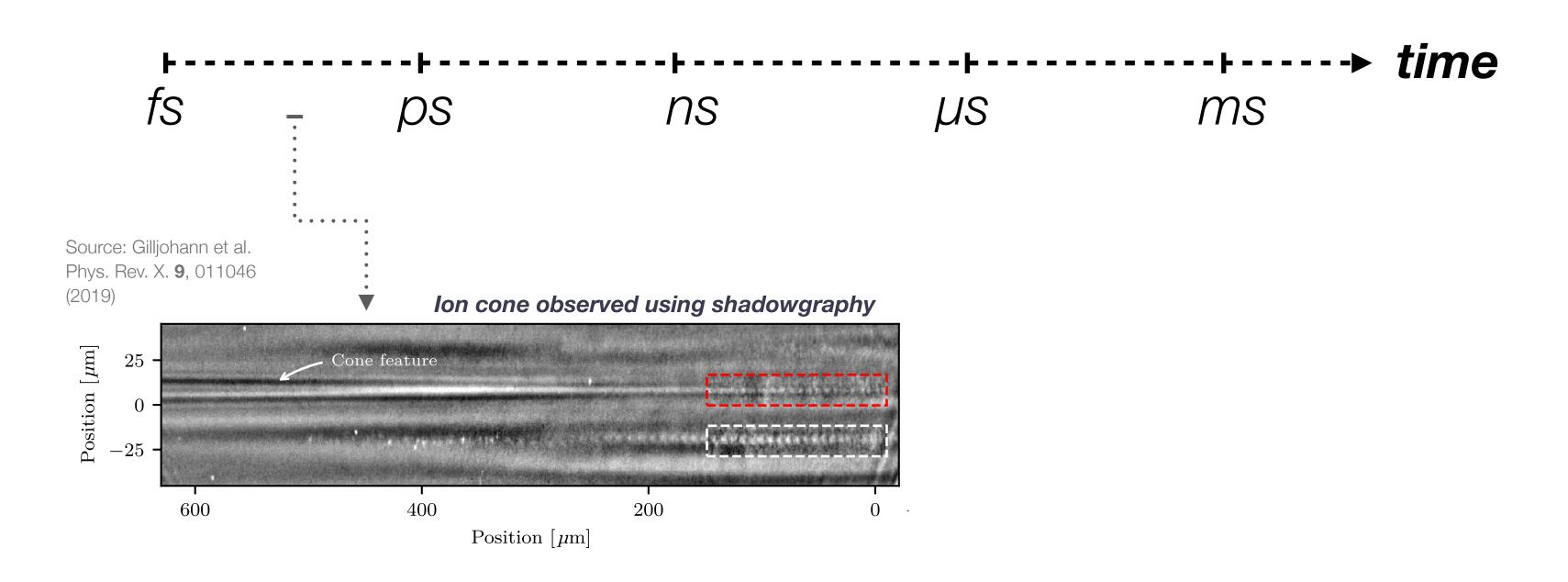
- > Runaway effect: Becomes more inefficient with increasing plasma temperature
- > **Obvious shortcomings:** No sinks in energy + no plasma evolution
 - > Hot and fast electrons carry a large portion of the initial wakefield energy
 - > Heated ions cause excitation and ionisation of background neutral atoms
 - > Plasma density and temperature distribution for each next bunch may not be uniform



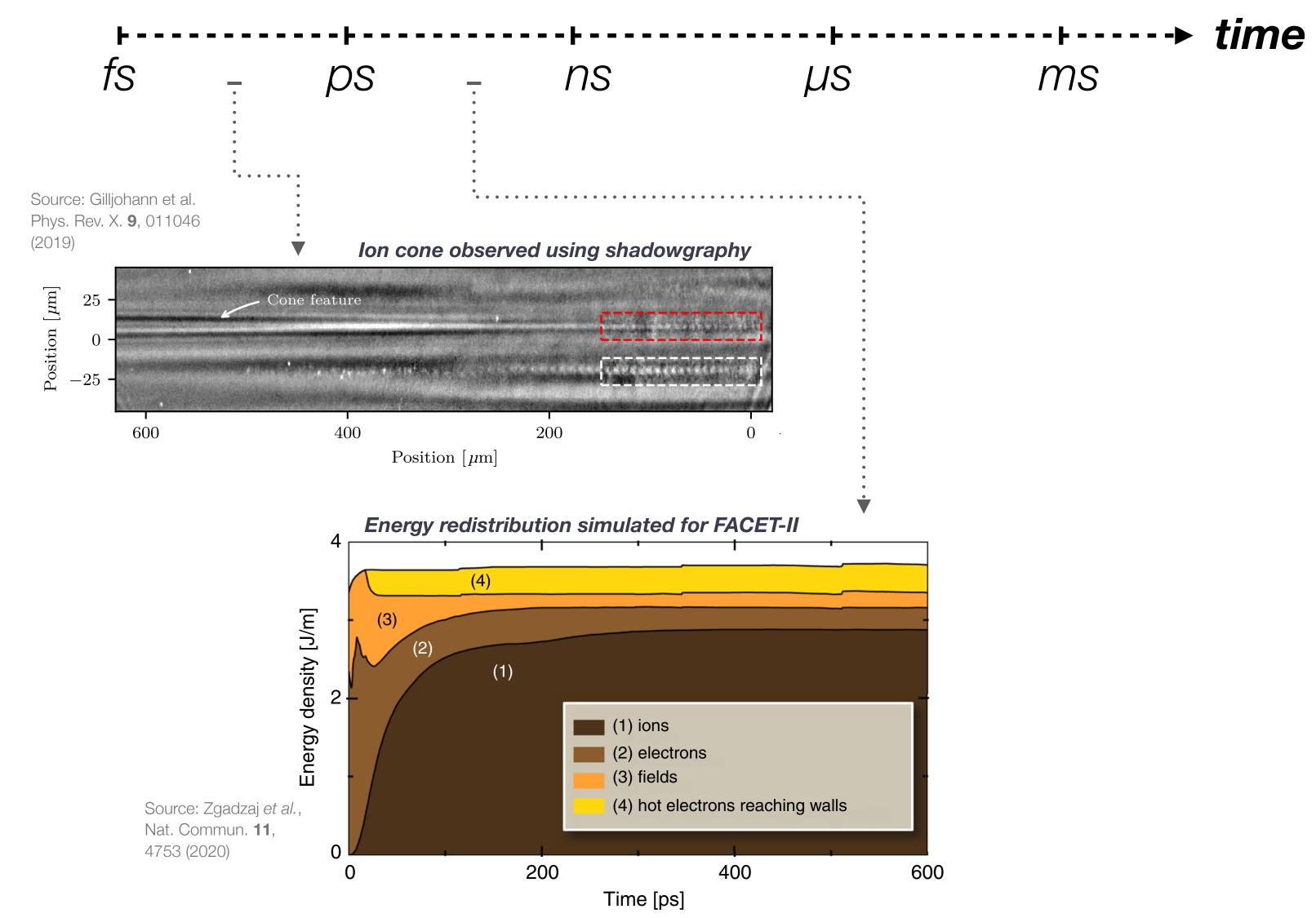
What are the effects of temperature on a plasma accelerator?

What are the relevant channels of energy transport that dictate plasma evolution between subsequent acceleration events?



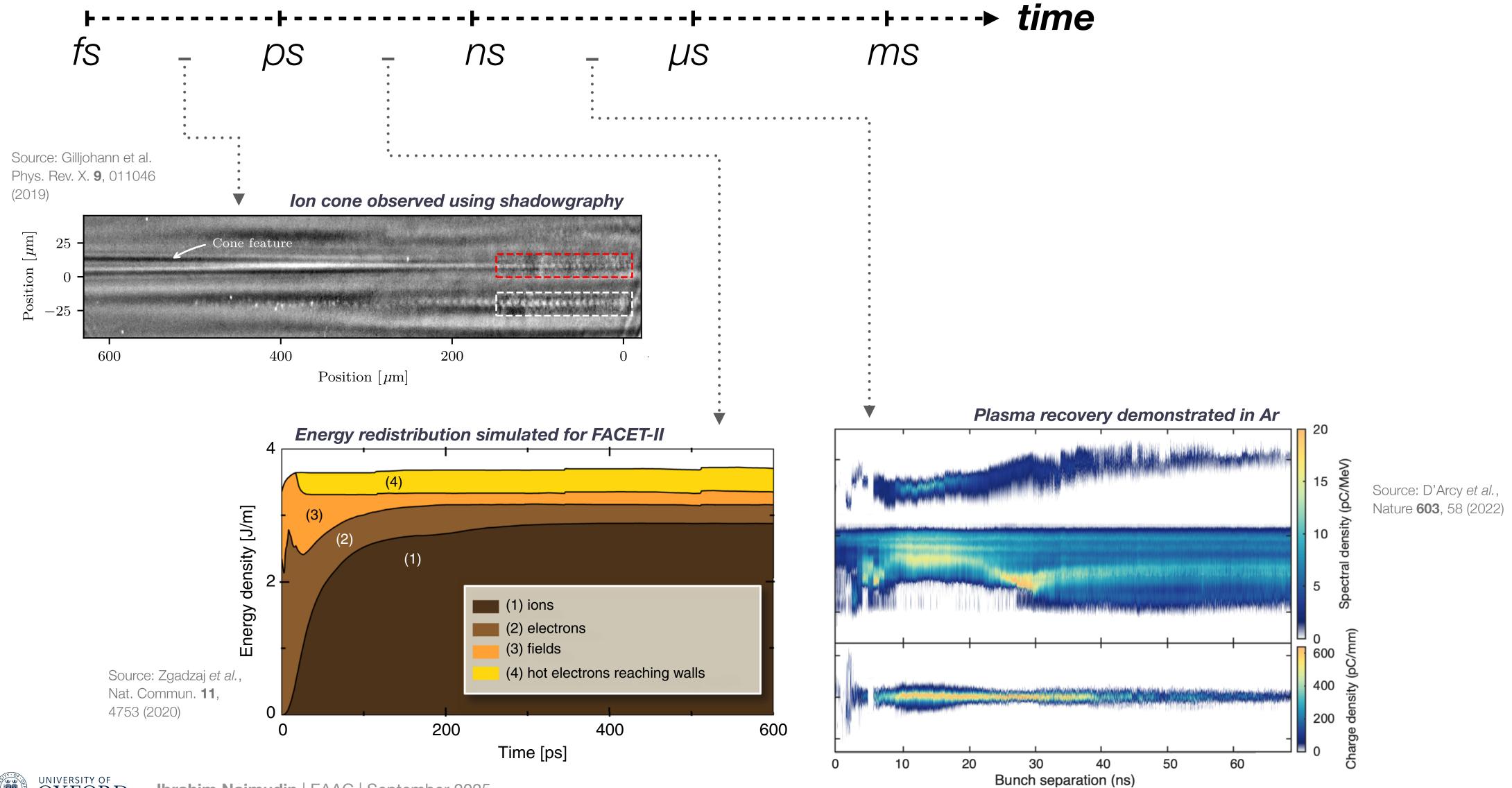




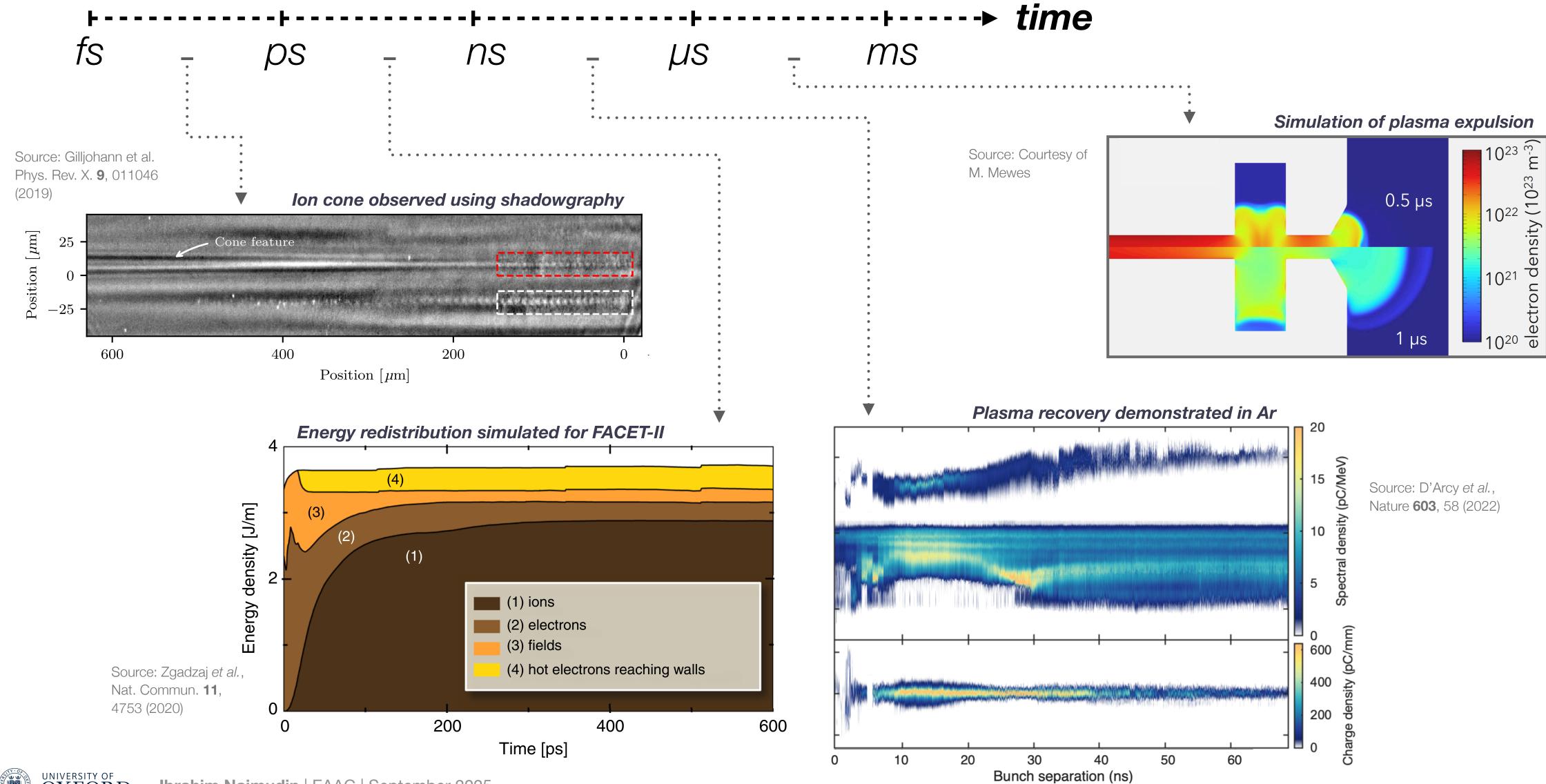




Ibrahim Najmudin | EAAC | September 2025



Ibrahim Najmudin | EAAC | September 2025



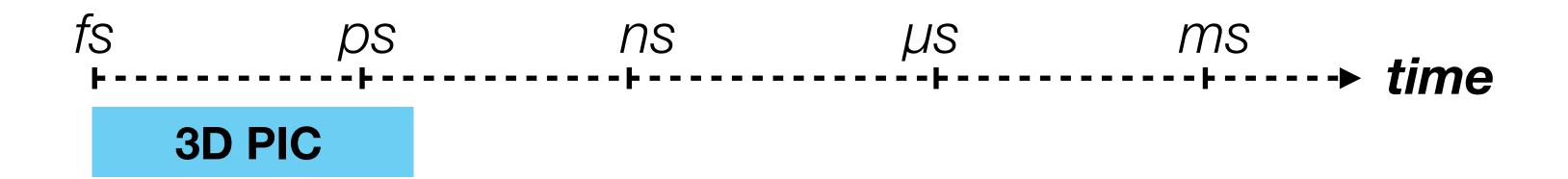
- > Goal: Simulate long timescales of plasma evolution after the intense & nonlinear perturbation from a driver
 - > Typically, PIC codes are used to guide experiments
 - > This is currently not the case for long timescale plasma studies
 - > Need a self-consistent framework to inform future experimental programs and facility designs (ad-hoc)



- > Goal: Simulate long timescales of plasma evolution after the intense & nonlinear perturbation from a driver
 - > Typically, PIC codes are used to guide experiments
 - > This is currently not the case for long timescale plasma studies
 - > Need a self-consistent framework to inform future experimental programs and facility designs (ad-hoc)

> How:

> 3D PIC codes are most suitable and most commonly used for short timescales involved in plasma wakefield excitation (fs-ps)



- > Goal: Simulate long timescales of plasma evolution after the intense & nonlinear perturbation from a driver
 - > Typically, PIC codes are used to guide experiments
 - > This is currently not the case for long timescale plasma studies
 - > Need a self-consistent framework to inform future experimental programs and facility designs (ad-hoc)

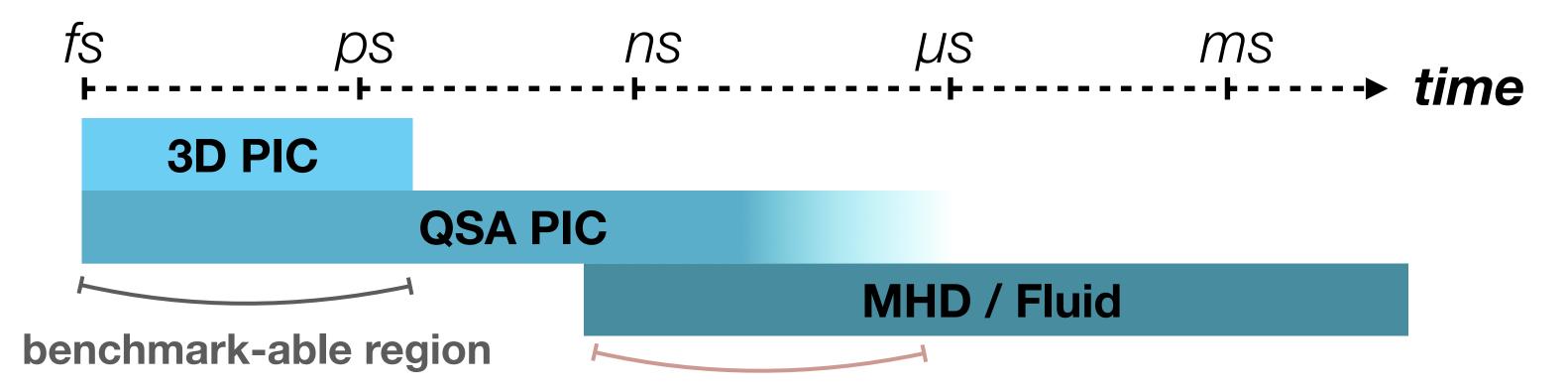
- > 3D PIC codes are most suitable and most commonly used for short timescales involved in plasma wakefield excitation (fs-ps)
- > MHD/Fluid codes are common to use for bulk plasma evolution effects on longer timescales ($ns-\mu s$)



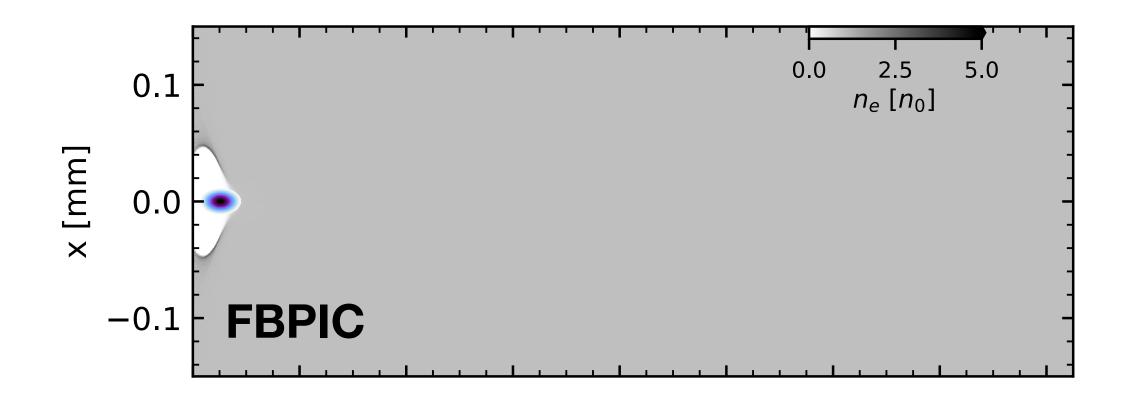


- > Goal: Simulate long timescales of plasma evolution after the intense & nonlinear perturbation from a driver
 - > Typically, PIC codes are used to guide experiments
 - > This is currently not the case for long timescale plasma studies
 - > Need a self-consistent framework to inform future experimental programs and facility designs (ad-hoc)

- > 3D PIC codes are most suitable and most commonly used for short timescales involved in plasma wakefield excitation (fs-ps)
- > MHD/Fluid codes are common to use for bulk plasma evolution effects on longer timescales ($ns-\mu s$)
- > Middle Ground: The Quasi-Static Approximation (QSA)
 - > Used to separate the timescales involved in plasma evolution and driver evolution
 - > A single timestep in driver evolution allows us to focus on the plasma evolution!

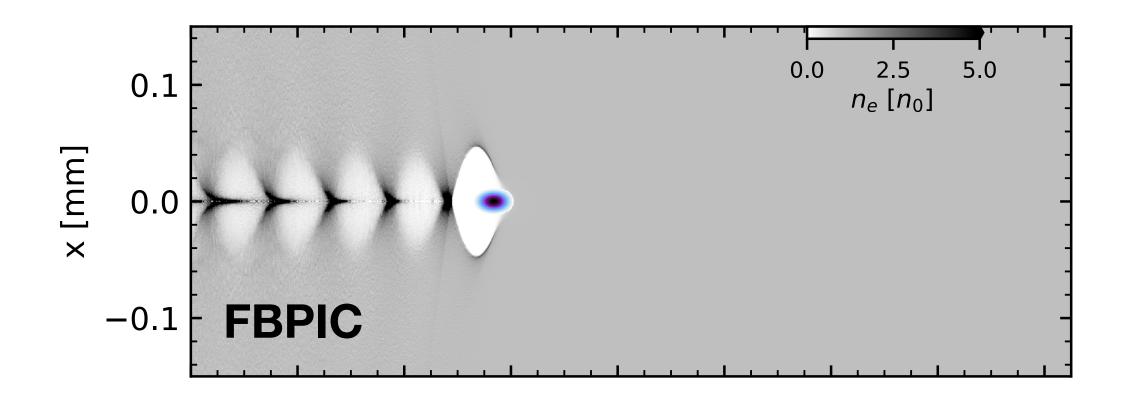






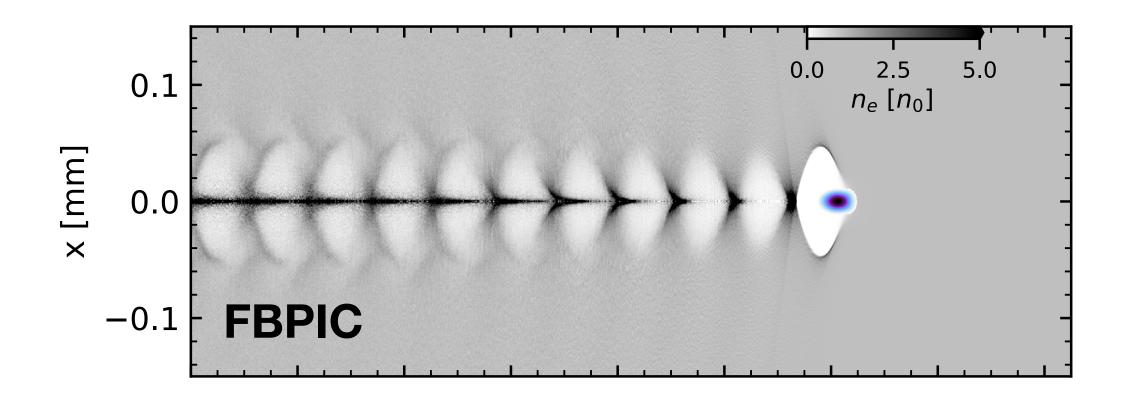
> How:





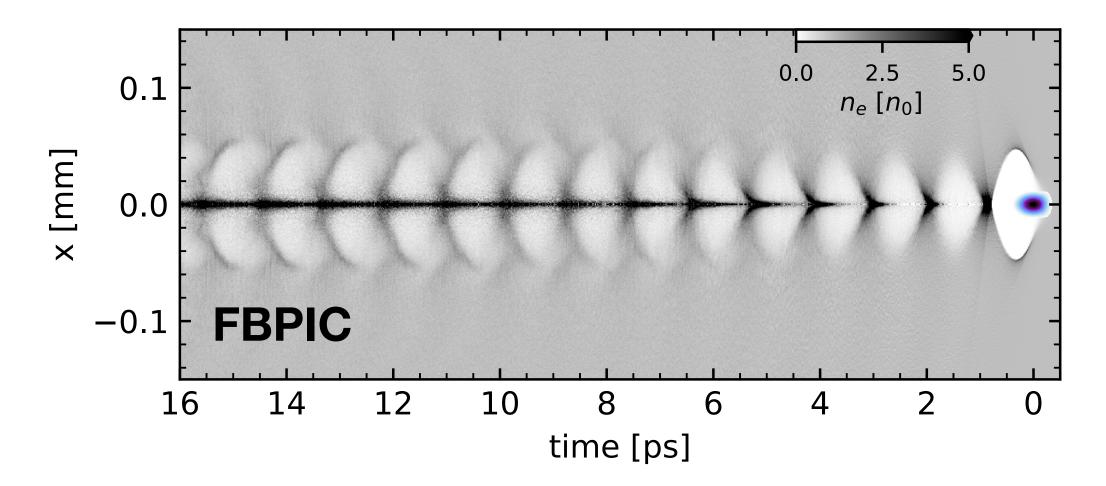
> How:





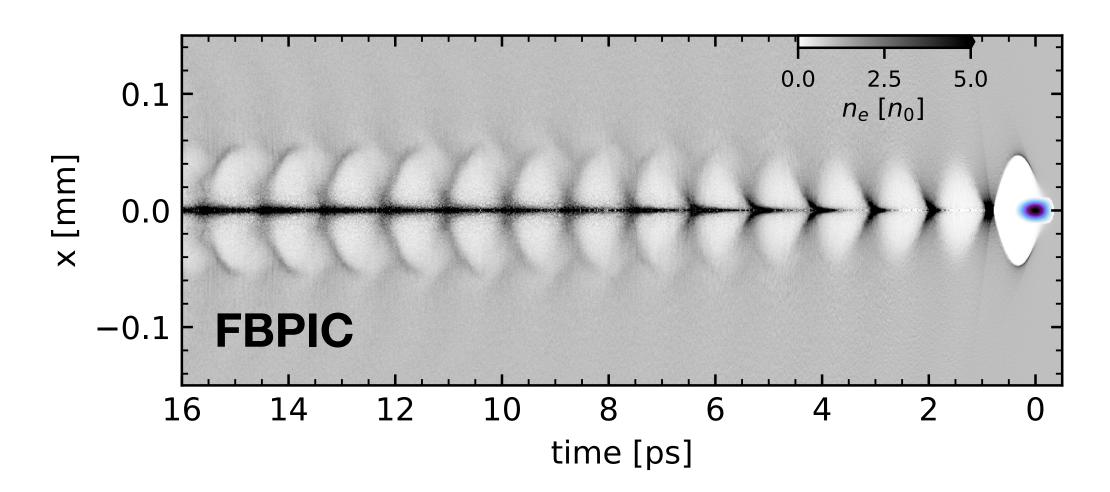
> How:

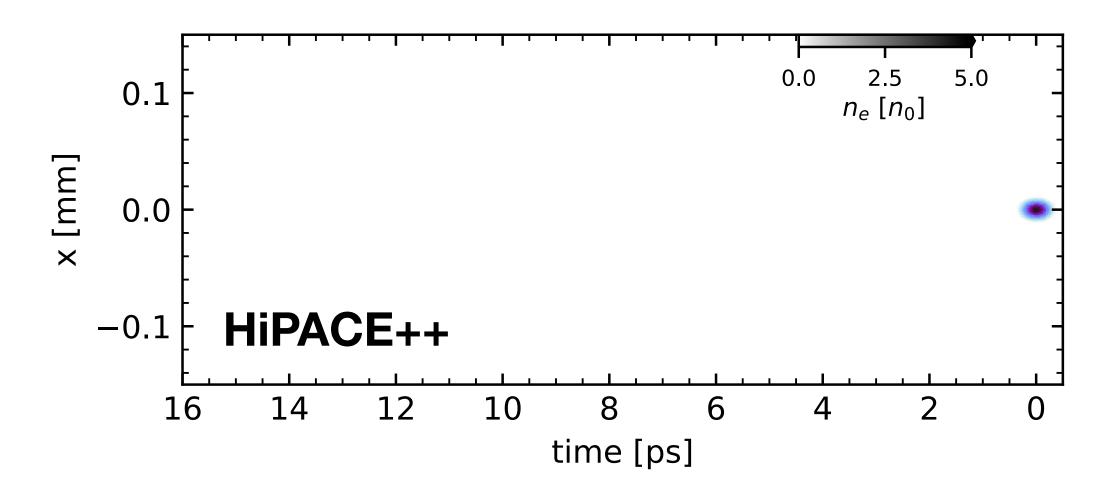




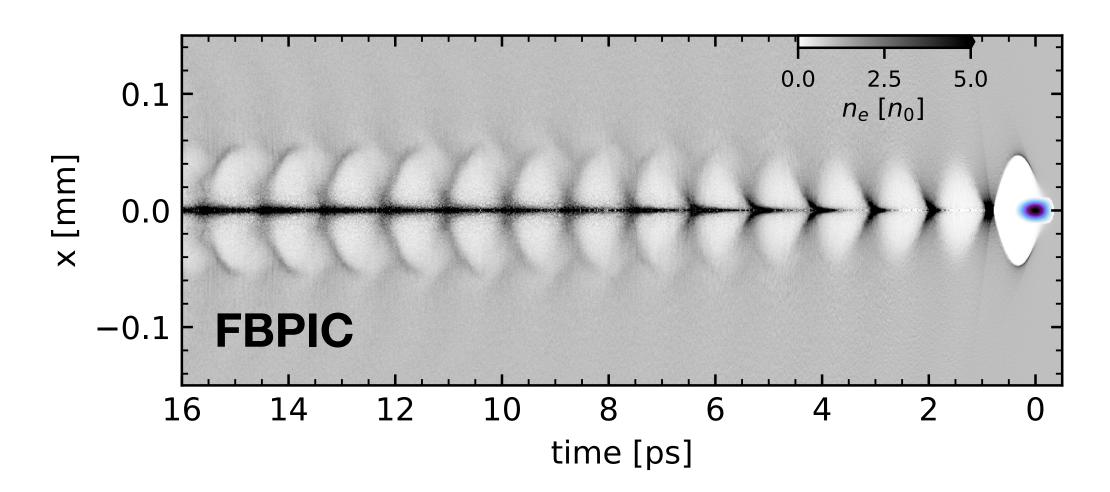
> How:

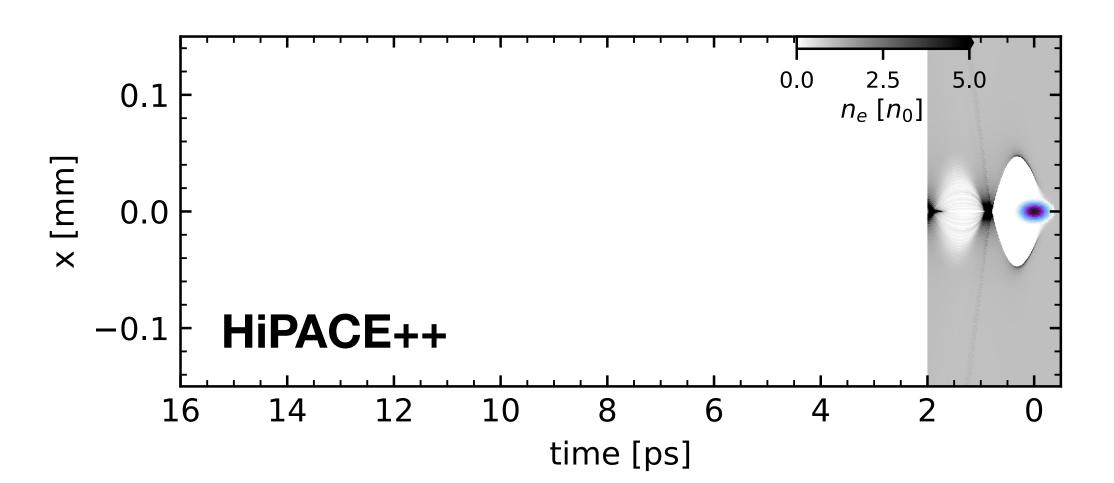




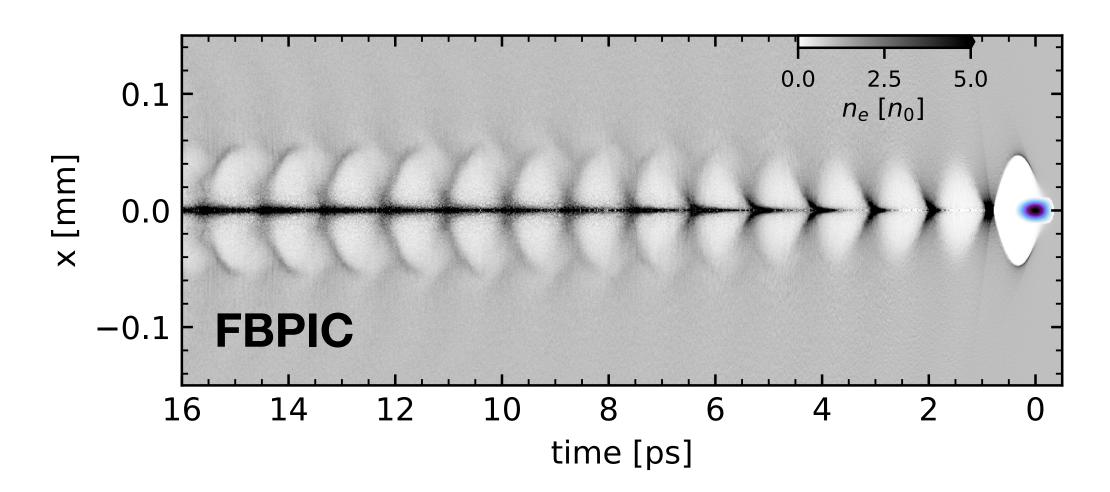


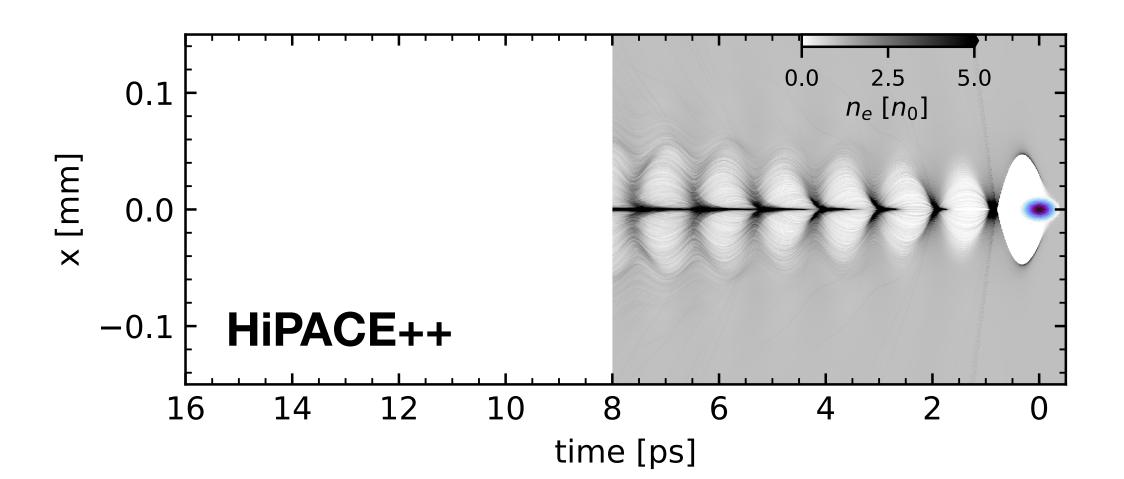
- > FBPIC : Propagate the beam forward in time through a stationary plasma
- > HiPACE++: Propagate the plasma backwards in time past a stationary beam



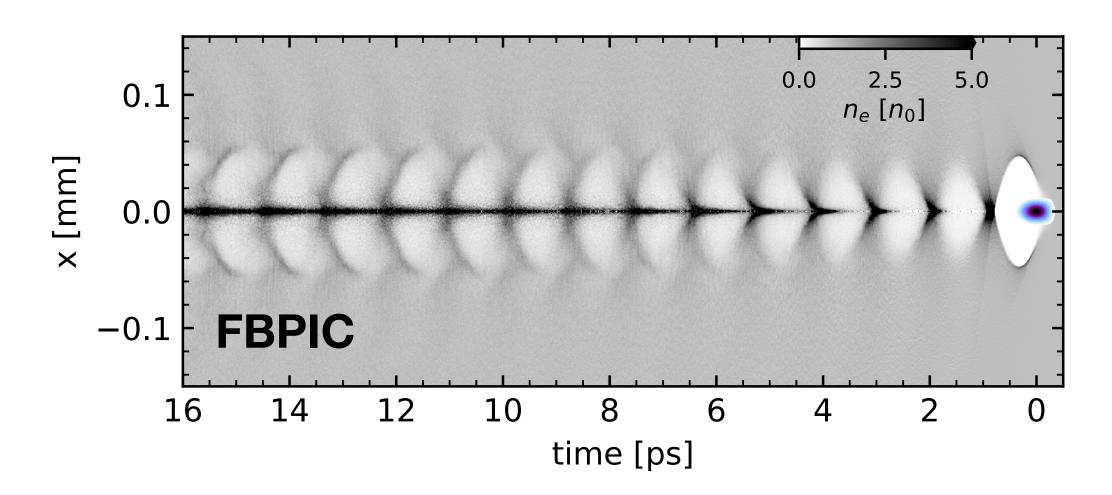


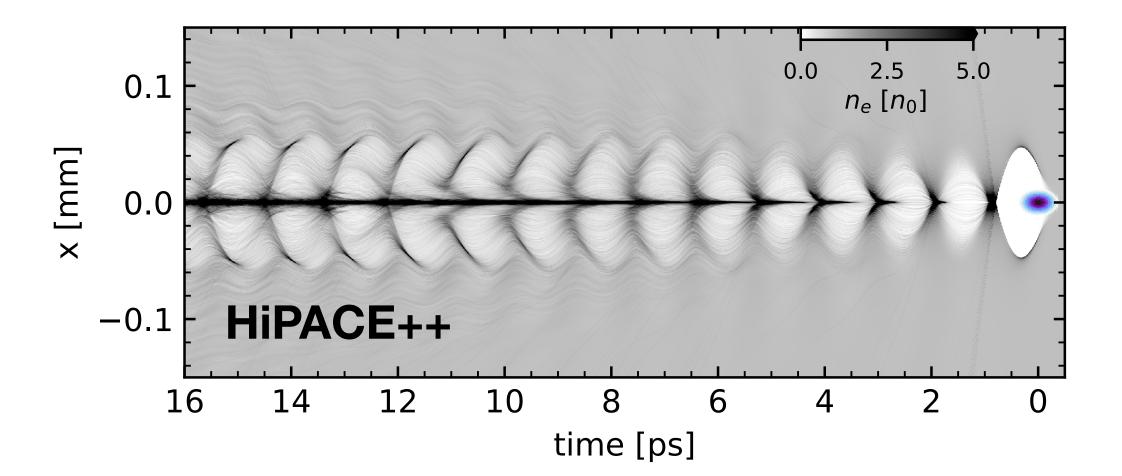
- > FBPIC : Propagate the beam forward in time through a stationary plasma
- > HiPACE++: Propagate the plasma backwards in time past a stationary beam



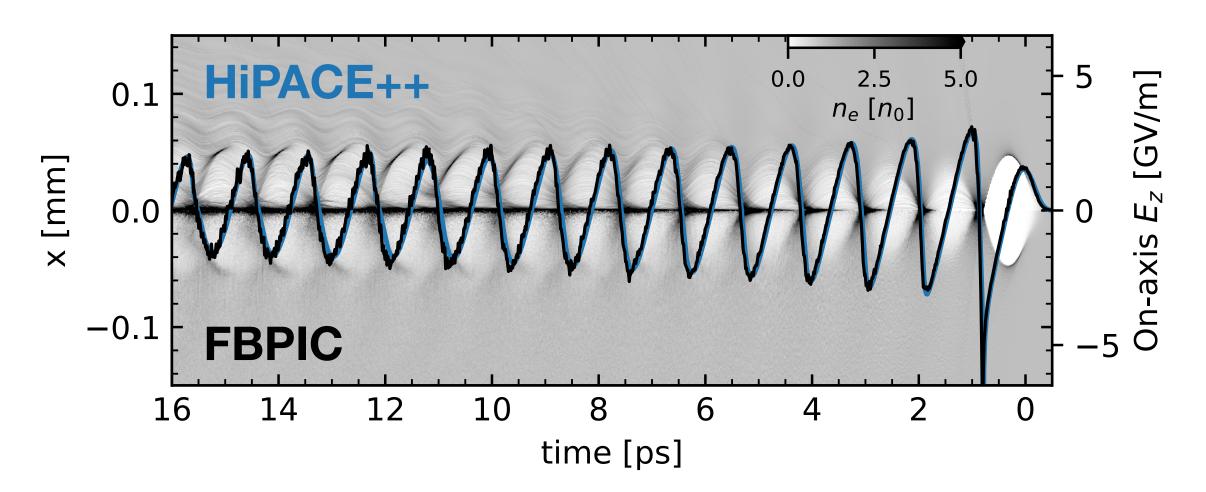


- > FBPIC : Propagate the beam *forward in time* through a stationary plasma
- > HiPACE++: Propagate the plasma backwards in time past a stationary beam



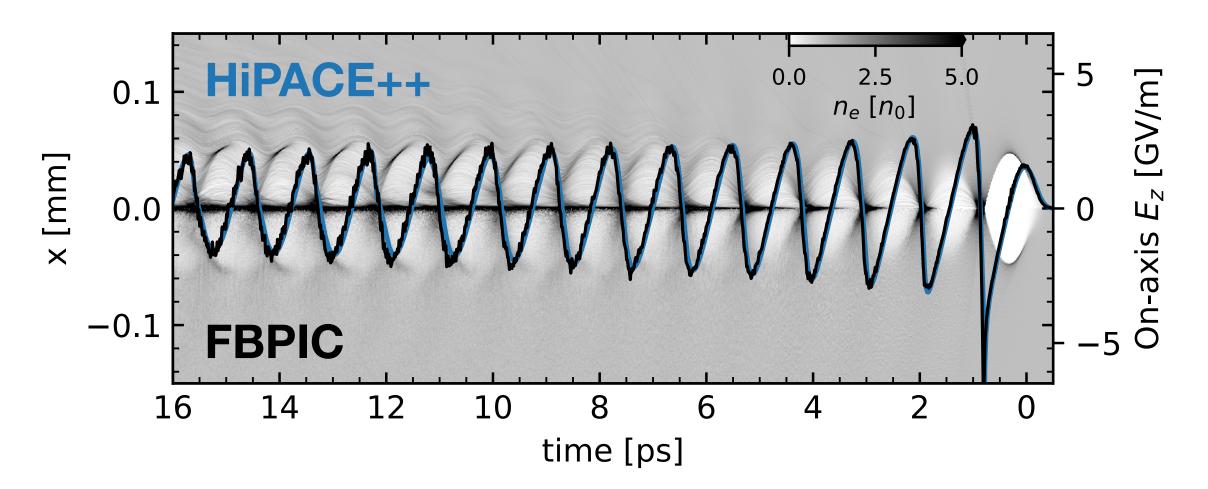


- > FBPIC : Propagate the beam forward in time through a stationary plasma
- > HiPACE++: Propagate the plasma backwards in time past a stationary beam



> Preliminary benchmarking beyond the first bubble:

- > Typically, comparisons of QSA-PIC codes ensure that correct accelerating fields are produced in the first bubble
- > Here we show that a QSA-PIC code can reliably describe plasma evolution at long timescales for cases where:
 - > Well-matched and/or slowly-evolving drivers are used $\sigma_r^2 = \epsilon_n/(\gamma k_\beta)$
 - > Plasma electron trapping is insignificant
- > Offers a boost in computational speedup + reduction in computational resource usage



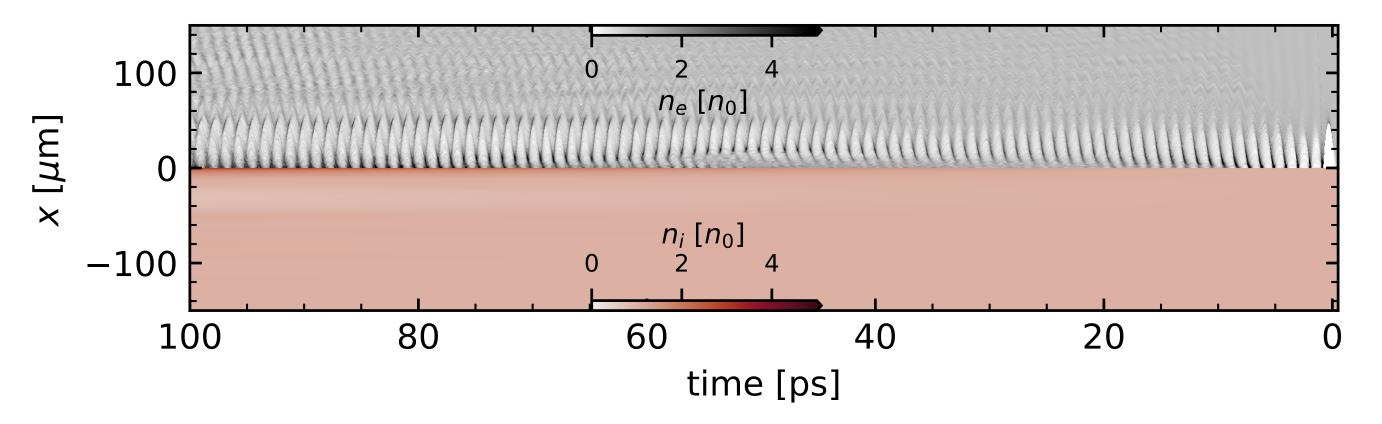
> Preliminary benchmarking beyond the first bubble:

- > Typically, comparisons of QSA-PIC codes ensure that correct accelerating fields are produced in the first bubble
- > Here we show that a QSA-PIC code can reliably describe plasma evolution at long timescales for cases where:
 - > Well-matched and/or slowly-evolving drivers are used $\sigma_r^2 = \epsilon_n/(\gamma k_\beta)$
 - > Plasma electron trapping is insignificant
- > Offers a boost in computational speedup + reduction in computational resource usage

> Advantages:

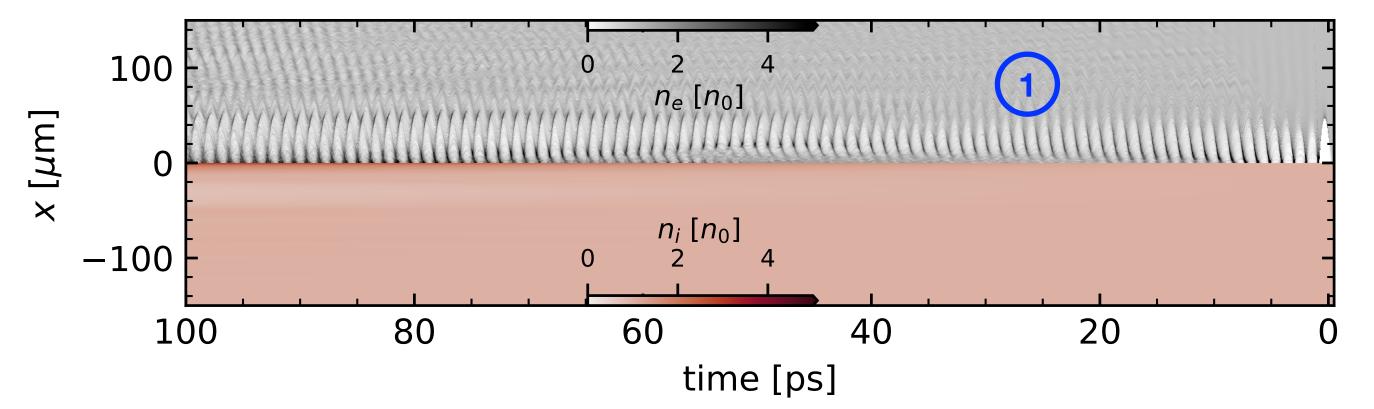
- > 1. Allows long timescales to be simulated feasibly
- > 2. Allows fast parameter scans of plasma evolution at short-medium timescales



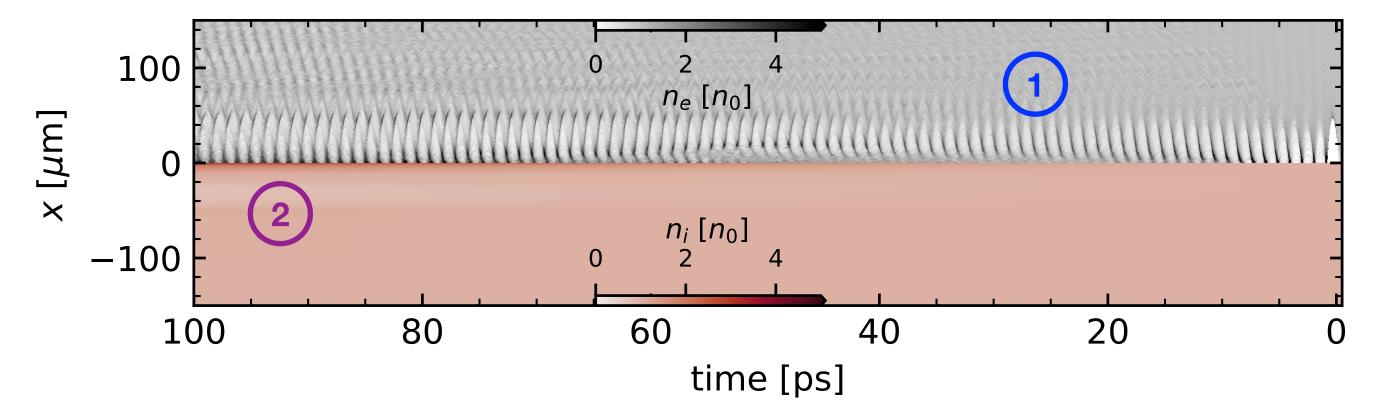




> 1. Prolonged plasma electron oscillations for many plasma periods

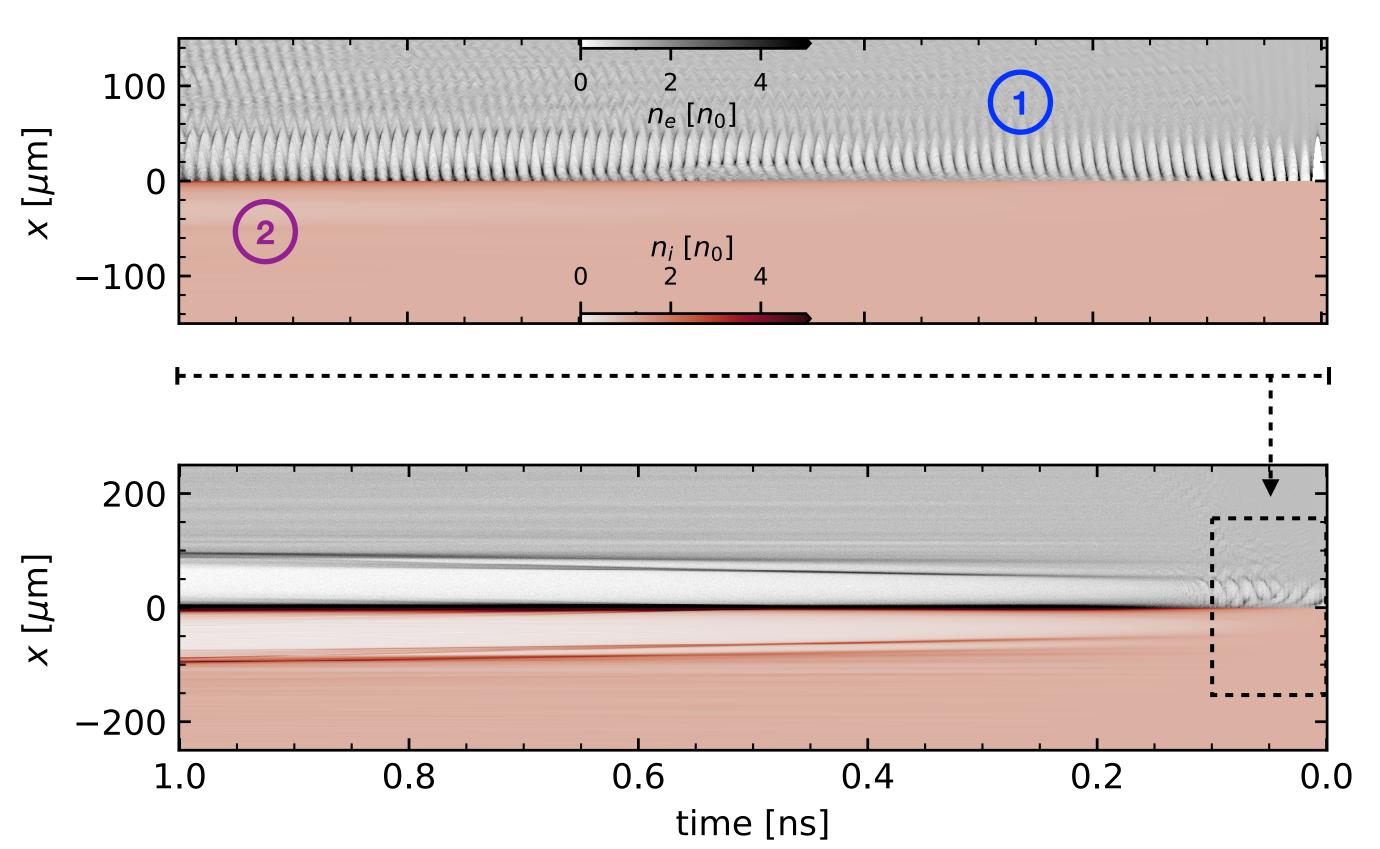


- > 1. Prolonged plasma electron oscillations for many plasma periods
- > 2. Facilitating a slow buildup of a central ion channel



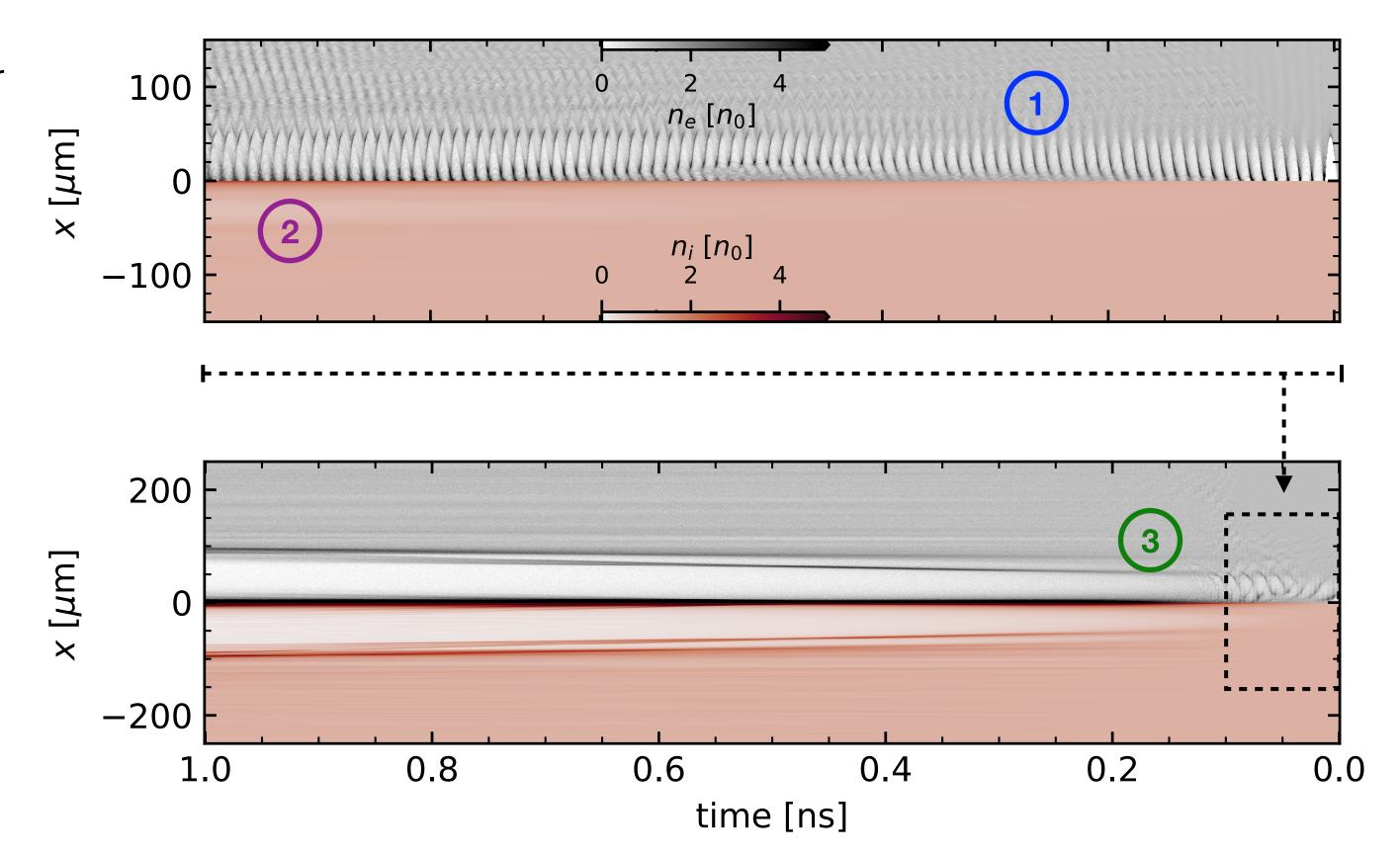


- > 1. Prolonged plasma electron oscillations for many plasma periods
- > 2. Facilitating a slow buildup of a central ion channel



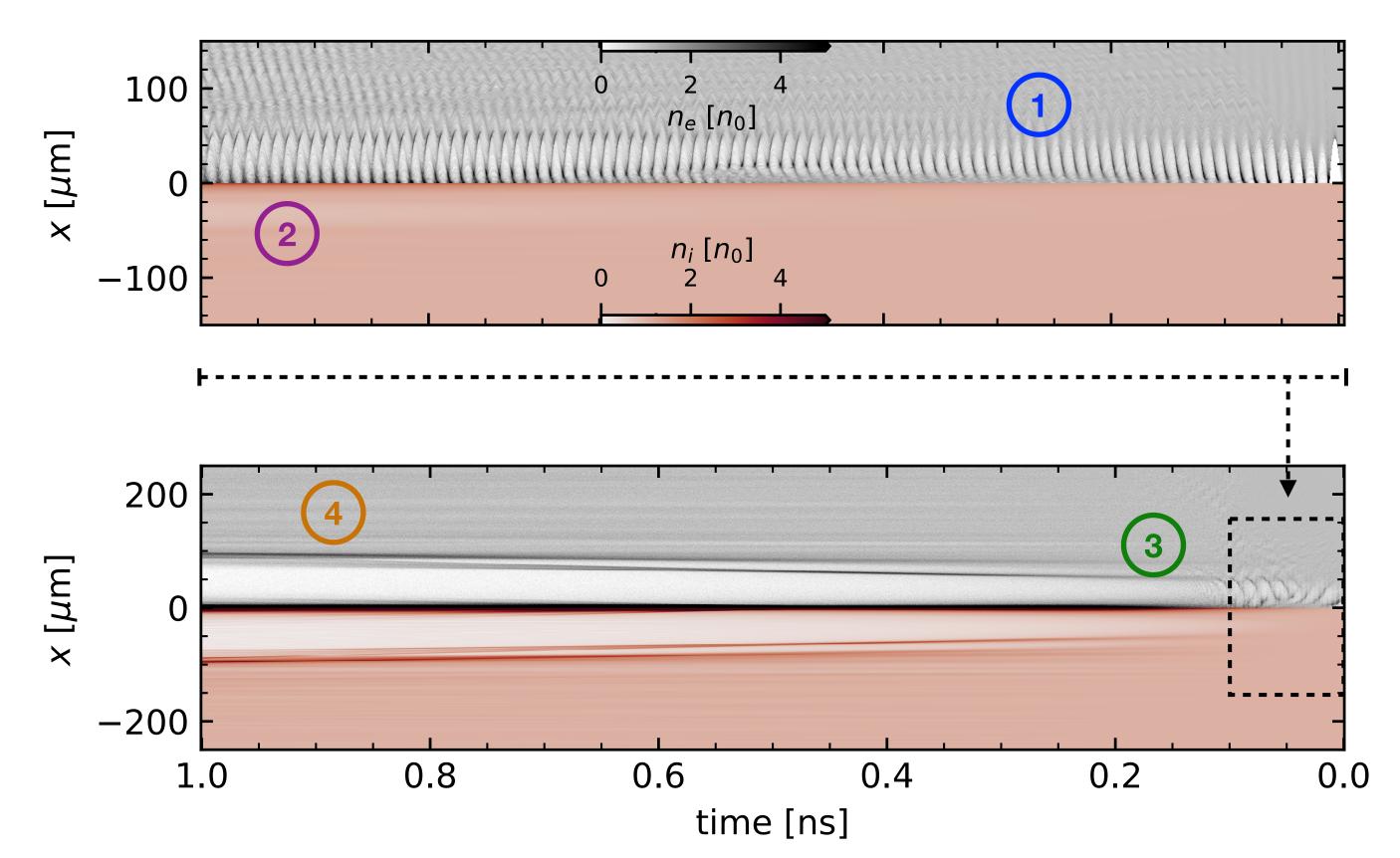


- > 1. Prolonged plasma electron oscillations for many plasma periods
- > 2. Facilitating a slow buildup of a central ion channel
- > 3. Electron oscillation decay energy from the fields in the plasma wave is gradually transferred to the electrons and ions in the channel

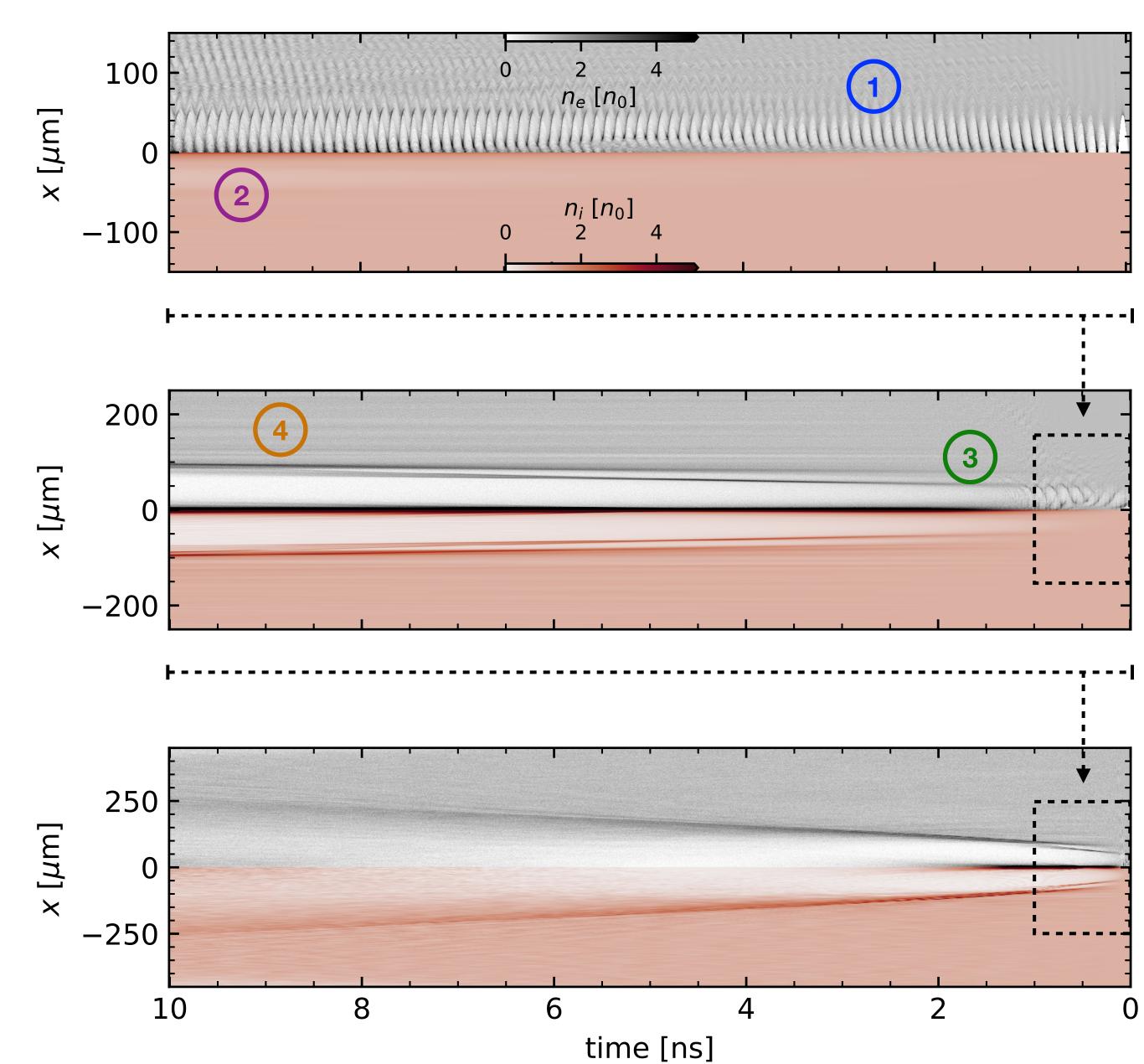




- > 1. Prolonged plasma electron oscillations for many plasma periods
- > 2. Facilitating a slow buildup of a central ion channel
- > 3. Electron oscillation decay energy from the fields in the plasma wave is gradually transferred to the electrons and ions in the channel
- > 4. The motion of the electrons follows the motion of the ions the plasma exhibits quasi-neutrality

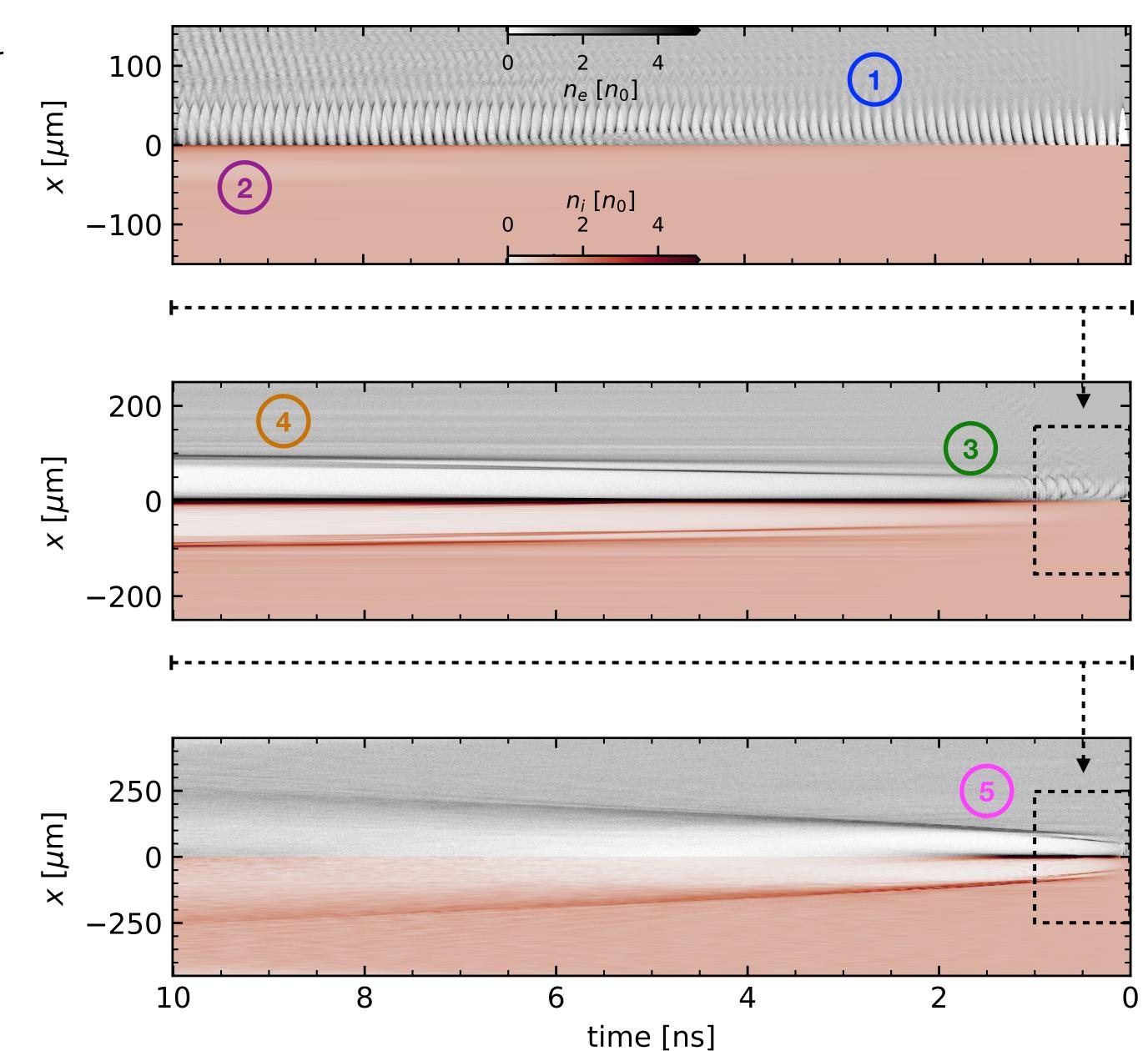


- > 1. Prolonged plasma electron oscillations for many plasma periods
- > 2. Facilitating a slow buildup of a central ion channel
- > 3. Electron oscillation decay energy from the fields in the plasma wave is gradually transferred to the electrons and ions in the channel
- > 4. The motion of the electrons follows the motion of the ions — the plasma exhibits quasi-neutrality



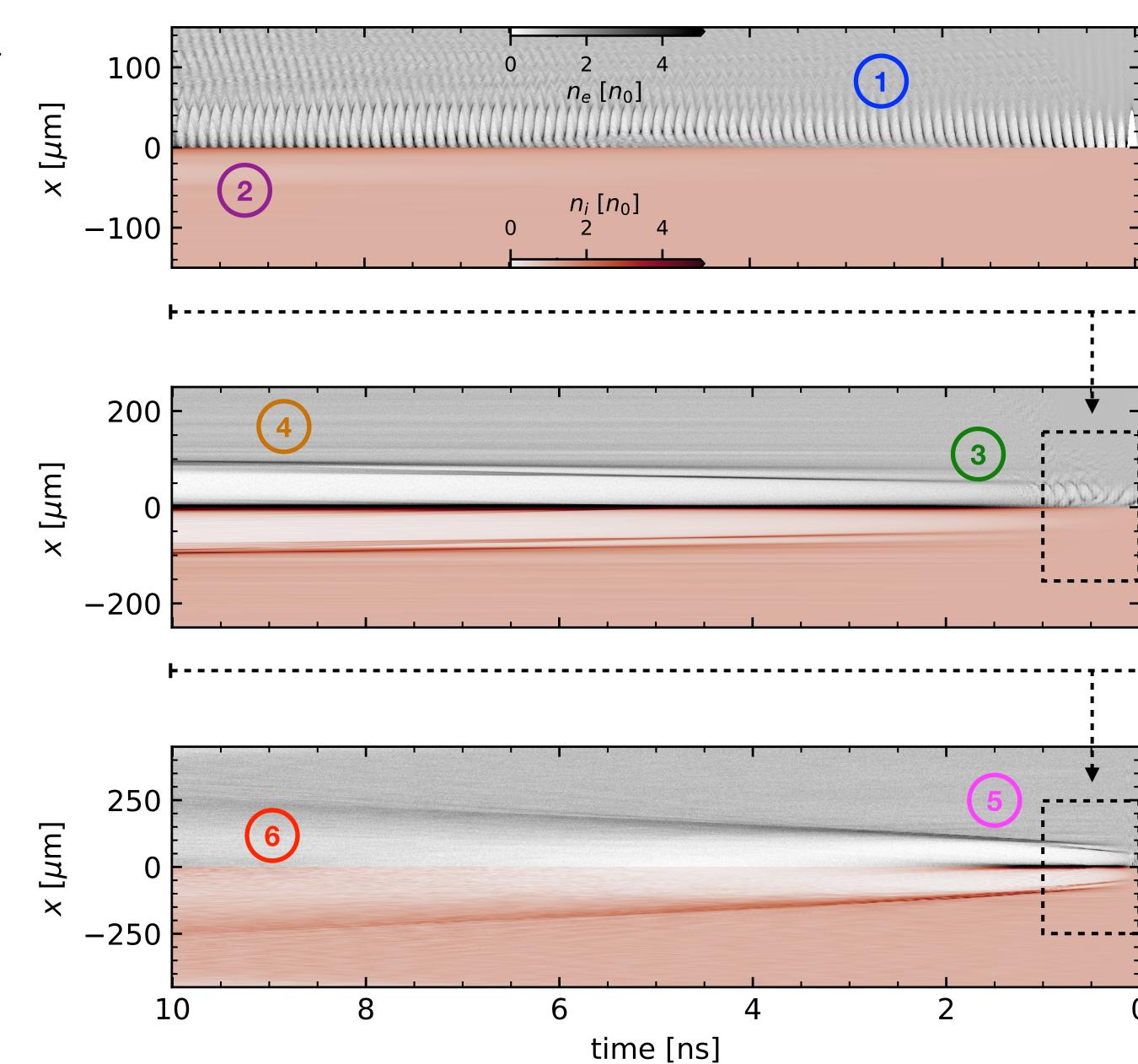


- > 1. Prolonged plasma electron oscillations for many plasma periods
- > 2. Facilitating a slow buildup of a central ion channel
- > 3. Electron oscillation decay energy from the fields in the plasma wave is gradually transferred to the electrons and ions in the channel
- > 4. The motion of the electrons follows the motion of the ions — the plasma exhibits quasi-neutrality
- > 5. Dense ion channel develops followed by a rapid shock expansion — depleting axial plasma density





- > 1. Prolonged plasma electron oscillations for many plasma periods
- > 2. Facilitating a slow buildup of a central ion channel
- > 3. Electron oscillation decay energy from the fields in the plasma wave is gradually transferred to the electrons and ions in the channel
- > 4. The motion of the electrons follows the motion of the ions the plasma exhibits quasi-neutrality
- 5. Dense ion channel develops followed by a rapid shock expansion — depleting axial plasma density
- > 6. At very late timescales, expect plasma to begin to thermalise and equilibrate





Enabling qualitative scalings with ion species

> Considerations of ion mass

- > FLASHForward demonstrated plasma recovery in 63 ns in **argon**
- > HALHF intends to use **helium**
- > How would plasma motion scale?
- > Dominated by the motion of the ions, so one may assume it would scale with the ion's plasma frequency:

timescales $\propto \sqrt{m_i}$

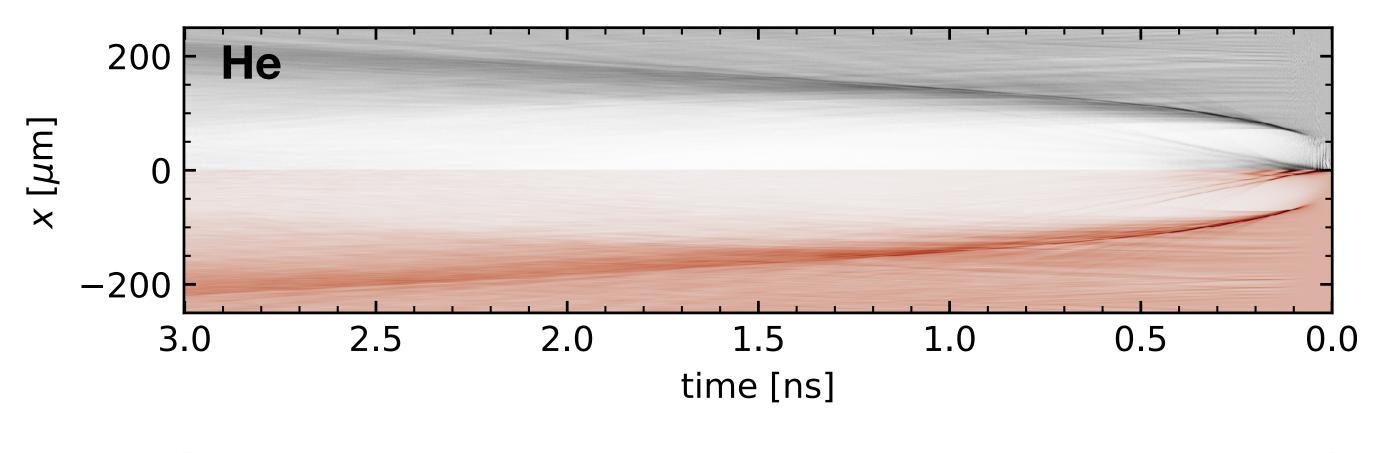


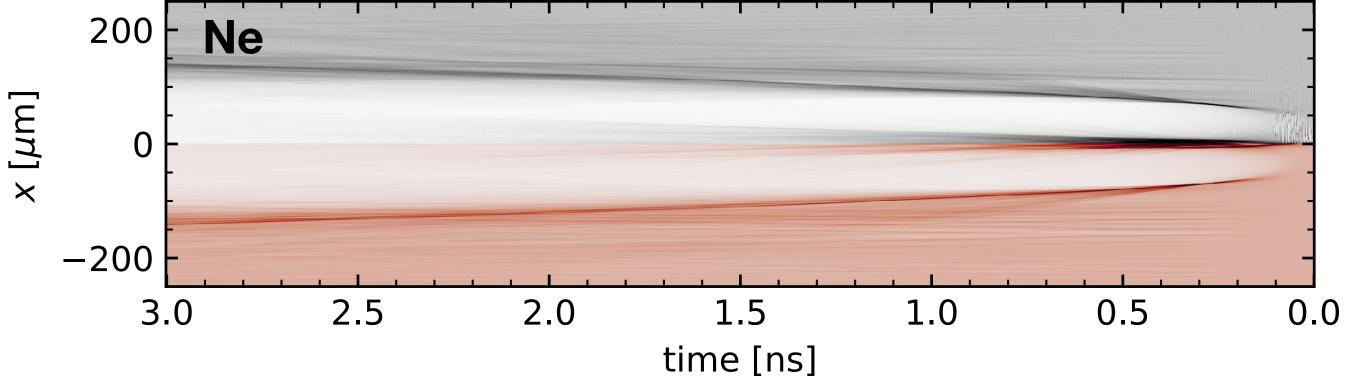
Enabling qualitative scalings with ion species

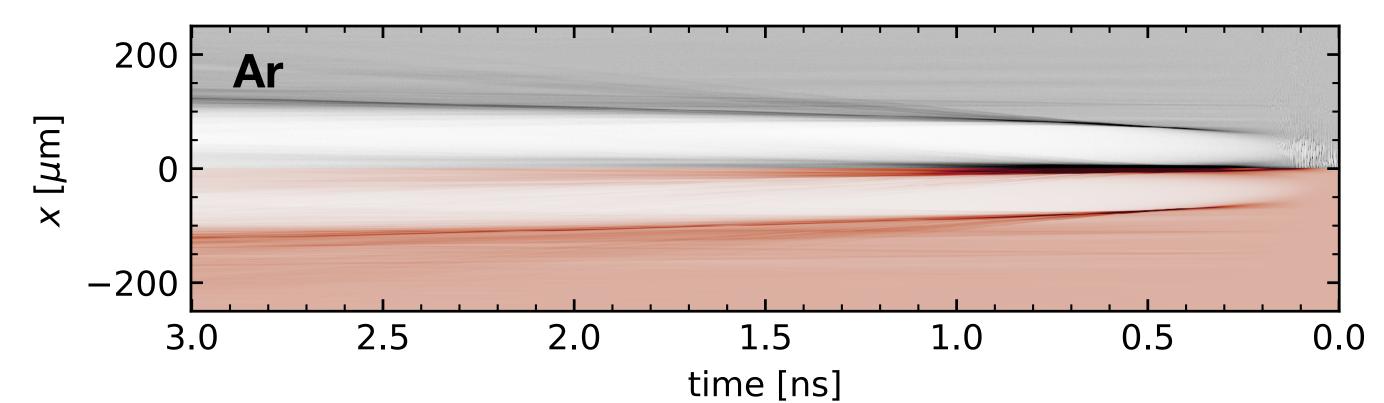
> Considerations of ion mass

- > FLASHForward demonstrated plasma recovery in 63 ns in **argon**
- > HALHF intends to use **helium**
- > How would plasma motion scale?
- > Dominated by the motion of the ions, so one may assume it would scale with the ion's plasma frequency:

timescales $\propto \sqrt{m_i}$







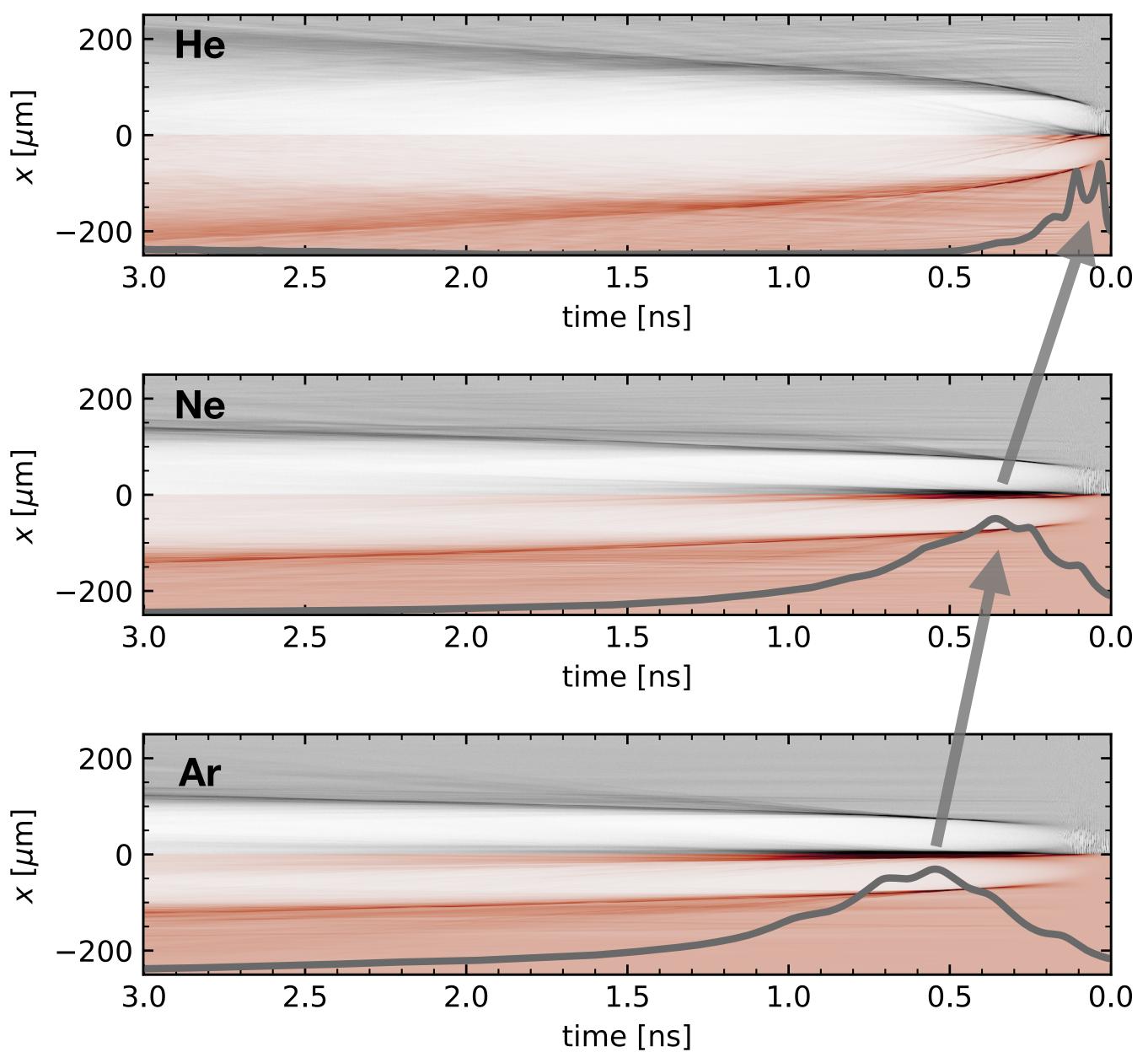


Enabling qualitative scalings with ion species

> Considerations of ion mass

- > FLASHForward demonstrated plasma recovery in 63 ns in **argon**
- > HALHF intends to use **helium**
- > How would plasma motion scale?
- > Dominated by the motion of the ions, so one may assume it would scale with the ion's plasma frequency:

timescales $\propto \sqrt{m_i}$

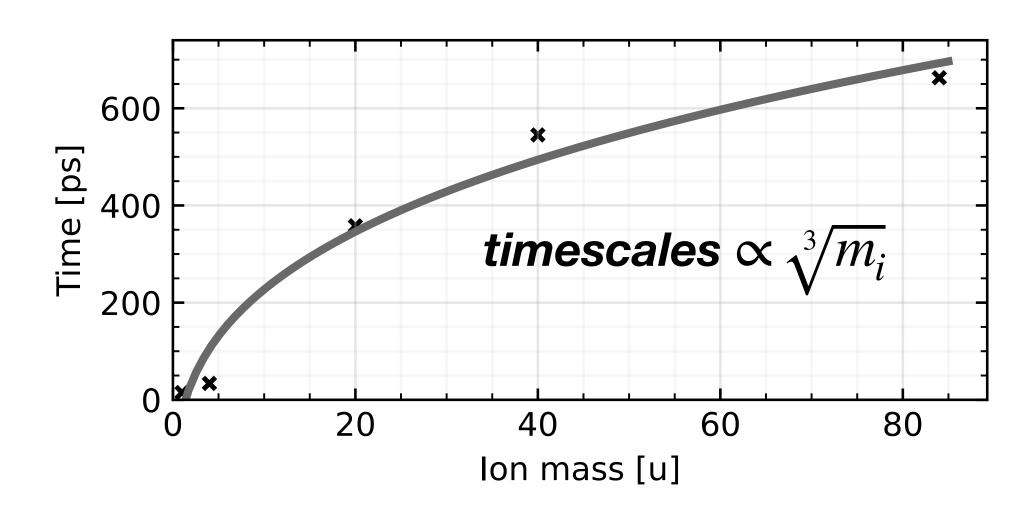


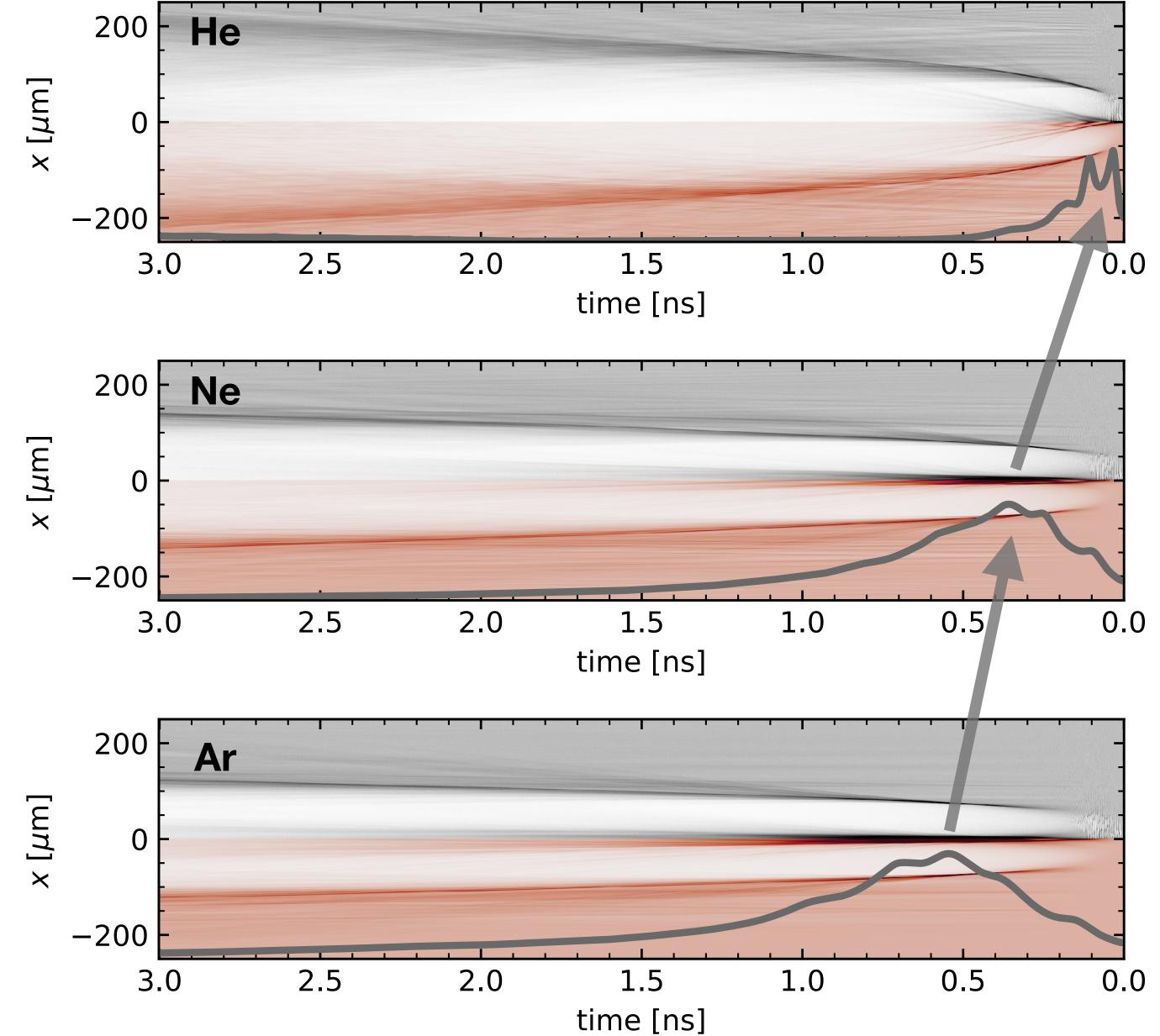
Enabling qualitative scalings with ion species

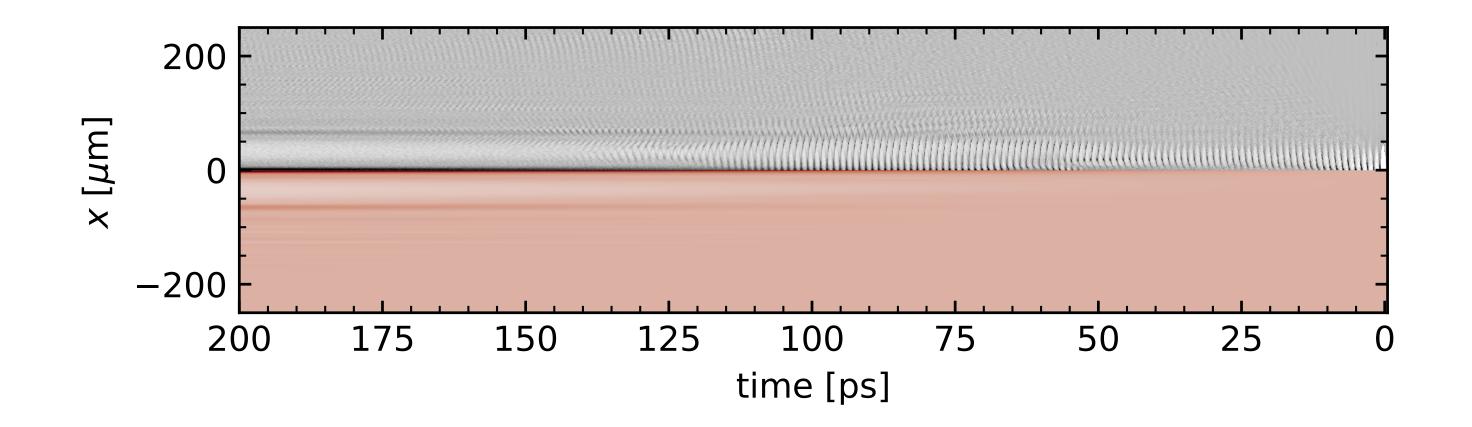
> Considerations of ion mass

- > FLASHForward demonstrated plasma recovery in 63 ns in **argon**
- > HALHF intends to use helium
- > How would plasma motion scale?
- > Dominated by the motion of the ions, so one may assume it would scale with the ion's plasma frequency:

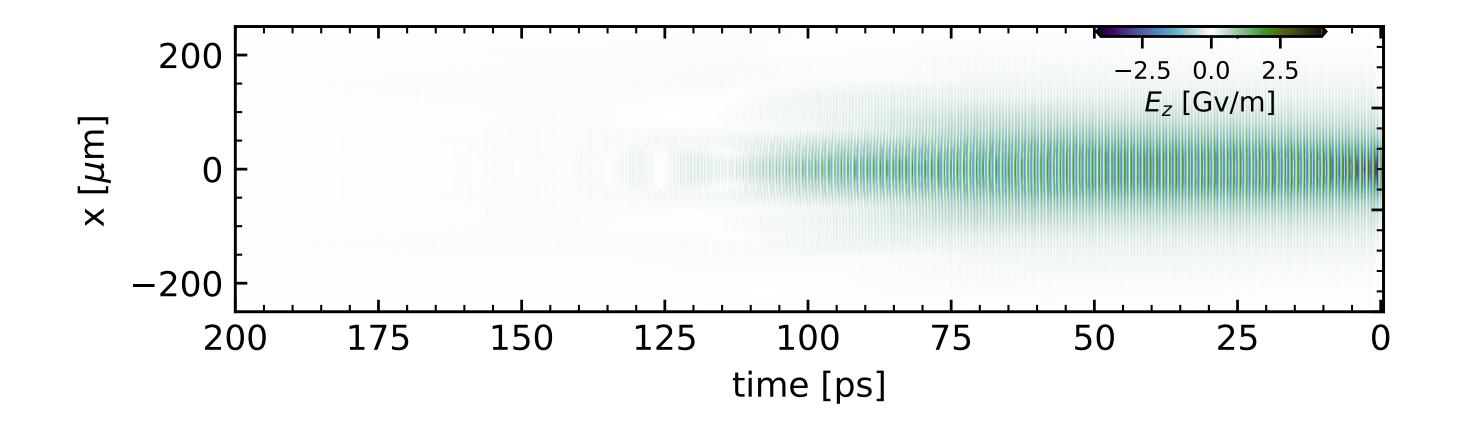
timescales
$$\propto \sqrt{m_i}$$



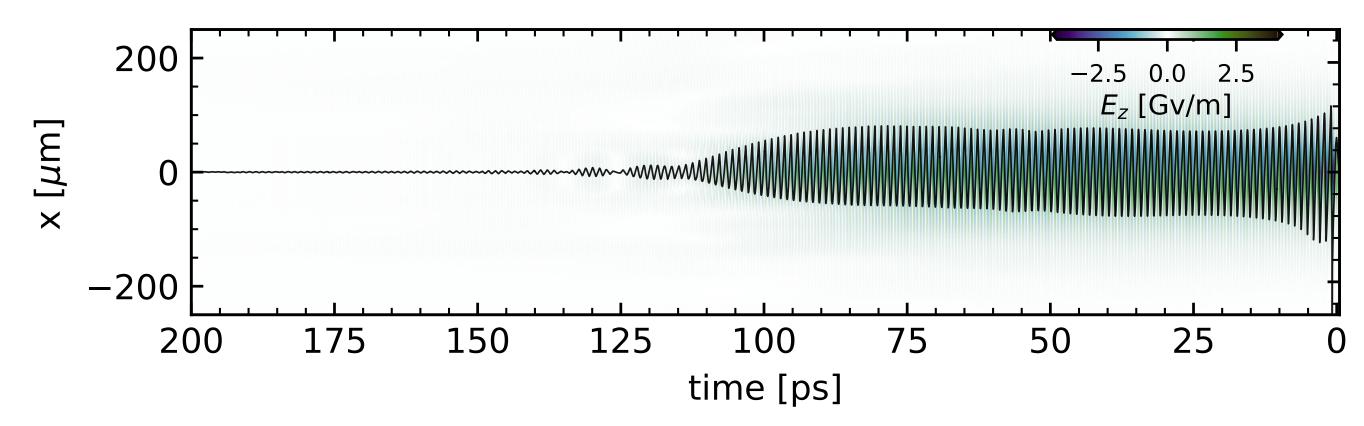






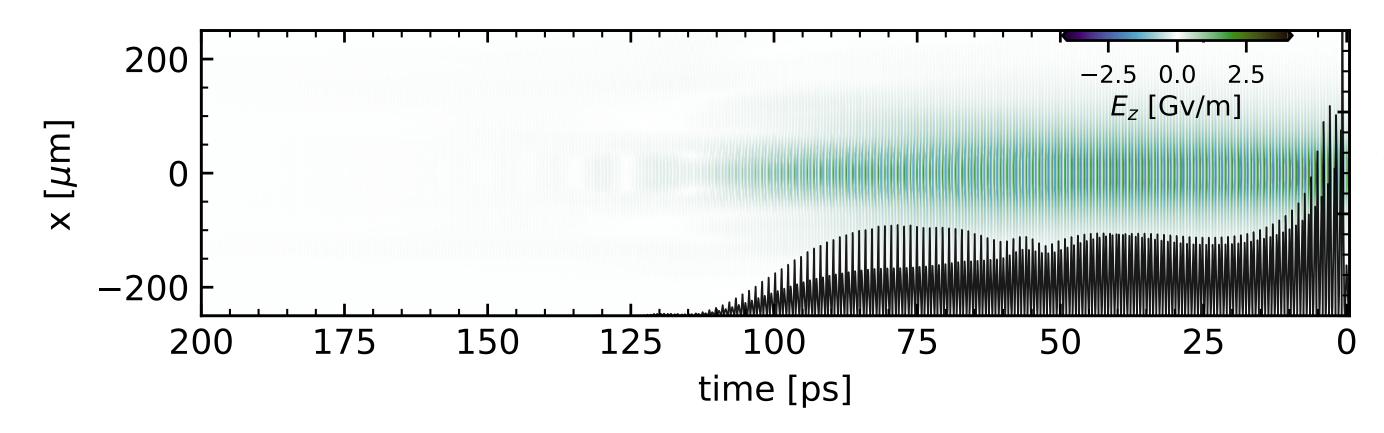




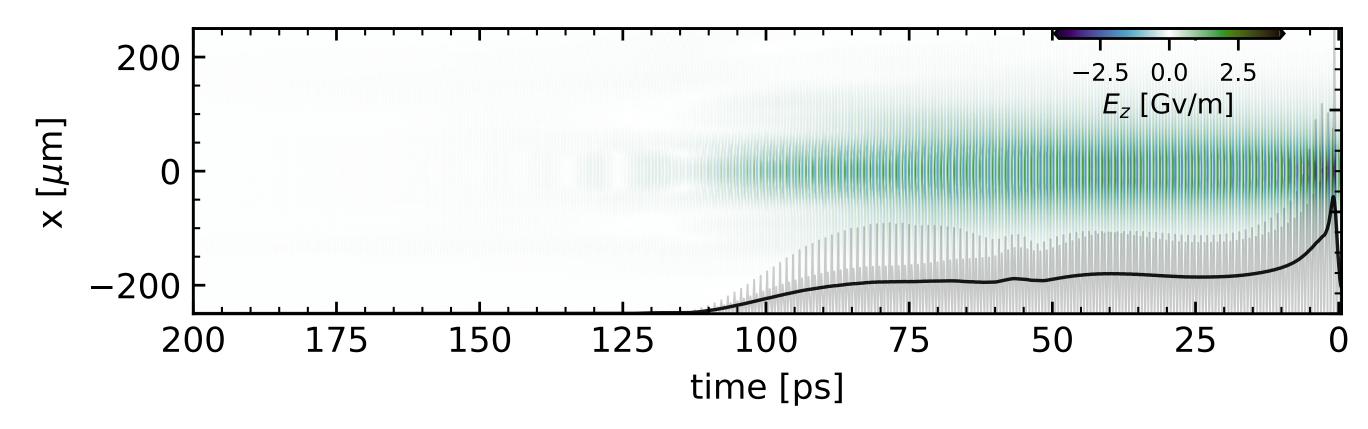


1. Look at axial longitudinal electric field



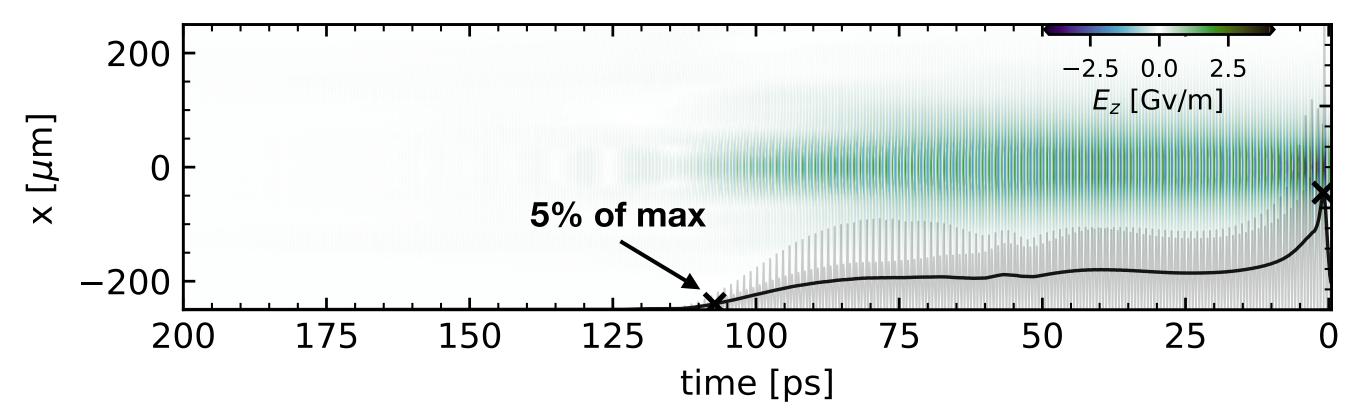


- 1. Look at axial longitudinal electric field
- 2. Look at axial field energy



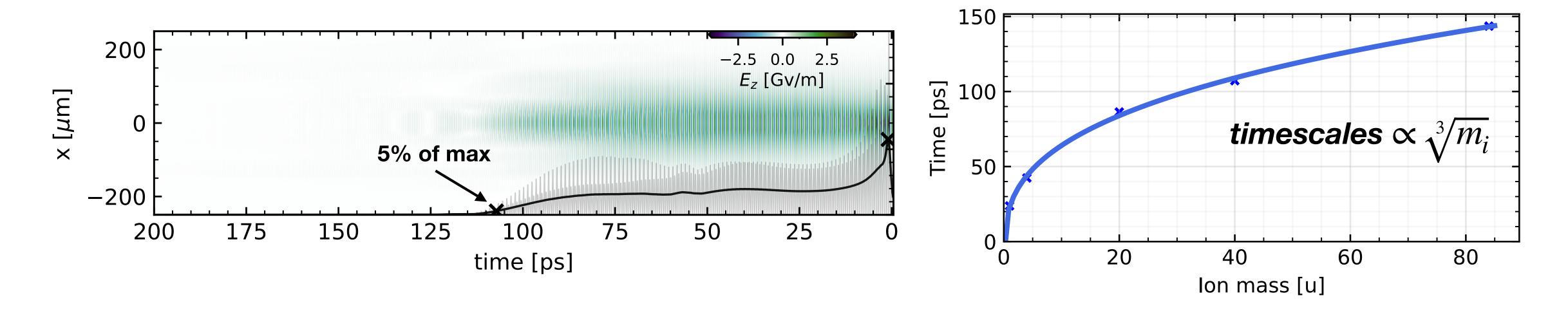
- 1. Look at axial longitudinal electric field
- 2. Look at axial field energy
- 3. Gaussian smoothing giving mean field energy



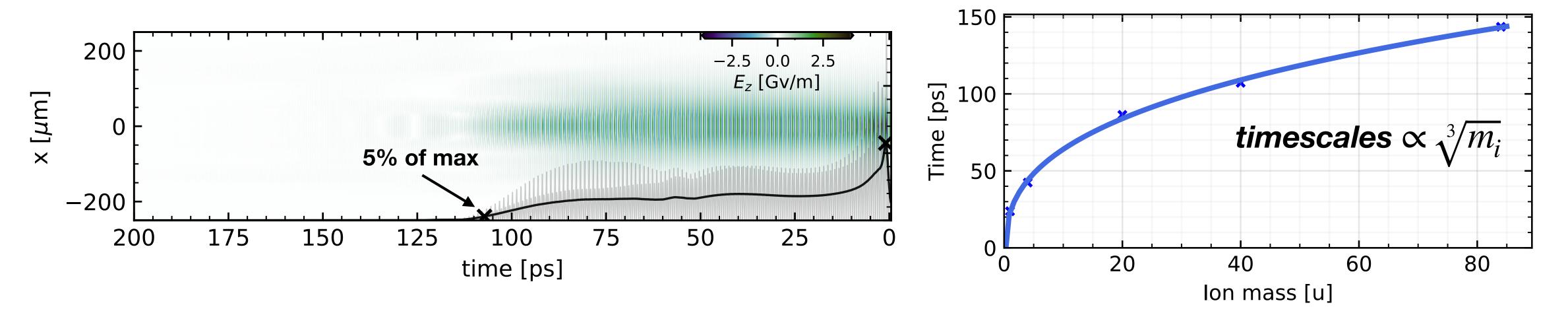


- 1. Look at axial longitudinal electric field
- 2. Look at axial field energy
- 3. Gaussian smoothing giving mean field energy
- 4. Pick low threshold of field energy to signify wave decay time



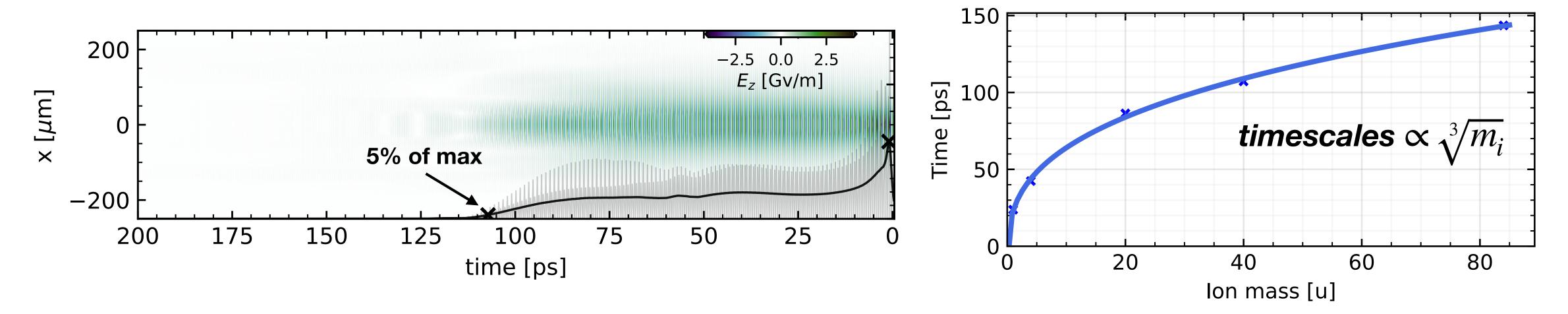






> Plasma evolution is a series of causes and effects:

- > Beam driver deposits energy into forming nonlinear plasma wave
- > lons slowly quench energy of the wave
- > This depletes the wave with a $\sqrt[3]{m_i}$ dependence which has been predicted and observed:
 - > Spitsyn et al. Phys. Plasmas 25, 103103 (2018) & M. Turner et al., Phys. Rev. Lett. 134, 155001 (2025)
 - > As a result, the ion spike peak ensues with the same $\sqrt[3]{m_i}$ dependence



> Plasma evolution is a series of causes and effects:

- > Beam driver deposits energy into forming nonlinear plasma wave
- > lons slowly quench energy of the wave
- > This depletes the wave with a $\sqrt[3]{m_i}$ dependence which has been predicted and observed:
 - > Spitsyn et al. Phys. Plasmas 25, 103103 (2018) & M. Turner et al., Phys. Rev. Lett. 134, 155001 (2025)
 - > As a result, the ion spike peak ensues with the same $\sqrt[3]{m_i}$ dependence
- > Plasma evolution and recovery relies on the interdependence of many consecutive interlinked effects
- > But so far, these were all only for a cold plasma...

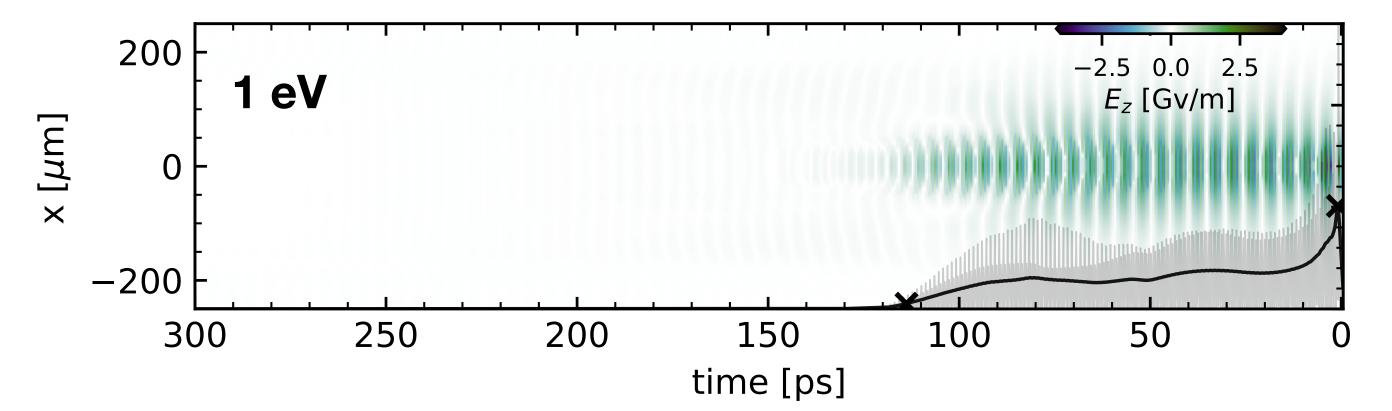


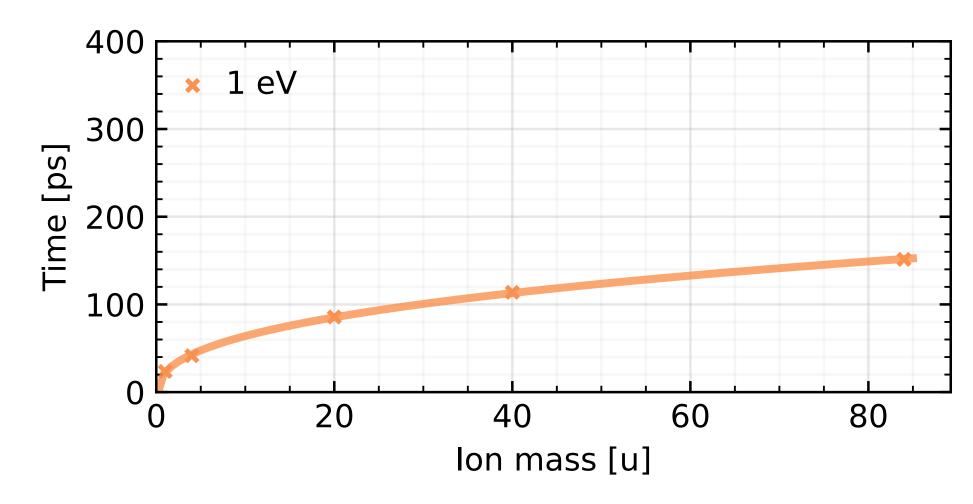
- > Expectations and effects
 - > At hot temperatures, expect random thermal motion to damp the wave faster



> Expectations and effects

> At hot temperatures, expect random thermal motion to damp the wave faster

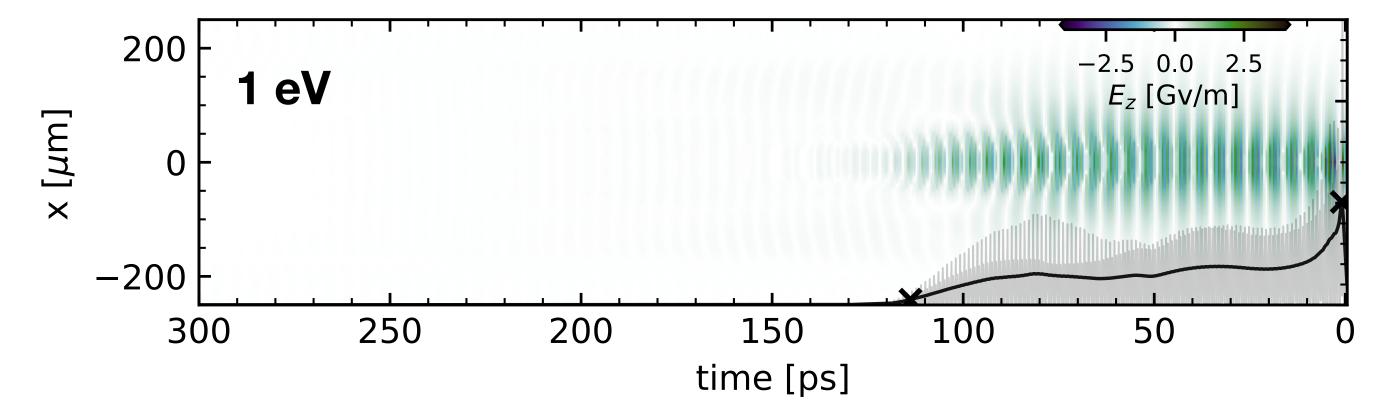


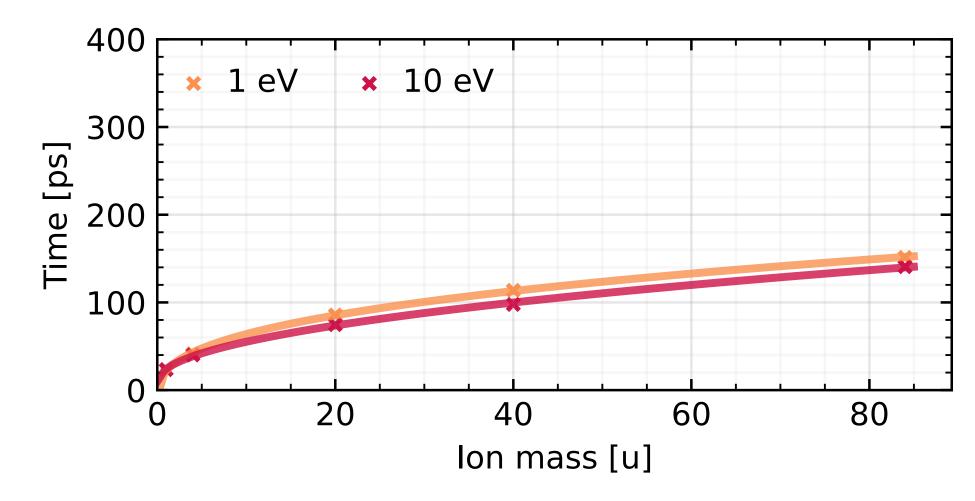




> Expectations and effects

- > At hot temperatures, expect random thermal motion to damp the wave faster
- > See little changes at low temperatures

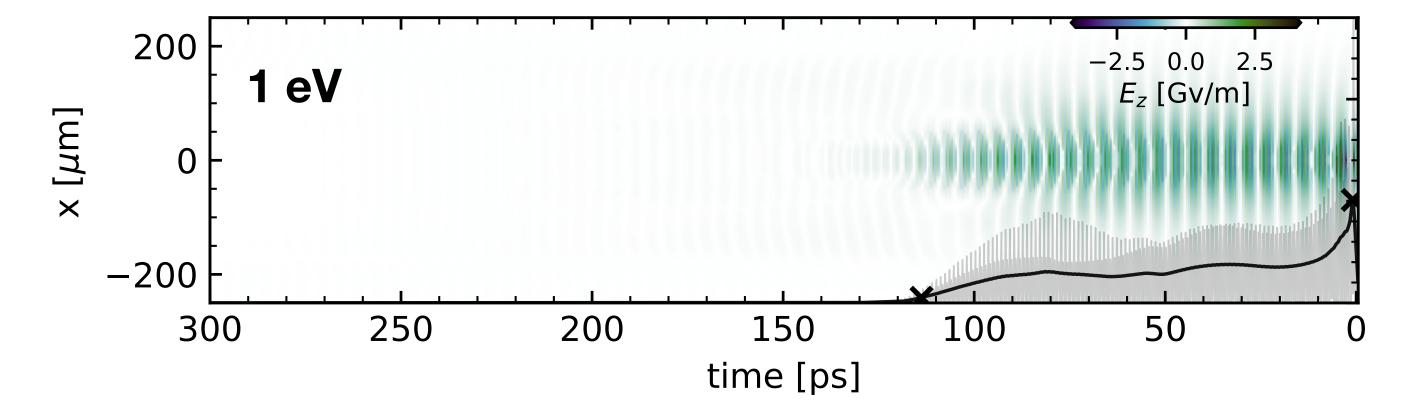


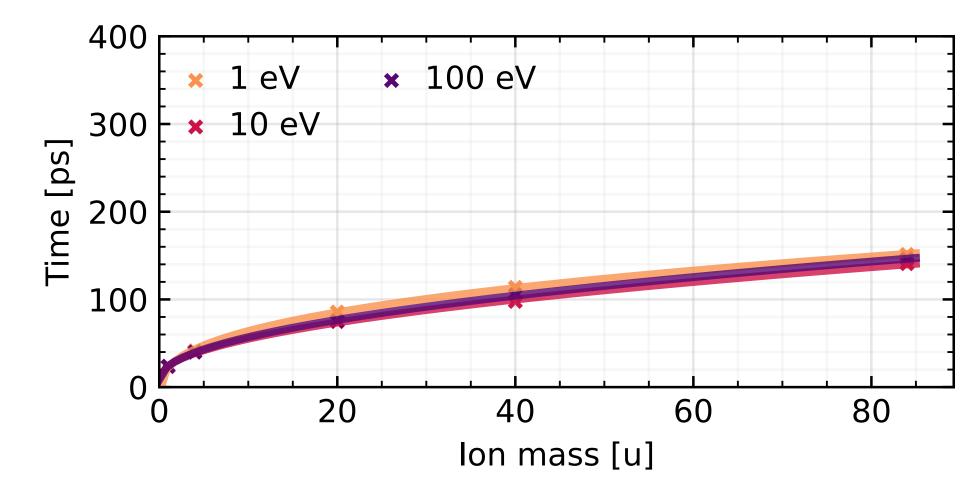




> Expectations and effects

- > At hot temperatures, expect random thermal motion to damp the wave faster
- > See little changes at low temperatures

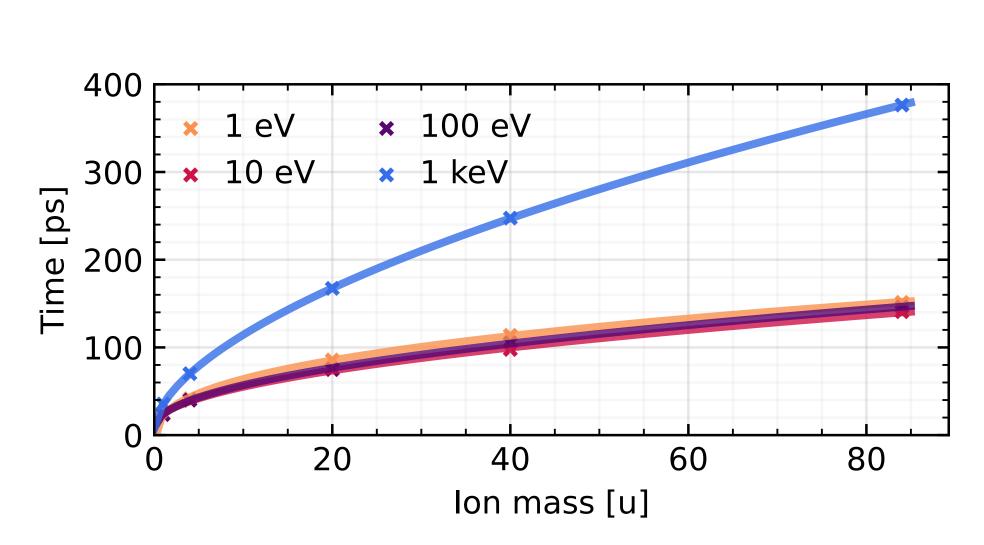


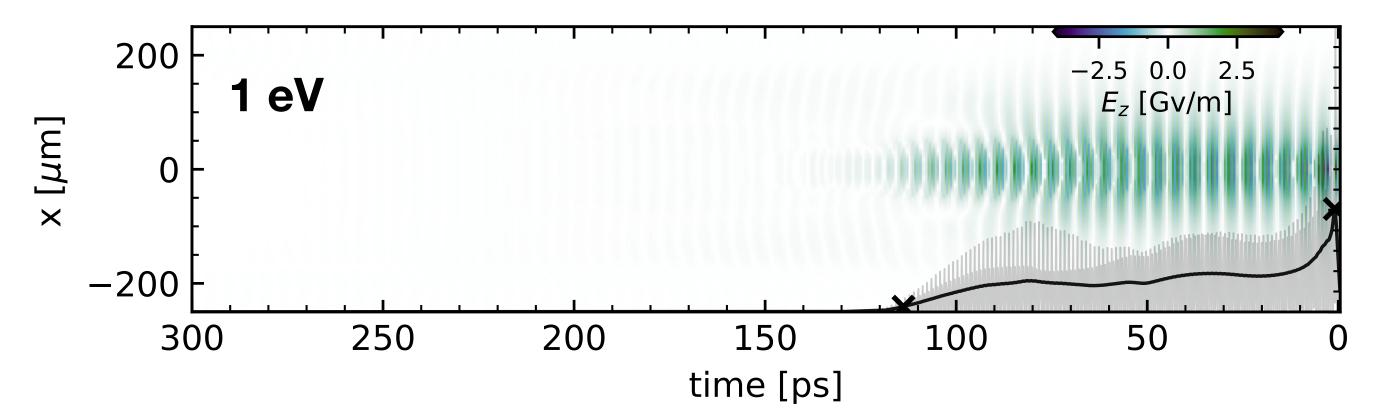


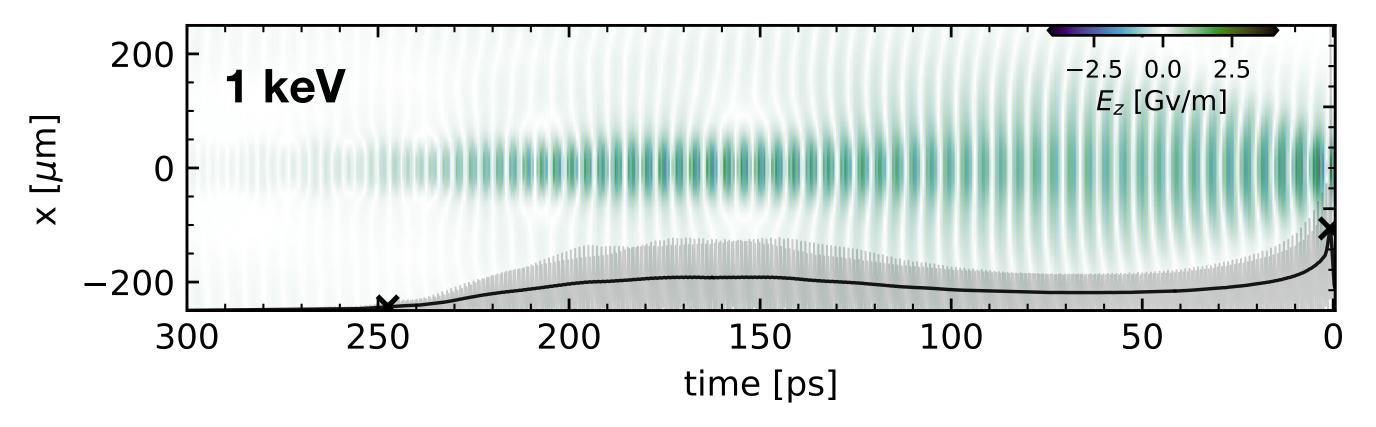


> Expectations and effects

- > At hot temperatures, expect random thermal motion to damp the wave faster
- > See little changes at low temperatures
- > See region of high temperature which prolongs the plasma wave

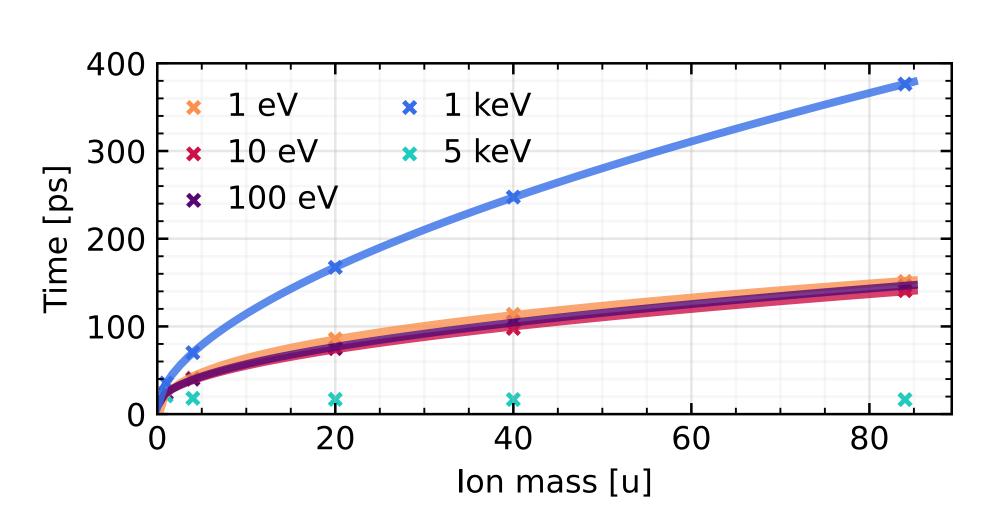


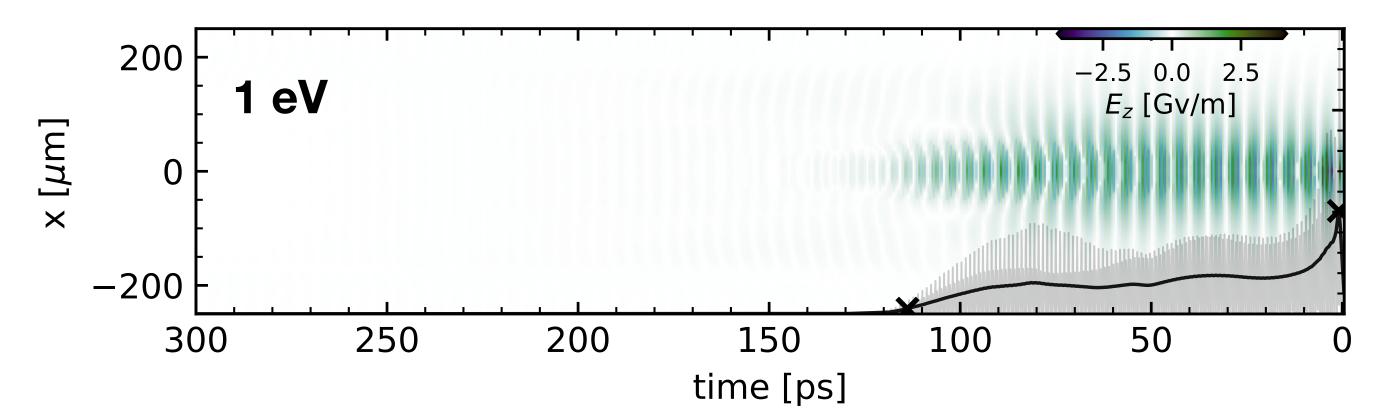


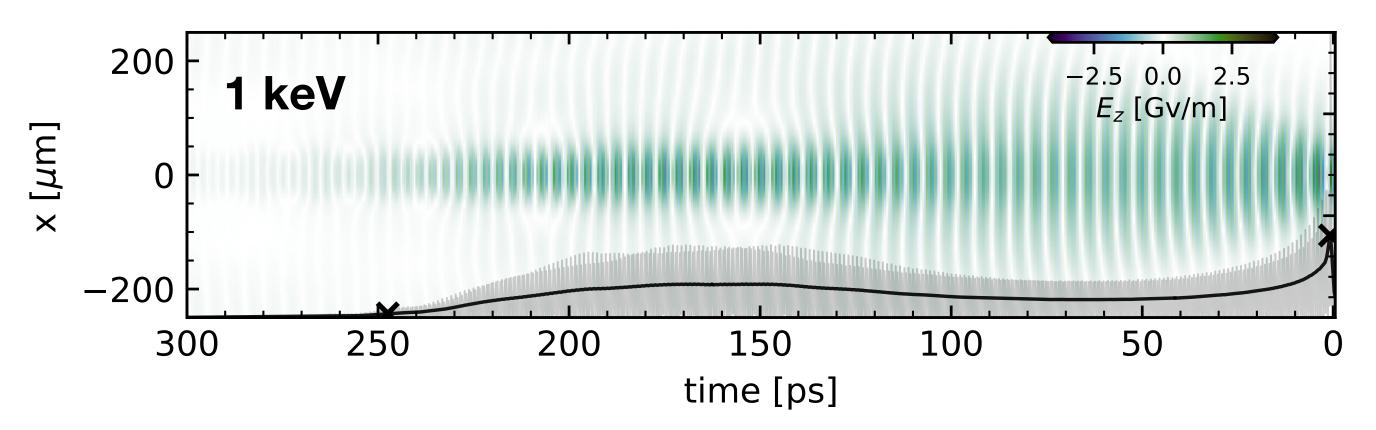


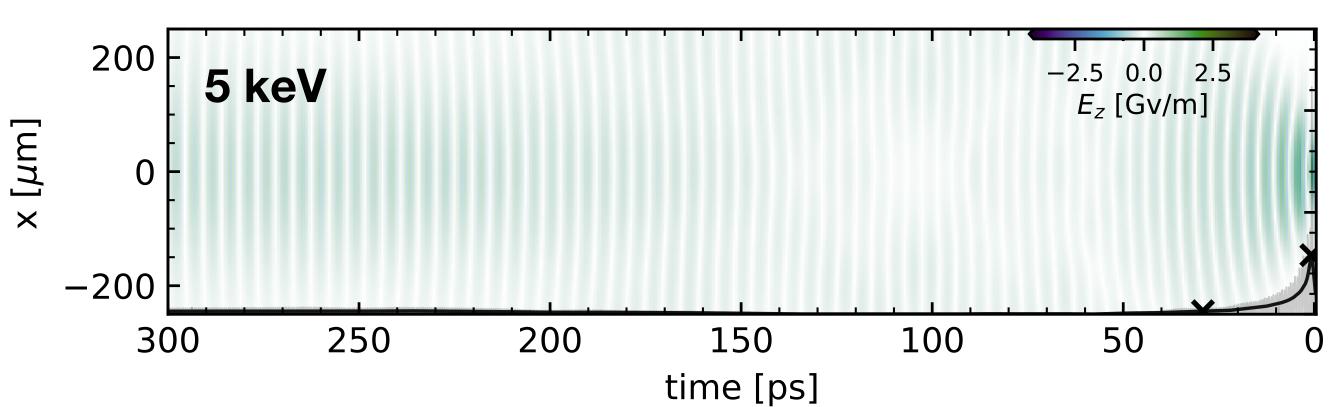
> Expectations and effects

- > At hot temperatures, expect random thermal motion to damp the wave faster
- > See little changes at low temperatures
- > See region of high temperature which prolongs the plasma wave
- See expected quick damping at extreme temperatures
- > Suggests use of extreme temperature plasma to be factored into designs?











Conclusions

> Summary:

- > Temperature effects are important for collider applications as they will deteriorate beam quality
- > Relevant timescales cannot be captured by PIC codes so introduced Quasi-Static Approximation
 - > QSA allows multi-nanosecond simulations extending beyond plasma quasi-neutrality
- > Demonstrated effects of ion species on plasma wave decay
 - > He could be used to run high-repetition-rate machines like HALHF if a $\sqrt[3]{m_i}$ scaling is followed
- > But plasma recovery timescales found to be dependent on initial plasma temperature

> Next steps:

- > Address several limitations in HiPACE++ relevant to long timescales
 - > Ionization, collisions, recombination etc.
- > Extend plasma evolution to longer timescales to reproduce experimental data of plasma recovery from FLASHForward
- > Apply these methods to inform on future experiments at FLASHForward
- > Apply these methods to inform on future baselines and modifications for the HALHF collider

Thank you!



Extra slide - simulation parameters & fits

Slide 2 simulations (HiPACE++)

Plasma: 1e16 /cm3, Singly ionised Ar

Driver parameters: 250 pC, 1 µm transverse rms, 40 µm longitudinal rms

Witness parameters: 65 pC, 1 µm transverse rms, 15 µm longitudinal rms

Slides 6-8 simulations (HiPACE++ & FBPIC)

Plasma density: 1e16 /cm3

Driver parameters: 600 pC, 5 µm transverse rms, 50 µm longitudinal rms

Slide 9-12 simulations (HiPACE++ varying 9. ion mass and 10. temperature in Ar)

Plasma density: 1e16 /cm3

Driver parameters: 567 pC (nb/n0=100), 5 µm transverse rms, 50 µm longitudinal rms

Fitting parameters

All fits of t vs m_ion are to: t = a*m_ion**b + c



Extra slide - simulation parameters & fits

Slide 2 simulations (HiPACE++)

Plasma: 1e16 /cm3, Singly ionised Ar

Driver parameters: 250 pC, 1 µm transverse rms, 40 µm longitudinal rms

Witness parameters: 65 pC, 1 µm transverse rms, 15 µm longitudinal rms

Slides 6-8 simulations (HiPACE++ & FBPIC)

Plasma density: 1e16 /cm3

Driver parameters: 600 pC, 5 µm transverse rms, 50 µm longitudinal rms

Slide 9-12 simulations (HiPACE++ varying 9. ion mass and 10. temperature in Ar)

Plasma density: 1e16 /cm3

Driver parameters: 567 pC (nb/n0=100), 5 µm transverse rms, 50 µm longitudinal rms

Fitting parameters

All fits of t vs m_ion are to: t = a*m_ion**b + c

Scan	b
Ion mass	0.31
Wave decay	0.33

T [eV]	b
1	0.37
10	0.49
100	0.49
1000	0.58