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Developing an ultra-short laser pulse for probing relativistic laser-plasma interactions

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Abstract

Ultra-short laser pulses are essential to resolve femtosecond-timescale dynamics in plasma-based particle accelerators.

Here we present the process behind a high-intensity hollow-core beam line designed to spectrally broaden an input pulse spectrum by a factor 5.58, with an output pulse energy of ~2 mJ. Our design enables Fourier limited compression of the pulse to ~6 fs and will be utilized to take crisp shadowgrams of plasma wave sub-structures in LWFAs and PWFAs and furthermore used as a photocathode injector laser in a PWFA.

The proposed experimental arrangement is based on the mechanism of self-phase modulation, which requires an intense, short laser pulse and a material with strong third-order nonlinearity. We plan to use noble gases, which are well suited to act as nonlinear media with their comparatively high ionization levels and sufficiently large nonlinear refractive indices.

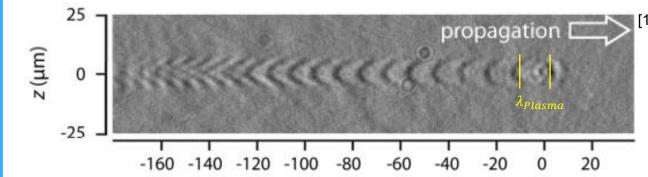
The required intensity for self-phase modulation is achieved by focusing a pre-compressed laser pulse down to a smaller beam diameter. A hollow-core fiber is used to maintain the intensity over a longer distance and clean the pulse profile for later use in shadowgraphy.

Challenges in designing such a hollow-core fiber beam line are spatial limitations, ionization defocusing, laser induced damages on optics and self-focusing effects due to the high pulse intensity.

Motivation

Ultra-short laser pulses with pulse durations of a few femtoseconds are required to resolve processes on timescales of tens of femtoseconds, such as the structure and motion of plasma waves in relativistic laser-plasma interactions.

Resolving a plasma wave requires exposure times of $t_{max} \leq \frac{\lambda_{Plasma}}{2*c_0}$.



Object size $\sim 10 \mu m \rightarrow t_{max} \sim 16 fs$

Figure 1: An exemplary shadowgram, indicating the size of the first plasma wave

A short pulse means a wide spectral range, though typical Ti:Sa lasers are limited in their spectral range. Self phase modulation is exploited to widen the spectral range of laser pulses and further reduce the Fourier limited pulse duration.

In addition to the short pulse duration, a clean pulse profile is required to ensure a good contrast on the shadowgrams.

Since hollow-core fibers act as spatial filters, they are preferred over other methods, such as gas-filled Heriott cells or thin plates.

Self phase modulation

- Direct consequence of the intensity dependent refractive index $n = n_0 + n_2 I$
- Maximum phase shift given by $\Delta \phi_{max} = \frac{\omega_0 n_2}{c_0} I_0 L_{eff}$ [2]
- Spectral broadening factor is $F = \frac{\Delta \omega_{out}}{\Delta \omega_{in}} = \sqrt{1 + \frac{4}{3\sqrt{3}}} \Delta \phi_{max}^2 = \frac{\Delta \tau_{in}}{\Delta \tau_{out}}$ [2]
- Spectral broadening imprints positive chirp on the pulse

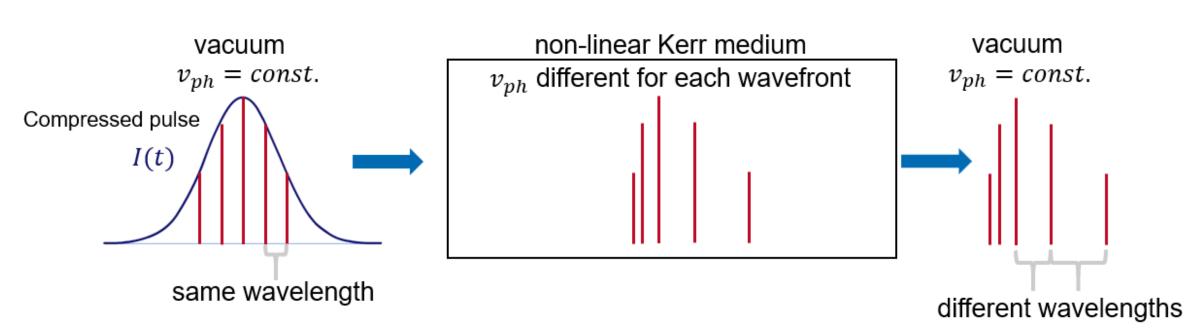


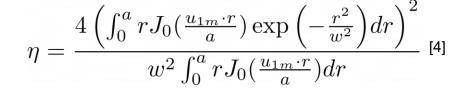
Figure 2: Simplified visualization of the physical process during self phase modulation.

sensitive optical component and compare with

the scaled damage threshold of the optic

Gaussian pulses and hollow-core fibers

- Pure Gaussian TEM modes can not exist inside the fiber, electronic hybrid modes are excited instead
- Coupling efficiency of excitation depends on focal spot size of laser and inner radius of fiber a
- Effective length of the nonlinear process depends on inner radius of fiber a, the fiber materials and the nonlinear mediums refractive index



$$L_{eff} = \frac{1 - \exp\left(-\frac{\alpha_m}{L}\right)}{\alpha_m}$$
 [2]

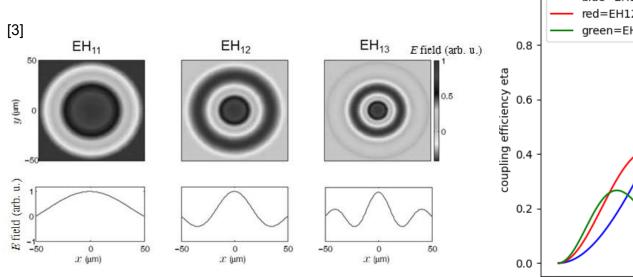
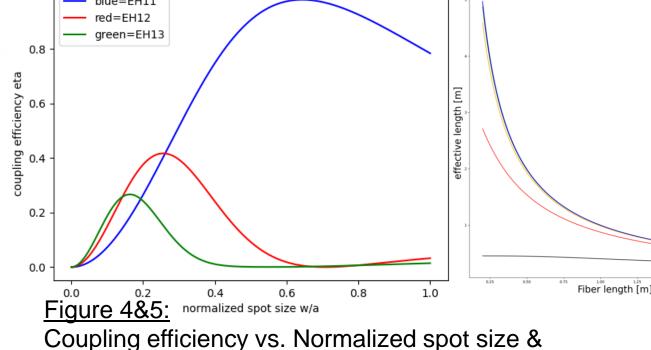
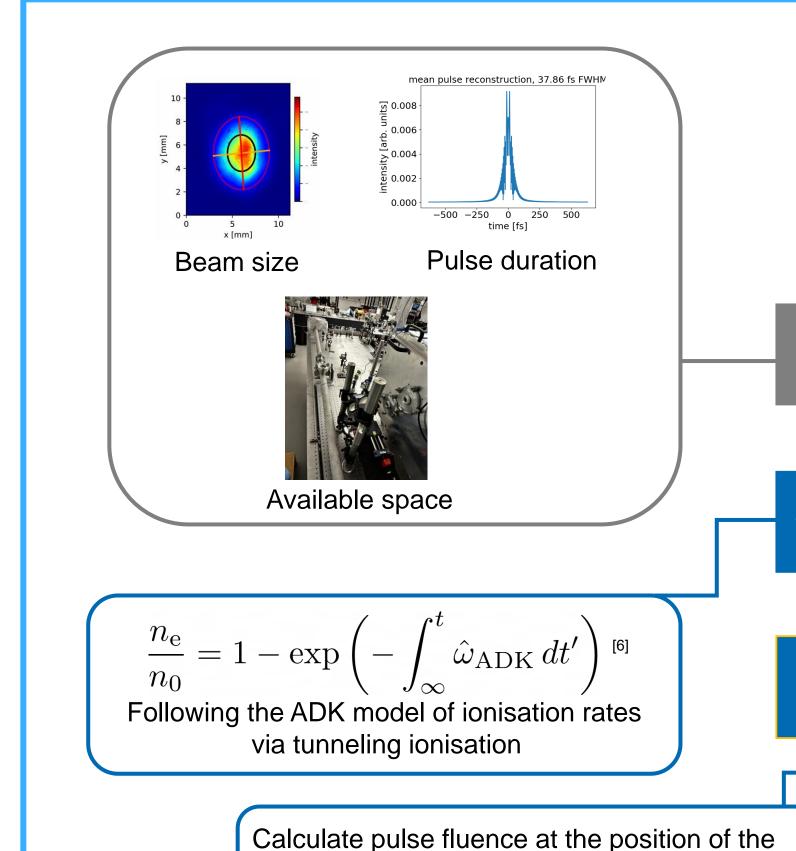


Figure 3: transverse profiles of

the first three EH_{1m} modes.





29.4 54.8 93.5 63.9 Nonlinear refractive indices of noble gases **Define boundary** Calculate the Calculate the conditions and broadening factor Choose nonlinear intensity required to for the desired desired beam broaden the medium spectrum via SPM pulse duration parameters the nonlinear medium is ionised Choose Calculate focal spot Check ionization Calculate focal commercially for selected focus size assuming a available focus length mirror beam energy mirror Optics get damaged or B-Integral is too large Check B-Integral in Determine fiber Check damage radius for most the window, find a thresholds of efficient coupling suitable window sensitive optics into the EH₁₁ mode position

Designing the Beamline

Effective lengths for different radii

 $I_0 = \frac{4\ln 2}{\sqrt{\pi}} \cdot \frac{E}{\tau \cdot \pi w^2} \quad ^{[7]}$

For a Gaussian pulse with a FWHM pulse duration auand a full pulse energy E

The design is an iterative process with three crucial check-points:

- The nonlinear medium is not ionised to prevent ionisation defocusing and energy loss by ionizing
- The pulse fluence is always lower than the optical damage threshold of the used optics such as mirrors, windows and the fiber itself
- The B-Integral is less or equal to 1 to ensure that nonlinear effects such as self focusing and beam breakdown are negligible

Sources

[2] C. Vozzi et. al., Applied Physics B 80, 285-289 (2005).

[3] B. Cros, CERN Accel. school proceedings 2014 [4] M. Nisoli, et. al., IEEE Journal of selected topics in quantum electronics 4,

[5] C. Bree et. al., IEEE Journal of Quantum Electronics 46, 433–437 (2010)

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Experimental arrangement and operation parameters

- 4mJ input pulse energy
- 2w = 8.6mm focused with f=4m to 239µm
- 975µm fiber diameter for 90% coupling efficiency

 $n_2(10^{-20} \text{cm}^2/\text{W})$ | Eq. (11)

- Expected broadening factor F=5.58
- Expected Fourier limited pulse duration $\Delta \tau$ =6.3fs, aimed to achieve with a dedicated chirped mirror compressor

