

EMISSION OF SPIN-2 PARTICLES FROM THE SUN

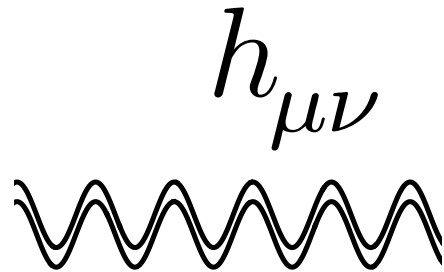
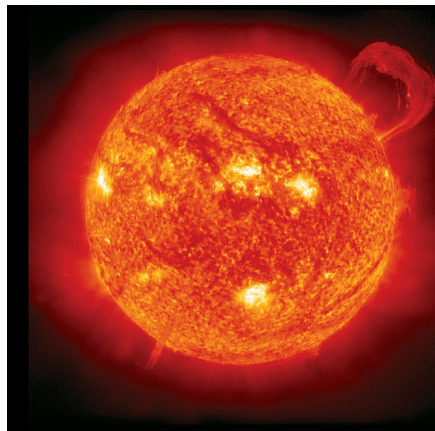
CAMILO GARCIA-CELY.
SEPTEMBER 9, 2025

COSMIC WISPers
Sofia, Bulgaria

In collaboration with Andreas Ringwald

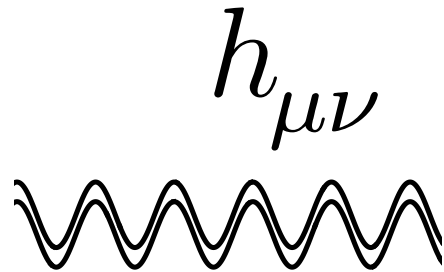
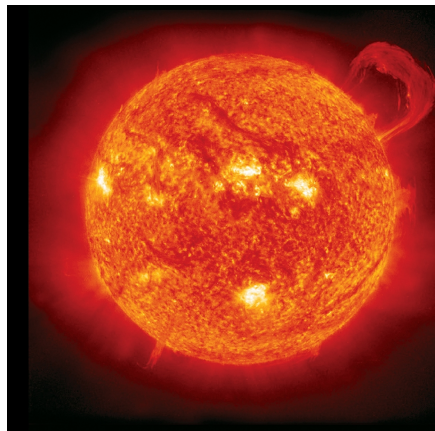


SPIN-2 PARTICLES



$$\mathcal{L} = \frac{1}{2} \partial_\rho h_{\mu\nu} \partial^\rho h^{\mu\nu} - \frac{1}{2} \partial_\rho h \partial^\rho h + \partial_\rho h \partial_\nu h^{\rho\nu} - \partial_\rho h_{\mu\nu} \partial^\nu h^{\mu\rho} \\ - m^2 (h_{\mu\nu} h^{\mu\nu} - h^2)/2 + (8\pi G')^{\frac{1}{2}} h_{\mu\nu} T^{\mu\nu}, \quad \text{Fierz-Pauli, 1939}$$

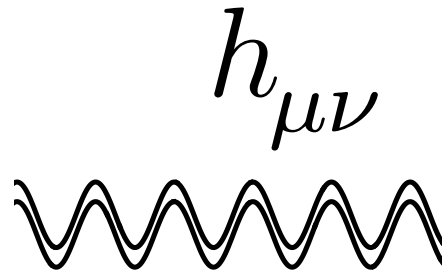
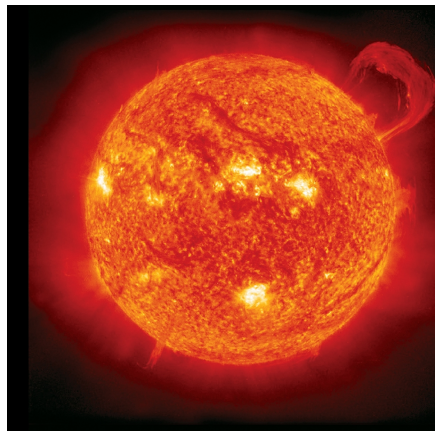
SPIN-2 PARTICLES



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— Mediators of dark Yukawa forces

SPIN-2 PARTICLES

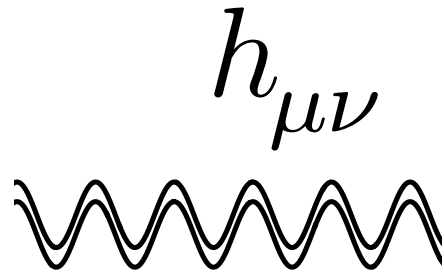
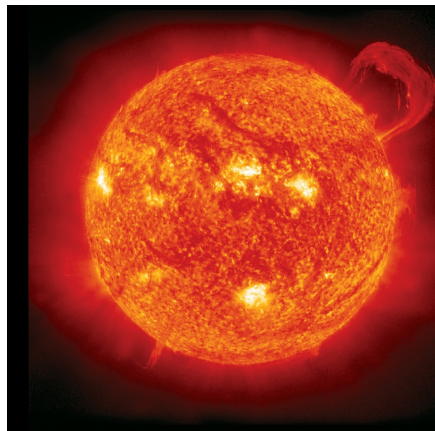


$$\mathcal{L} = \frac{1}{2} \partial_\rho h_{\mu\nu} \partial^\rho h^{\mu\nu} - \frac{1}{2} \partial_\rho h \partial^\rho h + \partial_\rho h \partial_\nu h^{\rho\nu} - \partial_\rho h_{\mu\nu} \partial^\nu h^{\mu\rho} - m^2 (h_{\mu\nu} h^{\mu\nu} - h^2)/2 + (8\pi G')^{\frac{1}{2}} h_{\mu\nu} T^{\mu\nu}, \quad \text{Fierz-Pauli, 1939}$$

— Mediators of dark Yukawa forces

— Spin -2 Dark Matter

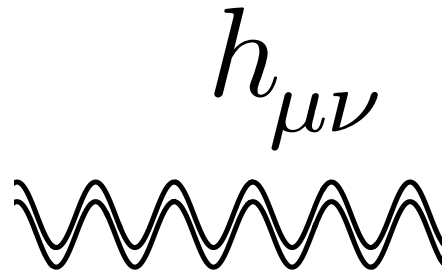
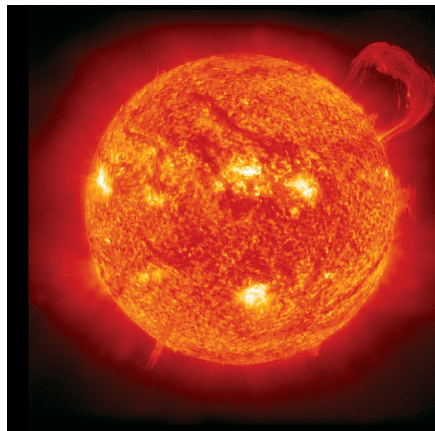
SPIN-2 PARTICLES



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- Mediators of dark Yukawa forces
- Spin -2 Dark Matter
- Bimetric theories

SPIN-2 PARTICLES



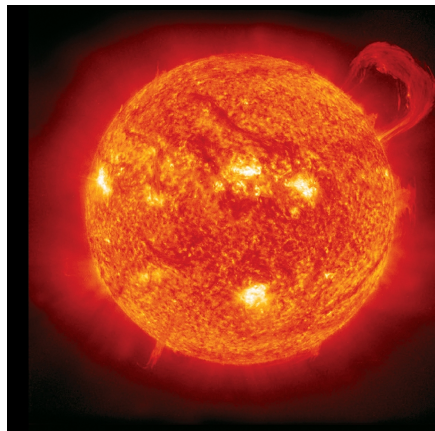
$$\mathcal{L} = \frac{1}{2} \partial_\rho h_{\mu\nu} \partial^\rho h^{\mu\nu} - \frac{1}{2} \partial_\rho h \partial^\rho h + \partial_\rho h \partial_\nu h^{\rho\nu} - \partial_\rho h_{\mu\nu} \partial^\nu h^{\mu\rho} - m^2 (h_{\mu\nu} h^{\mu\nu} - h^2)/2 + (8\pi G')^{\frac{1}{2}} h_{\mu\nu} T^{\mu\nu}, \quad \text{Fierz-Pauli, 1939}$$

- Mediators of dark Yukawa forces
- Spin -2 Dark Matter
- Bimetric theories

Naturally related to solar gravitational waves!

García-Cely, Ringwald, PRL 2025

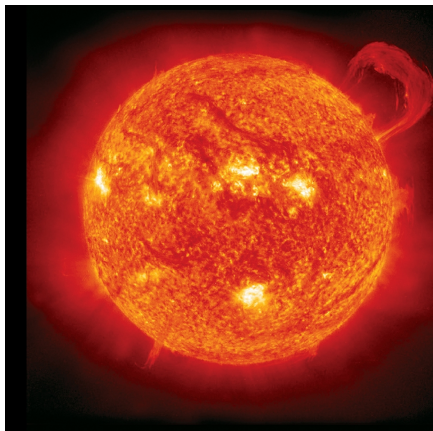
SOLAR EMISSION OF AXIONS



a



SOLAR EMISSION OF AXIONS



a



PHYSICAL REVIEW D

VOLUME 33, NUMBER 4

15 FEBRUARY 1986

Astrophysical axion bounds diminished by screening effects

Georg G. Raffelt

PHYSICAL REVIEW D

VOLUME 37, NUMBER 6

15 MARCH 1988

Plasmon decay into low-mass bosons in stars

Georg G. Raffelt

PHYSICAL REVIEW D **102**, 043019 (2020)

Axion helioscopes as solar magnetometers

Ciaran A. J. O'Hare^{1,*}, Andrea Caputo^{2,†}, Alexander J. Millar^{3,4,‡} and Edoardo Vitagliano^{5,§}

PHYSICAL REVIEW D **101**, 123004 (2020)

Revisiting longitudinal plasmon-axion conversion in external magnetic fields

Andrea Caputo^{1,*}, Alexander J. Millar^{2,3,†} and Edoardo Vitagliano^{4,‡}

PHYSICAL REVIEW D **102**, 123024 (2020)

**Production of axionlike particles from photon conversions
in large-scale solar magnetic fields**

Ersilia Guarini¹, Pierluca Carenza^{1,2}, Javier Galán³, Maurizio Giannotti⁴ and Alessandro Mirizzi^{1,2}

Axion emission by magnetic-field induced conversion of longitudinal plasmons

N. V. Mikheev

Department of Theoretical Physics, Yaroslavl State University, Sovetskaya 14, Yaroslavl 150000, Russia

G. Raffelt

Max-Planck-Institut für Physik (Werner-Heisenberg-Institut), Föhringer Ring 6, 80805 München, Germany

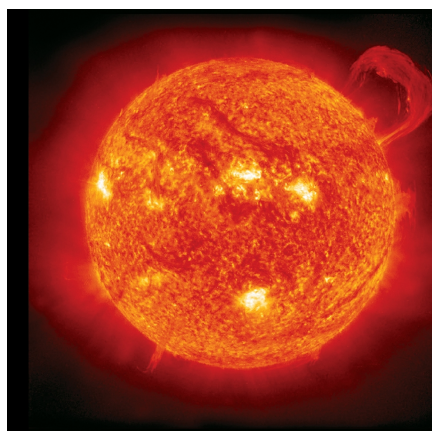
L. A. Vassilevskaya

Solar axion flux from the axion-electron coupling

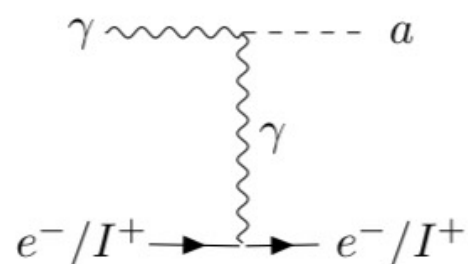
Javier Redondo (Munich U., ASC and Munich, Max Planck Inst.)

Oct 2, 2013

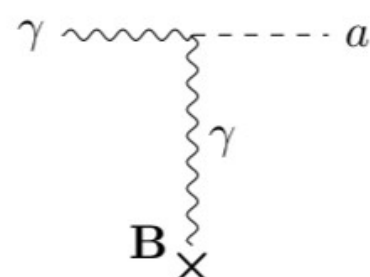
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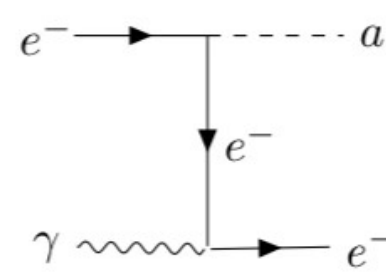
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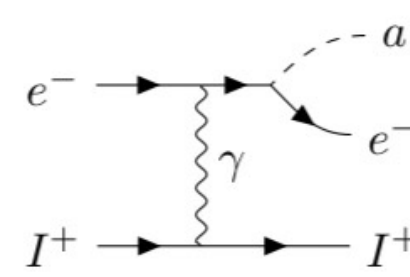
Primakoff (P)



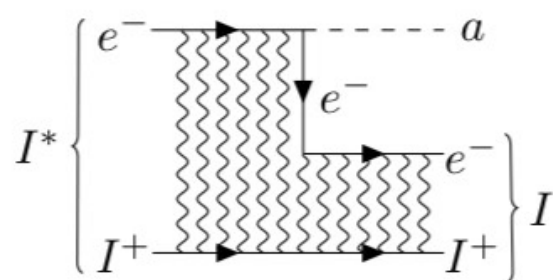
plasmon (LP, TP)



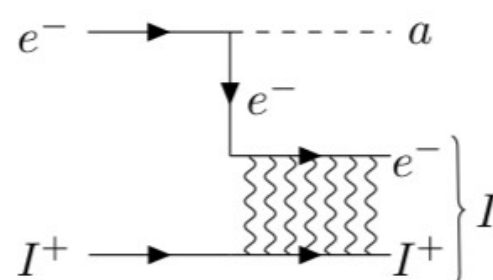
Compton (C)



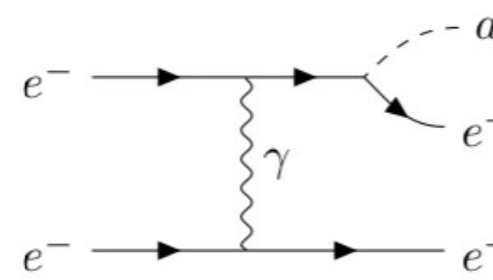
e - I bremsstrahlung (ff)



axio-deëxcitation (bb)

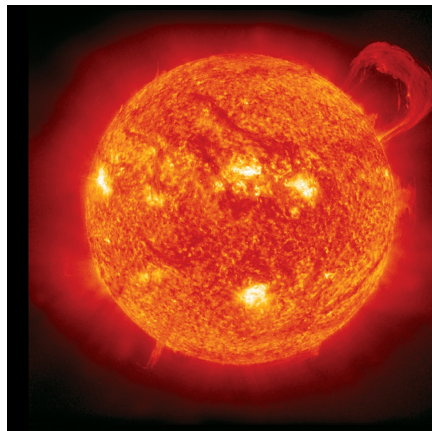


axio-recombination (fb)

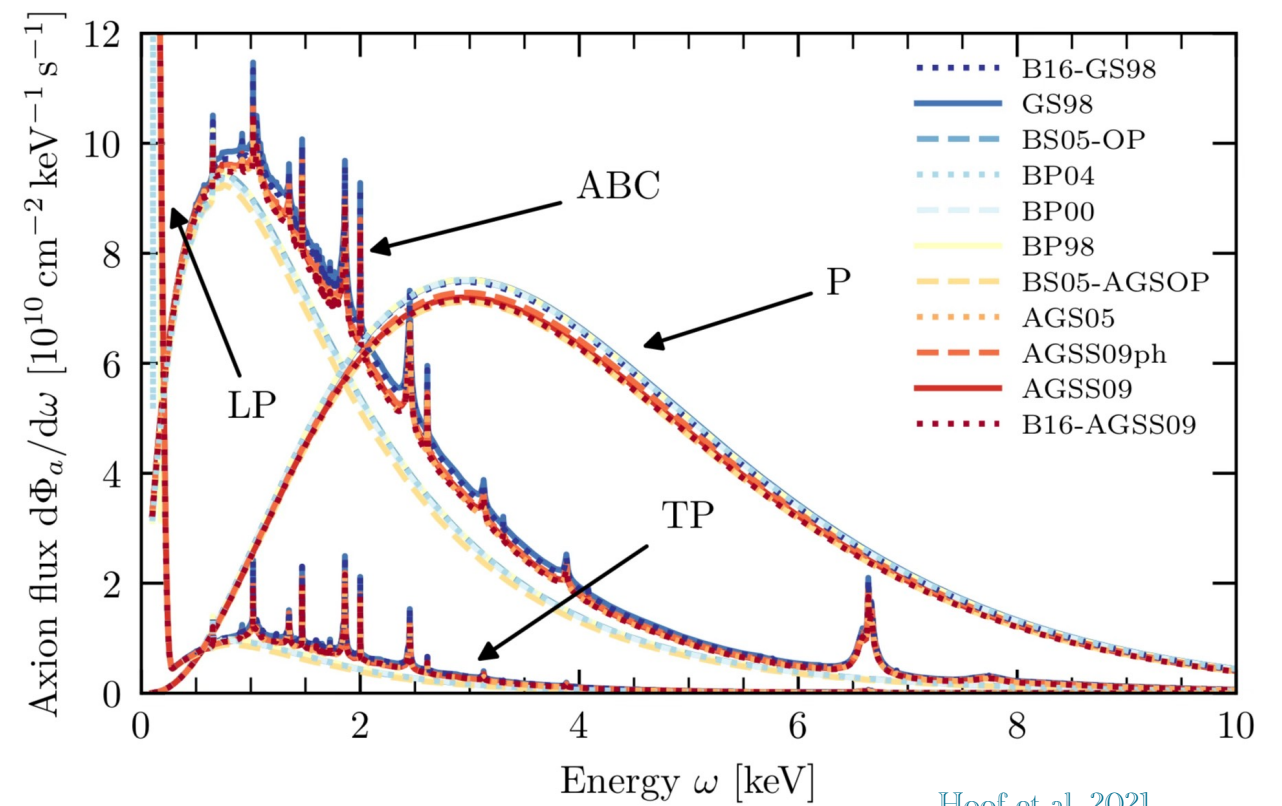


e - e bremsstrahlung (ee)

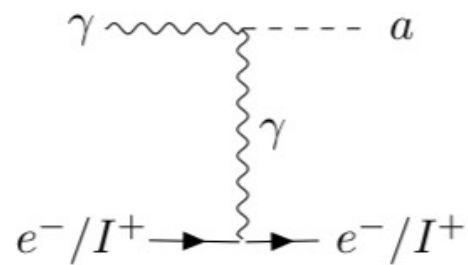
SOLAR EMISSION OF AXIONS



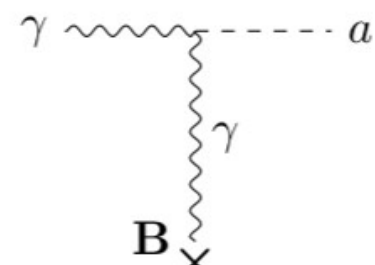
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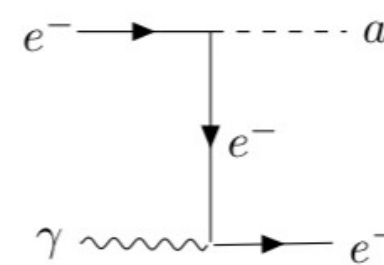
Hoof et al, 2021



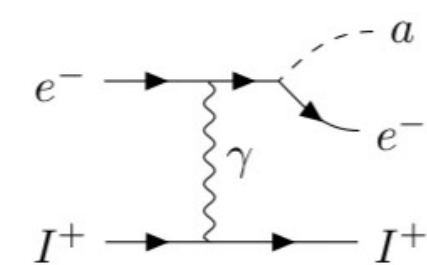
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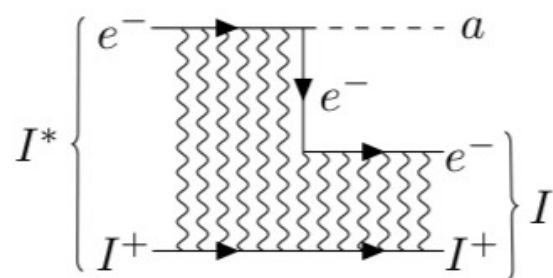
plasmon (LP, TP)



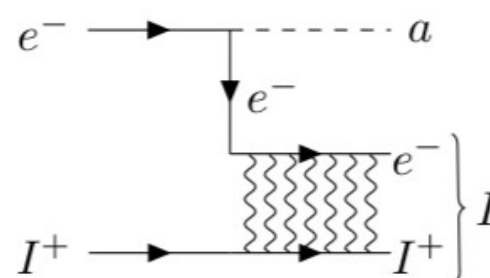
Compton (C)



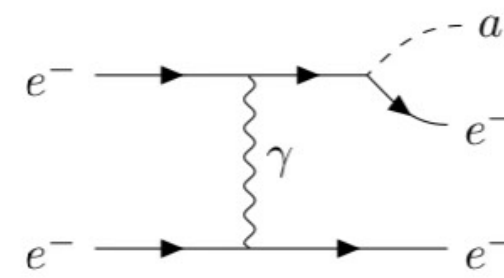
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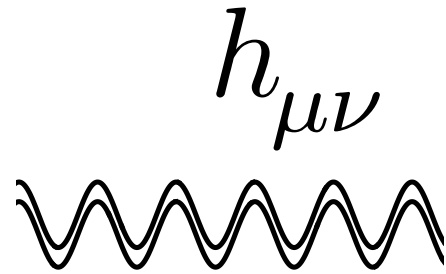
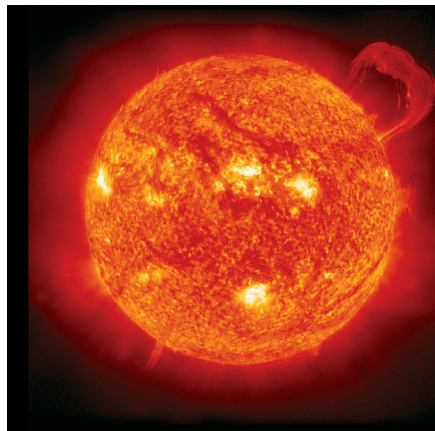


axio-recombination (fb)



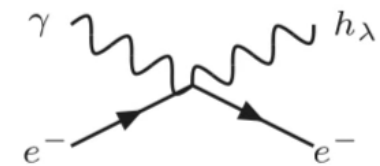
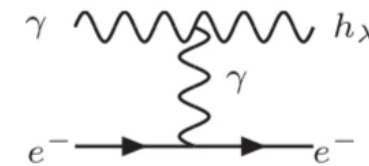
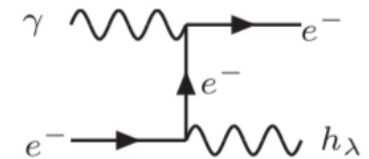
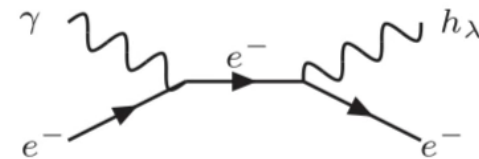
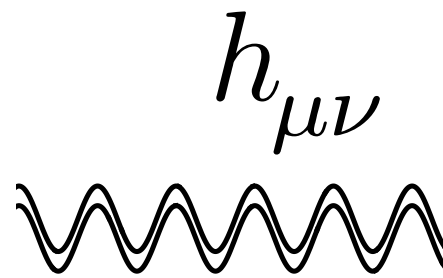
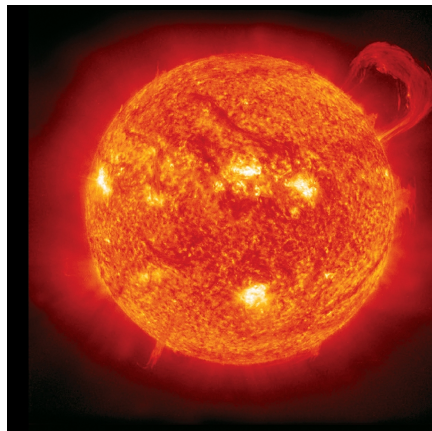
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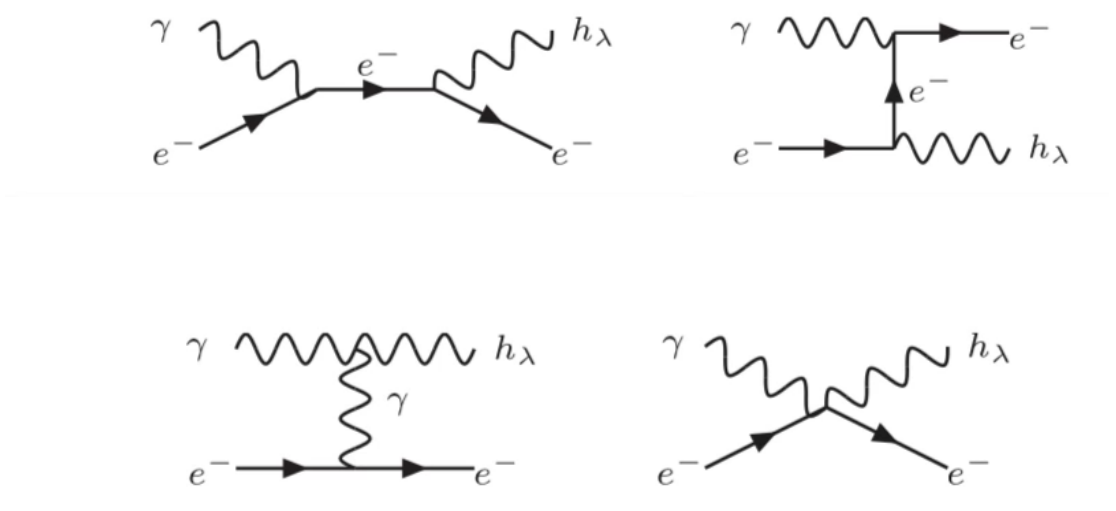
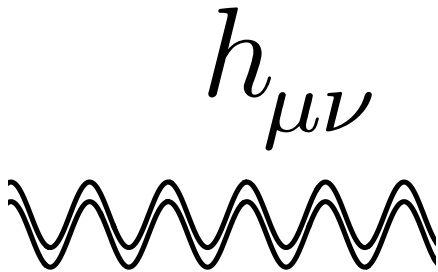
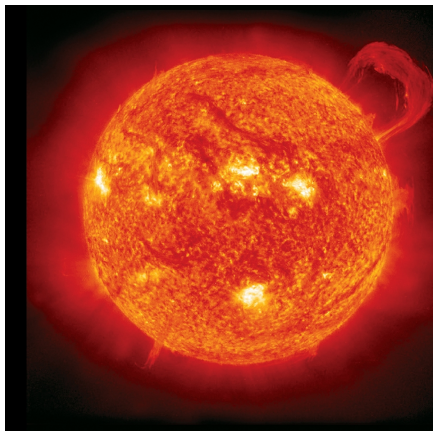
SOLAR EMISSION OF SPIN-2 PARTICLES

Photoproduction



SOLAR EMISSION OF SPIN-2 PARTICLES

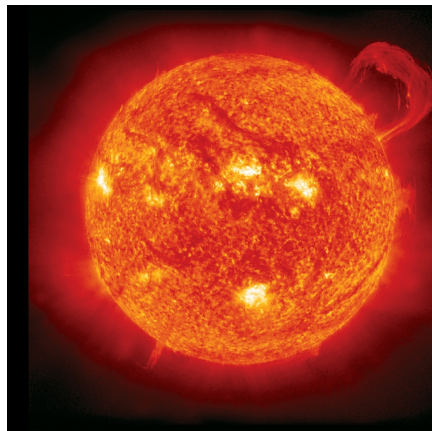
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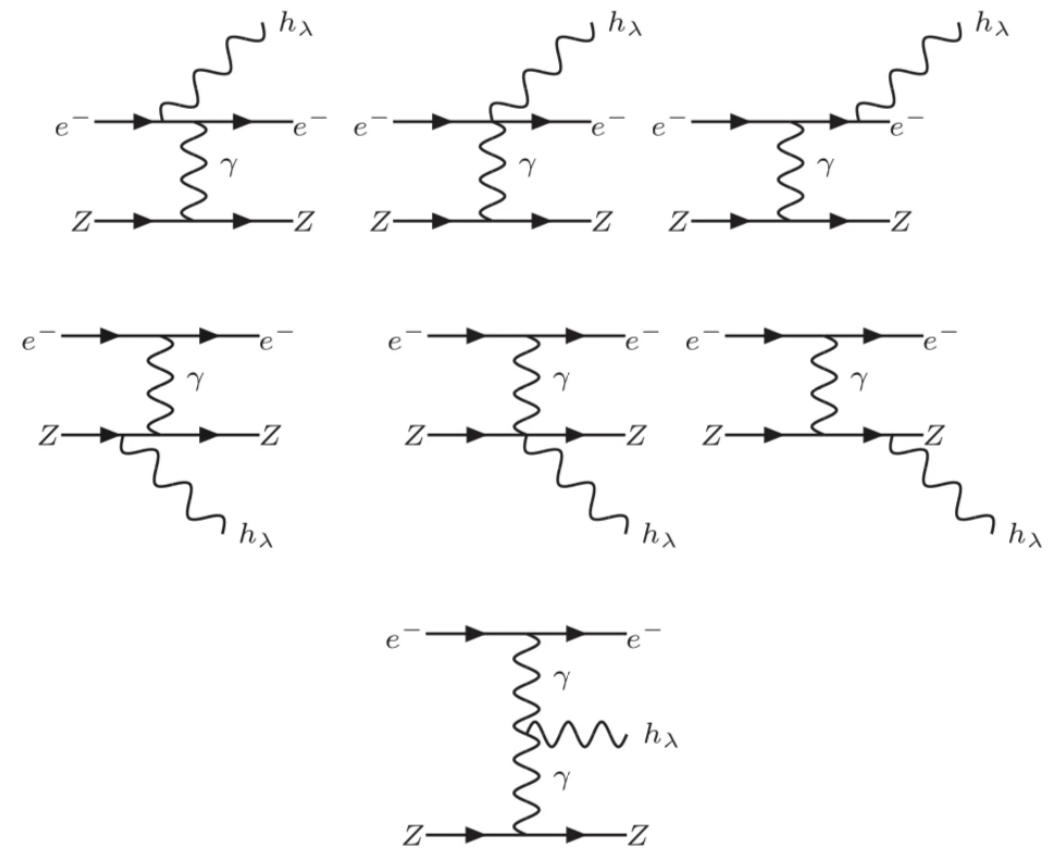
Collision	λ	$\frac{d\Gamma}{d\omega dV}$	
Photo-production $\gamma Z \rightarrow Z h_{\lambda}$	± 2 ± 1 0	$n_{\gamma} n_Z G' Z^2 \alpha \pi \delta(\omega - p_i) \int d\cos\theta \cot^2 \frac{\theta}{2} [1 + \cos^2 \theta] F(\theta)$ 0 $\frac{4}{3} n_{\gamma} n_Z G' Z^2 \alpha \pi \delta(\omega - p_i) \int d\cos\theta \cot^2 \frac{\theta}{2} \sin^4 \frac{\theta}{2} F(\theta)$	$F(\theta) = \frac{(2\omega \sin \frac{\theta}{2})^2}{\kappa^2 + (2\omega \sin \frac{\theta}{2})^2}$

SOLAR EMISSION OF SPIN-2 PARTICLES

Bremsstrahlung

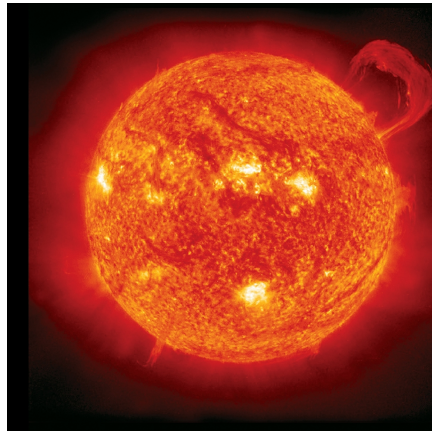


$$h_{\mu\nu}$$

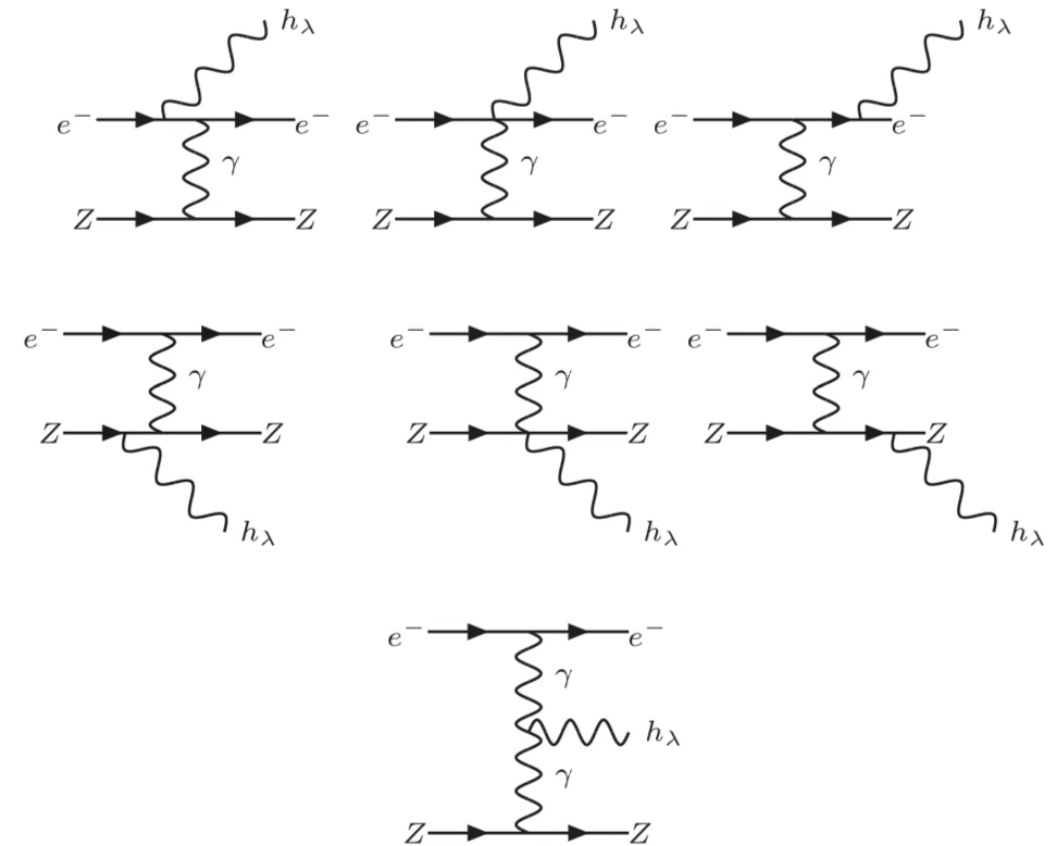


SOLAR EMISSION OF SPIN-2 PARTICLES

Bremsstrahlung

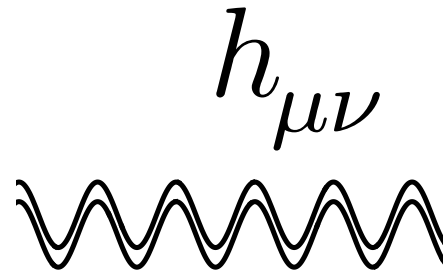
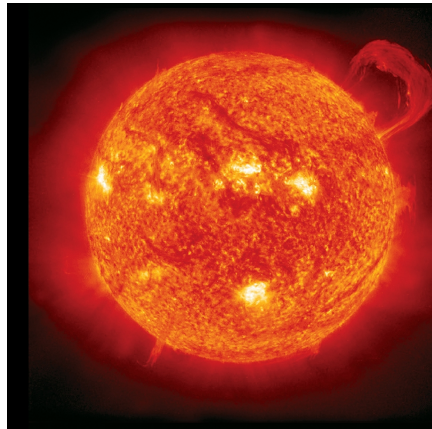


$$h_{\mu\nu}$$



Bremsstrahlung	± 2	$\frac{32n_e n_Z G' Z^2 \alpha^2 p_i}{15\omega} \left(\frac{1}{m_e} + \frac{1}{m_Z} \right) \left(3(1 + \xi^2)L + 10\xi + \mathcal{O}(\xi_s^2) \right)$	$\xi = \frac{p_f}{p_i}, \quad \xi_s = \frac{\kappa}{p_i}$ $\omega = E_i(1 - \xi^2)$ $L = \log \sqrt{\frac{(1+\xi)^2 + \xi_s^2}{(1-\xi)^2 + \xi_s^2}}$
$eZ \rightarrow eZ h_\lambda$	± 1	0	
	0	$\frac{16n_e n_Z G' Z^2 \alpha^2 p_i}{45\omega} \left(\frac{1}{m_e} + \frac{1}{m_Z} \right) \left((1 + \xi^2)L + 30\xi + \mathcal{O}(\xi_s^2) \right)$	
Bremsstrahlung	± 2	$\frac{16n_e^2 G' \alpha^2 p_i}{15\omega m_e} \left(\left(6(1 + \xi^2) - \frac{3(1-\xi^2)^4 + 7(1-\xi^4)^2}{2(1+\xi^2)^3} \right) L + 20\xi - \frac{6\xi(1+\xi^4)}{(1+\xi^2)^2} + \mathcal{O}(\xi_s^2) \right)$	
$ee \rightarrow ee h_\lambda$	± 1	0	
	0	$\frac{16n_e^2 G' \alpha^2 p_i}{15\omega m_e} \left(\left(\frac{1}{3}(1 + \xi^2) - \frac{(1-\xi^2)^4 + 29(1-\xi^4)^2}{12(1+\xi^2)^3} \right) L + \frac{29\xi}{3} + \frac{2\xi^3}{3(1+\xi^2)^2} + \mathcal{O}(\xi_s^2) \right) <$	

SOLAR EMISSION OF SPIN-2 PARTICLES

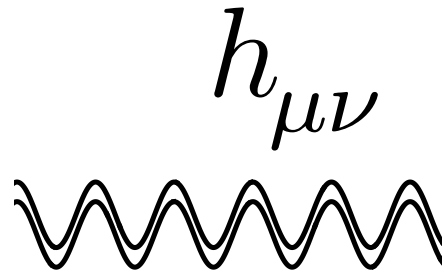
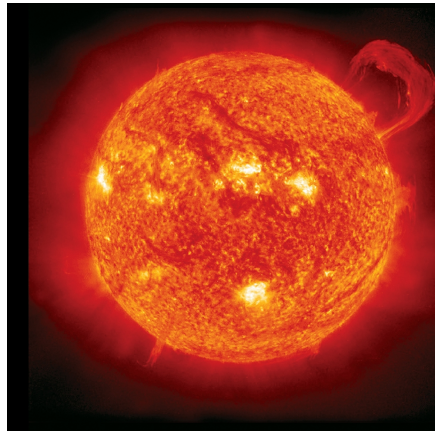


We thus verify the vDVZ discontinuity

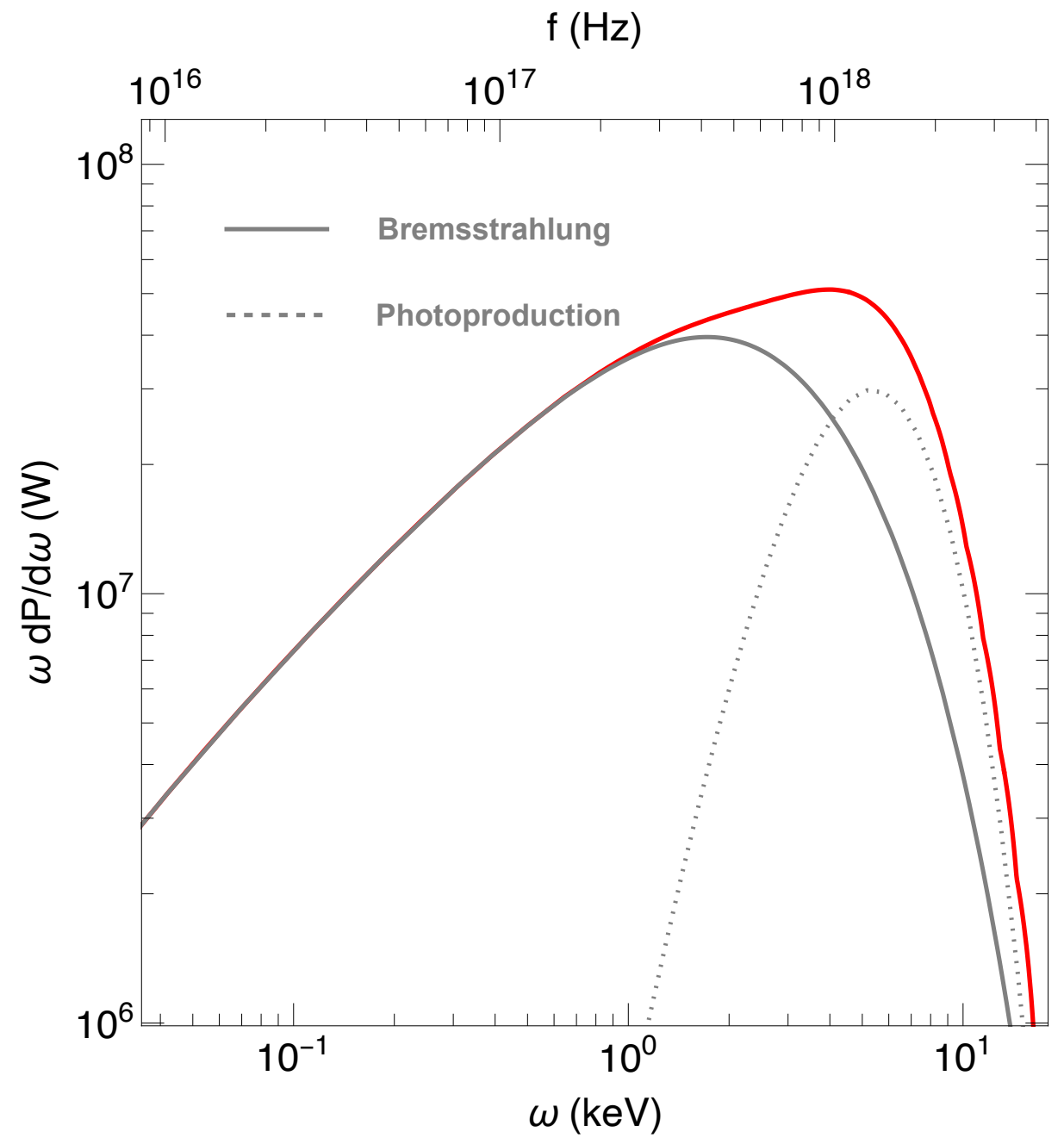
van Dam, Veltman, 1970; Zahkarov, 1970

Collision	λ	$\frac{d\Gamma}{d\omega dV}$	
Photo-production $\gamma Z \rightarrow Z h_\lambda$	± 2 ± 1 0	$\frac{n_\gamma n_Z G' Z^2 \alpha \pi \delta(\omega - p_i) \int d\cos\theta \cot^2 \frac{\theta}{2} [1 + \cos^2 \theta] F(\theta)}{0}$ $\frac{\frac{4}{3} n_\gamma n_Z G' Z^2 \alpha \pi \delta(\omega - p_i) \int d\cos\theta \cot^2 \frac{\theta}{2} \sin^4 \frac{\theta}{2} F(\theta)}{0}$	$F(\theta) = \frac{(2\omega \sin \frac{\theta}{2})^2}{\kappa^2 + (2\omega \sin \frac{\theta}{2})^2}$
Bremsstrahlung $eZ \rightarrow eZ h_\lambda$	± 2 ± 1 0	$\frac{32 n_e n_Z G' Z^2 \alpha^2 p_i}{15\omega} \left(\frac{1}{m_e} + \frac{1}{m_Z} \right) \left(3(1 + \xi^2) L + 10\xi + \mathcal{O}(\xi_s^2) \right)$ $\frac{16 n_e n_Z G' Z^2 \alpha^2 p_i}{45\omega} \left(\frac{1}{m_e} + \frac{1}{m_Z} \right) \left((1 + \xi^2) L + 30\xi + \mathcal{O}(\xi_s^2) \right)$	$\xi = \frac{p_f}{p_i}, \quad \xi_s = \frac{\kappa}{p_i}$ $\omega = E_i(1 - \xi^2)$
Bremsstrahlung $ee \rightarrow ee h_\lambda$	± 2 ± 1 0	$\frac{16 n_e^2 G' \alpha^2 p_i}{15\omega m_e} \left(\left(6(1 + \xi^2) - \frac{3(1 - \xi^2)^4 + 7(1 - \xi^4)^2}{2(1 + \xi^2)^3} \right) L + 20\xi - \frac{6\xi(1 + \xi^4)}{(1 + \xi^2)^2} + \mathcal{O}(\xi_s^2) \right)$ $\frac{16 n_e^2 G' \alpha^2 p_i}{15\omega m_e} \left(\left(\frac{1}{3}(1 + \xi^2) - \frac{(1 - \xi^2)^4 + 29(1 - \xi^4)^2}{12(1 + \xi^2)^3} \right) L + \frac{29\xi}{3} + \frac{2\xi^3}{3(1 + \xi^2)^2} + \mathcal{O}(\xi_s^2) \right) <$	$L = \log \sqrt{\frac{(1 + \xi)^2 + \xi_s^2}{(1 - \xi)^2 + \xi_s^2}}$

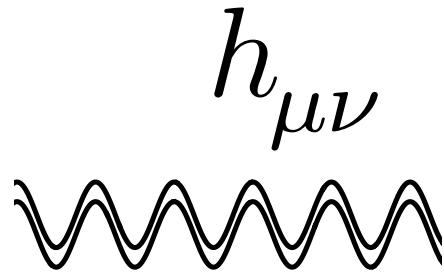
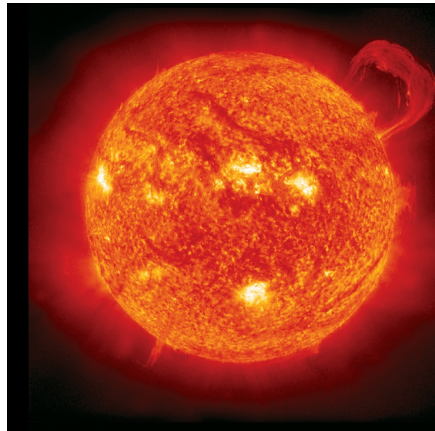
CASE $G=G'$



CGC, Ringwald, 2025 **PRELIMINARY**

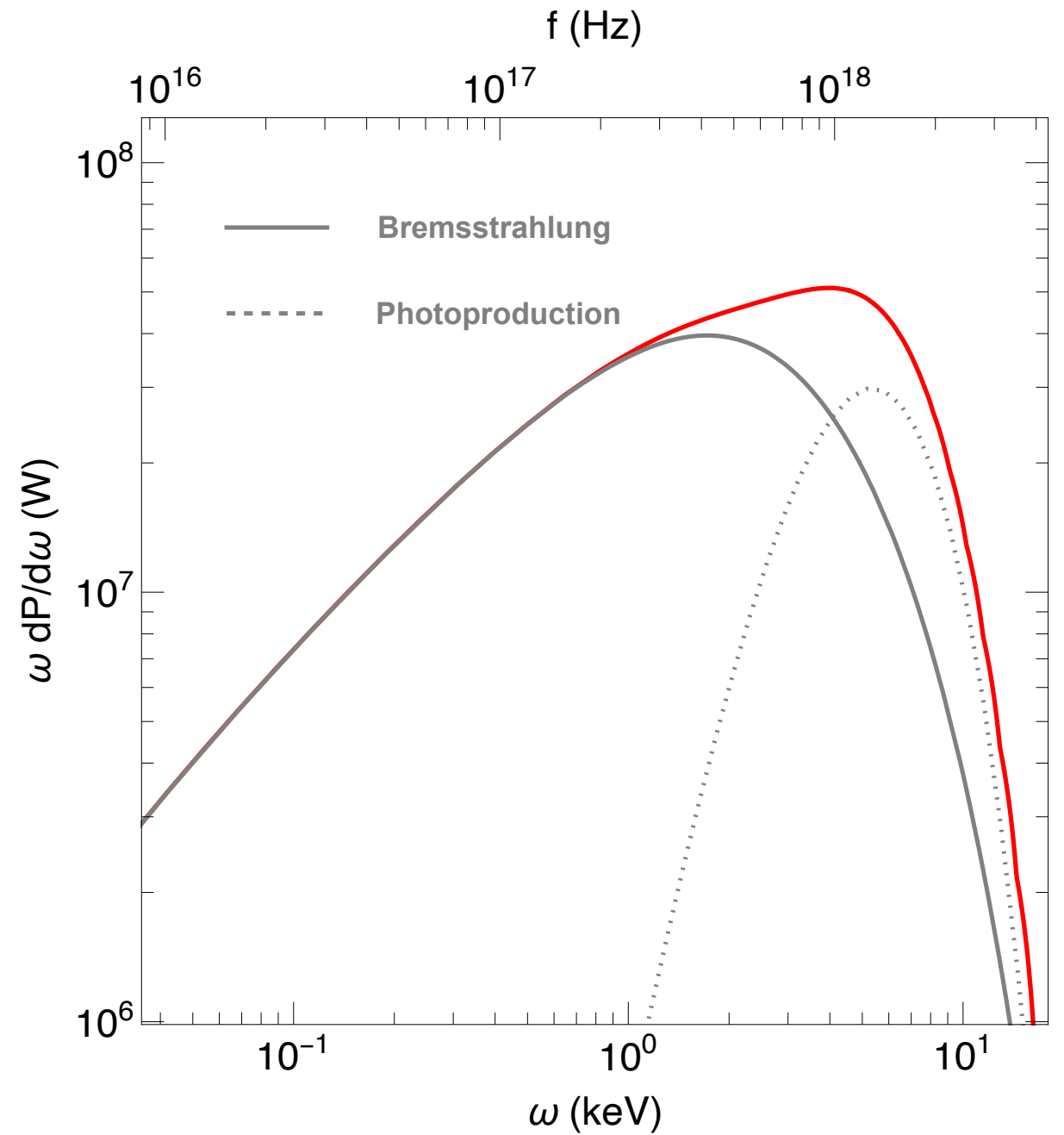


CASE $G=G'$



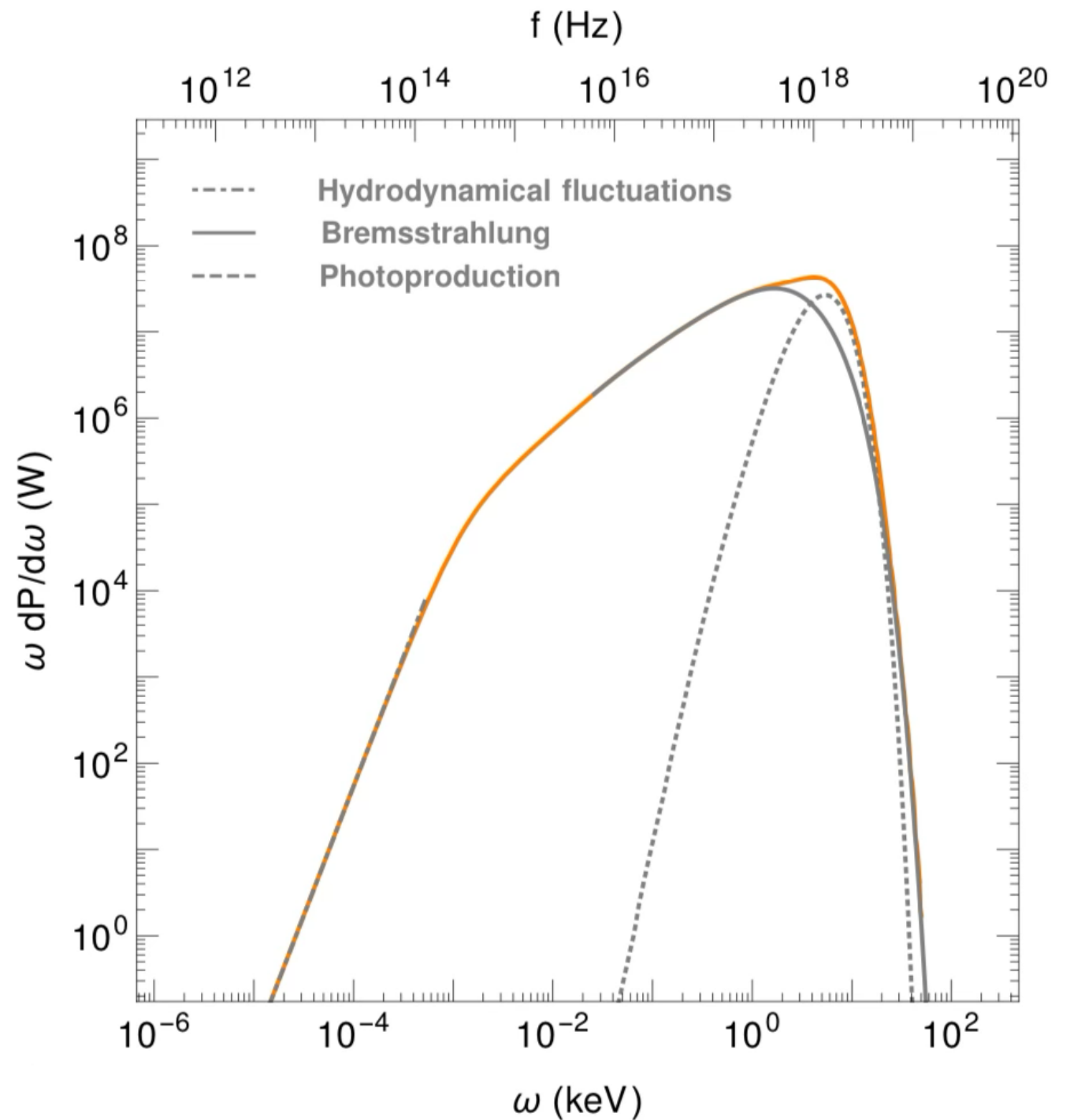
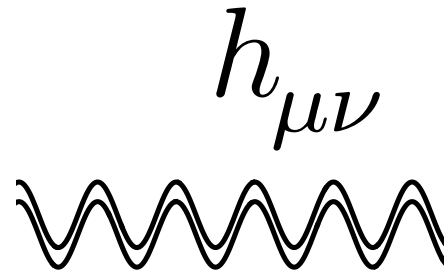
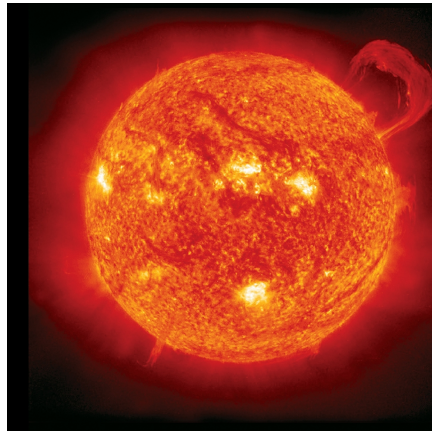
$$P_{\odot} = (1.5 \times 10^{15} \text{ erg/s}) \frac{G'}{G}$$

CGC, Ringwald, 2025 **PRELIMINARY**



CASE $G=G'$

Naturally related to solar gravitational waves!



$$P_{\odot} = (1.5 \times 10^{15} \text{ erg/s}) \frac{G'}{G}$$

CGC, Ringwald, 2025 **PRELIMINARY**

Infrared Photons and Gravitons*

STEVEN WEINBERG†

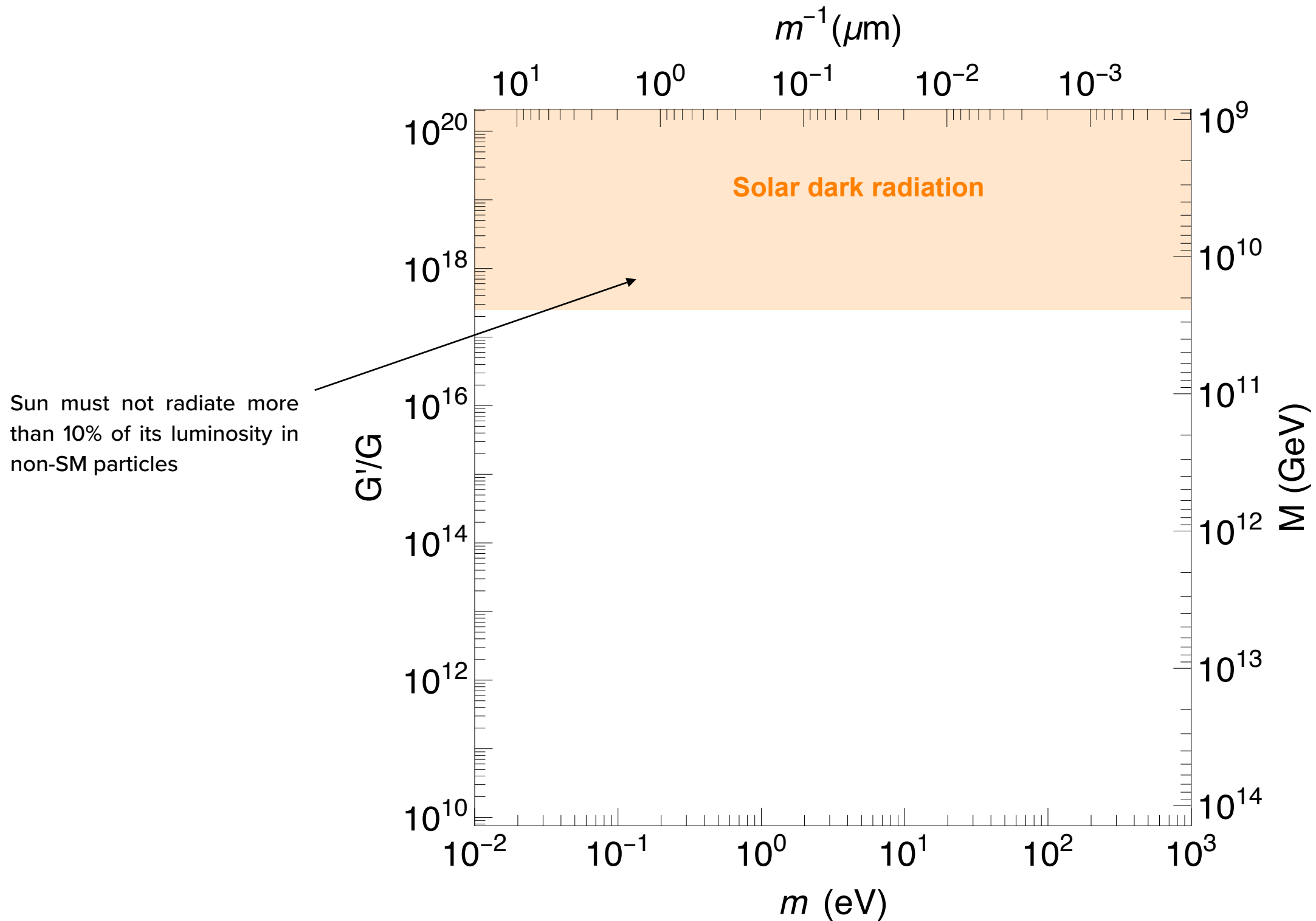
Department of Physics, University of California, Berkeley, California

(Received 1 June 1965)

The solar gravitational radiation power is then

$$P_{\odot} \simeq 6 \times 10^{14} \text{ erg/sec.} \quad (4.24)$$

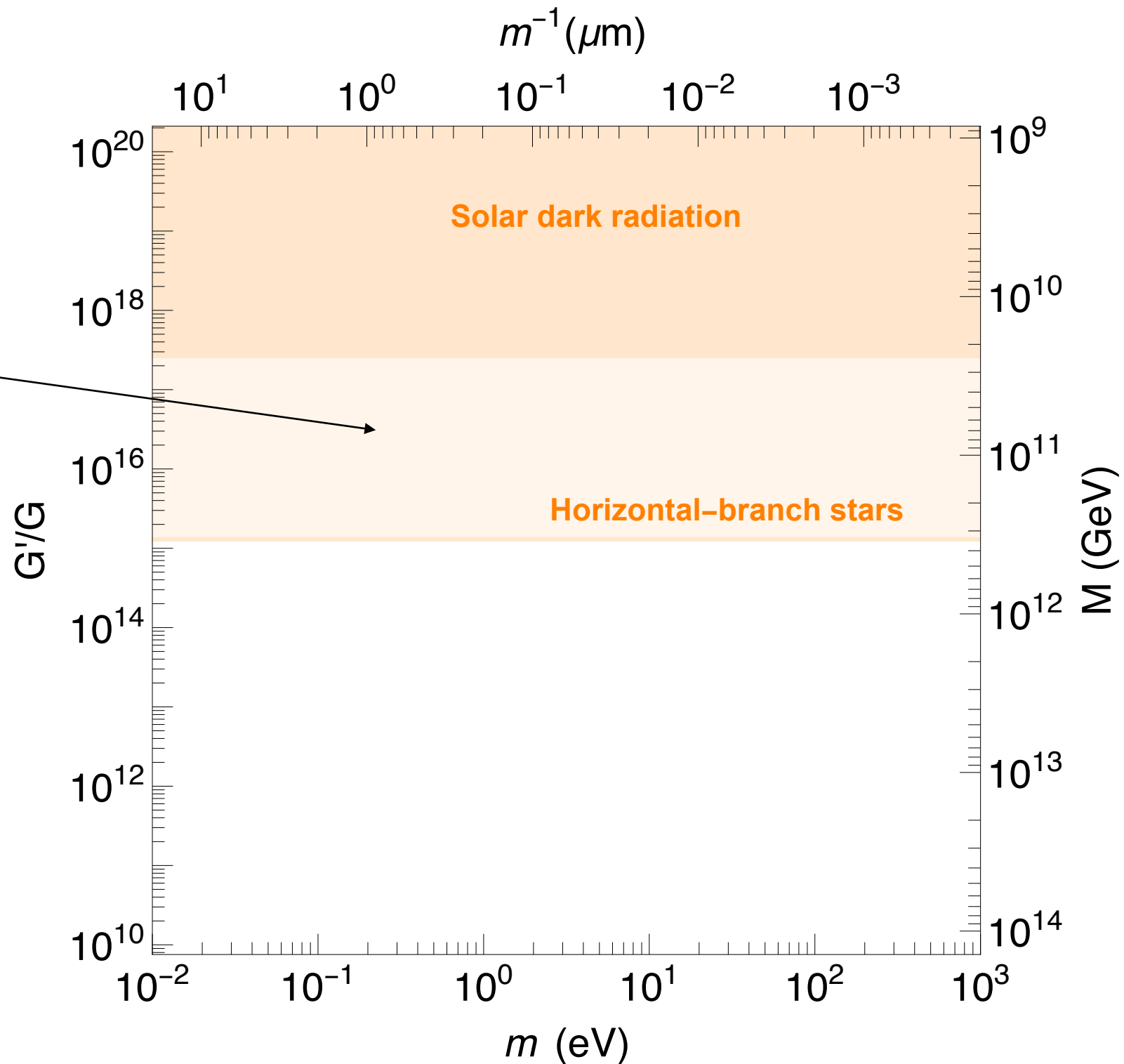
PARAMETER SPACE OF SPIN-2 PARTICLES



PARAMETER SPACE OF SPIN-2 PARTICLES

In horizontal branch stars, an extra energy loss will make the star burn Helium faster. To keep the helium-burning lifetime from shrinking by more than about ten percent.

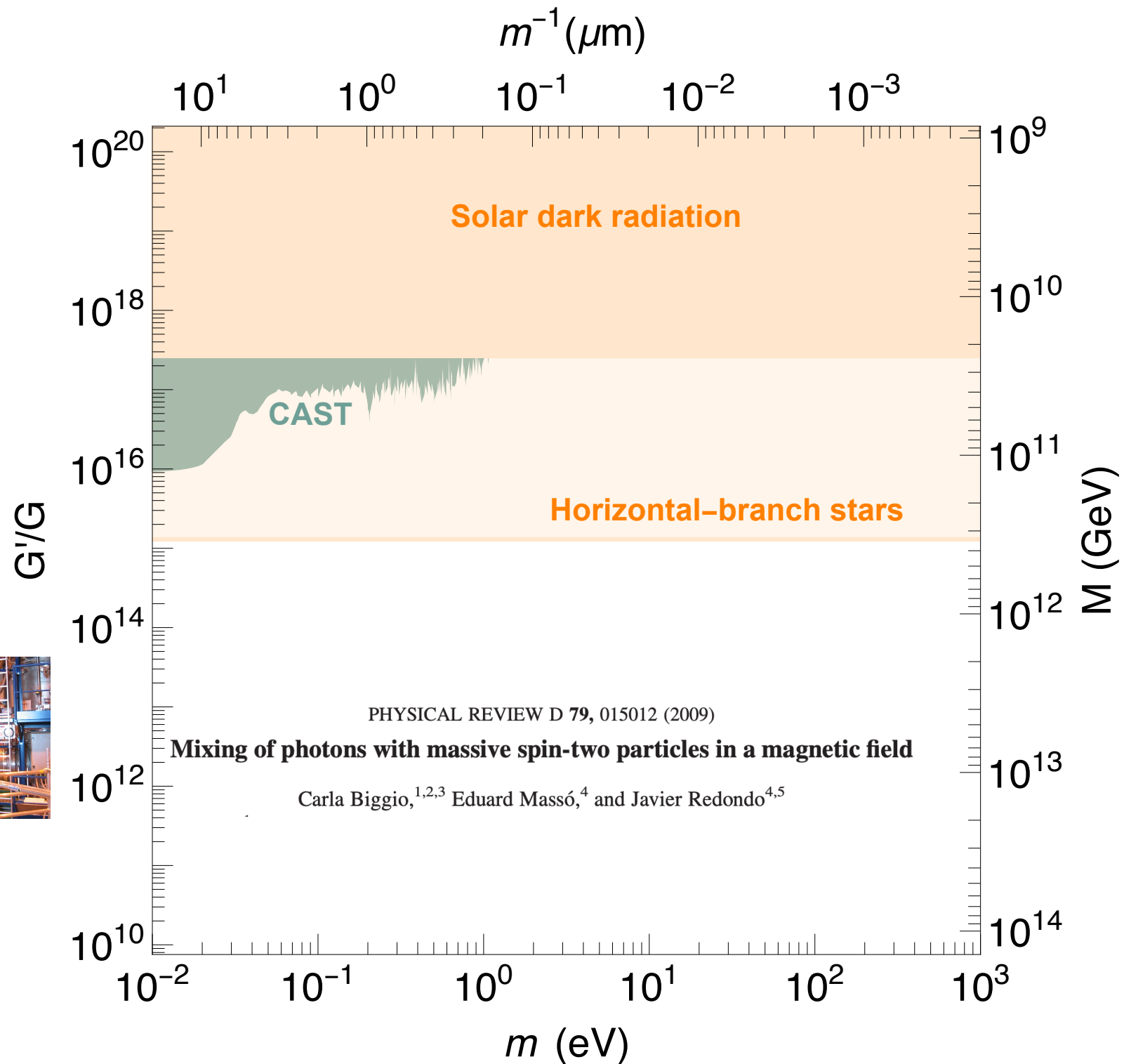
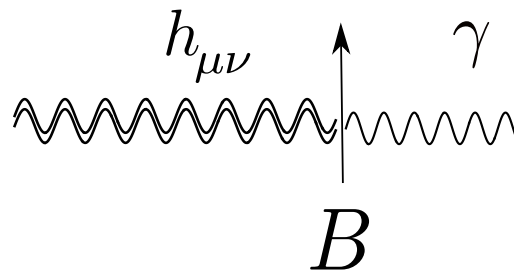
Raffelt, 1995



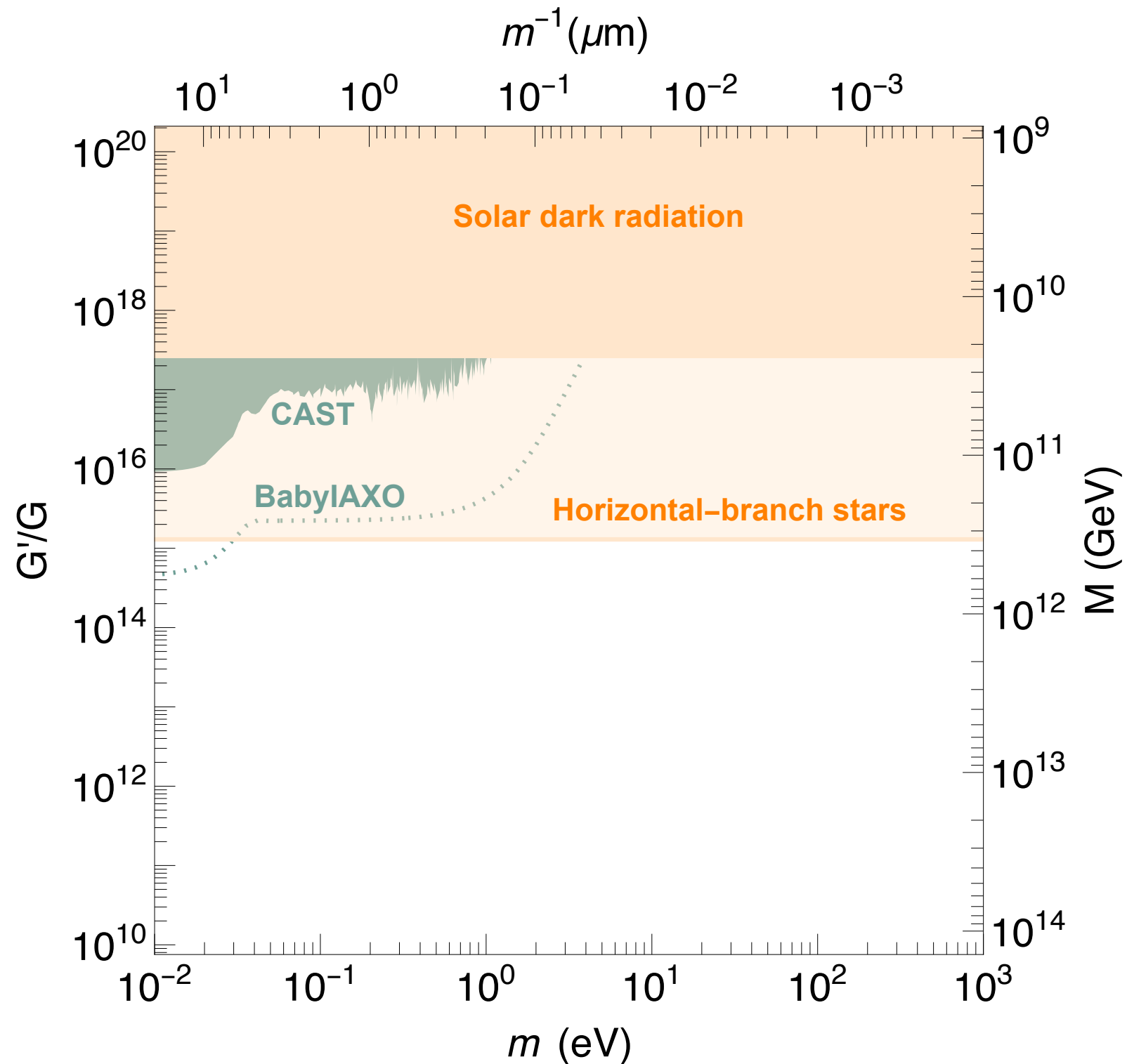
CGC, Ringwald, 2025 **PRELIMINARY**

Camilo García-Cely

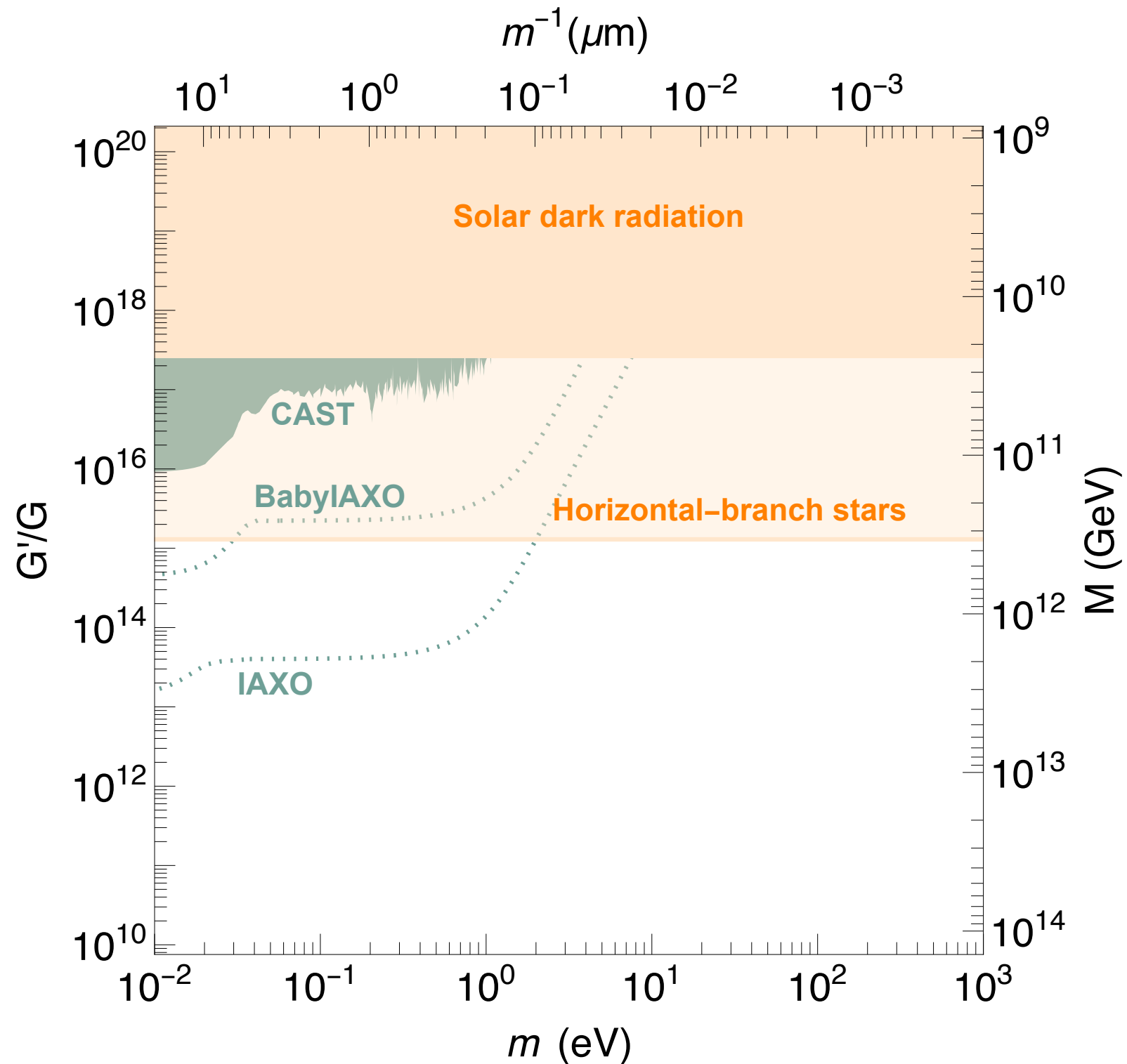
PARAMETER SPACE OF SPIN-2 PARTICLES



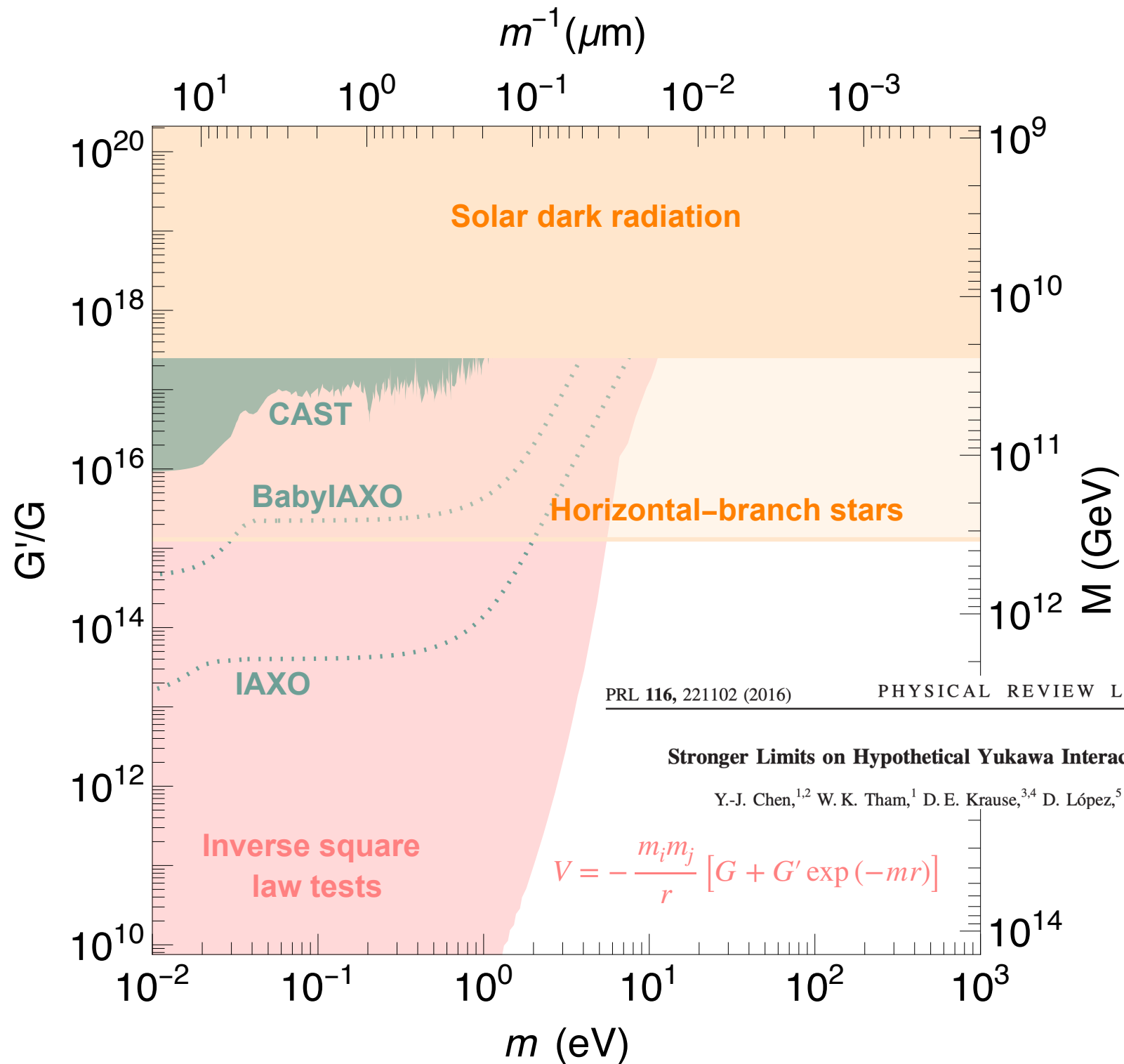
PARAMETER SPACE OF SPIN-2 PARTICLES



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PARAMETER SPACE OF SPIN-2 PARTICLES

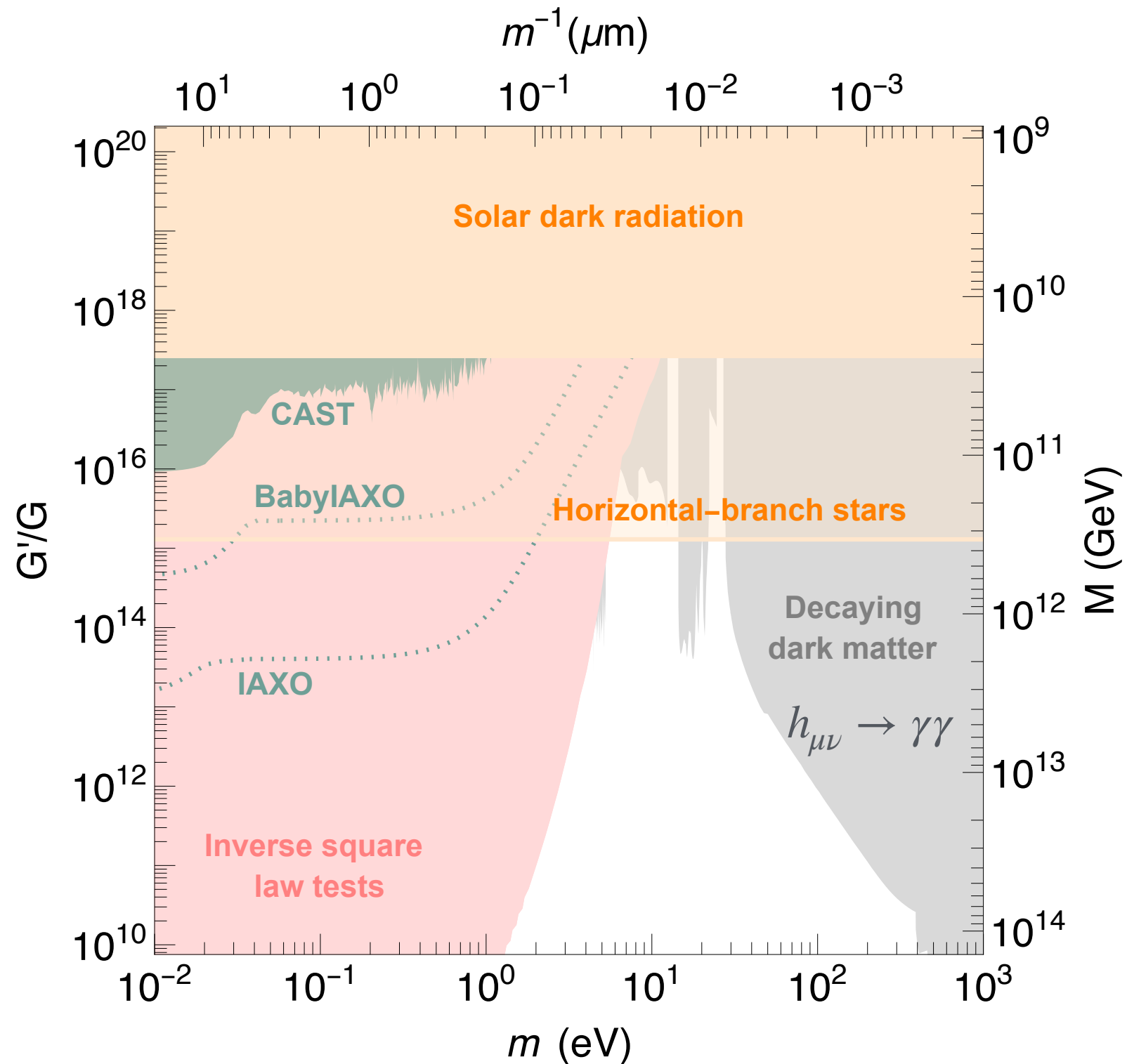


CGC, Ringwald, 2025 **PRELIMINARY**

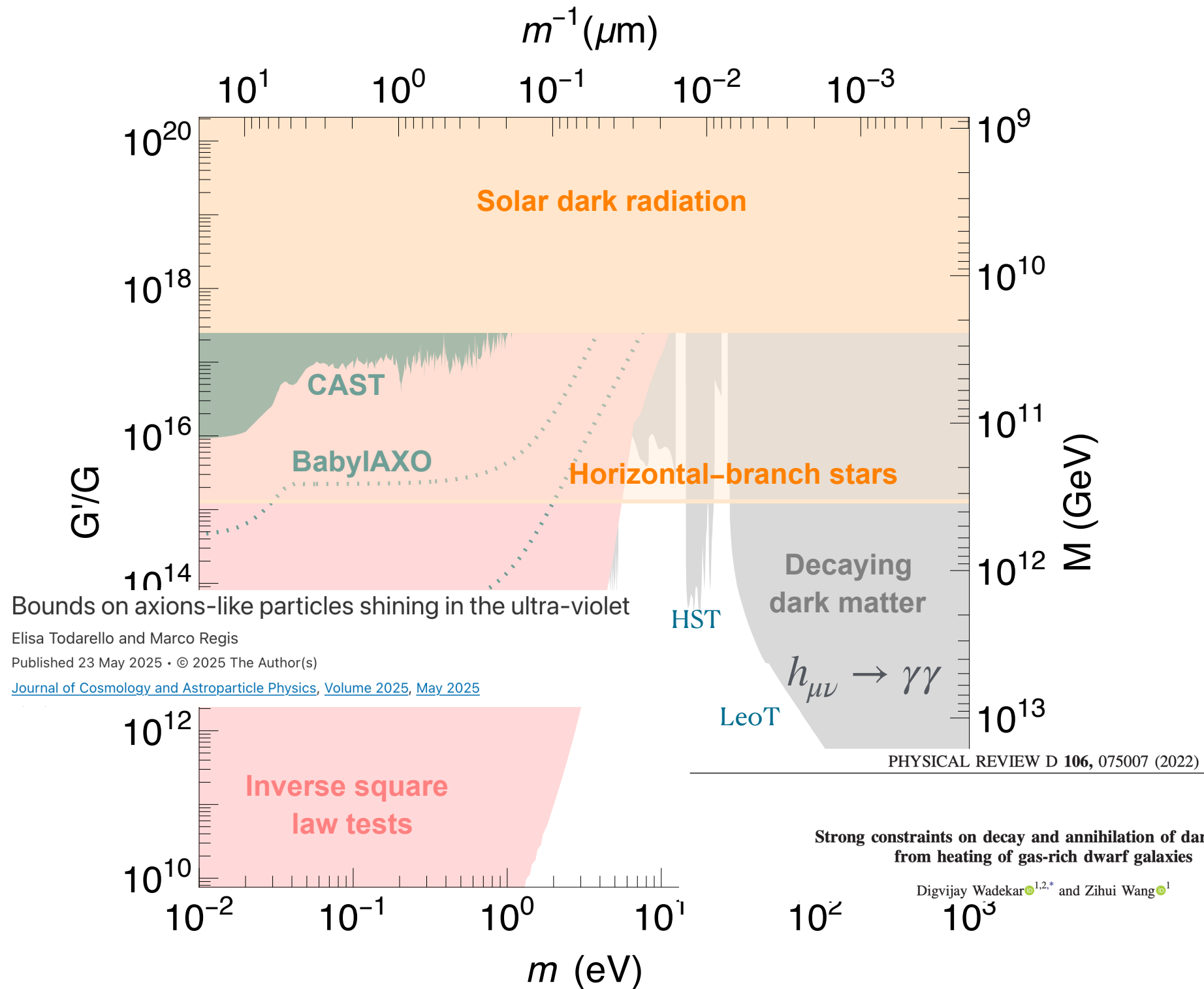
Camilo García-Cely



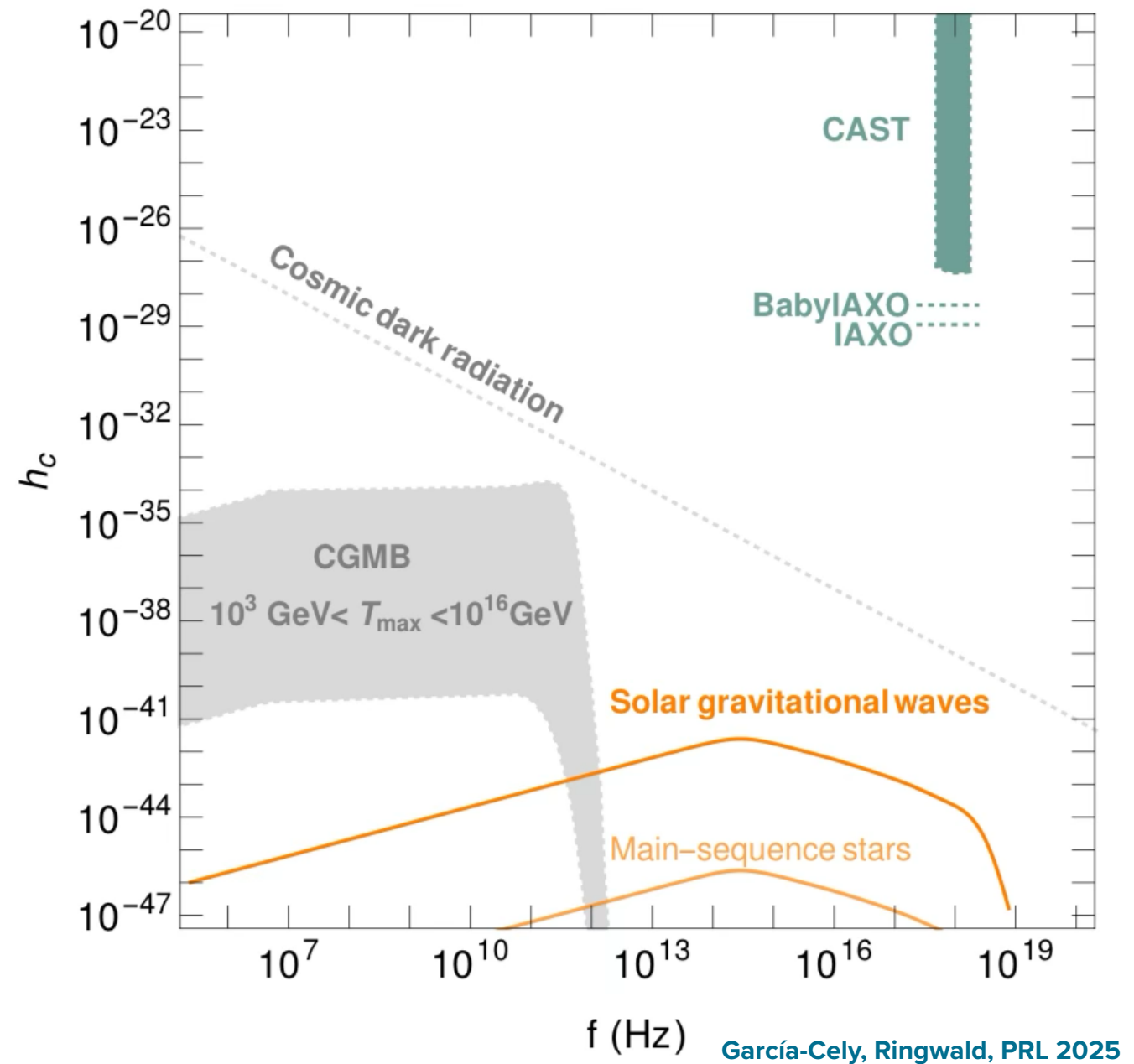
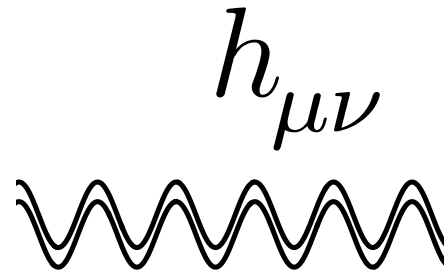
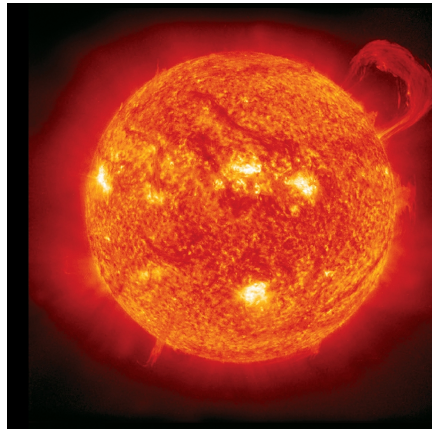
PARAMETER SPACE OF SPIN-2 PARTICLES



PARAMETER SPACE OF SPIN-2 PARTICLES



SOLAR GRAVITATIONAL WAVES



CONCLUSIONS

- Spin-2 particles can be emitted efficiently from the Sun through bremsstrahlung and photoproduction, providing a novel channel of “dark radiation.”
- Stellar evolution, especially the helium-burning phase of horizontal branch stars, sets strong limits on exotic energy losses, constraining the allowed parameter space of spin-2 particles.
- Solar luminosity arguments, combined with astrophysical and laboratory bounds, already probe large regions of parameter space of spin-2 particles.
- These results link solar physics, stellar evolution, and laboratory searches, showing that the Sun can serve as a laboratory for testing theories with new spin-2 states.
- There is a natural synergy with gravitational-wave physics.