

**3rd General Meeting of COST Action COSMIC
WISPerS (CA21106), Sofia, Bulgaria
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Dark Matter Freeze-in and Small-Scale Observables: Bounds and Viable Models

Alessandro Lenoci

The Hebrew University of Jerusalem

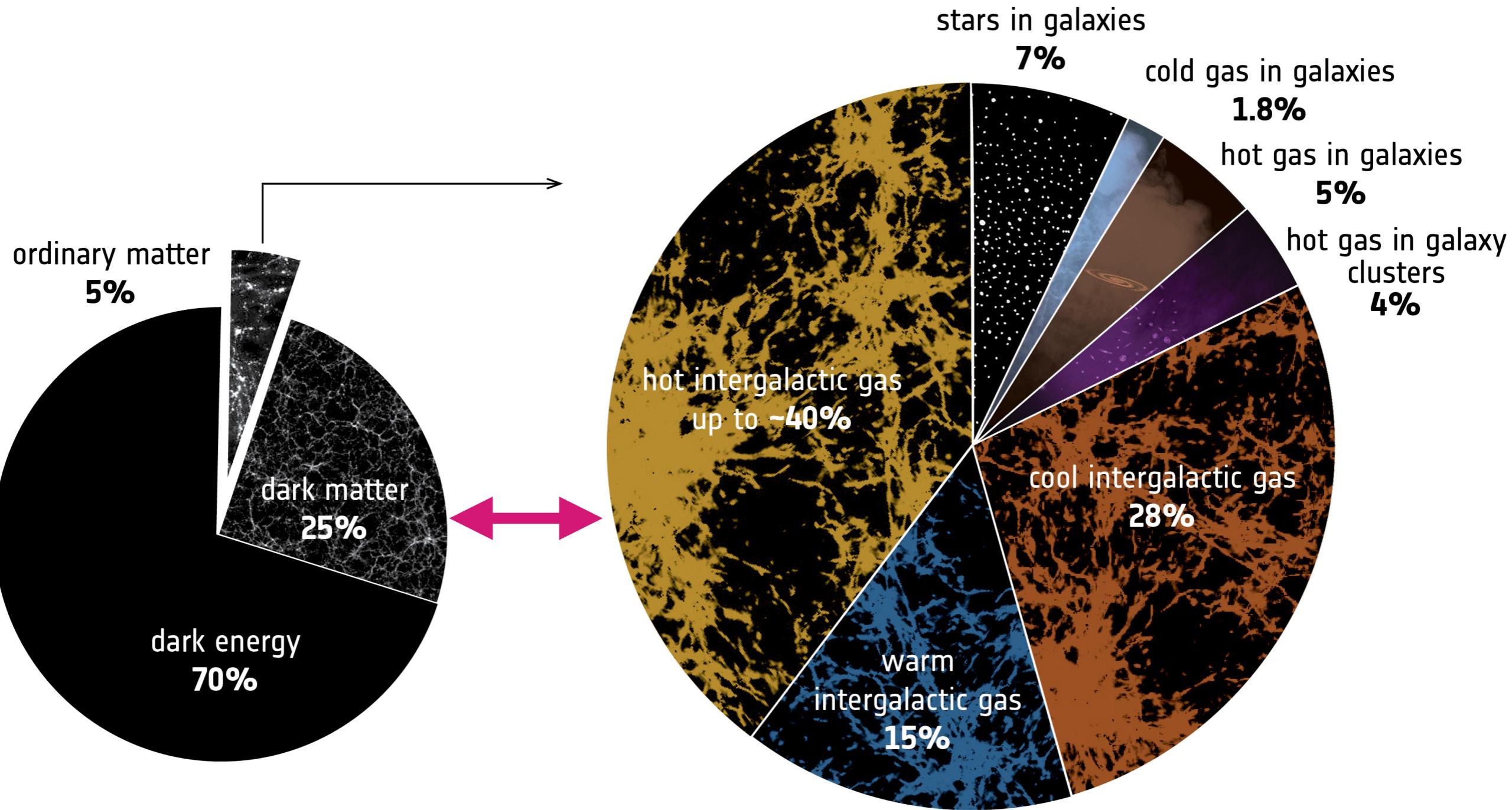
Based on

F. D'Eramo, AL, A. Dekker (2506.13864)



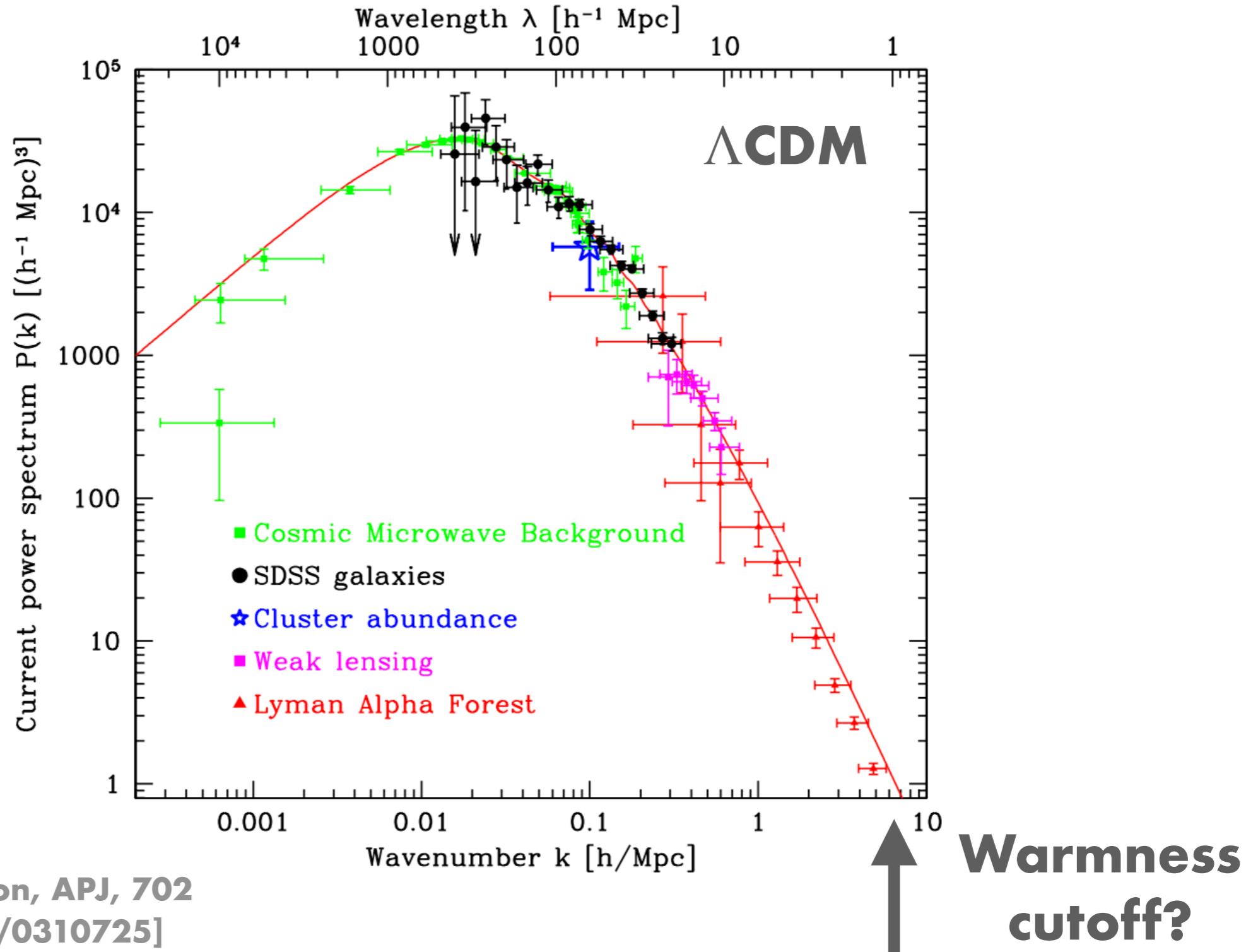
האוניברסיטה
העברית
בירושלים
THE HEBREW
UNIVERSITY
OF JERUSALEM

Dark and Visible Matter



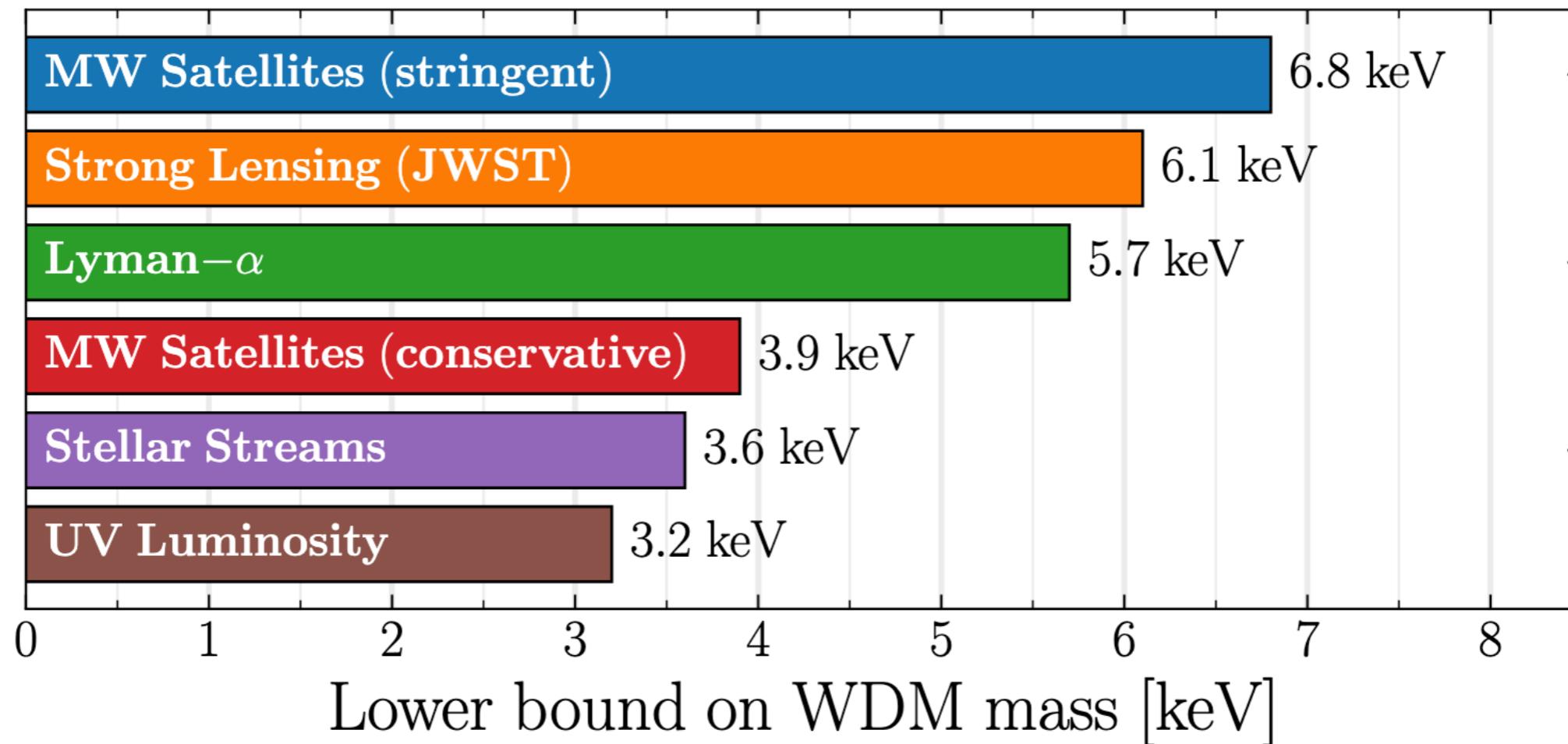
[ESA 2018]

Dark Matter Power Spectrum



[SDSS Collaboration, APJ, 702
(2004), astro-ph/0310725]

WDM constraints from observations



[R. E. Keeley+, MRAS 535, 1652 (2024), 2405.01620]

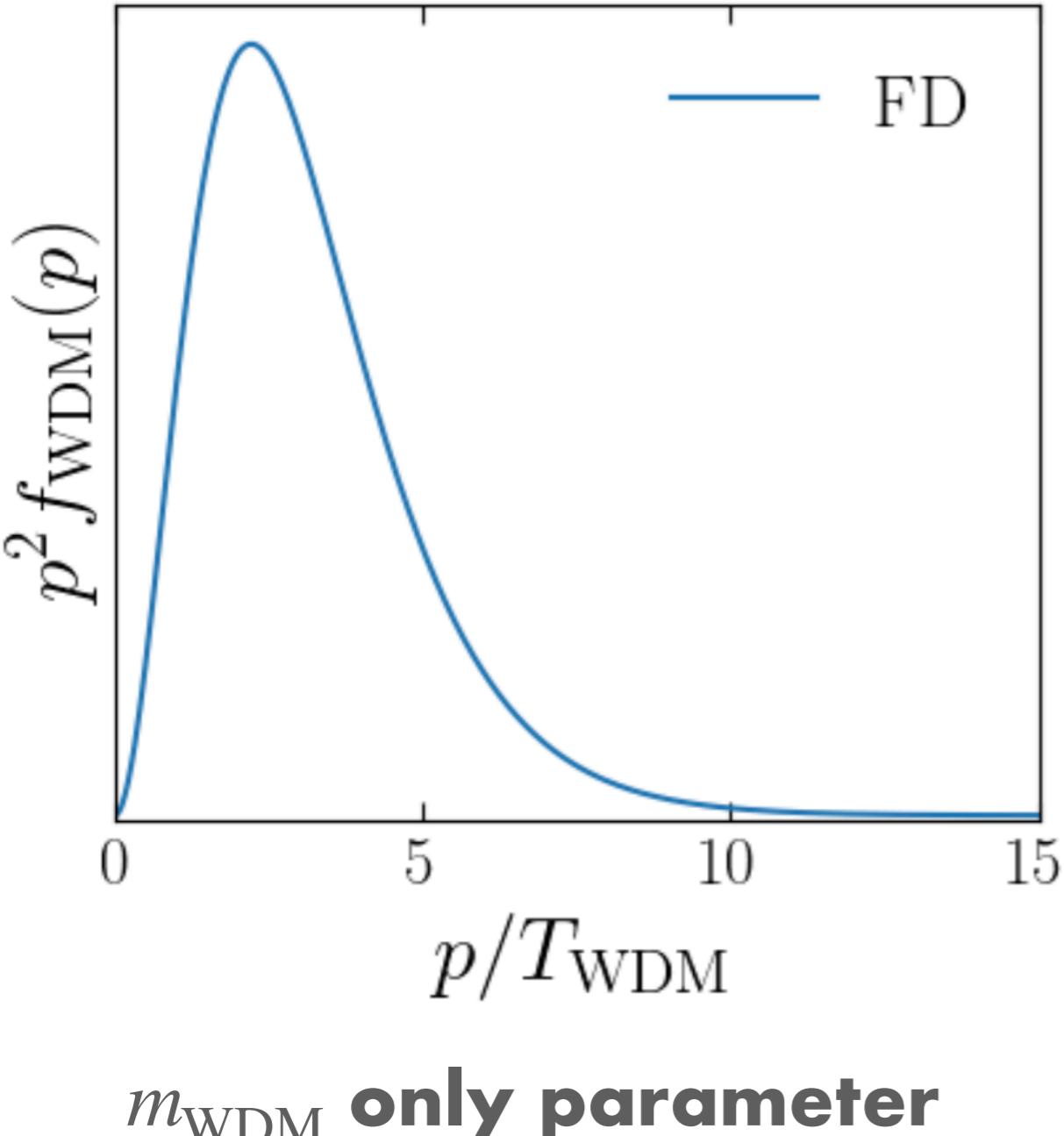
[V. Irsic+, PRD 109, 043511 (2024), 2309.04533]

[N. Banik+, MNRAS, 502, 2364 (2021), 1911.02662]

[B. Liu+, APJ 968, 79 (2024), 2404.13596]

Warm Dark Matter

The PSD is assumed to be thermal, e.g. with FD statistics $\sim \nu$



$$f_{\text{WDM}} \propto \frac{1}{\exp(p/T_{\text{WDM}}) + 1}$$

T_{WDM} set by the DM relic density

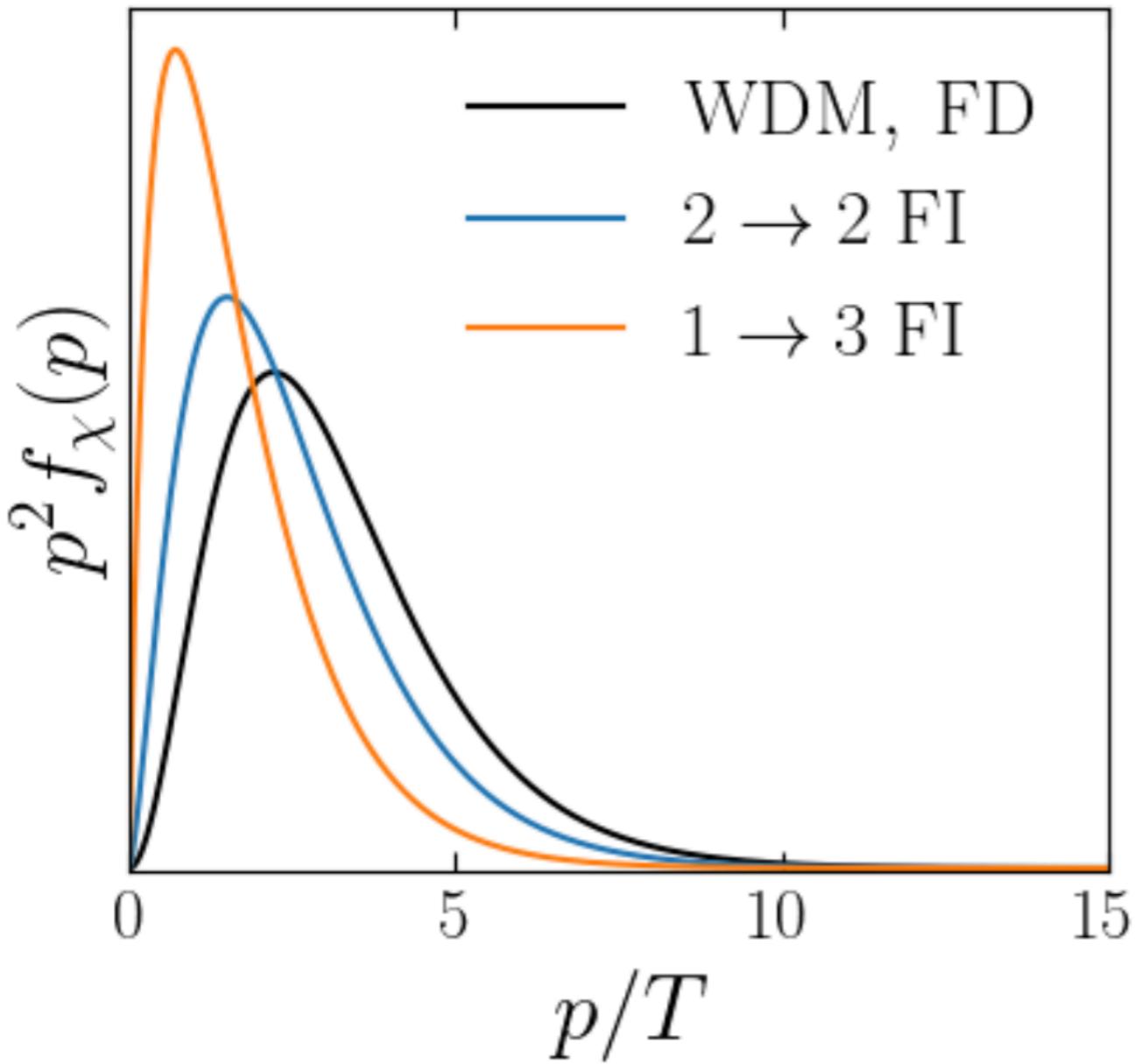
$$\rho_{\text{WDM}} \propto m_{\text{WDM}} T_{\text{WDM}}^3$$

RMS velocity

$$W_\chi = \frac{1}{m_{\text{WDM}}} \sqrt{\langle p^2 \rangle} = \sigma_{\text{WDM}} \frac{T_{\text{WDM}}}{m_{\text{WDM}}} \simeq 3.6$$

Quasi-thermal DM

Physical properties approximated by the very first few moments of the PSD



$$n = \int \frac{d^3 p}{(2\pi)^3} f_\chi(p)$$

$$\langle p \rangle = \int \frac{d^3 p}{(2\pi)^3} p f_\chi(p)$$

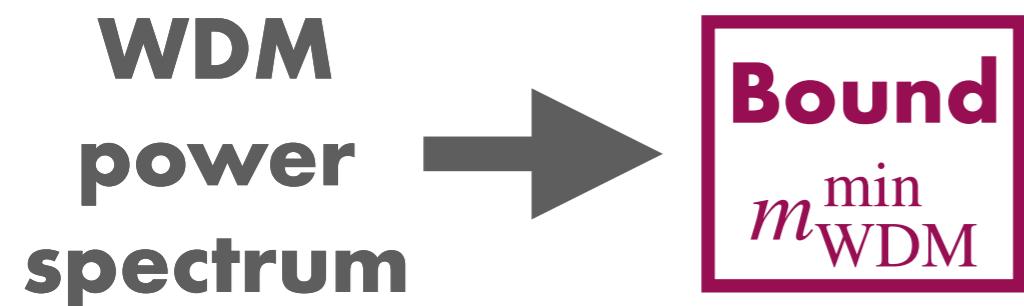
$$\sqrt{\langle p^2 \rangle} = \sqrt{\int \frac{d^3 p}{(2\pi)^3} p^2 f_\chi(p)}$$

Rms velocity

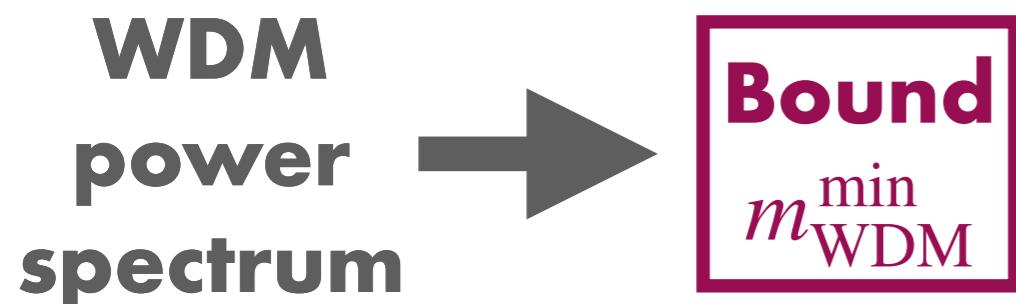
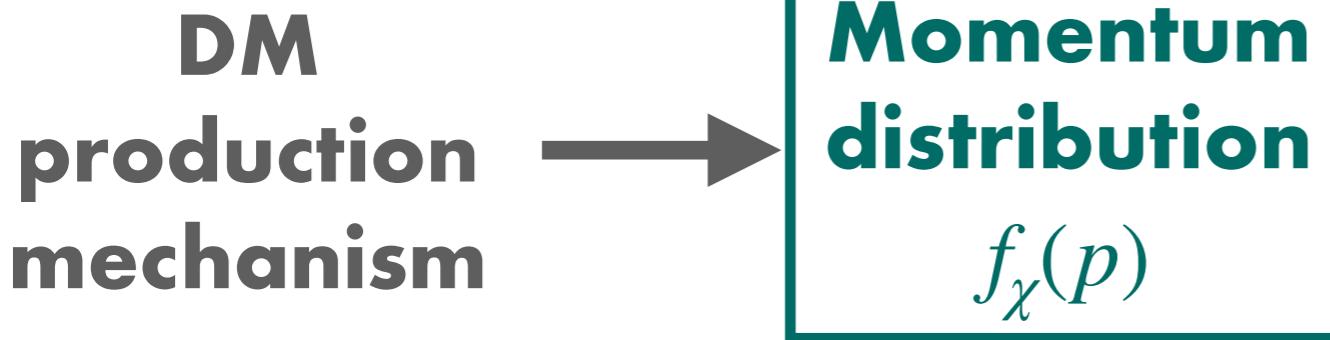
$$W_\chi = \frac{1}{m_\chi} \sqrt{\langle p^2 \rangle}$$

The general strategy

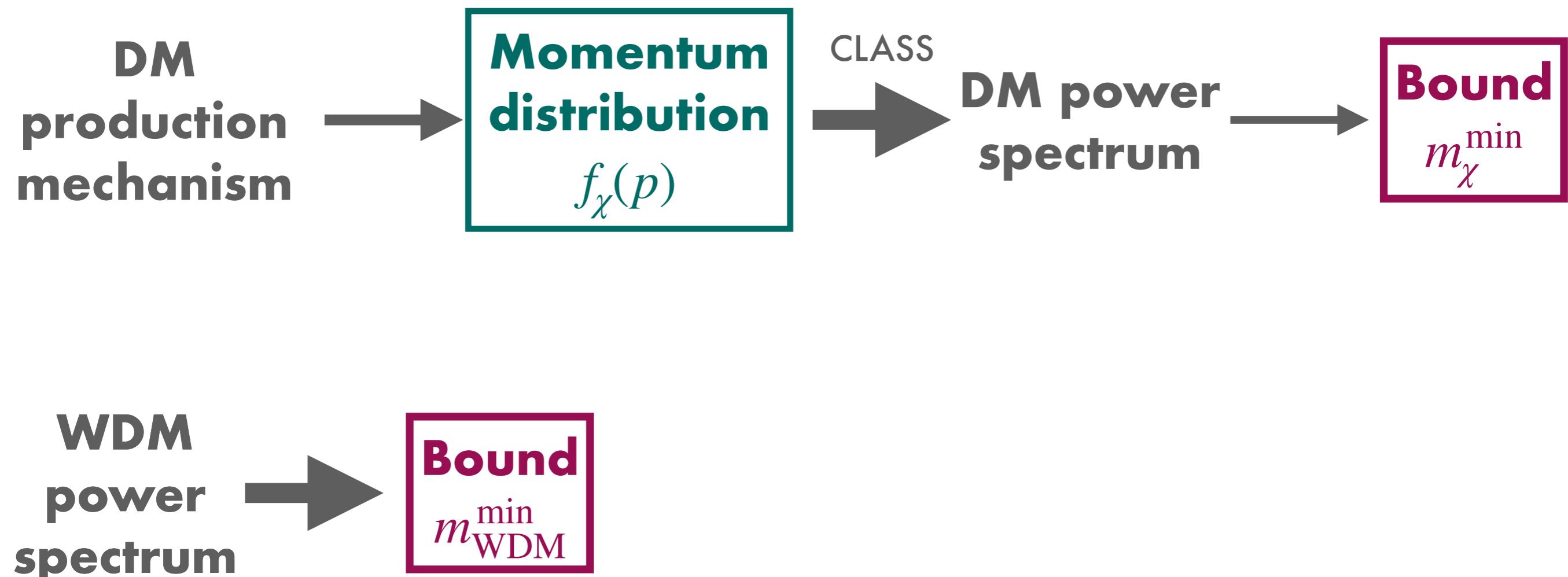
DM
production
mechanism



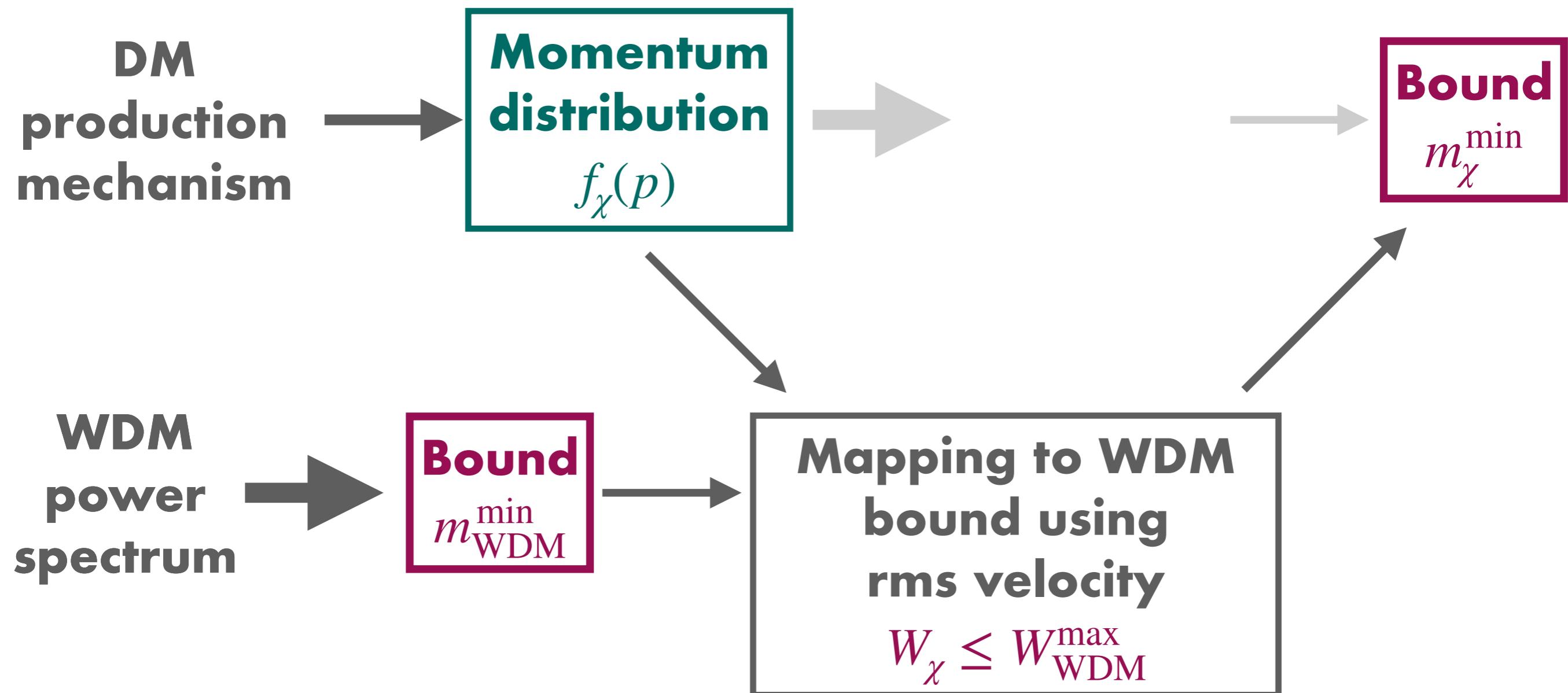
The general strategy



The general strategy



The general strategy



[F. D'Eramo, AL JCAP 10 (2021) 045, [2012.01446](#)]

The map to WDM

$$p \propto a^{-1}$$

After production, PSD is “frozen” and the DM free-streams

$$q = \frac{p}{T_P} \frac{a}{a(t_P)}$$

comoving momentum

The mapping requires

$$W_\chi \leq W_{\text{WDM}}(m_{\text{WDM}}^{\min})$$

$$m_\chi^{\min} = 22 \text{ keV} \left(\frac{m_{\text{WDM}}^{\min}}{6 \text{ keV}} \right)^{4/3} \left(\frac{\sigma_q}{3.6} \right) \left(\frac{106.75}{g_{\star s}(T_P)} \right)^{1/3}$$

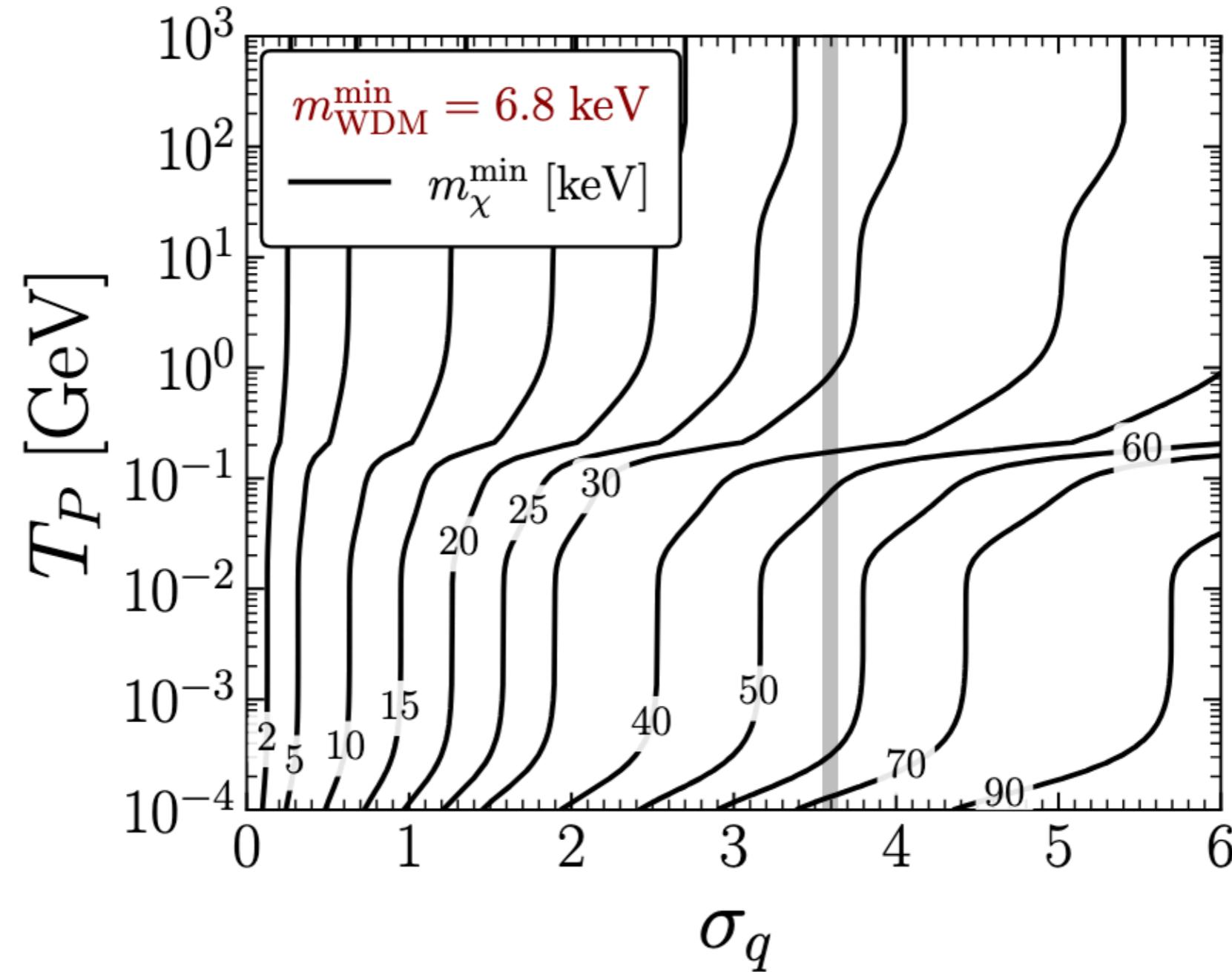
Reference
WDM bound from
observables

$$\sigma_q^2 = \frac{\int dq q^4 f_\chi}{\int dq q^2 f_\chi}$$

Reference
temperature,
determines
definition of σ_q

Mass bounds on quasi-thermal DM

Reference temperature, can be associated to the production energy scale



PSD width, depends on the production mechanism (PSD shape)

Quasi-thermal DM via freeze-in

$$\frac{df_\chi}{dt} = \frac{C(p, t)}{E}$$

Quasi-thermal DM via freeze-in

$$\frac{df_\chi}{dt} = \frac{C(p, t)}{E}$$

Two-body decays $\sigma_q \simeq 2.96\sqrt{1 - (m_2/m_1)^2}$

$$C_{1 \rightarrow 2\chi} = \frac{\ell}{2} \int d\Pi_1 d\Pi_2 (2\pi)^4 \delta^{(4)}(P_1 - P_2 - P) |\mathcal{M}|^2 f_1^{\text{eq}} .$$

Binary scatterings $\sigma_q \simeq 2.96$

$$C_{12 \rightarrow 3\chi} = \frac{\ell}{2} \int \prod_{i=1}^3 d\Pi_i (2\pi)^4 \delta^{(4)}(P_i - P_f) |\mathcal{M}|^2 f_1^{\text{eq}} f_2^{\text{eq}} .$$

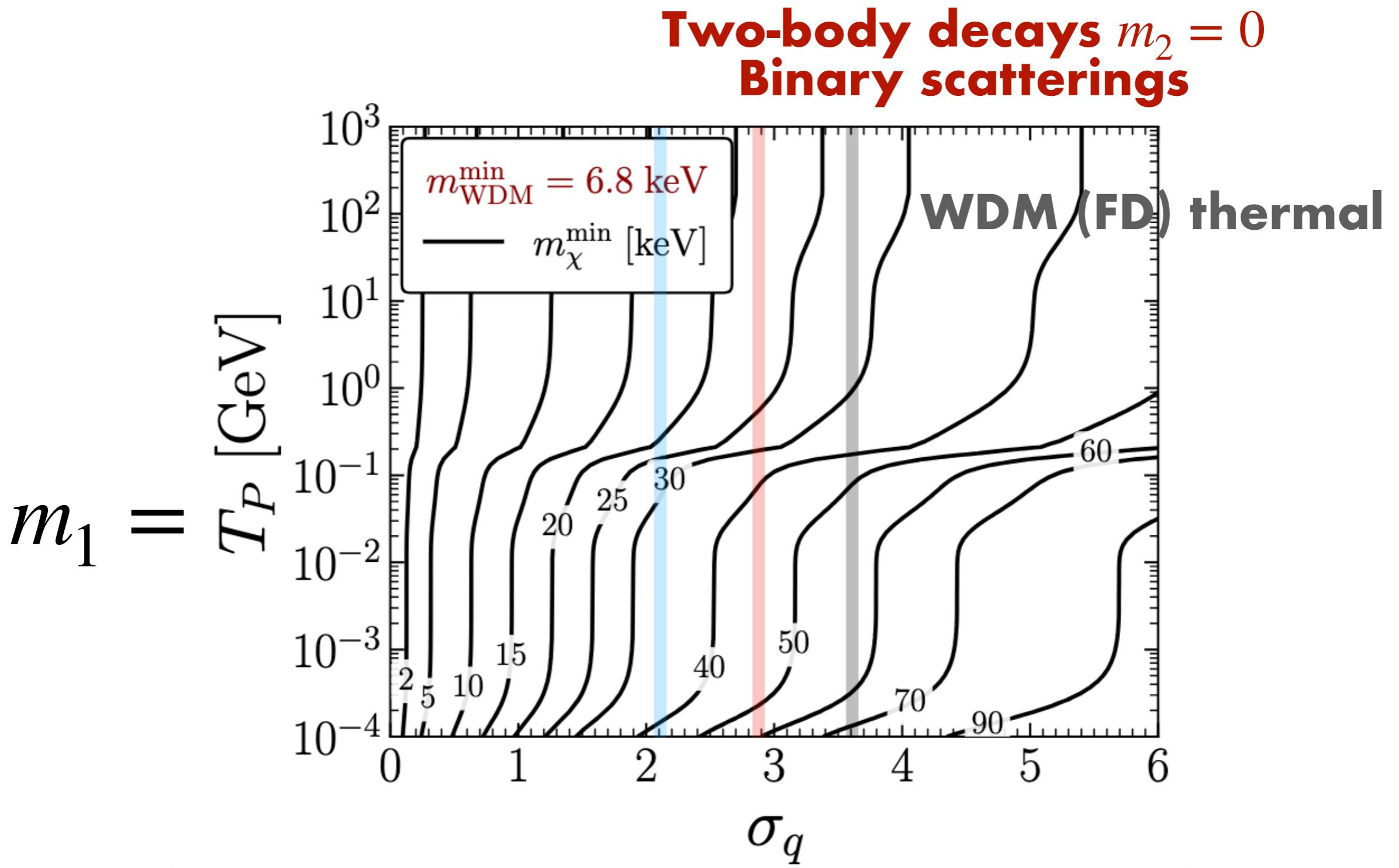
Three-body decays $\sigma_q \simeq 2.09\sqrt{1 - (m_2/m_1)^2}$

$$C_{1 \rightarrow 23\chi} = \frac{\ell}{2} \int \prod_{i=1}^3 d\Pi_i (2\pi)^4 \delta^{(4)}(P_i - P_f) |\mathcal{M}|^2 f_1^{\text{eq}} .$$

$$f_\chi \propto \frac{1}{\sqrt{q}} e^{-q(1-m_2^2/m_1^2)^{-1}}$$

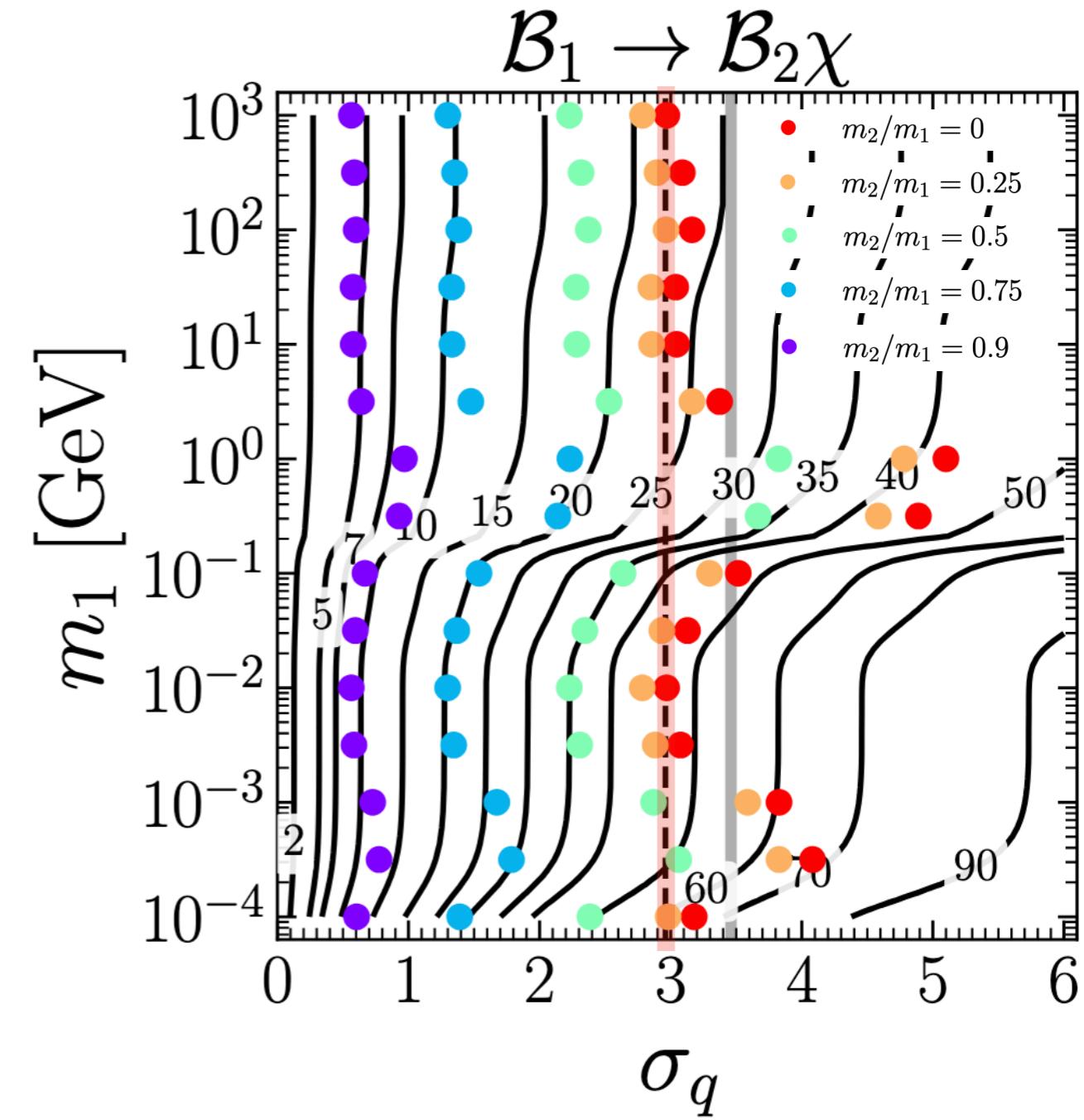
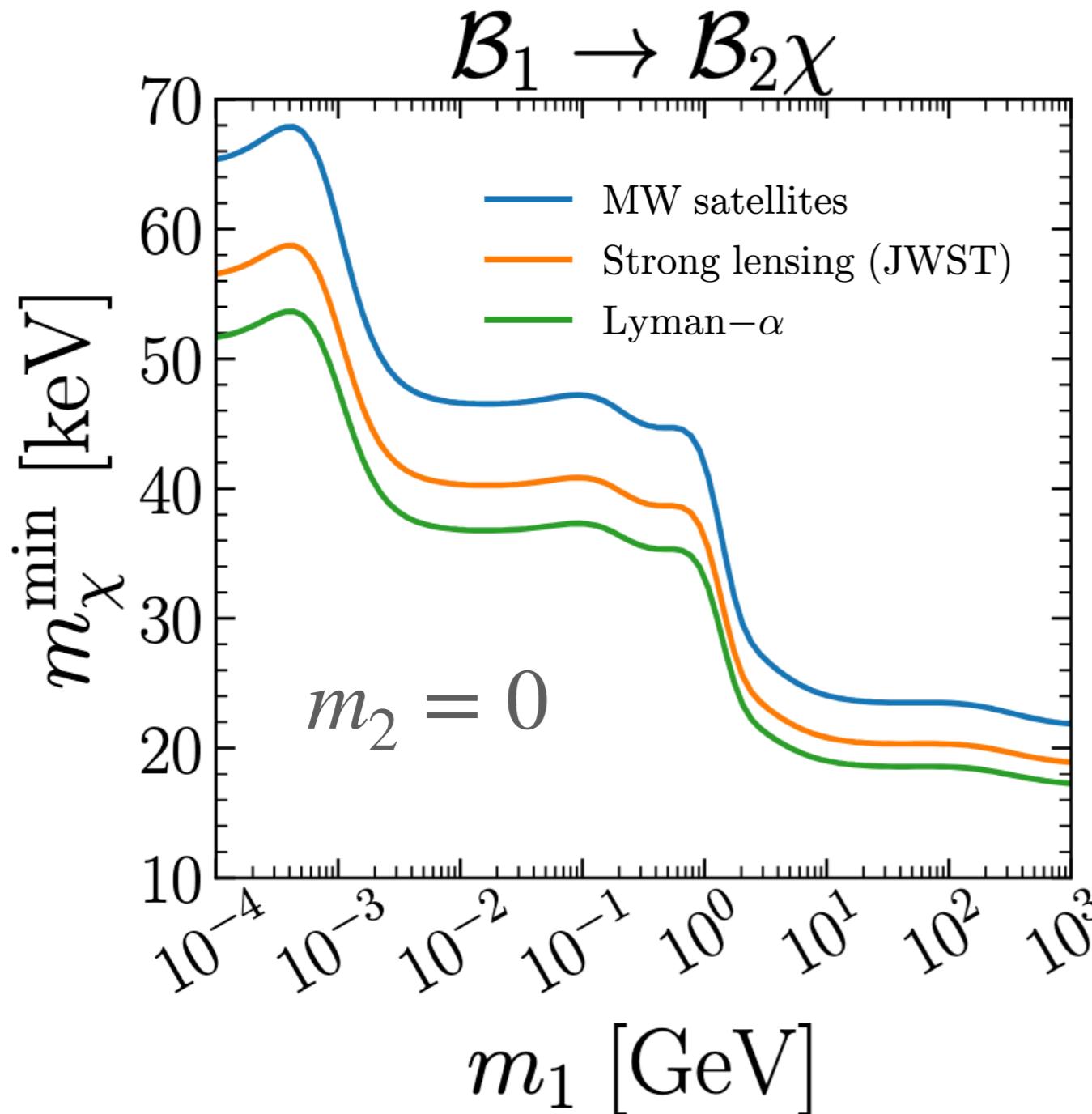
$$f_\chi \propto \frac{1}{\sqrt{q}} e^{-q}$$

$$f_\chi \propto \frac{1}{q} \operatorname{erfc} \left(\sqrt{\frac{q}{1 - m_2^2/m_1^2}} \right)$$



Three-body decays $m_2 = m_3 = 0$

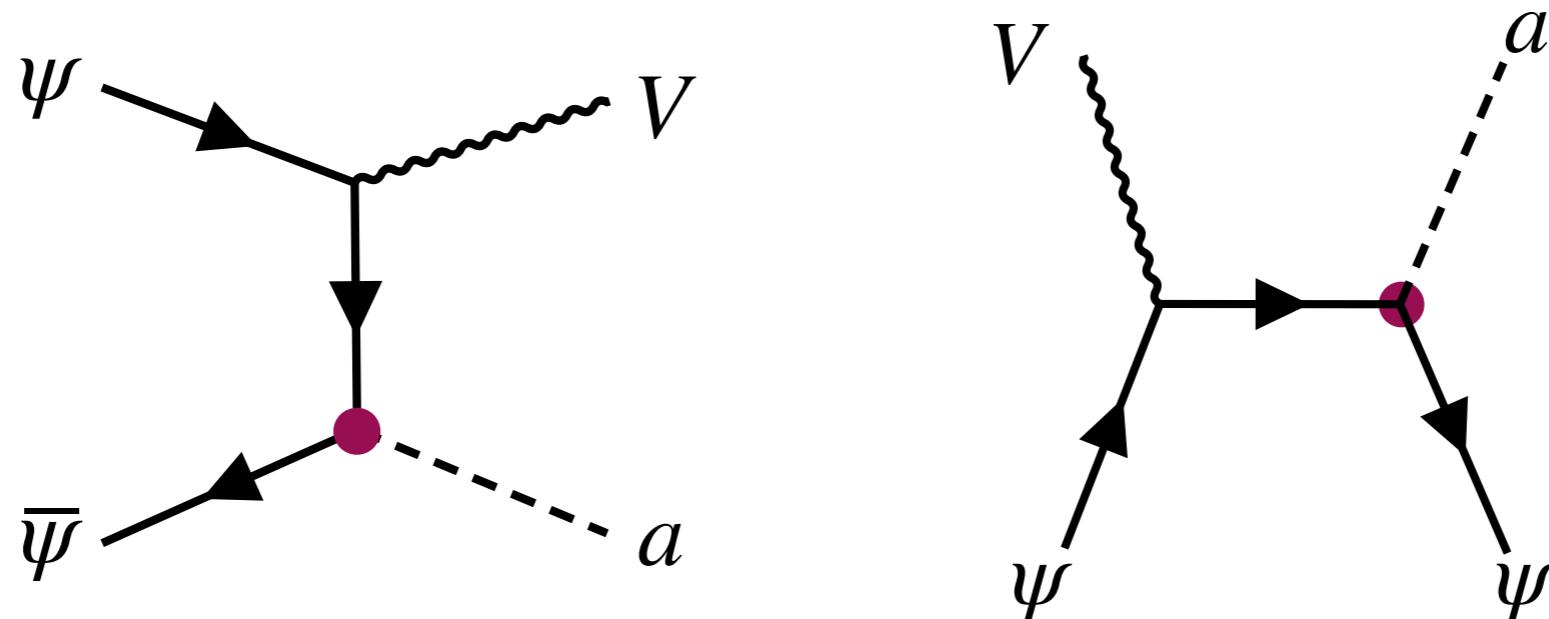
Freeze-in general bounds, decays



Axions coupled to SM fermions

$$\mathcal{L}_a = -\frac{\partial_\mu a}{2\Lambda_\psi} \sum_\psi c_\psi \bar{\psi} \gamma^\mu \gamma^5 \psi$$

ψ is a lepton or a quark
 V is a photon or gluon



IR-dominated freeze-in via binary scatterings

We work in the limit $m_a \ll m_\psi$ and $T < v_{EW}$

Lower mass bounds on quasi-thermal axion

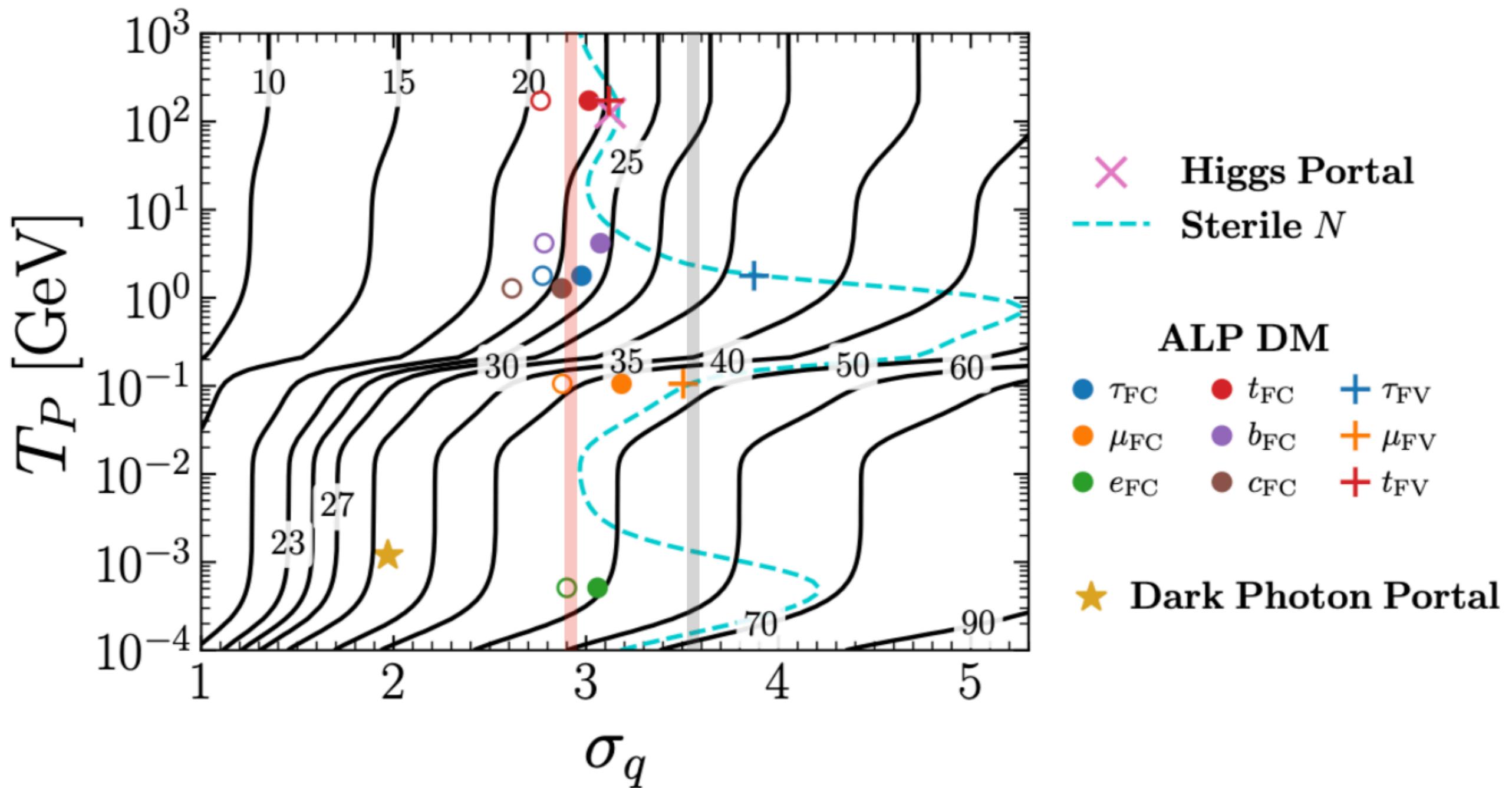
	m_a^{\min} (keV)			Relic density requirement and X-ray constraints			
ψ	Ly α	JWST	Satellite	$\Lambda_{\psi}^{\text{relic}}$ (GeV)	$-\Delta g_{a\gamma\gamma}$ (GeV $^{-1}$)	$\tau_{a \rightarrow \gamma\gamma}$ (sec)	Allowed?
electron	39	43	49	5.7×10^7	3.1×10^{-14}	1.2×10^{18}	✗
muon	33	37	43	3.6×10^8	5.6×10^{-20}	1.1×10^{30}	✓
tau	19	21	24	4.9×10^8	7.3×10^{-23}	1.8×10^{36}	✓
charm	19	20	23	4.5×10^9	1.8×10^{-23}	3.1×10^{37}	✓
bottom	19	21	24	7.3×10^9	3.2×10^{-25}	8.3×10^{40}	✓
top	18	19	22	2.8×10^{10}	1.5×10^{-28}	5.5×10^{47}	✓

**Minimum
ALP mass**

**Symmetry
breaking scale
(coupling) fixed
by relic density**

**1-loop
induced
coupling to
photons and
ALP lifetime**

Summary



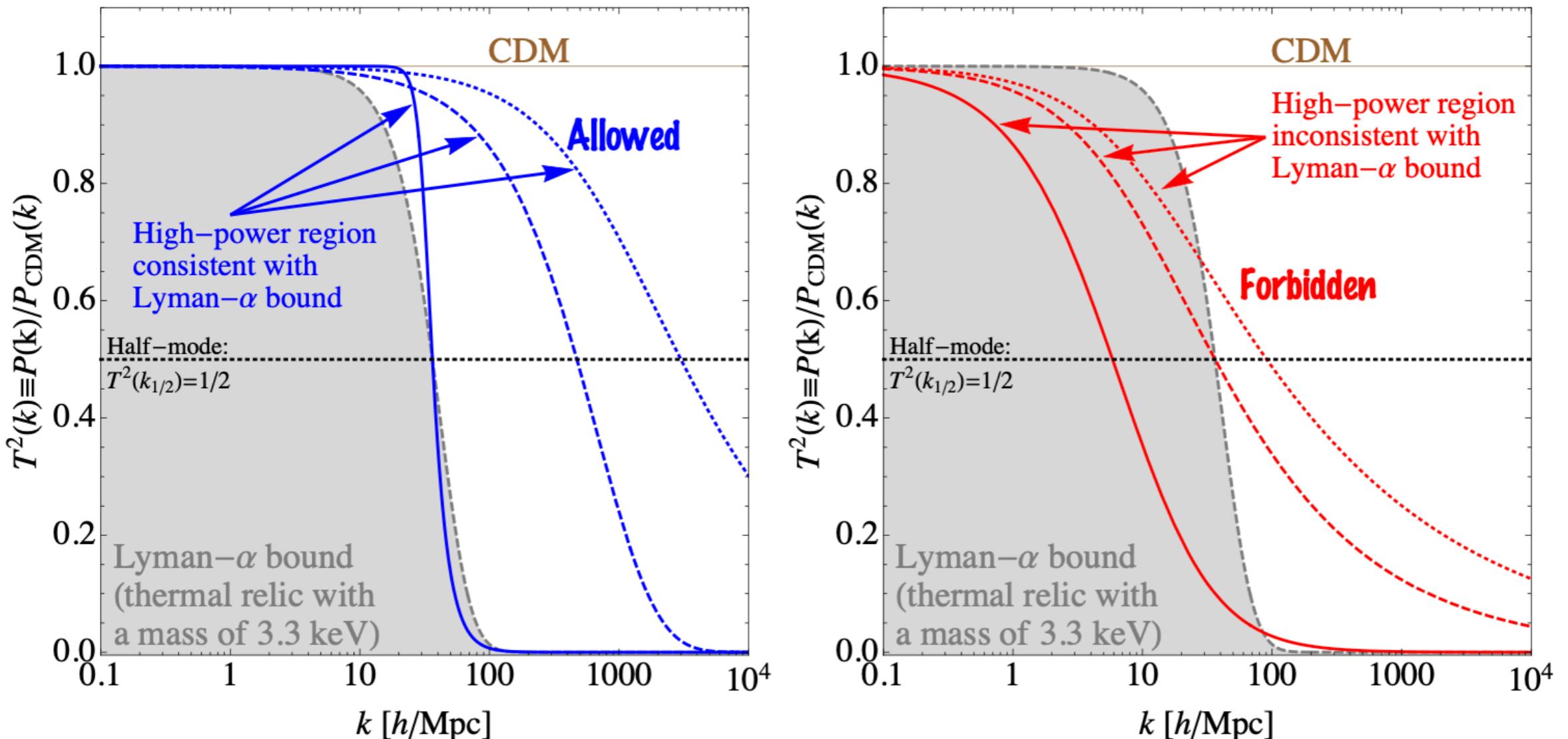
...and outlook

Much to do, different cosmologies,
mixed CDM+WDM scenarios, different models,
all in phase space!

Thank you!

alessandro.lenoci@mail.huji.ac.il

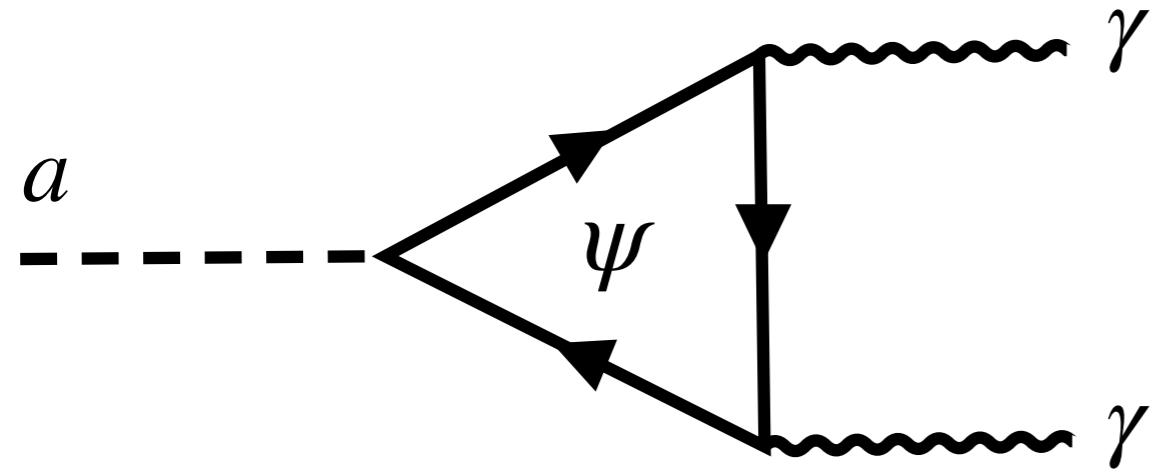
WDM power spectrum



WDM
power
spectrum → **Bound**
 m_{WDM}^{\min}

Constraints on keV axions

A coupling to photons is induced via a fermion loop



[Bauer et al 1708.00443]
[Bauer et al 2012.12272]

$$\Delta g_{a\gamma\gamma} \approx -\frac{\alpha_{\text{em}} N_\psi^c Q_\psi^2}{12\pi f_a} \left(\frac{m_a}{m_\psi}\right)^2$$

X-ray, NuStar, INTEGRAL, constrain

[Calore et al 2209.06299]
[eg. Roach et al 2207.04572]

$$g_{a\gamma\gamma} < 10^{-19} \text{ GeV}^{-1}$$

for axions in the 10 keV - 100 keV range.

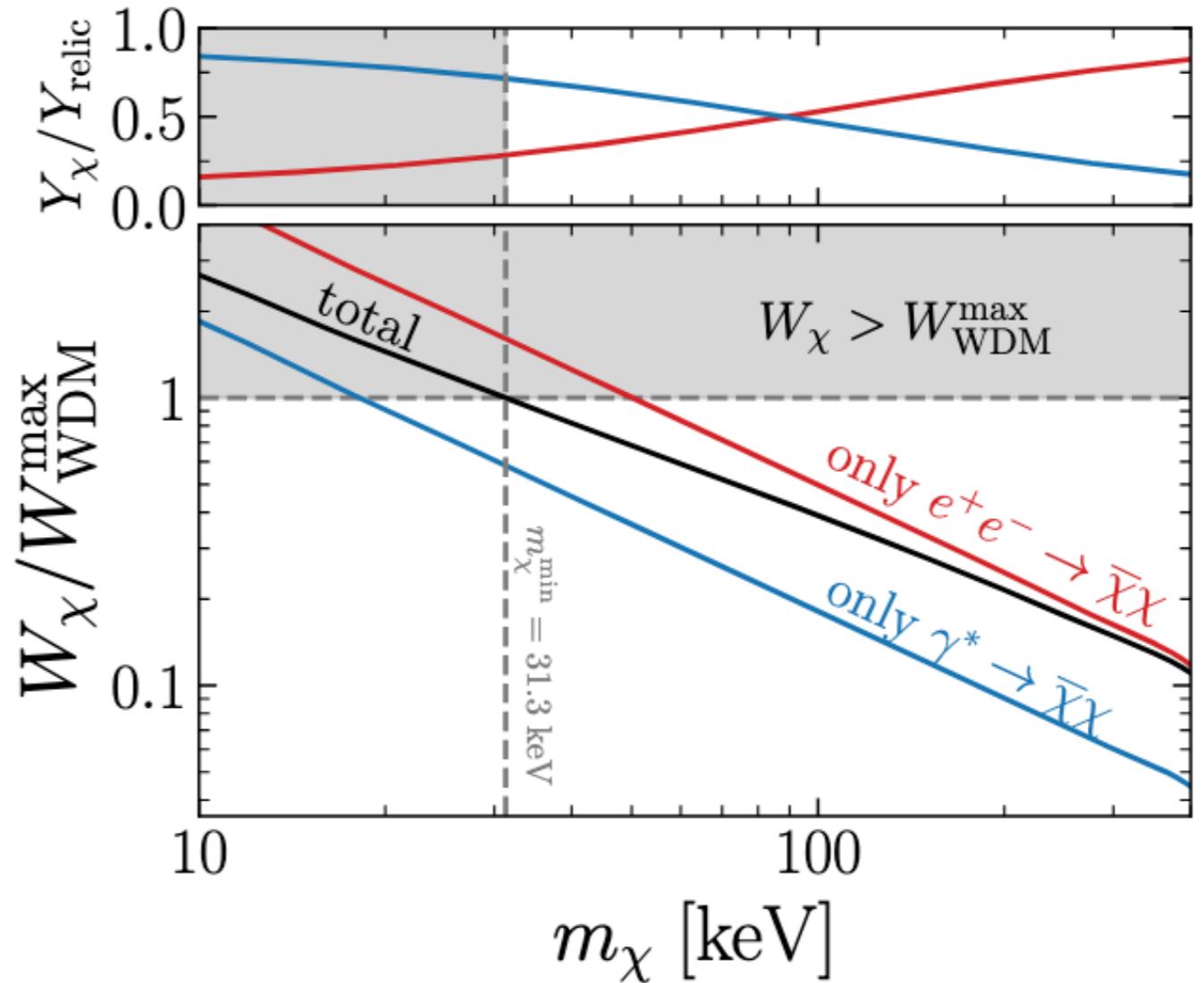
We check for $(m_a^{\min}, f_a^{\text{relic}})$ if the bound is respected or not

Dark photon portal DM

$$\begin{aligned} \mathcal{L}_{\text{DP}} = & -\frac{1}{4}F'_{\mu\nu}F'^{\mu\nu} + \frac{\kappa}{2}F'_{\mu\nu}F^{\mu\nu} + \frac{1}{2}m_{A'}^2 A'_\mu A'^\mu \\ & + eJ_{\text{EM}}^\mu A_\mu + \bar{\chi}\gamma^\mu(i\partial_\mu + e_D A'_\mu)\chi. \end{aligned}$$

Two processes:

- 1) Pair annihilation**
- 2) Plasmon decays**



Dark photon portal DM

$$\begin{aligned}\mathcal{L}_{\text{DP}} = & -\frac{1}{4}F'_{\mu\nu}F^{\prime\mu\nu} + \frac{\kappa}{2}F'_{\mu\nu}F^{\mu\nu} + \frac{1}{2}m_{A'}^2 A'_\mu A'^\mu \\ & + e J_{\text{EM}}^\mu A_\mu + \bar{\chi} \gamma^\mu (i\partial_\mu + e_D A'_\mu) \chi.\end{aligned}$$

