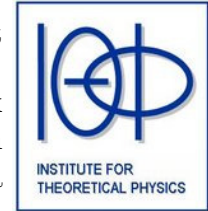




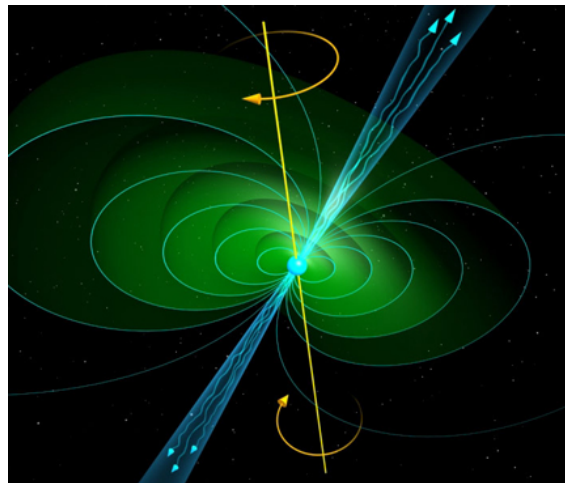
Andreas Schmitt

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1040 Vienna, Austria



Inverse magnetic catalysis in dense (holographic) matter

F. Preis, A. Rebhan, A. Schmitt, JHEP 1103, 033 (2011); JPG 39, 054006 (2012)



- **Magnetic catalysis**

K. G. Klimenko, Theor. Math. Phys. 89, 1161-1168 (1992)

V. P. Gusynin, V. A. Miransky, I. A. Shovkovy, PLB 349, 477-483 (1995)

- (massless) fermions in **Nambu-Jona-Lasinio (NJL)** model

$$\mathcal{L}_{\text{NJL}} = \bar{\psi} i\gamma^\mu \partial_\mu \psi + G \left[(\bar{\psi}\psi)^2 + (\bar{\psi} i\gamma^5 \psi)^2 \right]$$

- mean-field approximation: $\bar{\psi}\psi = \langle \bar{\psi}\psi \rangle + \text{fluctuations}$

Zero magnetic field:

dynamical fermion mass

$$M \propto \langle \bar{\psi}\psi \rangle \neq 0$$

for coupling $g > g_c = 1$

dimensionless coupling $g \equiv G\Lambda^2/\pi^2$

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K. G. Klimenko, Theor. Math. Phys. 89, 1161-1168 (1992)

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- mean-field approximation: $\bar{\psi}\psi = \langle \bar{\psi}\psi \rangle + \text{fluctuations}$

Nonzero magnetic field:

$M \neq 0$ for *arbitrarily small* g ,

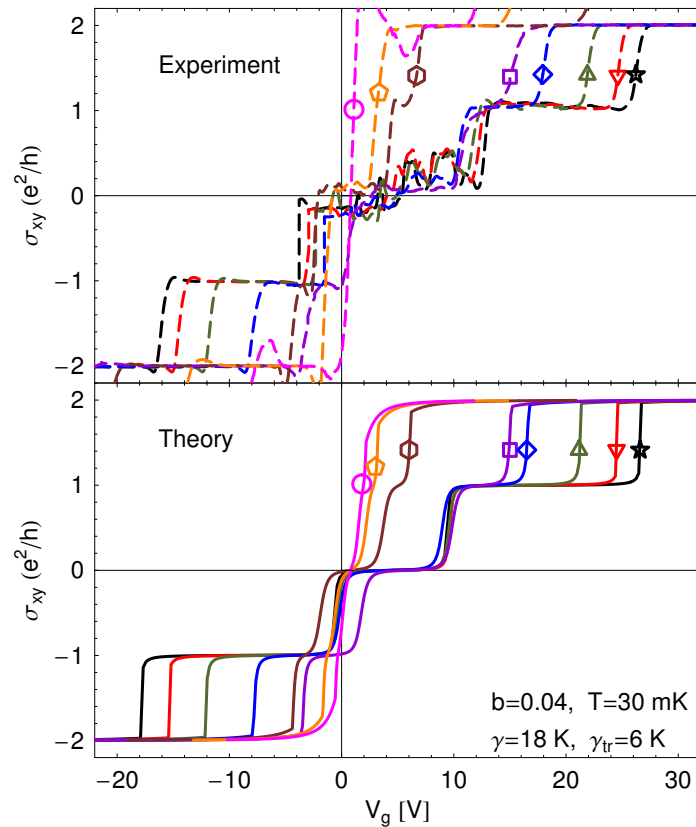
$$M \propto \sqrt{eB} e^{-\text{const.}/eBg}$$

at *weak coupling* $g \ll 1$

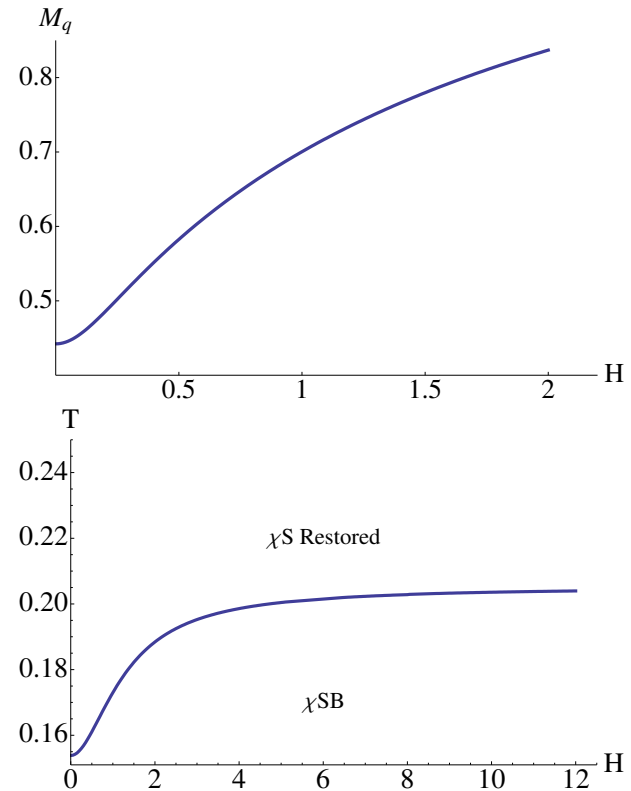


dimensionless coupling $g \equiv G\Lambda^2/\pi^2$

- **Magnetic catalysis in the real world and in holography**



V.P.Gusynin *et al.*, PRB 74, 195429 (2006)

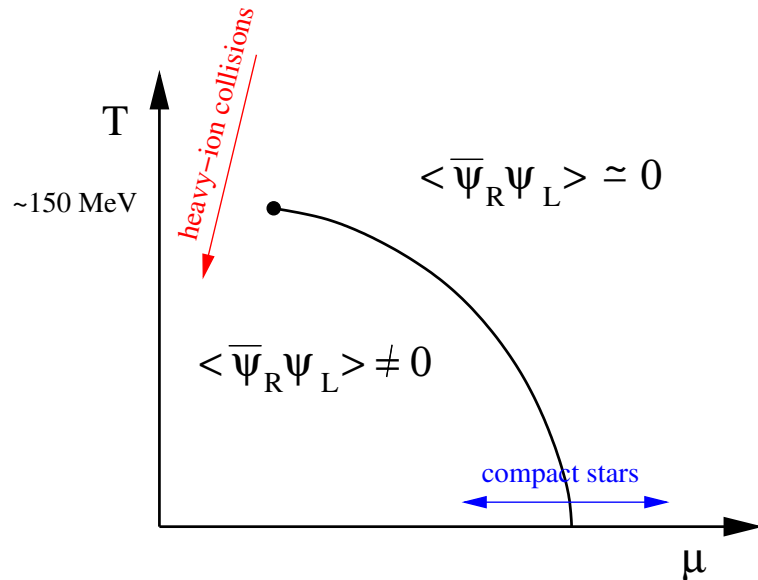


C.V.Johnson, A.Kundu, JHEP 0812, 053 (2008)

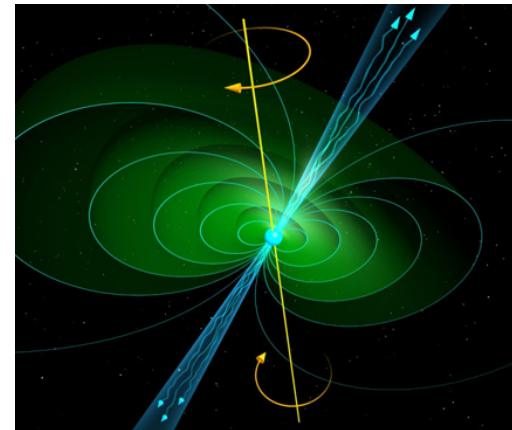
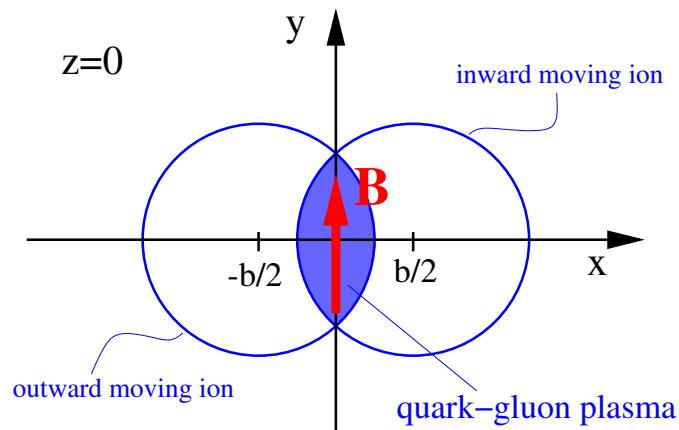
- **graphene**: appearance of additional plateaus in strong magnetic fields
[$B = 9$ T (pink), $B = 45$ T (black)]

- **holography**: magnetic field enhances dynamical mass M_q and critical temperature T_c

• (Chiral) magnetic catalysis in QCD?



- chiral transition probed in **heavy-ion collisions** and **compact stars**
- in both instances huge magnetic fields are present!



heavy-ion collisions: temporarily

$$B \sim 10^{18} \text{ G}$$

magnetars: at surface

$$B \sim 10^{15} \text{ G}$$

- **Chiral transition in the Sakai-Sugimoto model**

- original AdS/CFT correspondence

J. M. Maldacena, *Adv. Theor. Math. Phys.* 2, 231 (1998)

(supergravity limit of)
string theory on $AdS_5 \times S^5$

\Leftrightarrow

(strong coupling limit of)
 $\mathcal{N} = 4$ SYM theory on $\mathbb{R}^{3,1}$

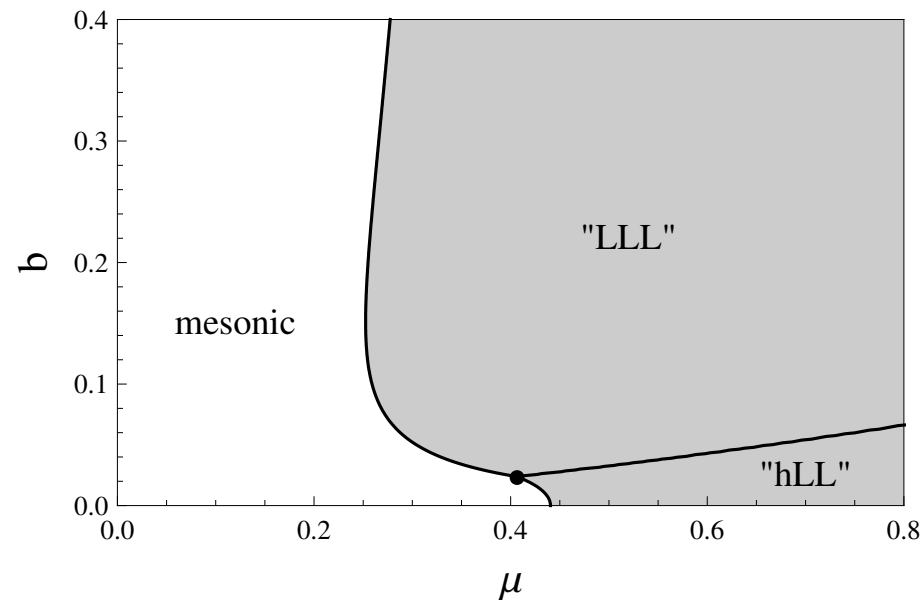
- Sakai-Sugimoto model

E. Witten, *Adv. Theor. Math. Phys.* 2, 505 (1998)

T. Sakai, S. Sugimoto, *Prog. Theor. Phys.* 113, 843-882 (2005)

- breaks supersymmetry completely
- accounts for confinement and chiral symmetry breaking
- reproduces meson/hadron spectra surprisingly well
- here: “NJL limit” (as opposed to large- N_c QCD limit)
 - small separation of flavor branes
 - deconfined phase
 - chiral transition sensitive to μ , B

- $T = 0$ phase diagram



- apparent Landau level transition

G. Lifschytz, M. Lippert, PRD 80, 066007 (2009)

(charge density exactly like in LLL for free fermions)

- non-monotonic behavior of critical μ

(in apparent contrast to magnetic catalysis)

- **Inverse magnetic catalysis**

Why does B *restore* chiral symmetry for certain μ ?

- free energy *gain* from $\bar{\psi} - \psi$ pairing increases with B (magnetic catalysis)
- μ induces free energy *cost* for pairing; this cost depends on B !

weak coupling (NJL):

E. V. Gorbar *et al.*, PRC 80, 032801 (2009)

$$\Delta\Omega \propto B[\mu^2 - M(B)^2/2]$$

just like Clogston limit $\delta\mu = \frac{\Delta}{\sqrt{2}}$

in superconductivity

A. Clogston, PRL 9, 266 (1962)

B. Chandrasekhar, APL 1, 7 (1962)

Sakai-Sugimoto:

large B :

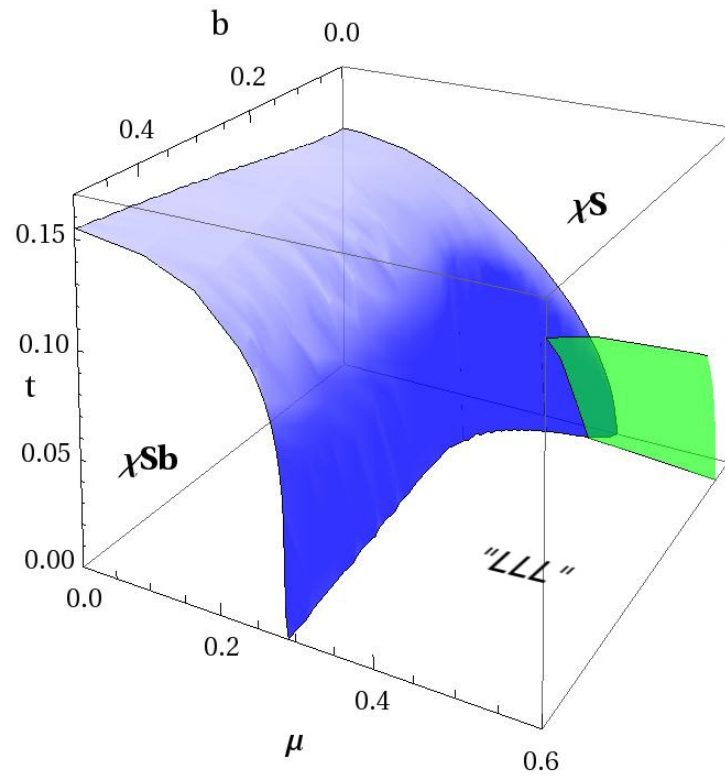
$$\Delta\Omega \propto B[\mu^2 - 0.12 M(B)^2]$$

small B :

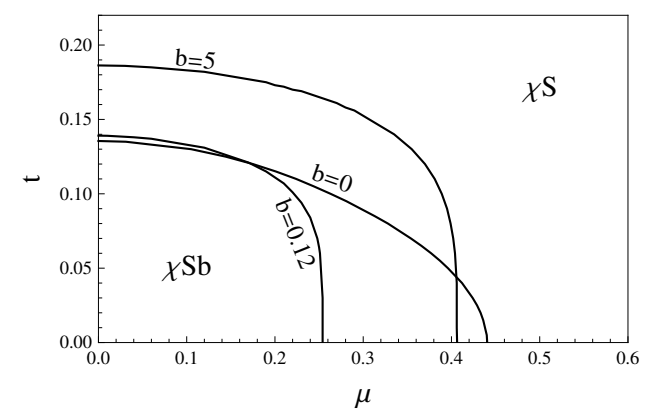
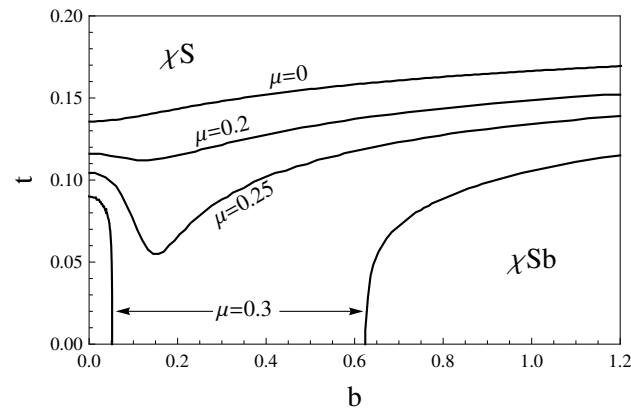
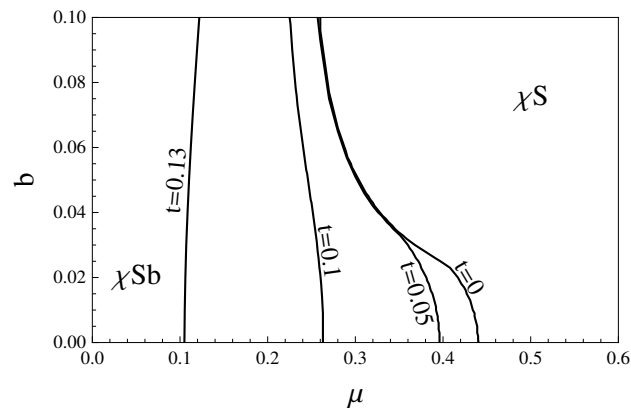
$$\Delta\Omega \propto \mu^2 B - \text{const} \times M(B)^{7/2}$$

“Inverse magnetic catalysis”

● Phase structure at nonzero temperature



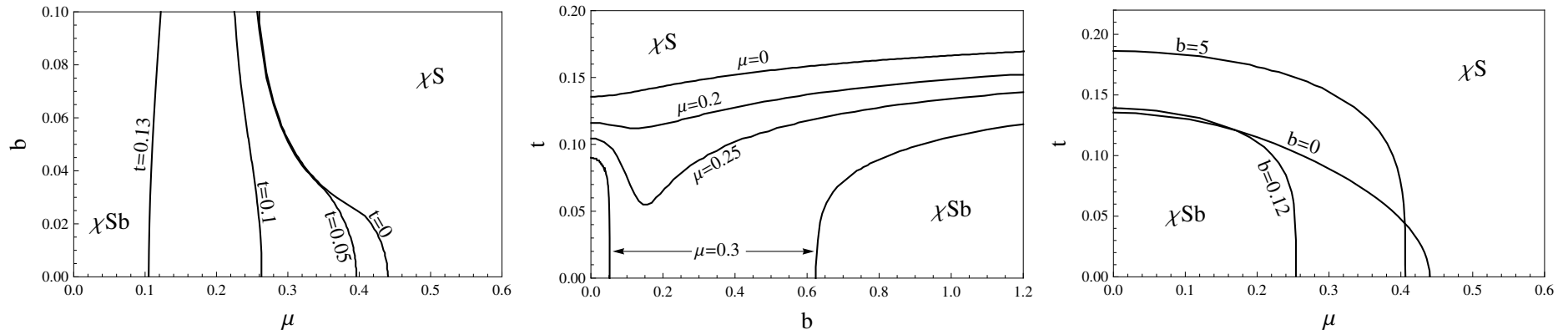
blue: chiral phase transition
green: "LLL" transition



- Agreement with NJL calculation

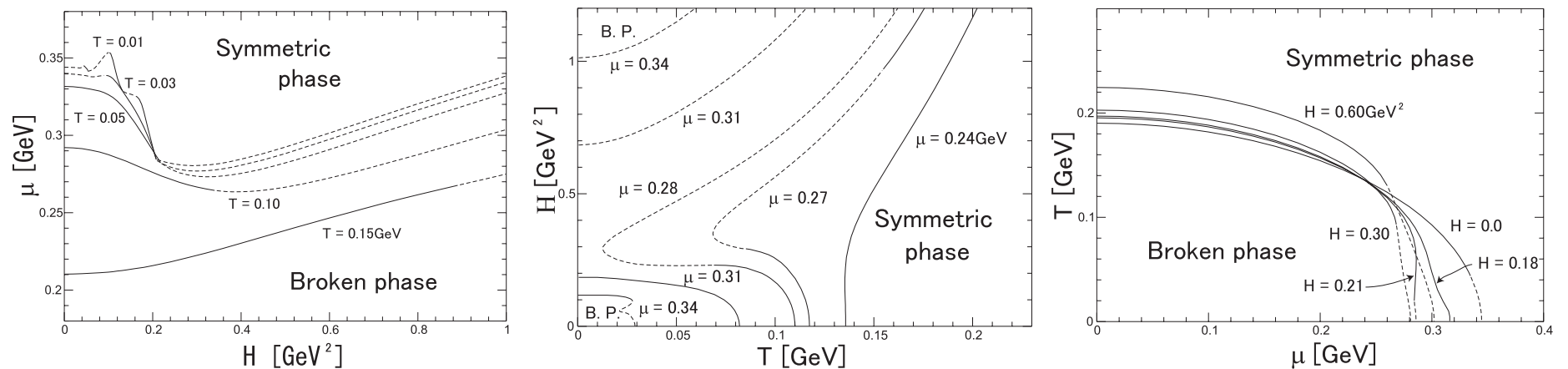
Sakai-Sugimoto:

F. Preis, A. Rebhan and A. Schmitt, JHEP 1103, 033 (2011)



NJL:

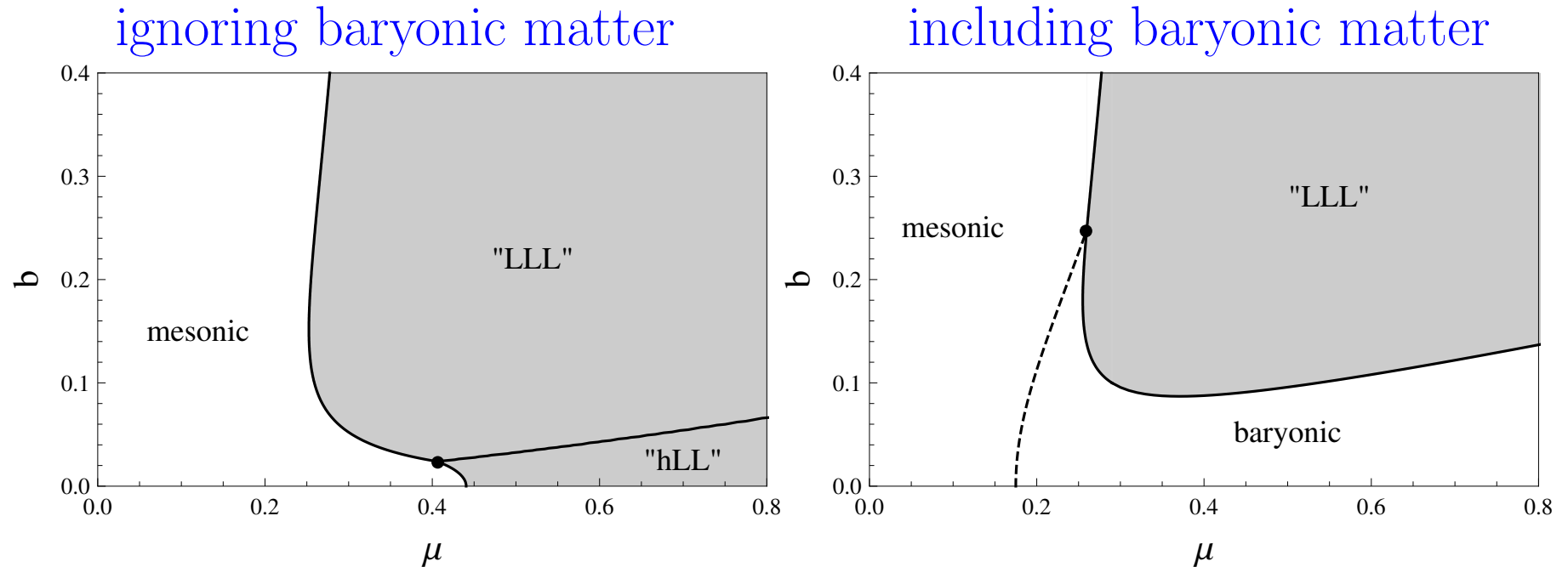
T. Inagaki, D. Kimura, T. Murata, Prog. Theor. Phys. 111, 371-386 (2004)



(IMC also in quark-meson model J. O. Andersen and A. Tranberg, arXiv:1204.3360 [hep-ph])

- **(Homogeneous) baryonic matter in Sakai-Sugimoto**
- baryons in AdS/CFT: wrapped D-branes with N_c strings
E. Witten, JHEP 9807, 006 (1998); D. J. Gross, H. Ooguri, PRD 58, 106002 (1998)
- baryons in Sakai-Sugimoto:
 - D4-branes wrapped on S^4
 - equivalently: instantons on D8-branes (\rightarrow skyrmions)
T. Sakai, S. Sugimoto, Prog. Theor. Phys. 113, 843-882 (2005)
H. Hata, T. Sakai, S. Sugimoto, S. Yamato, Prog. Theor. Phys. 117, 1157 (2007)
- pointlike approximation for $N_f = 1$
O. Bergman, G. Lifschytz, M. Lippert, JHEP 0711, 056 (2007)

- **Effect of baryons on $T = 0$ phase diagram**



- small b : baryonic matter prevents the system from restoring chiral symmetry
- baryon onset (2nd order!) intersects chiral phase transition
→ large b : mesonic matter superseded by quark matter
- with baryonic matter, IMC plays an even more prominent role in the phase diagram

- **Summary**
- dense matter in the “NJL limit” of the Sakai-Sugimoto model...
 - ... shows a transition reminiscent of LLL
 - ... shows a chiral phase transition with MC at large B and IMC at small B
- baryonic matter changes the phase diagram dramatically for small B
- agreement with NJL calculations only before including baryonic matter

- **Outlook**

- Sakai-Sugimoto calculation:

- understand chiral non-restoration with baryons
 - extend to $N_f = 2$
- more general: problem of large N_c vs. $N_c = 3!$

- inverse catalysis in QCD and NJL models

- how is IMC affected by the “chiral shift”?
 - E. V. Gorbar, V. A. Miransky and I. A. Shovkovy, PRC 80, 032801 (2009)
 - F. Preis, A. Rebhan, A. Schmitt, work in progress
- are recent lattice results for $T_c(B)$ related to IMC?
 - G. S. Bali *et al.*, JHEP 1202, 044 (2012)