

Subtracted dispersion relations for virtual Compton scattering off the proton

Matteo Ronchi^a, Barbara Pasquini^b, Marc Vanderhaeghen^a, Igor Danilkin^a.

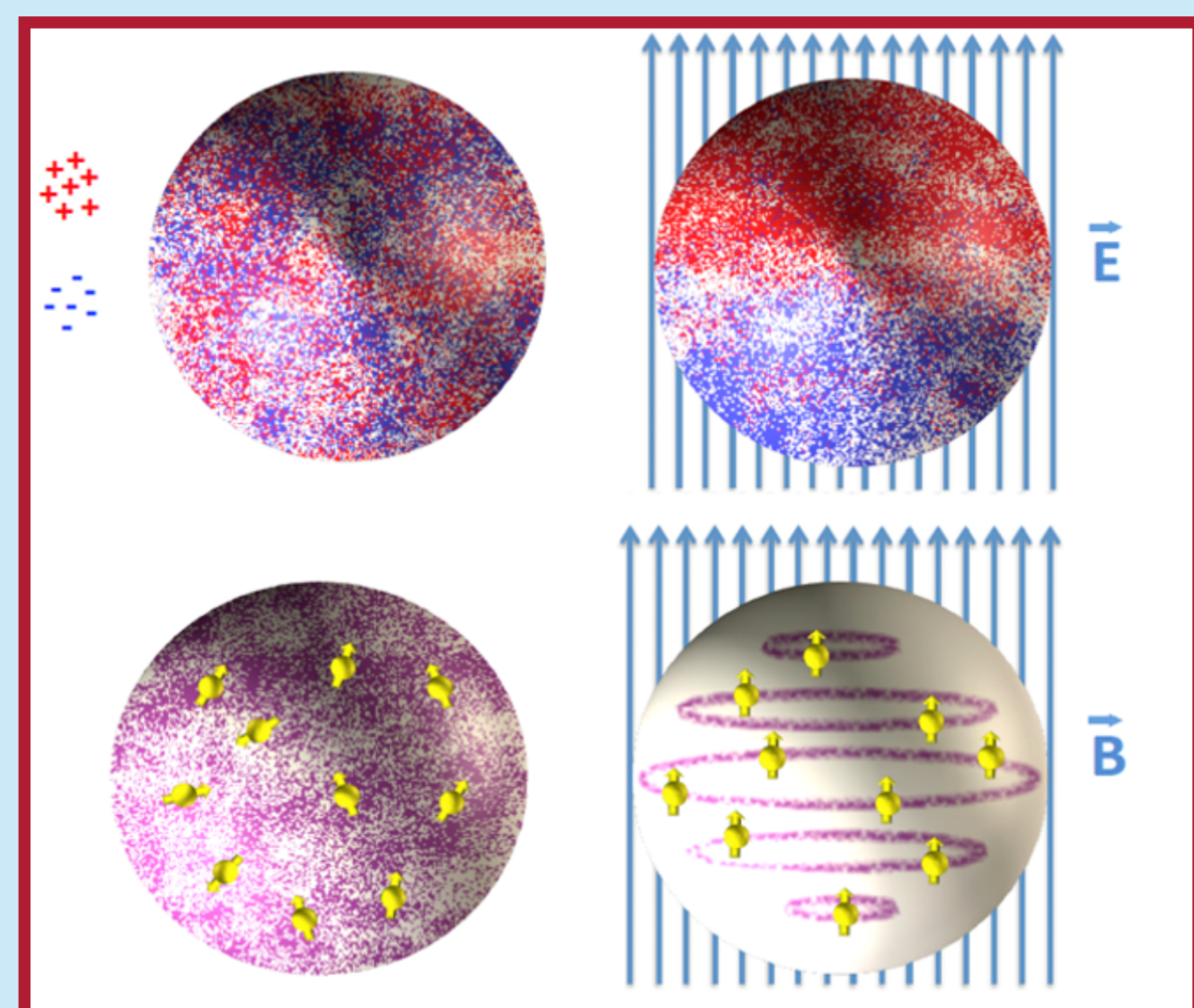
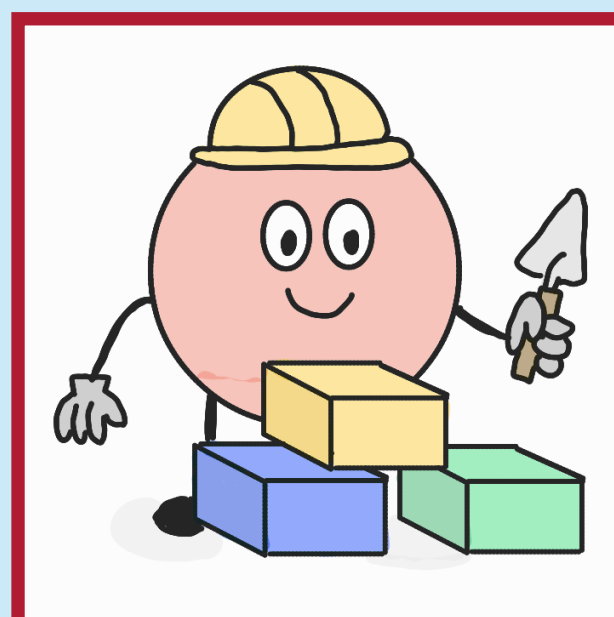
^aInstitut für Kernphysik, Johannes Gutenberg-Universität, Mainz

^bUniversità degli studi di Pavia, Pavia

The proton: what's still to learn?

The proton is the only stable composite particle in nature and the fundamental **building block** of matter

Studying the **proton's structure** is crucial for improving our understanding of the formation of matter through the strong interaction



Real Compton scattering (RCS)

Virtual Compton scattering (VCS)

$$\vec{D}_E \sim \alpha_{E1} \vec{E} \quad \vec{D}_E \sim \alpha_{E1}(Q^2) \vec{E}$$

$$\vec{D}_M \sim \beta_{M1} \vec{B} \quad \vec{D}_M \sim \beta_{M1}(Q^2) \vec{B}$$

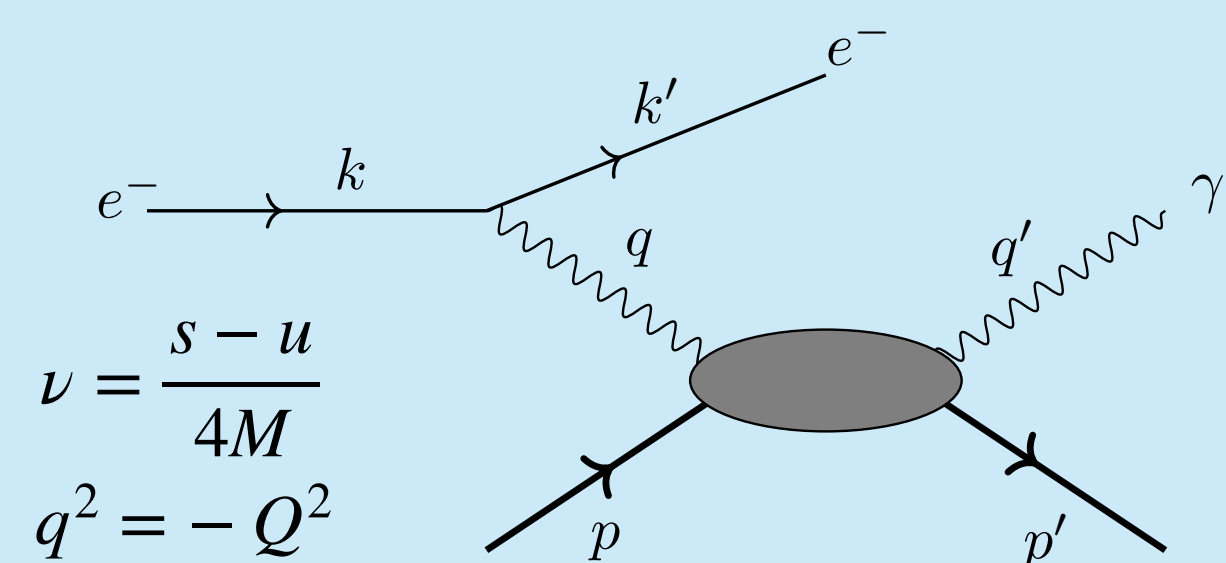
Q^2 : virtuality of the initial photon

An important property is its response to an external electromagnetic field, parametrized by static polarizabilities, α_{E1} and β_{M1} .

Generalized polarizabilities (GPs), $\alpha_{E1}(Q^2)$ and $\beta_{M1}(Q^2)$, map out the spatial distribution of the induced polarization.

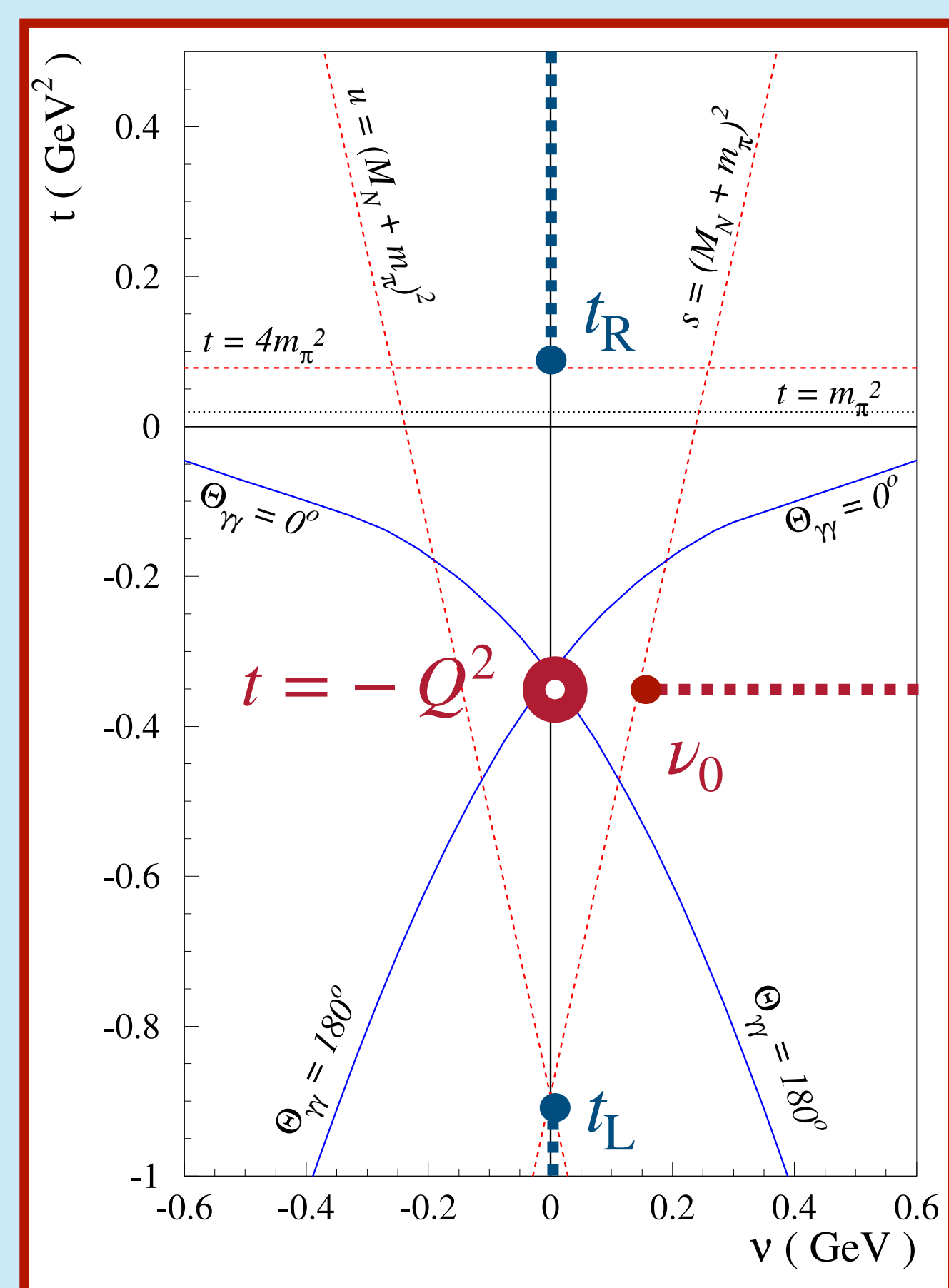
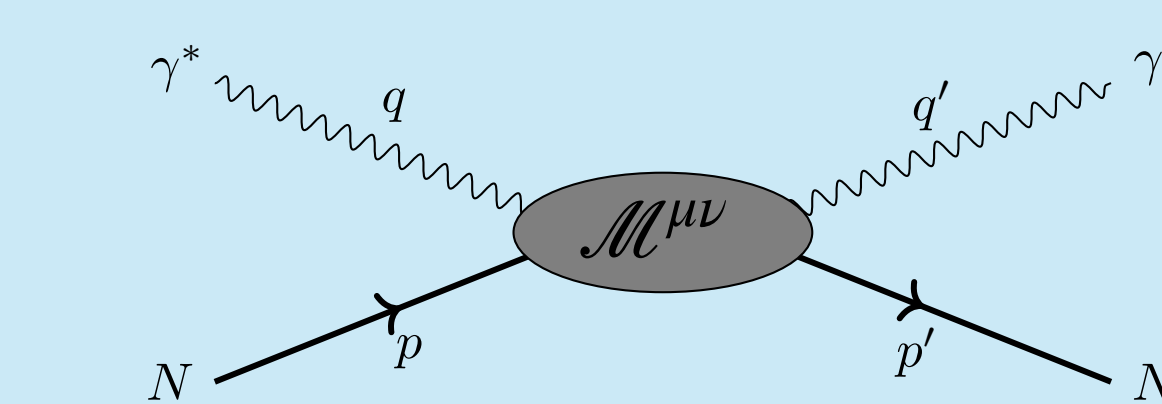
How can we extract GPs from VCS?

Through photon electroproduction off the proton, $ep \rightarrow ep\gamma$, we study the subprocess known as VCS ($\gamma^*p \rightarrow \gamma p$)



$$\mathcal{M}^{\mu\nu} = \sum_{i=1}^{12} F_i(Q^2, \nu, t) \tilde{\rho}_i^{\mu\nu} \rightarrow F_i(Q^2, \nu, t) = F_i(Q^2, -\nu, t), \quad \forall i$$

F_i scalar amplitudes



F_i are:
- free of kinematical singularities
- even in the variable ν

GPs are defined at the kinematical point $\nu = 0$ and $t = -Q^2$

$$F_i^{NB}(Q^2, 0, -Q^2)$$

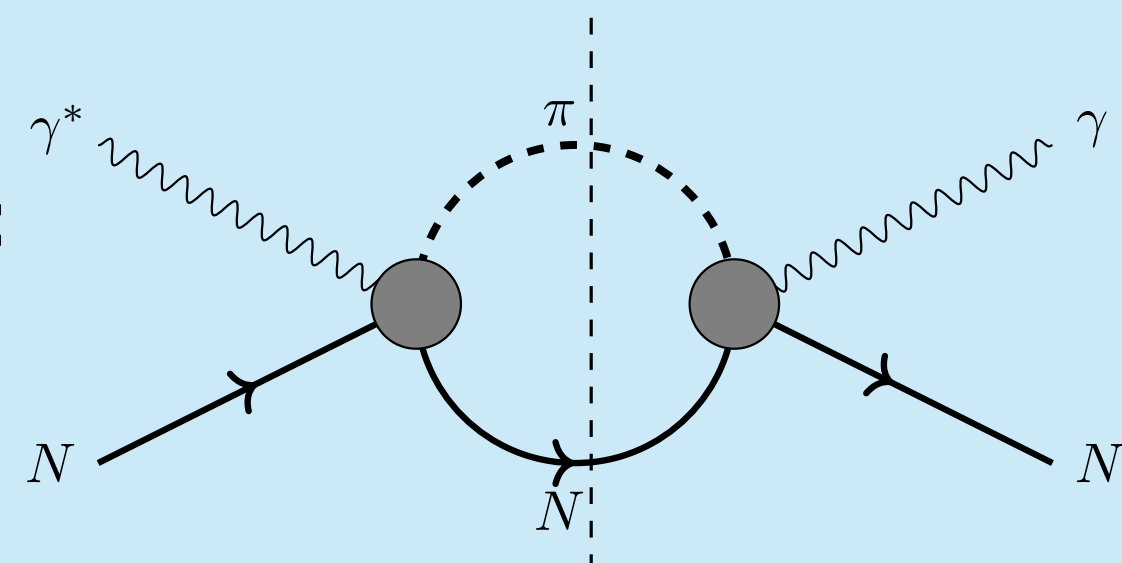
Assuming analyticity and appropriate high-energy behavior

We write unsubtracted **dispersion relations (DR)**

[B. Pasquini et al., EPJA 11, 185-208 (2001)]

$$\text{Re } F_i^{NB}(Q^2, \nu, t) = \frac{2}{\pi} \mathcal{P} \int_{\nu_0}^{\infty} d\nu' \frac{\nu' \text{Im}_s F_i(Q^2, \nu', t)}{\nu'^2 - \nu^2}$$

- **Low energies:** effective field theories
- **Intermediate region (~100 MeV to 1-2 GeV):** dominated by resonances, no complete theoretical description
- **High energies:** perturbative QCD



New way of extracting GPs

From Regge theory, F_1 and F_5 need a subtraction

Subtracted DR in ν at fixed t :

[B. Pasquini et al., Phys. Rev. C62 (2000)]

$$\text{Re } F_i^{NB}(Q^2, \nu, t) = F_i^{NB}(Q^2, 0, t) + \frac{2}{\pi} \nu^2 \mathcal{P} \int_{\nu_0}^{\infty} d\nu' \frac{\text{Im}_s F_i(Q^2, \nu', t)}{\nu'(\nu'^2 - \nu^2)}$$

"Price to pay" → subtraction function

$$F_i^{NB}(Q^2, 0, t) = F_i^{NB}(Q^2, 0, -Q^2) + \text{pole contribution}$$

GPs are free parameters

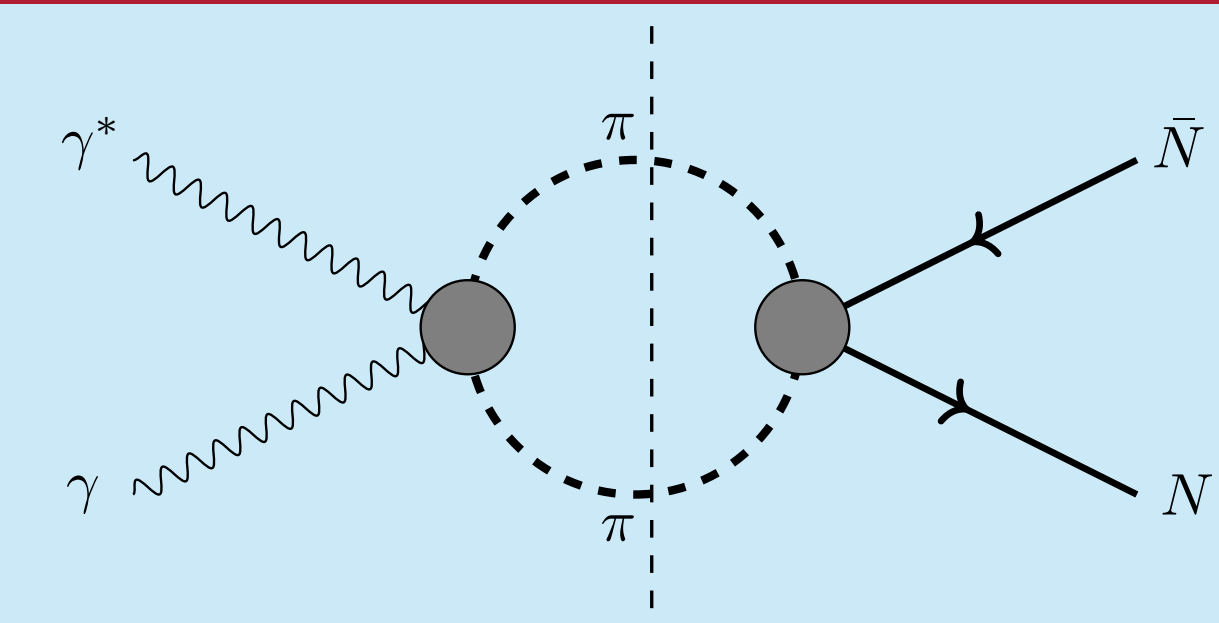
$$+ \frac{(t + Q^2)}{\pi} \mathcal{P} \int_{t_R}^{\infty} dt' \frac{\text{Im}_t F_i(Q^2, 0, t')}{(t' + Q^2)(t' - t)} \quad \text{:: right hand cut (RHC)}$$

$$+ \frac{(t + Q^2)}{\pi} \mathcal{P} \int_{-\infty}^{t_L} dt' \frac{\text{Im}_t F_i(Q^2, 0, t')}{(t' + Q^2)(t' - t)} \quad \text{:: left hand cut (LHC)}$$

Model dependence in the convergence of the integrals → reduced

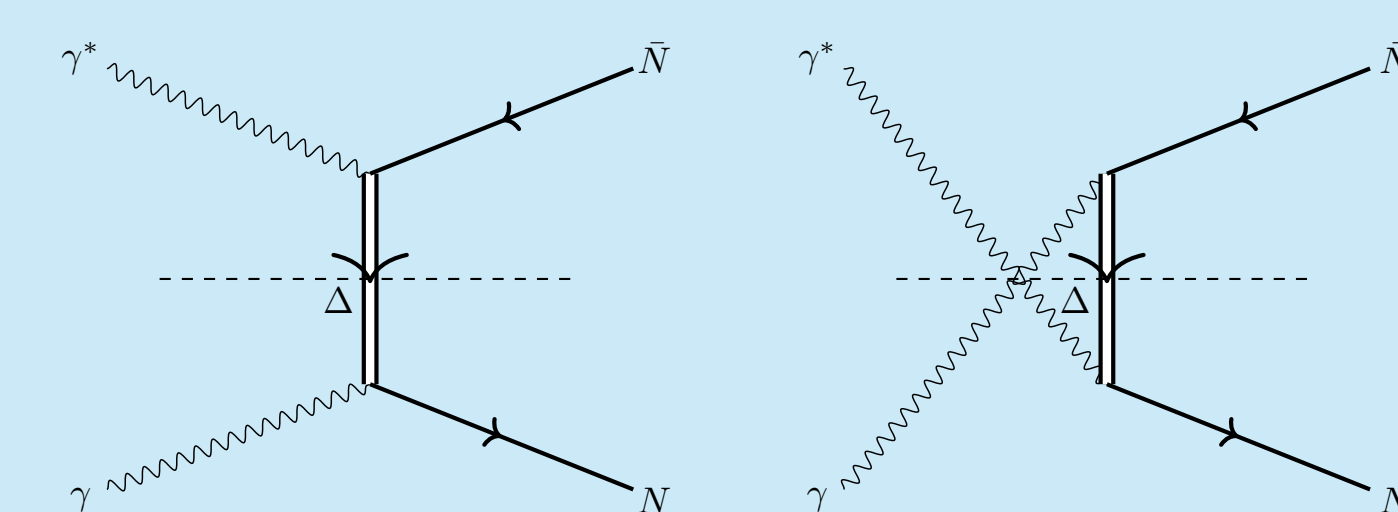
Evaluation of RHC and LHC

$$\text{RHC} :: \mathcal{P} \int_{t_R}^{\infty} dt' \frac{\text{Im}_t F_i(Q^2, 0, t')}{(t' - t)}$$



- 1) **Crossing symmetry:** integrand linked to t -channel VCS, $\gamma\gamma^* \rightarrow N\bar{N}$, helicity amplitudes
- 2) **Unitarity:** imaginary part expressed via the partial waves of the two subprocesses $\gamma\gamma^* \rightarrow \pi\pi$ and $\pi\pi \rightarrow N\bar{N}$

$$\text{LHC} :: \mathcal{P} \int_{-\infty}^{t_L} dt' \frac{\text{Im}_t F_i(Q^2, 0, t')}{(t' - t)}$$



Left-hand cut discontinuity modeled by the spectral function for the $\Delta(1232)$ -resonance excitation in s - and u - channel for the VCS process

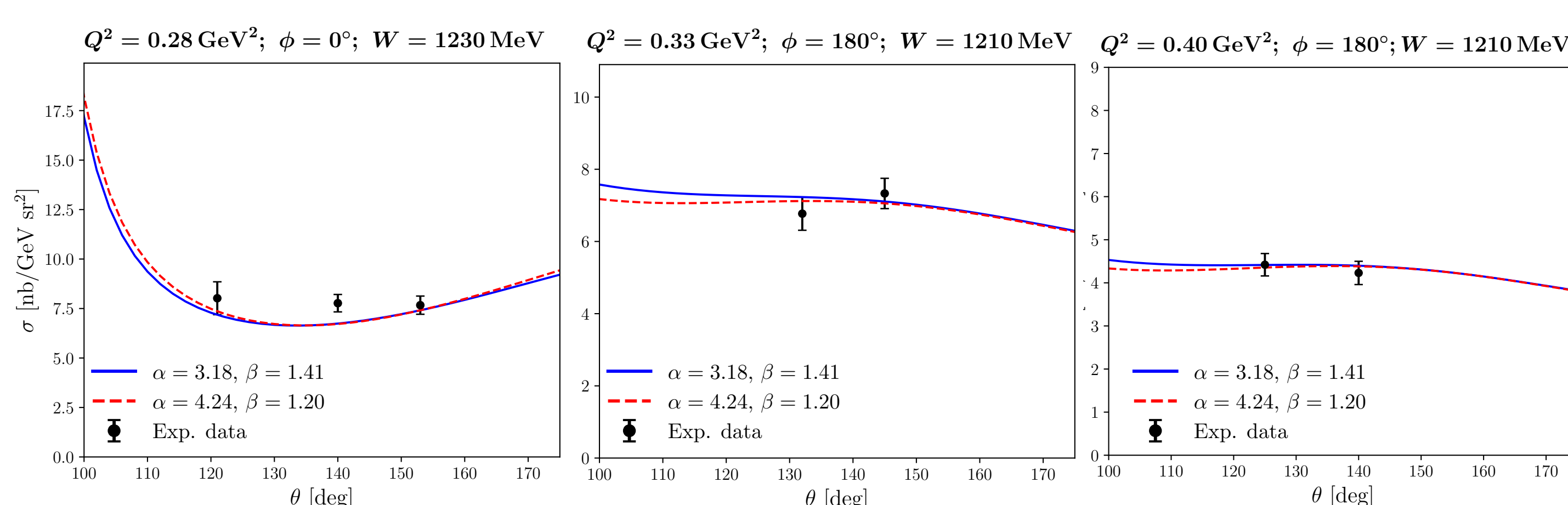
First results

We present the cross-section prediction as a function of the scattering angle θ at fixed Q^2, W and ϕ

W : total centre-of-mass energy of the (γ^*p) system

ϕ : azimuthal angle

At low energies, below $\Delta(1232)$ resonance, we have a good description of the data



In this approach the **GPs** are input parameters that need to be obtained from external sources; in the future they will be fitted to experimental data

Values of $\alpha_{E1}(Q^2)$ and $\beta_{M1}(Q^2)$ taken from [R. Li et al., Nature 611 (2022)]

Any further questions? Please contact me: ronchima@uni-mainz.de