

Classification of Modular UV Completions via cLFV observables

Adrián Moreno-Sánchez

with

A.Palavrić [2505.01535]



Junta de Andalucía

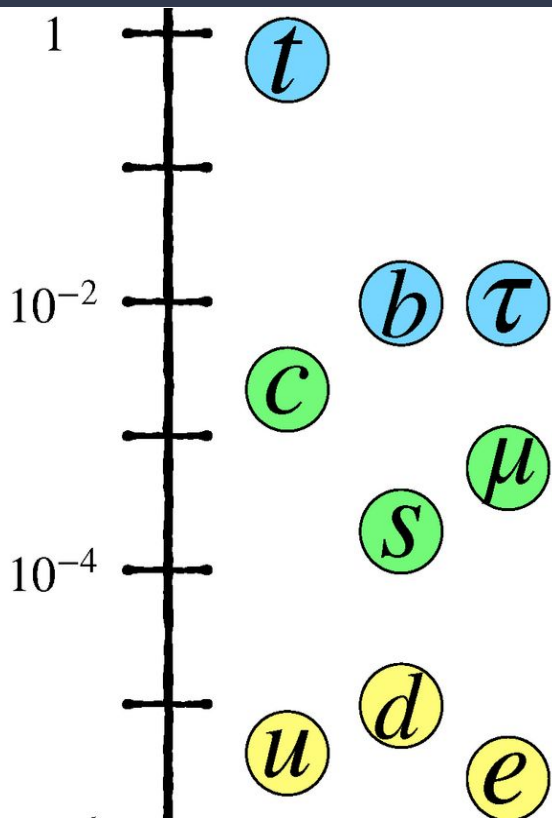
FTAE
High Energy Theory



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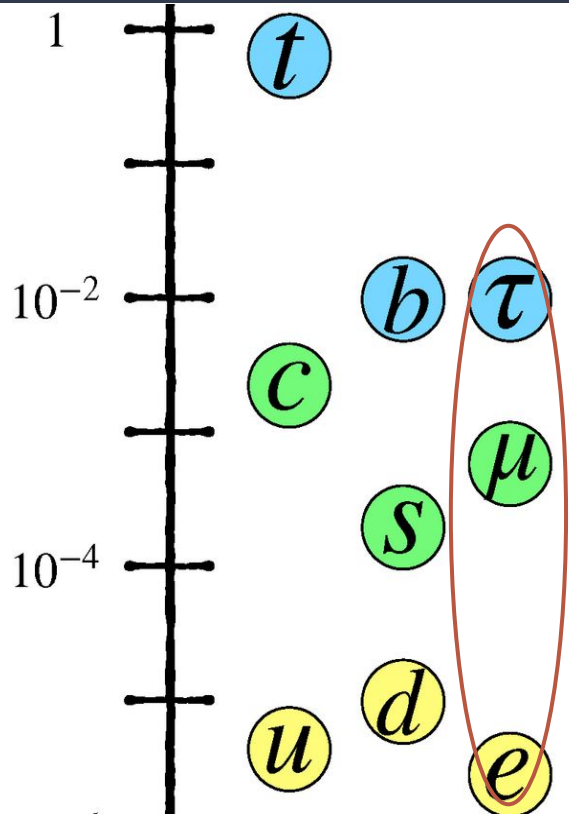
cLFV searches with Mu2e -LFN Workshop 2025

Introduction



- **Flavor Puzzle:** Three generations with no apparent explanation
- SM is flavor blind except for the yukawa sector
- Different UV explanations to reproduce the Yukawa matrix

Introduction



- **Flavor Puzzle:** Three generations with no apparent explanation
- SM is flavor blind except for the yukawa sector
- Different UV explanations to reproduce the Yukawa matrix
- This talk: just the **leptonic sector**

Effective Field Theories

Our best tool to parameterize extensions beyond the SM

- **Model-independent** agnostic framework
- Systematically capture low-energy effects of UV physics
- Constrain UV physics through **deviations** in our experimental measurements

Wilson
Coefficient

$$\mathcal{L}_{\text{SMEFT}} = \mathcal{L}_{\text{SM}} + \sum_{d \geq 5} \sum_i \frac{C_i}{\Lambda^{d-4}} \mathcal{O}_i^{(d)}$$

Flavor structure in SMEFT

Dimension 6: [Grzadkowski, Iskrzynski, Misiak, Rosiek]

- Single generation: 59 parameters
- Three generations: 2499 parameters

Proliferation arises from flavor: What if we assume flavor? [Greljo, Palavrić, Thomsen]

SMEFT $\mathcal{O}(1)$ terms (dim-6, $\Delta B = 0$)		Lepton sector															
		MFV _L		U(3) _V		U(2) ² × U(1) ²		U(2) ²		U(2) _V		U(1) ⁶		U(1) ³		No symm.	
Quark sector	MFV _Q	41	6	45	9	59	6	62	9	67	13	81	6	93	18	207	132
	U(2) ² × U(3) _d	72	10	78	15	95	10	100	15	107	21	122	10	140	28	281	169
	U(2) ³ × U(1) _{d3}	86	10	92	15	111	10	116	12	123	21	140	10	158	28	305	175
	U(2) ³	93	17	100	23	118	17	124	23	132	30	147	17	168	38	321	191
	No symmetry	703	570	734	600	756	591	786	621	818	652	813	612	906	705	1350	1149

Granada Dictionary

[de Blas, Criado, Pérez-Victoria, Santiago]

- All possible renormalizable **heavy** extensions in the leptonic sector
- Their matching to the SMEFT is automatized at tree-level and one-loop

UV Field	$-\mathcal{L}_{UV}^{(4)} \supset$
$S_1 \sim (\mathbf{1}, \mathbf{1})_1$	$[y_{S_1}]_{rij} S_{1r}^\dagger \bar{\ell}_i \sigma_2 \ell_j^c$
$S_2 \sim (\mathbf{1}, \mathbf{1})_2$	$[y_{S_2}]_{rij} S_{2r}^\dagger \bar{e}_i e_j^c$
$\varphi \sim (\mathbf{1}, \mathbf{2})_{\frac{1}{2}}$	$[y_\varphi]_{rij} \varphi_r \bar{\ell}_i e_j$
$\Xi_1 \sim (\mathbf{1}, \mathbf{3})_1$	$[y_{\Xi_1}]_{rij} \Xi_{1r}^{a\dagger} \bar{\ell}_i \sigma^a i \sigma_2 \ell_j^c$
$N \sim (\mathbf{1}, \mathbf{1})_0$	$[\lambda_N]_{ri} \bar{N}_{R,r} \tilde{\phi}^\dagger \ell_i$
$E \sim (\mathbf{1}, \mathbf{1})_{-1}$	$[\lambda_E]_{ri} \bar{E}_{R,r} \phi^\dagger \ell_i$
$\Delta_1 \sim (\mathbf{1}, \mathbf{2})_{-\frac{1}{2}}$	$[\lambda_{\Delta_1}]_{ri} \bar{\Delta}_{1L,r} \phi e_i$
$\Delta_3 \sim (\mathbf{1}, \mathbf{2})_{-\frac{3}{2}}$	$[\lambda_{\Delta_3}]_{ri} \bar{\Delta}_{3L,r} \tilde{\phi} e_i$
$\Sigma \sim (\mathbf{1}, \mathbf{3})_0$	$\frac{1}{2} [\lambda_\Sigma]_{ri} \bar{\Sigma}_{R,r}^a \tilde{\phi}^\dagger \sigma^a \ell_i$
$\Sigma_1 \sim (\mathbf{1}, \mathbf{3})_{-1}$	$\frac{1}{2} [\lambda_{\Sigma_1}]_{ri} \bar{\Sigma}_{1R,r}^a \phi^\dagger \sigma^a \ell_i$
$\mathcal{B} \sim (\mathbf{1}, \mathbf{1})_0$	$[g_{\mathcal{B}}^\ell]_{rij} \mathcal{B}_r^\mu \bar{\ell}_i \gamma_\mu \ell_j + [g_{\mathcal{B}}^e]_{rij} \mathcal{B}_r^\mu \bar{e}_i \gamma_\mu e_j$
$\mathcal{W} \sim (\mathbf{1}, \mathbf{3})_0$	$\frac{1}{2} [g_{\mathcal{W}}^\ell]_{rij} \mathcal{W}_r^{\mu a} \bar{\ell}_i \sigma^a \gamma_\mu \ell_j$
$\mathcal{L}_3 \sim (\mathbf{1}, \mathbf{2})_{-\frac{3}{2}}$	$[g_{\mathcal{L}_3}]_{rij} \mathcal{L}_{3r}^{\mu\dagger} \bar{e}_i^c \gamma_\mu \ell_j$

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- All possible renormalizable **heavy** extensions in the leptonic sector
- Their matching to the SMEFT is automatized at tree-level and one-loop
- **SMEFT directions**: Classification with a flavor assumption

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Discrete Symmetries

Broken flavor symmetry: A_4 [Ferreira, Morisi, Peinado, Holthausen,...]

- The leptonic sector can be very well accommodated in

$$\ell \equiv (\ell_1, \ell_2, \ell_3)^T \sim \mathbf{3},$$

$$(e_1, e_2, e_3) \sim (\mathbf{1}, \mathbf{1}', \mathbf{1}'')$$

- Yukawa interactions can be very well reproduced

Tensor Decomposition

$$\mathbf{1} \otimes \mathbf{1} = \mathbf{1}, \quad \mathbf{1}'' \otimes \mathbf{1}'' = \mathbf{1}',$$

$$\mathbf{1}' \otimes \mathbf{1}' = \mathbf{1}'', \quad \mathbf{1}' \otimes \mathbf{1}'' = \mathbf{1}.$$

$$\mathbf{3} \otimes \mathbf{3} = \mathbf{1} \oplus \mathbf{1}' \oplus \mathbf{1}'' \oplus \mathbf{3}_S \oplus \mathbf{3}_A$$

Modular Symmetries

[Feruglio, Criado , Asaka, Yoshida,...]

Breaking of A_4 dynamically controlled by VEV of a **modulus field** that transforms under the modulus group

$$\tau \rightarrow \gamma\tau = \frac{a\tau + b}{c\tau + d}$$

- Fields are modular forms

$$f(\tau) \rightarrow (c\tau + d)^k U f(\tau), \quad U \in \Gamma_3 \quad \begin{array}{l} k_\ell = 2 \\ k_e = 0 \end{array}$$

- Yukawa couplings too...

$$\begin{bmatrix} Y_1(\tau) \\ Y_2(\tau) \\ Y_3(\tau) \end{bmatrix} = \begin{bmatrix} 1 + 12q + 36q^2 + 12q^3 + \dots \\ -6q^{1/3}(1 + 7q + 8q^2 + \dots) \\ -18q^{2/3}(1 + 2q + 5q^2 + \dots) \end{bmatrix}, \quad q = e^{2i\pi\tau}$$

[Kobayashi,
Otsuka,
Tanimoto,
Yamamoto]

Modular Symmetries [Feruglio, Criado , Asaka, Yoshida,...]

Breaking of A4 dynamically controlled by VEV of a **modulus field** that transforms under the modulus group

$$\tau \rightarrow \gamma\tau = \frac{a\tau + b}{c\tau + d}$$

Yukawa couplings explained just with one parameter: τ

- Fields are modular forms

$$f(\tau) \rightarrow (c\tau + d)^k U f(\tau), \quad U \in \Gamma_3 \quad \begin{array}{l} k_l = 2 \\ k_e = 0 \end{array}$$

- Yukawa couplings too...

$$\begin{bmatrix} Y_1(\tau) \\ Y_2(\tau) \\ Y_3(\tau) \end{bmatrix} = \begin{bmatrix} 1 + 12q + 36q^2 + 12q^3 + \dots \\ -6q^{1/3}(1 + 7q + 8q^2 + \dots) \\ -18q^{2/3}(1 + 2q + 5q^2 + \dots) \end{bmatrix}, \quad q = e^{2i\pi\tau}$$

[Kobayashi,
Otsuka,
Tanimoto,
Yamamoto]

- Consider the UV interaction

$$[y_{\mathcal{S}_1}]_{rij} \mathcal{S}_{1r}^\dagger \bar{\ell}_i i \sigma_2 \ell_j^c$$

- Extract the flavor tensor

$$[y_{\mathcal{S}_1}]_{rij} = \frac{y_{\mathcal{S}_1}^{(1)}}{2} \left[\delta_{r1} (\delta_{i2} \delta_{j3} - \delta_{i3} \delta_{j2}) + \delta_{r2} (\delta_{i3} \delta_{j1} - \delta_{i1} \delta_{j3}) + \delta_{r3} (\delta_{i1} \delta_{j2} - \delta_{i2} \delta_{j1}) \right],$$

- Perform the matching to the SMEFT

$$\frac{1}{M_{\mathcal{S}_1}^2} [y_{\mathcal{S}_1}]_{rjl}^* [y_{\mathcal{S}_1}]_{rik} [\mathcal{O}_{ll}]_{ijkl}$$

- Compute the experimental bounds

$$y_{\mathcal{S}_1} \sim 1$$

Low Energy observables: Global Fit [Falkowski et al.]

Flavor Violating

- Radiative Decays [MEG, BaBar, Belle]
- 3 body leptonic decays [SINDRUM, Belle, Mu3e]
- Muon-to-electron conversion [SINDRUM, Mu2e]

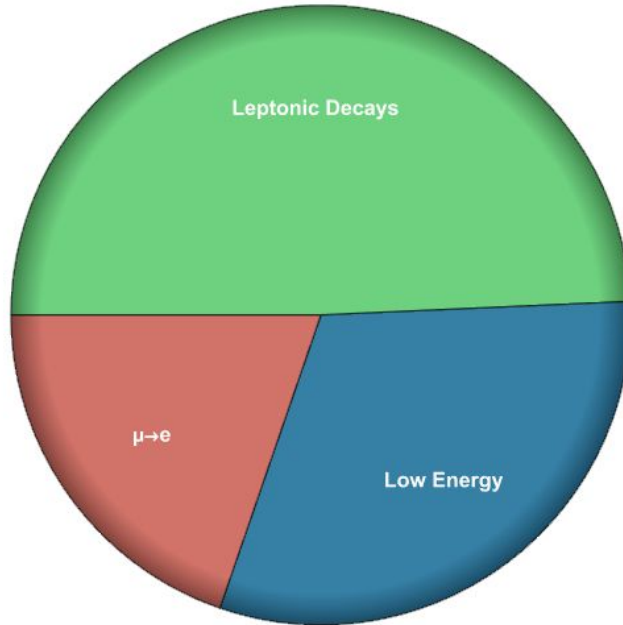
$$\text{BR}(\mu \rightarrow e, \gamma)$$

$$\text{BR}(\ell_i \rightarrow \ell_j \ell_k \ell_l)$$

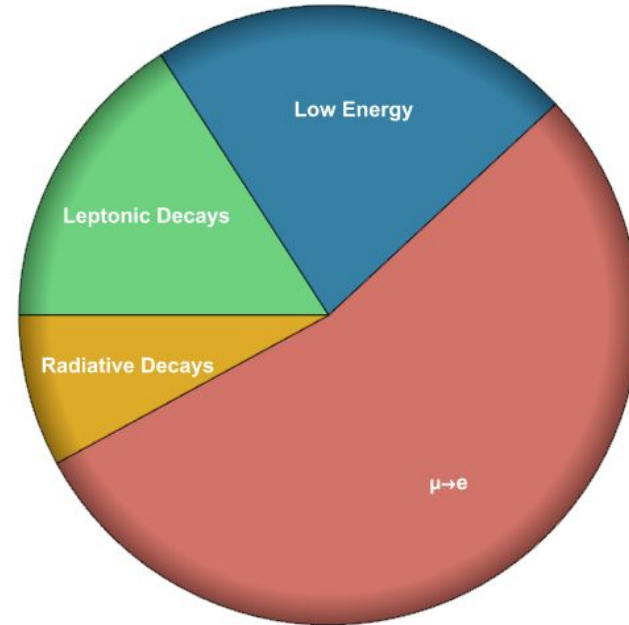
$$CR(\mu \rightarrow e, N)$$

Radiative Corrections

Tree Level



One Loop

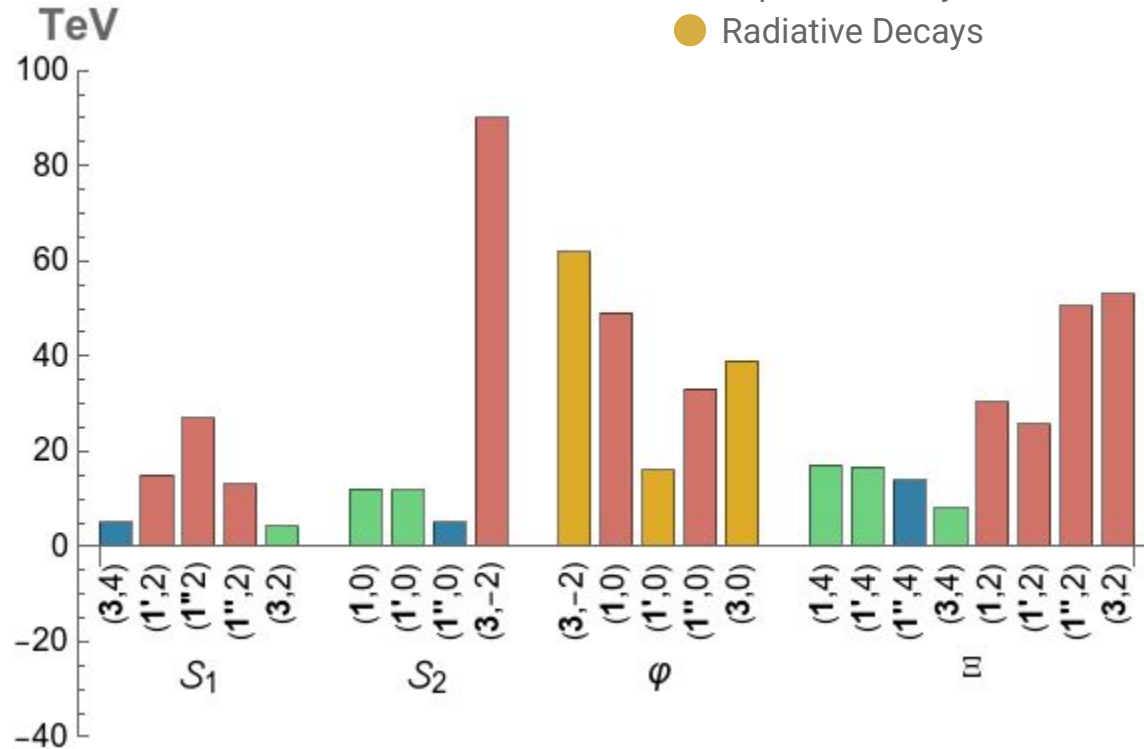


Scalars



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- Low Energy Observables
- Nuclei Conversion
- Leptonic Decays
- Radiative Decays

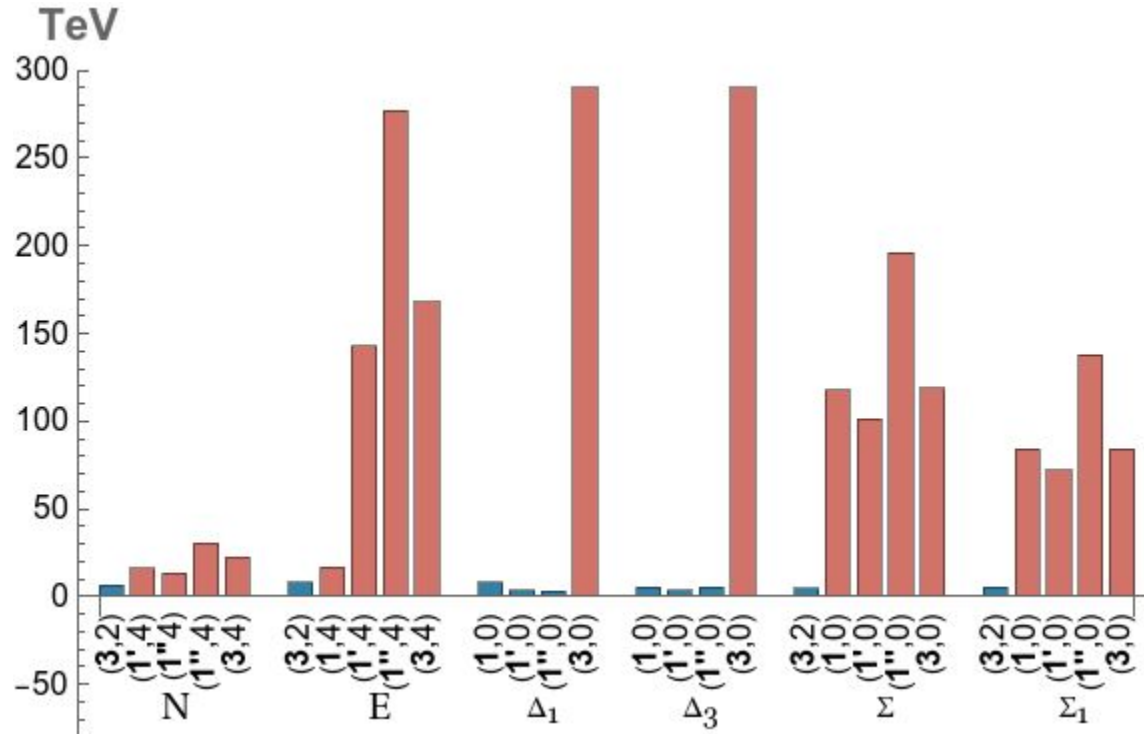


Fermions



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- Nuclei Conversion

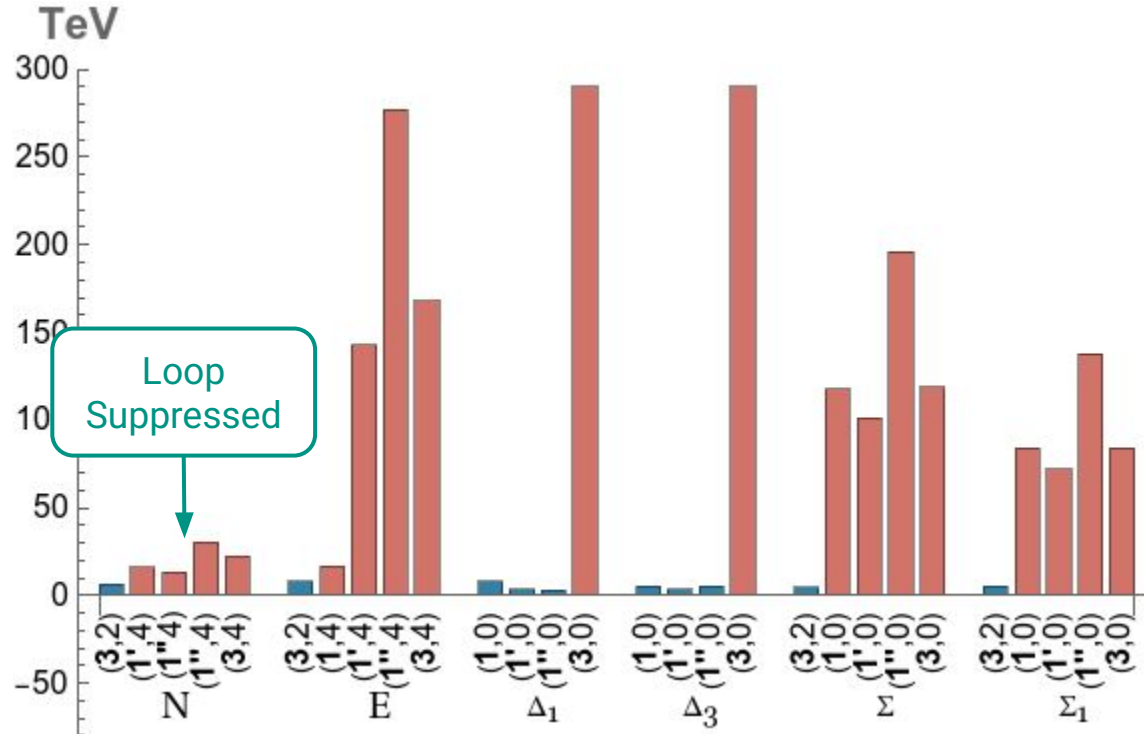


Fermions



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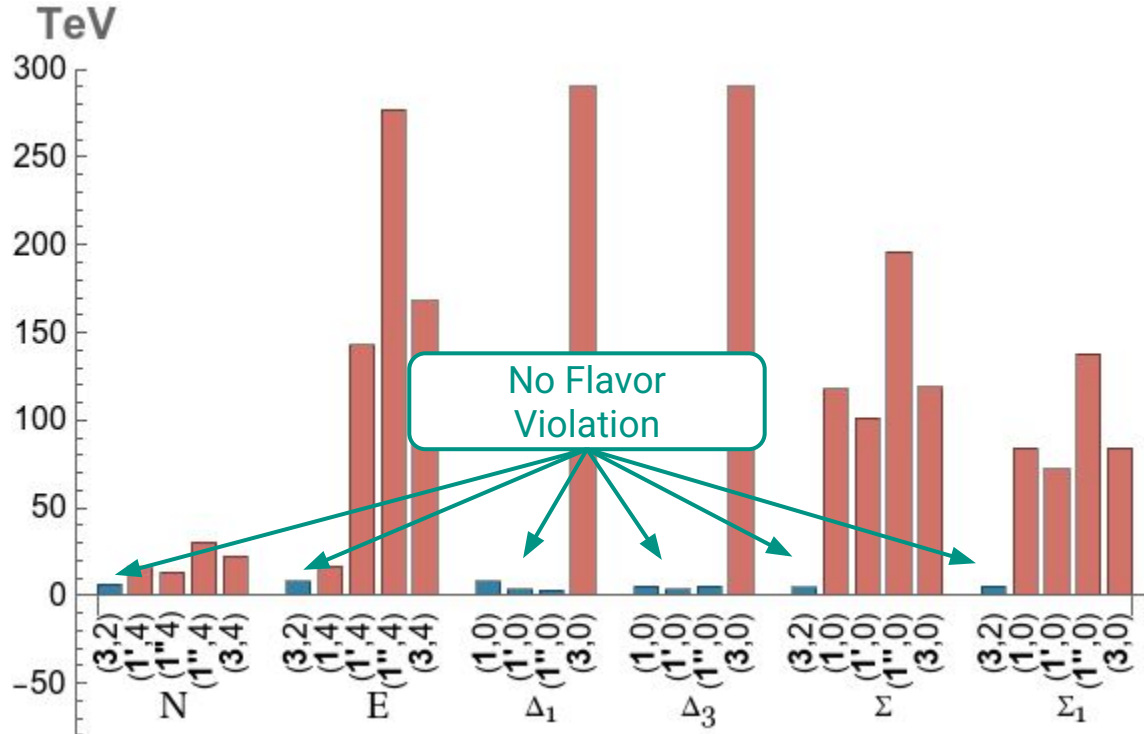
- Low Energy Observables
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Fermions



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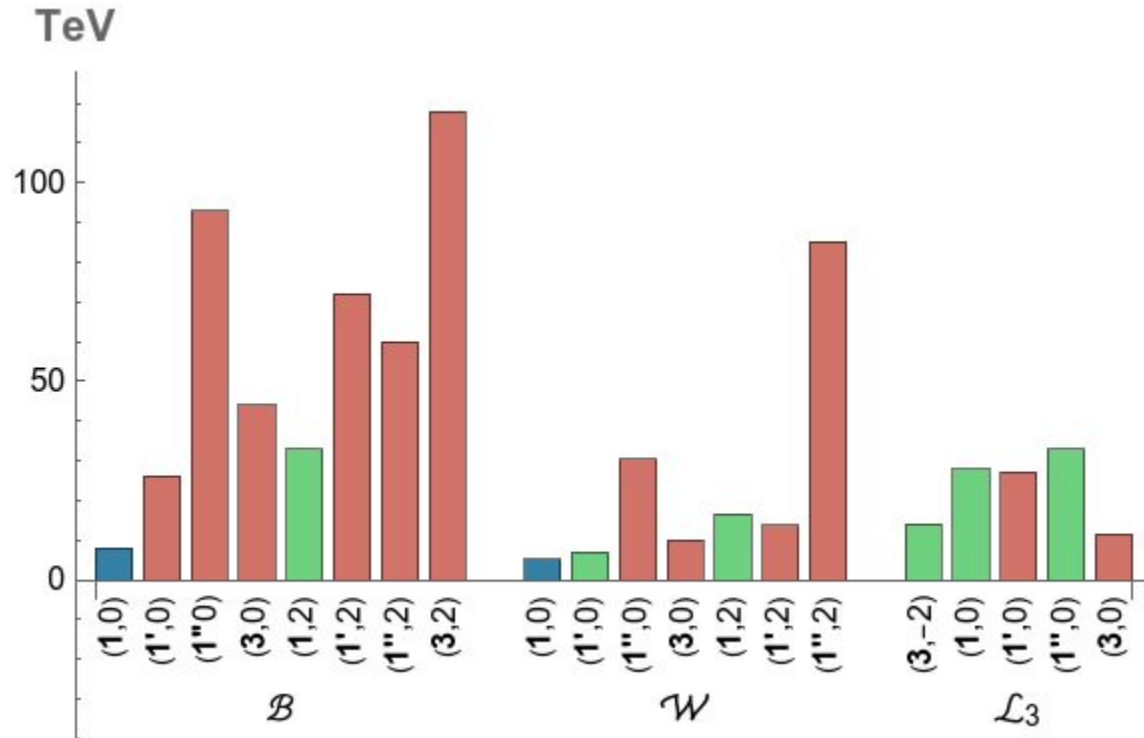


Vectors



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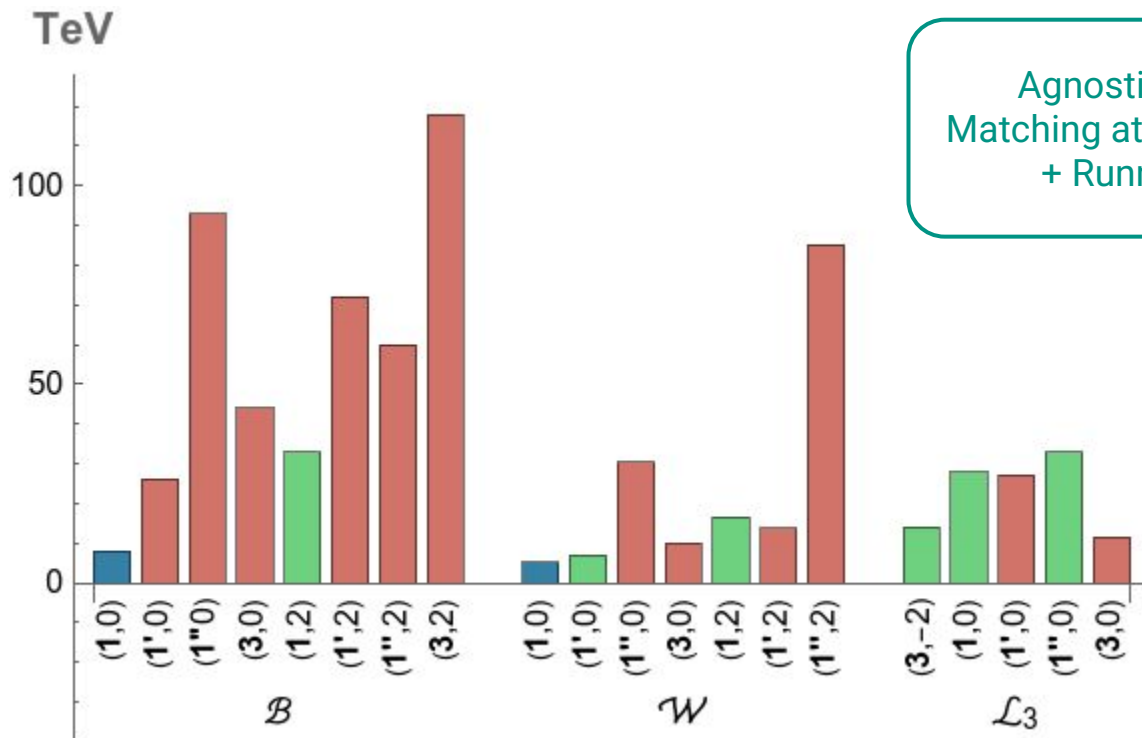
- Low Energy Observables
- Nuclei Conversion
- Leptonic Decays



Vectors



- Low Energy Observables
- Nuclei Conversion
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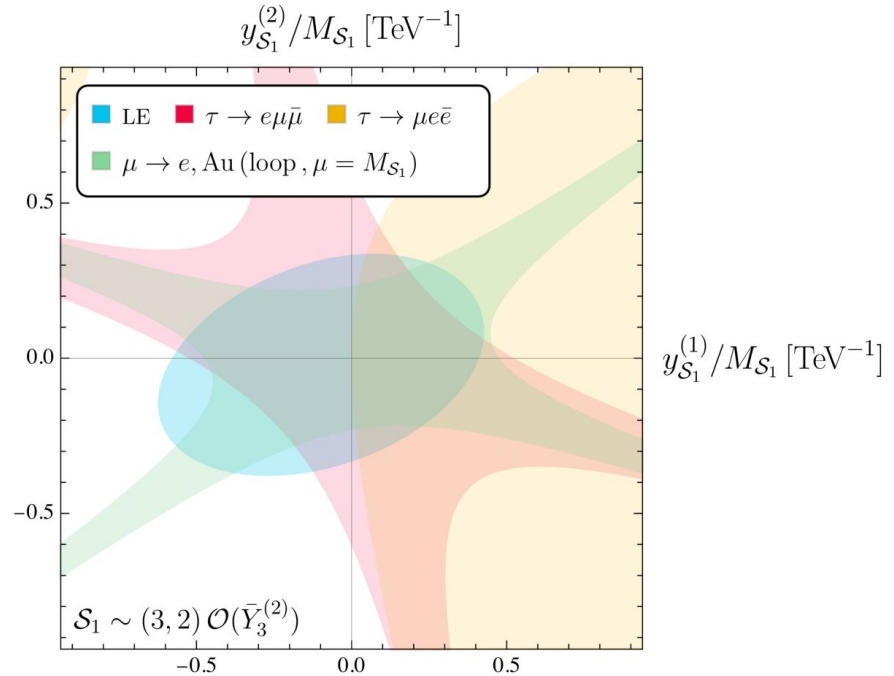


Agnostic UV:
Matching at tree-level
+ Running

Multi-Parameter

Flavor tensor with various different couplings

- Combined fit with correlations
- Projected two-dimensional likelihoods
- Blind lines



To take home

- Modular symmetries are a promising explanation for flavor
- UV modular physics is very constrained by cFLV
- Radiative correction are relevant
- $CR(\mu \rightarrow e, N)$ as the most relevant observable

Grazie mille!

