



From Flavor Bounds to Charming Higgs Discoveries

— 07 / 05 / 2025 —

MAURO VALLI

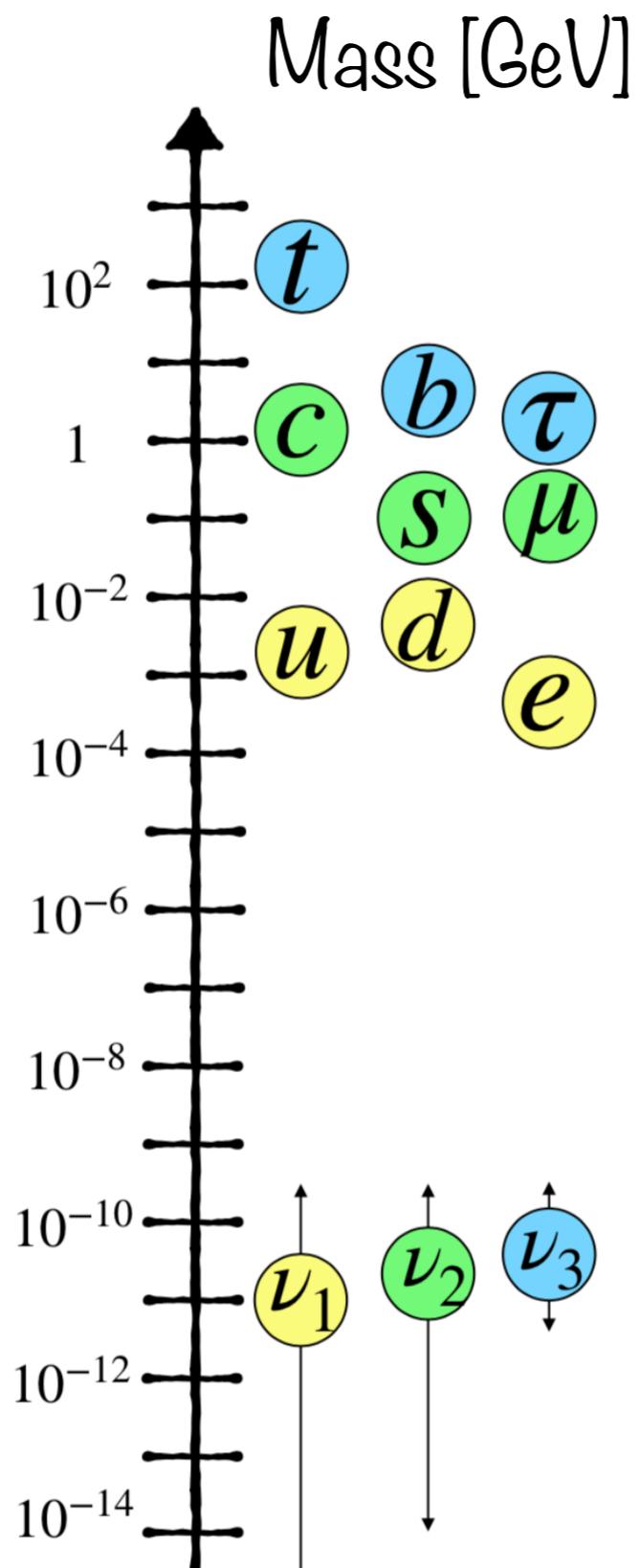
INFN Rome



MANY THANKS TO UTfit Collaboration + A.Giannakopoulou & P.Meade

The Standard Model

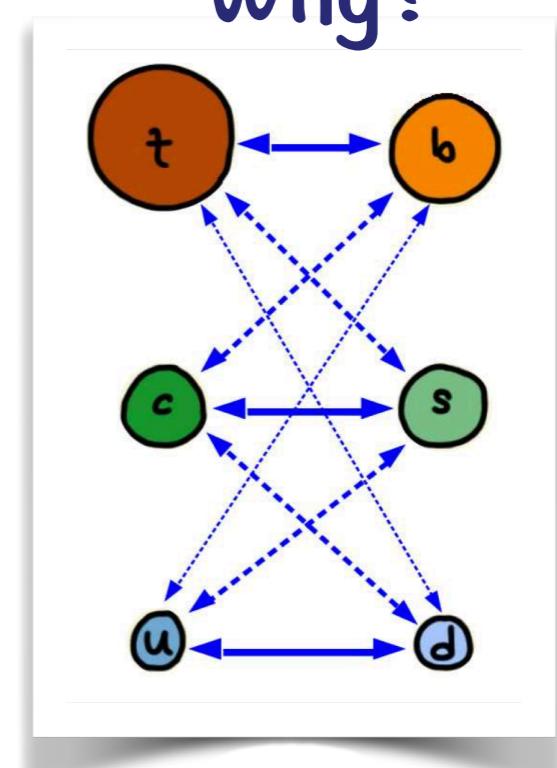
THE SM
FLAVOR
PUZZLE



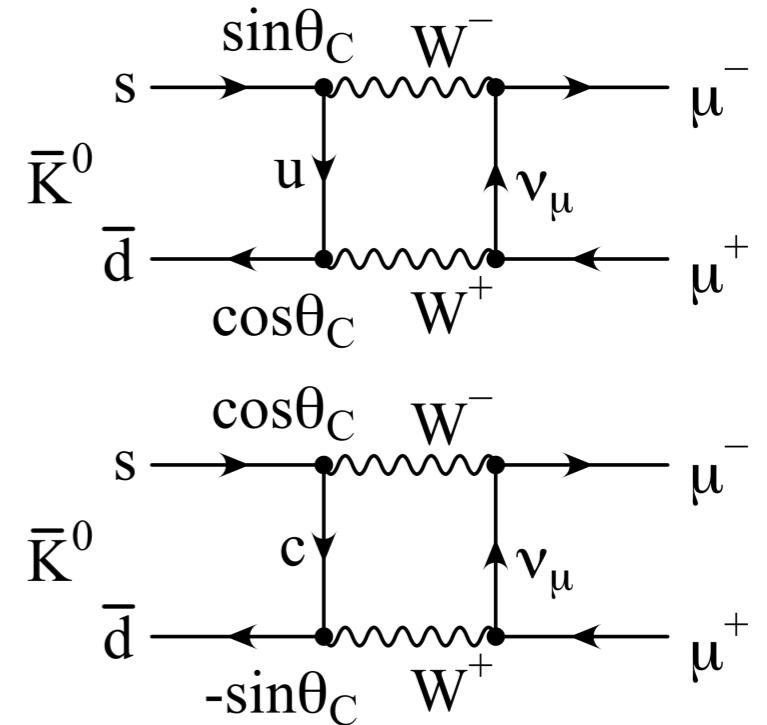
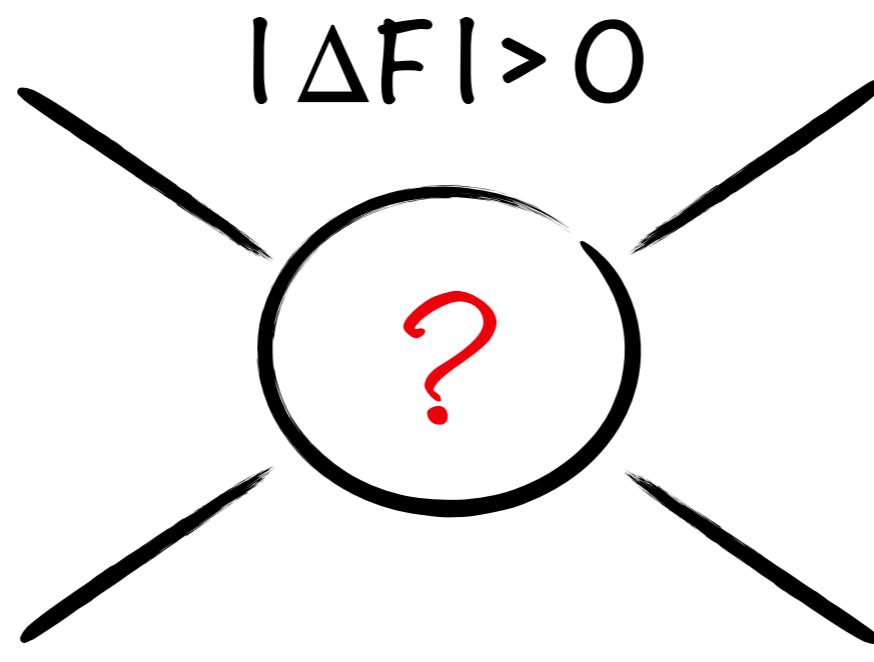
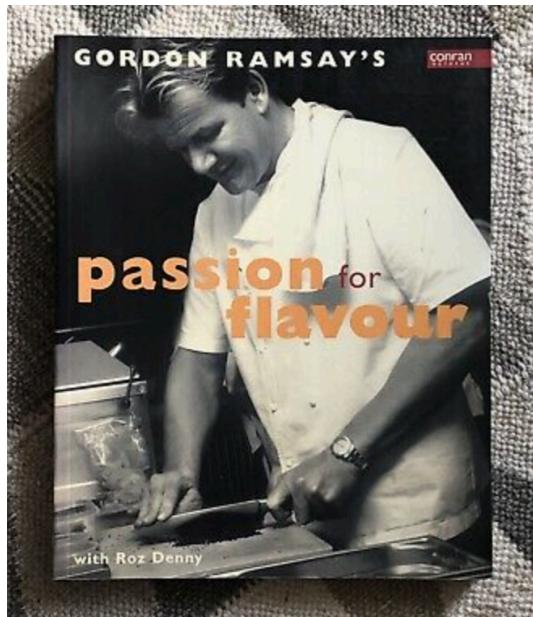
$$V_{\text{CKM}} = \begin{pmatrix} V_{ud} & V_{us} & V_{ub} \\ V_{cd} & V_{cs} & V_{cb} \\ V_{td} & V_{ts} & V_{tb} \end{pmatrix}$$

u
c
t
d
s
b

Why?



Precision Tests: Flavor



It led to “New Physics” (NP) !



<https://agenda.infn.it/event/41258>

- ▶ The 70th anniversary of CERN
 - ▶ The 50th anniversary of J/Ψ
- which marked the transition from
- MANY MODELS** to the **STANDARD THEORY**

2024

J.Iliopoulos @ The Rise of Particle Physics

Flavor Metrology:

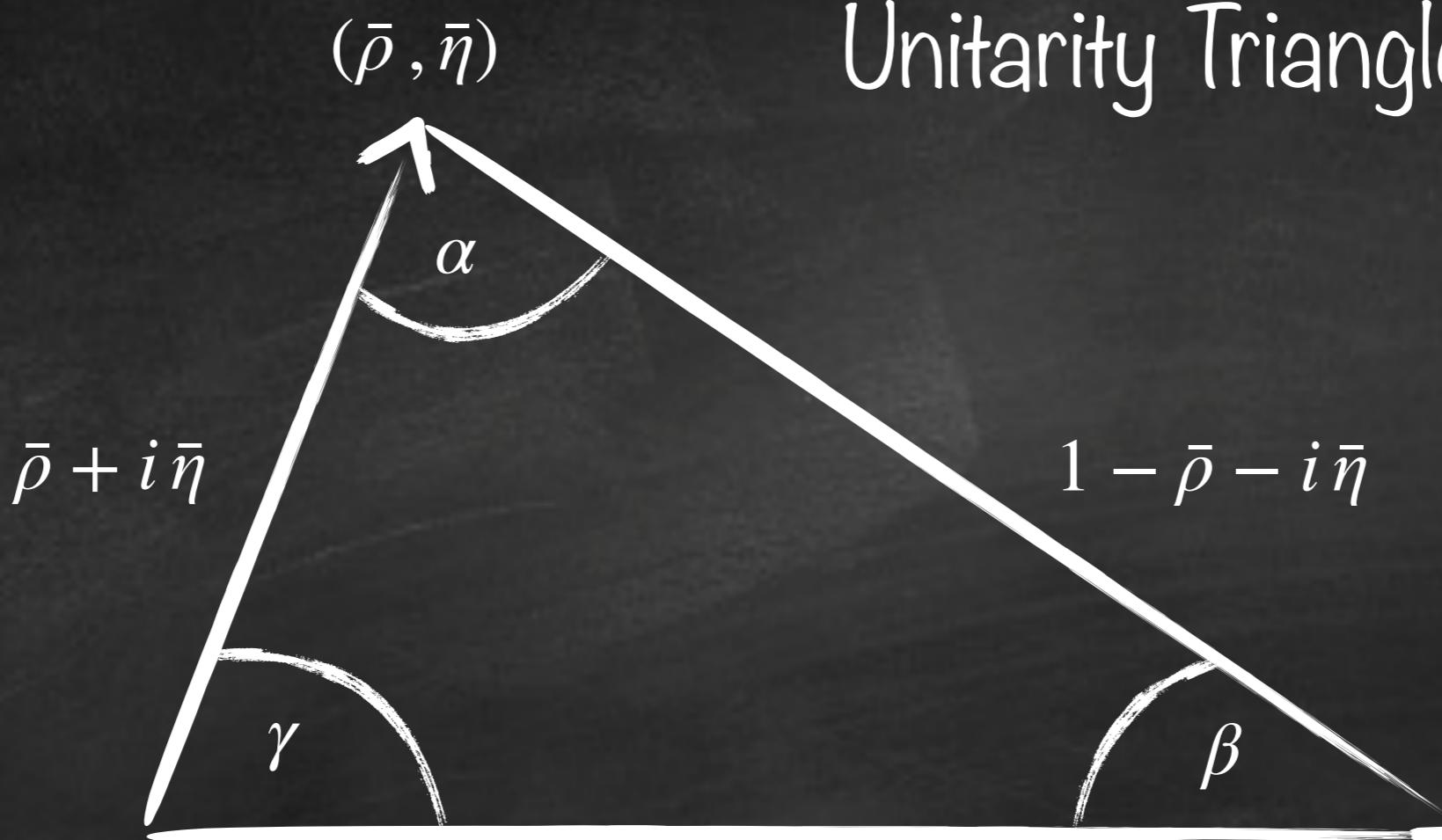


- Flavor violation in SM in charged weak-current $\longleftrightarrow V_{CKM}$
→ Flavor Changing Neutral Currents (FCNCs) **ONLY** @ one loop
- CKM matrix described by 4 params (3 angles + CP phase)

$$V_{CKM} = \begin{pmatrix} 1 - \lambda^2/2 & \lambda & A\lambda^3(\bar{\rho} - i\bar{\eta}) \\ -\lambda & 1 - \lambda^2/2 & A\lambda^2 \\ A\lambda^3(1 - \bar{\rho} - i\bar{\eta}) & -A\lambda^2 & 1 \end{pmatrix} + \mathcal{O}(\lambda^4)$$

$(\bar{\rho}, \bar{\eta})$ apex of $V_{ub}^* V_{ud} + V_{cb}^* V_{cd} + V_{tb}^* V_{td} = 0$

Unitarity Triangle (UT)



$$-\frac{V_{ud}V_{ub}^*}{V_{cd}V_{cb}^*} - \frac{V_{td}V_{tb}^*}{V_{cd}V_{cb}^*} = R_b e^{i\gamma} + R_t e^{-i\beta} = 1 \simeq (\bar{\rho} + i\bar{\eta}) + (1 - \bar{\rho} - i\bar{\eta})$$



www.utfit.org



M.Bona, M. Ciuchini, D. Derkach, F. Ferrari, E. Franco,
V. Lubicz, G. Martinelli, M. Pierini, L. Silvestrini, C.
Tarantino, V. Vagnoni, M. Valli, and L.Vittorio

— EXP
— TH



LINCEI CELEBRATIVE ESSAYS

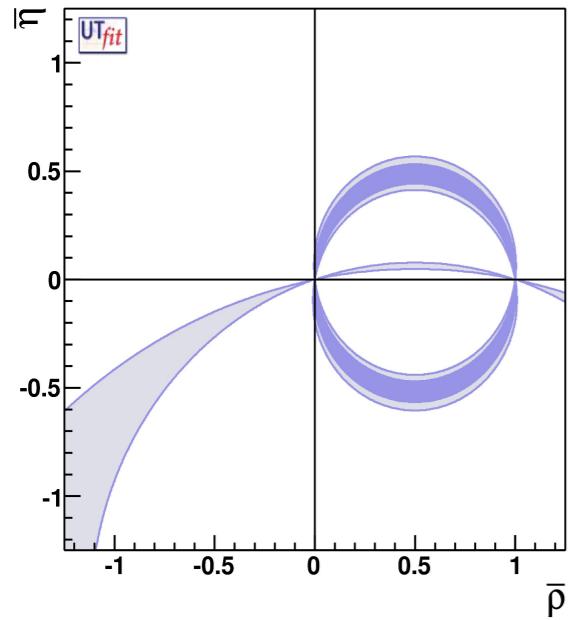
New **UTfit** analysis of the unitarity triangle
in the Cabibbo–Kobayashi–Maskawa scheme

arXiv: [2212.03894](https://arxiv.org/abs/2212.03894) — Rend.Lincei Sci.Fis.Nat. 34 (2023) 37-57

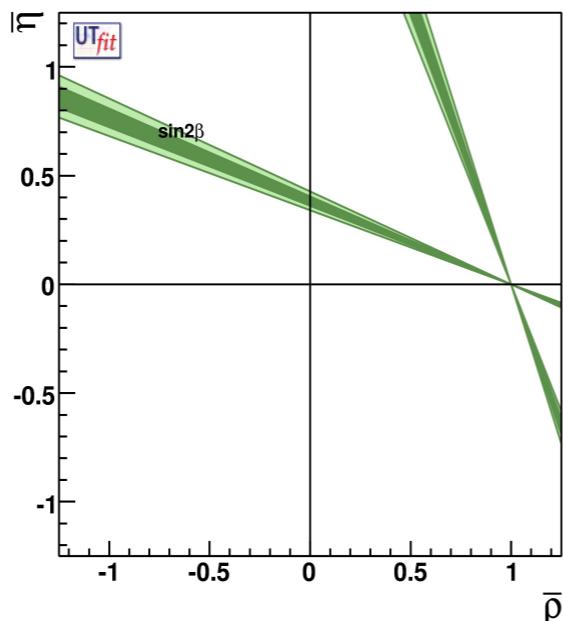
The Power of Redundancy

see, e.g., ***Les Houches Lect.Notes 108 (2020) - L.Silvestrini***

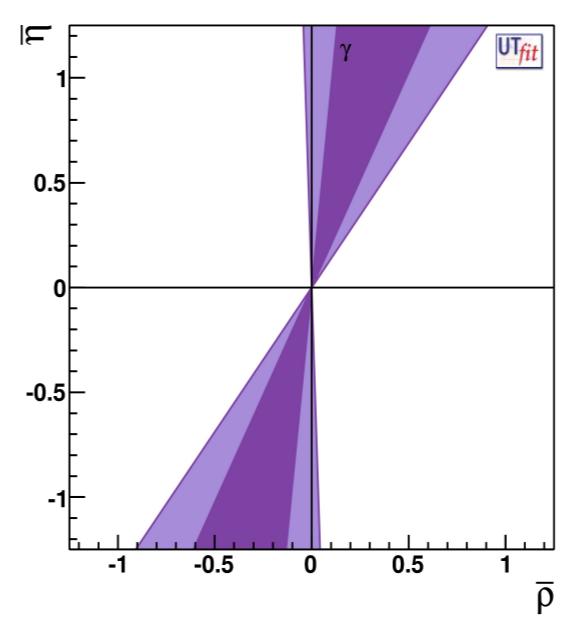
α ($B \rightarrow \pi\pi, \rho\rho$)



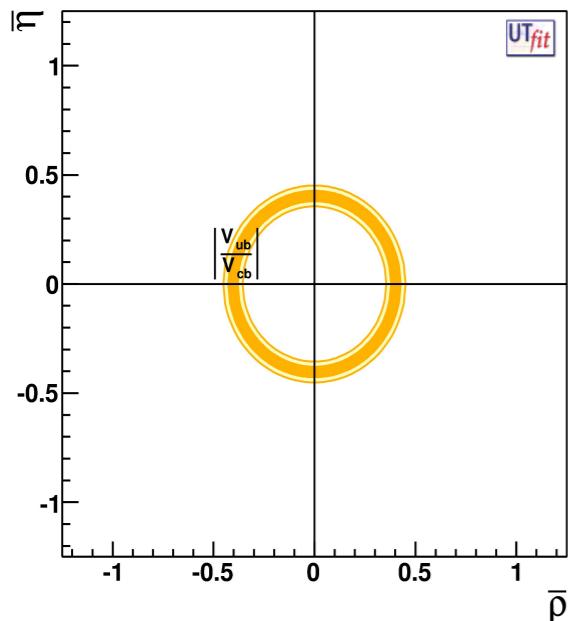
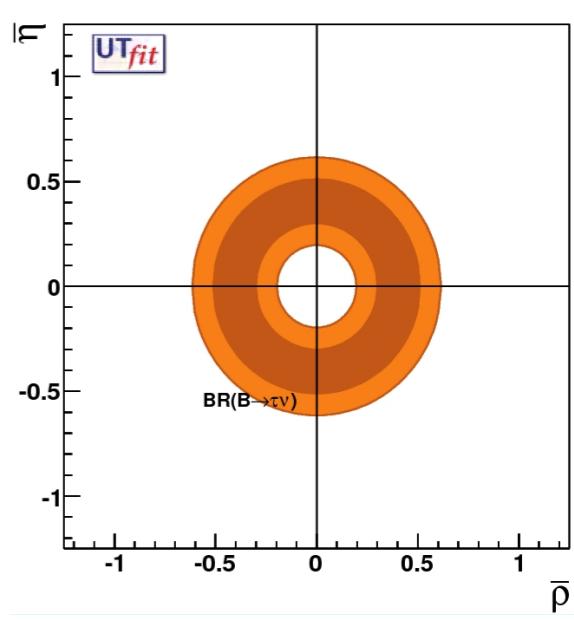
β ($B \rightarrow J/\psi K^{(*)}$)



γ ($B \rightarrow D^{(*)} K$)

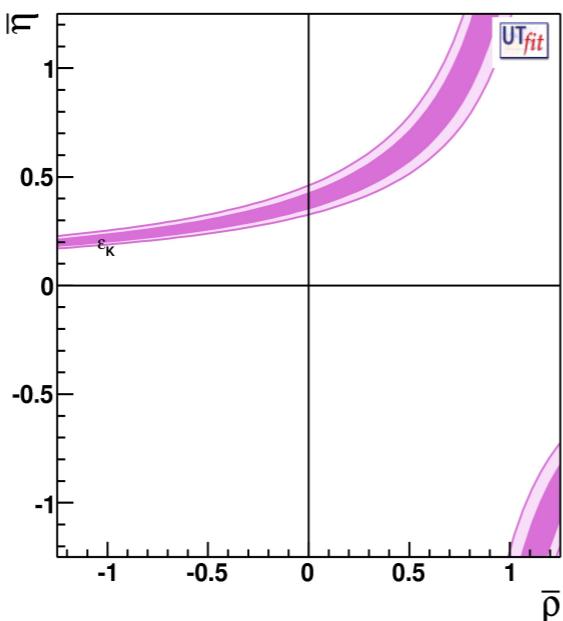


$\text{BR}(B \rightarrow \tau\nu)$



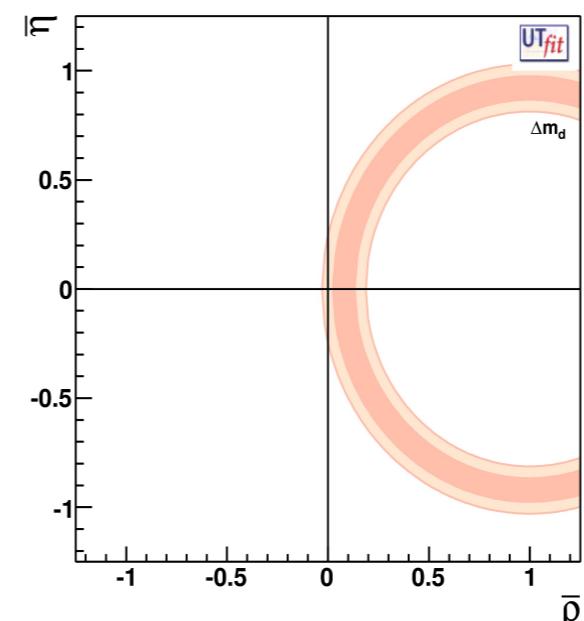
$|V_{ub}/V_{cb}|$

(semileptonic decays)



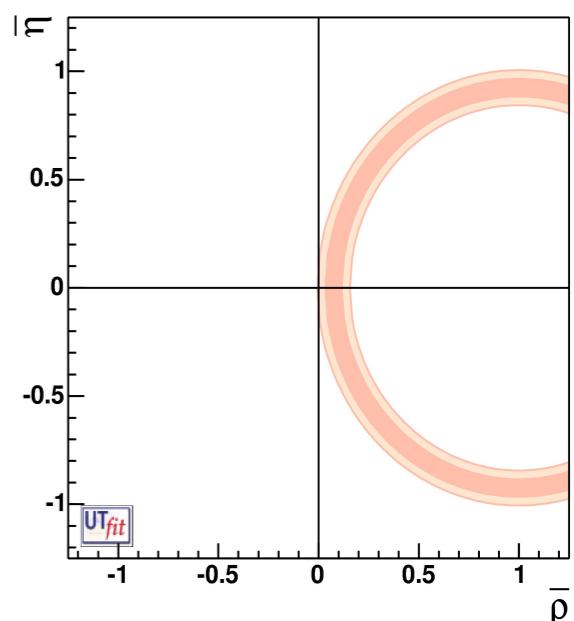
ϵ_K

(CPV in $K - \bar{K}$)



Δm_d

$(B_{d,s} - \bar{B}_{d,s})$



$\Delta m_d/\Delta m_s$

$$\mathcal{P}(\bar{\rho}, \bar{\eta}, \vec{p} \mid \vec{\mathcal{O}}) \sim \mathcal{P}(\vec{\mathcal{O}} \mid \bar{\rho}, \bar{\eta}, \vec{p}) \times \mathcal{P}_0(\bar{\rho}, \bar{\eta}, \vec{p})$$

posterior \sim likelihood \times prior

see, e.g., **JHEP 07 (2001) 013**

- $\vec{\mathcal{O}} = \{\epsilon_K, \Delta m_{d,s}, \dots\}$

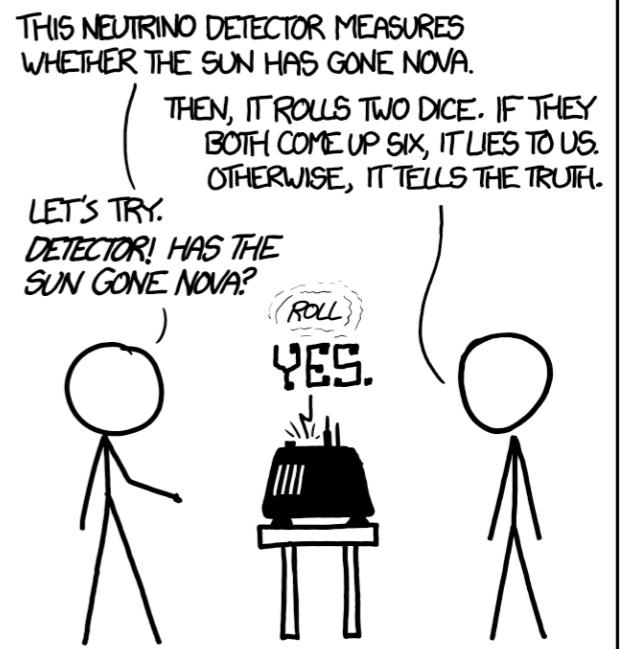
Observables \longleftrightarrow constraints in the fit

- $\vec{p} = \{f_{K,B}, B_{K,B}, \dots\}$

Parameters we can marginalize over

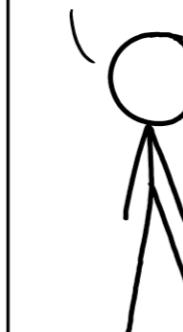
- $(\bar{\rho}, \bar{\eta}) \longleftrightarrow$ CKM pair to be inferred

DID THE SUN JUST EXPLODE?
(IT'S NIGHT, SO WE'RE NOT SURE.)



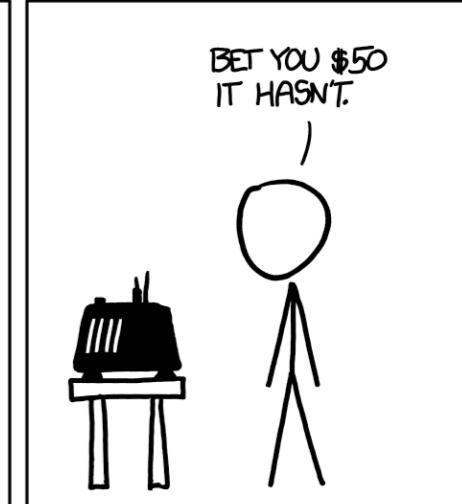
FREQUENTIST STATISTICIAN:

THE PROBABILITY OF THIS RESULT HAPPENING BY CHANCE IS $\frac{1}{36}=0.027$. SINCE $p < 0.05$, I CONCLUDE THAT THE SUN HAS EXPLODED.



BAYESIAN STATISTICIAN:

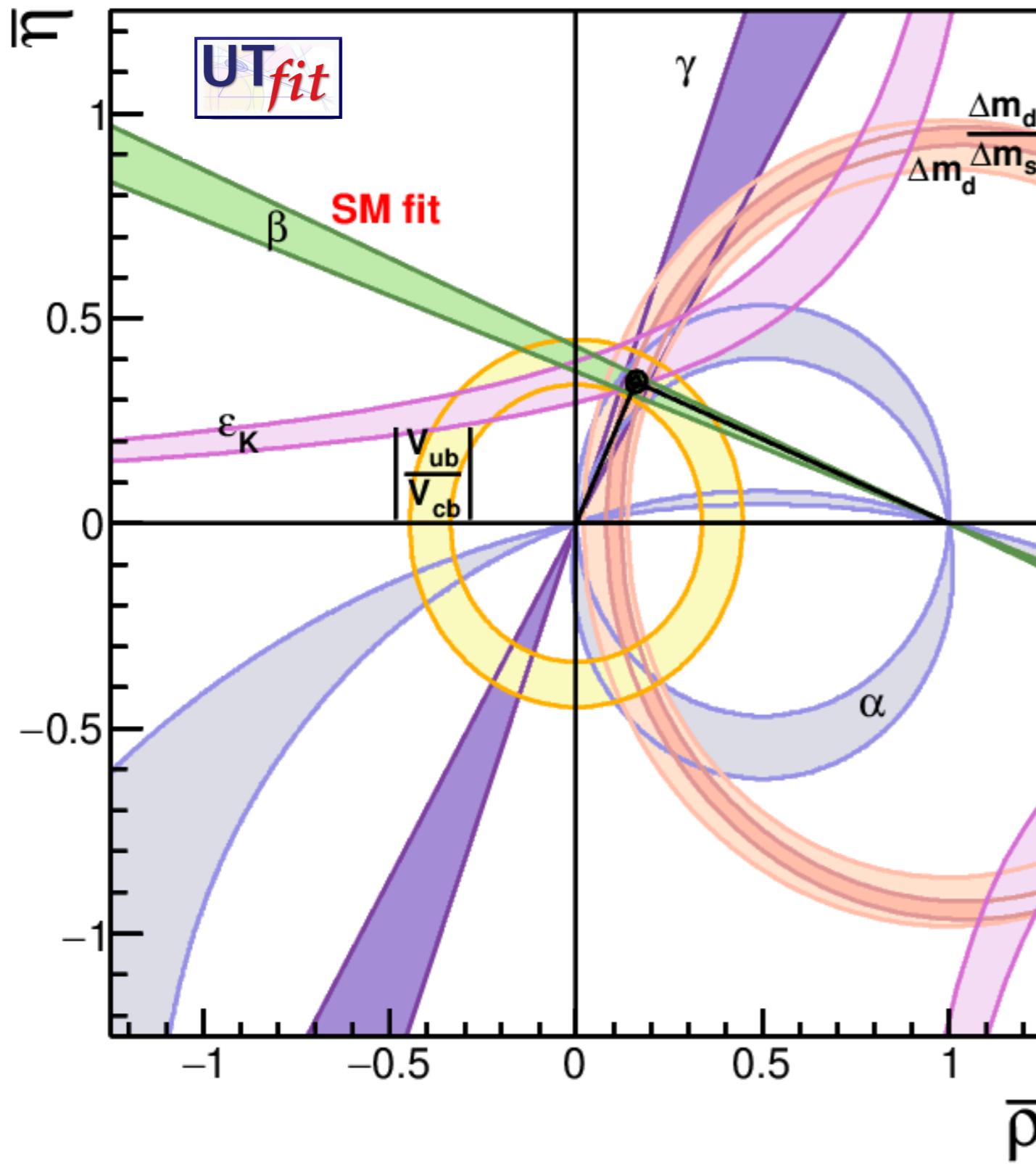
BET YOU \$50 IT HASN'T.



@ <https://xkcd.com/1132>

UTA: Unitarity Triangle Analysis

@ 95% prob

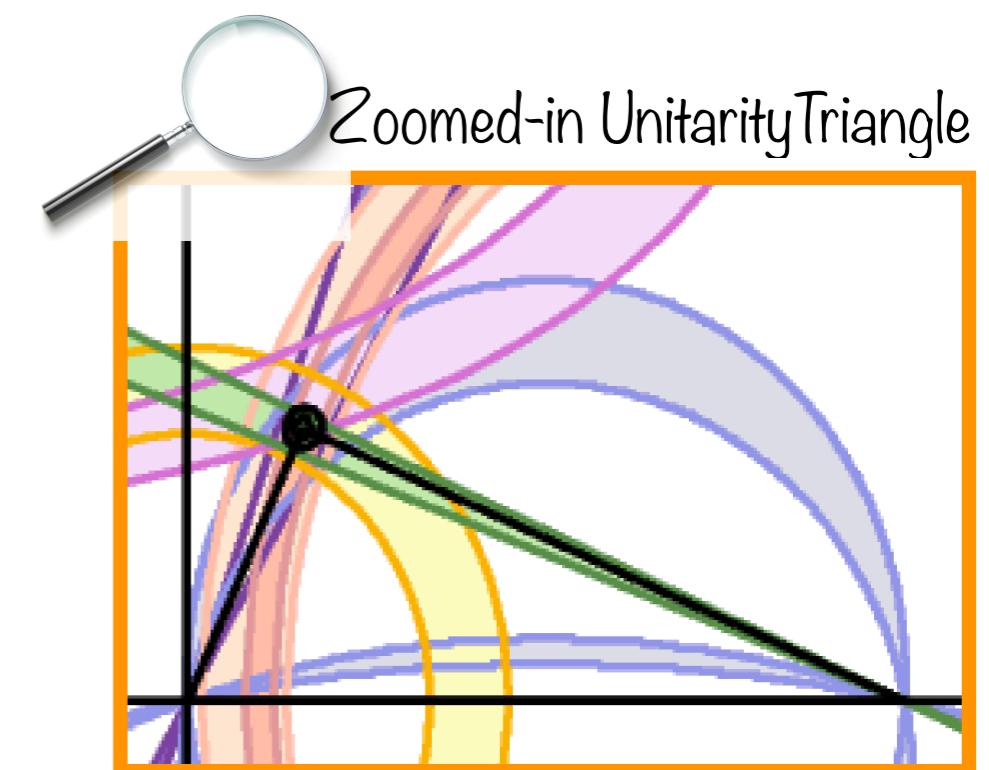


$$\bar{\rho} = 0.160 \pm 0.009 \sim 6\%$$

$$\bar{\eta} = 0.346 \pm 0.009 \sim 3\%$$

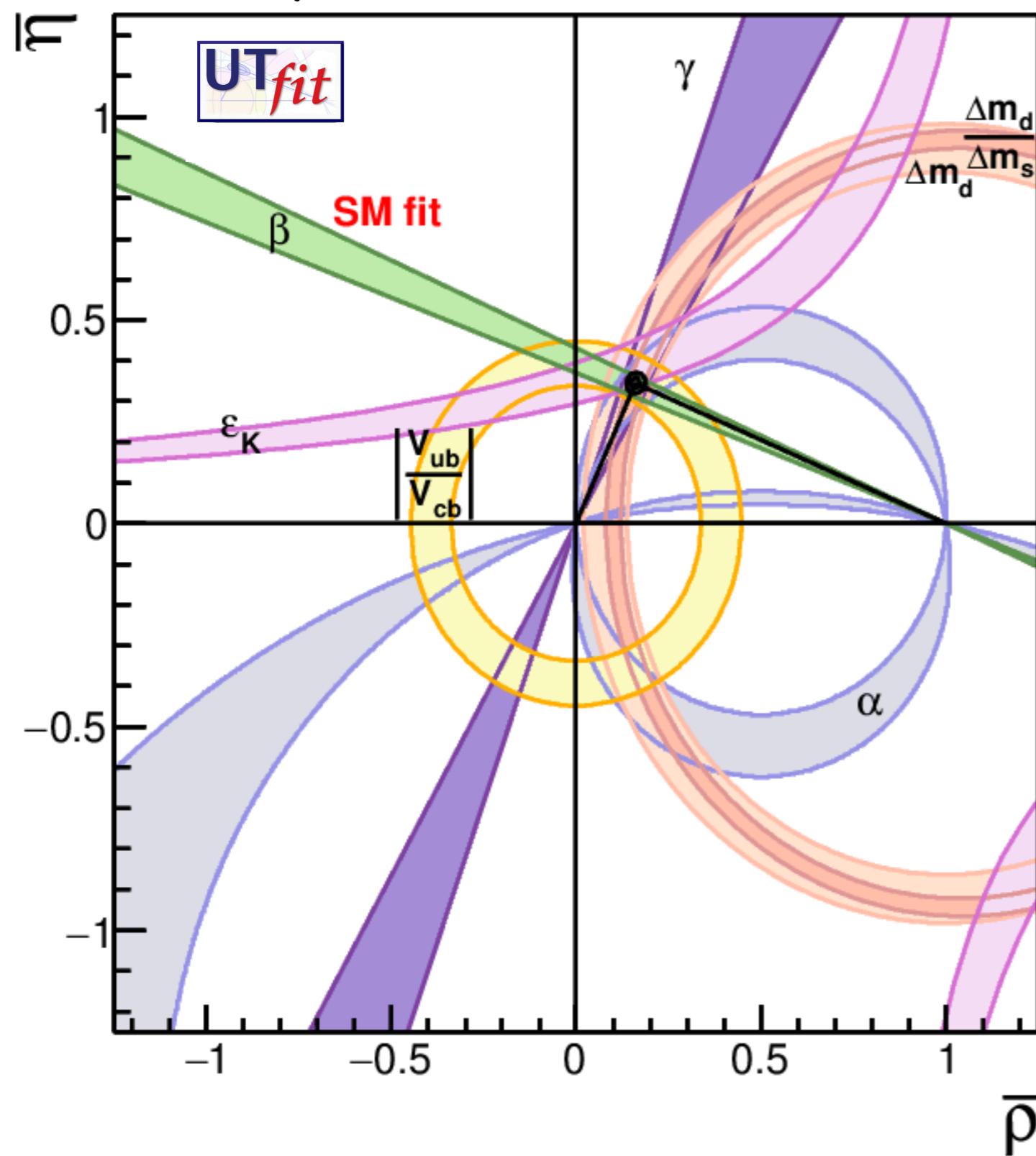
$$\lambda = 0.2251 \pm 0.0008$$

$$A = 0.827 \pm 0.010$$

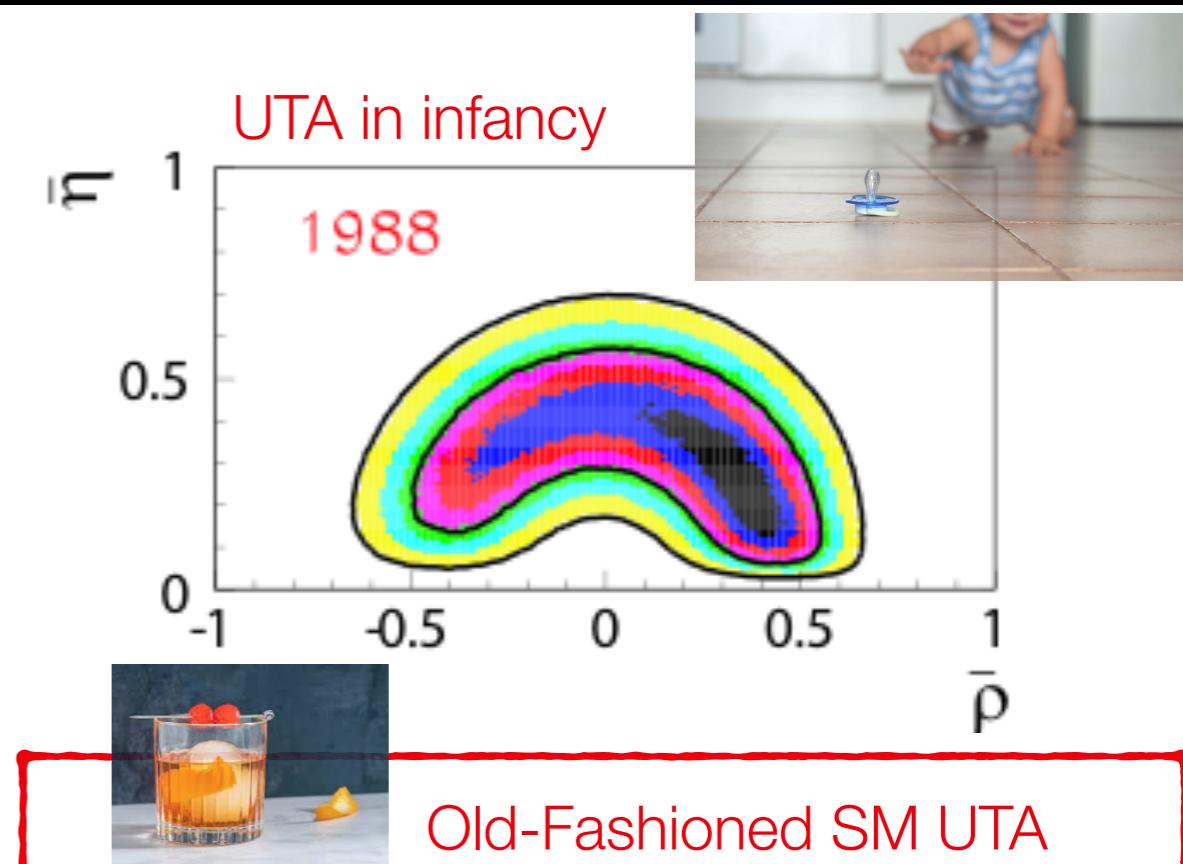


few % determination = decades of tremendous EXP + TH progress!

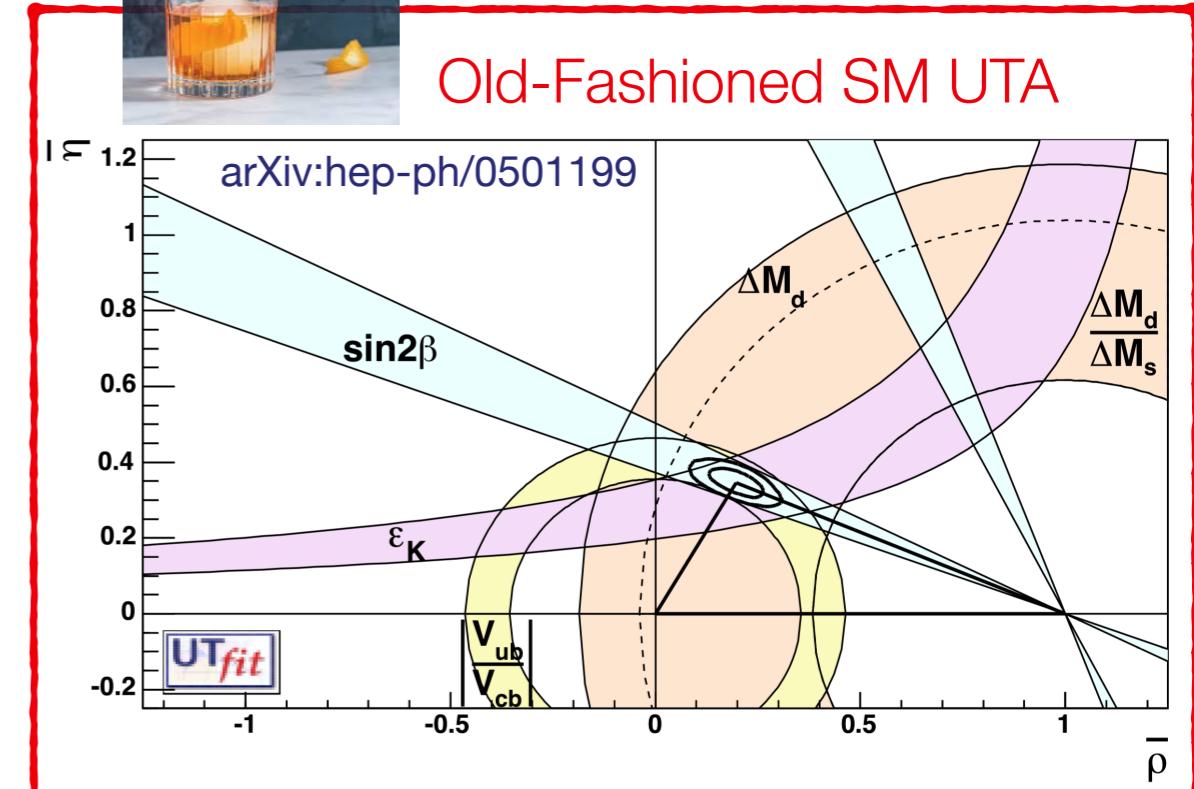
@ 95% prob



UTA in infancy



Old-Fashioned SM UTA



$$\bar{\rho} = 0.196 \pm 0.045 \sim 23\%$$

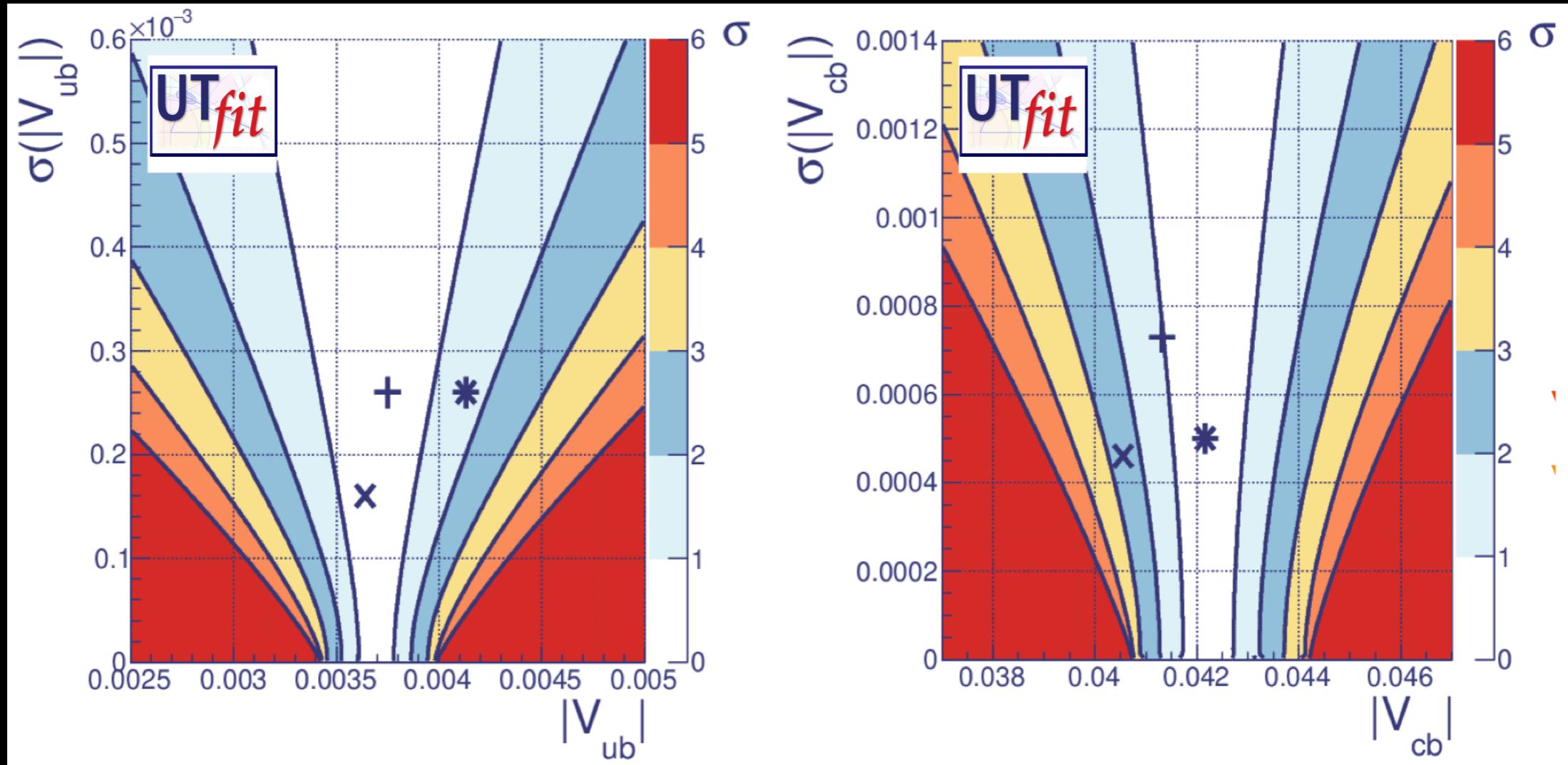
$$\bar{\eta} = 0.347 \pm 0.025 \sim 7\%$$

Compatibility plots

graphical pull of observables

Tensions in the fit

- $+$ \longleftrightarrow measurement
- x \longleftrightarrow exclusive
- $*$ \longleftrightarrow inclusive



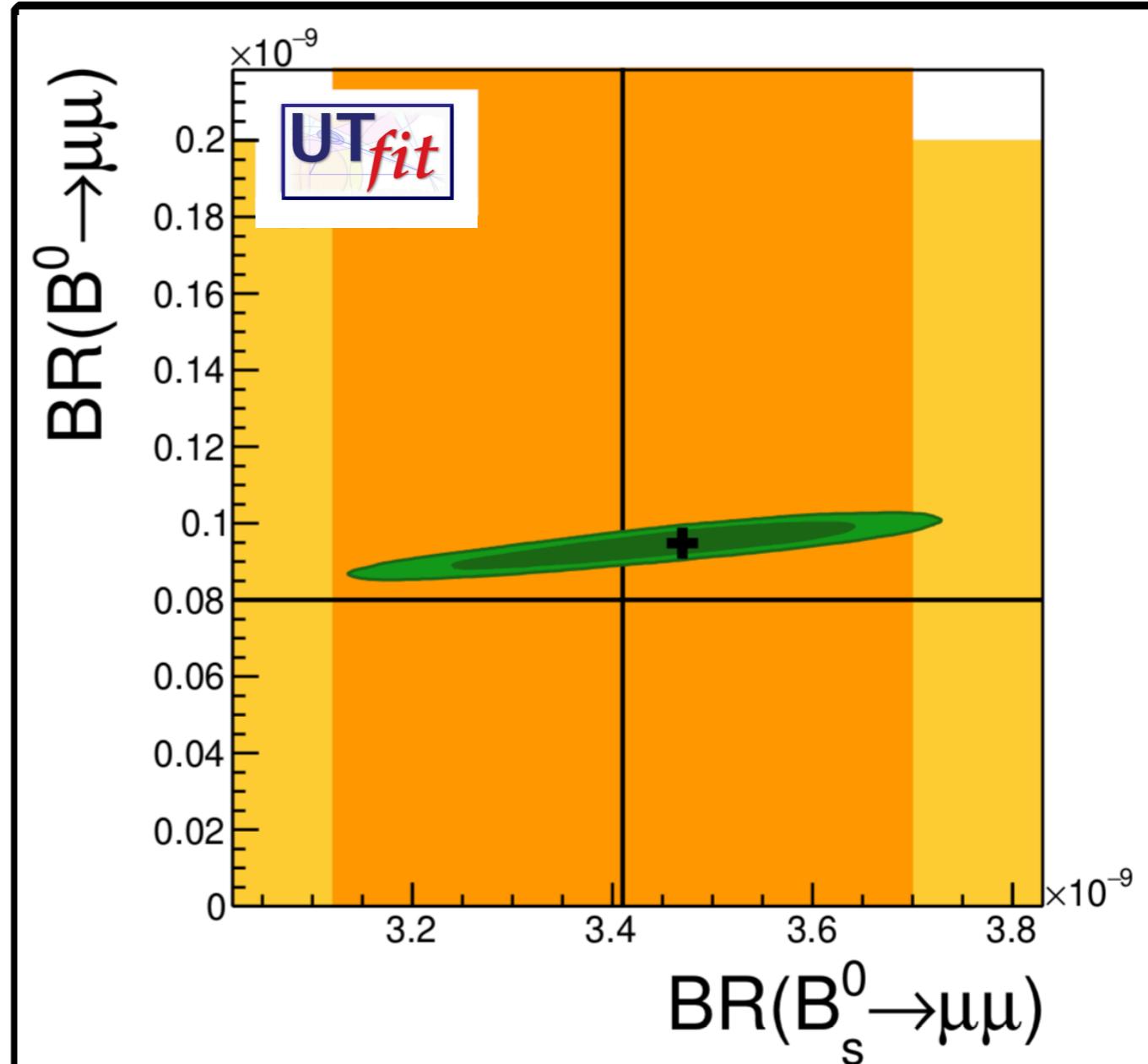
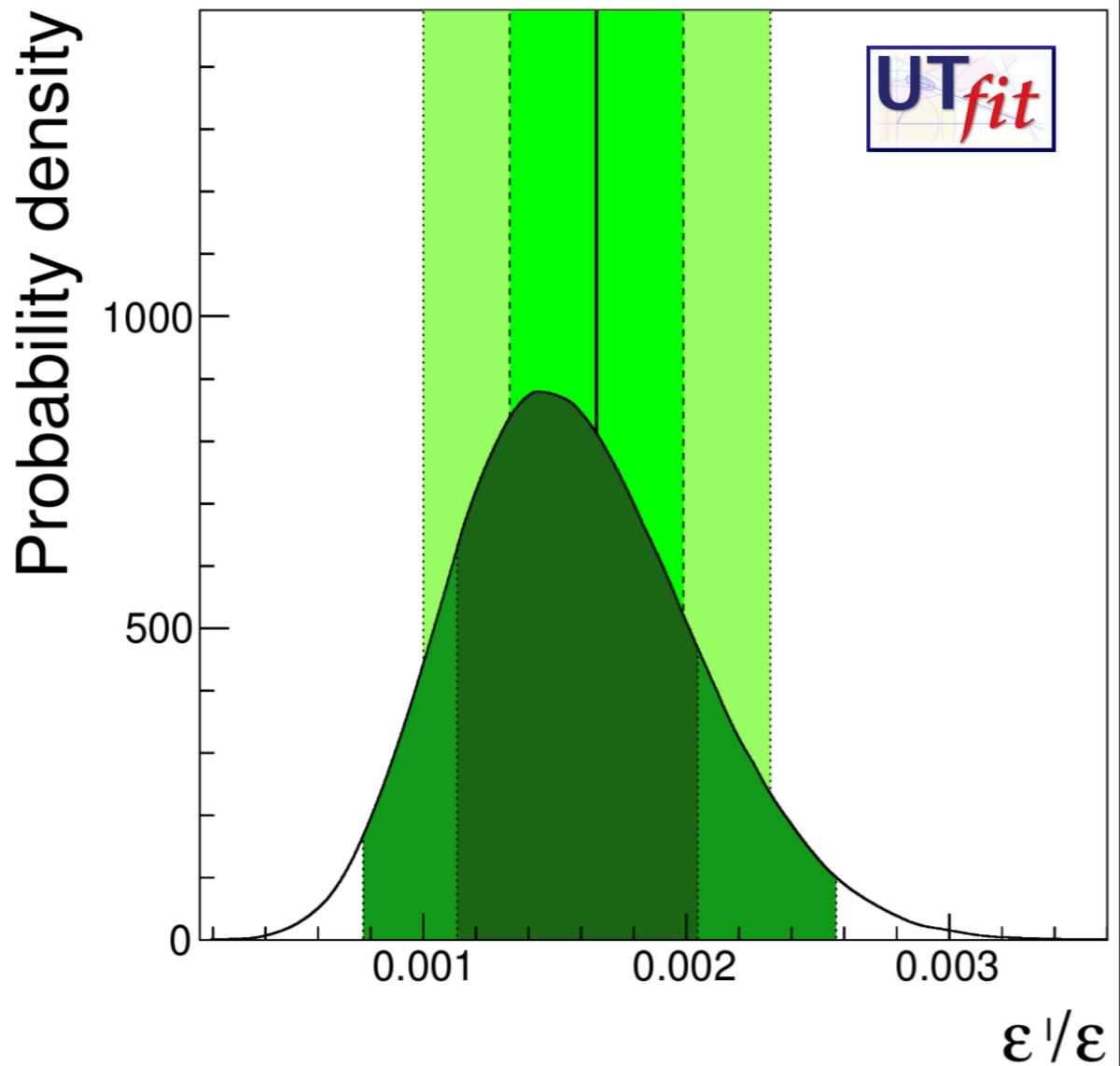
We find only a mild tension in the determination
of V_{cb} from exclusive modes ($< 3\sigma$).

UT Highlights: $|\Delta F| = 1$

UT now includes both indirect
and direct CP from $K \rightarrow \pi\pi$



to RBC/UKQCD Coll.
PRD 102 (2020) 5, 054509



$BR(s)$ in agreement w/ EXP data
→ disfavoring NP for B anomalies

PRD 107 (2023) 5, 055036

B ANOMALIES : A GLIMPSE

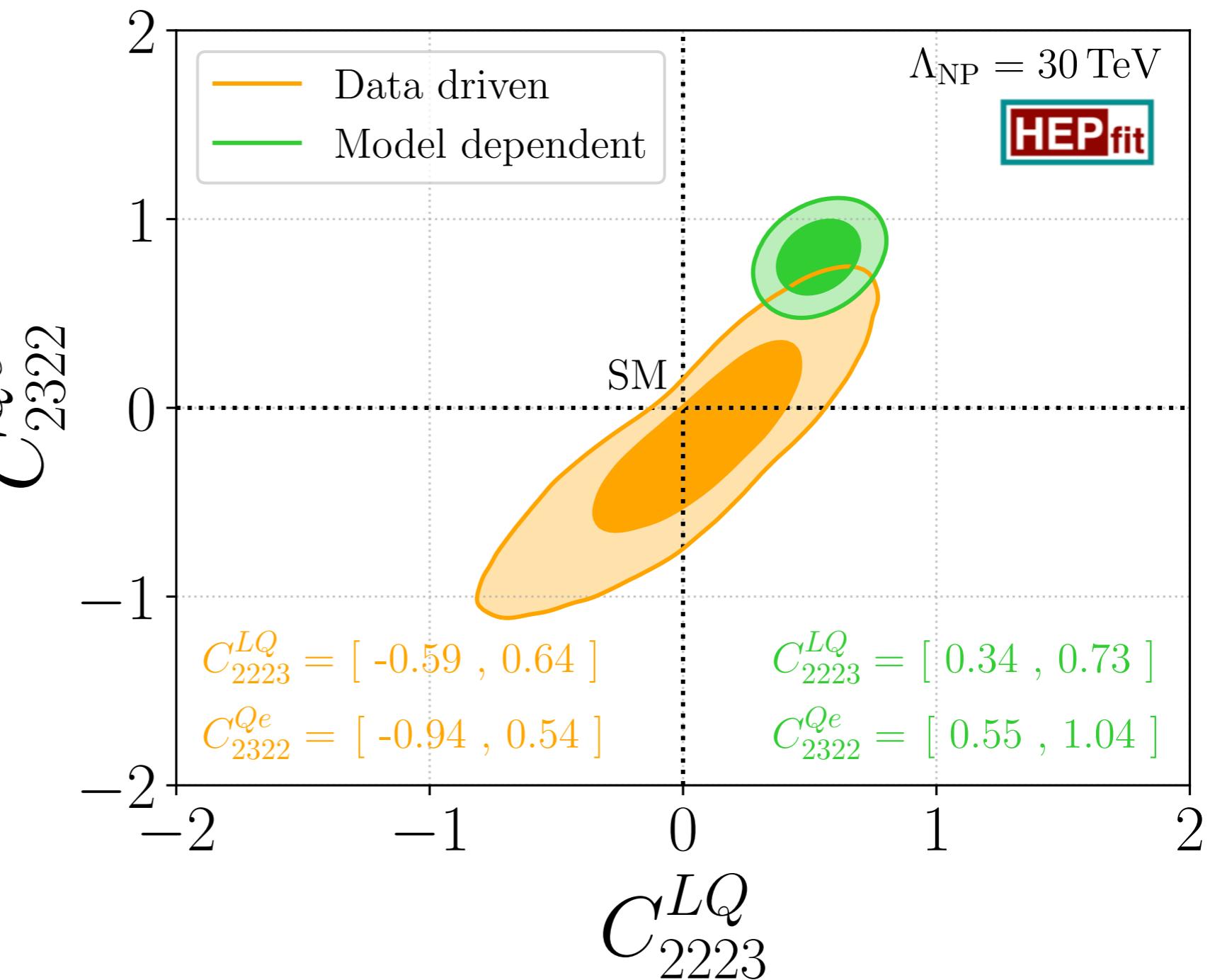
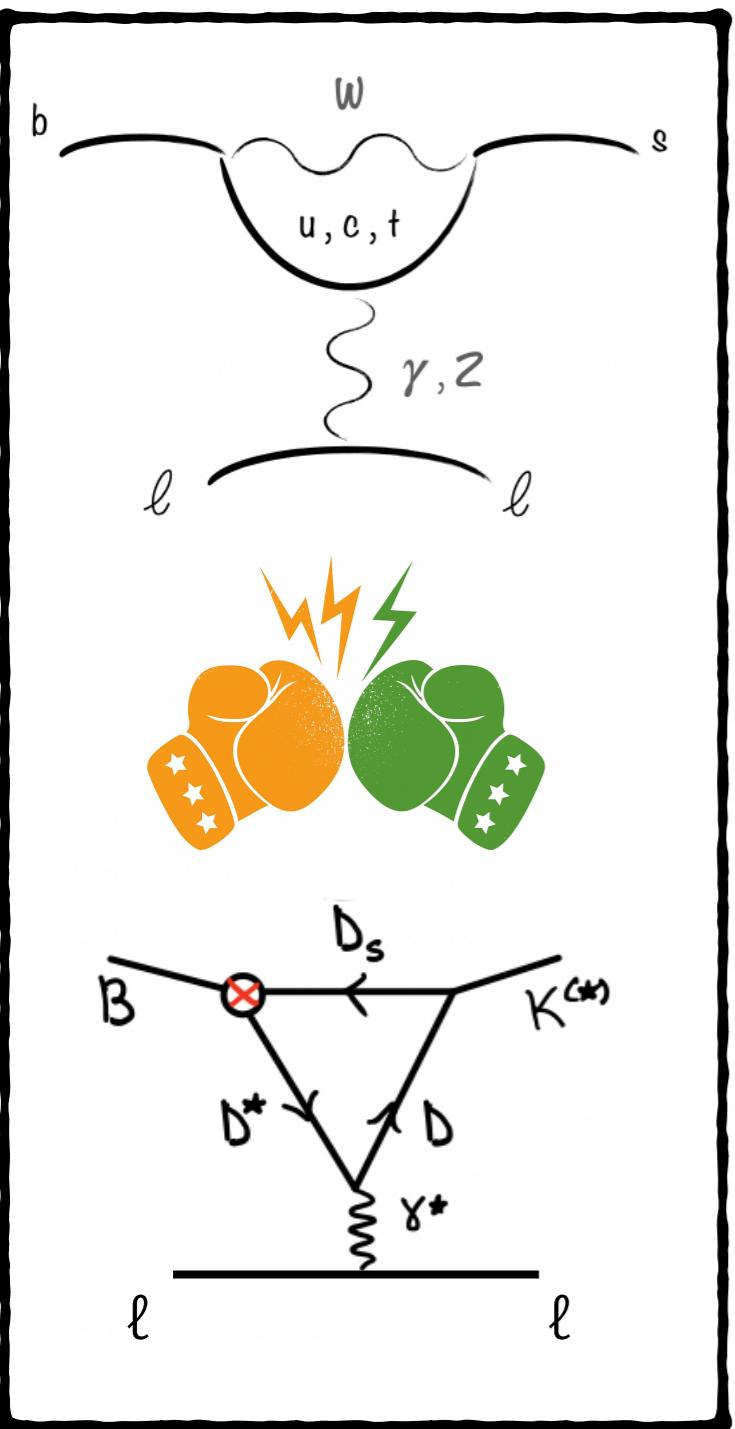
PRD 107 (2023) 5
SMEFT GLOBAL ANALYSIS:
KEY NP OPERATORS

$$O_{2223}^{LQ} = \bar{L}_2 \gamma_\mu L_2 \bar{Q}_2 \gamma^\mu Q_3$$

$$O_{2322}^{Qe} = \bar{Q}_2 \gamma_\mu Q_3 \bar{e}_2 \gamma^\mu e_2$$

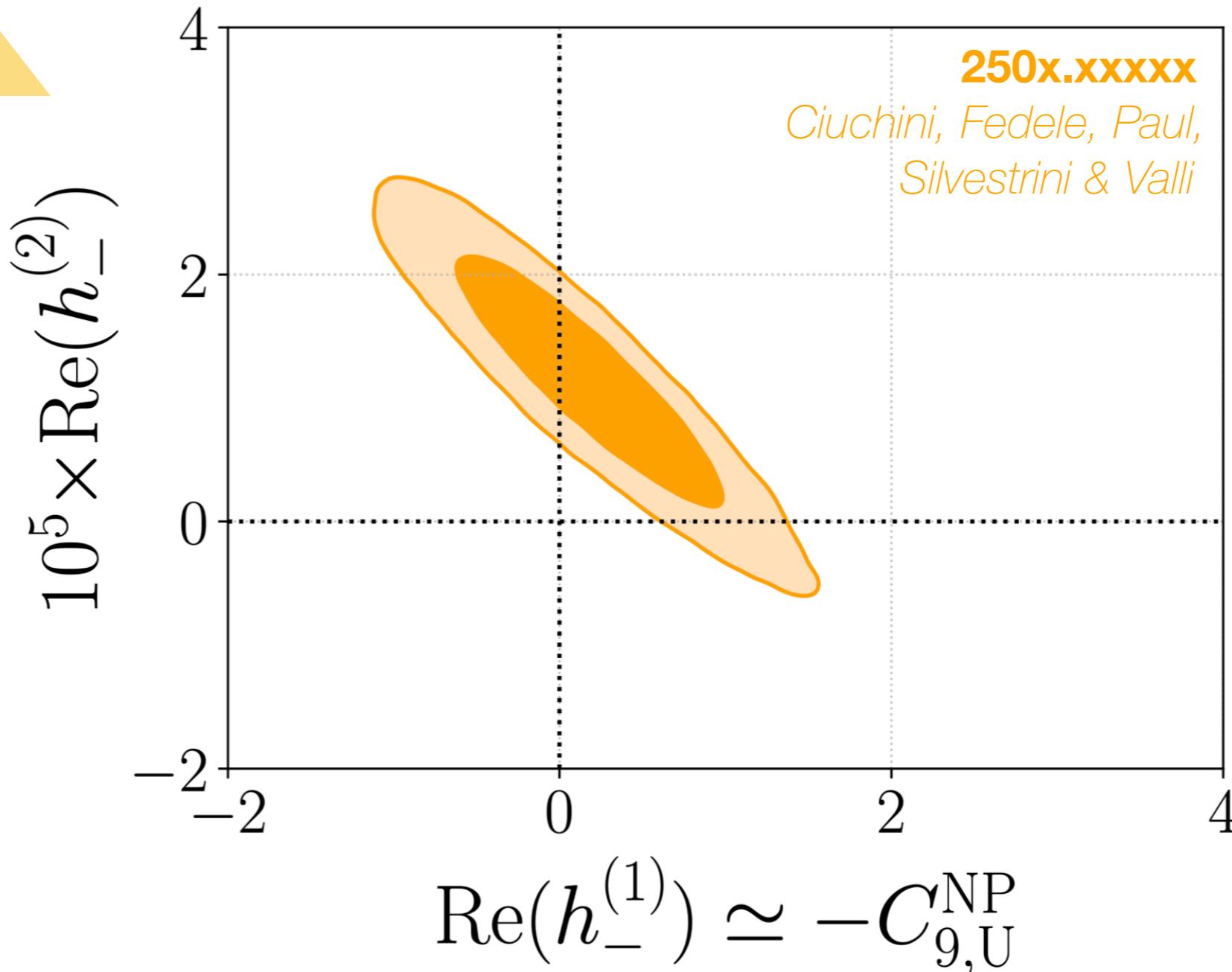
$$C_9 \propto C^{Qe} + C^{LQ}$$

$$C_{10} \propto C^{Qe} - C^{LQ}$$



B ANOMALIES CIRCA 2025

QCD ONLY



New data from
CMS and LHCb
start to point to
hadronic effects!

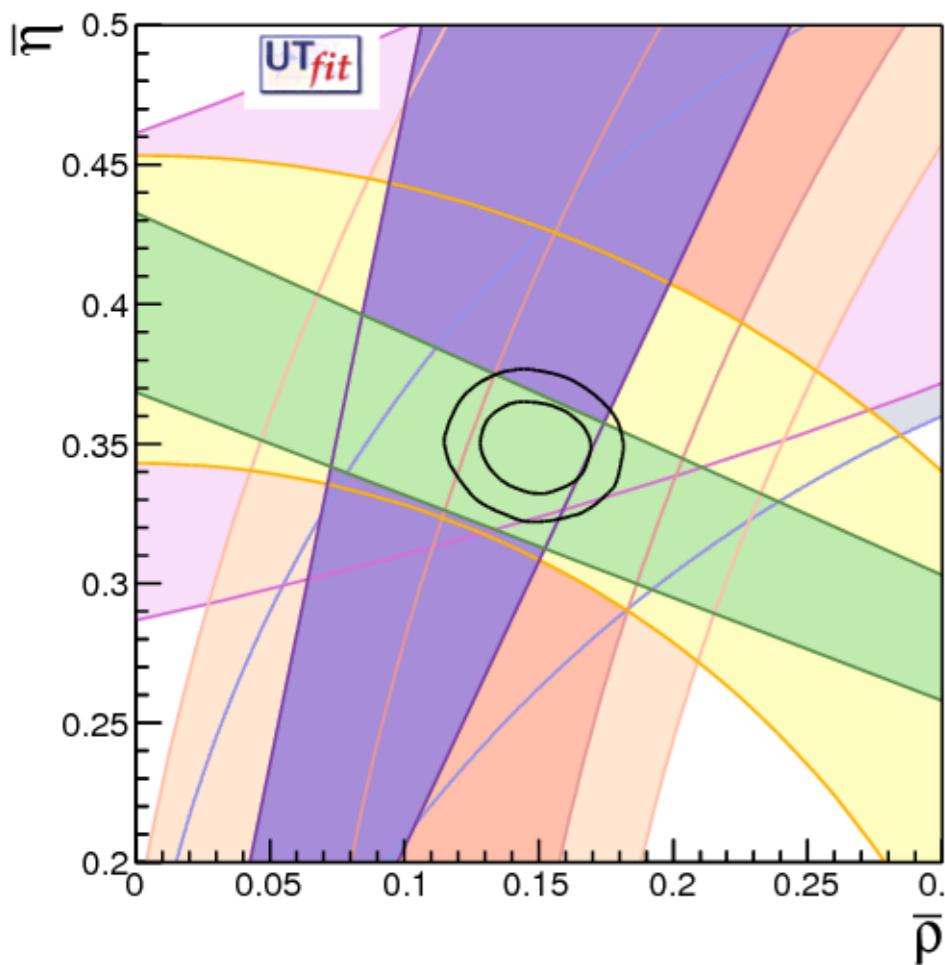
QCD ~ LEPTON UNIVERSAL NP

UT Highlights: HL-LHC

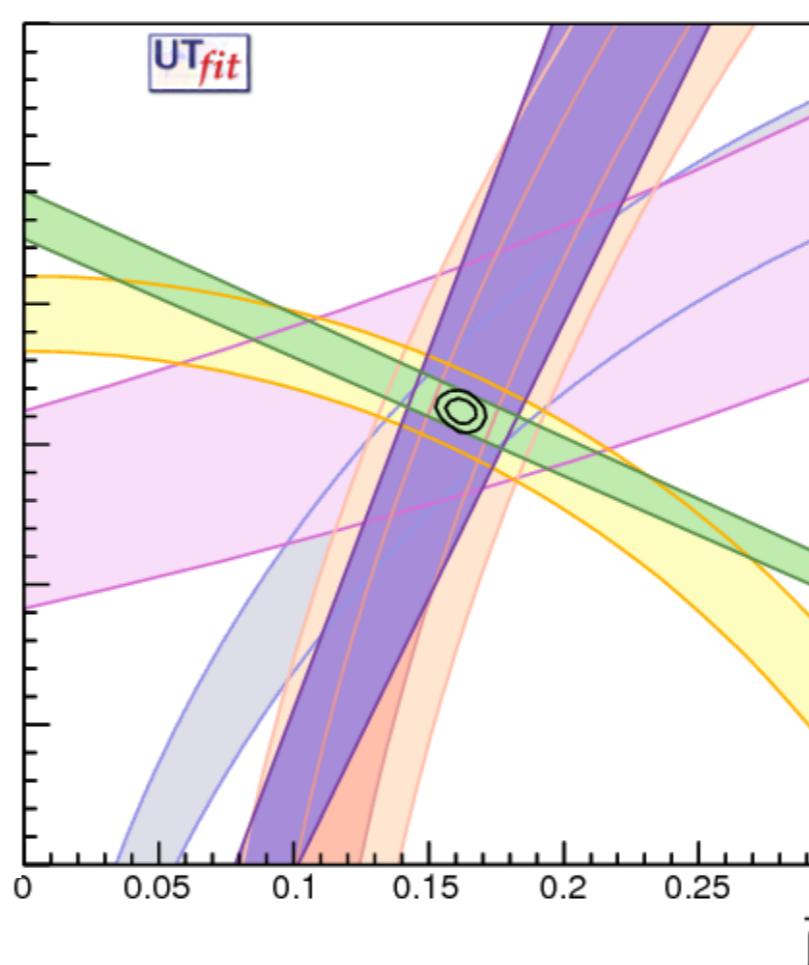
Ultimate upgrade of LHCb will possibly allow for $< 1\%$ in $(\bar{\rho}, \bar{\eta})$

— arXiv:1812.07638 —

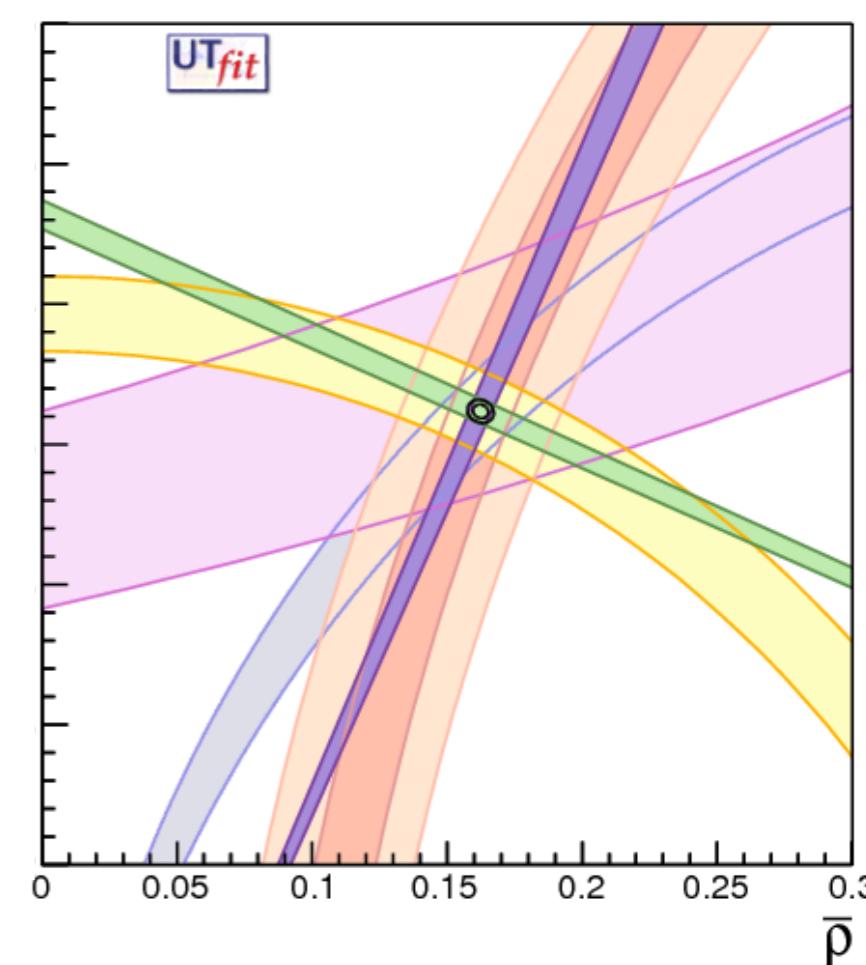
TODAY



RUN 3



HL-LHC

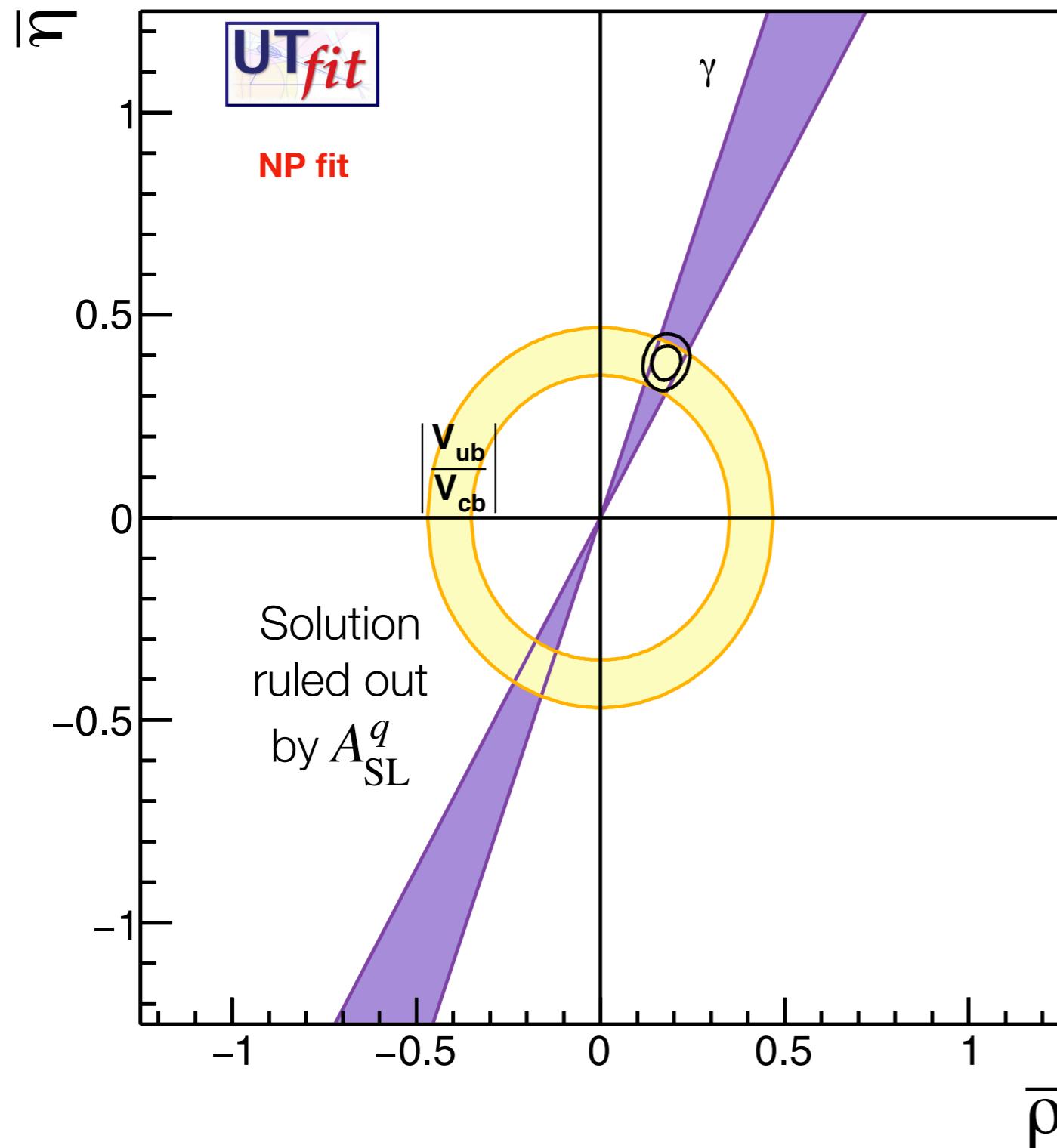


THIS IS W/O IMPACT OF BELLE II & RECENT B-PARKING OF CMS ...

HL-LHC EVEN MORE FLAVORFUL!

NP UT Analysis

*** Assumption: only FCNC amplitudes affected by NP ***



NP UT Analysis

$$\bar{\rho} = 0.167 \pm 0.025 \sim 15\%$$
$$\bar{\eta} = 0.361 \pm 0.027 \sim 7.5\%$$



$$\bar{\rho} = 0.160 \pm 0.009 \sim 6\%$$
$$\bar{\eta} = 0.346 \pm 0.009 \sim 3\%$$

SM UT Analysis

Flavour & BSM Physics

PoS WIFAI2023 (2024) 007

Most general $|\Delta F| = 2 H_{\text{eff}}$:

$$O_1^{q_i q_j} = \bar{q}_{jL}^\alpha \gamma_\mu q_{iL}^\alpha \bar{q}_{jL}^\beta \gamma^\mu q_{iL}^\beta$$

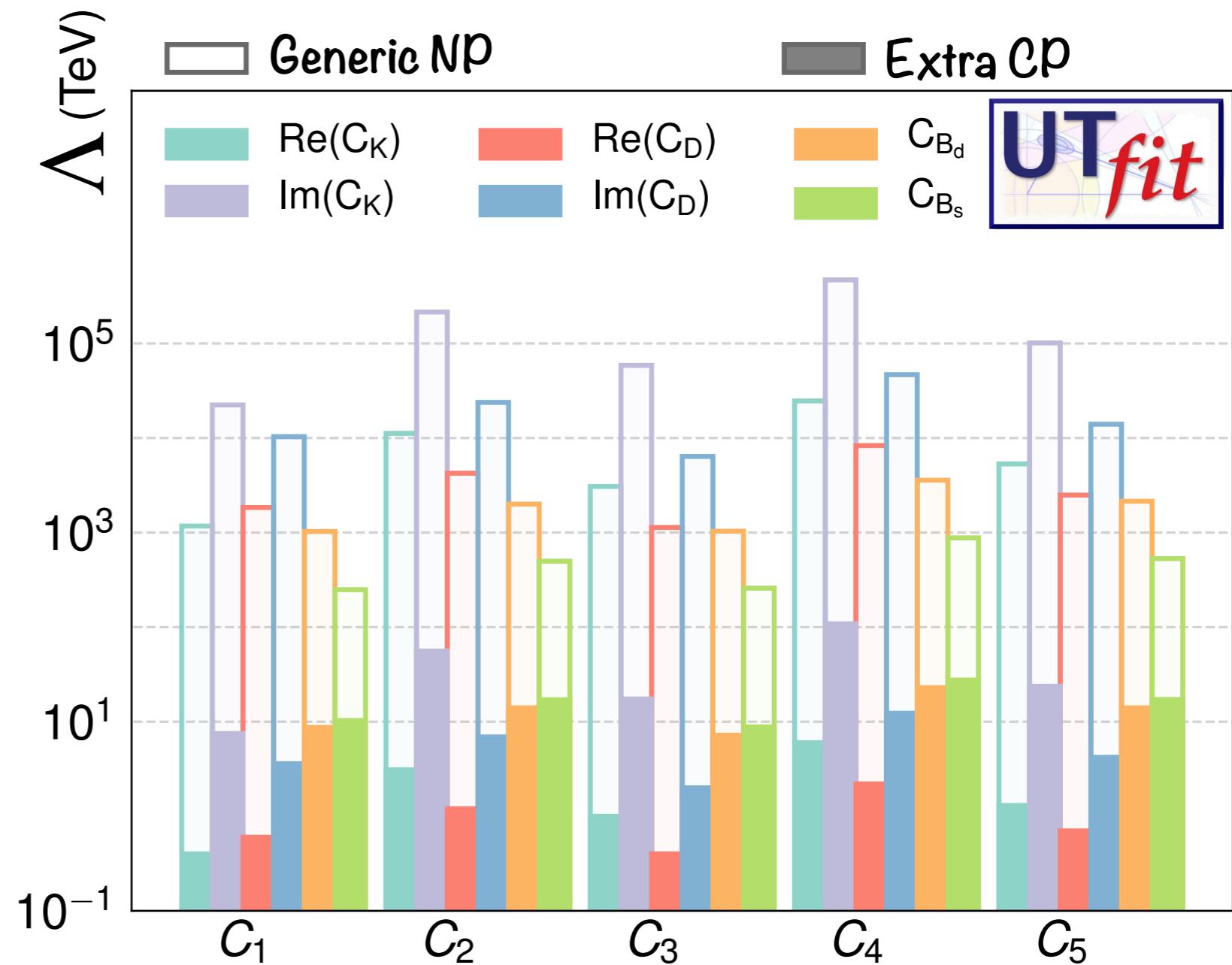
$$O_2^{q_i q_j} = \bar{q}_{jR}^\alpha q_{iL}^\alpha \bar{q}_{jR}^\beta q_{iL}^\beta \quad \text{SM}$$

$$O_3^{q_i q_j} = \bar{q}_{jR}^\alpha q_{iL}^\beta \bar{q}_{jR}^\beta q_{iL}^\alpha$$

$$O_4^{q_i q_j} = \bar{q}_{jR}^\alpha q_{iL}^\alpha \bar{q}_{jL}^\beta q_{iR}^\beta$$

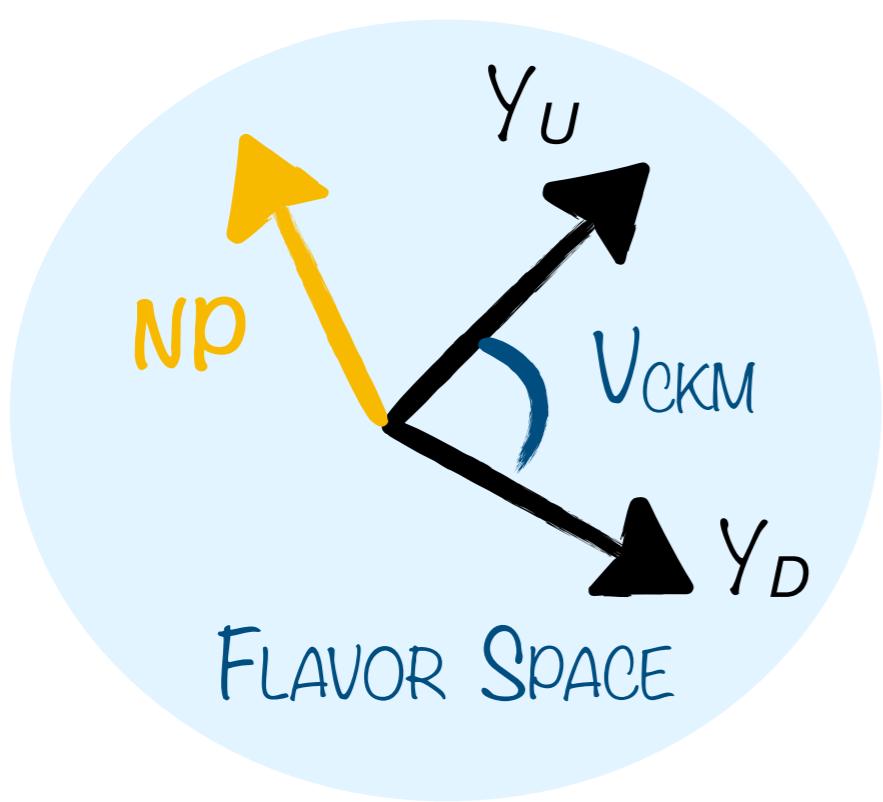
$$O_5^{q_i q_j} = \bar{q}_{jR}^\alpha q_{iL}^\beta \bar{q}_{jL}^\beta q_{iR}^\alpha$$

+ chirally flipped $\tilde{O}_{1,2,3}^{q_i q_j}$

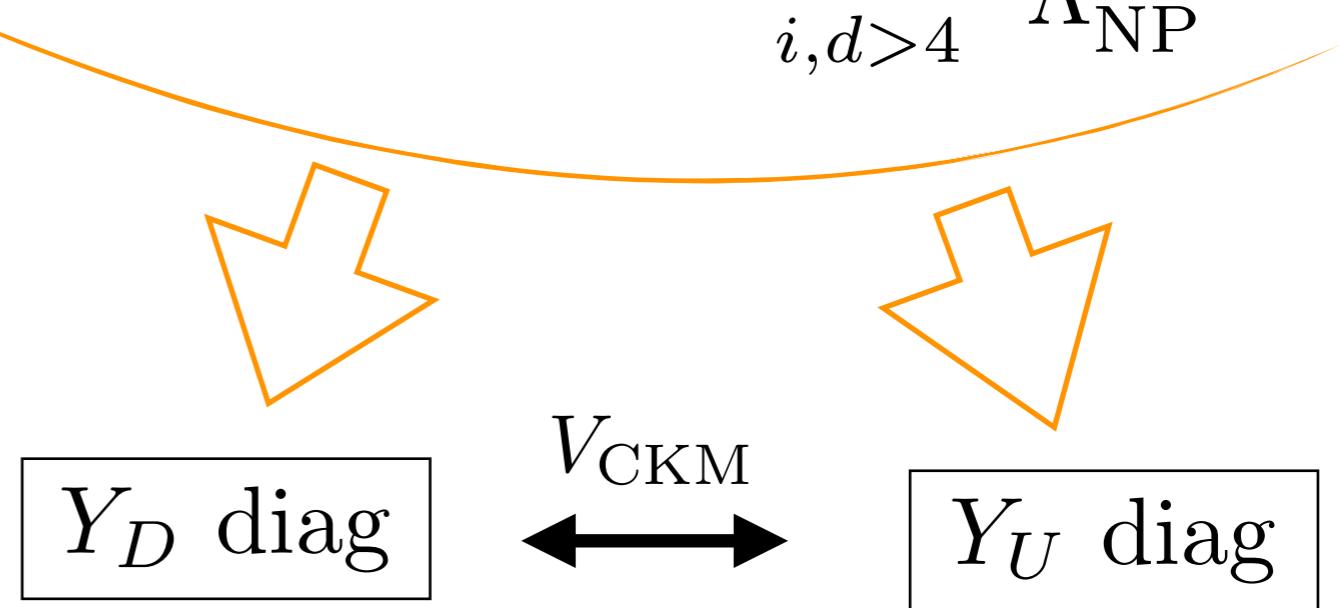


- **Generic NP** = no SM protection, i.e.: $C(\Lambda) \sim 1/\Lambda^2$ → $\Lambda > 4.7 \times 10^5 \text{ TeV}$
- **Extra CP** = SM-like protection but new $O(1)$ phases → $\Lambda > 108 \text{ TeV}$

Going Beyond the Weak EFT



$$\mathcal{L}_{\text{SMEFT}} = \mathcal{L}_{\text{SM}} + \sum_{i,d>4} \frac{C_i \mathcal{O}_i^{(d)}}{\Lambda_{\text{NP}}^{d-4}}$$



$|\Delta F| = 2$ in the SMEFT – **Silvestrini & Valli – Phys. Lett. B 799 (2019) 135062**

SMEFT RGE

$O_{jk}^{HQ(1[3])}$ $(H^\dagger i D_\mu^A H) (\bar{Q}_j \gamma^\mu [\tau^A] Q_k)$	O_{jjkl}^{LeQd} $(\bar{L}_j e_j) (\bar{d}_k Q_l)$	O_{jjkl}^{LeQu} $(\bar{L}_j e_j) i \tau^2 (\bar{Q}_k u_l)$	$O_{jklm}^{ud(1[8])}$ $(\bar{u}_j \gamma_\mu [T^a] u_k) (\bar{d}_l \gamma^\mu [T^a] d_m)$	$O_{jklm}^{QuQd(1[8])}$ $(\bar{Q}_j \gamma_\mu [T^a] u_k) i \tau^2 (\bar{Q}_l \gamma^\mu [T^a] d_m)$
$O_{jklm}^{QQ(1[3])}$ $(\bar{Q}_j \gamma_\mu [\tau^A] Q_k) (\bar{Q}_l \gamma^\mu [\tau^A] Q_m)$	O_{jklm}^{uu} $(\bar{u}_j \gamma_\mu u_k) (\bar{u}_l \gamma^\mu u_m)$	O_{jklm}^{dd} $(\bar{d}_j \gamma_\mu d_k) (\bar{d}_l \gamma^\mu d_m)$	$O_{jklm}^{Qd(1[8])}$ $(\bar{Q}_j \gamma_\mu [T^a] Q_k) (\bar{d}_l \gamma^\mu [T^a] d_m)$	$O_{jklm}^{Qu(1[8])}$ $(\bar{Q}_j \gamma_\mu [T^a] Q_k) (\bar{u}_l \gamma^\mu [T^a] u_m)$

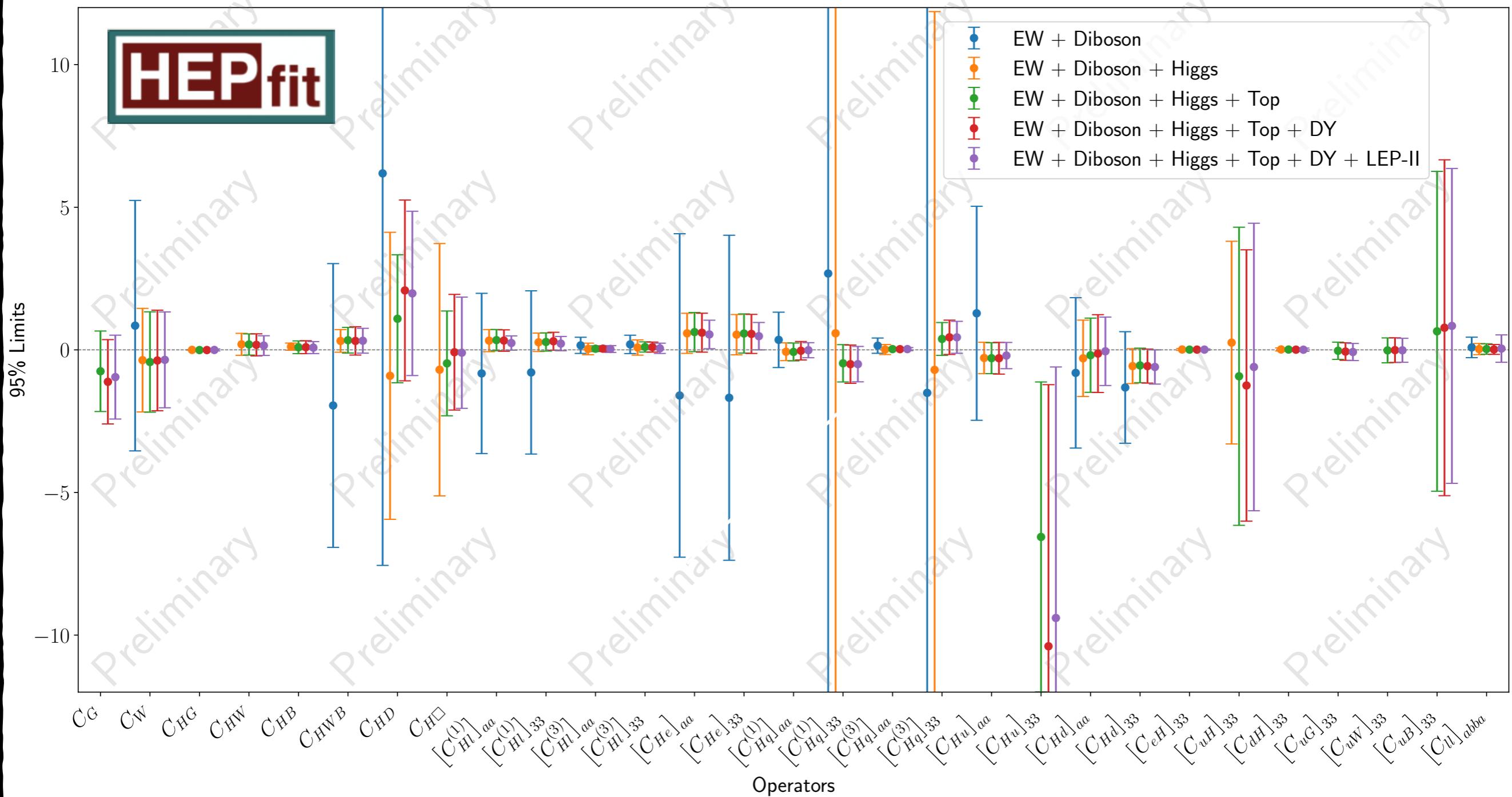
poorly constrained

FLAVOR MISALIGNMENT

UT ANALYSIS IN THE SMEFT: A LOT OF WORK YET TO BE DONE!

$U(2)^5$ flavour symmetry: 2-Fermion

Limits for WC at the scale $\Lambda_{UV} = 1$ TeV



Víctor Miralles

New physics constraints via global fits

Higgs and Effective Field Theory 2024



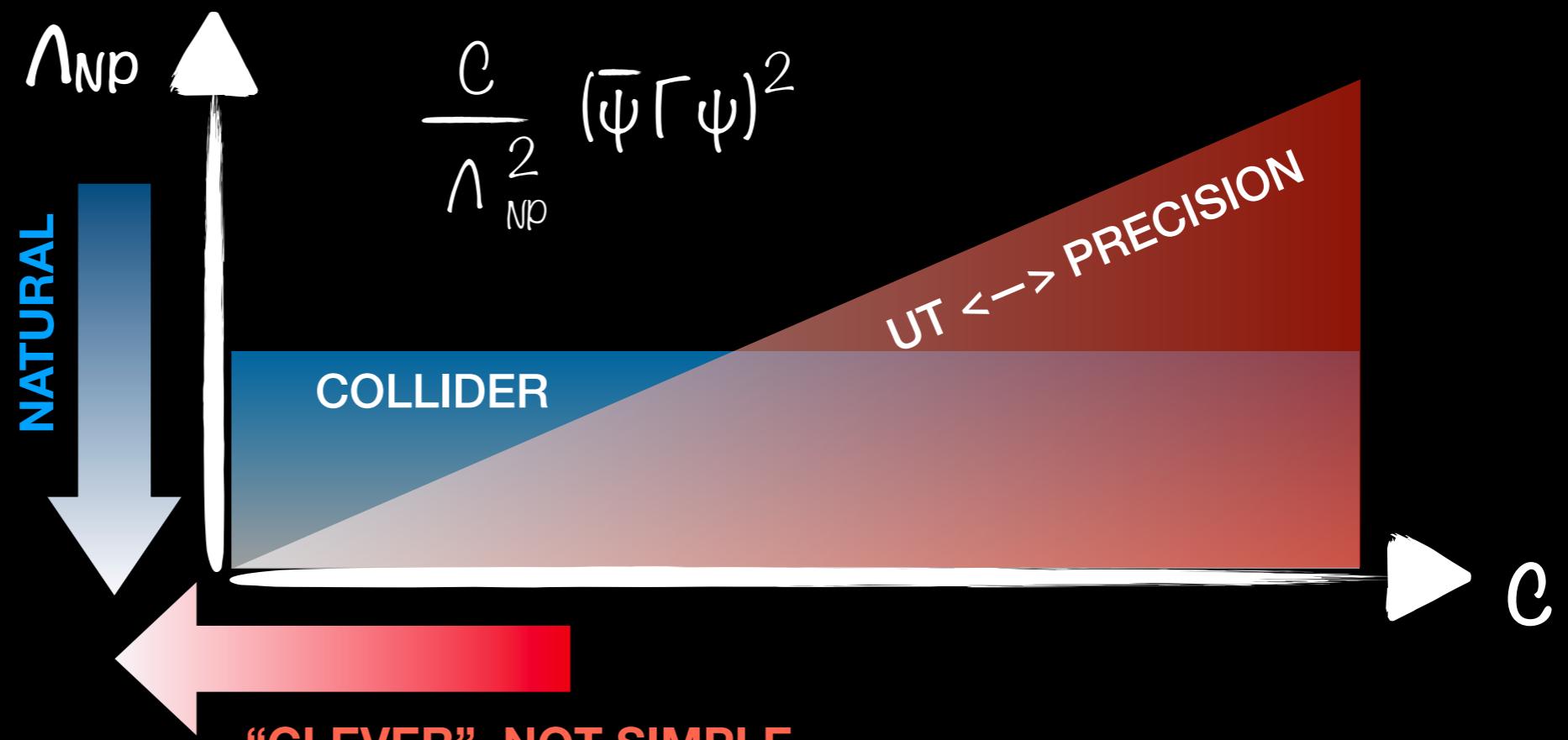
SMEFT vs Precision. Picture further constrained by Flavor...
[with J.de Blas, A.Goncalves, V.Miralles, L.Reina, L.Silvestrini]



Lessons from Precision

- **SM UT:** Towards % precision ... overall remarkable consistency!
—> *in the HL-LHC era we might aim at a permil test*

- **NP UT:**



**BOTTOM
LINE**

A Theory of Flavor is either **VERY CLEVER** or “**JUST**” UNNATURAL

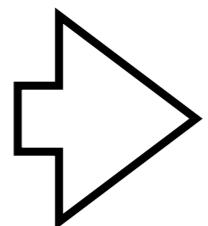
How To Be Clever ...

SEIBERGOLOGY: Yukawa are spurions breaking the flavor group

$$U(3)^5 = SU(3)_Q \times SU(3)_d \times SU(3)_u \times SU(3)_L \times SU(3)_e \times U(1)^5$$

$$U(1)^5 = U(1)_B \times U(1)_L \times U(1)_Y \times U(1)_{PQ} \times U(1)_e$$

E.g.: $-\mathcal{L}_{\text{SM}} \supset Y_{ij}^d \bar{Q}_i H d_j + Y_{ij}^u \bar{Q}_i \tilde{H} u_j + h.c.$



$$Y^d \sim (3, \bar{3}, 1) \quad Y^u \sim (3, 1, \bar{3})$$

under $SU(3)_Q \times SU(3)_d \times SU(3)_u$

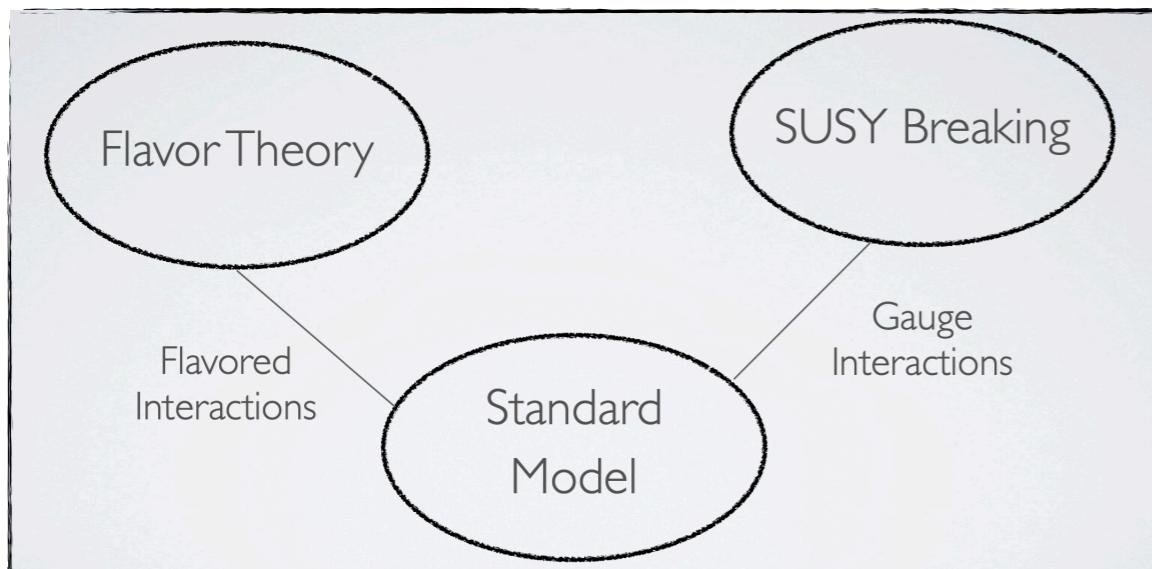
ASSUMPTION: only spurions that break flavor are **SM** Yukawas

→ **MINIMAL FLAVOR VIOLATION** [*Nucl.Phys.B* 645 (2002) 155]

MFV is an Ansatz

MFV allows to push a full-fledged theory of flavor to very high energies.

$$\Lambda_{\text{NP flavor}} \gg \Lambda_{\text{BSM}} \gtrsim v_{\text{EW}}$$



Example: **GAUGE MEDIATION**

In models with gauge-mediated SUSY breaking, the sfermion sector becomes flavorful under SM RGE.

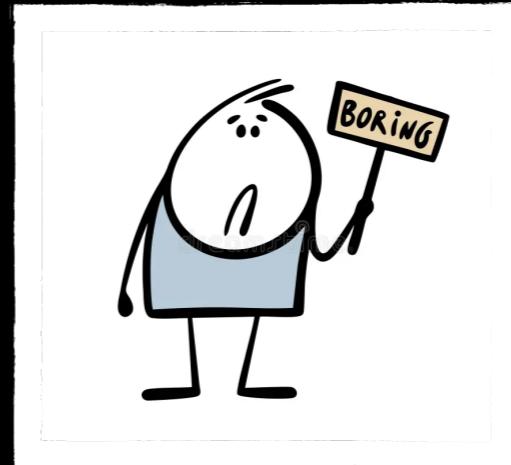
The theory of Flavor is decoupled from the SUSY messenger sector.

FLAVOR BOUNDS GREATLY RELAXED UNDER THE MFV ANSATZ.

“Minimal” for the # of spurions. Actually, **Maximal Flavor Conservation**.

MFV + SMEFT leaves room x **New Physics** just above a **few TeV**.

MFV \longleftrightarrow



Can we avoid large FCNCs
and have BSM with
radically different flavor



Alignment, a kind-of FCNC killer

Consider the SM Lagrangian in the up-quark diag basis:

$$-\mathcal{L}_{\text{SM}} \supset (V y^d)_{ij} \bar{Q}_i H d_j + y^u_{ij} \bar{Q}_i \tilde{H} u_j + h.c.$$

with $y^d = \text{diag}(y_d, y_s, y_b)$ and $y^u = \text{diag}(y_u, y_c, y_t)$

ASSUMPTION: new flavorful interactions **aligned** with y^u

E.g.: $\mathcal{L}_{\text{BSM}} \supset \lambda^u_{ij} \bar{Q}_i H_2 u_j$ with $\lambda^u = \text{diag}(\lambda_u, \lambda_c, \lambda_t)$



Should squarks be degenerate?

Yosef Nir *✉, Nathan Seiberg *✉

Physics Letters B

Volume 309, Issues 3–4, 15 July 1993, Pages 337-343

HOW TO FORMALLY JUSTIFY SUCH ASSUMPTION FOR BSM?

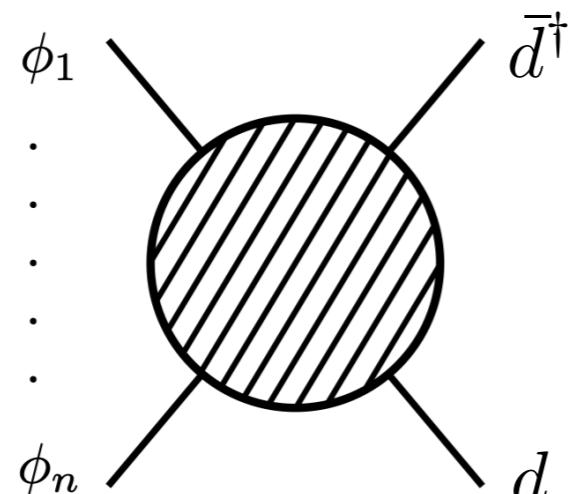
Another Way to Be Clever: Alignment!

— ANSATZ FOR ALIGNMENT IN THE UP-QUARK SECTOR —



- **NO** breaking of the family number $U(1)_f^3$ and of CP **other than** the wave-function renormalization Z^d of d
- **NO** fields / spurions transforming under $U(3)_d$ **but** d & Z^d

Mixing with
Spontaneously
Broken Flavor
Vacuum



$$\mathcal{L} \supset i Z_{ij}^d \bar{d}_i^\dagger \bar{\sigma}^\mu d_j + i \bar{u}_i^\dagger \bar{\sigma}^\mu u_i + i \bar{Q}_i^\dagger \bar{\sigma}^\mu Q_i$$

+ Yukawa terms all flavor diagonal

Canonical kinetic term implies: $(Z^d)^{-\frac{1}{2}} \sim V y^d$

The up-quark sector remains aligned.

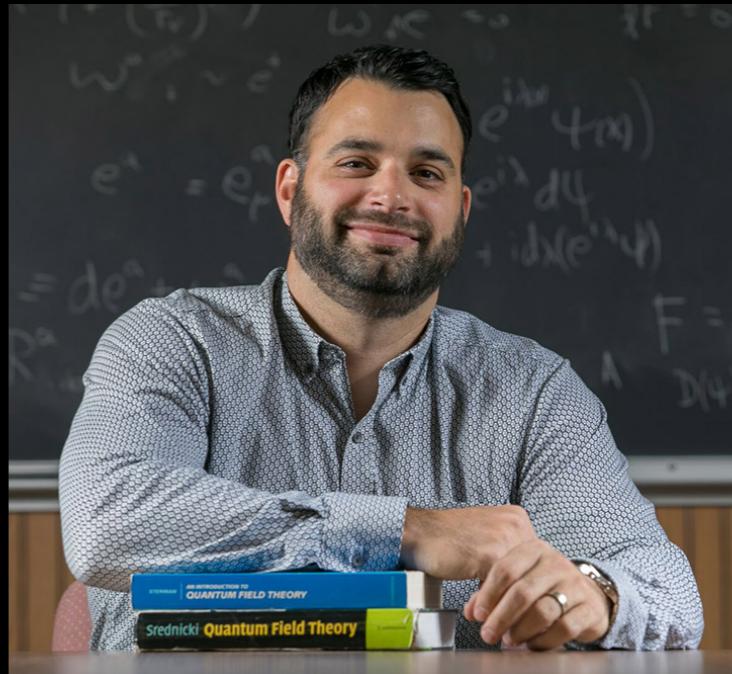
New spurions allowed: $\lambda^u = \text{diag}(\lambda_u, \lambda_c, \lambda_t)$

How charming can the Higgs be?

Artemis Sofia Giannakopoulou,^a Patrick Meade^a and Mauro Valli^b

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^b*INFN Sezione di Roma, Piazzale Aldo Moro 2, I-00185 Rome, Italy*



Artemis Giannakopoulou



SFV 2HDM

SFV physics may have generic family non-universal couplings to SM.

A simple framework: **SFV Two Higgs Doublet Models (2HDM).**

$$|D_\mu H_a|^2 - V(H_1, H_2) - \left(\mathcal{Y}_{aij}^u \bar{Q}_{Li} H_a U_{Rj} + \mathcal{Y}_{aij}^d \bar{Q}_{Li} H_a^c D_{Rj} + \mathcal{Y}_{aij}^\ell \bar{L}_{Li} H_a^c \ell_{Rj} + h.c. \right)$$

$$(a = 1, 2; i, j = 1, 2, 3)$$

2HDM features neutral 2 CP-even H:

$\cos(\beta - \alpha) \neq 0$ interesting!

$$h = H_1^0 \sin(\beta - \alpha) + H_2^0 \cos(\beta - \alpha)$$

$$H = H_1^0 \cos(\beta - \alpha) - H_2^0 \sin(\beta - \alpha)$$

+ a CP-odd and charged ones ... for simplicity:

$$m_H = m_{H^\pm} = m_A$$

DOWN-TYPE SFV ANSATZ ON 2HDM

$$\mathcal{Y}_1^u = y^u = \text{diag}(y_u, y_c, y_t)$$

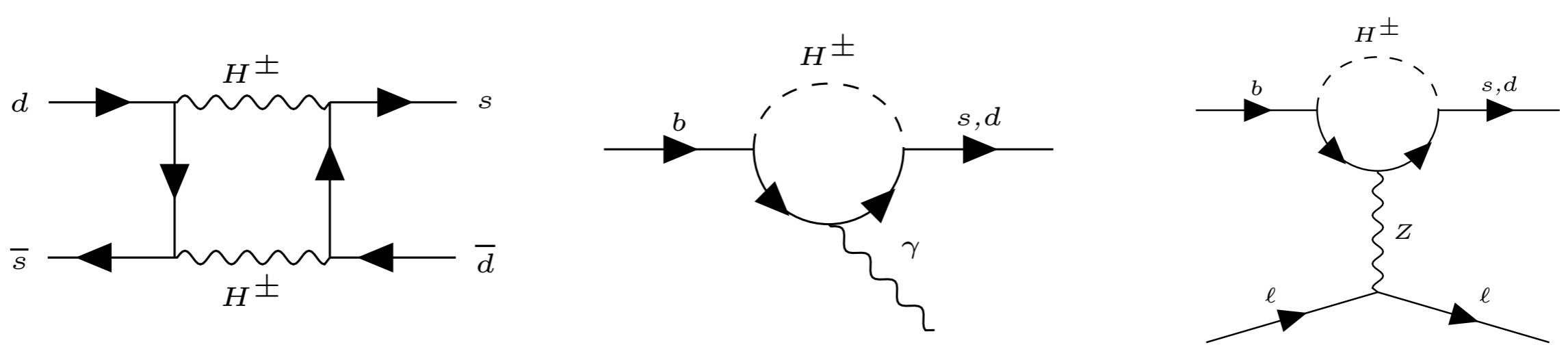
$$\mathcal{Y}_1^d = V y^d$$

$$\mathcal{Y}_1^\ell = y^\ell$$

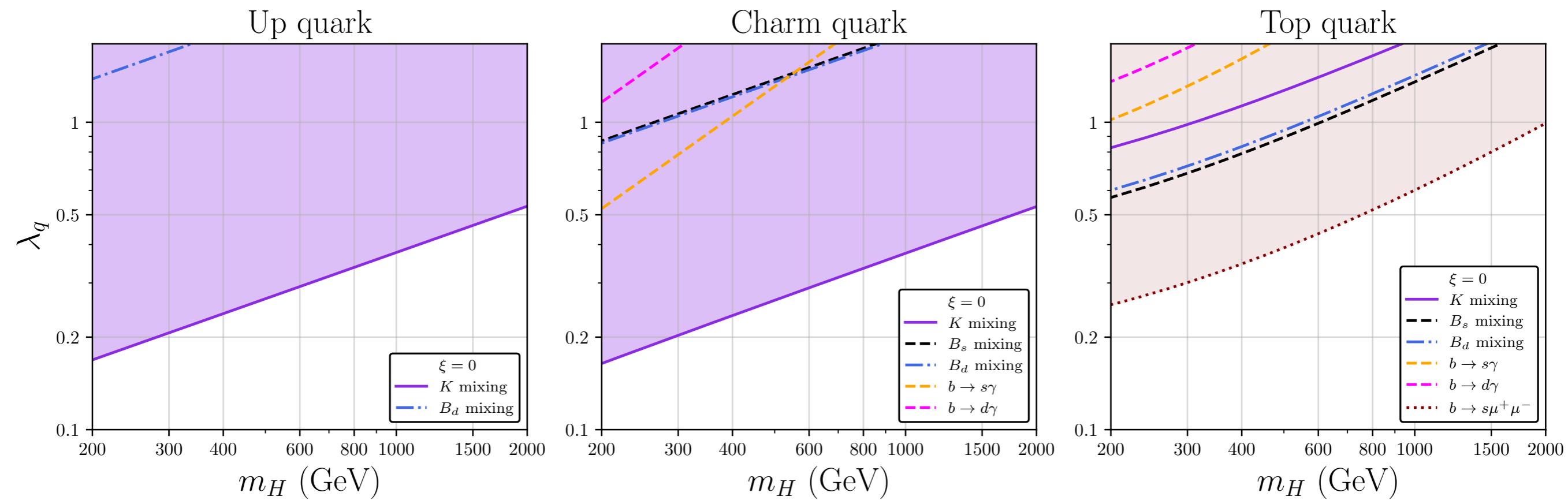
$$\mathcal{Y}_2^u = \lambda^u = \text{diag}(\lambda_u, \lambda_c, \lambda_t)$$

$$\mathcal{Y}_2^d = \xi V y^d$$

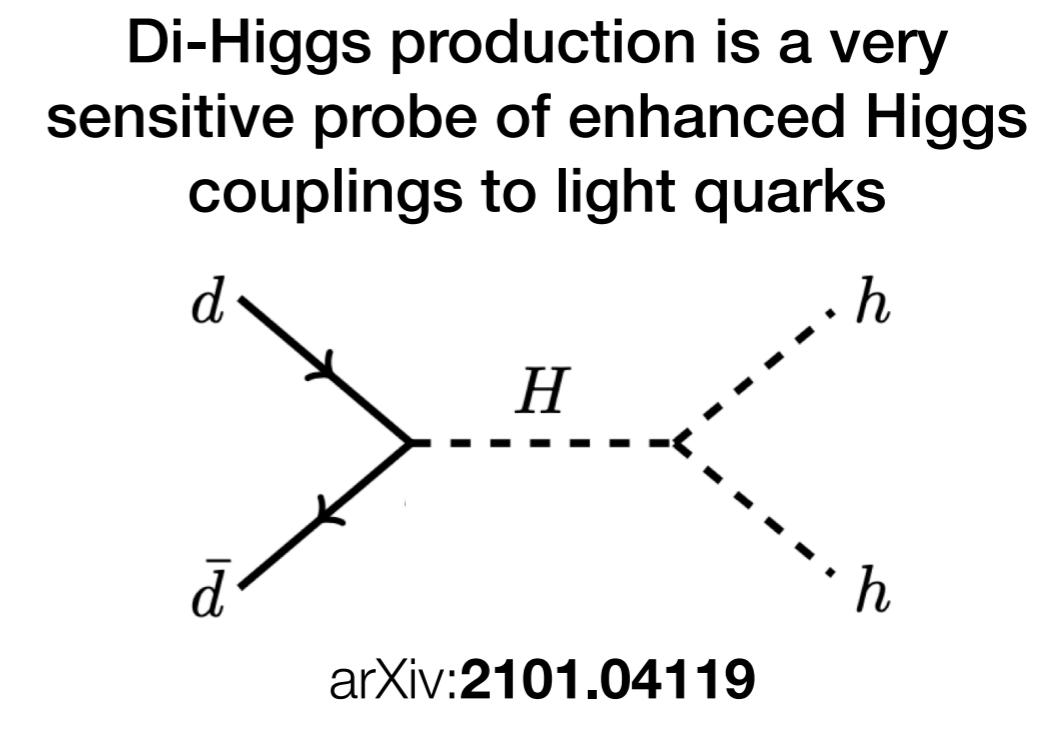
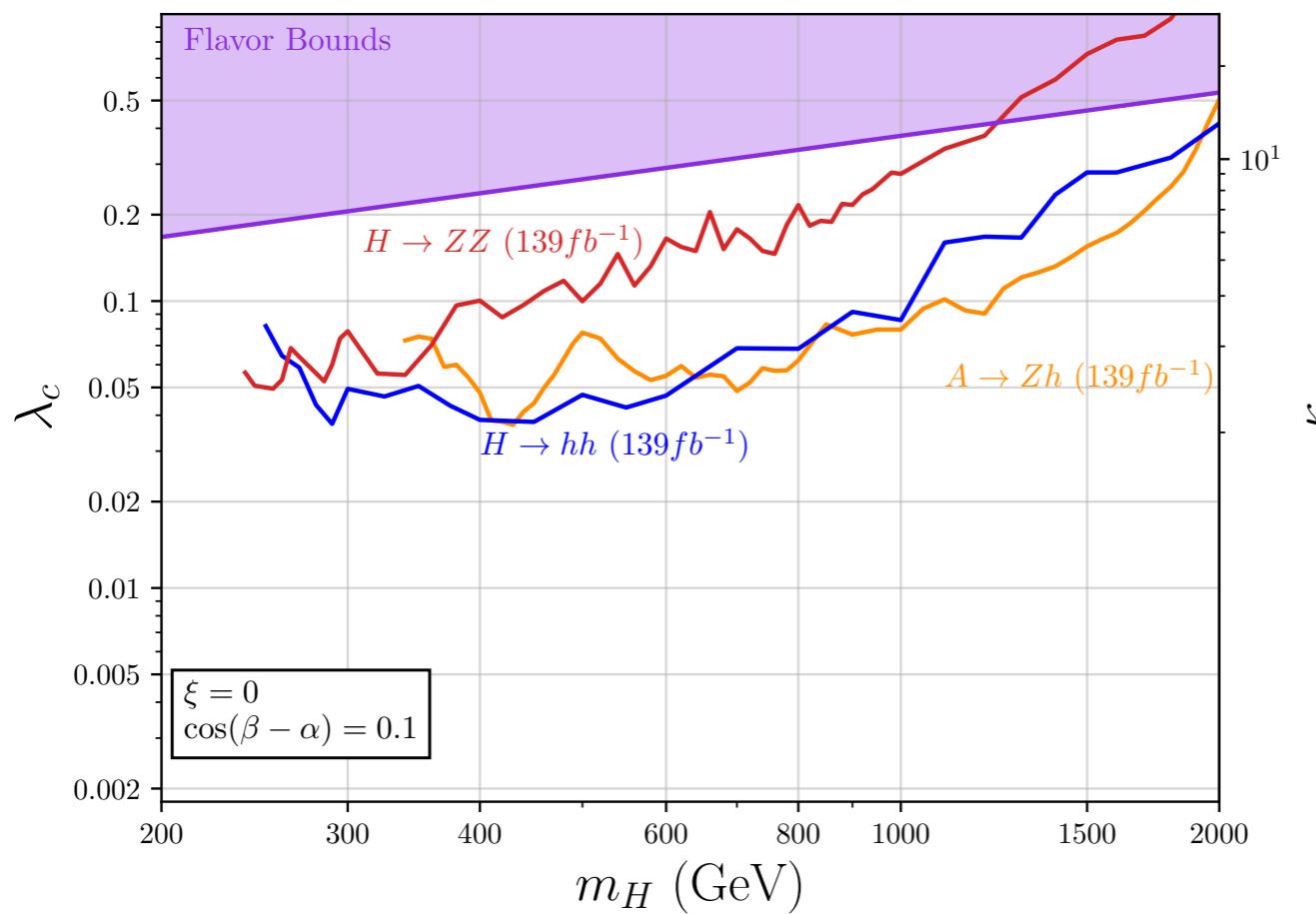
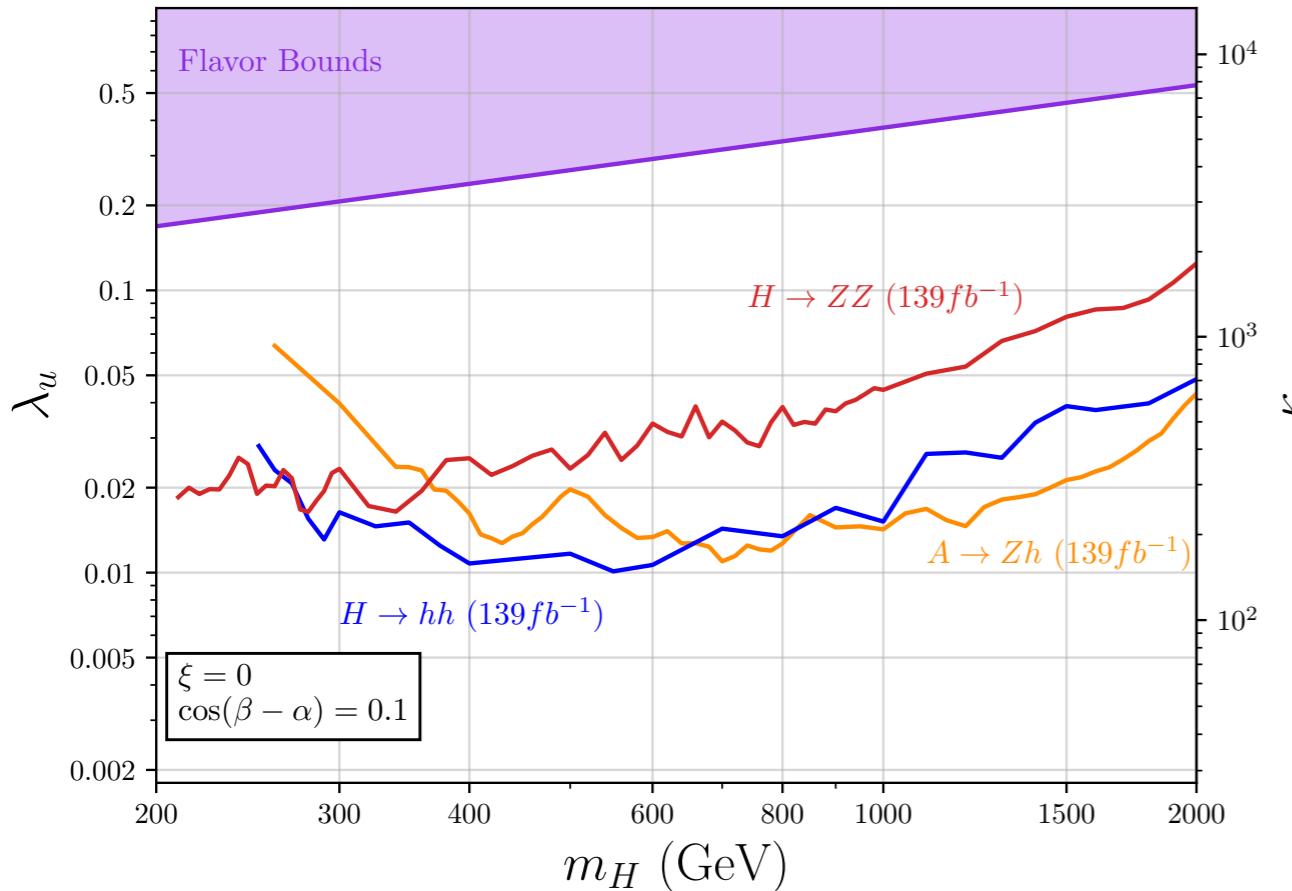
$$\mathcal{Y}_2^\ell = \xi^\ell y^\ell$$



FLAVOR PROTECTION IN DOWN-TYPE SFV ALLOWS NEW HIGGSES
 @ TEV TO COUPLE TO UP-TYPE QUARKS WITH ~ 0.1 STRENGTH.



Same outcome also for the up-type SFV 2HDM.



DOMINANT COLLIDER
BOUNDS FROM SEARCH
FOR NEW HIGGSES

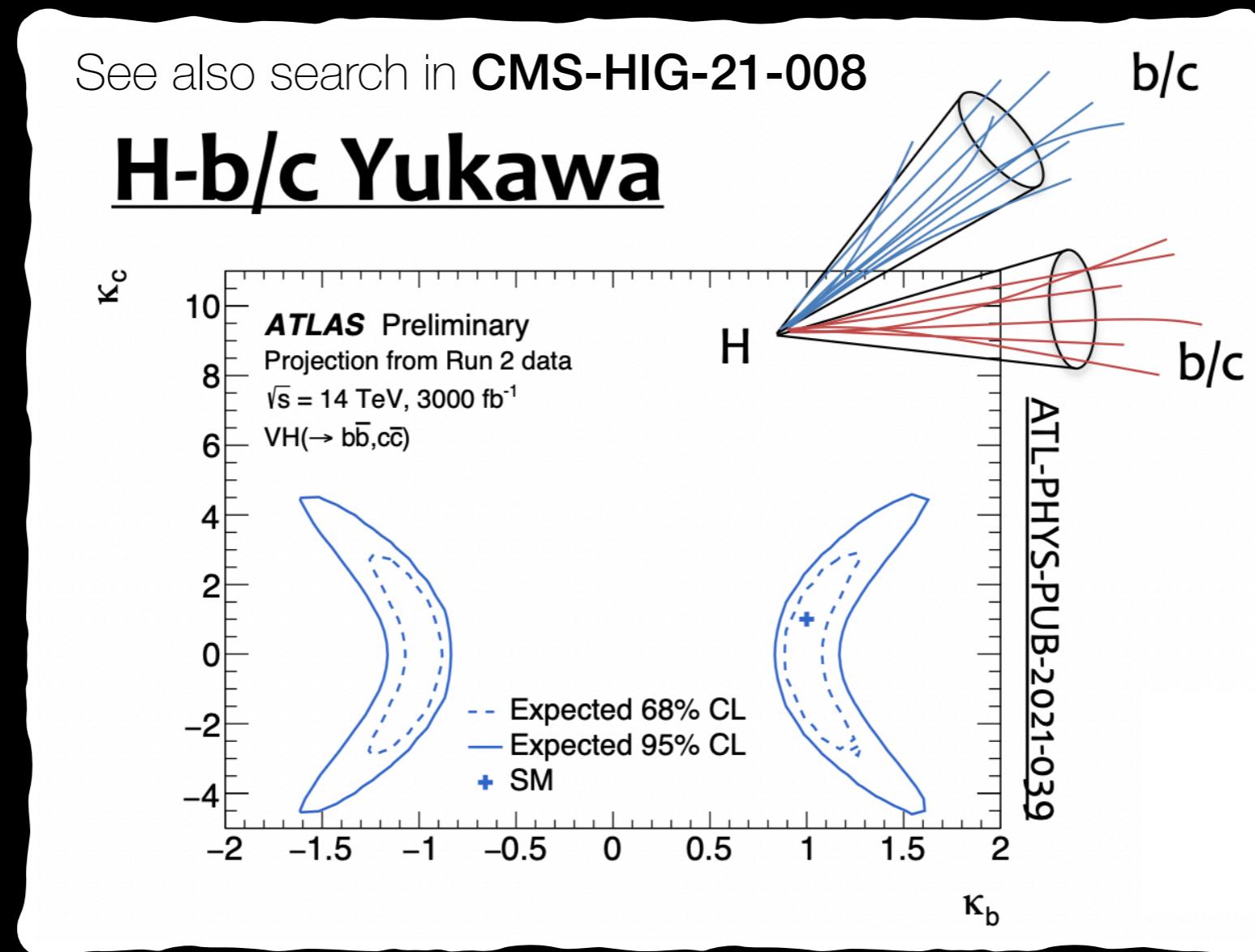
- Mainly from quark fusion, due to possibly large Yukawa
- Mostly into di-boson & di-Higgs due to non-zero mixing angle

OBS.

$\kappa_{u,c} \equiv y_{u,c}/y_{u,c}^{\text{SM}}$ is the Yukawa modifier

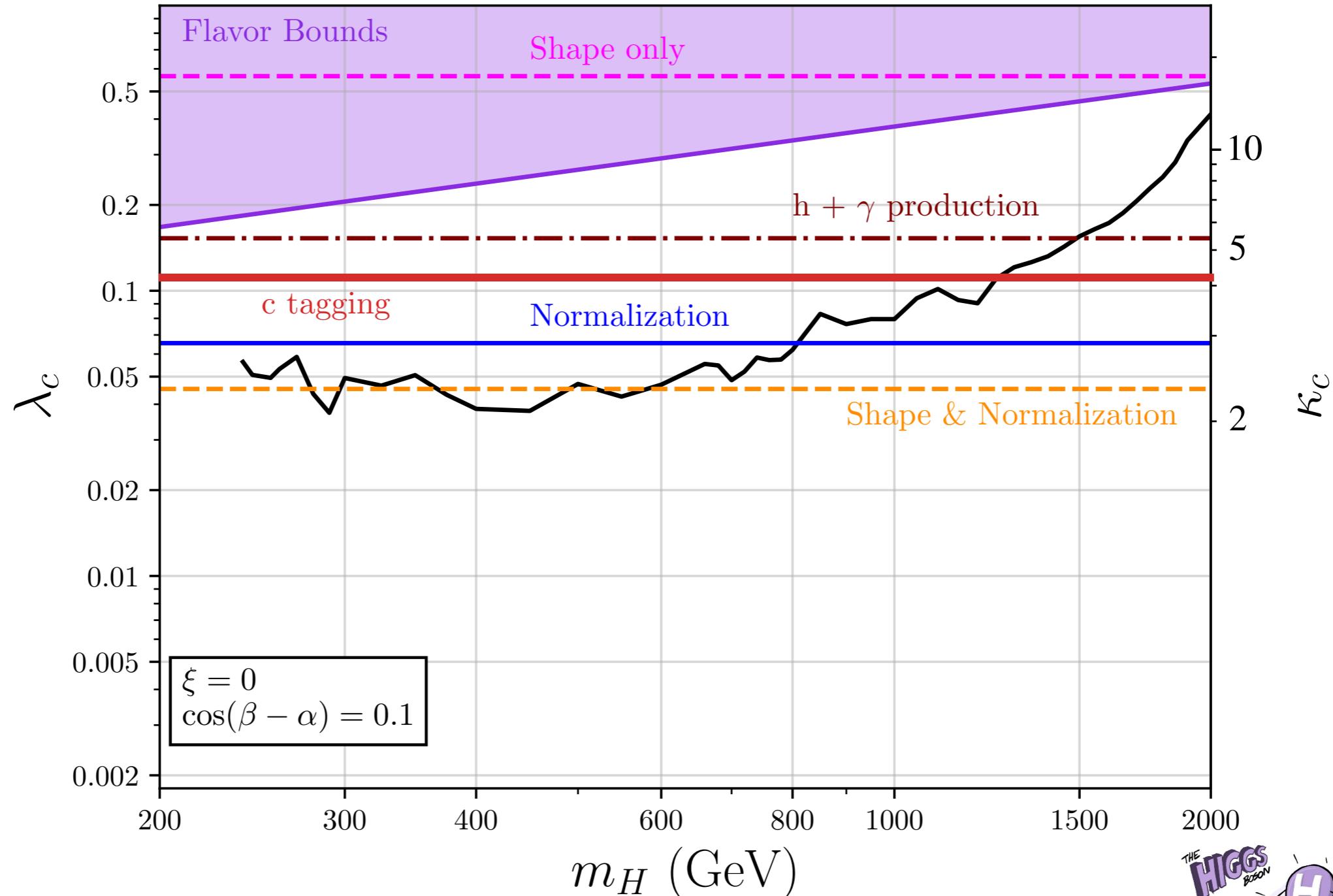
LHC probing Charm Yukawa!

Boosted topologies allow to disentangle rare signals from very large background ($H \rightarrow c\bar{c}$ 20x smaller than $H \rightarrow b\bar{b}$)



The advent of Deep Learning has been a game changer for ATLAS & CMS sensitivity to Yukawa coupling measurements.

Charming Higgs @ LHC



LHC **c tagging** turns out to be a remarkable probe of SFV ansatz



• Program of precision with flavor is still very rich,
but naively points to very high scales for NP.

No clear tension emerging from flavor data

• In absence of evidence for NP, the main direction
we pursued for BSM may need a paradigm shift.

E.g.: **MFV** hypothesis —> **SFV** ansatz

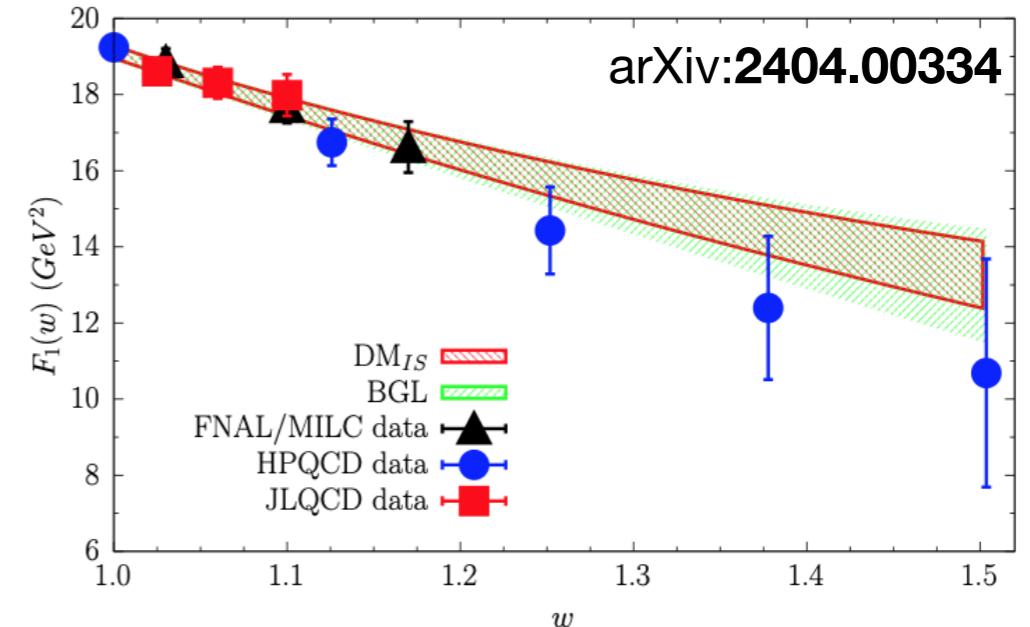
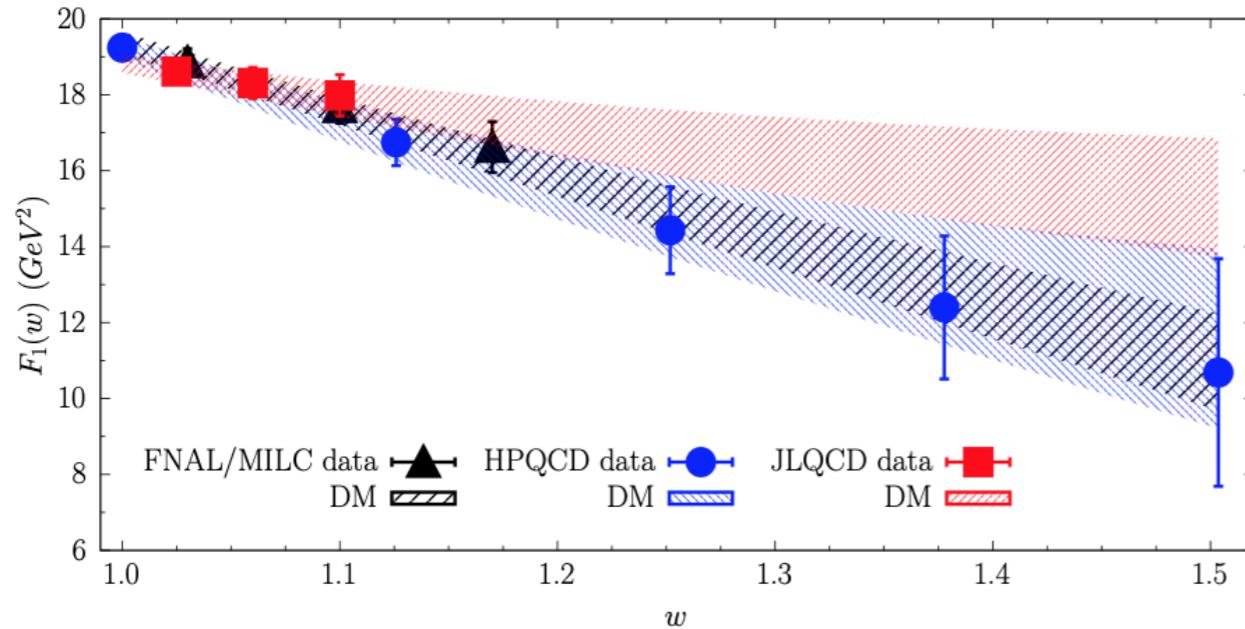
• Theory + EXP progress can lead to new targets.

Higgs Yukawa measurements insightful!

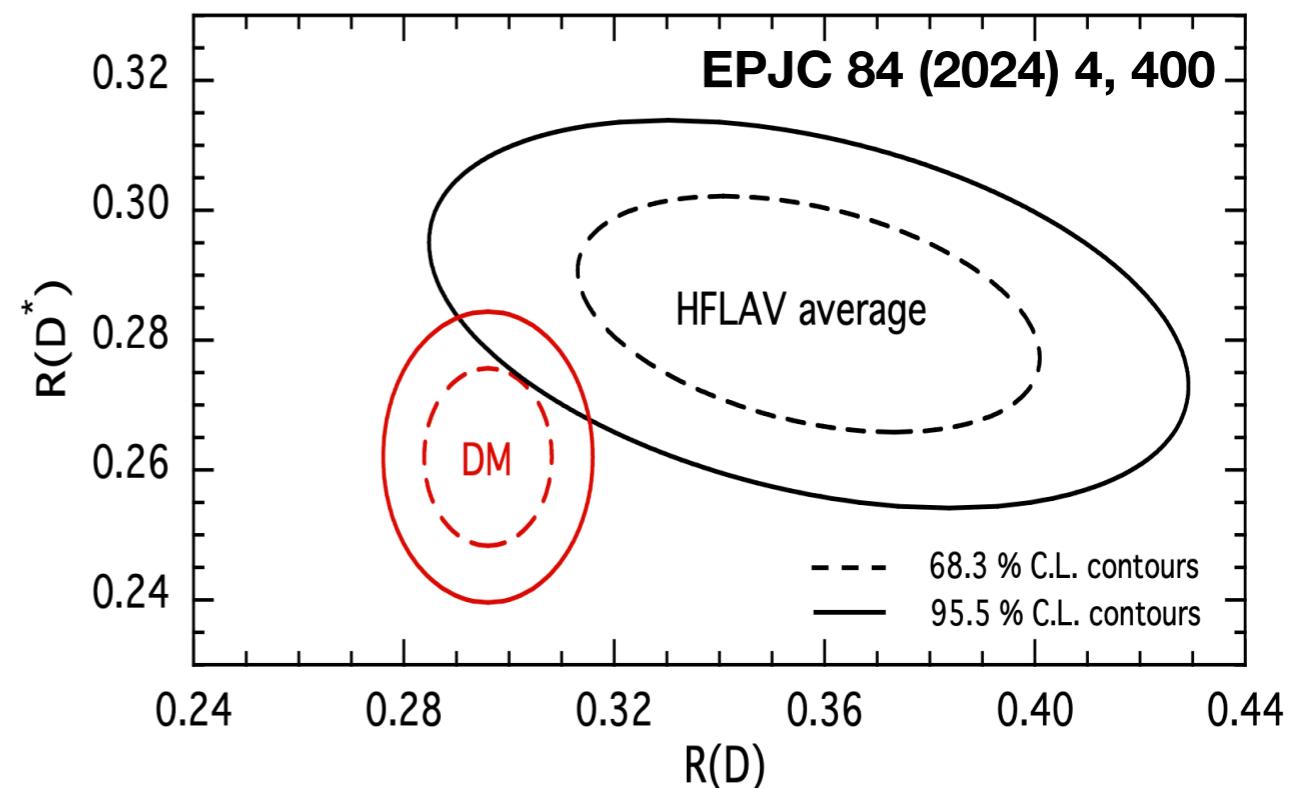
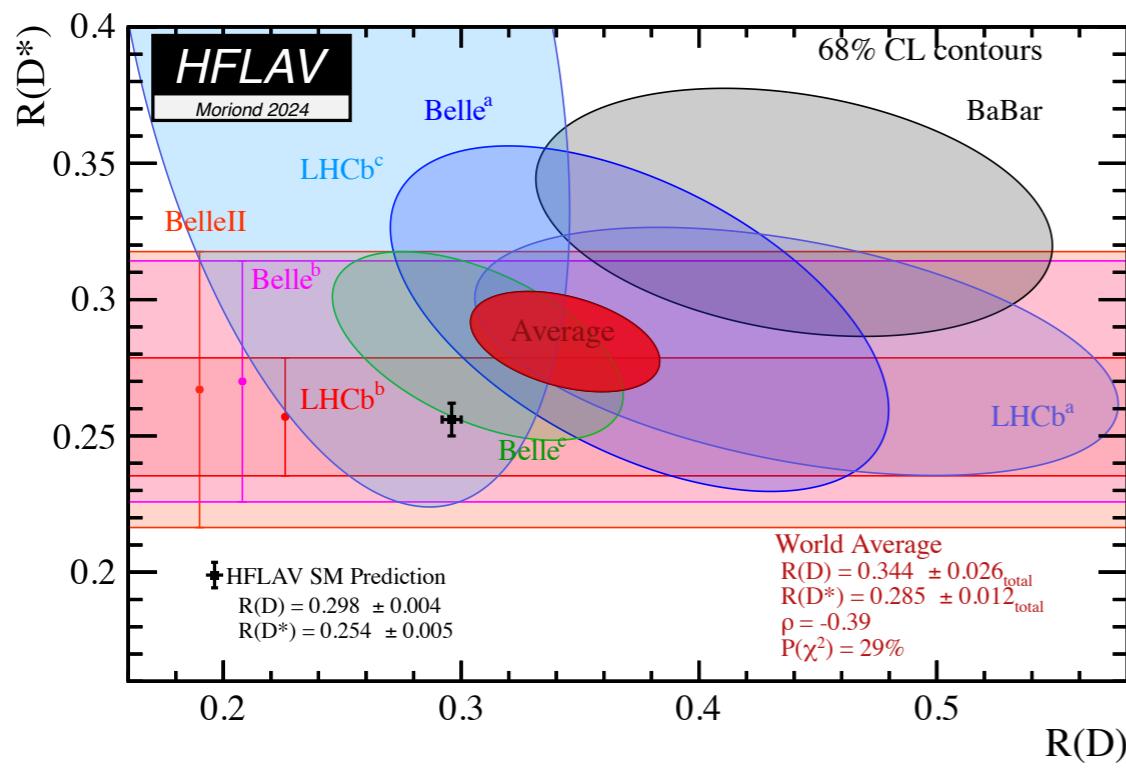
[See 2410.05236 and also 2410.08272]

BACKUP

A LOOK @ V_{cb}

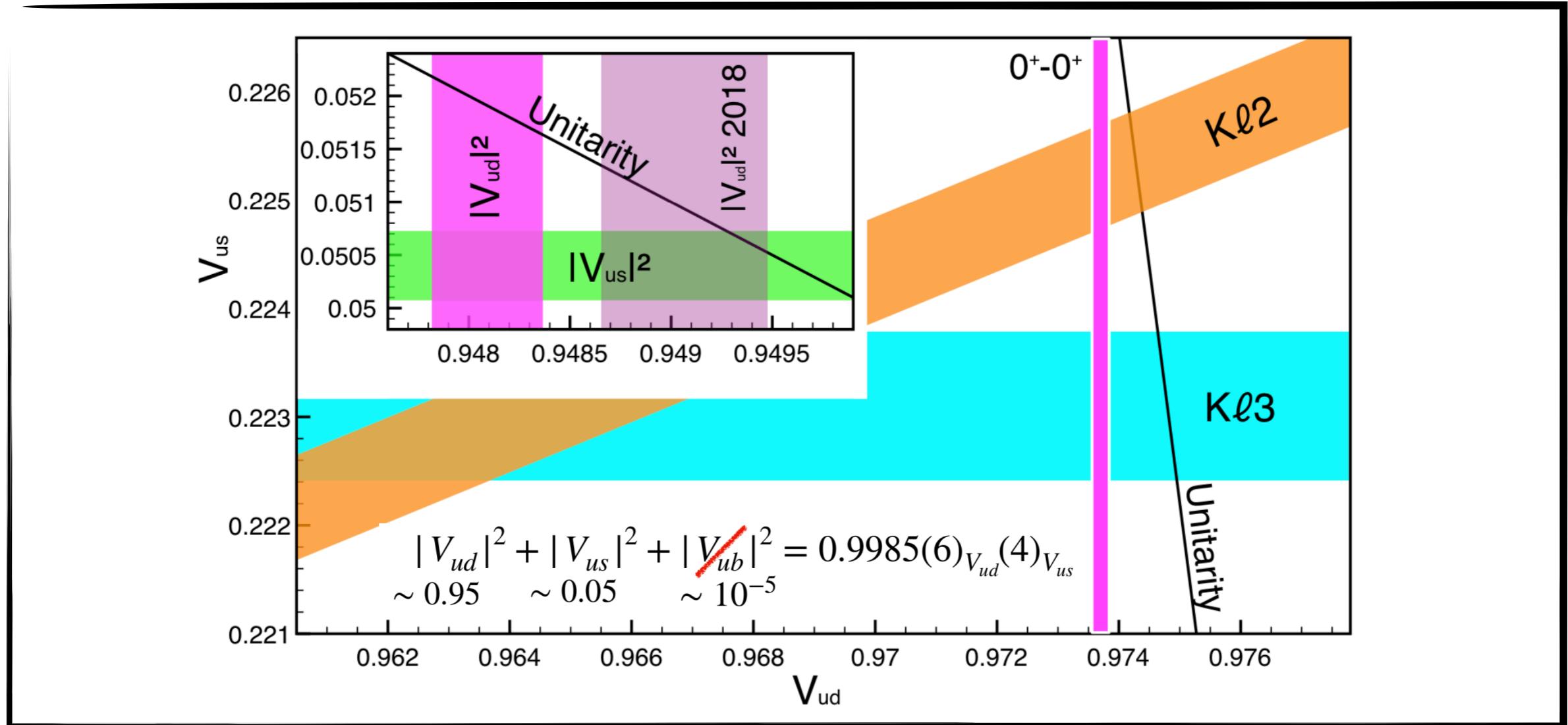


THE OUTCOME FOR THE OTHER 3 FORM FACTORS IMPACT ALSO $R_{D(\star)}$.



A LOOK @ 1st row

Misha Gorshteyn @ CKM 23



- LQCD on (semi)leptonic decays : Beyond % precision → **control of ΔI & QED**
See , e.g., *Phys.Rev.D 105 (2022) 11, 114507*
- 0⁺ → 0⁺ transitions “better” than neutron decay, **but $\pi^+ \rightarrow \pi^0 e^+ \nu$ cleanest though**
Interesting proposal: PIONEER – arXiv:2203.01981

UT NP Analysis — Wiki How

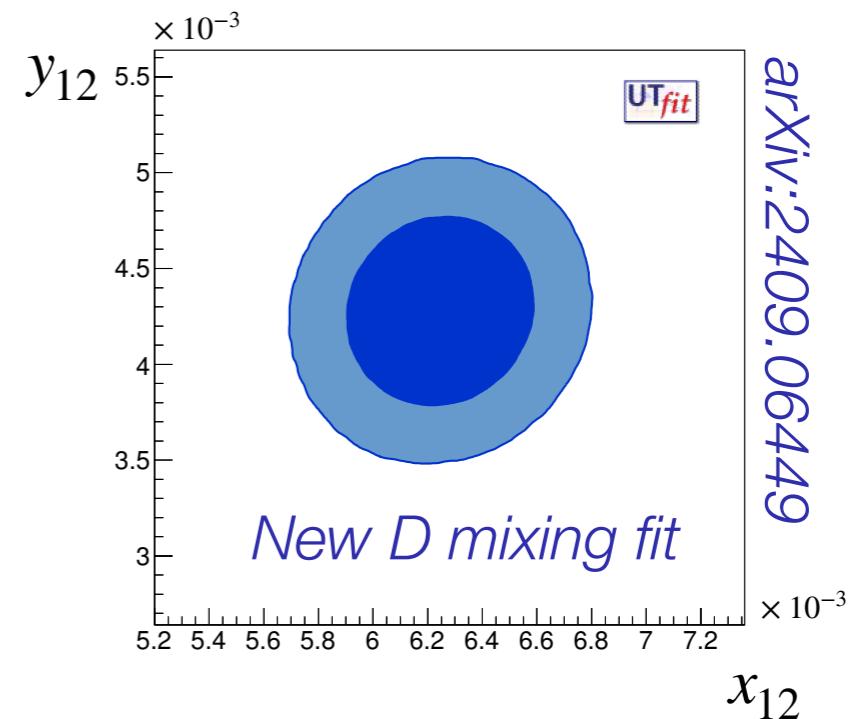
- Parametrize generic NP effects in $|\Delta F| = 2$ transitions:

$$A_q = C_{B_q} e^{2i\phi_{B_q}} A_q^{SM} e^{2i\phi_q^{SM}} = \left(1 + \frac{A_q^{NP}}{A_q^{SM}} e^{2i(\phi_q^{NP} - \phi_q^{SM})} \right) A_q^{SM} e^{2i\phi_q^{SM}}$$

- Include an extended list of observables to study also NP:

- same-side dilepton charge asymmetry
- semileptonic asymmetries in B^0 and B_s
- lifetime τ^{FS} in flavour-specific final states
- $\phi_s = 2\beta_s$ vs $\Delta\Gamma_s$ from $B_s \rightarrow J/\psi\phi$

+



- Fit simultaneously CKM & NP \rightarrow bound on NP scale
 $O(10)$ new parameters



: Vector-like quark

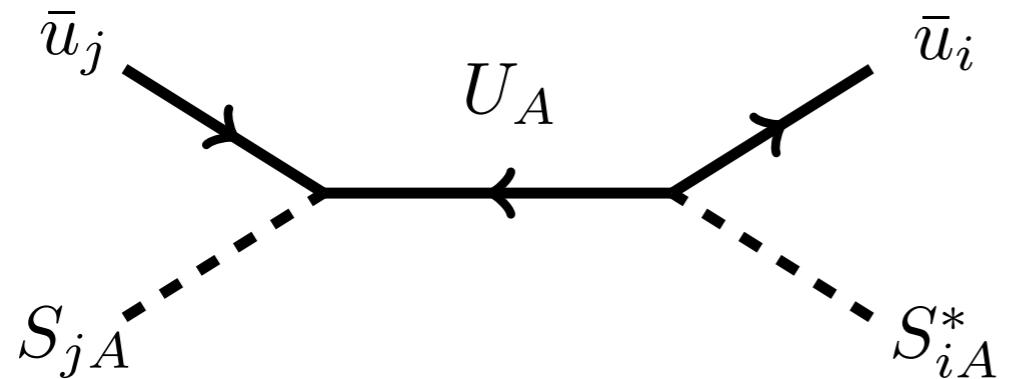
$$\mathcal{L} \supset M_{AB} U_A \bar{U}_B + \xi S_{iA} \bar{u}_i U_A$$

$$- [\eta_{ij}^u Q_i H \bar{u}_j - \eta_{ij}^d Q_i H^c \bar{d}_j + \text{h.c.}] + \mathcal{L}_{\text{BSM}}$$

No additional spurions/fields transforming under $U(3)_{\bar{u}}$

Introduce mixing between up-quark and heavy VLQs in a flavor breaking vacuum

	$U(3)_U$	$U(3)_{\bar{U}}$	$U(3)_{\bar{u}}$	$U(1)_B$	\mathbb{Z}_2
U	3			$1/3$	-1
\bar{U}		3		$-1/3$	-1
S	$\bar{3}$		$\bar{3}$		-1



$$Z_{ij}^u = \delta_{ij} + \frac{\xi^* \xi}{M_A^* M_A} S_{iA}^* S_{jA}$$

Integrating out heavy quarks leads to wave-function renormalization of the SM up-quarks

The source of all flavor-breaking!
CKM matrix arises from returning to canonical basis

S.Homiller@ppc2021

IN GENERAL, TWO CLASSES OF SFV: UP-TYPE AND DOWN-TYPE