

Probing Cosmology with Einstein Telescope

Angelo Ricciardone

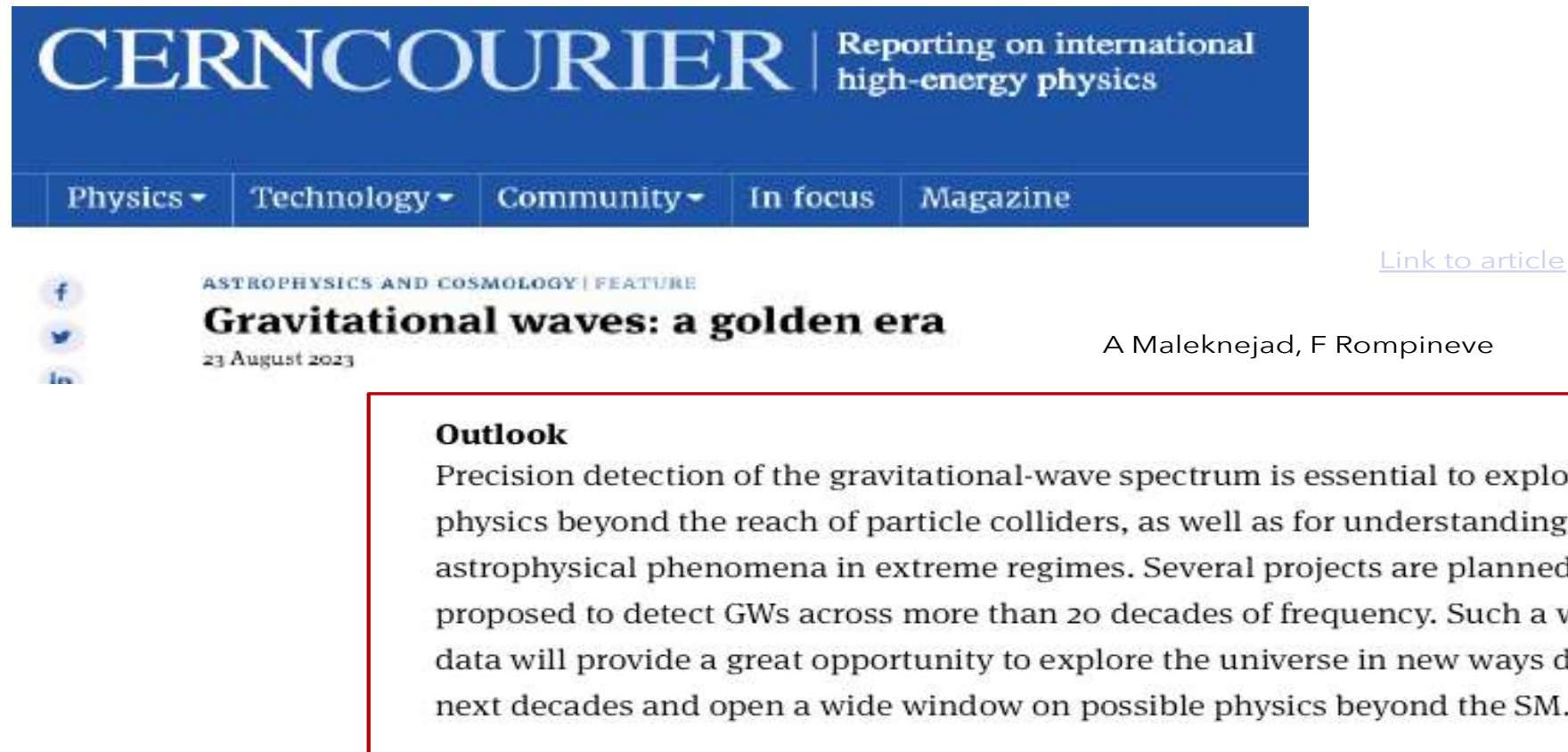
Department of Physics “E. Fermi”
University of Pisa

INFN - Sezione di Pisa



Latest Updates

- In June 2021, the European Strategy Forum on Research Infrastructures (ESFRI) decided to include the Einstein Telescope (ET) in the update of its roadmap for 2021.
- In June 2022 formal establishment of ET collaboration (today ~ 1400 members)
- In June 2023 Italian government present the Italian candidacy to host the Einstein Telescope
- On 25 January 2024 LISA has been adopted by ESA and construction will start in January 2025
- On 25 March 2024 ASI presented the LISA mission to the scientific community



The image is a screenshot of a CERN Courier article. The header is blue with 'CERN COURIER' in white and 'Reporting on international high-energy physics' in smaller white text. Below the header is a navigation bar with links: Physics, Technology, Community, In focus, and Magazine. The article title is 'Gravitational waves: a golden era' in bold black text, with 'ASTROPHYSICS AND COSMOLOGY | FEATURE' above it and '23 August 2023' below it. The authors 'A Maleknejad, F Rompineve' are listed to the right. A red box highlights the 'Outlook' section, which discusses the precision detection of the gravitational-wave spectrum and its importance for exploring particle physics and astrophysical phenomena.

CERN COURIER | Reporting on international high-energy physics

Physics ▾ Technology ▾ Community ▾ In focus Magazine

ASTROPHYSICS AND COSMOLOGY | FEATURE

Gravitational waves: a golden era

23 August 2023

A Maleknejad, F Rompineve

[Link to article](#)

Outlook

Precision detection of the gravitational-wave spectrum is essential to explore particle physics beyond the reach of particle colliders, as well as for understanding astrophysical phenomena in extreme regimes. Several projects are planned and proposed to detect GWs across more than 20 decades of frequency. Such a wealth of data will provide a great opportunity to explore the universe in new ways during the next decades and open a wide window on possible physics beyond the SM.

Where we are - LVK

The third observing run (B) from April 2019 to March 2020

Total number of gravitational waves observed to date (with probability of astrophysical origin > 0.5): ~ 90

(mostly BBHs, 2 BNS and 2 NS-BH)

GWTC-3 catalogue: arXiv:2111.03606

The run O4a May 2023 - January 2024 (GW230529)

The fourth run O4b+c has started in April with Virgo online

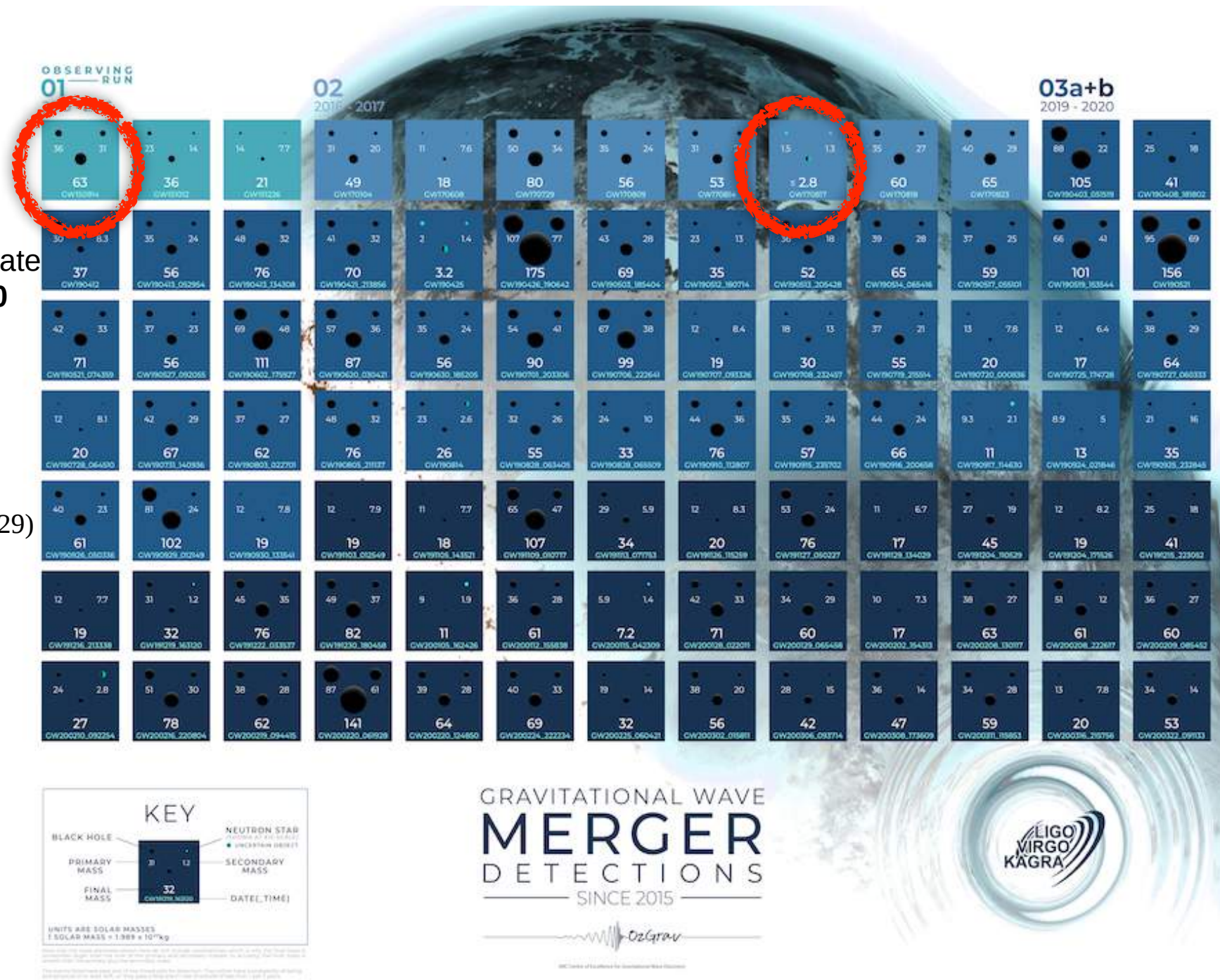
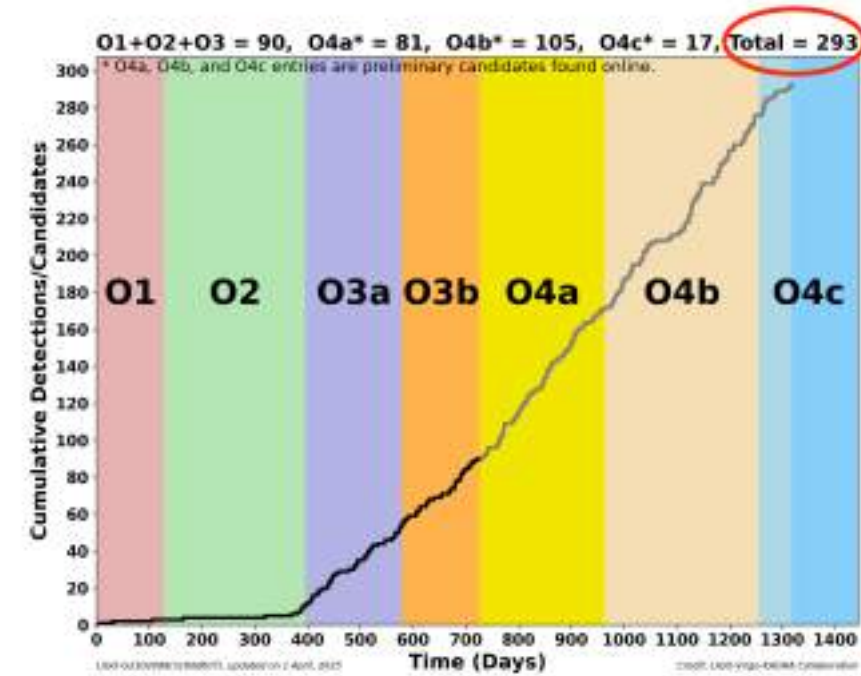
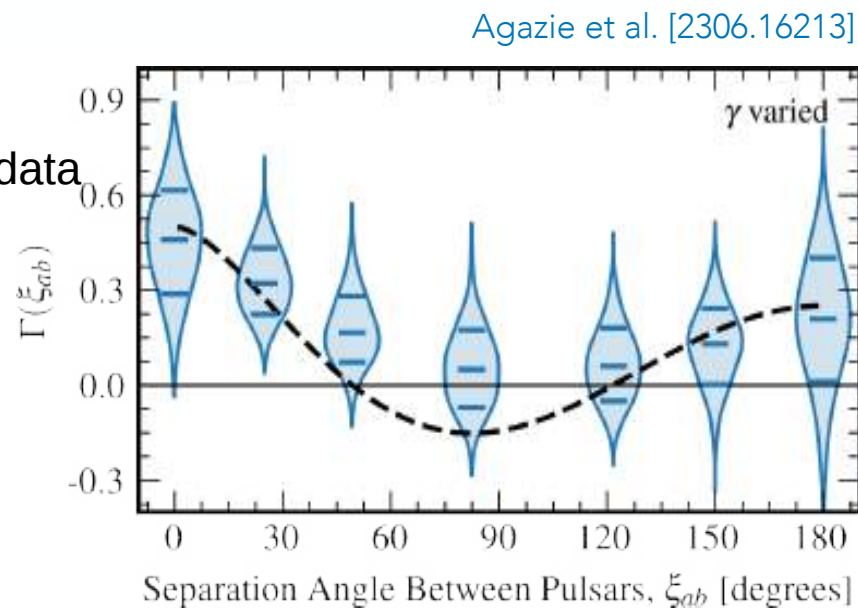


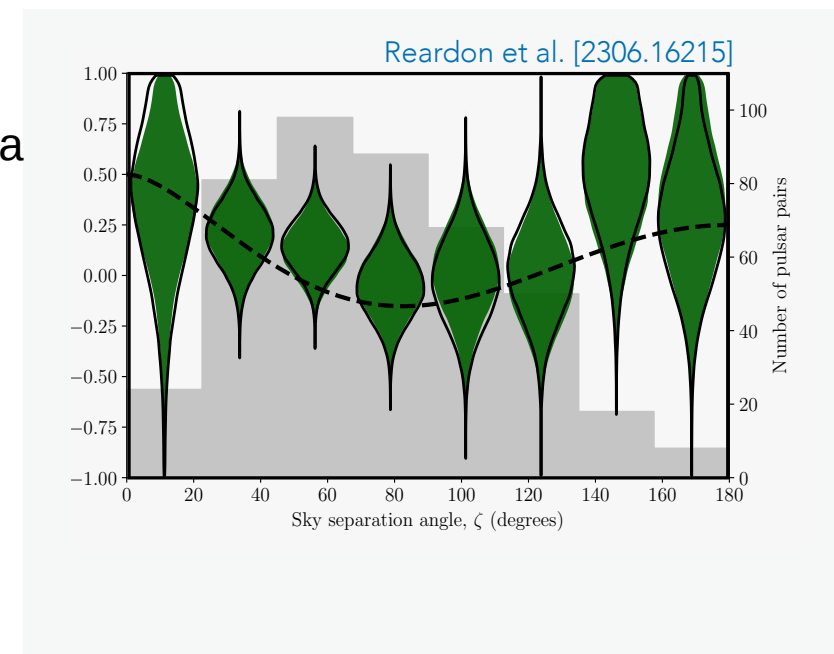
Image credit: LIGO / Virgo / KAGRA / C. Knox / H. Middleton

Where we are - PTA

NANOGRav:
68 pulsars, 16 yrs of data
 $\sim 3 - 4\sigma$ significance

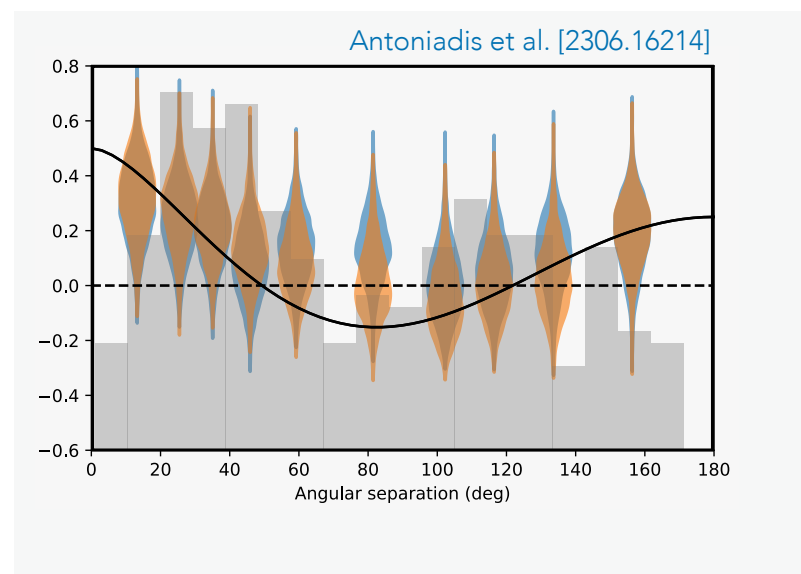


PPTA:
32 pulsars, 18 yrs of data
 $\sim 2\sigma$ significance

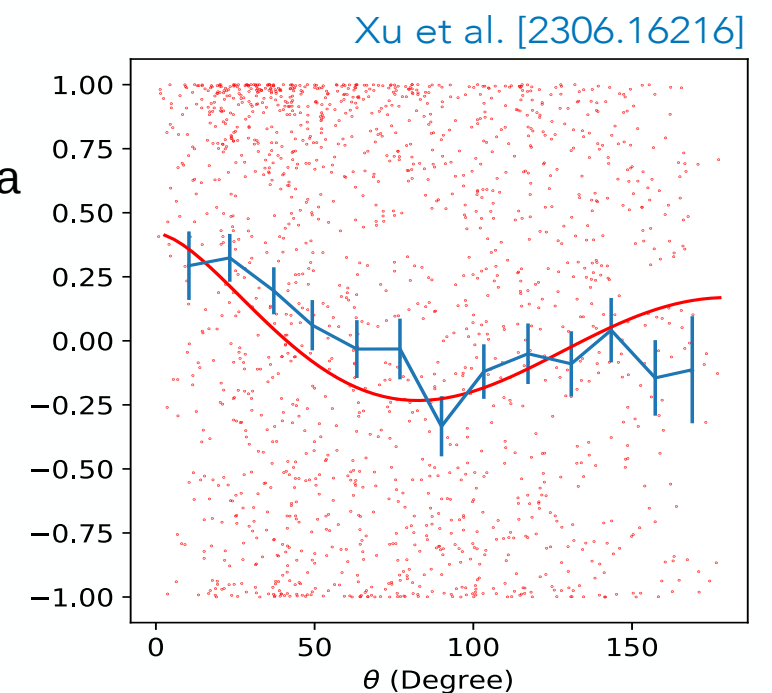


First Claim of Detection of a SGWB

EPTA+InPTA:
25 pulsars, 24 yrs of data
 $\sim 3\sigma$ significance

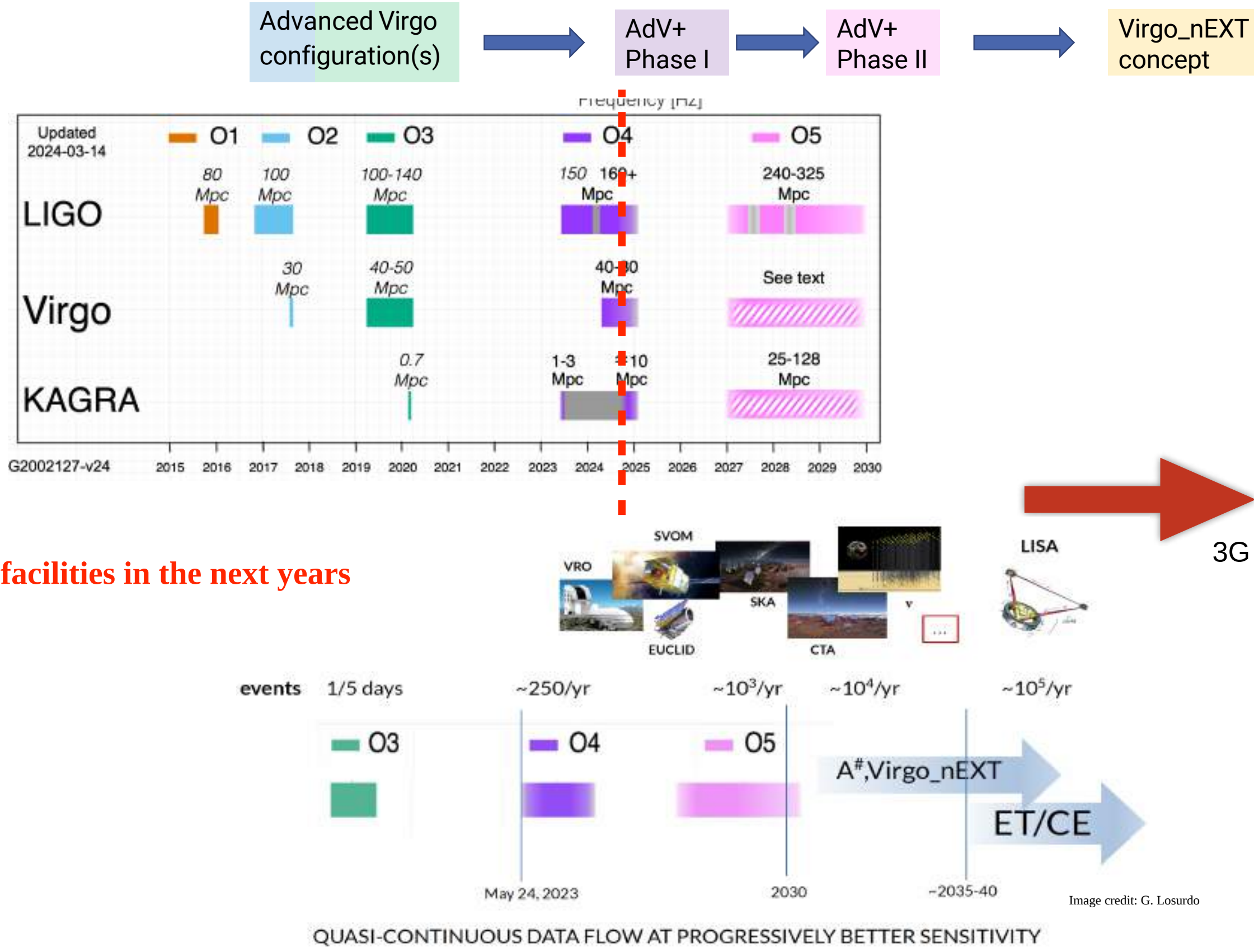


CPTA:
57 pulsars, 3 yrs of data
 $\sim 4.6\sigma$ significance

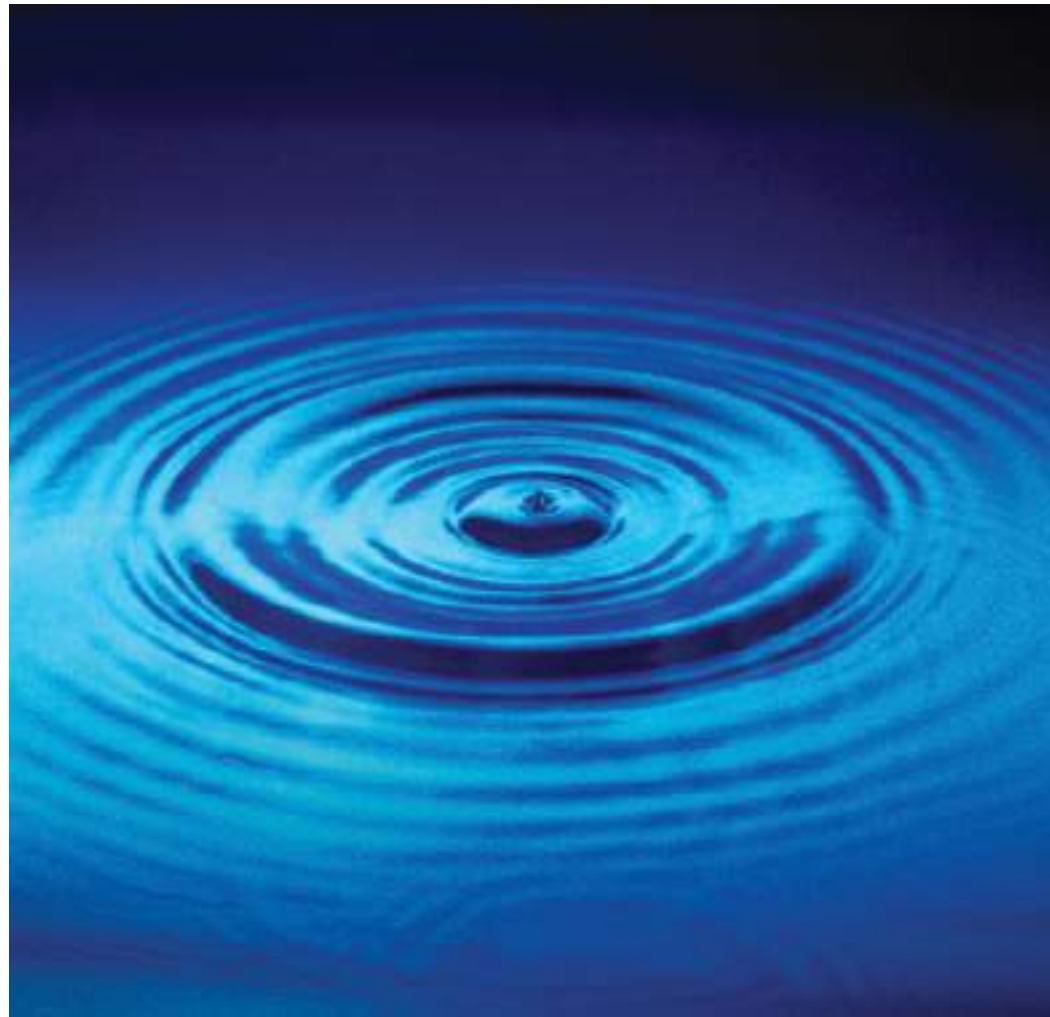


Bayesian reconstruction of normalized inter-pulsar correlations
Violins plot = marginal posterior densities (plus median and 68% credible values)

What are the plans for the LVK runs?



What are Gravitational Waves?



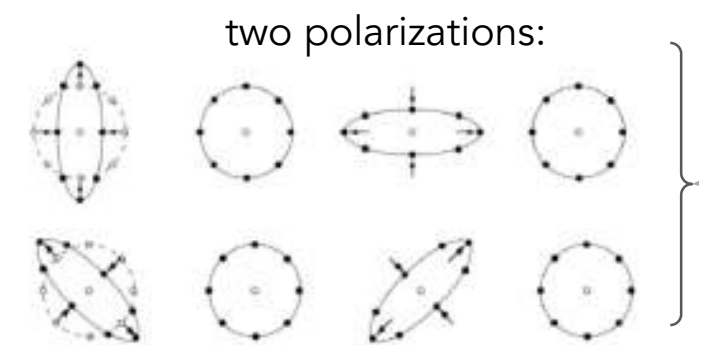
Rock in a pond

$$g_{\mu\nu} = \underbrace{\eta_{\mu\nu}}_{\text{flat}} + \underbrace{h_{\mu\nu}}_{\text{perturbation}}$$

linearized theory

$$\square \bar{h}_{\mu\nu} = \frac{16\pi G}{c^4} T_{\mu\nu}$$

wave equation for the perturbative metric term



What are Gravitational Waves?



Rock in a pond

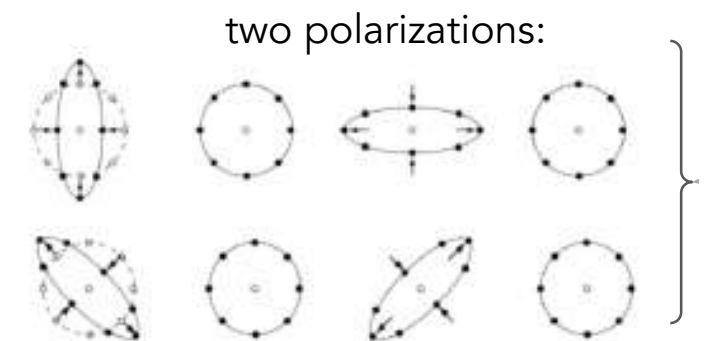
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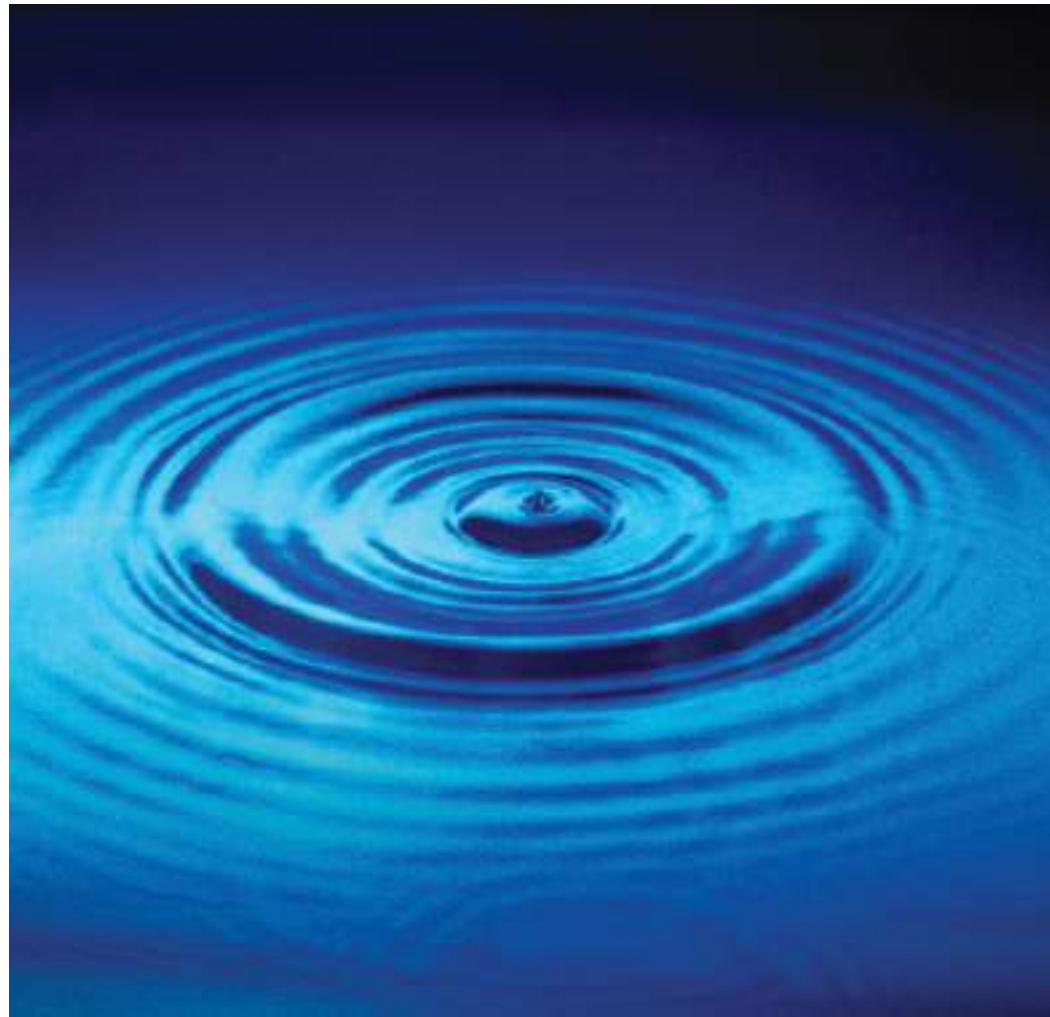
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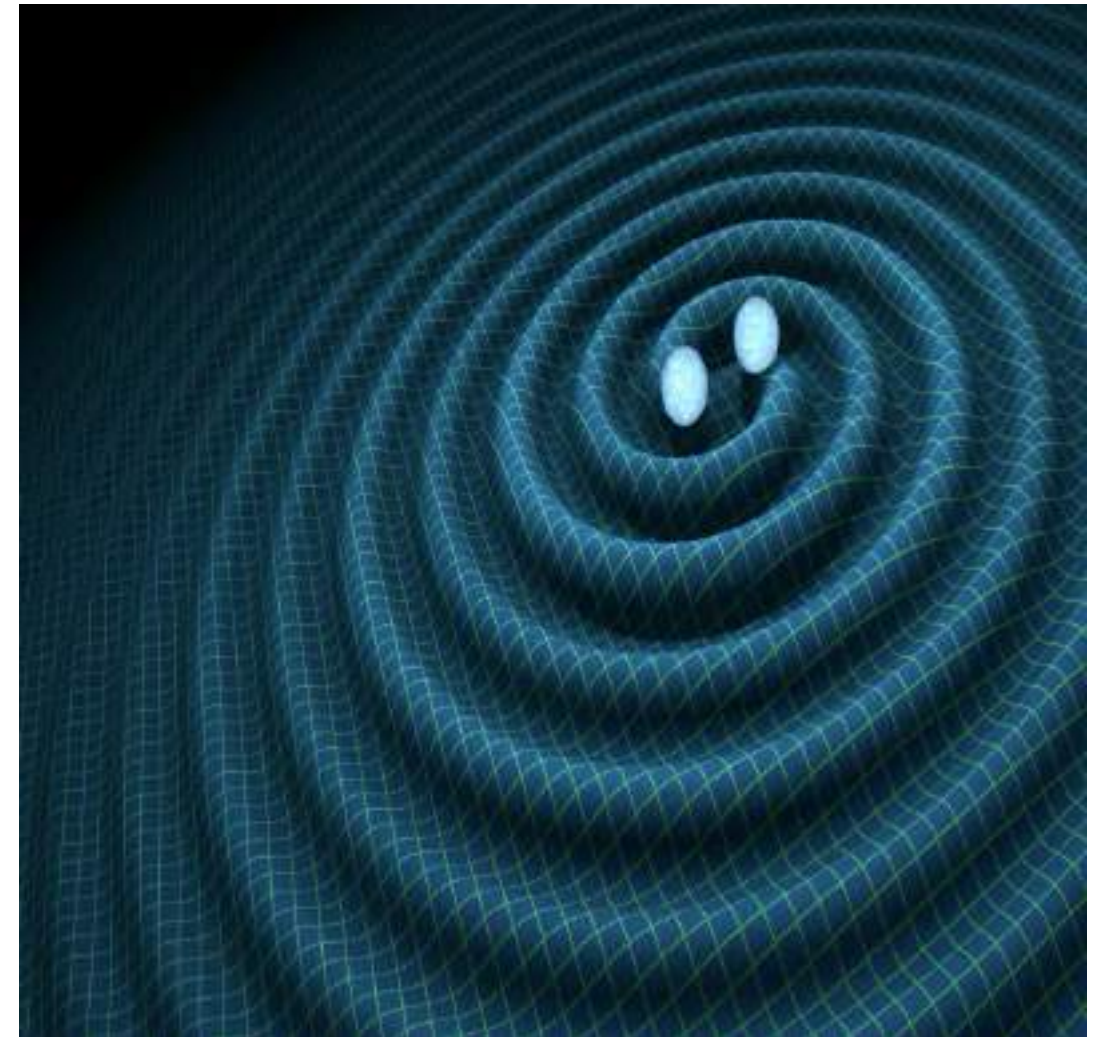
Merging COs in our universe



What are Gravitational Waves?



Rock in a pond



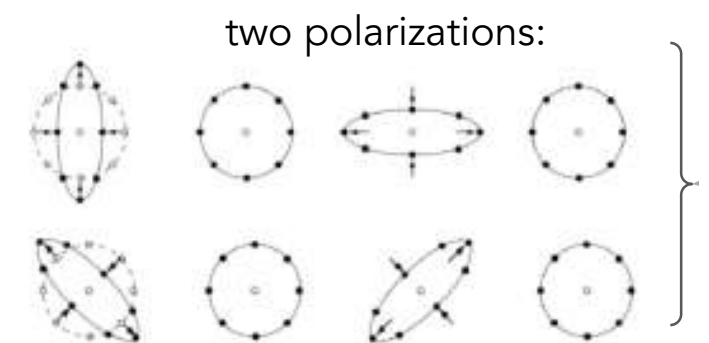
Merging COs in our universe

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linearized theory

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Gravitational Waves:

waves of space–time curvature
that accelerate free–falling particles



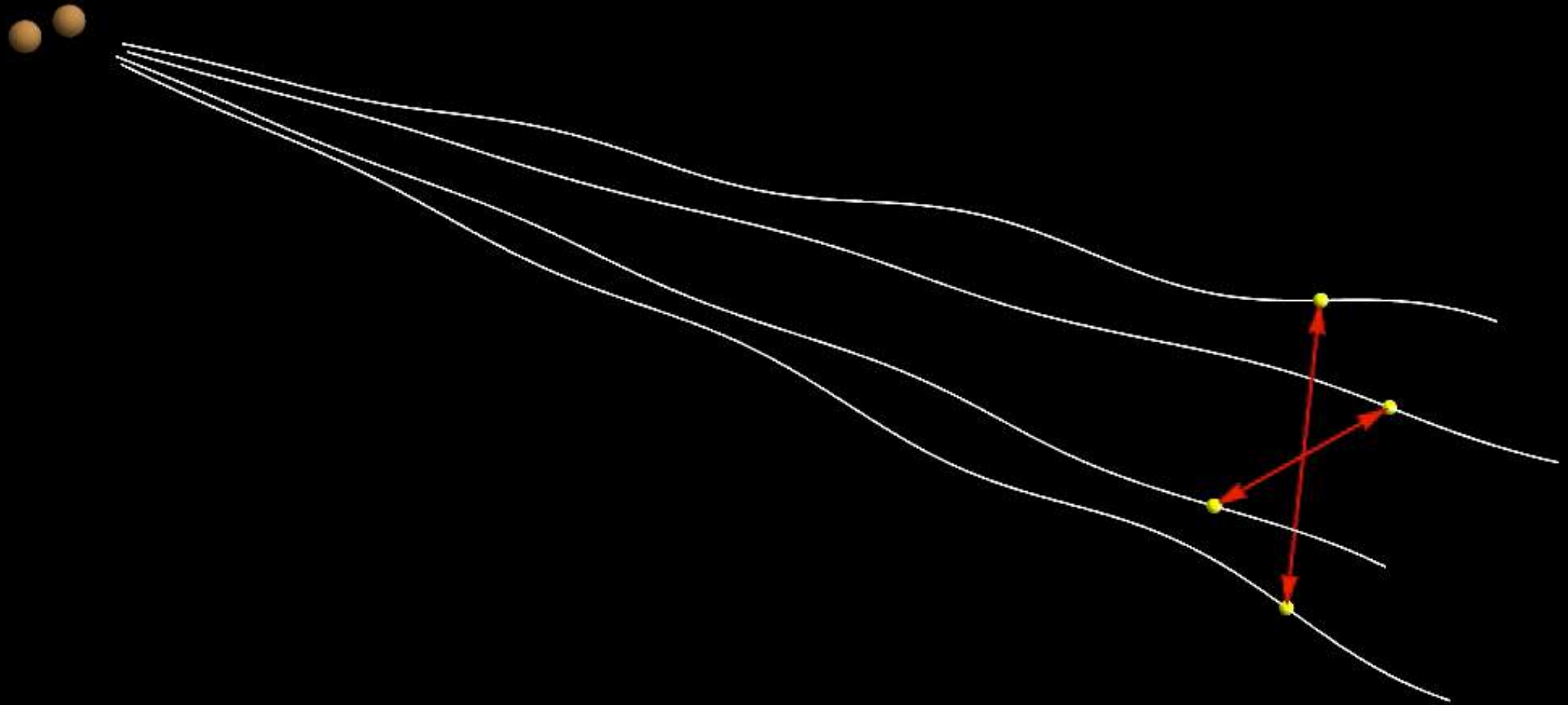
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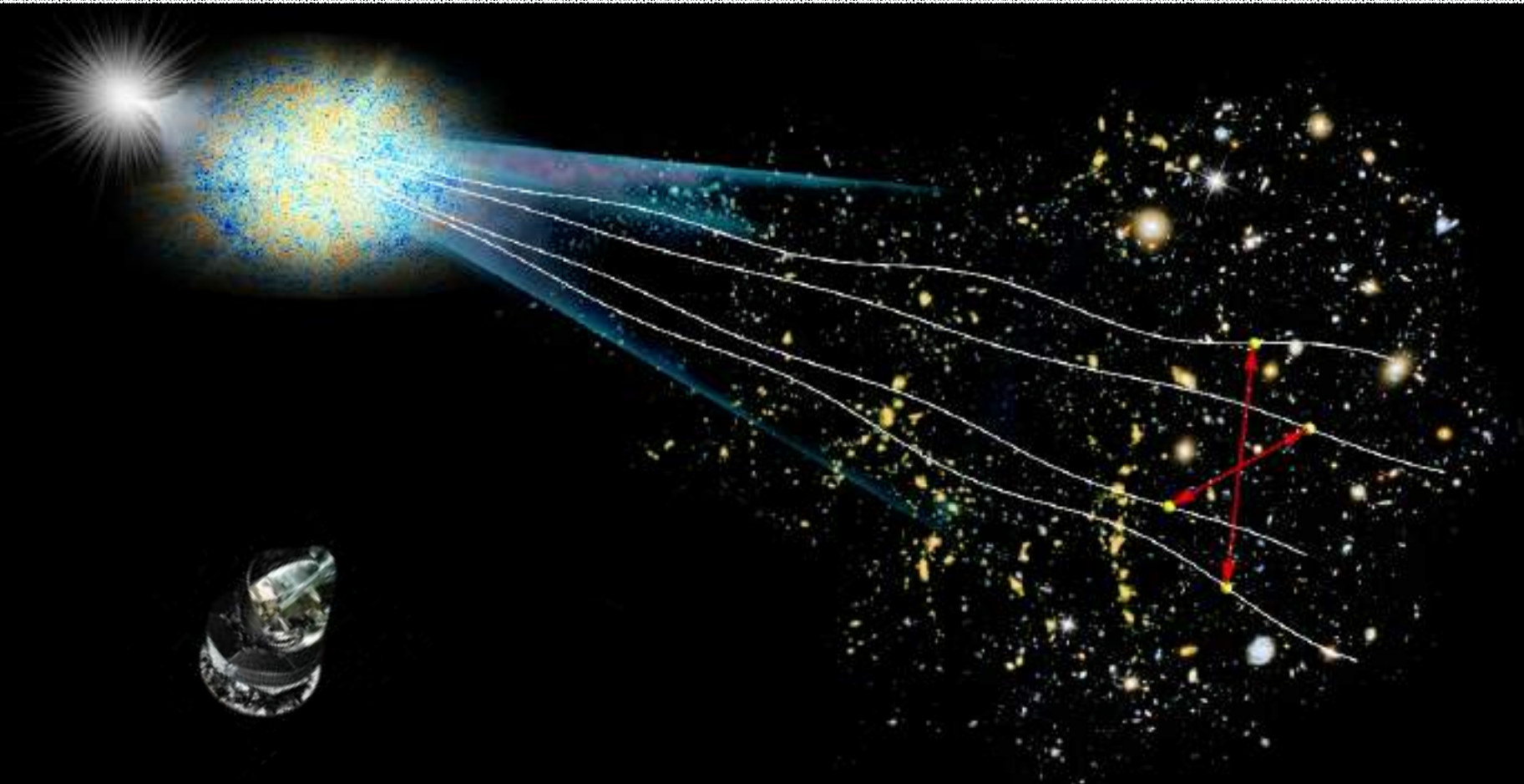


Gravitational Waves:

waves of space-time curvature
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“Indirect” vs Direct GW detection



**Polarization of CMB photons
through Thomson scattering
of electron and photon**

Only Tensor perturbations
can source B-mode

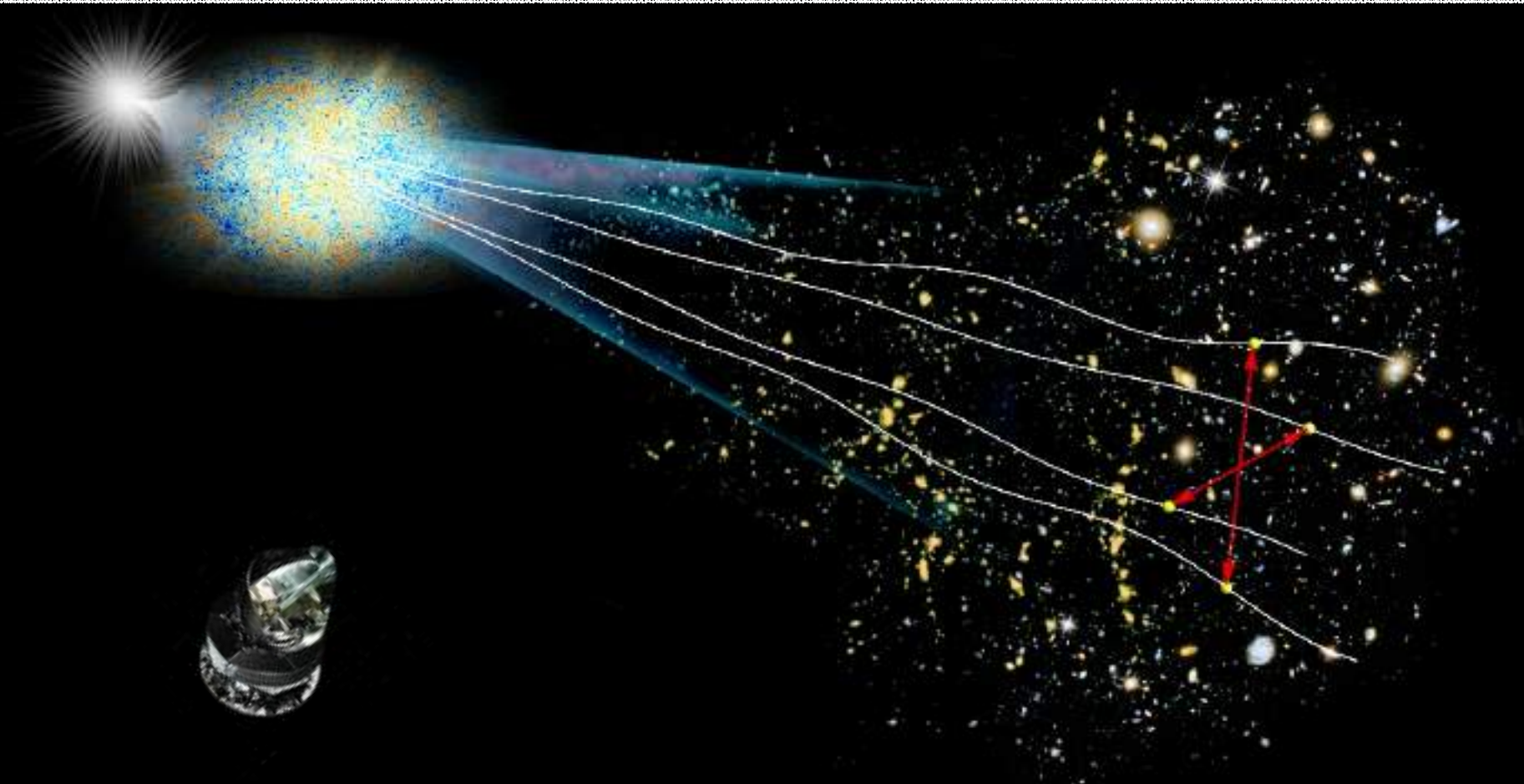
Poor and contaminated signal:

- foregrounds
- gravitational lensing (E→B at small scales)

**Distortion of space as GW
passes detector arms**

- ground-based
- space-based
- pulsar timing arrays

“Indirect” vs Direct GW detection



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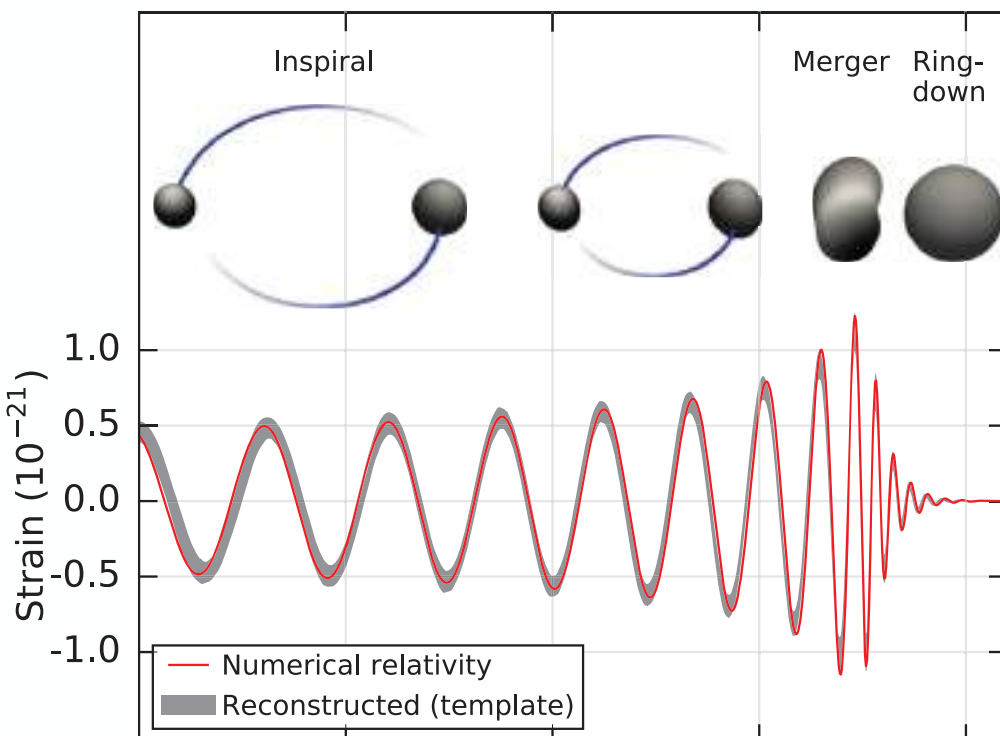
**Distortion of space as GW
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- space-based
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Sources of Gravitational Waves

Resolved Sources:

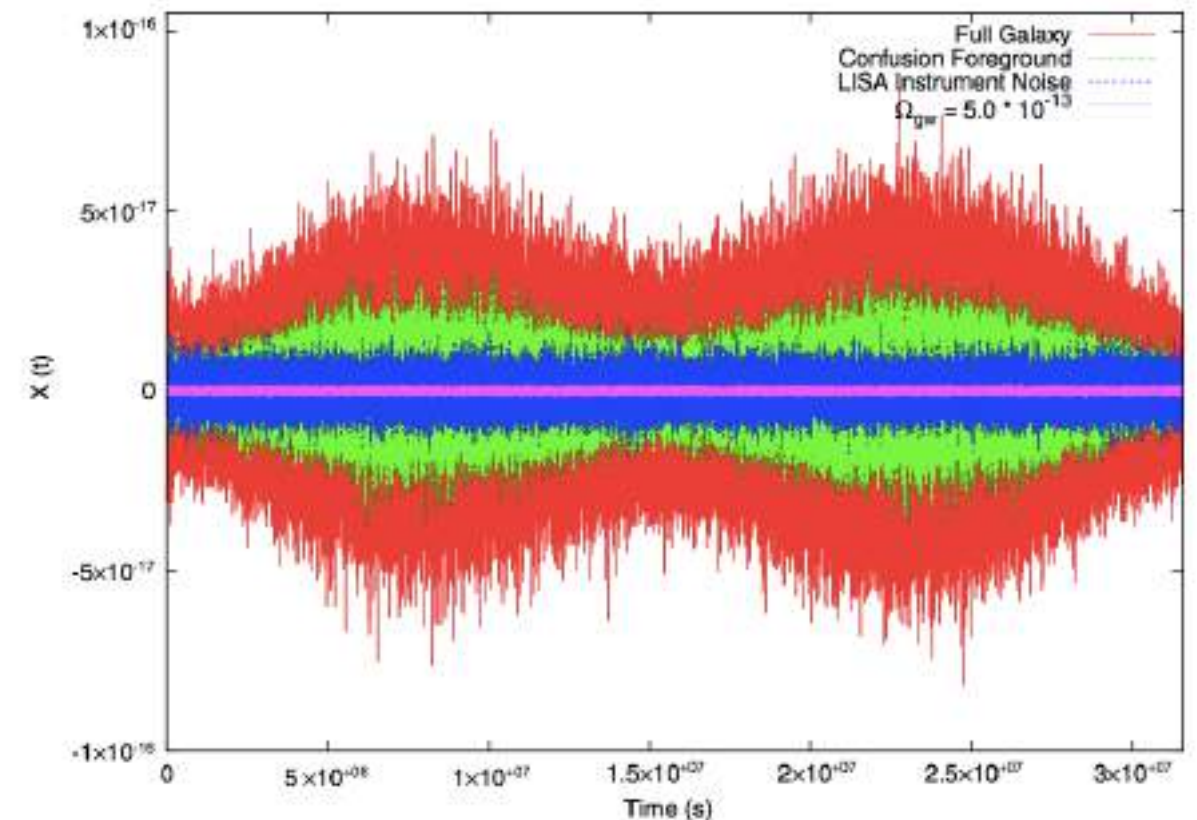
- Black Holes
- Neutron Stars
- White Dwarfs
- Supernovae
- ...



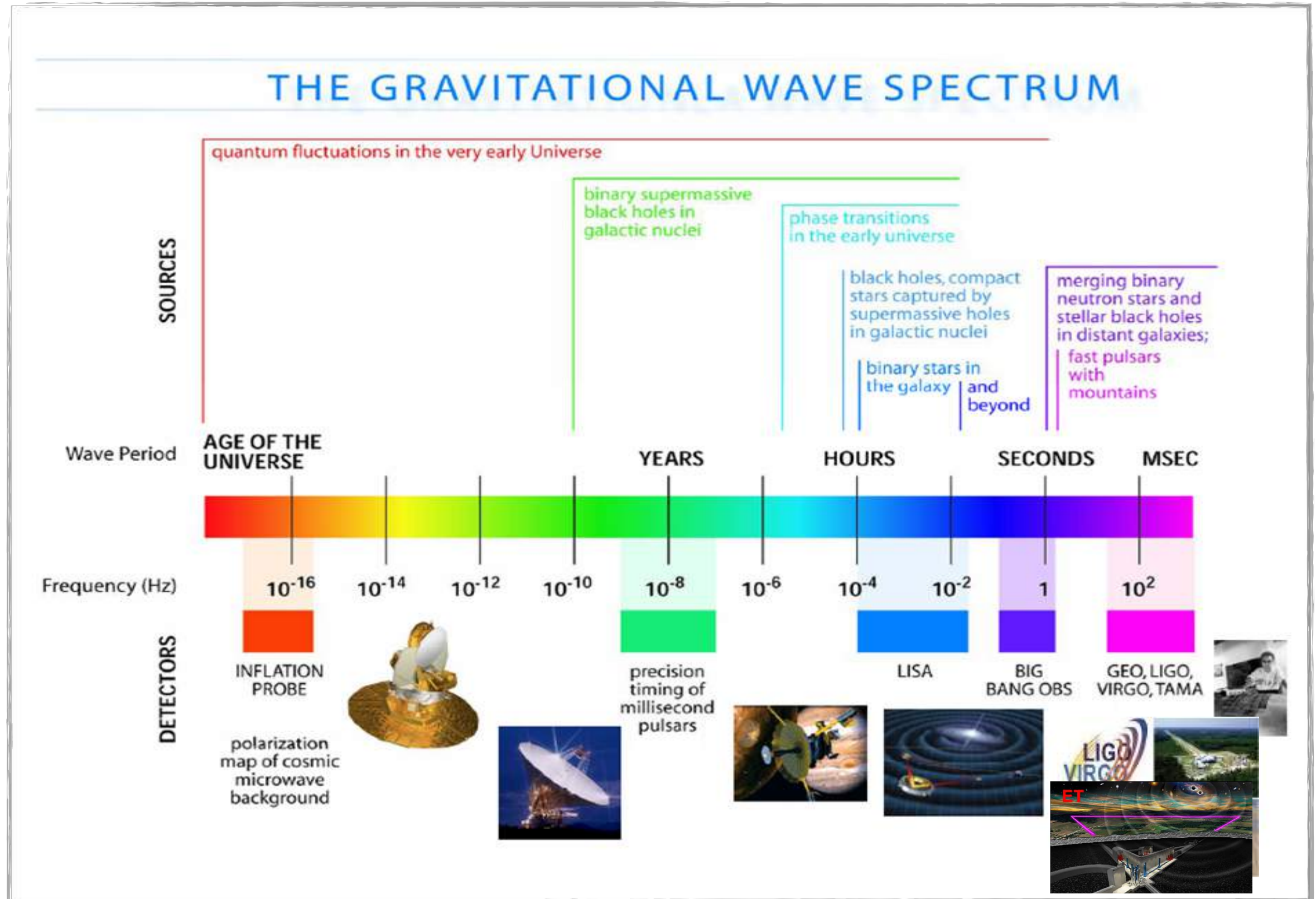
Unresolved Sources:

Stochastic Backgrounds

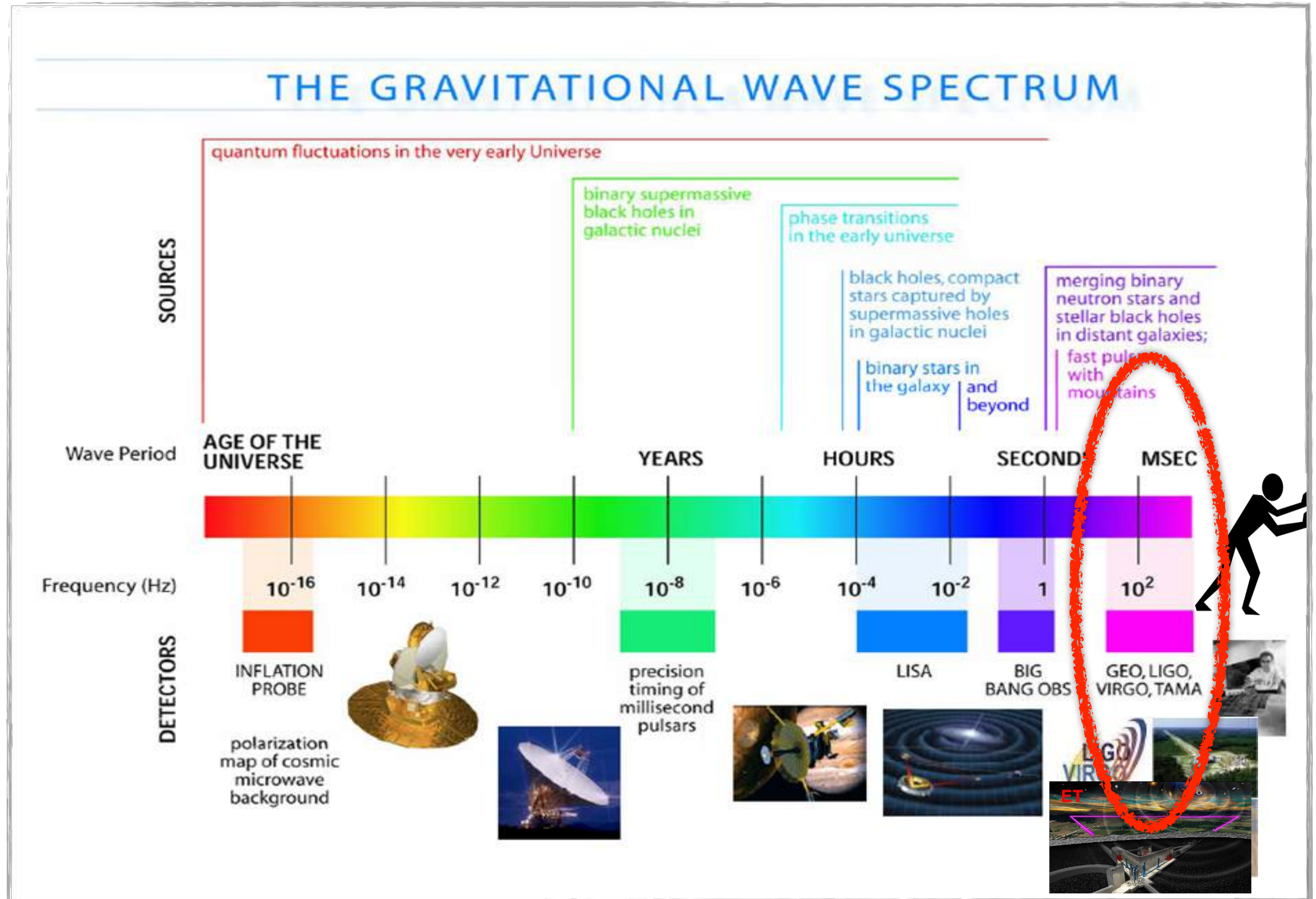
- Astrophysical
- Cosmological



The Gravitational Wave Spectrum

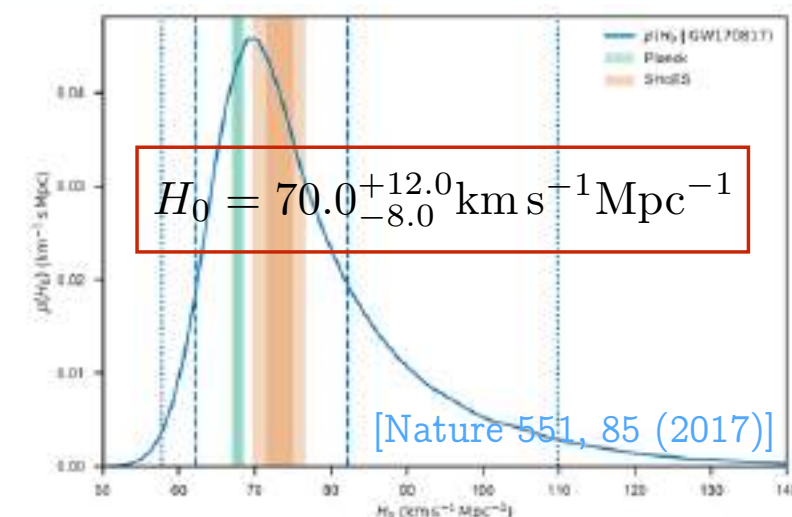


The Gravitational Wave Spectrum



What we have learned about Cosmology so far?

- We obtained the first measurement of the Hubble constant using GWs



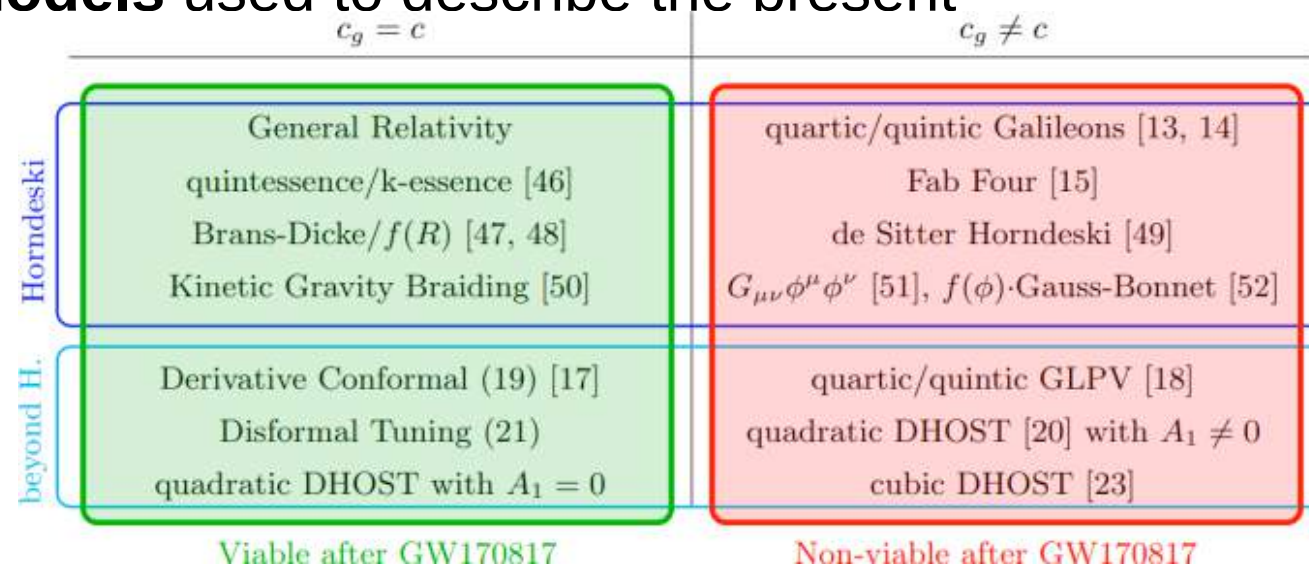
- We have tested and bounded deviations from GR (e.g., graviton mass, post-Newtonian coefficients, modified dispersion relations, etc.)

$$m_g \leq 1.27 \times 10^{-23} \text{ eV}/c^2$$

- The speed of GWs is the same as the speed of light

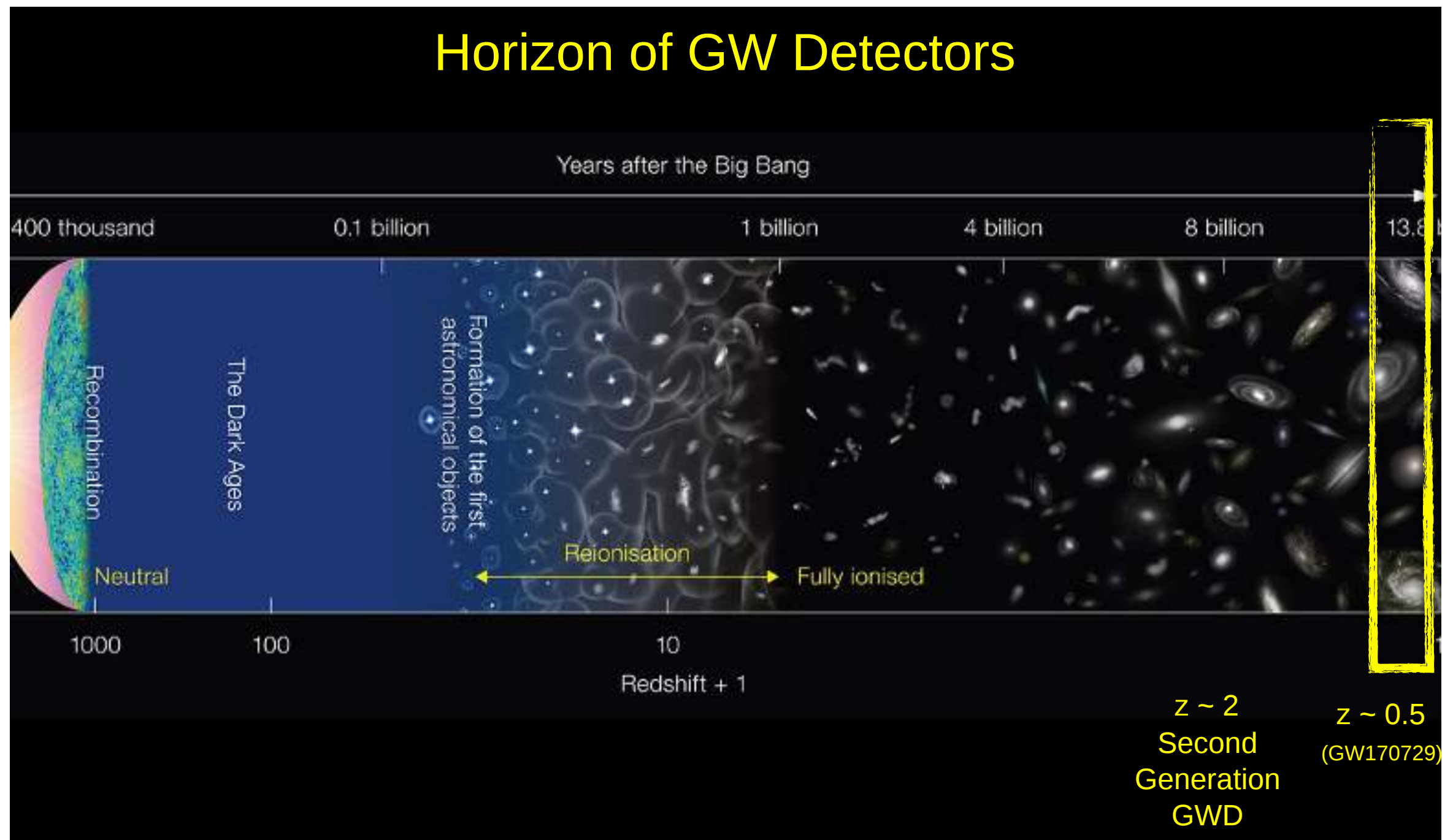
$$-3 \cdot 10^{-15} \leq c_g/c - 1 \leq 6 \cdot 10^{-16}$$

- We have ruled out many Modified Gravity models used to describe the present acceleration of our Universe



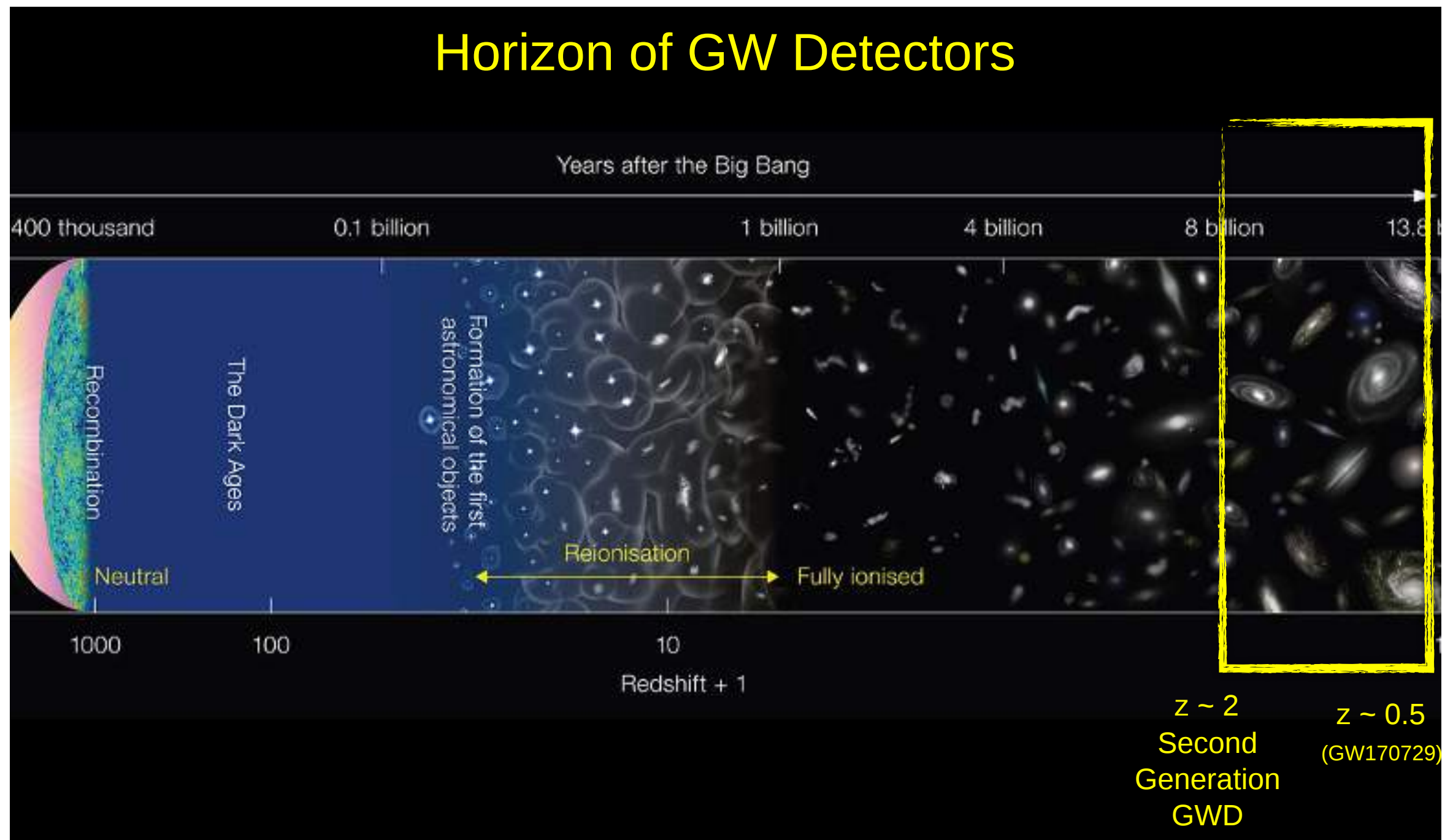
[LVK arxiv:2112.06861](https://arxiv.org/abs/2112.06861)

Where is the horizon?



Adv LIGO - Virgo - KAGRA:
BBHs only up to $z \sim 2$
BNSs in the very local Universe

Where is the horizon?



Adv LIGO - Virgo - KAGRA:
BBHs only up to $z \sim 2$
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Open questions in Cosmology

Early Universe

Primordial GWs

Cosmic Structure

Dark Matter

Primordial BHs?

Axion particles

Late Universe

Dark Energy

Modified Gravity?

Cosmic Tensions

H_0 measurements?

Isotropy

Testing
Cosmological
Principle

RESOLVED SOURCES

SGWB

Open questions in Cosmology

Early Universe

Primordial GWs

Late Universe

Dark Energy

Modified Gravity?

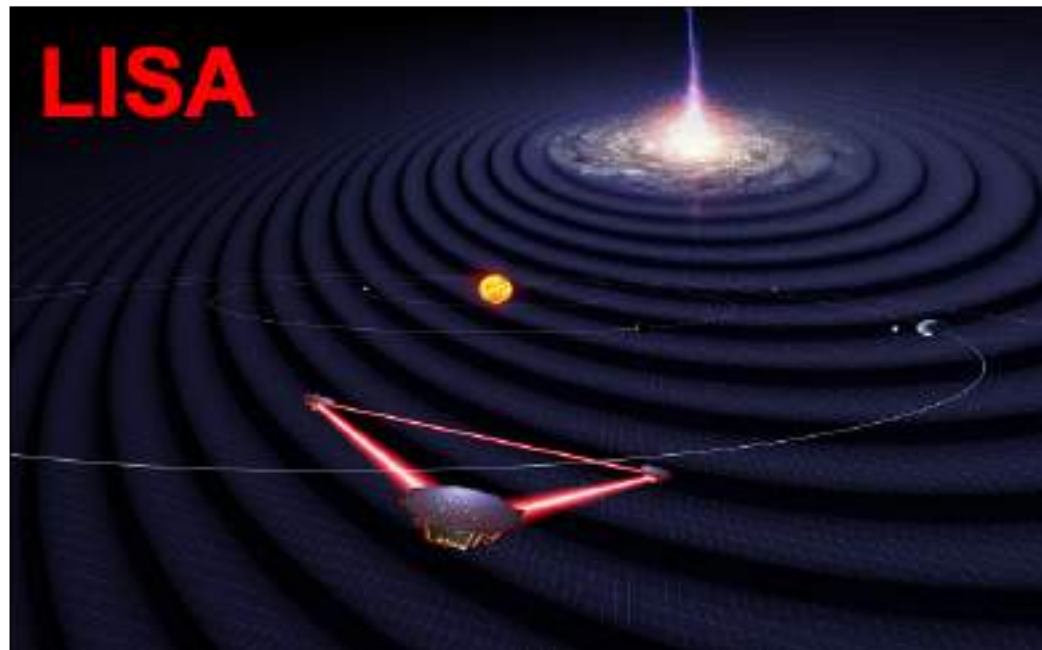
Cosmic Tensions

H_0 measurements?

RESOLVED SOURCES

SGWB

Prospects for next Generation GW Interferometers

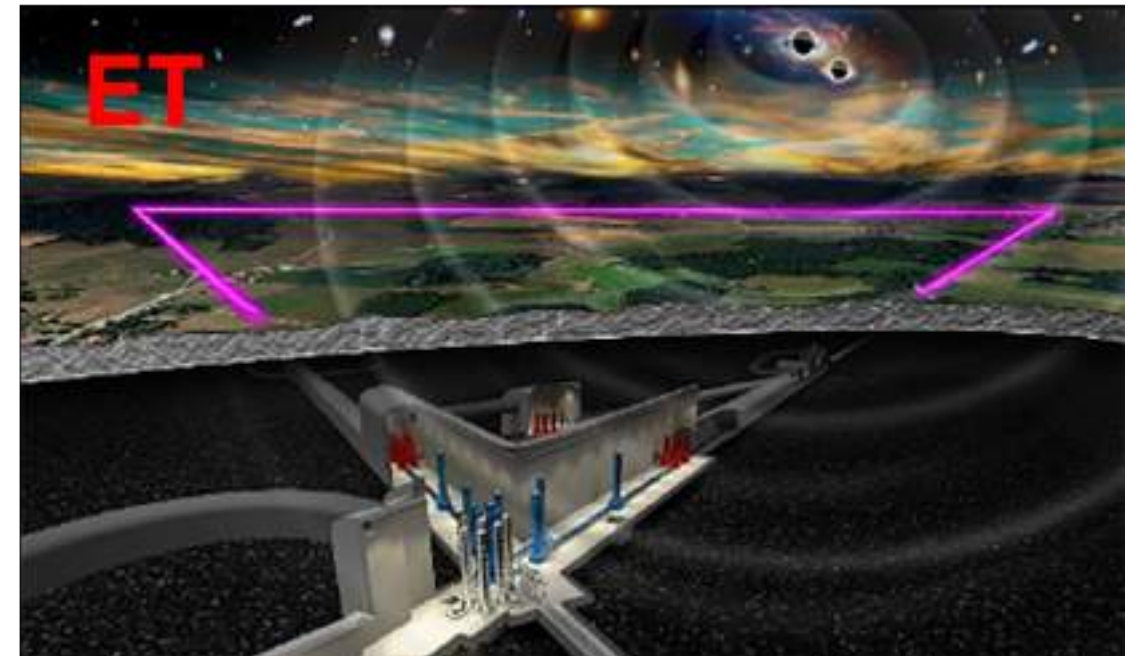


Geometry: **Constellation of 3 spacecraft in an equilateral configuration** (a giant interferometer)

Mission duration: **4 y science mission**
10 y nominal mission

Arm Length: **2.5 million km**

Expected Launch: 2034



Geometry: **Ground-based Triangular detector (HF+LF)**

Arm Length: **10 km**

Expected to be operative in: 2034

ET collaboration officially launched

+ CE, DECIGO, BBO, Taiji, TianQin, etc

AdV -> Virgo-nEXT

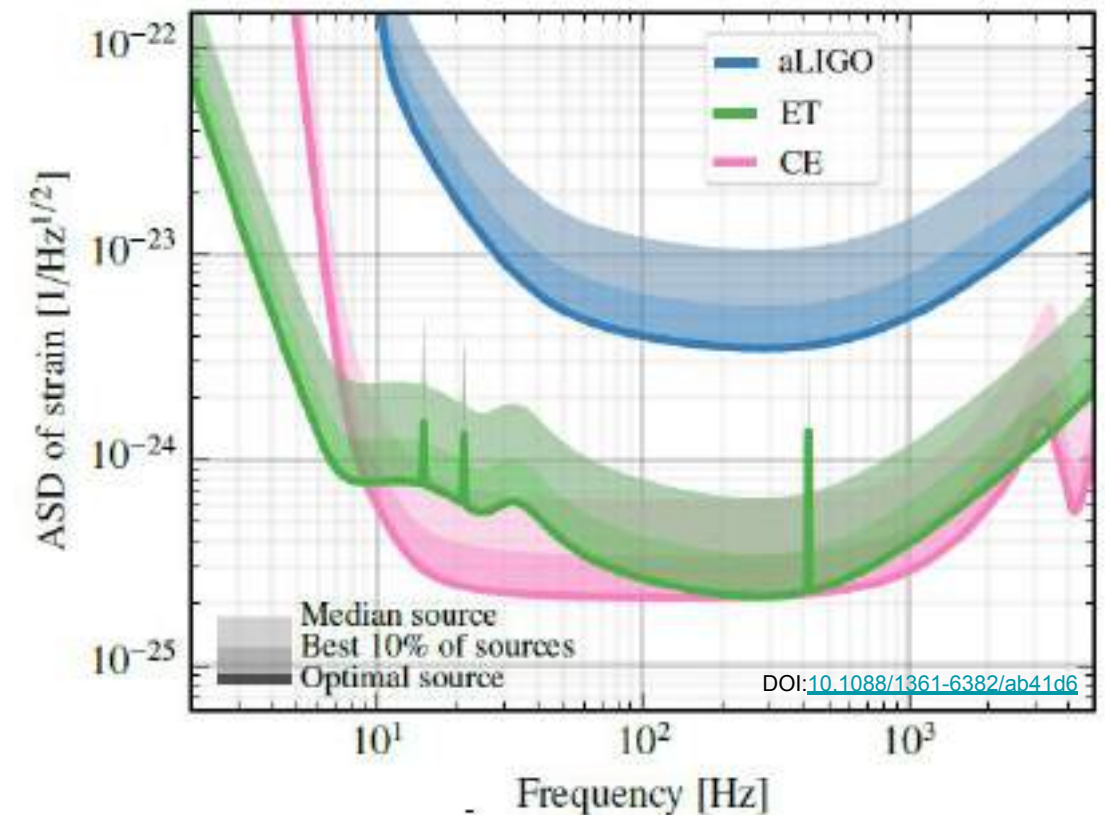
New PTA data

Einstein Telescope

But current LVK detectors have limitations: need to jump to next generation detectors:

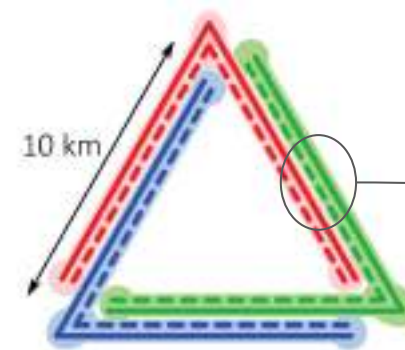
- **Einstein Telescope (ET):**
 - EU proposal for 3G observatory
 - Design study (baseline) triangle with arms of 10km
- **Cosmic Explorer (CE):**
 - US proposal for 3G observatory
 - L-shaped 40km interferometer

ET and CE will provide an improvement in sensitivity by one order of magnitude and a significant enlargement of the bandwidth



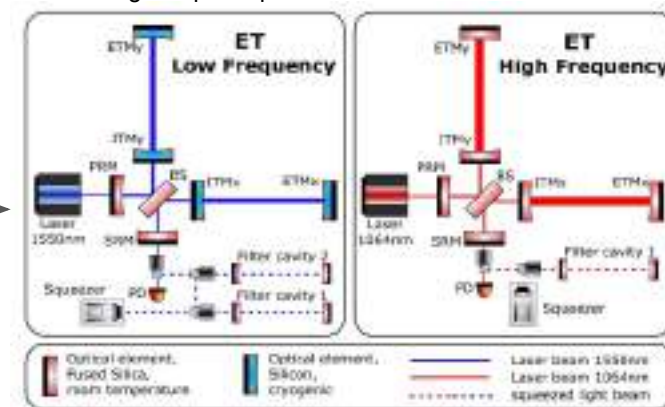
The current design of ET:

- **single site** located **200-300 meters underground** in order to significantly reduce seismic noise;
- **triangular shape**, consisting of three nested detectors
 - providing redundancy
- resolving the GW polarizations and a null stream
- **‘xylophone’** configuration: each detector consists of two interferometers
 - one tuned toward **high frequencies (HF)**, and using high laser power
 - one tuned toward **low-frequency (LF)**, working at cryogenic temperatures and low laser power



ET Design Report Update 2020

ET Design Report Update 2020



Einstein Telescope - possible designs

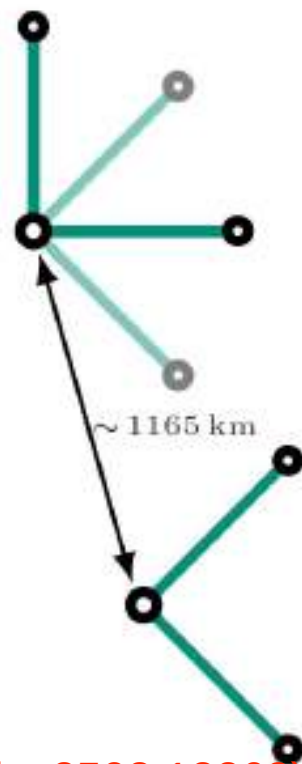
In the last years, proposals for **different designs** were made as they may bring **scientific advantages** with respect to the baseline design.

Different geometries are studied:

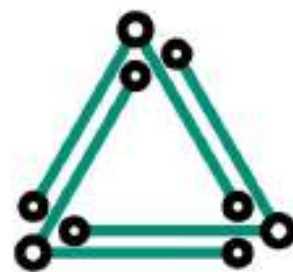
- **1 Triangular detector vs 2 L-shaped detector**
- **different arm length** (10, 15, 20 km)
- Cryo LF may be challenging to achieve: **HF+LF** vs **HF only**

For **2L-shape** two different orientations are proposed:

- parallel
- 45° angle



[arXiv:2303.15923](https://arxiv.org/abs/2303.15923)



PAPER • OPEN ACCESS

Science with the Einstein Telescope: a comparison of different designs

Marica Branchesi^{1,2}, Michele Maggiore^{3,4}, David Alonso⁵, Charles Badger⁶, Biswajit Banerjee^{1,2}, Freija Beirnaert⁷, Enis Belgacem^{3,4}, Swetha Bhagwat^{6,9}, Guillaume Boileau^{10,11}, Ssohrab Borhanian^{1,2}

[+ Show full author list](#)

Published 28 July 2023 • © 2023 The Author(s)

[Journal of Cosmology and Astroparticle Physics, Volume 2023, July 2023](#)

Citation Marica Branchesi et al JCAP07(2023)068

DOI 10.1088/1475-7516/2023/07/068

- 197 pages
- 75 authors

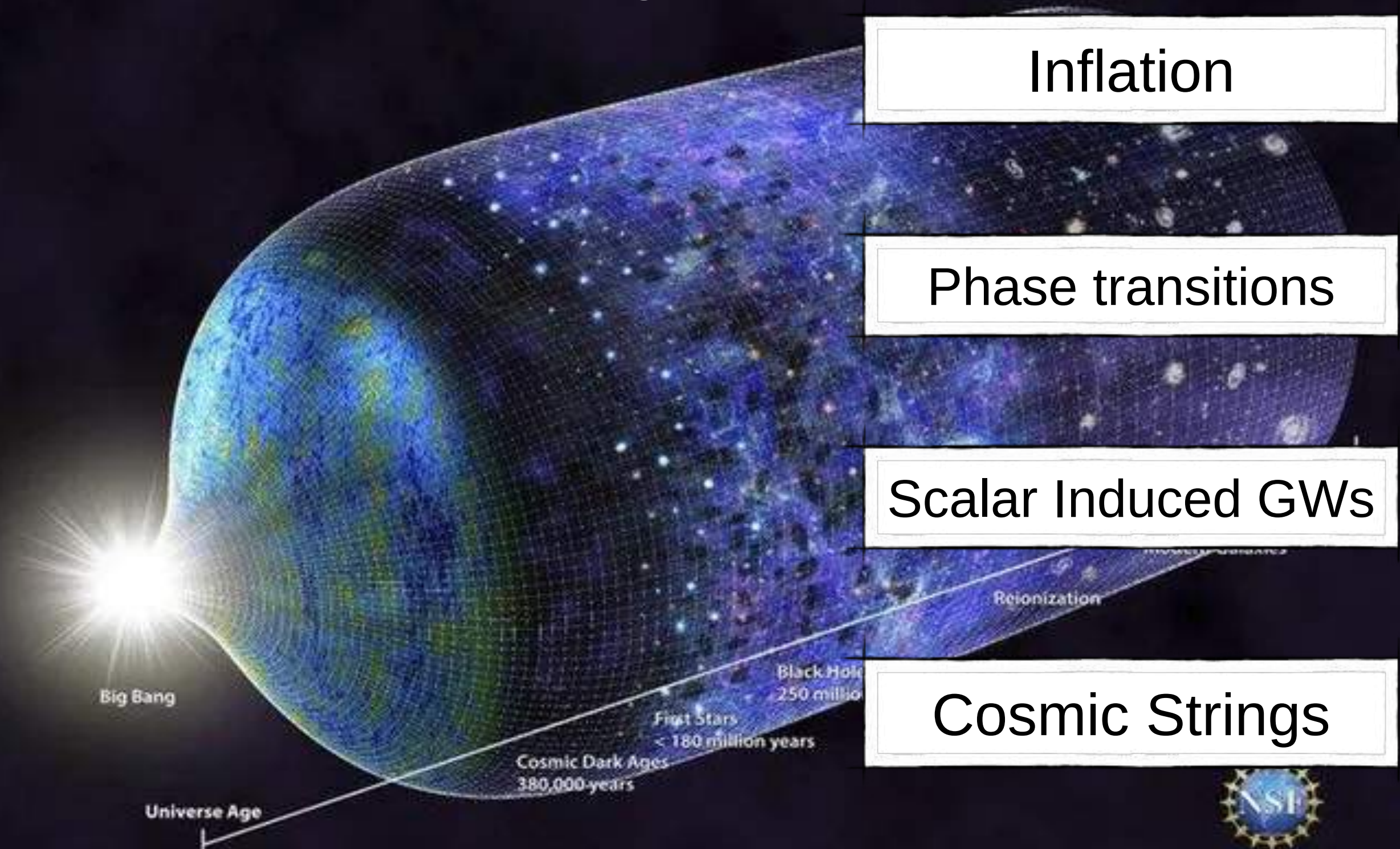


• More on the ET Blue Book (ArXiv: 2503.12263)

Early Universe Cosmology with Gravitational Waves

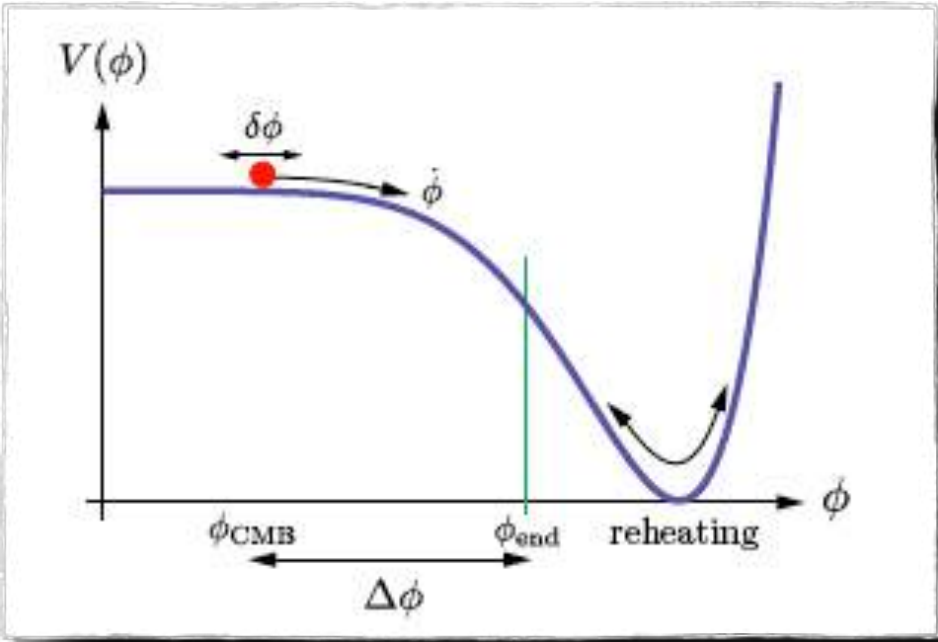
Probing the Early Universe

Cosmological Sources



Stochastic (i.e., *persistent, incoherent*) GWB of cosmological origin: **probe of the early Universe at energy scales above the ones achievable at current particle colliders**

Inflation and Primordial GWs

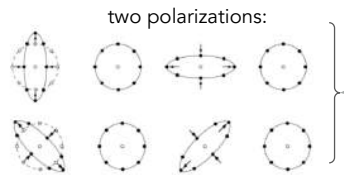


- Period of accelerated (exponential) expansion driven by a scalar field (inflaton) that rolls down on its flat potential

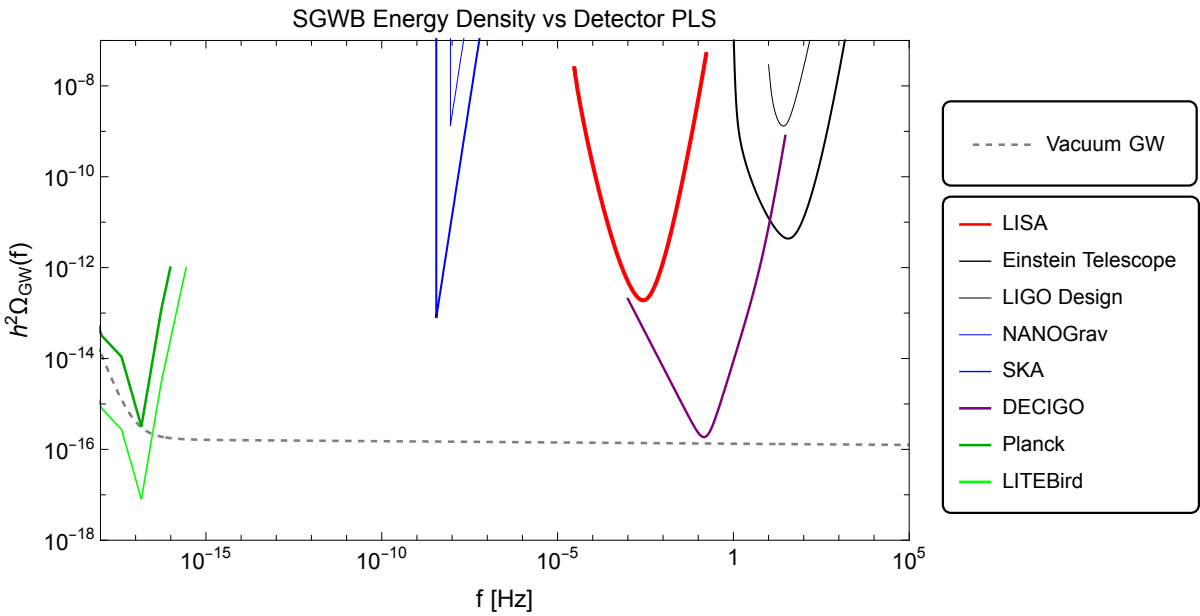
- Solve Standard Big-Bang shortcomings
- Generation of **TENSOR** and scalar perturbations
- Stretches the microphysics scales to super-horizon sizes

GW are represented by tensor perturbation h_{ij} of the FLRW metric

$$ds^2 = -dt^2 + a^2(t)(\delta_{ij} + h_{ij})dx^i dx^j$$

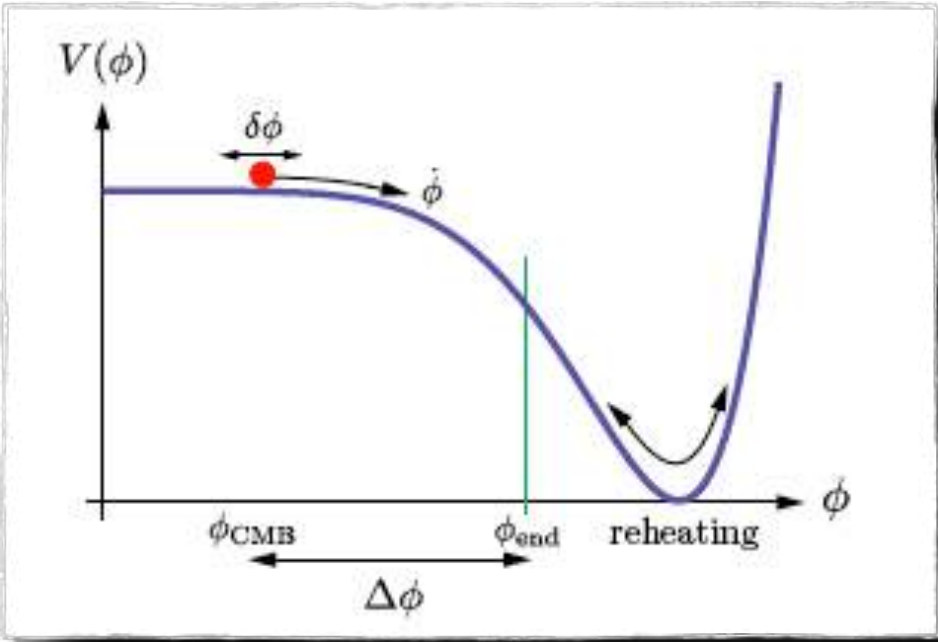


	$k \text{ [Mpc}^{-1}\text{]}$	$N_{\text{estim.}}$
CMB / LSS	$10^{-4} - 10^{-1}$	56 – 63
y - & μ -distortions	$10^{-1} - 10^4$	45 – 56
$P_\zeta \rightarrow \text{PBH} \rightarrow \text{GW @ PTA}$	$10^4 - 10^5$	41 – 44
$P_\zeta \rightarrow \text{PBH} \rightarrow \text{GW @ LISA}$	$10^5 - 10^7$	38 – 41
$P_\zeta \rightarrow \text{PBH} \rightarrow \text{GW @ ET}$	$10^7 - 10^8$	35 – 37
$P_{\delta g} \rightarrow \text{GW @ PTA}$	$10^6 - 10^8$	36 – 40
$P_{\delta g} \rightarrow \text{GW @ ET}$	$10^{11} - 10^{14}$	22 – 28
$P_{\delta g} \rightarrow \text{GW @ AdvLIGO}$	$10^{16} - 10^{17}$	15 – 17



LISA/ET → Possibility to test regions for which we have poor information

Inflation and Primordial GWs

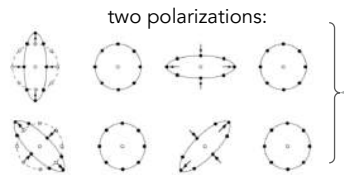


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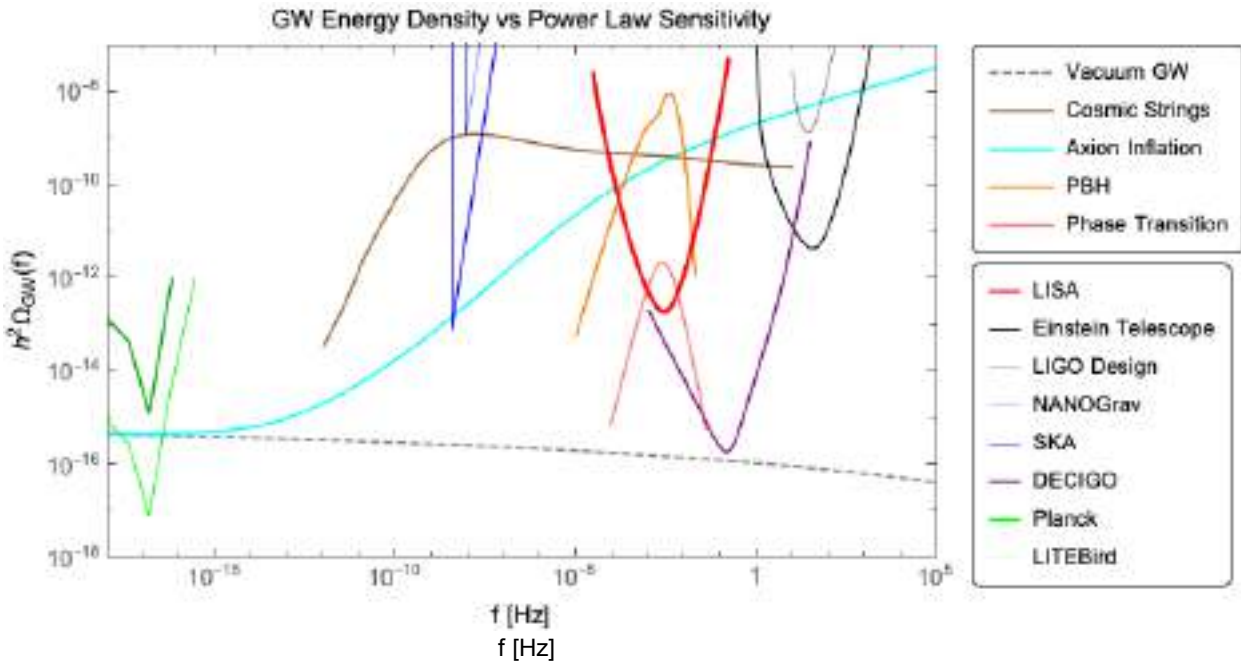
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$P_{\delta g} \rightarrow \text{GW @ AdvLIGO}$	$10^{16} - 10^{17}$	15 – 17



LISA/ET → Possibility to test regions for which we have poor information

Inflationary models producing such a signal

Axion inflation:

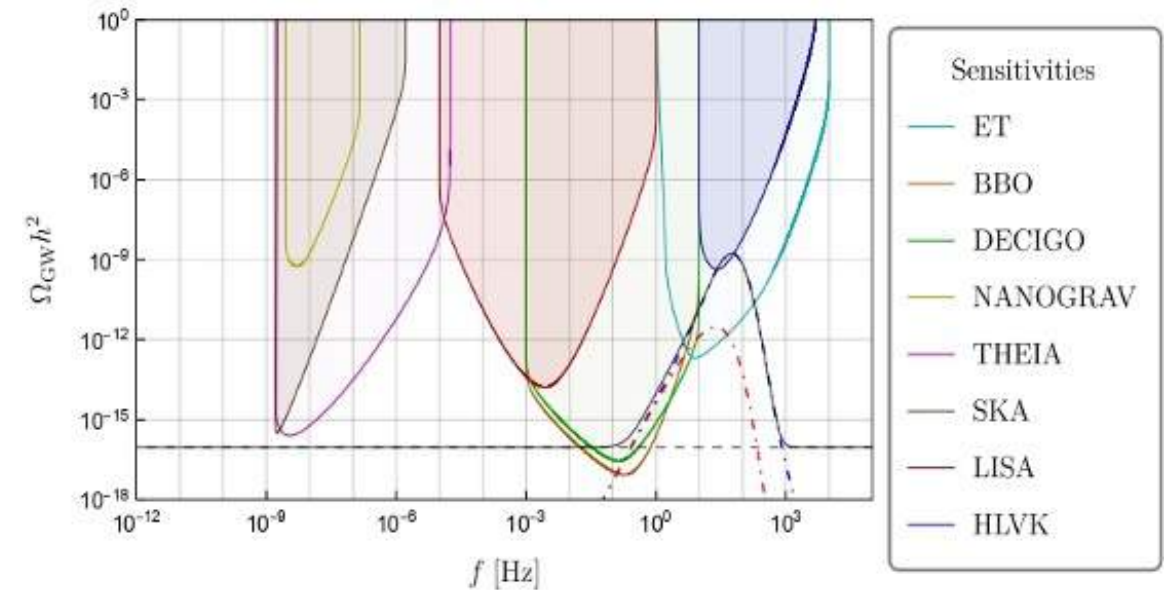
$$\mathcal{L} = -\frac{1}{2}\partial_\mu\phi\partial^\mu\phi - \frac{1}{4}F_{\mu\nu}^I F_I^{\mu\nu} - V(\phi) - \frac{\lambda^I}{4f}\phi F_{\mu\nu}^I \tilde{F}_I^{\mu\nu}$$

[Barnaby & Peloso 1011.1500]

[Sorbo 1101.1525]

The rolling axion strongly amplifies the gauge field, which in turn produces a strong SGWB.

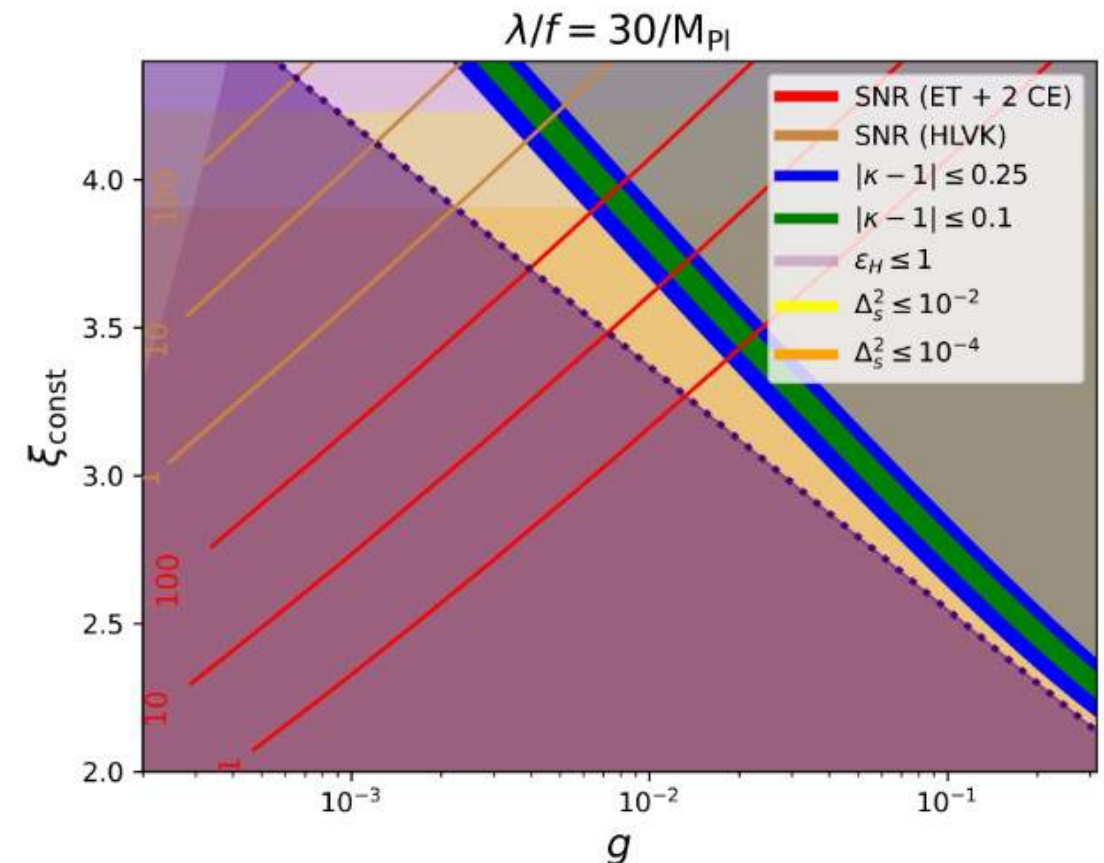
$$\xi \equiv \frac{\dot{\phi}}{2fH}$$



$$\Omega_{\text{GW}} \simeq \begin{cases} \frac{1}{12}\Omega_{R,0}\left(\frac{H^2}{\pi^2 M_{\text{Pl}}^2}\right)\left(1 + 4.3 \times 10^{-7} \frac{H^2}{M_{\text{Pl}}^2 \xi^6} e^{4\pi\xi}\right), & \text{for } U(1) \\ \frac{1}{12}\Omega_{R,0}\left(\frac{H^2}{\pi^2 M_{\text{Pl}}^2}\right)_{\xi=\xi_{cr}} \left[1 + \frac{1}{2}\xi_{cr}^6 \left(\frac{2^{7/4}H}{g\sqrt{\xi}} e^{(2-\sqrt{2})\pi\xi}\right)^2_{\xi=\xi_{ref}}\right], & \text{for } SU(2) \end{cases}$$

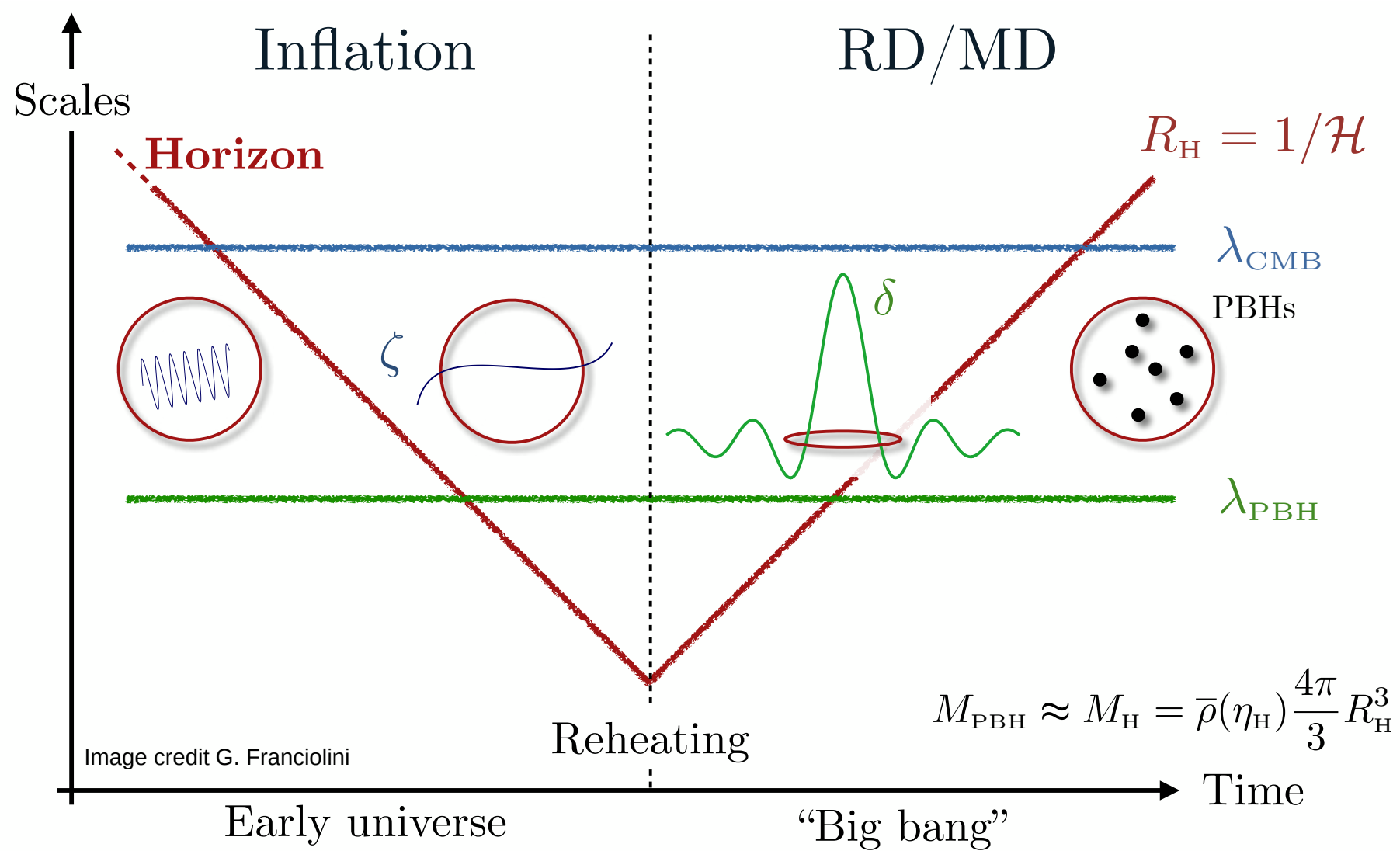
Presence of other fields (e.g. axions) in the Early Universe

Testing fundamental physics (e.g. axion decay constant) - related to high energy physics



GWs - Primordial Black Holes and Dark Matter

Large scalar perturbations source gravitational waves at 2nd order in perturbation theory when they re-enter the horizon during radiation era

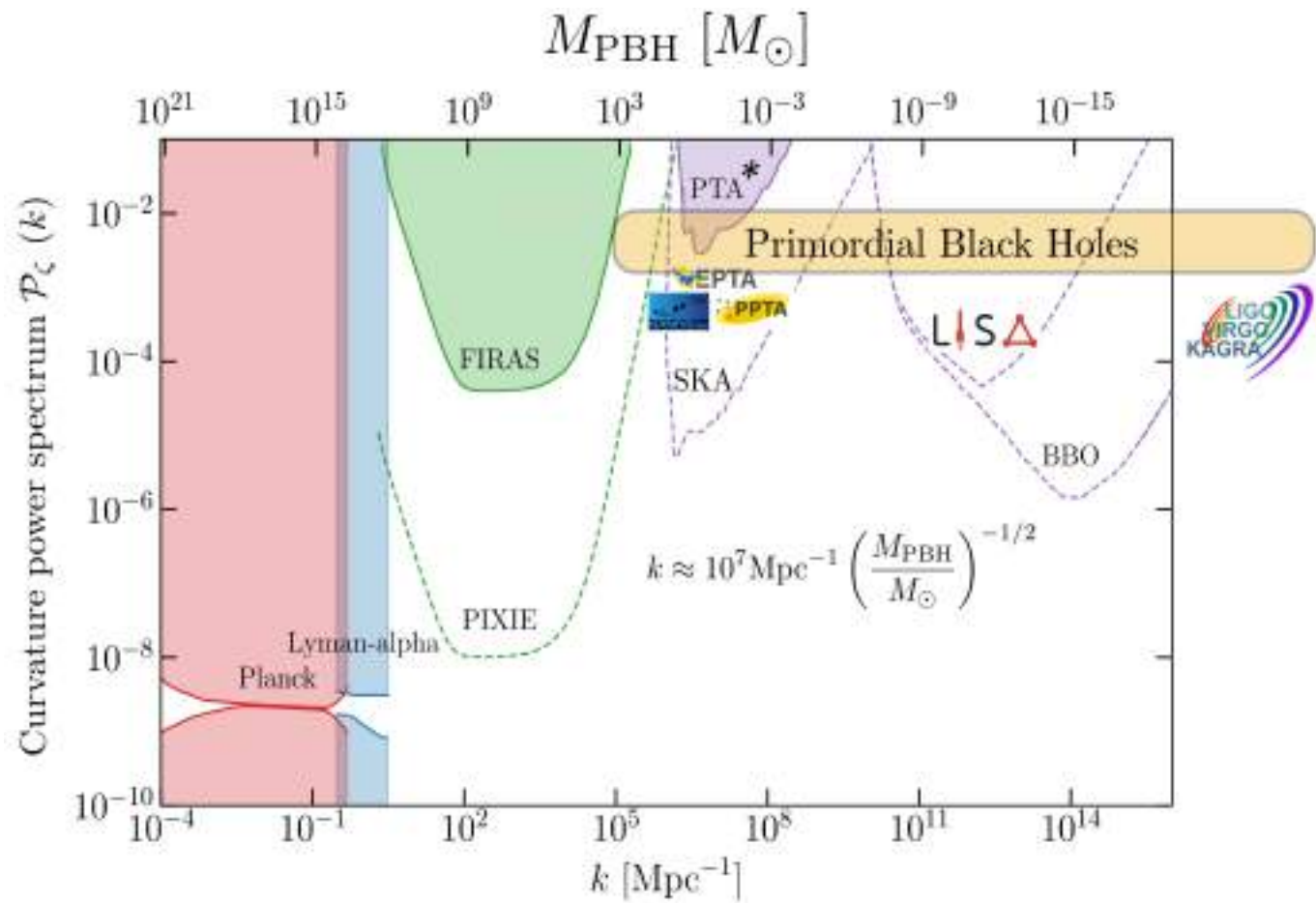


$$h''_{ij} + 2\mathcal{H}h'_{ij} - \nabla^2 h_{ij} = \mathcal{O}(\partial_i \zeta \partial_j \zeta)$$

[Tomita, K., 1967]
 [Matarrese, S., et al., 1993]
 [Domenech, G., review '21]

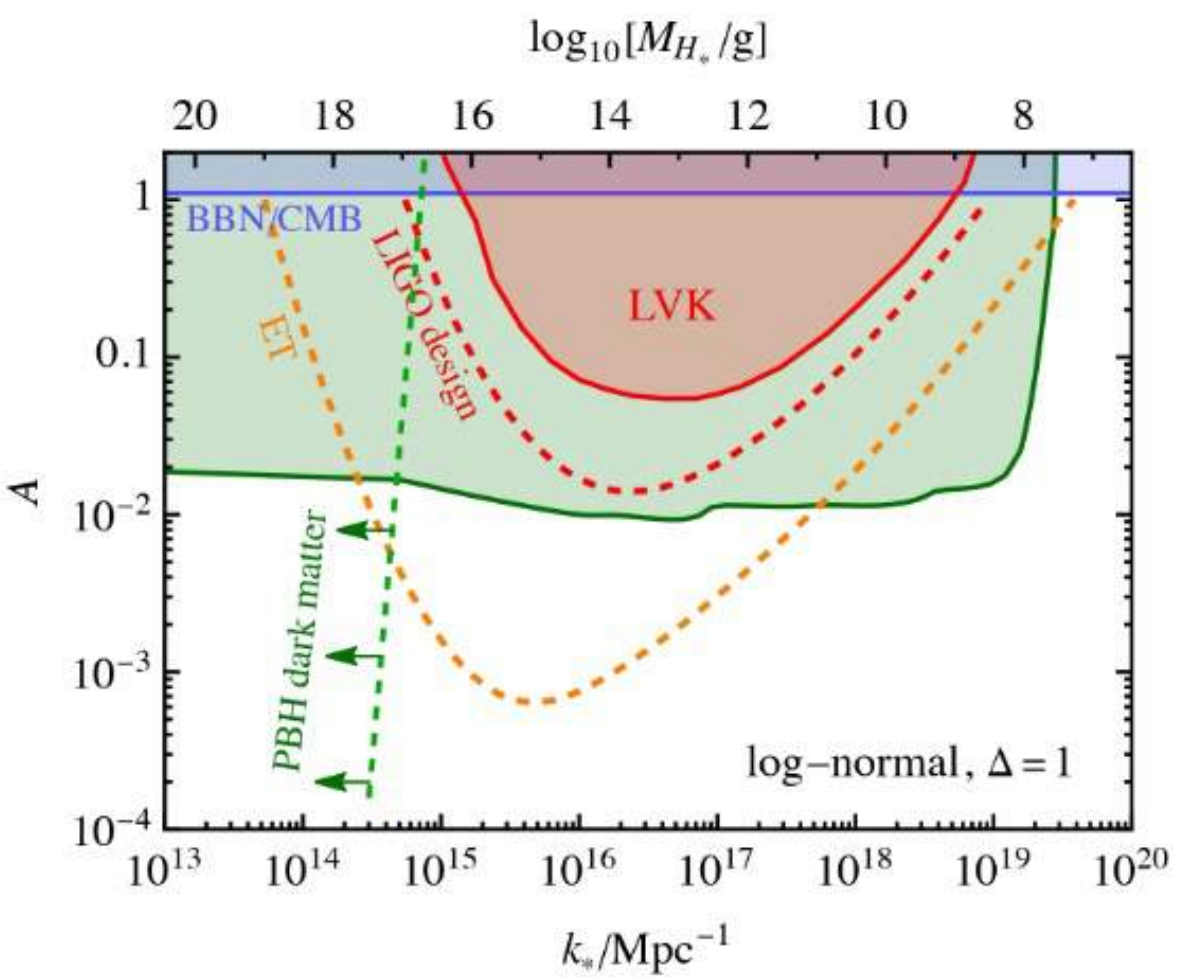
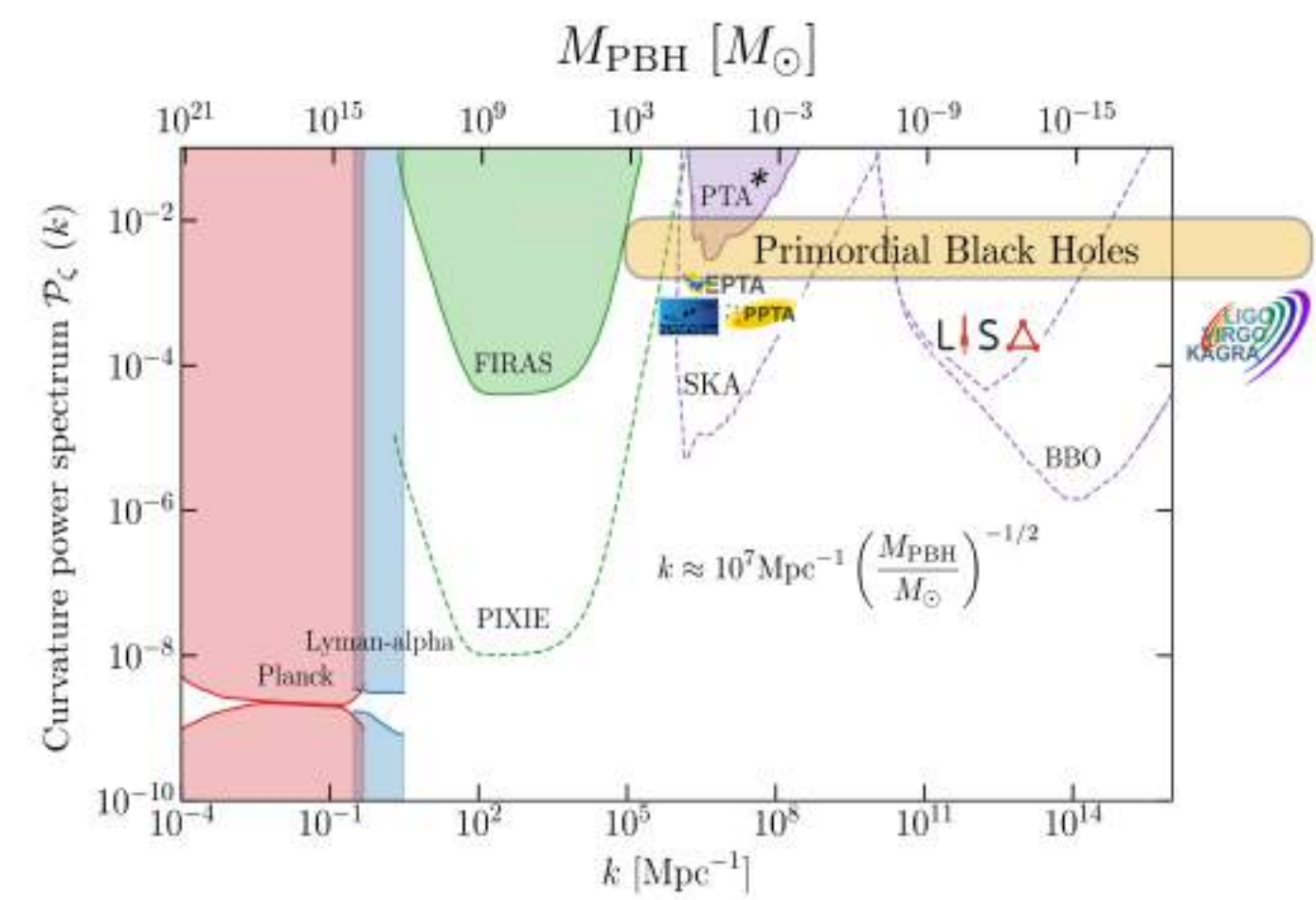
GWs - Primordial Black Holes and Dark Matter

$$\Omega_{\text{GW}} \simeq \Omega_{\text{rad}} A_{\zeta}^2$$



GWs - Primordial Black Holes and Dark Matter

$$\Omega_{\text{GW}} \simeq \Omega_{\text{rad}} A_{\zeta}^2$$

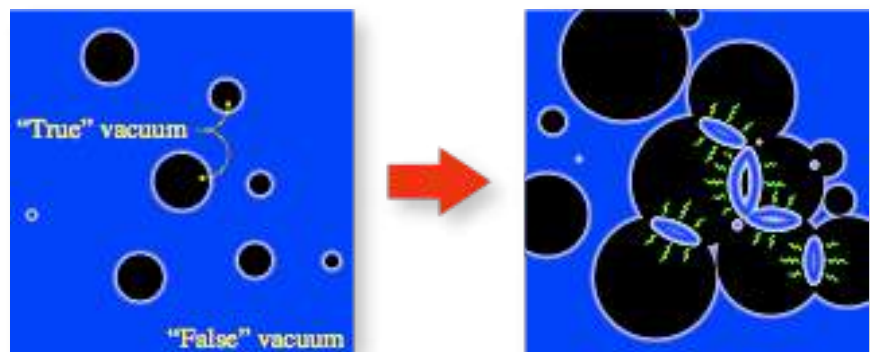


$$f_{\text{GW}} = 4.1 \text{ Hz} \left(\frac{\kappa}{2.51}\right) \left(\frac{g_{*}}{106.75}\right)^{-1/12} \left(\frac{M_H}{10^{-17} M_{\odot}}\right)^{-1/2}$$

Phase Transition in the Early Universe

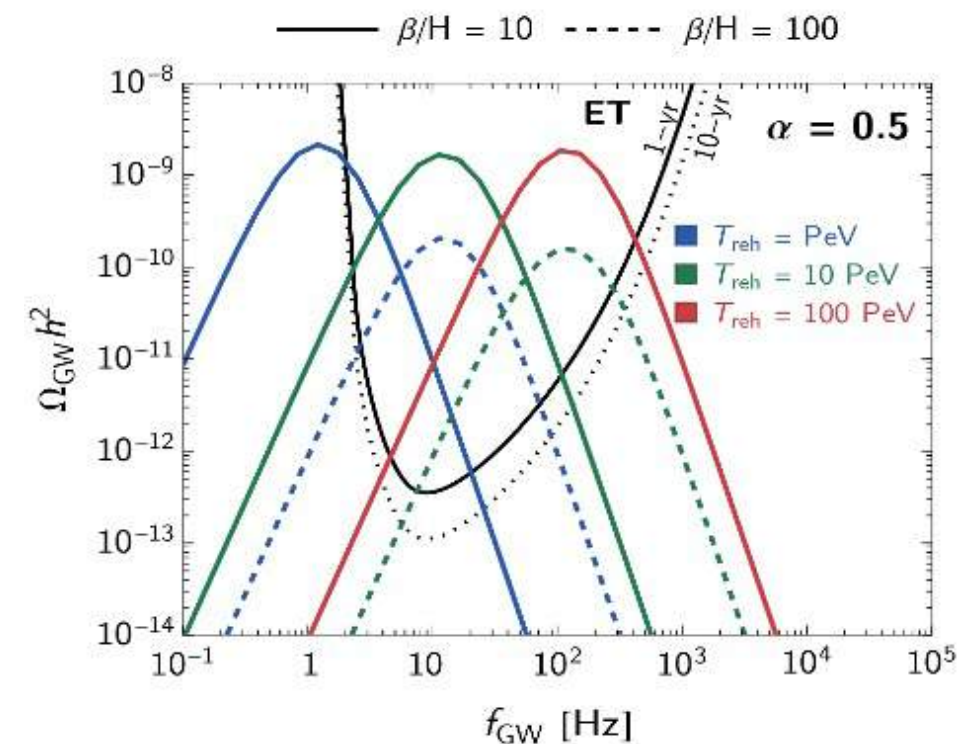
As the temperature in the very early universe decreases, there can be several **PTs: QCD, EW...**Beyond Standard Model?

If the **PT is first order**, the SGWB signal could be detectable by LISA



Processes

- Bubble collisions
- MHD Turbulence
- Sound Waves



$$\Omega_{\text{GW}}(f) \propto \frac{(a+b)^c f_{\text{peak}}^b f^a}{(b f_{\text{peak}}^{\frac{a+b}{c}} + a f^{\frac{a+b}{c}})^c}$$

$$f_{\text{peak}} \approx \mathcal{O}(1) \times 10^{-5} \frac{1}{H_{\text{reh}} R_{\text{coll}}} \frac{T_{\text{reh}}}{100 \text{ GeV}} \text{ Hz}$$

[Caprini C., et al '16, '19- LISA CosWG paper]

What we can learn from Phase Transition?

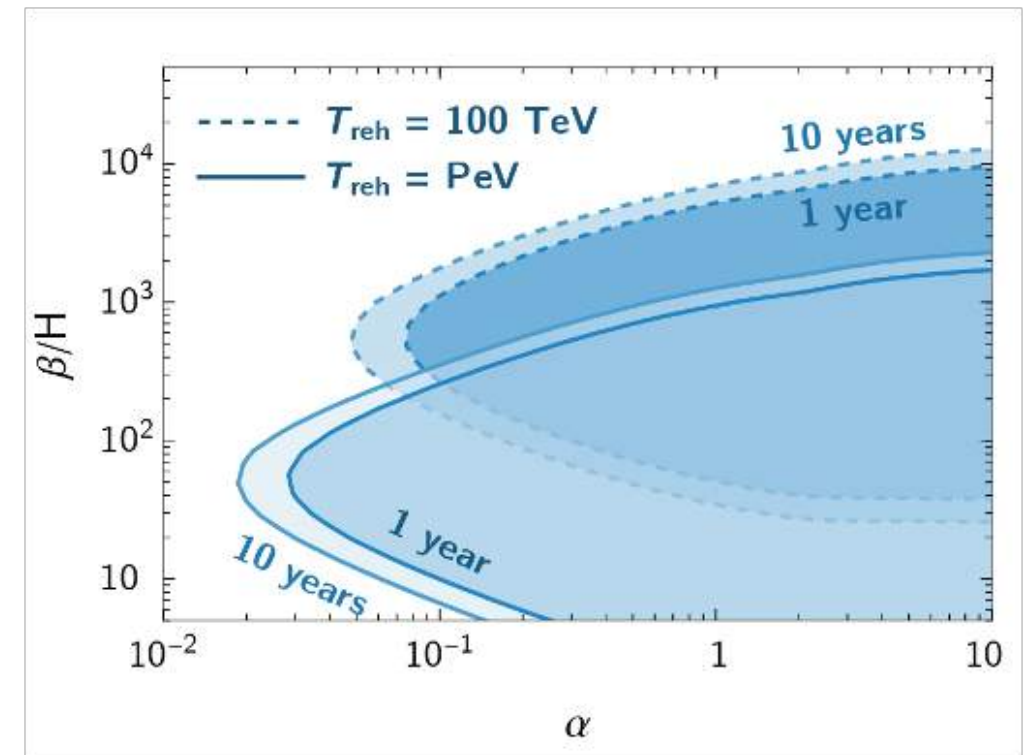
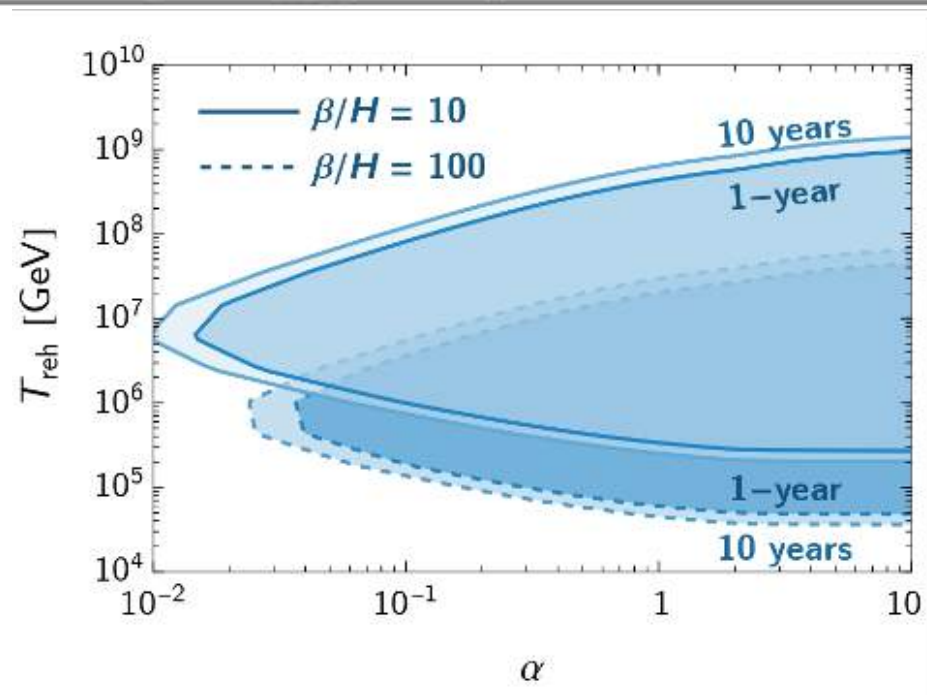
- ET could act as a **probe of Beyond Standard Model physics, complementary to colliders**

Simplest extensions of the SM

$$V_{\text{tree}}(\Phi, s) = -\mu_h^2 \Phi^\dagger \Phi + \lambda (\Phi^\dagger \Phi)^2 + \mu_s^2 \frac{s^2}{2} + \frac{\lambda_s}{4} s^4 + \frac{\lambda_{hs}}{2} s^2 \Phi^\dagger \Phi,$$

$$\text{BSM parameters}(\text{couplings, masses...}) \Rightarrow (\alpha, \beta, T_{\text{reh}}, v_w)$$

$$(\alpha, \beta, T_{\text{reh}}, v_w) \Rightarrow (f_{\text{peak}}, \Omega_{\text{peak}}, \text{GW spectrum}).$$



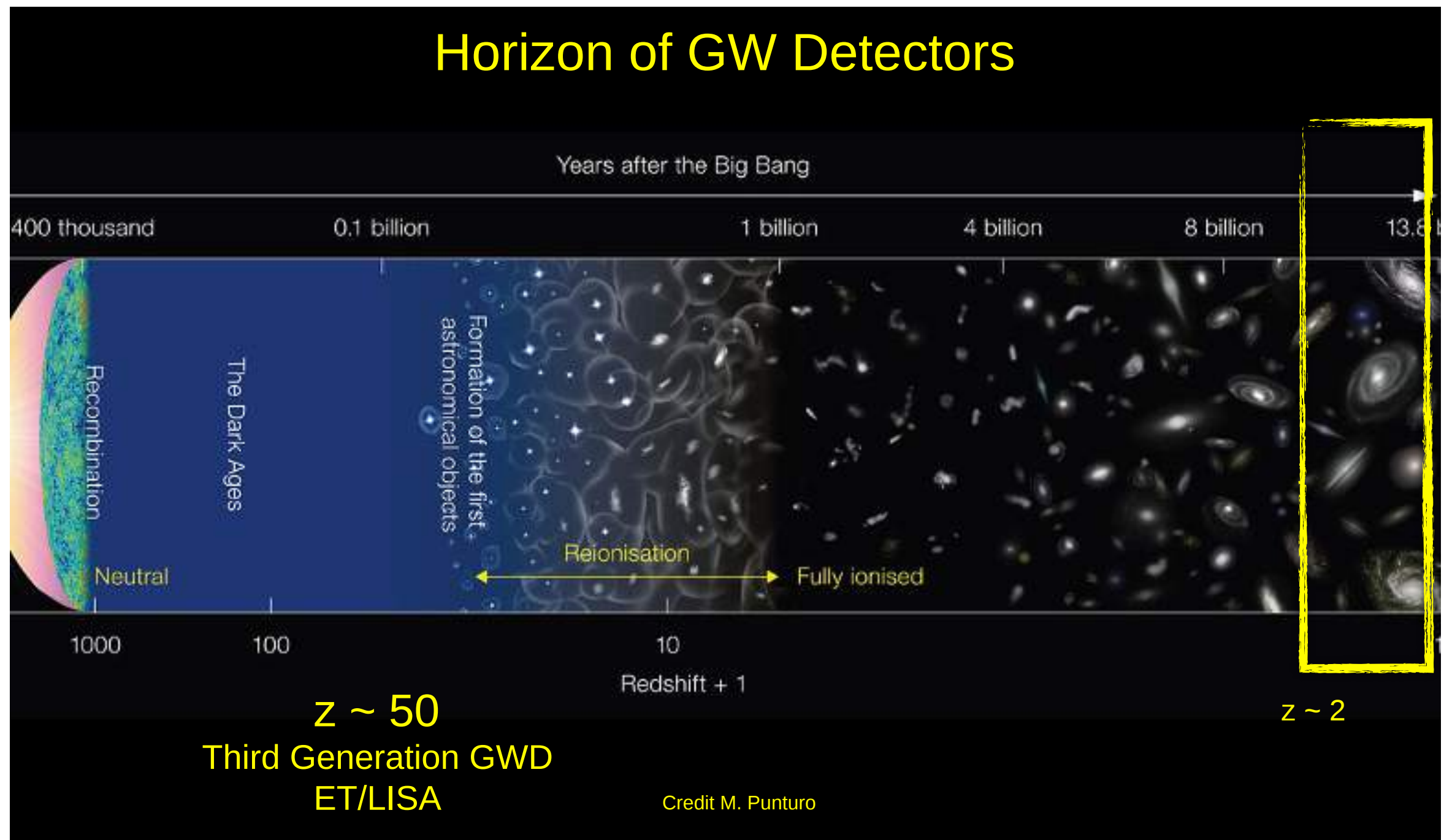
- ET could act as a **probe of Beyond Standard Model physics, complementary to colliders**
- In some BSM scenarios possible **joint detection at ET and LHC/FCC**

λ_{hs} Singlet coupling
 m_s Singlet mass

ET Blue book 2503.12263 (PT)

Late Universe Cosmology with Gravitational Waves

Where is the horizon for 3G detectors?

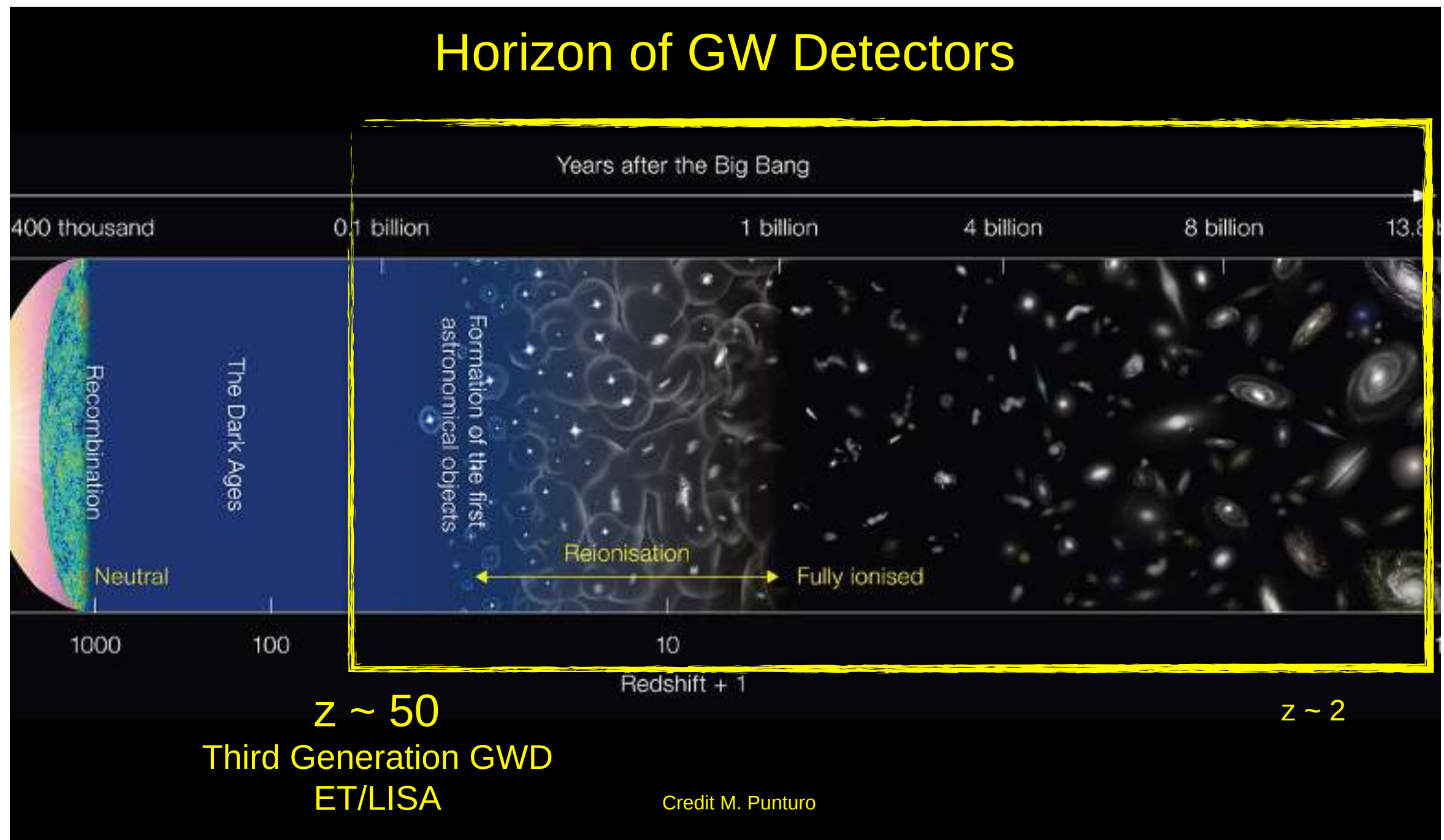


Einstein Telescope:

BBHs up to cosmic Dark Ages ($z > 30$)

BNSs up to cosmic Noon ($z \sim 2$)

Where is the horizon for 3G detectors?



Einstein Telescope:

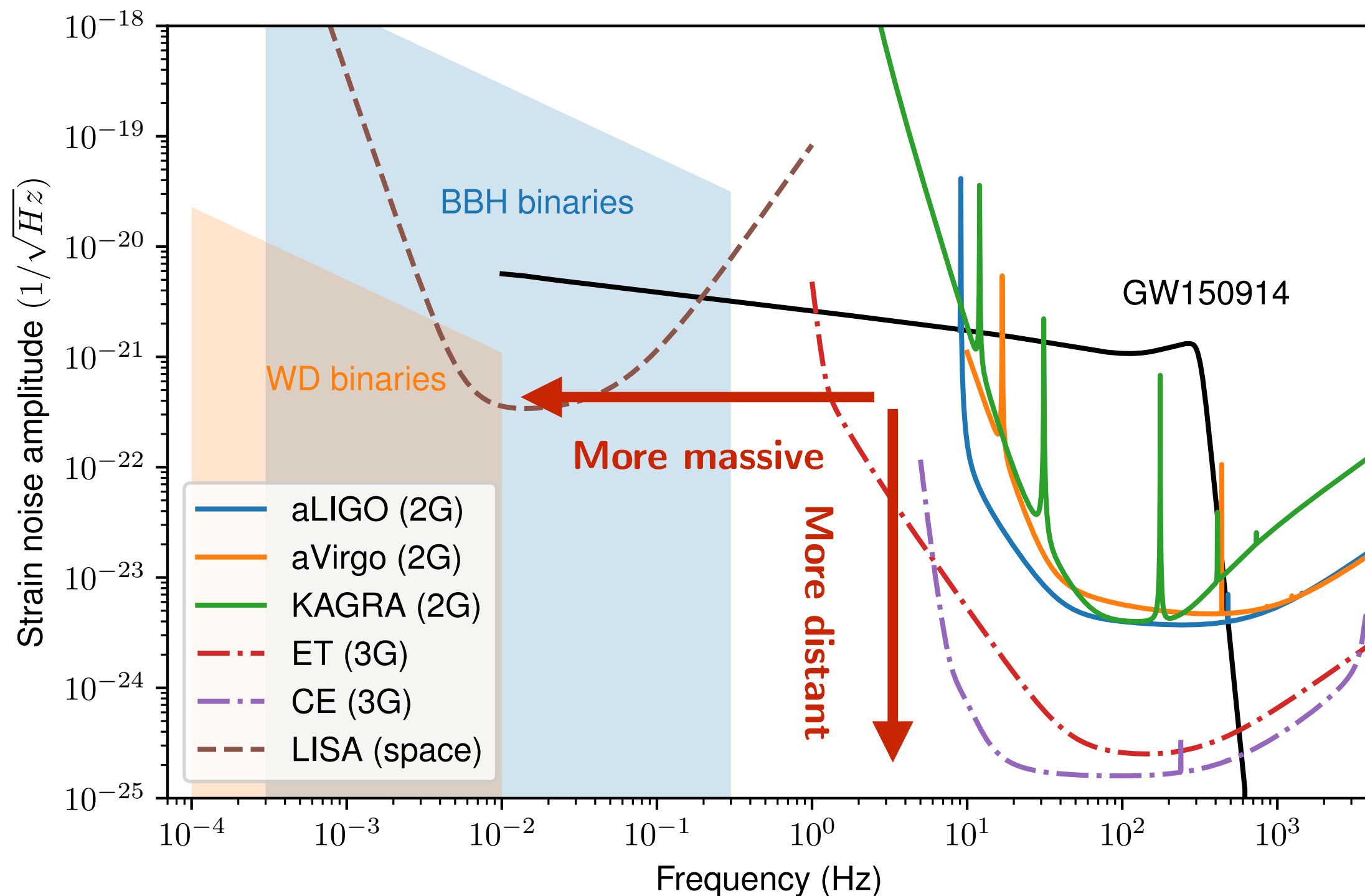
BBHs up to cosmic Dark Ages ($z > 30$)

BNSs up to cosmic Noon ($z \sim 2$)

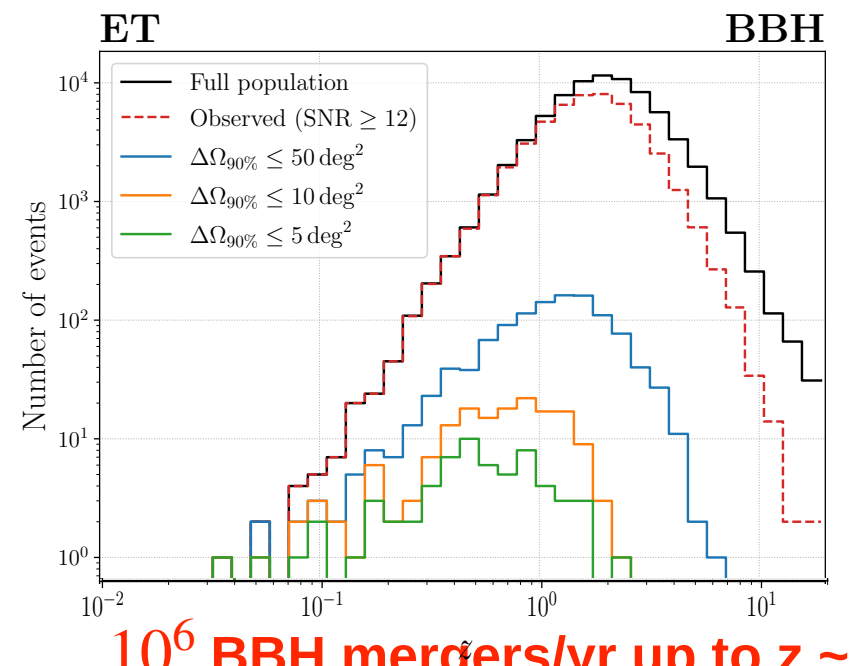
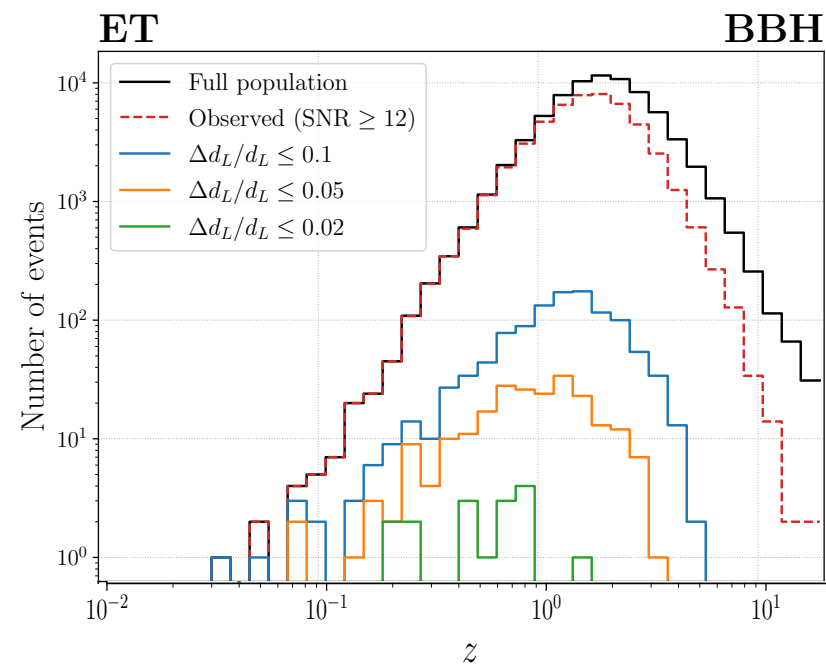
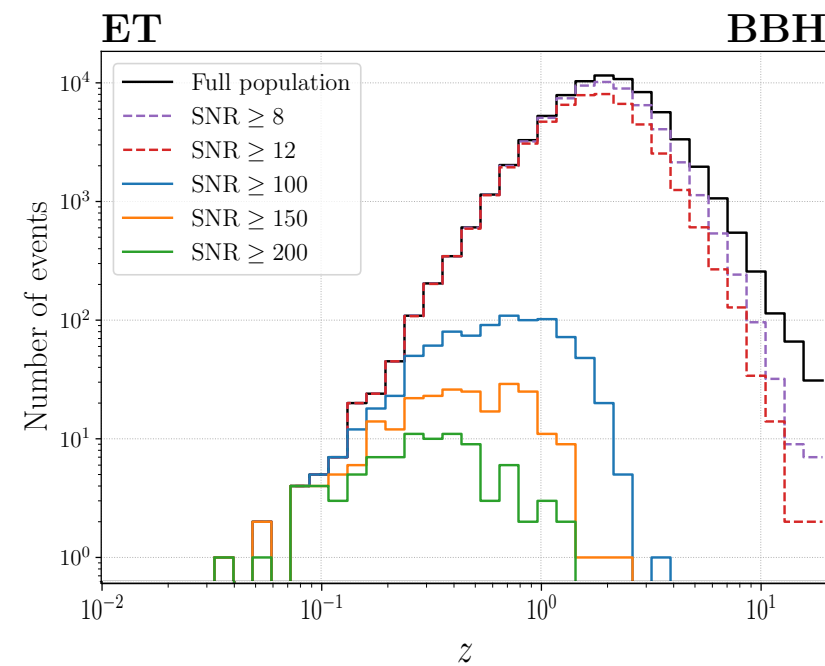
From 2G to 3G detectors: ET and LISA

$$h_{\text{gw}} \sim 1/d_L$$

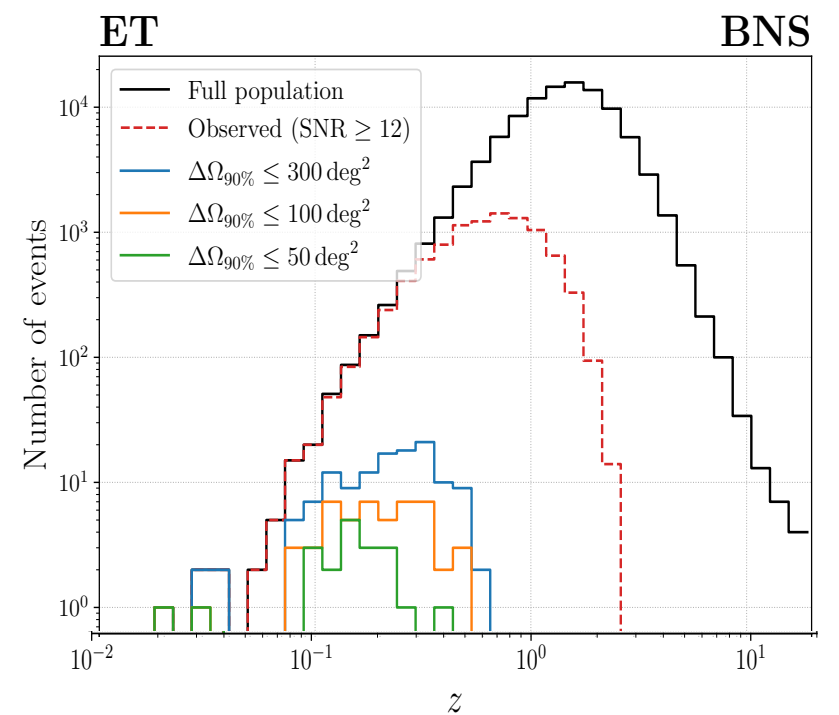
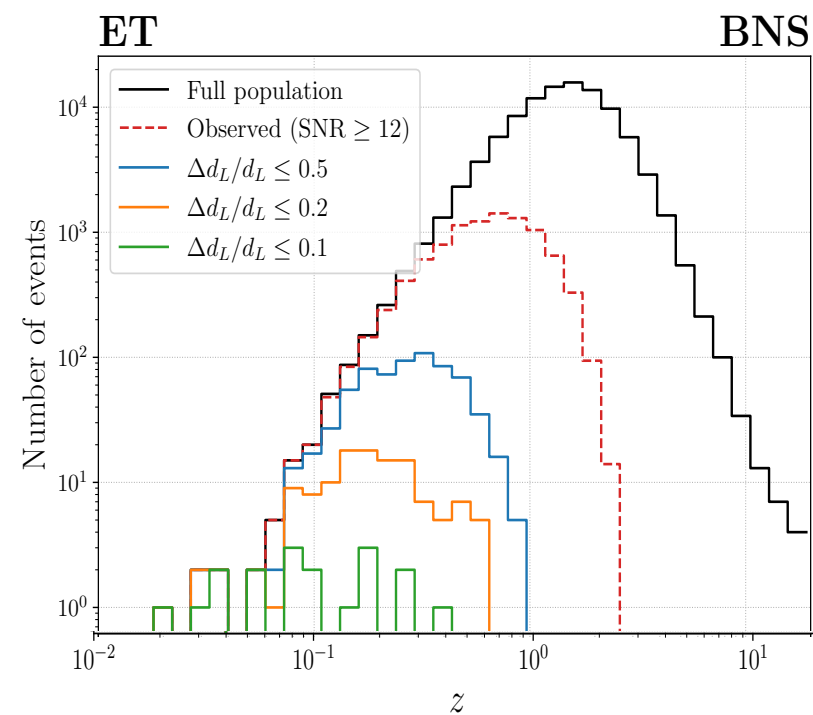
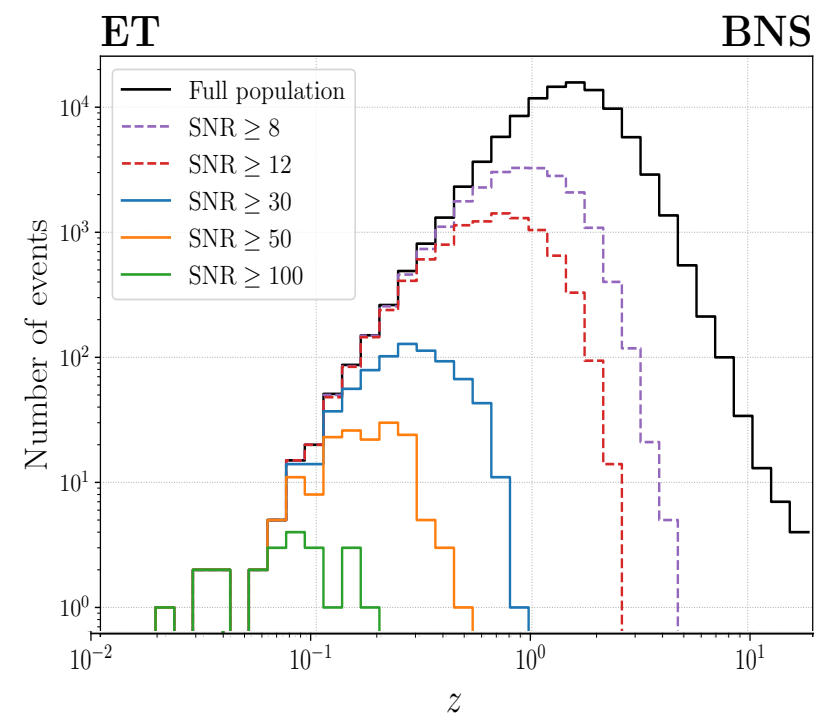
$$f_{\text{gw}} \sim 1\text{kHz} \left(\frac{10M_{\odot}}{M} \right)$$



Probing the Late Universe with ET



10^6 BBH mergers/yr up to $z \sim 50$



10^5 BNS mergers/yr up to $z \sim 2$

Many Golden Events

Using GWs as Standard Sirens

$$cz = H_0 \times d_L$$

Diagram illustrating the Hubble-Lemaître law equation $cz = H_0 \times d_L$. Arrows point from the labels to the corresponding variables in the equation:

- redshift points to z
- Hubble constant points to H_0
- Luminosity distance points to d_L

- three quantities: pick any two and infer the third.
- With standard sirens:

d_L from GW measurements;

z from, e.g. electromagnetic measurements (if have an optical counterpart, and know the host galaxy, can determine z).

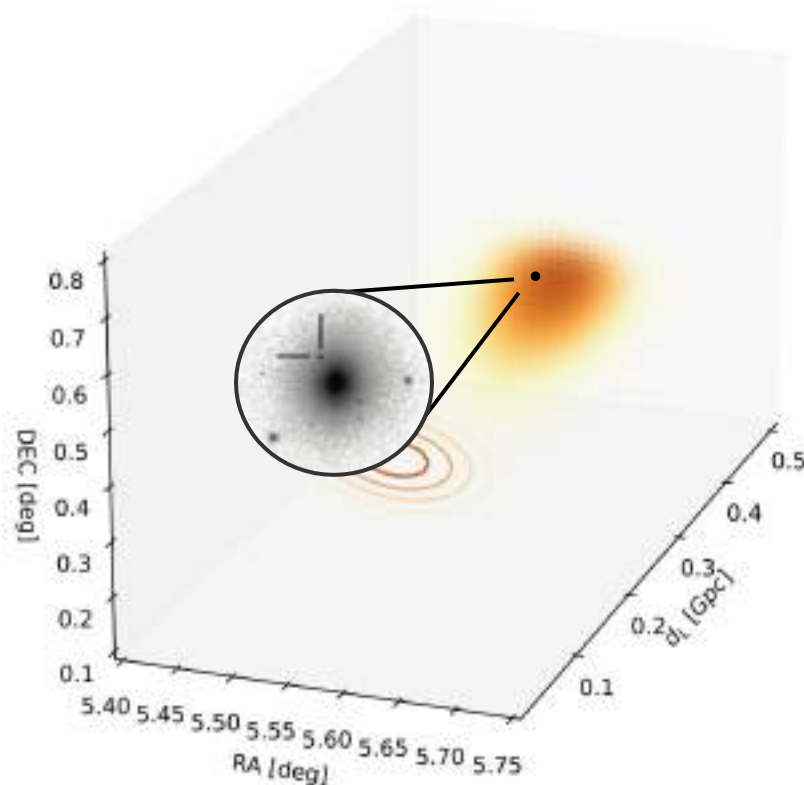
=> independent measure of H_0

Cosmology via the distance-redshift relation
... but no redshift measurement with GW data alone
(Degeneracy with masses)

Using GWs as Standard Sirens - Redshift Information

Bright Sirens

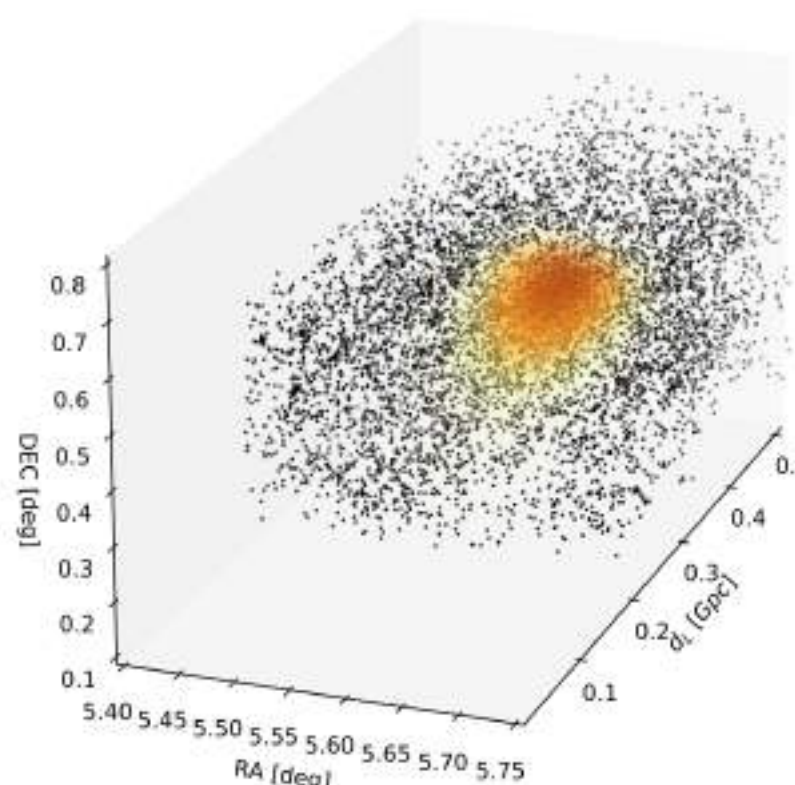
An **EM counterpart** is observed and used to obtain the host galaxy redshift.



(Holz & Hughes 2005)

Dark Sirens

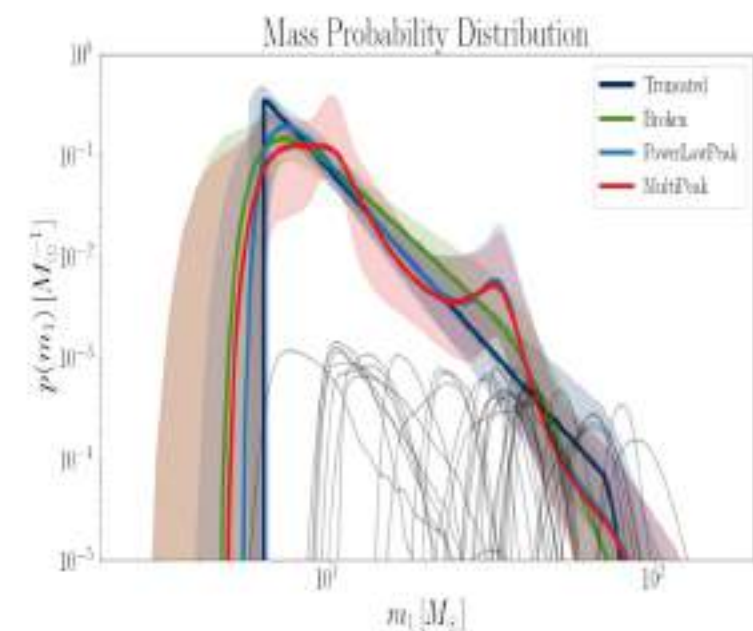
No EM counterpart observed. **Galaxy surveys** are used to provide redshift estimates for potential host galaxies.



(Schutz 1986, Del Pozzo 2012)

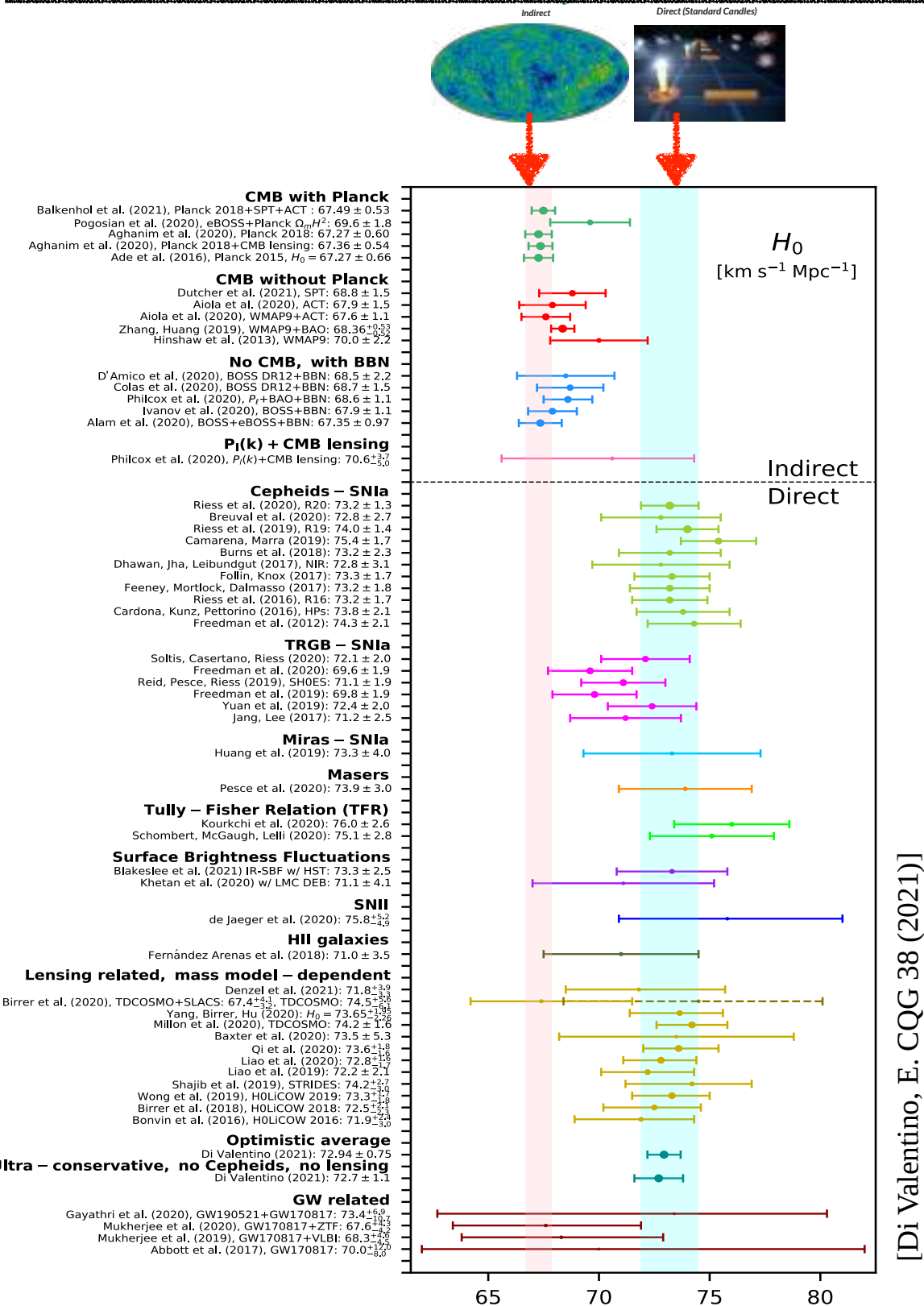
Spectral sirens

No EM counterpart or galaxy survey is used. Features in the **mass distribution** of the GW population break the mass-redshift degeneracy.

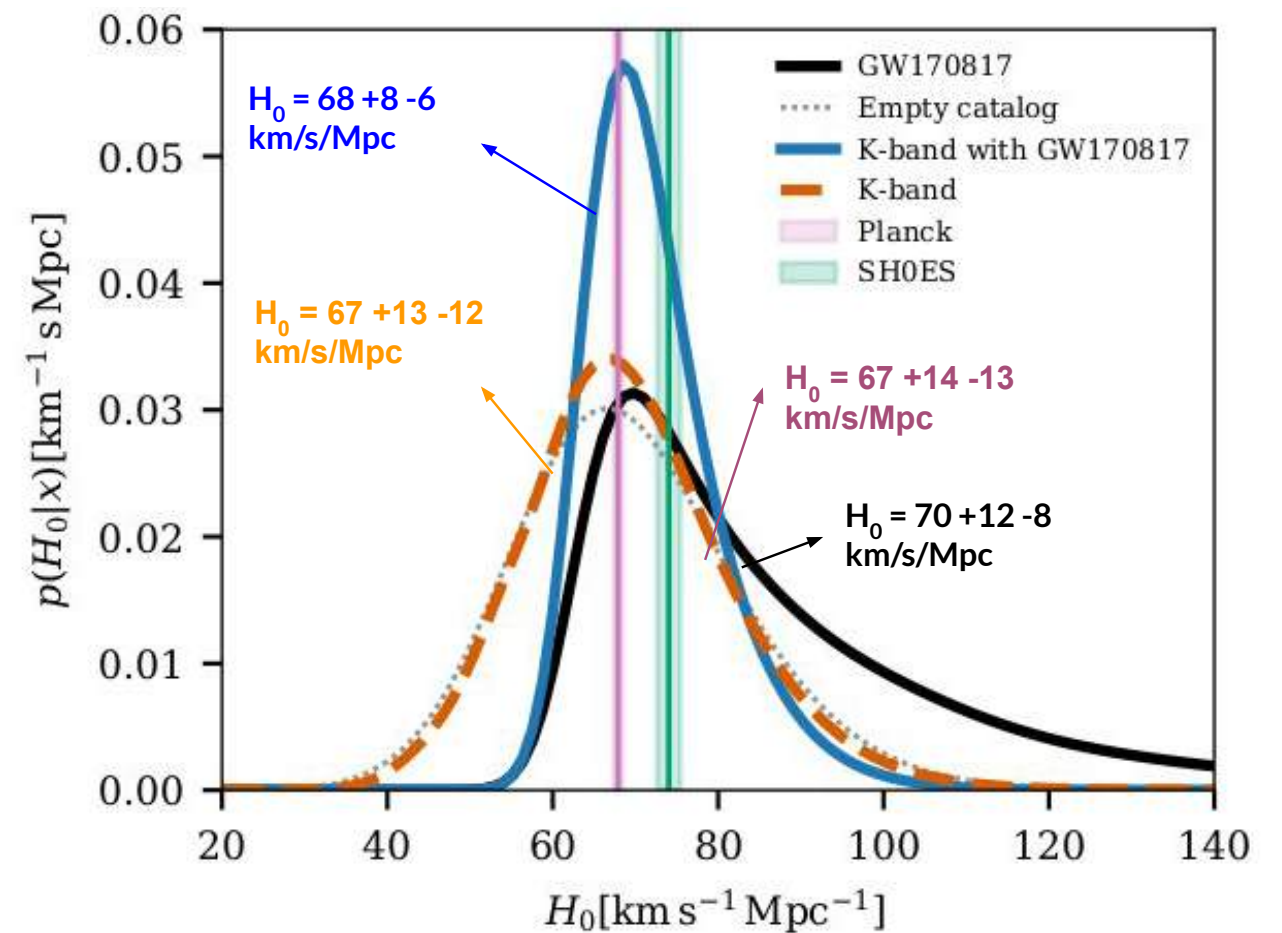


(Chernoff & Finn 1993)

H_0 - where we are: after GWTC-3 with Dark Sirens



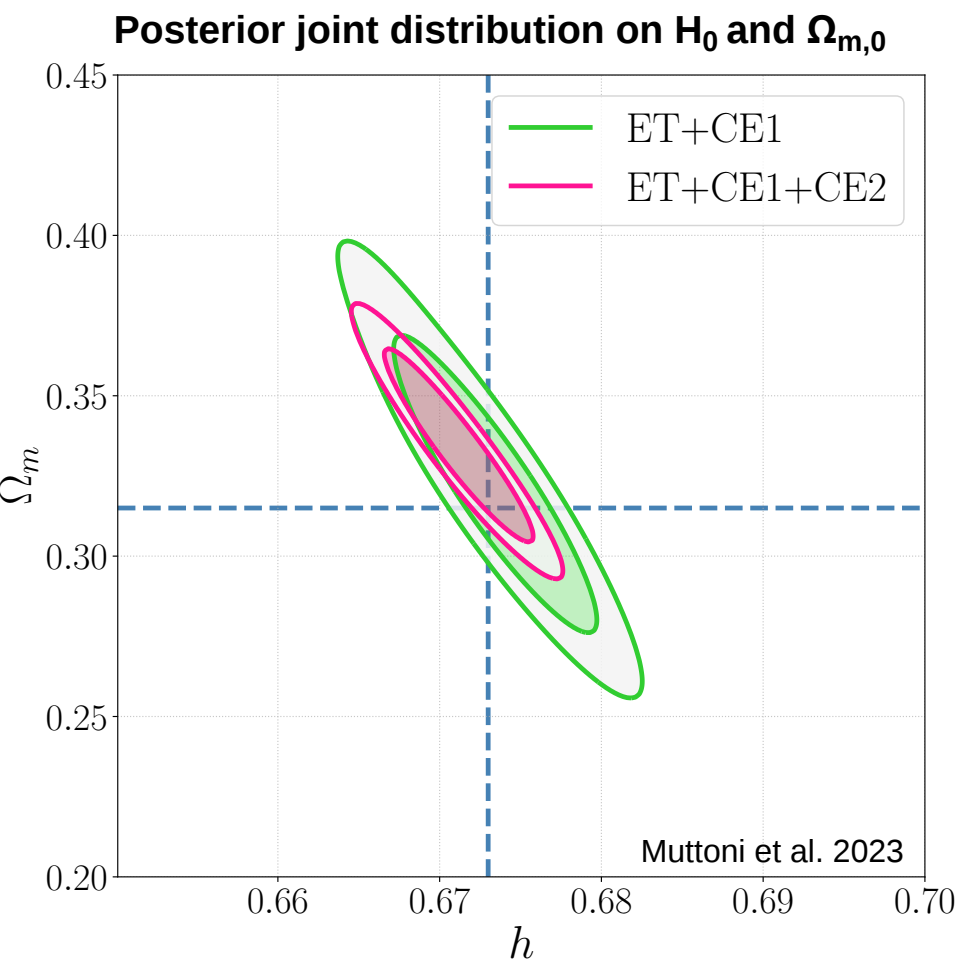
[GWCT3 - LVK+, ApJ 949 76 (2023)]



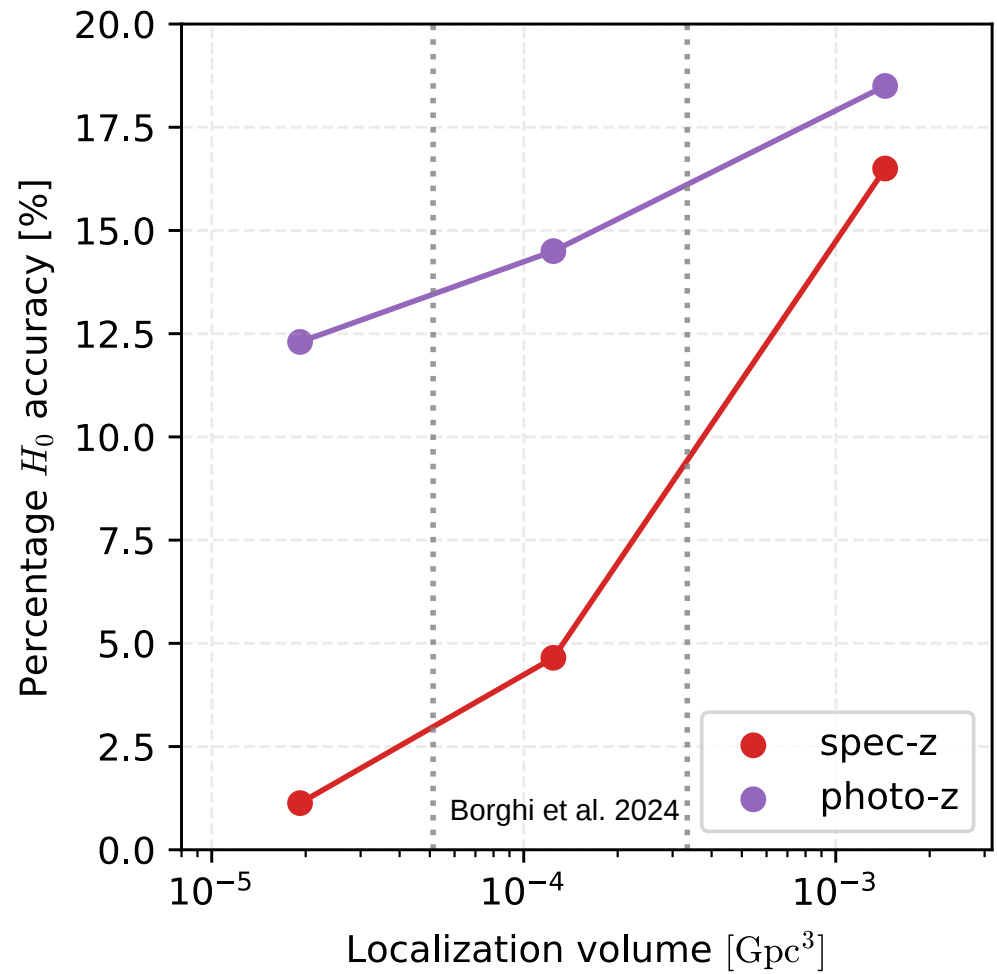
K-band for the luminosities of galaxies and the preferred mass model (Power Law+Gaussian peak)

~ 5 sigma tension between low and high redshift measurements of the Hubble parameter

H_0 - where we will be with ET? DARK SIRENS method



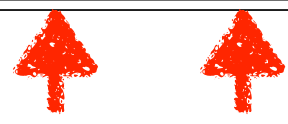
H_0 accuracy dependence VS redshift error of the galaxy catalogue



Fiducial scenario

- 1 year of observation
- Full duty cycle
- Loudest GW event localization
- Complete galaxy catalogue (generated using L-Galaxies running SAM on the Millennium simulation)
- Halo mass resolution ($2 \times 10^{10} M_\odot$)
- Box size (685 Mpc)

Network	N		$\Delta h/h$ (%)				$\Delta \Omega_m/\Omega_m$ (%)		
	$z < 1$	$z < 3$	$z < 1$	$z < 1$ fixed Ω_m	$z < 1$ single-host	$z < 3$	$z < 1$	$z < 1$ single-host	$z < 3$
ET+CE1	207	248	0.6 (1.1)	0.2 (0.4)	3.3-7.1 (5.6-11.2)	0.7 (1.1)	8.8 (14.4)	-	8.8 (14.6)
ET+CE1+CE2	278	348	0.5 (0.8)	0.2 (0.3)	1.7-2.1 (2.7-3.3)	0.4 (0.7)	6.1 (10.0)	-	5.3 (8.7)



Expected cosmological constraints at the 68% (90%) CI for multiyear 3G observations estimated from the 1 year fiducial results

Using GWs to Constrain Dark Energy

If the Dark Energy equation of states evolves in time

$$p_{\text{DE}}(z) = w_{\text{DE}}(z) \rho_{\text{DE}}(z)$$

$$\frac{\rho_{\text{DE}}(z)}{\rho_0} = \Omega_{\text{DE}} \exp \left\{ 3 \int_0^z \frac{d\tilde{z}}{1 + \tilde{z}} [1 + w_{\text{DE}}(\tilde{z})] \right\}$$

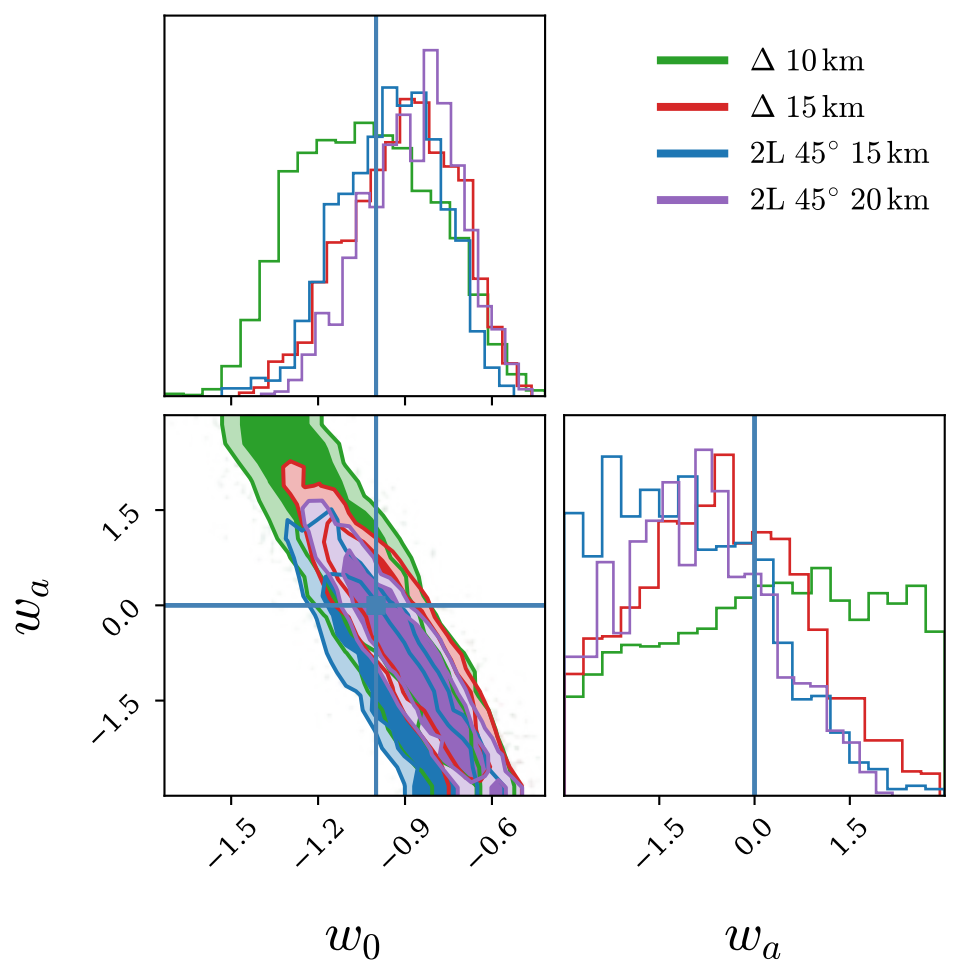
$$d_L(z) = \frac{c}{H_0} (1+z) \int_0^z \frac{d\tilde{z}}{\sqrt{\Omega_M(1+\tilde{z})^3 + \Omega_R(1+\tilde{z})^4 + \Omega_\Lambda}} \quad \rightarrow \quad d_L(z) = \frac{c}{H_0} (1+z) \int_0^z \frac{d\tilde{z}}{\sqrt{\Omega_M(1+\tilde{z})^3 + \Omega_R(1+\tilde{z})^4 + \rho_{\text{DE}}(\tilde{z})/\rho_0}}$$

$$w_{\text{DE}}(z) = w_0 + \frac{z}{1+z} w_a$$

Linder parameterization

Configuration	Δw_0	Δw_a
Δ -10km	0.49	3.81
Δ -15km	0.40	2.65
2L-15km-45°	0.35	2.55
2L-20km-45°	0.34	2.40

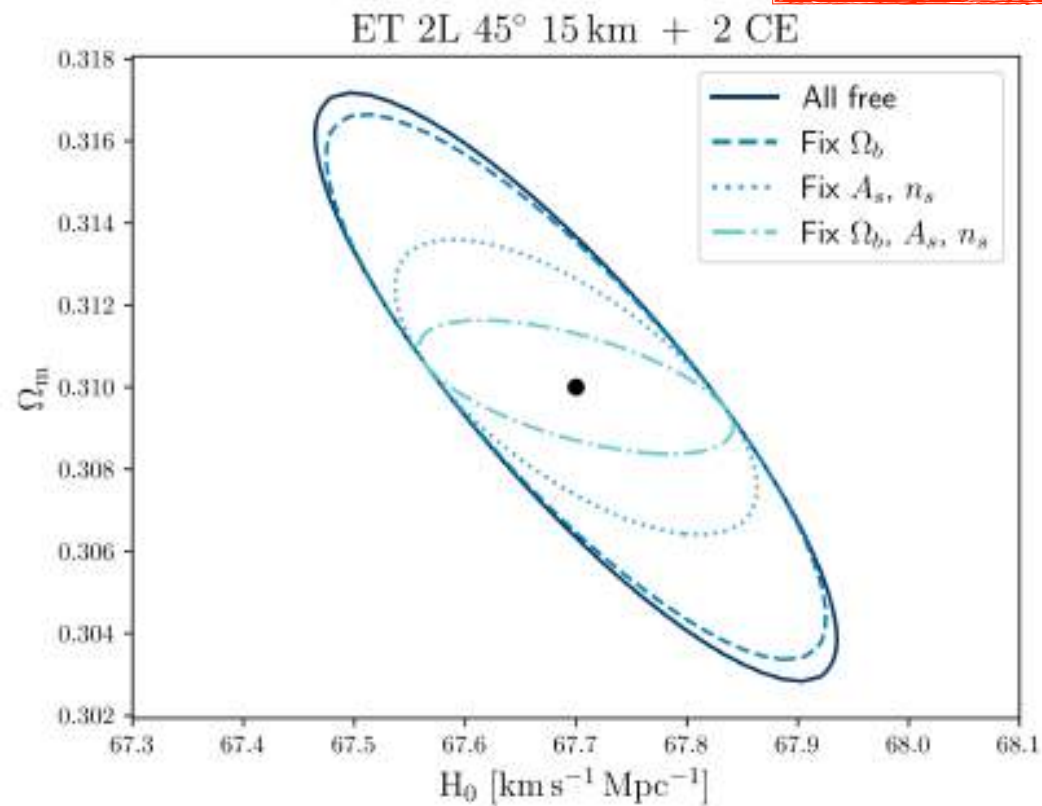
Einstein Telescope with 10³ bright sirens



joint GW+EM events obtained with ET+THESEUS, 5 yrs observation

Cross-correlation between GW and LSS (ETxEuclid)

$$C_{\ell}^{XY}(x_i, x_j) = \frac{2}{\pi} \int dk k^2 P(k) \Delta_{\ell}^{X, x_i}(k) \Delta_{\ell}^{Y, x_j}(k)$$



Parameter	All nuisance parameters free	Fix Ω_b	Fix A_s, n_s	Fix Ω_b, A_s, n_s
ET 2L 45° 15 km + 2CE				
H_0	0.8	0.75	0.55	0.47
Ω_m	5.3	4.9	2.6	1.2
ET Δ 10 km + 2CE				
H_0	0.8	0.79	0.6	0.53
Ω_m	5.3	4.9	2.6	1.2

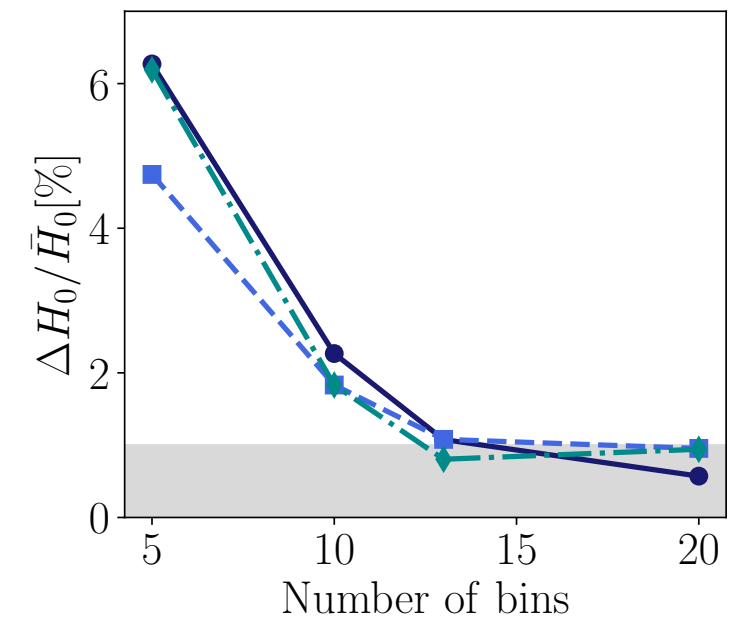


~ % level precision with one year of observation

(sky coverage of 15,000 square degrees and $\sim 1.6 \times 10^9$ observed galaxies)

Photometric sample of the Euclid satellite

Euclid Photometric Euclid Spectroscopic SKA



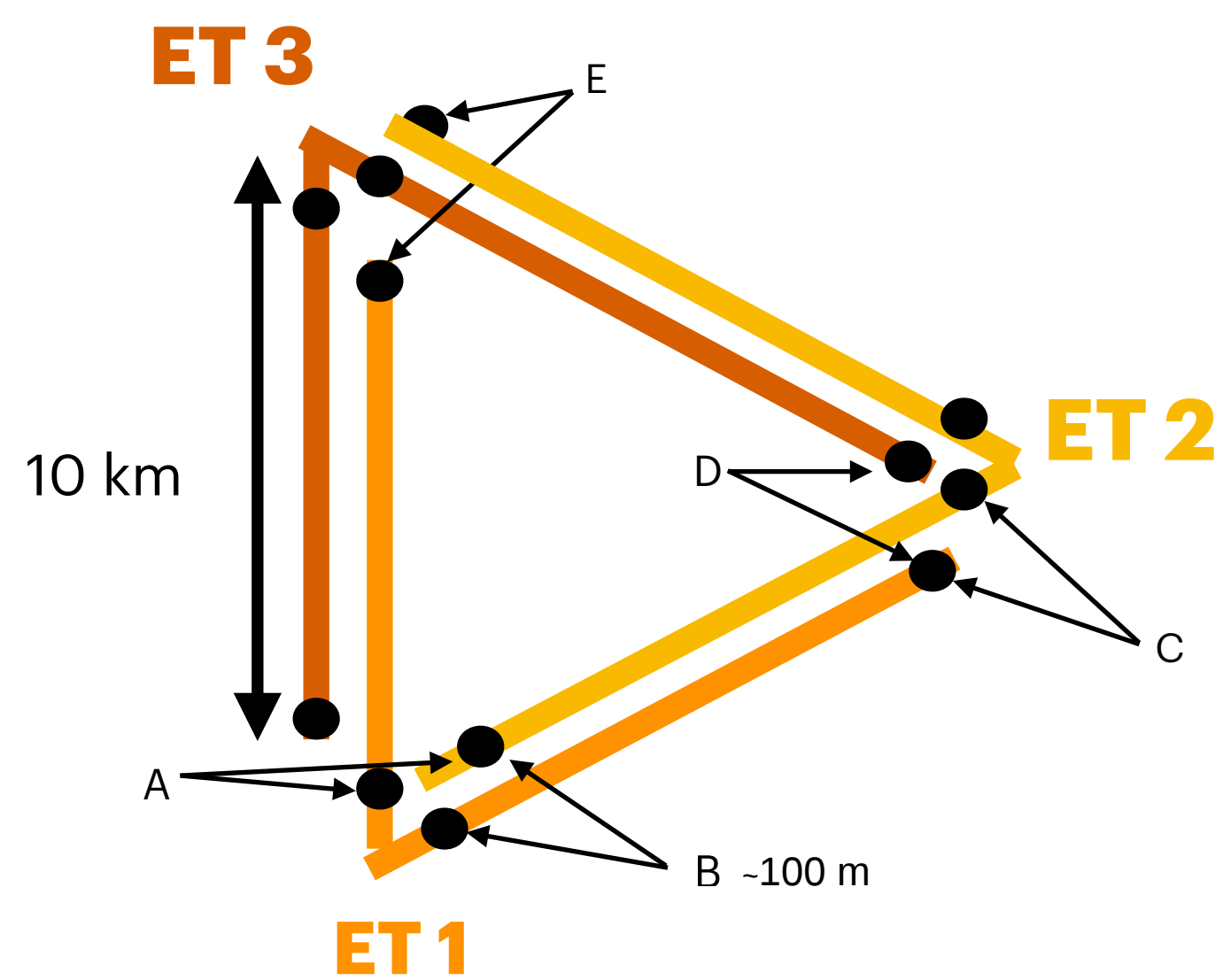
Impact of ET Correlated Noise on SGWB searches

Based on ArXiv: 2501.09057

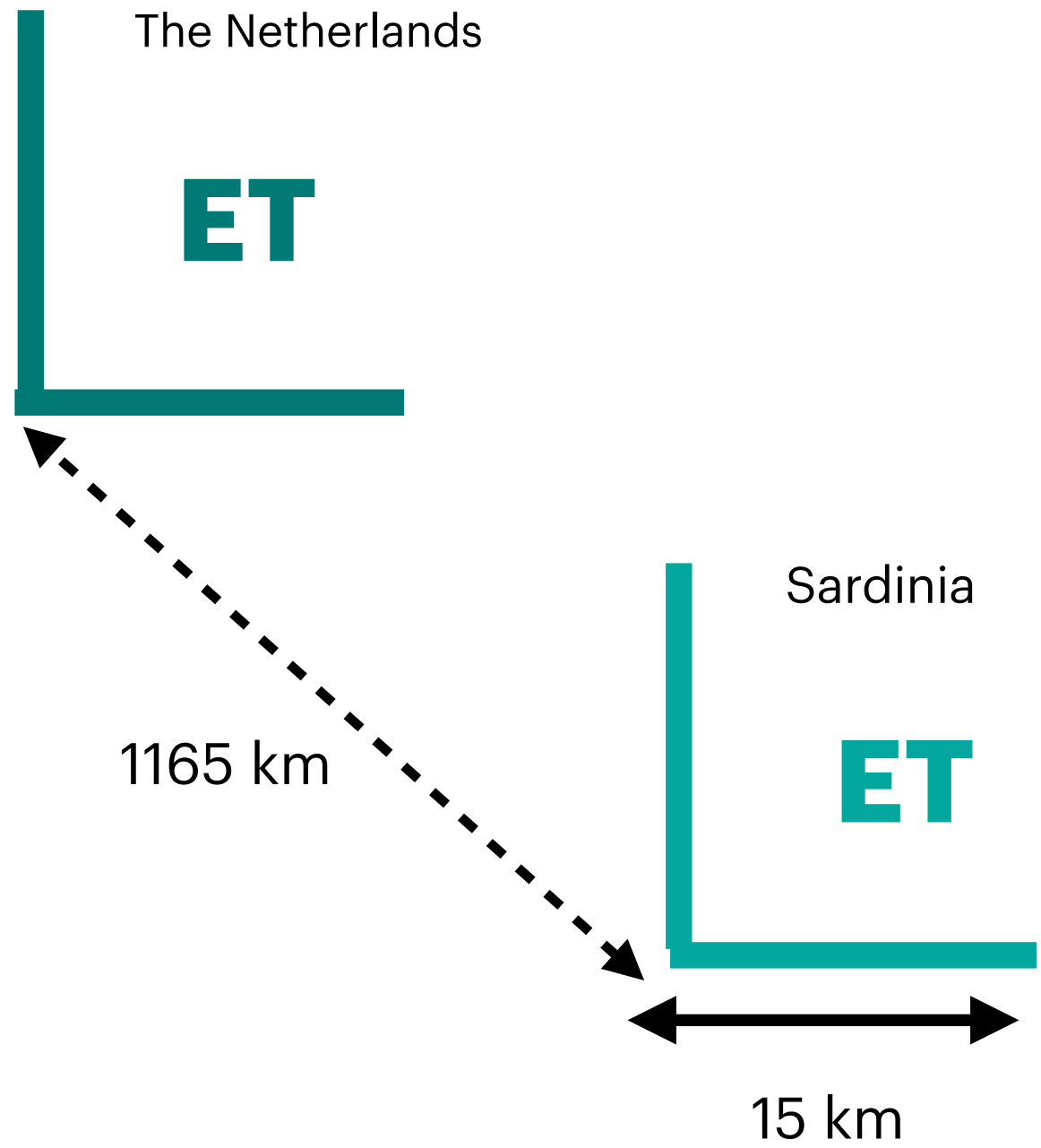
(I. Caporali, G. Capurri, W. Del Pozzo, AR, L. Valbusa Dall'Armi)

Einstein Telescope - possible designs

Triangular Configuration



2L Configuration

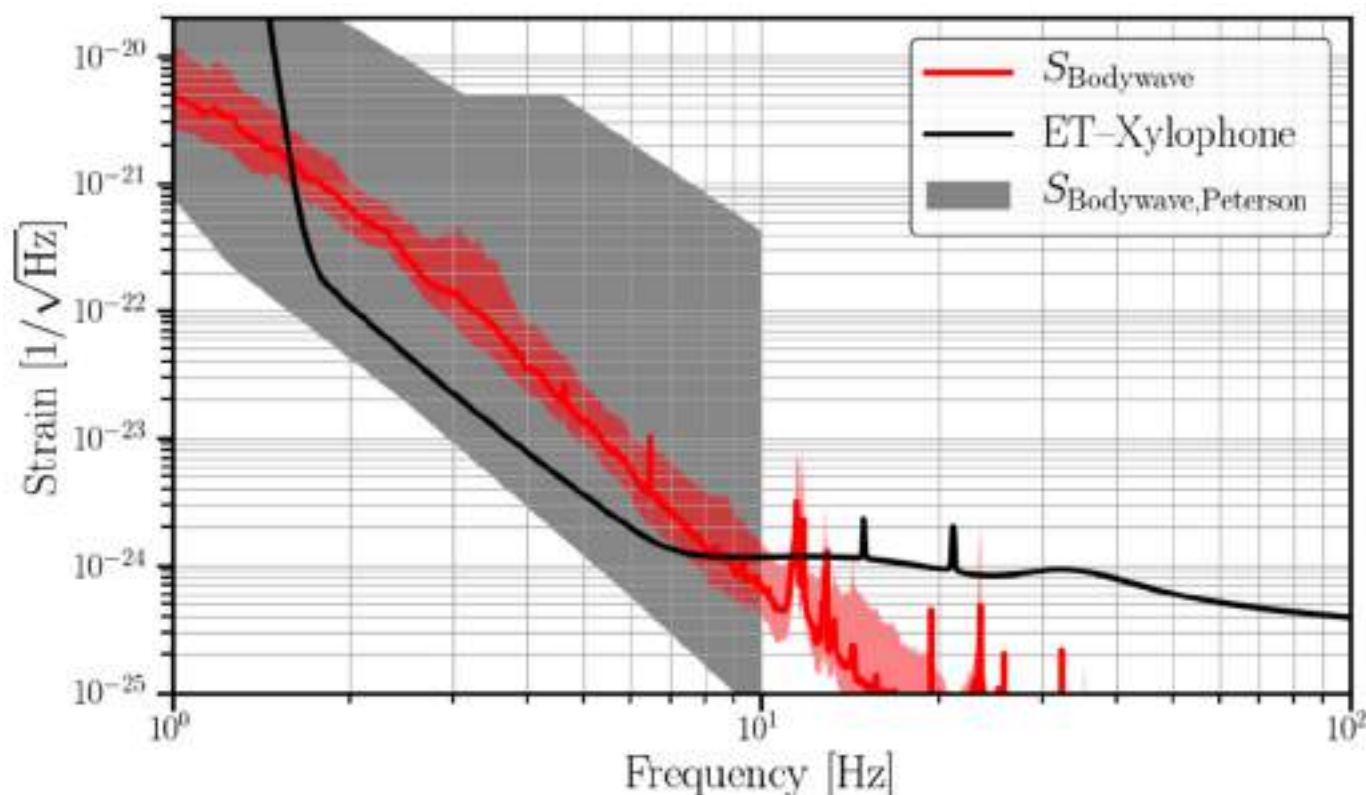


[ET CoBa paper JCAP, vol. 07,2023]

Correlated Newtonian Noise

Newtonian noise is a disturbance produced in GW detectors by local fluctuations in the Gravitational field

Local fluctuations generated by **changes** in density of rocks



Originated by passing seismic waves

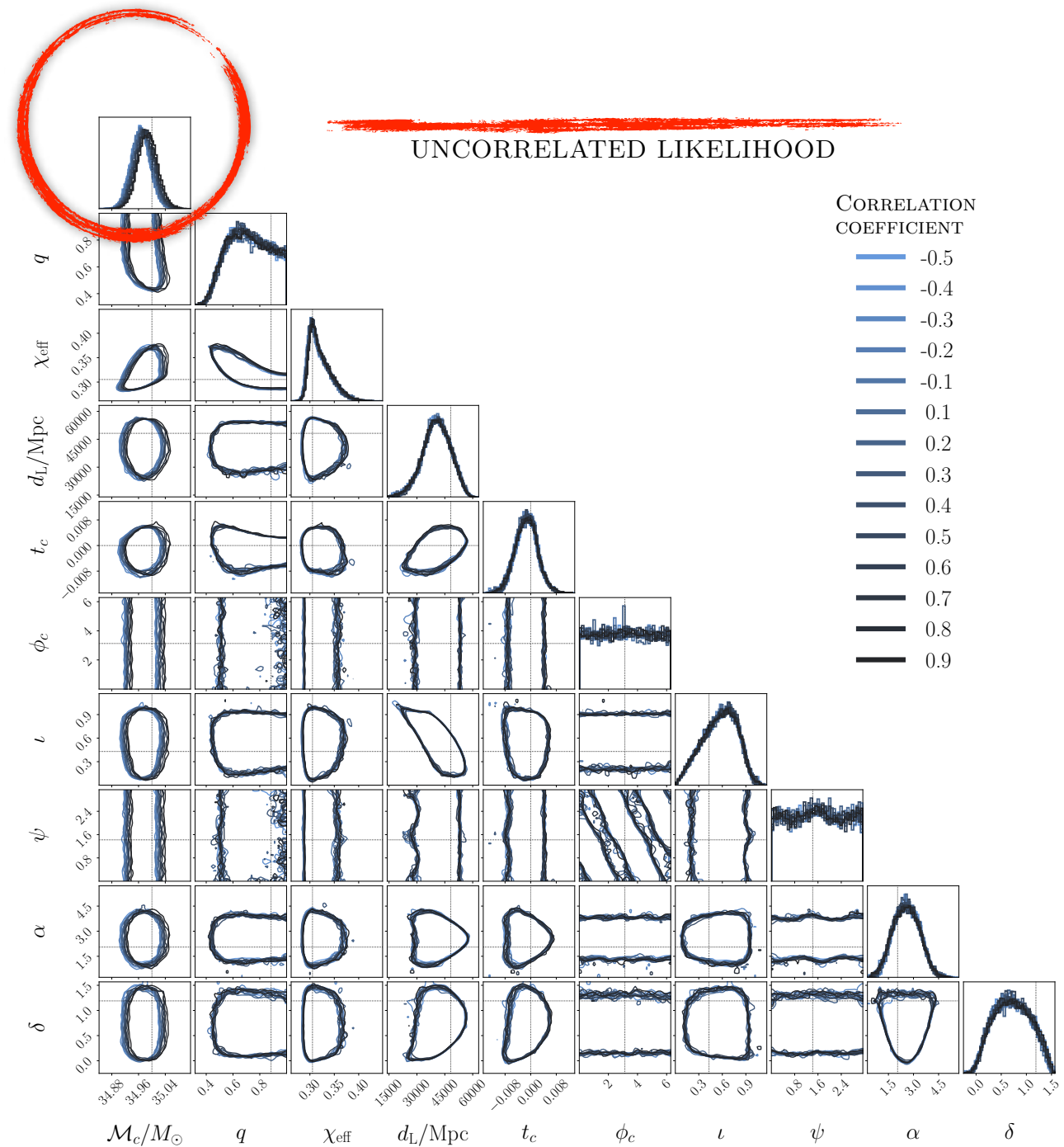
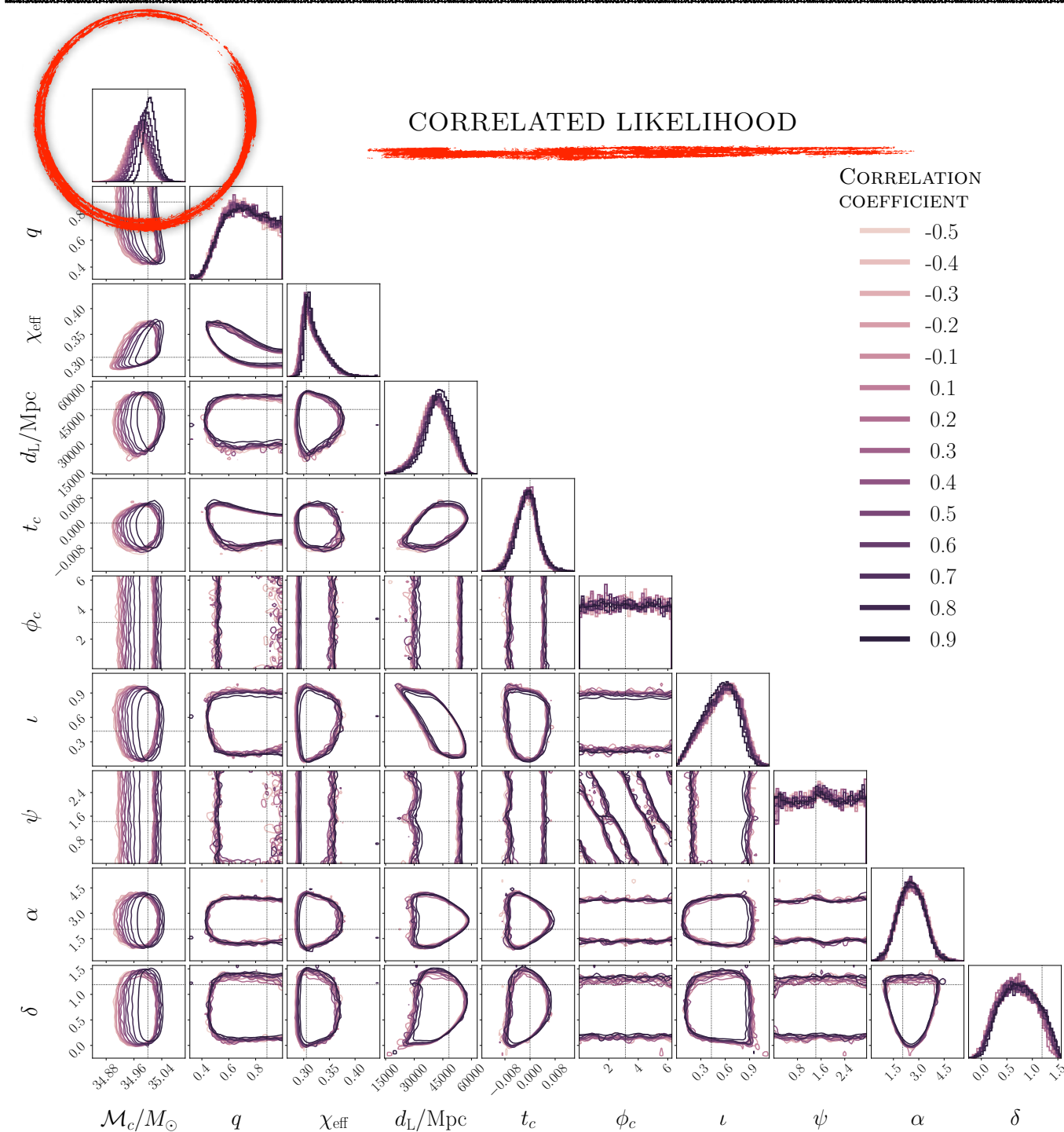
Body Waves

Surface Waves
(Rayleigh Waves)

Underground colocated interferometers will only be affected by this effect

[Janssens et al. arXiv:2206.06809]

Impact of Correlated Noise on PE Resolved Events



Neglecting these correlations may significantly reduce the accuracy of the chirp mass reconstruction.

[F.Cireddu et al. arXiv:2312.14614]

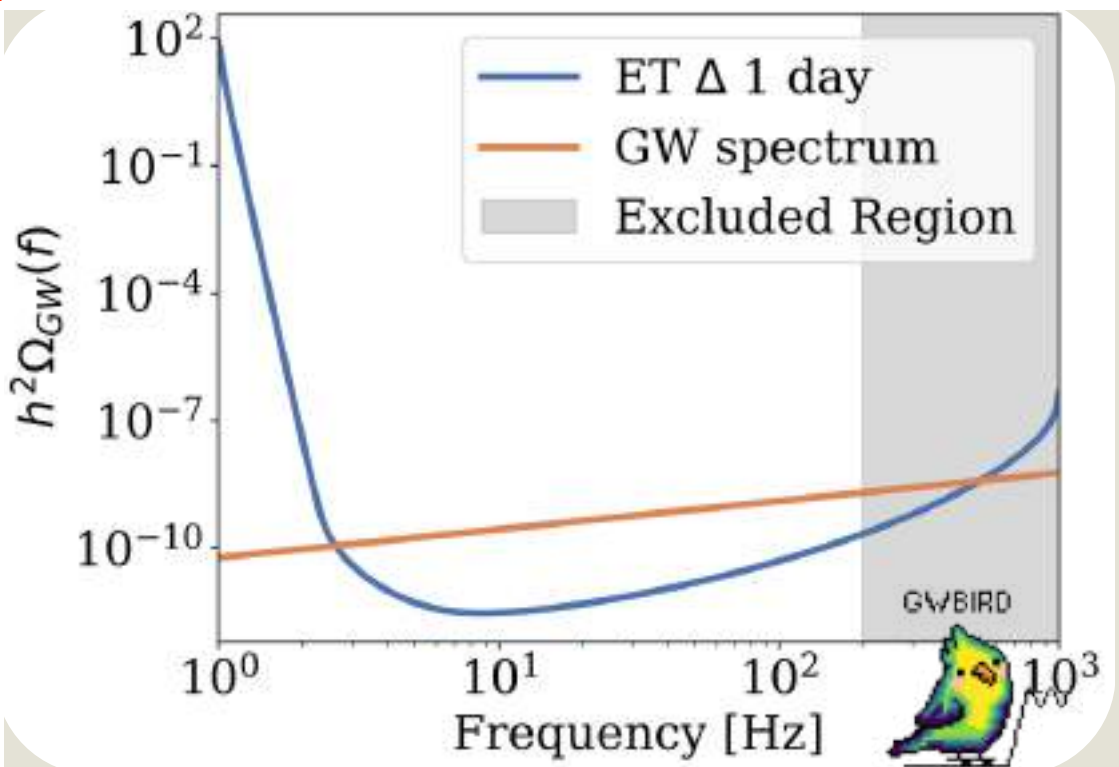
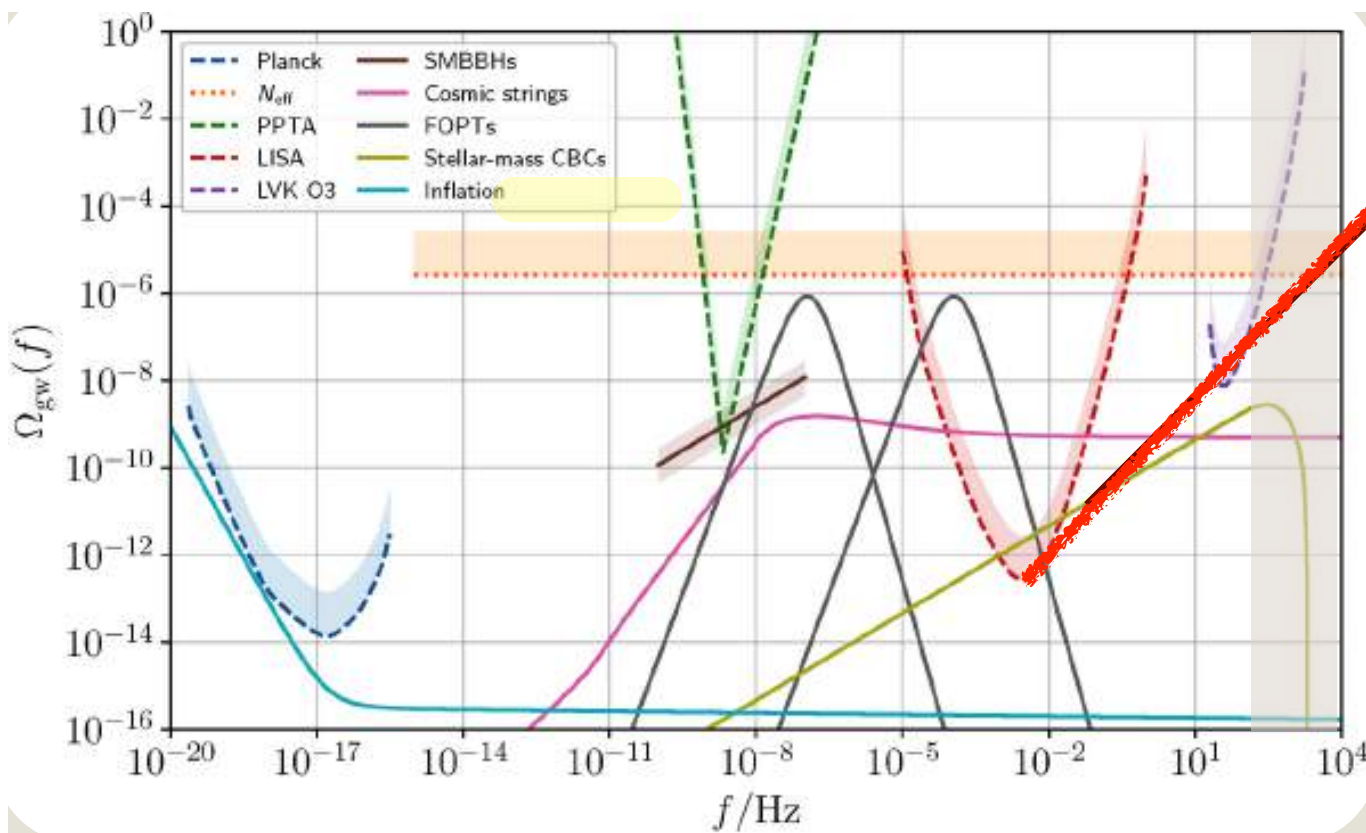
Impact of Correlated Noise SGWB

Correlated noise is a threat to all the possible **Astrophysical and Cosmological sources of SGWB**

$$\Omega_{\text{GW}} = \frac{1}{\rho_c} \frac{d\rho_{\text{GW}}}{d \log f}$$

$$\Omega_{\text{GW}}(f) = A_{\text{GW}} \left(\frac{f}{f_*} \right)^{n_{\text{GW}}}$$

Power Law Behavior



[I.Caporali, A. Ricciardone to appear]

Noise Covariance Triangular configuration

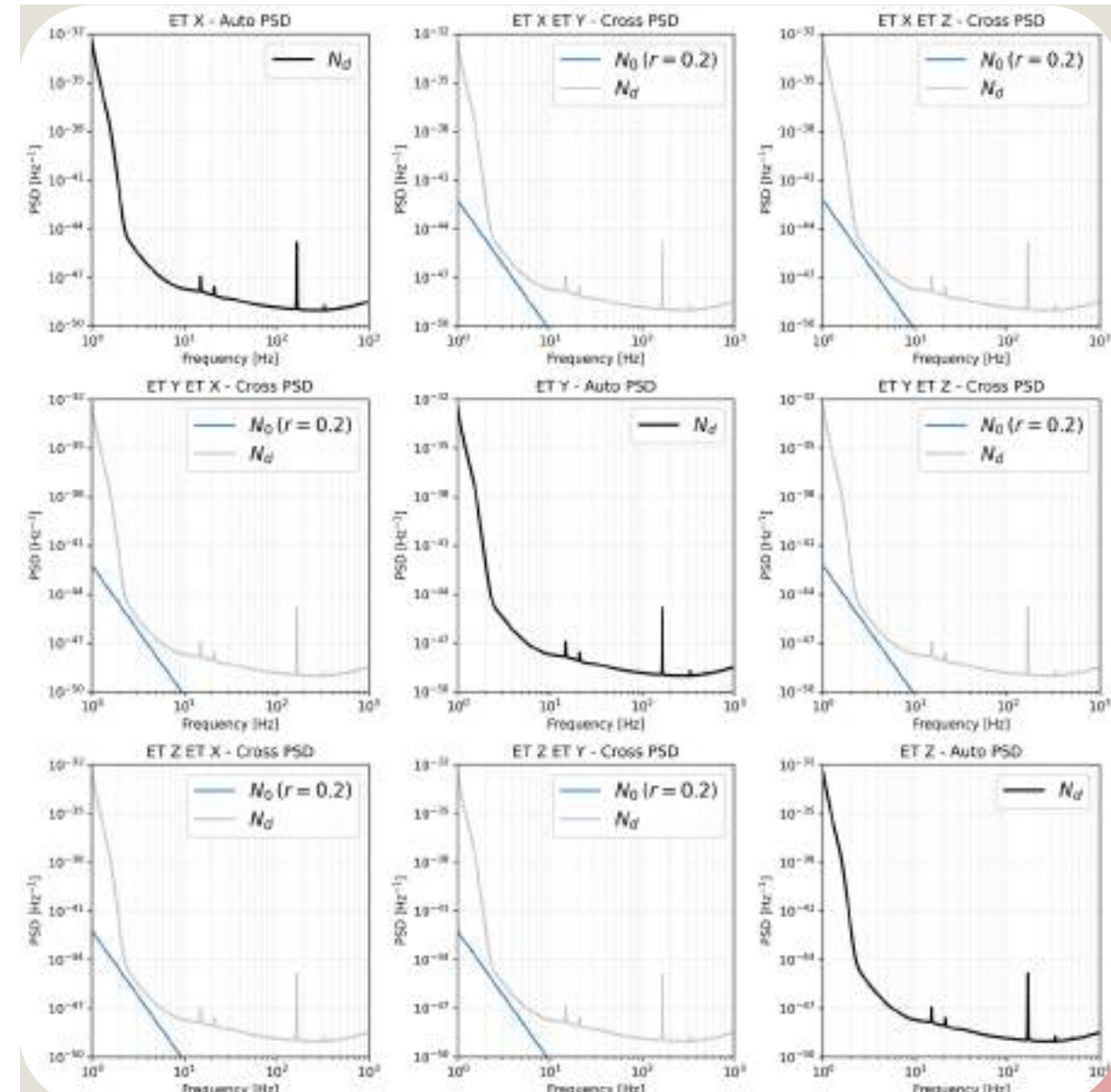
$$\langle \tilde{n}_I(f) \tilde{n}_J^*(f') \rangle \equiv \frac{\delta(f - f')}{2} N_{IJ}(f)$$

$$N_{IJ}(f) \equiv \begin{pmatrix} N_d(f) & N_o(f) & N_o(f) \\ N_o(f) & N_d(f) & N_o(f) \\ N_o(f) & N_o(f) & N_d(f) \end{pmatrix}$$

$$-1/2 \leq N_o(f)/N_d(f) \leq 1$$

Correlated noise template

$$N_o(f) = N_d(2.75 \text{ Hz}) r \left(\frac{f}{2.75 \text{ Hz}} \right)^{n_{\text{noise}}}$$

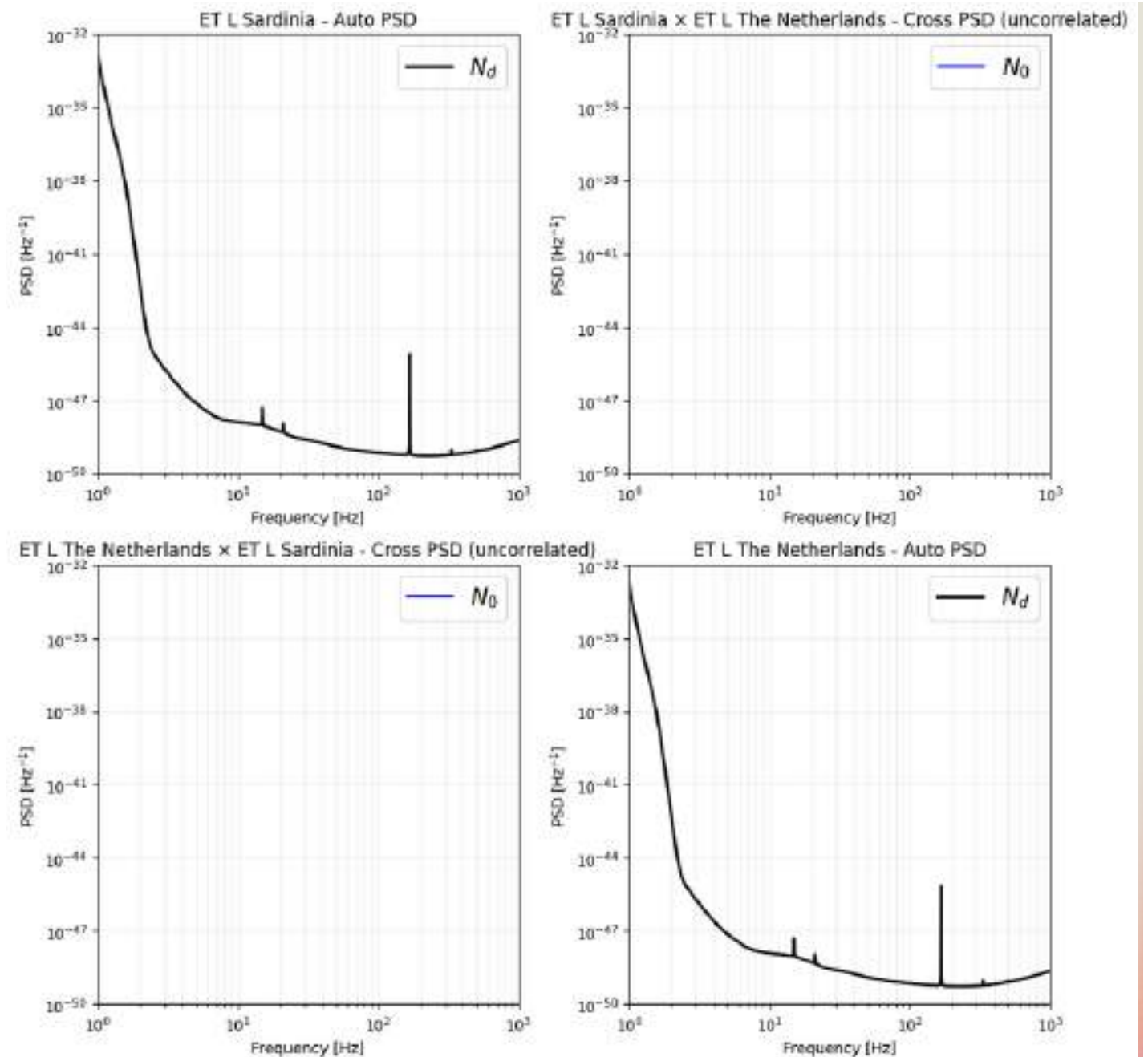


Noise Covariance 2L config

$$\langle \tilde{n}_I(f) \tilde{n}_J^*(f') \rangle \equiv \frac{\delta(f - f')}{2} N_{IJ}(f)$$

$$N_{IJ}(f) \equiv \begin{pmatrix} N(f) & 0 \\ 0 & N(f) \end{pmatrix}$$

Neglecting correlated magnetic noise



Likelihood for correlated and uncorrelated configurations

We introduce a time-averaged **estimator** of a **quadratic combination of the data**

$$\hat{C}_{IJ}(f) \equiv \sum_t \frac{2}{T_{\text{obs}} S_0(f)} \Re[\tilde{s}_I(t, f) \tilde{s}_J^*(t, f)]$$

[Abbott et al. 1903.02886]

[Abbott et al. 2101.12130]

t = time segments

$$T_{\text{seg}} = 4s$$

$$T_{\text{obs}} = 1\text{day}$$

$$S_0(f) = \frac{3H_0^2}{10\pi^2 f^3}$$

Averaging over many time segments allows us to treat the estimator of the SGWB as a **Gaussian random variable**, due to the central limit theorem

Likelihood for correlated and uncorrelated configurations

Average of the estimator

$$\bar{C}_{IJ}(f) \equiv \langle \hat{C}_{IJ}(f) \rangle = \gamma_{IJ}(f) \Omega_{\text{GW}}(f) + \frac{N_{IJ}(f)}{S_0(f)}$$

Covariance of estimator

$$\Sigma_{IJ}(f) \equiv \left\langle \left[\hat{C}_{IJ}(f) - \bar{C}_{IJ}(f) \right]^2 \right\rangle = \frac{\bar{C}_{II} \bar{C}_{JJ} + \bar{C}_{IJ}^2}{2N_{\text{seg}}}$$

(Gaussian) Likelihood

$$\mathcal{L} = \prod_{I,J} \frac{1}{\sqrt{2\pi \Sigma_{IJ}}} \exp \left[-\frac{1}{2} \frac{(\hat{C}_{IJ} - \bar{C}_{IJ})^2}{\Sigma_{IJ}} \right]$$

Triangular Configuration
(XYZ to the AET basis)
(I, J) = {(A,A), (E,E)}

2L Configuration
(I, J) = (ET1, ET2)

SGWB Data Generation

Noise and SGWB in different time segments are sampled from Nseg independent Gaussian distributions.

Data in different channels are correlated, implying they are generated from a multivariate Gaussian distribution.

Signal

$$h_{\text{ampl}}(f) = \sqrt{\frac{T_{\text{obs}}}{2} \frac{3H_0^2}{10\pi^2 f^3} \Omega_{\text{GW}}(f) \gamma_{ij}(f)}$$

$$h(f) = h_{\text{ampl}} \left(\mathcal{N}(0, 1/\sqrt{2}) + i\mathcal{N}(0, 1/\sqrt{2}) \right)$$

Noise

$$n_{\text{ampl}}(f) = \sqrt{\frac{T_{\text{obs}}}{2} N_{ij}(f)}$$

$$n(f) = n_{\text{ampl}} \left(\mathcal{N}(0, 1/\sqrt{2}) + i\mathcal{N}(0, 1/\sqrt{2}) \right)$$

Data stream

$$s(f) = h(f) + n(f)$$

Assuming that resolved sources are perfectly subtracted and neglecting time-frequency induced correlation

PE Outcomes

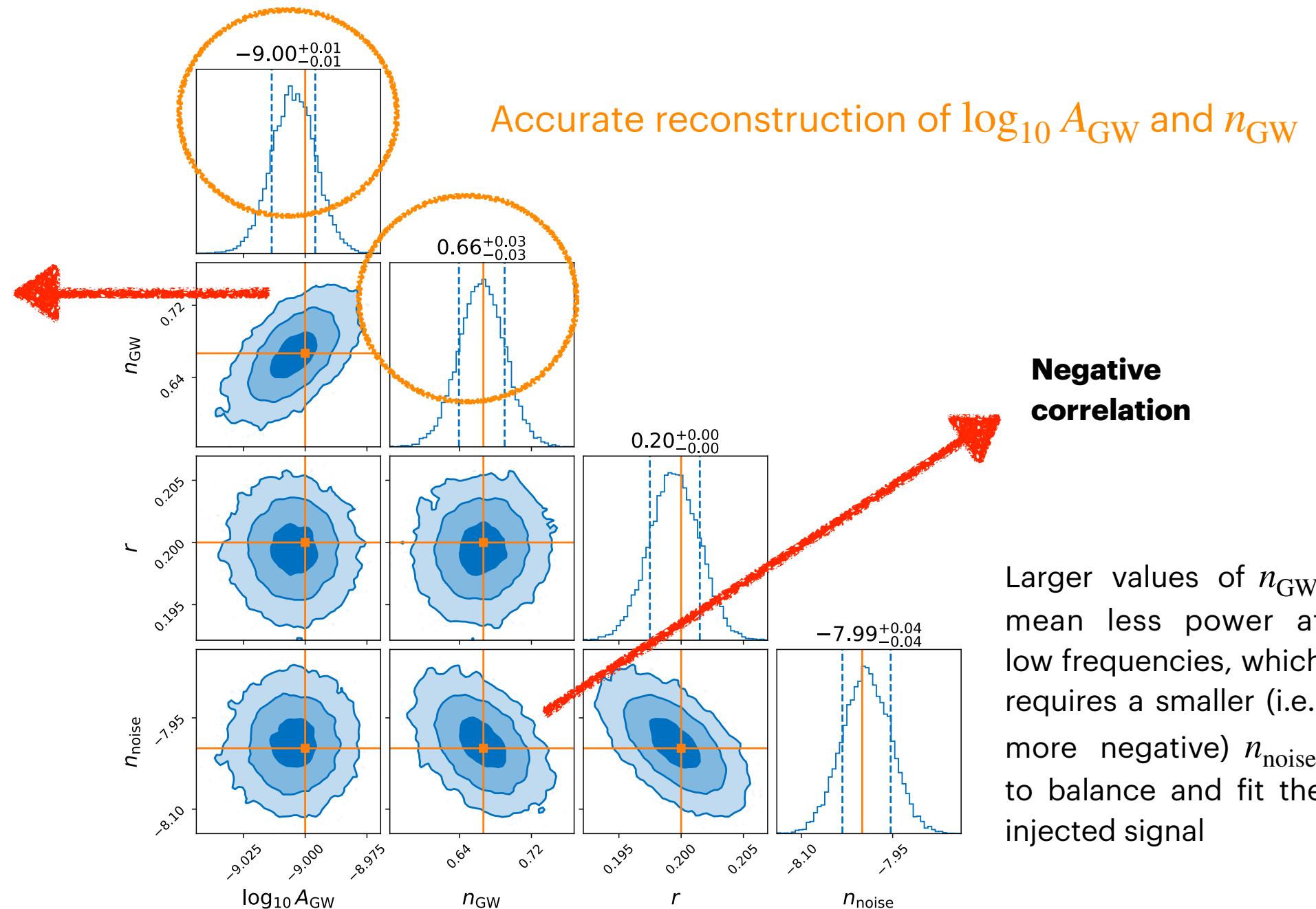
We implement our likelihood in **Bilby** and reconstruct the posterior using the **Dynesty** sampler

Positive correlation

Since the chosen pivot frequency for the SGWB, 25 Hz, is higher than the frequencies at which ET is most sensitive to the SGWB (~10 Hz), a larger retrieved value of n_{GW} requires a larger value of $\log_{10} A_{\text{GW}}$ to fit the injected signal

Injected parameters:

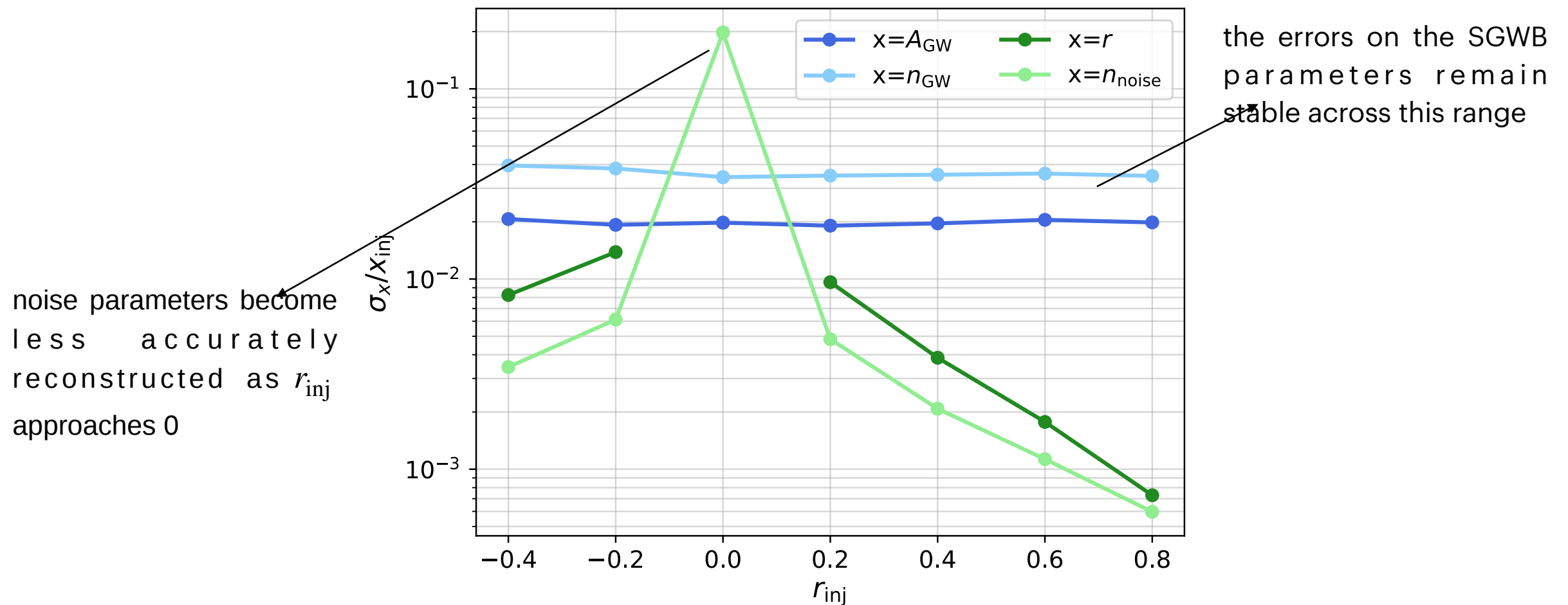
- $\log_{10} A_{\text{GW}} = -9$
- $n_{\text{GW}} = 2/3$
- $r = 0.2$
- $n_{\text{noise}} = -8$



PE Accuracy

Accurate reconstruction of the SGWB

(Having a model for the noise)

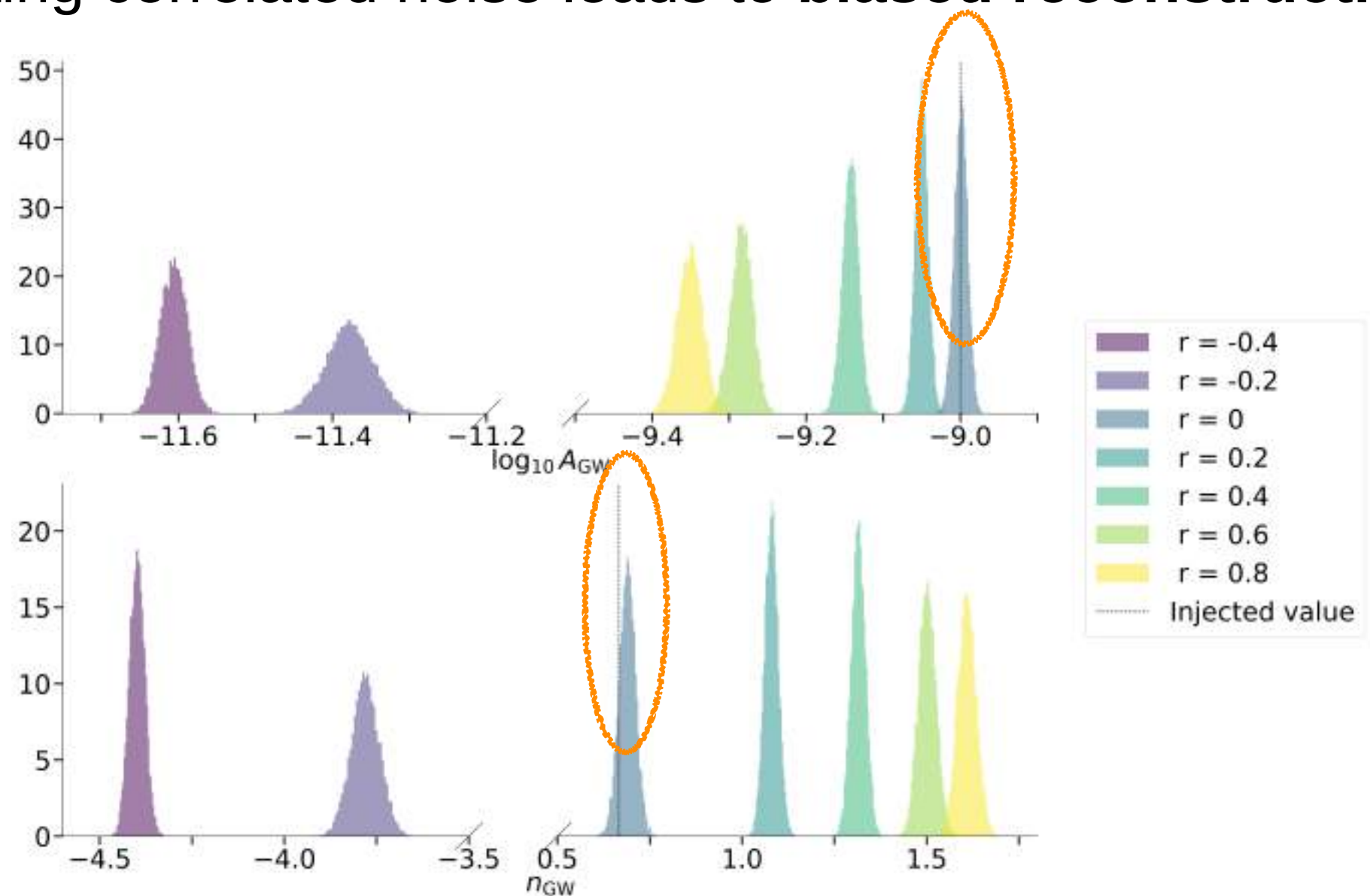


Main message: correlated noise with smaller power is inherently harder to reconstruct.

(To be tested also with noise agnostic searches)

PE Biased Results

Neglecting correlated noise leads to **biased reconstruction**



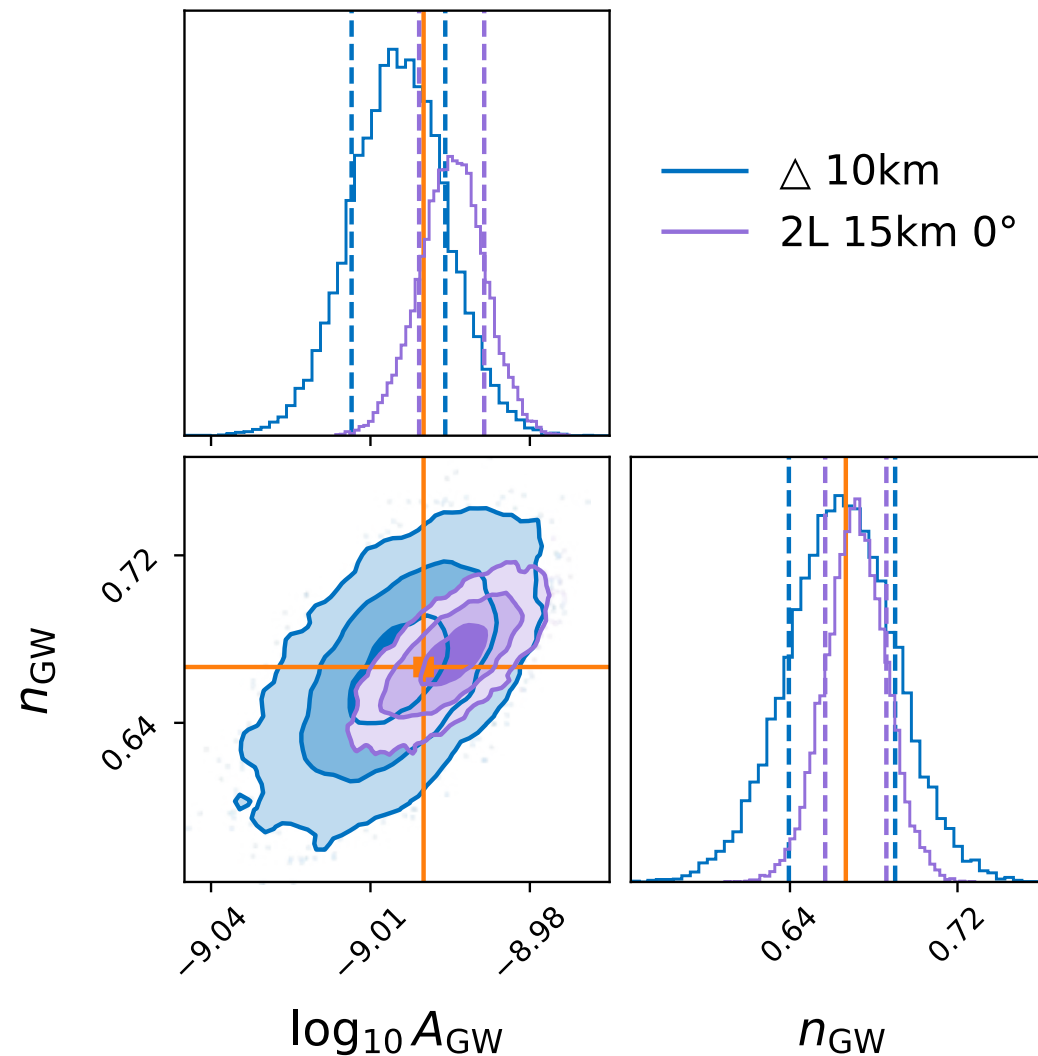
The injected parameters are **correctly retrieved** only for $r_{\text{inj}} = 0$, since there is **no correlated noise** in this case and the likelihood that neglects it is accurate.

14

The presence of **correlated noise** completely disrupts the signal interpretation when using the incorrect likelihood that does not account for it

Triangle vs 2L performance

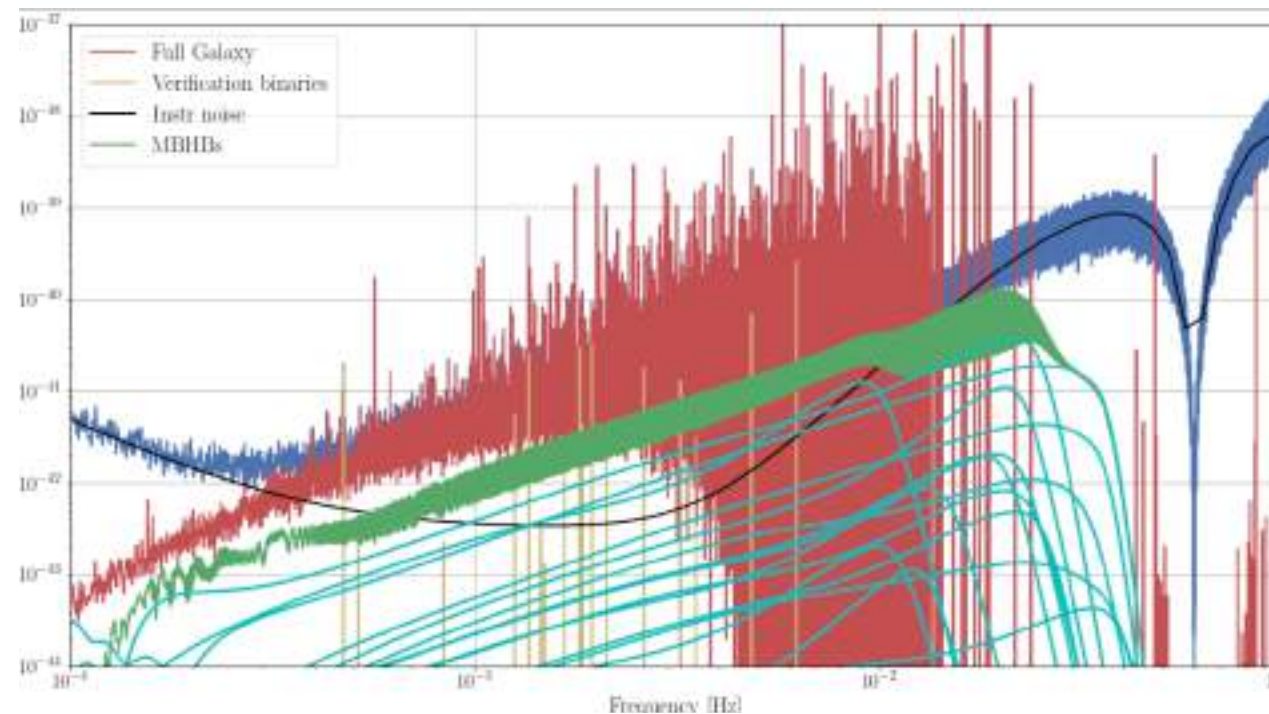
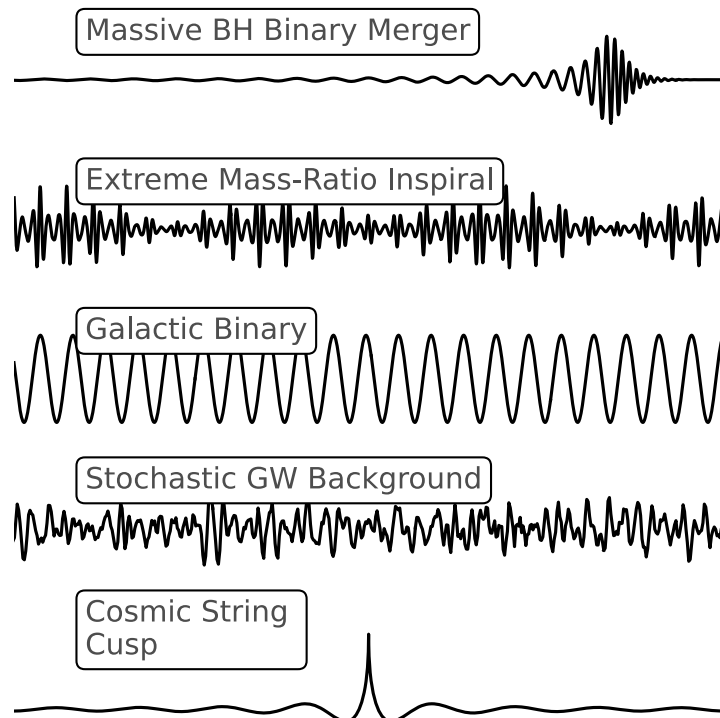
Triangular configuration is **competitive** to the 2L
(If noise is properly modeled)



Why the 2L performs slightly better?

Mainly because the **15 km arms** of the **2L** provide better sensitivity to the SGWB compared to the **triangle** with **10 km arms**.

How to approach the data analysis - Global Fit



- Large number of overlapping sources
- Residuals from sources subtraction
- Confusion from unresolved sources
- One doesn't cross-correlate like LVK
- Prediction of the instrumental noise?
- Artefacts: gaps in the data, glitches...

LISA Ground Segment - Italia

Develop a GLOBAL FIT pipeline for LISA

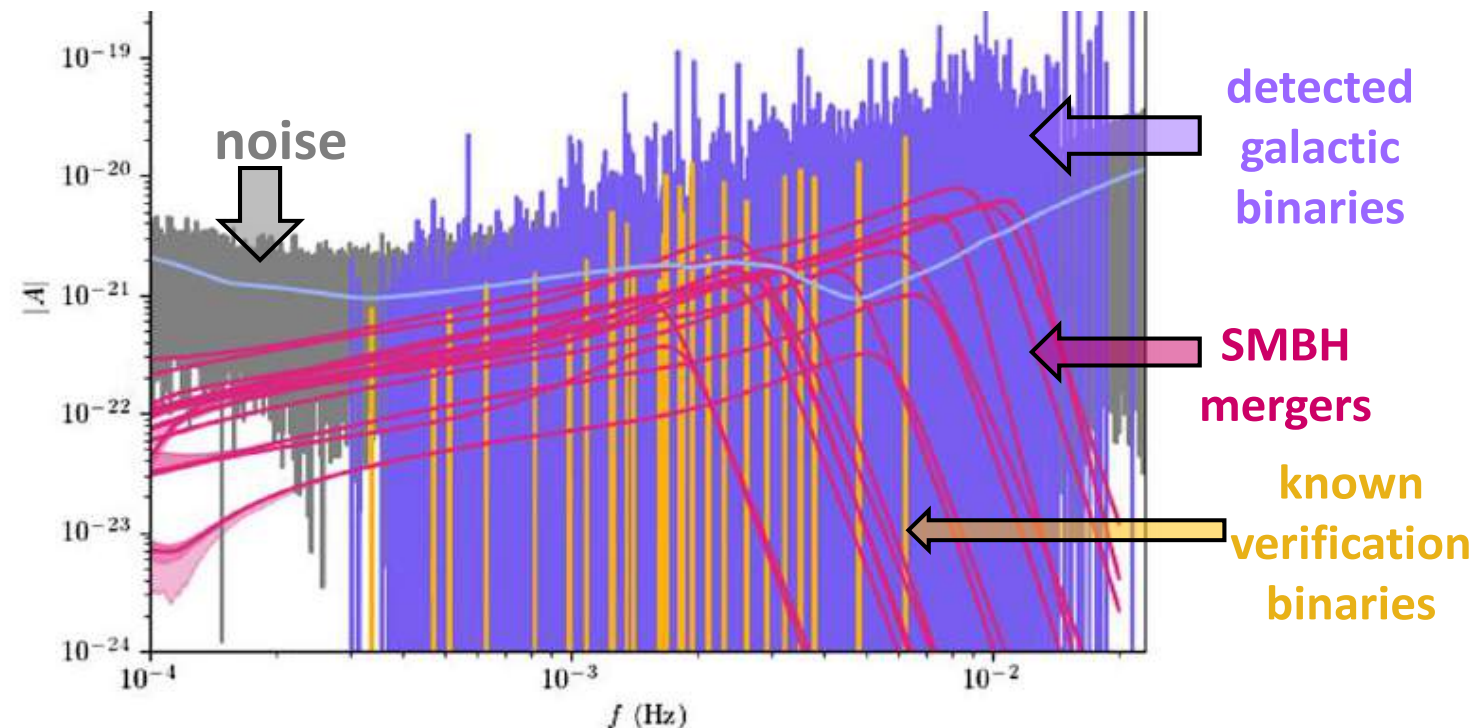
Pisa node

(Similar approach for ET)

Cornish, Littenberg, 23

Marsat+,24

Katz, Karnesis et al, '23



Cornish & Littenberg 2023

Conclusions

- All **GW interferometers** have been initially **conceived as GW astrophysics observatories**, they have **not been designed to do cosmology**
- **However**, they can provide new information on a variety of scales: from the Galaxy to Hubble scales, from the present time to the **very early universe** -> therefore they can be used as a **cosmological observatory** as well
- **We can have access to energy scales not accessible in any collider**
- We can test the ***late-time universe*** through the observation of the GW emission from compact binaries, and **constrain cosmological parameters**.
- Importance of including all detector layout effect in **SGWB parameter estimation**
- A lot of cosmological expertise in this community to push forward on GW topics

Thank you!