

From data to the analytic S -matrix: A Bootstrap Fit of the $\pi\pi \rightarrow \pi\pi$ amplitude

Based on arXiv: 2410.23333

with K. Häring, N. Su

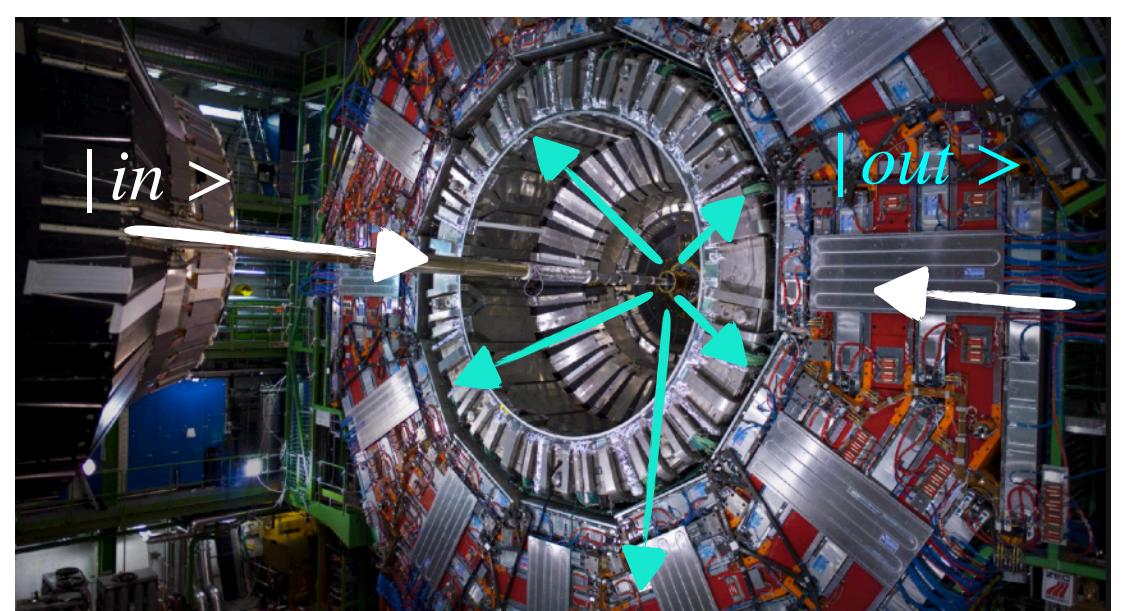
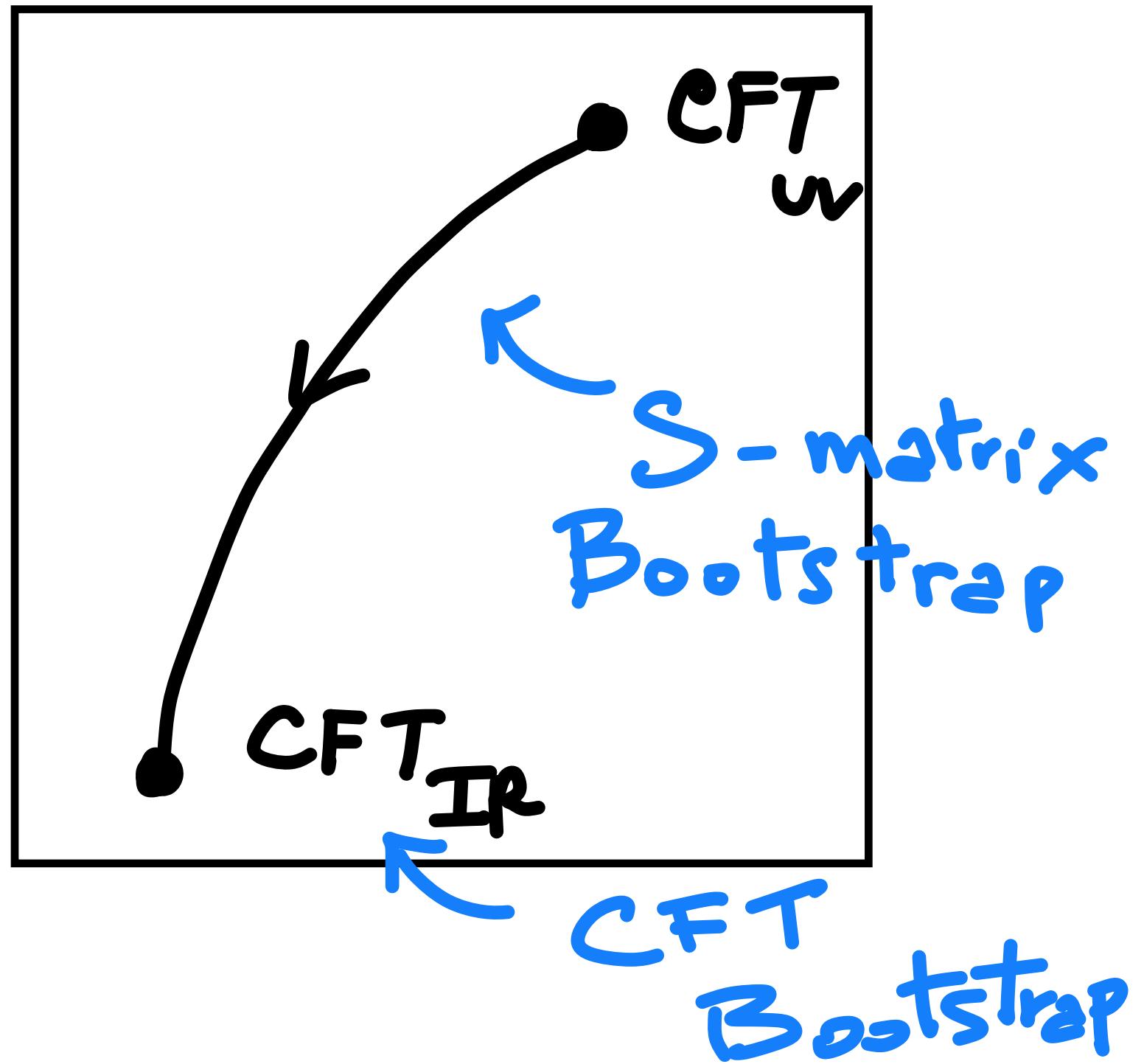
23/04/25

Bootstrap

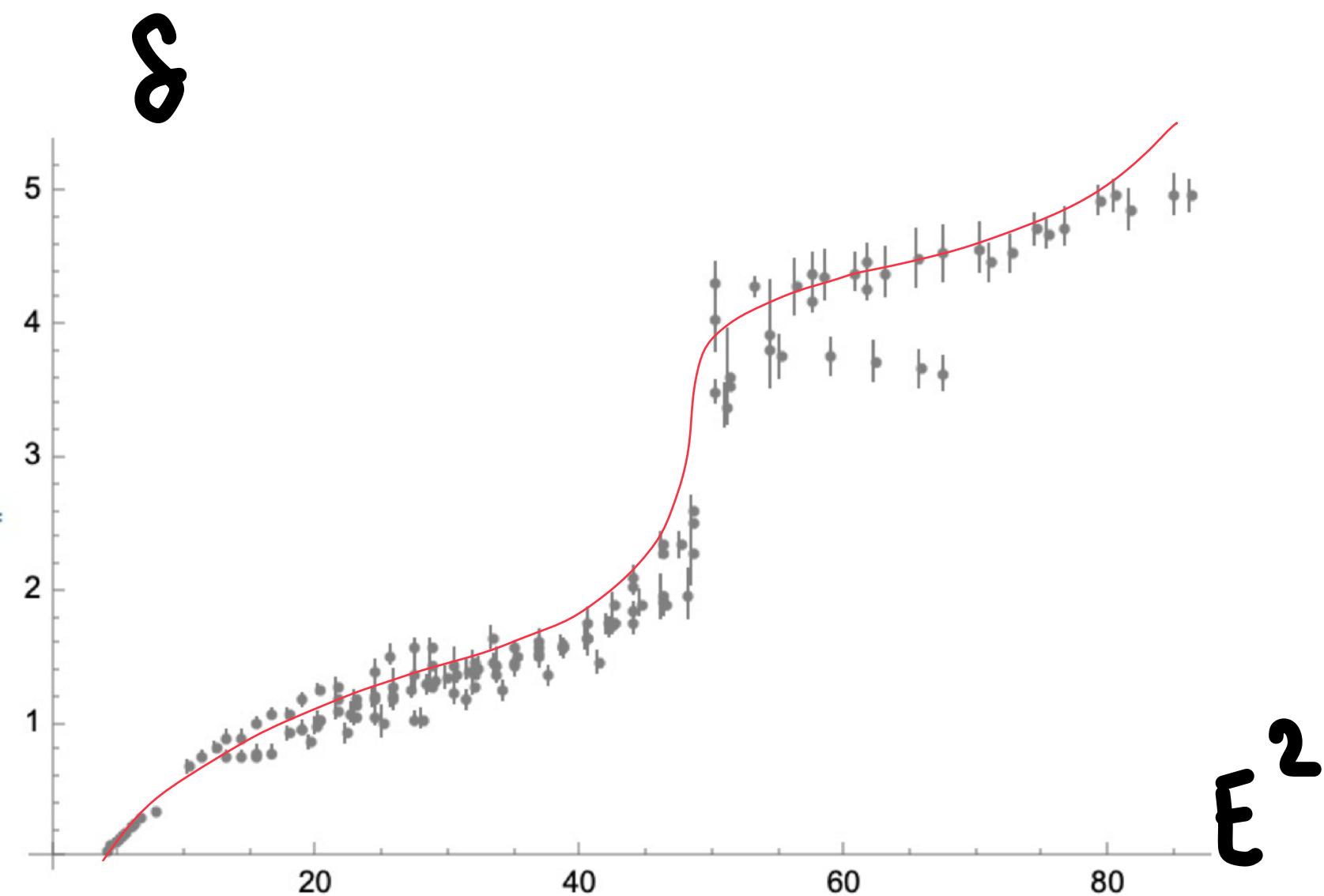
Map out the space of QFT_S
↳ Map out the space of Observables

CFT correlators
↳ critical exponents,
OPE coefficients , ...

QFT S -matrix / correlators
↑
this talk
↳ spectrum, couplings, ...



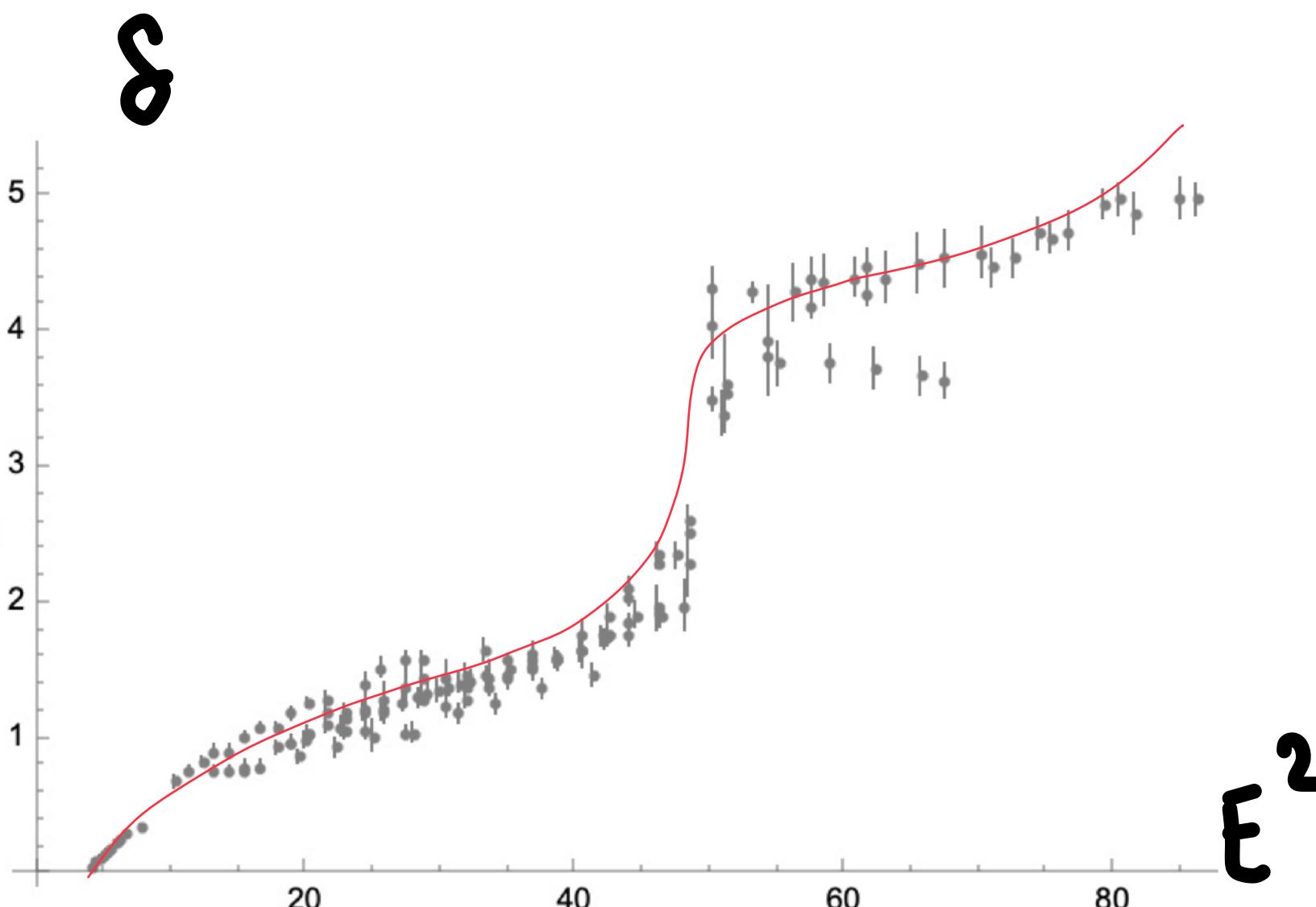
General Idea (I)



MODEL

Spectrum
Couplings
low Energy constants
...

General Idea (I)



MODEL

- 1) Isobar Models
(background + resonance)
- 2) Unitarized χ -PT approaches
- 3) Roy Equations [Caprini, Colangelo, Gasser, Leutwiler]

Spectrum

Couplings

low Energy constants

...

4) Can we use the Bootstrap to construct a better Model?

General Idea (II)

The ideal model is the QFT scattering amplitude

- 1) Causality \leftrightarrow Analyticity
- 2) Crossing Symmetry
- 3) Unitarity

Plan of the talk is to explain how to use the Bootstrap to construct such a model.

(Roy Equations close to satisfy all properties,
but limitations: low energy, low spins)

Pions trivia

1) lightest particles in the QCD spectrum

$$m_\pi \approx 140 \text{ MeV}$$

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(3 pions)
moderately weakly coupled @ low energy

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 $\tau_{\pi^\pm} \approx 2.6 \times 10^{-8} \text{ sec}$
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- 4) Approximate Isospin Symmetry $\frac{m_{\pi^+} - m_{\pi^0}}{\bar{m}_\pi} \approx 0.03$

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 ~~$\tau_{\pi^0} \approx 8.4 \times 10^{-17} \text{ sec}$~~ $\rightarrow \infty$
- 4) ~~Approximate Isospin Symmetry~~ $\frac{m_{\pi^+} - m_{\pi^0}}{m_\pi} \approx 0.03 = 0$

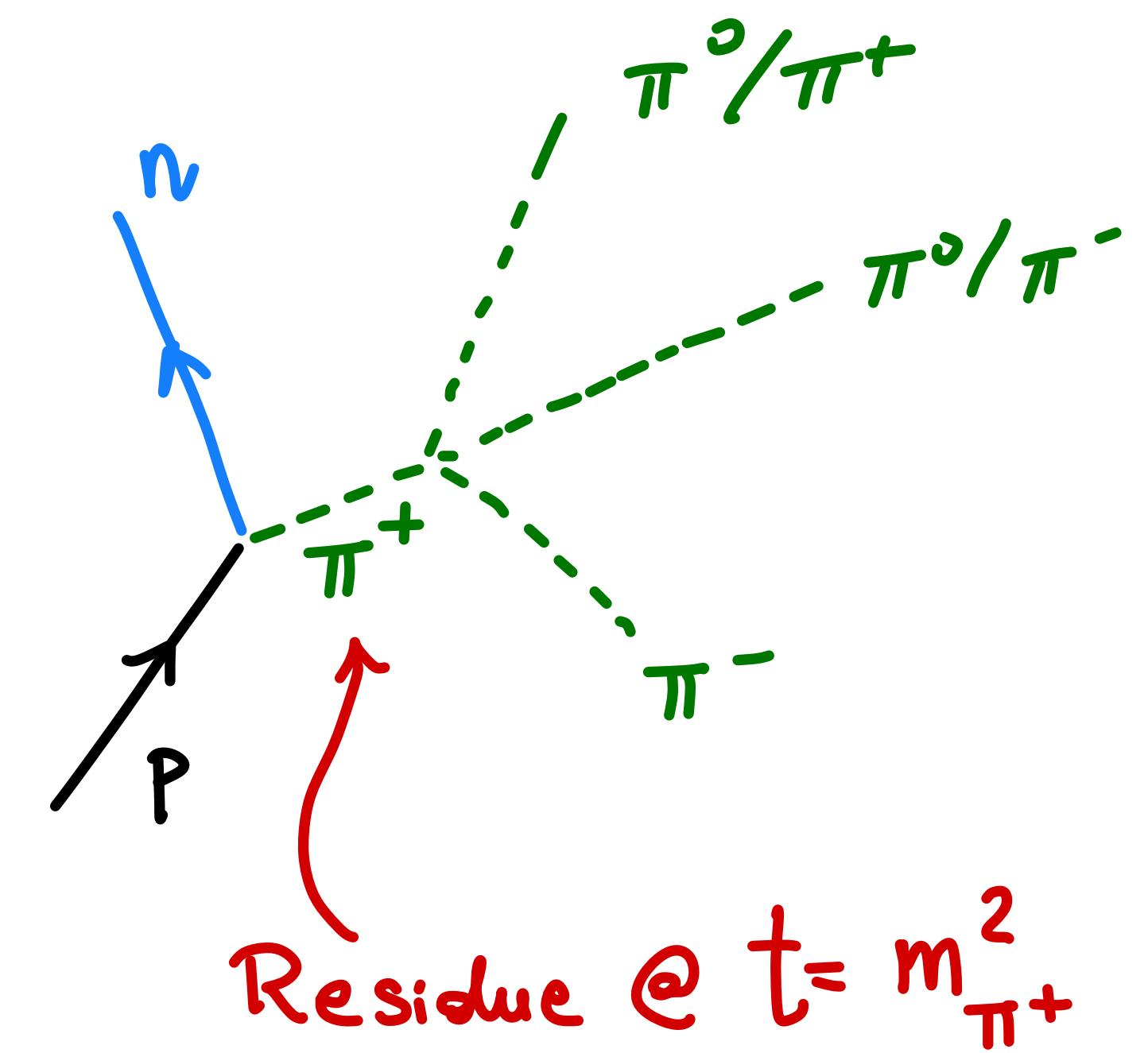
Pion S-matrix \cong S-matrix of $SO(3)$ scalars

$\pi\pi$ scattering Trivia

In Experiments we have to deal with 3), 4)

- CERN-MUNICH Coll. 60's-70's

$$\pi N \rightarrow \pi\pi N$$



$\pi\pi$ scattering Trivia

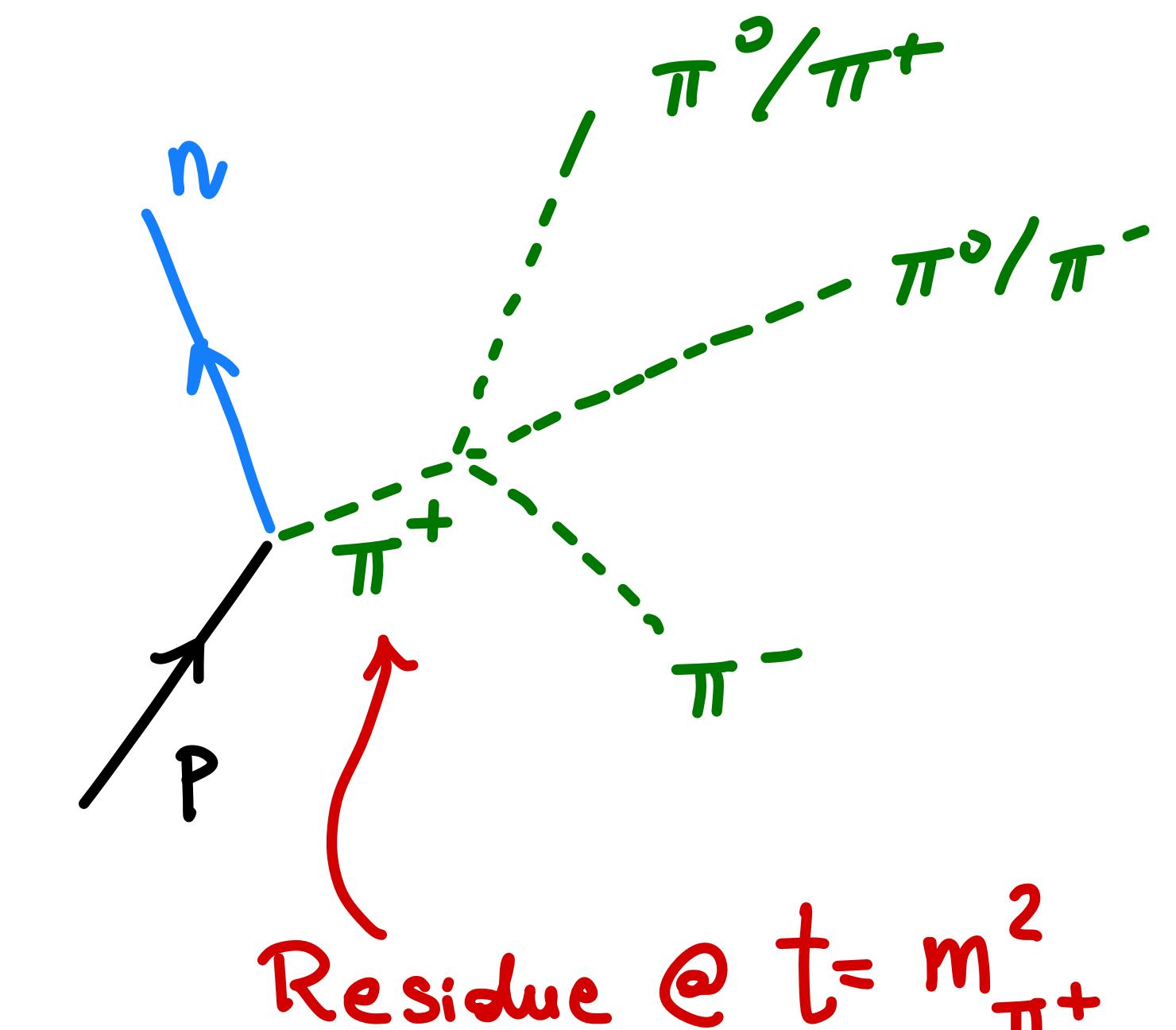
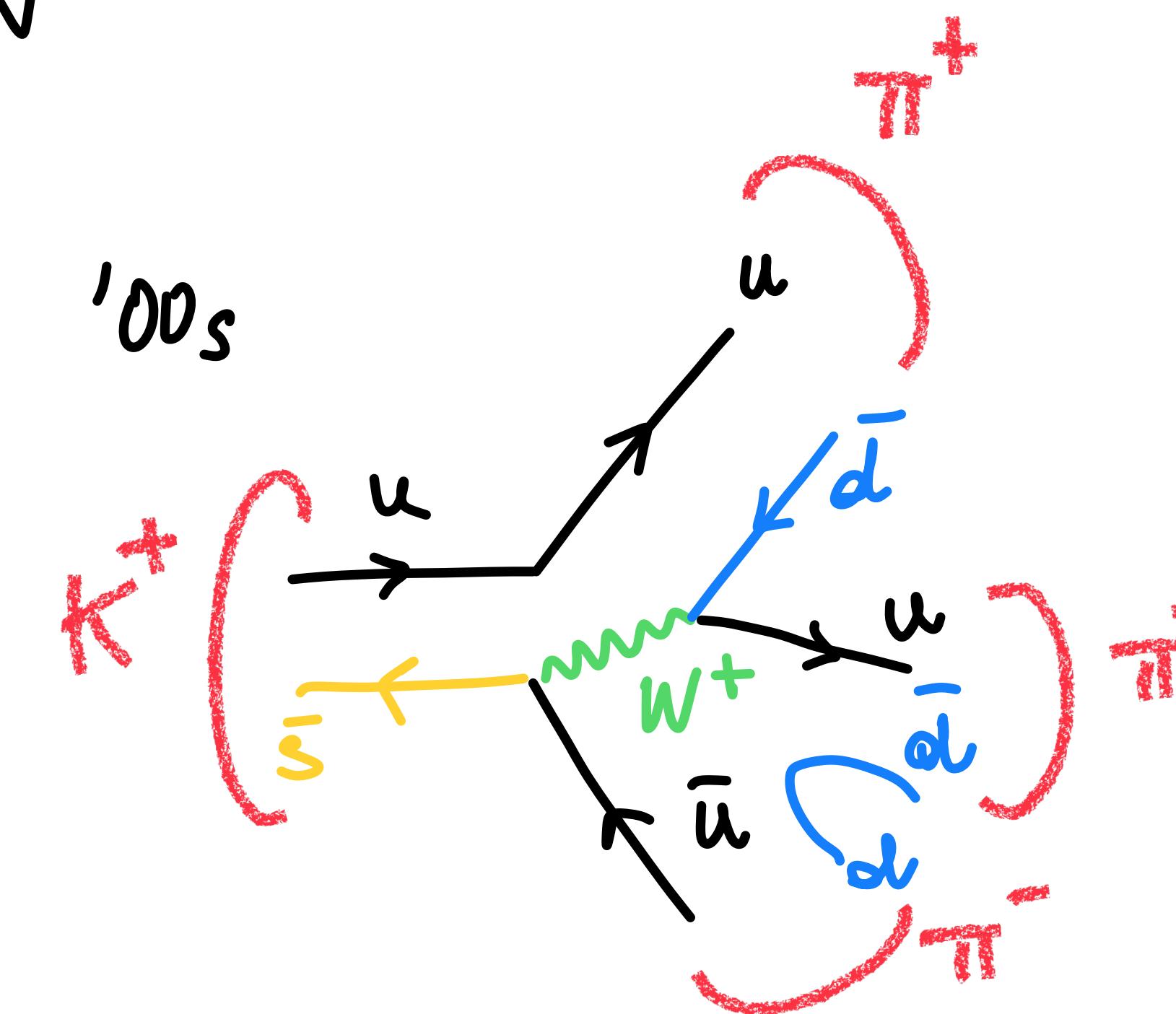
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- NA48 Experiment '00s

$$K \rightarrow \pi\pi\pi \\ (80 \text{ MeV})$$



Residue @ $t = m_{\pi^+}^2$

$\pi\pi$ scattering Trivia

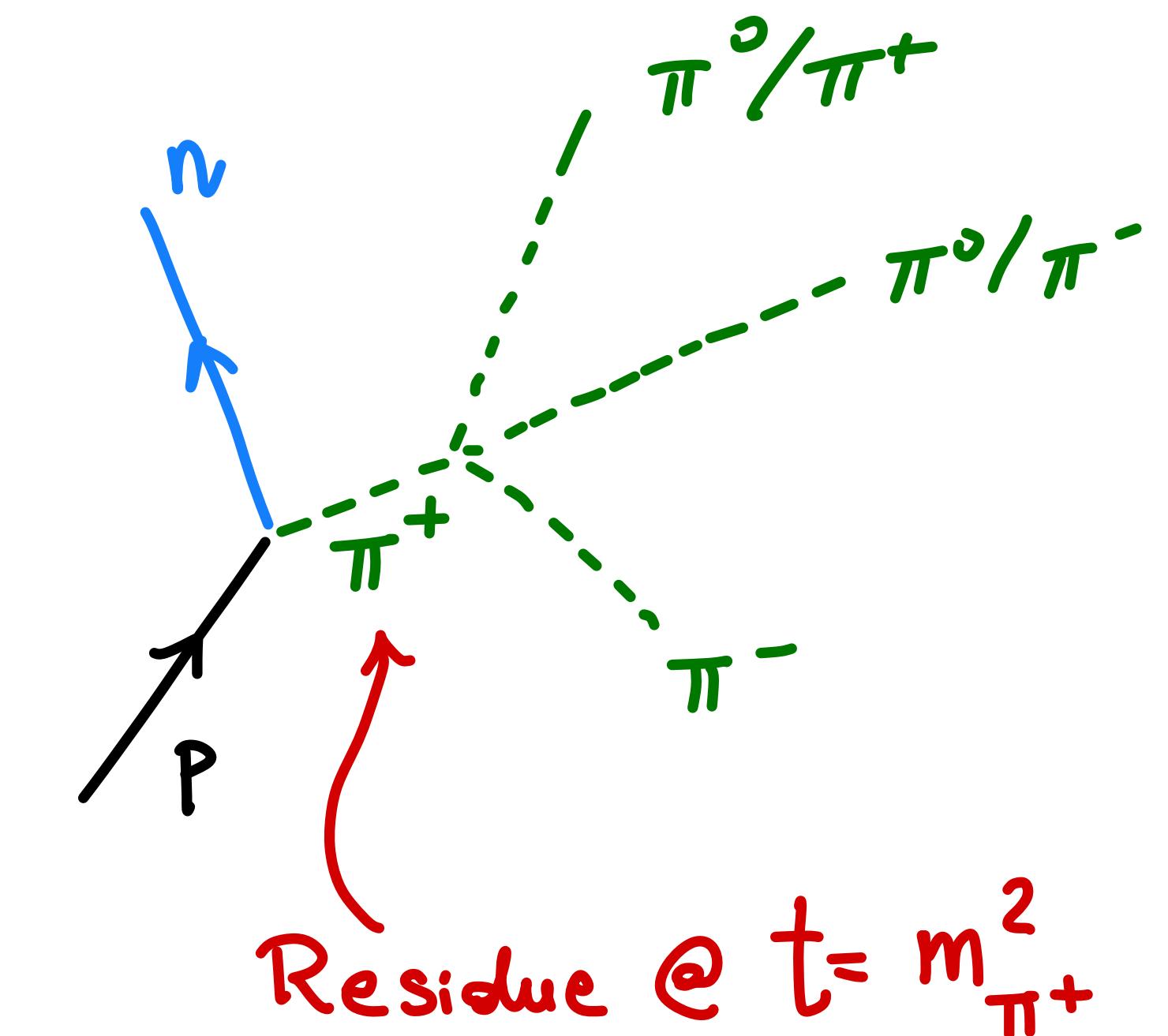
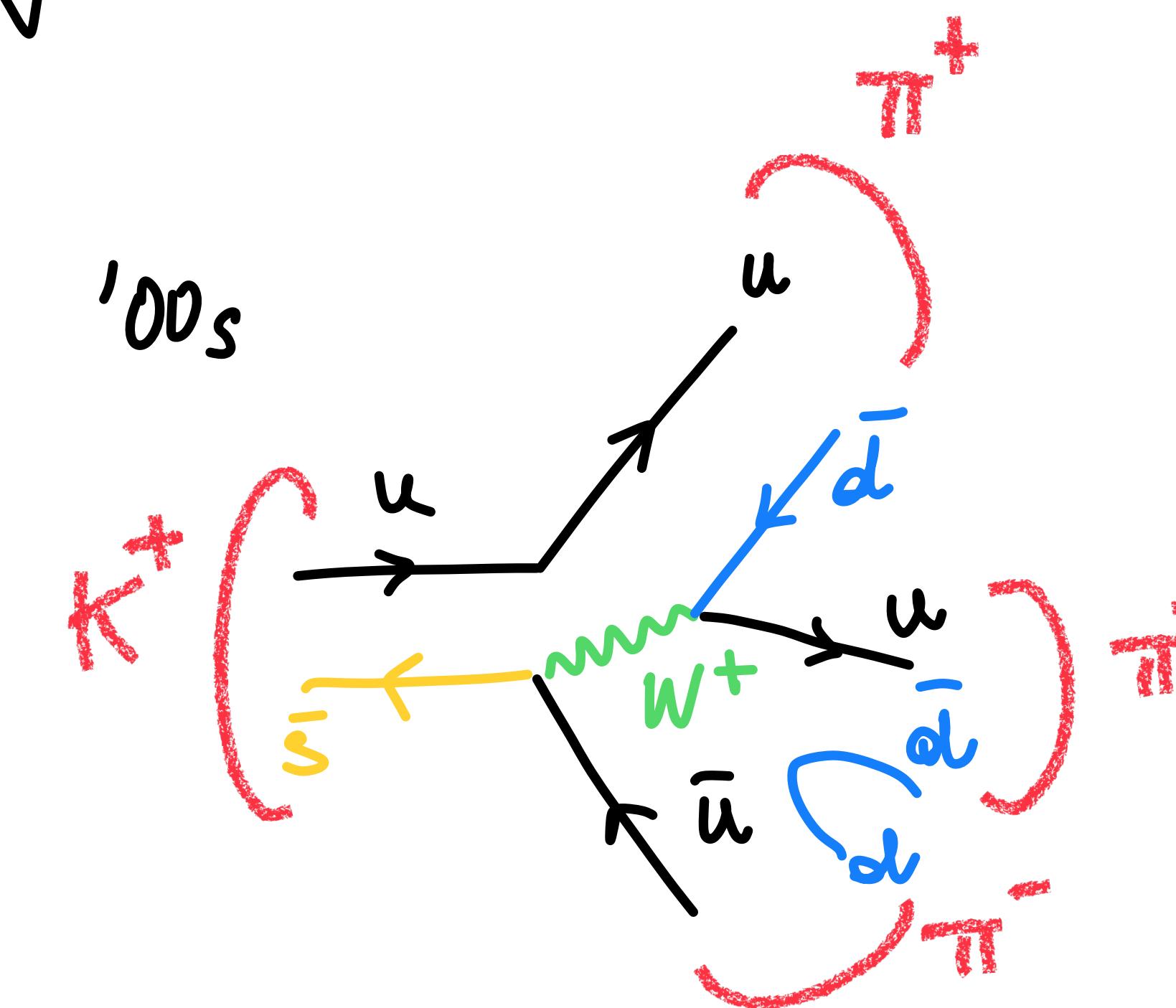
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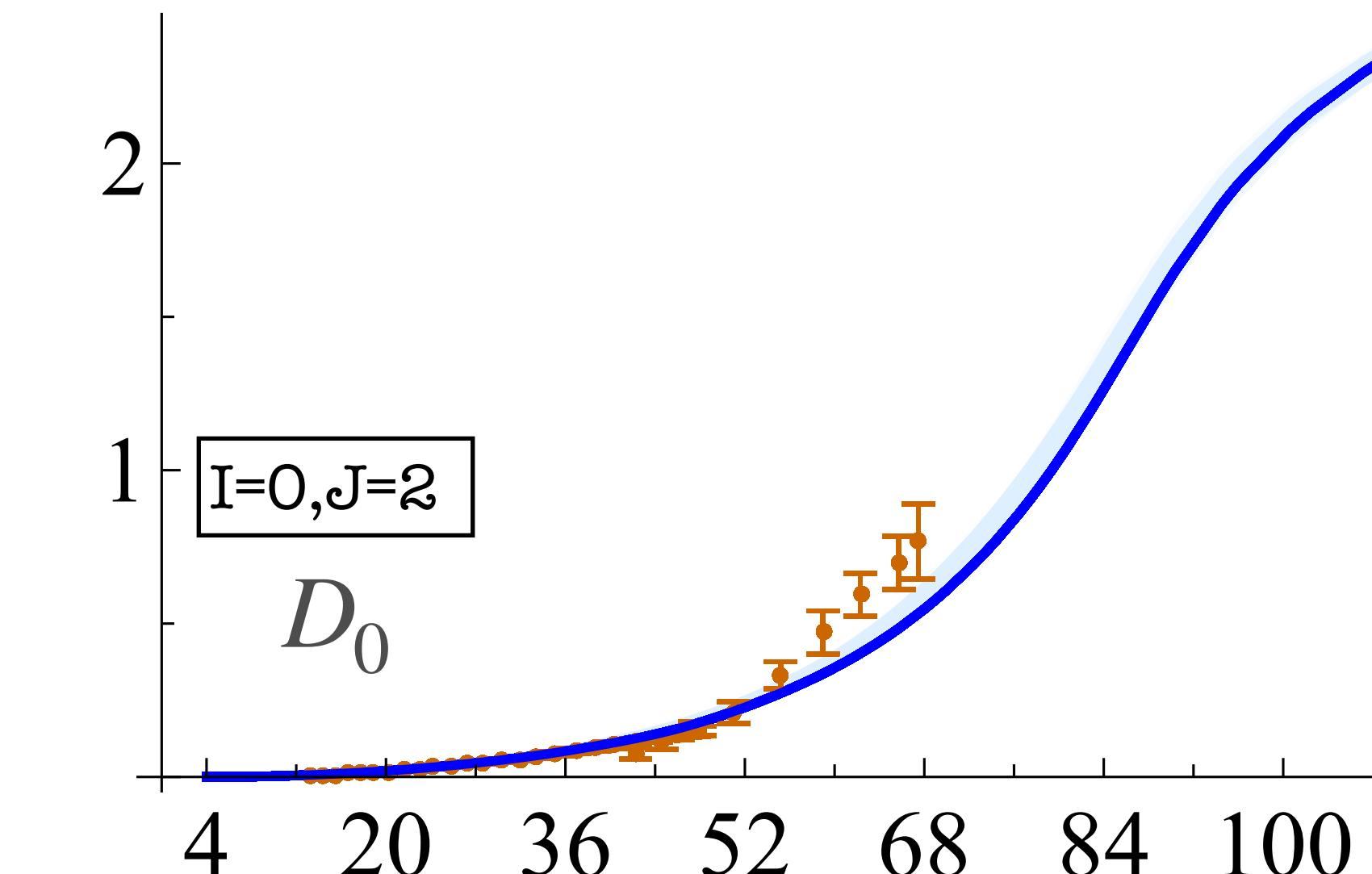
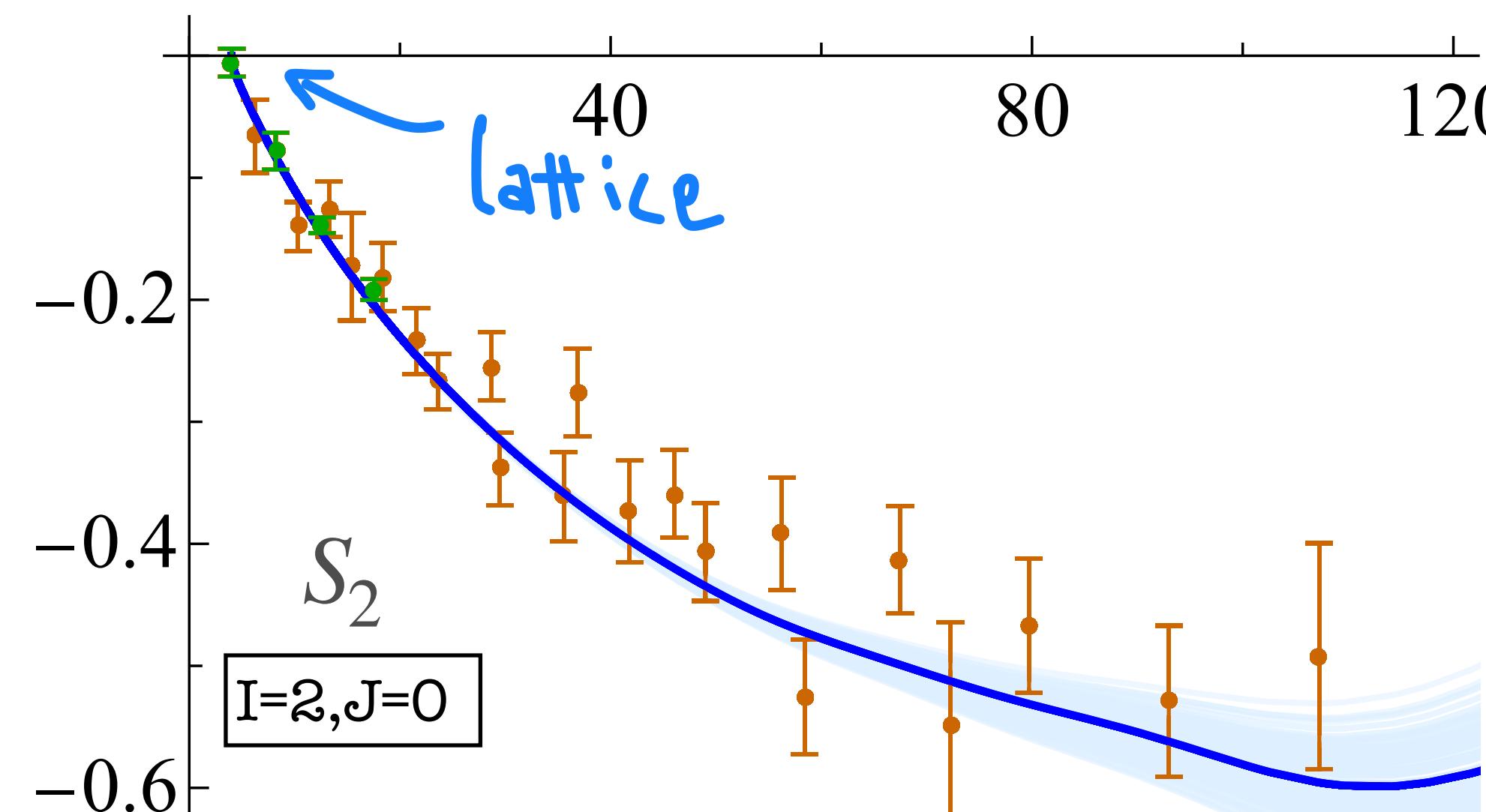
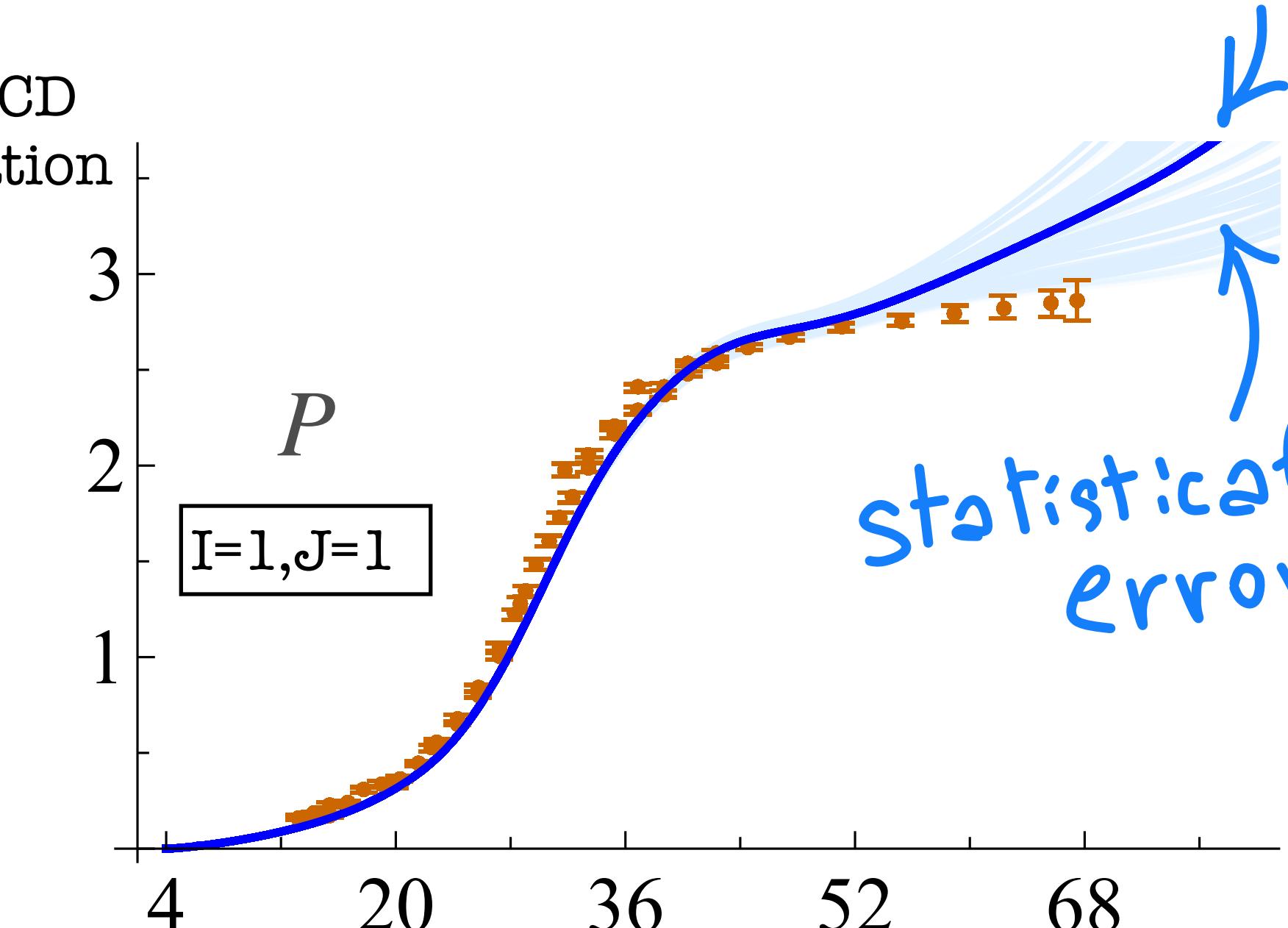
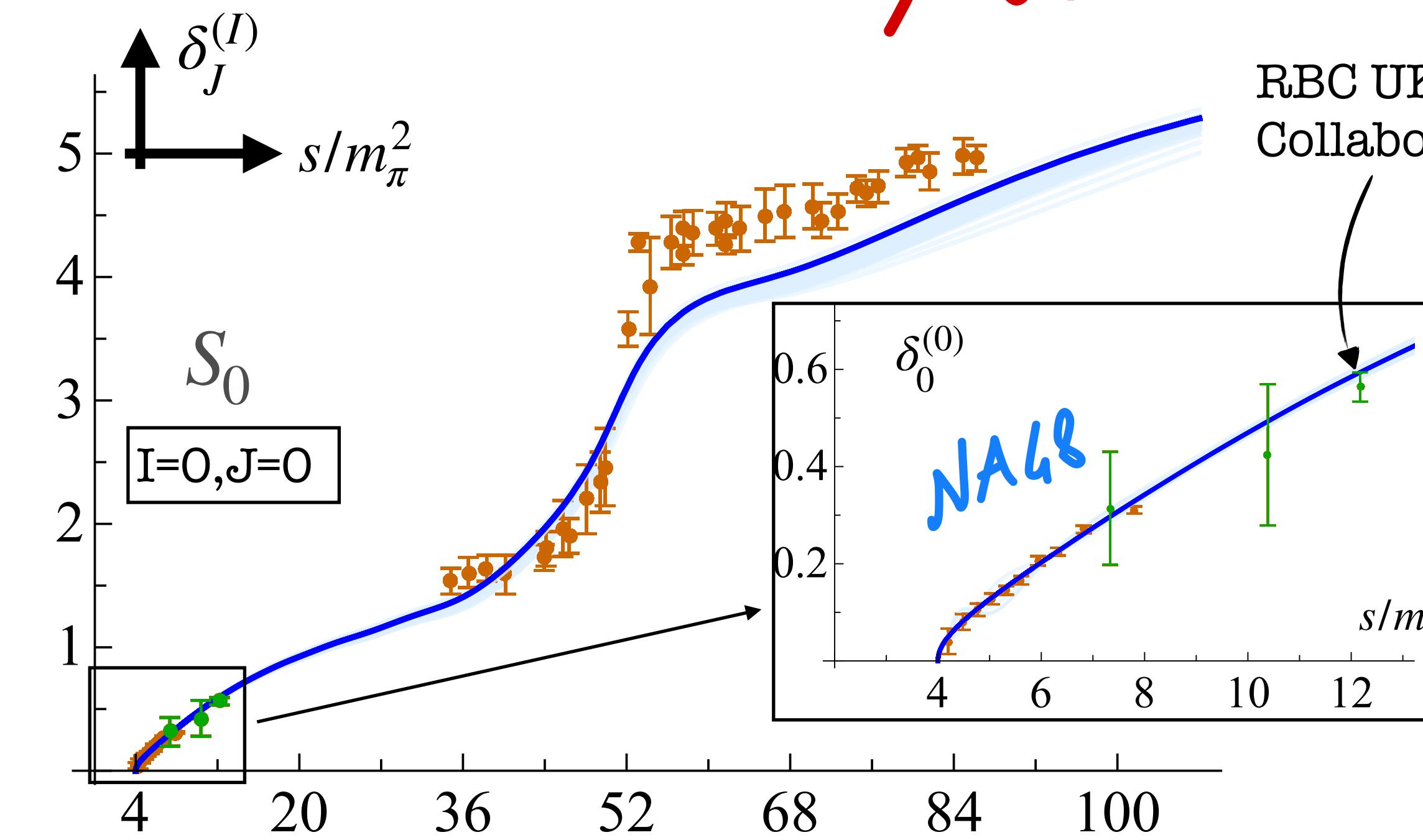
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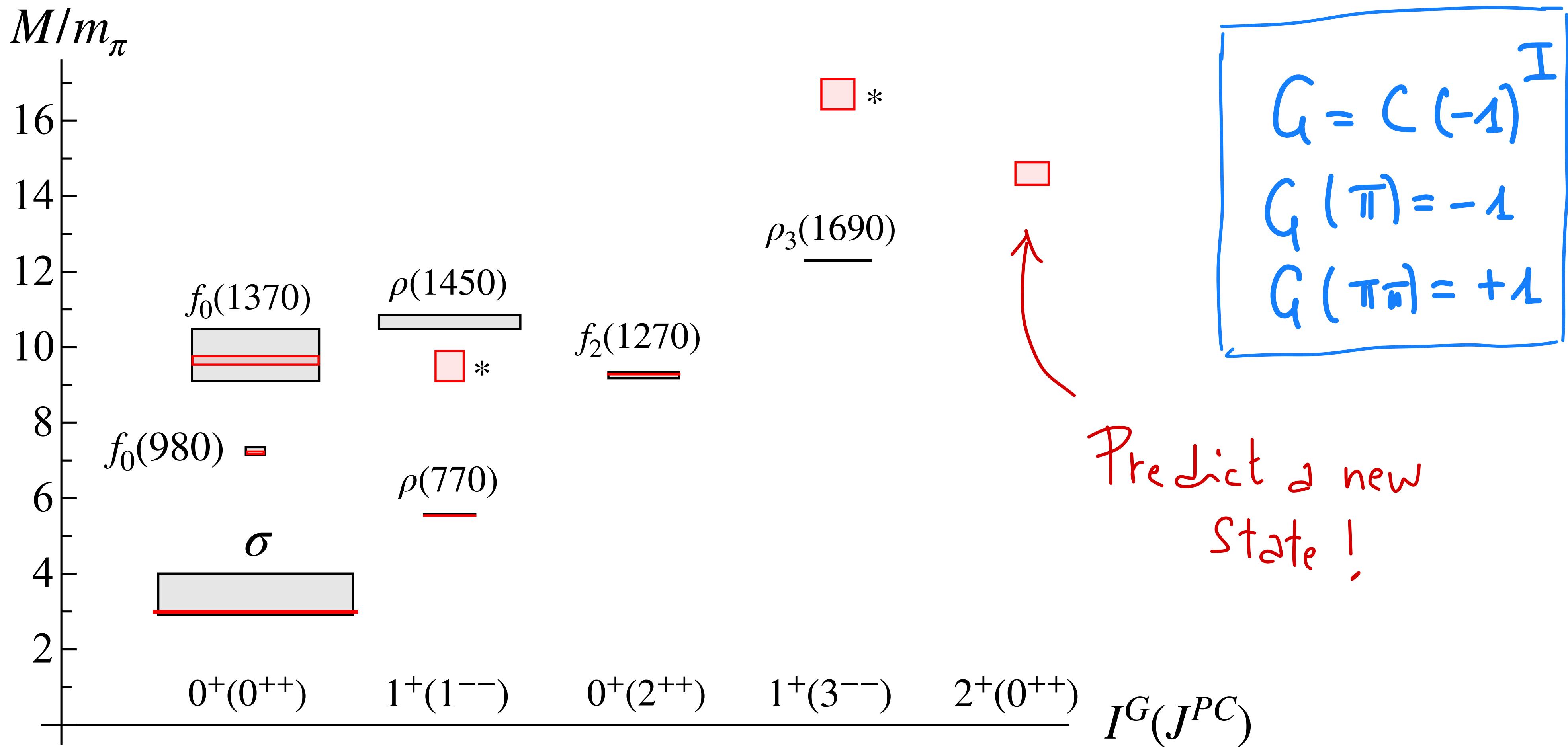
Lattice QCD (the cleanest): few data points, but important!

$$(S_J^{(I)} = \gamma_J^{(I)} e^{\zeta_J \delta_J^{(I)}})$$

Main Result



Spectrum with $G=+1$



The ρ -Expansion

$$M_{ab}^{cd}(s,t,u) =$$

$$t = (p_1 - p_3)^2$$

$$s = (p_1 + p_2)^2$$

$$u = (p_1 - p_4)^2 = 4m_\pi^2 - s - t$$

$$=$$

O(N) Bootstrap:

Cordova, Vieira, He, Kruczenski (2D) '18

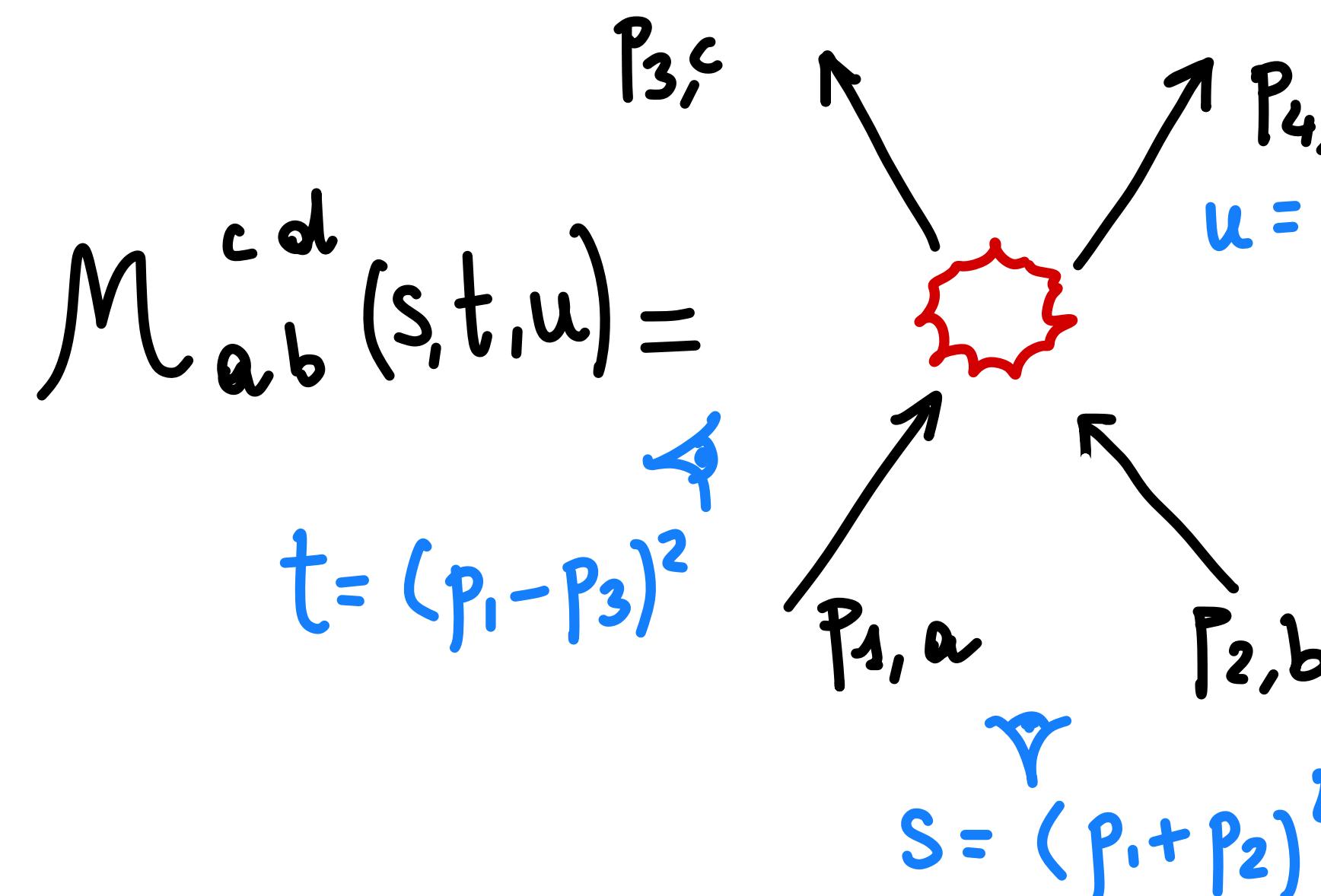
AG, Vieira, Penedones (4D + pions) '18

Sinha et al (pions) '19, '20

Elias-Mirri, AG, Guimaraes (O(4) + Higgs) '23

He, Kruczenski (pions + form factors) '23, '24

The ρ -Expansion

$M_{ab}^{cd}(s,t,u) =$


 $t = (p_1 - p_3)^2$
 $s = (p_1 + p_2)^2$
 $u = (p_1 - p_4)^2 = 4m_\pi^2 - s - t$

$$= \delta_{ab}\delta^{cd} A_s + \delta_a^c\delta_b^d A_t + \delta_a^d\delta_b^c A_u$$

where $A_s = A(s|t,u)$

Building blocks are $t-u$ symmetric

The ρ -Expansion

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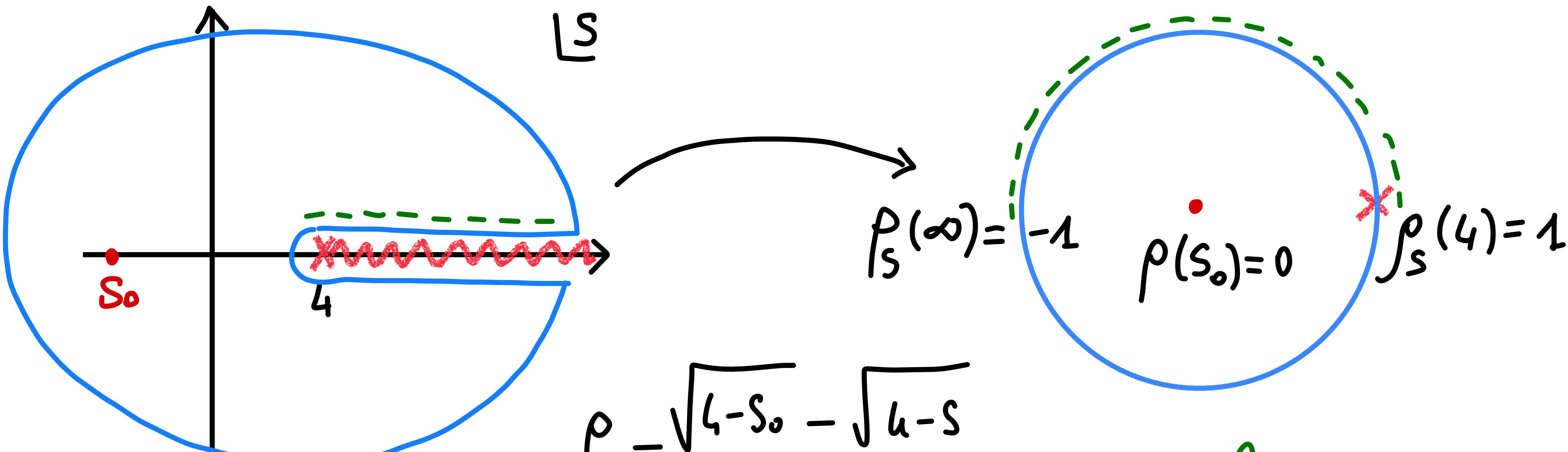
Building blocks are $t-u$ symmetric

$$A_s = \sum_{n+m \leq N} \alpha_{n,m} \rho_s^n \left(\rho_t^m + \rho_u^m \right) + \sum_{n+m \leq N} \beta_{(n,m)} \rho_t^n \rho_u^m$$

Crossing Symmetry ✓

The ρ -Expansion

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$$\rho_s = \frac{\sqrt{4-s_0} - \sqrt{4-s}}{\sqrt{4-s_0} + \sqrt{4-s}}$$

Analyticity ✓

Unitarity

1) Change of basis to flavour Irreps

$$(I=1) \times (I=1) = \begin{matrix} 0 \\ \text{Singlet} \end{matrix} + \begin{matrix} 1 \\ \text{Antisymmetric} \end{matrix} + \begin{matrix} 2 \\ \text{Symmetric} \\ \text{traceless} \end{matrix}$$

$$\left\{ A^{(0)} = N_f^{\geq 3} A_S + A_t + A_u, \quad A^{(1)} = A_t - A_u, \quad A^{(2)} = A_t + A_u \right\}$$

Unitarity

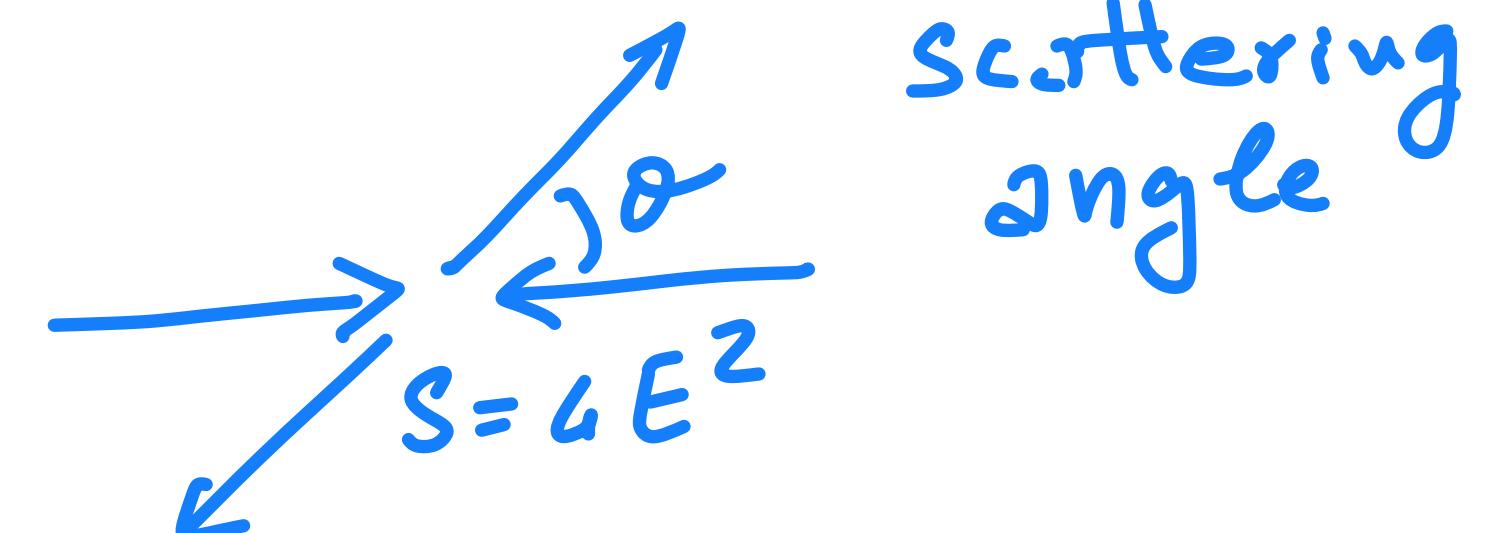
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2) Change of basis to Angular Momentum Irreps

$$f_\ell^{(I)}(s) = \frac{1}{32\pi} \int_{-1}^1 d\cos\theta P_\ell(\cos\theta) A^{(I)}(s, t = \frac{4-s}{2}(1+\cos\theta), u = \frac{4-s}{2}(1-\cos\theta))$$



Unitarity

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↑ linear functions of $\{\alpha_{n,m}; \beta_{n,m}\}$

Unitarity is diagonal

$$S_\ell^{(I)}(s) = 1 + i \sqrt{\frac{s-4}{s}} f_\ell^{(I)}(s)$$

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$$|S_\ell(s)|^2 \leq 1 \iff \ell = 0, 1, 2, \dots, \infty; \quad s \geq 4$$

$$\begin{pmatrix} 1 - \text{Im } S_\ell^{(I)} & \text{Re } S_\ell^{(I)} \\ \text{Re } S_\ell^{(I)} & 1 + \text{Im } S_\ell^{(I)} \end{pmatrix} \succcurlyeq 0$$

SDP

Pion Soft theorem

If $m_u = m_d = 0$, then $m_\pi^2 = 0$ and pions are Goldstones

Soft-theorem $M_{ab}^{cd}(s, t, u) \xrightarrow{P_i \rightarrow 0} 0$

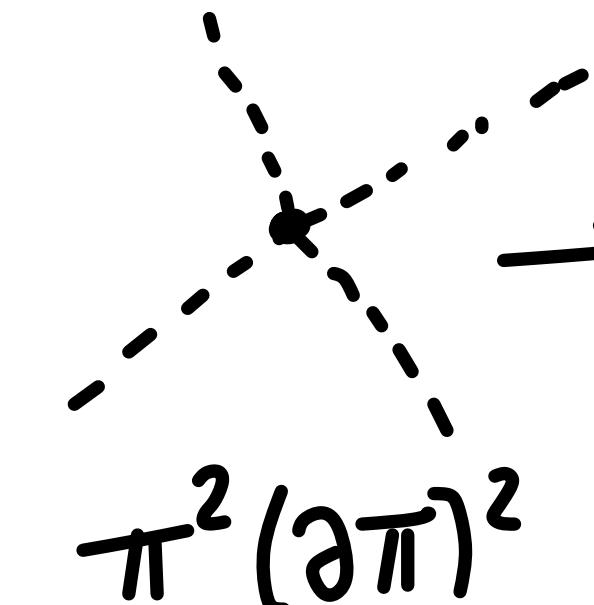
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Since $m_\pi^2 \neq 0$, we use χ -EFT (Weinberg)

$$\mathcal{L} = \frac{1}{2} (\partial_\mu \vec{\pi})^2 - \frac{1}{2} m_\pi^2 \vec{\pi}^2 + \dots$$



$$A_S = \frac{s - m_\pi^2}{f_\pi^2}$$

↑ $s=t=u=m_\pi^2$
Adler zero

↑ f_π
pion decay constant

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$$f_0^{(0)} = \frac{2S - m_\pi^2}{32\pi f_\pi^2} \rightarrow f_0^{(0)}\left(\frac{m_\pi^2}{2}\right) = 0$$

$$f_1^{(1)} = \frac{S - 4m_\pi^2}{96\pi f_\pi^2} \rightarrow f_1^{(1)}(4m_\pi^2) = 0$$

$$f_0^{(2)} = \frac{2m_\pi^2 - S}{16\pi f_\pi^2} \rightarrow f_0^{(2)}(2m_\pi^2) = 0$$

Non-perturbatively

$$f_0^{(0)}(z_0) = 0, \quad f_0^{(2)}(z_0) = 0$$

$$0 \leq z_0, z_2 \leq 4M_\pi^2$$

pion decay
constant

Spectrum Assumptions

QCD contains many resonances that couple to $\pi\pi$

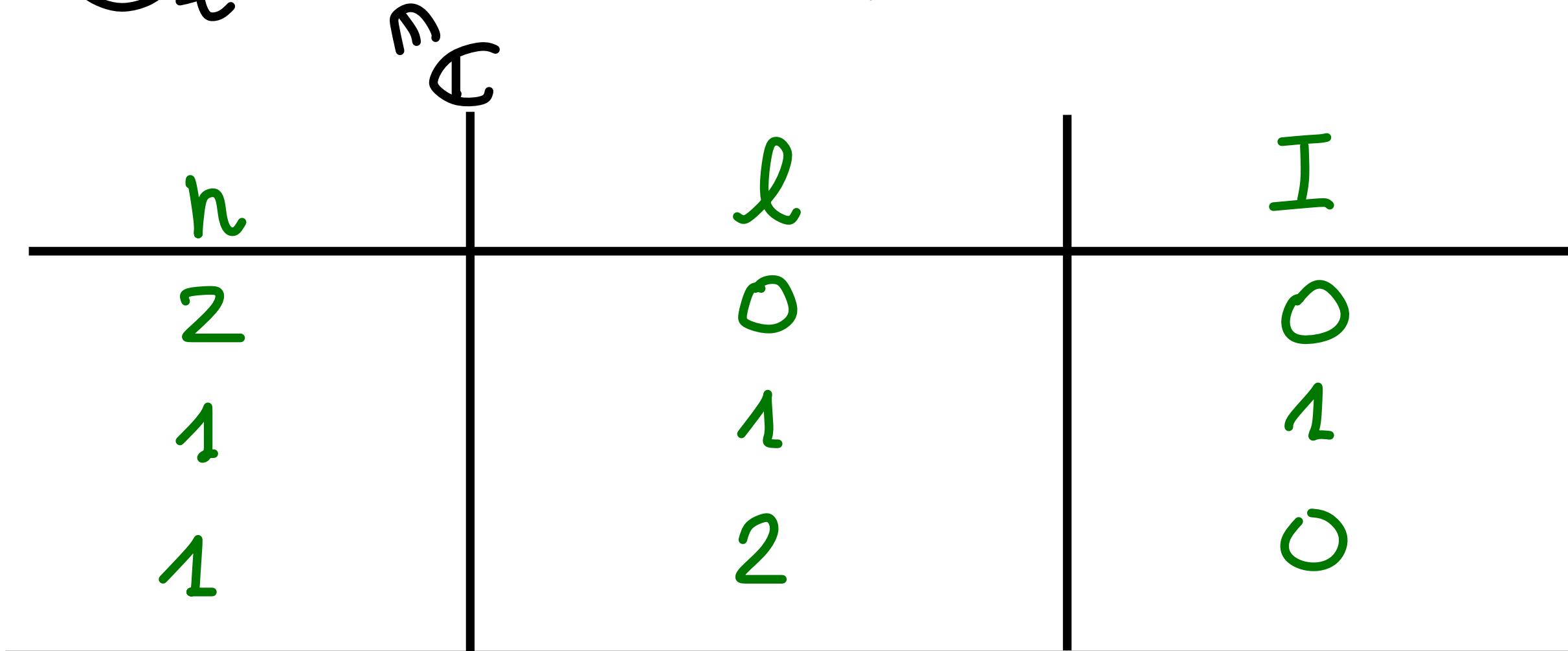
(Not necessary) We can help the model imposing the existence
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$$S_l^{(I)}(m_R^2) = 0 \quad (\text{pole in the II sheet})$$



$$m_R^2 = 4 \times 2$$

2 additional parameters

Summary

Crossing +
Analyticity

$$A_s = \sum_{n+m \leq N} a_{n,m} \rho_s^n (\rho_t^m + \rho_u^m) + \sum_{n+m \leq N} \beta_{(n,m)} \rho_t^n \rho_u^m$$

Unitarity

$$|S_\ell^{(I)}(s)|^2 \leq 1$$

$$\ell = 0, 1, 2, \dots, \cancel{\infty}$$

$N \rightarrow \infty$ any
analytic function

$$s \geq 4 \quad \text{grid} \rightarrow \infty \quad L \text{max} \rightarrow \infty$$

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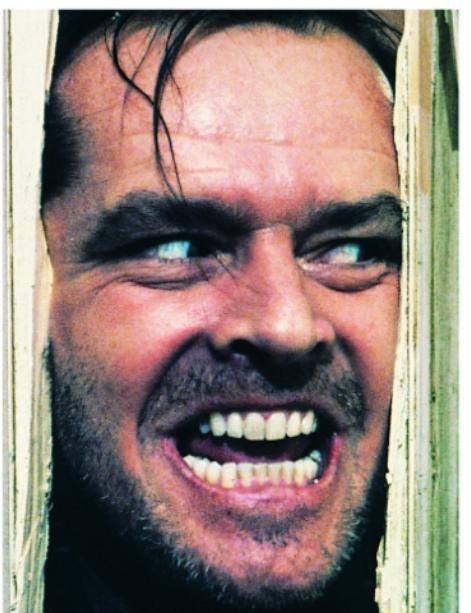
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$$\lim_{\text{grid} \rightarrow \infty} \quad \lim_{L_{\max} \rightarrow \infty} \quad \lim_{N \rightarrow \infty}$$



Triple extrapolation

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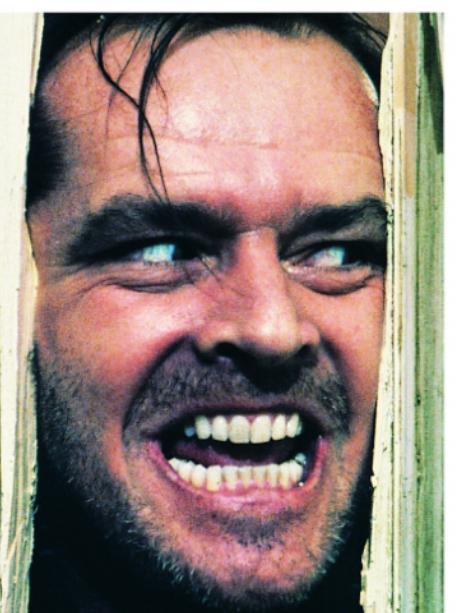
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Triple extrapolation Chebyshev grid

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Triple extrapolation Chebyshev grid

$$\operatorname{Im} A^{(I)}(s, 0 \leq t < h) \geq 0$$

positive sum rules for $\ell > L_{\max}$

$$\lim_{\text{grid} \rightarrow \infty} \quad \lim_{L_{\max} \rightarrow \infty} \quad \lim_{N \rightarrow \infty}$$

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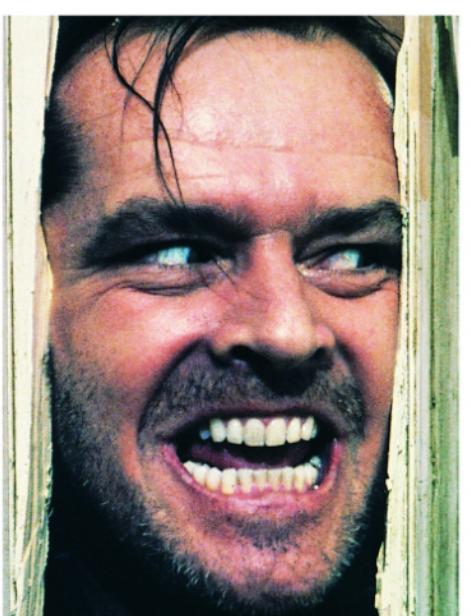
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Triple extrapolation Chebyshev grid

$$\operatorname{Im} A^{(I)}(s, 0 \leq t < 4) \geq 0 \quad \begin{matrix} \text{positive sum} \\ \text{rules for } \ell > L_{\max} \end{matrix}$$

$$\lim_{\text{grid} \rightarrow \infty}$$

$$\lim_{L_{\max} \rightarrow \infty}$$

$$\lim_{N \rightarrow \infty}$$

$$\left(\rho_s = \frac{\sqrt{4-s_0} - \sqrt{4-s}}{\sqrt{4-s_0} + \sqrt{4-s}} \right)$$

moving s_0
(Wavelets, multi
foliation)

The Model (I)

β -expansion with parameters $\{\alpha_{n,m}; \beta_{n,m}\}$ - 8 of them

Chiral zeroes

$$f_0^{(0)}(z_0) = 0, f_0^{(2)}(z_2) = 0$$

Spectrum Assumptions

$$S_e^{(I)}(m_R^2) = 0$$

In principle,
we want to tune
 α 's, β 's, z_0 , z_2 , m_R^2 to
fit the data

The Model (I)

ρ -expansion with parameters $\{\alpha_{n,m}; \beta_{n,m}\}$ - 8 of them

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Spectrum Assumptions

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ρ -expansion typically $\mathcal{O}(1000)$ parameters \rightarrow overfitting

Procedure to choose few of them and fix the rest

EFT

Bootstrap

The Model (II)

$\{\alpha_{n,m}; \beta_{(n,m)}\}$ have no direct physical interpretation

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$\{ \alpha_{n,m}; \beta_{(n,m)} \}$ have no direct physical interpretation

Better consider low-energy constants

Universal expansion

$$\operatorname{Re} f_e^{(I)}(s) = (s - 4)^l \left[a_e^{(I)} + b_e^{(I)} (s - 4) + \dots \right]$$

↑ Scattering lengths ↑ effective range

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↑ ↑
scattering lengths effective range

Choose e.g. $\{a_0^{(0)}, a_1^{(1)}, a_0^{(2)}, \dots\}$, and to fix the rest ...

"Go to the boundary of the allowed region"

Why the boundary

β -expansion
 $+ z_0, z_2, m_R^2$
 $+ \text{unitarity}$

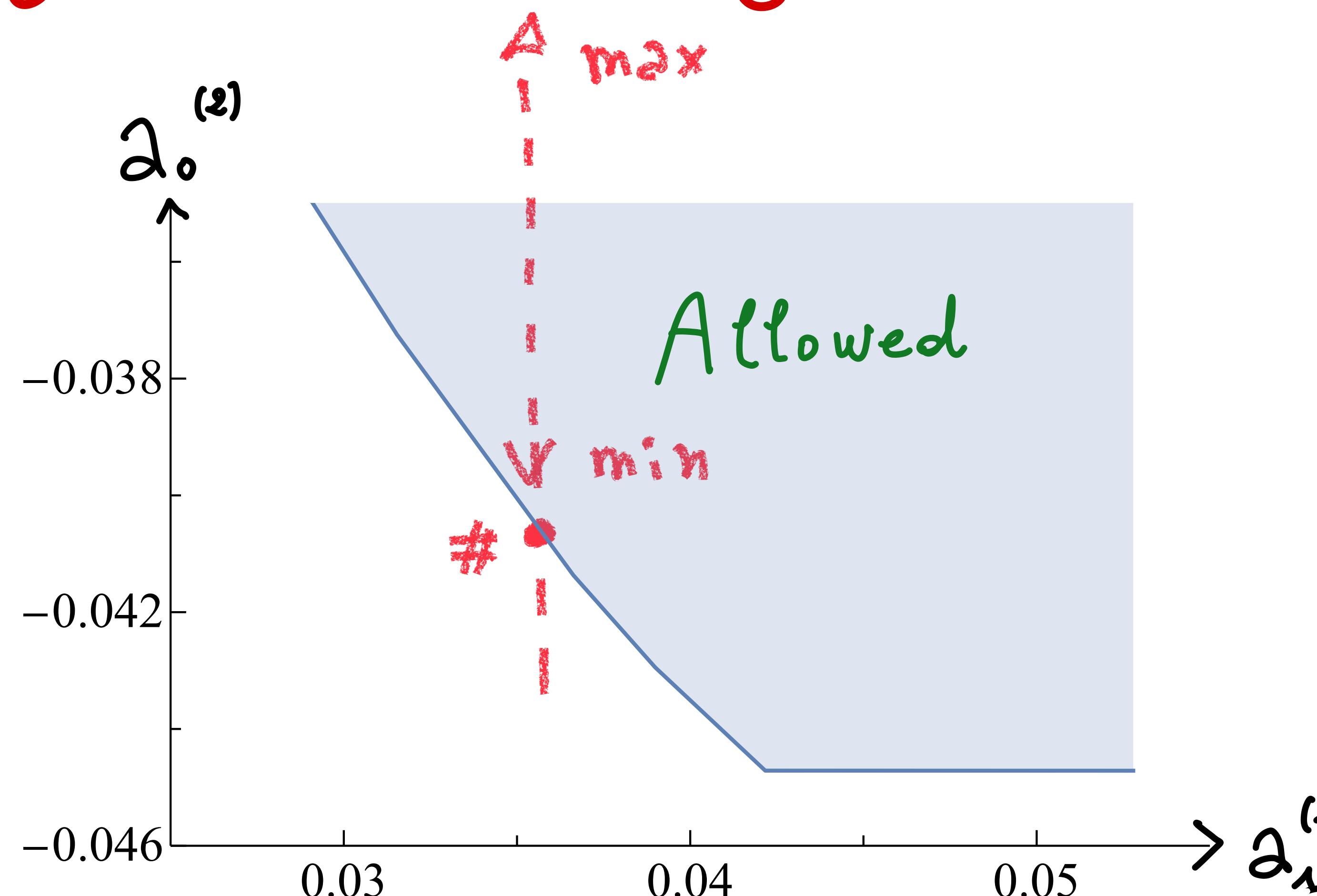
Fix $z_0^{(0)}, z_1^{(1)}$

$$\max z_0^{(2)} = \infty$$

$$\min z_0^{(2)} = \#$$

At the max/min
 all remaining

$\{\alpha_{n,m}; \beta_{n,m}\}$ are fixed ← The amplitude @ the boundary is unique!

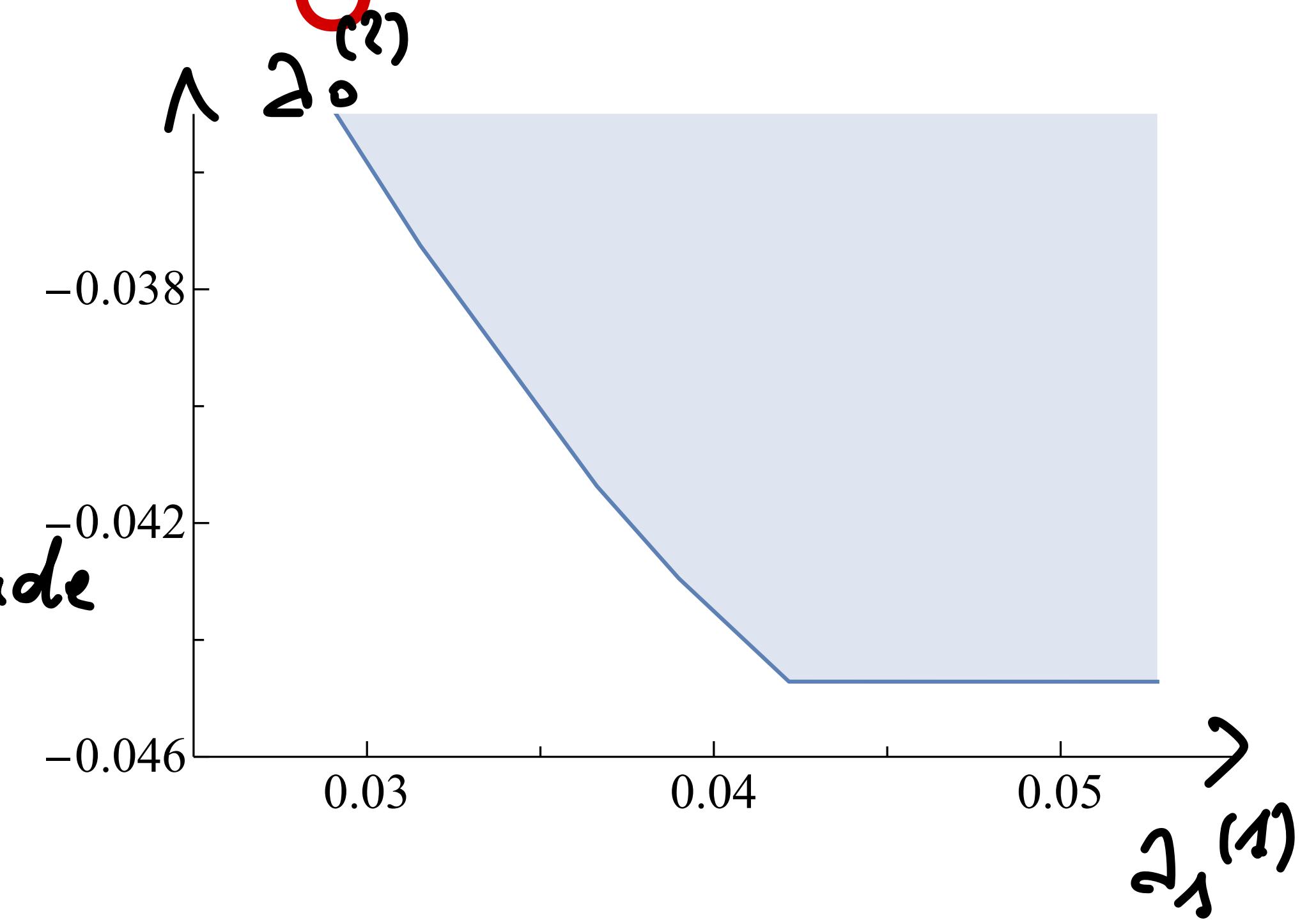


Why the boundary

ρ -expansion
+ z_0, z_2, m_R^2
+ unitarity

} \nwarrow boundary defined by
the constraints

If $|S_\ell^{(I)}(s)|^2 < 1 \rightarrow$ boundary amplitude
is elastic

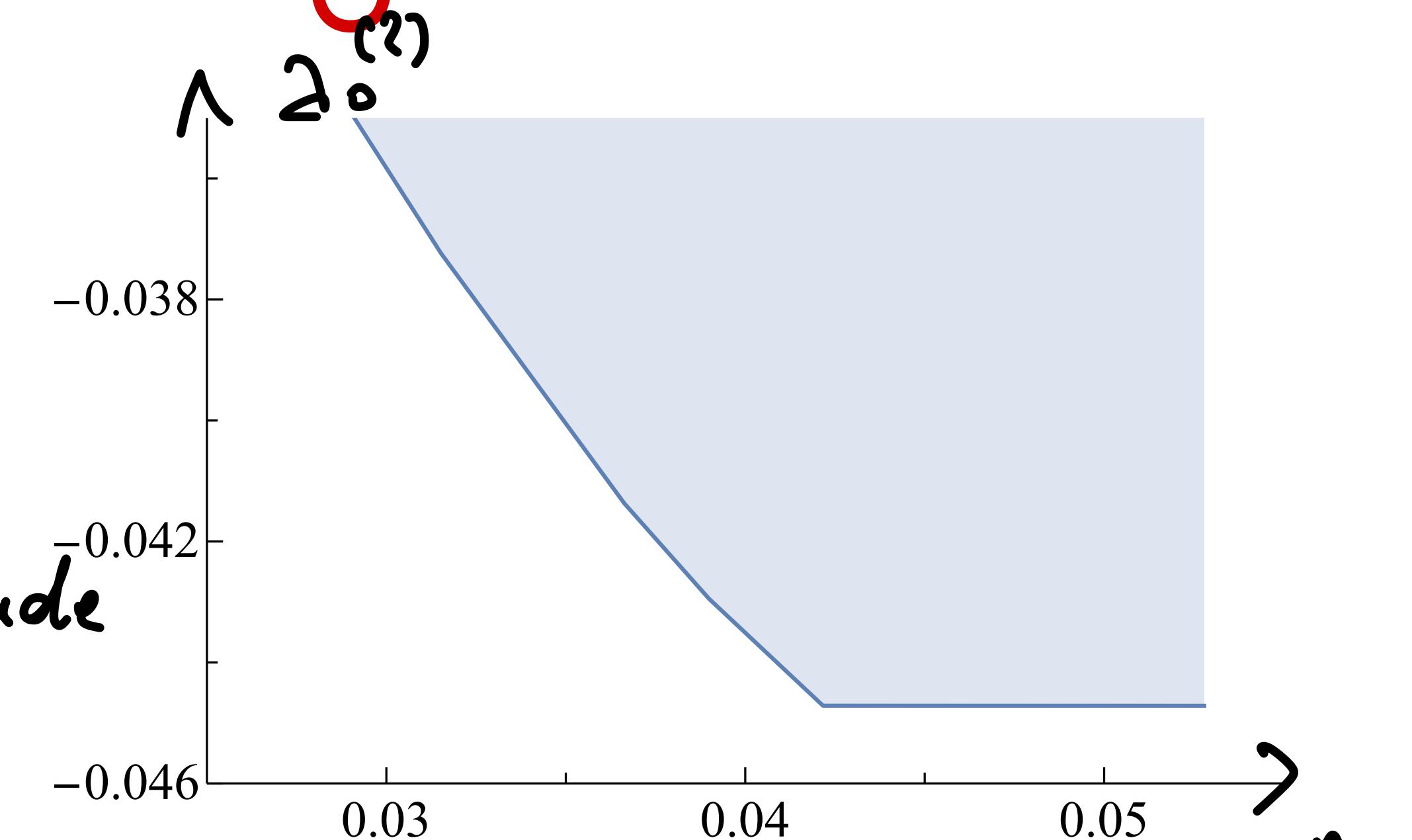


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- We impose the experimental inelasticity
 \hookrightarrow bndry amplitude has correct production

$$|S_\ell^{(I)}(s)|^2 \leq (\gamma_e^{(I)}(s))^2$$

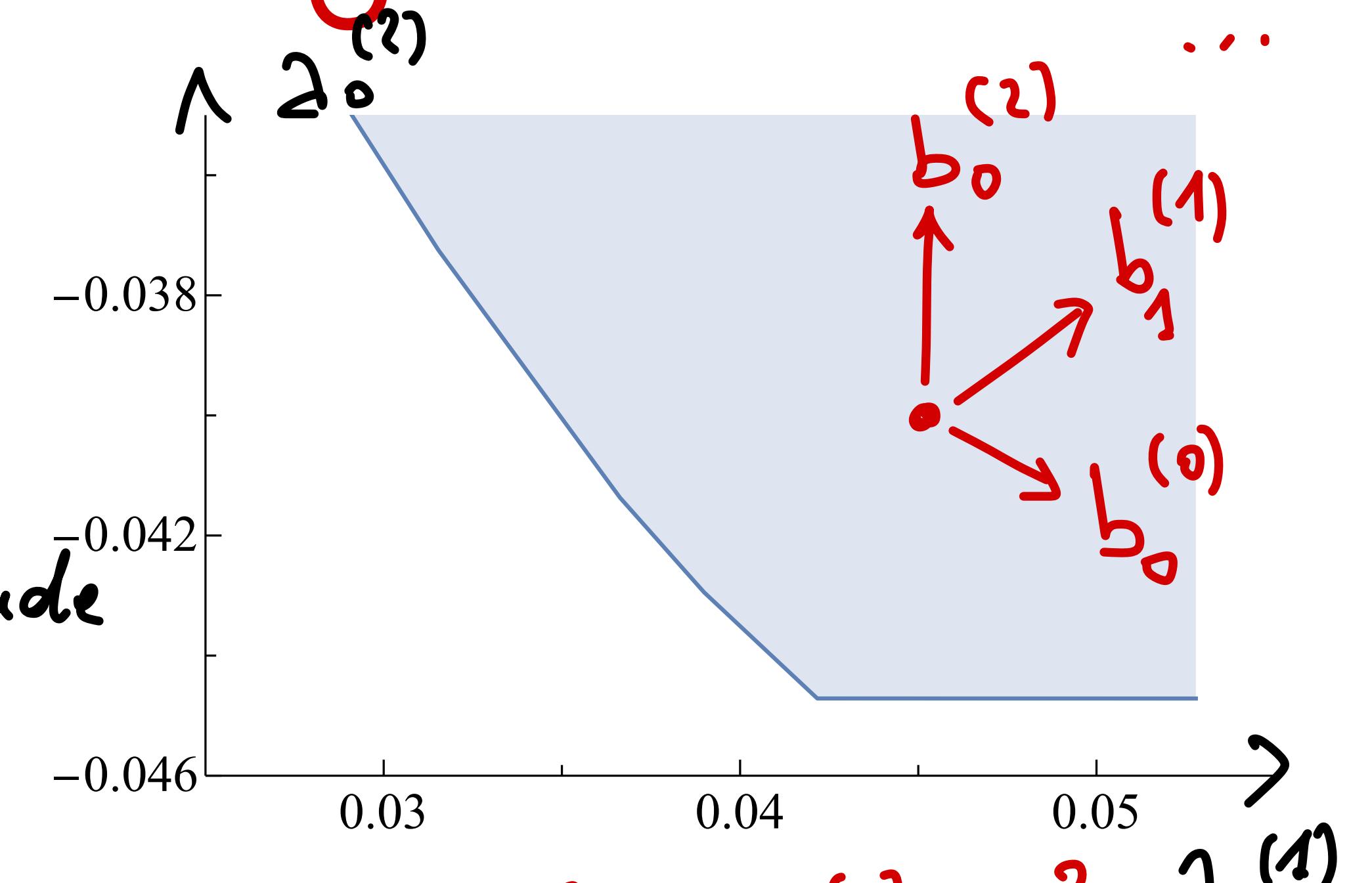
↑ taken from
Pelizz et al

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If $|S_\ell^{(I)}(s)|^2 < 1 \rightarrow$ boundary amplitude
is elastic



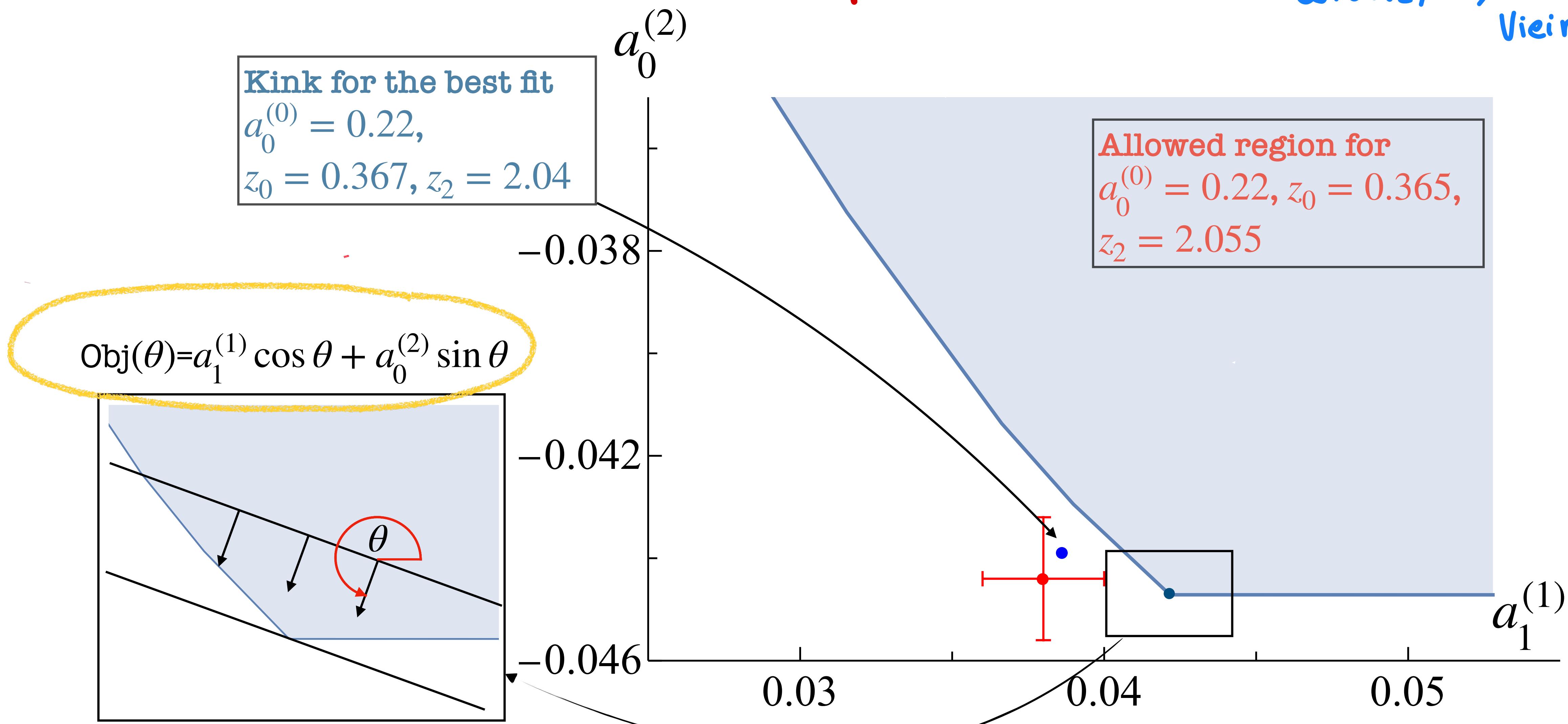
- We impose the experimental inelasticity
 \hookrightarrow bndry amplitude has correct production

- We tune more coefficients, e.g. fix $\{z_1^{(1)}, z_0^{(2)}\}$

$$|S_\ell^{(I)}(s)|^2 \leq (\gamma_e^{(I)}(s))^2$$

Normal Optimization

← Cordova, He, Kruczak, Vieira



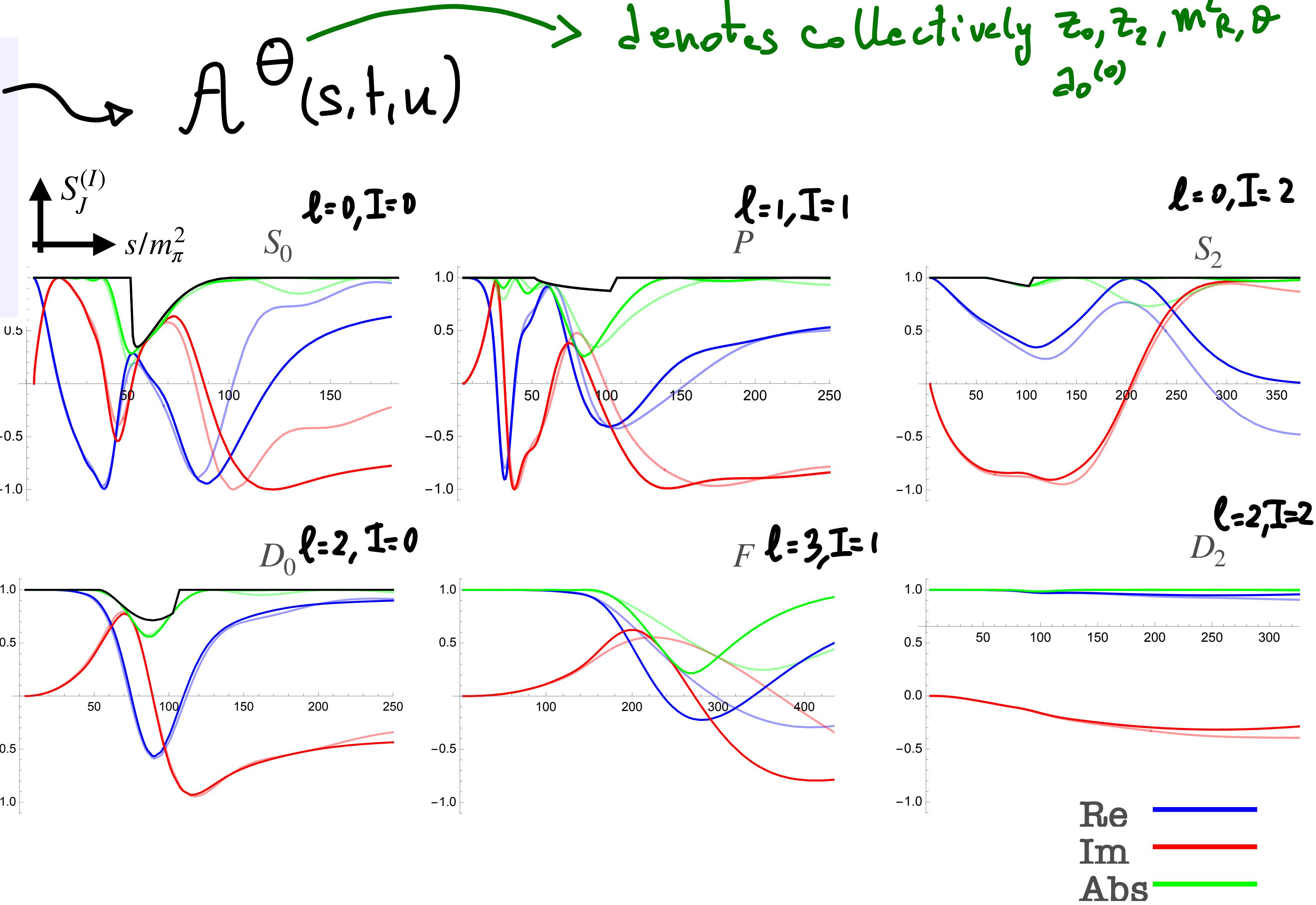
Bootstrap Solution

Given Fit Ansatz
 $\Theta = \{\theta, a_0^{(0)}, z_0, z_2, m_\rho^2, m_{f_0}^2, m_{f'_0}^2, m_{f_2}^2\}$

Maximize Obj(θ)
 in $\mathcal{A}^{\text{ansatz}}(s, t, u)$

constr. by $t_0^{(0)}(4) = 2a_0^{(0)}$, $t_0^{(0)}(z_0) = 0$, $t_0^{(2)}(z_2) = 0$
 $S_1^{(1)}(m_\rho^2) = 0$, $S_0^{(0)}(m_{f_0}^2) = 0$
 $S_0^{(0)}(m_{f'_0}^2) = 0$, $S_2^{(0)}(m_{f_2}^2) = 0$

$s \geq 4$ $\mathcal{U}_\ell^{(I)} \geq 0$ for $\ell \in \mathbb{N}, I = 0, 1, 2$ (9)



A non-linear / non-convex problem

To find the best fit we must tune Θ to minimize

$$\chi^2_{\Theta} = \sum_{i=1}^{\text{# data points}} \left(\left(S_{\ell_i}^{I_i}(s_i) \right)^{\text{exp}} - \left(S_{\ell_i}^{I_i}(s_i) \right)^{\text{Bootstrap}} \right)^2 / \left(S_{\ell_i}^{I_i}(s_i) \right)^2$$

$$S_{\ell}^{(I)}(s) = \gamma_{\ell}^{(I)}(s) e^{2i \delta_{\ell}^{(I)}(s)}$$

phase shift

Bottleneck:

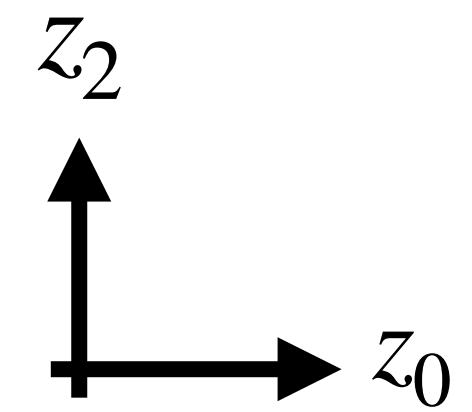
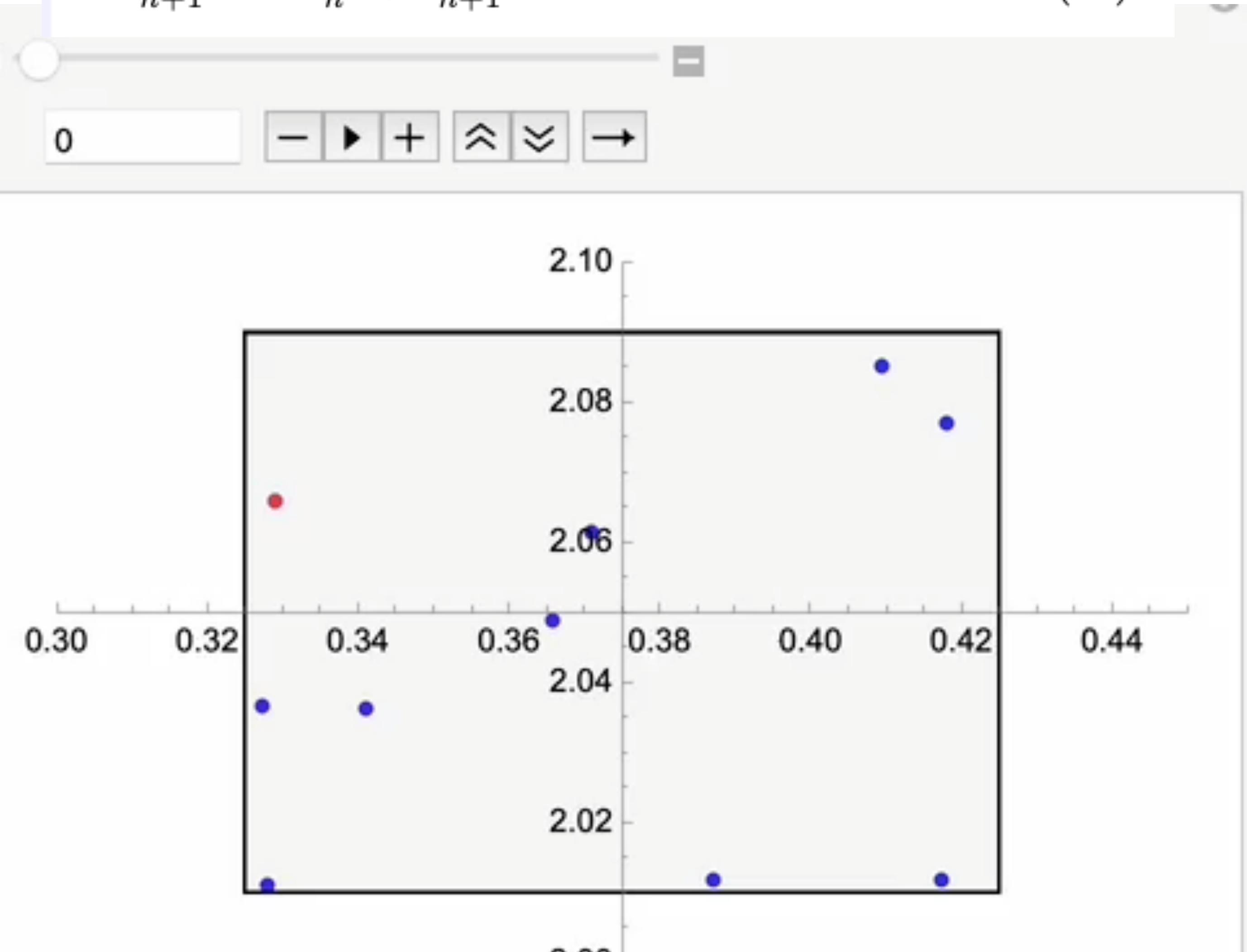
Evaluate $\chi^2_{\Theta} \leftrightarrow$ Solve a full Bootstrap problem ($\sim 2/3 h$)

$\partial_{\Theta} \chi^2_{\Theta}$ even longer \leftarrow Navigator
Skydiving Algorithm

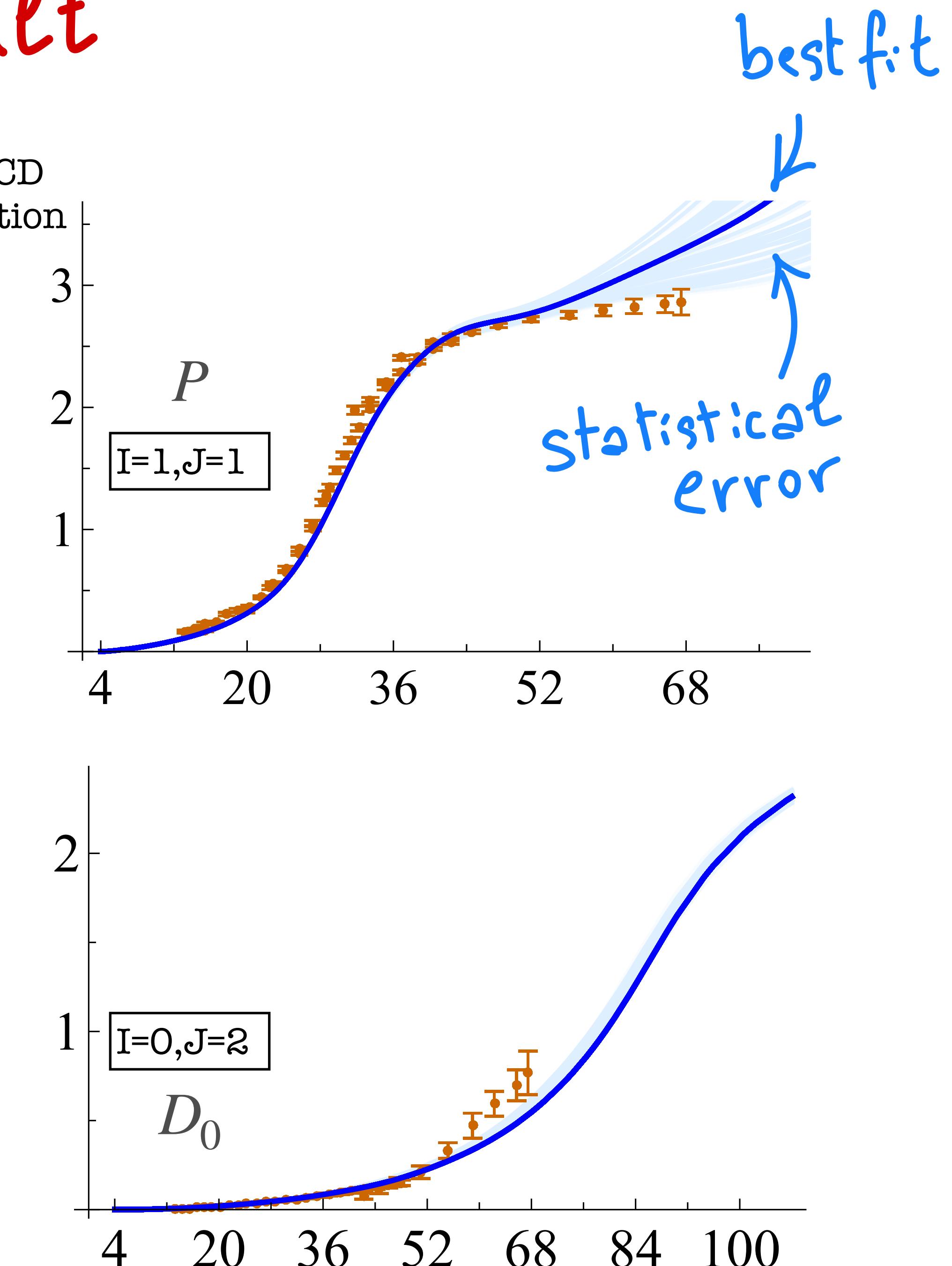
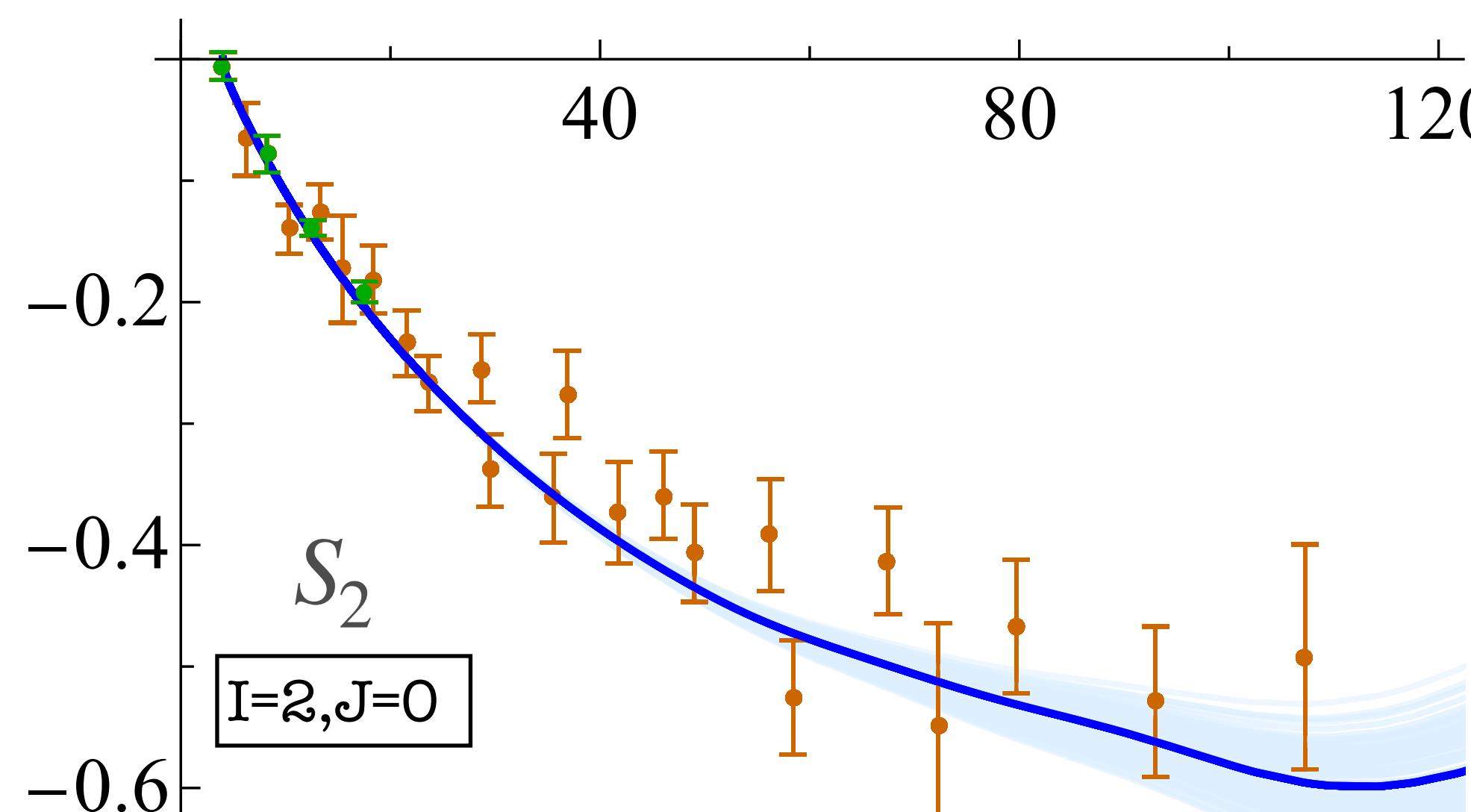
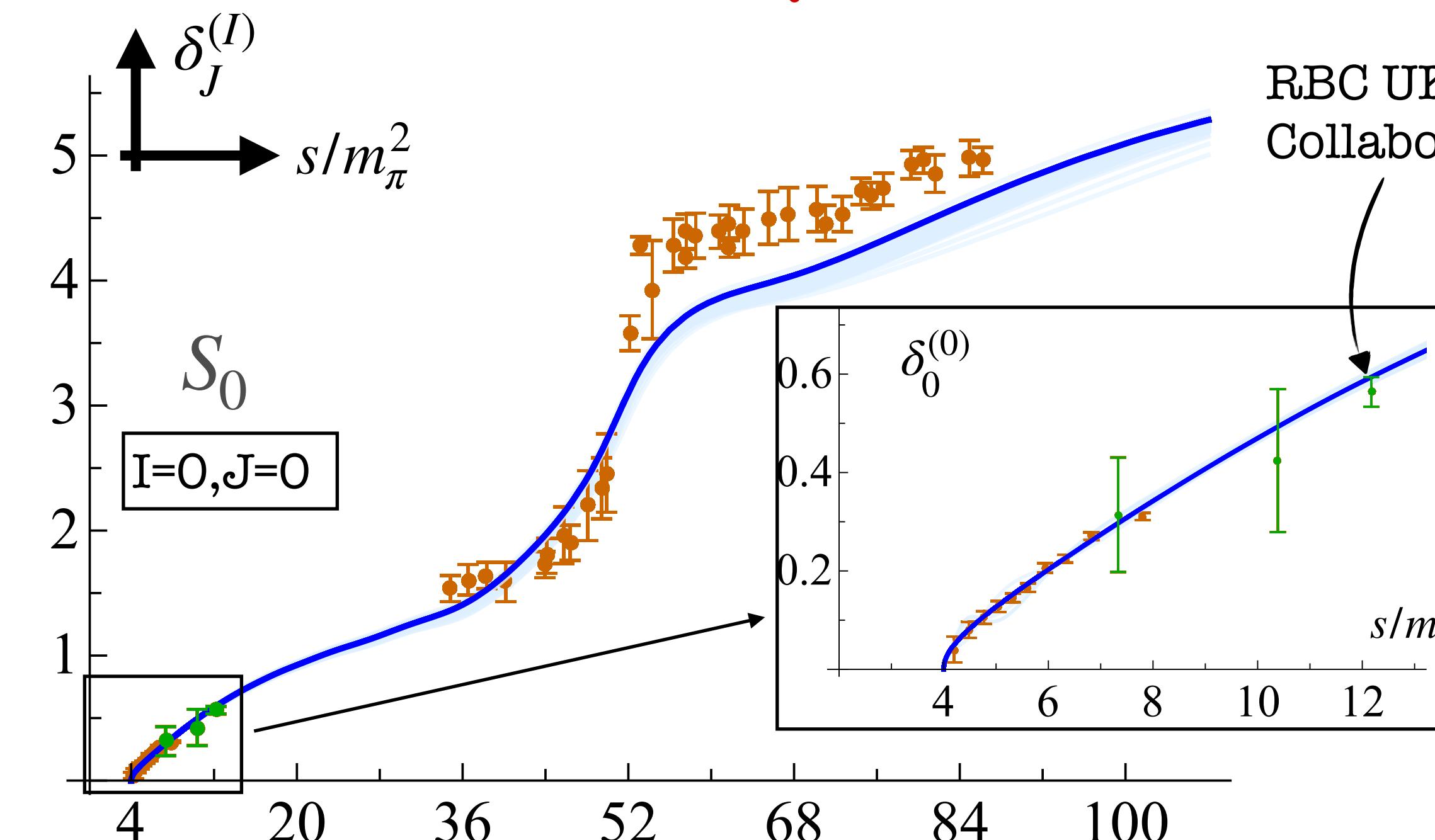
Particle Swarm Optimization

Basic Model
(swarm of birds dynamics)

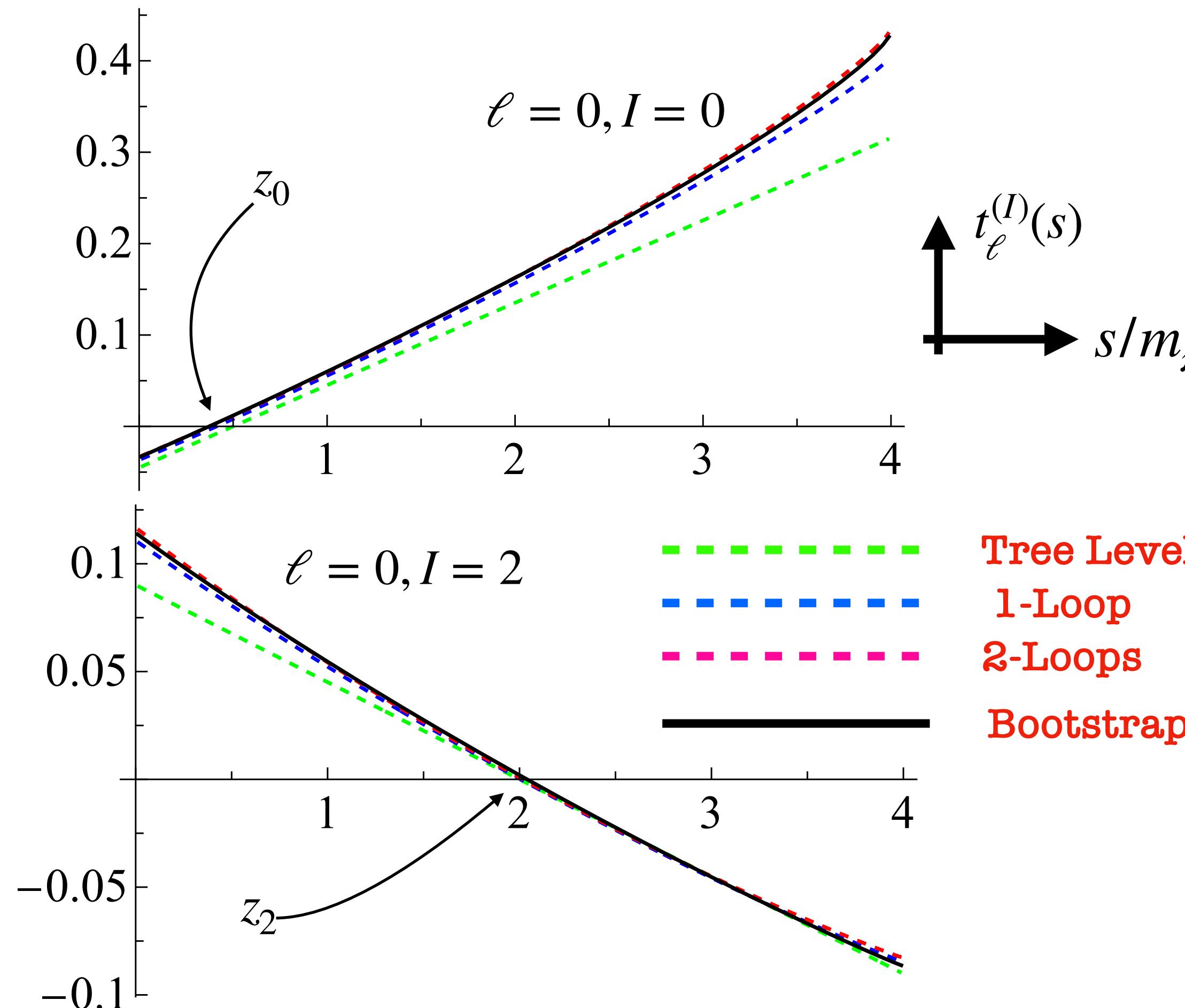
$$\begin{aligned} v_{n+1}^{(i)} &= \omega v_n^{(i)} + c_1 r_1 (\Theta_n^{(i)} - X_n^{(i)}) + c_2 r_2 (\Theta_n^{(i)} - Y_n), \\ \Theta_{n+1}^{(i)} &= \Theta_n^{(i)} + v_{n+1}^{(i)}. \end{aligned} \quad (11)$$



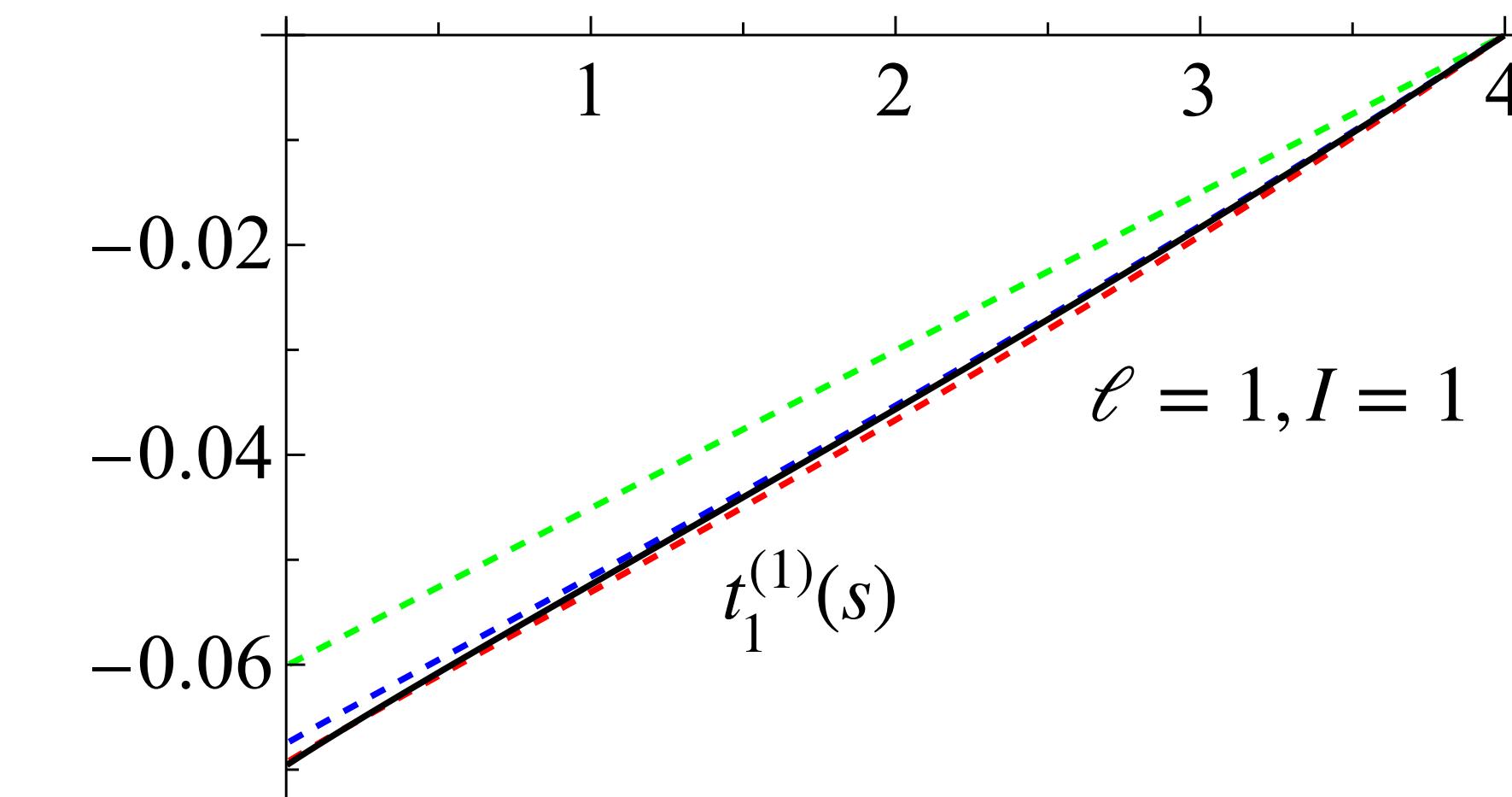
Main Result



Checks @ low energy



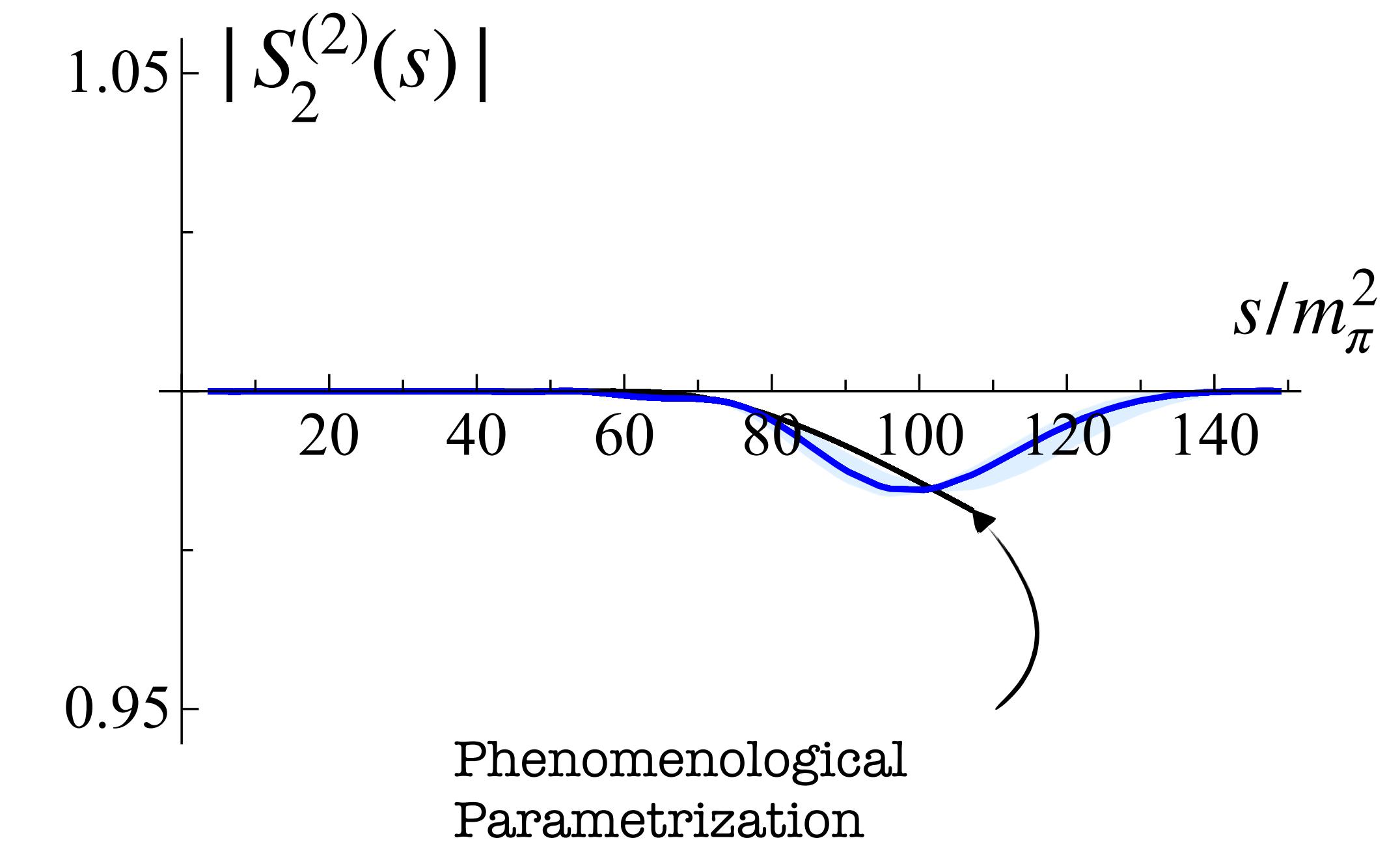
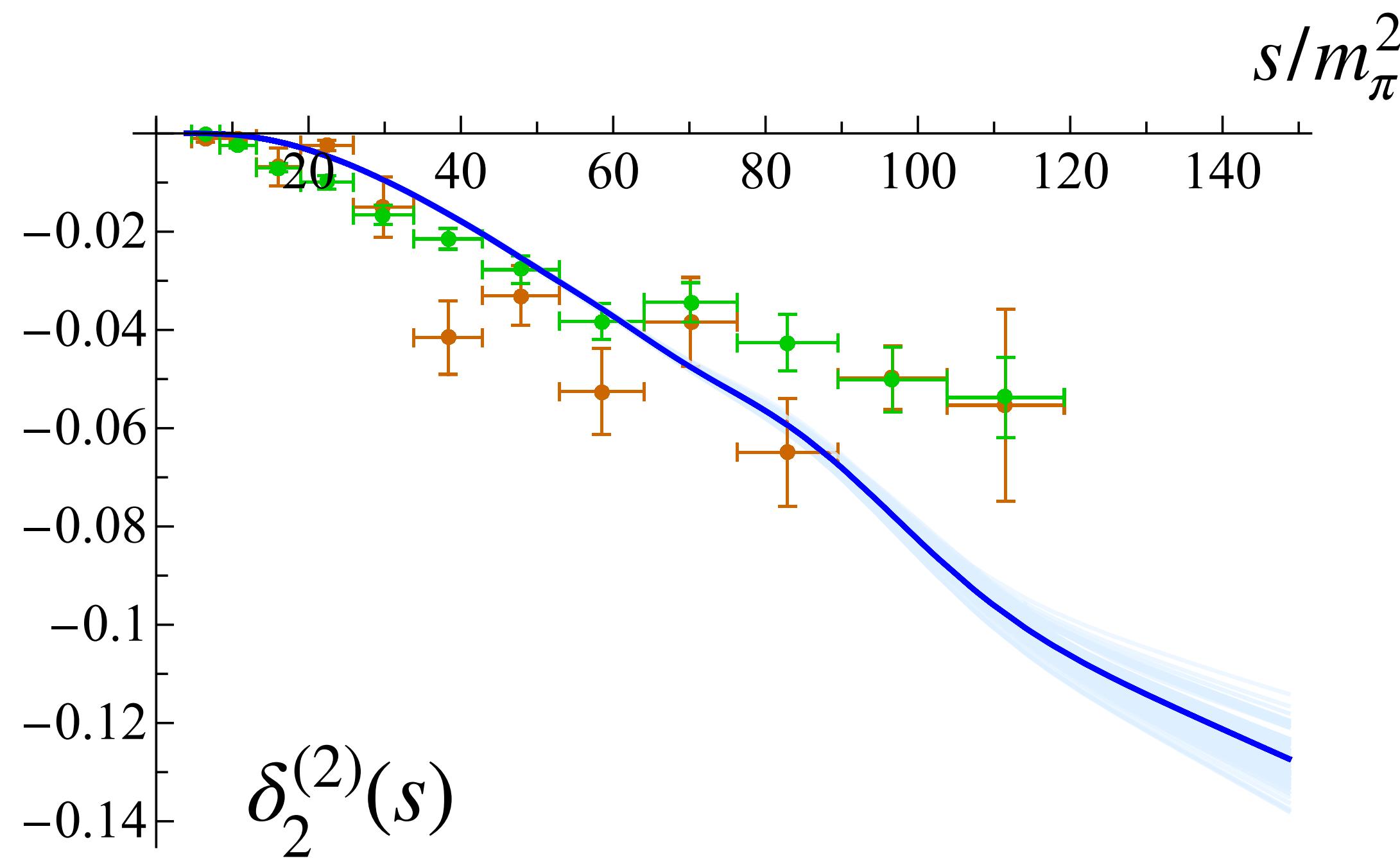
2-loops from
Colangelo, Gasser, Leutwyler



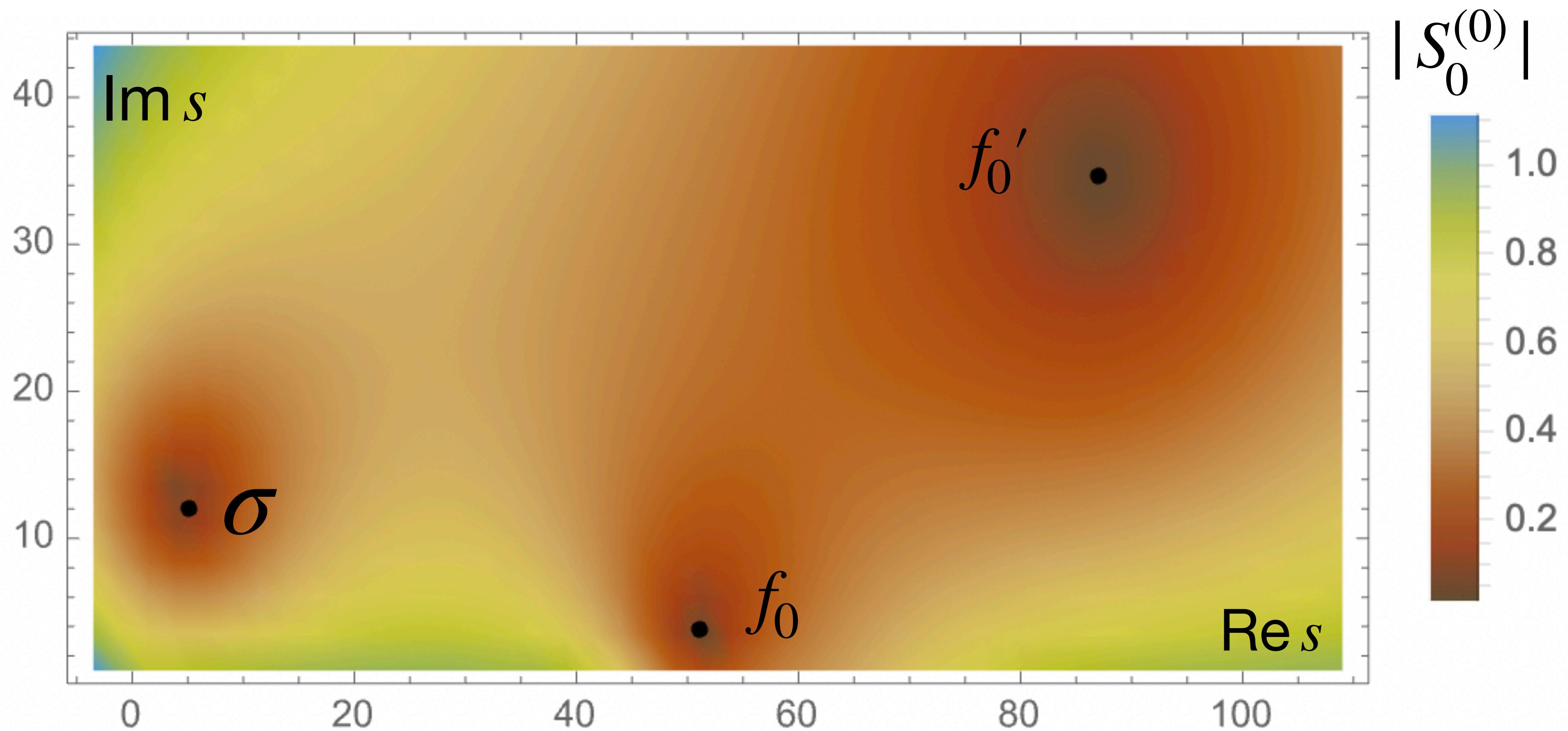
	Bootstrap Fit	Literature
$a_0^{(2)}$	$(-0.432 \pm 0.001) \times 10^{-1}$	$(-0.444 \pm 0.012) \times 10^{-1}$
$a_1^{(1)}$	$(0.380 \pm 0.002) \times 10^{-1}$	$(0.379 \pm 0.05) \times 10^{-1}$
$b_0^{(0)}$	0.265 ± 0.030	0.276 ± 0.006
$b_0^{(2)}$	$(-0.797 \pm 0.002) \times 10^{-1}$	$(-0.803 \pm 0.012) \times 10^{-1}$
$b_1^{(1)}$	$(0.61 \pm 0.02) \times 10^{-2}$	$(0.57 \pm 0.01) \times 10^{-2}$
$a_2^{(0)}$	$(0.53 \pm 0.11) \times 10^{-2}$	$(0.175 \pm 0.003) \times 10^{-2}$
$a_2^{(2)}$	$(0.51 \pm 0.18) \times 10^{-3}$	$(0.170 \pm 0.013) \times 10^{-3}$
$a_1^{(3)}$	$(1.5 \pm 0.4) \times 10^{-4}$	$(0.56 \pm 0.02) \times 10^{-4}$



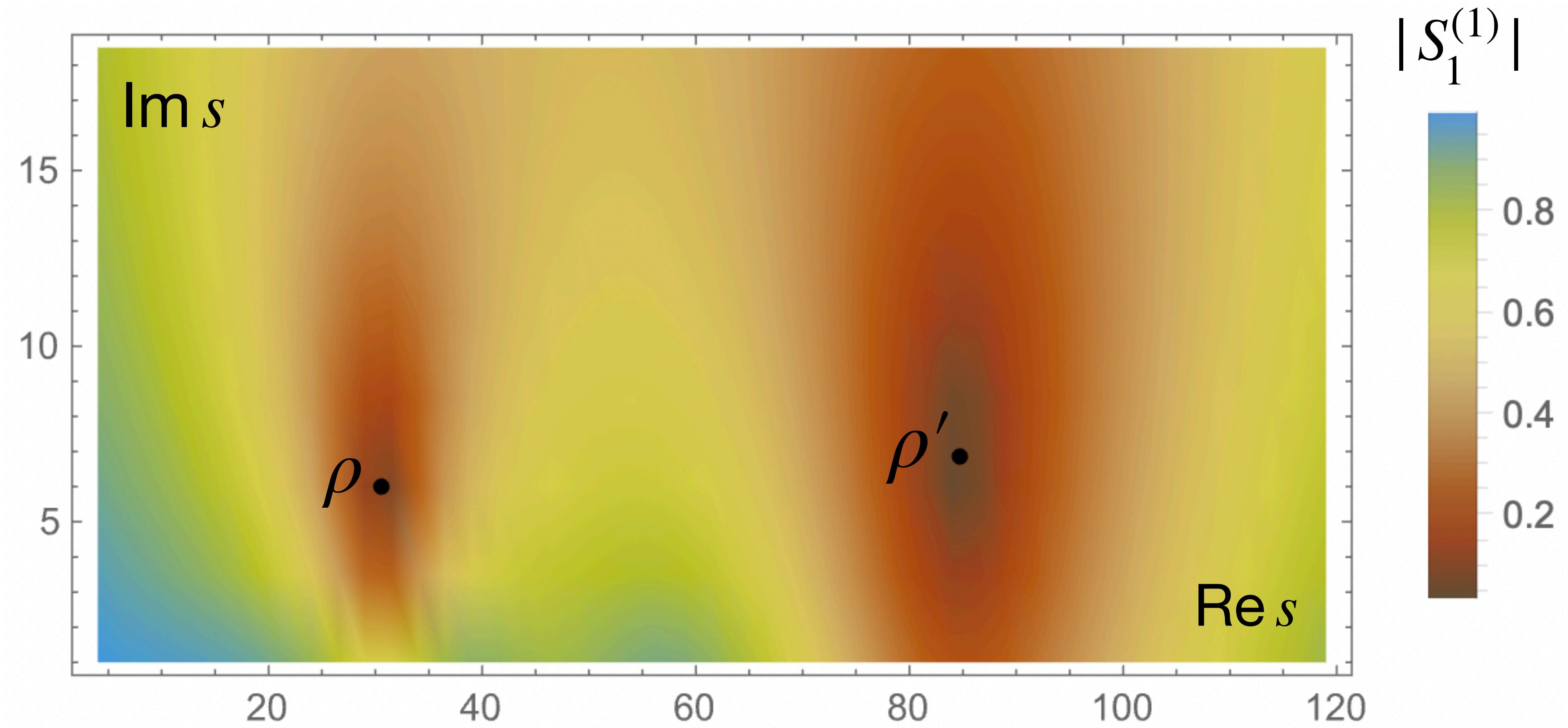
Predictions for $I=2, \ell=2$



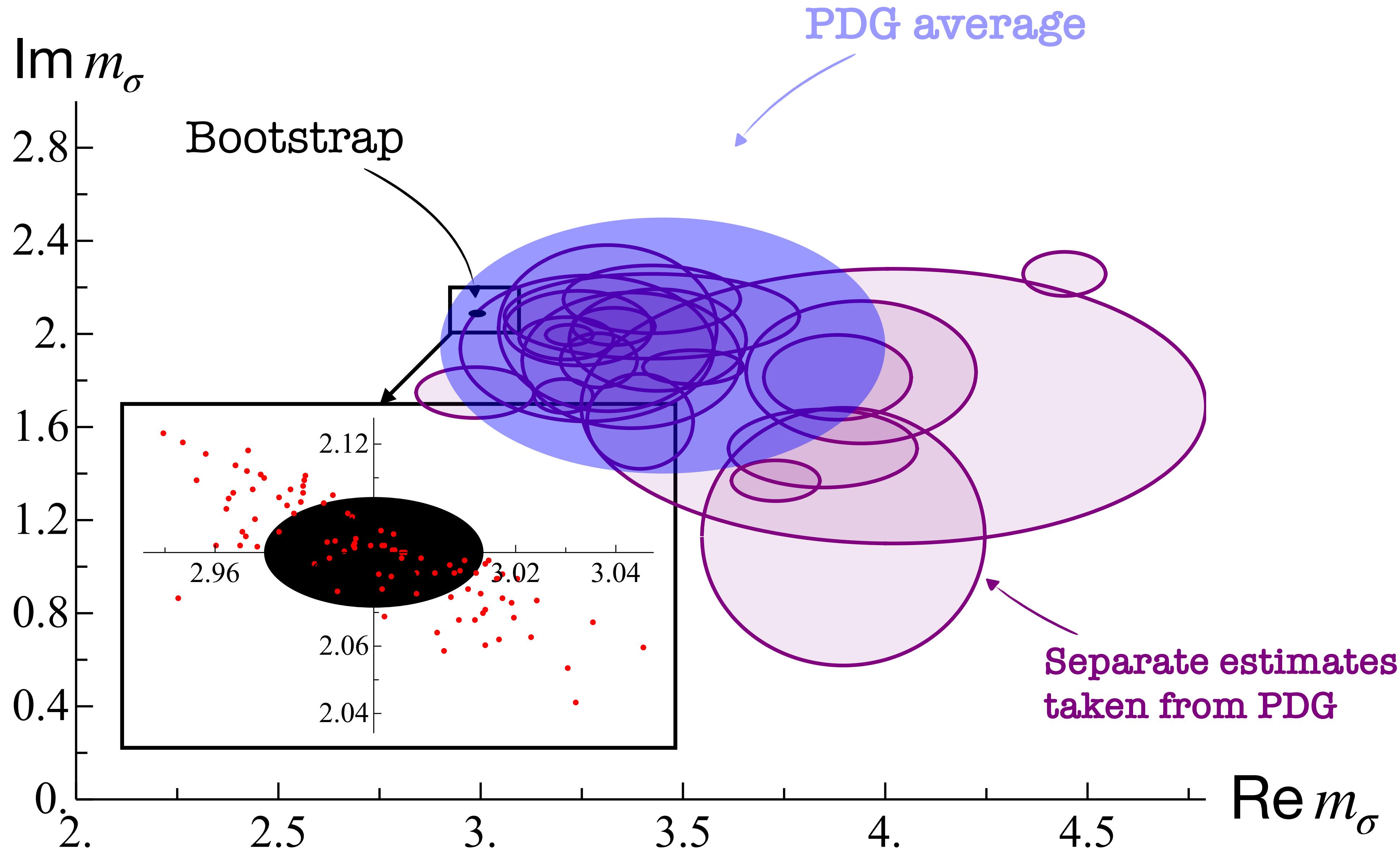
Spectrum for I=0, J=0



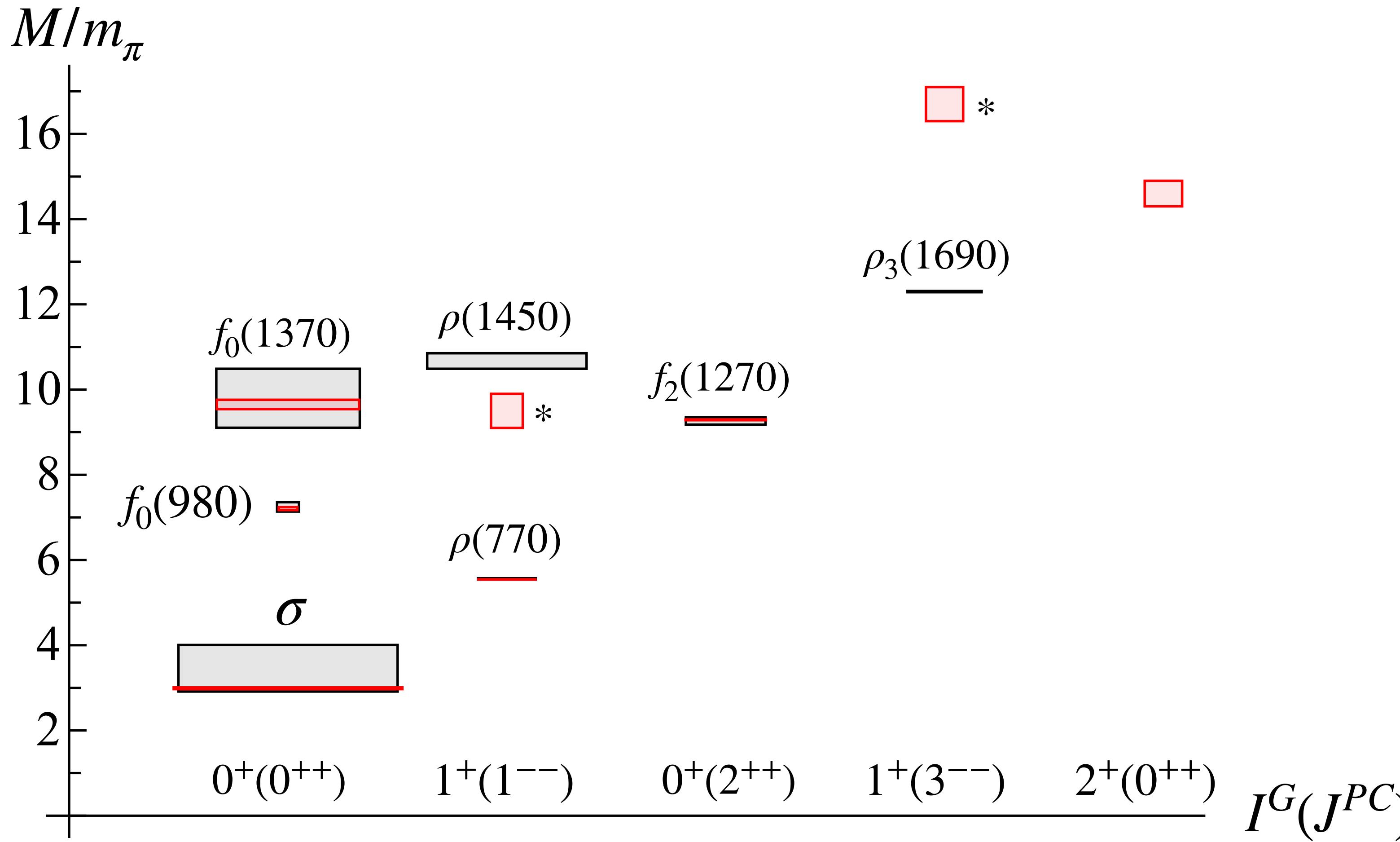
Spectrum for I=1, J=1



Sigma parameters determination

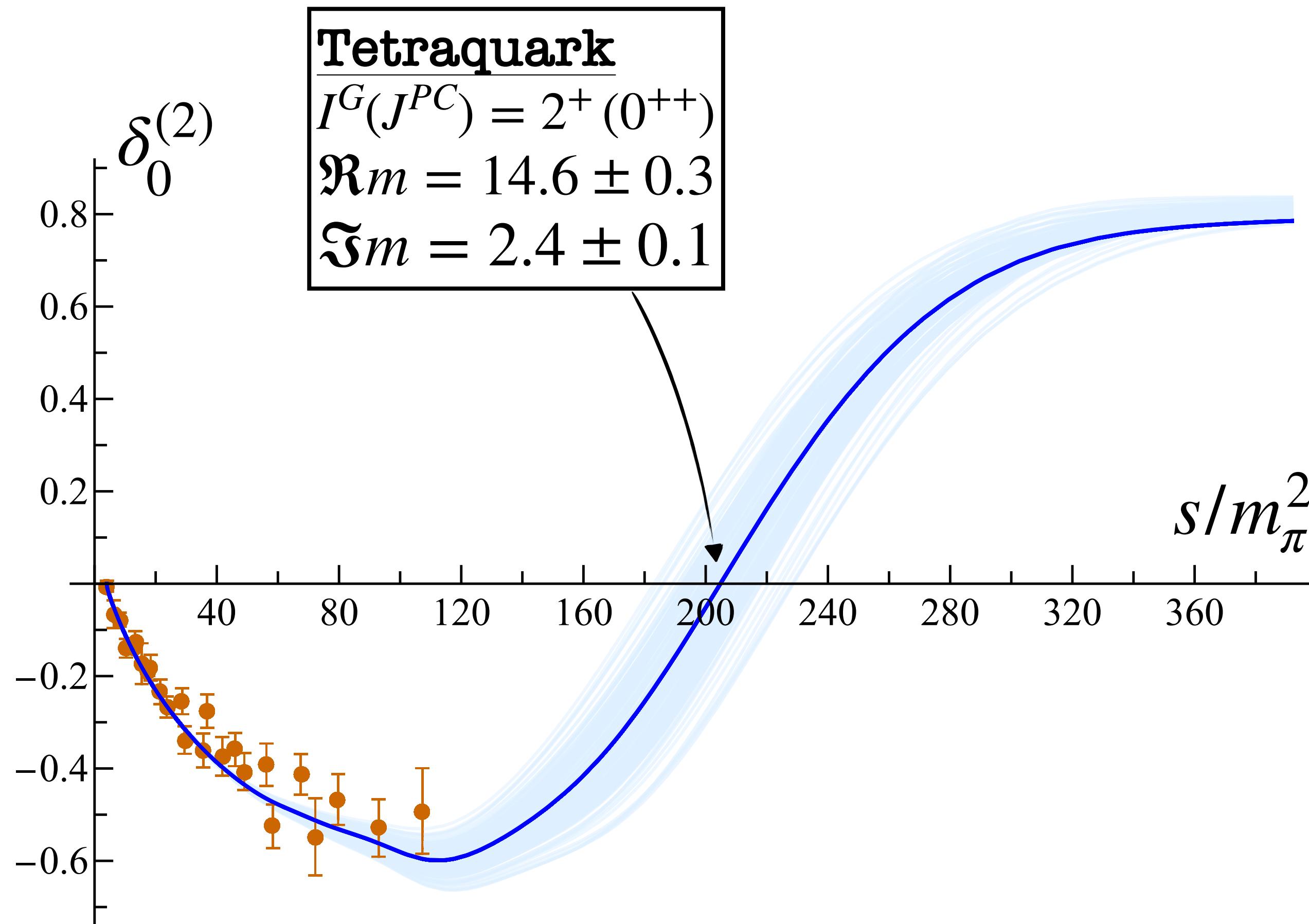


Full Spectrum < 1.4 GeV, with G-parity +1



The Tetraquark

- Each quark carries $I = \frac{1}{2}$
- 4 quarks to make $I = 2$
- Charge 2 state



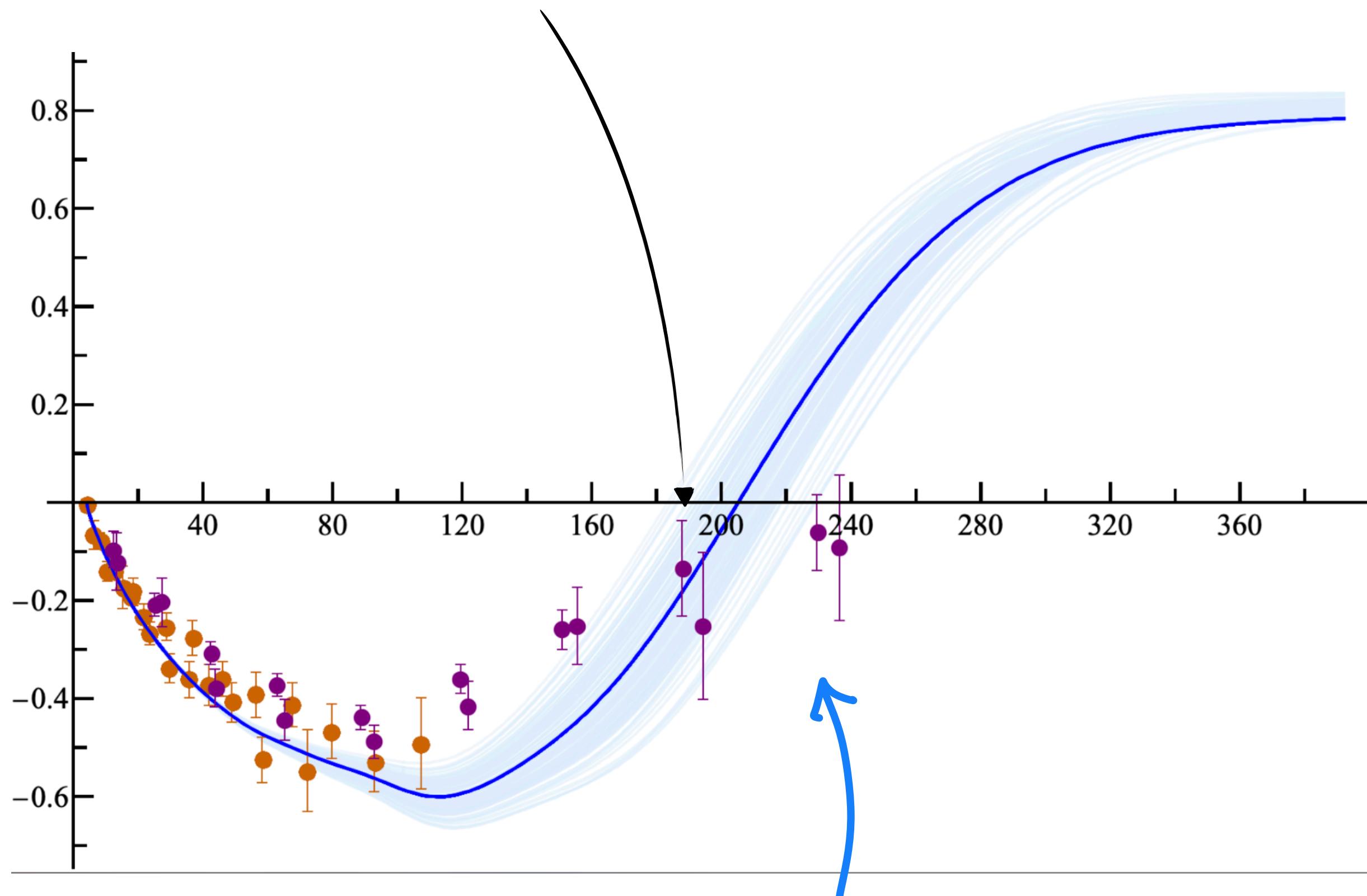
LHCb looking into $B^+ \rightarrow \pi^- \pi^+ \pi^+$!

(Cannot show results yet !)

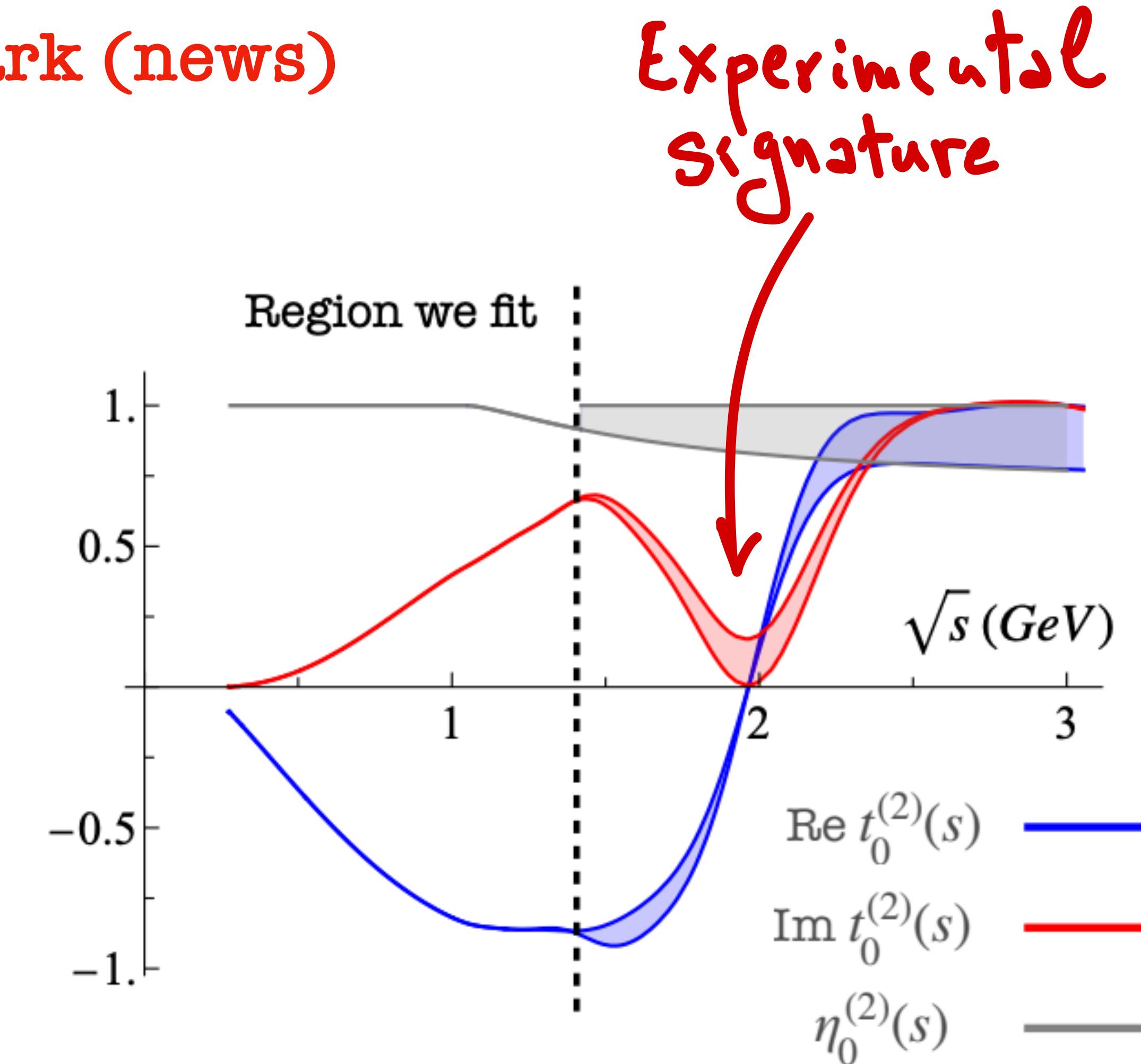
$M \sim 2\text{GeV}$,
 $\Gamma \sim 600\text{MeV}$

The Tetraquark (news)

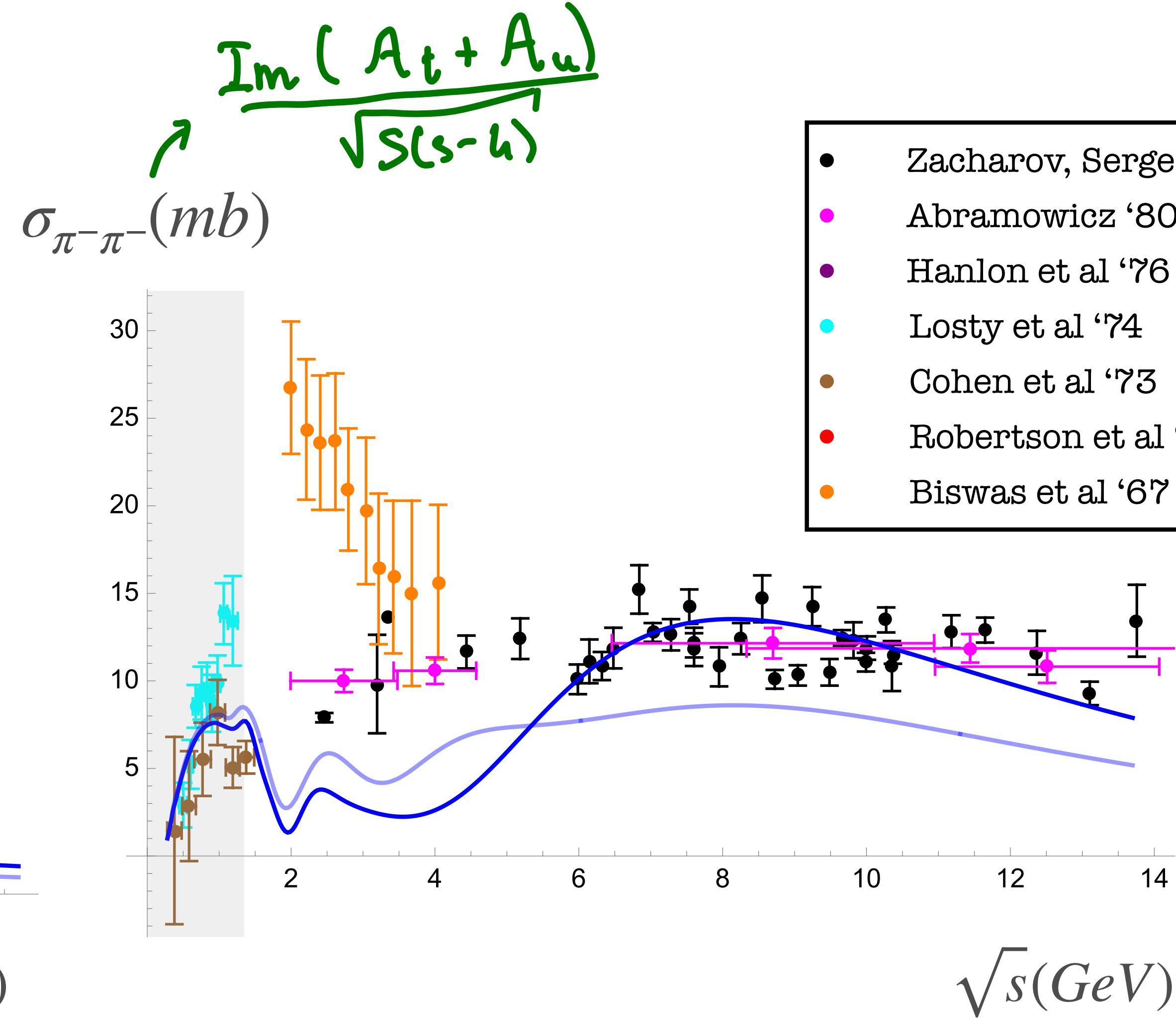
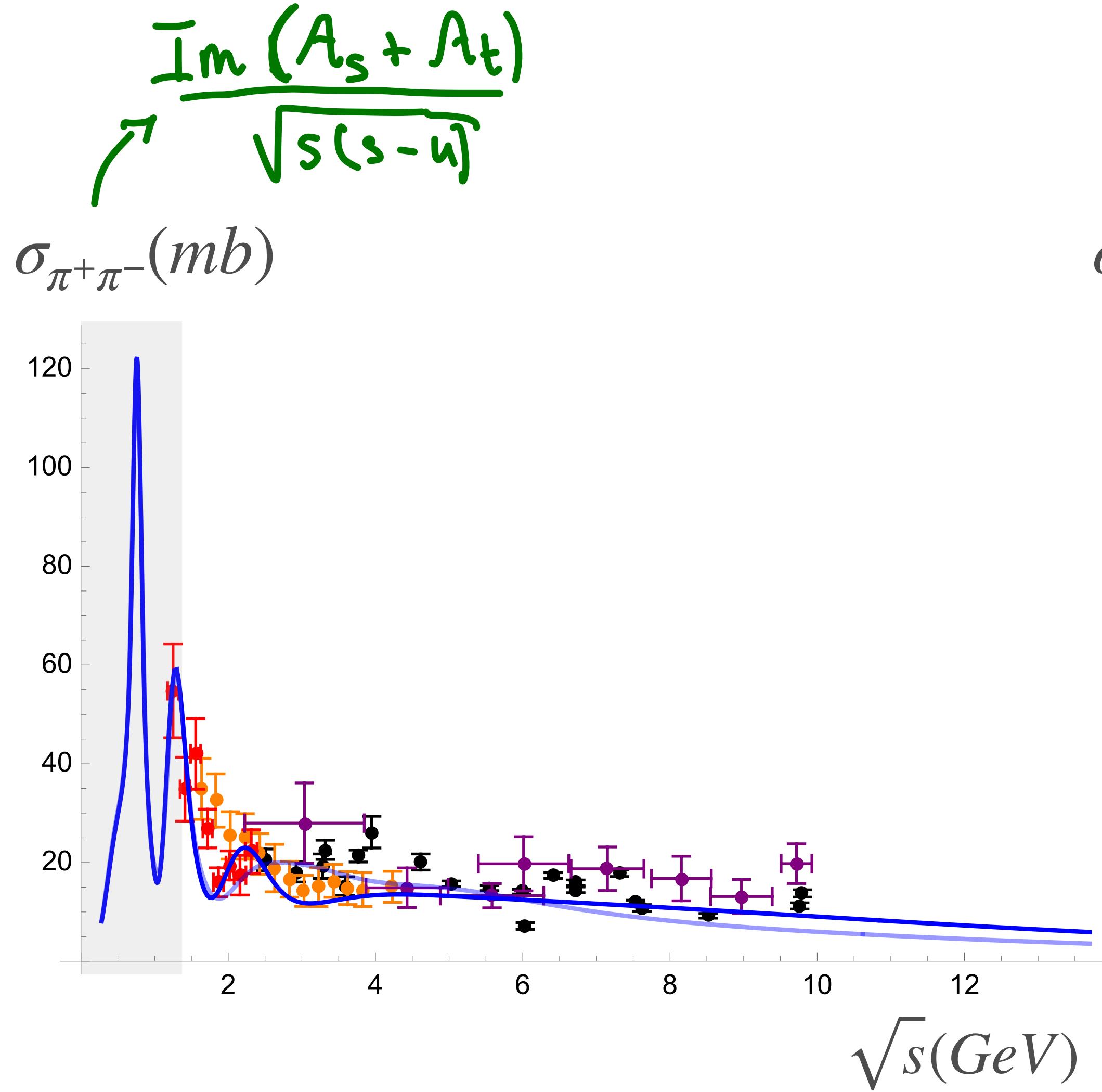
Data we had no idea about



Thanks to Luiz Vale Silva



Check @ High Energy



- Zacharov, Sergeev '84
- Abramowicz '80
- Hanlon et al '76
- Losty et al '74
- Cohen et al '73
- Robertson et al '73
- Biswas et al '67

Overview

Input

- ★ Experimental phase shifts data for S_0, S_2, P, D_0 waves
- ★ Inelasticity model for S_0, S_2, P, D_0 waves
- ★ \exists chiral zeros
- ★ \exists resonances $\rho(770), f_0(980), f_0(1370), f_2(1270)$

Output

- ♣ Fit for the phase shifts, chiral zeros, resonances positions
- ♣ Scattering lengths and effective ranges for any isospin and spin $\ell < 2$
- ♣ S_0, S_2, P waves for $0 < s < 4$ compatible with χ PT
- ♣ D_2 phase shift and inelasticity compatible with experiments
- ♣ Dynamical generation of $\sigma, \rho(1450), \rho_3$ resonances, plus a tetra quark
- ♣ $\sigma_{\pi^+\pi^-}$ and $\sigma_{\pi^-\pi^-}$ cross sections

Outlook

- 1) Improve numerics, go to higher energy, extract more resonances
- 2) $\pi\pi \rightarrow K\bar{K}$, $K\bar{K} \rightarrow K\bar{K}$ into the mix to fit inelasticity and study resonances with $S=1$.
- 3) Add form factors to fit data and constrain the Vacuum hadronic polarization contribution.

To appear soon

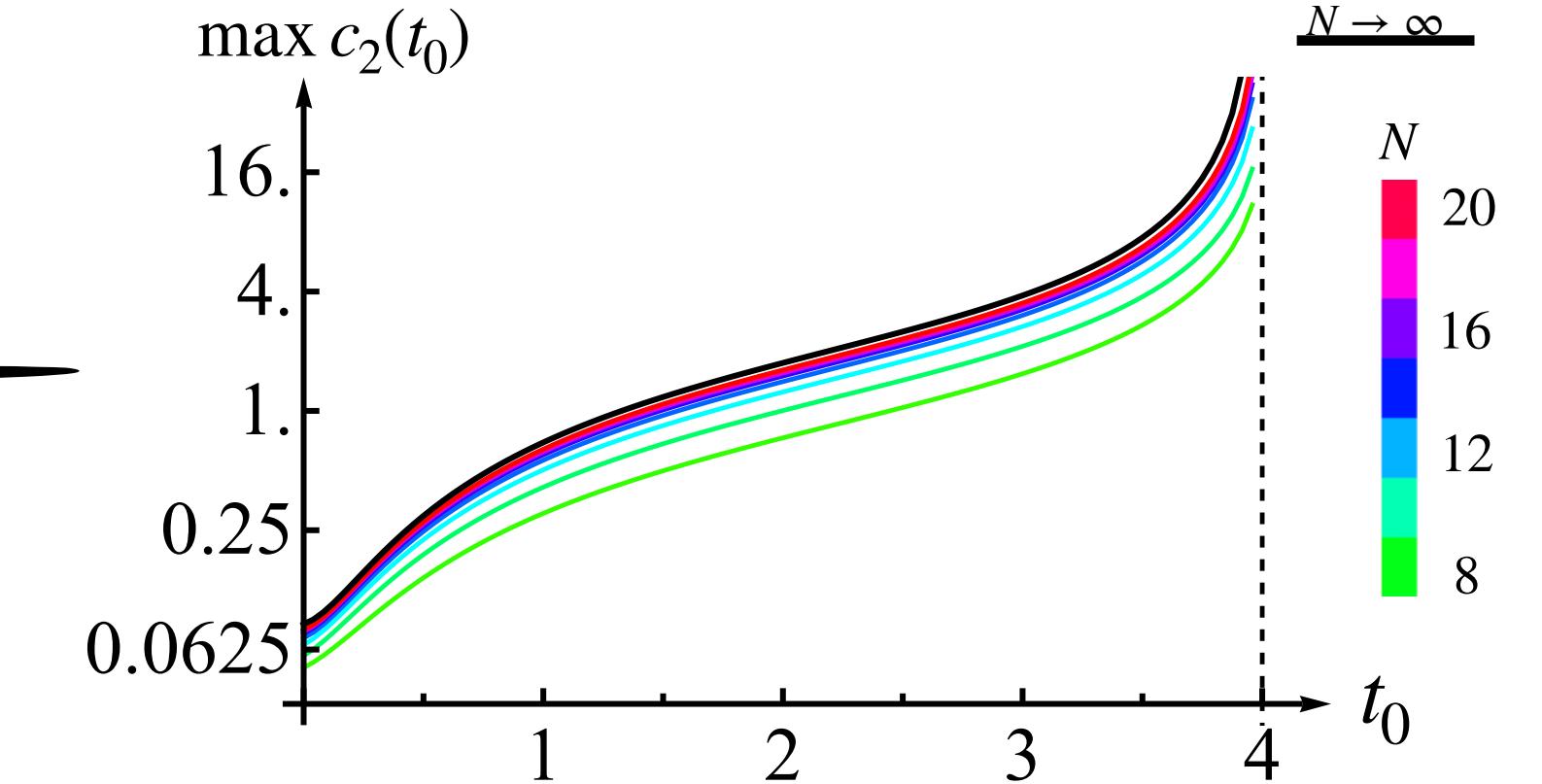
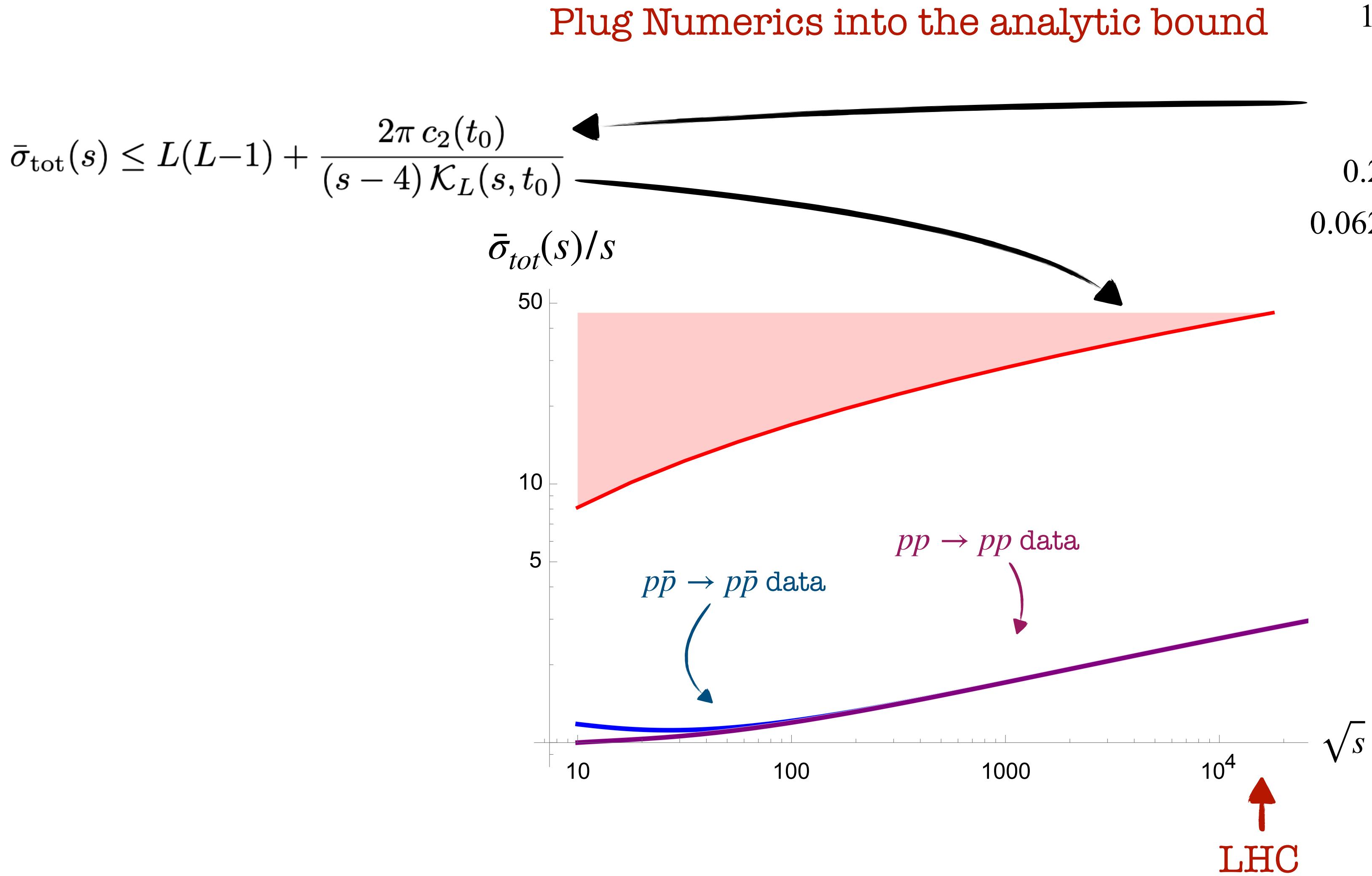
- 1) Maximizing the interaction strength: Cross-section Bootstrap with A. Georgoudis, and M. Correia
- 2) Geometry of crossing symmetric dispersion relations with J. Elias-Miro', M. Gumus, A. Zahed

← NEW APPROACH TO
THE FROISSART BOUND
"BEST TOY MODEL FOR P-P
SCATTERING"
(HAVE PLOTS
IF YOU WANT)

Connection Algebraic Geometry
and Bootstrap

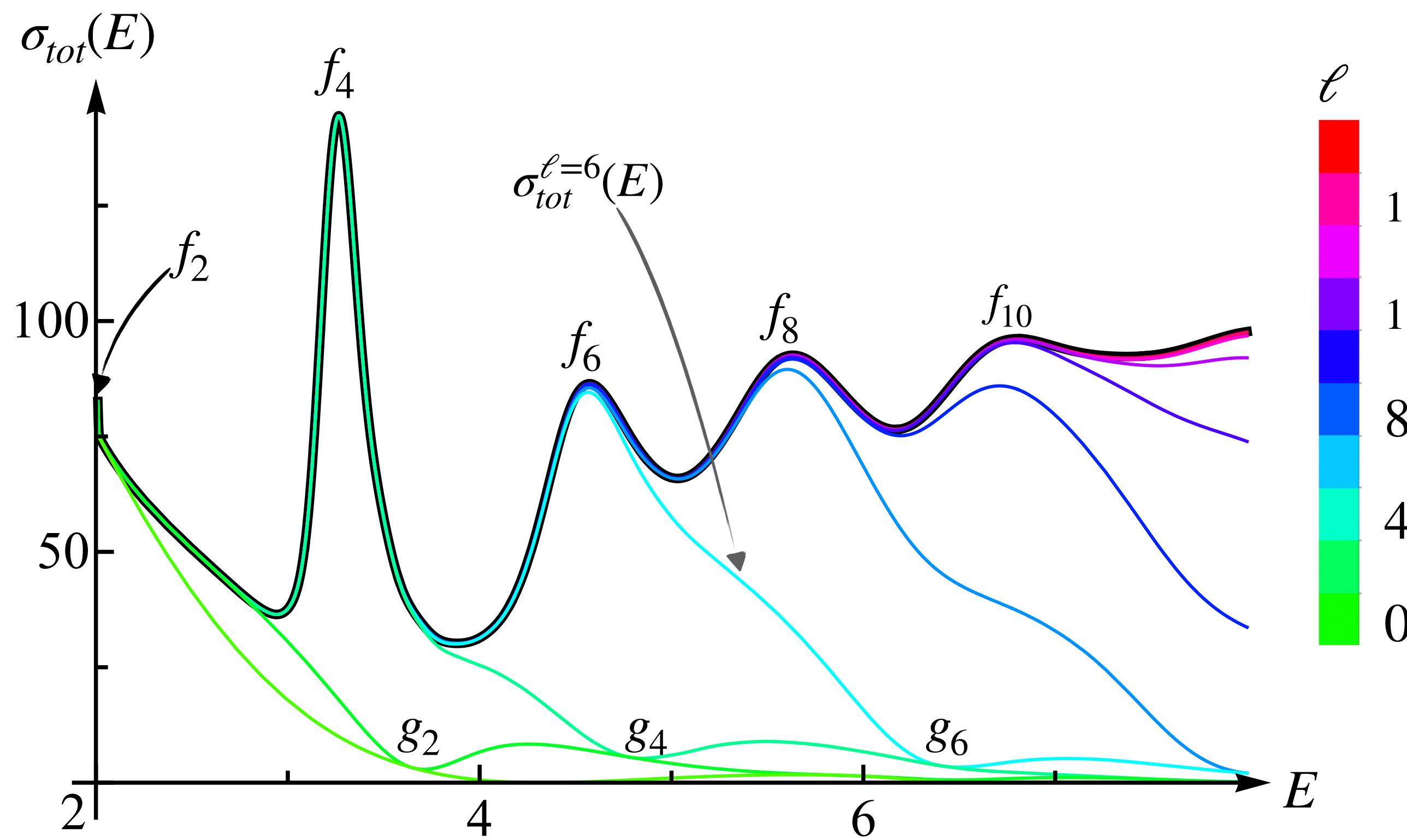
Backup

$$\bar{\sigma}_{\text{tot}}(s) \equiv \frac{1}{16\pi} \int_{4m^2}^s \frac{s' - 4m^2}{s - 4m^2} \sigma_{\text{tot}}(s') ds'$$

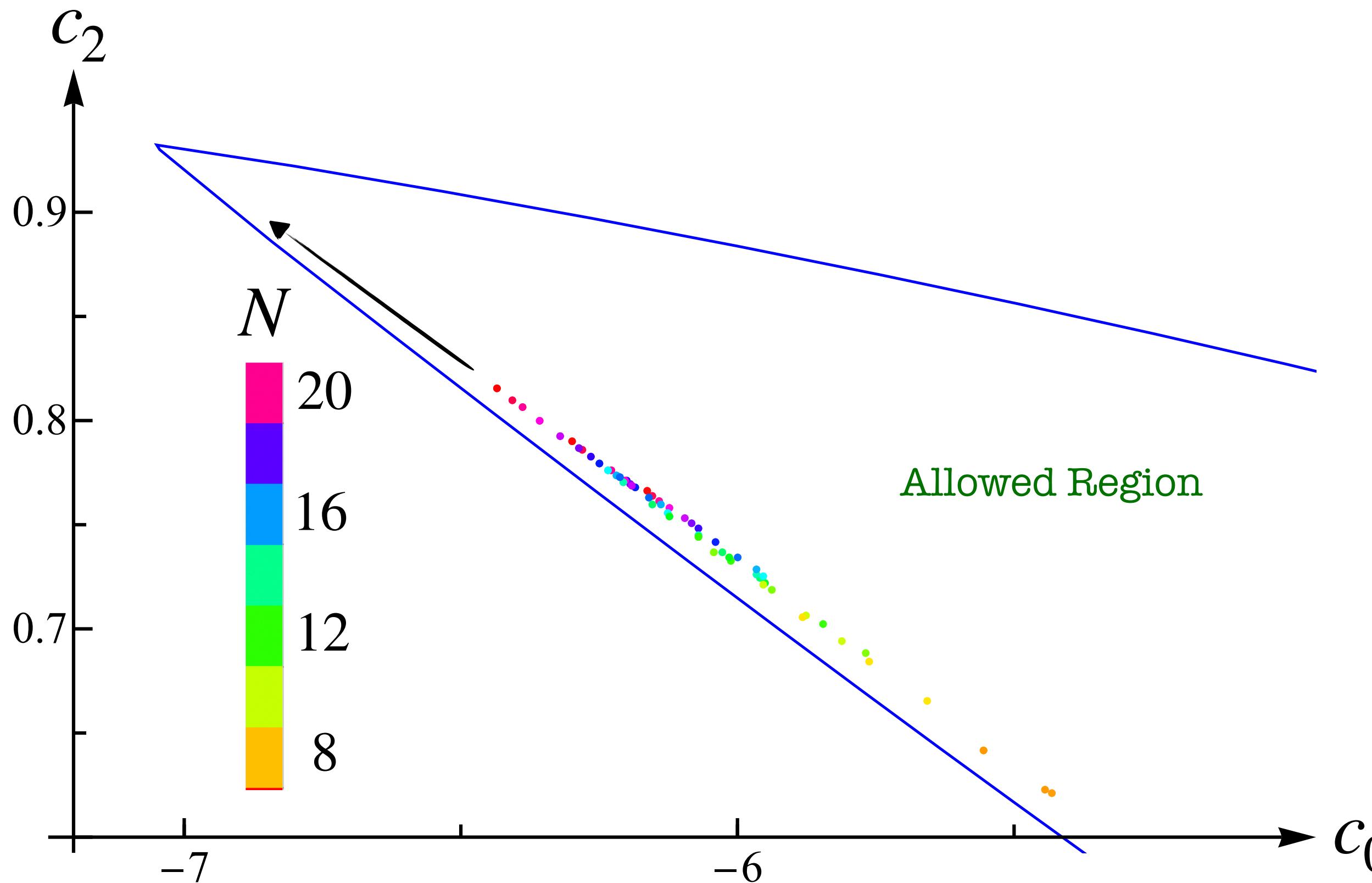


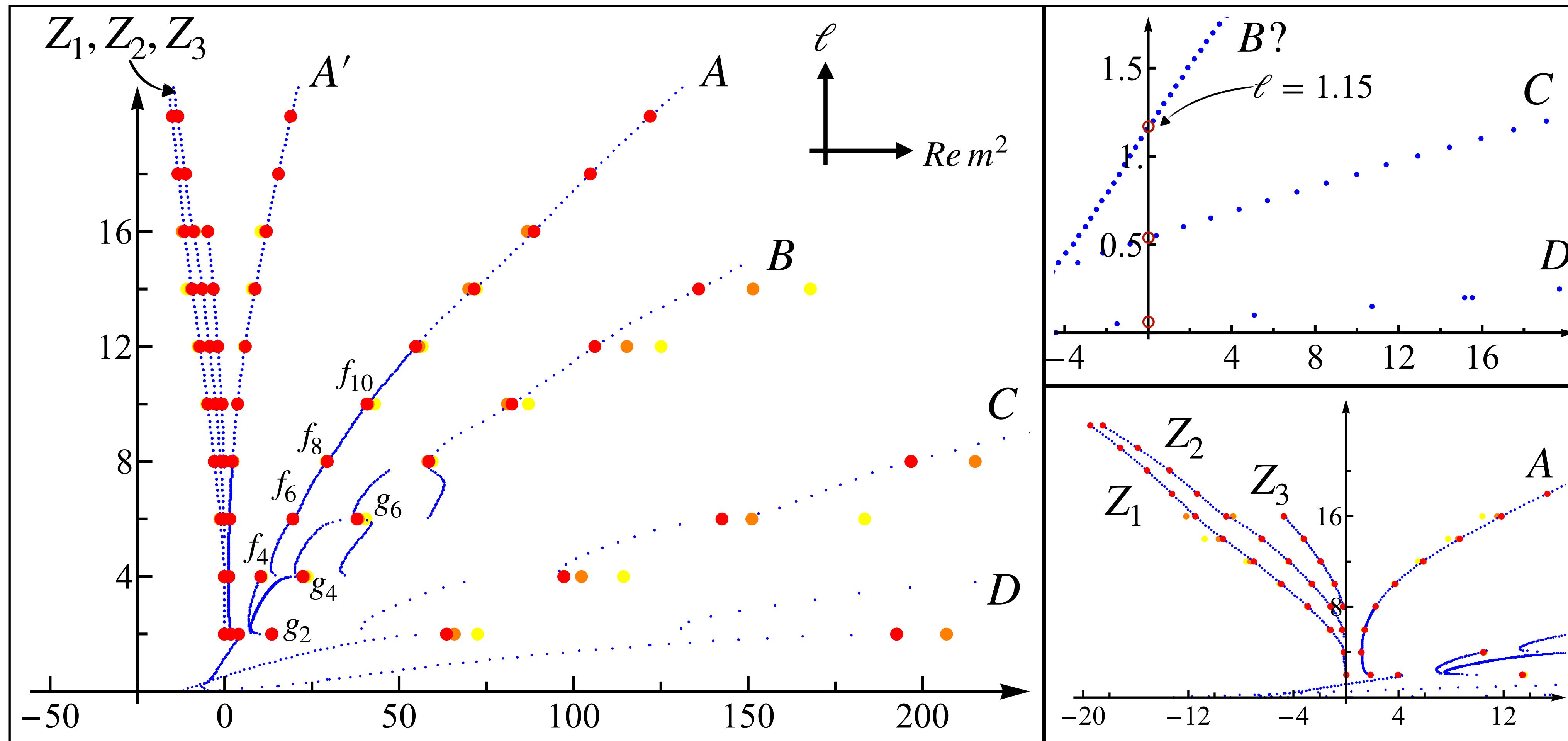
$$\lim_{s \rightarrow \infty} \max \bar{\sigma}_{tot} = \infty$$

The amplitude saturating the bound has a finite limit: Froissart Amplitude

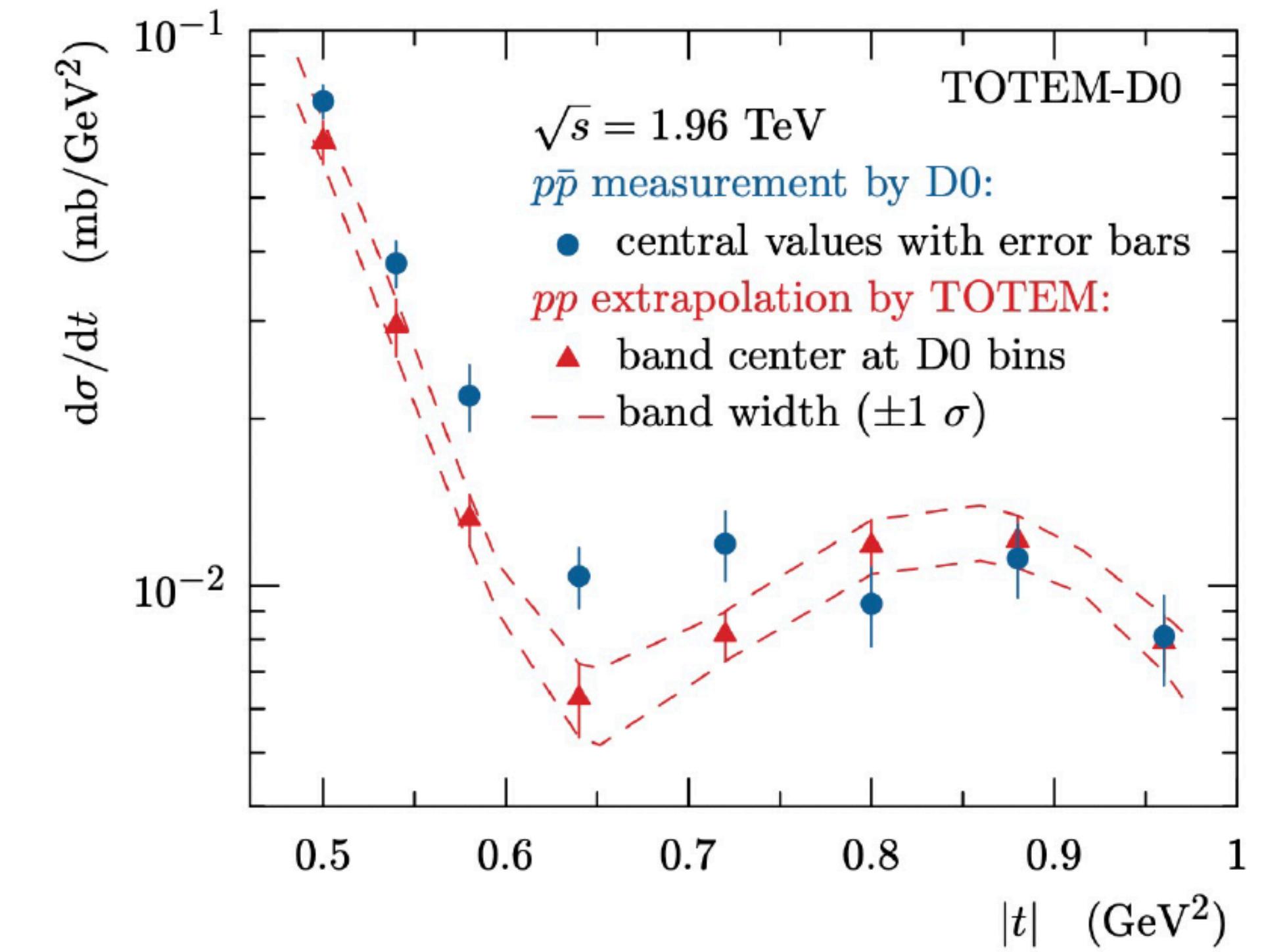
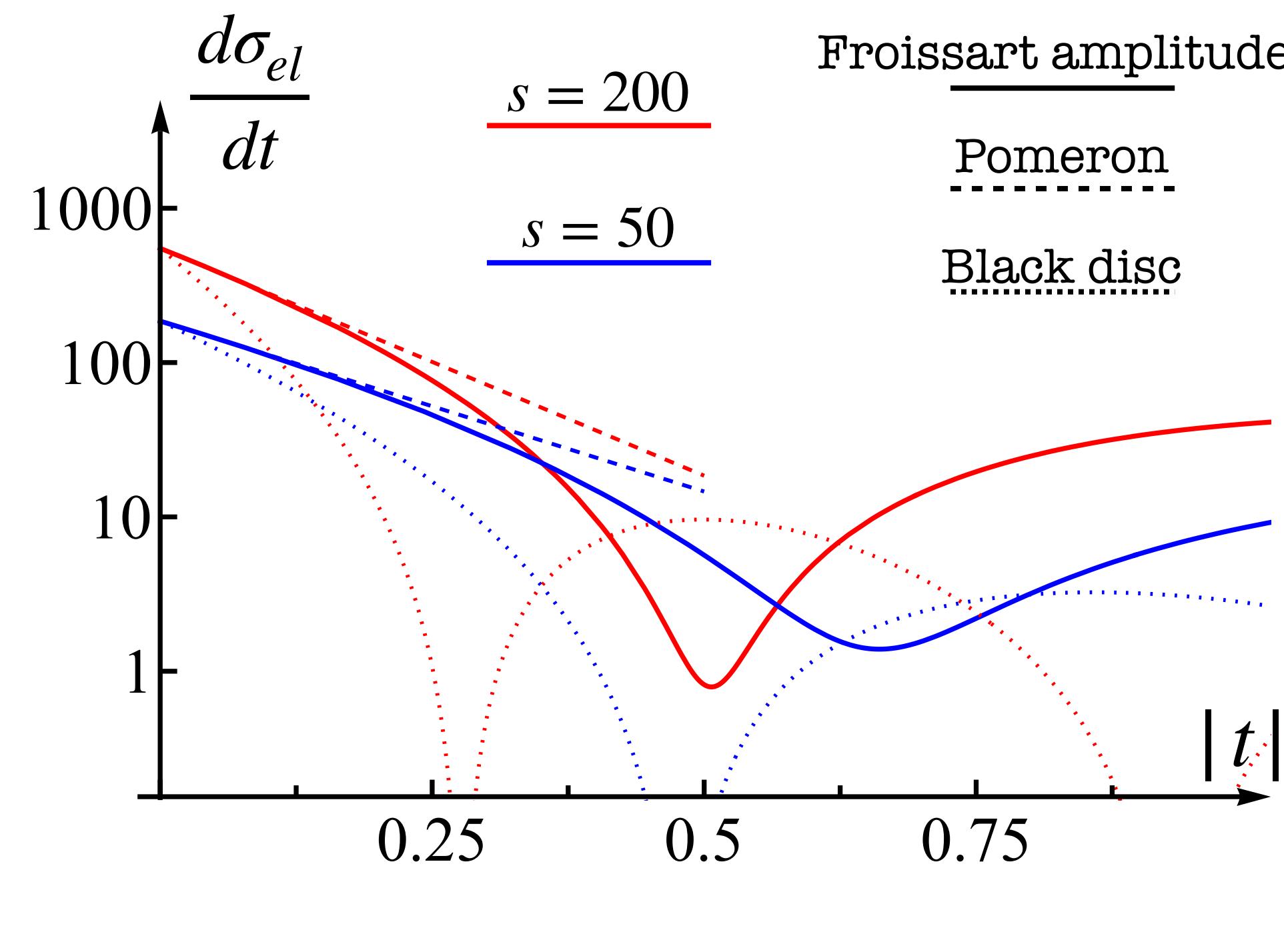


The Froissart Amplitude lives at an infinite dimensional kink





Similar Pattern for the diffractive cone



For $s \rightarrow \infty$, and fixed- t , Regge theory implies $T(s, t) \sim f(t)s^{\alpha(t)}$

Leading intercept should be $s^{1.15}$, but the ρ ansatz goes to a constant!

