## Lattice Progress on $\varepsilon^{\prime} / \varepsilon$

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- Introduction and reminder about quenched $\varepsilon^{\prime} / \varepsilon$ results
- Progress with $2+1$ flavor dynamical QCD simulations
- Chiral perturbation theory and lattice data for $\mathrm{m}_{\pi}$ and $\mathrm{f}_{\pi}$
- The denominator of $\varepsilon^{\prime} / \varepsilon: \mathrm{B}_{\mathrm{K}}$ and $\varepsilon$
- Preliminary results for $\Delta S=1$ matrix elements for $\varepsilon^{\prime} / \varepsilon$
- Conclusions


## CP Violation in the Kaon System

- Two amplitudes determine $\epsilon$ and $\epsilon^{\prime}$

$$
\eta_{+-}=\frac{A\left(K_{L}^{0} \rightarrow \pi^{+} \pi^{-}\right)}{A\left(K_{S}^{0} \rightarrow \pi^{+} \pi^{-}\right)}=\epsilon+\epsilon^{\prime} \quad \eta_{00}=\frac{A\left(K_{L}^{0} \rightarrow \pi^{0} \pi^{0}\right)}{A\left(K_{S}^{0} \rightarrow \pi^{0} \pi^{0}\right)}=\epsilon-2 \epsilon^{\prime}
$$

- SM: $\bar{K}^{0}-K^{0}$ mixing via $Q^{(\Delta S=2)}=\left(\bar{s}_{\alpha} d_{\alpha}\right)_{V-A}\left(\bar{s}_{\beta} d_{\beta}\right)_{V-A}$ defines $B_{K}$ as;

$$
\left.\left\langle\bar{K}^{0}\right| Q^{(\Delta S=2)}(\mu)\left|K^{0}\right|\right\rangle \equiv \frac{8}{3} B_{K}(\mu) f_{K}^{2} m_{K}^{2}
$$

- RGI parameter $\hat{B}_{K} \equiv B_{K}(\mu)\left[\alpha_{s}^{(3)}(\mu)\right]^{-2 / 9}\left[1+\frac{\alpha_{s}^{(3)}(\mu)}{4 \pi} J_{3}\right]$ relates SM and $\epsilon$

$$
\epsilon=\hat{B}_{K} \operatorname{Im} \lambda_{t} \frac{G_{F}^{2} f_{K}^{2} m_{K} M_{W}^{2}}{12 \sqrt{2} \pi^{2} \Delta M_{K}}\left\{\operatorname{Re} \lambda_{c}\left[\eta_{1} S_{0}\left(x_{c}\right)-\eta_{3} S_{0}\left(x_{c}, x_{t}\right)\right]-\operatorname{Re} \lambda_{t} \eta_{2} S_{0}\left(x_{t}\right)\right\} \exp (i \pi / 4)
$$

- Defining $A\left(K^{0} \rightarrow \pi \pi(I)\right) \equiv A_{I} e^{\left(i \delta_{I}\right)}, P_{2} \equiv \operatorname{Im} A_{2} / \operatorname{ReA}_{2}, P_{0} \equiv \operatorname{Im} A_{0} / \operatorname{ReA}_{0}$ :

$$
\epsilon^{\prime}=\frac{i e^{i\left(\delta_{2}-\delta_{0}\right)}}{\sqrt{2}}\left(\frac{\operatorname{Re} A_{2}}{\operatorname{Re} A_{0}}\right)\left(\frac{\operatorname{Im} A_{2}}{\operatorname{Re} A_{2}}-\frac{\operatorname{Im} A_{0}}{\operatorname{Re} A_{0}}\right) \quad w \equiv \frac{\operatorname{Re} A_{0}}{\operatorname{Re} A_{2}} \approx 22
$$

## Low Energy Standard Model Diagrams for $\mathrm{B}_{\mathrm{K}}$



Electroweak process at high energy scales reduce to a single 4-fermion operator at low energies

Need correctly normalized value of the $\mathrm{Q}^{\Delta \mathrm{S}=2}$ operator in kaon states.

## Standard Model Diagrams for $\varepsilon^{\prime} / \varepsilon$



## $K \rightarrow \pi \pi$ in 3-flavor Effective Theory

- Hamiltonian for 3-flavor effective theory: only 7 of 10 operators independent

$$
\mathcal{H}^{(\Delta S=1)}=\frac{G_{F}}{\sqrt{2}} V_{u d} V_{u s}^{*}\left\{\sum_{i=1}^{10}\left[z_{i}(\mu)+\tau y_{i}(\mu)\right] Q_{i}\right\}
$$

- $K \rightarrow \pi \pi$ from lattice calculations and LO chiral perturbation theory.

| Irrep | Isospin | $K^{+} \rightarrow \pi^{+}$ | $K^{0} \rightarrow \pi^{+} \pi^{-}$ |
| :---: | :--- | :---: | :---: |
| $(27,1)$ | $1 / 2,3 / 2$ | $-\frac{4 m_{M}^{2}}{f^{2}} \alpha^{(27,1)}$ | $-\frac{4 i}{f^{3}} m_{K^{0}}^{2} \alpha^{(27,1)}$ |
| $(8,8)$ | $1 / 2,3 / 2$ | $-\frac{12}{f^{2}} \alpha^{(8,8)}$ | $-\frac{12 i}{f^{3}} \alpha^{(8,8)}$ |
| $(8,1)$ | $1 / 2$ | $\frac{4 m_{M}^{2}}{f^{2}}\left(\alpha_{1}^{(8,1)}-\alpha_{2}^{(8,1)}\right)$ | $\frac{4 i}{f^{3}} m_{K^{0}}^{2} \alpha_{1}^{(8,1)}$ |

- $(8,1)$ coefficient $\alpha_{2}^{(8,1)}$ is power divergent, $\mathcal{O}\left(1 / a^{2}\right)$. Determine from $K \rightarrow|0\rangle$


## Quenched Chiral Extrapolations $(27,1)$ and $(8,1)$

Fit with known continuum chiral logarithm for quenched theory

$$
1-\frac{6 m_{M}^{2}}{(4 \pi f)^{2}} \ln \left(m_{M}^{2} / \Lambda^{2}\right)
$$

Good description of data, but $400 \mathrm{MeV} \leq m_{\mathrm{PS}} \leq 800 \mathrm{MeV}$


Only slope relevant in subtracted ME


## Real $K \rightarrow \pi \pi$ Amplitudes from Quenched QCD and $\chi \mathbf{P T}$



## $\epsilon^{\prime} / \epsilon$ from Quenched QCD and $\chi \mathbf{P T}$

- Dominant contribution: $Q_{2}$ to $\operatorname{Re} A_{2}$ and $\operatorname{Re} A_{0}, Q_{6}$ to $\operatorname{Im} A_{0}, Q_{8}$ to $\operatorname{Im} A_{2}$.
- Contributions depend on renormalization scale GeV
- Schematic formula for $\epsilon^{\prime} / \epsilon$

$$
\operatorname{Re}\left(\epsilon^{\prime} / \epsilon\right) \approx\left(\frac{\omega}{\sqrt{2}|\epsilon|}\right)_{\exp }\left\{\left[\frac{\alpha_{\mathrm{W}} \alpha_{8}}{\alpha_{\mathrm{W}} \alpha_{8}+\alpha_{2} m_{K^{0}}^{2} \xi}\right]^{(3 / 2)}-\left[\frac{\alpha_{\mathrm{W}} \alpha_{8}+\alpha_{\mathrm{S}} \alpha_{6} m_{K^{0}}^{2} \xi}{\alpha_{\mathrm{W}} \alpha_{8}+\alpha_{2} m_{K^{0}}^{2} \xi}\right]^{(1 / 2)}\right\}
$$



## Achieving Accurate Kaon Physics on the Lattice

| Issue | Current status |
| :--- | :--- |
| Quenched approximation | $2+1$ flavor DWF and ASQTAD |
| Chiral symmetry breaking | Staggered fermions <br> Twisted mass Wilson fermions <br> Domain wall fermion $\checkmark$ <br> Overlap fermions $\checkmark$ |
| Heavy pions | ASQTAD: one pion has $\mathrm{m}_{1}=\mathrm{m}_{\mathrm{s}} / 10$ <br> DWF: correct light pions with $\mathrm{m}_{1}=\mathrm{m}_{\mathrm{s}} / 7$ |
| Operator Renormalization | Non-perturbative renormalization (NPR) <br> Schrodinger functional methods |
| Extrapolation to chiral limit | Chiral perturbation theory: <br> DWF - continuum like <br> ASQTAD - include taste breaking |
| Multiparticle final states | 1) Avoid via ChPT <br> 2) Use finite volume effects |
| More computing speed | Many sustained Teraflops currently |



25,000 nodes at Brookhaven RBRC and USDOE machines

14,000 nodes at the University of Edinburgh


## Collaboration Members

RBC members:

Y. Aoki, C. Aubin, T. Blum, M. Cheng, N. Christ, S. Cohen, C.<br>Dawson, T. Doi, K. Hashimoto, T. Ishikawa, T. Izubuchi, C. Jung, M. Li, S. Li, M. Lightman, H. Lin, M. Lin, O. Loktik, R. Mawhinney, S. Ohta, S. Sasaki, E. Scholz, A. Soni, T. Yamazaki

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## Zero Temperature Ensembles

| Volume | $a^{-1}(\mathrm{GeV})$ | $\left(m_{l} m_{s}\right)$ | $m_{\text {res }}$ | MD time units |
| :---: | :---: | :---: | :---: | :---: |
| $16^{3} \times 32 \times 12$ | 1.69(5) |  | 0.00137(5) | $\begin{aligned} & 2680.5 \\ & 3097.5 \\ & 3252.5 \\ & \hline \end{aligned}$ |
| $16^{3} \times 32 \times 8$ | 1.8(1) | $\begin{array}{\|c} \hline(0.02,0.04) \\ (0.04,0.04) \\ \hline \end{array}$ | 0.0107(1) | $\begin{aligned} & 1797.5 \\ & 1797.5 \\ & \hline \end{aligned}$ |
| $16^{3} \times 32 \times 16$ | 1.62(4) | $\begin{gathered} \hline(0.01,0.04) \\ (0.02,0.04) \\ (0.03,0.04) \\ \hline \end{gathered}$ | 0.00308(4) | 4015 <br> 4045 <br> $4020+3580$ |
| $24^{3} \times 64 \times 16$ | 1.6-1.7 | $\begin{gathered} \hline(0.005,0.04) \\ (0.01,0.04) \\ (0.02,0.04) \\ (0.03,0.04) \\ \hline \end{gathered}$ | 0.0031 | $\begin{aligned} & 4500 \\ & 3785 \\ & 2850 \\ & 2813 \\ & \hline \end{aligned}$ |
| $32^{3} \times 64 \times 16$ | 2.1-2.2 | $\begin{aligned} & (0.004,0.03) \\ & (0.006,0.03) \\ & \hline \end{aligned}$ | $\approx 0.0005$ | $\begin{aligned} & 500 \\ & 892 \end{aligned}$ |

First row is with DBW2 gauge action, all others use the Iwasaki action.

## Partially Quenched NLO ChPT for $\mathrm{m}_{\pi}$ and $\mathrm{f}_{\pi}$




| $(3 \mathrm{fm})^{3}$ <br> volume |  |
| :---: | :---: |
| $\mathrm{m}_{\text {dyn }}$ <br> $(\mathrm{MeV})$ | $\mathrm{m}_{\text {val }}$ <br> $(\mathrm{MeV})$ |
| 550 | 550 |
| 400 | 400 |
| 320 | 320 |
|  | 220 |
|  |  |
|  |  |
| 400 | 400 |
| 320 | 320 |
|  | 220 |

Meifeng Lin
DWF@ 10

## $B_{K}$ Plateau

 Saul Cohen

$B$ Plateau on $16^{3}, m_{s}^{\text {sea }}=0.04, m_{l}^{\text {sea }}=0.02$

$24^{3}$ volume




## Fitting $\mathrm{B}_{\mathrm{PS}}$ to NLO Partially Quenched ChPT

Both $16^{3}$ and $24^{3}$ volumes are fit to same range of masses, 400 to 750

MeV.





NLO formula are reasonable interpolations of our data, but fail to go through light quark mass points.

## Unitary Extrapolation for $\mathrm{B}_{\mathrm{K}}$

Compare simple unitary extrapolation to result with partially quenched ChPT.

Get same result within statistical errors.


## Final Result for $\mathrm{B}_{\mathrm{K}}$


$\mathrm{B}_{\mathrm{K}}{ }^{(\overline{\mathrm{MS}})}(2 \mathrm{GeV})=0.557(12)(29)$ extrapolated to continuum

## $\mathrm{B}_{\mathrm{K}}$ Comparison



# Graph from Chris Dawson Lattice 2005 

## $\mathrm{B}_{\mathrm{K}}{ }^{(\overline{\mathrm{MS}})}(2 \mathrm{GeV})=0.557(12)(29)$ extrapolated to continuum RBC and UKQCD Collaborations

$$
\mathrm{B}_{\mathrm{K}}{ }^{\overline{\mathrm{MS}})}(2 \mathrm{GeV})=0.58(3)(6) \text { world average }
$$

Chris Dawson, Lattice 2005, PoS(LAT2005) 007

## $\mathrm{B}_{\mathrm{K}}$ in the Chiral Limit: $\mathrm{B}_{0}$

- Use data from (3 fm $)^{3}$ volume
- Only pseudoscalars with mass $\leq 400 \mathrm{MeV}$
- 12 data points used in fits
- Preliminary result: $\mathrm{B}_{0}{ }^{(\overline{\mathrm{MS}})}(2 \mathrm{GeV})=0.34(5)$



## $\varepsilon^{\prime} / \varepsilon$ on $2+1$ flavor, $(3 \mathrm{fm})^{3}$ ensembles

- Valence masses $0.001,0.005,0.01,0.02,0.03,0.04\left(\mathrm{~m}_{\mathrm{s}} / 10\right.$ to $\left.\mathrm{m}_{\mathrm{s}}\right)$
- Concentrating on 0.005/0.04 and 0.01/0.04 ensembles
- Large contributions by Tom Blum, Saul Cohen, Sam Li.
- 0.005/0.04 ensemble: 40 configurations separated by 80 MD time units. 0.01/0.04 ensemble: 30 configurations separated by 80 MD time units
- Concentrating on lighter quark masses where NLO chiralperturbation theory should be reasonable.
- Coulomb gauge fixed wall sources at $\mathrm{t}=5$ and 59
- Random noise source of length 40 for pupil calculations
- $1 / 2$ of time in wall source calculations, the other $1 / 2$ in pupils


## $\Delta \mathrm{I}=3 / 2$ Plateau



## $\Delta \mathrm{I}=3 / 2$ Plateau Comparisons

## Previous Quenched

## New 2+1 Flavor




## $\Delta \mathrm{I}=3 / 2$ Plateau



## $\Delta \mathrm{I}=1 / 2$ Plateau - unsubtracted $\mathrm{Q}_{6}$



## $\Delta \mathrm{I}=1 / 2$ Plateau $-\overline{\mathrm{s}} \mathrm{d}$



## Comparing $\Delta \mathrm{I}=1 / 2$ Plateau

Common fluctuations in $\mathrm{Q}_{6}$ and $\overline{\text { sd }}$


## $\Delta \mathrm{I}=1 / 2$ Subtraction



## $\Delta \mathrm{I}=1 / 2$ Subtraction Coefficient Comparisons

$$
\frac{\langle 0| Q_{i, \text { lat }}\left|K^{0}\right\rangle}{\langle 0|\left(\bar{s} \gamma_{5} d\right)_{\text {lat }}\left|K^{0}\right\rangle}=\eta_{0, i}+\eta_{1, i}\left(m_{s}-m_{d}\right)
$$

## Previous Quenched

New 2+1 Flavor



## Subtracting $\mathrm{Q}_{6}$



## Subtracting $\mathrm{Q}_{6}$



## Subtracted $\mathrm{Q}_{6}$



## Subtracted $\mathrm{Q}_{6}$ Comparison

$$
\left\langle\pi^{+}\right| Q_{i, \mathrm{lat}}^{(1 / 2)}\left|K^{+}\right\rangle+\eta_{1, i}\left(m_{s}+m_{d}\right)\left\langle\pi^{+}\right|(\bar{s} d)_{\text {lat }}\left|K^{+}\right\rangle
$$

## Previous Quenched

## New 2+1 Flavor




## $\mathrm{m}_{\text {res }}$ and $\Delta \mathrm{S}=1$ matrix elements

- Spurion field $\Omega$ at midpoint represents residual $\chi \mathrm{SB}$
- Transforms as $(\overline{3}, 3)$ under chiral symmetry
- For low energy observables, $\Omega$ goes to $\mathrm{m}_{\text {res }}$
- For divergent quantities, new parameters enter which are $\mathrm{O}\left(\mathrm{m}_{\text {res }}\right)$
- Due to unsupressed modes in 5-d, two powers of $\Omega$ can enter with the same size as a single power of $\Omega$
- Higher order terms are a few percent effect and can be subtracted
- Discussed by Christ and Sharpe at DWF@ 10 meeting at BNL

$$
\left\langle K^{+}\right| Q_{6}\left|\pi^{+}\right\rangle \sim\left\{\frac{(m+\Omega)}{a^{3}}+\frac{(m+\Omega)^{3}}{a^{3}}\right\}\left\langle K^{+}\right| \bar{s} d\left|\pi^{+}\right\rangle+\frac{\Omega}{a}\left\langle K^{+}\right| \bar{s} \sigma \cdot F d\left|\pi^{+}\right\rangle
$$

## Chiral Perturbation Theory and $\varepsilon^{\prime} / \varepsilon$

- Simulations will have a fixed, dynamical strange quark mass, which may be outside range of utility of NLO ChPT
- Lightest quark mass, $\mathrm{m}_{\mathrm{s}} / 10$, may need finite volume corrections added to ChPT formula.
- 2+1 flavor partially quenched ChPT being done by Aubin, Laiho and Li
- $(8,8)$ and $(27,1)$ operators complete. $(8,1)$ operators are well underway


## $\varepsilon^{\prime} / \varepsilon$ Summary

- Have summarized RBC-UKQCD calculation of NLO coefficients from K -> $\pi$ and K -> vacuum
- Work on $\mathrm{K}->\pi \pi$ at unphysical kinematics underway
- Lee and Sharpe are using ASQTAD staggered fermions, with smearings, to calculate $\mathrm{B}_{\mathrm{K}}, \mathrm{B}_{7}$ and $\mathrm{B}_{8}$. Testing to see how much smearing can help with operator mixing
- Hernandez, et. al. are working in the epsilon regime (quenched) to explore $\Delta \mathrm{I}=1 / 2$ rule
- Lellouch, et. al. are using overlap fermions on lattices generated with 2 flavors of Wilson fermions to look at $\mathrm{B}_{\mathrm{K}}$ and $\varepsilon^{\prime} \varepsilon$.


## Conclusions

- 2+1 flavor DWF QCD simulations well underway
$(3 \mathrm{fm})^{3}$ volumes at two lattice scales
$m_{l}=m_{s} / 5$ on $a^{-1}=1.6 \mathrm{GeV}$ lattices
$m_{l}=m_{s} / 7$ on $a^{-l}=2.1 \mathrm{GeV}$ lattices
- $\Delta \mathrm{S}=1$ matrix elements appear to be benefitting from large spatial volume, giving reduced statistical errors.
- From comparison of ChPT to data, we are investigating range of pseudoscalar masses where NLO ChPT is accurate to, say $10 \%$.
- For $\varepsilon^{\prime} / \varepsilon$, NLO fits for $\Delta \mathrm{I}=3 / 2$ amplitudes should work.
- For $\Delta \mathrm{I}=1 / 2$ amplitudes, statistical errors will likely limit NLO fits
- $\mathrm{K}->\pi \pi$, tests underway to get needed constant.
- Major systematic in final result likely ChPT
- Multiparticle final states a few years away...

