

KAON'07

Laboratori Nazionali di Frascati dell'INFN

May 21-25, 2007

Rare Kaon Decays beyond the Standard Model

- A Brief Introduction (Formalism, NP Models, Main Features of $K_L \rightarrow \pi^0 \nu \bar{\nu}$, $K^+ \rightarrow \pi^+ \nu \bar{\nu}$, $K_L \rightarrow \pi^0 e^+ e^-$, $K_L \rightarrow \pi^0 \mu^+ \mu^-$)
- NP Effects within Minimal Flavour Violation (MFV)
- NP Effects beyond MFV

Class of NP models with NO new sources of flavour-violation w.r.t. the SM (motivated by the SM success in Flavour Physics)

Cecilia Tarantino
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Master Formula for Weak Decays

A.J. Buras, hep-ph/0101336
hep-ph/0109197

(General) Minimal Flavour Violation

(MFV) G.D'Ambrosio et al., hep-ph/0207036

Constrained Minimal Flavour Violation

(CMFV) A.J.Buras et al., hep-ph/0007085

$$\begin{aligned}
 \mathbf{A}(\text{Decay}) = & \mathbf{V}_{\text{CKM}}^i \left\{ \mathbf{B}_i \eta_{\text{QCD}}^i \left[\mathbf{F}_{\text{SM}}^i + \mathbf{F}_{\text{CMFV}}^i \right] + \mathbf{B}_i^{\text{MFV}} \left(\eta_{\text{QCD}}^i \right)^{\text{MFV}} \mathbf{F}_{\text{MFV}}^i \right\} \\
 & + \mathbf{V}_{\text{NEW}}^i \mathbf{B}_i^{\text{NEW}} \left(\eta_{\text{QCD}}^i \right)^{\text{NEW}} \mathbf{F}_{\text{NEW}}^i
 \end{aligned}$$

} Most General NP
(beyond MFV)

\mathbf{F}_i^i 's: short-distance Loop Functions (Penguins, Boxes) [from perturbation theory]

(η_{QCD}^i) 's: QCD corrections [from RG improved perturbation theory]

\mathbf{B}_i 's: long-distance Parameters [the most uncertain: from experiments, Lattice, ...]

CMFV: (pragmatic approach) a CMFV Model has to satisfy two constraints:
the only source of flavour-violation is the CKM, the only operators are the SM ones

[A.J. Buras et al., hep-ph/0007085]

MFV: (top-down approach) in building a MFV Model the SM Yukawa couplings
are the only *building-blocks* of flavour violation [G.D'Ambrosio et al., hep-ph/0207036]

General NP beyond MFV: new sources of flavour violation ($\mathbf{V}_{\text{NEW}}^i$) can appear

The Most Appealing Models

CMFV

Littlest Higgs (LH) [N.Arkani-Hamed, A.G.Cohen, E.Katz, A.E.Nelson, hep-ph/0206021]

(Higgs= Goldstone boson of a global symmetry $SU(5) \xrightarrow{f} SO(5)$,

“little hierarchy problem” cured by heavy partners with the same spin-statistics

(W_H, Z_H, A_H, T, Φ)

f is constrained to be large ($>2-3$ TeV) by electroweak (ew) precision tests)

Appelquist-Cheng-Dobrescu (ACD) [T.Appelquist, H.C.Cheng, B.A.Dobrescu, hep-ph/0012100]

(1 universal extra-dimension \rightarrow SM particles propagate in 5 dimensions,

in 4 dimensions Kaluza-Klein modes appear, KK—Parity conservation,

only one additional free parameter: $1/R > 200-300$ GeV from ew precision tests)

(General) MFV

MSSM [P.Fayet (1976,1977,1979), G.R.Farrar, P.Fayet(1978),...] with MFV

(boson-fermion symmetry, minimal spectrum (sparticles, 2 Higgs doublets, $\tan\beta=v_u/v_d$), R-parity)

• low $\tan\beta \rightarrow$ **CMFV** (true only if μ is small) [W.Altmannshofer, A.J.Buras, D.Guadagnoli, hep-ph/0703200]

• large $\tan\beta \rightarrow$ new operators are generated by Higgs contributions

Beyond MFV

MSSM [P.Fayet (1976,1977,1979), G.R.Farrar, P.Fayet(1978),...] with a general Flavour Structure
(boson-fermion symmetry, minimal spectrum (sparticles, 2 Higgs doublets),
R-parity,
27 new Flavour Changing couplings along the squark lines)

Littlest Higgs with T-Parity (LHT) [H.C.Cheng, I.Low, hep-ph/0308199, hep-ph/0405243]
(LH + a discrete symmetry (T-Parity)
→heavy particles cannot appear at tree-level
→ew precision tests are less constraining ($f > 500\text{GeV}$),
in addition to W_H, Z_H, A_H, T, Φ ,
new mirror fermions with new flavour interactions)

Minimal 3-3-1 Model [P.H.Frampton(1992), F.Pisano, V.Pleitez(1992), R.Foot et al.(1993)]
(extended gauge group $SU(3)_c \times SU(3)_L \times U(1)_X$, 3 generations required for anomaly
cancellation and QCD asymptotic freedom, new heavy quarks and gauge bosons,
**FCNC's can be transmitted at tree-level by the new neutral
gauge boson Z' and involve a new mixing matrix)**

Main Features of Rare Kaon Decays See talk by Uli for details and complete bibliography

$$\mathbf{K_L \rightarrow \pi^0 \nu \bar{\nu} \text{ and } K^+ \rightarrow \pi^+ \nu \bar{\nu}}$$

Golden modes [intrinsic theoretical uncertainty <5%]

- Short-distance dominated
- Operator matrix elements extracted from $K/3 \longrightarrow$ Beyond LO in ChPT

F.Mescia and C.Smith, hep-ph/0705.2025 (last week!)

Strongly suppressed within the SM

K_L

$$K_L \cong \frac{1}{\sqrt{2}} (K^0 + \bar{K}^0)$$

$$\langle \pi^0 | (\bar{d}s)_{V-A} | \bar{K}^0 \rangle = -\langle \pi^0 | (\bar{d}s)_{V-A} | K^0 \rangle$$

$$A(K_L \rightarrow \pi^0 \nu \bar{\nu}) \propto (\lambda_t - \lambda_t^*) X(x_t) + (\lambda_c - \lambda_c^*) X(x_c)$$

$$\cong \text{Im}(\lambda_t) X(x_t) \rightarrow \mathbf{O(\lambda^5)}$$

$X(x) \approx x$
 $\text{Im}(\lambda_t) \approx \text{Im}(\lambda_c) \approx \lambda^5$

K^+

$$A(K^+ \rightarrow \pi^+ \nu \bar{\nu}) : \text{Im}(\lambda_t) X(x_t), \text{Re}(\lambda_t) X(x_t) \rightarrow \mathbf{O(\lambda^5)}$$

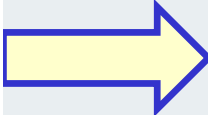
and $\text{Re}(\lambda_c) X(x_c)$ suppressed as well by x_c

\longrightarrow At NNLO

A.J.Buras et al., hep-ph/0508165

\cong
 λ + long-distance and dimension-8 operator effects

G.Isidori, F.Mescia, C.Smith, hep-ph/0503107



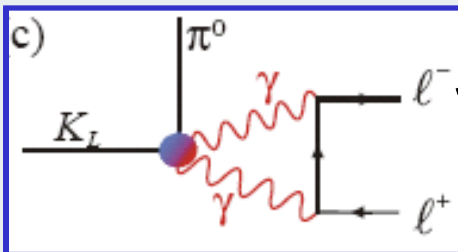
highly sensitive to NP, where theoretical cleanness is not spoiled!

$$\mathbf{K_L \rightarrow \pi^0 e^+ e^- \text{ and } K_L \rightarrow \pi^0 \mu^+ \mu^-}$$

Promising modes

Three contributions of comparable size:

- a) direct CP-violating (short-distance \rightarrow clean),
- b) indirect CP-violating (\leftrightarrow exp. observation of $K_S \rightarrow \pi^0 l^+ l^-$),
- c) long-distance CP-conserving (\leftrightarrow exp. studies of $K_L \rightarrow \pi^0 \gamma \gamma$)



Recently, theoretically
under control

G.Buchalla et al., hep-ph/0308008,
G.Isidori et al., hep-ph/0404127

The two-photon CP-conserving contribution $K_L \rightarrow \pi^0 (\gamma \gamma)_{J=0,2} \rightarrow \pi^0 (l^+ l^-)_{J=0,2}$ represents the main difference between muon and electron channels: $J^{CP}=0^{++}$ is helicity suppressed and can be neglected in the electron mode

Different NP effects can be disentangled by measuring both rates

F.Mescia, C.Smith, S.Trine, hep-ph/0606081

See talk by Christopher

SM Predictions and Experimental Measurements

Golden Modes	SM	Experiment
$K^+ \rightarrow \pi^+ \nu \bar{\nu}$	$(7.8 \pm 0.8) \cdot 10^{-11}$ F.Mescia and C.Smith, hep-ph/0705.2025	$(14.7^{+13}_{-8.9}) \cdot 10^{-11}$ E787 E949
$K_L \rightarrow \pi^0 \nu \bar{\nu}$	$(2.5 \pm 0.4) \cdot 10^{-11}$ F.Mescia and C.Smith, hep-ph/0705.2025	$< 2.1 \cdot 10^{-7}$ E391a
$K_L \rightarrow \pi^0 e^+ e^-$	$(3.5^{+1.0}_{-0.9}) \cdot 10^{-11}$ F.Mescia et al. hep-ph/0606081	$< 2.8 \cdot 10^{-10}$ KTeV
$K_L \rightarrow \pi^0 \mu^+ \mu^-$	$(1.4 \pm 0.3) \cdot 10^{-11}$ F.Mescia et al. hep-ph/0606081	$< 3.8 \cdot 10^{-10}$ KTeV

Very clean

**Still a lot of
room for NP**



**Rare Kaon Decays
within **Constrained-**
and **General-MFV****

Imposing available information from the Universal Unitary Triangle Analysis and from the measurements of $\text{Br}(B \rightarrow X_s \gamma)$, $\text{Br}(B \rightarrow X_s t \bar{t})$ and $\text{Br}(K^+ \rightarrow \pi^+ \nu \bar{\nu})$

- Determination of the allowed range for the short-distance function $C(\text{Z-penguin})$
- Evaluation of the upper bounds for Rare Decays

Branching Ratios	CMFV (95%)	SM (95%)	SM (68%)	Exp
$\text{Br}(K^+ \rightarrow \pi^+ \nu \bar{\nu}) \cdot 10^{11}$	<11.9	<10.9	8.3 ± 1.2	$14.7^{+13.0}_{-8.9}$
$\text{Br}(K_L \rightarrow \pi^0 \nu \bar{\nu}) \cdot 10^{11}$	<4.6	<4.2	3.1 ± 0.6	$<5.9 \cdot 10^4$

NP effects in box and gluon-penguin diagrams are assumed negligible
 $\chi^{\text{NP}} \approx C^{\text{NP}}$

2007 update by U.Haisch and A.Weiler, including constraints from R_b^0 , A_b , $A_{\text{FB}}^{0,b}$:

- The negative (non-SM) sign of C is now excluded
- Br's are also bounded from below

• Existing constraints forbid significant enhancements (<30%)
 • Rare Kaon Decays are potentially the most constraining
 [Next Decade Scenario with ~5% exp. uncertainty is looked forward]

Branching Ratios	CMFV (95%)	SM (95%)	Exp.
$\text{Br}(K^+ \rightarrow \pi^+ \nu \bar{\nu}) \cdot 10^{11}$	[4.3, 10.7]	[5.4, 9.1]	$14.7^{+13.0}_{-8.9}$
$\text{Br}(K_L \rightarrow \pi^0 \nu \bar{\nu}) \cdot 10^{11}$	[1.6, 4.4]	[2.2, 3.5]	$<2.1 \cdot 10^4$

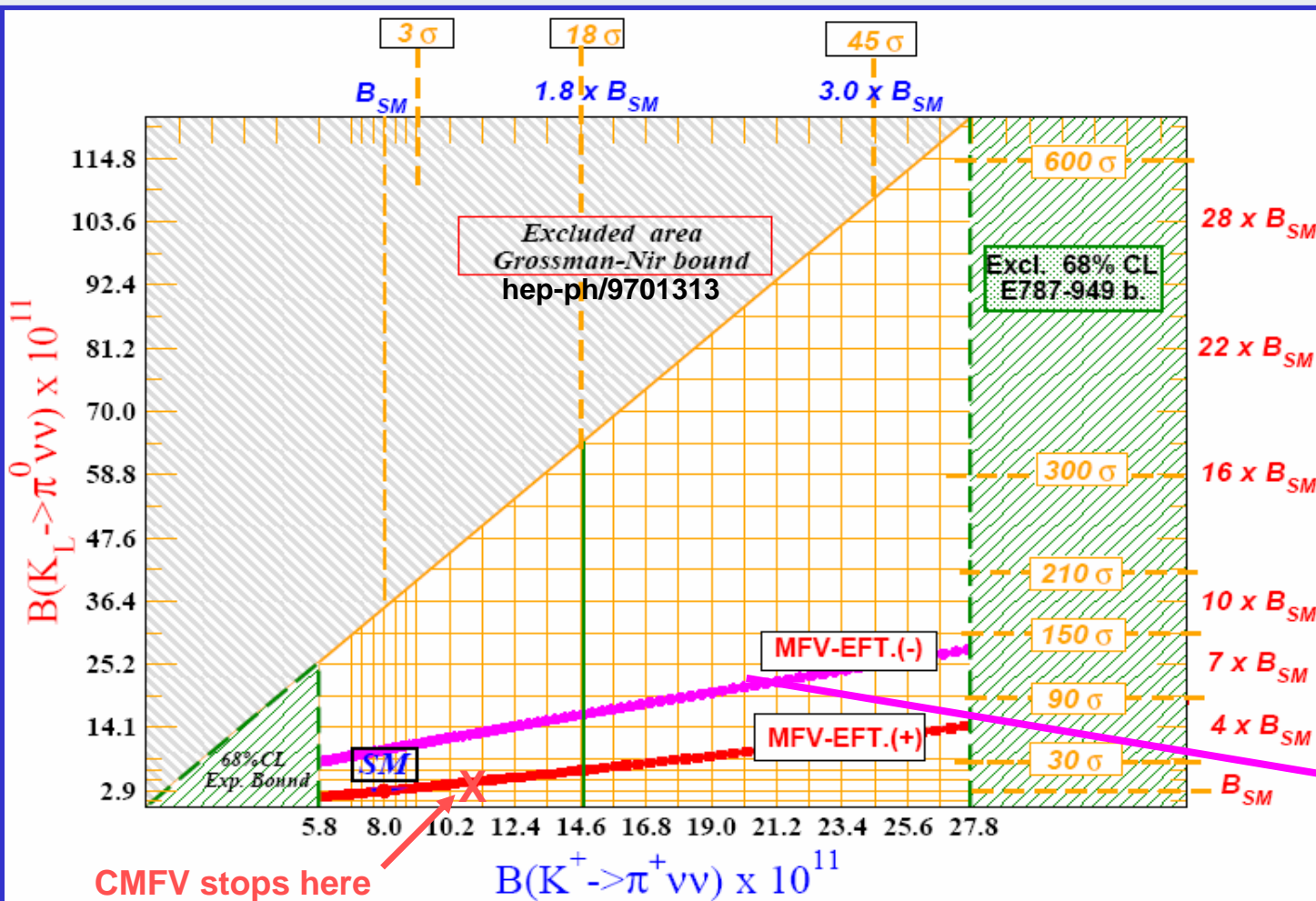
$$\left. \frac{\text{Br}(K_L \rightarrow \pi^0 \nu \bar{\nu})}{\text{Br}(K^+ \rightarrow \pi^+ \nu \bar{\nu})} \right|_{\text{CMFV}} \approx \left. \frac{\text{Br}(K_L \rightarrow \pi^0 \nu \bar{\nu})}{\text{Br}(K^+ \rightarrow \pi^+ \nu \bar{\nu})} \right|_{\text{SM}}$$

In agreement with the MFV prediction:

G.D'Ambrosio et al., hep-ph/0207036
A.J.Buras, R.Fleischer, hep-ph/0104238

$$\left. \frac{\text{Br}(K_L \rightarrow \pi^0 \nu \bar{\nu})}{\text{Br}(K^+ \rightarrow \pi^+ \nu \bar{\nu})} \right|_{\text{MFV}} = \left. \frac{\text{Br}(K_L \rightarrow \pi^0 \nu \bar{\nu})}{\text{Br}(K^+ \rightarrow \pi^+ \nu \bar{\nu})} \right|_{\text{SM}} \times [1 + \epsilon_c \text{sign}(X)]$$

Federico Mescia [CKM06]



The "minus" sign branch is now ruled out within CMFV [2007 update by U.Haisch, A.Weiler]

In Specific CMFV Models:

Littlest Higgs (LH) [without T-parity]

Ew precision tests constrain the NP scale to be large: $f > 2\text{-}3 \text{ TeV}$
→ The enhancements of $\text{Br}(K \rightarrow \pi \nu \bar{\nu})$ amount to at most 15%

A.J.Buras et al., hep-ph/0607189

Appelquist-Cheng-Dobrescu [1 universal extra-dimension]

- For $1/R=200 \text{ GeV}$, Br's are enhanced by 16% (K^+) and 17% (K_L)
- For $1/R=300 \text{ GeV}$ (from ew tests if $m_H < 250 \text{ GeV}$) the enhancements are twice smaller

A.J.Buras, M.Spranger, A.Weiler, hep-ph/0212143

• Even smaller enhancements would be found with the new NNLO SM prediction for $b \rightarrow s \gamma$ by M.Misiak et al., hep-ph/0609232

MSSM with MFV and moderate (<30) $\tan\beta$

Y.Nir, M.P.Worah, hep-ph/9711215

A.J.Buras, A.Romanino, L.Silvestrini, hep-ph/9712398

G.Colangelo, G.Isidori, hep-ph/9808487

The deviations from the SM in the two $\text{Br}(K \rightarrow \pi \nu \bar{\nu})$ are naturally small ($\sim 10\%$) but could saturate the model-independent bound [G.Isidori et al., hep-ph/0604074]

See talk by Christopher

- At large $\tan\beta$, charged-Higgs loops yield the new operator $(\bar{s}d)_{V+A} (\bar{\nu}\nu)_{V-A}$ not suppressed by $m_{d,s}$
- However, its impact on rare K decays is significant only beyond MFV [G.Isidori, P.Paradisi, hep-ph/0601094]

What about $K_L \rightarrow \pi^0 \ell^+ \ell^-$ within CMFV?

- In general, models without new operators (like CMFV models) induce **smaller effects** (<10%) for $K_L \rightarrow \pi^0 \ell^+ \ell^-$ than for $K \rightarrow \pi \nu \bar{\nu}$
- Still, **NP effects** in vector (Q_{7V}) and axial-vector (Q_{7A}) can be **disentangled** by measuring both electron and muon modes



Rare Kaon Decays
beyond MFV

MSSM with a general Flavour Structure

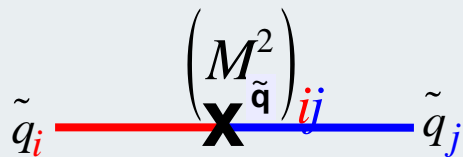
To obtain simple analytic formulae for $s \rightarrow d$ transitions, it is useful to work in a basis where $d_L^i - \tilde{d}_L^j - N_n$ and $d_L^i - \tilde{u}_L^j - \chi_n$ are flavour diagonal and to expand (non-diagonal) sfermion mass matrices

Mass Insertion Approximation

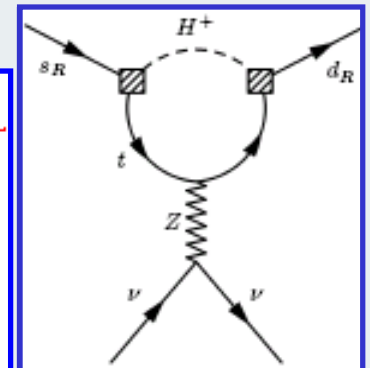
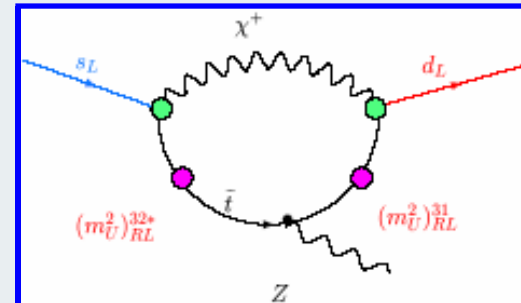
$$M_{\tilde{u}}^2 = \begin{pmatrix} (m_U^2)_{LL} & (m_U^2)_{LR} \\ (m_U^2)_{LR}^\dagger & (m_U^2)_{RR} \end{pmatrix}$$

3x3 non-diagonal flavour matrices expanded in small off-diagonal entries: e.g., $(\delta_{LL}^U)_{ij} \equiv (m_U^2)_{ij}^{LL} / \tilde{m}^2$

Flavour non-diagonality is brought by squark propagators



- Y.Nir, M.P.Worah, hep-ph/9711215
- A.J.Buras, A.Romanino, L.Silvestrini, hep-ph/9712398
- G.Colangelo, G.Isidori, hep-ph/9808487
- A.J.Buras et al., hep-ph/9908371
- A.J.Buras et al., hep-ph/0408142
- G.Isidori, P.Paradisi, hep-ph/0601094



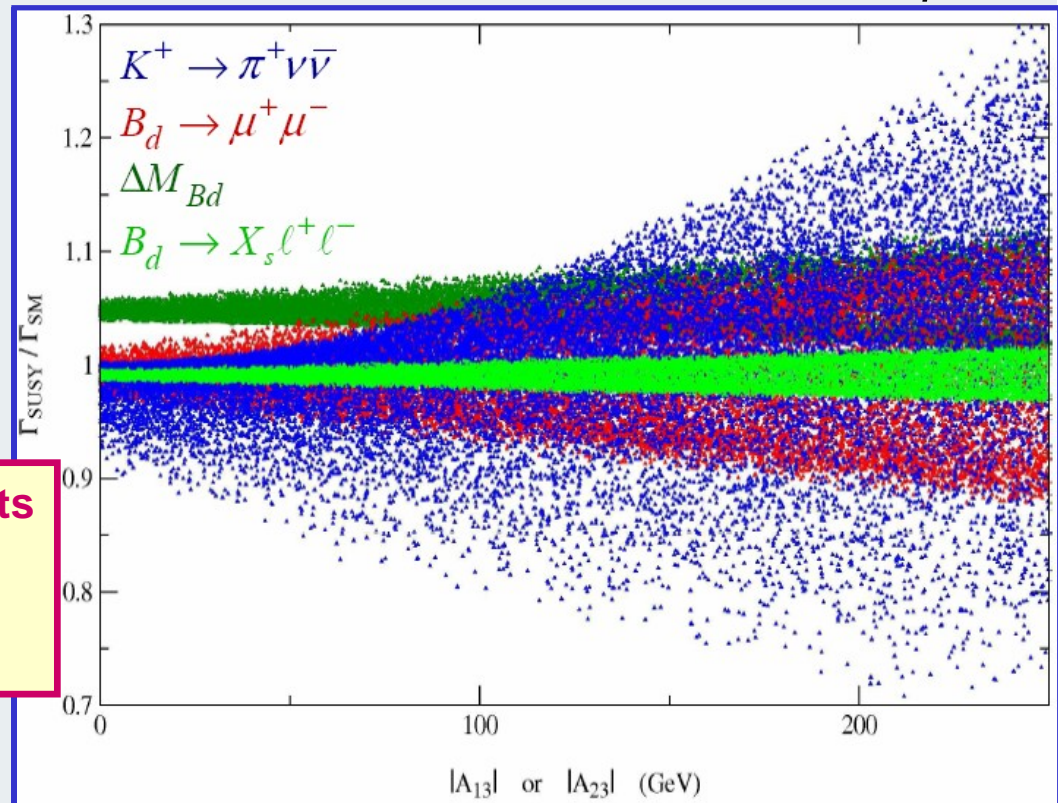
The dominant NP effects in rare kaon decays come from **chargino up-squark loops** (non-MFV terms in the LR up sector, double ins. not suppressed by $O(M_W/M_{SUSY})$ nor CKM) and **charged-Higgs top-quark loops** (non-MFV terms in the RR down sector, not suppressed by $m_{d,s}$ and enhanced by $\tan^4\beta$)

$\tan\beta < 30$

δ_{LR}^u mass insertions
are not strongly constrained
by other B- and K-observables

$$\delta_{LR}^u \leftrightarrow \left(M_u^2 \right)_{iL3R} \approx m_t A_{i3}$$

- $K \rightarrow \pi \nu \bar{\nu}$ would be the best measurements to determine/constrain δ_{LR}^u ($\leftrightarrow A_{i3}$)
- Large ($O(1)$) departures from the SM are allowed in $K \rightarrow \pi \nu \bar{\nu}$



G.Isidori, P.Paradisi, hep-ph/0601094

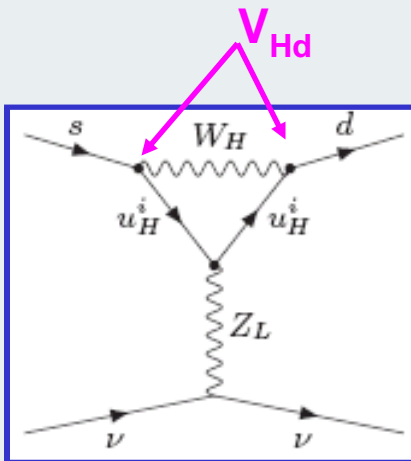
- Precise measurements of $K \rightarrow \pi \nu \bar{\nu}$ rates could provide the most stringent bounds on charged Higgs contributions (δ_{RR}^d)

What about $K_L \rightarrow \pi^0 \ell^+ \ell^-$?

- At large $\tan\beta$ (pseudo)-scalar operators $(\bar{s}d)(\bar{\ell}\ell)$ and $(\bar{s}d)(\bar{\ell}\gamma_5 \ell)$ are enhanced
- Being helicity suppressed, they affect only the muon mode and can lead to a clear signal

Littlest Higgs Model with T-Parity (LHT)

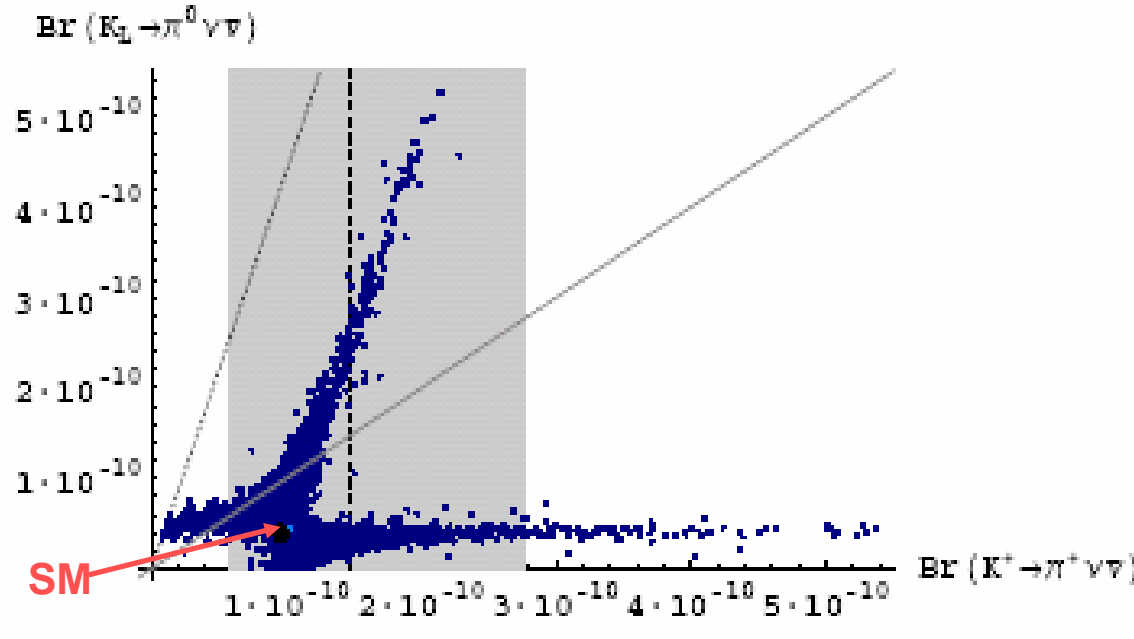
NP scale
 $\approx 1\text{TeV}$



- Mirror fermions introduce new flavour interactions governed by the new mixing matrix V_{Hd}
- Large effects are expected, in particular in rare K decays where the CKM λ_t suppression has in general no analogy in V_{Hd}

Blanke, Buras, Poschenrieder, Recksiegel, C.T., Uhlig, Weiler, hep-ph/0610298

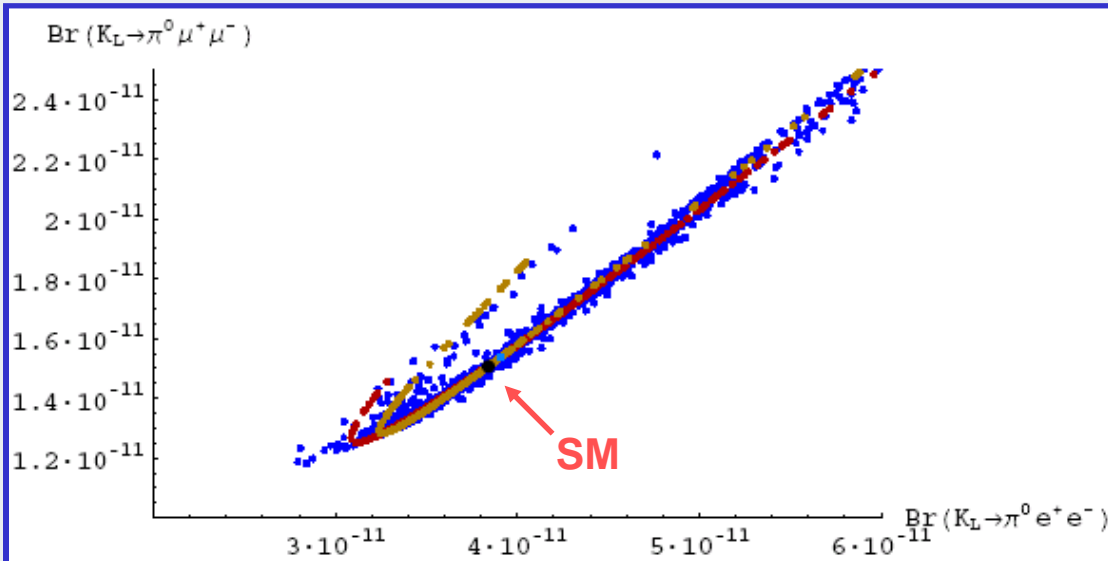
K_L vs. K^+



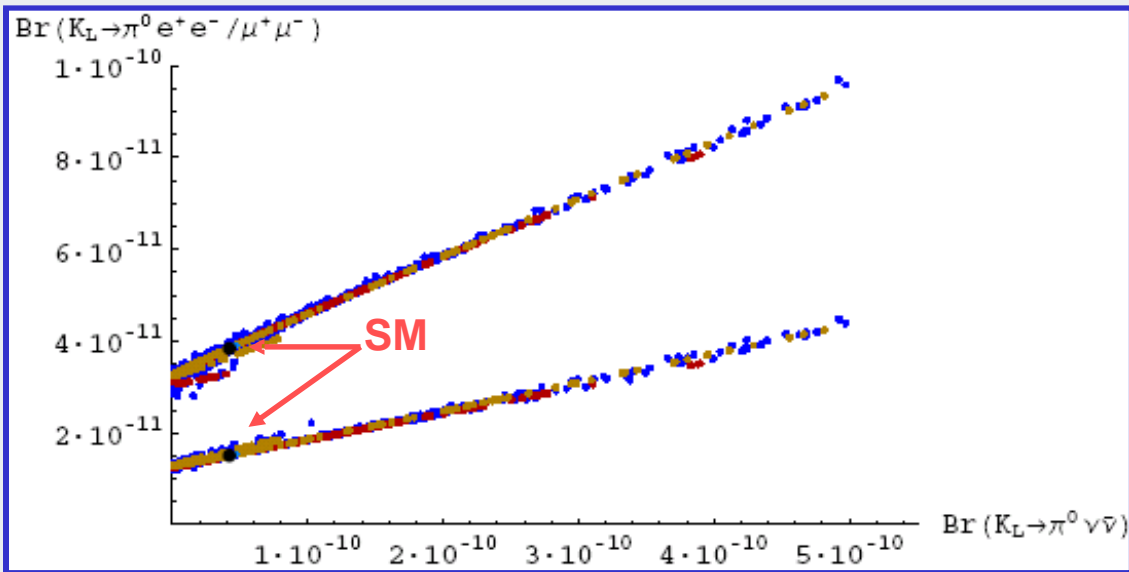
The strongest constraint comes from ε_K , which selects two branches:

- $\text{Br}(K^+)$ can be enhanced by a factor 5 with $\text{Br}(K_L) \sim \text{Br}(K_L)_{\text{SM}}$
- $\text{Br}(K_L)$ can be enhanced by a factor 10 with $\text{Br}(K^+) < 2-3 \cdot \text{Br}(K^+)_{\text{SM}}$

$K_L \rightarrow \pi^0 e^+ e^-$ and $K_L \rightarrow \pi^0 \mu^+ \mu^-$



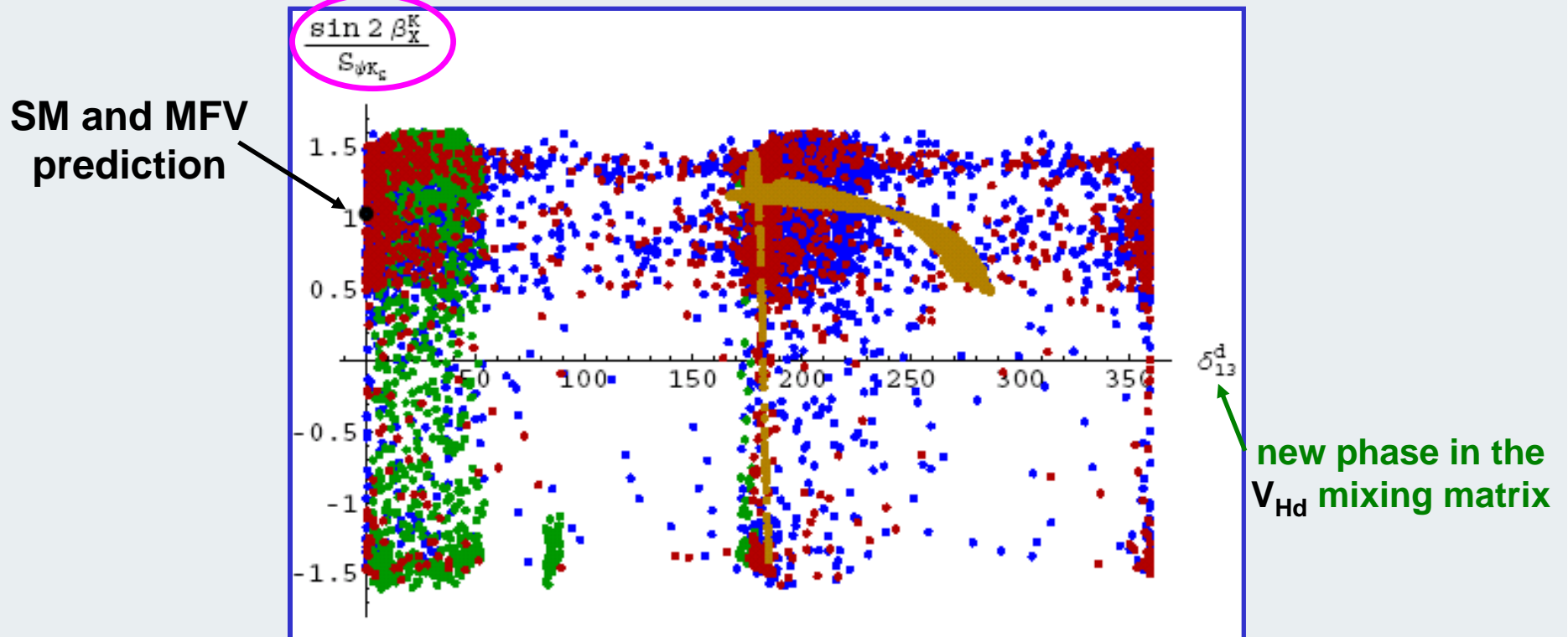
Blanke, Buras, Poschenrieder, Recksiegel, C.T., Uhlig, Weiler, hep-ph/0610298



- Strong correlations between muon and electron modes and with $K_L \rightarrow \pi^0 \nu \bar{\nu}$
- Enhancements up to a factor 2 w.r.t. the SM

- Different values of β from $S_{\psi K_S}$ and $K_L \rightarrow \pi^0 \nu \bar{\nu}$ would be a clear signal of NP beyond MFV
- In LHT, the universality between B and K systems can be strongly violated

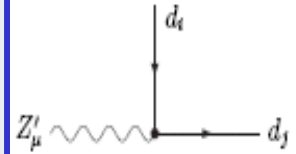
Blanke,Buras,Poschenrieder,Recksiegel,C.T.,Uhlig,Weiler, hep-ph/0610298



Minimal 3-3-1 Model

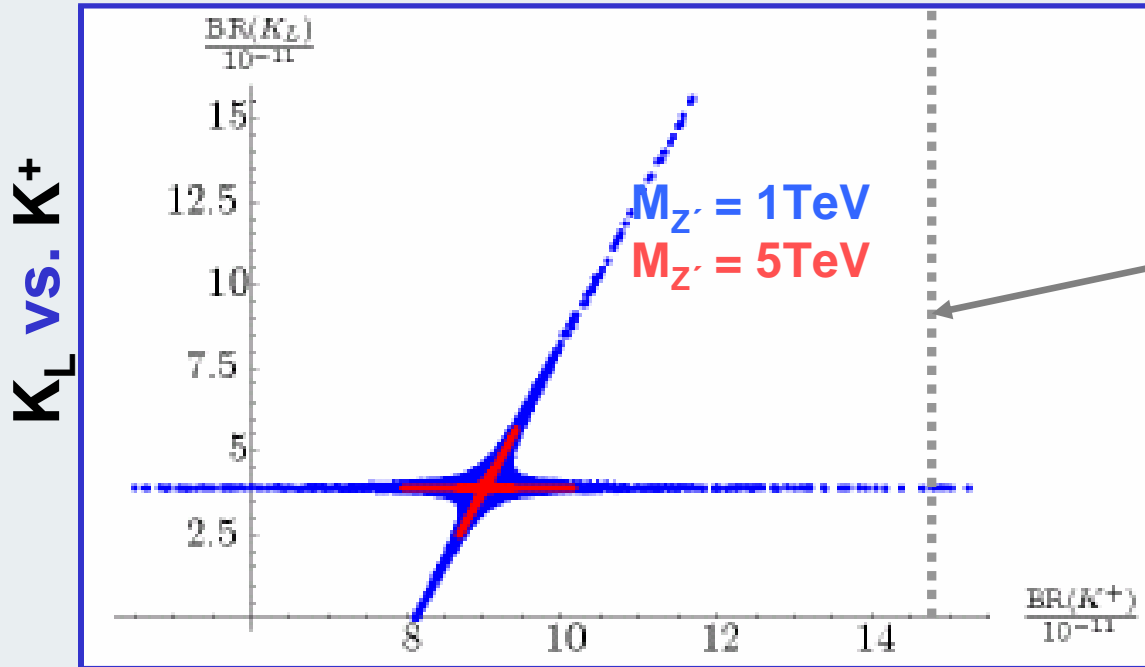
$M_{Z'} \approx 1-5 \text{ TeV}$

FCNC's can be transmitted at tree-level by the new heavy gauge boson Z' and involve a new mixing matrix



$$i \frac{g_C w}{\sqrt{3} \sqrt{1 - 4s_w^2}} V_{3d_i}^* V_{3d_j} P_L \quad i, j = 1, 2, 3, i \neq j$$

C.Promberger, S.Schatt, F.Schwab, hep-ph/0702169



The present experimental central value for $\text{Br}(K^+)$ can be reached only for a rather low $M_{Z'} \approx 1 \text{ TeV}$

$$K_L \rightarrow \pi^0 e^+ e^- \text{ and } K_L \rightarrow \pi^0 \mu^+ \mu^-$$

- The 2-branch pattern is similar to LHT
- Due to the leptophobic nature of the $\bar{l}lZ'$ coupling (suppressed by $\sqrt{1 - 4s_w^2}$), smaller enhancements, w.r.t. LHT, are allowed: by a factor 2 (K^+) and 4 (K_L)

• Strong correlations between muon and electron modes and with $K_L \rightarrow \pi^0 \nu \bar{\nu}$

• Enhancements up to a factor 3 w.r.t. the SM are allowed

Role of ε'/ε in constraining rare kaon decays

direct CP-violation
in the K-system

$$\frac{\varepsilon'}{\varepsilon} = (16.7 \pm 1.6) \cdot 10^{-3}$$

NA48, KTeV

- ε'/ε and rare kaon decays are strongly correlated as both involve the short-distance functions X_K, Y_K, Z_K
- ε'/ε is unfortunately affected by a very large theoretical uncertainty (non-perturbative matrix element of the gluon-penguin operator $\leftrightarrow R_6$)

See talks by Chris Sachrajda and Bob Mawhinney

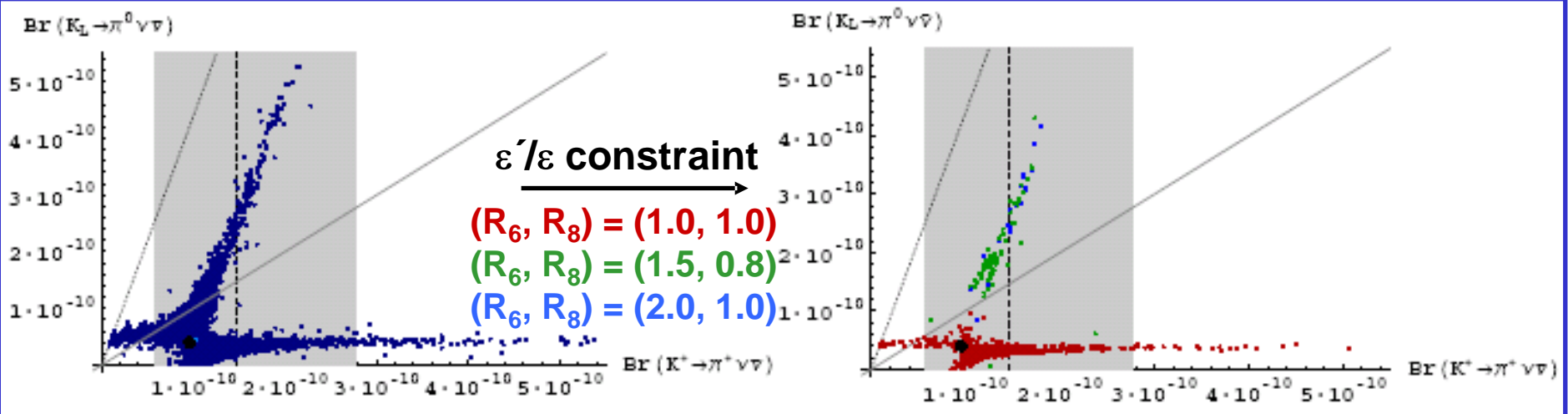
MSSM beyond MFV

A.J.Buras, L.Silvestrini, hep-ph/9811471; A.J.Buras et al., hep-ph/9908371

- Data on ε'/ε may constraint considerably the double LR insertions and the corresponding enhancement of rare kaon decays
- Significant enhancements are still possible

LHT

M.Blanke, A.J.Buras, S.Recksiegel, C.T., S.Uhlig, hep-ph/0704.3329



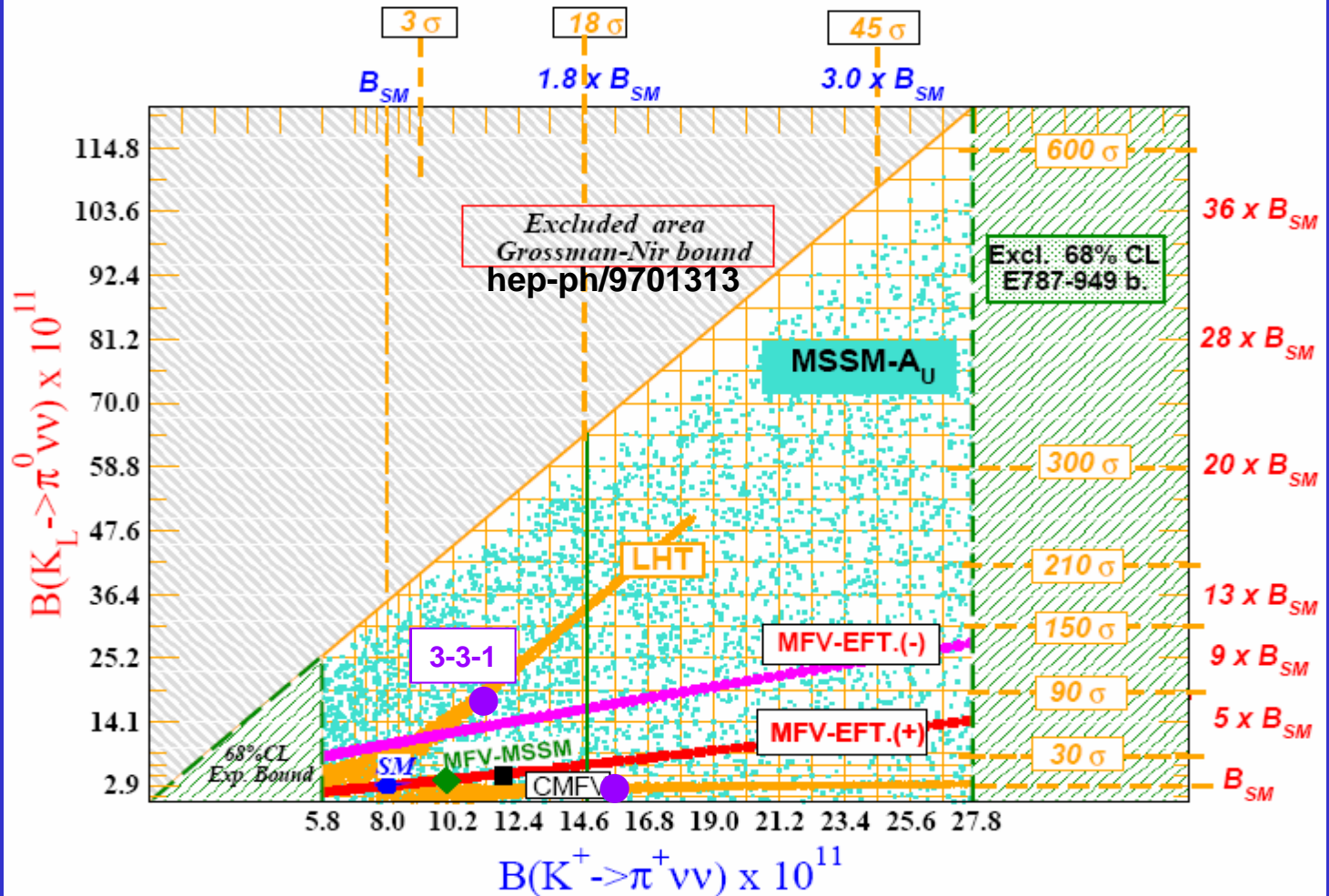
- With an improved theoretical accuracy, ε'/ε could strongly constrain rare K decays
- If the large-N values $(R_6, R_8) = (1.0, 1.0)$ are confirmed, O(10)-enhancements of K_L are excluded

Conclusive Messages

$$\begin{aligned} K_L &\rightarrow \pi^0 \nu \bar{\nu} \\ K^+ &\rightarrow \pi^+ \nu \bar{\nu} \end{aligned}$$

Rare K decays are excellent probes of NP, for their theoretical cleanness and strong suppression within the SM

Federico Mescia [CKM06]



- Large NP signals could be seen in rare K decays (without large effects in B-systems)
- In particular, in MSSM or LHT, K_L can be enhanced by an order of magnitude
- Specific correlations could help in discriminating NP models

$$K_L \rightarrow \pi^0 e^+ e^- \text{ and } K_L \rightarrow \pi^0 \mu^+ \mu^-$$

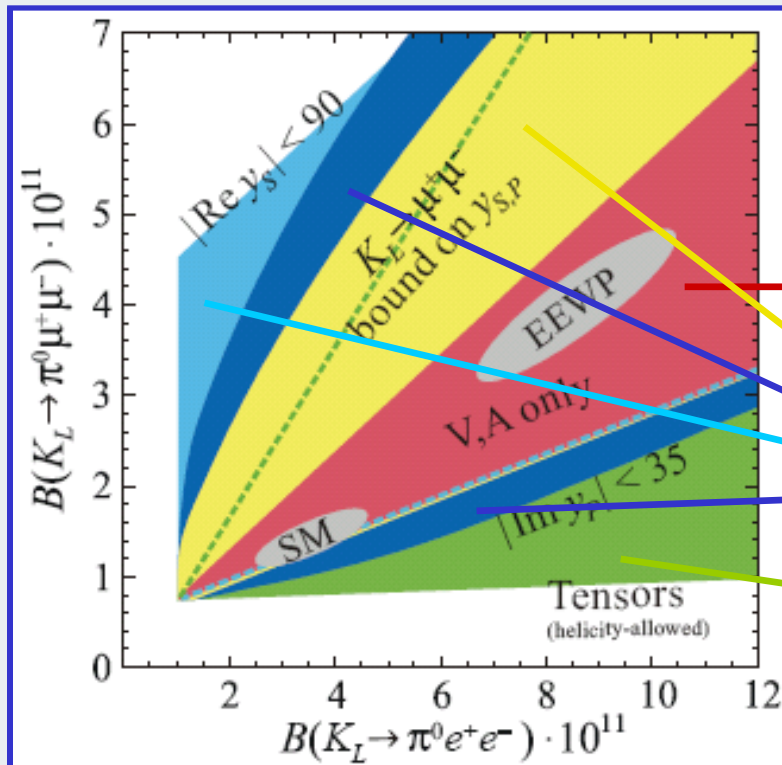
- The recent theoretical control over the long-distance components makes $K_L \rightarrow \pi^0 \ell^+ \ell^-$ decays very clean and highly sensitive to NP

[G.D'Ambrosio et al., hep-ph/9808289, G.Buchalla et al., hep-ph/0308008, G.Isidori et al., hep-ph/0404127]

- Different μ and e masses allow to probe helicity-suppressed effects
- NP enhancements by a factor 2-5 are possible (MSSM, LHT, 3-3-1)

F.Mescia, C.Smith, S.Trine, hep-ph/0606081

Model-independent analysis of the impact of $\Delta S=1$ four-fermion operators $(\bar{s} \Gamma d)(\bar{\ell} \Gamma \ell)$



NP Models with no new operators
(MSSM at low $\tan\beta$ [and small μ], LHT, 3-3-1)

NP Models with new (helicity-suppressed) operators (MSSM at large $\tan\beta$)

Other NP Models not discussed in this talk
(e.g. with tree-level leptoquark interactions)

See next talk, by Christopher...

Rare K Decays are unique probes of NP



↓
Accurate experimental
measurements are
looked forward!

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