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Quantum and Classical Dynamics with Random Permutation Circuits

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School of Physics and Astronomy
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Motivation: quantum many-body dynamics and thermalization

not simple in eigenbasis of H
short-range correlated

Quantum quench: $|\Psi_0\rangle \rightarrow |\Psi_t\rangle = e^{-it\overbrace{H}^{\text{local}}} |\Psi_0\rangle$

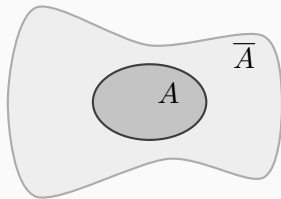
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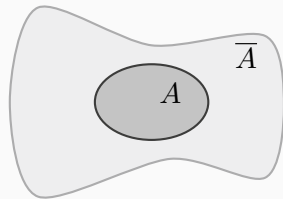
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Many subtleties: conservation laws, localisation, scars, ...

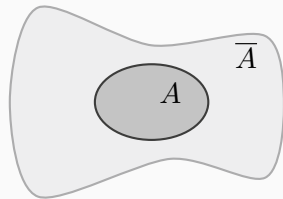
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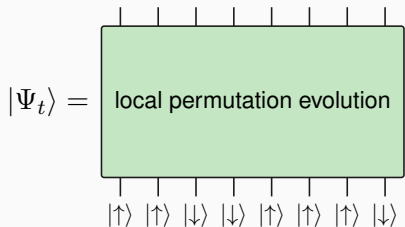
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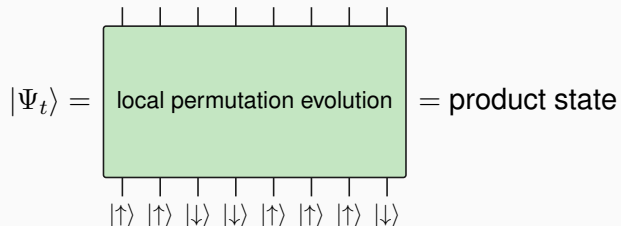
Q: What happens under classical deterministic dynamics?

$\Leftarrow \exists$ a basis : $|\underline{s}\rangle \rightarrow |\underline{s}'\rangle$

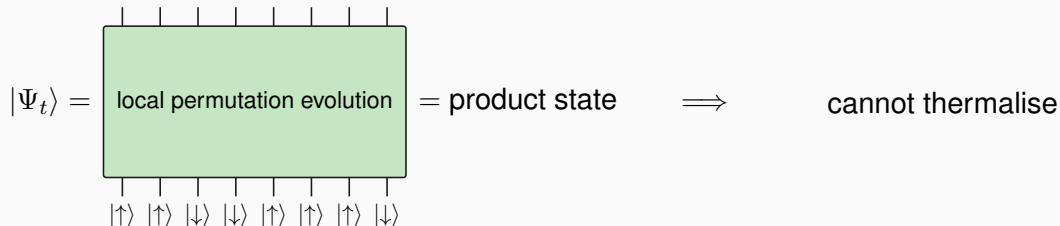
Classical deterministic dynamics



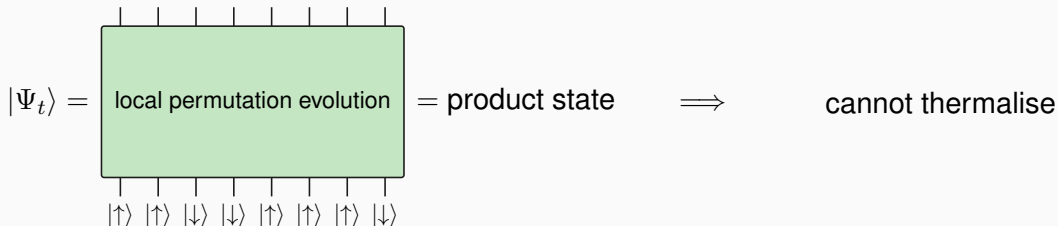
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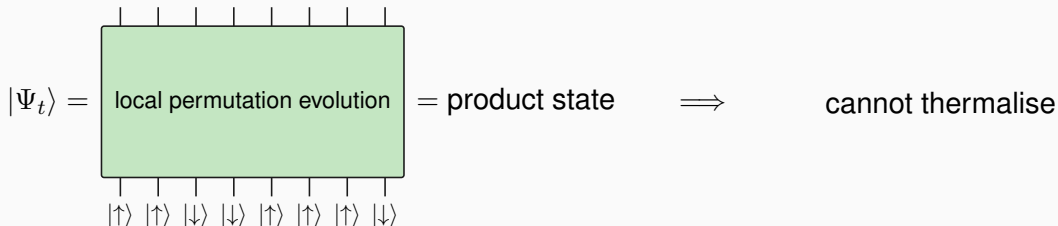


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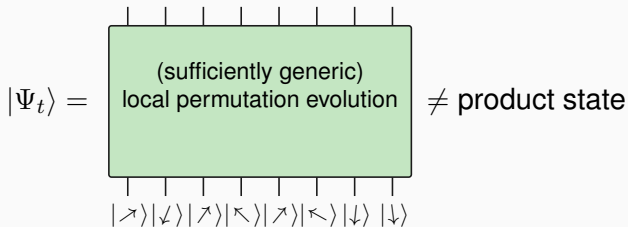


This statement is basis dependent!

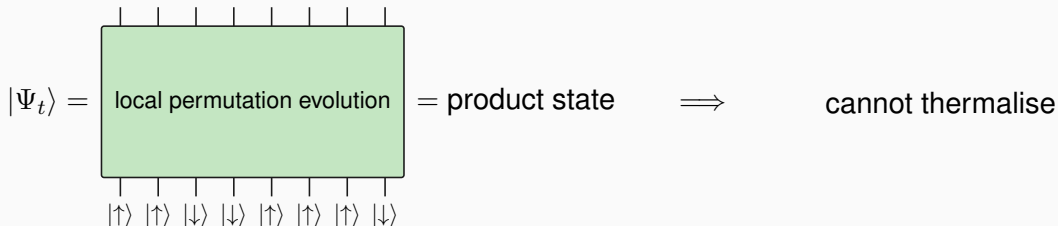
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Is there a simple way to distinguish it from generic quantum dynamics?

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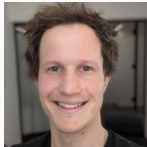
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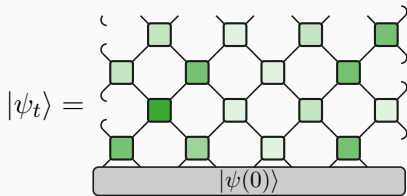
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In analogy with random-unitary circuits:

chain of L qudits with locally applied randomly chosen permutations



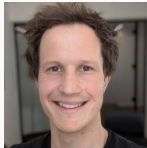
$$\text{green square} \in S(d^2)$$



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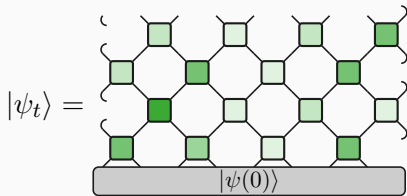
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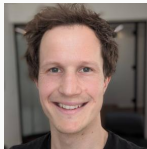
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How do the averaged quantities (e.g. OTOCs and purity) behave?

Technical tool: Averages over the permutation group I

Basic building block:

$$\mathcal{Q}^{(m)} := \frac{1}{d^2!} \sum_{U \in S(d^2)} \underbrace{U \otimes U \otimes U \otimes \dots \otimes U}_m = \text{[Diagram: a green square with 'm' inside, four lines extending from the corners]}$$

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Detour: Uniform averages over $d \times d$ unitary matrices

$$\text{[Diagram: A red square with a diagonal line from top-left to bottom-right, labeled '2n' inside, with four lines extending from the corners.]} := \int_{U(d)} dU \underbrace{U \otimes U^* \otimes U \otimes \dots \otimes U^*}_{2n}$$

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- Symmetry under reshuffling of replicas: $|\sigma\rangle = \sum_{s_1, \dots, s_n=0}^{d-1} |s_1 s_{\sigma(1)} s_2 s_{\sigma(2)} \dots s_n s_{\sigma(n)}\rangle \quad \forall \sigma \in S(n)$

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Invariant states given by **partitions** of $\mathbb{Z}_m$

$$\pi = \left\{ \{i_{1,1}, i_{1,2}, \dots, i_{1,l_1}\}, \dots, \{i_{k,1}, \dots, i_{k,l_k}\} \right\}$$

$$\sum_{j=1}^k l_j = m, \quad \bigcup_{j=1}^k \{i_{j,1}, \dots, i_{j,l_j}\} = \mathbb{Z}_m$$

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Example: $2n = m = 4$

$B_4 = 15$ partition states (and $2! = 2$ permutation states)

$$\pi_1 = \{\{0, 1, 2, 3\}\}$$

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$$\pi_3 = \{\{0, 1, 3\}, \{2\}\}$$

$$\pi_4 = \{\{0, 2, 3\}, \{1\}\}$$

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$$|\bigcirc\rangle = \frac{1}{d} \sum_{a,b=0}^{d-1} |aabb\rangle$$

$$|\square\rangle = \frac{1}{d} \sum_{a,b=0}^{d-1} |abba\rangle$$

$$|\times\rangle = \frac{1}{\sqrt{d}} \sum_{a=0}^{d-1} |aaaa\rangle$$

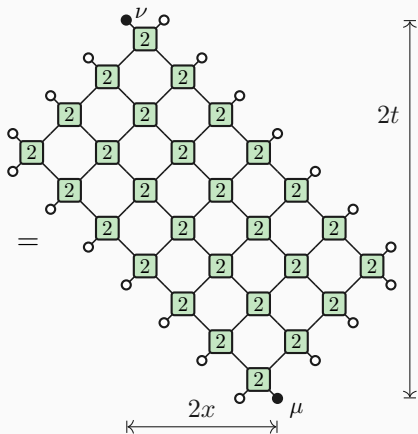
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Averaged infinite-temperature two-point correlation function I

$$C_{\mu\nu}(x, t) = \frac{1}{\text{tr} \mathbb{1}} \langle \text{tr}[\mathcal{O}_\mu(x, t) \mathcal{O}_\nu(0, 0)] \rangle_{\mathcal{Q}}$$

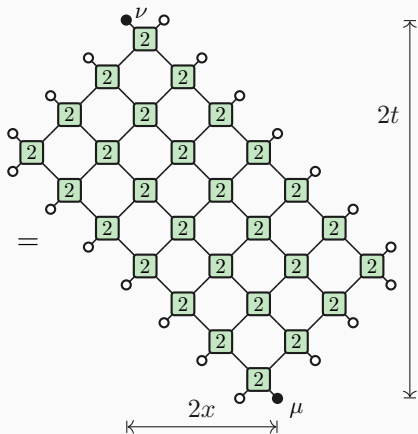
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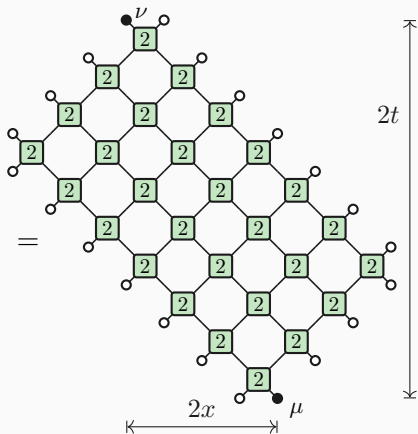
$$\boxed{2} = Q^{(2)}$$

$$|\mathcal{O}\rangle = \frac{1}{\sqrt{d}} \sum_{s=0}^{d-1} |ss\rangle \propto |\pi_1\rangle$$

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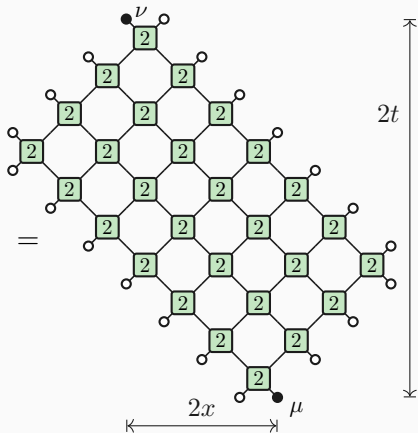
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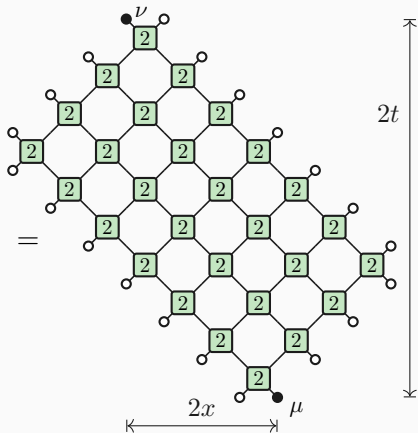
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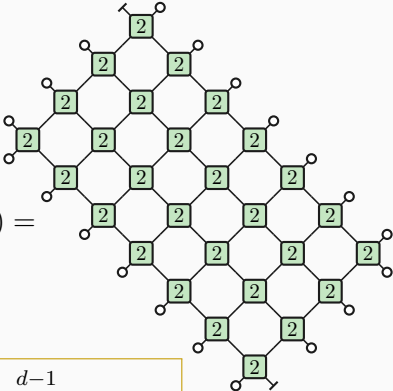
First step: project $|\bullet_{\mu,\nu}\rangle$ to $\{|\mathcal{O}\rangle, |-\rangle\}$

Averaged infinite-temperature two-point correlation function II

$$C_{\mu\nu}(x, t) = \frac{o_\mu o_\nu}{(1 - 1/d)^2} \left[C_{--}(x, t) - \frac{1}{d} \right]$$

Averaged infinite-temperature two-point correlation function II

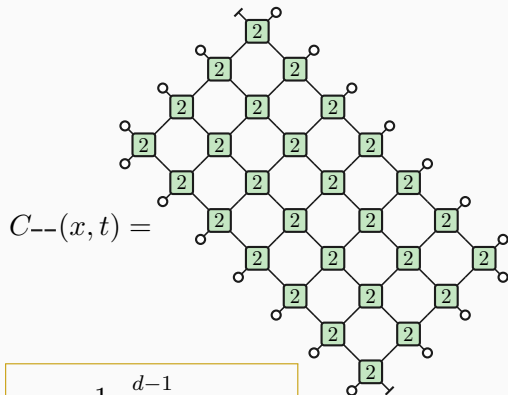
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$$C_{--}(x, t) =$$


$$o_\mu = \frac{1}{d^{\frac{3}{2}}} \sum_{i,j=0}^{d-1} \langle i | \mathcal{O}_\mu | j \rangle$$

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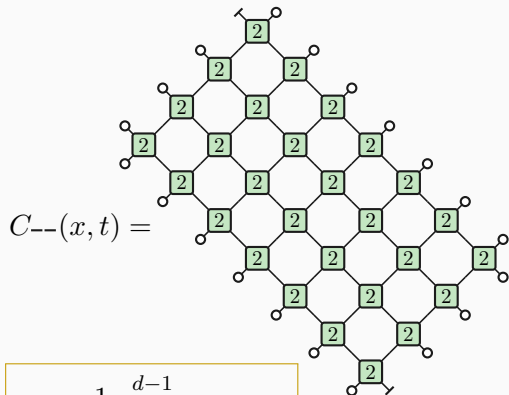


$$o_\mu = \frac{1}{d^{\frac{3}{2}}} \sum_{i,j=0}^{d-1} \langle i | \mathcal{O}_\mu | j \rangle$$

$$\begin{aligned} \text{Square with 2 top legs and 2 bottom legs, each with a small circle} &= \text{Two vertical lines} \\ \text{Square with 2 top legs and 2 bottom legs, each with a small circle} &= \text{Two vertical lines} \\ \text{Square with 2 top legs and 2 bottom legs, each with a small circle} &= \frac{\sqrt{d}}{1+d} (\text{Two vertical lines} + \text{Two vertical lines}) \end{aligned}$$

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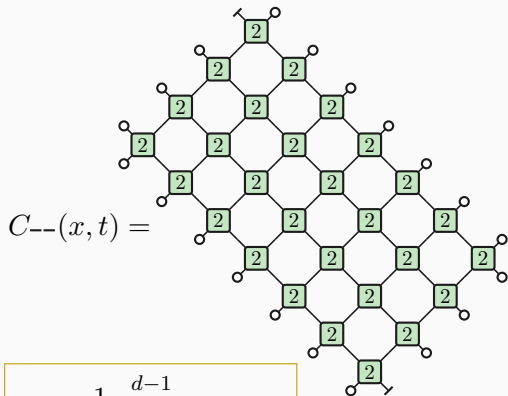
Diagrammatic equations for the two-point correlation function. The first equation shows a green square with '2' and two external lines (one horizontal, one vertical) equal to two separate vertical lines. The second equation shows a green square with '2' and two external lines (one horizontal, one vertical) equal to two separate horizontal lines. The third equation shows a green square with '2' and two external lines (one horizontal, one vertical) equal to the sum of two diagrams: one with two vertical lines and one with two horizontal lines, multiplied by a factor of $\frac{\sqrt{d}}{1+d}$.

Analogous to $2n = 4$ for random unitaries:

$$\frac{C_{\mu\nu}(x, t)}{o_\mu o_\nu} \Big|_{d \rightarrow d^2} = \left\langle \left(\frac{\text{tr}[O_\mu(x, t) O_\nu(0, 0)]}{\text{tr} \mathbb{1}} \right)^2 \right\rangle_{\text{U}}$$

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Diagrammatic rules for the 2-point correlation function:

- A green square with '2' and two external lines (one incoming, one outgoing) is equal to two separate vertical lines: $\text{green square with 2 and two lines} = \text{two vertical lines}$.
- A green square with '2' and two external lines (one incoming, one outgoing) is equal to two separate vertical lines: $\text{green square with 2 and two lines} = \text{two vertical lines}$.
- A green square with '2' and two external lines (one incoming, one outgoing) is equal to a green square with '2' and two external lines (one incoming, one outgoing) multiplied by $\frac{\sqrt{d}}{1+d}$ times the sum of two vertical lines: $\text{green square with 2 and two lines} = \text{green square with 2 and two lines} \cdot \frac{\sqrt{d}}{1+d} (\text{two vertical lines} + \text{two vertical lines})$.

Analogous to $2n = 4$ for random unitaries:

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Asymptotic scaling:

$$C_{\mu\nu}(x, t) \simeq \frac{o_\mu o_\nu c(d)}{t^2} \left(\frac{4/d}{(1 + 1/d)^2} \right)^{2t} e^{-2x^2/t}$$

Out-of-time-ordered correlation functions I

$$O_{\mu\nu}(x, t) = \frac{1}{2} \frac{\left\langle \text{tr} \left[\left| [\mathcal{O}_\mu(x, t), \mathcal{O}_\nu(0, 0)] \right|^2 \right] \right\rangle_{\mathbb{Q}}}{\text{tr } \mathbb{1}}$$

Out-of-time-ordered correlation functions I

$$O_{\mu\nu}(x, t) = \underbrace{\frac{\langle \text{tr}[\mathcal{O}_\mu^2(x, t) \mathcal{O}_\nu^2(0, 0)] \rangle_Q}{\text{tr} \mathbb{1}}}_{\tilde{O}_{\mu\nu}^{(1)}(x, t) = 1 + \text{subleading terms}} - \underbrace{\frac{\langle \text{tr}[\mathcal{O}_\mu(x, t) \mathcal{O}_\nu(0, 0) \mathcal{O}_\mu(x, t) \mathcal{O}_\nu(0, 0)] \rangle_Q}{\text{tr} \mathbb{1}}}_{\tilde{O}_{\mu\nu}^{(2)}(x, t)}$$

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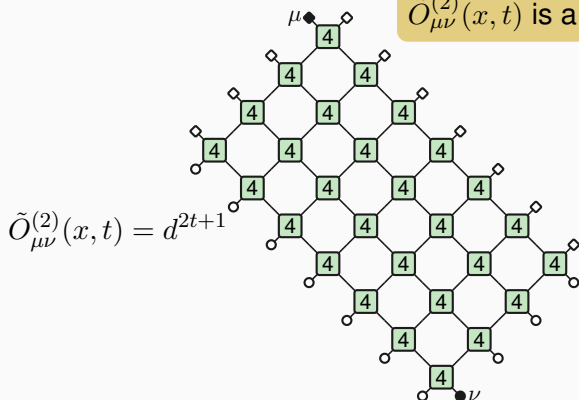
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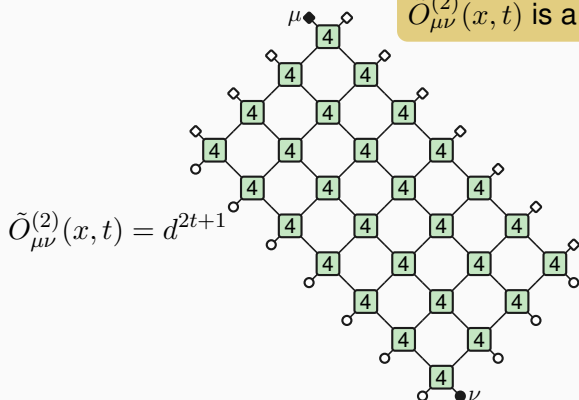
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$B_4 = 15$ states!



$$\tilde{O}_{\mu\nu}^{(2)}(x,t) = d^{2t+1}$$

$$|\bigcirc\rangle = \frac{1}{d} \sum_{a,b=0}^{d-1} |aabb\rangle$$

$$|\square\rangle = \frac{1}{d} \sum_{a,b=0}^{d-1} |abba\rangle$$

$$|\times\rangle = \frac{1}{\sqrt{d}} \sum_{a=0}^{d-1} |aaaa\rangle$$

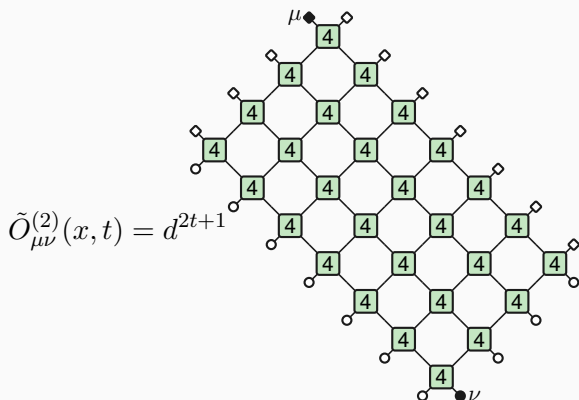
$$|-\rangle = \frac{1}{d^2} \sum_{a,b,c,e=0}^{d-1} |abce\rangle$$

$$|\bullet_\nu\rangle = (\mathcal{O}_\nu \otimes \mathbb{1})^{\otimes 2} |\bigcirc\rangle$$

$$|\blacksquare_\mu\rangle = (\mathcal{O}_\mu \otimes \mathbb{1})^{\otimes 2} |\square\rangle$$

Out-of-time-ordered correlation functions II

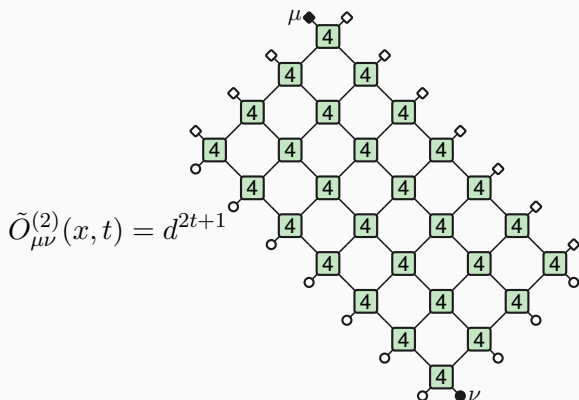
Problem: generic $|\blacksquare_\mu\rangle, |\bullet_\nu\rangle$ are projected to all 15 states $\Rightarrow \tilde{O}^{(2)}(x, t)$ hard to evaluate



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Additional structure for $\mathcal{O}_\nu \rightarrow \mathcal{O}_d$ diagonal

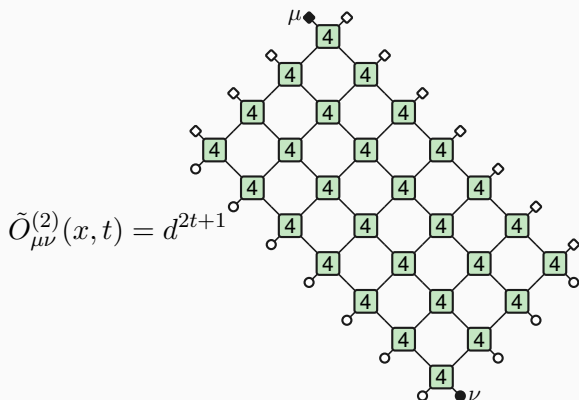


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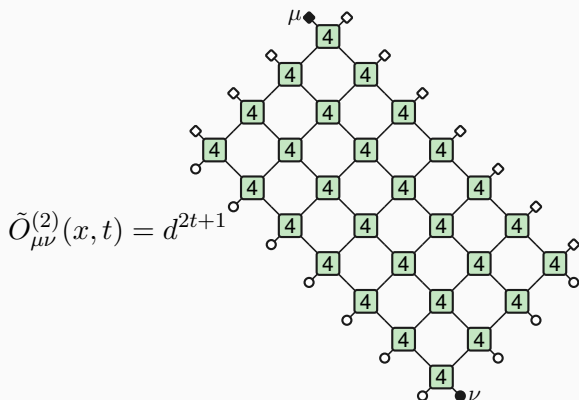
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Action of $\mathcal{Q}^{(4)}$ on these states closes:

$$\begin{aligned} \begin{array}{c} \diagup \diagdown \\ \text{4} \\ \diagdown \diagup \\ \circ \quad \circ \end{array} &= \begin{array}{c} | \\ | \\ \circ \quad \circ \end{array} & \begin{array}{c} \diagup \diagdown \\ \text{4} \\ \diagdown \diagup \\ \times \quad \times \end{array} &= \begin{array}{c} | \\ | \\ \times \quad \times \end{array} \\ \begin{array}{c} \diagup \diagdown \\ \text{4} \\ \diagdown \diagup \\ \times \quad \circ \end{array} &= \begin{array}{c} \diagup \diagdown \\ \text{4} \\ \diagdown \diagup \\ \circ \quad \times \end{array} &= \frac{\sqrt{d}}{1+d} \left(\begin{array}{c} | \\ | \\ \circ \quad \circ \end{array} + \begin{array}{c} | \\ | \\ \times \quad \times \end{array} \right) \end{aligned}$$



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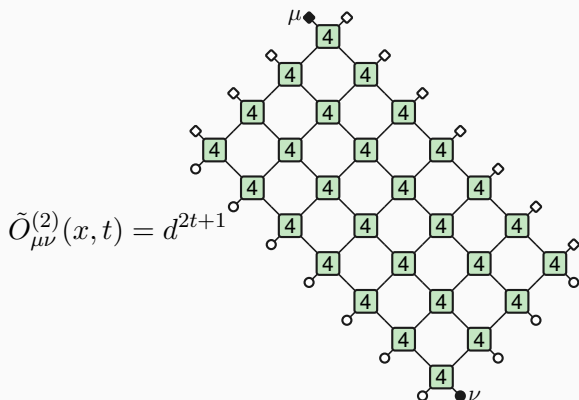
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We can obtain the closed-form expression



Out-of-time-ordered correlation functions III

$$C_{\mu\nu}(x, t) = \frac{d^2 o_\mu o_\nu}{d-1} f(x, t) \quad O_{\mu d}(x, t) = \frac{1 - o_{\mu,2}}{d-1} [f(x, t)]_{d \rightarrow \frac{1}{d}}$$

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$$O^{(\text{RU})}(x, t) = O_{\mu d}(x, t)|_{o_{\mu,2} \mapsto 0, d \mapsto d^2}$$

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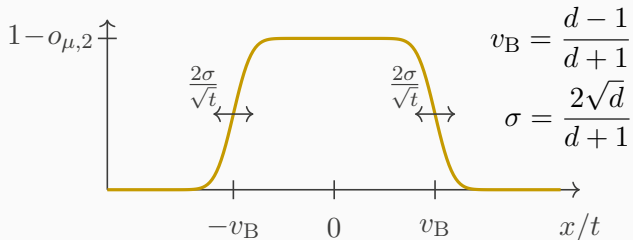
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Asymptotic scaling:



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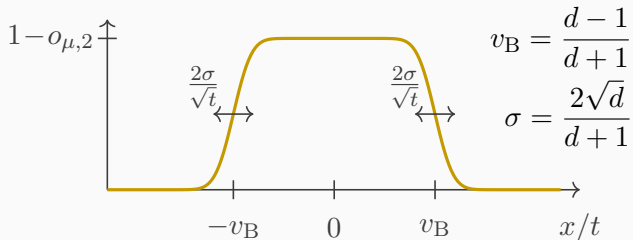
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Asymptotic scaling:



Non-generic value inside the light-cone!

Out-of-time-ordered correlation functions III

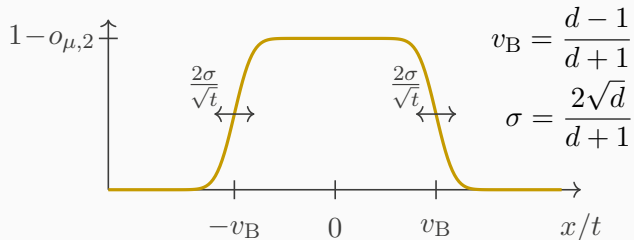
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Dynamics preserves diagonal operators

$$[\mathcal{O}_{d,1}(0, 0), \mathcal{O}_{d,2}(x, t)] = 0$$

Out-of-time-ordered correlation functions III

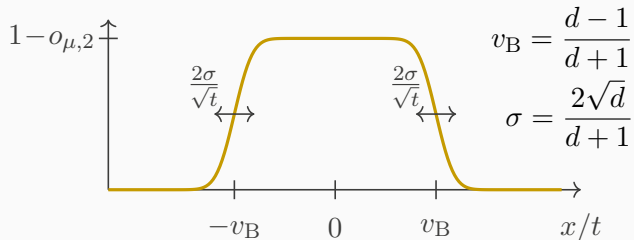
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Expectation: $O_{\mu\nu}(0, t) \rightarrow 1 - o_{\mu,2} o_{\nu,2}$

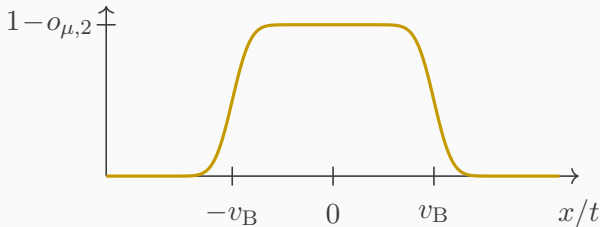
Summary and outlook

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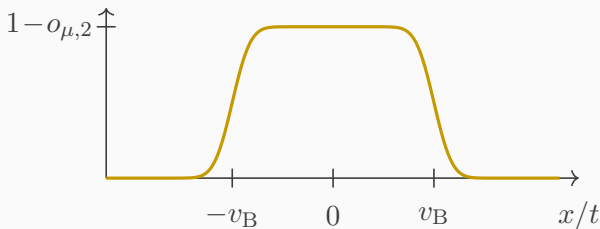
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Summary and outlook

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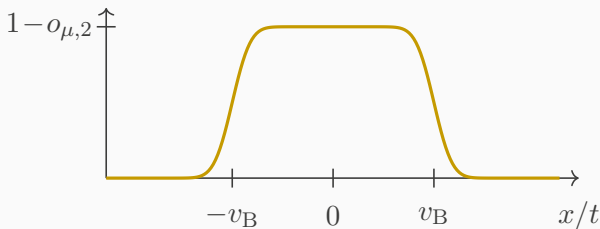
What about entanglement? ← cf. Michele's talk

D. Szász-Schagrín, M. Mazzoni, B. Bertini, KK, L. Piroli, 2025, arXiv:2505.06158

Summary and outlook

Motivating question: what are the generic features of many-body dynamics in the case of purely classical deterministic time-evolution?

- OTOCs have generic shape and *nontrivial* value
- Dependence on choice of local operators/states



What about entanglement? ← cf. Michele's talk

D. Szász-Schagrín, M. Mazzoni, B. Bertini, KK, L. Piroli, 2025, arXiv:2505.06158

Higher replica quantities: local operator entanglement

B. Bertini, KK, P. Kos, D. Malz, 2025, arXiv:2508.10890