

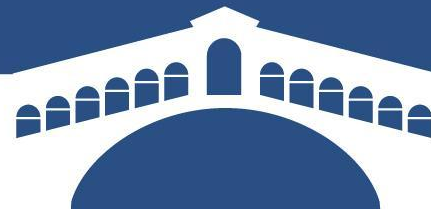


PARALLEL 1 / **ELECTROWEAK PHYSICS**

Electroweak inputs

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23-27 JUNE 2025 Lido di Venezia



Linear vs. Circular

- Z-pole ($\sqrt{s} \sim 90$ GeV)
- WW threshold ($\sqrt{s} \sim 160$ GeV)
- ZH threshold ($\sqrt{s} \sim 250$ GeV)



Circular

FCC-ee

- Huge luminosity
- 4IP
- absolute beam energy calibration up to 160 GeV
- 6×10^{12} Z (5 years)
- 2.4×10^8 WW (2 years)

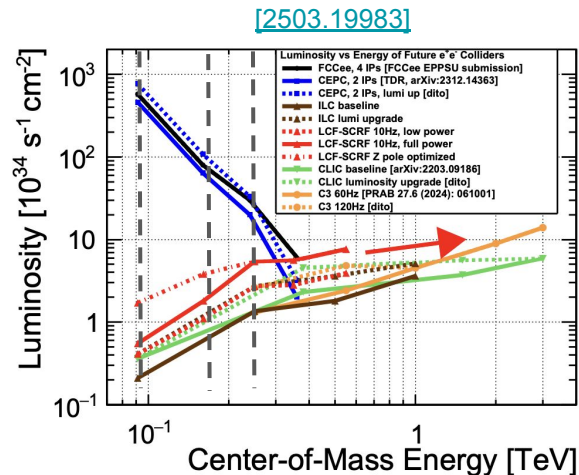
LEP3

- Luminosity = FCC-ee/3.5
- 2IP
- absolute beam energy calibration up to 90 GeV
- 2×10^{12} Z (6 years)
- 10^8 WW (4 years)

Linear

Linear Collider Facility (LCF)

- Luminosity / 1000
- 1,2 IP
- longitudinally polarised beams
- 5×10^9 Z (2 years)
- 5×10^7 WW (8 years)



Why EWK precision at intensity frontier?

- Potential to test the consistency of the SM via loop corrections

- $\sin^2\theta_W^{\text{eff}}$, m_W and other EWPOs parametrically depend on:

- $\alpha_{\text{QED}}(m_Z), m_{\text{top}}, m_Z, \alpha_S(m_Z), \dots$

- great precision is needed on these parameters to interpret potential deviations in terms of new physics
- impacts Higgs precision physics
- best to have many EWPOs to explore nature of BSM 'signals'

- Programmes

- Z lineshape:

- $m_Z, \Gamma_Z, A_f(\sin\theta_W^{\text{eff}}), R_\ell = \Gamma_{\text{had}}/\Gamma_\ell, R_q = \Gamma_{q\bar{q}}/\Gamma_{\text{had}}, \sigma_{\text{had}}^0$

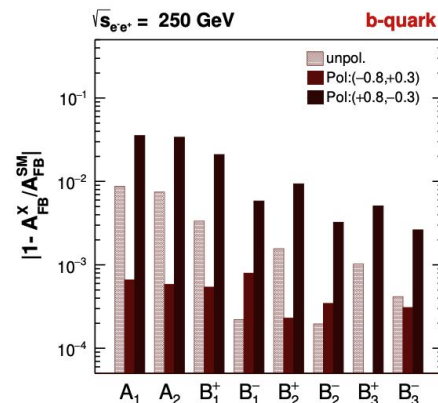
- $\alpha_{\text{QED}}(m_Z), \alpha_S(m_Z)$

- LCF: from lattice
- FCC-ee, LEP3: direct measurement + Lattice

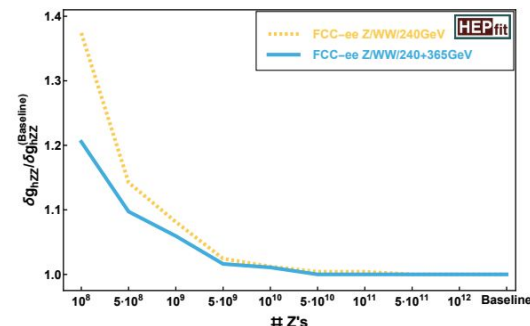
- WW threshold (and above)

- $m_W, \Gamma_W, \text{BR}(W \rightarrow \ell\nu), \alpha_S(m_Z), \text{triple/quartic gauge couplings}$

[2306.11413]



[2505.00272]



Caveats

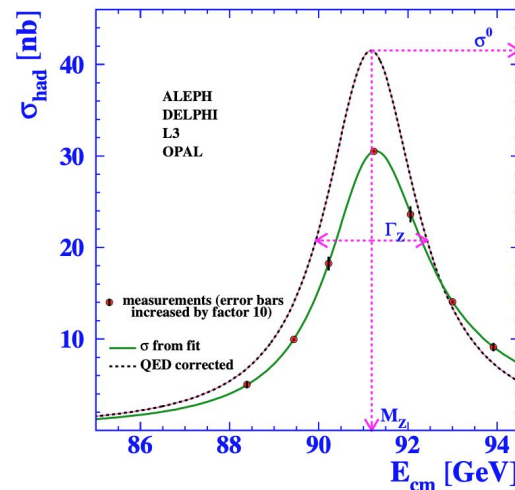
- **EWK factories will deliver $10^9 - 10^{12}$ Z**
 - vs. LEP/SLD 10-100 reduction in statistical uncertainties
 - statistical accuracies on some of the measurements down to 10^{-6}
 - systematic projections should be intended as targets for detectors/th. calculations
 - improvement of 10-1000X vs LEP
 - some will be challenging to reach, HARD work is needed
 - several systematics are de-facto of statistical nature
 - some of the projections are contentious and being challenged
- **Theory uncertainties and tools are NOT discussed here (see this afternoon for more in-depth discussion) lots of work to do!**
 - missing higher orders, multi-photon emission
 - MC modelling
 - NP QCD (hadronization) / parton showers
 - will be determined + heavily constrained by data

Z lineshape
(the most challenging)

Z mass / Width

- **Z Mass** is a parametric input to xsec, width, BR, ...
 - Current uncertainty $\Delta m_Z \sim 2 \text{ MeV}$ (LEP)
 - Statistical uncertainty scales as $\sim \Gamma_Z / 2\sqrt{N_Z^{\text{off-peak}}}$
 - 4 (7) keV at FCC-ee (LEP3)
 - 20 keV at LCF
 - Dominant systematic:
 - FCC-ee (LEP3): absolute beam energy calibration:
 - resonant depolarisation $\Delta\sqrt{s} \sim \Delta m_Z \sim 100 \text{ keV}$
 - achieved at LEP, ongoing effort to improve
 - LCF: absolute momentum scale
 - $\Delta p \sim \Delta m_Z \sim 200 \text{ keV}$ using J/ψ mass ($K_S \rightarrow \pi\pi$)
 - absolute limit 2 ppm (is 10 ppm more feasible?)
- **Z total width** is sensitive to fermion couplings and to BSM
 - Current $\Delta\Gamma_Z \sim 2 \text{ MeV}$
 - Dominant systematic:
 - relative absolute beam energy calibration
 - point-to-point $\Delta\sqrt{s}_{\text{p.t.p}}$ (uncorrelated component)
 - can be measured in-situ with $\mu\mu$ events
 - $\Delta\Gamma_Z \sim 12$ (25) (125) keV at FCC-ee (LEP3) (LCF)

[0509008]



$$\Gamma_Z = \Gamma_{ee} + \Gamma_{\mu\mu} + \Gamma_{\tau\tau} + \Gamma_{\text{had}} + \Gamma_{\text{inv}}.$$

$$\sigma_{\text{had}}^0 \equiv \frac{12\pi}{m_Z^2} \frac{\Gamma_{ee}\Gamma_{\text{had}}}{\Gamma_Z^2}$$

$$R_\ell^0 \equiv \Gamma_{\text{had}}/\Gamma_{\ell\ell}$$

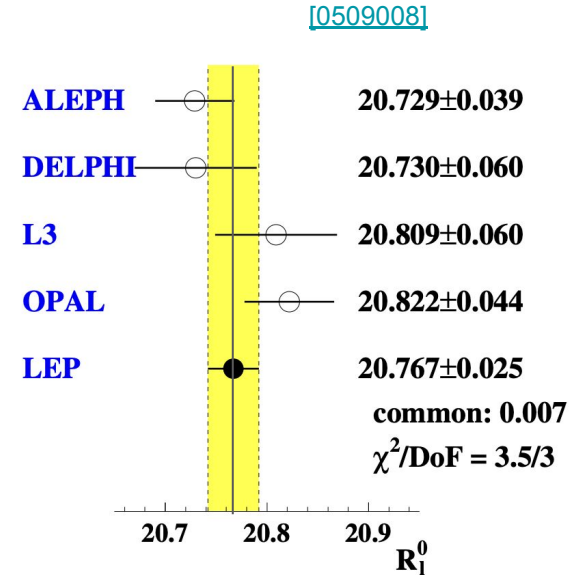
lepton universality:

$$R_{\text{inv}}^0 = \left(\frac{12\pi R_\ell^0}{\sigma_{\text{had}}^0 m_Z^2} \right)^{\frac{1}{2}} - R_\ell^0 - (3 + \delta_\tau)$$

$$R_{\text{inv}}^0 = N_\nu \left(\frac{\Gamma_{\nu\bar{\nu}}}{\Gamma_{\ell\ell}} \right)_{\text{SM}}$$

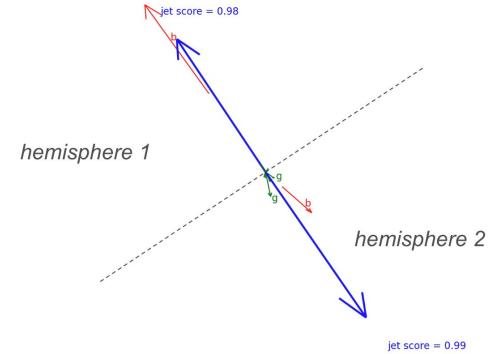
Leptonic Branching Ratios

- Defined as $R_\ell = \Gamma_{\text{had}}/\Gamma_\ell$ with $\ell = e, \mu, \tau$ measure individual axial-vector couplings to the Z
 - provides α_s measurement through radiation QCD corrections
- Current (relative) uncertainties $\sim 10^{-3}$
- Ratio, not affected to uncertainties luminosity, production cross-section etc ..
 - Dominant systematics:
 - acceptance determination at low angle for $e, \mu, (\tau)$ (negligible for hadronic modes)
 - can be translated into detector requirement:
 - absolute polar angles (and radial, longitudinal)
 - 5-10 μm at $10^\circ(20^\circ)$ to match stat. accuracy at FCC-ee (3×10^{-6})
 - 10-20 μm at $10^\circ(20^\circ)$ to match stat. accuracy at LEP3 (6×10^{-6})
 - easily achievable at LCF (10^{-4})
 - Low angle Bhabha subtraction increases uncertainty for $e+e$ -channel

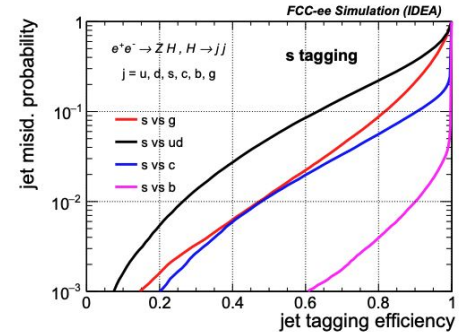


Hadronic Branching Ratios

- Defined as $R_q = \Gamma_q / \Gamma_{\text{had}}$ with $q = b, c, s$ measure individual chiral couplings to the Z
 $\rightarrow (\sim g_L^2 + g_R^2)$
- Current (relative) uncertainties $\sim 10^{-3}$
- Dramatic improvements compared to LEP are expected, driven by:
 - Reduced beam-spot sizes, light beam-pipe
 - Light and precise vertex detectors (few μm single point resolution)
 - Particle ID allowing strange tagging (K+ identification up to 30-40 GeV)
 - and NEW measurement of R_s
 - Advanced AI flavor tagging algorithms
 - pure b, c and strange jets $\rightarrow$ background contamination negligible
- Dominant systematics:
 - Hemisphere correlations
 - mainly driven by QCD (gluon emissions, $g \rightarrow b\bar{b}/c\bar{c}$, etc ..)
 - can be positive (negative) for hard (soft) emissions
 - can be reduced with (acoplanarity) cuts
 - measured directly in data
 - 10^9 (10^6) gluon splitting samples in FCC-ee (LCF)
- Projections:
 - 2(b) - 10(s) (4-10) $\times 10^{-6}$ for FCC-ee (LEP3) and 50-200 $\times 10^{-6}$ for LCF**



[2202.03285]



Asymmetries (LCF)

- A_f measures the asymmetry between Left (L) and Right (R) handed **Zff** couplings ($\sin^2\theta_w$):

$$A_f = \frac{g_{Lf}^2 - g_{Rf}^2}{g_{Lf}^2 + g_{Rf}^2}$$

- Linear colliders can measure A_e directly via LR asymmetry, using longitudinal beam polarisation combinations

- very clean experimentally (measure total hadronic cross-section)
 - independent of Z decay mode
- requires excellent control of initial state polarisation
 - determined by Blondel scheme
- provided best measurement of $\sin^2\theta_w$ at SLC

$$P_{e^-} \approx \pm 0.80, P_{e^+} \approx \pm 0.30$$

$$P_{\text{eff}} = \frac{P_{e^-} - P_{e^+}}{1 - P_{e^-}P_{e^+}} \approx 0.89$$

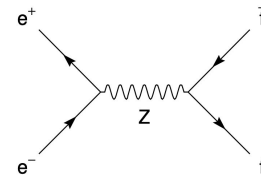
$$A_{\text{LR}} = \frac{1}{P_{\text{eff}}} \frac{N_L - N_R}{N_L + N_R} \approx A_e$$

- provides a measurement of $\Delta A_e / A_e \sim 2 \times 10^{-4}$
 - $\Delta \sin^2\theta_w \sim 4 \times 10^{-6}$ (**100x** vs. today world average)

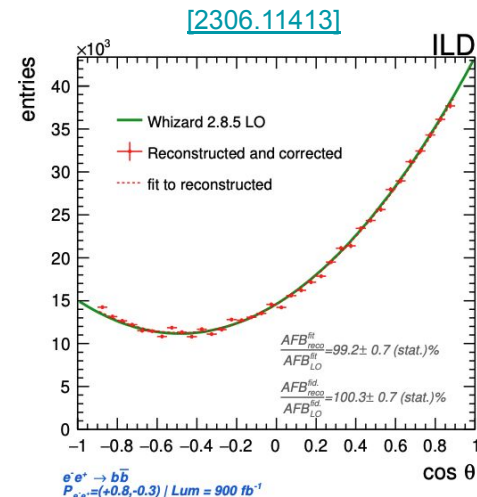
Other fermionic symmetries (LCF)

- A_f can then be measured as the Left Right - Forward Backward asymmetry:

$$A_{\text{LR,FB}}^f = \frac{1}{P_{\text{eff}}} \frac{(N^F - N^B)_L - (N^F - N^B)_R}{(N^F + N^B)_L + (N^F + N^B)_R} \propto A_f$$



- with a statistical uncertainty inflated by $1/\sqrt{B(Z \rightarrow ff)}$
- systematics:
 - Modelling uncertainties (symmetric) cancel
 - initial state polarisation (from A_{LR})
 - Measured at few 10^{-4} (depending on the fermion species)
 - 20-100X better than LEP

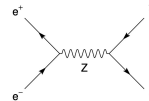


Asymmetries (FCC-ee/LEP3)

- At circular machines A_f can be measured through forward backward asymmetries in $ee \rightarrow ff$:

$$\frac{d\sigma}{d\cos\theta} = \frac{3}{8} \sigma_{\text{tot}} \left[1 + \cos^2\theta + \frac{8}{3} A_{\text{FB}} \cos\theta \right]$$

$$A_{\text{FB}} = \frac{3}{4} A_e A_f$$



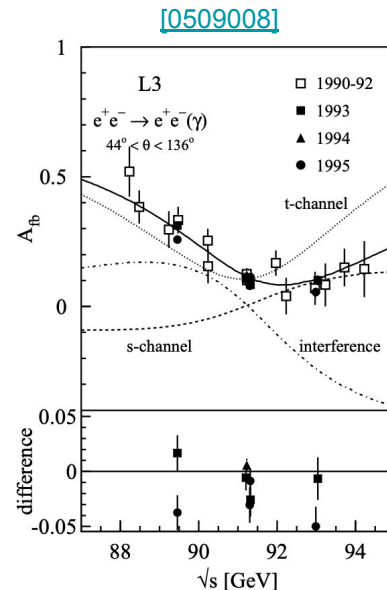
- To infer A_f , one has to measure A_e by other means:

- In $ee \rightarrow ee$, measure forward-backward asymmetry: $A_{\text{FB}}^e = \frac{3}{4} A_e^2$

Challenges:

- Requires t-channel Bhabha + interference subtraction
 - can be controlled off-peak
- Dominant systematics:
 - A_{FB}^e is measured off- and on-peak, and $A_{\text{FB}}^{0,e}$ is extracted as the value at m_Z corrected of ISR, interference effects
 - $\Delta\sqrt{s}_{\text{p.t.p}} \sim 20 \text{ keV}$ (uncorrelated component)

- $\Delta A_e \sim 1.3 (2.6) \times 10^{-5}$ at FCC-ee (LEP3)

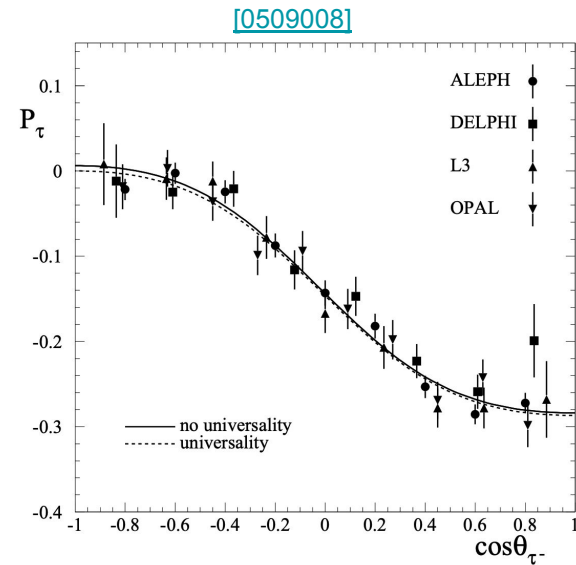


Asymmetries: A_e from τ -polarisation

- Without longitudinally polarised beams, one can access final helicity in $ee \rightarrow ff$ using τ -leptons
 - exploit that ν_τ flies in the same (opposite) direction as τ for L(R) in the τ center-of-mass
- Can construct polarisation of final state as:

$$\mathcal{P}_f = \frac{d(\sigma_r - \sigma_l)}{d \cos \theta} \bigg/ \frac{d(\sigma_r + \sigma_l)}{d \cos \theta} = - \frac{\mathcal{A}_f(1 + \cos^2 \theta) + 2\mathcal{A}_e \cos \theta}{(1 + \cos^2 \theta) + 2\mathcal{A}_f \cos \theta}$$

- A_e can then be accessed as the forward backward tau polarisation asymmetry:
 - Very clean measurement since minimally affected by potential data/MC mis-modelling, as long as they are FB symmetric
 - Example: hadronization tau decays
 - Non-tau backgrounds (bhabha, $Z \rightarrow \mu\mu$), enough to assume 10x vs LEP
 - Big improvement to come with IP tau lifetime
 - $\Delta A_e / A_e \sim 2(4) \times 10^{-5}$ at FCC-ee (LEP3)



Other fermion asymmetries (FCC-ee/LEP3)

$$A_{\text{FB}}^{0,f} = \frac{3}{4} \mathcal{A}_e \mathcal{A}_f$$

- Once $\mathbf{A}_e$ is known, $\mathbf{A}_f$ can then be measured from Forward-Backward asymmetries

- dominant systematics:

- $A_{\text{FB}}^{\mu,\tau}$: point-to-point energy calibration

- $\Delta A_{\text{FB}}^{\mu} \approx 2.3 [2.4] \times 10^{-6} \rightarrow \Delta A_{\mu}/A_{\mu} \sim 3(6) \times 10^{-5}$
 - $\Delta A_{\text{FB}}^{\tau} \approx 2.7 [2.4] \times 10^{-6} \rightarrow \Delta A_{\tau}/A_{\tau} \sim 4(7) \times 10^{-5}$

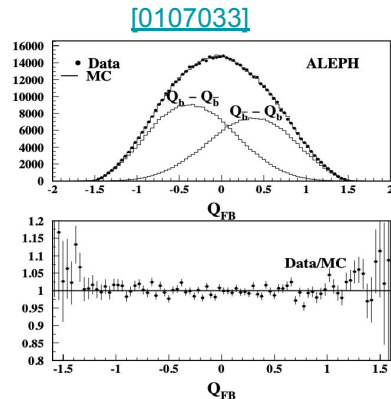
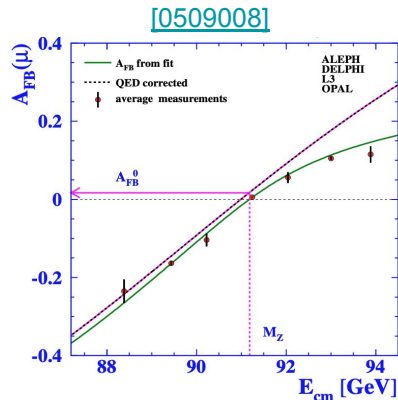
- $\mathbf{A}_q$ for heavy flavors (q=b,c,s) requires the determination of the jet (hemisphere) charge

- Measure simultaneously:

- average FB hemisphere charge difference AND sum
 - Measure charge separation
 - $\rightarrow \Delta A_{\text{FB}}^q \sim 5 \times 10^{-6} \rightarrow \Delta A_q \sim 2 \times 10^{-4}$

- dominant systematics:

- Asymmetry in detector material (nuclear interactions produce excess of positive charge)
 - can be monitored with conversions
 - Background contamination
 - tagger purities much larger than LEP



$$\alpha_{\text{QED}}(m_Z), \sigma_{\text{had}}^0$$

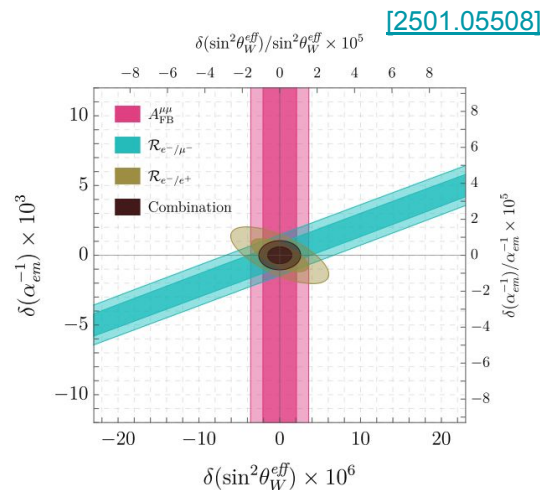
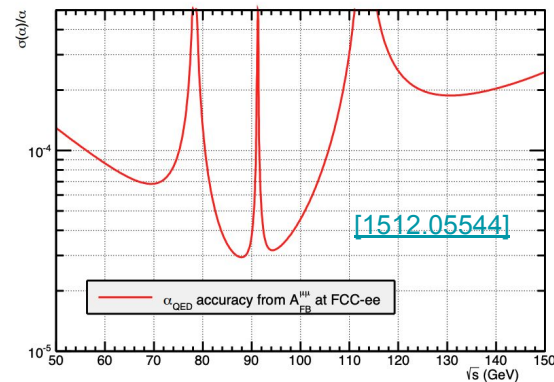
$$\cos^2 \theta_{\text{eff}}^f \sin^2 \theta_{\text{eff}}^f = \frac{\pi \alpha(0)}{\sqrt{2} m_Z^2 G_F} \frac{1}{1 - \Delta r^f}$$

• Electro-magnetic constant

- Dominant parametric uncertainty in EW precision ($\sin \vartheta_W^{\text{eff}}$ and m_W) fit:
 - Current uncertainty $\delta\alpha/\alpha = 1.4 \times 10^{-4}$
- LCF has to assume required precision from Lattice
- FCC-ee (LEP3) can directly measure it:
 - from off-peak FB asymmetry (interference with γ^*) in $\mu\mu$ events ($\delta\alpha/\alpha = 3 \times 10^{-5}$)
 - small experimental uncertainty, stat dominated
 - Z-pole energy points chosen to optimize measurement !
 - from $R_{e^+e^-}, R_{e^-/\mu^-}$ ($\delta\alpha/\alpha = 0.6 \times 10^{-5}$)
 - e^+/e^- efficiency control (charge mis-id), material budget
 - e^-/μ^- acceptance difference (to be determined from 10^{11} lepton pairs)
 - Can then provide comparison with Lattice calculation

• Peak cross-section

- Sensitive to Zee coupling and total width, provides N_ν
- limited by luminosity determination, $\delta\mathcal{L}/\mathcal{L}$
 - Bhabha $\sim 10^{-4}$ (FCC-ee, LEP3, LCF)
 - $e^+e^- \rightarrow \gamma\gamma \sim 2(4) \cdot 10^{-5}$ (FCC-ee, LEP3)



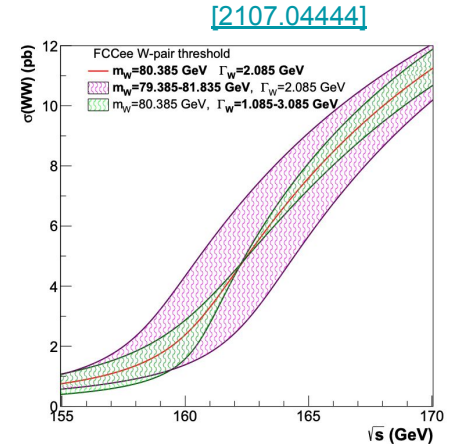
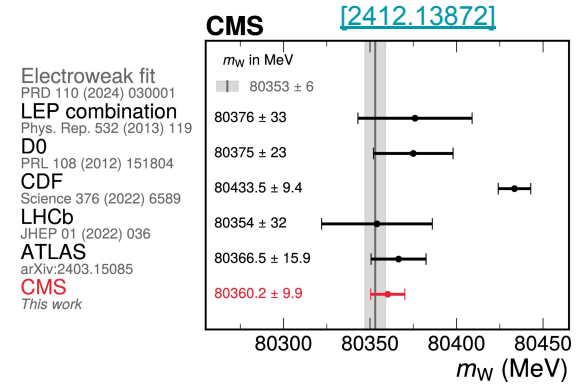
WW threshold and above

W mass / width (threshold scan) - Circular

- **W mass** is a crucial input to test SM consistency
 - Current uncertainty $\Delta m_W \sim 10 \text{ MeV}$ (LHC)
- At lepton colliders, the simplest way is to measure it, is via a threshold scan, i.e cross-section vs. beam energy:
- With 2 or more energy points, m_W and Γ_W can be extracted simultaneously

$$\Delta m_W(\text{stat}) = \left(\frac{d\sigma_{WW}}{dm_W} \right)^{-1} \frac{\sqrt{\sigma_{WW}}}{\sqrt{\mathcal{L}}}$$

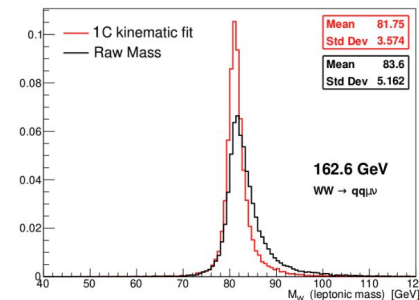
- $\Delta m_W^{\text{scan}}(\text{stat}) \sim 400 \text{ keV}$ (FCC-ee)
- $\Delta \Gamma_W^{\text{scan}}(\text{stat}) \sim 1 \text{ MeV}$ (FCC-ee)
- Dominant systematics, beam energy calibration
 - Absolute energy calibration
 - $\Delta m_W(\text{syst}) \sim 150 \text{ keV}$ (FCC-ee) (300 keV on $\sqrt{s}$)
 - via resonant depolarisation (**only at FCC-ee**)
 - $\Delta m_W(\text{syst}) \sim 700 \text{ keV}$ (LEP3)
 - via radiative return Z (**at LEP3**)
- Through polarized scan, $\Delta m_W(\text{LCF}) \sim 2 \text{ MeV}$ (not in official run plan)



W mass, kinematic fit

- **W mass** can also be accessed through direct reconstruction of decay products: a kinematic fit (done at LEP2)
 - can be performed at threshold (~ 160 GeV) $\rightarrow$ FCC-ee, LEP3
 - in the “boosted” regime (~ 240 GeV) $\rightarrow$ LCF, FCC-ee, LEP3
- Major sensitivity from the **qq ℓ v** channel (1C) (4 - 3)
 - Avoid issues with color reconnection
 - Use full 4-momentum constraints
 - Allows to fully resolve the neutrino kinematics $\rightarrow 4-3=1$ constraint
 - Additional constraint that W masses are equal (2C)
 - Dominant systematics
 - absolute $\sqrt{s}$ and BES calibration
 - FCC-ee : resonant depol. + di-muon
 - LEP3, ILC: radiative Z returns
 - hadronization
 - more important where WW system is boosted (since non isotropic)

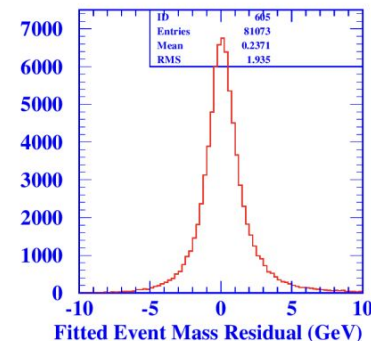
[M.Beguin](#)



$$\Delta m_W^{\text{kin}}(\text{stat}) \sim \mathbf{210} \text{ (400) keV} - \text{FCC-ee(LEP3)}$$

$$\Delta m_W^{\text{kin}}(\sqrt{s}) \sim \mathbf{160} \text{ (700) keV} - \text{FCC-ee (LEP3)}$$

[G. Wilson](#)



$$\Delta m_W^{\text{kin}}(\text{stat}) \sim \mathbf{600} \text{ keV (LCF)}$$

$$\Delta m_W^{\text{kin}}(\text{had.}) \sim \mathbf{1} \text{ MeV}$$

W properties

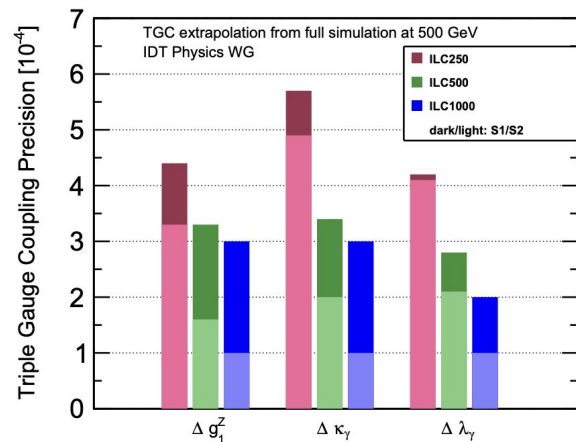
- Large **W** production rate can be exploited to constrain leptonic BRs

- absolute measurement will be limited by luminosity determination
 - from $\gamma\gamma / \text{BhaBha}$, $\delta\mathcal{L}/\mathcal{L} \sim 10^{-4}$
- FCC-ee / LEP3 / LCF $\sim 10^{-4} / 2 \cdot 10^{-4} / 2 \cdot 10^{-4}$

- Triple/Quartic Gauge couplings

- through WW/VBS production
- optimal observable analysis exploiting full 5D angular information of WW decays in semilep decay
- LCF has powerful constraints thanks to:
 - Initial state polarisation
 - Higher energy
 - energy growth of relevant operators
 - higher luminosity
- few 10^{-4} in reach at LCF

$\mathcal{B}(W \rightarrow e\nu_e) \times 10^4$	1071	$\pm$	16	0.13	0.10	From WW and ZH threshold luminosity
$\mathcal{B}(W \rightarrow \mu\nu_\mu) \times 10^4$	1063	$\pm$	15	0.13	0.10	From WW and ZH threshold luminosity
$\mathcal{B}(W \rightarrow \tau\nu_\tau) \times 10^4$	1138	$\pm$	21	0.13	0.15	From WW scan ZH threshold luminosity



Conclusions

- For Z-pole precision measurements:
 - Linear Collider has the advantage of longitudinal polarisation
 - Circular colliders have the advantage of luminosity (300-1000)
- Per unit of luminosity, Chiral observables (asymmetries) are measured better at linear colliders ,
 - by $\sim 1/A_e^2 \sim 40$
- However FCC-ee (LEP3) have 1000(300) more luminosity at the Z pole, so overall measurements are more precise by factors of:
 - $\sin^2\theta_W$ → **5 : 2 : 1** for FCC-ee : LEP3 : LCF
 - lepton universality/couplings → **30 : 15 : 1**
 - electromagnetic and strong coupling only at FCC-ee (LEP3)
 - m_Z → **2 : 2 : 1**
 - Γ_Z → **10 : 5 : 1**
- For the W mass, difference is milder, but resonant depol. makes a difference for FCC-ee vs LEP3/LCF
 - m_W → **4 : 1 : 1**
- Di-fermion production and multi-boson production can be uniquely probed by
 - linear colliders (ILC, CLIC, ...) up to 500-3 TeV
 - muon collider up to 10 TeV
 - FCC-hh up to 40-50 TeV

Backup

Lepton universality tests and $\sin^2\theta_W$

Table 6: Projected ABSOLUTE uncertainties on value of $\sin^2\theta_W^{\text{eff}}$ obtained at the Z-pole from FCC-ee, LCvision, and LEP3. When two numbers are given, the first [second] corresponds to the projected statistical [systematic]; total uncertainty.

FCC		
quantity	uncertainty	$\sin^2\theta_W^{\text{eff}}$
$A_{\text{FB}}^{0,\ell}(10^{-6})$	1.4 [2.7]	0.75 [1.44] ; 1.6
$\mathcal{A}_e$ from $A_{\text{FB}}^{\text{pol}(\tau)}$	7 [20]	0.9 [2.5] ; 2.8
$A_{\text{FB}}^{0,b}(10^{-6})$	4 [4]	0.74 [0.74] ; 1.03
$A_{\text{FB}}^{0,c}(10^{-6})$	5 [5]	1.3 [1.3] ; 1.8
$A_{\text{FB}}^{0,\tau}(10^{-6})$	7.4 [7.4]	1.4 [1.4] ; 1.9
FCC combination of $\sin^2\theta_W^{\text{eff}}$		0.4 [0.5] ; 0.7
LEP3		
quantity	uncertainty	$\sin^2\theta_W^{\text{eff}}$
$A_{\text{FB}}^{0,\ell}(10^{-6})$	2.6 [4.6]	1.4 [2.7] ; 3.0
$\mathcal{A}_e$ from $A_{\text{FB}}^{\text{pol}(\tau)}$	13 [39]	1.7 [4.7] ; 5.2
$A_{\text{FB}}^{0,b}(10^{-6})$	7.8 [7.8]	1.4 [1.4] ; 2
$A_{\text{FB}}^{0,c}(10^{-6})$	9.5 [9.5]	2.4 [2.4] ; 3.4
$A_{\text{FB}}^{0,\tau}(10^{-6})$	14 [14]	2.6 [2.6] ; 3.6
LEP3 combination of $\sin^2\theta_W^{\text{eff}}$		0.75 [0.95] ; 1.35
LCvision		
quantity	uncertainty	$\sin^2\theta_W^{\text{eff}}$
$\mathcal{A}_e$ from $A_{\text{LR}}(10^{-6})$	23 [19]	2.9 [2.4] ; 4
$\mathcal{A}_\mu$ from $A_{\text{LRFB}}^{(\mu)}(10^{-6})$	100 [50]	12.6 [6.3] ; 14
$\mathcal{A}_\tau$ from $A_{\text{LRFB}}^{(\tau)}(10^{-6})$	100 [50]	12.6 [6.3] ; 14
LCvision combination of $\sin^2\theta_W^{\text{eff}}$		2.75 [2.4] ; 3.7

Table 5: Projected RELATIVE uncertainties on partial-width ratios obtained on peak at the Z-pole from FCC-ee, LCvision, and LEP3. When two numbers are given, the first [second] corresponds to the projected statistical [systematic] uncertainty. .

Observable (value)	present uncertainty	FCC-ee Stat. [Syst.]	LCF Stat. [Syst.]	LEP3 Stat. [Syst.]	Comment and leading uncertainty
$R_\ell \equiv \frac{\Gamma_{\text{had}}}{\Gamma_\ell}$	20.767 ± 0.025				
relative uncertainty (10^{-6})	1200	1.5 [2.3]		2.8 [2.3]	Low angle acceptance
$R_e \equiv \frac{\Gamma_{\text{had}}}{\Gamma_e}$	20.804 ± 0.050				
relative uncertainty (10^{-6})	2400	3.4 [2.3]	200[500]	6.4 [2.3]	Low angle acceptance
$R_\mu \equiv \frac{\Gamma_{\text{had}}}{\Gamma_\mu}$	20.785 ± 0.033				
relative uncertainty (10^{-6})	1600	2.4 [2.3]	200 [200]	4.5 [2.3]	Low angle acceptance
$R_\tau \equiv \frac{\Gamma_{\text{had}}}{\Gamma_\tau}$	20.764 ± 0.045				
relative uncertainty (10^{-6})	2100	2.7 [2.3]	200 [200]	5.1 [2.3]	Low angle acceptance
LEPTON UNIVERSALITY					
Axial vector couplings					
R_e/R_ℓ	rel. (10^{-6})	3.1	600	5.8	
R_μ/R_ℓ	rel. (10^{-6})	2.8	300	5.3	
R_τ/R_ℓ	rel. (10^{-6})	3.6	300	6.8	
Vector couplings					
$A_{\text{FB}}^{0,e} / A_{\text{FB}}^{0,\ell}$	rel. (10^{-6})	210	N.A.	390	
$A_{\text{FB}}^{0,\mu} / A_{\text{FB}}^{0,\ell}$ or $\mathcal{A}_\mu/\mathcal{A}_e$	rel. (10^{-6})	157	700	295	
$A_{\text{FB}}^{0,\tau} / A_{\text{FB}}^{0,\ell}$ or $\mathcal{A}_\tau/\mathcal{A}_e$	rel. (10^{-6})	183	700	345	
QUARK FLAVOURS					
$R_b \equiv \frac{\Gamma_b}{\Gamma_{\text{had}}}$	0.21629 ± 0.00066				
relative uncertainty (10^{-6})	3300	1.2 [1.6]	20 [60]	2.2 [3.0]	multiple flavour tags
$R_c \equiv \frac{\Gamma_c}{\Gamma_{\text{had}}}$	0.1721 ± 0.0030				
relative uncertainty (10^{-6})	17000	1.4 [2.2]	100[250]	2.6 [4.2]	multiple flavour tags
NEW! $R_s \equiv \frac{\Gamma_s}{\Gamma_{\text{had}}}$	relative (10^{-6})	N.A.			
relative uncertainty (10^{-6})	N.A.	2.5 [11]	–	4.7 [21]	multiple flavour tags