



PARALLEL SESSION / **FLAVOUR PHYSICS**

Theoretical landscape in flavour physics

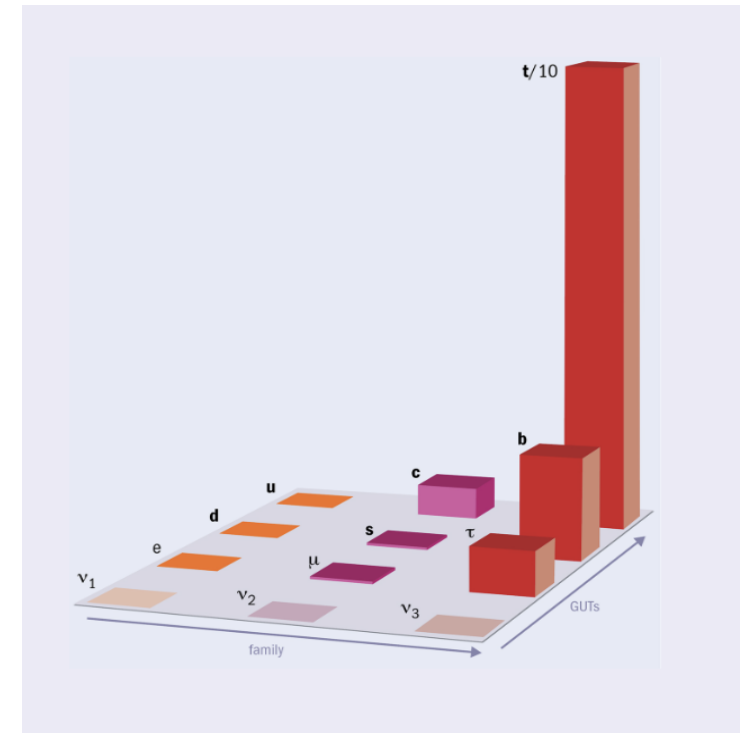
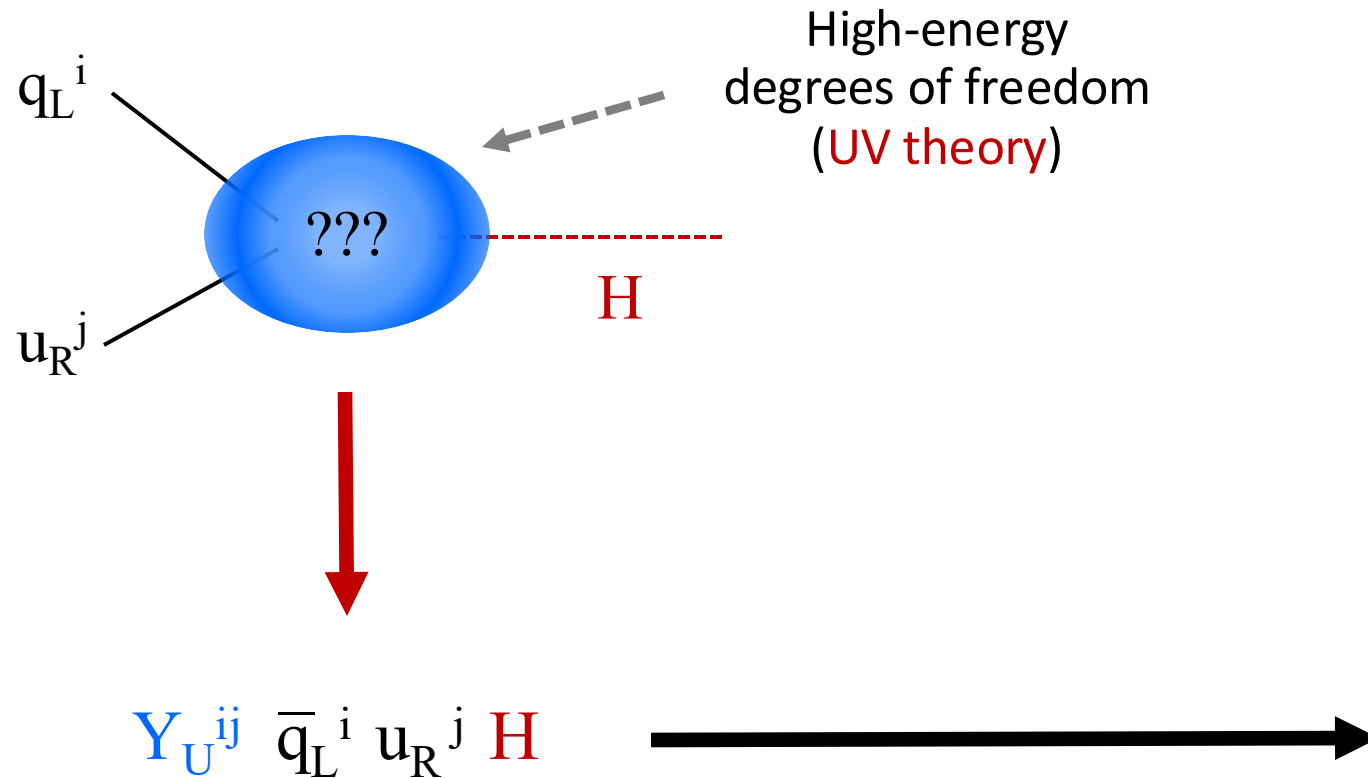
Gino Isidori (University of Zürich)

23-27 JUNE 2025 Lido di Venezia



The open questions: *the SM flavour puzzle*

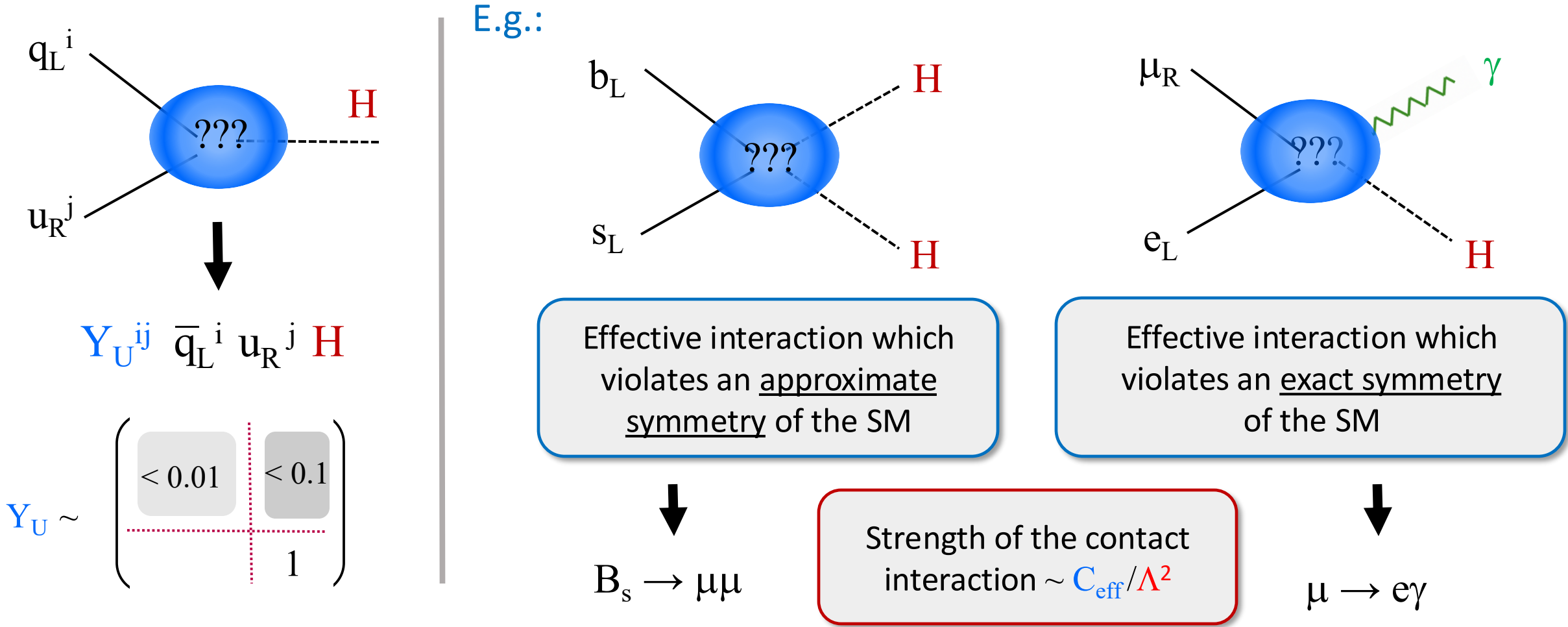
Fermion masses are the results of the Yukawa interactions → Inescapable link between Higgs and flavor, whose origin can be addressed only **beyond the SM**...



“Message from the UV”
that we need to “decode”

The open questions: *the NP flavour puzzle*

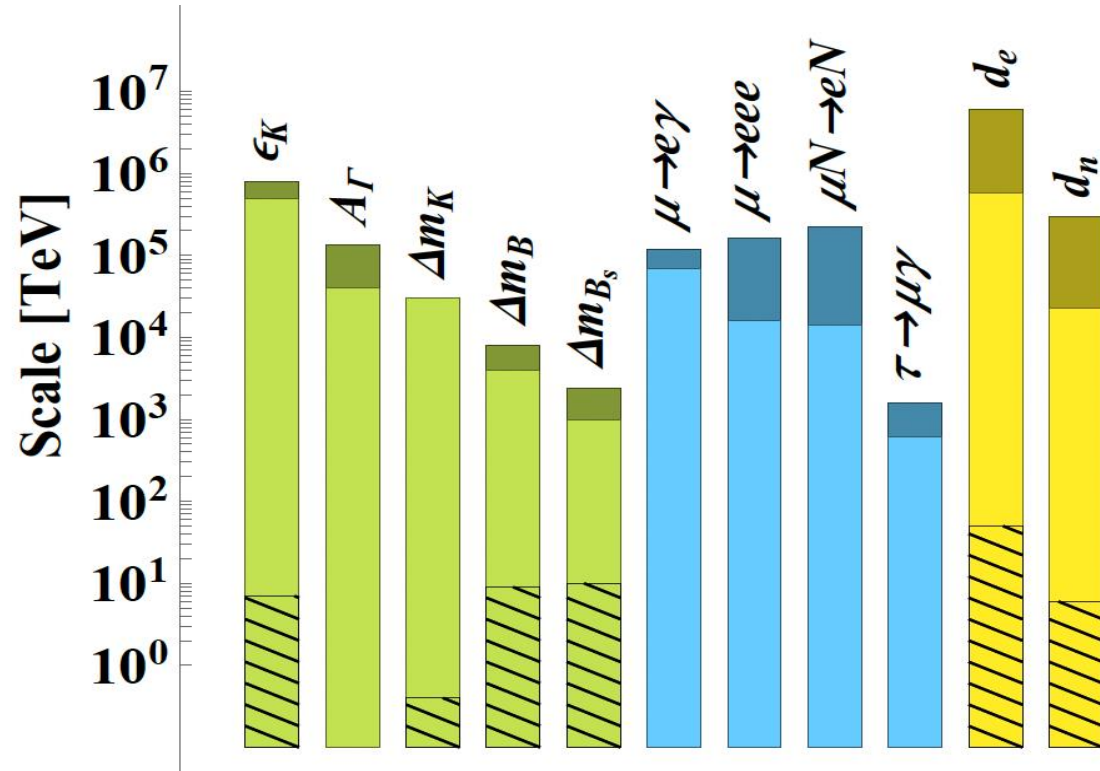
On general grounds, the flavor non-universal dynamics in the UV should give rise to **effective contact interactions** → **violations of exact & approximate symmetries of the SM**:



The open questions: *the NP flavour puzzle*

On general grounds, the flavor non-universal dynamics in the UV should give rise to **effective contact interactions** → **violations of exact & approximate symmetries of the SM**.

At present we do not see any of these effects, and this is often expressed in terms of strong bounds on effective NP scales:



N.B.: These bounds are obtained assuming unit couplings for the effective coefficients

E.g.:

$$\frac{C_{\mu e \gamma}}{\Lambda^2} \rightarrow \frac{1}{\Lambda^2}$$

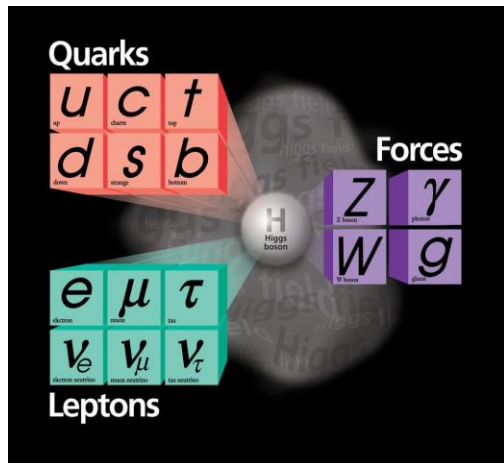
Useful to show the potential,
but rather misleading...

How to address the challenges: *Flavour in the SM EFT*

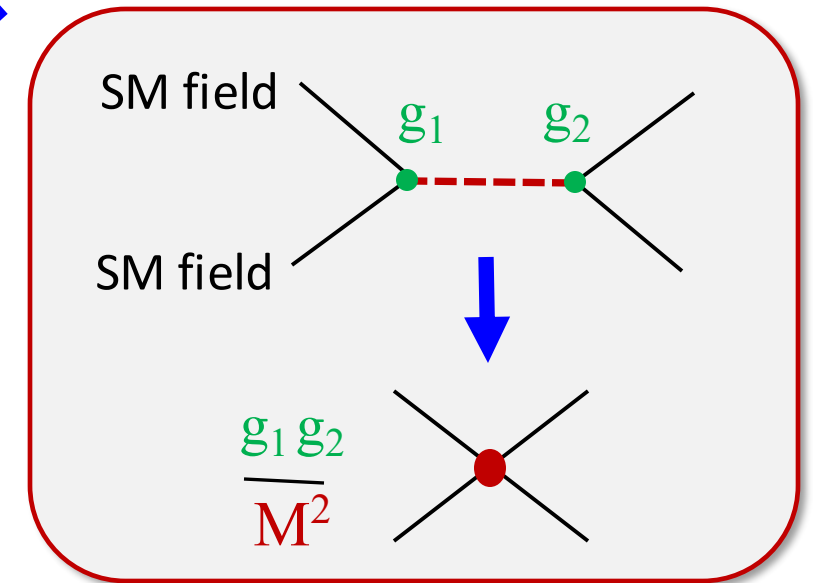
No matter if NP is at the TeV scale or well above, it is always very heavy compared to the energy scale of low-energy flavour experiments → systematic description via the “SM EFT”

$$\mathcal{L}_{\text{SM-EFT}} = \underbrace{\mathcal{L}_{\text{gauge}} + \mathcal{L}_{\text{Higgs}}}_{\text{Long-range structure of the theory}} + \sum_{d > 4} \frac{c_i^{[d]}}{\Lambda^{d-4}} \mathcal{O}_i^{[d]}(\psi_{\text{SM}})$$

Long-range structure
of the theory



Most general
description of heavy NP
that we are not able to
directly “excite”



How to address the challenges: *Flavour in the SM EFT*

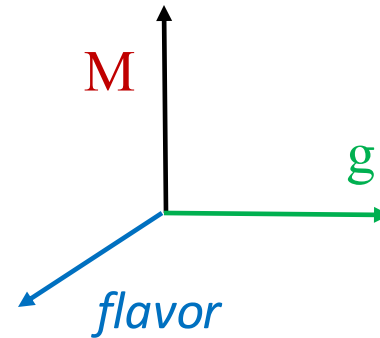
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Understanding the flavour structure of the SM EFT is a question of pivotal importance.

N. of independent couplings:

- No flavor symmetry 2499 [@d=6]
- Maximal flavor symmetry 47 [@d=6]



Don't be scared by the large number of independent couplings of eff. operators!

They do not indicate model parameters, but rather the number of observables we can probe (*whose correlations depend on the unknown underlying model*)

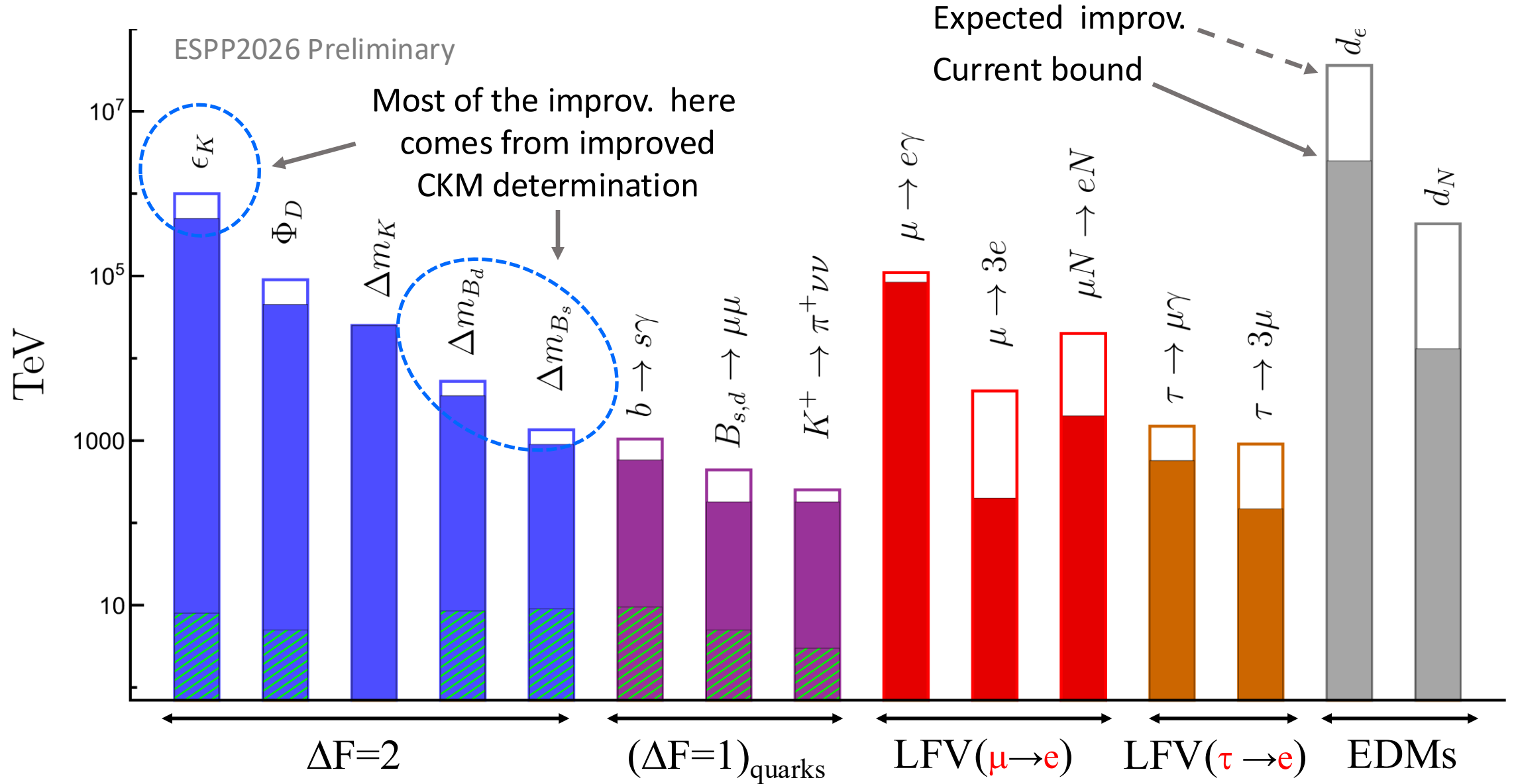
How to address the challenges @ large-scale facilities [b,c,τ]

	γ	quarks	μ	e	τ	ν
$3_q \rightarrow 2_q$	$b \rightarrow s \gamma$	$B_s \leftrightarrow \bar{B}_s$	$b \rightarrow s \mu\mu$ $b \rightarrow c \mu\nu$	$R_{bs}(e/\mu)$ $R_{bc}(e/\mu)$	$b \rightarrow s \tau\tau$ $b \rightarrow c \tau\nu$	$b \rightarrow s \nu\nu$
$3_q \rightarrow 1_q$	$b \rightarrow$	$B(B_s \rightarrow \mu\mu),$ $B(B \rightarrow K^* \mu\mu),$ $A_{FB}(B \rightarrow K^* \mu\mu),$ $\square$	$b \rightarrow d \mu\mu$ $b \rightarrow u \mu\nu$	$R^{\tau/\mu}(D),$ $R^{\tau/\mu}(D^*), \dots$	$b \rightarrow d \tau\tau$ $b \rightarrow u \tau\nu$	$B(B \rightarrow K \nu\nu),$ $B(B \rightarrow K^* \nu\nu),$ $\square$
$2_q \rightarrow 2_q$	$c \rightarrow u \gamma$	$D \leftrightarrow \bar{D}$	$c \rightarrow u \mu\mu$ $c \rightarrow s \mu\nu$ $c \rightarrow d \mu\nu$	$R_{cu}(e/\mu)$ $R_{cs}(e/\mu)$ $R_{cd}(e/\mu)$		$c \rightarrow u \nu\nu$
	N.B.: several observables for each transition					
$3_l \rightarrow 2_l$	$\tau \rightarrow \mu \gamma$ $\tau \rightarrow e \gamma$	$\tau \rightarrow \mu qq$ $\tau \rightarrow e qq$	$\tau \rightarrow \mu \mu\mu$ $\tau \rightarrow \mu ee$	$\tau \rightarrow \mu ee$ $\tau \rightarrow e ee$		$\tau \rightarrow \mu \nu\nu$ $\tau \rightarrow e \nu\nu$

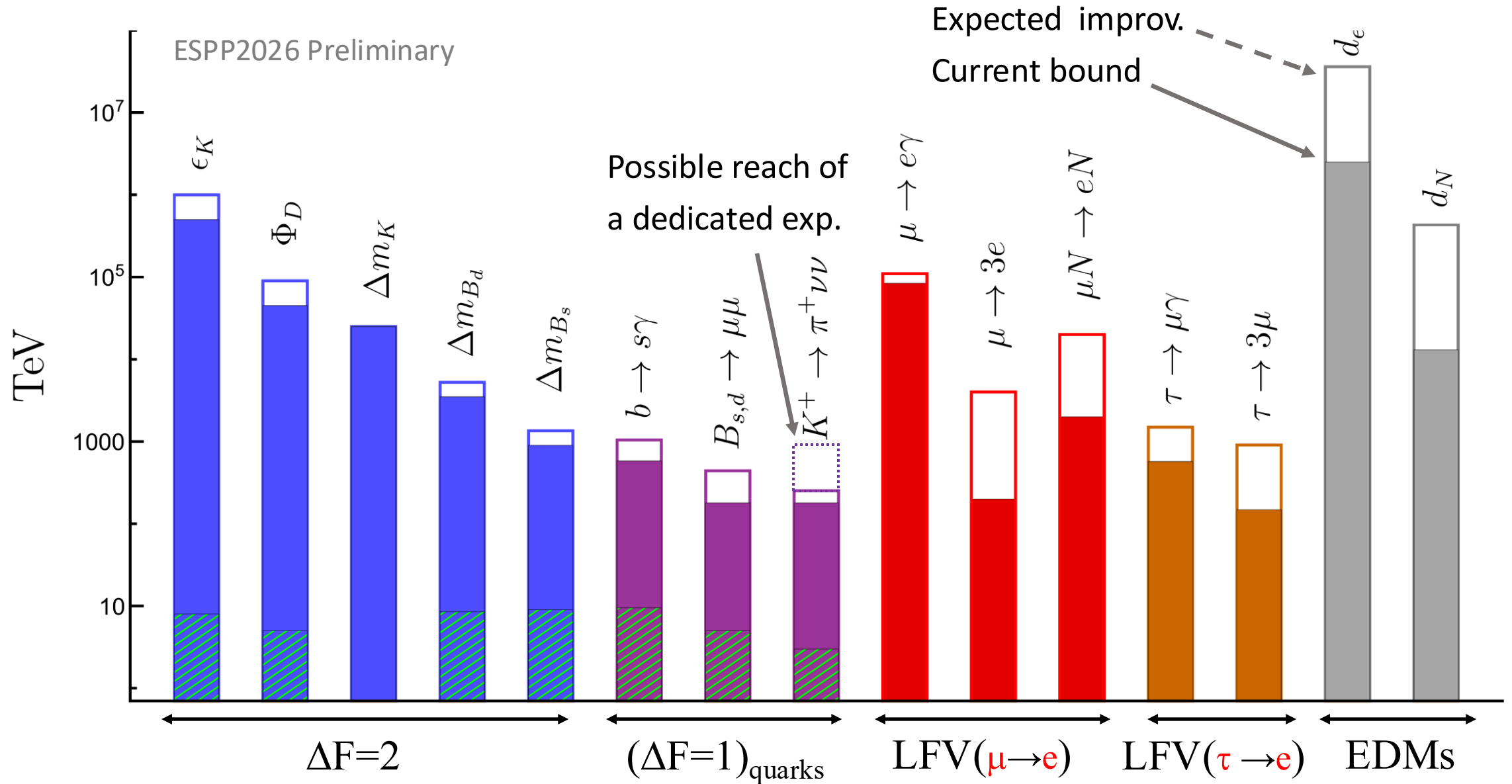
How to address the challenges @ dedicated facilities $[K, \pi, \mu, \text{EDMs}]$

	γ	quarks	μ	e	τ	ν
$2_q \rightarrow 1_q$	$s \rightarrow d \gamma$	$K \leftrightarrow \bar{K}$	$s \rightarrow d \mu \mu$ $s \rightarrow u \mu \nu$ $d \rightarrow u \mu \nu$	$R_{sd}(e/\mu)$ $R_{su}(e/\mu)$ $R_{du}(e/\mu)$		$s \rightarrow d \nu \nu$
$2_l \rightarrow 1_l$	$\mu \rightarrow e \gamma$	$\mu N \rightarrow e N$		$\mu \rightarrow e e e$		$\mu \rightarrow e \nu \nu$
$2_l \rightarrow 2_l$	$g_\mu \quad d_\mu$					
$1_l \rightarrow 1_l$	d_e					
$1_q \rightarrow 1_q$	d_N					

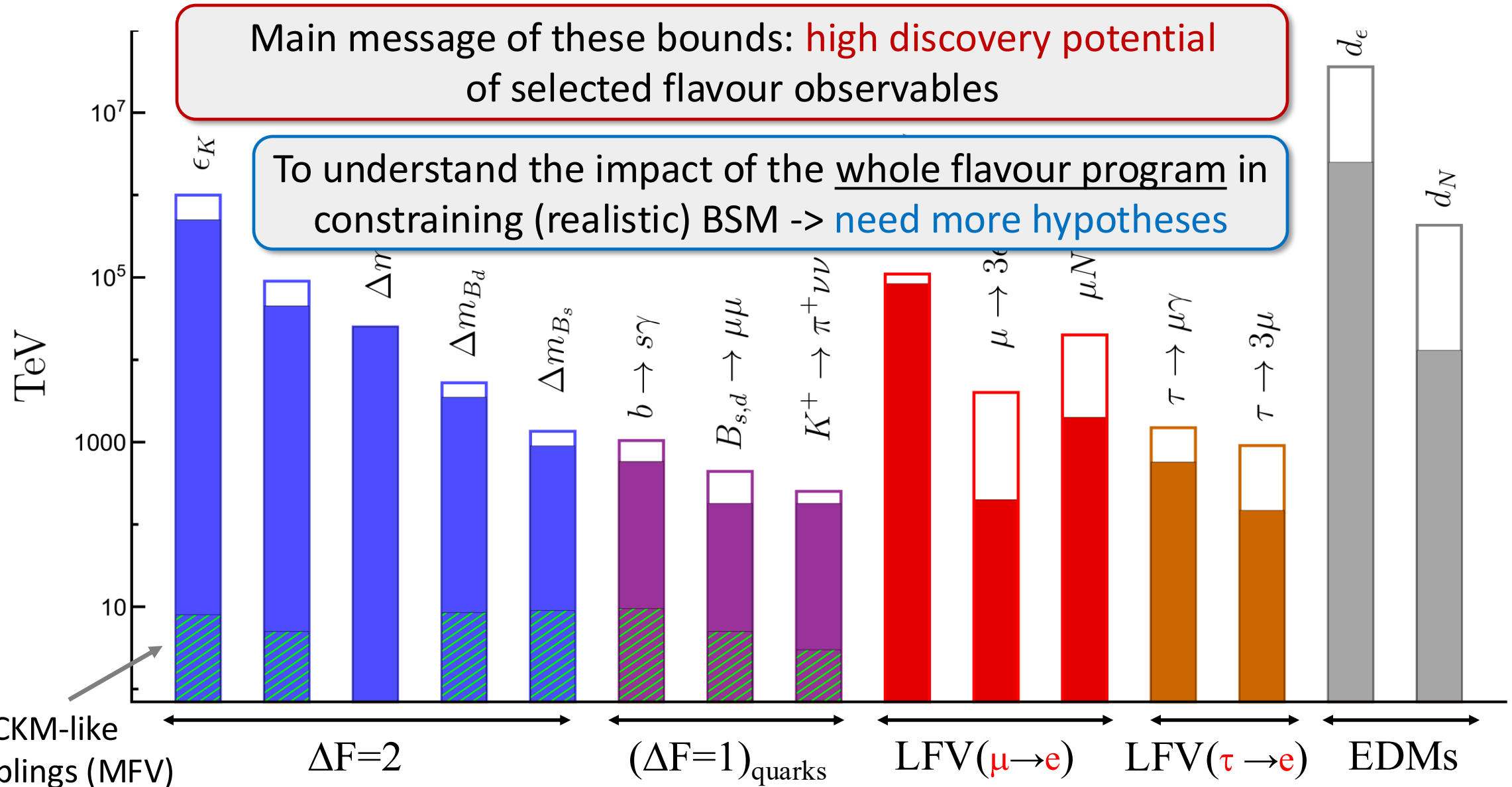
I. Bounds on effective scales for generic (unit) couplings



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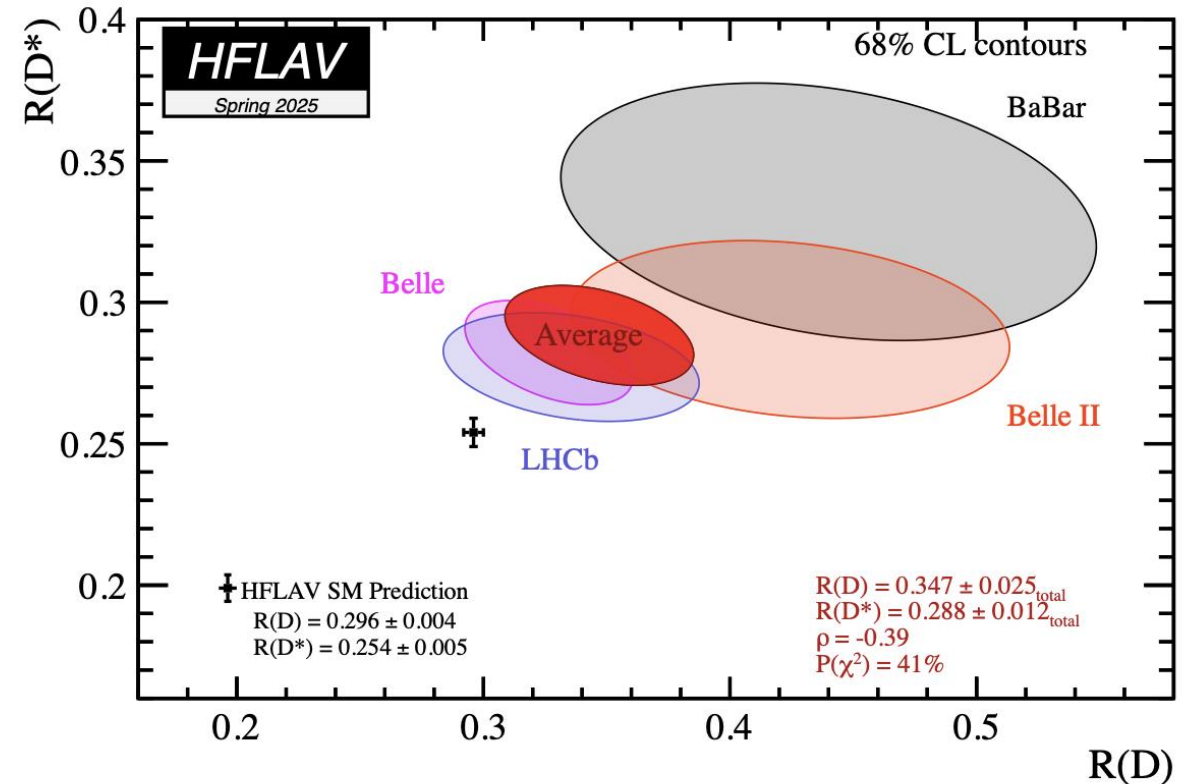


II. An explicit example of flavour – EW interplay (*assuming NP*)

Looking at bounds from individual observables misses the main point of indirect NP searches, which is the interplay of a large number of observables in “decoding” the underlying theory

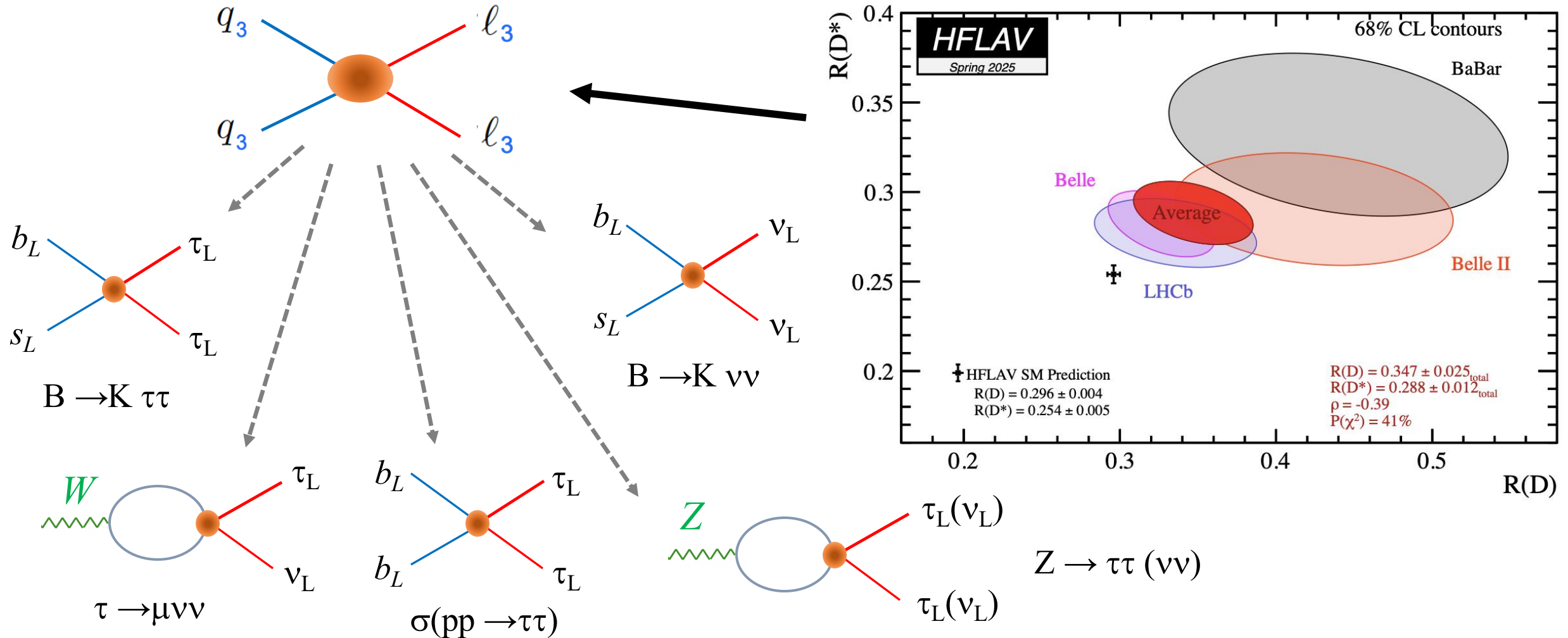
To illustrate this point worth to consider an explicit example assuming a NP signal

$$R_{D^{(*)}} = \frac{\mathcal{B}(B \rightarrow D^{(*)} \tau \nu)}{\mathcal{B}(B \rightarrow D^{(*)} \mu \nu)}$$



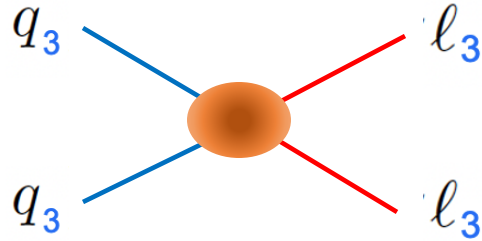
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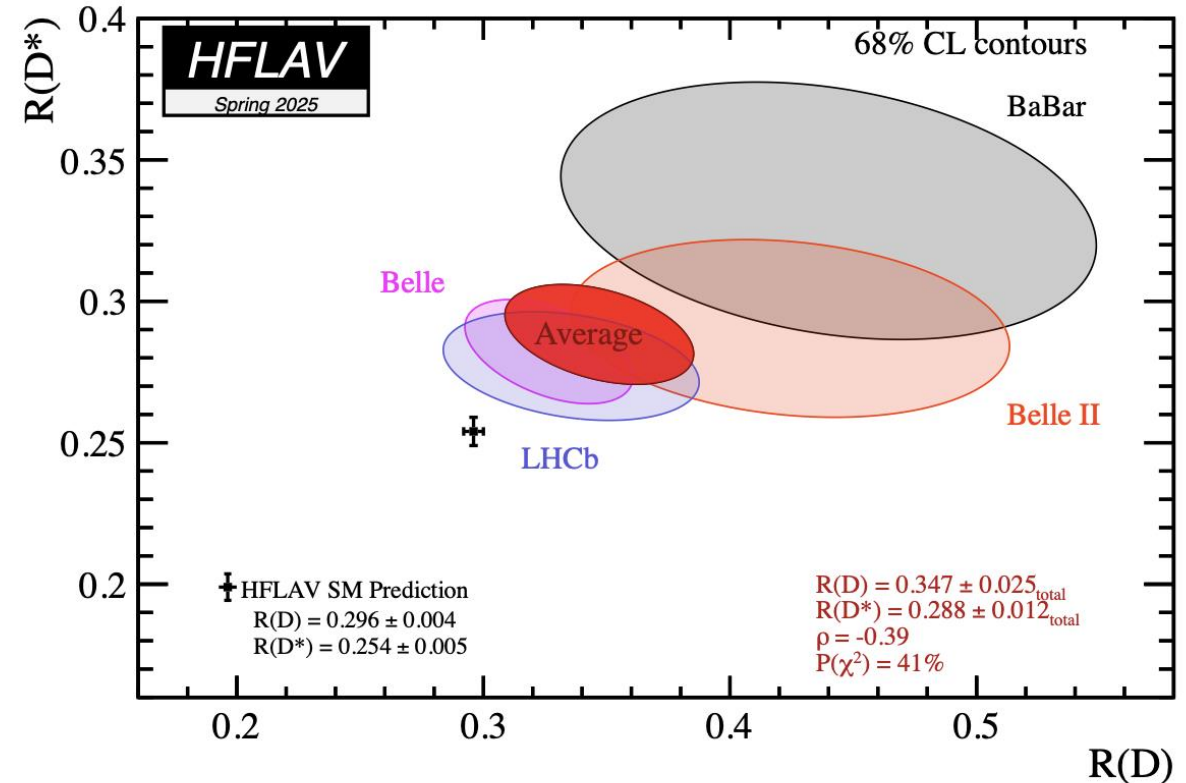
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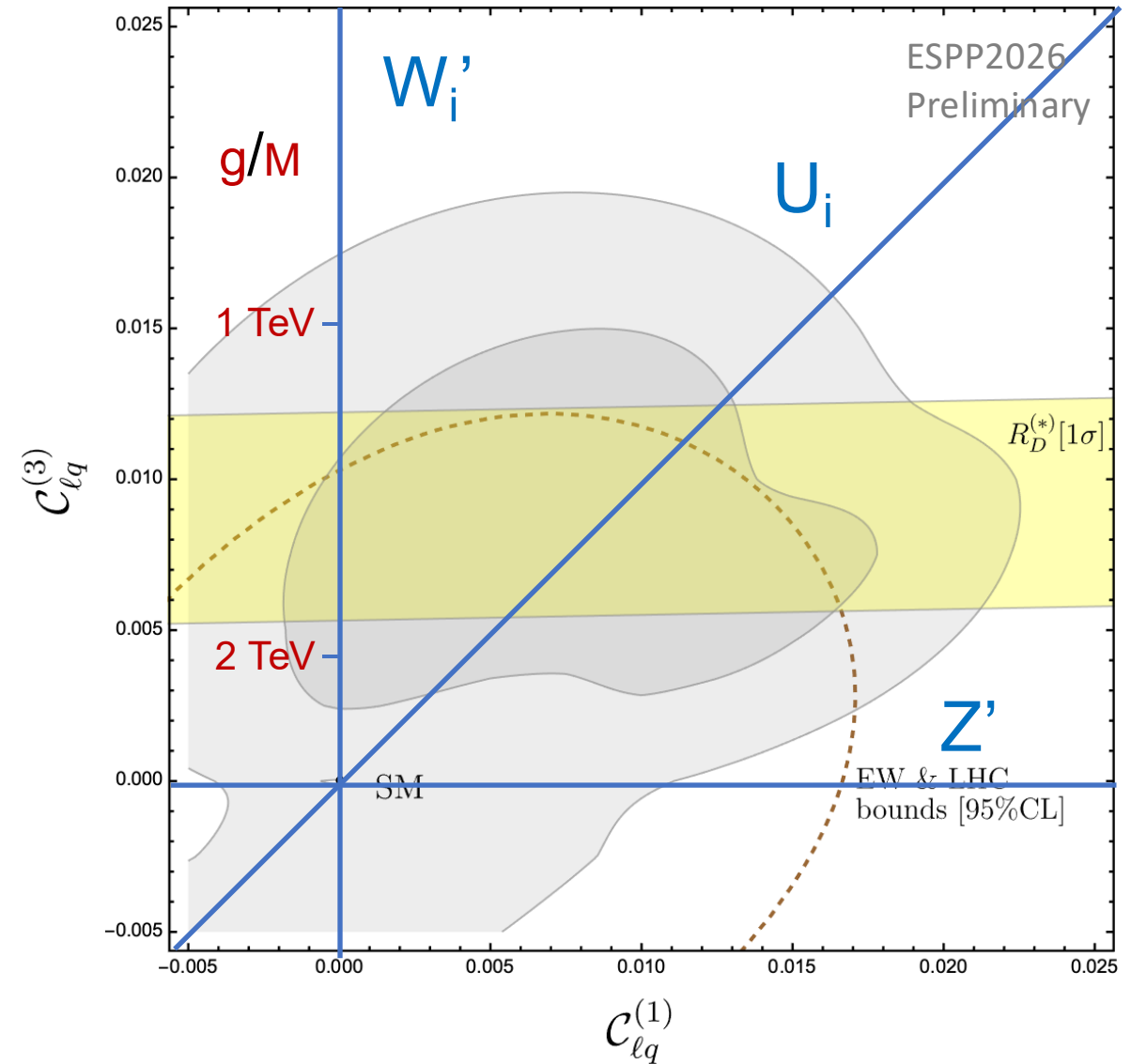
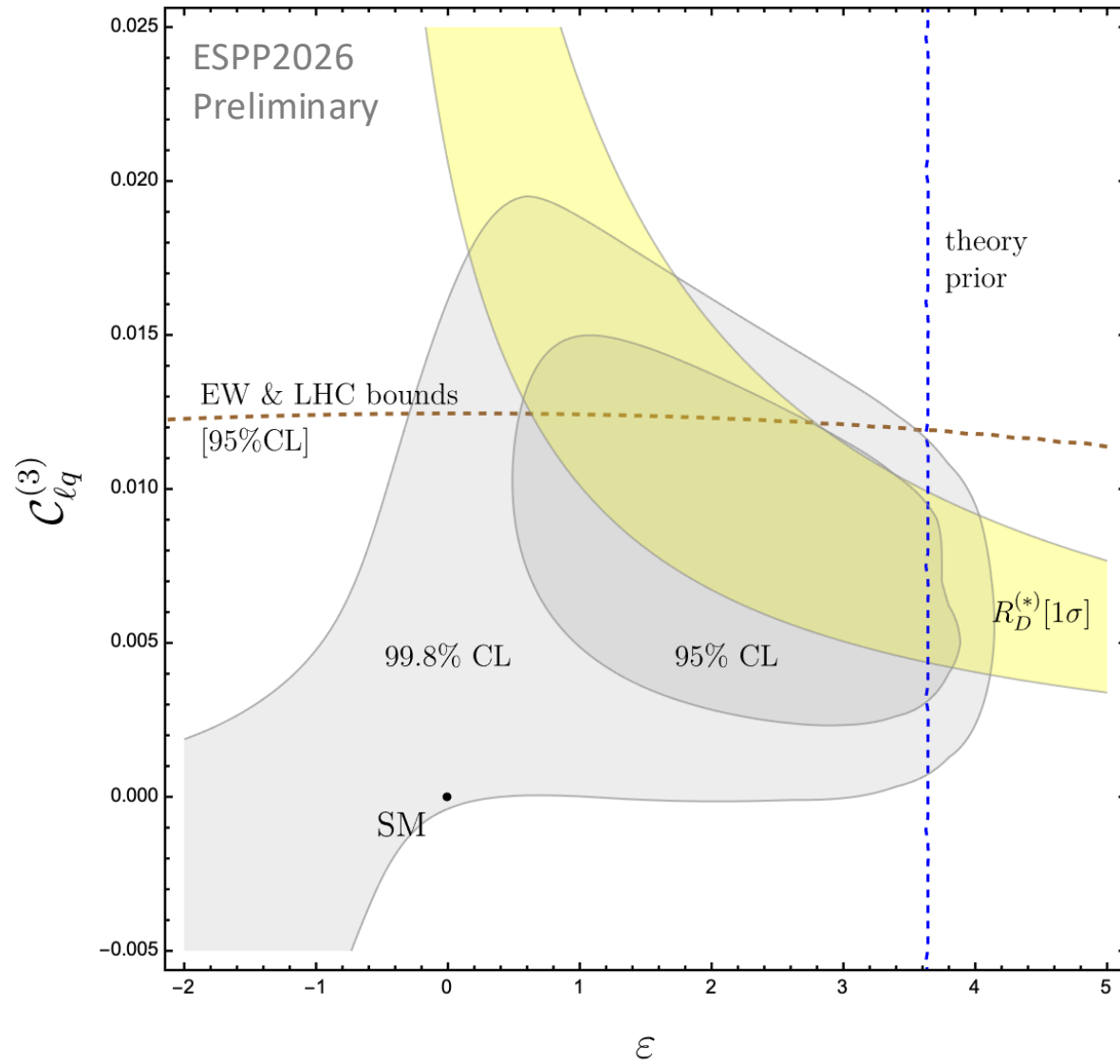
A simple, consistent, description of the system involves (*at least*) 3 parameters:

■ Two Wilson coefficients, $C_{\ell q}^{(1)}$ $C_{\ell q}^{(3)}$

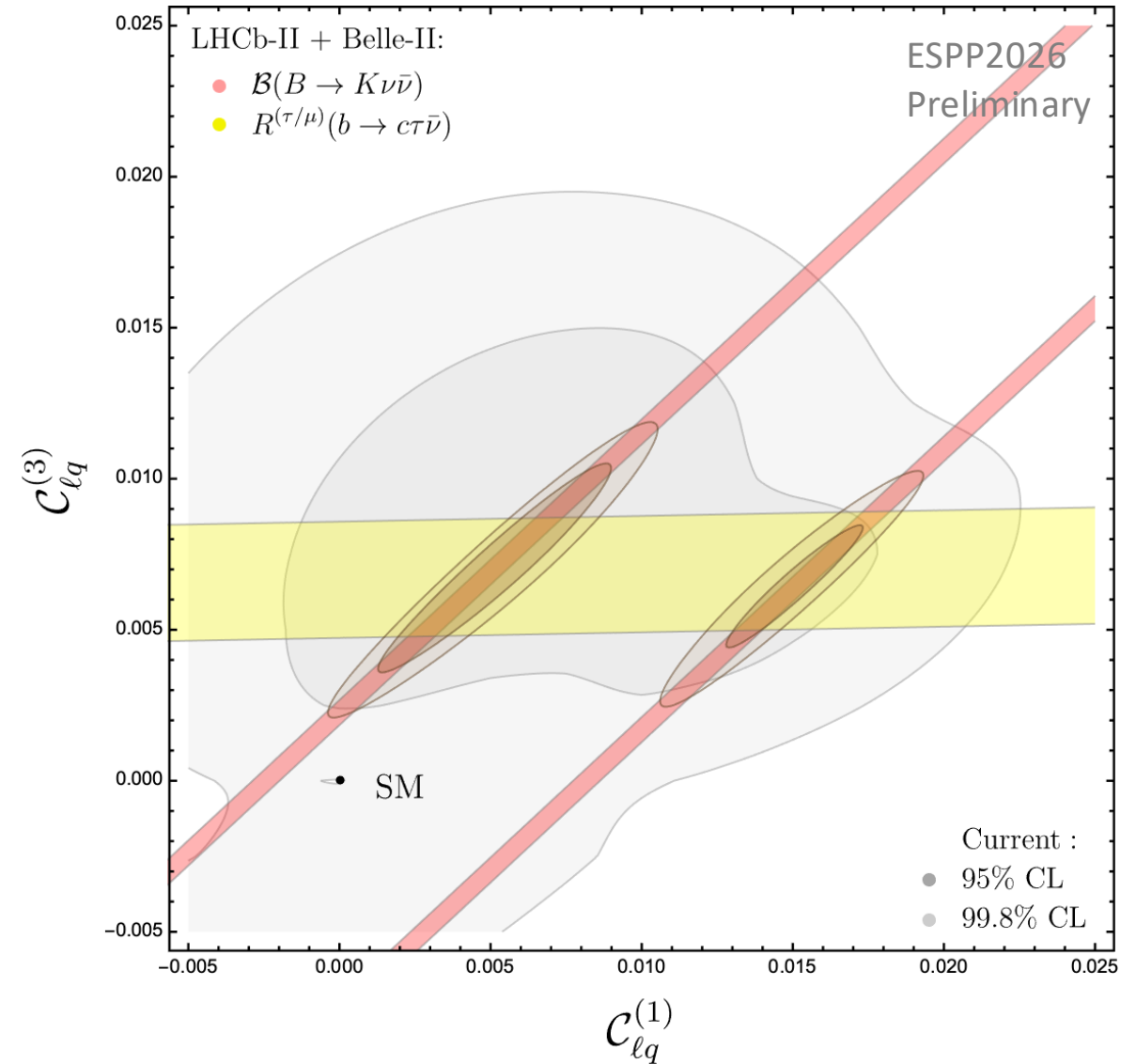
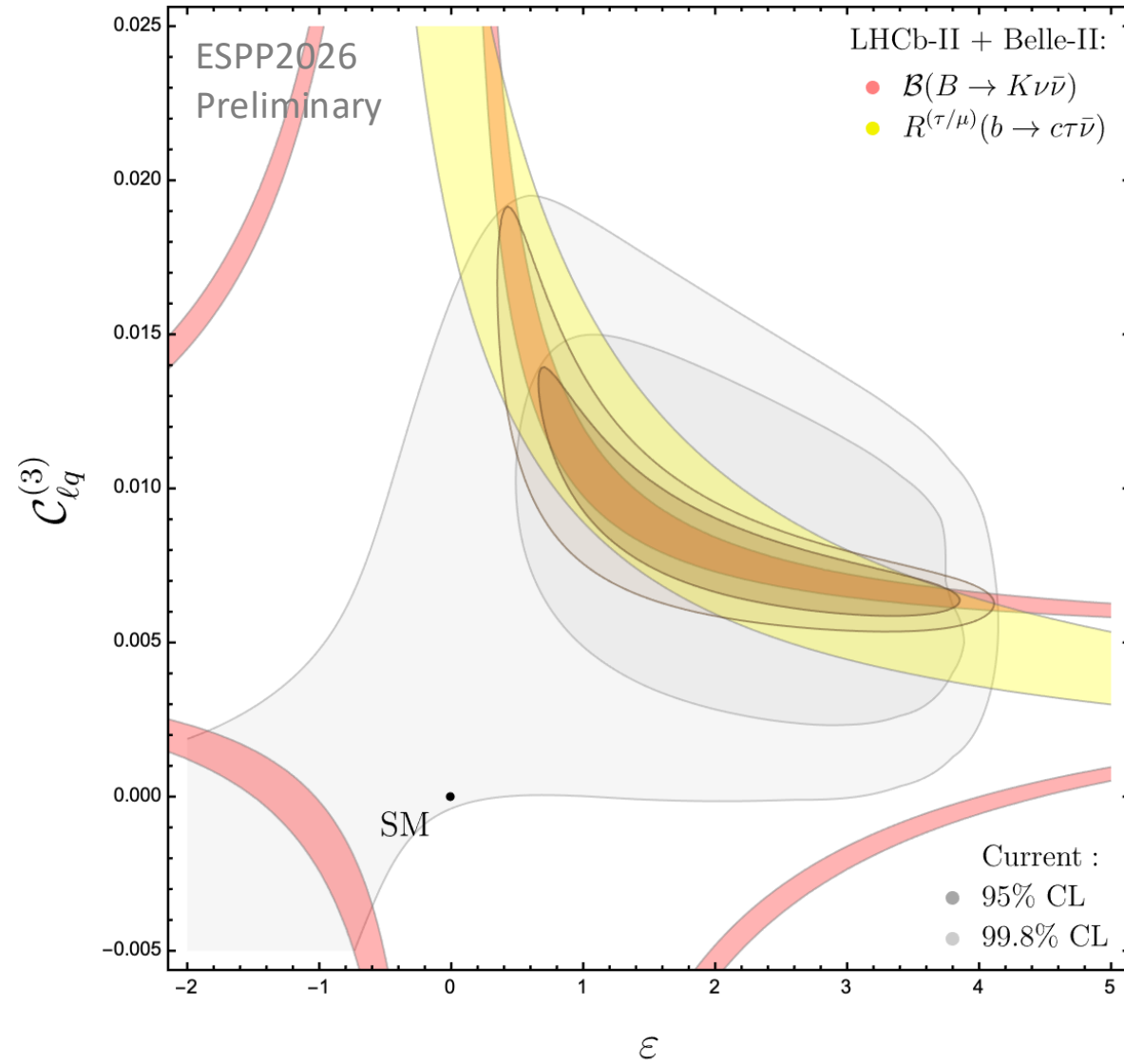
■ One flavour-mixing parameter, $\epsilon = O(1)$: $q_3 = b_L + \epsilon |V_{ts}| s_L$



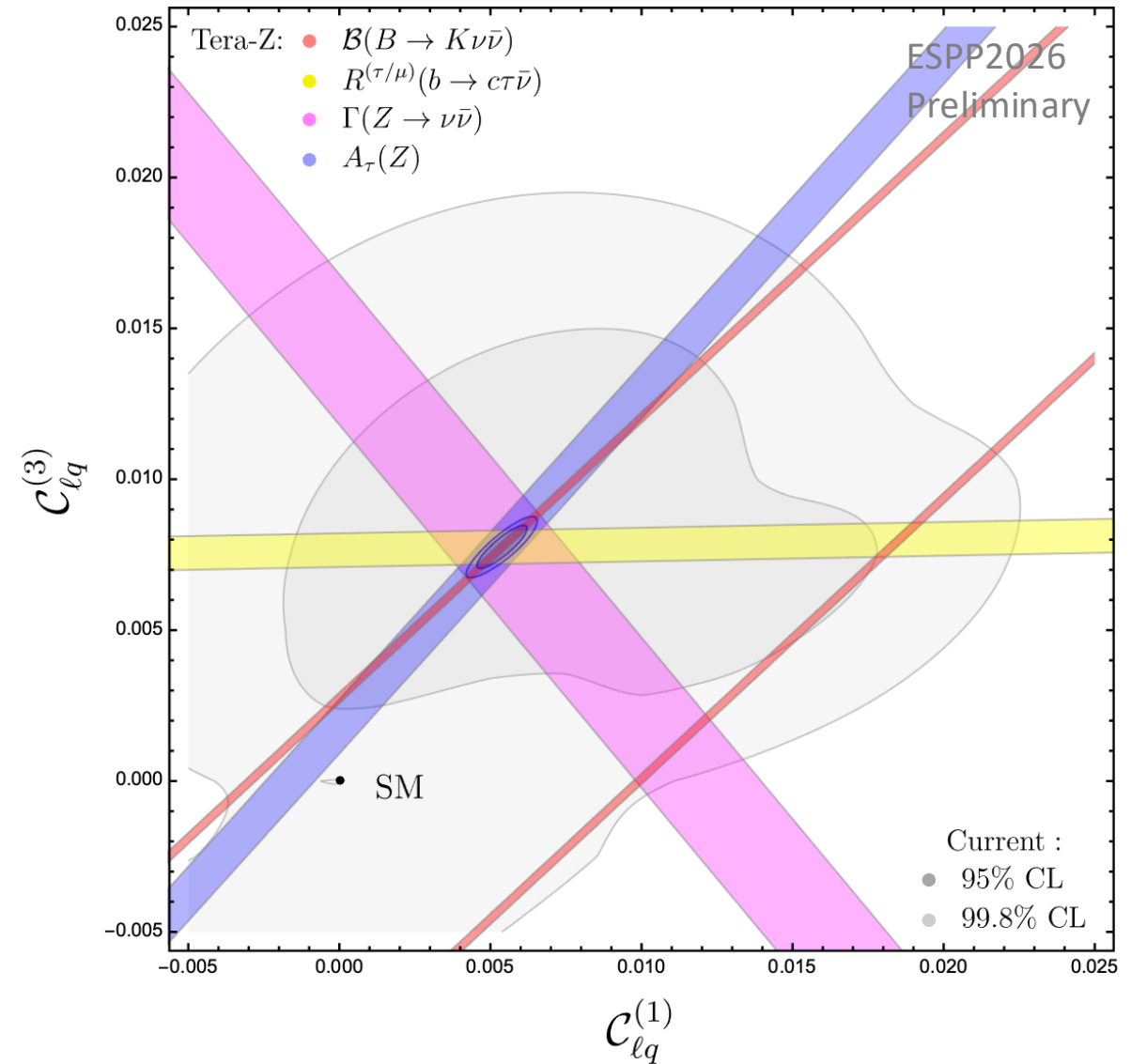
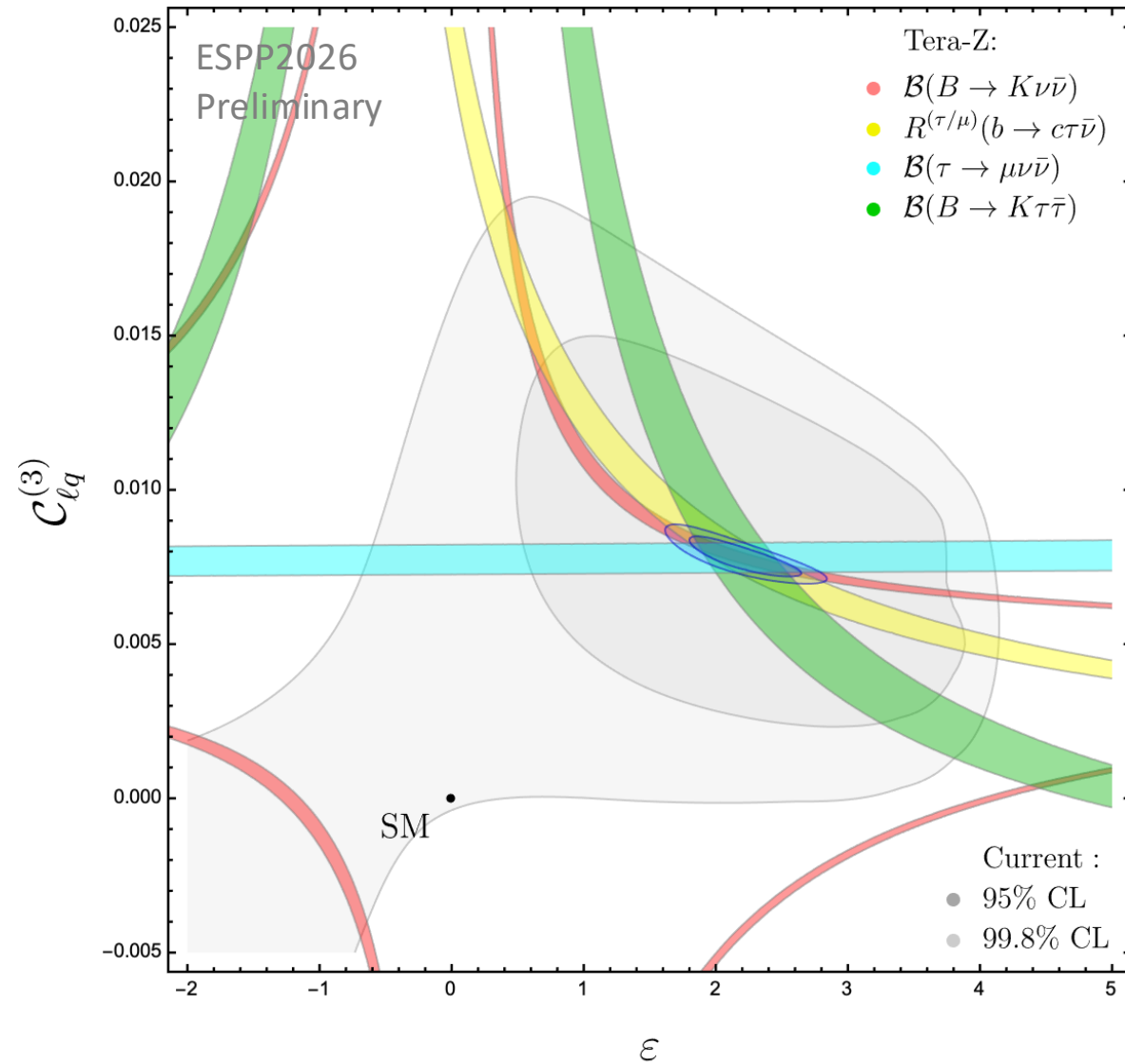
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About flavour, EW, and Higgs physics interplay

What I have illustrated is only an example, but the conclusions are very general:

The interplay of a large number of observables is a key ingredient if we aim to “**decode**” the underlying theory via indirect measurements and/or “**validate**” indirect BSM evidences (*in concrete models there are no hundreds parameters, but likely more than three...*) and flavour physics represent a particularly “rich” domain of NP sensitive observables.

Precision **Higgs**, **electroweak**, and **flavour physics** are three faces of the same strategy:

- for all of them the primary objective is indirect discovery
- in realistic models, with TeV-scale NP, they provide very complementary information