Bayes-ic Interactions -Bayesian analysis of nucleon-nucleon scattering data in pionless effective field theory

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Outline

- Motivation/Bayesian EFT Model Calibration
- BUQEYE Formalism
- Interaction Choice
- Results



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 Motivation/Bayesian EFT Model Calibration • **BUQEYE Formalism** Interaction Choice • Results





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nucleon-nucleon interactions.



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• Constrain only in relavent regimes











Microscopic nuclear interactions can be formulated using EFT.

LO

NLO (p/Λ_b)

 $N^{2}LO$ (p/Λ_b)

N³LO (p/Λ_b)

N⁴LO (p/Λ_b)

 $N^{5}LO$ $\left(p/\Lambda_b\right)^6$





Microscopic nuclear interactions can be formulated using EFT.

Potentials represented in expansions of $Q \equiv p/\Lambda_b < 1$

LO $\left(p/\Lambda_b\right)^0$

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Two nucleon



3N



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 <math>
 \left(p/\Lambda_b \right)^c$

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 ${
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Potentials represented in expansions of $Q \equiv p/\Lambda_b < 1$

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A specific organization is known as a power counting.

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$$\sum_{i=0}^{\infty} y^{(i)}Q^i$$

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where $y^{(i)}$ is calculated using a potential at order *i*, and *Q* is the parameter that orders the interaction.



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orders the interaction.

Truncating this tells us what we are neglecting: $y_k \neq 0$



$$\sum_{i=0}^{\infty} y^{(i)}Q^i$$











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What do I mean by this?



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 $y = y_{th}(\mathbf{a}) + \delta y$







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Quantity we $y = y_{th}(\mathbf{a}) + \delta y$ want to describe







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Quantity we **y** want to describe Model output







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Using Bayesian statistics:

Probability of a given data y

posterior



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data \mathbf{y} Probability of data \mathbf{y} given \mathbf{a}
likelihood

Probability of **a** given data **y** Pr

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Posterior



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Likelihood Prior

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How do we calculate the likelihood?



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$$\operatorname{pr}(\vec{y} \mid \vec{a}) \propto \exp\left[\left(\vec{y}_{th} - \vec{y}_{exp}\right)^T \Sigma^{-1} \left(\vec{y}_{th} - \vec{y}_{exp}\right)\right]$$



 $pr(\vec{a} | \vec{y}, I) \propto pr(\vec{y} | \vec{a}) pr(\vec{a} | I)$

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X +++



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$$y_{\text{th}}(x) = y_{\text{ref}}(x) \sum_{n=0}^{k} c_n Q^n + y_r$$







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$\delta y_{\text{th}}^{(k)}(x) + \delta y_{\text{th}}^{(k)}(x)$ < th `` *n=k*+1 Prediction

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$\operatorname{ref}(X)$ (\mathcal{X}) th *n*=*k*+1 Uncertainty Prediction

Discrepancy Model



Discrepancy Model

We can find an estimate for the effects of

$$\delta y_{\rm th}^{(k)}(x) =$$

by making a good set of assumptions.



11

Freets of $y_{ref}(x) \sum_{n=k+1}^{\infty} c_n Q^n$

Discrepancy Model

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• Statistics on c_n 's: $\mathbb{E}[c_n] = 0$, $\mathbb{E}[c_n] = 0$



11

$y_{ref}(x) \sum c_n Q^n$ n=k+1

$$[c_n c_m] = \bar{c}^2 \delta_{mn}$$
Discrepancy Model

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by making a good set of assumptions.

• Statistics on c_n 's: $\mathbb{E}[c_n] = 0$, $\mathbb{E}[c_n c_m] = \bar{c}^2 \delta_{mn}$

$$\Sigma_{ij}^{\text{th}} = \mathbb{E}\left[\delta_{\text{th},i}^{(k)}\delta_{\text{th},j}^{(k)}\right] - \mathbb{E}\left[\delta_{\text{th},i}^{(k)}\right]\mathbb{E}$$



 $\delta y_{\text{th}}^{(k)}(x) = y_{\text{ref}}(x) \sum_{n=1}^{\infty} c_n Q^n$ n=k+1

 $\left[\delta_{\text{th},j}^{(k)}\right] = \frac{\left(y_{\text{ref},i}\bar{c}Q_i^{k+1}\right)\left(y_{\text{ref},j}\bar{c}Q_j^{k+1}\right)}{1}$ $1 - Q_i Q_j$





For our models, we want to estimate the total probability distribution



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$\operatorname{pr}(\mathbf{a}, \overline{c}^2, \Lambda_b | \mathbf{y}, I) \propto \operatorname{pr}(\mathbf{y} | \mathbf{a}, \overline{c}^2, \Lambda_b, I) \operatorname{pr}(\mathbf{a} | I) \operatorname{pr}(\overline{c}^2 | \Lambda_b, \mathbf{a}, I) \operatorname{pr}(\Lambda_b | \mathbf{a}, I)$





pr(
$$\mathbf{a}, \bar{c}^2, \Lambda_b | \mathbf{y}, I$$
) $\propto \text{pr}(\mathbf{y} | \mathbf{a}, \bar{c}^2, \Lambda)$
 $\mathbf{x} e^{(\mathbf{y}_{exp} - \mathbf{y}_{th})^T \Sigma^{-1} (\mathbf{y}_{exp} - \mathbf{y}_{th})}$



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$\Lambda_b, I) \operatorname{pr}(\mathbf{a} | I) \operatorname{pr}(\overline{c}^2 | \Lambda_b, \mathbf{a}, I) \operatorname{pr}(\Lambda_b | \mathbf{a}, I)$





pr(**a**,
$$\bar{c}^2$$
, $\Lambda_b | \mathbf{y}, I$) \propto pr(**y** | **a**, \bar{c}^2 , Λ_b
× $e^{(\mathbf{y}_{exp} - \mathbf{y}_{th})^T \Sigma^{-1} (\mathbf{y}_{exp} - \mathbf{y}_{th})}$
Multivariate Gaussian



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$\Lambda_b, I) \operatorname{pr}(\mathbf{a} | I) \operatorname{pr}(\overline{c}^2 | \Lambda_b, \mathbf{a}, I) \operatorname{pr}(\Lambda_b | \mathbf{a}, I)$





 $\propto e^{(\mathbf{y}_{exp}-\mathbf{y}_{th})^{\mathrm{T}}\Sigma^{-1}(\mathbf{y}_{exp}-\mathbf{y}_{th})}$ Multivariate Gaussian ~ $\mathcal{N}(\mu, \Sigma_{\text{prior}})$



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 $\mathbf{x} e^{(\mathbf{y}_{exp} - \mathbf{y}_{th})^{\mathrm{T}} \Sigma^{-1} (\mathbf{y}_{exp} - \mathbf{y}_{th})}$ Multivariate Gaussian $(\mu, \Sigma_{\text{priot}})$



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 $\mathbf{x} e^{(\mathbf{y}_{exp} - \mathbf{y}_{th})^{\mathrm{T}} \Sigma^{-1} (\mathbf{y}_{exp} - \mathbf{y}_{th})}$ Multivariate Gaussian



For our models, we want to estimate the total probability distribution











 $\mathbf{x} e^{(\mathbf{y}_{exp} - \mathbf{y}_{th})^{\mathrm{T}} \Sigma^{-1} (\mathbf{y}_{exp} - \mathbf{y}_{th})}$ **Multivariate Gaussian**



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$\operatorname{pr}(\mathbf{a}, \bar{c}^2, \Lambda_b | \mathbf{y}, I) \propto \operatorname{pr}(\mathbf{y} | \mathbf{a}, \bar{c}^2, \Lambda_b, I) \operatorname{pr}(\mathbf{a} | I) \operatorname{pr}(\bar{c}^2 | \Lambda_b, \mathbf{a}, I) \operatorname{pr}(\Lambda_b | \mathbf{a}, I)$ $(\nu, \tau^2) \propto \frac{\operatorname{pr}(\Lambda_b | I)}{1}$ $\sim \chi^{-2}$ Hyperparameters found in order-by-order calculations









 $\mathbf{x} e^{(\mathbf{y}_{exp} - \mathbf{y}_{th})^{\mathrm{T}} \Sigma^{-1} (\mathbf{y}_{exp} - \mathbf{y}_{th})}$ **Multivariate Gaussian**



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The hyperparameters come out of the order-by-order analysis:

 $c_{n,i} =$



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$$\nu_0 + N_{\rm obs} n_c$$



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$$\nu_0 \tau_0 + \sum_{n,i} c_{n,i}^2$$



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Results







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- Degrees of freedom
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- Representation
- Regularization scheme





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$y_{\text{th}}(x) = y_{\text{ref}}(x) \sum c_{2n}(x)Q^{2n}(x)$ n=0





$y_{\text{th}}(x) = \underbrace{\overline{y_{\text{ref}}(x)}}_{y_{\text{exp}}} \sum_{n=0}^{\infty} c_{2n}(x) Q^{2n}(x)$

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Our interaction takes the form:



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$$v_{\text{NLO}}^{\text{CI}}(\vec{k},\vec{K}) = C_1 k^2 + C_2 k^2 \sigma_1 \cdot \sigma_2 + C_3 S_{12}(k) + C_4 k^2 \tau_1 \cdot \tau_2$$
$$+ i C_5 \vec{S} \cdot (\vec{K} \times \vec{k}) + C_6 k^2 \tau_1 \cdot \tau_2 \sigma_1 \cdot \sigma_2 + C_7 S_{12}(k) \tau_2 \cdot \tau_2$$



$$y_{\text{th}}(x) = \underbrace{y_{\text{ref}}(x)}_{\text{ref}} \sum_{n=0}^{\infty} c_{2n}(x) Q^{2n}(x)$$
$$y_{\text{exp}}(x) \quad n=0$$

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$$+ i C_5 \vec{S} \cdot (\vec{K} \times \vec{k}) + C_6 k^2 \tau_1 \cdot \tau_2 \sigma_1 \cdot \sigma_2 + C_7 S_{12}(k) \tau_2 \cdot \tau_2$$

 $v_{\text{NLO}}^{\text{CD}} = C_0^{\text{TT}} T_{12} + C_0^{\text{TV}} (\tau_{1z} + \tau_{2z})$



$$y_{\text{th}}(x) = \underbrace{y_{\text{ref}}(x)}_{\text{ref}} \sum_{n=0}^{\infty} c_{2n}(x) Q^{2n}(x)$$
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We employ a Gaussian cutoff in coordinate space, which smears δ -functions upon Fourier transformation

 $f(r) = -\frac{1}{\pi^2}$



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$$\frac{1}{\frac{3}{2}R_s^3}e^{-\left(\frac{r}{R_s}\right)^2}$$

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We employ a Gaussian cutoff in coordinate space, which smears δ -functions upon Fourier transformation

 $f(r) = -\frac{\pi^2}{\pi^2}$

We choose $R_s \in [1.5, 2.0, 2.5]$ fm which are $\sim \frac{400}{R_s}$ MeV in momentum space.



To use these interactions, they must be regularized in some fashion and may be

$$\frac{1}{3/2R_s^3}e^{-\left(\frac{r}{R_s}\right)^2}$$





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MeV + deuteron binding energy + nn scattering length.



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- Our choice of data is the pp and np Granada database (251 differential cross sections, 133 total cross sections, 4 polarized cross sections) up to 5
- We then use Markov Chain Monte Carlo (MCMC) to sample the posteriors at LO (Q^0) , NLO (Q^2) , and N3LO (Q^4) , allowing for the order-by-order convergence analysis for LO \rightarrow NLO and NLO \rightarrow N3LO to estimate \bar{c} and Λ_{h} .



Prior Choices

 $\mathbf{pr}(\vec{a} \mid I) \sim \mathcal{N}\left(\vec{a}_{p.s.}^{MAP}, \overrightarrow{10^2}\right)$

- $pr(\Lambda_h | I) \sim \mathcal{N}(500 \text{ MeV}, 1000^2)$
- $\operatorname{pr}(\bar{c}^2 | \mathbf{I}) \sim \chi^{-2}(\nu_0 = 1.5, \tau_0^2 = 1.5)$
- $r(x_i, x_j; \vec{l}) = e^{|p_i p_j|/2l_p} e^{|\theta_i \theta_j|/2l_\theta} \delta_{\text{type}_i, \text{type}_j}, \quad l_p = 0.3 \text{ MeV}, \ l_\theta = 20^\circ$
- $p_{\rm soft} = \begin{cases} p_d \sim 45 \; {\rm MeV}/c, & {\rm for} \; np \; {\rm scattering} \\ 1/{}^1 a_{\rm pp} \sim 25 \; {\rm MeV}, & {\rm for} \; pp \; {\rm scattering} \, . \end{cases}$



$$MeV^2$$
)

Outline

- Motivation/Bayesian EFT Model Calibration
- BUQEYE Formalism
- Interaction Choice
- Results











More figures







(S, T) = (0, 1)

More figures





(S, T) = (1, 0)

More figures





More figures

2.5 fm \bar{c} and Λ_b Posteriors







2.0 fm \bar{c} and Λ_b Posteriors









1.5 fm \bar{c} and Λ_b Posteriors



















Why is there dependence on the order?







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• Power-counting?







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- Power-counting?
- Flawed assumption of geometric series?







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For our analysis, we used naïve dimensional analysis for the powercounting, which is known to be problematic.







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JB et al. Phys. Rev. C **111**, 034005 (2025)







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Ekström and Platter, Phys.Lett.B 860 139207 (2025)







$$pr(\mathbf{y}_{th} | \mathbf{y}, \mathbf{x}, I) = \int d\mathbf{a} \, d\bar{c}^2 \, d\Lambda_b \, \mathcal{N} \left(\mathbf{y}_{th} (\mathbf{a}, \mathbf{x}), \Sigma_{th} \left(\bar{c}^2, \Lambda_b \right) \right) \times pr(\mathbf{a} | \mathbf{y}, I) pr(\bar{c}^2 | \Lambda_b, \mathbf{a}, I) pr(\Lambda_b | \mathbf{a}, I),$$







95 % Confidence interval $\rightarrow \sim 150$ % uncertainty $pr(\mathbf{y}_{th} | \mathbf{y}, \mathbf{x}, I) = \int d\mathbf{a} \, d\bar{c}^2 \, d\Lambda_b \, \mathcal{N} \left(\mathbf{y}_{th}(\mathbf{a}, \mathbf{x}), \Sigma_{th}\left(\bar{c}^2, \Lambda_b\right) \right) \times$ $pr(\mathbf{a} | \mathbf{y}, I) pr(\bar{c}^2 | \Lambda_b, \mathbf{a}, I) pr(\Lambda_b | \mathbf{a}, I),$







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> • Poorly constrained *d*-waves • 2b corrections at $O(Q^5)$ <u>
> → SHOULD BE CONSISTENT</u> <u>WITH 0</u>



Applications to Few-Body Systems





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- What is the momentum scale in bound systems?
 - Easy for chiral interactions: $p \approx m_{\pi}$
 - Not so easy for pionless interactions
- three-nucleon interactions are important, and we do not have uncertainty quantification for these







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interaction.

• Re-counting the index of interactions





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- How do we fix this?



Summary

- EFT interactions provide a systematic framework for microscopic nuclear interactions.
 - Enable rigorous uncertainty quantification.
- Bayesian analysis is a useful tool in EFT contexts.
 - Uncertainty quantification
 - Extract information of physics from data
 - E.g., inconsistencies in EFT power counting

"Bayesian analysis is like having a discussion with data."

What can the data tell us about the physics we want to explore?



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