



# The need for accurate nuclear cross sections for solar (and stellar) modeling

Marciana 2025 – Lepton Interactions with Nucleons & Nuclei  
Marciana Marina – June 22<sup>nd</sup>-27<sup>th</sup> - 2025

A. Serenelli



This project has received funding  
from the European Union's  
Horizon 2020 research and  
innovation programme under  
grant agreement No 101008324  
(ChETEC-INFRA).

The Sun as a cornerstone for astrophysics

(Which?) Solar composition

a view from helioseismology

a view from solar neutrinos

the need for accurate & precise nuclear cross sections

Solar-like stars

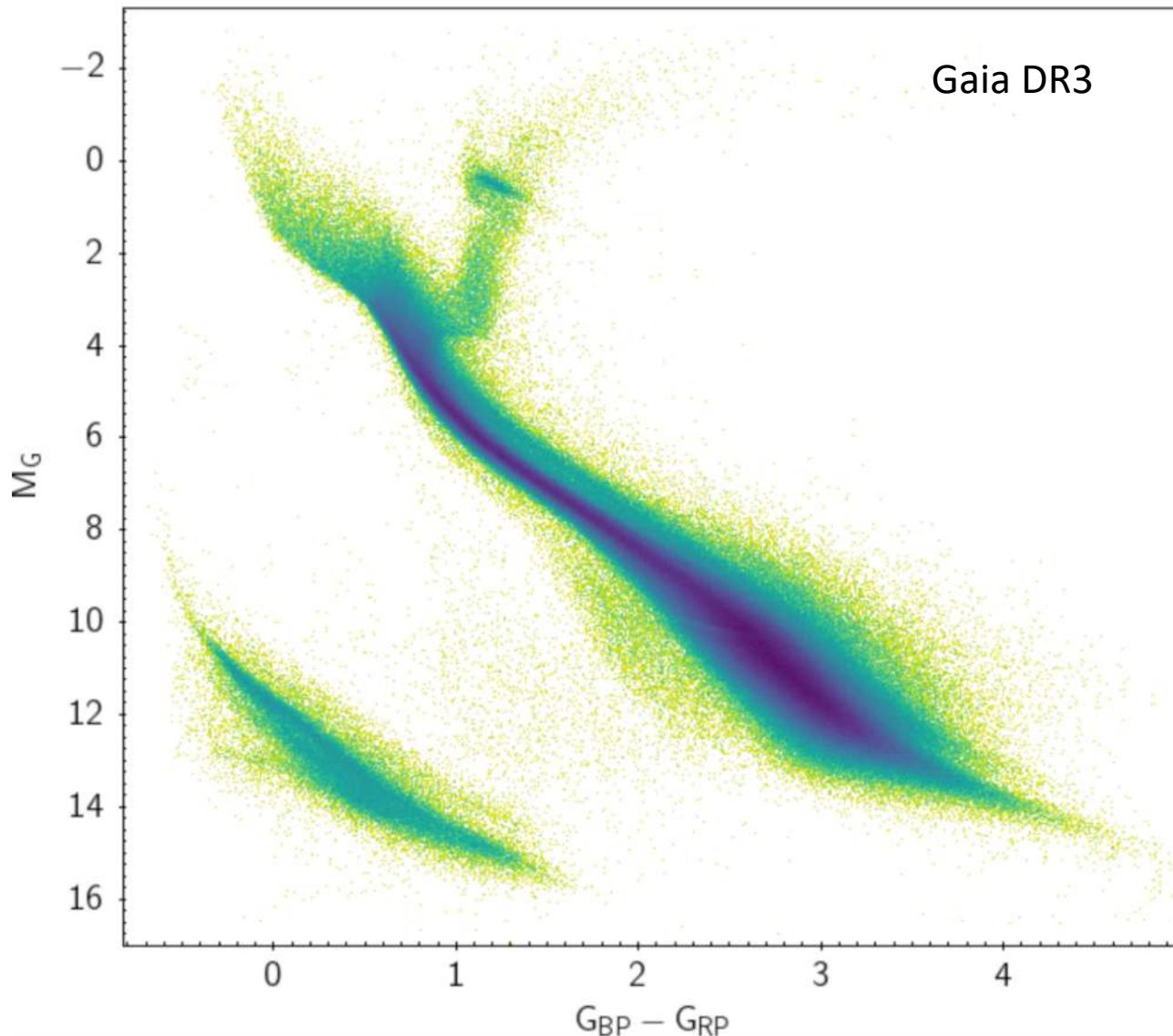
convective cores in  $1.1 - 1.5 M_{\odot}$  stars in the context of asteroseismic missions (PLATO)

the need for accurate & precise  $^{14}\text{N}+\text{p}$  cross section

# Why the Sun? It is “foundation” science



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~ $10^9$  individual stars with measurements  
colors, temperature, luminosity, (composition)

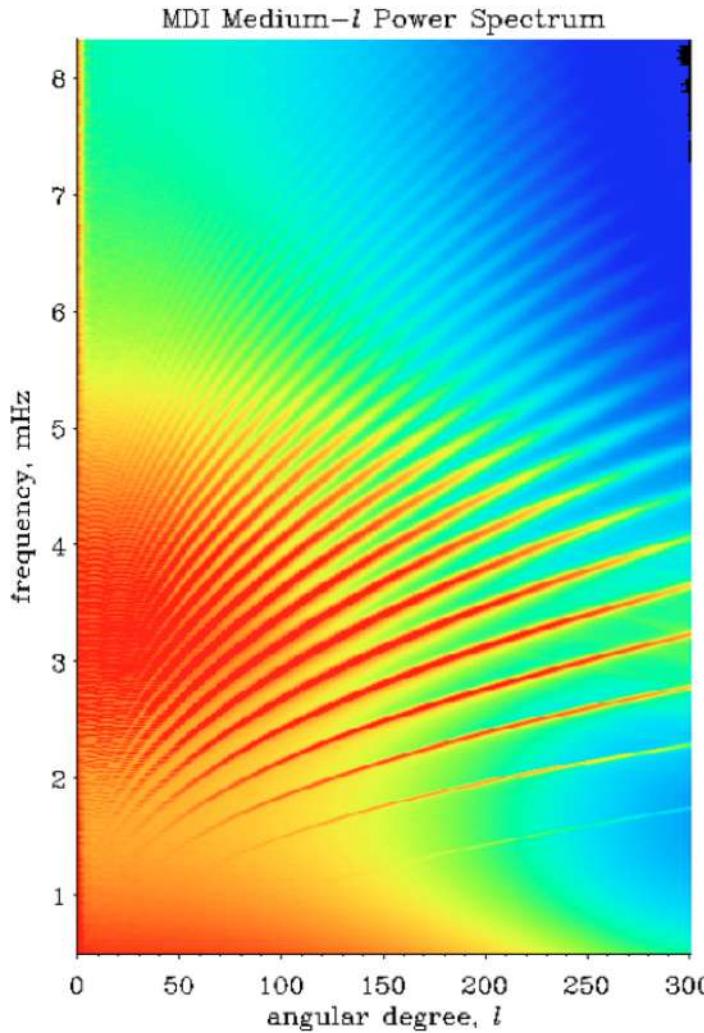
~ $10^3$  with accurate, precise, (model) independent  
mass determinations  
selective club: eclipsing binaries

**1 star with accurate, precise, (model) independent  
age determination**  
**meteoritic dating**  
**+ highly accurate radius & mass**

# Why the Sun? It is “foundation” science



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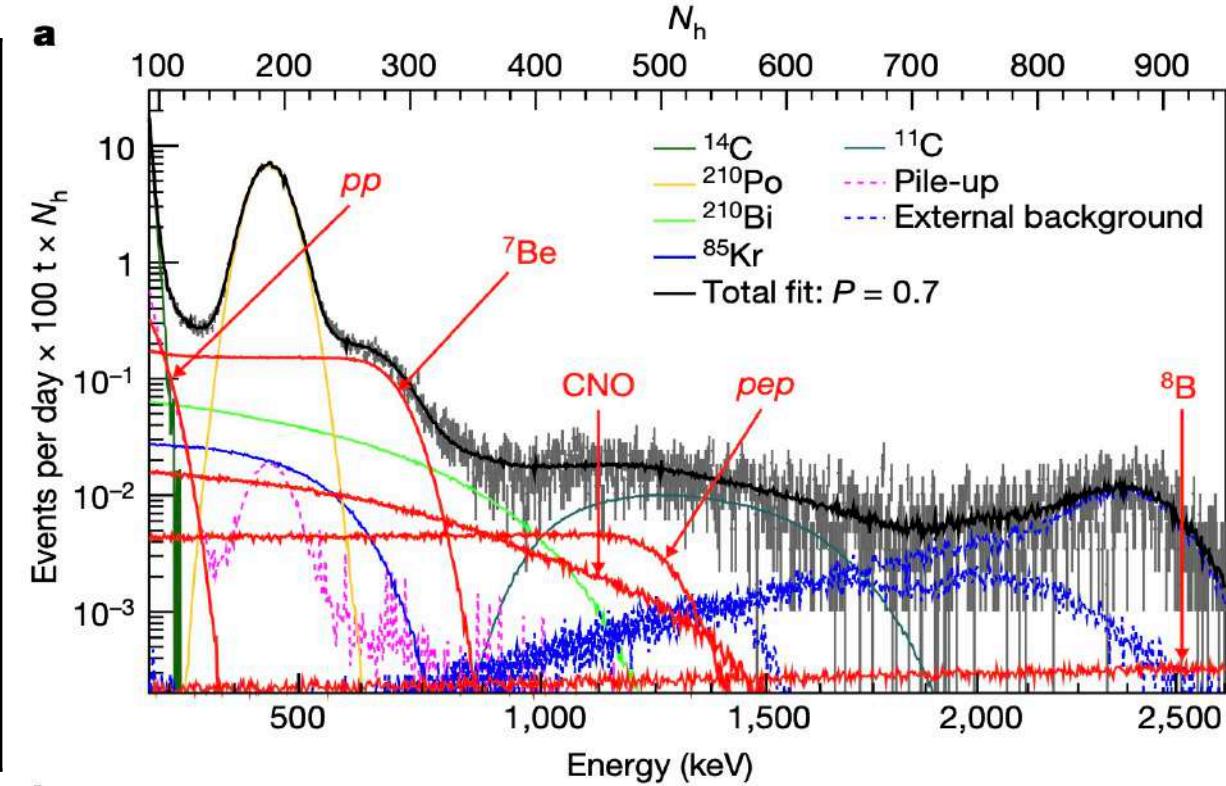
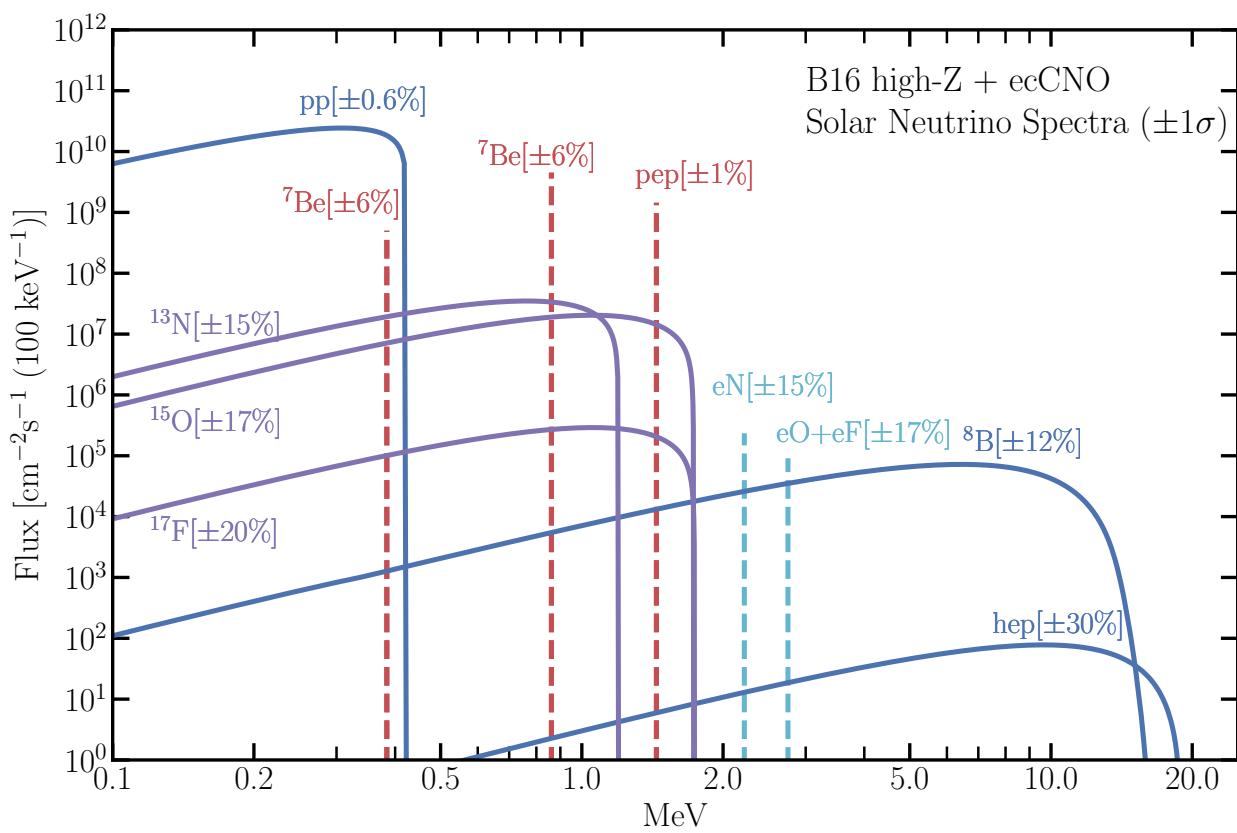
## Helioseismology

- >10<sup>5</sup> eigenmodes → inversion of internal structure:  
sound speed, density, adiabatic index (EoS)
- global quantities:  
surface helium, depth of convective envelope
- beyond standard solar models:  
internal rotation profile (depth and latitude)

**Allows testing theory of stellar evolution by looking at internal structure**

# Why the Sun? It is "foundation" science

**Solar neutrinos → information on solar core, nuclear physics**

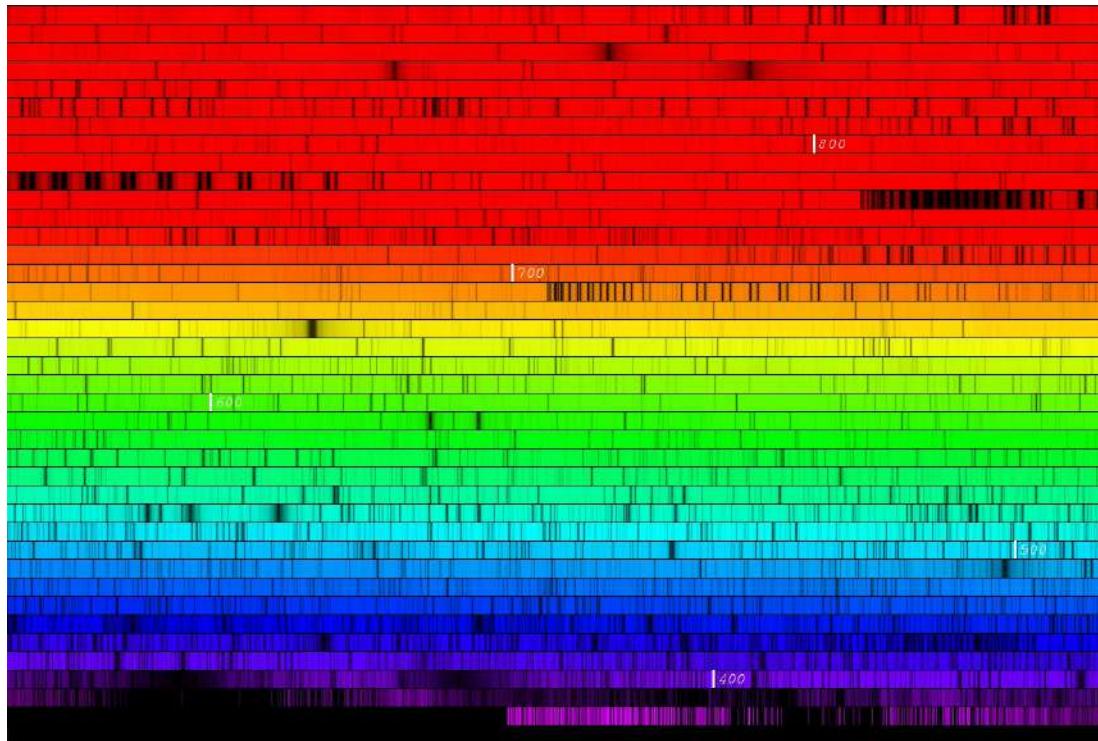


Borexino solar neutrino spectrum  
and identified solar fluxes

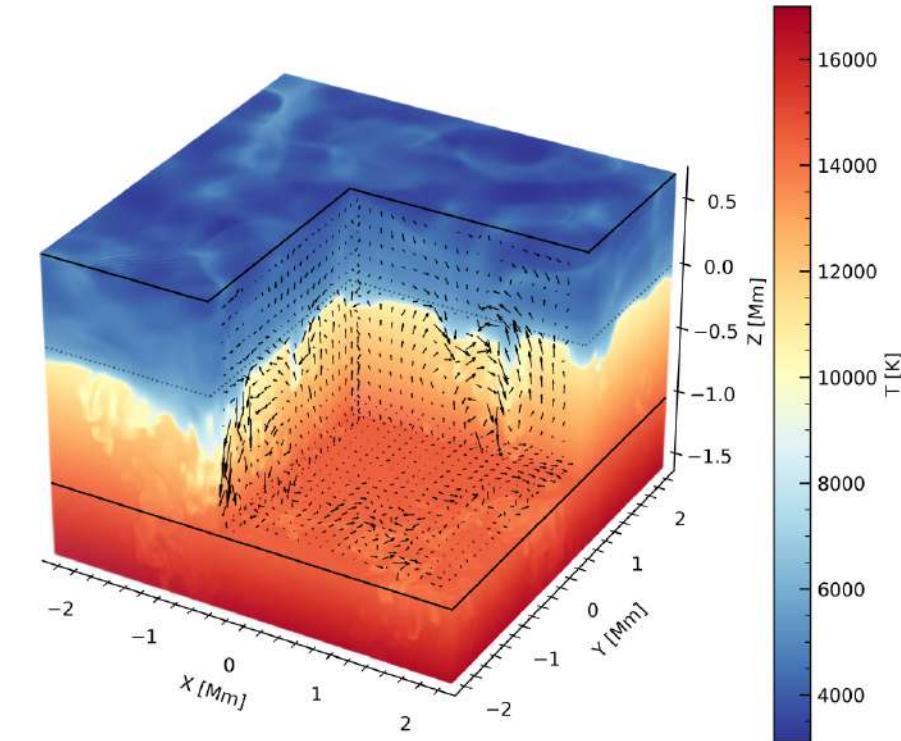
# Foundation science: Solar spectrum & abundances



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Solar envelope is convective  
→ hydrodynamic models  
→ 3D atmosphere model



Model atmosphere  
→ detailed radiative transfer  
→ synthetic spectrum to compare with observed one  
→ determination of abundances

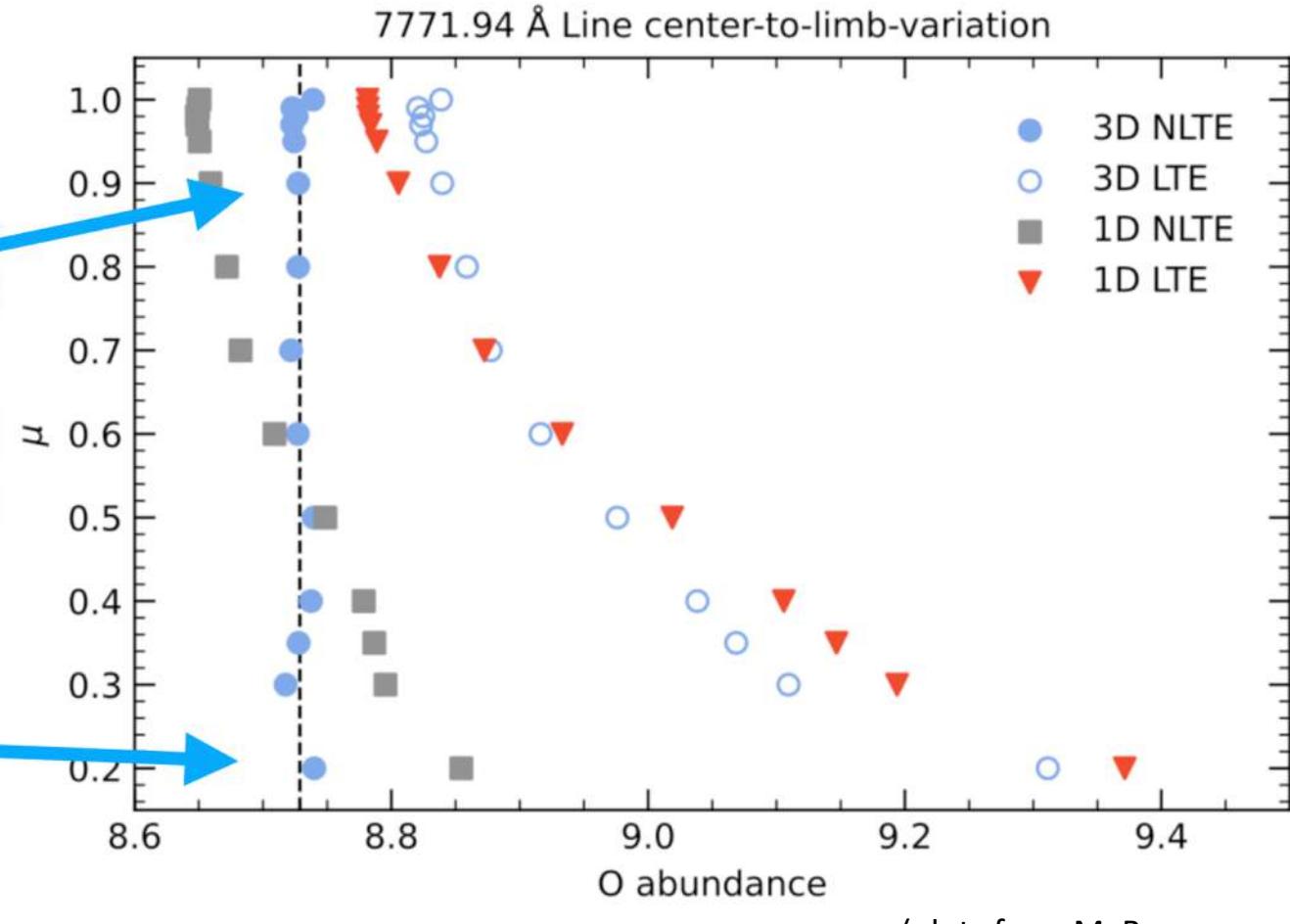
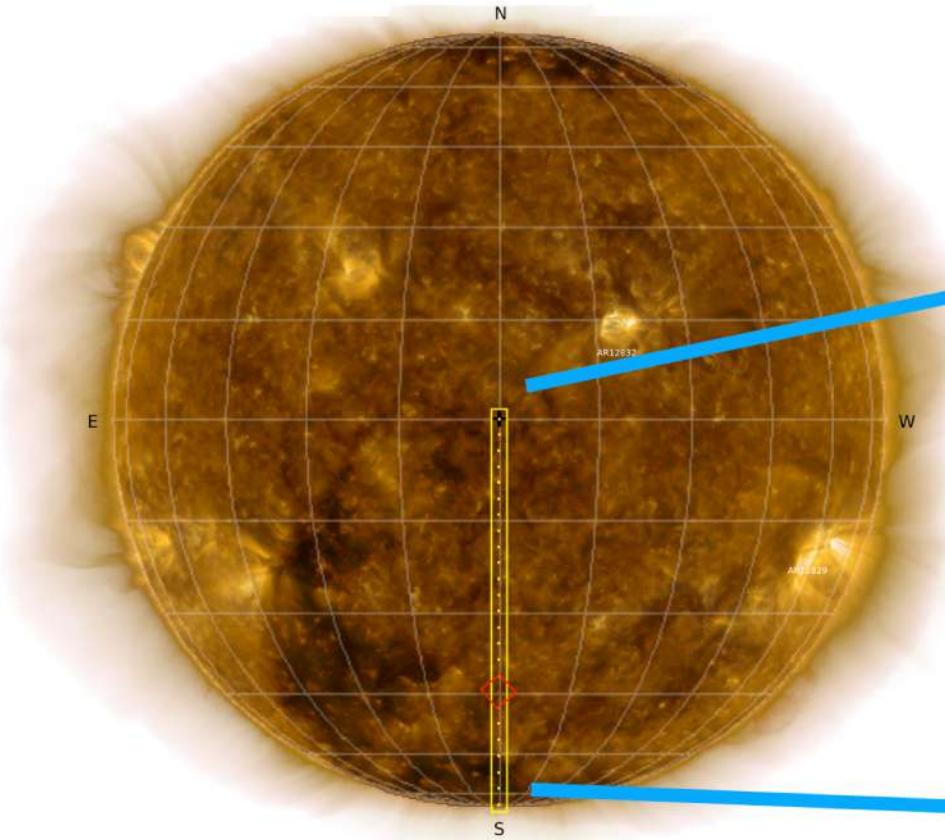
# Foundation science: Solar spectrum & abundances



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Only star that allows detailed tests, e.g. center-to-limb variations

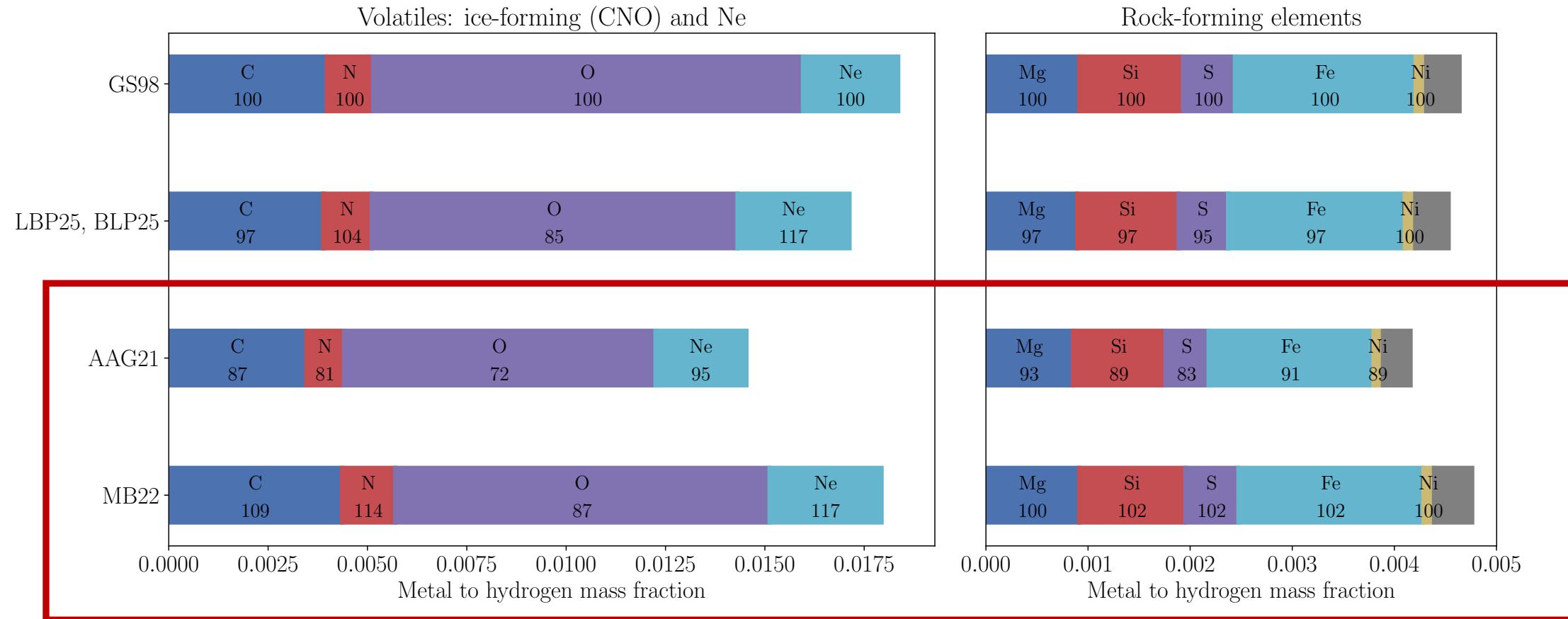


Pietrow, Hoppe, Bergemann et al. 2023

Bergemann, Hoppe, et al. 2021

(plots from M. Bergemann)

# Which solar composition?



GS98: Grevesse & Sauval 1998

LBP25/BLP25: Lodders, Bergemann, Palme 2025

AAG21: Asplund et al. 2021,

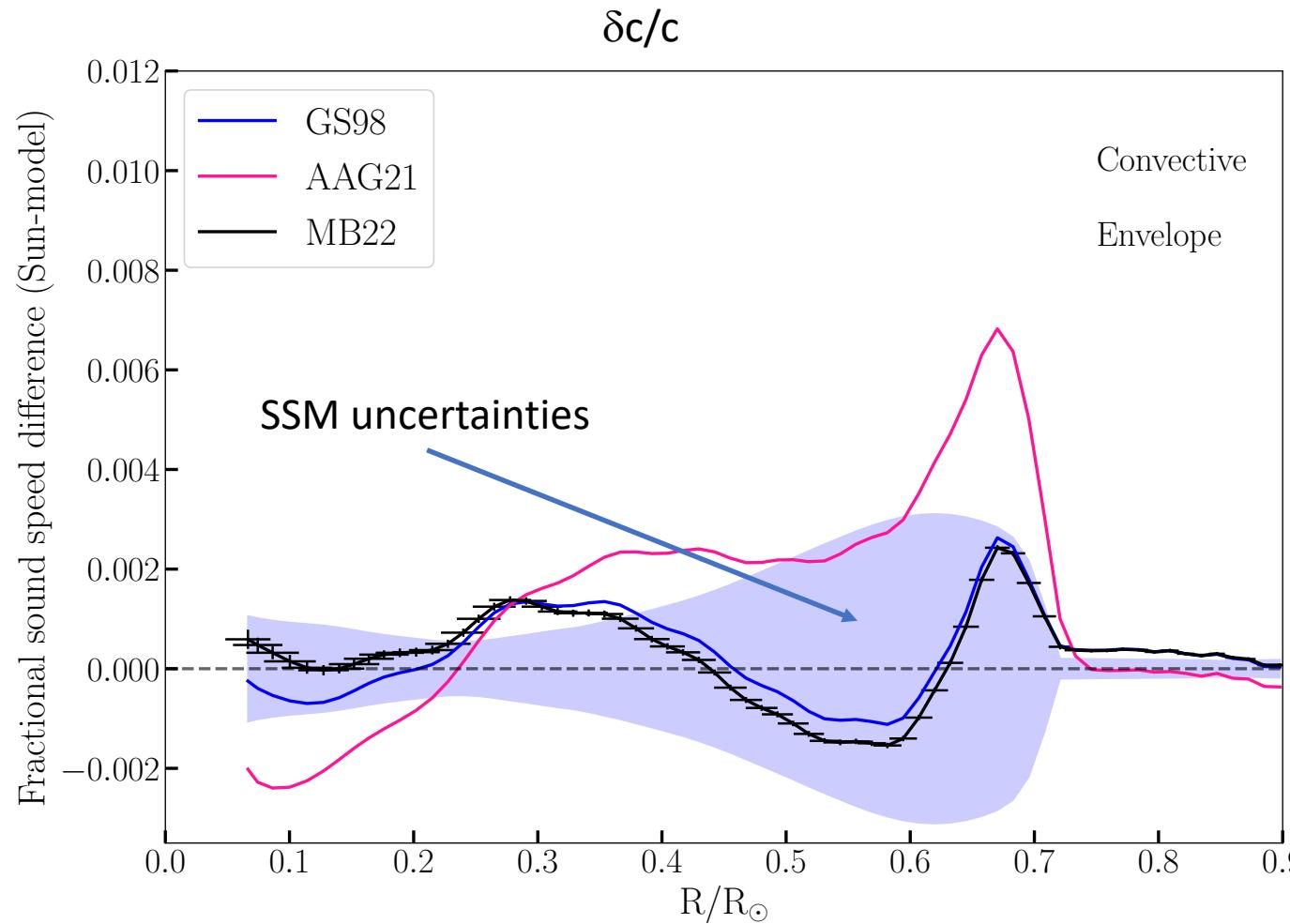
MB22: Magg et al. 2022

**Chemical abundances are a constraint, not a prediction, of (non-) standard solar models**

# What helioseismology tells us



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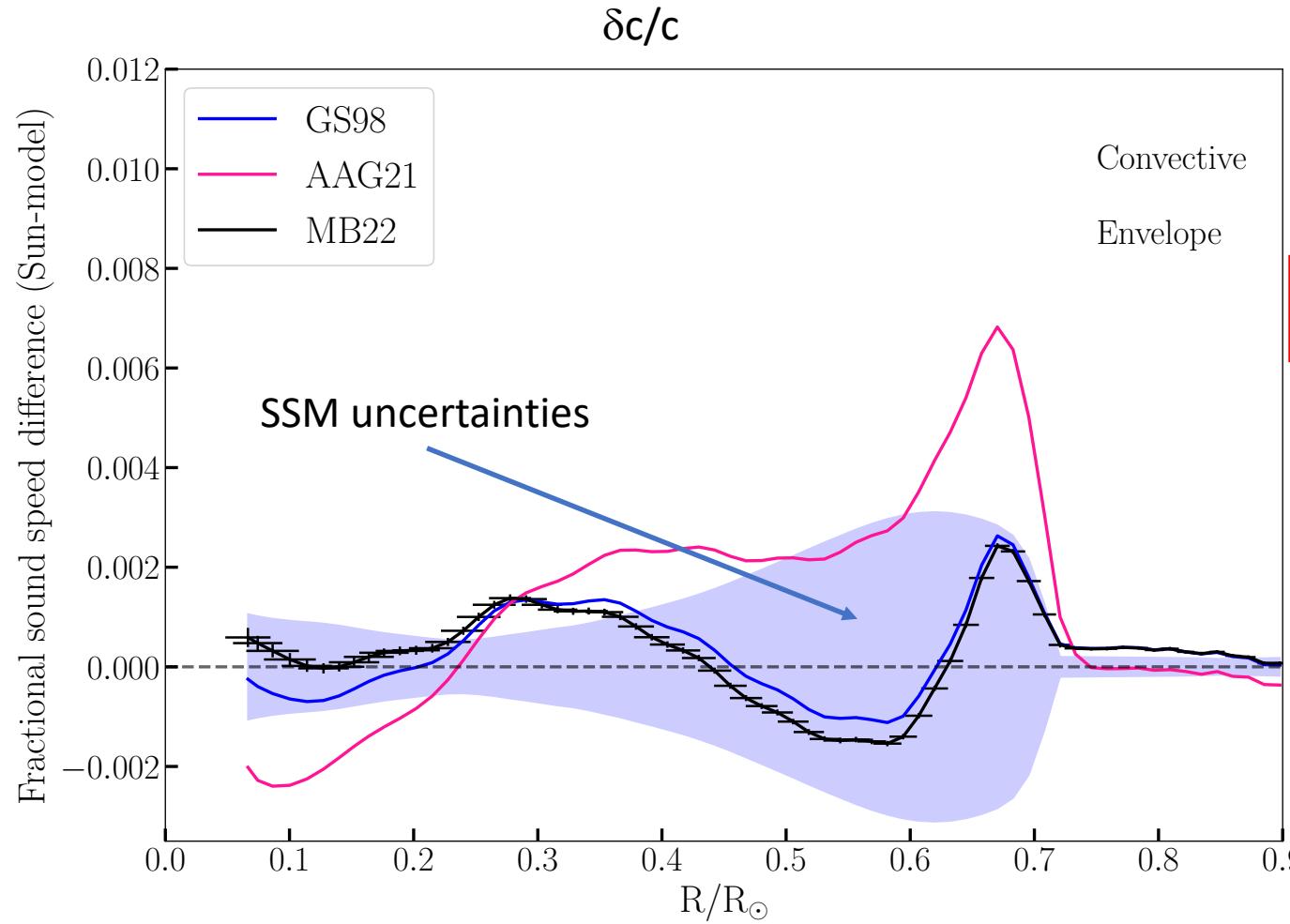


Model	$R_{\text{CZ}}/R_\odot$	$Y_S$	$\langle \delta c/c \rangle$	$Y_{\text{ini}}$	$Z_{\text{ini}}$
MB22	0.7123	0.2439	0.0010	0.2734	0.0176
AAG21	0.7197	0.2343	0.0027	0.2638	0.0155
GS98	0.7122	0.2425	0.0010	0.2718	0.0187
Solar	<b>0.713</b> <b><math>\pm 0.001</math></b>	<b>0.2485</b> <b><math>\pm 0.0035</math></b>			

# What helioseismology tells us



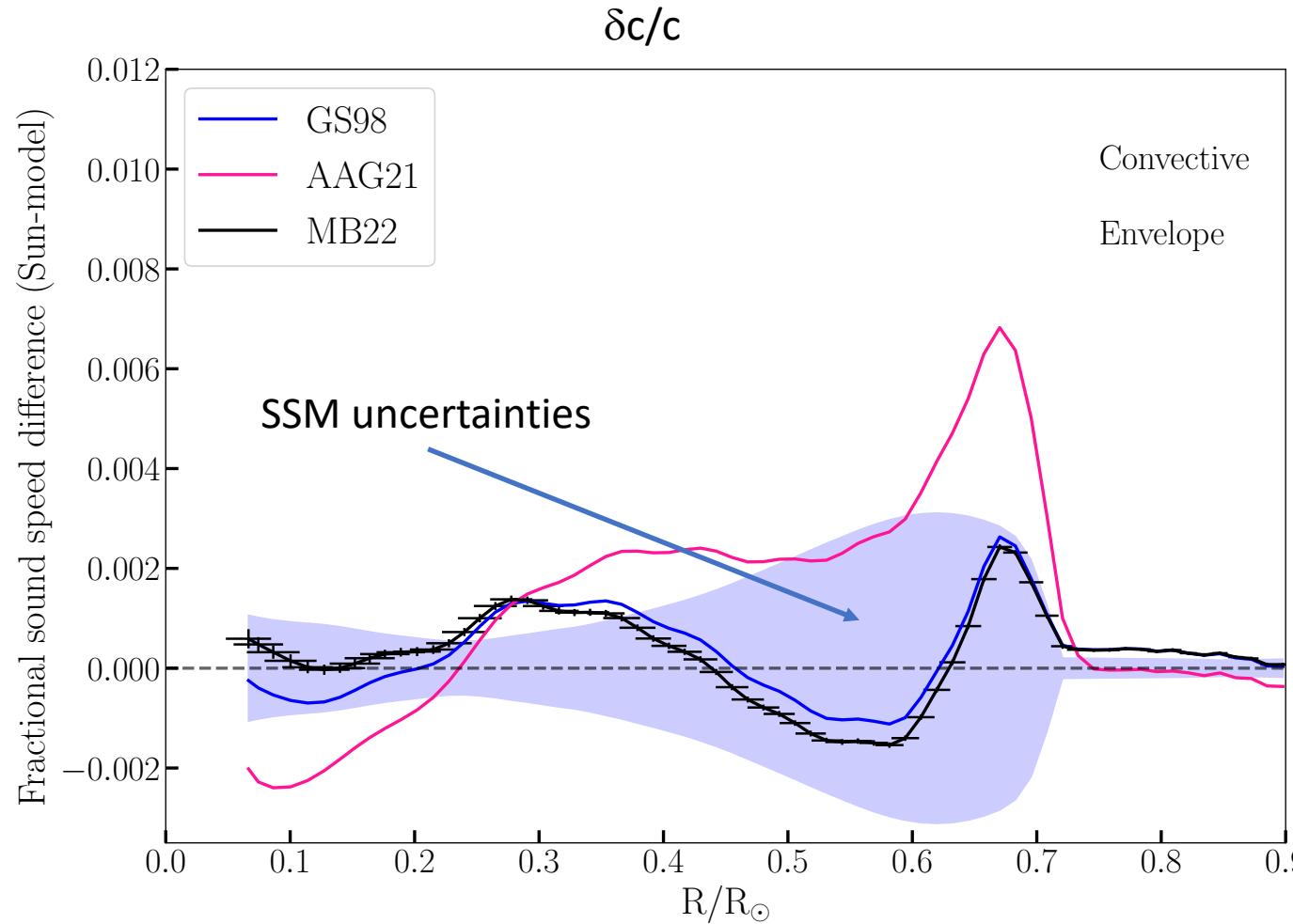
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# What helioseismology tells us



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Results sensitive to thermal structure  
because sound speed scales with  $T^{1/2}$

$$\nabla T^4 \propto \kappa$$

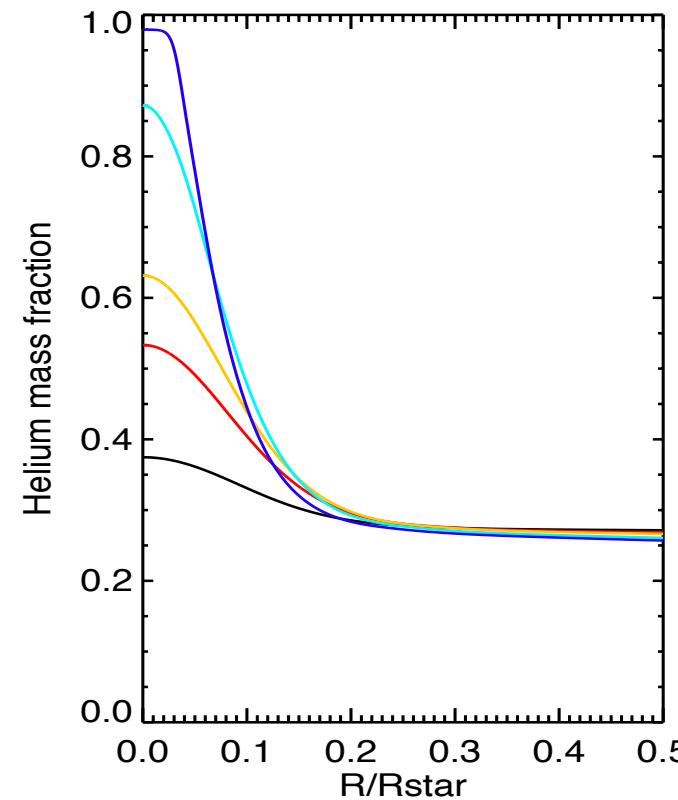
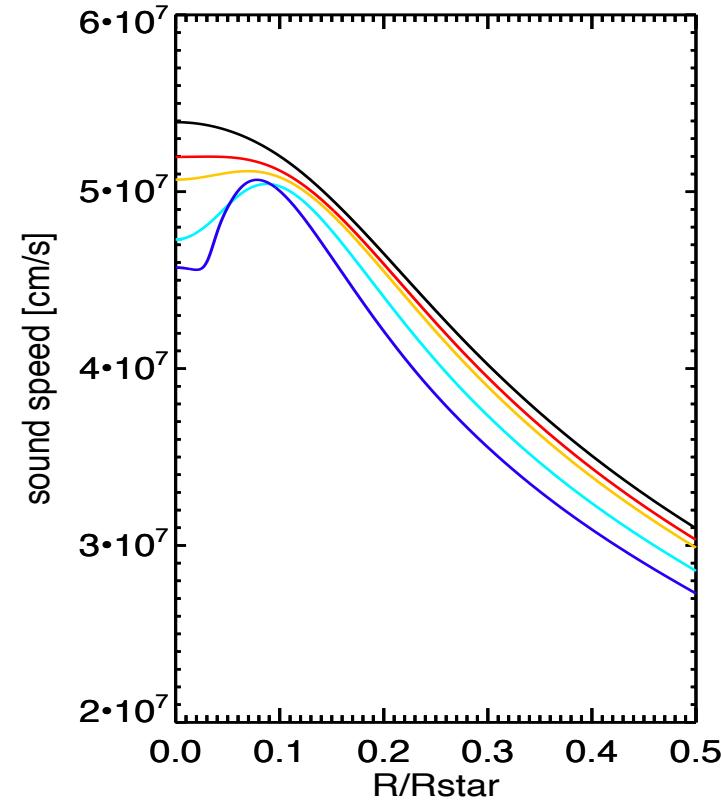
→ (composition + radiative opacities)

$$c^2 \propto \frac{T}{\mu}$$

# Dating the Sun "as a star"

Cancellation effects limit modes to  $l=0, 1, 2, (3)$  for other stars (e.g. Kepler, TESS, PLATO)

$$\nu_{n,\ell} - \nu_{n-1,\ell+2} \propto \frac{1}{4\pi\nu_{n,\ell}} \int_0^R \frac{dc}{dr} dr \quad \longrightarrow \text{age diagnostics}$$

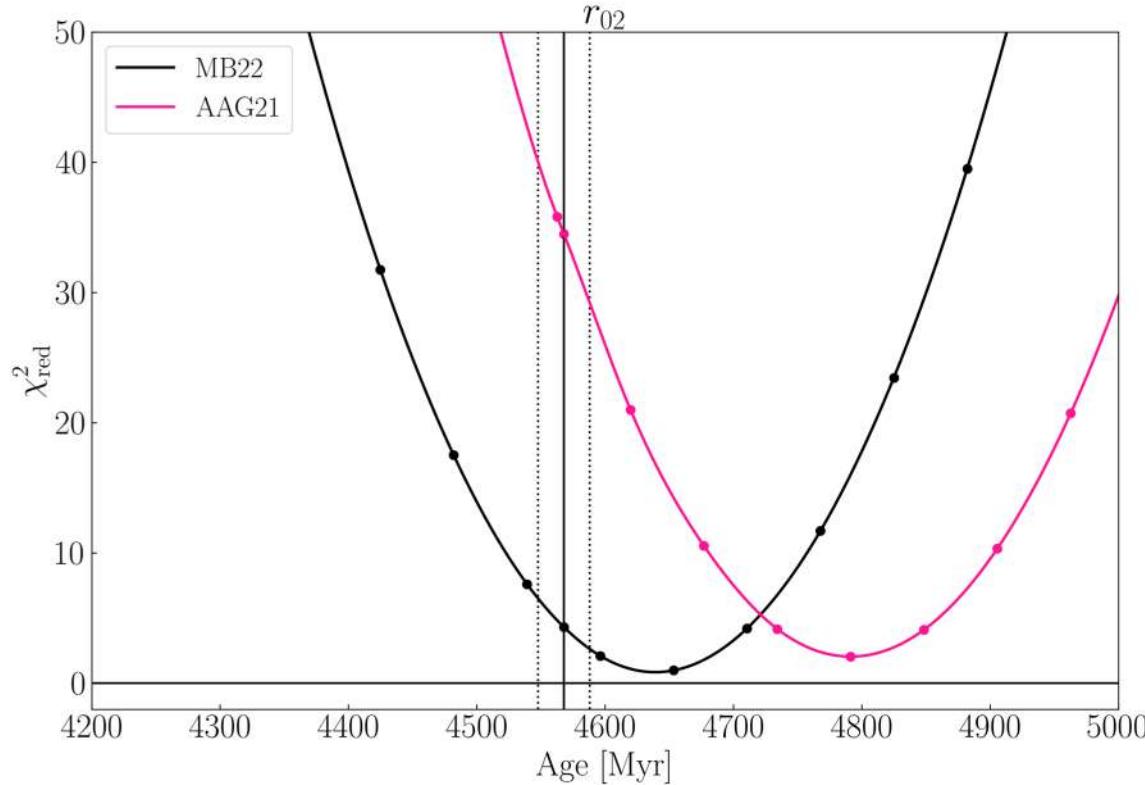


$$c^2 \propto \frac{T}{\mu}$$

# The Sun from afar

No independent age for other stars

$$\nu_{n,\ell} - \nu_{n-1,\ell+2} \propto \frac{1}{4\pi\nu_{n,\ell}} \int_0^R \frac{dc}{dr} \frac{dr}{r}$$



	Solar age (Gyr)	$\chi^2$ (33 dofs)
Sun	<b><math>4.568 \pm 0.020</math></b>	---
AAG21	$4.755 \pm 0.034$	76.6
MB22	$4.611 \pm 0.032$	38.4

**Composition introduces a systematic effect on age determination of about 250 Myr (5%)**

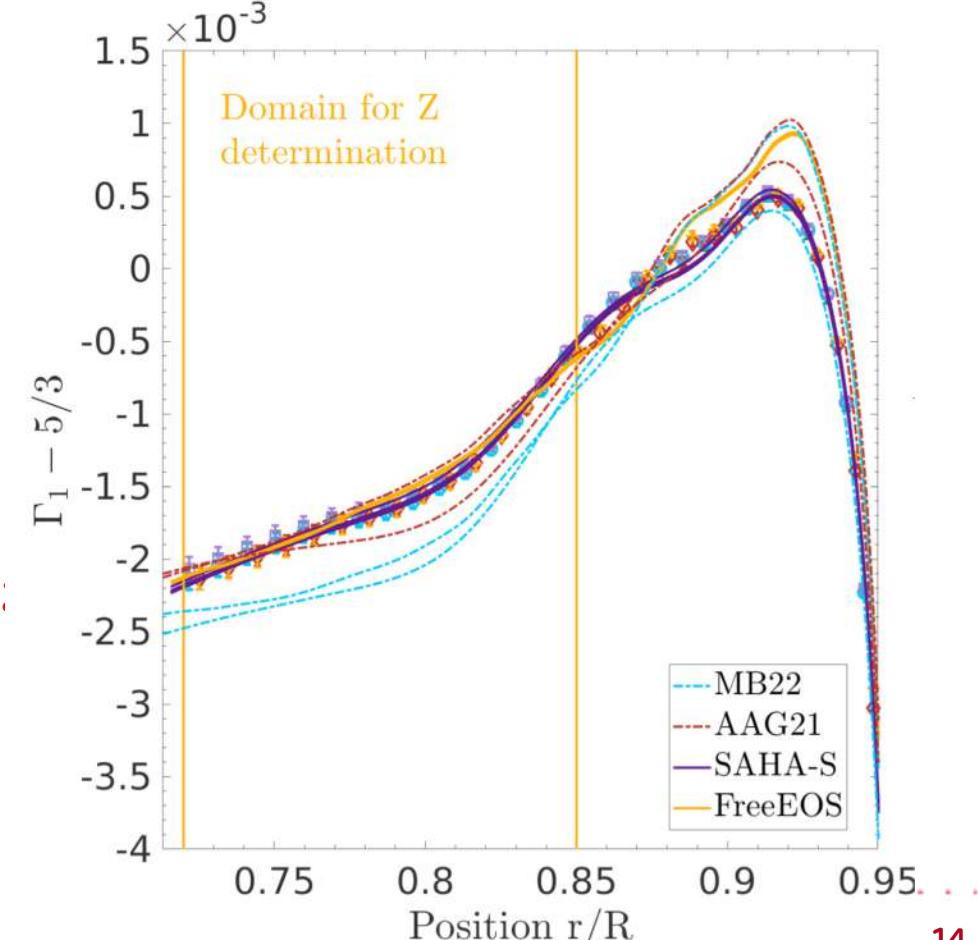
However,

Propagation of sound waves carry information about composition through adiabatic index:

$$\Gamma_1 = \left( \frac{\partial \ln P}{\partial \ln \rho} \right)_{\text{ad}} = 5/3 \text{ (for fully ionized gas)} < 5/3 \text{ in partial ionization regions}$$

- It can be determined through inversion of solar oscillations and compared to solar models.
- Only sensitive to total Z (not individual elements)
- Results are degenerate with equation of state

**Results indicate agreement with AAG21 (low-Z) rather than MB22**



# Consensus (as of end of 2022) H-burning x-sections

INT WORKSHOP INT-22-82W

## Solar Fusion Cross Sections III

July 26, 2022 - July 29, 2022



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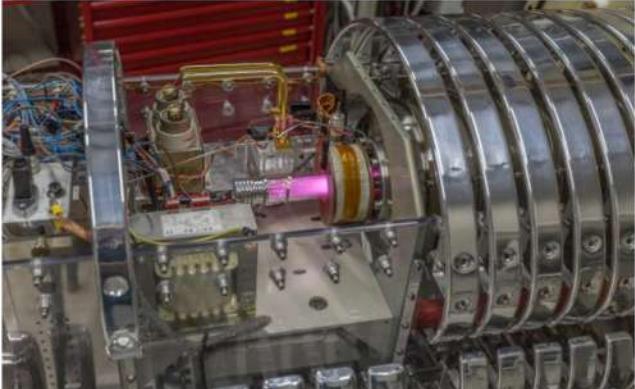
### ORGANIZERS

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Note to applicants: This workshop will be held in the David Brower Center near the UC Berkeley campus in Berkeley, CA.

WORKING GROUP AND PRESENTATIONS WEBSITE



Draft, February 26, 2024

## Solar fusion III: New data and theory for hydrogen-burning stars.

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**arXiv: 2405.06470**  
**RMP coming (soon)**

# Consensus (as of end of 2022) H-burning x-sections

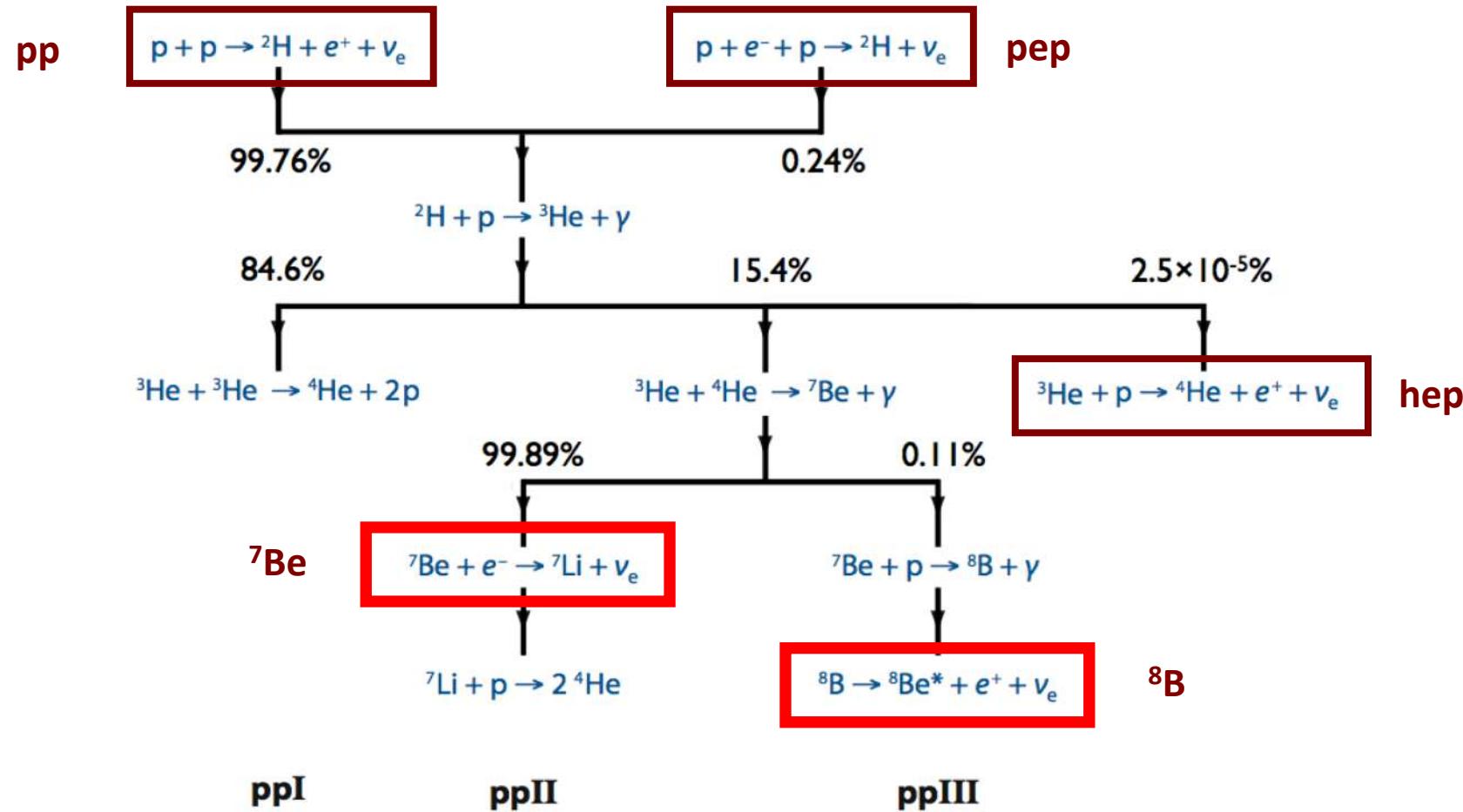


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Reaction	1 $\sigma$ error (SFIII)	1 $\sigma$ error (SFII, 2010)	
p+p	1.5%	0.9%	GA/GV
$^3\text{He} + ^3\text{He}$	6.6%	5.2%	uncert. in protons spectral shape
$^3\text{He} + ^4\text{He}$	5.1%	5%	
$^7\text{Be} + \text{p}$	3.6%	7.5%	halo EFT & Fitting
$^{14}\text{N} + \text{p}$	8.4%	7.2%	R-matrix and data tension for ground state transition

# What solar neutrinos from pp-chain tell us



# What solar neutrinos from pp-chain tell us



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$\text{cm}^{-2}\text{s}^{-1}$	AAG21	MB22	Sun
pp ( $10^{10}$ )	6.00 (0.6%)	5.95 (0.6%)	5.94 (0.4%)
pep ( $10^8$ )	1.45 (1.1%)	1.42 (1.1%)	1.42 (1.6%)
hep ( $10^3$ )	8.16 (30%)	7.92 (30%)	30 (33%)
${}^7\text{Be}$ ( $10^9$ )	4.52 (7.4%)	4.90 (7.6%)	4.93 (2%)
${}^8\text{B}$ ( $10^6$ )	4.31 (12.6%)	5.13 (13.1%)	5.20 (1.9%)



"Sun": experimental results from Gonzalez-García et al. 2024

**Model uncertainties >> experimental ones**

**x-sections ( $S_{17}, S_{34}, S_{11}$ )**

**6% for  ${}^7\text{Be}$**

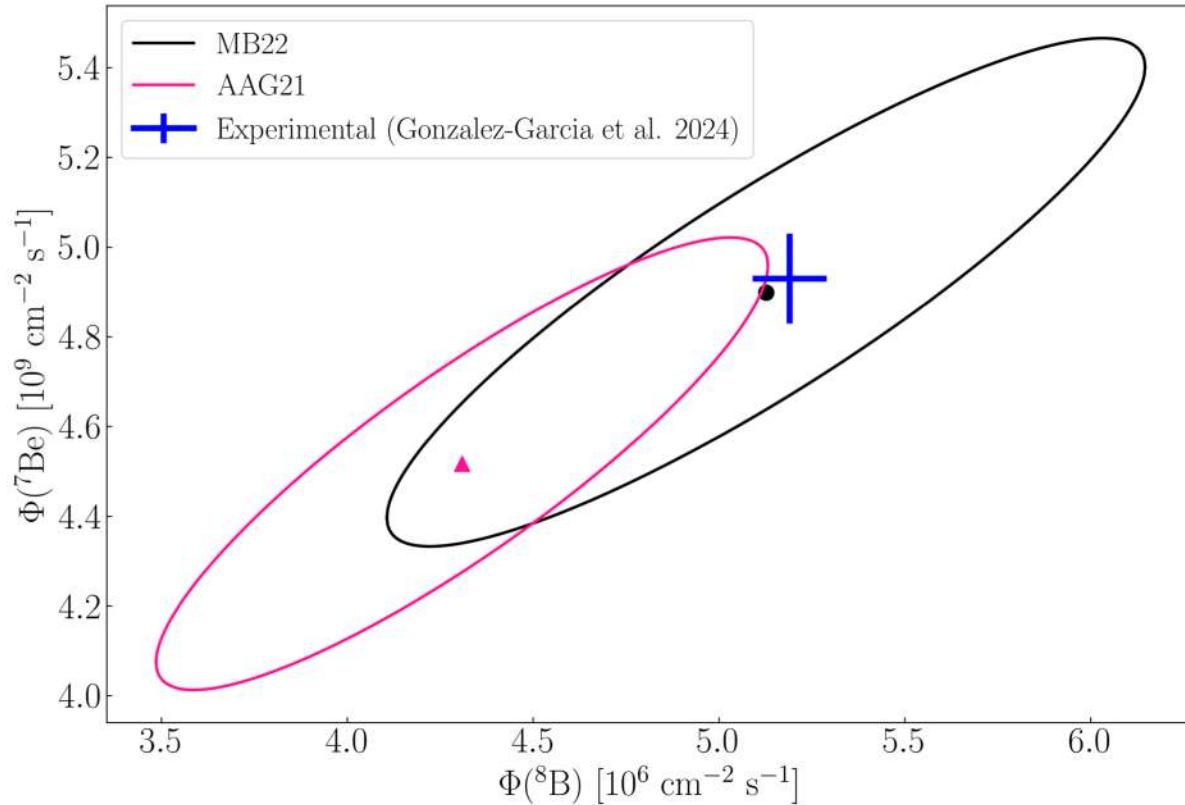
**8% for  ${}^8\text{B}$**

**radiative opacity**



# What solar neutrinos from pp-chain tell us

Mostly a temperature sequence with slope determined by nuclear reaction rates



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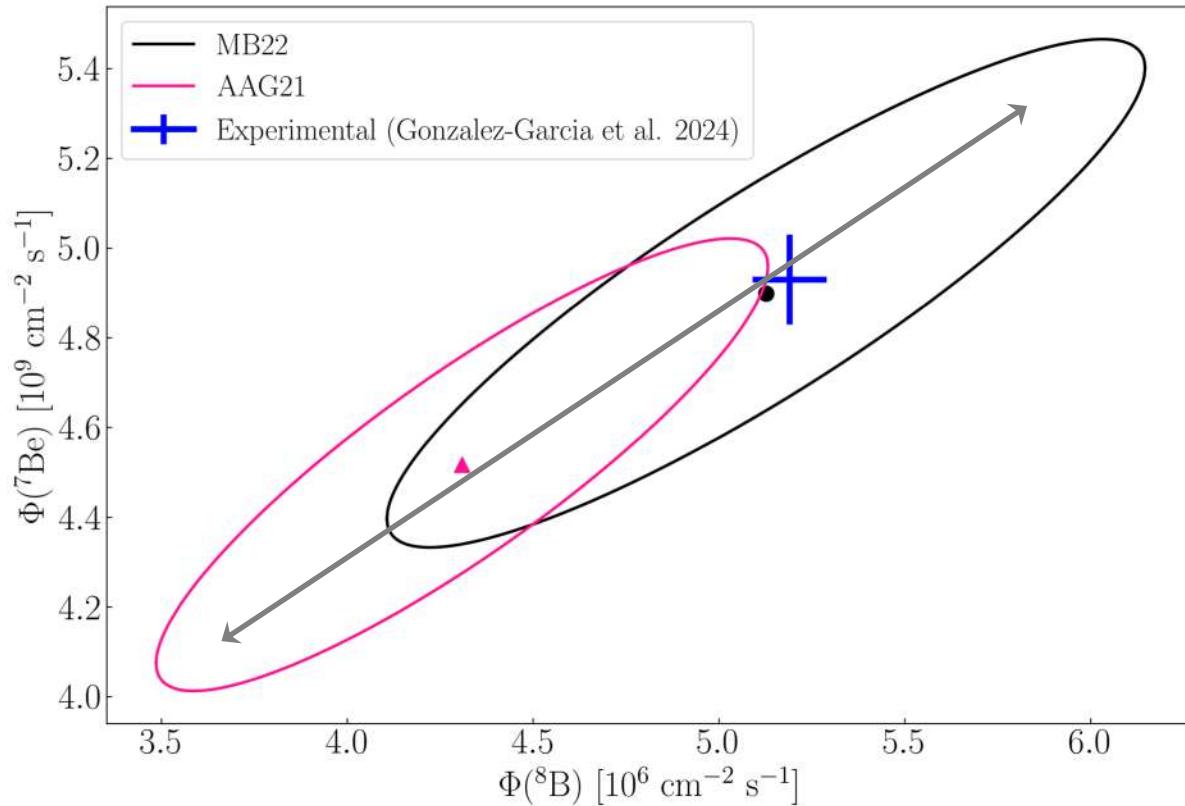
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Mostly a temperature sequence with slope determined by nuclear reaction rates



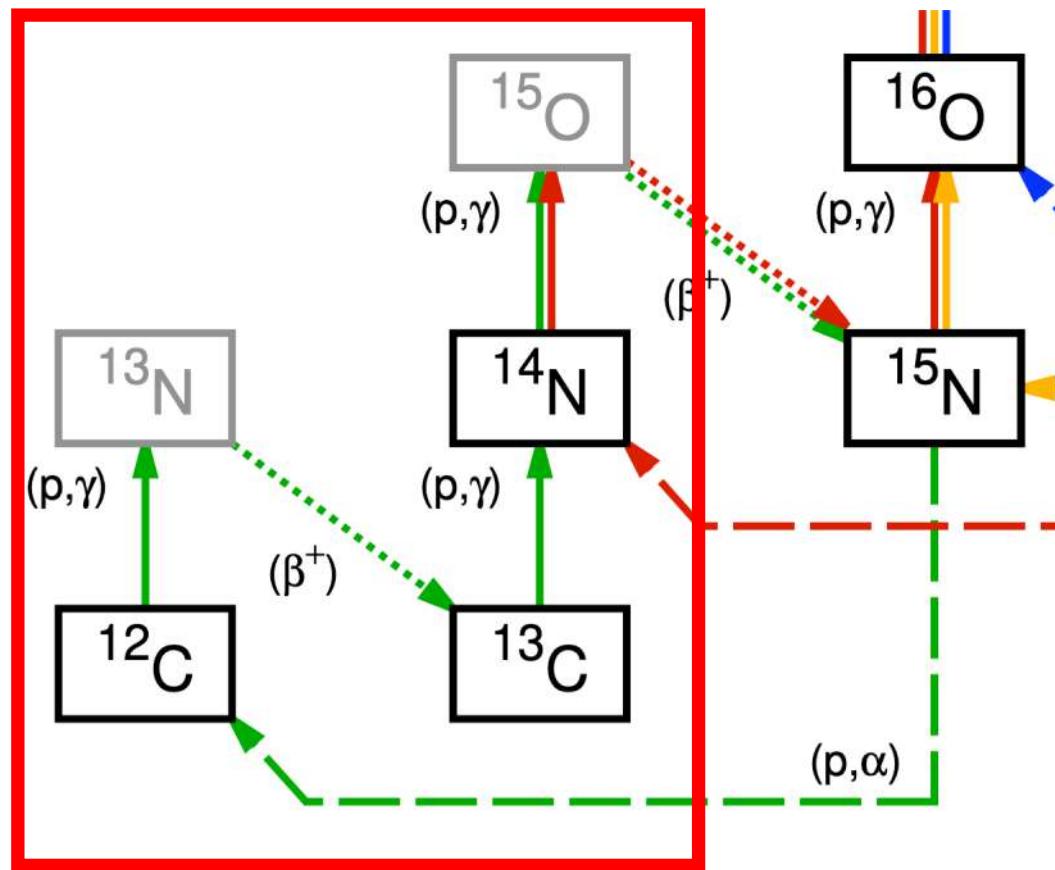
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"Sun": experimental results from Gonzalez-Garcia et al. 2024

Combination of **composition + radiative opacities** →  
→ core temperature consistent with higher opacity (Z?) models

# CN-cycle is a trace contribution to solar structure

CN operates against a “fixed” structure determined by pp-chains

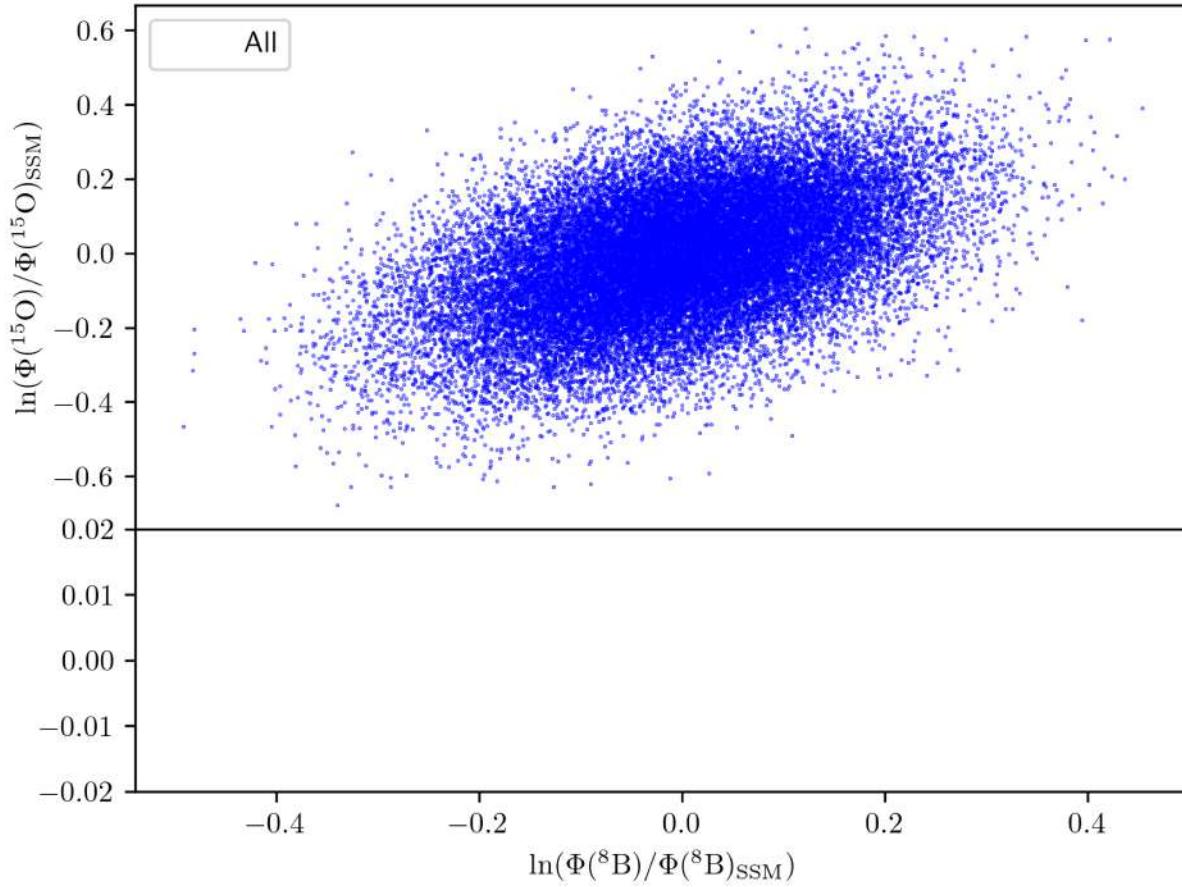


Changes in physics affecting CN do not change structure, **i.e. core temperature**,  
→ retain explicit dependences:

- e.g. linear response to bottleneck nuclear reaction  $^{14}\text{N}(\text{p},\gamma)^{15}\text{O}$
- **linear dependence on abundance of catalysts in solar core: C+N**
- **one-to-one relation between neutrino fluxes and CN abundance**

# What CN solar neutrinos tell us

$^8\text{B}$  as a thermometer

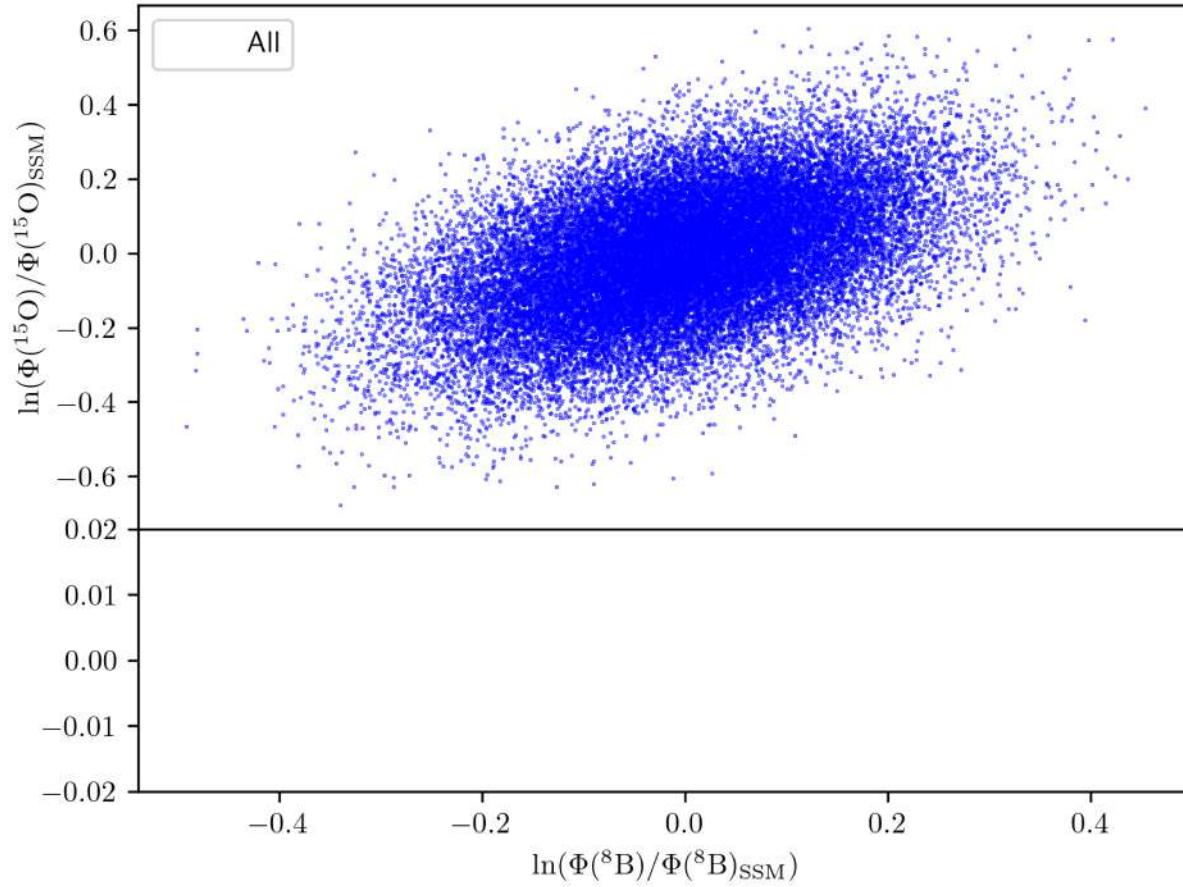


Neutrino fluxes depend on:

- **solar core temperature – environmental quantities**
  - opacity
  - heavy elements (Si, Mg, Fe)
  - luminosity, age
- uncertainties in these quantities affect n-fluxes in a fully correlated way
- **nuclear reaction rates**
  - specific dependence for specific fluxes
  - (e.g.  $^{14}\text{N}(\text{p},\text{g})^{15}\text{O}$  does not affect pp-chain)
- **catalyzing effect of abundances**
  - C & N abundance in the solar core → CN-cycle

# What CN solar neutrinos tell us

$^8\text{B}$  as a thermometer



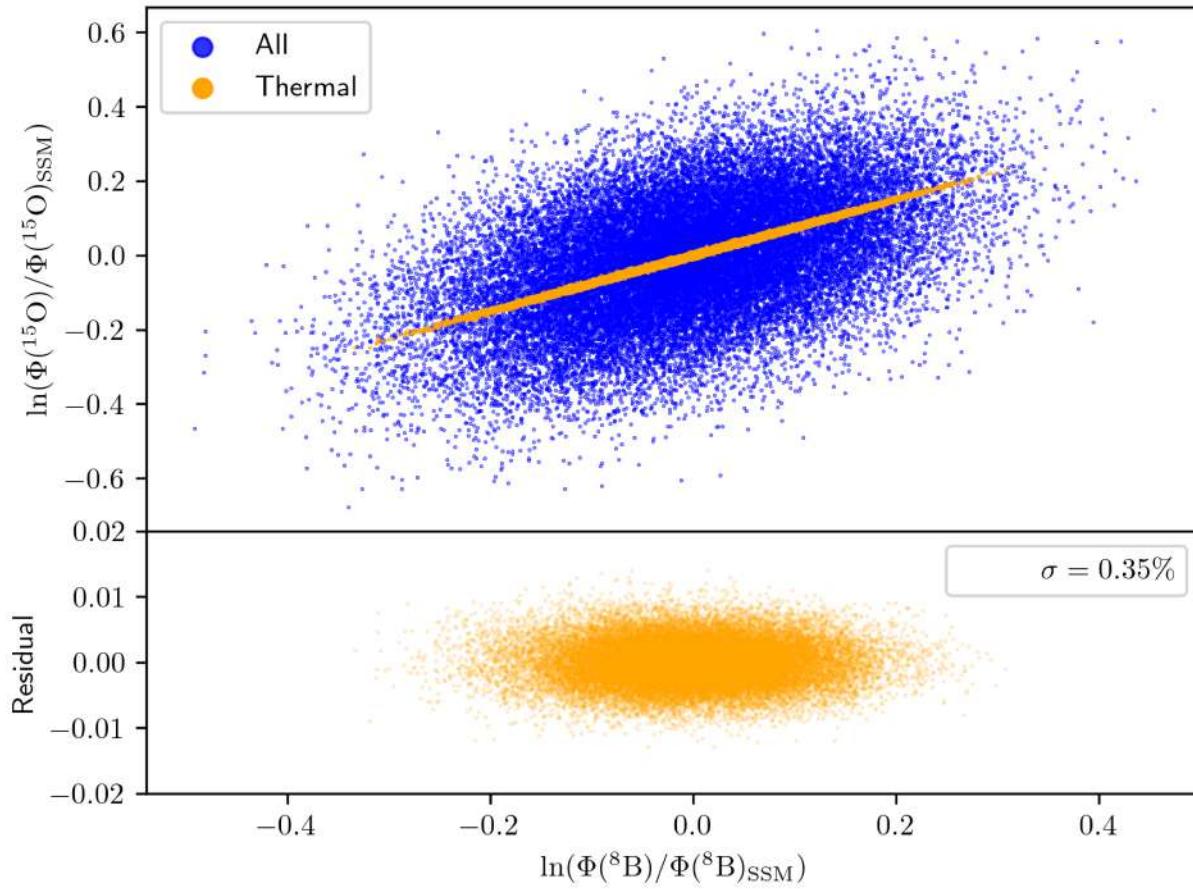
Neutrino fluxes as power-laws:

$$\frac{\phi(^{15}\text{O})}{\phi(^{15}\text{O})^{\text{SSM}}} = [L_{\odot}^{5.942} O^{2.034} A^{1.364} D^{0.382}] \\ \times [S_{11}^{-2.912} S_{33}^{0.024} S_{34}^{-0.052} S_{17}^{0.0} S_{e7}^{0.0} S_{114}^{1.00}] \\ \times [x_C^{0.815} x_N^{0.217} x_O^{0.112} x_{\text{Ne}}^{0.081} x_{\text{Mg}}^{0.069} x_{\text{Si}}^{0.150} x_S^{0.109} x_{\text{Ar}}^{0.028} x_{\text{Fe}}^{0.397}]$$

$$\frac{\phi(^8\text{B})}{\phi(^8\text{B})^{\text{SSM}}} = [L_{\odot}^{6.966} O^{2.734} A^{1.319} D^{0.278}] \\ \times [S_{11}^{-2.665} S_{33}^{-0.419} S_{34}^{0.831} S_{17}^{1.028} S_{e7}^{-1} S_{114}^{0.00}] \\ \times [x_C^{0.022} x_N^{0.007} x_O^{0.128} x_{\text{Ne}}^{0.102} x_{\text{Mg}}^{0.092} x_{\text{Si}}^{0.198} x_S^{0.138} x_{\text{Ar}}^{0.034} x_{\text{Fe}}^{0.498}]$$

# What CN solar neutrinos tell us

${}^8\text{B}$  as a thermometer

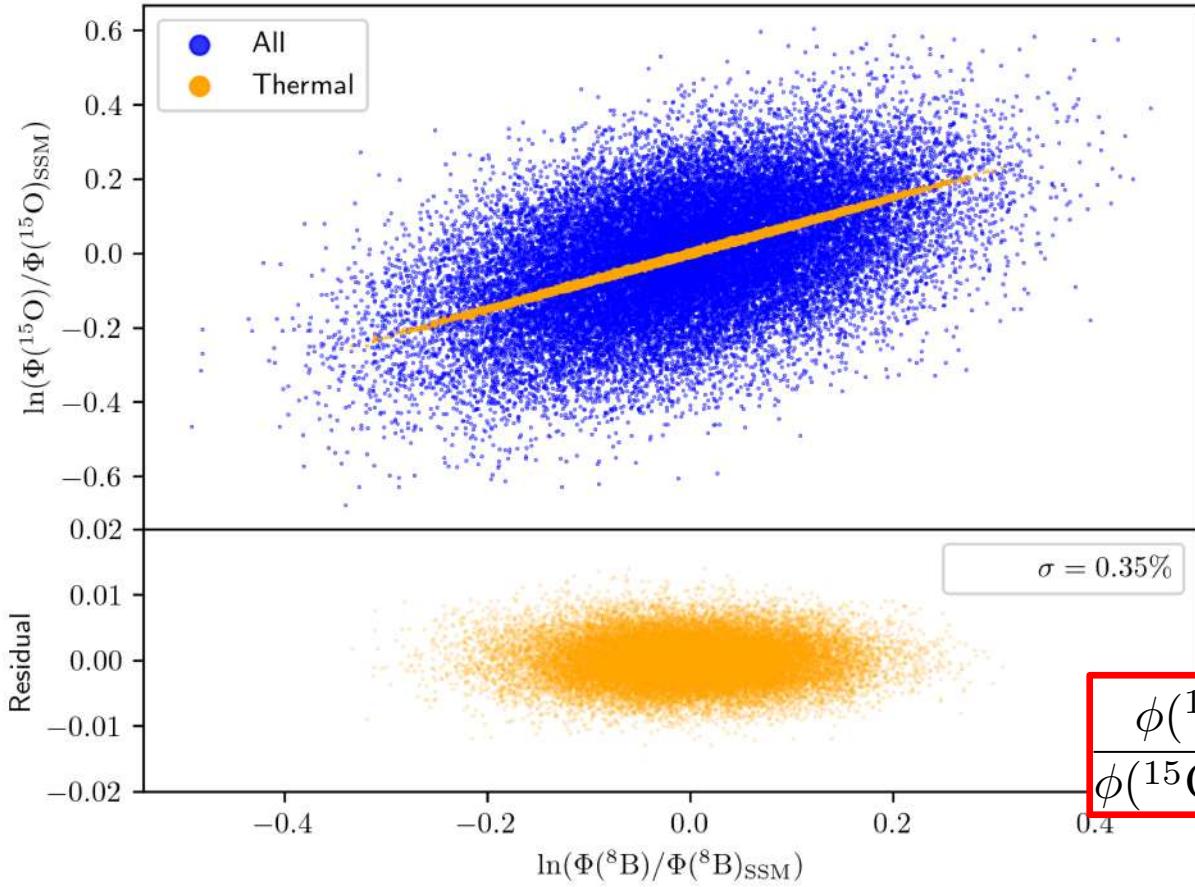


**Thermal uncertainties are cancelled out**, absorbed by a  ${}^8\text{B}$  experimental measurement, down to 0.3%

$$\frac{\phi({}^{15}\text{O})}{\phi({}^{15}\text{O})_{\text{SSM}}} \Big/ \left[ \frac{\phi({}^8\text{B})}{\phi_{\text{SSM}}({}^8\text{B})} \right]^{0.785} = x_C^{0.794} x_N^{0.212} D^{0.172} \\ \times [L_{\odot}^{0.515} O^{-0.016} A^{0.308}] \\ \times [S_{11}^{-0.831} S_{33}^{0.342} S_{34}^{-0.685} S_{17}^{-0.785} S_{e7}^{0.785} S_{114}^{0.995}] \\ \times [x_O^{0.003} x_{\text{Ne}}^{-0.005} x_{\text{Mg}}^{-0.003} x_{\text{Si}}^{-0.001} x_S^{-0.001} x_{\text{Ar}}^{0.001} x_{\text{Fe}}^{0.003}]$$

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$^8\text{B}$  as a thermometer



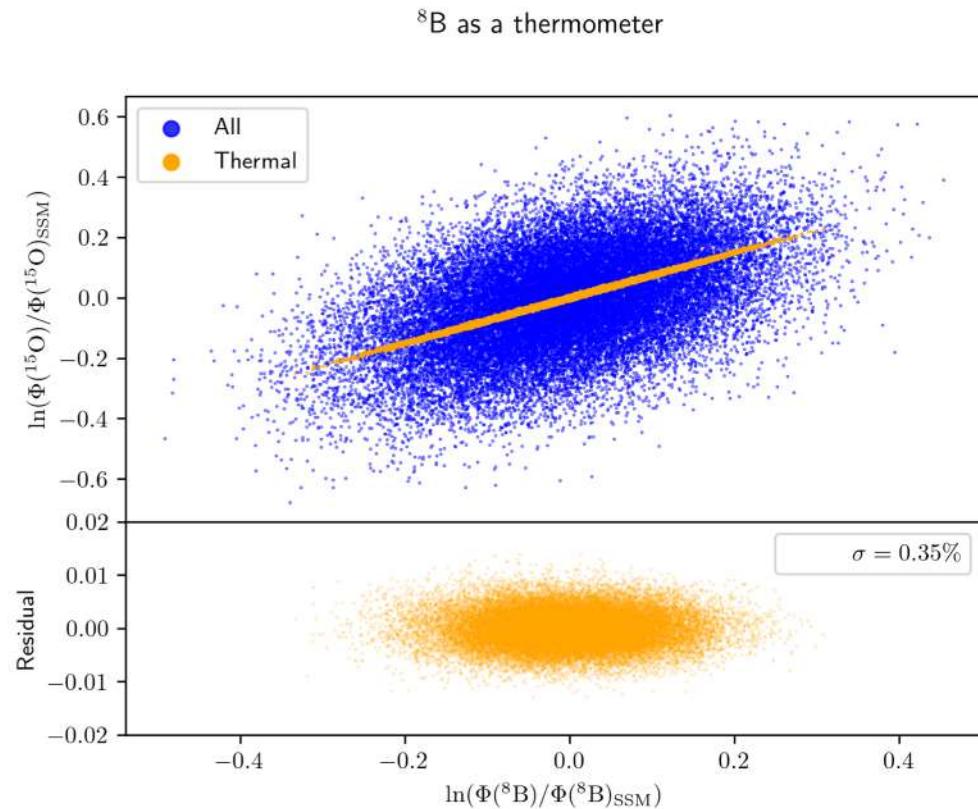
$$\frac{\phi(^{15}\text{O})}{\phi(^{15}\text{O})_{\text{SSM}}} \approx \left( \frac{\phi(^8\text{B})}{\phi(^8\text{B})_{\text{SSM}}} \right)^{0.785} (x_{\text{C+N}})[1 \pm 0.10(\text{nuc}) + 0.03(\text{D})]$$

Linear dependence  
on C+N

$$\frac{\phi(^{15}\text{O})}{\phi(^{15}\text{O})_{\text{SSM}}} / \left[ \frac{\phi(^8\text{B})}{\phi_{\text{SSM}}(^8\text{B})} \right]^{0.785} = x_C^{0.794} x_N^{0.212} D^{0.172} \\ \times [L_\odot^{0.515} O^{-0.016} A^{0.308}] \\ \times [S_{11}^{-0.831} S_{33}^{0.342} S_{34}^{-0.685} S_{17}^{-0.785} S_{e7}^{0.785} S_{114}^{0.995}] \\ \times [x_O^{0.003} x_{\text{Ne}}^{-0.1} x_{\text{Fe}}^{0.003}]$$

Nuclear reaction rates

# What CN solar neutrinos tell us



SuperKamiokande, SNO, others → 2% measurement of <sup>8</sup>B flux

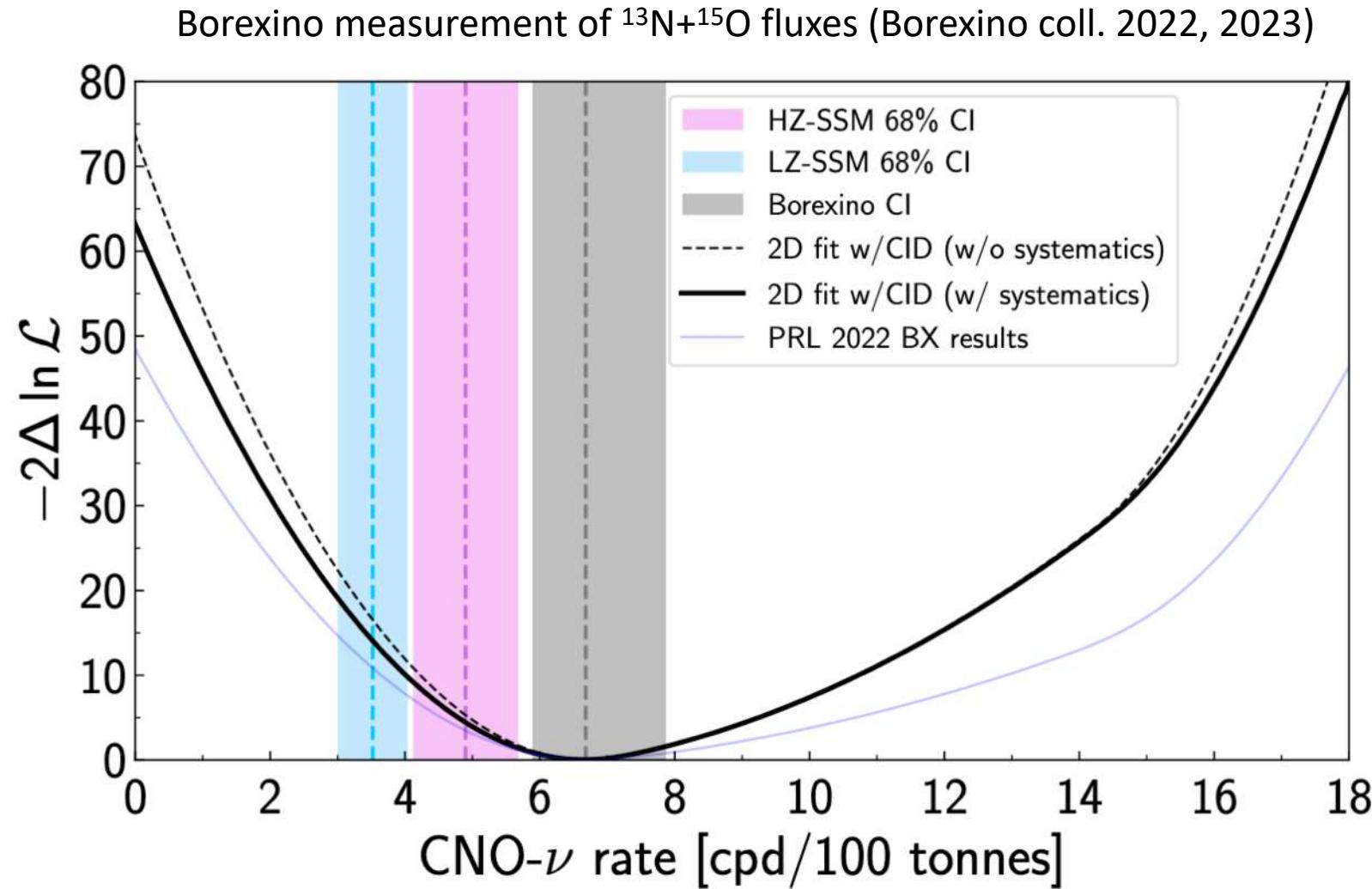
$$\frac{\phi(^{15}\text{O})}{\phi(^{15}\text{O})^{\text{SSM}}} \simeq \left( \frac{\phi(^8\text{B})}{\phi(^8\text{B})^{\text{SSM}}} \right)^{0.785} (x_{\text{C+N}})[1 \pm 0.10(\text{nuc}) + 0.03(\text{D})]$$

C+N abundance if <sup>15</sup>O flux is measured  
(or any combination of <sup>13</sup>N & <sup>15</sup>O fluxes)

**10% uncertainty (nuclear rates)**  
**experimental uncertainty in CN-ν fluxes dominated by**

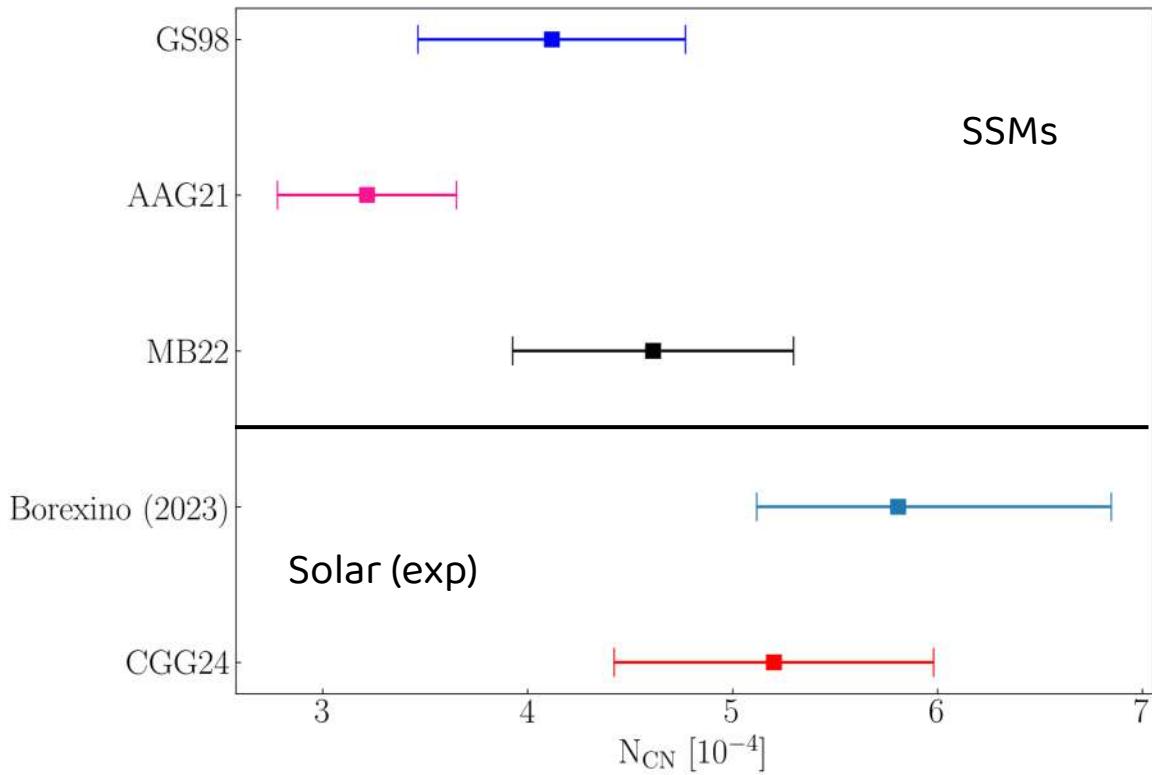
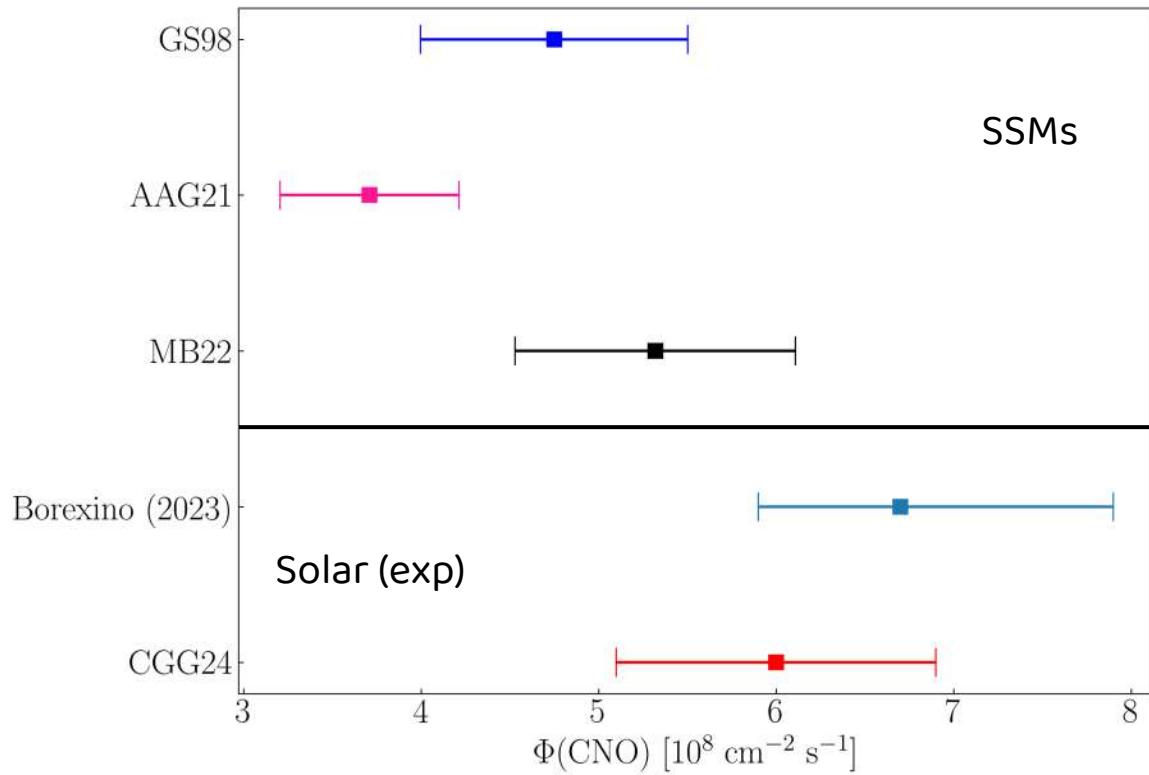
- <sup>14</sup>N(p,γ)<sup>15</sup>O (8.5%)
- <sup>3</sup>He(<sup>4</sup>He,γ)<sup>7</sup>Be (5%)

# CN measurement by Borexino



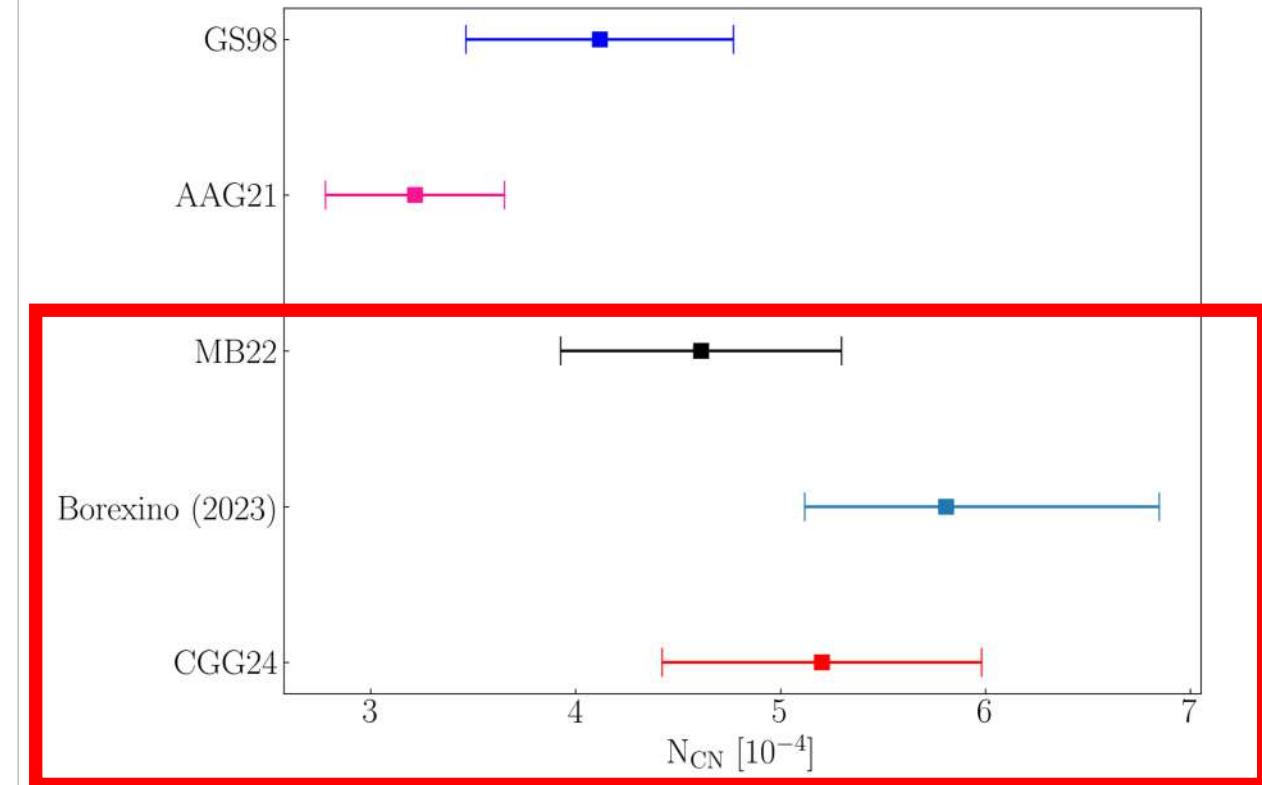
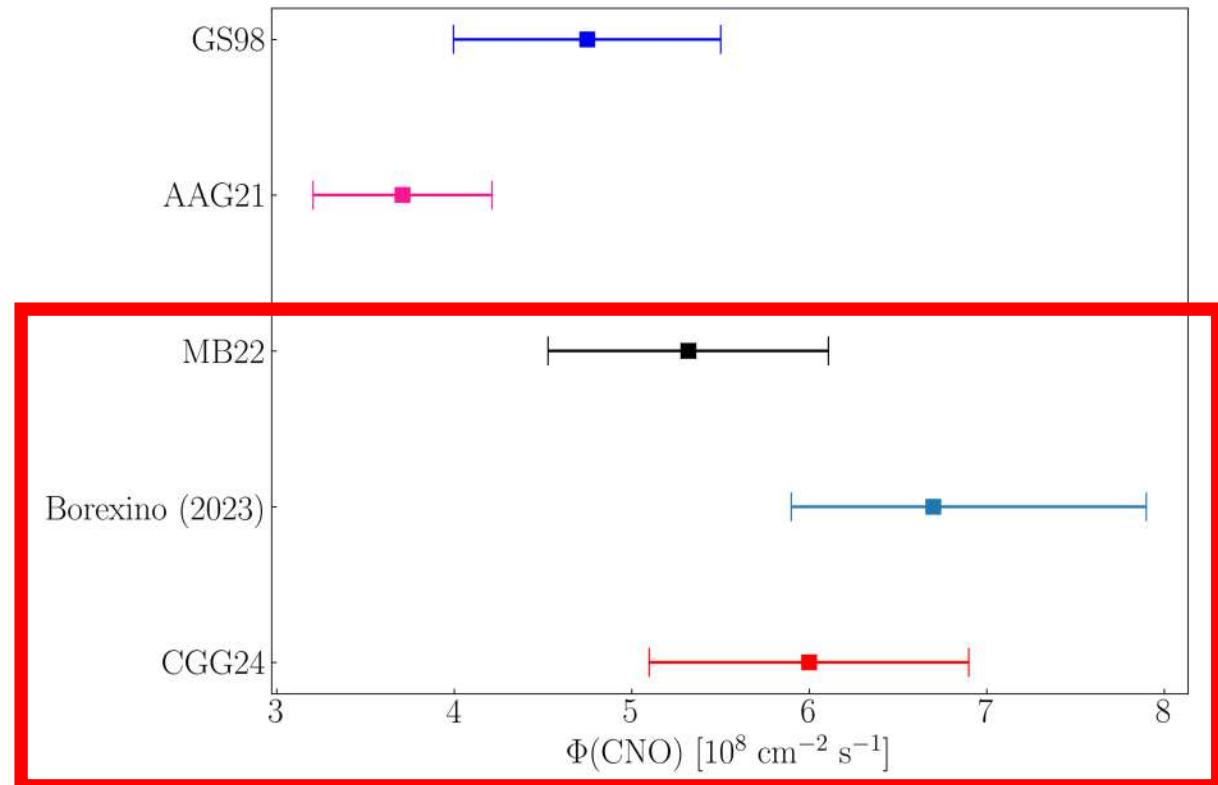
# What CN solar neutrinos tell us

$$\frac{\phi(^{15}\text{O})}{\phi(^{15}\text{O})^{\text{SSM}}} \simeq \left( \frac{\phi(^8\text{B})}{\phi(^8\text{B})^{\text{SSM}}} \right)^{0.785} (x_{\text{C+N}})[1 \pm 0.10(\text{nuc}) + 0.03(\text{D})]$$



# What CN solar neutrinos tell us

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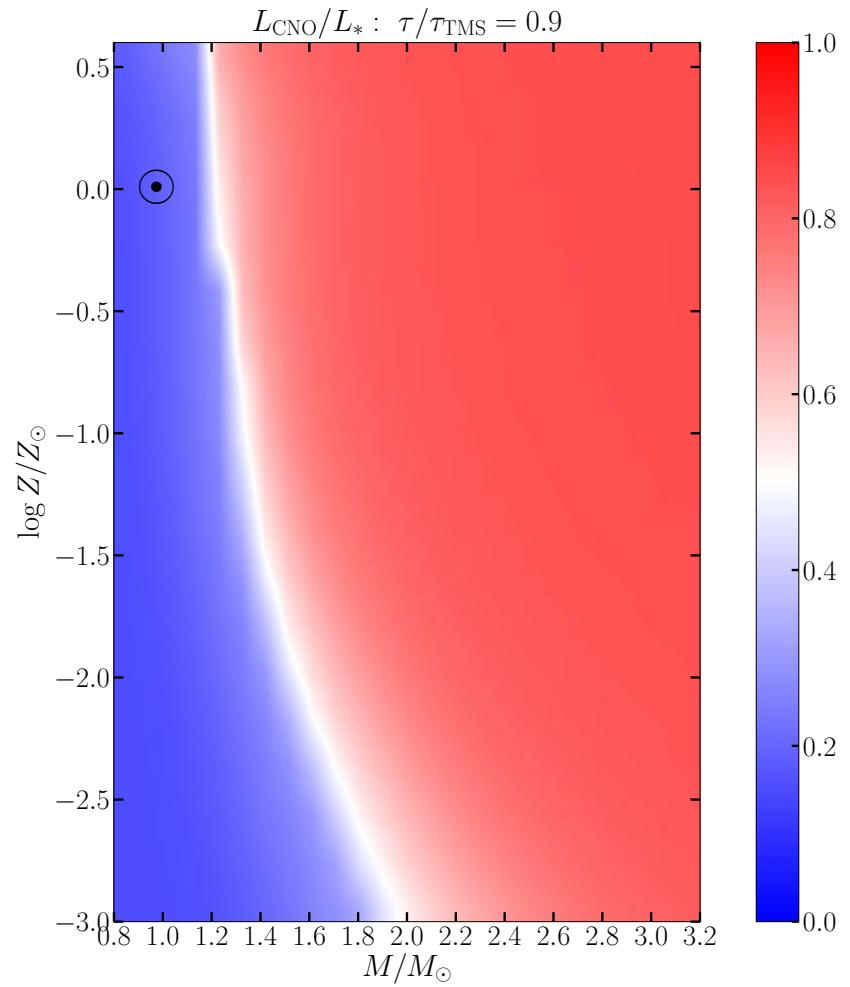


CN neutrinos break the degeneracy between composition and opacity  
Favor large CN abundance

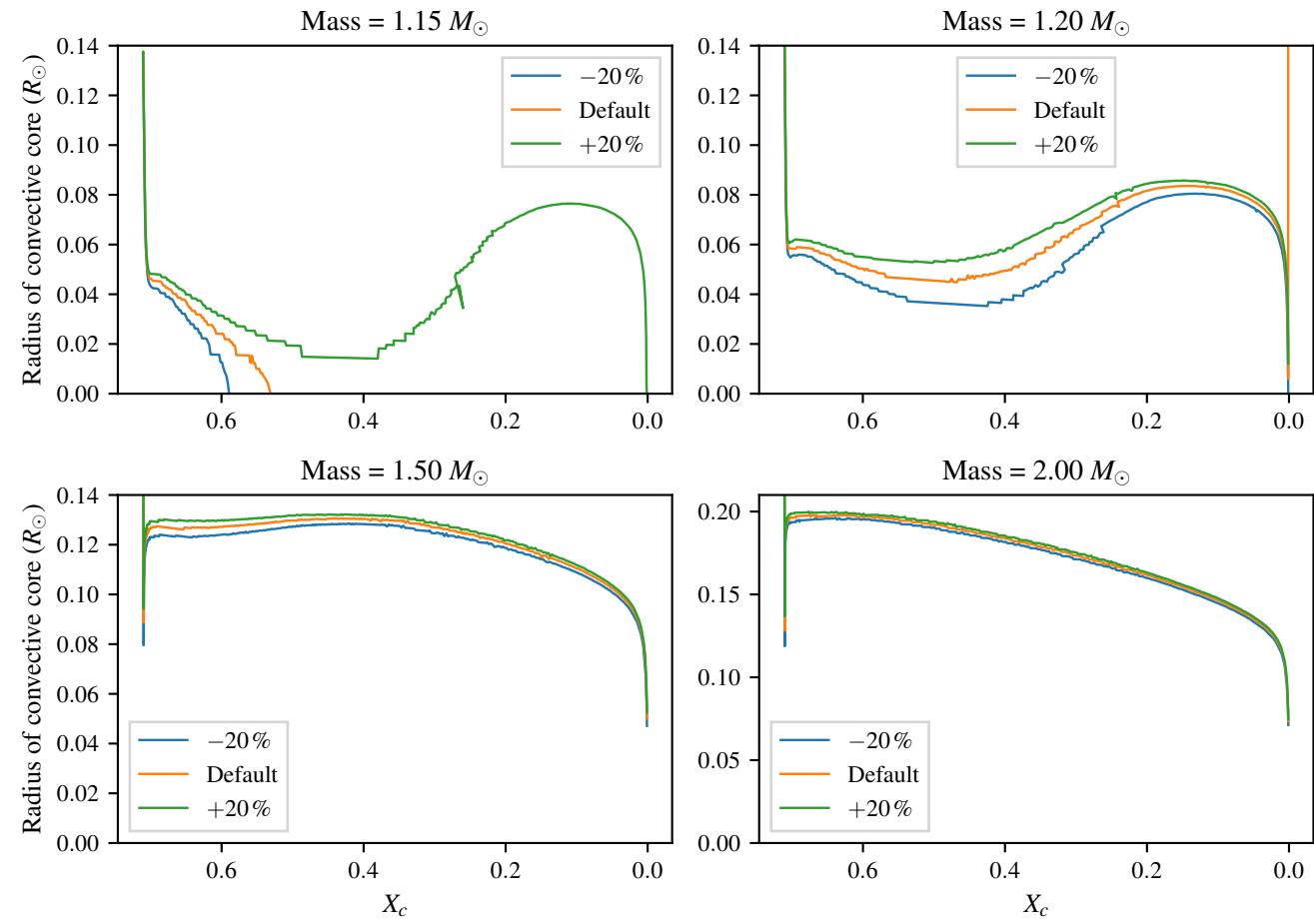
Nuclear rates largest source of uncertainty, but one we can control

# CNO vs mass and metallicity

## Stellar luminosity: CNO vs pp

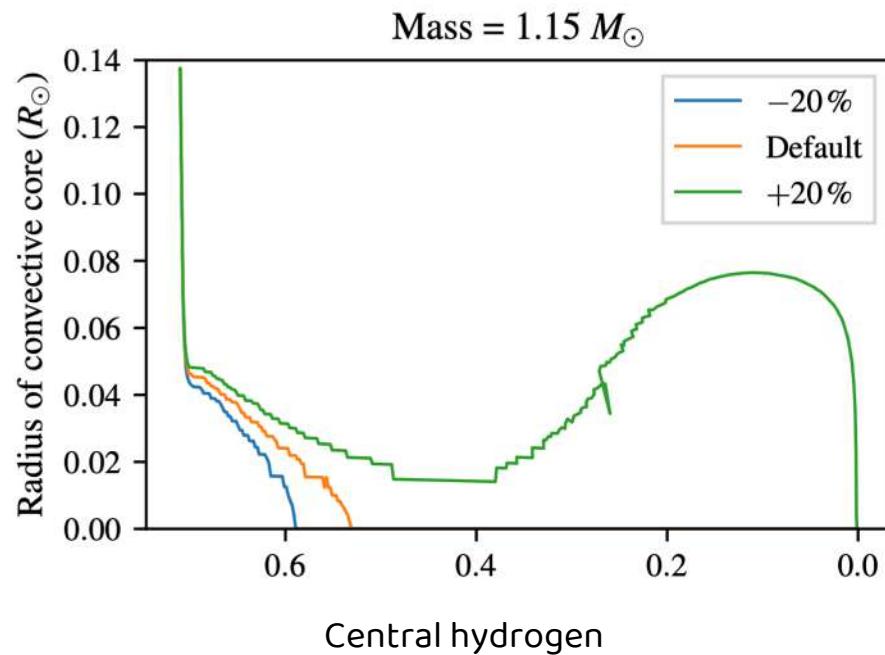


## Dependence on $^{14}\text{N} + \text{p}$ rate

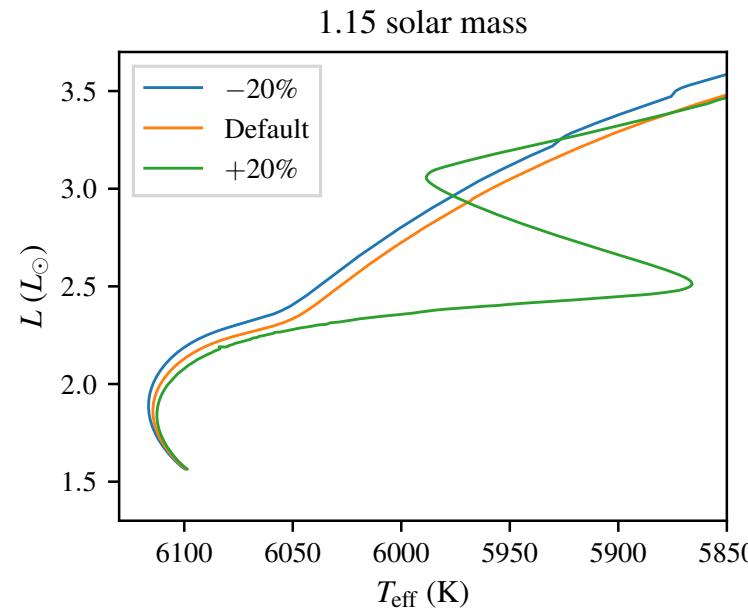


# Convective core: to be or not to be

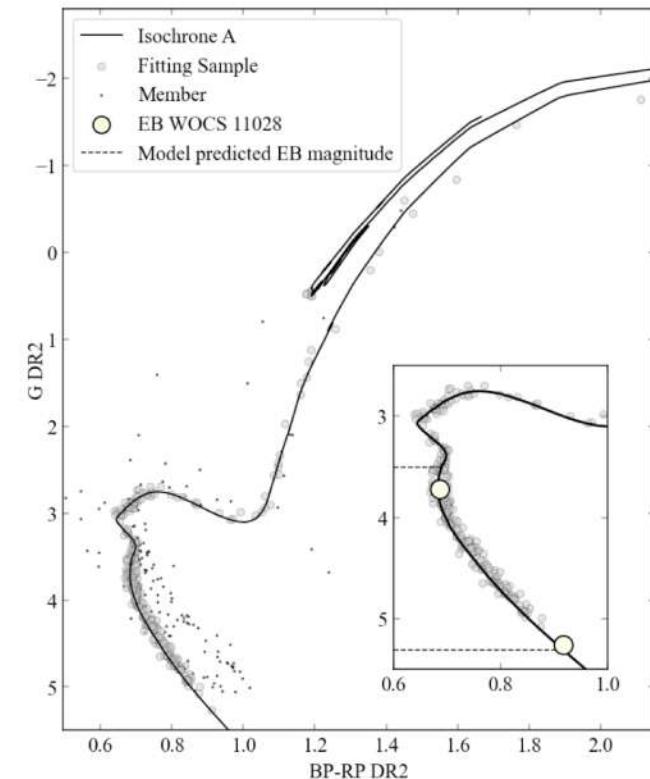
Size of convective core



Hertzsprung-Russell Diagram



M67 Color-magnitude Diagram



Stellar mass @ which convective core develops  
depends on N14+p reaction rate  
Qualitative change in HRD morphology

Reyes et al. 2024

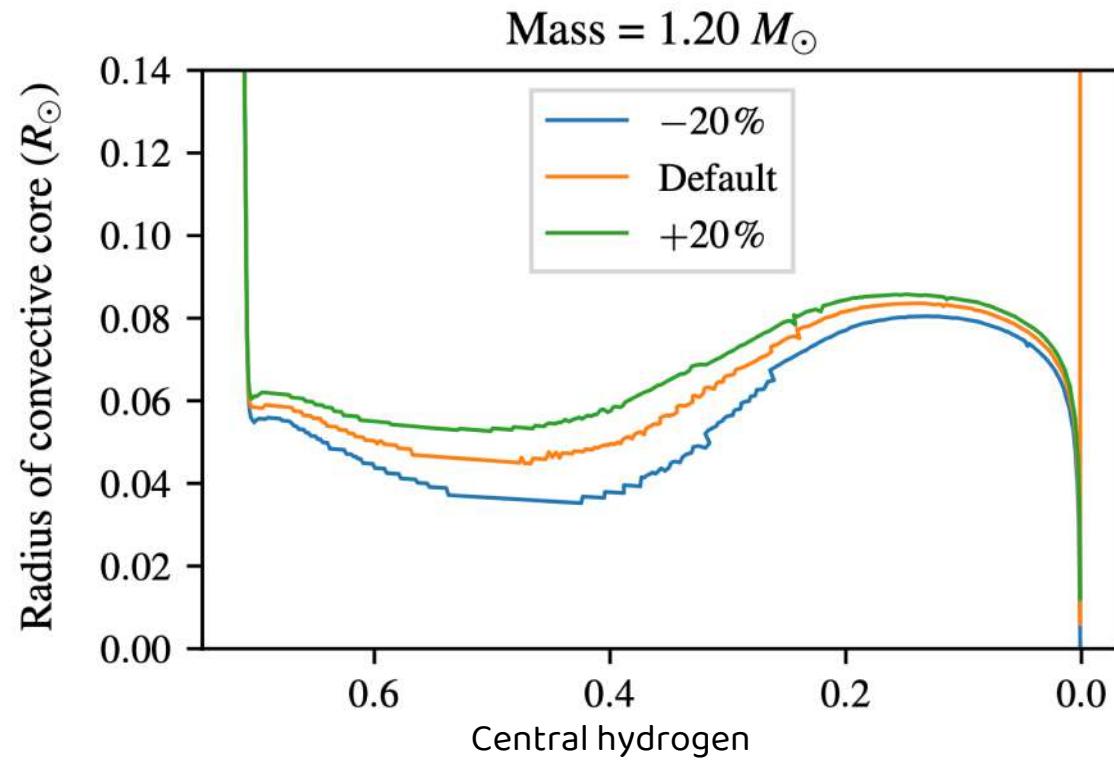
# Size of mixed core: convective boundary + overshooting

Size of chemically homogenous (mixed) core is formed by

+ truly convective core

+ **overshooting region** (parametrized,  $\alpha_{ov}$ , no 1st principles model)

It can be measured with astereoseismology



If  $^{14}\text{N}+\text{p}$  is not controlled, no way to separate the true CC boundary and the OV region

OV is the largest uncertain in stellar modeling

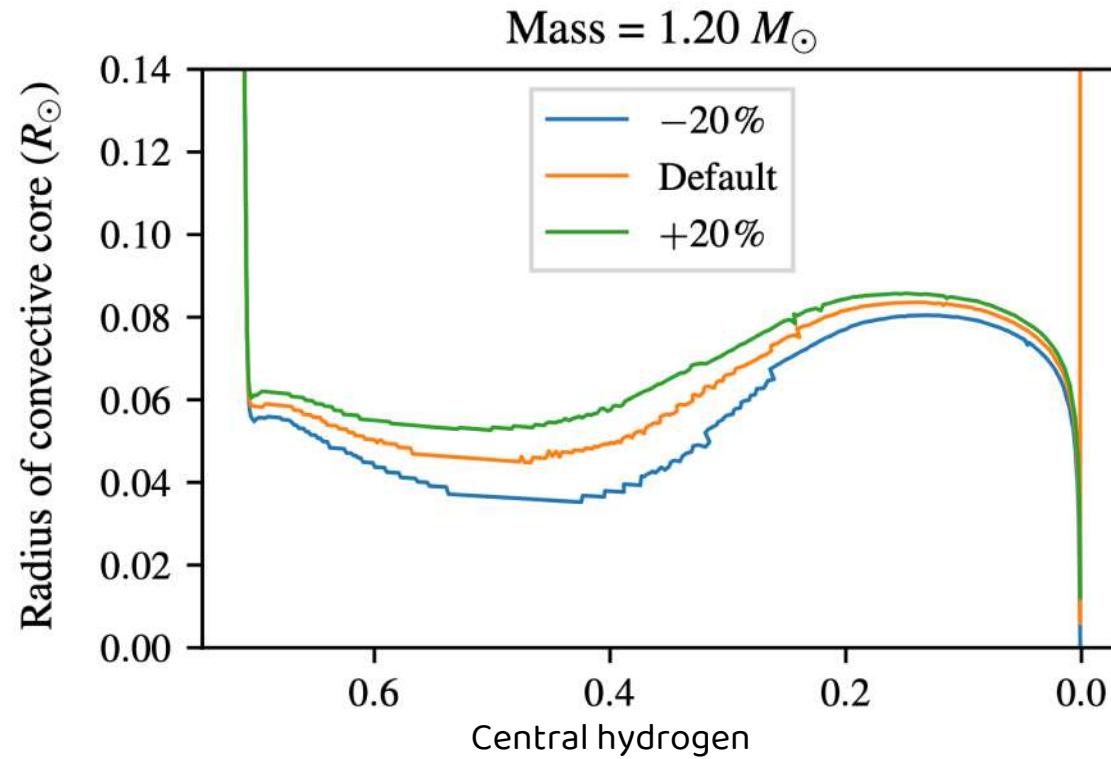
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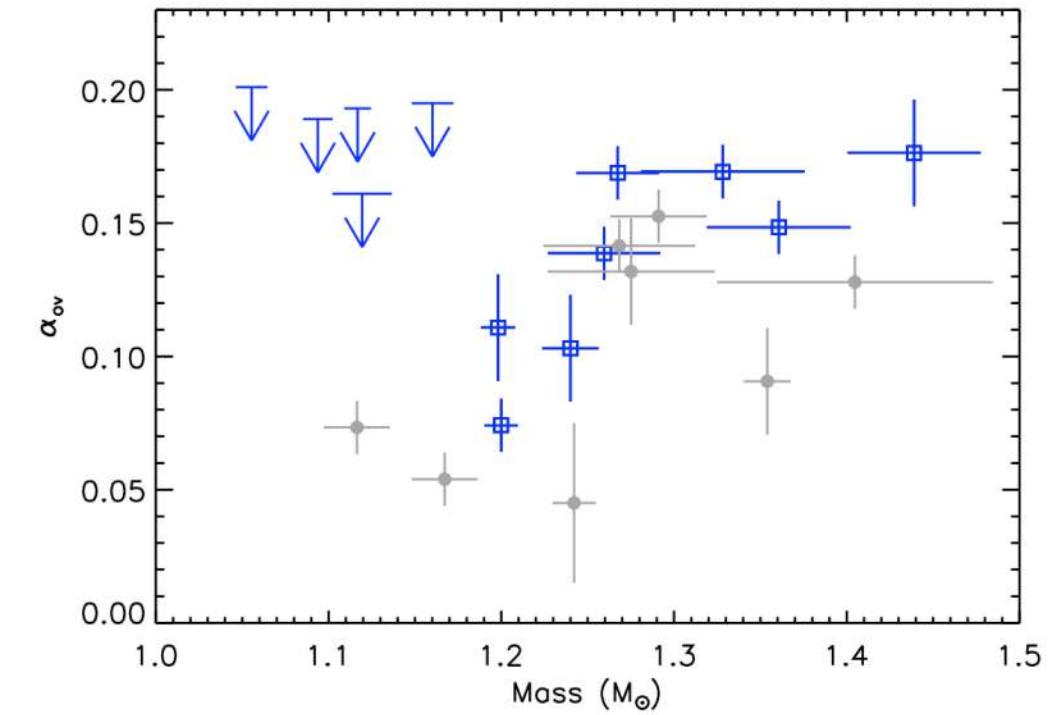
+ truly convective core

+ **overshooting region** (parametrized,  $\alpha_{ov}$ , no 1st principles model)

It can be measured with astereoseismology



Calibrating OV in stars  
Only if N14+p well known



Deheuvels et al. 2016

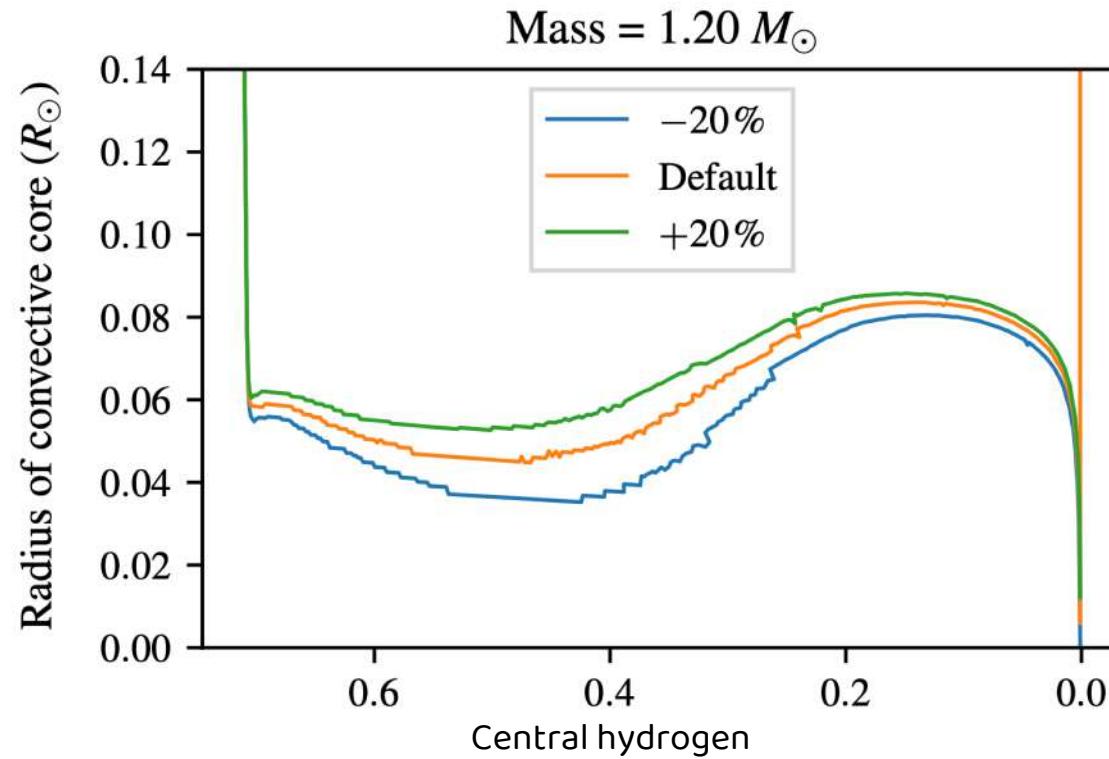
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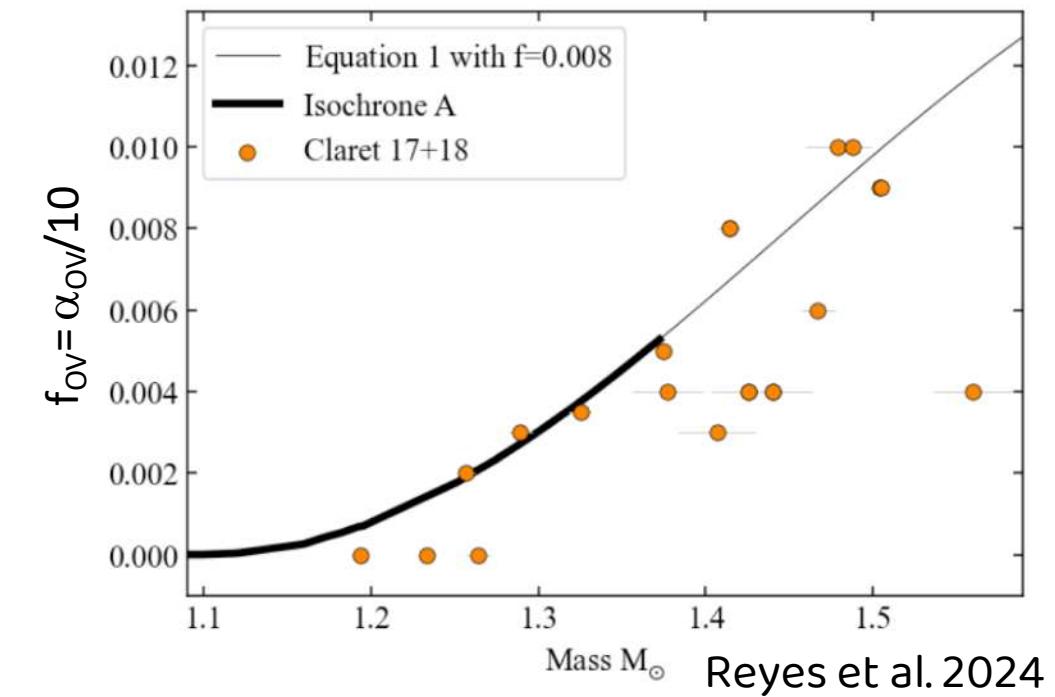
+ truly convective core

+ **overshooting region** (parametrized,  $\alpha_{ov}$ , no 1st principles model)

It can be measured with astereoseismology



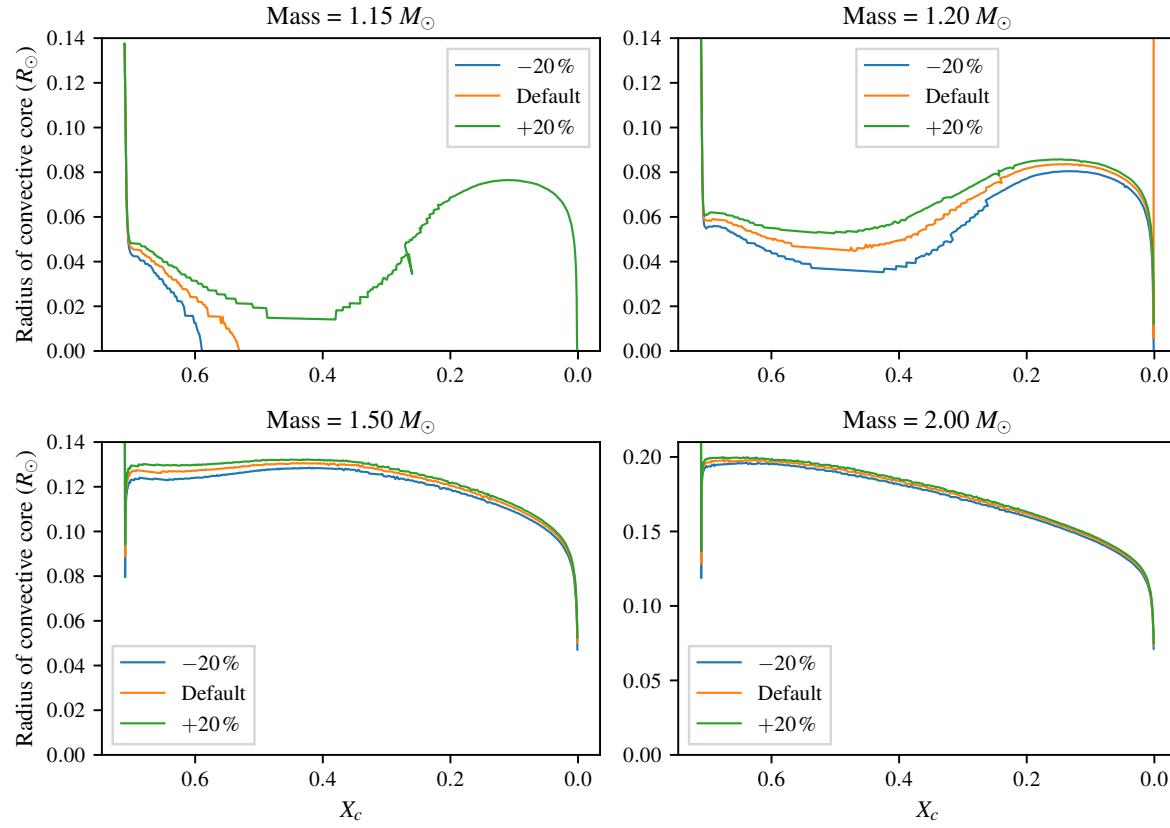
Calibrating OV in M67  
Only if N14+p well known



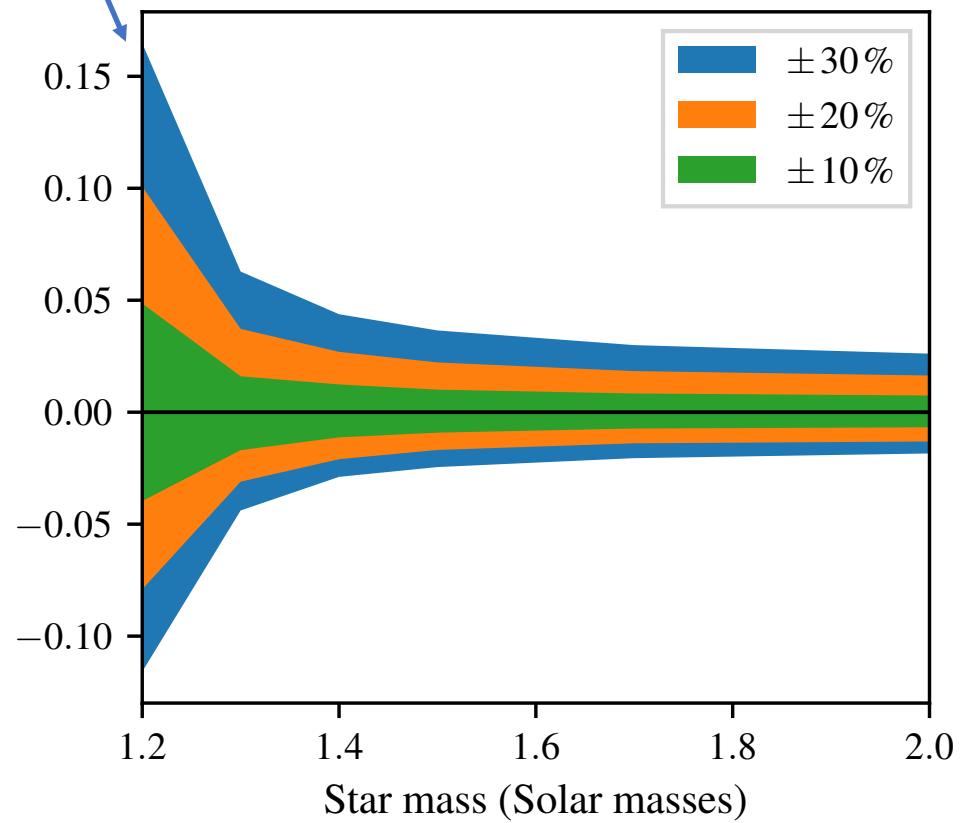
# Size of mixed core: convective boundary + overshooting

Parametrization of overshooting for stars in the  $1.1\text{-}1.4 M_{\odot}$  range requires improved  $^{14}\text{N}+\text{p}$

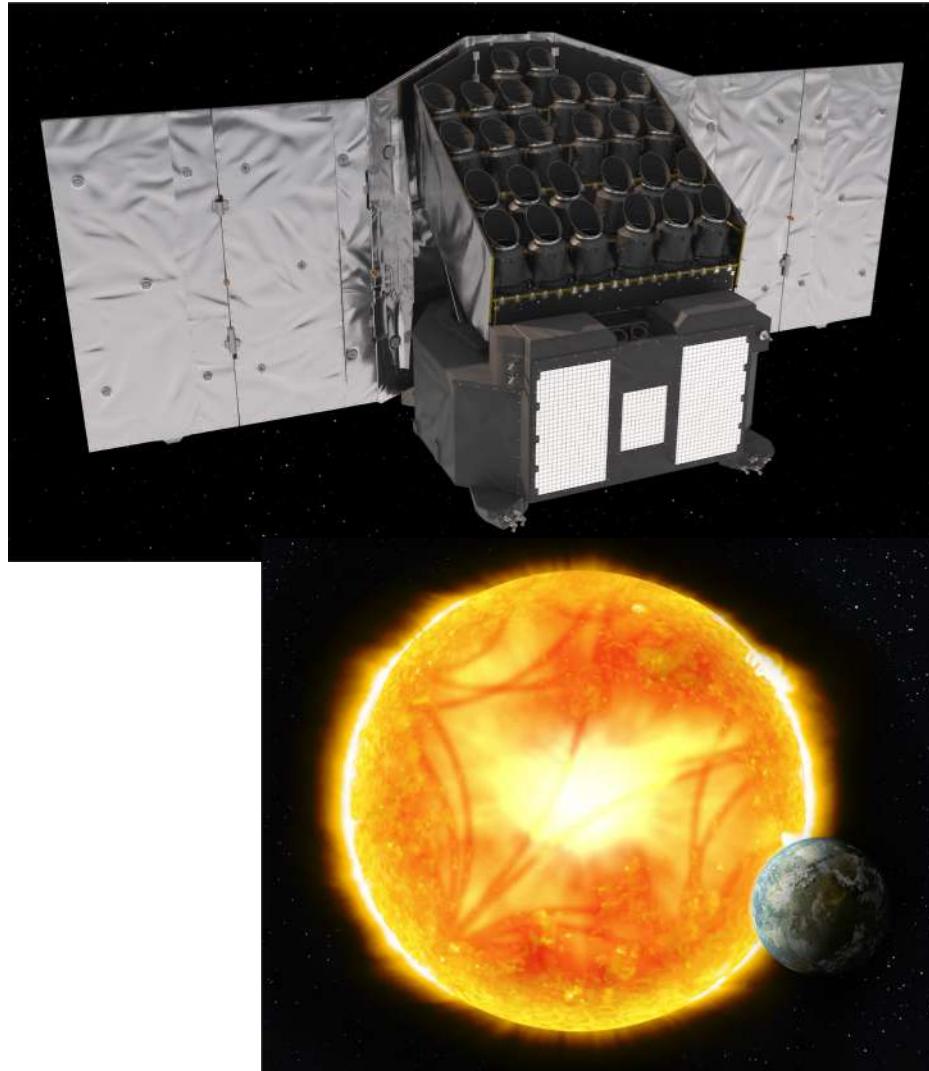
Typical OV values are 0.2-0.25 → at low masses, uncertainty is very large



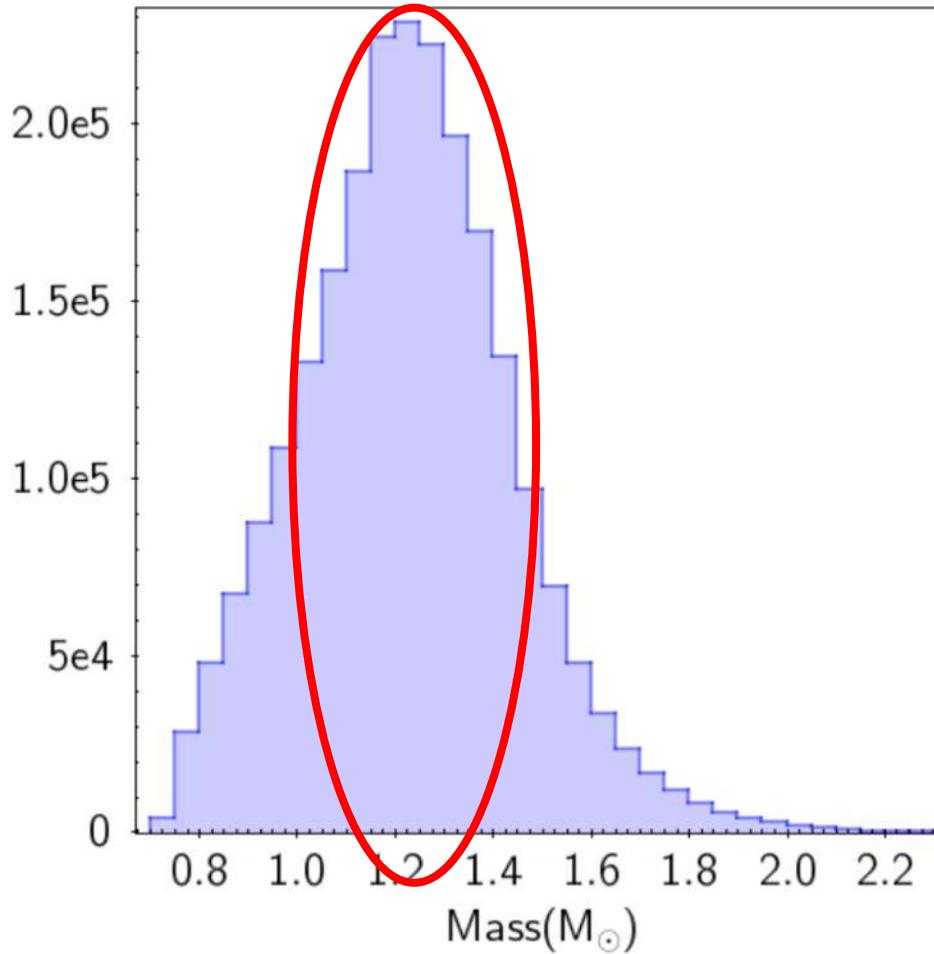
Systematic uncertainties on overshoot



# Critical for PLATO science, ESA's planet hunter & asteroseismology mission (launch December 2026)



Estimated mass distribution for FGK PLATO Sample



Constraints to compute solar models – composition is the big uncertainty:

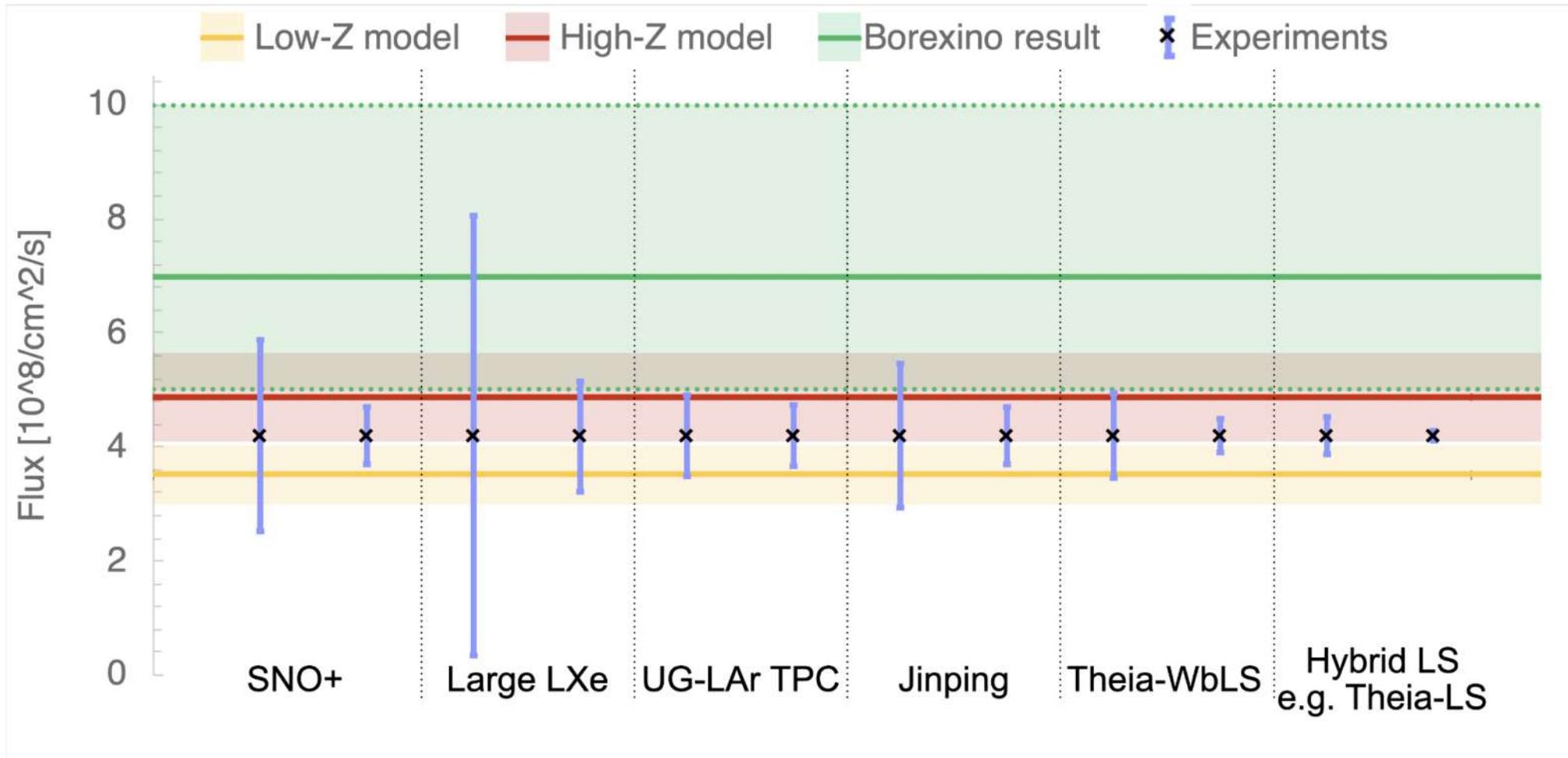
- High solar metallicity is favored by (degenerate with opacity):  
sound speed, surface helium, depth of convective envelope  
pp-chain solar neutrinos
- but lower solar metallicity is favored by (degenerate with equation of state):  
adiabatic index
- CNO neutrino break degeneracy with opacity → nuclear reactions main (and controllable) uncertainty
- Uncertainty in composition can be tamed by better CN neutrinos (models and experiments)

Science – only possible in the Sun

For solar-like oscillators:

- main target of PLATO missions – finding an Earth analog around solar-like stars
- $^{14}\text{N} + \text{p}$  is fundamental to model evolution of stars in range  $1.1$  to  $1.5 M_{\odot}$
- empirical determinations of overshooting – the largest uncertainty in stellar modeling – are contingent to our knowledge of  $^{14}\text{N} + \text{p}$  rate





Orebi-Gann et al. 2021

# Future: measurement of diffusive processes in the Sun?



Institute of  
Space Sciences



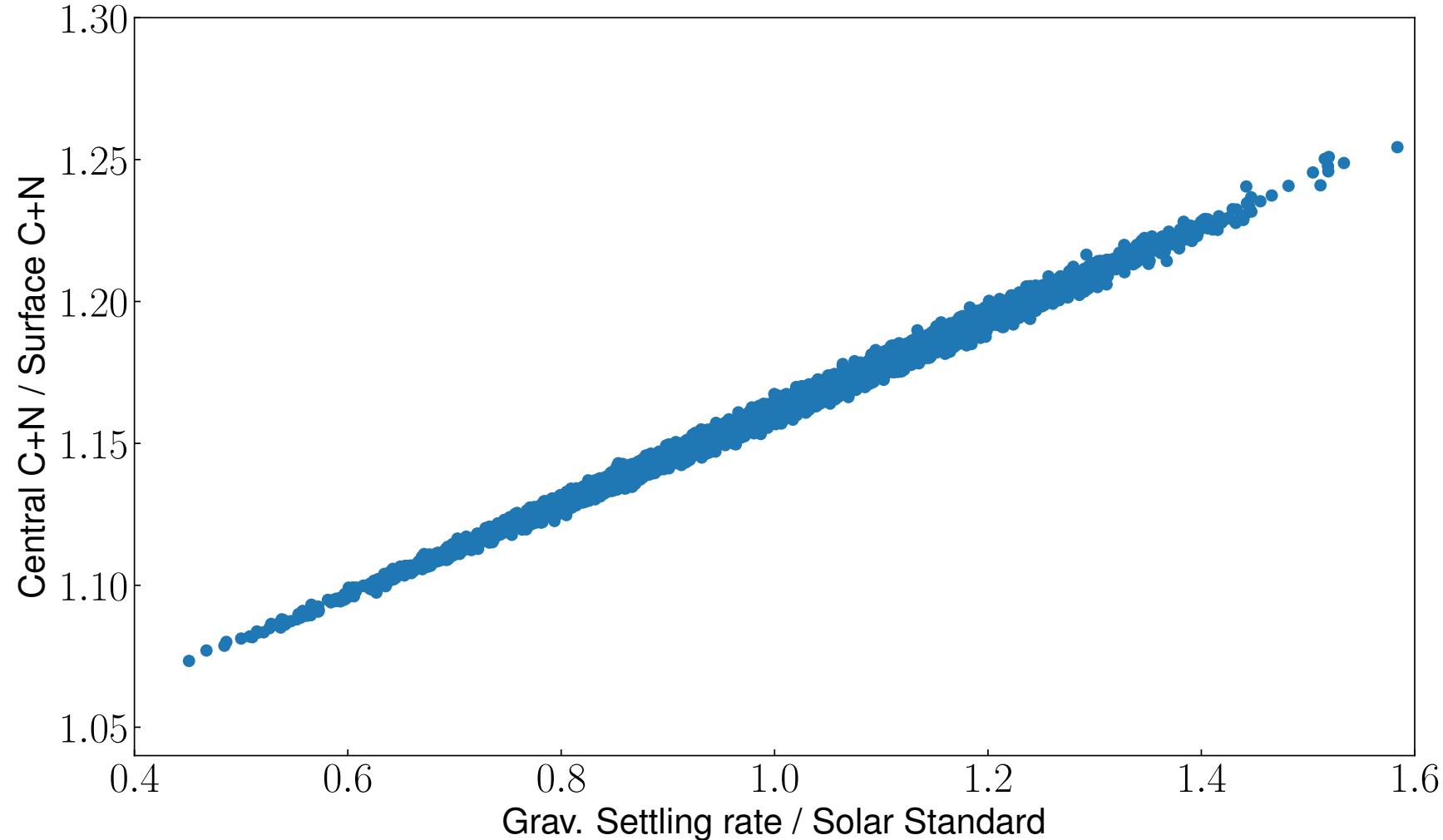
Current errors too large

- 10-15% spectroscopy
- 20% neutrinos

Errors of a few % would  
be needed

Possible in the (distant?)  
future

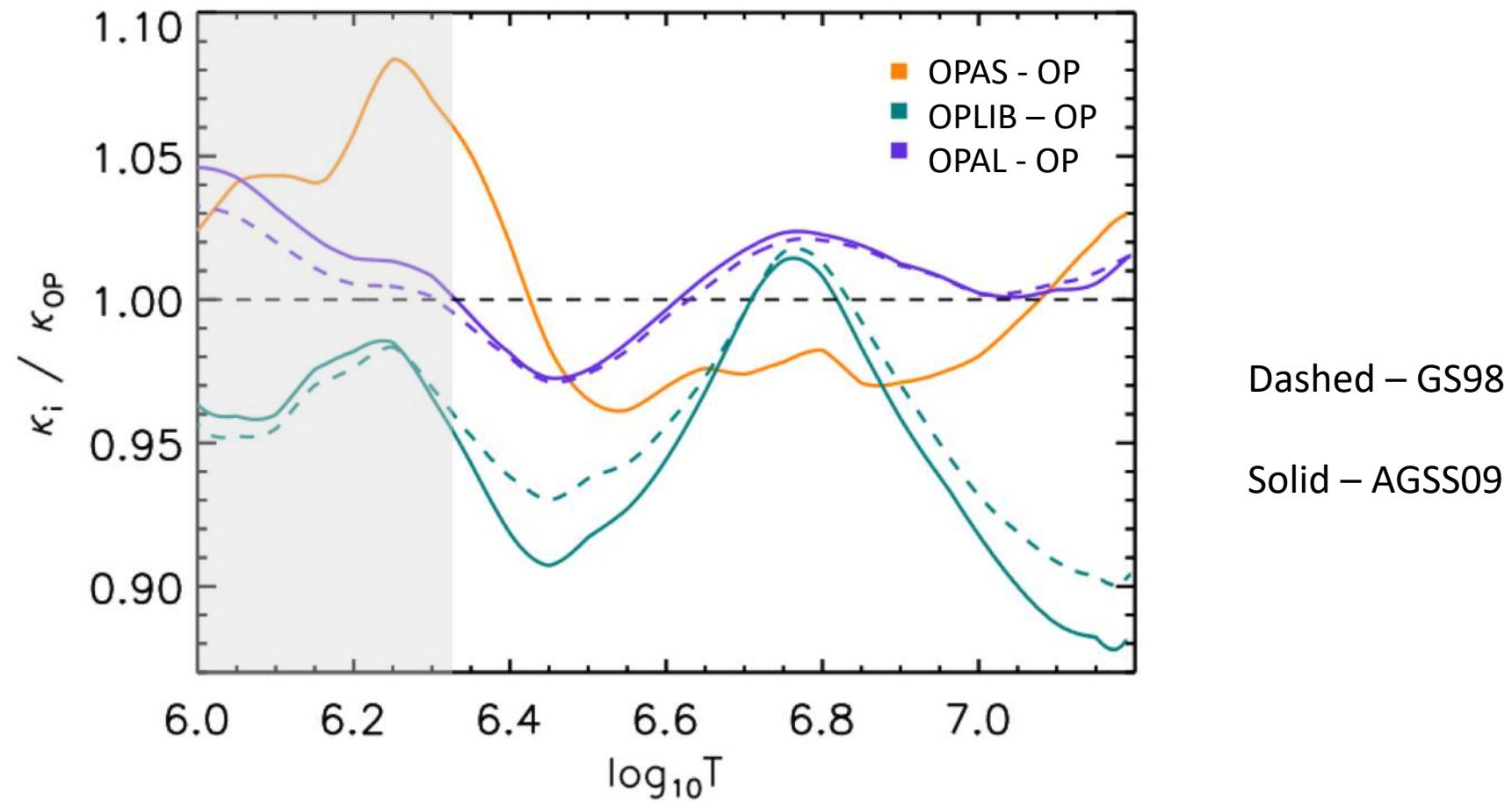
Ratio of central to surface C+N abundances



# Solar (stellar) radiative opacities



Institute of  
Space Sciences



# Solar (stellar) opacities

Bailey et al. 2015



Institute of  
Space Sciences

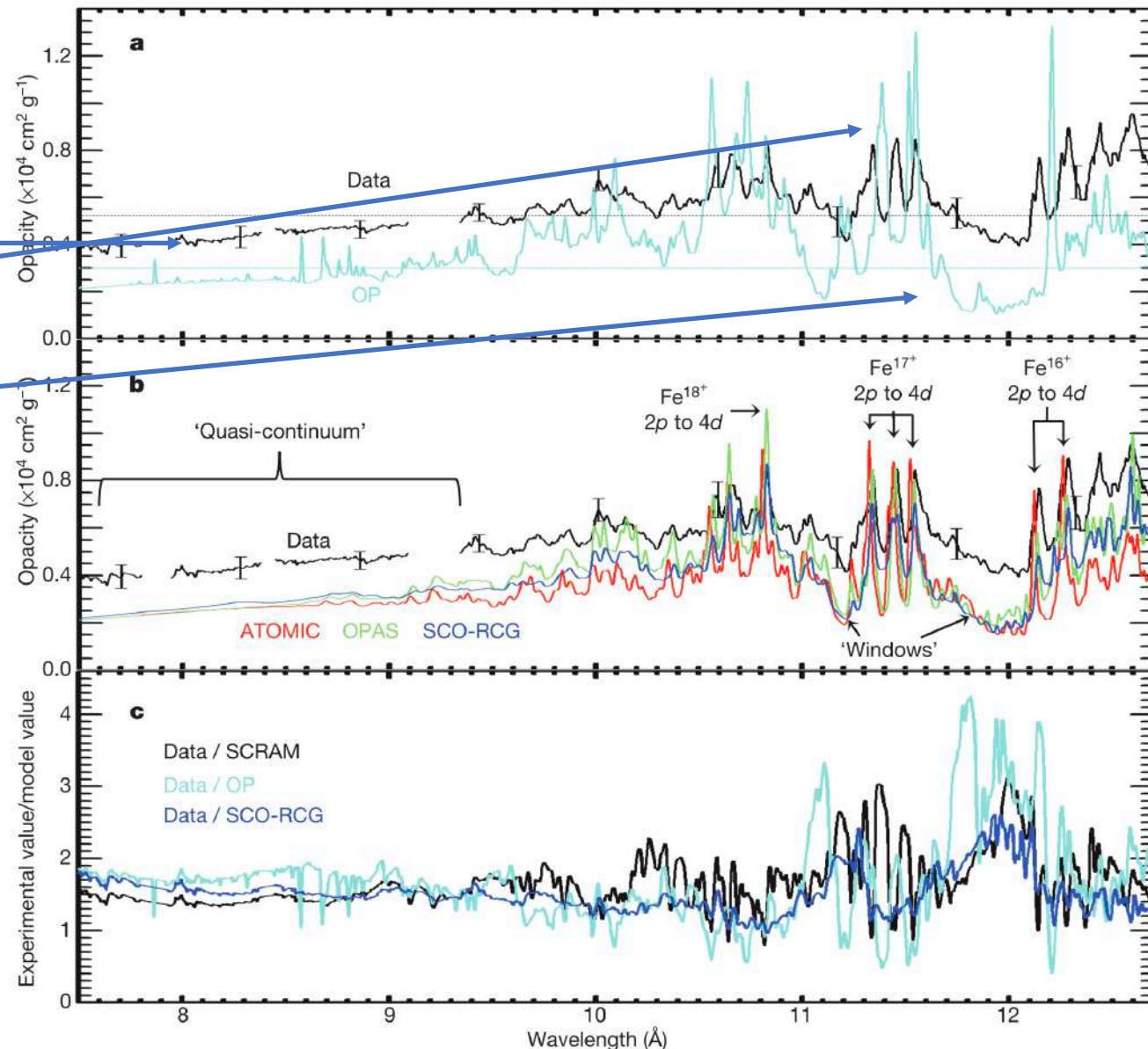


Calculations have:

1) Lower quasicontinuum

2) Narrower lines

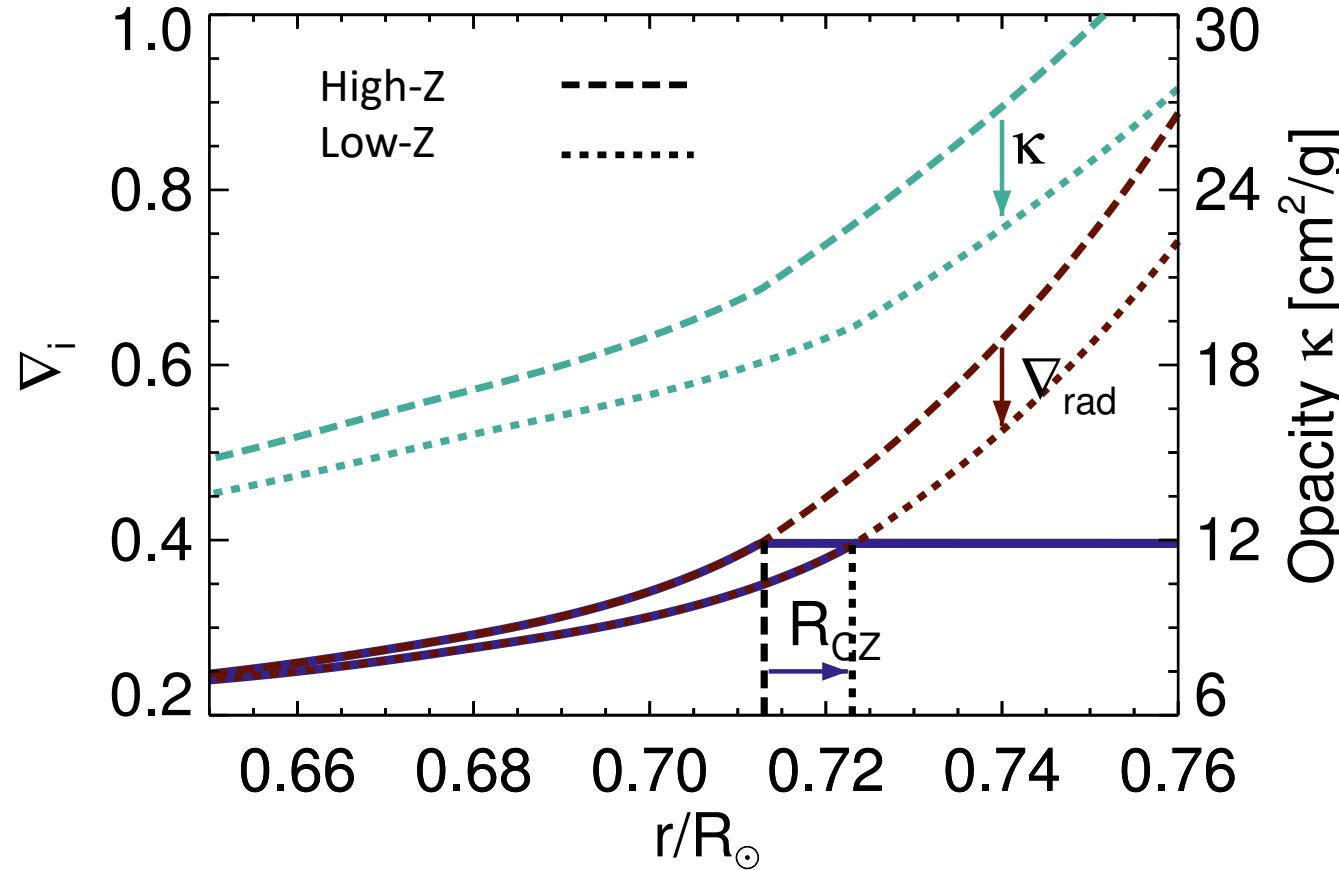
3) Deeper opacity windows



Experimental hint of higher opacity than theoretical calculations predict  
but situation unclear because of large differences in continuum



# Impact of metallicity



Solar model with low-Z has overall lower opacity

- flatter temperature profile
- slightly lower internal temperature
- affects helioseismology
- pp-chain neutrinos

Degeneracy between metals and opacity very difficult to break

**Opacities are the worst known fundamental piece of physics in solar/stellar modeling**