



Enrico Viguzzi INFN Milano

Inelastic and transfer reactions of light weakly bound nuclei

In collaboration with

Francisco Barranco University of Sevilla

Gregory Potel

Thanks to

Yassid Ayyad University of Santiago

Jie Chen Southern University of Science and Technology

NUSDAF Meeting , Catania, 22-24 Ottobre 2025

International Symposium Commemorating the 40th Anniversary of the Halo Nuclei

October 12-18, 2025 | Beijing



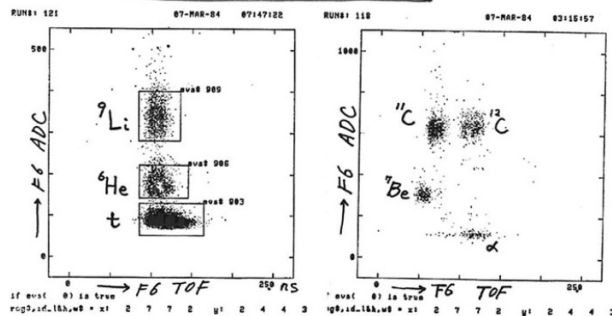
March 7, 1984

HOT!!

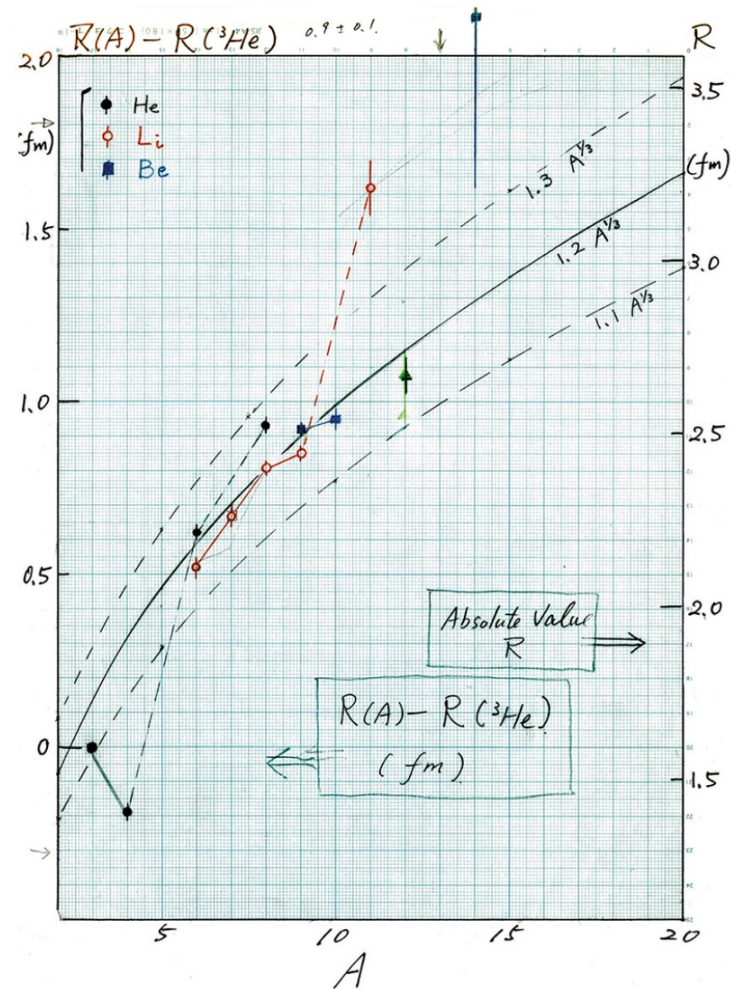
News from LBL.

We have got unstable nuclear
Beams.

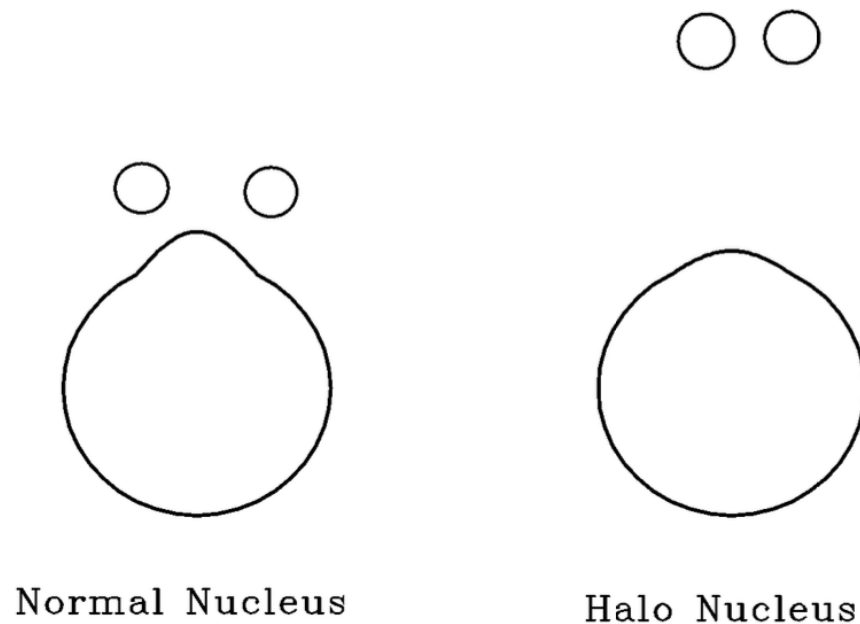
^{12}C Beam of ~ 800 A-MeV



小林. Tanihata, O. Hashimoto, 山崎



In 1997 Kuo et al predicted that in halo nuclei core polarization would be suppressed, and that the fundamental nucleon-nucleon interaction could be probed in a clearer and more direct way in halo nuclei than in ordinary nuclei....



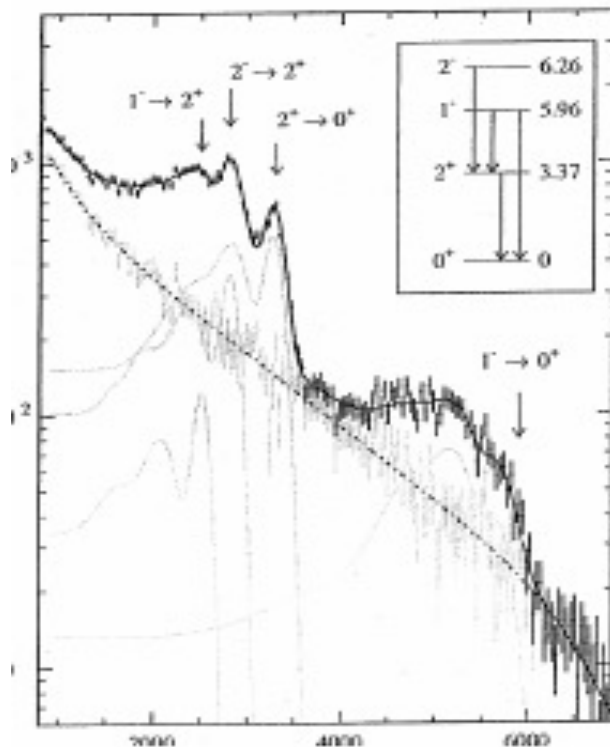
T.T.S. Kuo et al,
PRL 78 (1997) 2708

... But experiments demonstrated that the core dynamics plays an important role...

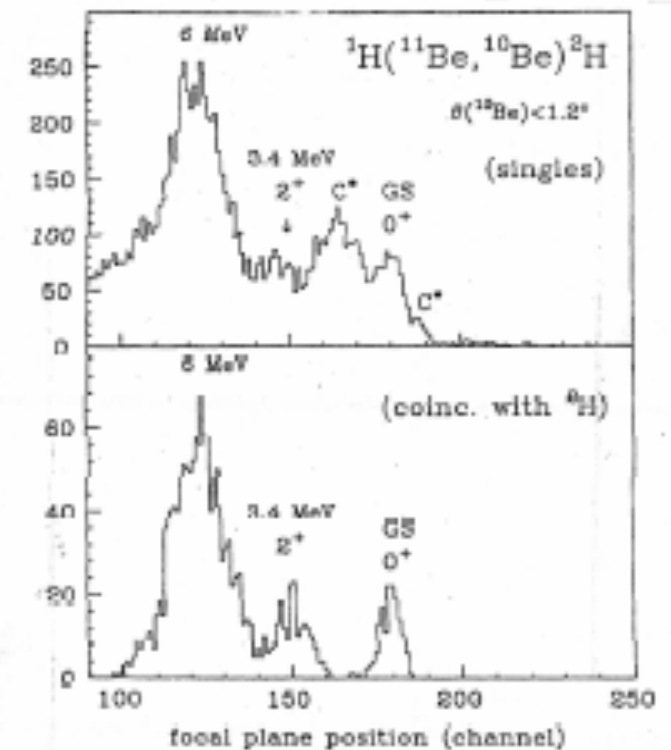
The admixture of $d_{5/2} \times 2^+$ configuration in the $1/2^+$ g.s. of ^{11}Be is about 15%

$p(^{11}\text{Be}, ^{10}\text{Be})d$

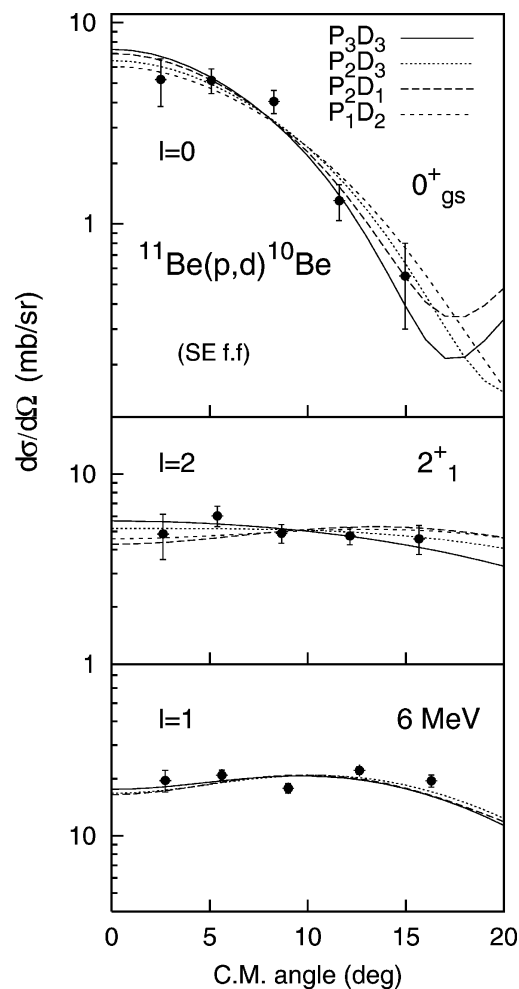
$9\text{Be}(^{11}\text{Be}, ^{10}\text{Be}+\gamma)\text{X}$



T. Aumann et al. PRL 84 (2000) 35



S. Fortier et al. Phys. Lett. B461 (1999)22
J.S. Winfield et al., Nucl.Phys. A683 (2001)48



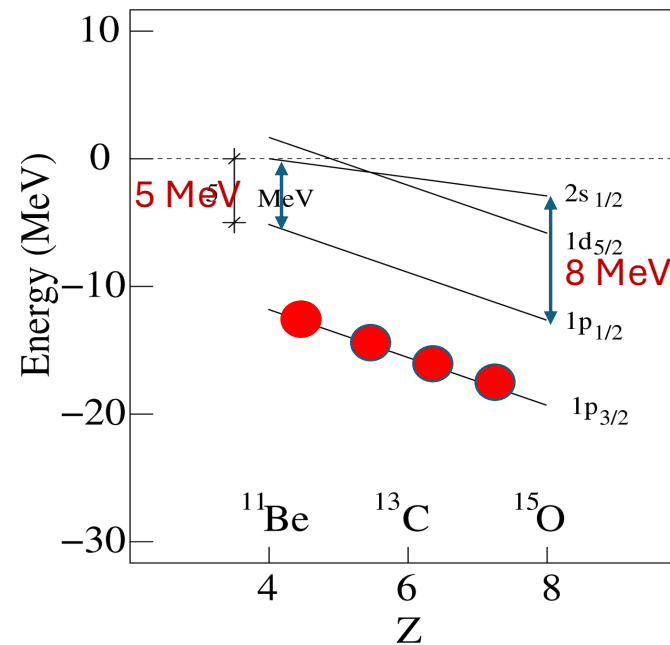
A careful analysis of transfer reactions is needed to estimate phonon admixtures in the wavefunctions

$$|^{11}\text{Be}_{gs}\rangle = \alpha |^{10}\text{Be}(0^+) \otimes 2s\rangle + \beta |^{10}\text{Be}(2^+) \otimes 1d\rangle$$

Good agreement with 2+ cross sections is obtained in DWBA with $\beta^2 = 0.17$ considering the coupling effects on the transfer form factor; using β as a simple spectroscopic factor one finds $\beta^2 = 0.28$

J.S. Winfield et al.,
Nucl.Phys. A683 (2001)48

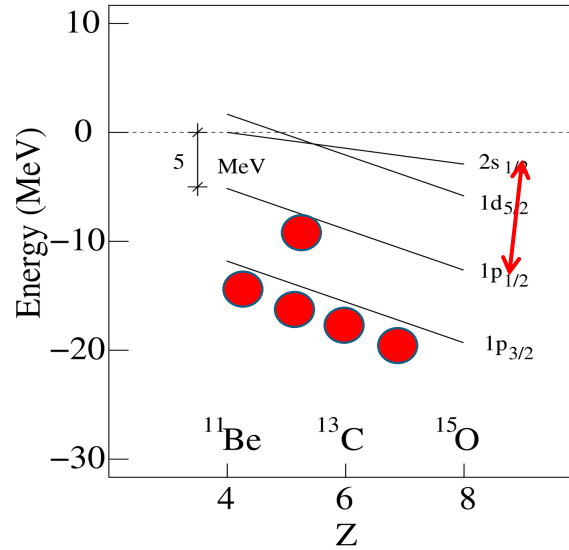
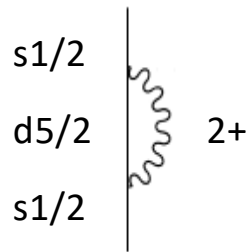
Parity inversion in N=7 isotones is not reproduced by spherical mean field obtained from non relativistic energy density functionals.



Typical spherical
mean-field results
with Skyrme forces
(Sagawa, Brown, Esbensen
PLB 309(93)1)

^{11}Be

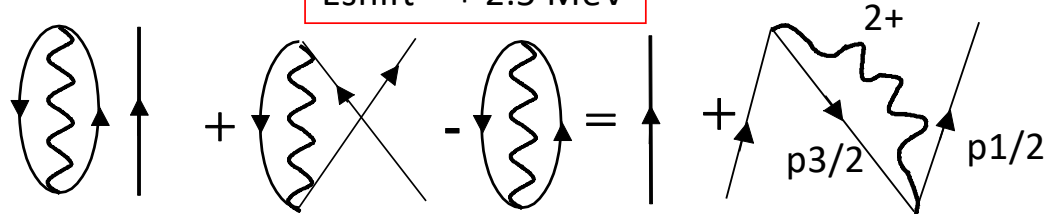
Eshift = - 2.5 MeV



Self-energy

+

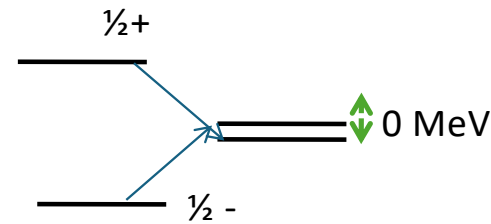
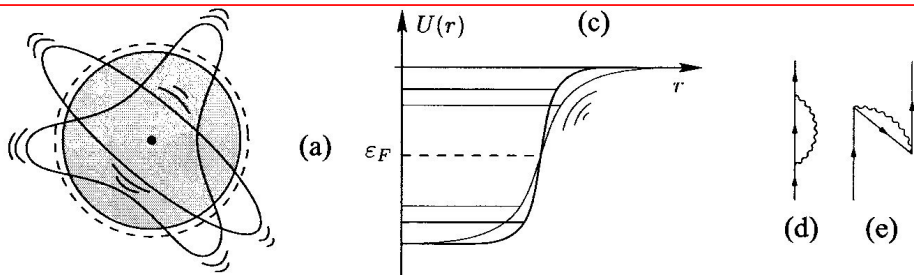
Eshift = + 2.5 MeV

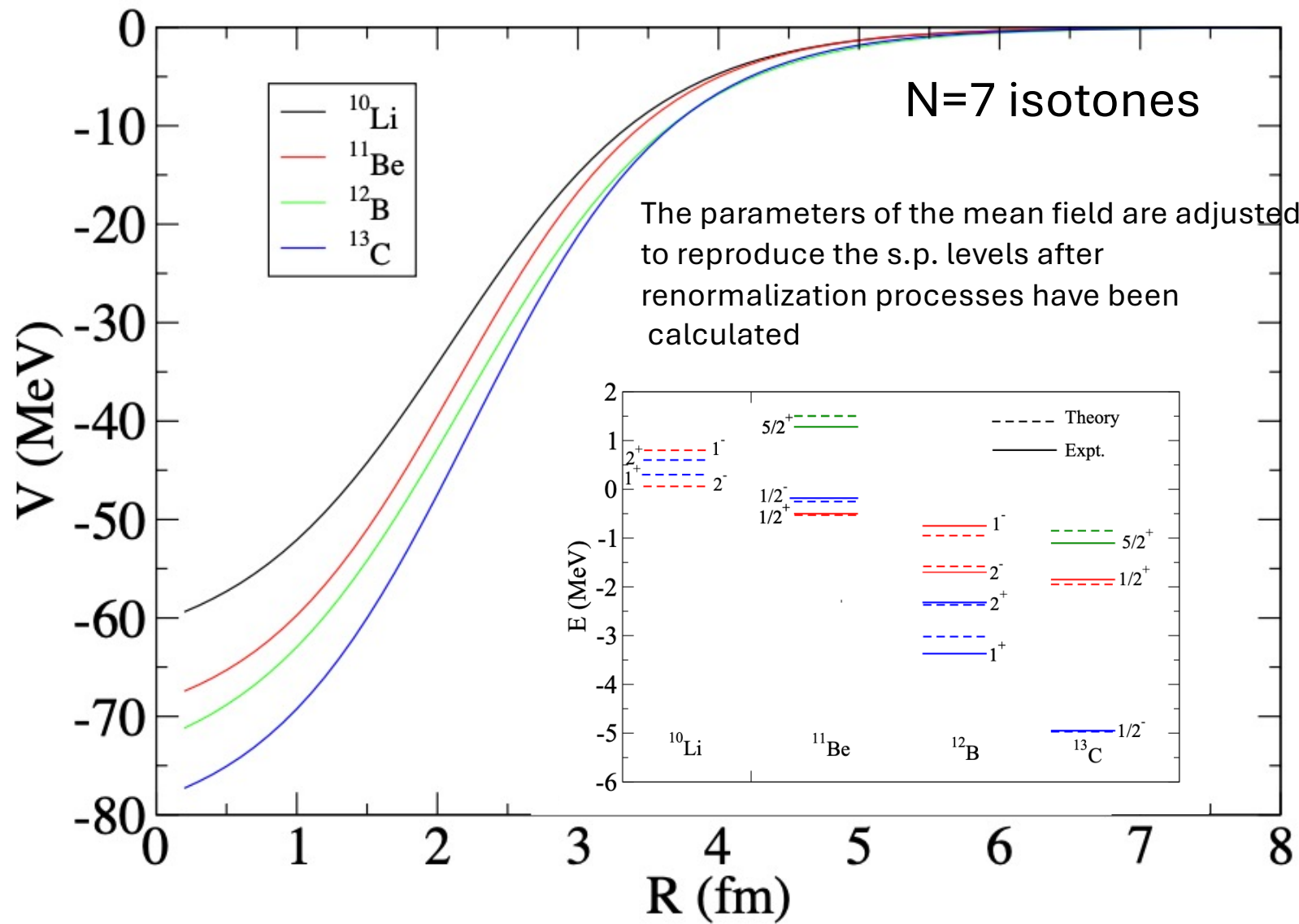


Pauli blocking of core ground state correlations



Level inversion

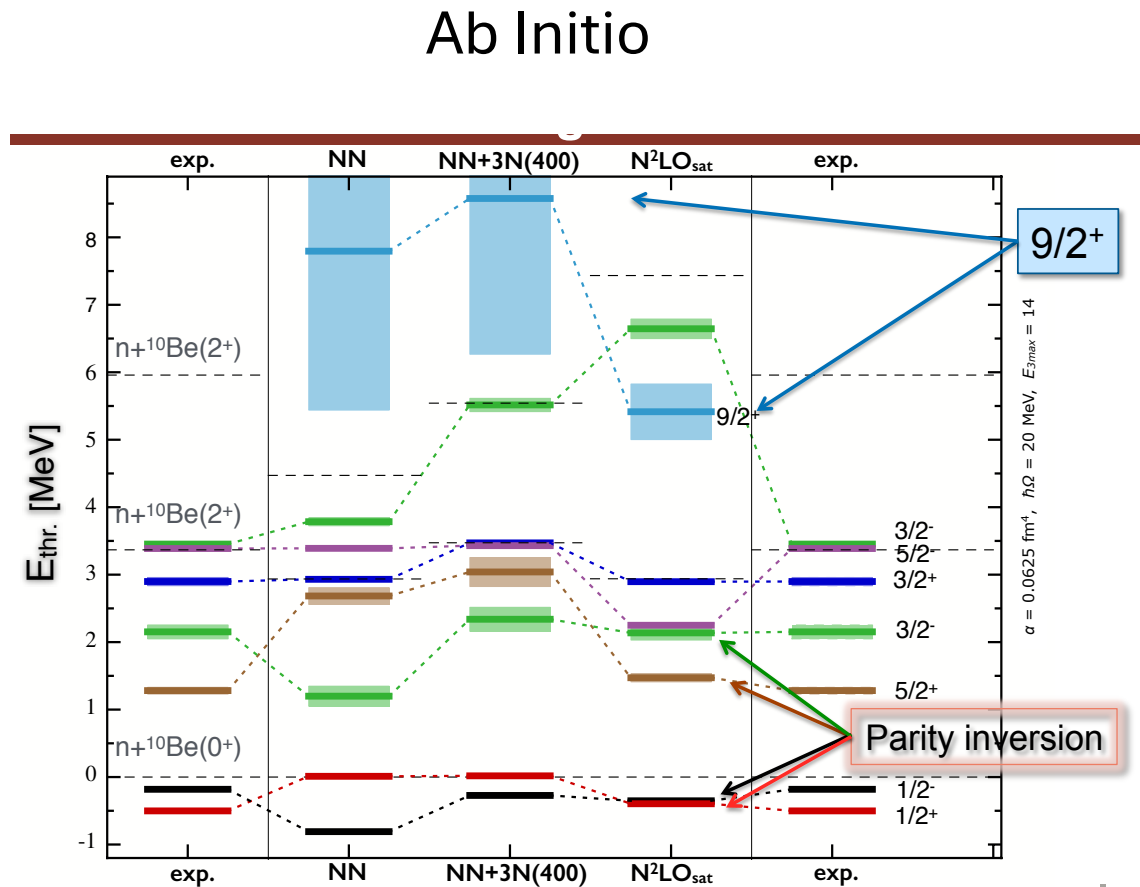
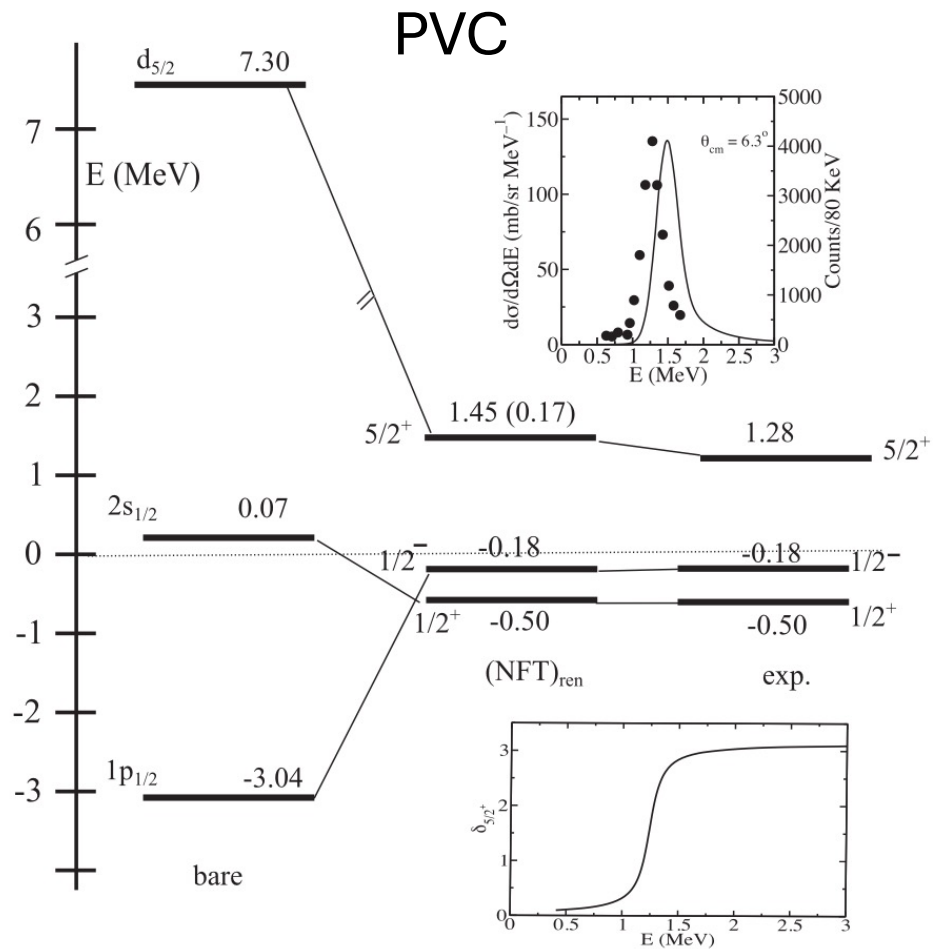




F. Barranco et al.,
 PRL 119 (2017)
 082501
 PRC 101 (2020)
 031305(R)

The description of the experimental results from complementary approaches can be of great interest

^{11}Be



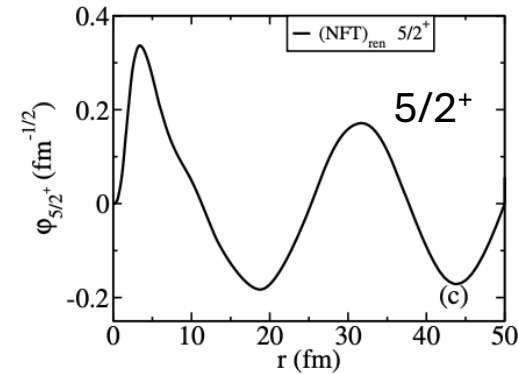
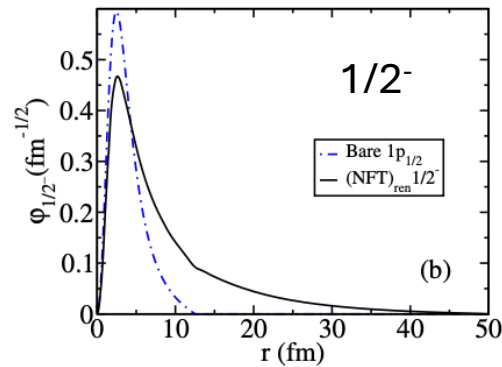
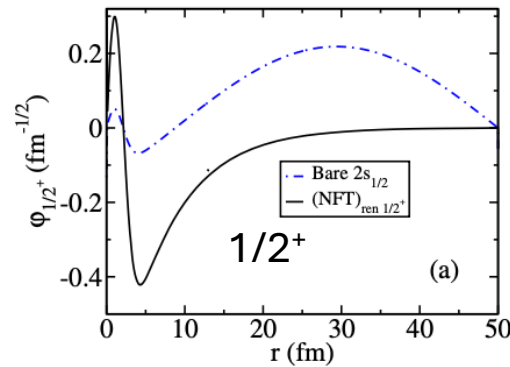
F. Barranco et al., PRL 119 (2017) 082501

A. Calci, P. Navrátil, R. Roth, J. Dohet-Eraly, S. Quaglioni, G. Hupin, PRL 117, 242501 (2016)

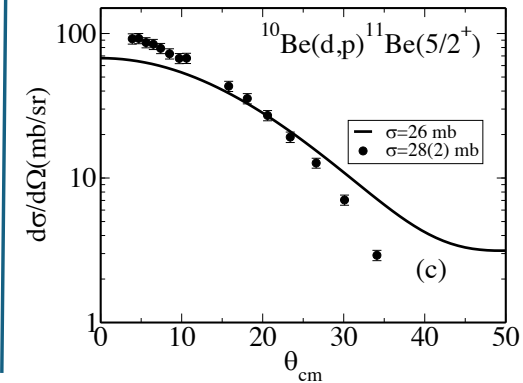
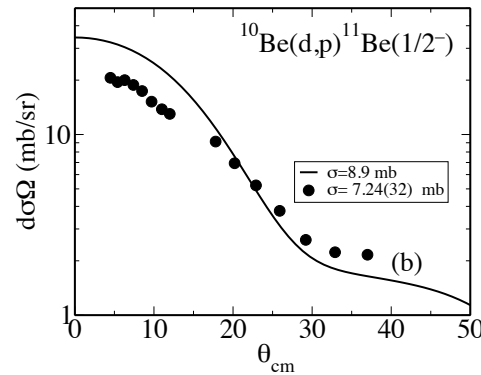
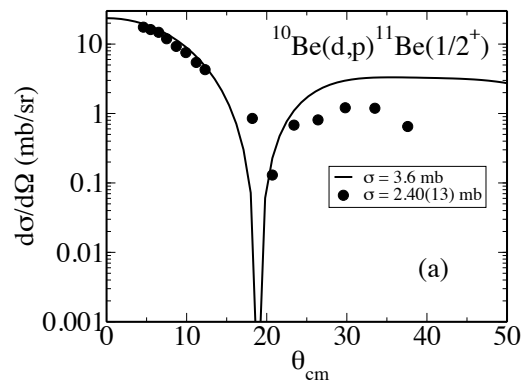
$^{10}\text{Be}(d,p)^{11}\text{Be}$ at $E_d = 21.4$ MeV

Test of the single-particle component of the many-body wavefunction

Form factors

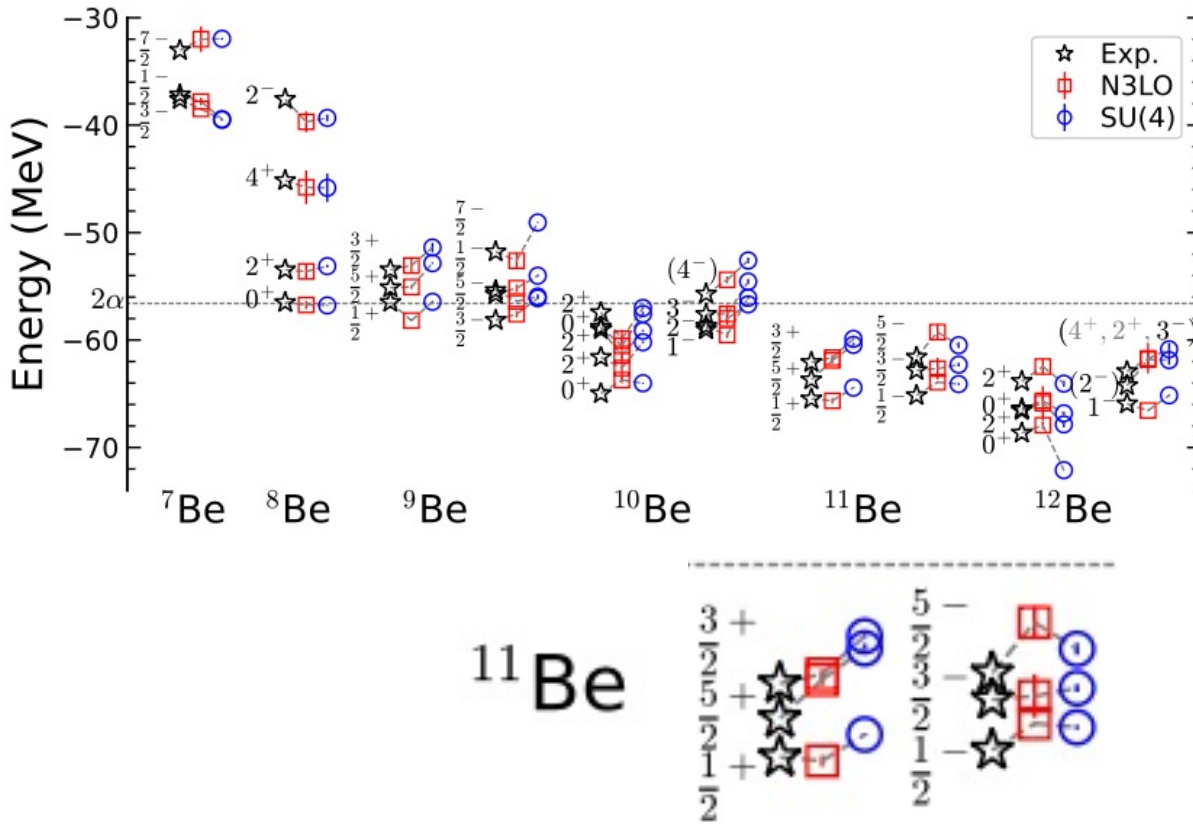


Cross sections



K.T. Schmitt et al., PRC88 (2012) 064612

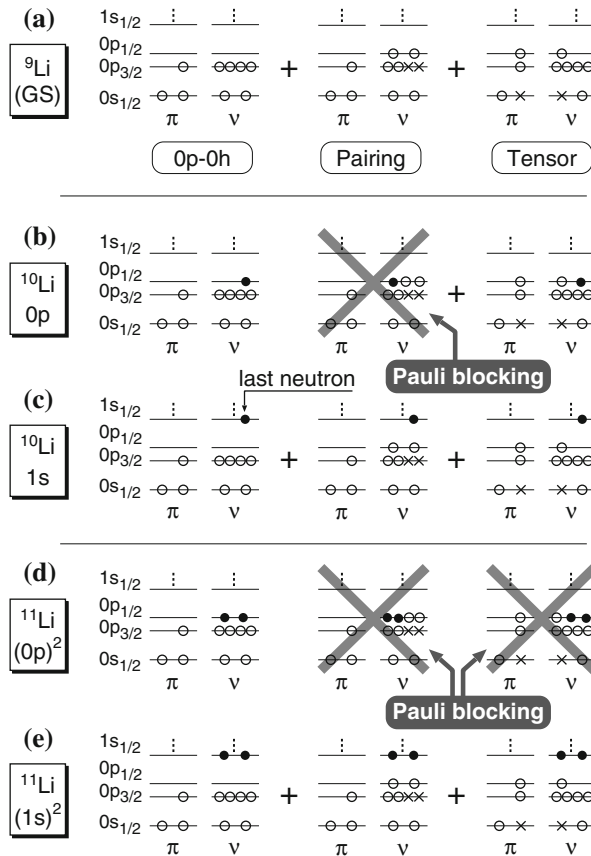
Lattice Effective field theory



		SU(4)	N ³ LO	Exp.
⁷ Be	$E2, \frac{3}{2}^- \rightarrow \frac{1}{2}^-$	16.0(2)	15.2(5)	26(6)(3) [91]
⁹ Be	$Q(\frac{3}{2}^-)$	7.3(1)	7.4(1.0)	5.29(4) [92]
	$E1, \frac{1}{2}^+ \rightarrow \frac{3}{2}^-$	0.131(3)	0.060(15)	0.136(2) [93]
	$E1, \frac{5}{2}^+ \rightarrow \frac{3}{2}^-$	0.045(14)	0.049(5)	0.010(8) [67]
	$E2, \frac{5}{2}^- \rightarrow \frac{3}{2}^-$	35.7(1.8)	27.8(1.9)	27.1(2.0) [67]
	$E2, \frac{7}{2}^- \rightarrow \frac{3}{2}^-$	11.6(2.5)	5.3(8)	9.5(4.1) [67]
¹⁰ Be	$E1, 3_1^- \rightarrow 2_1^+$	0.026(2)	0.004(3)	0.009(1) [67]
	$E2, 2_1^+ \rightarrow 0_1^+$	10.6(4)	8.5(9)	9.2(3) [31]
¹¹ Be	$E1, \frac{1}{2}^- \rightarrow \frac{1}{2}^+$	0.023(3)	0.038(3)	0.102(2) [94]
¹² Be	$E1, 0_1^+ \rightarrow 1_1^-$	0.049(2)	0.056(26)	0.051(13) [95]
	$E2, 2_1^+ \rightarrow 0_1^+$	7.8(1.1)	9.0(3.1)	14.2(1.0)(2.0) [6]

S. Shen et al, 134, 162503 (2025)

Comparison with the model by Ikeda, Myo et al.

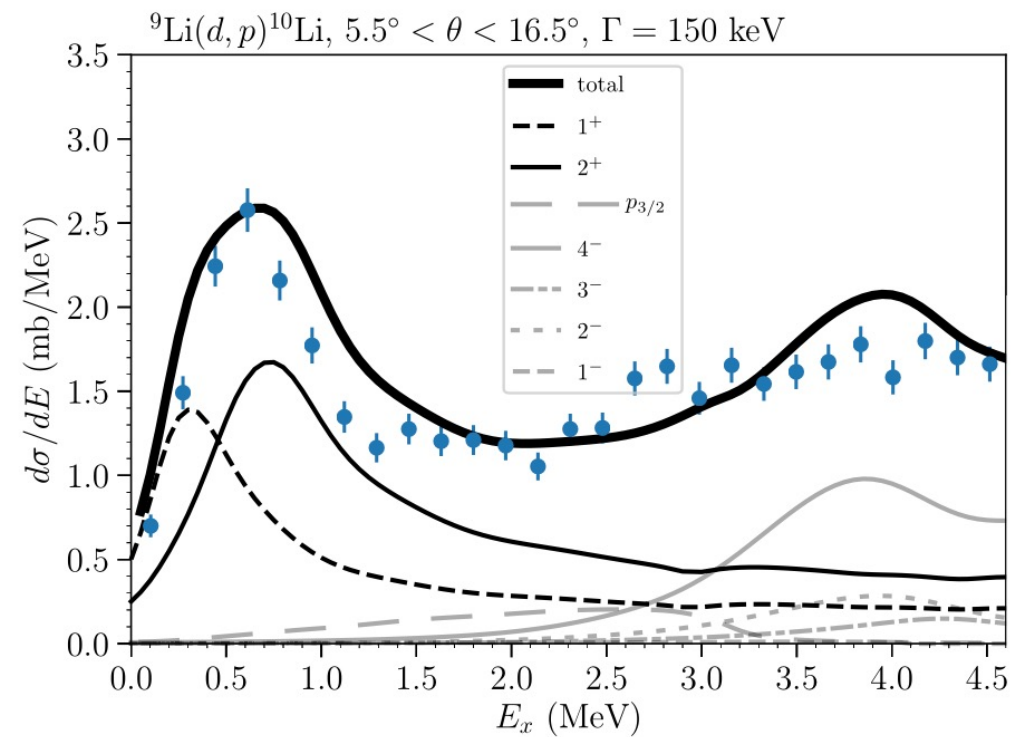
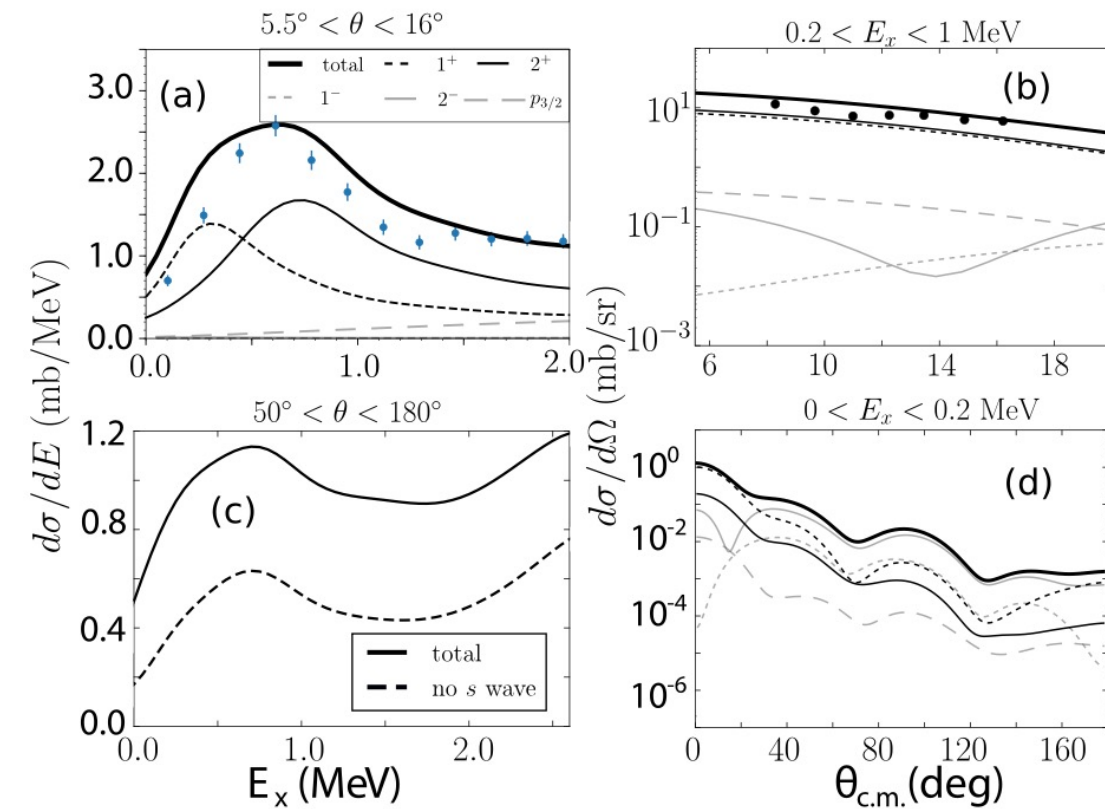


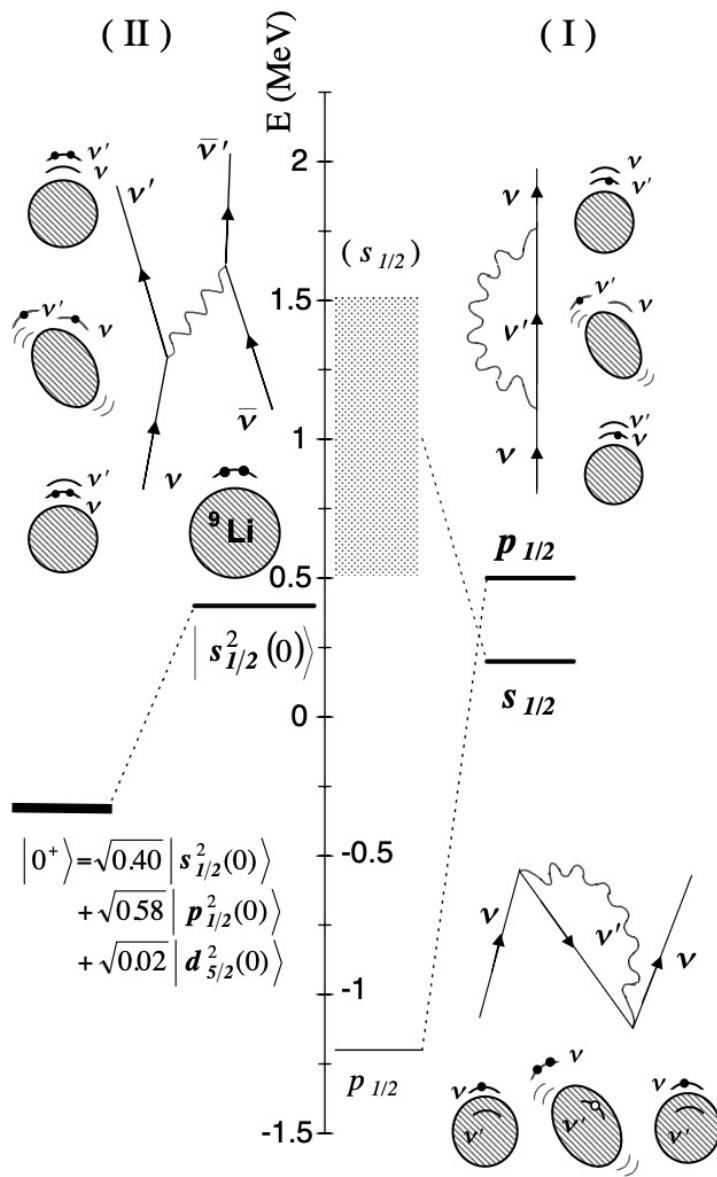
The $p_{1/2}$ orbit is pushed up by pairing correlations and tensor force. Only 3/2- configurations are included: coupling to core vibrations (1/2-) is not considered. Binding energy is given as input. 50%(s²)-50%(p²) wavefunction is obtained

K. Ikeda et al,
Lect. Notes in Physics 818 (2010)

^{10}Li and the $^9\text{Li}(d,p)^{10}\text{Li}$ reaction

F. Barranco et al,
Phys. Rev. C 101 (2020)
031305(R)





^{11}Li correlated wave function

$$|\tilde{0}\rangle = |0\rangle + 0.7|(ps)_{1^-} \otimes 1^-; 0\rangle + 0.1|(sd)_{2^+} \otimes 2^+; 0\rangle$$

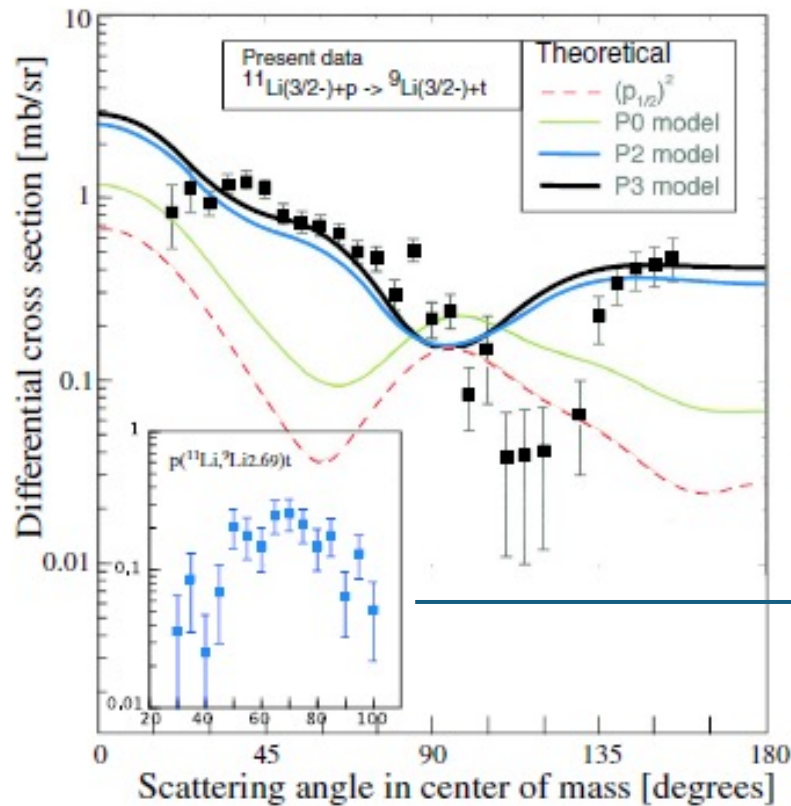
$$|0\rangle = 0.45|s_{1/2}^2(0)\rangle + 0.55|p_{1/2}^2(0)\rangle + 0.04|d_{5/2}^2(0)\rangle$$

^{11}Li

F. Barranco et al. EPJ A11 (2001) 385

Probing ^{11}Li halo-neutrons correlations via (p,t) reaction

I. Tanihata et al, PRL 100 (2008) 192502

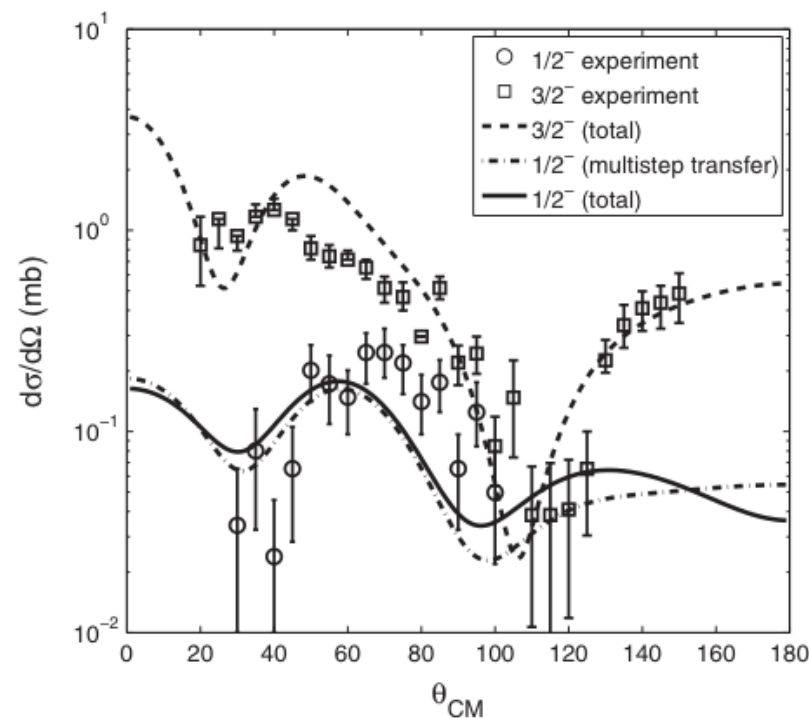


Unique information:
angular distribution associated with
the excited state of the ^9Li core

TABLE I. Optical potential parameters used for the present calculations.

	V MeV	r_V fm	a_V fm	W MeV	W_D MeV	r_W fm	a_W fm	V_{so} MeV	r_{so} fm	a_{so} fm
$p + ^{11}\text{Li}$ [10]	54.06	1.17	0.75	2.37	16.87	1.32	0.82	6.2	1.01	0.75
$d + ^{10}\text{Li}$ [11]	85.8	1.17	0.76	1.117	11.863	1.325	0.731	0		
$t + ^9\text{Li}$ [12]	1.42	1.16	0.78	28.2	0	1.88	0.61	0		

Calculation of absolute two-nucleon transfer cross section
by finite-range DWBA calculation



$$|\bar{0}\rangle = |0\rangle + 0.7|(ps)_{1-} \otimes 1^-; 0\rangle + 0.1|(sd)_{2+} \otimes 2^+; 0\rangle$$

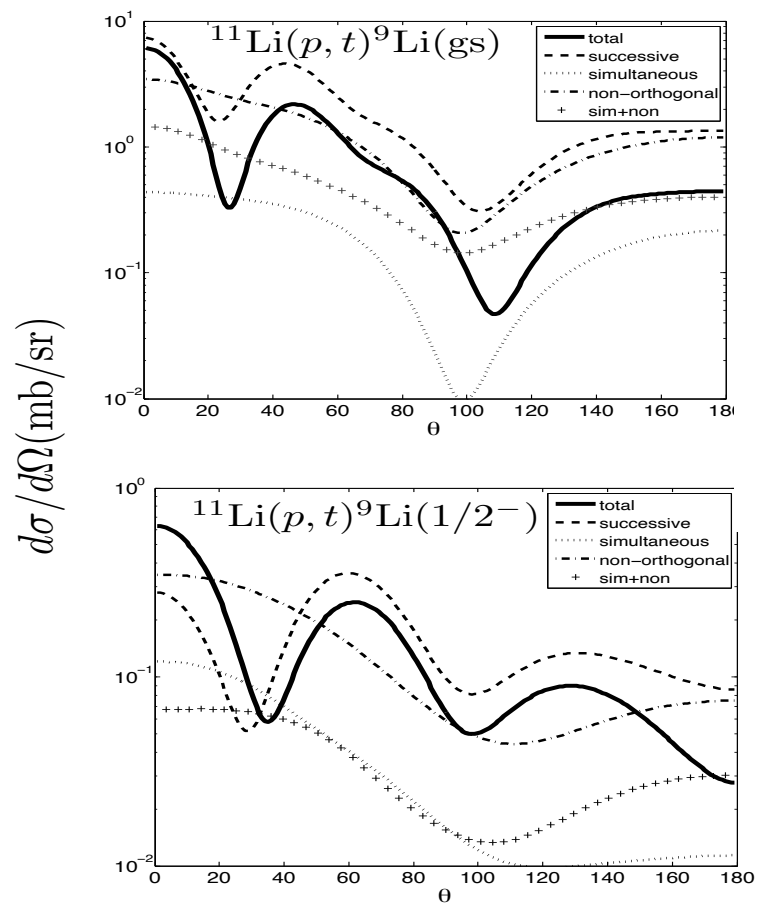
Essential component to populate
the excited state of the core

This calculation should be
repeated by other groups!

$\sigma(^{11}\text{Li}(\text{gs}) \rightarrow ^9\text{Li}(\text{i})) \text{ (mb)}$			
i	ΔL	Theory	Experiment
gs ($3/2^-$)	0	6.1	5.7 ± 0.9
2.69 MeV ($1/2^-$)	2	0.5	1.0 ± 0.36

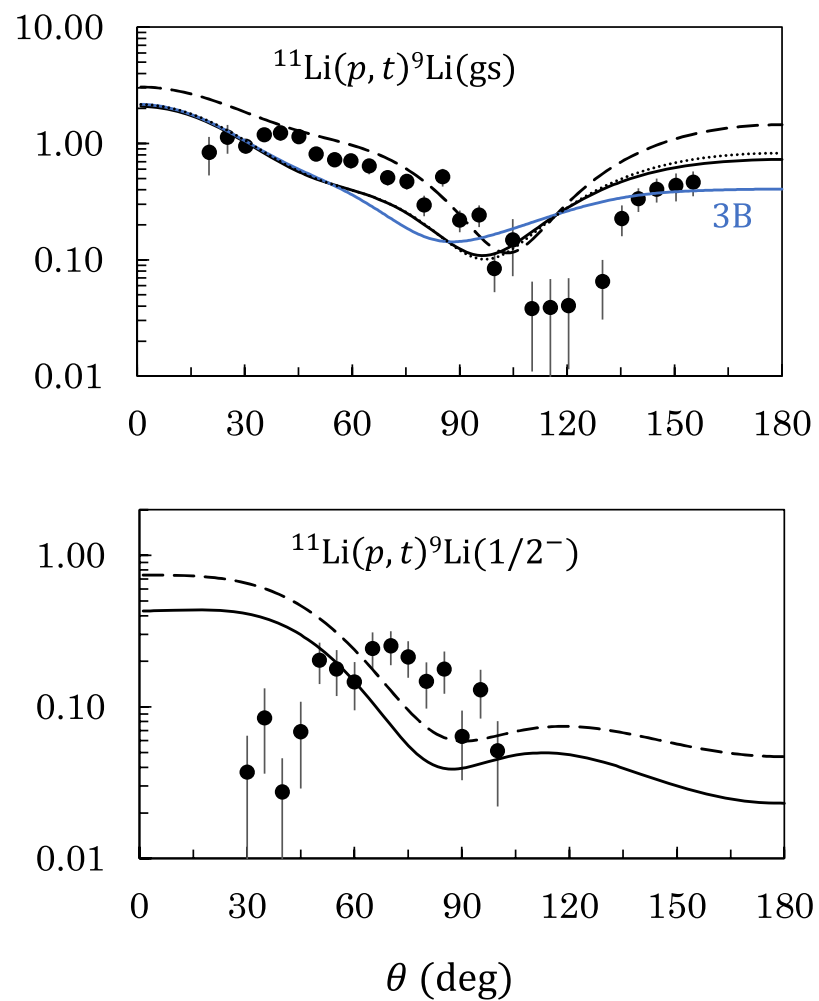
G. Potel et al., PRL 105 (2010) 172502

^{11}Li w.f.: 50% 1- admixture
1% 2+ admixture

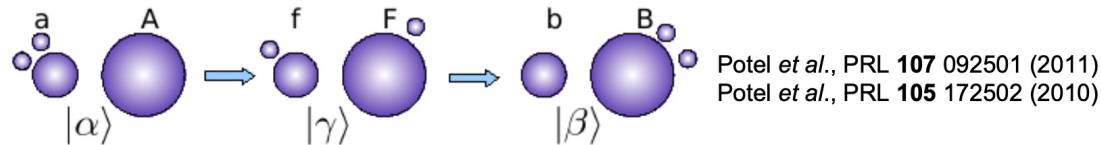


^{11}Li w.f. : 10% 2+ admixture

P. Descouvemont, PLB **862** (2025) 139356



Second order DWBA



$$\sigma \sim |T_{2NT}|^2 = \left| \sum_{j_f, j_i} B_{j_f} B_{j_i} \left(T^{(1)} + T_{succ}^{(2)} + T_{NO}^{(2)} \right) \right|^2 \Rightarrow \text{coherent sum of many terms}$$

structure input:
coefficients of the initial and final
wavefunctions (e.g., from a BCS
calculation)

Successive transfer

$$\begin{aligned} T_{succ}^{(2)}(j_i, j_f) = & 2 \sum_{K, M} \sum_{\substack{\sigma_1 \sigma_2 \\ \sigma'_1 \sigma'_2}} \int d\mathbf{r}_{fF} d\mathbf{r}_{b1} d\mathbf{r}_{A2} [\psi^{j_f}(\mathbf{r}_{A1}, \sigma_1) \psi^{j_f}(\mathbf{r}_{A2}, \sigma_2)]_0^{0*} \\ & \times \chi_{bB}^{(-)*}(\mathbf{r}_{bB}) v(\mathbf{r}_{b1}) [\psi^{j_f}(\mathbf{r}_{A2}, \sigma_2) \psi^{j_i}(\mathbf{r}_{b1}, \sigma_1)]_M^K \\ & \times \int d\mathbf{r}'_{fF} d\mathbf{r}'_{b1} d\mathbf{r}'_{A2} G(\mathbf{r}_{fF}, \mathbf{r}'_{fF}) [\psi^{j_f}(\mathbf{r}'_{A2}, \sigma'_2) \psi^{j_i}(\mathbf{r}'_{b1}, \sigma'_1)]_M^K \\ & \times \frac{2\mu_{fF}}{\hbar^2} v(\mathbf{r}'_{f2}) [\psi^{j_i}(\mathbf{r}'_{b2}, \sigma'_2) \psi^{j_i}(\mathbf{r}'_{b1}, \sigma'_1)]_0^{0*} \chi_{aA}^{(+)}(\mathbf{r}'_{aA}) \end{aligned}$$

Study of (p,t) reaction on ^{11}Li at 6 MeV/nucleon

X. Wang,^{1,2} J. L. Ferreira,³ J. Lubian,³ J. Tanaka,² R. Kanungo,^{4,5} M. Alcorta,⁵ N. Aoi,^{2,6} H. Bidaman,⁷ C. Burbadge,⁷ G. Christian,⁵ S. Cruz,⁵ B. Davids,⁵ A. Diaz Varela,⁷ J. Even,^{5,8} G. Hackman,⁵ J. Henderson,⁵ S. Ishimoto,⁹ S. Kaur,⁴ M. Keefe,⁴ R. Krücken,^{5,10} K.G. Leach,⁵ J. Lighthall,⁵ E. Padilla Rodal,^{5,11} J. S. Randhawa,⁴ P. Ruotsalainen,⁵ A. Sanetullaevd,⁵ J. K. Smith,⁵ O. Workman,⁴ and I. Tanihata¹

¹*Department of Physics, Beihang University, Beijing 100191, China*

²*RCNP, Osaka University, 10 - 1, Mihogaoka, Ibaraki, Osaka 567 - 0047, Japan*

³*Instituto de Física, Universidade Federal Fluminense, Niterói, Rio de Janeiro, 24210 - 340, Brazil*

⁴*Astronomy and Physics Department, Saint Mary's University, Halifax, Nova Scotia B3H3C3, Canada*

⁵*TRIUMF, Vancouver, V6T2A3, Canada*

⁶*CNS, University of Tokyo, Bunkyo - ku, Tokyo 113 - 0033, Japan*

⁷*Department of Physics, University of Guelph, Guelph N1G2W1, Canada*

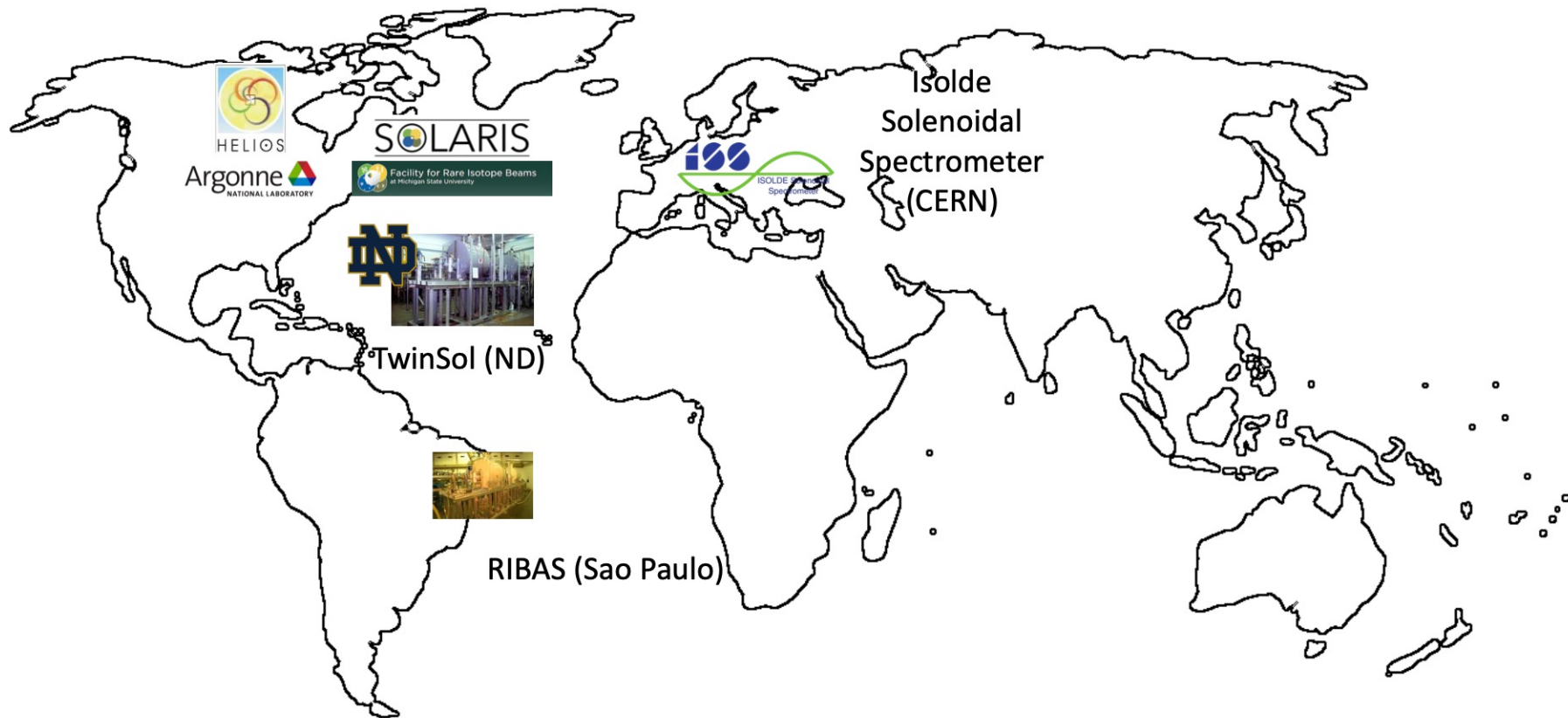
⁸*KVI-CART, University of Groningen, 9747 AA Groningen, The Netherlands*

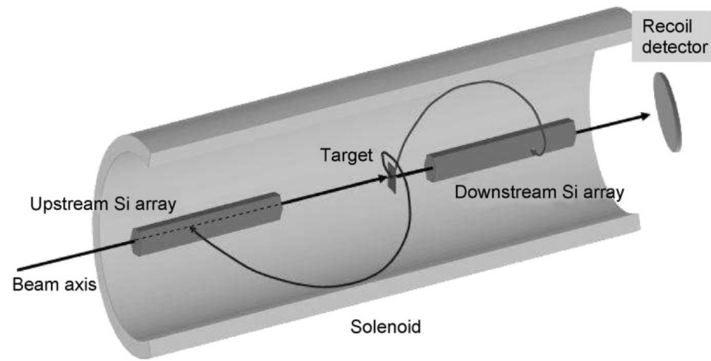
⁹*High Energy Accelerator Research Organization (KEK), Ibaraki 305 - 0801, Japan*

¹⁰*Department of Physics and Astronomy, University of British Columbia, Vancouver, BC V6T1Z1, Canada*

¹¹*Instituto de Ciencias Nucleares, UNAM, Mexico City 04510, Mexico*

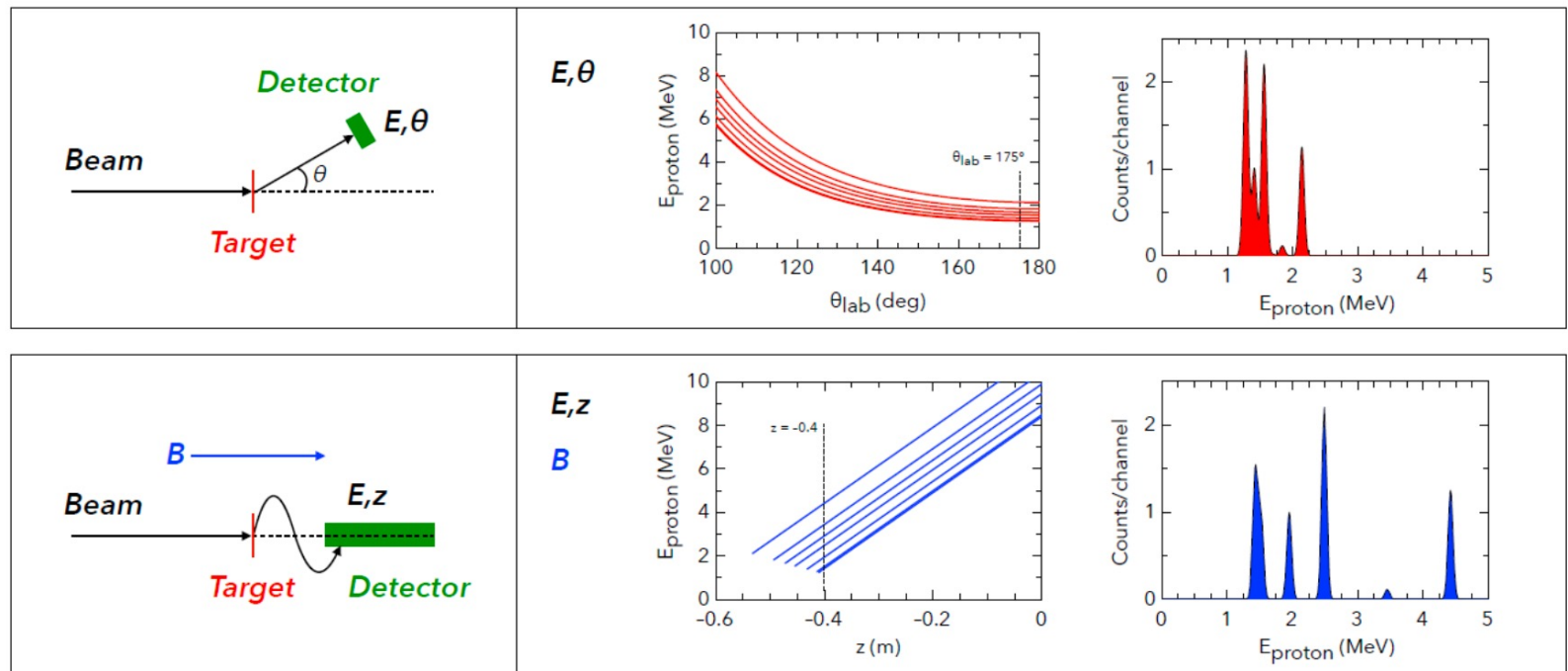
Solenoidal spectrometers around the world





$$T_{\text{cyc}} = \frac{2\pi r}{v_{\perp}} = \frac{2\pi m}{\mathcal{B} q e}.$$

The solenoidal-spectrometer technique



Cluster structure in ^{10}Be and dipole excitations

➤ Cluster structure in ^{10}Be g.s. has been confirmed by cluster knockout reaction.

P. J. Li *et al.*, Phys. Rev. Lett. 131, 212501 (2023)

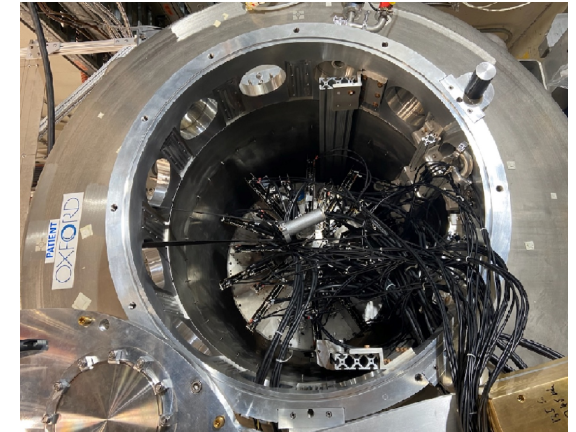
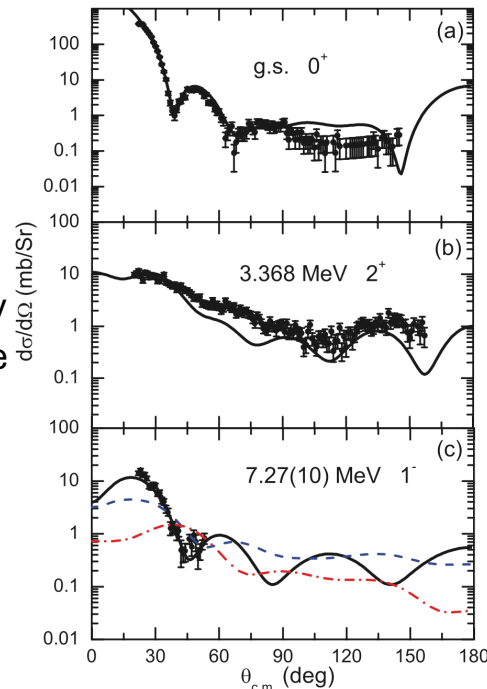
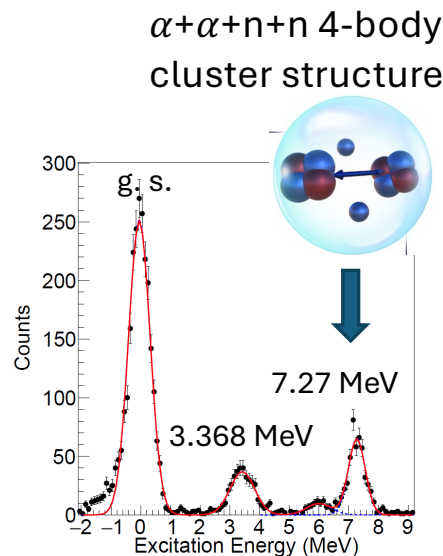
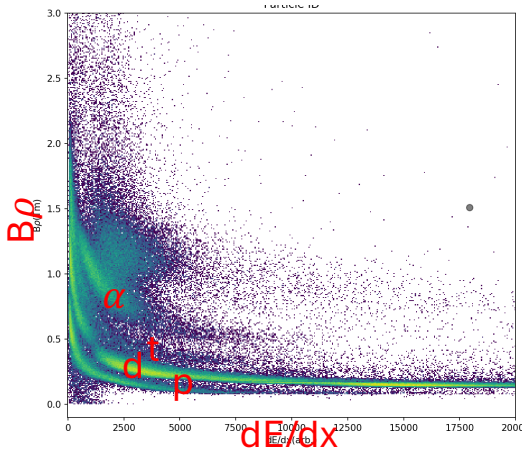
➤ $^{10}\text{Be}(d,d')^{10}\text{Be}$

AT-TPC Filled with D_2 : 600 Torr

^{10}Be beams @ 9.1 MeV/u

2000 pps

SOLARIS field: $B=3\text{ T}$



- Energy resolution $\sim 700\text{ keV}$
- Simultaneous measurement of (d,p) (d,t) and (d,d') reactions with $\sim 100\text{ pps}$ beam within one week!
- Open wide opportunities in the future with weak exotic beams

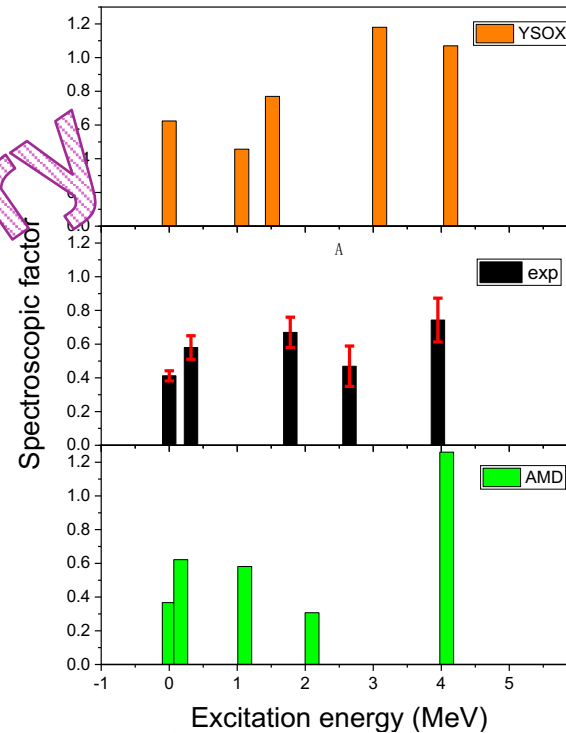
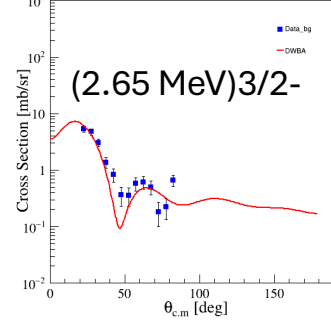
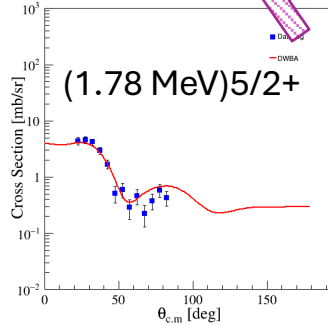
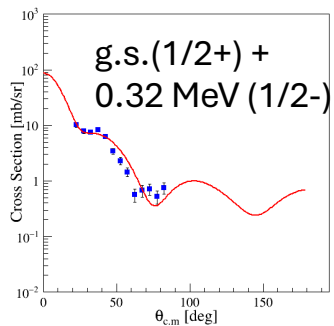
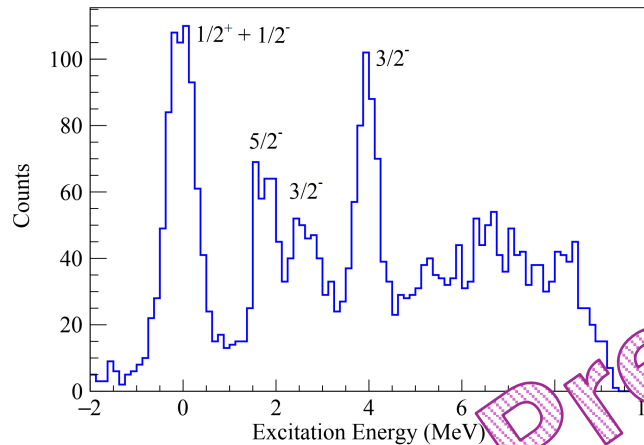
J. Chen, Y. Ayyad *et al.* Phys. Rev. Lett. 134, 012502 (2025).

One-neutron removal reaction on ^{12}Be

- $^{12}\text{Be}(p,d)^{11}\text{Be}$ measurement
- Configuration mixing in the psd shell
- Occupancy in the $1s_{1/2}, 0p_{1/2}, 0d_{5/2}$ of the ^{12}Be g.s.

AT-TPC Filled with H_2 : 300 Torr
HELIOS @ ATLAS: B=2.5 T

Neutron removal reaction $^{12}\text{Be}(p,d)^{11}\text{Be}$

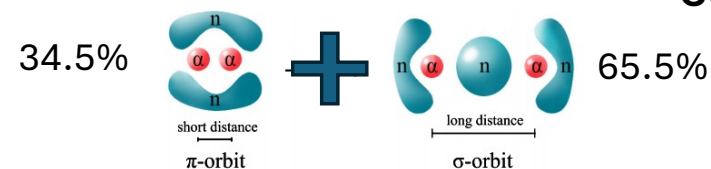


shell model →
Single-particle
structure

experiment

AMD model →
Cluster structure

Courtesy: Qing Zhao

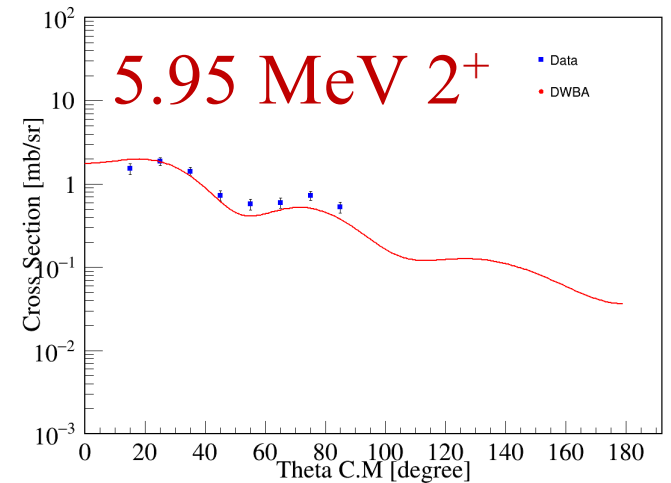
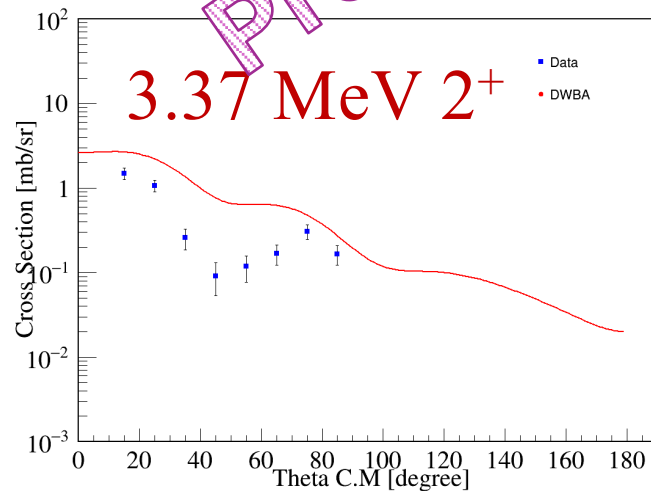
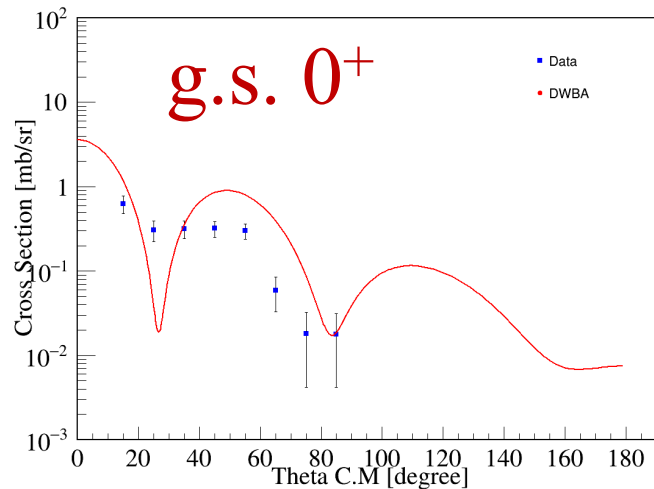
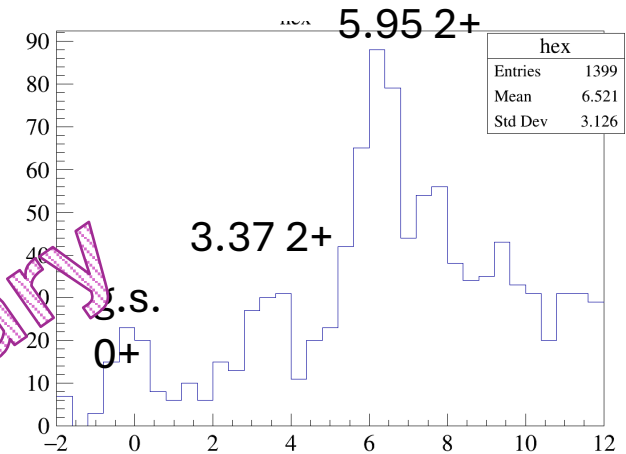


Two-neutron removal reaction on ^{12}Be

-- How much core excitations in ^{12}Be g.s.?

$^{12}\text{Be}(p,t)$ excite the g.s. and low-lying state of ^{10}Be .

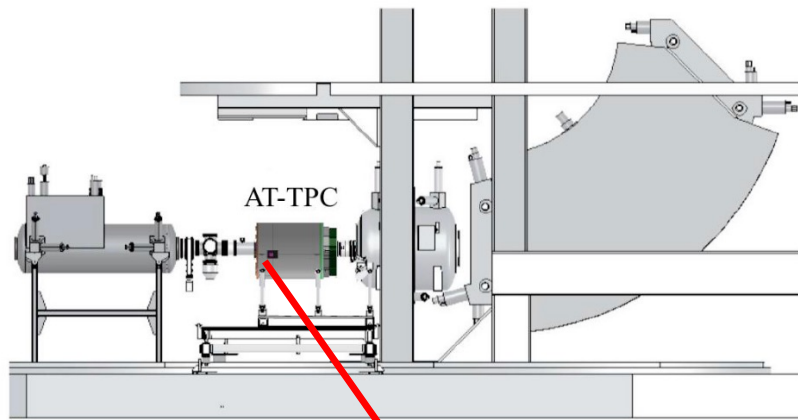
- The g.s. and first $2+$ state was populated.
- The second $2+$ state at ~ 6 MeV was strongly populated.



New measurement of dipole resonance of ^{11}Li

-- Proton inelastic scattering of ^{11}Li with AT-TPC+S800

The S800 spectrograph beamline at FRIB



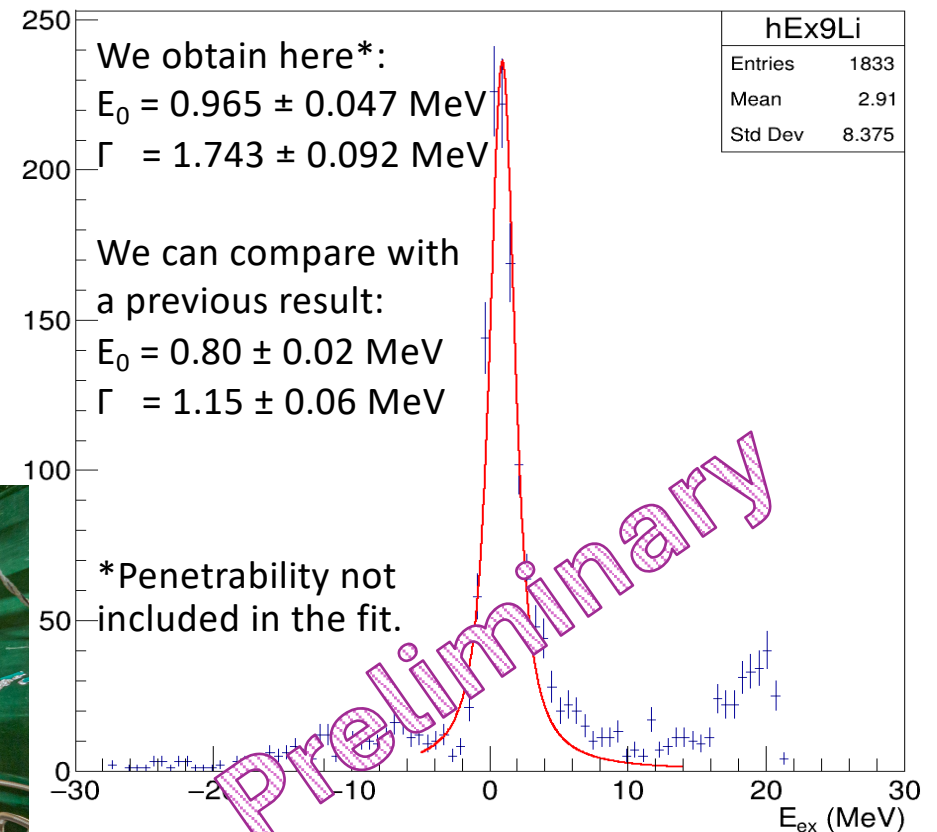
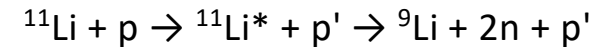
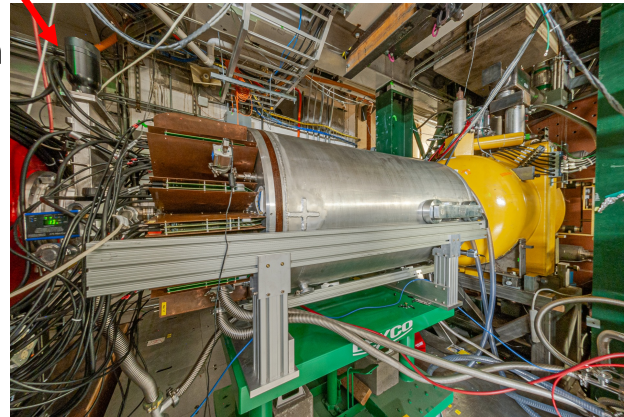
$1\text{E}+4$ pps, ~ 53.4 MeV/u, ^{11}Li beam

Target: AT-TPC filled with:

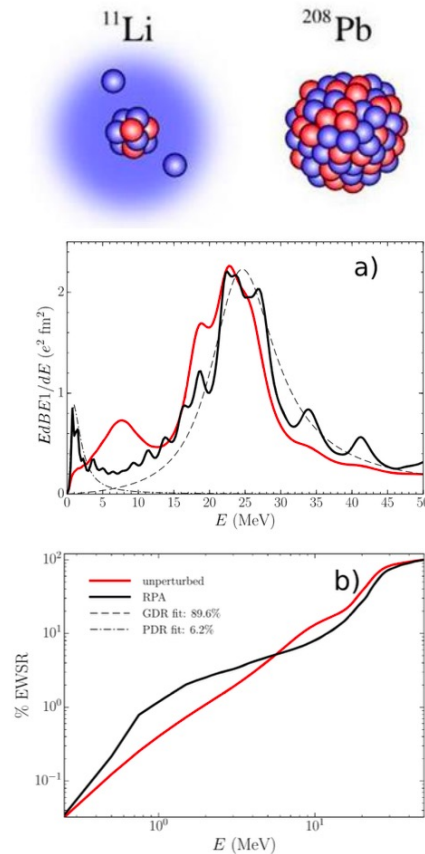
A) H_2 @600Torr

B) 90% H_2 +10% CF_4 @700Torr

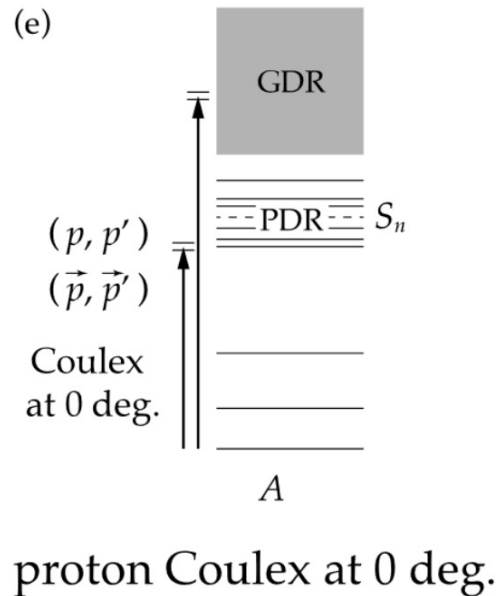
Courtesy: Yassid Ayyad



Electromagnetic response of two-neutron halo nucleus ^{11}Li



Phys. Scr. 94 (2019) 114002



Progress in Particle and Nuclear Physics 106 (2019) 360–433

- Scattered protons: Inclusive measurement relevant to the full photo-absorption cross section containing all the decay channels.
- Broad excitation energy region.
- Polarization transfer observables measured.
- Absolute photo-absorption cross sections extracted from the measured proton scattering cross section.
- Branching ratios to decaying channels.

Tamii and Kobayashi AAPPS Bulletin (2024) 34:7
A. Tamii et al. Phys. Rev. Lett. 107, 062502 (2011)

Increase TPC 's rate capabilities and use rare/hazardous gases

One and two nucleon transfer reactions (**t,p**) and (**³He,d**) and (**³He,p**)

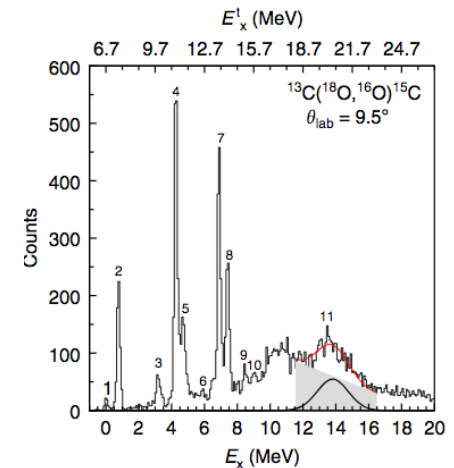
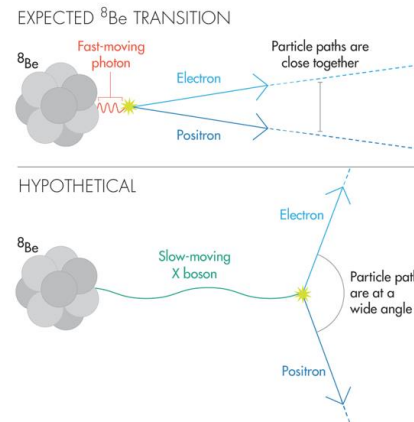
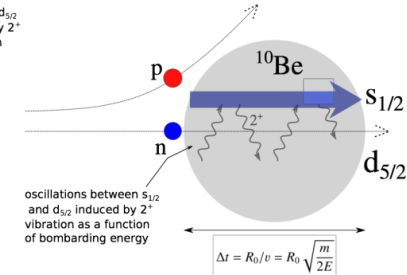
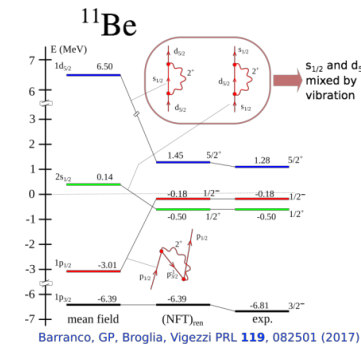
- Isoscalar np pairing
- Giant Pairing Vibrations (GPV)
- Pairing correlations
- Shape co-existence
- Astrophysics

Efficient determination of excitation functions

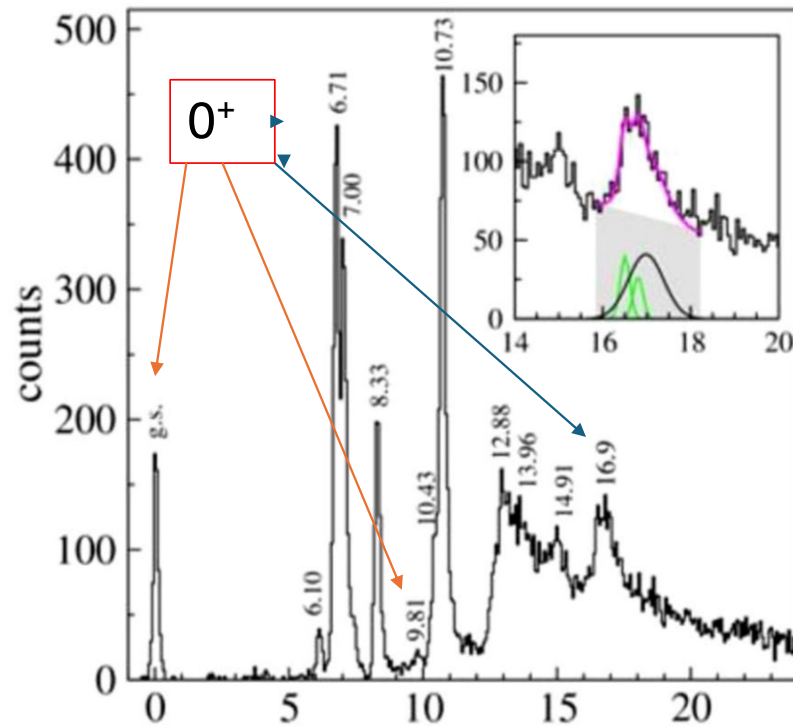
- Resonant scattering
- Transfer reactions as a function of bombarding energy

Charge exchange reactions (t,³He)

Search for exotic BSM standard physics.

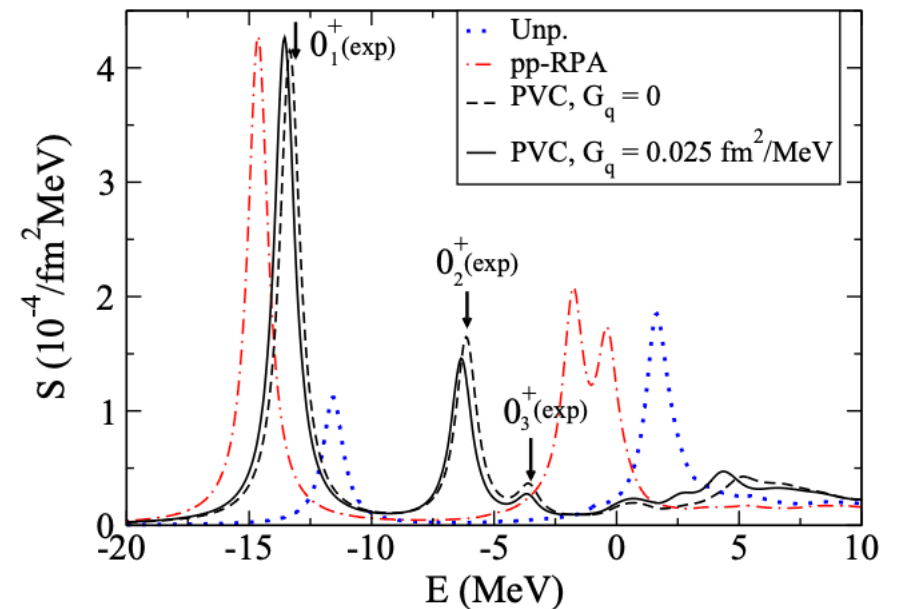


Several unsuccessful experimental searches have been carried out over the years , but recently a bump has been detected at $E^* \approx 16$ MeV in the reaction $^{12}\text{C}(^{18}\text{O}, ^{16}\text{O})^{14}\text{C}$ at $E_{\text{lab}} = 84$ and 275 MeV and interpreted as a signature of GPV



F. Cappuzzello et al., Nat. Comm. 6 (2015) 6743

A new theoretical calculation of the pair strength function (no cross sections yet!)



F. Barranco et al., PRL 134 (2025) 062501

We cannot compute two-neutron transfer cross sections, and for the moment we limit ourselves to pairing strengths

$$S_k(R_{box}) = \left| \sum_{pp'} X(k)_{pp'} \int dr \psi_p(r) \psi_{p'}(r) 1/V_0 dV/dr(r) < j_p || Y_0 || j_{p'} > + \right. \\ \left. \sum_{hh'} Y(k)_{hh'} \int dr \psi_h(r) \psi_{h'}(r) 1/V_0 dV/dr < j_h || Y_0 || j_{h'} > \right|^2$$

$$S(E; R_{box}) = \sum_k S_k(R_{box}) \times \mathcal{L}(E - E_k; W)$$

$$S(E) = \frac{\sum_{i=1, N_{box}} S(E; R_{box, i})}{N_{box}}$$