

Large scale simulations of photosynthetic antenna systems: interplay of cooperativity and disorder

Main author: **Alessia Valzelli**

3rd year PhD student

DINFO, Dept. of Physics and INFN of Florence

Co-authors:

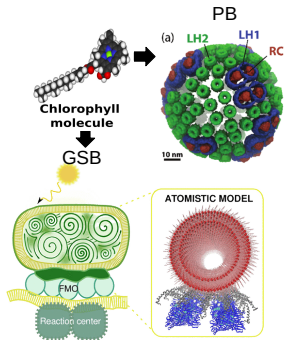
A. Boschetti (INRiM Turin, LENS and U. Florence), F. Mattiotti (U. Strasbourg), A. Kargol (Loyola U. New Orleans), C. Green (Loyola U. New Orleans), F. Borgonovi (UCSC and INFN Milan) and **G. L. Celardo** (U. of Florence, INFN and LENS)

Florence Theory Group Day
March 19th 2025

Simmetry and cooperativity in Biology

QUANTUM BIOLOGY:

1. Role of quantum mechanics in biological systems
2. Applications: PHOTON SENSORS
BIO-INSPIRED SUNLIGHT PUMPED LASER
SOLAR CELLS



Light harvesting complexes:
high degree of symmetry, order and complexity

Symmetry favors the emergence of quantum cooperative properties
robust to noise: **SUPERRADIANCE** and **SUPER-TRANSFER**

How does cooperativity donate robustness?
Fast time scales (faster than thermal relaxation) → high internal efficiency

Related papers:

J. Am. Chem. Soc. 2014, 136, 5, 2048–2057
J. Phys. Chem. Lett. 2012, 3, 4, 536–542
Marco Gullì et al 2019 New J. Phys. 21 013019

MOTIVATIONS

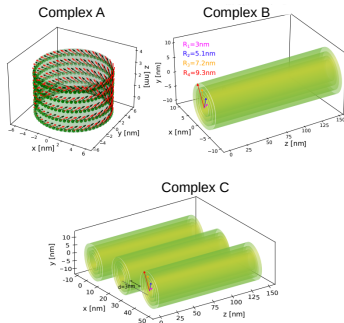
- Relationship between structure and functionality.
- Cooperativity and superradiance in GSB and PB light-harvesting systems.
- Robustness to static and thermal noise.

Ongoing collaborations:

- Study of the EET in GSB and PB antennae.
- Applications: bio-inspired sunlight pumped lasers.

1. Hierarchical organization in GSB and PB complexes

GSB complexes A,B,C



$$L_{max} = 148.57 \text{ nm}$$

$$N_{max} = 132840 \text{ Bchl}c \text{ (complex C)}$$

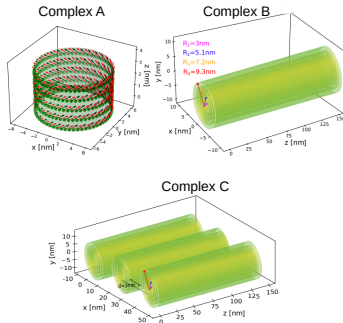
$$e_0 = 15390 \text{ cm}^{-1}$$

$$\lambda_0 \sim 650 \text{ nm}$$

$$|\vec{\mu}| = \sqrt{30} \text{ D}$$

1. Hierarchical organization in GSB and PB complexes

GSB complexes A,B,C



$$L_{max} = 148.57 \text{ nm}$$

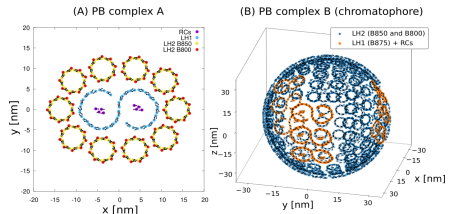
$$N_{max} = 132840 \text{ Bchl } c \text{ (complex C)}$$

$$\epsilon_0 = 15390 \text{ cm}^{-1}$$

$$\lambda_0 \sim 650 \text{ nm}$$

$$|\vec{\mu}| = \sqrt{30} \text{ D}$$

PB complexes A,B



$$N = 334 \text{ Bchl } a$$

$$\bar{\epsilon}_0 = 12500 \text{ cm}^{-1}$$

$$\lambda_0 \sim 800 \text{ nm}$$

$$|\vec{\mu}| = 10.157 \text{ D}$$

$$R = 30 \text{ nm}$$

$$N = 4113 \text{ Bchl } a$$

$$\bar{\epsilon}_0 = 12500 \text{ cm}^{-1}$$

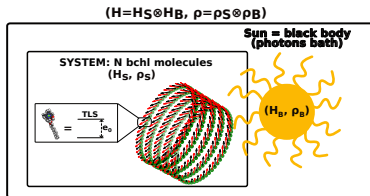
$$\lambda_0 \sim 800 \text{ nm}$$

$$|\vec{\mu}| = 10.157 \text{ D}$$

2. The model: why the Non-Hermitian Hamiltonian approach?

Bchl molecule = TLS ($e_0 = \hbar\omega_0, \vec{\mu}_0, \gamma$)

OPEN QUANTUM SYSTEM



$$\frac{d\rho_S}{dt} = -\frac{i}{\hbar} [H_0 + \Delta, \rho_S] + \frac{1}{\hbar} \sum_{i,j} \left(\sigma_j^- \rho_S \sigma_i^+ - \frac{1}{2} \left\{ \sigma_i^+ \sigma_j^-, \rho_S \right\} \right)$$

- 1 Born-Markov and secular approx. (Lindblad approach);
- 2 EMF: black body at 0 K;
- 3 Sunlight very dilute: single excitation manifold.

Radiative non-Hermitian Hamiltonian (NHH)

$$\hat{H}_{\text{eff}} = \sum_j \left(\omega_0 - i\frac{\gamma}{2} \right) |j\rangle \langle j| + \sum_{\substack{i,j \\ i \neq j}} \left(\Delta_{ij} - \frac{i}{2} \gamma_{ij} \right) |i\rangle \langle j|$$

Spontaneous emission of an ensemble of TLS and interactions mediated by the vacuum fluctuations of the EMF.

Complex eigenvalues:

$$\epsilon_n = E_n - \frac{i}{2} \Gamma_n$$

- $\Gamma_n \gg \gamma$: superradiant eigenstate
- $\Gamma_n \ll \gamma$: subradiant eigenstate

Related paper: Grad J., Hernandez G., Mukamel S., PRA 37, 3835 (1988)

Radiative non-Hermitian Hamiltonian (NHH)

$$\hat{H}_{eff} = \sum_j \left(\omega_0 - i\frac{\gamma}{2} \right) |j\rangle \langle j| + \sum_{\substack{i,j \\ i \neq j}} \left(\Delta_{ij} - \frac{i}{2} \gamma_{ij} \right) |i\rangle \langle j|$$

Spontaneous emission of an ensemble of TLS and interactions mediated by the vacuum fluctuations of the EMF.

Complex eigenvalues:

$$\epsilon_n = E_n - \frac{i}{2} \Gamma_n$$

- $\Gamma_n \gg \gamma$: superradiant eigenstate
- $\Gamma_n \ll \gamma$: subradiant eigenstate

Related paper: Grad J., Hernandez G., Mukamel S., PRA 37, 3835 (1988)

The Hermitian Hamiltonian and the Dipole strength

HH: Hermitian Hamiltonian

Resonance overlap criterion: $\Gamma_n < \delta$

$\delta = \frac{E_{\max} - E_{\min}}{N}$ (average energy mean level spacing)

$$H_H = \sum_{i=1}^N \epsilon_0 |i\rangle \langle i| + \sum_{i \neq j} \Delta_{ij} |i\rangle \langle j|$$

Dipole strength:

$$\vec{D}_n = \sum_{i=1}^N \langle i | E_n \rangle \hat{\mu}_i.$$

where $|E_n\rangle = \sum_{i=1}^N C_{ni} |i\rangle$

If resonances do not overlap: $|\vec{D}_n|^2 \approx \Gamma_n / \gamma$.

The Hermitian Hamiltonian and the Dipole strength

HH: Hermitian Hamiltonian

Resonance overlap criterion: $\Gamma_n < \delta$

$\delta = \frac{E_{max} - E_{min}}{N}$ (average energy mean level spacing)

$$H_H = \sum_{i=1}^N \epsilon_0 |i\rangle \langle i| + \sum_{i \neq j} \Delta_{ij} |i\rangle \langle j|$$

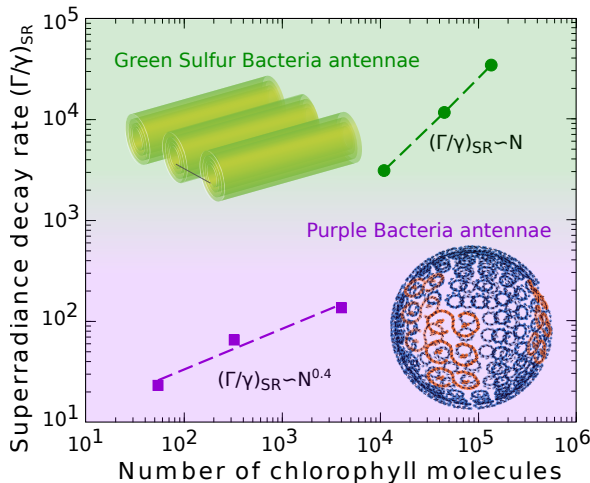
Dipole strength:

$$\vec{D}_n = \sum_{i=1}^N \langle i | E_n \rangle \hat{\mu}_i.$$

where $|E_n\rangle = \sum_{i=1}^N C_{ni} |i\rangle$

If resonances do not overlap: $|\vec{D}_n|^2 \approx \Gamma_n / \gamma$.

3. The emergence of Superradiance in GSB and PB



J. Phys. Chem. B 2024, 128, 9643-9655

4. Robustness to thermal noise and static disorder

THERMAL NOISE

Density matrix of a state at the **canonical equilibrium**:

$$\rho_{ij} = \sum_n \frac{e^{-\beta E_n}}{\text{Tr}(e^{-\beta \hat{H}})} \langle i | E_n \rangle \langle E_n | j \rangle,$$

where $\beta = 1/k_B T$ at room temperature.

The **thermal coherence length** L_ρ is defined by:

$$L_\rho = \frac{1}{N} \frac{\left(\sum_{ij} |\rho_{ij}| \right)^2}{\sum_{ij} |\rho_{ij}|^2}$$

L_ρ measures how much a single excitation is spread coherently over the molecules of the aggregate: $\frac{1}{N} \leq L_\rho \leq N$

4. Robustness to thermal noise and static disorder

THERMAL NOISE

Density matrix of a state at the **canonical equilibrium**:

$$\rho_{ij} = \sum_n \frac{e^{-\beta E_n}}{\text{Tr}(e^{-\beta \hat{H}})} \langle i | E_n \rangle \langle E_n | j \rangle,$$

where $\beta = 1/k_B T$ at room temperature.

The **thermal coherence length** L_ρ is defined by:

$$L_\rho = \frac{1}{N} \frac{\left(\sum_{ij} |\rho_{ij}| \right)^2}{\sum_{ij} |\rho_{ij}|^2}$$

L_ρ measures how much a single excitation is spread coherently over the molecules of the aggregate: $\frac{1}{N} \leq L_\rho \leq N$

Static and thermal noise in GSB

STATIC DISORDER W [cm^{-1}]: strength of the energy fluctuations

$$H_{ij} = (e_0 + \Delta e)|i\rangle\langle i|$$

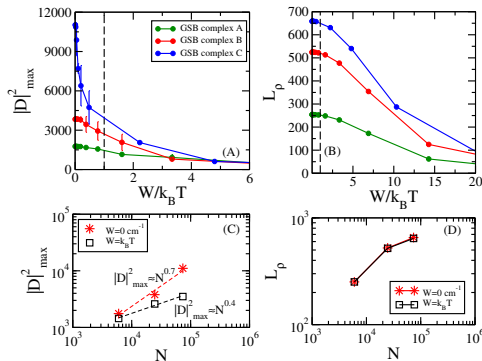
with $-\frac{W}{2} \leq \Delta e \leq \frac{W}{2}$.

Static and thermal noise in GSB

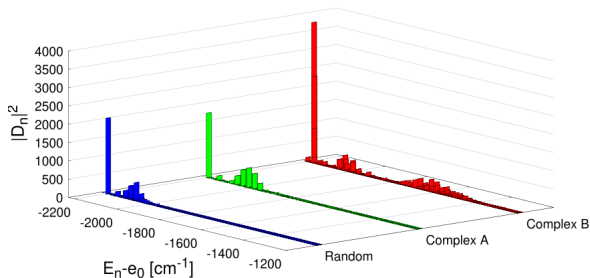
STATIC DISORDER W [cm^{-1}]: strength of the energy fluctuations

$$H_{ij} = (e_0 + \Delta e)|i\rangle\langle i|$$

with $-\frac{W}{2} \leq \Delta e \leq \frac{W}{2}$.

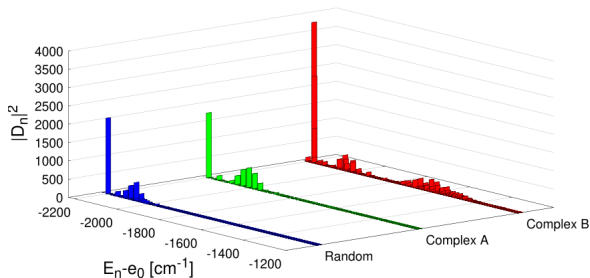


5. Robustness of Superradiance with respect to random dipole orientation in GSB



Increasing the system size by adding the molecules in the same positions but with randomized dipole directions: **no superradiant enhancement**.

5. Robustness of Superradiance with respect to random dipole orientation in GSB



Increasing the system size by adding the molecules in the same positions but with randomized dipole directions: **no superradiant enhancement**.

Large Scale Simulations of Photosynthetic Antenna Systems: Interplay of Cooperativity and Disorder

Alessia Valzelli,[†] Alice Boschetti, Francesco Mattiotti, Armin Kargol, Coleman Green, Fausto Borgonovi,
and G. Luca Celardo

Cite This: *J. Phys. Chem. B* 2024, 128, 9643–9655



New Journal of Physics

The open access journal at the forefront of physics



Published in partnership
with: Deutsche Physikalische
Gesellschaft and the Institute
of Physics



OPEN ACCESS

PAPER

Macroscopic coherence as an emergent property in molecular nanotubes

Marco Gulli[†], Alessia Valzelli[†], Francesco Mattiotti[†], Mattia Angeli[†], Fausto Borgonovi[†] and
Giuseppe Luca Celardo[†]

RECEIVED
30 July 2018

REVISED
18 October 2018



6. Conclusions and outlook

CONCLUSIONS

- 1 **Superradiance** in the whole antenna complexes.
- 2 Cooperative effects robust even with **disorder** and **noise** levels comparable with ambient conditions.
- 3 Cooperative effects strictly related to the **geometry of the system**.

OUTLOOK

- 1 Experimental validation of cooperative response in LHC.
- 2 Artificial devices for light-harvesting and clean energy production.

6. Conclusions and outlook

CONCLUSIONS

- 1 **Superradiance** in the whole antenna complexes.
- 2 Cooperative effects robust even with **disorder** and **noise** levels comparable with ambient conditions.
- 3 Cooperative effects strictly related to the **geometry of the system**.

OUTLOOK

- 1 Experimental validation of cooperative response in LHC.
- 2 Artificial devices for light-harvesting and clean energy production.

6. Conclusions and outlook

CONCLUSIONS

- 1 **Superradiance** in the whole antenna complexes.
- 2 Cooperative effects robust even with **disorder** and **noise** levels comparable with ambient conditions.
- 3 Cooperative effects strictly related to the **geometry of the system**.

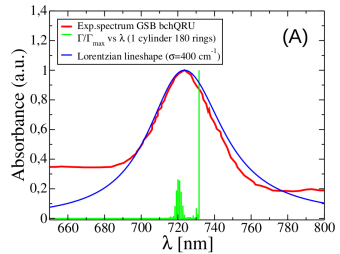
OUTLOOK

- 1 Experimental validation of cooperative response in LHC.
- 2 Artificial devices for light-harvesting and clean energy production.

6. Conclusions and outlook

CONCLUSIONS

- 1 **Superradiance** in the whole antenna complexes.
- 2 Cooperative effects robust even with **disorder** and **noise** levels comparable with ambient conditions.
- 3 Cooperative effects strictly related to the **geometry of the system**.



OUTLOOK

- 1 Experimental validation of cooperative response in LHC.
- 2 Artificial devices for light-harvesting and clean energy production.

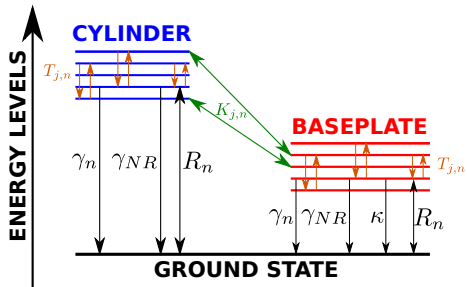
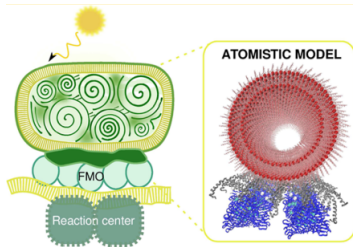
II. ONGOING COLLABORATIONS

**Study of EET in PB and GSB
light-harvesting units**

Bio-inspired sunlight pumped lasers

Optimization of EET in GSB light-harvesting complexes

- ANTENNA SYSTEM (chlorosome-baseplate-FMO trimer-RCs).
- Incoherent rate equations to model EET.
- MC-FRET: transition rate between chlorosome and baseplate
- Solar radiation: black body at Sun temperature $T = 5800$ K ($R_n =$).
- k average transfer rate from the FMO to the RC.



Related papers:

Phys. Chem. Chemical Phys. 18 7459 (2016)

J. Am. Chem. Soc. 2014, 136, 5, 2048–2057

Strong collaboration with DINFO, prof. D. Fanelli and F. Bagnoli

Bio-inspired sunlight pumped lasers: how to turn unconcentrated thermal sunlight into a coherent laser beam with LHC

TOWARDS A BIO-MIMETIC SUNLIGHT PUMPED LASER BASED ON PHOTOSYNTHETIC ANTENNA COMPLEXES



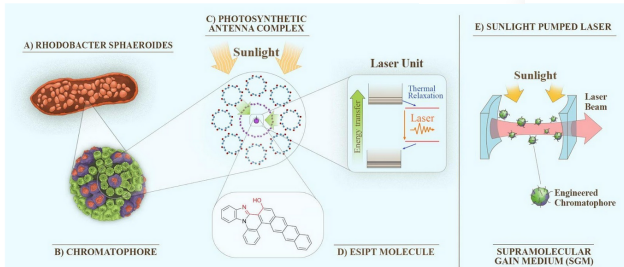
Consiglio Nazionale
delle Ricerche



UNIVERSITÀ
DI PARMA



MAX-PLANCK-GESELLSCHAFT



Related paper: Francesco Mattiotti et al 2021 New J. Phys. 23 103015

THANK YOU

FOR THE

ATTENTION!

Appendix 1: Radiative non-Hermitian Hamiltonian

Bchl molecule = TLS ($e_0 = \hbar\omega_0, \vec{\mu}, \gamma$)

NHH: radiative non-Hermitian-Hamiltonian

$$\hat{H}_{\text{eff}} = \sum_j \left(\omega_0 - i\frac{\gamma}{2} \right) |j\rangle \langle j| + \sum_{\substack{i,j \\ i \neq j}} \left(\Delta_{ij} - \frac{i}{2} \gamma_{ij} \right) |i\rangle \langle j|$$

$$\epsilon_n = E_n - \frac{i}{2} \Gamma_n$$

$$\begin{aligned} \Delta_{ij} = & \frac{3\gamma}{4} \left[\left(-\frac{\cos(k_0 r_{ij})}{(k_0 r_{ij})} + \frac{\sin(k_0 r_{ij})}{(k_0 r_{ij})^2} + \frac{\cos(k_0 r_{ij})}{(k_0 r_{ij})^3} \right) \hat{\mu}_i \cdot \hat{\mu}_j + \right. \\ & \left. - \left(-\frac{\cos(k_0 r_{ij})}{(k_0 r_{ij})} + 3\frac{\sin(k_0 r_{ij})}{(k_0 r_{ij})^2} + 3\frac{\cos(k_0 r_{ij})}{(k_0 r_{ij})^3} \right) (\hat{\mu}_i \cdot \hat{r}_{ij}) (\hat{\mu}_j \cdot \hat{r}_{ij}) \right], \\ \gamma_{ij} = & \frac{3\gamma}{2} \left[\left(\frac{\sin(k_0 r_{ij})}{(k_0 r_{ij})} + \frac{\cos(k_0 r_{ij})}{(k_0 r_{ij})^2} - \frac{\sin(k_0 r_{ij})}{(k_0 r_{ij})^3} \right) \hat{\mu}_i \cdot \hat{\mu}_j + \right. \\ & \left. - \left(\frac{\sin(k_0 r_{ij})}{(k_0 r_{ij})} + 3\frac{\cos(k_0 r_{ij})}{(k_0 r_{ij})^2} - 3\frac{\sin(k_0 r_{ij})}{(k_0 r_{ij})^3} \right) (\hat{\mu}_i \cdot \hat{r}_{ij}) (\hat{\mu}_j \cdot \hat{r}_{ij}) \right]. \end{aligned}$$

Appendix 2: Radiative Hamiltonian and the Dipole Approximation

HH: Hermitian Hamiltonian

Resonance overlap criterion: $\Gamma_n < \delta$

$\delta = \frac{E_{max} - E_{min}}{N}$ (average energy mean level spacing)

$$H_H = \sum_{i=1}^N e_0 |i\rangle \langle i| + \sum_{i \neq j} \Delta_{ij} |i\rangle \langle j|$$

DH: Dipole Hamiltonian (resonance overlap criterion and $k_0 r_{ij} \ll 1$)

$$H_{dip} = \sum_{i=1}^N e_0 |i\rangle \langle i| + \sum_{i \neq j} \frac{\vec{\mu}_i \cdot \vec{\mu}_j - 3(\vec{\mu}_i \cdot \hat{r}_{ij})(\vec{\mu}_j \cdot \hat{r}_{ij})}{r_{ij}^3} |i\rangle \langle j|$$

Appendix 2: Radiative Hamiltonian and the Dipole Approximation

HH: Hermitian Hamiltonian

Resonance overlap criterion: $\Gamma_n < \delta$

$\delta = \frac{E_{max} - E_{min}}{N}$ (average energy mean level spacing)

$$H_H = \sum_{i=1}^N \mathbf{e}_0 |i\rangle \langle i| + \sum_{i \neq j} \Delta_{ij} |i\rangle \langle j|$$

DH: Dipole Hamiltonian (resonance overlap criterion and $k_0 r_{ij} \ll 1$)

$$H_{dip} = \sum_{i=1}^N \mathbf{e}_0 |i\rangle \langle i| + \sum_{i \neq j} \frac{\vec{\mu}_i \cdot \vec{\mu}_j - 3(\vec{\mu}_i \cdot \hat{r}_{ij})(\vec{\mu}_j \cdot \hat{r}_{ij})}{r_{ij}^3} |i\rangle \langle j|$$

Appendix 3: Dipole strength

For HH and DH models we compute the **dipole strength**:

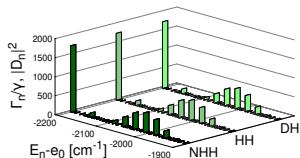
$$\vec{D}_n = \sum_{i=1}^N \langle i | E_n \rangle \hat{\mu}_i.$$

where $|E_n\rangle = \sum_{i=1}^N C_{ni} |i\rangle$

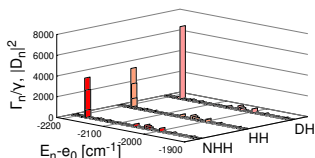
If resonances do not overlap: $|\vec{D}_n|^2 \approx \Gamma_n/\gamma$.

Appendix 4a: Results in GSB

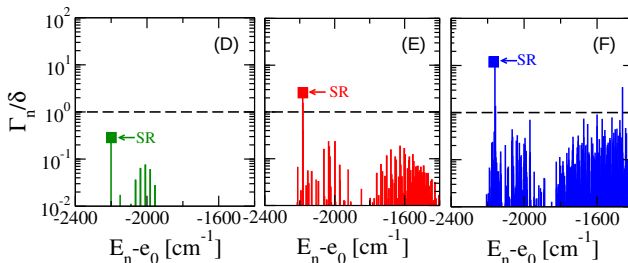
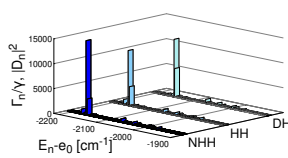
(A) GSB complex A



(B) GSB complex B

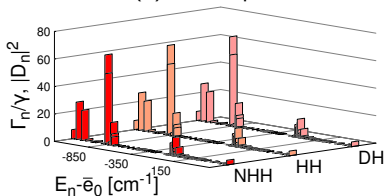


(C) GSB complex C



Appendix 4b: Results in PB

(A) PB complex A



(B) PB complex B

