







# LHCspin: a polarized fixed-target experiment at the LHC



L. L. Pappalardo (pappalardo@fe.infn.it)

#### Sar WorS 2025 – 4° Sardinian Workshop on Spin

- World top energies
- high luminosity
- p and ion beams
- highly sophisticated state-of-the-art detectors



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#### LHCb fixed-target setup

- Since 2015 LHCb can also be operated as a fixed-target experiment with the SMOG system, by injecting low pressure noble gases (He, Ne, Ar) into the VELO vessel.
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#### FT kinem. with $E_p = 6.8$ TeV:

gas (He, Ne, Ar)

•  $\sqrt{s_{NN}} \approx 110 \ GeV$ 

 $\sqrt{s_{NN}} = 72 \; {
m GeV}$ 

- $-3.0 \le y_{CM} \le 0$
- $x_F < 0$
- intermediate-large *x*<sub>B</sub>
- intermediate  $Q^2$



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#### Many interesting published analyses:

- Antiproton prod. cross section in p-He
- Charmonium production in p-Ne and Pb-Ne
- Open charm production in p-Ne and Pb-Ne

#### The SMOG2 upgrade

[SMOG2 TDR]



- > 20 cm storage cell for the target gas installed upstream of the VELO
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- more gas species:  $H_2$ ,  $D_2$ , He,  $N_2$ ,  $O_2$ , Ne, Ar (Kr and Xe to be tested)
- target density increased by large factor (up to 30)
- precise density (luminosity) determination
- negligible impact on LHC and LHCb performance
- can run un parallel with collider mode!



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- LHCb is now the first (unique) LHC experiment with two simultaneous interaction regions!



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- > Large beam-gas samples already collected in 2024 with all available gases!





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#### The LHCspin project

The LHCspin project represents the natural evolution of SMOG2 and will allow for the first time to perform spin physics measurements at the LHC through the implementation of a new-generation polarized gaseous target in the LHCb spectrometer.



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Study multi-dimensional nucleon structure at unique kinematic conditions (backward CM region, poorly explored large-x region at intermediate  $Q^2$ )

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#### **Physics Motivations**

Study multi-dimensional nucleon structure at unique kinematic conditions (backward CM region, poorly explored large-x region at intermediate  $Q^2$ )

#### **Points of strenght**

- ✓ use of well-established polarized gas target technology (HERMES @ DESY, ANKE @ COSY,...)
- ✓ marginal impact on LHC beam lifetime and LHCb mainstream physics program and performances
- ✓ can run in parallel with collider mode (well displaced interaction regions)
- $\checkmark\,$  can benefit from both protons and ion beams
- ✓ allows also injection of non-polarized gases (a-la SMOG2): H<sub>2</sub>, D<sub>2</sub>, He, N<sub>2</sub>, O<sub>2</sub>, Ne, Ar, Kr, Xe
- ✓ broad and unique physics program (next slides)

# Nucleon tomography in momentum space: TMDs



- Describe spin-orbit correlations of the form  $\vec{S} \cdot (\vec{p}_1 \times \vec{p}_2)$
- generate distorsions of the parton densities in transverse momentum plane (e.g. Sivers effect)
- can provide sensitivity to unknown parton OAM!

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#### Quark and gluon TMDs

	quark pol.				
		U	L	Т	
Indiana por	U	$f_1$		$h_1^\perp$	
	L		$g_{1L}$	$h_{1L}^{\perp}$	
	Т	$f_{1T}^{\perp}$	<b>g</b> <sub>1T</sub>	$h_1, h_{1T}^\perp$	

- 8 independent quark TMDs at leading-twist
- significant experimental progress in the last 20 years!
- main results from SIDIS (HERMES, COMPASS, JLAB,  $\rightarrow$  EIC)
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- Similar notation, but important differences!
- different naïve-time-reversal properties
- Experimental access still very limited!

	T-even	T-odd
q	$\mathbf{h}_{1}^{\mathbf{q}}$	$\mathbf{h_1^{\perp q}}$
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#### Polarized hadronic collisions with LHCspin offer a complementary approach!



- Measure experimental observables sensitive to both quarks and gluons TMDs
- Make use of new probes (charmed and beauty mesons)
- Test non-trivial process dependence of quarks and (especially) gluons TMDs

nucleon pol.

	quark pol.				
		U	L	Т	
leon pol.	U	$f_1$		$h_1^\perp$	
	L		$g_{1L}$	$h_{1L}^{\perp}$	
nuc	Т	$f_{1T}^{\perp}$	$g_{1T}$	$h_1, \ egin{smallmatrix} m{h}_{1T}^\perp \ m{h}_{1$	

Unpolarized Drell-Yan



- Theoretically cleanest hard h-h scattering process
- LHCb has excellent  $\mu$ -ID & reconstruction for  $\mu^+\mu^-$
- dominant:  $\overline{q}(x_{beam}) + q(x_{target}) \rightarrow \mu^+ \mu^-$
- beam sea quarks probed at small *x*
- target valence quarks probed at large x



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Expected mass resolution  $\sim 10 \ MeV$  in  $2 < M_{\mu\mu} < 5 \ GeV$  $_{450} = 2 < M_{\mu\mu} < 5 \text{ GeV}$ 13570 Entries Mean -0.01197 Std Dev 12.44 400 E Constant 444.5 0.4011 350 300 250 150 Sigma 10.97 50Ē 40 M<sub>nen</sub>-M<sub>reco</sub> [MeV]



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Sensitive to unpol. and BM TMDs

 $d\sigma_{UU}^{DY} \propto f_1^{\bar{q}} \otimes f_1^q + \cos 2\phi \ h_1^{\perp,\bar{q}} \otimes h_1^{\perp,q}$ 



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Sensitive to quark TMDs through TSSAs

$$A_N^{DY} = \frac{1}{P} \frac{\sigma_{DY}^{\uparrow} - \sigma_{DY}^{\downarrow}}{\sigma_{DY}^{\uparrow} + \sigma_{DY}^{\downarrow}} \implies A_{UT}^{sin\phi_S} \sim \frac{f_1^q \otimes f_{1T}^{\perp q}}{f_1^q \otimes f_1^q}, \quad A_{UT}^{sin(2\phi-\phi_S)} \sim \frac{h_1^{\perp q} \otimes h_1^q}{f_1^q \otimes f_1^q}, \dots$$



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- In DY extraction of qTMDs does not require knowledge of FF
- Verify sign change of Sivers func. wrt SIDIS  $f_{1T}^{\perp}|_{DY} = -f_{1T}^{\perp}|_{SIDIS}$
- Test flavour sensitivity using both H and D targets

#### Gluon TMDs

gluon pol.					
	U	Circularly	Linearly		
U	$f_1^g$		$h_1^{\perp g}$		
L		$g^g_{1L}$	$h_{1L}^{\perp g}$		
Т	$f_{1T}^{\perp g}$	$g_{1T}^g$	$h_1^g,h_{1T}^{\perp g}$		

In high-energy hadron collisions, heavy quarks are dominantly produced through gg fusion:



the most efficient way to access the gluon dynamics inside the proton at LHC is to **measure heavy-quark observables** 

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Polarized gTMDs can be accessed through TSSAs in **inclusive heavy meson production** 

$$A_N = \frac{1}{P} \frac{\sigma^{\uparrow} - \sigma^{\downarrow}}{\sigma^{\uparrow} + \sigma^{\downarrow}} \propto \left[ f_{1T}^{\perp g}(x_a, k_{\perp a}) \otimes f_g(x_b, k_{\perp b}) \right] \sin \phi_S + \cdots$$



c, b

 $\overline{c}, \overline{b}$ 

g

g

## Gluon TMDs



**Gluon Sivers function:** 

- Sheds light on spin-orbit correlations of unpol. gluons inside a transv. pol. proton
- is sensitive to gluon OAM

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c, b

 $\overline{c}, \overline{b}$ 

 $\eta_c, \chi_{c0}, \dots$ 

 $\eta_{b}, \chi_{b0}, ...$  $J/\psi, \psi', ...$ 

Y(nS)

TMD factorization requires  $q_T(Q) \ll M_Q$ . Can look at **associate quarkonia production**, where only the relative  $q_T$  needs to be small, e.g.:  $pp^{(\uparrow)} \rightarrow J/\psi + J/\psi + X$ 

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$$\frac{\mathrm{d}\sigma}{\mathrm{d}M_{QQ}\mathrm{d}Y_{QQ}\mathrm{d}^{2}P_{QQr}\mathrm{d}\Omega} = \frac{\sqrt{M_{QQ}^{2} - 4M_{Q}^{2}}}{(2\pi)^{28s}M_{QQ}^{2}} \left\{ F_{1}(M_{QQ},\theta_{\mathrm{CS}}) \mathcal{C}\left[f_{1}^{g}f_{1}^{g}\right](x_{1,2}, P_{QQr}) + F_{2}(M_{QQ},\theta_{\mathrm{CS}}) \mathcal{C}\left[w_{2}h_{1}^{\perp g}h_{1}^{\perp g}\right](x_{1,2}, P_{QQr}) \right\} \\ + \left(F_{3}(M_{QQ},\theta_{\mathrm{CS}}) \mathcal{C}\left[w_{3}f_{1}^{g}h_{1}^{\perp g}\right](x_{1,2}, P_{QQr}) + F_{3}'(M_{QQ},\theta_{\mathrm{CS}}) \mathcal{C}\left[w_{3}'h_{1}^{\perp g}f_{1}^{g}\right](x_{1,2}, P_{QQr})\right) \cos 2\phi_{\mathrm{CS}} + F_{4}(M_{QQ},\theta_{\mathrm{CS}}) \mathcal{C}\left[w_{4}h_{1}^{\perp g}h_{1}^{\perp g}\right](x_{1,2}, P_{QQr}) \cos 4\phi_{\mathrm{CS}}\right\} \\ \frac{\mathrm{gluon pol.}}{\mathrm{gluon pol.}}$$

nucleon pol.

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$f_{1T}^{\perp g}$	$g_{1T}^g$	$h_1^g,h_{1T}^{\perp g}$

L

Т

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$$F_1(M_{\mathcal{Q}\mathcal{Q}},\theta_{\mathrm{CS}}) \, \mathcal{C}\Big[f_1^g f_1^g\Big](x_{1,2}, \boldsymbol{P}_{\mathcal{Q}\mathcal{Q}^T}) + F_2(M_{\mathcal{Q}\mathcal{Q}},\theta_{\mathrm{CS}}) \, \mathcal{C}\Big[w_2 h_1^{\perp g} h_1^{\perp g}\Big](x_{1,2}, \boldsymbol{P}_{\mathcal{Q}\mathcal{Q}^T})$$

 $+ \left( F_3(M_{\mathcal{Q}\mathcal{Q}},\theta_{\mathrm{CS}}) \mathcal{C} \left[ w_3 f_1^g h_1^{\perp g} \right](x_{1,2}, \boldsymbol{P}_{\mathcal{Q}\mathcal{Q}_T}) + F_3'(M_{\mathcal{Q}\mathcal{Q}},\theta_{\mathrm{CS}}) \mathcal{C} \left[ w_3' h_1^{\perp g} f_1^g \right](x_{1,2}, \boldsymbol{P}_{\mathcal{Q}\mathcal{Q}_T}) \right) \cos 2\phi_{\mathrm{CS}} + F_4(M_{\mathcal{Q}\mathcal{Q}},\theta_{\mathrm{CS}}) \mathcal{C} \left[ w_4 h_1^{\perp g} h_1^{\perp g} \right](x_{1,2}, \boldsymbol{P}_{\mathcal{Q}\mathcal{Q}_T}) \cos 4\phi_{\mathrm{CS}} \right\}$ 

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 $\langle \cos 2\phi_{\rm CS} \rangle = -0.029 \pm 0.050 \text{ (stat)} \pm 0.009 \text{ (syst)}$  $\langle \cos 4\phi_{\rm CS} \rangle = -0.087 \pm 0.052 \text{ (stat)} \pm 0.013 \text{ (syst)}$ 

- azimuthal amplitudes consistent with zero
- a few-% asymmetry cannot be excluded
- uncertainties statistically dominated
- But very challenging at fixed-target kinematics

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 $s 4\phi_{\rm C}$ 

# GPDs: a complementary approach to the nucleon tomography



Courtesy QuantOm Collaboration



quark pol.ULTU $f_1$  $h_1^{\perp}$ U $f_1$  $h_1^{\perp}$ L $g_{1L}$  $h_{1L}^{\perp}$ T $f_{1T}^{\perp}$  $g_{1T}$ h\_1,  $h_{1T}^{\perp}$ 

#### Gluon GPDs and UPC

GPD	U	L	T
U	H		$\mathcal{E}_T$
		$\tilde{H}$	$ ilde{E}_T$
T	E	$ ilde{E}$	$H_T, \  ilde{H}_T$

Gluon GPDs can be accessed at LHC in **Ultra-Peripheral collisions (UPC)** where a quasi-real photon is emitted by the relativistic beam particle [PRD 85 (2012), 051502]

At LHC energies, these photons are energetic enough to trigger the production of hard dileptons and charmonia and bottomonia.


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Gluon GPDs can be accessed at LHC in **Ultra-Peripheral collisions (UPC)** where a quasi-real photon is emitted by the relativistic beam particle [PRD 85 (2012), 051502]

At LHC energies, these photons are energetic enough to trigger the production of hard dileptons and charmonia and bottomonia.



L. L. Pappalardo

#### Gluon GPDs and UPC

GPD	U	L	T
U	H		$\mathcal{E}_T$
L		$\tilde{H}$	$ ilde{E}_T$
T	E	$ ilde{E}$	$H_T, \  ilde{H}_T$

With LHCspin exclusive photo-production of  $J/\psi$  in UPC of proton (or lead) beams with polarized  $H^{\uparrow}$  target can be studied, providing constraints to the essentially unknown gluon GPD  $E_g$  which plays a crucial role in the Ji sum rule:

$$J^{g} = \frac{1}{2} \int_{0}^{1} dx \Big( H^{g}(x,\xi,0) + E^{g}(x,\xi,0) \Big)$$



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$$A_{N} = \frac{\sigma^{h_{A}h_{B}^{\downarrow}} - \sigma^{h_{A}h_{B}^{\uparrow}}}{\sigma^{h_{A}h_{B}^{\downarrow}} + \sigma^{h_{A}h_{B}^{\uparrow}}} = \frac{\int dk \frac{dn_{A}}{dk} A_{N}^{\gamma} \sigma^{\gamma h_{B}}}{\int dk \left[\frac{dn_{A}}{dk} \sigma^{\gamma h_{B}} + \frac{dn_{B}}{dk} \sigma^{\gamma h_{A}}\right]}$$

The hadronic STSA  $A_N$  can be parametrized in terms of the photonic STSA  $A_N^{\gamma}$  which incorporates the GPDs  $H^g$  and  $E^g$  through their gluonic CFFs  $\mathcal{H}^g$  and  $\mathcal{E}^g$ 



- Extraction based on models for the GPD H<sup>g</sup> (Goloskokov-Kroll) and E<sup>g</sup> (PRD 85, 051502 (2012))
- AFTER model-dependent predictions
   very promising for pH<sup>↑</sup> UPC

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- probe collective phenomena in heavy-light systems through ultrarelativistic collisions of heavy nuclei with trasv. pol. deuterons
- polarized light target nuclei offer a unique opportunity to control the orientation of the formed fireball by measuring the elliptic flow relative to the polarization axis (ellipticity).



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Unpol. deuterons: the fireball is azimuthally symmetric  $\rightarrow v_2 \approx 0$ .

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- Compact dipole magnet for static transverse field to maintain polarization inside the cell
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- Need to modify main flange of VELO vessel (inward)
- No need for additional detectors!
- Possibility to switch from dipole magnet to solenoid to realize a Longitudinal polarized target in a future phase

#### Kinematic coverage

$$\sqrt{s} = \sqrt{2m_N E_p} = 115 \text{ GeV}$$

$$x_F = 2 E_T / \sqrt{s} \sinh(y^*) \quad E_T = \sqrt{M^2 + p_T^2}$$

$$Q^2 = p_T^2 + m^2 \quad x = Q e^{-y_{CM}/\sqrt{s}}$$

The kinematic coverage **depends on the cell position**:

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10-1

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- Probable position of LHCspin cell: [-670,-470] mm



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 $J/\Psi \rightarrow \mu^+\mu^-PV X$  track reconstruction efficiency



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#### Target

- $I_0 = 6.5 \cdot 10^{16} s^{-1}$  (HERMES)  $C_{tot} = 17.4$  l/s (20 cm cell)  $\theta = 3.7 \cdot 10^{13}$  atoms/cm<sup>2</sup>

#### Beam (Run4)

- $2.2 \cdot 10^{11}$  p/bunch •
- 2760 bunches •
- $I_{beam} = 6.8 \cdot 10^{18} \ p/s$

$$\mathcal{L}_{pH} \approx 2.5 \cdot 10^{32} \,\mathrm{cm}^{-2} \mathrm{s}^{-1}$$
$$L_{pH}(Run) \approx 5 \,f \,b^{-1}$$

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Channel	Events / week	Total yield
$J/\psi  o \mu^+\mu^-$	$2.6  imes 10^7$	$3.1 \times 10^9$
$D^0 \to K^- \pi^+$	$1.3  imes 10^8$	$1.6  imes 10^{10}$
$\psi(2S) \rightarrow \mu^+ \mu^-$	$4.6  imes 10^5$	$5.5  imes 10^7$
$J/\psi J/\psi \to \mu^+ \mu^- \mu^+ \mu^-$ (DPS)	$1.7 imes10^1$	$2.1  imes 10^3$
$J/\psi J/\psi \to \mu^+ \mu^- \mu^+ \mu^-$ (SPS)	$5.1  imes 10^1$	$6.1  imes 10^3$
Drell Yan (5 $< M_{\mu\mu} < 9 \text{ GeV}$ )	$1.5  imes 10^4$	$1.8  imes 10^6$
$\Upsilon  o \mu^+ \mu^-$	$1.1  imes 10^4$	$1.3  imes 10^6$
$\Lambda_c^+ \to p K^- \pi^+$	$2.6  imes 10^6$	$3.1 \times 10^8$

fully reconstructed and selected events (after cuts)!

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Will be much higher in Run5 (HL-LHC)!!

### The jet target option

Alternative solution with **jet target** also under evaluation:

- lower density (~ $10^{12}$  atoms/ $cm^2$ )  $\rightarrow$  about a factor of 40 smaller
- higher polarization (up to 90%)
- lower systematics in P measurement (virtually close to 0)
- Compatible with SMOG2 setup



Necessary pre-requisites for approval of the project at LHCb (Run5)

- R&D campaign for the apparatus towards the final setup for LHCb
- feasibility studies in a dedicated exp. area served by LHC beams
- Develop and test a RHIC-like absolute polarimeter



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- Develop and test a RHIC-like absolute polarimeter

#### **IR4** is the ideal place for our R&D:

- Lots of free space for our instrumentation
- Rails, cables and racks already available in-situ



ABS and BR-polarimeter have already been moved to Ferrara for first tests and optimization prior to full installation at IR4.





LS3 (2027-30):

- Installation at IR4 of existing setup (ABS + BR polarimeter)
- Implementation of a RHIC-like absolute polarimeter: exploits Coulomb-Nuclear Interference in beam-target scattering to measure of molecular polarization (to be calibrated with the BRP)
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- In-beam polarimetry studies
- proof-of-principle prototype experiment: first pol. meas. at the LHC (SSAs in inclusive light hadron production)

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Apr 2025

22

[hep-ex]

arXiv:2504.16034v1

April 23, 2025

#### LHCspin: a Polarized Gas Target for LHC

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https://arxiv.org/abs/2504.16034



Sar Wors 2025 - 4th Sardinian Workshop on Spin - Pula (Cagliari) 11-13 June 2025









#### Conclusions

- > The Fixed-Target program at LHCb is active since Run 2, now greatly enriched with SMOG2
- LHCspin is the natural evolution: a polarized fixed target at LHCb will bring spin-physics for the first time at the LHC and will open the way to a broad and unique physics program!
- > Novel approaches and reactions will be exploited to study the 3D nucleon structure
- First insights into the yet unknown gluon TMDs (such as the GSF) will be possible thanks to the excellent capabilities of LHCb in reconstructing quarkonia states and heavy mesons.
- The approval process of the first phase (R&D at IR4) is in progress.



#### Absolute polarimetry

Under **Coulomb-Nuclear Interference** (CNI) conditions one can measure L-R asymmetries in elastic  $pH^{\uparrow}$  scattering:

RHIC with p-C scattering

$$A_N(t) = \frac{\mu_p - 1}{m_p} \sqrt{-3t_e} \frac{(t/t_e)^{3/2}}{3(t/t_e)^2 + 1} \qquad t_e = -\frac{8\pi\sqrt{3}}{\sigma_e}$$





Analyzing power is maximal (4-5%) for  $t = t_e$ 

#### For a 7 TeV proton beam:

- $\sigma_{tot} \approx 47 \ mb$
- proton recoil energies: 1.7 4.6 MeV
- proton recoil angles  $87^{\circ} < \theta_{lab} < 89^{\circ}$
- Find more here: <u>https://www.maths.tcd.ie/~nhb/talks/2019\_07\_16\_nhb.pdf</u>

#### Physics measurements at IR4

Even though the focus will be on polarimetry and beam interactions, we performed preliminary calculations to determine if a simple detector could meet our needs



we can achieve a resolution  $\delta p/p < 1\%$  within a few meters of lever arm (depending on space constraints) for momenta up to a few GeV and with N = 10 hit measurements

with  $\delta p/p \sim 1\%$  we have  $\delta m \sim 40$  MeV, excellent for any other measurement

it is even possible to have a ToF PID @ $3\sigma$  level for  $\pi - K$  $p \sim 1 \ GeV \rightarrow \sigma_T \mathcal{O}(100) \ ps$ 



# A synergic attack to gTMDs

#### [D. Boer: Few-body Systems 58, 32 (2017)]

	DIS	DY	SIDIS	$pA \to \gamma \operatorname{jet} X$	$e p \to e' Q \overline{Q} X$ $e p \to e' j_1 j_2 X$	$pp \to \eta_{c,b} X$ $pp \to H X$	$\begin{array}{c} pp \rightarrow J/\psi \ \gamma \ X \\ pp \rightarrow \Upsilon \ \gamma \ X \end{array}$
$f_1^{g[+,+]}$ (WW)	×	×	×	×	$\checkmark$	$\checkmark$	$\checkmark$
$f_1^{g[+,-]}$ (DP)	$\checkmark$	$\checkmark$	$\checkmark$	$\checkmark$	×	Х	×

Can be measured at the EIC



Can be measured at RHIC & LHC (including LHCb+SMOG2)

	$pp \to \gamma \gamma X$	$pA \to \gamma^* \text{ jet } X$	$e p \to e' Q \overline{Q} X$ $e p \to e' j_1 j_2 X$	$pp \to \eta_{c,b} X$ $pp \to H X$	$\begin{array}{c} pp \to J/\psi \gamma X \\ pp \to \Upsilon \gamma X \end{array}$
$h_1^{\perp g  [+,+]}  (WW)$	$\checkmark$	×	$\checkmark$	$\checkmark$	$\checkmark$
$h_1^{\perp g [+,-]}$ (DP)	×	$\checkmark$	×	×	×

	DY	SIDIS	$p^{\uparrow} A \to h X$	$p^{\uparrow}A \to \gamma^{(*)} \text{ jet } X$	$ \begin{array}{c} p^{\uparrow}p \rightarrow \gamma  \gamma  X \\ p^{\uparrow}p \rightarrow J/\psi  \gamma  X \\ p^{\uparrow}p \rightarrow J/\psi  J/\psi  X \end{array} $	$e p^{\uparrow} \rightarrow e' Q \overline{Q} X$ $e p^{\uparrow} \rightarrow e' j_1 j_2 X$
$f_{1T}^{\perp g  [+,+]}  (WW)$	×	×	×	×	$\checkmark$	$\checkmark$
$f_{1T}^{\perp g [+,-]}$ (DP)	$\checkmark$	$\checkmark$	$\checkmark$	$\checkmark$	×	×

Can be measured at RHIC and LHCb+LHCspin
# gluon TMDs



- Depending on their combinations, there are 2 independent versions of each gTMD that can be probed in different processes and can have different magnitude and widths and different x and k<sub>T</sub> dependencies!
- E.g. there are 2 types of  $f_1^g$  and  $h_1^{\perp g}$ : [++] = [--] Weizsacker-Williams (WW) ; [+-] = [-+] DiPole (DP)
- 2 indep. GSF:  $f_{1T}^{\perp g[+,+]}$  "f-type"  $\rightarrow$  antisymm. colour structure ;  $f_{1T}^{\perp g[+,-]}$  "d-type"  $\rightarrow$  symm. colour structure

### A preliminary analysis tool for pseudo-data

A pseudo-data set based on a Transversely Pol. H target has been generated to study the interplay between statistical and systematic uncertainties (due to the measurement of the polarization).

Similar approach used at HERMES (Appendix C of [JHEP, 12:010, 2020]):

- Use official LHCb MC data for inclusive production of  $J/\psi \rightarrow \mu^+\mu^-$  in fixed-target configuration (PYTHIA8 + EPOS)
- Introduce a spin-dependence in the simulation: assign to each simulated event a target polarization state (↑ or ↓) using a random extraction modulated with a model for the cross section
- The model assumes a dominant sin φ modulation (e.g. sensitive to the gluon Sivers) plus a suppressed sin 2φ modulation (to account e.g. for possible higher-twist contributions). Both terms depend mildly on the kinematics (x, p<sub>T</sub>):

$$p = \frac{1}{2} \left[ 1 + \left( a_1 + a_2 \frac{x - \overline{x}}{x_{max}} + a_3 \frac{p_T - \overline{p_T}}{p_{T \ max}} \right) \sin \phi + \left( b_1 + b_2 \frac{x - \overline{x}}{x_{max}} + b_3 \frac{p_T - \overline{p_T}}{p_{T \ max}} \right) \sin 2\phi \right]$$

• Using these pseudo-data the TSSA is computed in the usual way:

$$A_N = \frac{1}{P} \frac{N^{\uparrow} - N^{\downarrow}}{N^{\uparrow} + N^{\downarrow}}$$

and the uncertainties on  $N^{\uparrow(\downarrow)}$  (Poisson) and P (systematic) propagated accordingly.

## A preliminary analysis tool for pseudo-data

• The data points are binned in  $x_F$  and  $p_T$  (2D binning), represented vs.  $\phi$  and fitted with  $f = a_1 \sin \phi + a_2 \sin 2\phi$  where the free parameters  $a_1$  and  $a_2$  represent the amplitude of the corresponding azimuthal modulation



- The extracted parameters  $a_1$  and  $a_2$  are consistent with those used to generate the model (no bias is observed)
- With the available MC statistics (corresponding to 2 weeks of data-taking) there is no sensitivity for the  $\sin 2\phi$  term
- The amplitudes  $a_1$  are reported vs.  $x_F$  in bins of  $p_T$  (and vice-versa)
- A mild kinematic dependence is observed consistent with the model

#### Statistical vs Systematics uncertainties

• The analysis tool described above allows to study the interplay between statistical uncertainties and systematic uncertainties (due to the measurement of the polarization) under different data-taking scenarios

$p_T ~({ m MeV})$	$x_F$	$a_1 \ (\Delta P = 0\%)$	$a_1 \ (\Delta P = 5\%)$	$a_1 \ (\Delta P = 20\%)$	$a_1 \ (\Delta P = 50\%)$
[0, 1500]	[-0.70, -0.09]	$0.090\pm0.013$	$0.089 \pm 0.013$	$0.087 \pm 0.014$	$0.087 \pm 0.022$
[0, 1500]	[-0.09, -0.06]	$0.104 \pm 0.011$	$0.104 \pm 0.012$	$0.103 \pm 0.016$	$0.100 \pm 0.027$
[0, 1500]	[-0.06, -0.04]	$0.098 \pm 0.012$	$0.098 \pm 0.013$	$0.097 \pm 0.016$	$0.094 \pm 0.027$
[0, 1500]	[-0.04, 0.05]	$0.118 \pm 0.014$	$0.117 \pm 0.014$	$0.114\pm0.017$	$0.113 \pm 0.030$
$[1500,\!6000]$	[-0.70, -0.09]	$0.093 \pm 0.010$	$0.092\pm0.010$	$0.090\pm0.013$	$0.089 \pm 0.023$
$[1500,\!6000]$	[-0.09, -0.06]	$0.108 \pm 0.011$	$0.108 \pm 0.011$	$0.108 \pm 0.015$	$0.107 \pm 0.027$
$[1500,\!6000]$	[-0.06, -0.04]	$0.105\pm0.012$	$0.105\pm0.012$	$0.104\pm0.015$	$0.103 \pm 0.026$
$[1500,\!6000]$	[-0.04, 0.05]	$0.105\pm0.011$	$0.105\pm0.012$	$0.102\pm0.015$	$0.102\pm0.026$

- A 5% systematic uncertainty on P has no impact on the total uncertainty on  $a_1$
- For  $\Delta P = 20\%$  the systematic uncertainty amounts to 30-40% of the statistical uncertainty
- For  $\Delta P = 50\%$  the systematic uncertainty approximately equals the statistical uncertainty
- We expect  $\Delta P \approx 10-15\%$  for the storage cell hypothesis (and close to 0 for the jet target hypothesis )

#### The LHCb detector

- ► LHCb is a general-purpose single-arm spectrometer, fully instrumented in  $2 < \eta < 5$  and optimised for detection of charmed and beauty hadrons [JINST 3 (2008) S08005] [IJMPA 30 (2015)1530022]
- ➤ Excellent particle identification and momentum resolution:  $\sigma_p/p \le 1.0$  % ( $p \in [2,200]$  GeV)
- Precise primary and secondary vertex reconstruction (VELO)



# Types of collisions (Collider mode):protonsprotonslead ionslead ions $\overbrace{\sqrt{s}} = 13.6 \text{ TeV}$ $\overbrace{\sqrt{s_{NN}}} = 8.2 \text{ TeV}$ lead ions



#### More physics reach with unpolarized FT reactions

- Intrinsic heavy-quark [S.J. Brodsky et al., Adv. High Energy Phys. 2015 (2015) 231547]
  - 5-quark Fock state of the proton may contribute at high x!
  - charm PDFs at large x could be larger than obtained from conventional fits
- pA collisions (using unpolarized gas: He, N, Ne, Ar, Kr, Xe)
  - constraints on nPDFs (e.g. on poorly understood gluon antishadowing at high x)
  - studies of parton energy-loss and absorption phenomena in the cold medium
  - reactions of interest for cosmic-ray physics and DM searches
- PbA collisions at √s<sub>NN</sub> ≈ 72 GeV (using unpolarized gas: He, N, Ne, Ar, Kr, Xe)
   Study of QGP formation (search for predicted sequential quarkonium suppression)







 $c\overline{c}$  states:  $J/\psi$ ,  $\chi_c$ ,  $\psi'$ ,... Different binding energies, different dissociation temperatures  $\rightarrow$  **medium thermometer** 

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