# **Description of the full** $q_T$ **spectrum** of low Q Drell Yan production

**IV Sardinian Workshop on Spin** 

In collaboration with S. Camarda and G. Ferrera



Istituto Nazionale di Fisica Nucleare

#### Lorenzo Rossi

June 12th

UNIVERSITÀ **DEGLI STUDI DI MILANO** 







#### Drell Yan production $h_1(p_1) + h_2(p_2) \rightarrow V + X \rightarrow \ell_1 + \ell_2 + X$





#### **Drell Yan production** $h_1(p_1) + h_2(p_2) \rightarrow V + X \rightarrow \ell_1 + \ell_2 + X$







#### **Drell Yan production** $h_1(p_1) + h_2(p_2) \rightarrow V + X \rightarrow \ell_1 + \ell_2 + X$

 $\frac{d\sigma}{d^2\mathbf{q_T} dM^2 dy d\Omega} = \sum_{a,b} \int_0^1 dx_1 \int_0^1 dx_2 f_{a,h_1}(x_1,\mu_F^2) f_{b,h_2}(x_2,\mu_F^2) \frac{d\hat{\sigma}_{ab}}{d^2\mathbf{q_T} dM^2 d\hat{y} d\Omega} (\alpha_S(\mu_R^2),\mu_R^2,\mu_F^2)$ 

























#### When $q_T \ll M$ :







**Short Recap** 

When  $q_T \ll M$ :  $\int_{-}^{q_T^2} d\bar{q}_T^2 \frac{d\hat{\sigma}_{q\bar{q}}}{d\bar{q}_T^2} \sim 1 + \alpha_S [c_{12}L_{q_T}^2 + c_{11}L_{q_T} + \dots]$  $+\alpha_{S}^{2}[c_{24}L_{q_{T}}^{4}+\ldots+c_{21}L_{q_{T}}+\ldots]+\mathcal{O}(\alpha_{S}^{3})$ whit  $\alpha_S^n L_{q_T}^m = \alpha_S^n \log^m \left(\frac{M^2}{a_T^2}\right) \gg 1$ ΙI







JHEP 05 (2001) 025, Phys.Lett.B 564 (2003) 65-72, Nucl.Phys.B 737 (2006) 73-120

# $q_T$ resummation in QCD



# $q_T$ resummation in QCD $\frac{d\hat{\sigma}}{d^2\mathbf{q}_{\mathrm{T}}} = \frac{d\hat{\sigma}^{(1)}}{d^2}$

JHEP 05 (2001) 025, Phys.Lett.B 564 (2003) 65-72, Nucl.Phys.B 737 (2006) 73-120

$$d\hat{\sigma}^{(\text{res})} + \frac{d\hat{\sigma}^{(\text{fin})}}{d^2 \mathbf{q_T}}$$



In impact parameter space:

 $\frac{d\hat{\sigma}^{(\text{res})}}{d^2\mathbf{q_T}} = \frac{M^2}{\hat{s}} \int \frac{d^2\mathbf{b_T}}{4\pi} e^{i\mathbf{b_T}\cdot\mathbf{q_T}} W(b, M)$ 

JHEP 05 (2001) 025, Phys.Lett.B 564 (2003) 65-72, Nucl.Phys.B 737 (2006) 73-120

### $q_T$ resummation in QCD







In impact parameter space:



In Mellin space:

JHEP 05 (2001) 025, Phys.Lett.B 564 (2003) 65-72, Nucl.Phys.B 737 (2006) 73-120

 $\frac{d\hat{\sigma}^{(\text{res})}}{d^2\mathbf{q_T}} = \frac{M^2}{\hat{s}} \int \frac{d^2\mathbf{b_T}}{4\pi} e^{i\mathbf{b_T}\cdot\mathbf{q_T}} W(b, M)$ 

#### $W_{\mathcal{N}}(b, M) = \hat{\sigma}^{(0)} H_{\mathcal{N}}(\alpha_{S}) \times \exp\{G_{\mathcal{N}}(\alpha_{S}, L)\}$



Nucl.Phys.B 261 (1985) 104-142



Nucl.Phys.B 261 (1985) 104-142



5



5



5







Eur.Phys.J.C 80 (2020) 3, 251





Eur.Phys.J.C 80 (2020) 3, 251



# Study of the full $q_T$ spectrum of high energy data







Eur.Phys.J.C 80 (2020) 3, 251



Study of the full  $q_T$  spectrum of high energy data



Study of the Z and  $W^{\pm}$  boson masses







Eur.Phys.J.C 80 (2020) 3, 251



Study of the full  $q_T$  spectrum of high energy data



Study of the Z and  $W^{\pm}$  boson masses

#### And at low invariant mass?









Only the resummed part...



#### Only the resummed part...



 $\equiv$  README.md

Nanga Parbat is a fitting framework aimed at the determination of the non-perturbative component of TMD distributions.

#### Download

You can obtain NangaParbat directly from the github repository:

https://github.com/MapCollaboration/NangaParbat

For the last development branch you can clone the master code:

git clone git@github.com:MapCollaboration/NangaParbat.git



#### Only the resummed part...



 $\equiv$  README.md

Nanga Parbat is a fitting framework aimed at the determination of the non-perturbative component of TMD distributions.

#### Download

You can obtain NangaParbat directly from the github repository:

https://github.com/MapCollaboration/NangaParbat

For the last development branch you can clone the master code:

git clone git@github.com:MapCollaboration/NangaParbat.git

#### arTeMiDe









<u>Recent version/release can be found in repository</u>

#### Articles, presentations & supplementary materials



Extra pictures for the paper arXiv:1902.08474

Seminar of A.Vladimirov in Pavia 2018 on TMD evolution.

Link to the text in Inspire.

Archive of older links/news.

#### About us & Contacts



If you have found mistakes, or have suggestions/questions, please, contact us.

Some extra materials can be found on <u>Alexey's web-page</u>

Alexey Vladimirov Alexey.Vladimirov@physik.uni-regensburg.de

Ignazio Scimemi ignazios@fis.ucm.es





Or only the finite part...



#### Or only the finite part...

**PHENIX**  $\sqrt{s} = 200 \text{GeV}$ 4.8 GeV < Q < 8.2 GeV 1.2 < |y| < 2.2Central scale:  $\mu_R = \mu_F = E_T^{\gamma}$ 10<sup>1</sup> 100 NNLOJET at LO  $10^{-1}$ NNLOJET at NLO NNLOJET at NNLO  $10^{-2}$ PHENIX 12 0  $p_T$  (GeV)

Physics Letters B 829 (2022) 137111



#### Or only the finite part...

**PHENIX**  $\sqrt{s} = 200 \text{GeV}$ 4.8 GeV < Q < 8.2 GeV 1.2 < |y| < 2.2Central scale:  $\mu_R = \mu_F = E_T^{\gamma}$ 10<sup>1</sup> 10<sup>0</sup> NNLOJET at LO  $10^{-1}$ NNLOJET at NLO NNLOJET at NNLO  $10^{-2}$ PHENIX 2 0  $p_T$  (GeV)

Physics Letters B 829 (2022) 137111



Phys.Rev.D 100 (2019) 1, 014018



#### Or only the finite part...

**PHENIX**  $\sqrt{s} = 200 \text{GeV}$ 4.8 GeV < Q < 8.2 GeV 1.2 < |y| < 2.2Central scale:  $\mu_R = \mu_F = E_T^{\gamma}$ 10<sup>1</sup> 100 NNLOJET at LO  $10^{-1}$ NNLOJET at NLO NNLOJET at NNLO  $10^{-2}$ PHENIX 2 0  $p_T$  (GeV)

Physics Letters B 829 (2022) 137111



Phys.Rev.D 100 (2019) 1, 014018





Finite part

and the second and the





and the second second and the second second second





•  $N^4LL$  perturbative accuracy, i.e. up to exp( ~  $\alpha_S^n L^{n-3}$ )



#### **Resummed part**



Finite part

- $N^4LL$  perturbative accuracy, i.e. up to exp
- Matching with NNLO corrections (i.e. up to  $\mathcal{O}(\alpha_S^3)$ ) at large  $q_T$



$$o(\sim \alpha_S^n L^{n-3})$$



Finite part

- $N^4LL$  perturbative accuracy, i.e. up to exp
- Matching with NNLO corrections (i.e. up to  $\mathcal{O}(\alpha_S^3)$ ) at large  $q_T$ Now implemented in DYTurbo



$$o(\sim \alpha_S^n L^{n-3})$$


Finite part

- $N^4LL$  perturbative accuracy, i.e. up to exp
- Matching with NNLO corrections (i.e. up to  $\mathcal{O}(\alpha_S^3)$ ) at large  $q_T$



$$p(\sim \alpha_S^n L^{n-3})$$

Now implemented in DYTurbo Benchmark with MCFM Phys. Rev. D 107 (2023) 1, L011506





Finite part

- $N^4LL$  perturbative accuracy, i.e. up to exp
- Matching with NNLO corrections (i.e. up to  $\mathcal{O}(\alpha_S^3)$ ) at large  $q_T$ Now implemented in DYTurbo



$$p(\sim \alpha_S^n L^{n-3})$$

Benchmark with MCFM Phys.Rev.D 107 (2023) 1, L011506 Benchmark with inclusive  $N^3LO$  cross section JHEP 12 (2022) 066





Finite part

- $N^4LL$  perturbative accuracy, i.e. up to exp
- Matching with NNLO corrections (i.e. up to  $\mathcal{O}(\alpha_S^3)$ ) at large  $q_T$

NLL QED corrections, i.e. up to  $\exp(\sim \alpha_{em}^n L^{n-1})$  JHEP 07 (2023) 104



$$p(\sim \alpha_S^n L^{n-3})$$

Now implemented in DYTurbo Benchmark with MCFM Phys.Rev.D 107 (2023) 1, L011506 Benchmark with inclusive  $N^3LO$  cross section JHEP 12 (2022) 066





Finite part

- $N^4LL$  perturbative accuracy, i.e. up to exp
- Matching with NNLO corrections (i.e. up to  $\mathcal{O}(\alpha_S^3)$ ) at large  $q_T$

- NLL QED corrections, i.e. up to  $\exp(\sim \alpha_{em}^n L^{n-1})$  JHEP 07 (2023) 104
- LL mixed QCD-QED accuracy, i.e. up to  $\exp(\sim \alpha_S^n \alpha_{em}^n L^{2n})$  JHEP 07 (2023) 104



$$p(\sim \alpha_S^n L^{n-3})$$

Now implemented in DYTurbo Benchmark with MCFM Phys.Rev.D 107 (2023) 1, L011506 Benchmark with inclusive  $N^3LO$  cross section JHEP 12 (2022) 066







Experiment	Observable	$\sqrt{s}  [\text{GeV}]$	$Q \; [{ m GeV}]$	$y \text{ or } x_F$
E605	$Ed^{3}\sigma/d^{3}\boldsymbol{q}$	38.8	7 - 18	$x_{F} = 0.1$
$E288 \ 200 \ GeV$	$Ed^{3}\sigma/d^{3}q$	19.4	4 - 9	y = 0.40
E288 300 $\mathrm{GeV}$	$Ed^{3}\sigma/d^{3}oldsymbol{q}$	23.8	4 - 12	y = 0.21
$E288 \ 400 \ GeV$	$Ed^{3}\sigma/d^{3}q$	27.4	5 - 14	y = 0.03
ATLAS 8 TeV on-peak	$(1/\sigma)d\sigma/d \mathbf{q}_T $	8000	66 - 116	$\begin{array}{l}  y  < 0.4 \\ 0.4 <  y  < 0.8 \\ 0.8 <  y  < 1.2 \\ 1.2 <  y  < 1.6 \\ 1.6 <  y  < 2 \\ 2 <  y  < 2.4 \end{array}$
CDF Run II	$d\sigma/d oldsymbol{q}_T $	1960	66 - 116	Inclusive



Experiment	Observable	$\sqrt{s} \; [\text{GeV}]$	$Q \; [{ m GeV}]$	$y \text{ or } x_F$
E605	$Ed^{3}\sigma/d^{3}q$	38.8	7 - 18	$x_{F} = 0.1$
E288 200 $\mathrm{GeV}$	$Ed^{3}\sigma/d^{3}oldsymbol{q}$	19.4	4 - 9	y = 0.40
E288 300 $\mathrm{GeV}$	$Ed^{3}\sigma/d^{3}oldsymbol{q}$	23.8	4 - 12	y = 0.21
E288 400 $\mathrm{GeV}$	$Ed^{3}\sigma/d^{3}oldsymbol{q}$	27.4	5 - 14	y = 0.03
ATLAS 8 TeV on-peak	$(1/\sigma)d\sigma/d \mathbf{q}_T $	8000	66 - 116	$\begin{array}{l}  y  < 0.4 \\ 0.4 <  y  < 0.8 \\ 0.8 <  y  < 1.2 \\ 1.2 <  y  < 1.6 \\ 1.6 <  y  < 2 \\ 2 <  y  < 2.4 \end{array}$
CDF Run II	$d\sigma/d oldsymbol{q}_T $	1960	66 - 116	Inclusive



Experiment	Observable	$\sqrt{s} \; [\text{GeV}]$	$Q \; [{ m GeV}]$	$y \text{ or } x_F$
E605	$Ed^{3}\sigma/d^{3}q$	38.8	7 - 18	$x_{F} = 0.1$
E288 200 $\mathrm{GeV}$	$Ed^{3}\sigma/d^{3}oldsymbol{q}$	19.4	4 - 9	y = 0.40
E288 300 $\mathrm{GeV}$	$Ed^{3}\sigma/d^{3}oldsymbol{q}$	23.8	4 - 12	y = 0.21
E288 400 $\mathrm{GeV}$	$Ed^{3}\sigma/d^{3}oldsymbol{q}$	27.4	5 - 14	y = 0.03
ATLAS 8 TeV on-peak	$(1/\sigma)d\sigma/d \mathbf{q}_T $	8000	66 - 116	$\begin{array}{l}  y  < 0.4 \\ 0.4 <  y  < 0.8 \\ 0.8 <  y  < 1.2 \\ 1.2 <  y  < 1.6 \\ 1.6 <  y  < 2 \\ 2 <  y  < 2.4 \end{array}$
CDF Run II	$d\sigma/d oldsymbol{q}_T $	1960	66 - 116	Inclusive



#### Already well described



Experiment	Observable	$\sqrt{s} \; [\text{GeV}]$	$Q \; [{ m GeV}]$	$y \text{ or } x_F$
E605	$Ed^{3}\sigma/d^{3}q$	38.8	7 - 18	$x_{F} = 0.1$
E288 200 $GeV$	$Ed^{3}\sigma/d^{3}oldsymbol{q}$	19.4	4 - 9	y = 0.40
E288 300 $\mathrm{GeV}$	$Ed^{3}\sigma/d^{3}oldsymbol{q}$	23.8	4 - 12	y = 0.21
$E288 \ 400 \ GeV$	$Ed^{3}\sigma/d^{3}oldsymbol{q}$	27.4	5 - 14	y = 0.03
ATLAS 8 TeV on-peak	$(1/\sigma)d\sigma/d \mathbf{q}_T $	8000	66 - 116	$\begin{array}{l}  y  < 0.4 \\ 0.4 <  y  < 0.8 \\ 0.8 <  y  < 1.2 \\ 1.2 <  y  < 1.6 \\ 1.6 <  y  < 2 \\ 2 <  y  < 2.4 \end{array}$
CDF Run II	$d\sigma/d m{q}_T $	1960	66 - 116	Inclusive



#### Some stuff...

#### Some stuff...

#### • Landau pole:

#### Some stuff...

Landau pole:  $b_T \rightarrow b_T \qquad b_{\pm}(b_T) = \frac{b_T}{\sqrt{1 + \frac{b_T^2}{b_{max}^2}}}$  Nucl. Phys. B250 (1985)199



#### Landau pole:

(We tried also the minimal prescription)

 $b_T \rightarrow b_{\star}(b_T) = \frac{b_T}{\sqrt{1 + \frac{b_T^2}{b_{max}^2}}}$  Nucl. Phys. B250 (1985)199



# Landau pole:

(We tried also the minimal prescription)

Unitary condition:

 $b_T \rightarrow b_{\star}(b_T) = \frac{b_T}{\sqrt{1 + \frac{b_T^2}{b_{max}^2}}}$  Nucl.Phys.B250 (1985)199



- Landau pole:  $b_T \rightarrow b_{\star}(b_T) = \frac{b_T}{\sqrt{1 + \frac{b_T^2}{b_{max}^2}}}$  Nucl. Phys. B250 (1985)199
- Unitary condition:  $L = log(M^2b^2)$

#### Some stuff...

- (We tried also the minimal prescription)
- Unitary condition:  $L = log(M^2b^2)$   $\tilde{L} = log(M^2b^2 + 1)_{Nucl. Phys. B 407 (1993) 3-42}$

Landau pole:  $b_T \rightarrow b_{\pm}(b_T) = \frac{b_T}{\sqrt{1 + \frac{b_T^2}{b_{max}^2}}}$  Nucl. Phys. B250 (1985) 199

#### Some stuff...

- (We tried also the minimal prescription)
- Unitary condition:  $L = log(M^2b^2)$   $\tilde{L} = log(M^2b^2 + 1)_{Nucl. Phys. B 407 (1993) 3-42}$

Landau pole:  $b_T \rightarrow b_{\pm}(b_T) = \frac{b_T}{\sqrt{1 + \frac{b_T^2}{b_{max}^2}}}$  Nucl. Phys. B250 (1985)199

 $\exp\{\alpha_S^n \tilde{L}^k\}\Big|_{b_T=0} = 1$ 

#### Some stuff...

- (We tried also the minimal prescription)
- Unitary condition:  $L = log(M^2b^2)$   $\tilde{L} = log(M^2b^2 + 1)_{Nucl. Phys. B 407 (1993) 3-42}$

$$\int_0^\infty d^2 \mathbf{q}_{\mathbf{T}} \, \frac{d\sigma}{d^2 \mathbf{q}_{\mathbf{T}}} = \sigma^{tot}$$

Landau pole:  $b_T \rightarrow b_T (b_T) = \frac{b_T}{\sqrt{1 + \frac{b_T^2}{b_{max}^2}}}$  Nucl. Phys. B250 (1985)199

 $\exp\{\alpha_S^n \tilde{L}^k\}\Big|_{b_T=0} = 1$ 









#### • Non perturbative parametrisation:







Non perturbative parametrisation:  $S_{NP}(b_T) = \exp\left[-g_j(b_T) - g_k(b_T)\log\frac{M^2}{O_0^2}\right]$ 

Phys.Rev.D 91 (2015) 7, 074020







$$g_j(b_T) = \frac{g_1 b_T^2}{1 + \lambda b_T^2} + \operatorname{sign}(q) \left(1\right)$$

Non perturbative parametrisation:  $S_{NP}(b_T) = \exp\left[-g_j(b_T) - g_k(b_T) \log \frac{M^2}{O_0^2}\right]$ 

Phys.Rev.D 91 (2015) 7, 074020

 $-\exp[-|q|b_T^4]\Big)$ 





$$g_{j}(b_{T}) = \frac{g_{1}b_{T}^{2}}{1+\lambda b_{T}^{2}} + \operatorname{sign}(q) \left(1 - \exp[-|q|b_{T}^{4}]\right)$$
$$g_{K}(b_{T}) = g_{0} \left(1 - \exp\left[-\frac{C_{F}\alpha_{S}(\mu_{b_{\star}})b_{T}^{2}}{\pi g_{0}b_{max}^{2}}\right]\right)$$

Non perturbative parametrisation:  $S_{NP}(b_T) = \exp\left[-g_j(b_T) - g_k(b_T)\log\frac{M^2}{Q_0^2}\right]$ 

Phys.Rev.D 91 (2015) 7, 074020

 $-\exp[-|q|b_T^4]$ 







Non perturbative parametrisation:  $S_{NP}(b_T) = \exp\left[-g_j(b_T) - g_k(b_T)\log\frac{M^2}{Q_0^2}\right]$ 

Phys.Rev.D 91 (2015) 7, 074020

180<sup>°</sup> max





Non perturbative parametrisation:



 $S_{NP}(b_T) = \exp\left[-g_j(b_T) - g_k(b_T) \log \frac{M^2}{O_0^2}\right]$ 

Phys.Rev.D 91 (2015) 7, 074020









Eur.Phys.J.C 80 (2020) 3, 251





Eur.Phys.J.C 80 (2020) 3, 251







Eur.Phys.J.C 80 (2020) 3, 251

![](_page_66_Picture_3.jpeg)

![](_page_66_Picture_4.jpeg)

![](_page_66_Picture_5.jpeg)

![](_page_67_Picture_1.jpeg)

Eur.Phys.J.C 80 (2020) 3, 251

#### Fit at central scale $\mu_F = \mu_R = \mu_O = M$

#### Including PDFs uncertainty (MSHT20\_an3lo)

![](_page_67_Picture_5.jpeg)

![](_page_67_Picture_7.jpeg)

![](_page_68_Picture_1.jpeg)

Eur.Phys.J.C 80 (2020) 3, 251

#### Fit at central scale $\mu_F = \mu_R = \mu_O = M$

#### Including PDFs uncertainty (MSHT20\_an3lo)

#### 378 datapoints

![](_page_68_Picture_6.jpeg)

![](_page_68_Picture_8.jpeg)

![](_page_69_Picture_1.jpeg)

Eur.Phys.J.C 80 (2020) 3, 251

#### Fit at central scale $\mu_F = \mu_R = \mu_Q = M$

- Including PDFs uncertainty (MSHT20\_an3lo)
- 378 datapoints  $\chi^2 / N_{d.o.f.} = 1.28$

![](_page_69_Picture_6.jpeg)

![](_page_69_Picture_9.jpeg)

![](_page_70_Picture_1.jpeg)

Eur.Phys.J.C 80 (2020) 3, 251

#### Fit at central scale $\mu_F = \mu_R = \mu_O = M$

- Including PDFs uncertainty (MSHT20\_an3lo)
- 378 datapoints  $\gamma^2/N_d$ **u.0**.**j**.

![](_page_70_Picture_6.jpeg)

#### Not to compare with the

(unpolarized) TMD fit

![](_page_70_Picture_12.jpeg)

![](_page_71_Picture_1.jpeg)

Eur.Phys.J.C 80 (2020) 3, 251

#### Fit at central scale $\mu_F = \mu_R = \mu_O = M$

- Including PDFs uncertainty (MSHT20\_an3lo)
- 378 datapoints  $\gamma^2/N_{do}$ *a.0*.*j*.

![](_page_71_Picture_6.jpeg)

#### Not to compare with the

(unpolarized) TMD fit

We have high- $q_T$  data!

.28

![](_page_71_Picture_14.jpeg)


### E288 300 GeV, 5 < Q < 6





### E288 300 GeV, 5 < Q < 6





### E288 300 GeV, 5 < Q < 6



### E605, 11.5 < Q < 13.5





### E288 300 GeV, 5 < Q < 6



### E605, 11.5 < Q < 13.5



 $\chi^2/N_{data} = 0.79$ 





Phys.Rev.D 100 (2019) 1, 014018



Phys.Rev.D 100 (2019) 1, 014018





Phys.Rev.D 100 (2019) 1, 014018

### Starting point





Phys.Rev.D 100 (2019) 1, 014018

### Starting point



Arrival point









# High invariant mass data remain well descrbed



































We are working on the choice of the non perturbative (hyper)parameters...





# • We found a very good description of by using the "full" formalism

We found a very good description of the low invariant mass  $q_T$  spectrum data



- We found a very good description of by using the "full" formalism
- We found a good consistency in the operation invariant mass data

We found a very good description of the low invariant mass  $q_T$  spectrum data



- by using the "full" formalism
- invariant mass data

We extracted a Collins Soper kernel compatible with the literature

We found a very good description of the low invariant mass  $q_T$  spectrum data



- by using the "full" formalism
- invariant mass data

We extracted a Collins Soper kernel compatible with the literature

Refine the non perturbative effects

We found a very good description of the low invariant mass  $q_T$  spectrum data



- by using the "full" formalism
- invariant mass data

We extracted a Collins Soper kernel compatible with the literature

Refine the non perturbative effects

Study of the available data-sets

We found a very good description of the low invariant mass  $q_T$  spectrum data



### Resummation, Evolution, Factorization 2025

### 13–17 Oct 2025 Physics Department, Milan University

Europe/Rome timezone

#### Overview

Workshop venue

Registration

Registration #2: fee payment

Call for Abstracts

Participant List

Timetable

Accommodation

Workshop dinner (Wednesday 15/10)

Videoconference information



#### Contact

Enter your search term



Q