

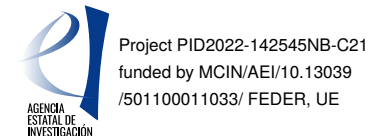
Present status of three-family global fits to neutrino oscillation data

Michele Maltoni

Instituto de Física Teórica UAM/CSIC

2nd workshop on “Multi-Aspect Young-ORiented Advanced Neutrino Academy (MAYORANA)”

Palazzo Grimaldi, Modica, Italy – June 17th, 2025



SM with ν masses: general three-neutrino framework

- Equation of motion: **6 parameters** (including **Dirac** and neglecting **Majorana** phases):

$$i \frac{d\vec{\nu}}{dt} = H \vec{\nu}; \quad H = U_{\text{vac}} \cdot D_{\text{vac}} \cdot U_{\text{vac}}^\dagger \pm V_{\text{mat}};$$

$$U_{\text{vac}} = \begin{pmatrix} 1 & 0 & 0 \\ 0 & c_{23} & s_{23} \\ 0 & -s_{23} & c_{23} \end{pmatrix} \cdot \begin{pmatrix} c_{13} & 0 & s_{13} e^{-i\delta_{\text{CP}}} \\ 0 & 1 & 0 \\ -s_{13} e^{i\delta_{\text{CP}}} & 0 & c_{13} \end{pmatrix} \cdot \begin{pmatrix} c_{12} & s_{12} & 0 \\ -s_{12} & c_{12} & 0 \\ 0 & 0 & 1 \end{pmatrix} \cdot \begin{pmatrix} e^{i\eta_1} & 0 & 0 \\ 0 & e^{i\eta_2} & 0 \\ 0 & 0 & 1 \end{pmatrix}, \quad \vec{\nu} = \begin{pmatrix} \nu_e \\ \nu_\mu \\ \nu_\tau \end{pmatrix};$$

$$D_{\text{vac}} = \frac{1}{2E_\nu} \left[\text{diag} (0, \Delta m_{21}^2, \Delta m_{31}^2) + \cancel{m_1^2 I} \right]; \quad V_{\text{mat}} = \sqrt{2} G_F N_e \text{diag} (1, 0, 0).$$

6 parameters $\iff$ 6 types of experiments

- ν_e disapp: $\left\{ \begin{array}{l} \text{-- solar experiments (mainly SNO)} \rightarrow \theta_{12} \\ \text{-- reactor LBL (KamLAND)} \rightarrow \Delta m_{21}^2 \\ \text{-- reactor MBL (Double-Chooz, Daya-Bay, Reno)} \rightarrow \theta_{13} [\Delta m_{31}^2] \end{array} \right\} \text{ MeV sources}$
- ν_μ disapp: $\left\{ \begin{array}{l} \text{-- atmospheric experiments (SK, IC)} \rightarrow \theta_{23} [\Delta m_{31}^2] \\ \text{-- accelerator long-baseline (T2K, NOvA)} \rightarrow |\Delta m_{31}^2| [\theta_{23}] \end{array} \right\} \text{ GeV sources}$
- ν_e appear: $\left\{ \begin{array}{l} \text{-- atmospheric (SK) \& accelerator (T2K, NOvA)} \rightarrow \delta_{\text{CP}}, \pm \Delta m_{31}^2 \end{array} \right\}$

Solar and reactor neutrinos

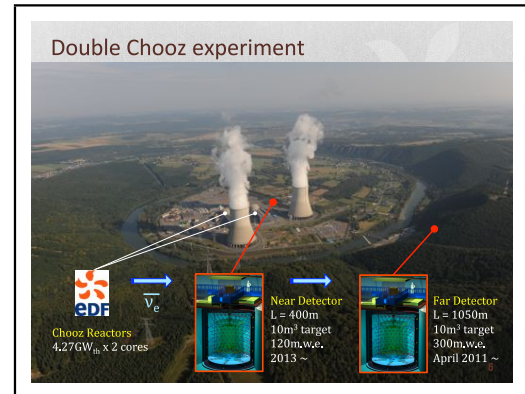
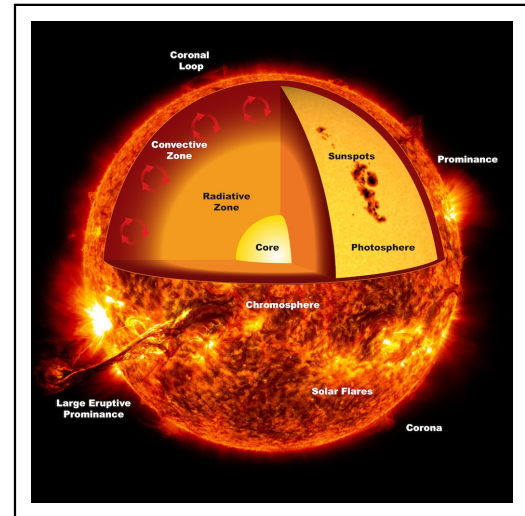
- ν_e from **nuclear** reactions $\Rightarrow$ energy in the MeV range;
- ν_μ and ν_τ indistinguishable $\Rightarrow$ no sensitivity to θ_{23} or δ_{CP} .

Reactor neutrinos

- $\bar{\nu}_e$ produced by nuclear **fission** in reactor's core;
- detection: inverse beta decay ($\bar{\nu}_e + p \rightarrow e^+ + n$), both e^+ and n observed in coincidence;
- negligible matter effects $\Rightarrow$ mostly vacuum params.

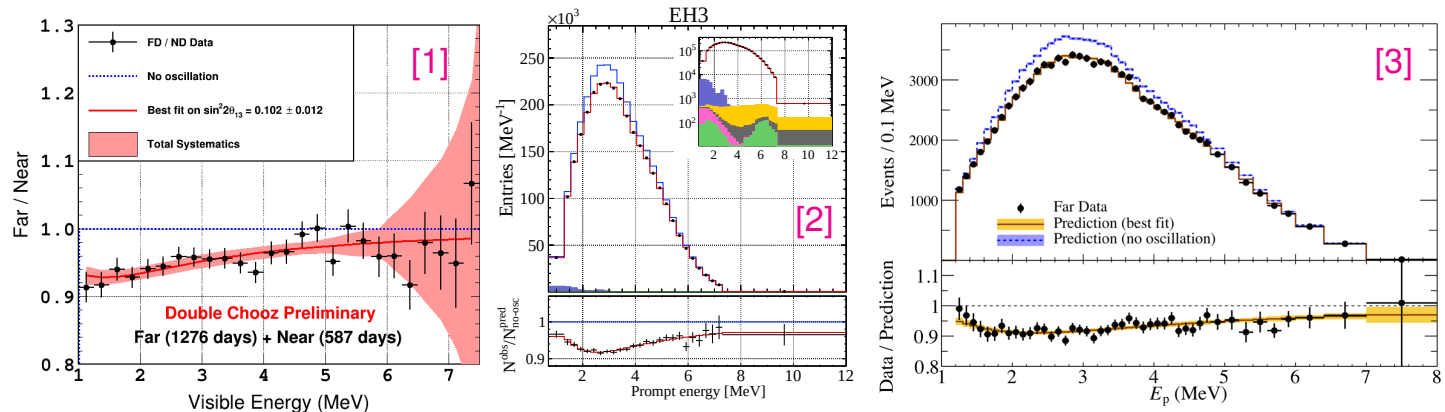
Solar neutrinos

- ν_e produced by nuclear **fusion** in the core of the Sun;
- two different mechanisms at work: **p-p chain** and **CNO cycle**. Both give $4p \rightarrow {}^4\text{He} + 2e^+ + 2\nu_e + \gamma \Rightarrow$ solar **light** and **neutrinos** in well-defined mutual proportions;
- detection: various processes (**CC- ν_e** , **NC**, **ES**);
- matter effects very important (MSW effect).



Medium-baseline reactor neutrino disappearance and θ_{13}

- Positive $\bar{\nu}_e$ disappearance (≈ 1 km) in DOUBLE-CHOOZ [1], DAYA-BAY [2], RENO [3];
- experimental results are mutually consistent $\Rightarrow$ it is now a firmly established fact that $\theta_{13} \neq 0$
 $\Rightarrow$ full 3ν oscillation phenomenology;
- all these experiments have spectral capabilities and detector units placed at different baselines
 $\Rightarrow$ uncertainties in the reactor flux predictions do **not** affect the results.



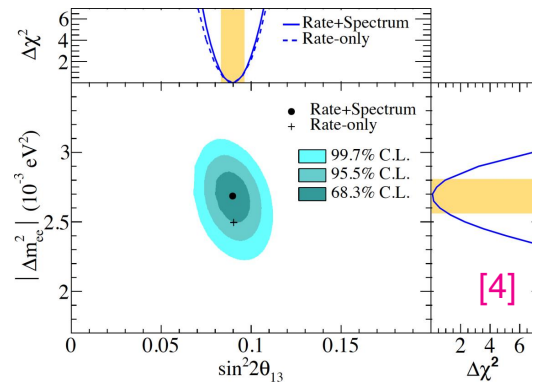
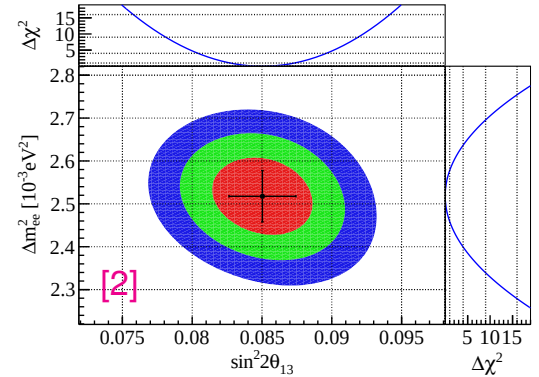
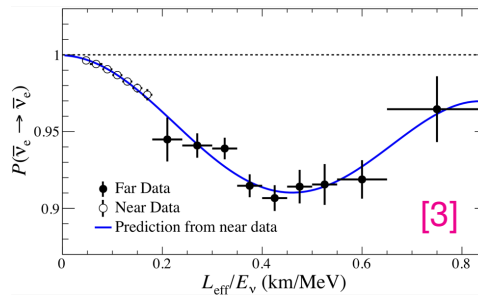
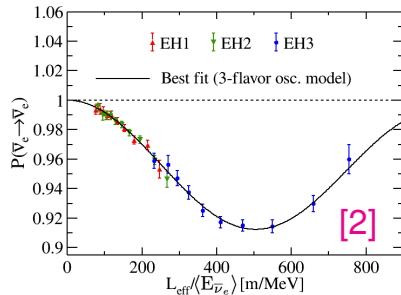
[1] T. Bezerra [DOUBLE-CHOOZ], online talk at Neutrino 2020, Fermilab, USA, June 22–July 2, 2020.

[2] F.P. An *et al.* [DAYA BAY], Phys. Rev. Lett. **130** (2023) 161802 [arXiv:2211.14988].

[3] J. Yoo [RENO], online talk presented at Neutrino 2020, Fermilab, USA, June 22–July 2, 2020.

Measuring θ_{13} and Δm_{31}^2 from reactor data

- FAR/NEAR spectral ratio $\Rightarrow$ flux shape irrelevant;
- spectral information from Double-Chooz, Daya-Bay and Reno $\Rightarrow$ oscillation pattern clearly visible $\Rightarrow \theta_{13}$ and Δm_{31}^2 accurately determined by reactor data;
- accuracy from **reactor** $\nu_e \rightarrow \nu_e$ comparable with **LBL** $\nu_\mu \rightarrow \nu_\mu$, but oscillation channel is different $\Rightarrow$ important **complementary** information available.



[2] F.P. An *et al.* [DAYA BAY], Phys. Rev. Lett. **130** (2023) 161802 [arXiv:2211.14988].

[3] J. Yoo [RENO], online talk at Neutrino 2020, Fermilab, USA, 22/06–2/07/2020.

[4] K.K. Joo [RENO], online talk at Neutrino 2022, Virtual Seoul, Korea, 30/05–4/06/2022.

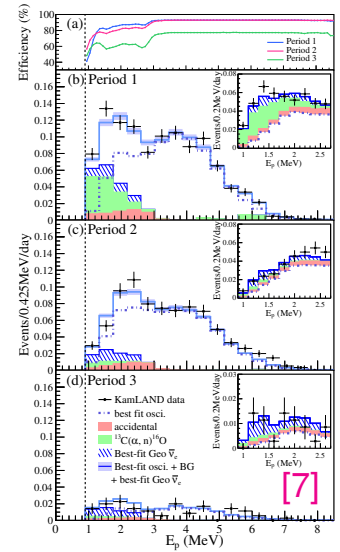
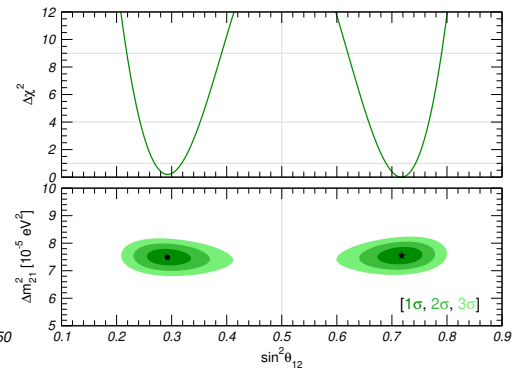
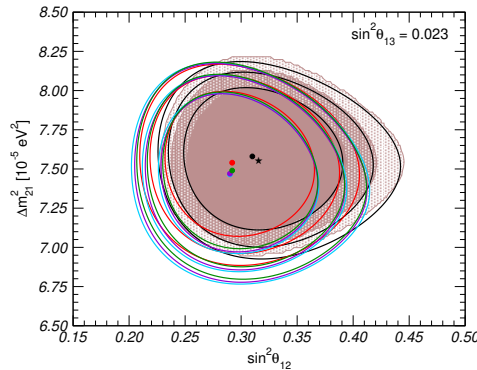
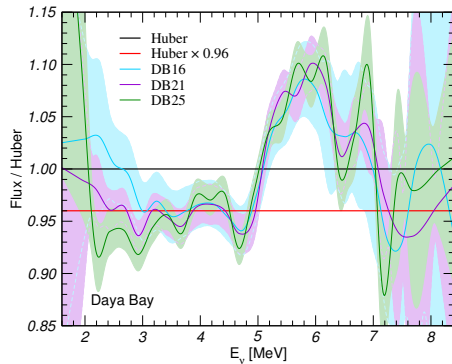
Measuring θ_{12} and Δm_{21}^2 with KamLAND data

- Much longer baseline (≈ 180 km) $\Rightarrow$ sensitive to θ_{12} and Δm_{21}^2 ;
- lack of a near detector $\Rightarrow$ spectral distortions may be an issue;
- problem discussed in [5] $\Rightarrow$ impact on Δm_{21}^2 found to be small;
- solution: bind KamLAND spectrum to Daya-Bay measurement [6];
- marginal matter contributions $\Rightarrow$ both θ_{12} octants allowed by data.

[5] F. Capozzi *et al.*, NPB **908** (2016) 218 [arXiv:1601.07777].

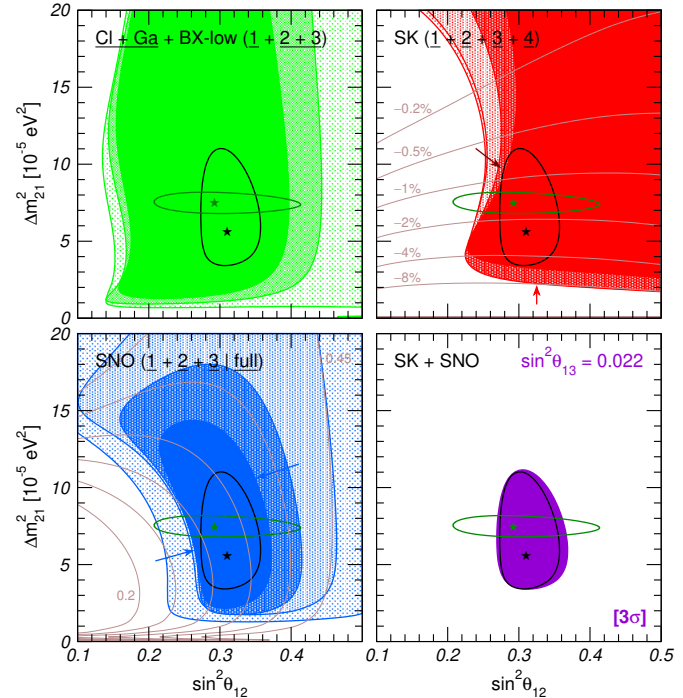
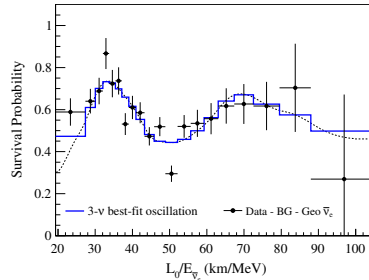
[6] F.P. An *et al.* [Daya Bay], PRL **134** (2025) 201802 [arXiv:2501.00746].

[7] A. Gando *et al.* [KamLAND], PRD **88** (2013) 033001 [arXiv:1303.4667].



Determination of θ_{12} and Δm_{21}^2 from solar neutrino data

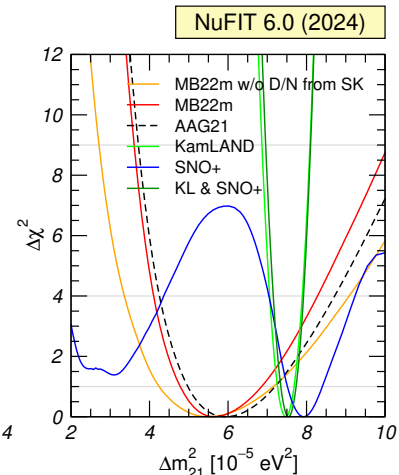
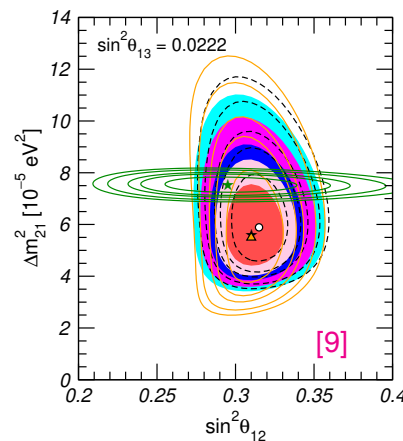
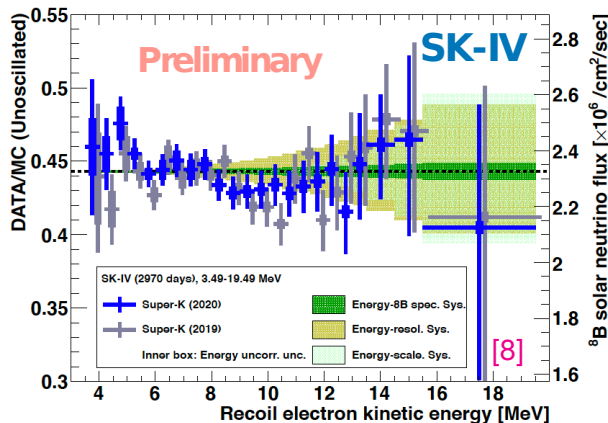
- $P_{ee} = c_{13}^4 P_{\text{eff}} + s_{13}^4$, $i \frac{d\vec{v}}{dt} = \left[\frac{\Delta m_{21}^2}{4E_\nu} \begin{pmatrix} -\cos 2\theta_{12} & \sin 2\theta_{12} \\ \sin 2\theta_{12} & \cos 2\theta_{12} \end{pmatrix} \pm \sqrt{2} G_F N_e \begin{pmatrix} c_{13}^2 & 0 \\ 0 & 0 \end{pmatrix} \right] \vec{v}$, $\vec{v} = \begin{pmatrix} \nu_e \\ \nu_a \end{pmatrix}$;
 - $\nu_\mu \equiv \nu_\tau \Rightarrow$ no sensitivity to θ_{23} and δ_{CP} ;
 - $\Delta m_{31}^2 \approx \infty \Rightarrow$ specific Δm_{31}^2 value irrelevant;
- $\Rightarrow$ data only depend on Δm_{21}^2 , θ_{12} and θ_{13} ;
- param's: $\begin{cases} \theta_{12} \text{ dominated by SNO;} \\ \Delta m_{21}^2 \text{ dominated by KamLAND;} \end{cases}$
 - region given by high-E data, low-E marginal;
 - strong matter effects $\Rightarrow$ fixes $\theta_{12} < 45^\circ$ octant;
 - SNO-NC measurement confirms SSM;
 - KamLAND precisely determines the oscillation pattern.



Comparison between solar and KamLAND measurements

- Long-standing weak tension on preferred Δm_{21}^2 from solar and KamLAND data;
- choice of the assumed solar model (GS, AGSS, MB22, ...) has little impact on the issue;
- cause:
 - too much D/N asymmetry in SK
 - no indication of low-E turn-up
- new data [8]:
 - D/N: 3.6% $\rightarrow$ 2.1%,
 - “hints” of turn-up.

$\Rightarrow$ tension considerably reduced after Neutrino 2020 conference.



[8] Y. Nakajima [SK], online talk at Neutrino 2020, Fermilab, USA, June 22–July 2, 2020.

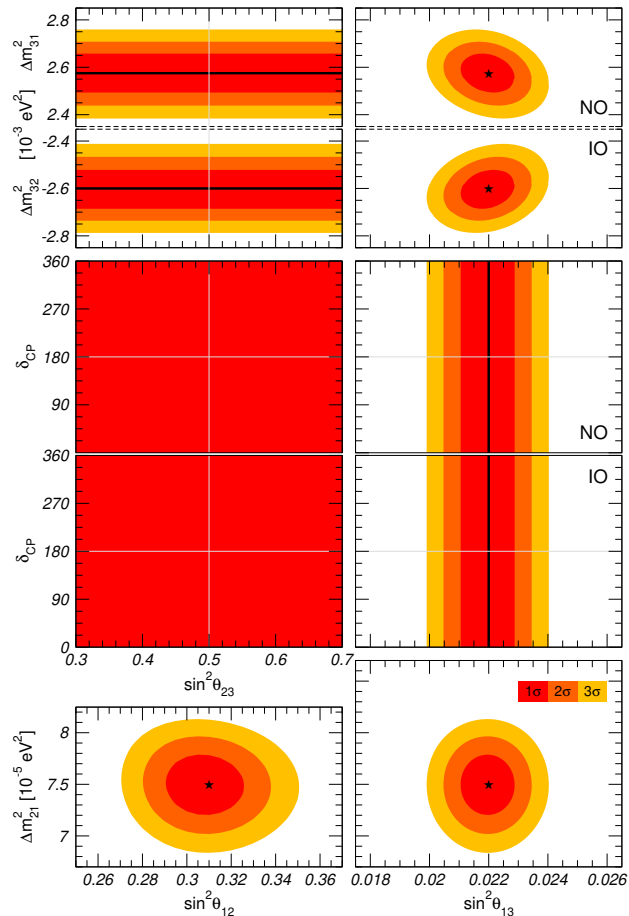
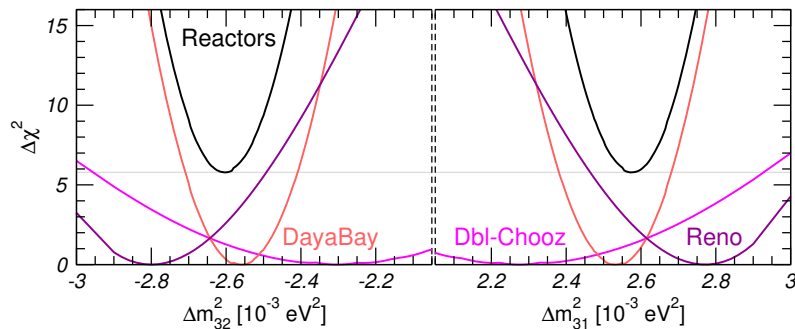
[9] I. Esteban *et al.*, JHEP **12** (2024) 216 [arXiv:2410.05380] & NuFIT 6.0 [http://www.nu-fit.org].

Global ν_e disappearance bounds

- Both $(\theta_{12}, \Delta m_{21}^2)$ and $(\theta_{13}, |\Delta m_{31}^2|)$ determined;
- all MBL reactors sensitive to Δm_{ij}^2 only through the same effective combination $|\Delta m_{ee}^2|$, with [10]:

$$\Delta m_{ee}^2 \equiv \frac{|U_{e1}|^2 \Delta m_{31}^2 + |U_{e2}|^2 \Delta m_{32}^2}{|U_{e1}|^2 + |U_{e2}|^2};$$

- consequently, tensions among MBL reactors are the same for **NO** and **IO** $\Rightarrow$ no info on ordering.



[10] Nunokawa *et al.*, Phys. Rev. D [hep-ph/0503283].

Atmospheric and accelerator neutrinos

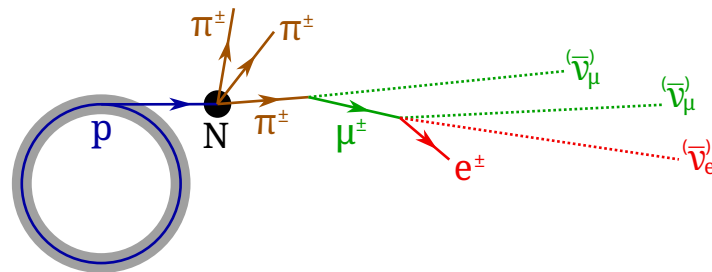
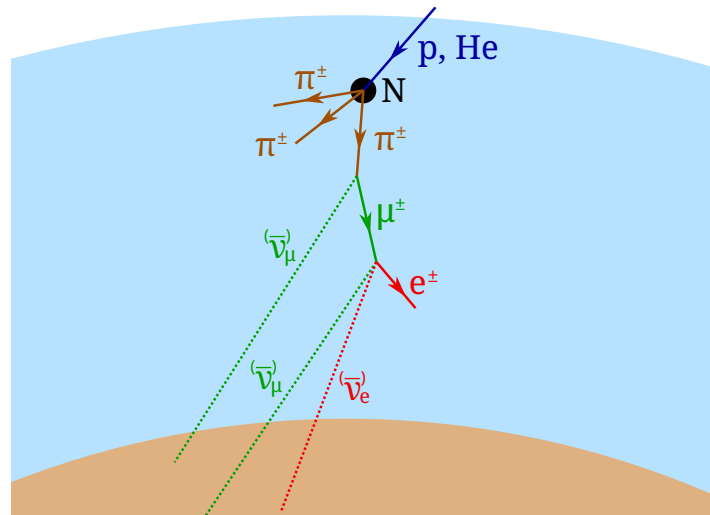
- Atmospheric (accelerator) neutrinos are produced by the interaction of *cosmic rays* (*protons*) with the *Earth's atmosphere* (*target*):

$$1 \quad A_{\text{in}} + A_{\text{tgt}} \rightarrow \pi^\pm, K^\pm, K^0, \dots$$

$$2 \quad \pi^\pm \rightarrow \mu^\pm + \nu_\mu,$$

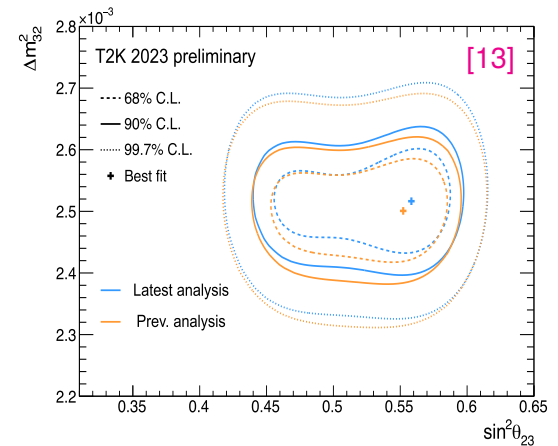
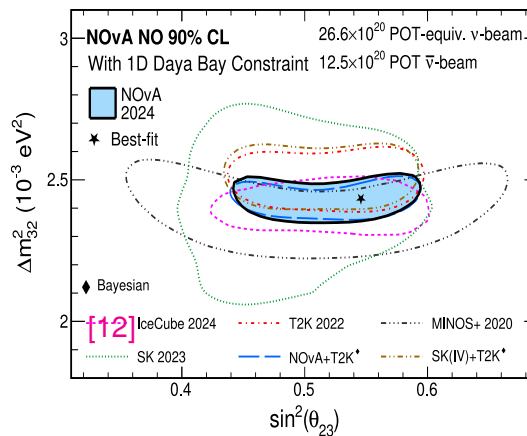
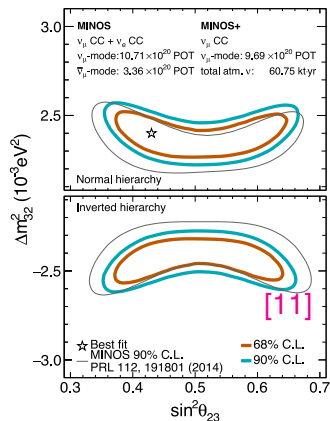
$$3 \quad \mu^\pm \rightarrow e^\pm + \nu_e + \nu_\mu;$$

- at the detector, some ν interacts and produces a **charged lepton**, which is observed;
- atmospheric**: fluxes of ν_μ and ν_e are known with poor precision ($\approx 20\%$), but the accuracy on the ν_μ/ν_e ratio is better ($\approx 5\%$);
- accelerator**: a **near detector** allow to characterize the unoscillated ν flux.



Consistency of Δm_{31}^2 and θ_{23} from accelerator data

- Δm_{31}^2 & θ_{23} dominated by LBL disappearance ($\nu_\mu \rightarrow \nu_\mu$) data;
- Δm_{21}^2 & θ_{12} subleading contributions to LBL appearance ($\nu_\mu \rightarrow \nu_e$) $\Rightarrow$ relevant for δ_{CP} ;
- reasonably good agreement between all experiments in the allowed regions, although some small differences are visible.



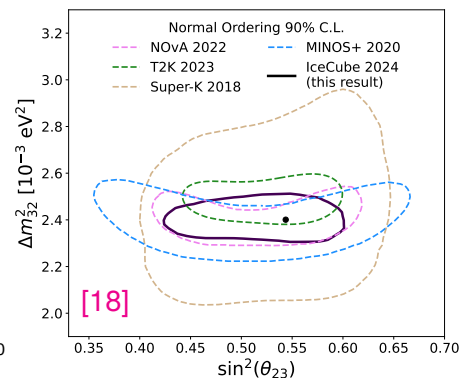
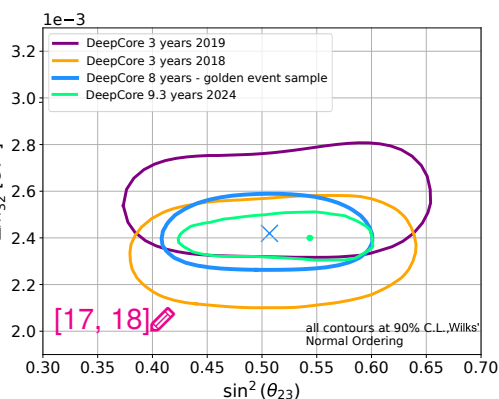
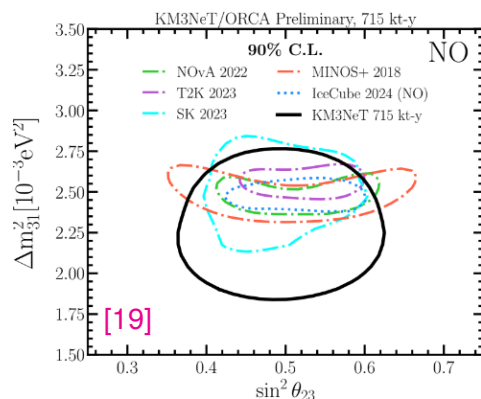
[11] T. Carroll [MINOS], online talk at Neutrino 2020, Fermilab, USA, June 22–July 2, 2020.

[12] J. Wolcott [NOvA], talk at Neutrino 2024, Milan, Italy, June 16–22, 2024.

[13] D. Carabadjac [T2K], talk at ICHEP 2024, Prague, Czech Republic, July 17–24, 2024.

The contribution of neutrino telescopes through atmospheric data

- IceCUBE/DeepCore: after many 3-year “calibration” fits [14, 15, 16], updated 8-year [17] and 9.3-year [18] results presented (but not “released”) $\Rightarrow$ competitive with reactors and LBL;
- Km3NET/ORCA: under deployment (23/115 strings so far [19]), fit (ORCA6+11) catching up.



[14] M.G. Aartsen *et al.* [ICECUBE], PRD **91** (2015) 072004 [arXiv:1410.7227], updated Oct. 2016. [IC16]

[15] M.G. Aartsen *et al.* [ICECUBE], PRL **120** (2018) 071801 [arXiv:1707.07081].

[16] M.G. Aartsen *et al.* [ICECUBE], PRD **99** (2019) 032007 [arXiv:1901.05366]. [IC19]

[17] R. Abbasi *et al.* [ICECUBE], PRD **108** (2023) 012014 [arXiv:2304.12236].

[18] R. Abbasi *et al.* [ICECUBE], PRL **134** (2025) 091801 [arXiv:2405.02163]. [IC24]

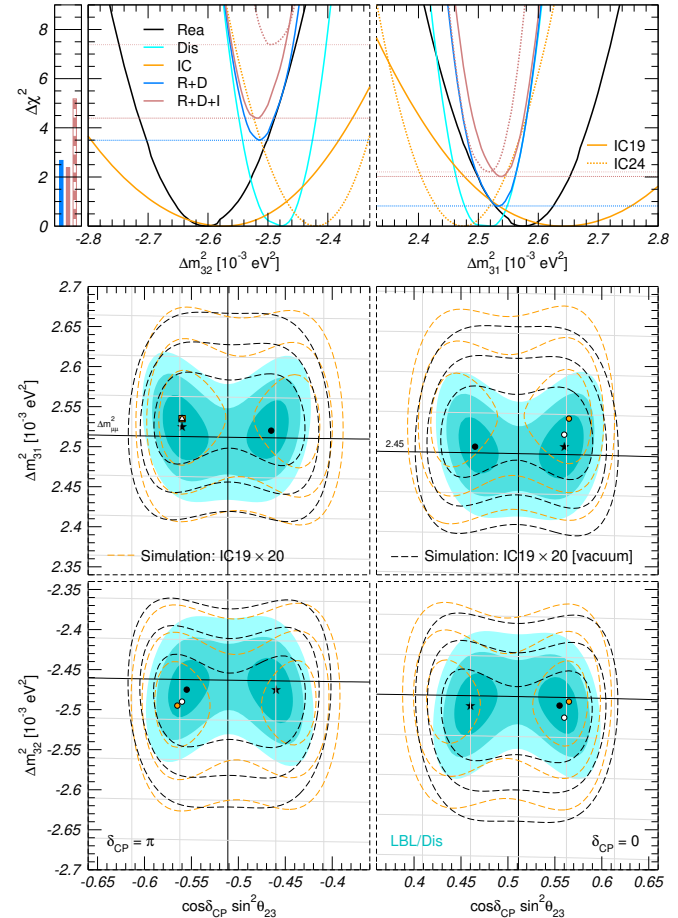
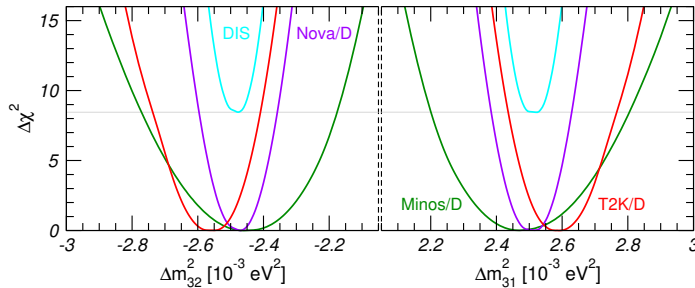
[19] J. Coelho [KM3NET], talk at Neutrino 2024, Milan, Italy, June 16–22, 2024.

Δm_{31}^2 and mass ordering

- All **LBL-dis** data effectively measure [10]:

$$\Delta m_{\mu\mu}^2 \equiv \frac{|U_{\mu 1}|^2 \Delta m_{31}^2 + |U_{\mu 2}|^2 \Delta m_{32}^2}{|U_{\mu 1}|^2 + |U_{\mu 2}|^2};$$

- neither **reactor- ν_e** nor **accelerator- ν_μ** data are sensitive to the ordering, but their **combination** is: $\Delta m_{ee}^2 \neq \Delta m_{\mu\mu}^2 \Rightarrow$ better agreement for **NO**;
- IC**: matter effects sizable $\Rightarrow \nu_\mu$ data do **not** follow $\Delta m_{\mu\mu}^2 \Rightarrow$ complementary to **LBL-dis**.



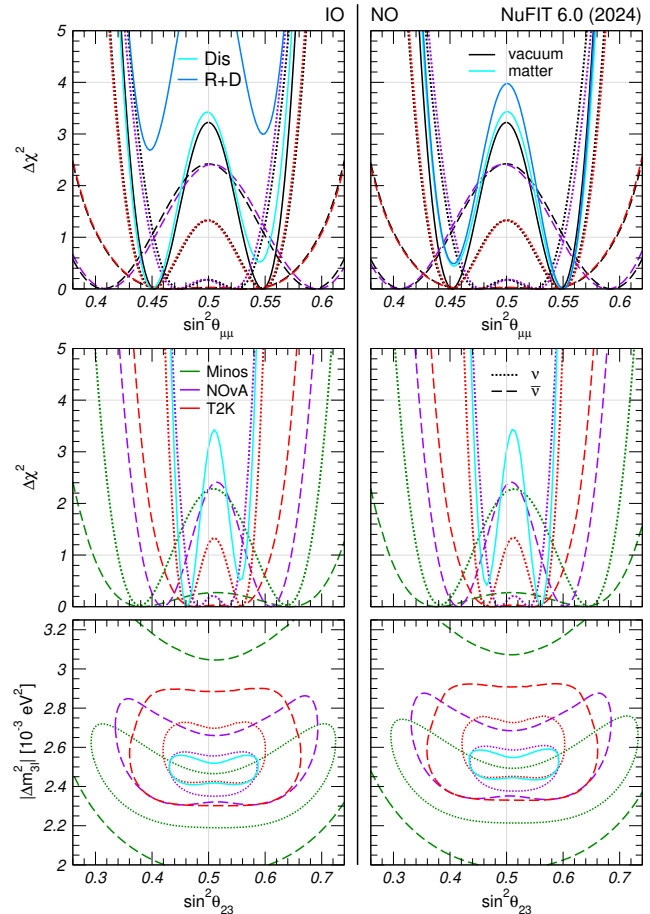
[10] Nunokawa *et al.*, PRD [hep-ph/0503283].

θ_{23} mixing and octant

- Effective LBL-disappearance mixing angle:

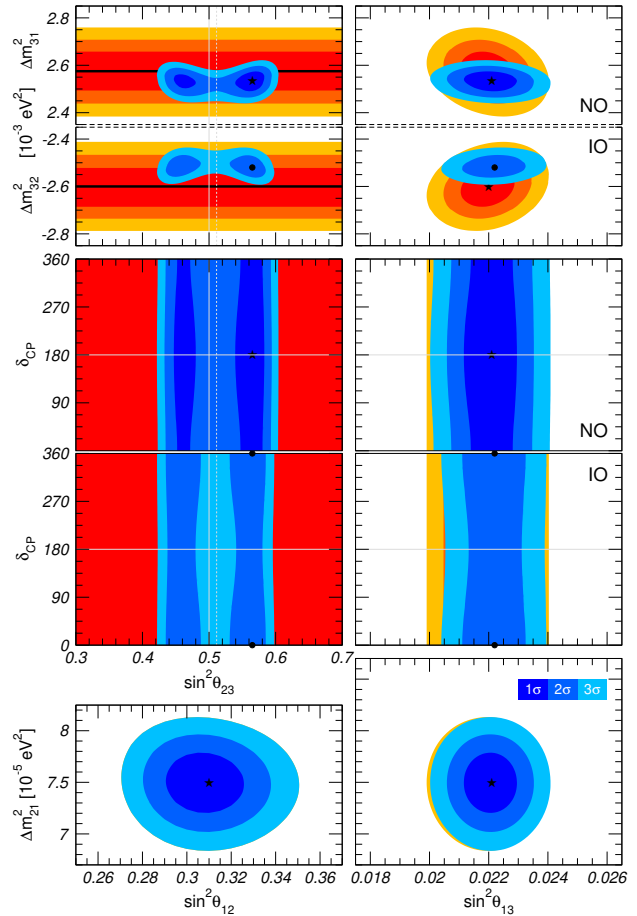
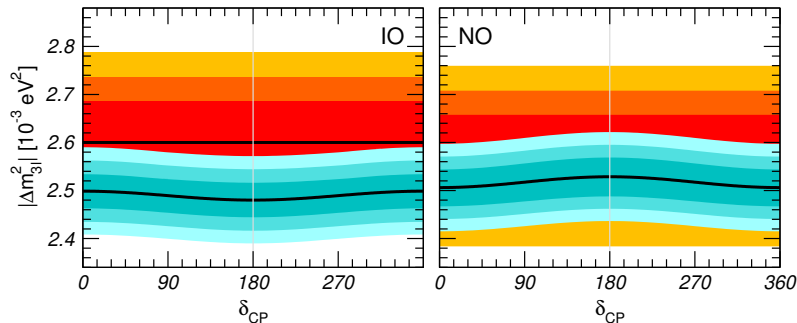
$$\sin^2 \theta_{\mu\mu} \equiv |U_{\mu 3}|^2 = \sin^2 \theta_{23} \cos^2 \theta_{13};$$

- each individual LBL-dis ν or $\bar{\nu}$ data slightly favor deviation from maximal $\theta_{\mu\mu}$ mixing, but without any preference for a given octant;
- matter effects (subleading) affect $\theta_{\mu\mu}$ preferred value, but not its significance. However:
 - they have opposite sign for ν and $\bar{\nu}$ samples;
 - they impact **NOvA** (& **Minos**) more than **T2K**;
- hence, **overall** LBL show weak octant preference (**NO** $> 45^\circ$ & **IO** $< 45^\circ$ at $\Delta\chi^2 \sim 0.4$);
- maximal $\theta_{\mu\mu}$ mixing **disfavored** at $\Delta\chi^2 \sim 3.4$;
- since $\Delta m_{\mu\mu}^2$ depends on θ_{23} but Δm_{ee}^2 does not, **combining** LBL+REA impacts θ_{23} range.



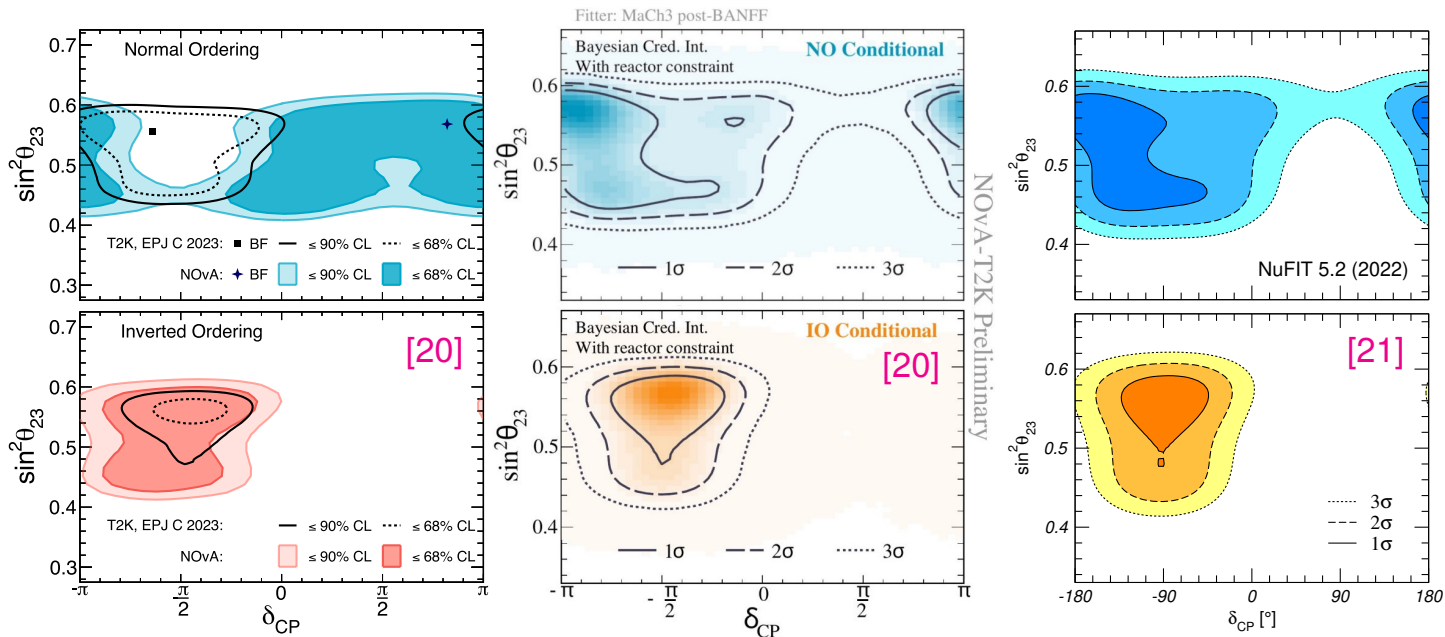
Global $\nu_e + \nu_\mu$ disappearance bounds

- Intrinsic contributions of ν_μ -disapp data sample:
 - uniquely determines the θ_{23} mixing angle;
 - sizably reduces the allowed range of Δm_{31}^2 ;
- LBL-disapp alone insensitive to: $\left\{ \begin{array}{l} \Delta m^2 \text{ ordering,} \\ \text{CP violation;} \end{array} \right.$
- but combination with ν_e -disapp (& atm-dis) yields:
 - preference for **NO**: 1.5σ (IC19) $\rightarrow$ 2.3σ (IC24);
 - non-maximal θ_{23} : 2.2σ (IC19) $\rightarrow$ 1.9σ (IC24);
 - a weak dependence on the δ_{CP} phase.



Tension between NOvA and T2K data: summer 2020

- Neutrino 2020: tension on δ_{CP} between T2K and NOvA for **NO** (no problem for **IO**);
- official joint T2K/NOvA analysis finally presented [20], results very similar to estimates [21].

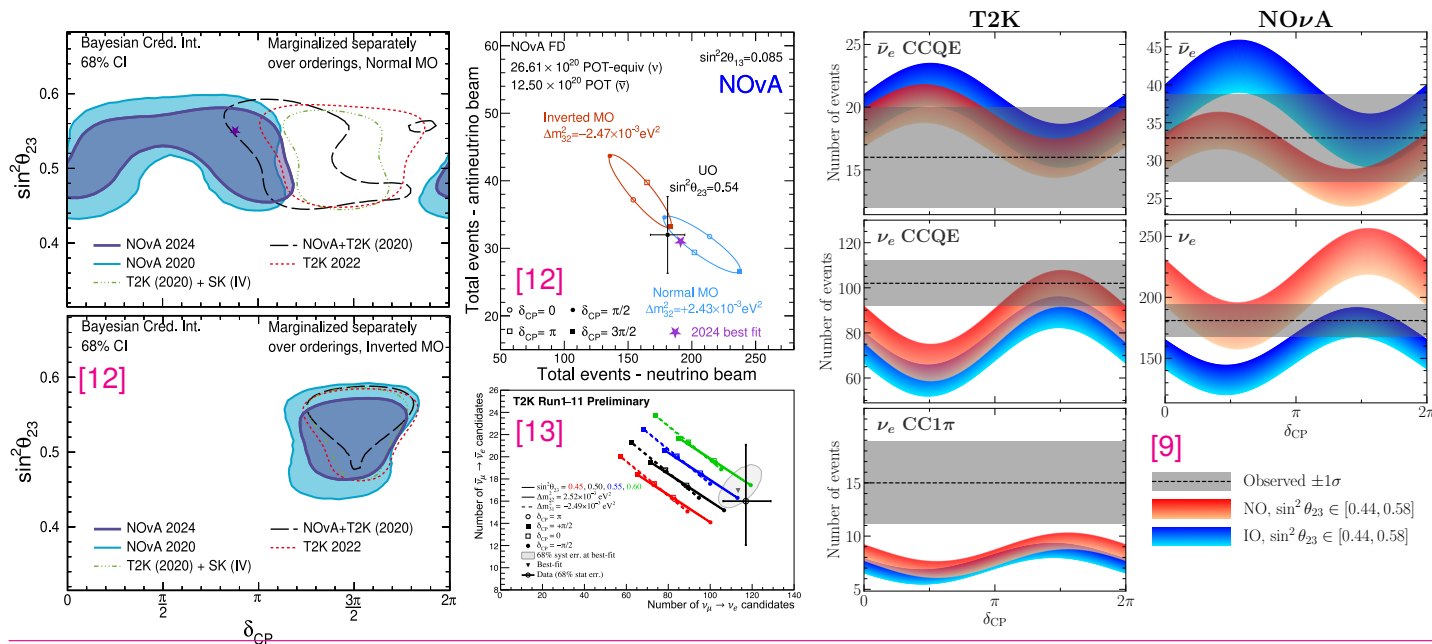


[20] M. Sanchez [NOvA], talk at Moriond-EW 2024, La Thuile, Italy, March 24–31, 2024.

[21] I. Esteban *et al.*, JHEP **09** (2020) 178 [arXiv:2007.14792].

Tension between NOvA and T2K data: today

- Neutrino 2024: NOvA substantially increased ν statistics, but no qualitative change on δ_{CP} .



[9] I. Esteban *et al.*, JHEP **12** (2024) 216 [arXiv:2410.05380] & NuFIT 6.0 [http://www.nu-fit.org].

[12] J. Wolcott [NOvA], talk at Neutrino 2024, Milan, Italy, June 16–22, 2024.

[13] D. Carabadjac [T2K], talk at ICHEP 2024, Prague, Czech Republic, July 17–24, 2024.

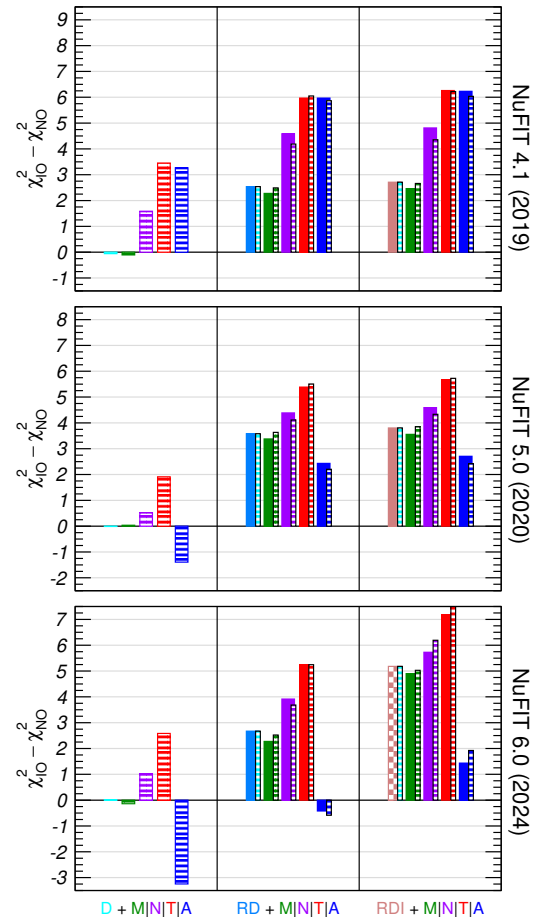
Impact on the neutrino mass ordering

- when taken by themselves, both **T2K** and **NOvA** appearance data exhibit a long-standing preference for **NO**;
- but the δ_{CP} tension among them (which has become more severe over time) implies that a **joint LBL** fit prefers **IO**;
- negative** $\Delta\chi^2_{\text{IO-NO}}$ contribution from ν_e appearance data is largely independent from the **positive** one stemming from ν_μ disappearance (both **w/o** and **with** IC atmos data):

$$\text{IO} - \text{NO}: \begin{cases} \Delta\chi^2_{\text{Rea+Dis+App}} \approx \Delta\chi^2_{\text{Rea+Dis}} + \Delta\chi^2_{\text{Dis+App}}, \\ \Delta\chi^2_{\text{Rea+Dis+IC+App}} \approx \Delta\chi^2_{\text{Rea+Dis+IC}} + \Delta\chi^2_{\text{Dis+App}} \end{cases}$$

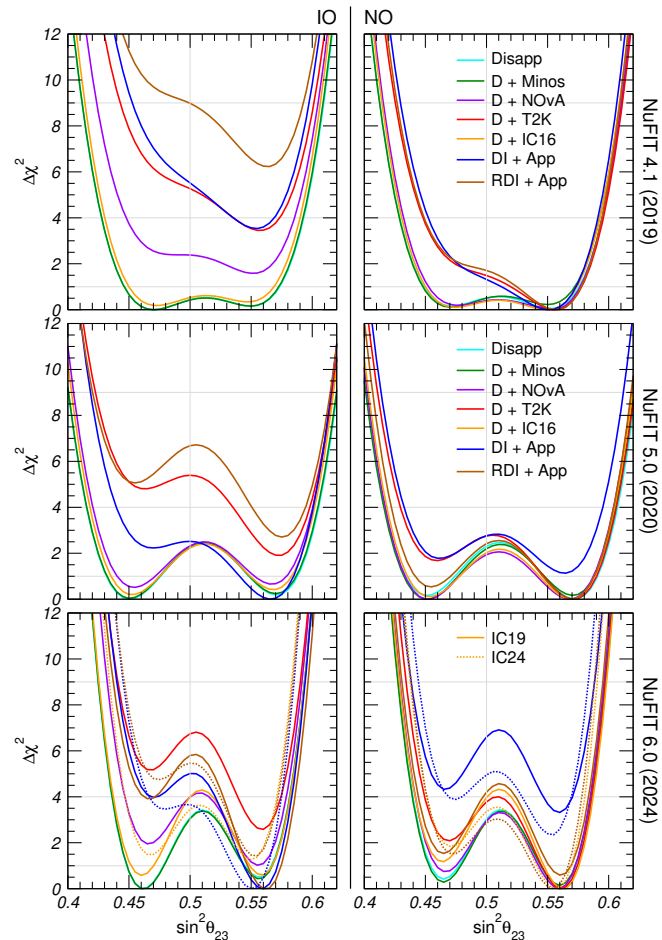
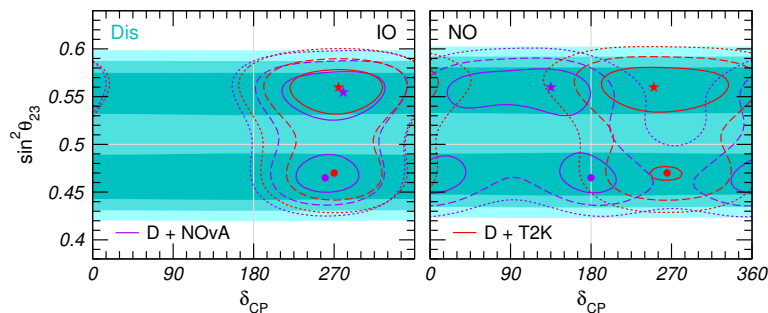
for $\text{App} \in \{\text{Minos}, \text{NOvA}, \text{T2K}, \text{All-LBL}\}$;

- these *contrasting preferences* between ν_e -app and ν_μ -dis mutually cancel, leading to no indication ($\lesssim 1\sigma$);
- inclusion of Super-K atmos. data (multiple ν_e -dis, ν_μ -dis, $\nu_\mu \leftrightarrow \nu_e$ channels) yields a $\Delta\chi^2 = 5.7$ push for **NO**.



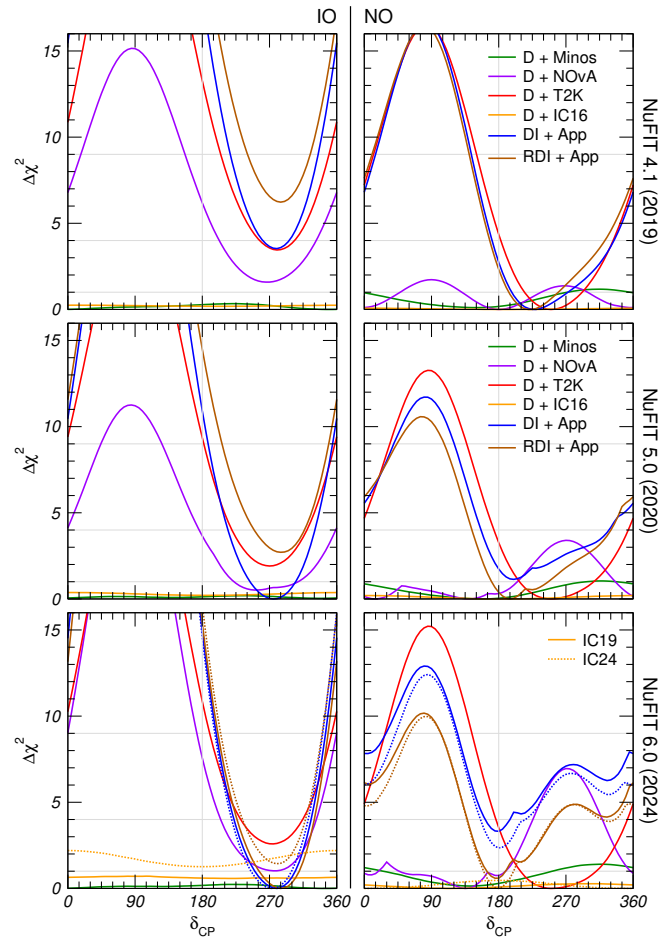
Impact of LBL-app on the θ_{23} octant

- **Minos** and **IC16** contribution marginal; **IC19/24** get some sensitivity from “cascade” (ν_e) events;
- **T2K** (more) and **NOvA** (less) both push for $\theta_{23} > 45^\circ$, irrespective of the mass ordering;
- combination of LBL-appearance data:
 - **IO**: preference for $\theta_{23} > 45^\circ$ is confirmed;
 - **NO**: the δ_{CP} tension between **T2K** and **NOvA** is less severe for $\theta_{23} < 45^\circ$, hence it cancels this hint and leads to similar minima.



Status of the CP phase

- **T2K** data show a clear preference for maximal CP violation ($\delta_{\text{CP}} \simeq 270^\circ$), irrespective of the assumed mass ordering;
- **NOvA** data also favor such value for **IO**, but for **NO** it disfavors it, preferring instead the CP conserving value $\delta_{\text{CP}} \simeq 180^\circ$;
- **NOvA** rejection of $\delta_{\text{CP}} \simeq 270^\circ$ has steadily increased over time: $1.2\sigma \rightarrow 1.8\sigma \rightarrow 2.6\sigma$ from the analysis of NuFIT 4.1 $\rightarrow$ 5.0 $\rightarrow$ 6.0 data;
- **Minos** & **IceCube** have practically no sensitivity to δ_{CP} and give negligible contribution;
- combined **LBL+IC** experiments indicate $\delta_{\text{CP}} \simeq \pi$ for **NO**, thus dominated by **NOvA**. Further inclusion of **reactors** does not change this picture.



Neutrino oscillations: where we are

- Global 6-parameter fit (including δ_{CP}):
 - Solar**: Cl + Ga + SK(1–4) + SNO-full (I+II+III) + BX(1–3);
 - Atmospheric**: IC19 | IC24 + SK(1–5);
 - Reactor**: KamLAND + SNOplus + IC + DB + Reno;
 - Accelerator**: Minos + T2K + NOvA;

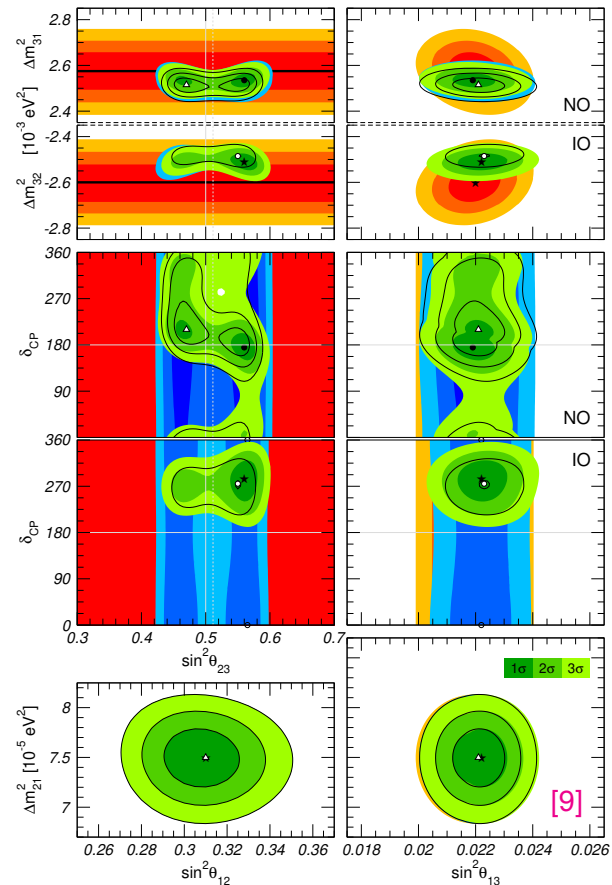
- best-fit point and 1σ (3σ) ranges:

$$\begin{aligned} \theta_{12} &= 33.68^{+0.73}_{-0.70} \left({}^{+2.27}_{-2.05} \right), & \Delta m_{21}^2 &= 7.49^{+0.19}_{-0.19} \left({}^{+0.56}_{-0.57} \right) \times 10^{-5} \text{ eV}^2, \\ \theta_{23} &= \begin{cases} 48.5^{+0.7}_{-0.9} \left({}^{+2.0}_{-7.6} \right), \\ 48.6^{+0.7}_{-0.9} \left({}^{+2.0}_{-7.2} \right), \end{cases} & \Delta m_{31}^2 &= \begin{cases} +2.534^{+0.025}_{-0.023} \left({}^{+0.072}_{-0.071} \right) \times 10^{-3} \text{ eV}^2, \\ -2.510^{+0.024}_{-0.025} \left({}^{+0.072}_{-0.073} \right) \times 10^{-3} \text{ eV}^2, \end{cases} \\ \theta_{13} &= 8.58^{+0.11}_{-0.13} \left({}^{+0.33}_{-0.39} \right), & \delta_{\text{CP}} &= 285^{+25}_{-28} \left({}^{+129}_{-182} \right); \end{aligned}$$

- neutrino mixing matrix:

$$|U|_{3\sigma} = \begin{pmatrix} 0.801 \rightarrow 0.842 & 0.519 \rightarrow 0.580 & 0.142 \rightarrow 0.155 \\ 0.248 \rightarrow 0.505 & 0.473 \rightarrow 0.682 & 0.649 \rightarrow 0.764 \\ 0.270 \rightarrow 0.521 & 0.483 \rightarrow 0.690 & 0.628 \rightarrow 0.746 \end{pmatrix};$$

- ordering: $\Delta\chi^2_{\text{IO-NO}} = -0.6$ (IC19) | $+6.1$ (IC24 + SK).



[9] I. Esteban *et al.*, JHEP 12 (2024) 216 [arXiv:2410.05380] & NuFIT 6.0 [http://www.nu-fit.org].

- Most of the present data from **solar**, **atmospheric**, **reactor** and **accelerator** experiments are well explained by the 3ν oscillation hypothesis. The three-neutrino scenario is robust;
- the long-standing “hints” concerning the **mass ordering**, with **NO** favored over **IO** at the $2\sigma \div 3\sigma$ level, are cancelled by the T2K/NOvA tension;
- the discovery of large θ_{13} opened the road to searches for **CP violation**. However, results on this topic need further clarifications;
- deviation from **maximal θ_{23} mixing** is also still an open issue. The region $\theta_{23} > 45^\circ$ seems to be slightly preferred, especially for **IO**;
- synergies between different experiments will be crucial to increase the sensitivity.

