



The $U(1)_{L_\mu + L_\tau}$ Model Meets the $(g - 2)_\mu$ Data and Muon Neutrino Trident Scattering Again

arXiv: 2506.05511

In collaboration with **Xiao-Gang He, Andrew Cheek and Xinhui Chu**

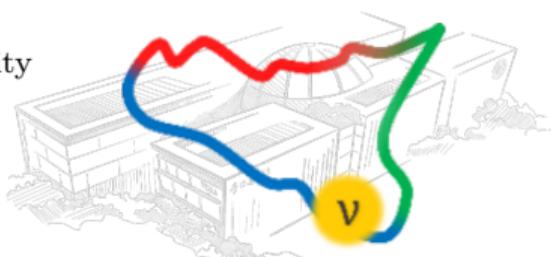
Speaker: **Ming-Wei Li**

limw2021@sjtu.edu.cn



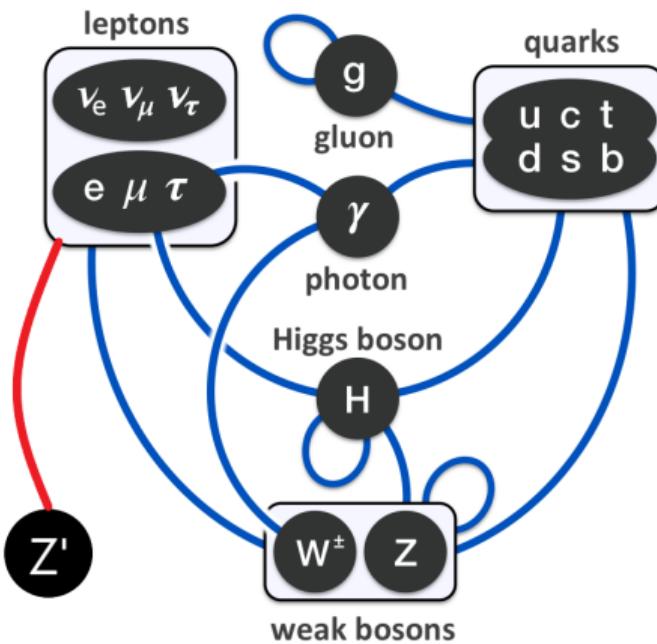
T.D. Lee Institute, Shanghai Jiao Tong University

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$U(1)$: The Simplest Extensions to the SM

Figure 1: A possible fifth fundamental force.



One of the simplest extensions is the $U(1)_{L_\mu - L_\tau}$ model, (He et al., 1990)

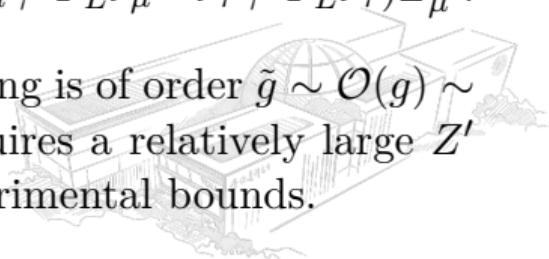
$$\begin{aligned}\mathcal{L}_{Z'} = & -\frac{1}{4} (Z')_{\mu\nu} (Z')^{\mu\nu} + \frac{1}{2} m_{Z'}^2 Z'_\mu Z'^\mu \\ & - \tilde{g} (\bar{L}_\mu \gamma^\mu L_\mu - \bar{L}_\tau \gamma^\mu L_\tau + \bar{\mu}_R \gamma^\mu \mu_R - \bar{\tau}_R \gamma^\mu \tau_R) Z'_\mu.\end{aligned}$$

where $L_\alpha = (\nu_\alpha, \ell_\alpha)^T$ is the $SU(2)_L$ doublet.

The new interaction terms:

$$-\tilde{g} (\bar{\mu} \gamma^\mu \mu - \bar{\tau} \gamma^\mu \tau + \bar{\nu}_\mu \gamma^\mu P_L \nu_\mu - \bar{\nu}_\tau \gamma^\mu P_L \nu_\tau) Z'_\mu .$$

Typically, the coupling is of order $\tilde{g} \sim \mathcal{O}(g) \sim \mathcal{O}(10^{-1})$, which requires a relatively large Z' mass to satisfy experimental bounds.



Non-Standard Interaction Z' Model



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$$-\tilde{g}(\bar{\mu}\gamma^\mu\mu - \bar{\tau}\gamma^\mu\tau + \bar{\nu}_\mu\gamma^\mu P_L\nu_\mu - \bar{\nu}_\tau\gamma^\mu P_L\nu_\tau)Z'_\mu.$$

$U(1)_{L_\mu-L_\tau}$ model involves only the μ and τ flavors,

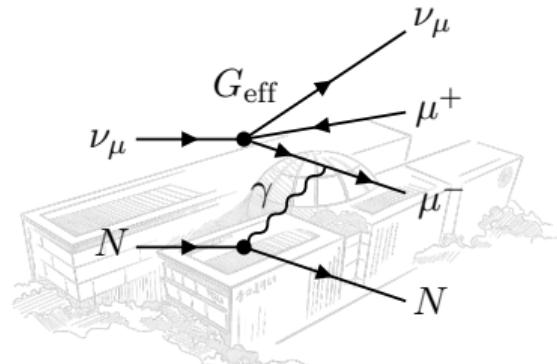
① Muon magnetic moment

$$\Rightarrow \Delta a_\mu^{Z'} = \frac{\tilde{g}^2}{8\pi^2} \frac{m_\mu^2}{m_{Z'}^2} \int_0^1 \frac{2x^2(1-x)dx}{1-x+(m_\mu^2/m_{Z'}^2)x^2} \approx \frac{\tilde{g}^2}{12\pi^2} \frac{m_\mu^2}{m_{Z'}^2} + \mathcal{O}\left(\frac{m_\mu^4}{m_{Z'}^4}\right)$$

② Muon neutrino trident scattering

$$\mathcal{L}_{\text{SM}}^{\text{eff}} = -\frac{g}{8m_W^2} (\bar{\nu}_\mu \gamma^\mu P_L \nu_\mu) (\bar{\mu} \gamma_\mu (1 + 4 \sin^2 \theta_W - \gamma^5) \mu)$$

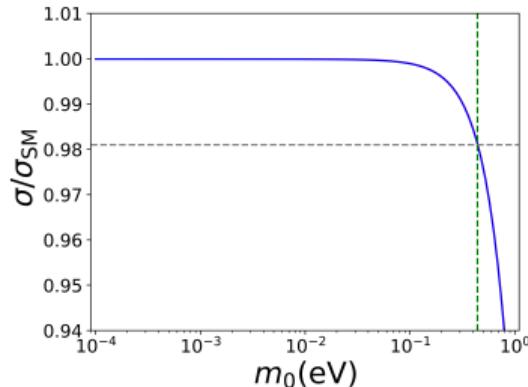
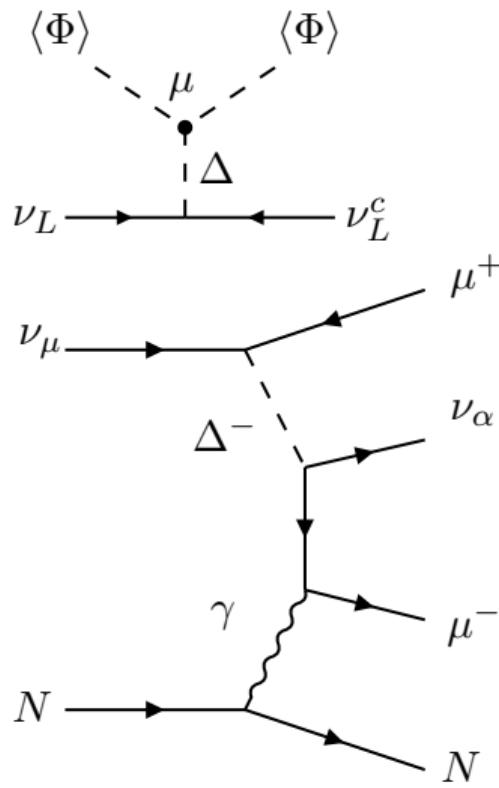
$$\mathcal{L}_{Z'}^{\text{eff}} = -\frac{\tilde{g}^2}{m_{Z'}^2} (\bar{\nu}_\mu \gamma^\mu P_L \nu_\mu) (\bar{\mu} \gamma_\mu \mu)$$



Muon Neutrino Trident Scattering and BSM



arXiv: 2204.05031, Phys.Lett.B 831 (2022) 137218.



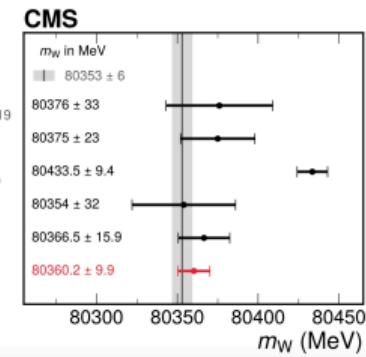
$$\rho \equiv \frac{m_W^2}{m_Z^2 \cos^2 \theta_W}, \quad \rho_{SM} = 1$$

$$\rho_\Delta = 1 - \frac{2v_\Delta^2}{v^2 + 4v_\Delta^2} < 0$$

$$\rho_{CDF} = 1.0019 > 0$$

arXiv: 2412.13872
(Also Commented by
Francesco Vissani)

Electroweak fit
PRD 110 (2024) 030001
LEP combination
Phys. Rep. 532 (2013) 119
D0
PRL 108 (2012) 151804
CDF
Science 376 (2022) 6589
LHCb
JHEP 01 (2022) 036
ATLAS
arXiv:2403.15085
CMS
This work



Constraints from $(g - 2)_\mu$



In the heavy mass limits $\Delta a_\mu^{Z'} \approx \frac{\tilde{g}^2}{12\pi^2} \frac{m_\mu^2}{m_{Z'}^2}$, (arXiv: 2104.03281, 2506.03069)

$$\begin{cases} \Delta a_\mu^{2021} = (251 \pm 59) \times 10^{-11} \\ \Rightarrow \frac{\tilde{g}^2}{m_{Z'}^2} = (2.66 \pm 0.63) \times 10^{-5} \text{ GeV}^{-2} \end{cases}, \quad \begin{cases} \Delta a_\mu^{2025} = (39 \pm 64) \times 10^{-11} \\ \Rightarrow \frac{\tilde{g}^2}{m_{Z'}^2} = (4.14 \pm 6.79) \times 10^{-6} \text{ GeV}^{-2} \end{cases}$$

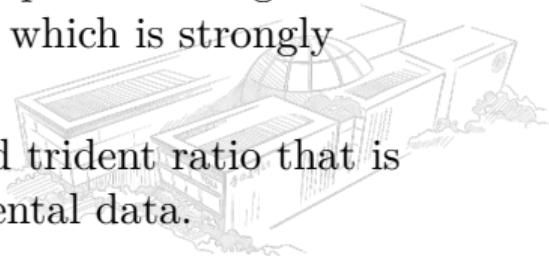
$$R_{\text{trident}} \equiv \frac{\sigma_{\text{SM}+Z'}}{\sigma_{\text{SM}}} = \frac{1 + \left(1 + 4 \sin^2 \theta_W + \frac{\sqrt{2} G_{Z'}}{G_F}\right)^2}{1 + \left(1 + 4 \sin^2 \theta_W\right)^2}, \quad \text{where } G_F \equiv \frac{\sqrt{2}}{8} \frac{g^2}{m_W^2}, \quad G_{Z'} \equiv \frac{\tilde{g}^2}{m_{Z'}^2}$$

$$\Rightarrow R_{\text{trident}}^{2021} = 5.94 \pm 1.08, \quad R_{\text{trident}}^{2025} = 1.47 \pm 0.87. \quad G_F = 1.166 \times 10^{-5} \text{ GeV}^{-2}$$

$$\begin{cases} \frac{\sigma_{\text{CHARM-II}}}{\sigma_{\text{SM}}} = 1.58 \pm 0.57 \\ \frac{\sigma_{\text{CCFR}}}{\sigma_{\text{SM}}} = 0.82 \pm 0.28 \\ \frac{\sigma_{\text{NuTeV}}}{\sigma_{\text{SM}}} = 0.72^{+1.73}_{-0.72} \end{cases}$$

$$\Rightarrow \sigma_{\text{exp}}/\sigma_{\text{SM}} = 0.96 \pm 0.25,$$

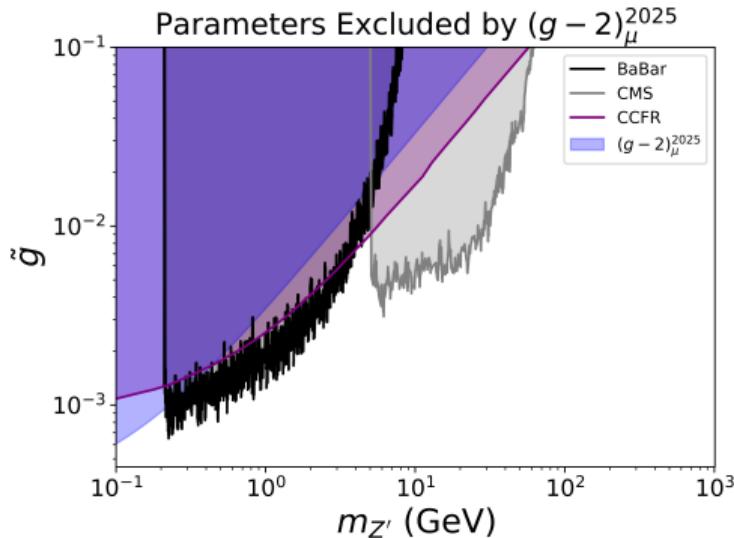
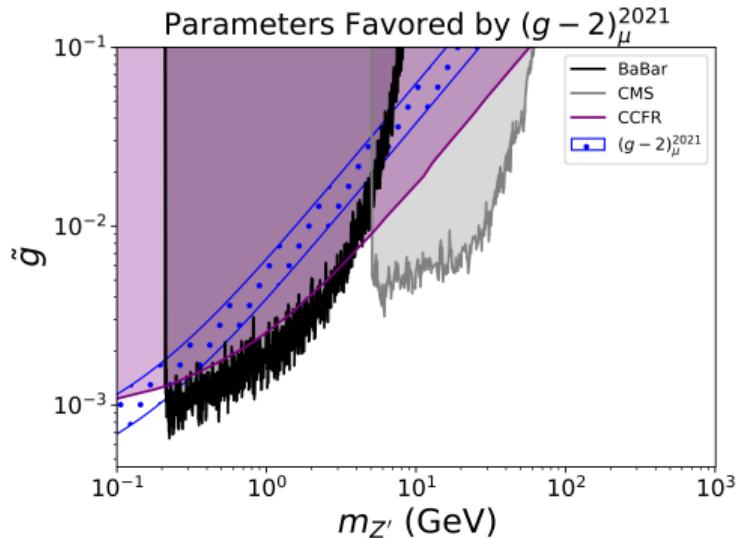
- ▶ From Δa_μ^{2021} , $U(1)_{L_\mu - L_\tau}$ predicts a significant enhancement in R_{trident} , which is strongly disfavored.
- ▶ Δa_μ^{2025} yields a predicted trident ratio that is consistent with experimental data.



Constraints from $(g - 2)_\mu$



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The new $(g - 2)_\mu$ data reopens the parameter space for large $m_{Z'}$.



Muon Collider Signatures

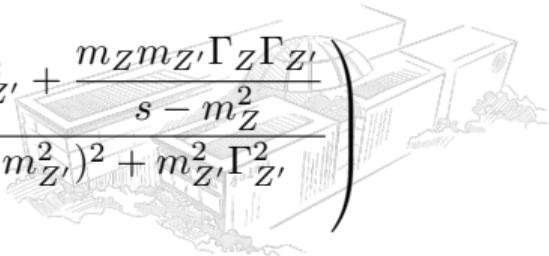


For the $U(1)_{L_\mu - L_\tau}$ model, $\mu^- \mu^+ \rightarrow \tau^- \tau^+$ will be the most consequential,

$$\begin{aligned} \mathcal{M} = & \left(\frac{4\pi\alpha_{\text{em}}}{s} - \frac{G_{Z'} m_{Z'}^2}{q^2 - m_{Z'}^2 + im_{Z'}\Gamma_{Z'}} \right) (\bar{\mu}\gamma^\mu\mu)(\bar{\tau}\gamma_\mu\tau) \\ & + \frac{\sqrt{2}G_F m_Z^2}{q^2 - m_Z^2 + im_Z\Gamma_Z} (\bar{\mu}\gamma^\mu(g_V - g_A\gamma_5)\mu)(\bar{\tau}\gamma_\mu(g_V - g_A\gamma_5)\tau) \end{aligned}$$

where $g_V = -\frac{1}{2} + 2\sin^2\theta_W$ and $g_A = -\frac{1}{2}$ and $\Gamma_{Z'} = \frac{G_{Z'} m_{Z'}^3}{4\pi}$.

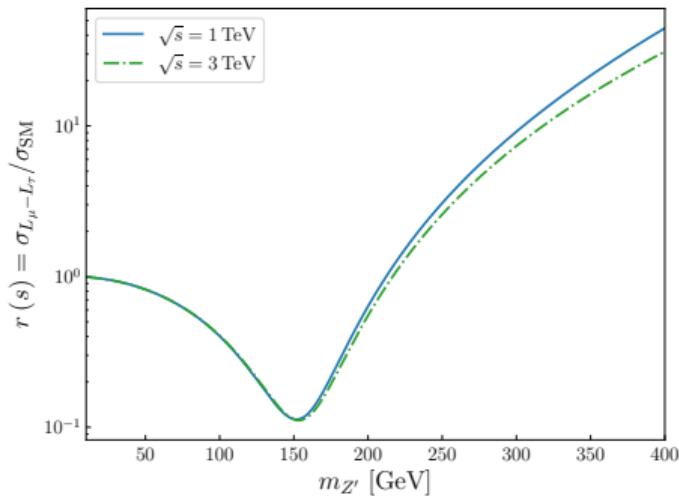
$$\begin{aligned} \sigma_{\text{SM}+Z'} = & \frac{4\pi\alpha_{\text{em}}^2}{3s} + \frac{G_{Z'} m_{Z'}^2 \left((G_{Z'} m_{Z'}^2 - 8\pi\alpha_{\text{em}})s + 8\pi\alpha_{\text{em}} m_{Z'}^2 \right)}{12\pi \left((s - m_{Z'}^2)^2 + m_{Z'}^2 \Gamma_{Z'}^2 \right)} + \frac{G_F^2 m_Z^4 s}{6\pi} \frac{(g_V^2 + g_A^2)^2}{(s - m_Z^2)^2 + m_Z^2 \Gamma_Z^2} \\ & + \frac{2\sqrt{2}\alpha_{\text{em}} G_F m_Z^2 g_V^2 (s - m_Z^2)}{3 \left((s - m_Z^2)^2 + m_Z^2 \Gamma_Z^2 \right)} \left(1 - \frac{G_{Z'} m_{Z'}^2 s}{4\pi\alpha_{\text{em}}} \frac{s - m_{Z'}^2 + \frac{m_Z m_{Z'} \Gamma_Z \Gamma_{Z'}}{s - m_Z^2}}{(s - m_{Z'}^2)^2 + m_{Z'}^2 \Gamma_{Z'}^2} \right) \end{aligned}$$



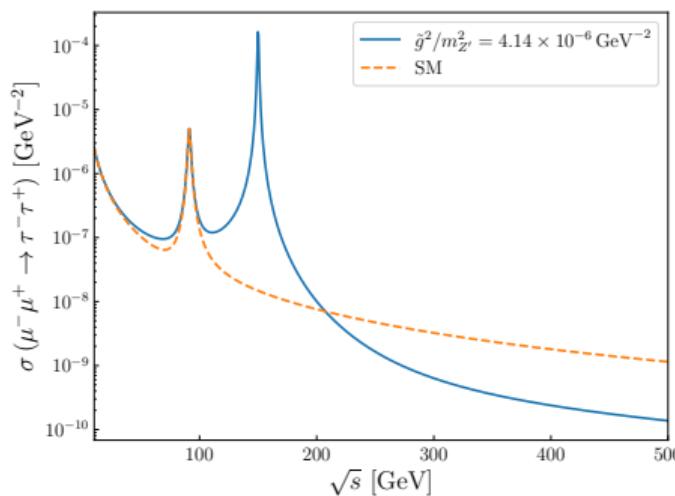
Muon Collider Signatures



Using the central value of Δa_μ^{2025} , ie. $G_{Z'}^{2025} \equiv \frac{\tilde{g}^2}{m_{Z'}^2} = 4.14 \times 10^{-6} \text{ GeV}^{-2}$,



(a) $\sigma_{\mu-\tau}^{\text{SM}+Z'}/\sigma_{\mu-\tau}^{\text{SM}}$.



(b) $\sigma_{\mu-\tau}$ for $m_{Z'} = 150 \text{ GeV}$.

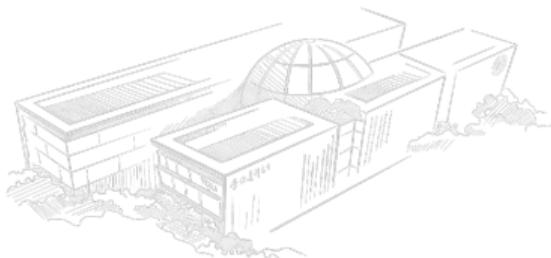


Summary



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- ① One of the simplest extensions, the $U(1)_{L_\mu-L_\tau}$ model, introduces non-standard interactions (NSIs) that couple only to the μ and τ flavors.
- ② It contributes to both $(g - 2)_\mu$ and muon neutrino trident (MNT) scattering. Recent results allow for heavier Z' masses, relaxing previous bounds.
- ③ A future muon collider could provide a promising opportunity to directly probe the $U(1)_{L_\mu-L_\tau}$ model.



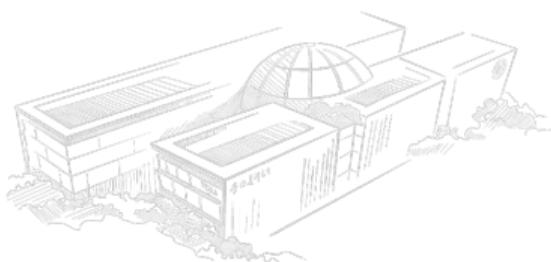
Thanks



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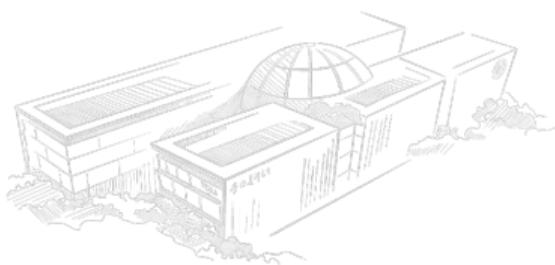
Thanks for
Your Attention



Backup



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