



Inelastic neutrino scattering on argon

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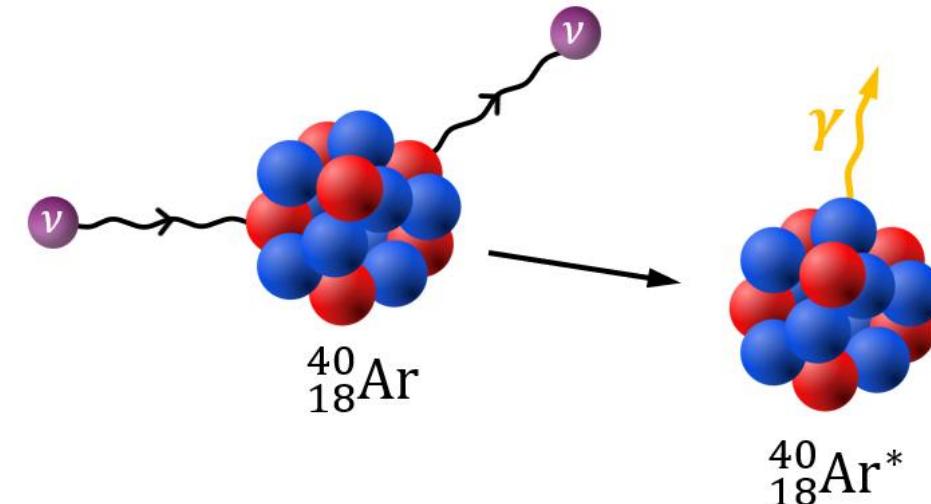
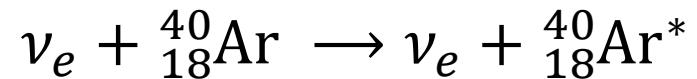
MAYORANA School
Modica, 19-25 June 2025



Neutrino interactions on ^{40}Ar at low energies ($\lesssim 100$ MeV)

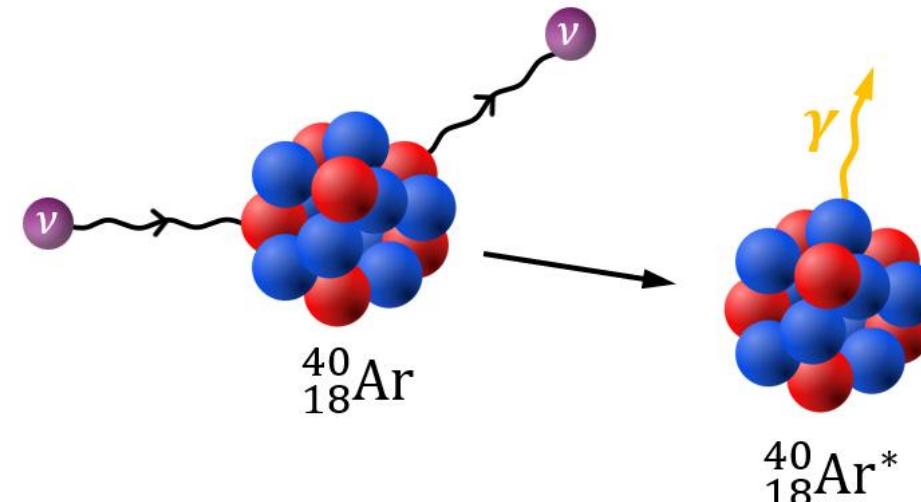
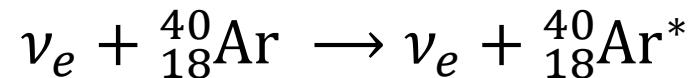
Neutrino interactions on $^{40}_{18}\text{Ar}$ at low energies ($\lesssim 100$ MeV)

Neutral Current (NC)

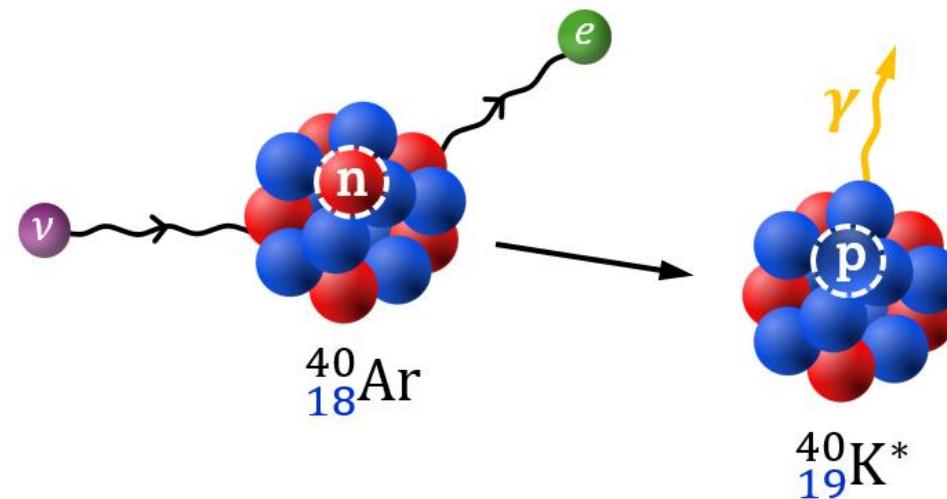
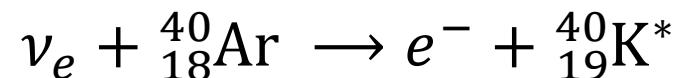


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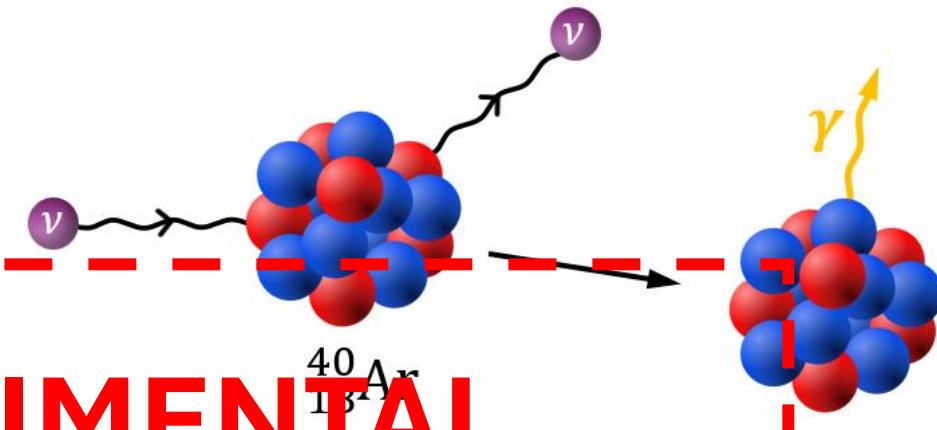
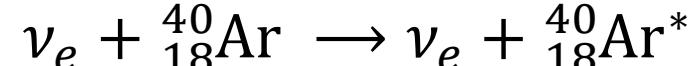


Charged Current (CC)



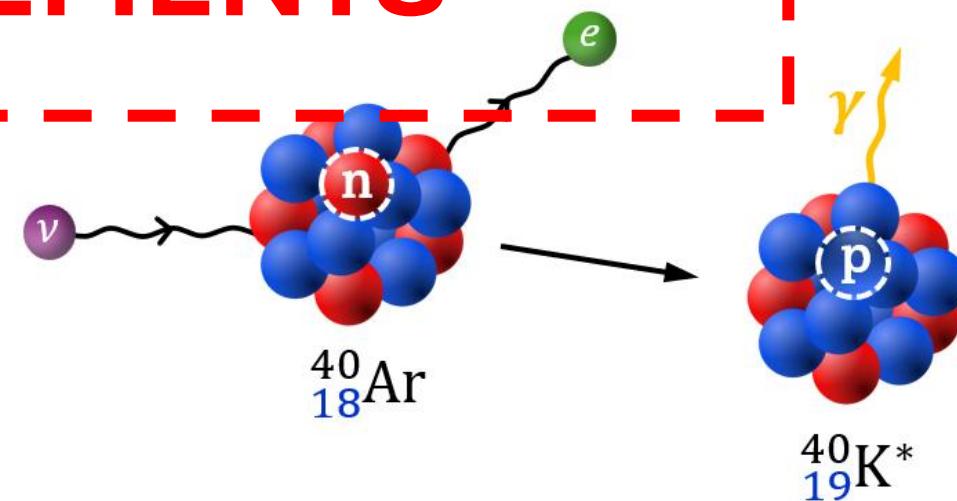
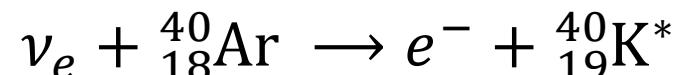
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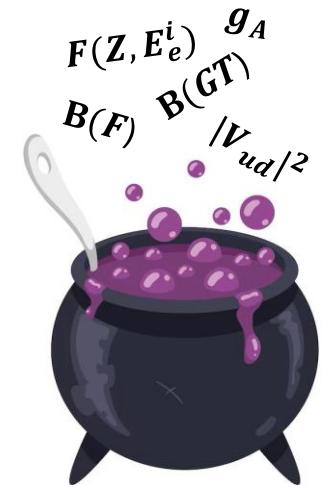


NO EXPERIMENTAL
MEASUREMENTS

Charged Current (CC)



Inelastic cross sections recipe



Inelastic cross sections recipe



NC

$$\sigma_{GT}^{(NC)} = \sum_i \frac{G_F^2 g_A^2}{\pi} (E_v - \omega)^2 B_i(GT)$$

Nuclear excitation energy

Nuclear transition probabilities

The equation shows the calculation of the inelastic cross section for a ground-state transition. It consists of a sum over states i of a term involving the Fermi coupling constant G_F^2 , the axial coupling constant g_A^2 , the ratio of nucleon form factors π , the nuclear excitation energy $(E_v - \omega)$, and the nuclear transition probability $B_i(GT)$. Two boxes with arrows point to the terms: one labeled 'Nuclear excitation energy' points to $(E_v - \omega)^2$, and another labeled 'Nuclear transition probabilities' points to $B_i(GT)$.

Inelastic cross sections recipe



NC

$$\sigma_{GT}^{(NC)} = \sum_i \frac{G_F^2 g_A^2}{\pi} (E_v - \omega)^2 \mathbf{B}_i(GT)$$

Nuclear excitation energy

Nuclear transition probabilities

CC

$$\sigma_{GT+F}^{(CC)} = \sum_i \frac{G_F^2 |V_{ud}|^2}{\pi} E_e^i p_e^i [\mathbf{B}_i(F) + \mathbf{B}_i(GT)] F(Z, E_e^i)$$

CKM matrix element

Fermi function

Quantifies the Coulomb interaction between the outgoing electron and protons of the ${}^{40}K$ nucleus

Inelastic cross sections recipe



NC

$$\sigma_{GT}^{(NC)} = \sum_i \frac{G_F^2 g_A^2}{\pi} (E_v - \omega)^2 B_i(GT)$$

Nuclear excitation energy

Nuclear transition probabilities

No direct experimental measurements
No reliable theoretical predictions



CC

$$\sigma_{GT+F}^{(CC)} = \sum_i \frac{G_F^2 |V_{ud}|^2}{\pi} E_e^i p_e^i [B_i(F) + B_i(GT)] F(Z, E_e^i)$$

CKM matrix element

Fermi function

Quantifies the Coulomb interaction between the outgoing electron and protons of the ${}^{40}K$ nucleus

HOW? Estimation of nuclear matrix elements

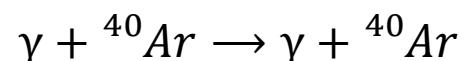


1 Theoretical approach

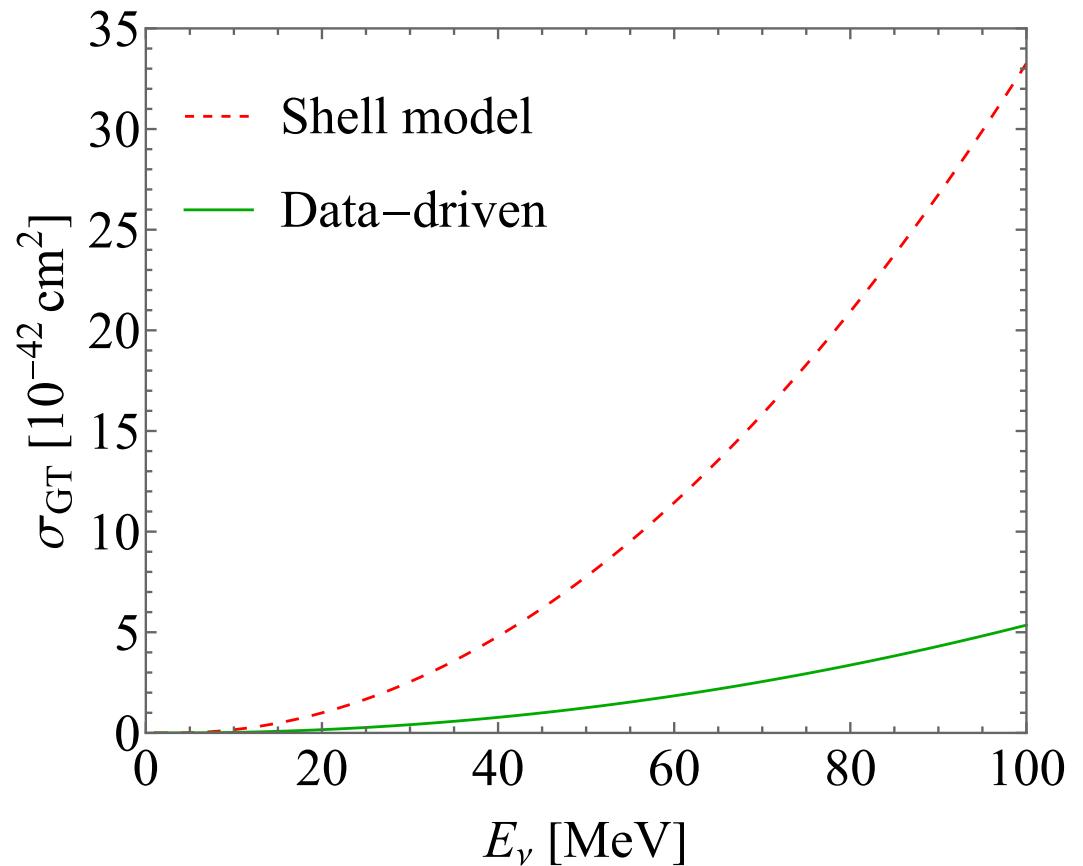
Transition probabilities are computed through BIGSTICK, a **nuclear shell model** code.

2 Data-driven approach

We consider **experimental measurements** of the magnetic dipole amplitudes $B(M1)$ for the process:



At low energies: $B(GT)^{exp} = \frac{B(M1)}{g_A^2(2.2993 \mu_N)^2}$



HOW? Estimation of nuclear matrix elements

CC

- 1 (p,n) scattering on ^{40}Ar : $p + {}^{40}\text{Ar} \rightarrow n + {}^{40}\text{K}^*$

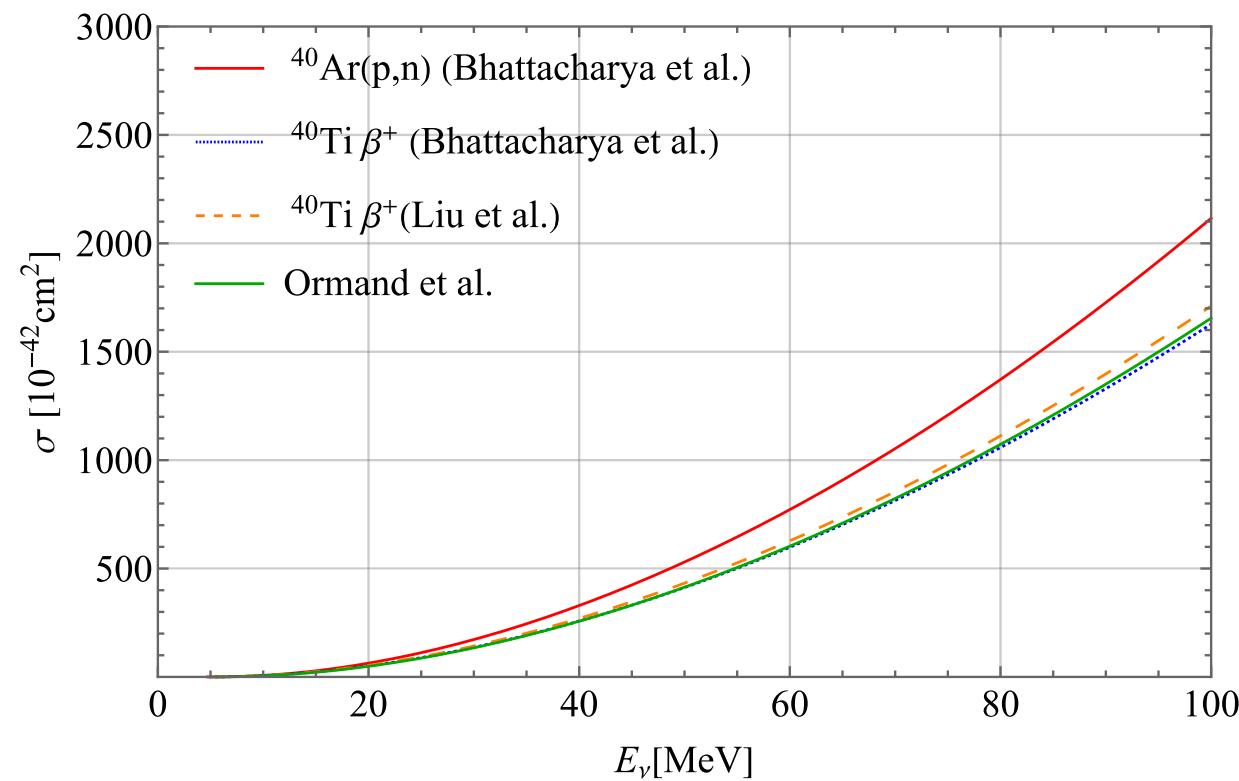
= Bhattacharya et al., Phys. Rev. C 80, 055501 (2009).

- 2 β decay of the mirror nucleus: ${}^{40}\text{Ti} \rightarrow {}^{40}\text{Sc}^* + e^+ + \nu_e$

= Bhattacharya et al., Phys. Rev. C 58, 36773687 (1998)
Liu et al., Phys. Rev. C 58, 26772688 (1998)

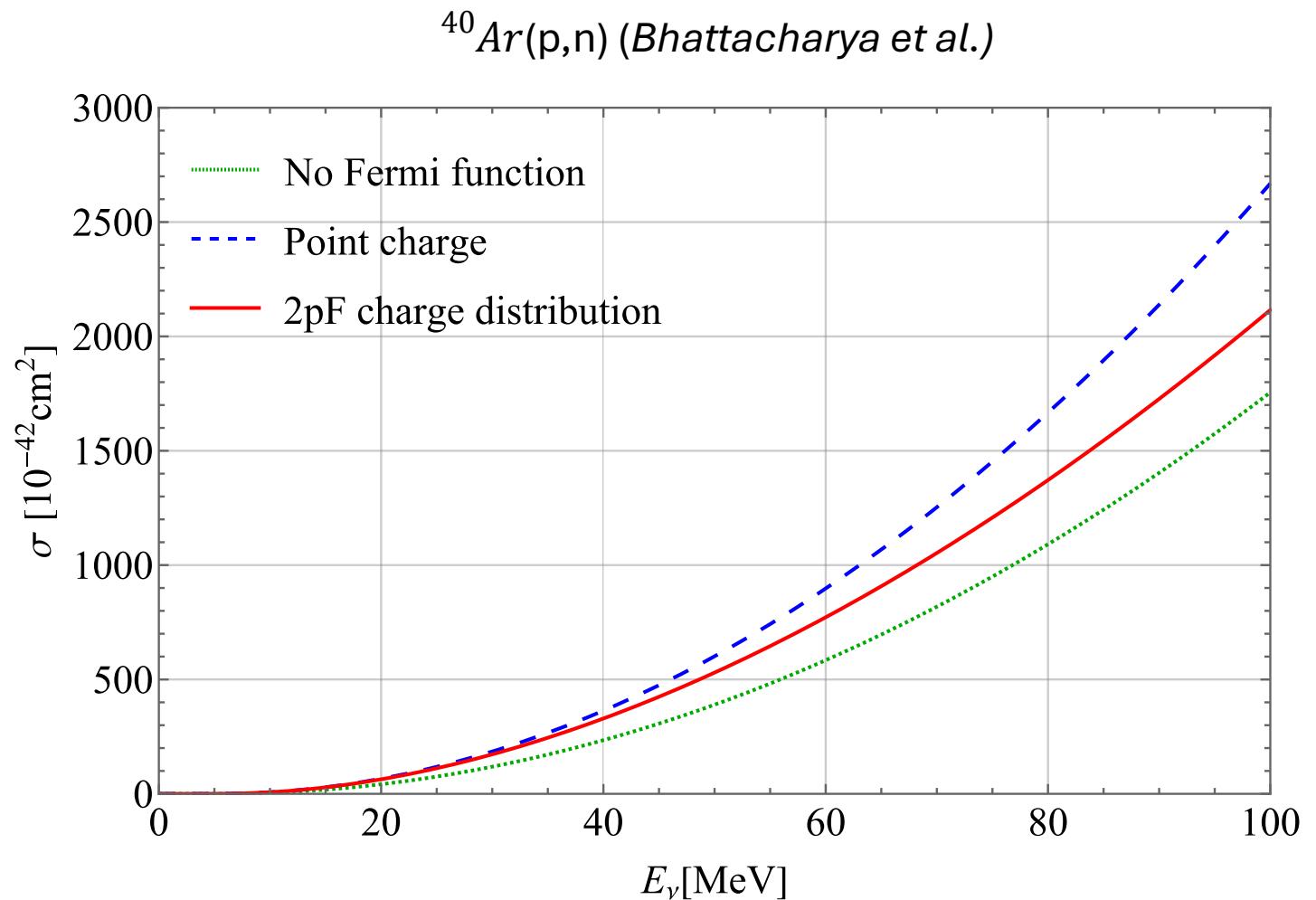
- 3 Theoretical shell model prediction

= W. E. Ormand et al., Phys. Lett. B345, 343 (1995)



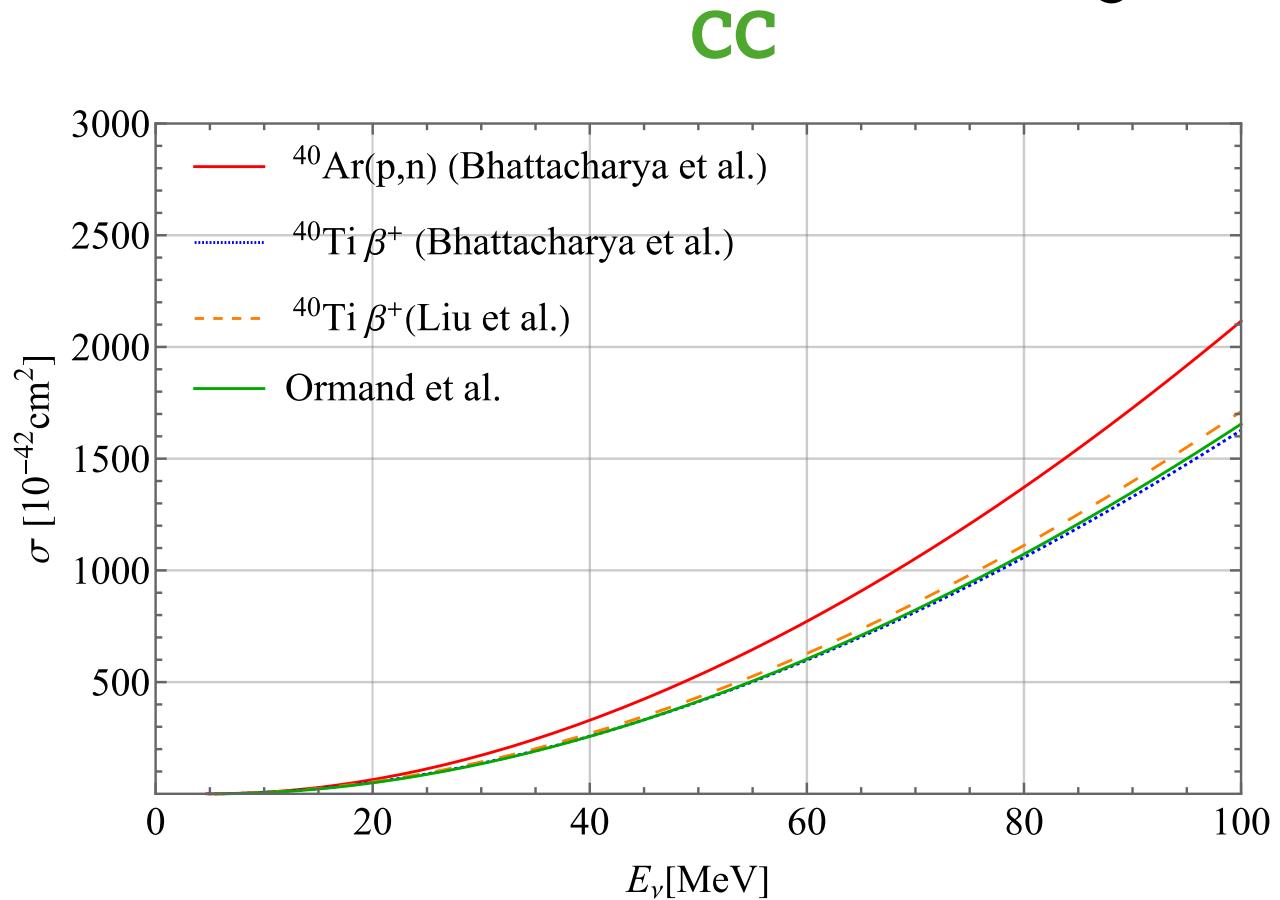
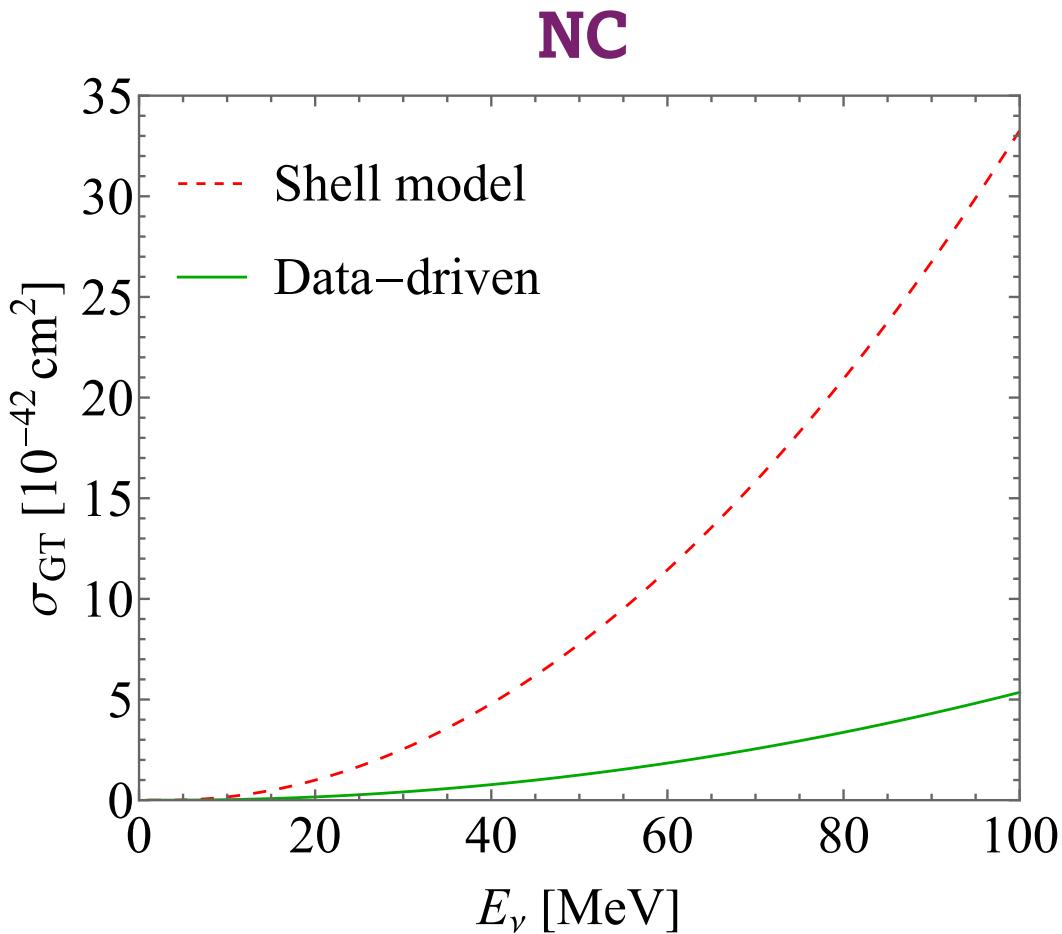
HOW? Estimation of nuclear matrix elements + Fermi function

CC



Summary: preliminary results

NC + CC





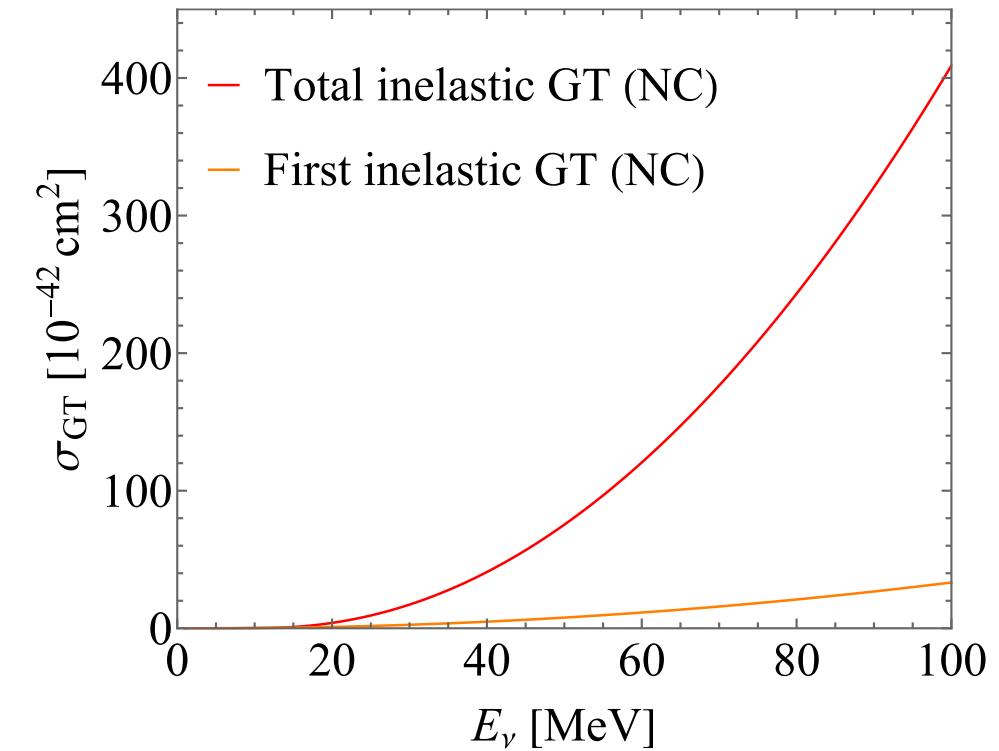
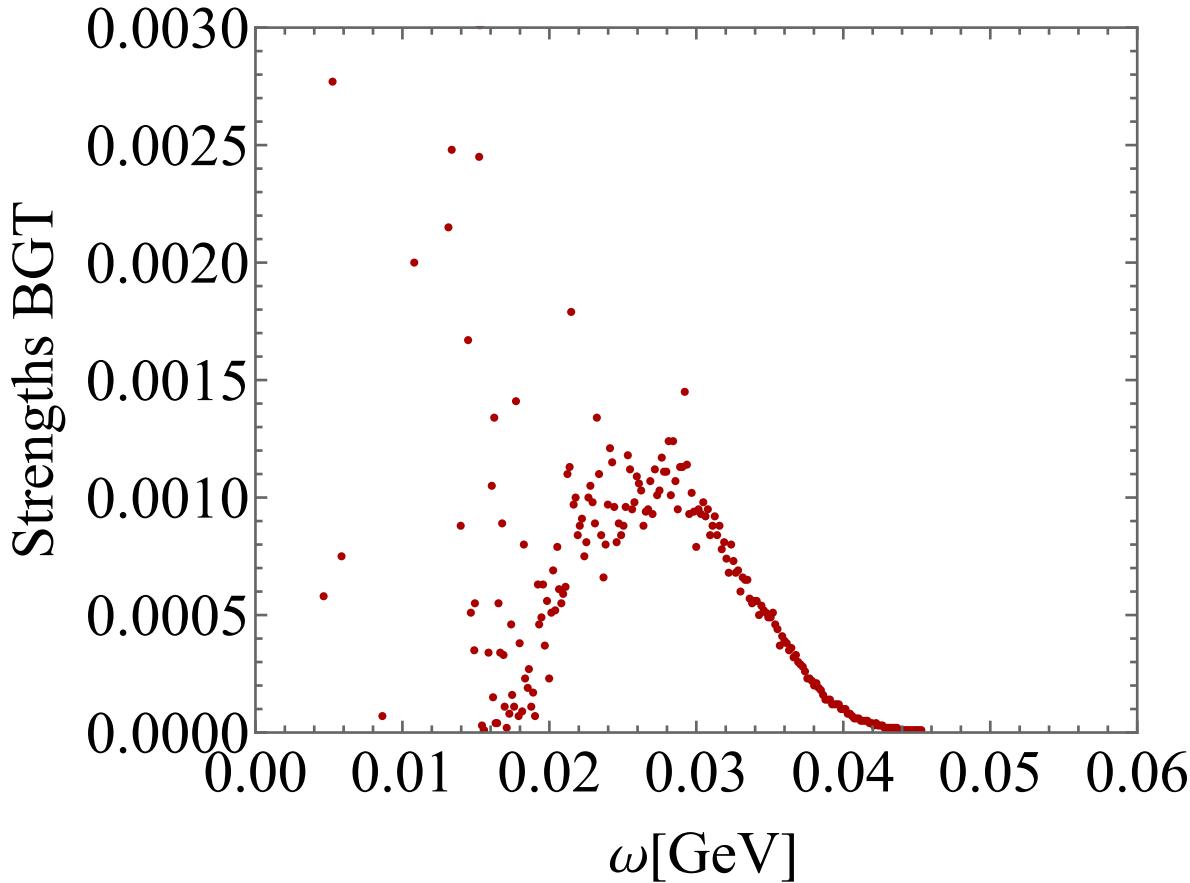
Thanks for your attention!

Estimating $B(GT)$: theoretical shell model prediction



C. W. Johnson et al. (2018), arXiv:1801.08432v1 [physics.comp-ph]

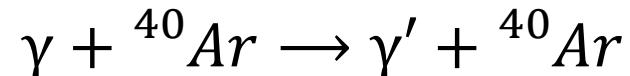
- 1 $B(GT)$ matrix elements are estimated through **BIGSTICK**, a nuclear shell model code used to compute transition probabilities.



E_ν [MeV]	$\sigma_{GT}[10^{-42}\text{cm}^2]$	$\sigma_{GT}^{\text{first}}[10^{-42}\text{cm}^2]$
20	3.9	1
50	75	7.8
80	243	21

Estimating B(GT): data-driven approach

- 2 We determine the B(GT) amplitudes from the experimental measurements of the magnetic dipole amplitudes **B(M1)**, conducted at the Triangle Universities Nuclear Laboratory (TUNL) for the process:



W. Tornov et al., Phys. Lett. B 835 (2022)

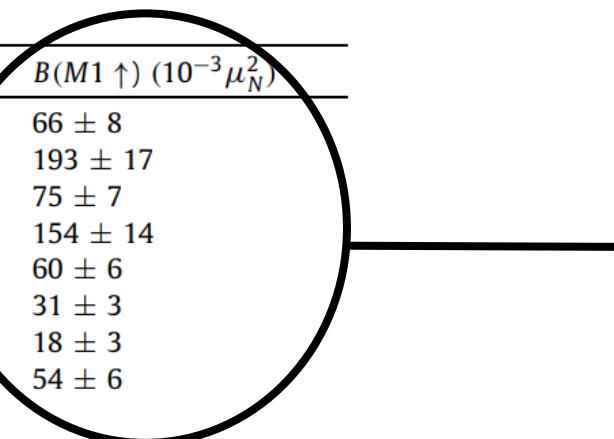
At low energies, these amplitudes are related to the BGT amplitudes through the relation:

$$B(GT)^{exp} = \frac{B(M1)}{g_A^2(2.2993 \mu_N)^2}$$



K. Langanke et al., Phys. Rev. Lett. 93 (2004)

ω (keV)	J^π	Γ_0 (meV)	$B(M1 \uparrow) (10^{-3} \mu_N^2)$
9697.5 \pm 1.4	1 ⁺	233 \pm 27	66 \pm 8
9758.3 \pm 1.1	1 ⁺	692 \pm 60	193 \pm 17
9805.6 \pm 1.3	1 ⁺	272 \pm 26	75 \pm 7
9841.3 \pm 1.3	1 ⁺	566 \pm 51	154 \pm 14
9871.7 \pm 1.2	1 ⁺	223 \pm 21	60 \pm 6
9893.9 \pm 1.4	1 ⁺	116 \pm 12	31 \pm 3
10020.5 \pm 1.7	1 ⁺	71 \pm 12	18 \pm 3
10033.9 \pm 1.4	1 ⁺	210 \pm 25	54 \pm 6



$$B(M1)_{tot} = 651(98) \times 10^{-3} \mu_N^2$$

Fermi function

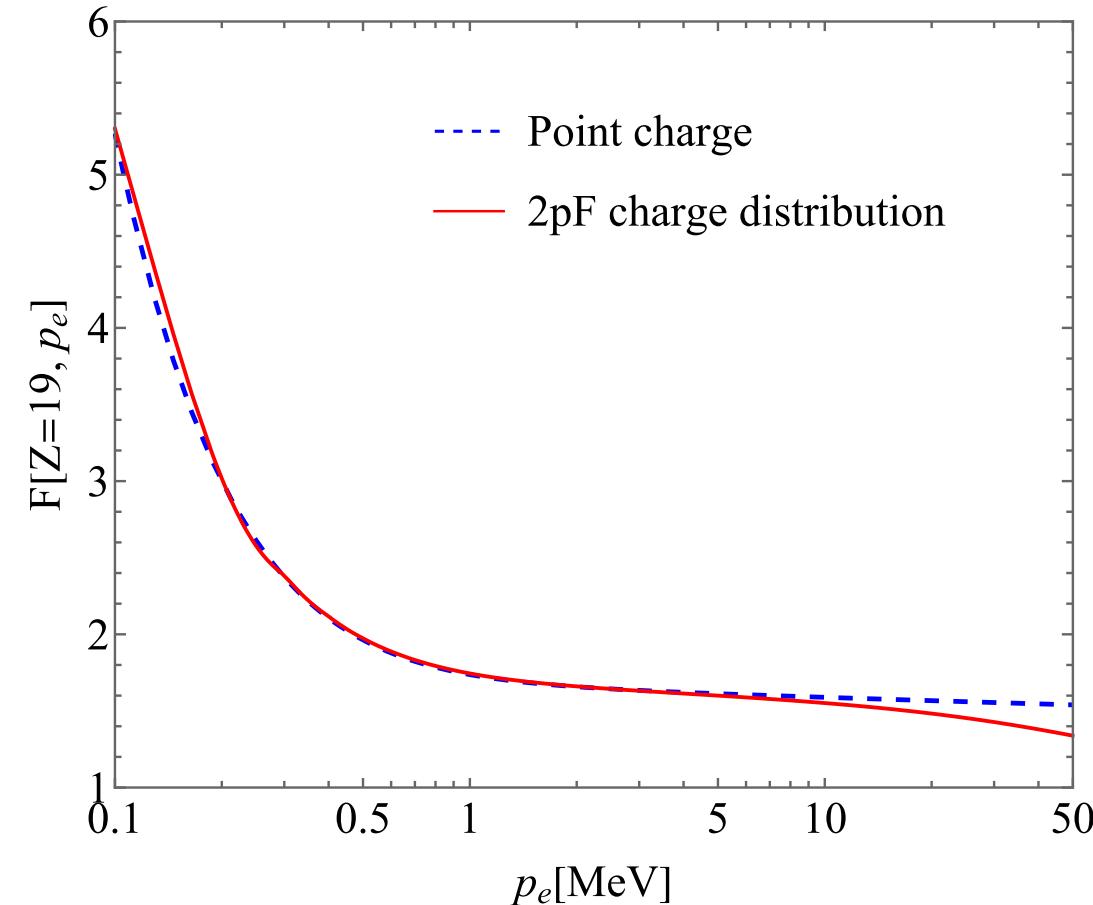
Fermi function: quantifies the Coulomb interaction between the outgoing electron and a proton of the ${}^{40}K$ nucleus

$$F(Z, p_e) = \frac{|\phi_e(0)_{coulomb}|^2}{|\phi_e(0)|^2}$$

→ Electron in a Coulomb field
→ Free electron

- 1 **Point charge:** analytical solution

- 2 **Charge distribution:** different nuclear charge distributions (uniform, 2pF, 3pF,...)
only numerical solution



Elastic and inelastic cross sections

