

# *Resonant $U_1$ leptoquark production @ LHC*

in collaboration with Peter Krack (NIKHEF) & Nudžeim Selimović (INFN Padova)

Eur. Phys. J. C 84, 304 (2024)

arXiv: 2311.13635

*Val di Luce, 2024*



Arman Korajac (INFN Pisa)

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Arman **Coraiazz** (INFN Pisa)

# OUTLINE

- **UV model for  $U_1$**
- **NLO QCD + QED results**
- **Simulation details & exclusion limits**
- **Conclusions**

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- **Conclusions**

# 4321 for U1

- Passing step in many UV flavor theories
- Decoupling the 3rd gen and the light fermions in the UV, effectively starting out with a  $U(2)^5$  flavor symmetry
- Motivated by solving the  $R(D^*)$  anomaly

$$\begin{array}{c}
 [SU(4) \times SU(2)_L \times SU(2)_R]^{[3]} \times [SU(3) \times SU(2)_L \times U(1)_{B-L} \times U(1)_R]^{[12]} \\
 \downarrow \langle \phi \rangle \\
 [SU(4) \times SU(2)_L \times SU(2)_R]^{[3]} \times [SU(3) \times SU(2)_L \times U(1)_Y]^{[12]} \\
 \begin{array}{c} \langle \Sigma_L \rangle \downarrow \langle \Sigma_R \rangle \end{array} \\
 SU(4)^{[3]} \times SU(3)^{[12]} \times SU(2)_L \times U(1)_X \\
 \begin{array}{c} \langle \Omega_3 \rangle \downarrow \langle \Omega_1 \rangle \end{array} \\
 SU(3)_c \times SU(2)_L \times U(1)_Y ,
 \end{array}$$

[Courtesy of N. Selimović]

[L. Di Luzio, A. Greljo and M. Nardecchia, arXiv: 1708.08450]

[L. Di Luzio, J. Fuentes-Martin, A. Greljo, M. Nardecchia, S. Renner, arXiv: 1808.00942]

[M. Bordone, C. Cornella, J. Fuentes-Martin, G. Isidori, arXiv: 1712.01368, 1805.09328]

[H. Georgi, Y. Nakai, arXiv: 1606.05865]

[Fuentes-Martin et al, arXiv: 1910.13474, 2006.16250, 2009.11296]

[D. Guadagnoli, M. Reboud, P. Stangl, arXiv: 2005.10117] ...

# 4321 for U1

- Using the minimal set-up, enough for gauge boson dynamics:

$$\mathcal{L} \supset -\frac{1}{4}H_{\mu\nu}^a H_{\mu\nu}^a + \sum_{i=1,3} (D_\mu \Omega_i)^\dagger (D_\mu \Omega_i) + \sum_{f=L,R^\pm} i\bar{\psi}_f \not{D}\psi_f + V(\Omega_i)$$

$$- Y_{\text{light,u}}^{ij} \bar{q}_L^i \tilde{H} u_R^j - Y_{\text{light,d}}^{ij} \bar{q}_L^i H d_R^j - y_{LR}^- \bar{\psi}_L H \psi_R^- - y_{LR}^+ \bar{\psi}_L \tilde{H} \psi_R^+ + \text{h.c.}$$

Field	$SU(4)_3$	$SU(3)_{12}$	$SU(2)_L$	$U(1)_X$
$\psi_L$	<b>4</b>	<b>1</b>	<b>2</b>	0
$\psi_R^+$	<b>4</b>	<b>1</b>	<b>1</b>	1/2
$\psi_R^-$	<b>4</b>	<b>1</b>	<b>1</b>	-1/2
$q_L^i$	<b>1</b>	<b>3</b>	<b>2</b>	1/6
$\ell_L^i$	<b>1</b>	<b>1</b>	<b>2</b>	-1/2
$u_R^i$	<b>1</b>	<b>3</b>	<b>1</b>	2/3
$d_R^i$	<b>1</b>	<b>3</b>	<b>1</b>	-1/3
$e_R^i$	<b>1</b>	<b>1</b>	<b>1</b>	-1
$H$	<b>1</b>	<b>1</b>	<b>2</b>	1/2
$\Omega_3$	$\bar{\mathbf{4}}$	<b>3</b>	<b>1</b>	1/6
$\Omega_1$	$\bar{\mathbf{4}}$	<b>1</b>	<b>1</b>	-1/2

$i = 1, 2$

$$\psi_L \equiv (q_L^3 \ell_L^3)^\top$$

$$\psi_R^- \equiv (d_R^3 e_R^3)^\top \quad \psi_R^+ \equiv (u_R^3 \nu_R^3)^\top$$

# 4321 for U1

- Using the minimal set-up, enough for gauge boson dynamics:

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$\ell_L^i$	<b>1</b>	<b>1</b>	<b>2</b>	-1/2
$u_R^i$	<b>1</b>	<b>3</b>	<b>1</b>	2/3
$d_R^i$	<b>1</b>	<b>3</b>	<b>1</b>	-1/3
$e_R^i$	<b>1</b>	<b>1</b>	<b>1</b>	-1
$H$	<b>1</b>	<b>1</b>	<b>2</b>	1/2
$\Omega_3$	$\bar{\mathbf{4}}$	<b>3</b>	<b>1</b>	1/6
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$$\psi_R^- \equiv (d_R^3 e_R^3)^\top \quad \psi_R^+ \equiv (u_R^3 \nu_R^3)^\top$$



# 4321 for U1

$$\mathcal{L} \supset \left[ -\frac{1}{4} H_{\mu\nu}^a H_{\mu\nu}^a + \sum_{i=1,3} (D_\mu \Omega_i)^\dagger (D_\mu \Omega_i) + \sum_{f=L,R^\pm} i \bar{\psi}_f \not{D} \psi_f + V(\Omega_i) \right]$$

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SSB :  $\mathcal{G}_{4321} \rightarrow \mathcal{G}_{\text{SM}}$

$$G' \sim (\mathbf{8}, \mathbf{1}, 0), \quad U_1 \sim (\mathbf{3}, \mathbf{1}, 2/3), \quad Z' \sim (\mathbf{1}, \mathbf{1}, 0)$$

**SM fields**

$$G_\mu^a = s_3 H_\mu^a + c_3 C_\mu^a, \quad B_\mu = s_1 H_\mu^{15} + c_1 X_\mu$$

**Massive gauge bosons**

$$G_\mu^{\prime a} = c_3 H_\mu^a - s_3 C_\mu^a, \quad Z'_\mu = c_1 H_\mu^{15} - s_1 X_\mu$$

$$U_\mu^{1,2,3} = \frac{1}{\sqrt{2}} (H_\mu^{9,11,13} - i H_\mu^{10,12,14})$$

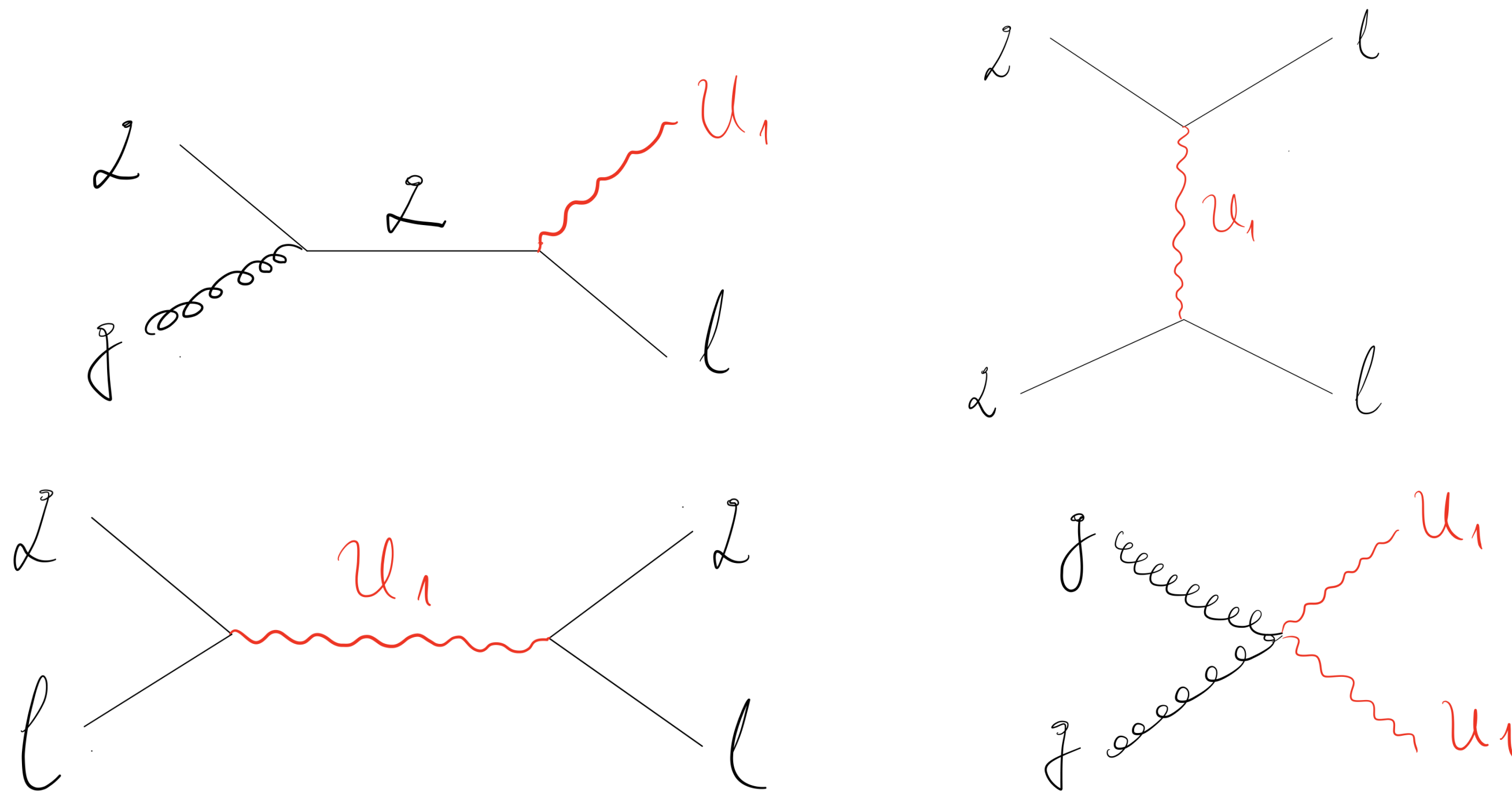


# 4321 for U1

$$\mathcal{L}_{U_1} \supset \frac{g_4}{\sqrt{2}} U_\mu (\beta_L^{33} \bar{q}_L^3 \gamma^\mu \tau_L + \beta_R^{33} \bar{b}_R \gamma^\mu \tau_R + \text{h.c.})$$

$$- ig_s U_\mu^\dagger T^a U_\nu G^{a,\mu\nu} - \frac{2}{3} ie U_\mu^\dagger U_\nu F^{\mu\nu}$$

Can be extended to couplings with light-gen fermions; additional fields required



- Signature can be seen through t-channel production, LQ + lepton or double LQ production
- **Goal: Examine the resonant s-channel production for U1**
- Already done, *however*:
  - no lepton PDFs at that moment
  - no experimental results at that time
  - @ LO.

[U. Haisch, G. Polesello, arXiv:2012.11474]

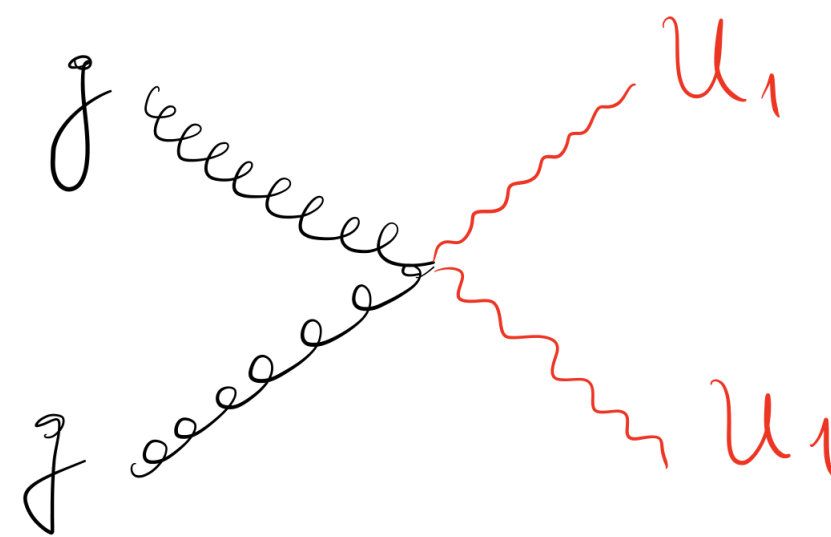
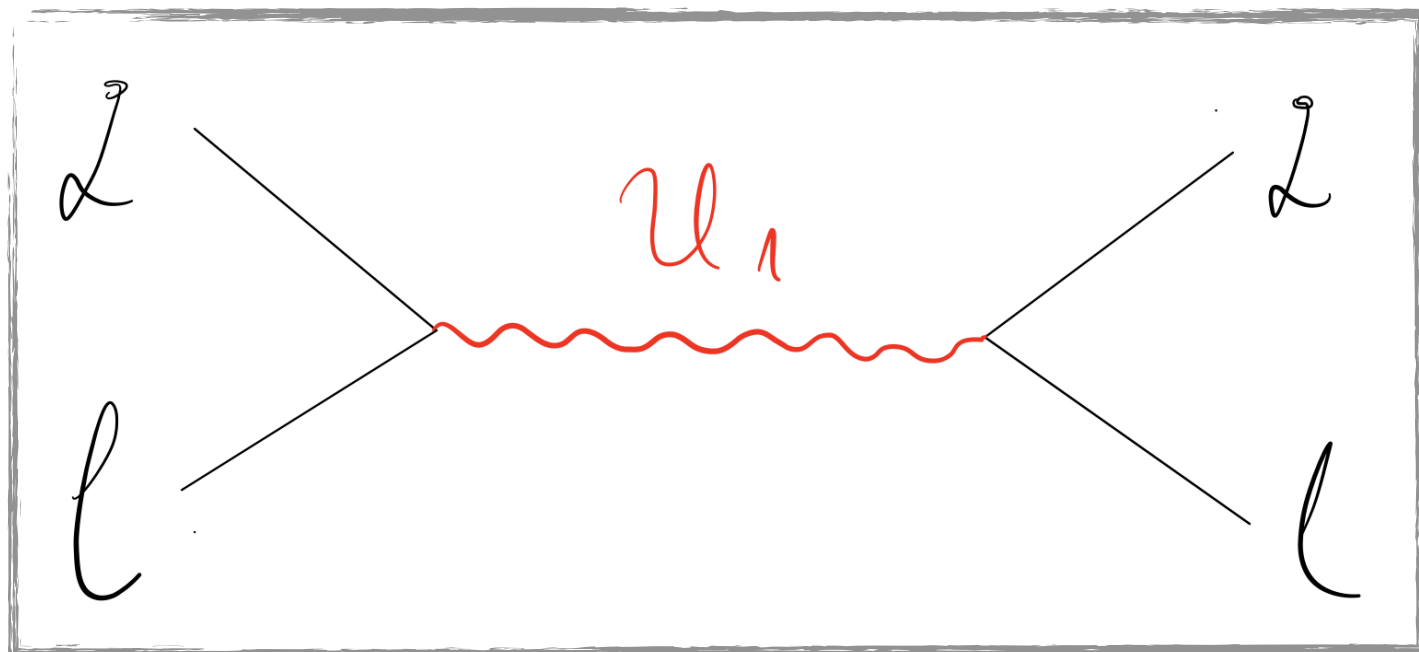
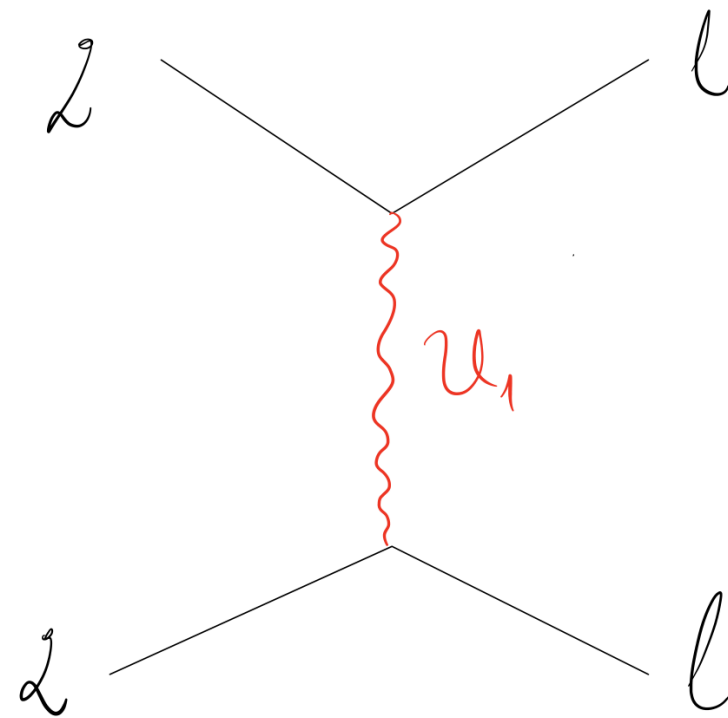
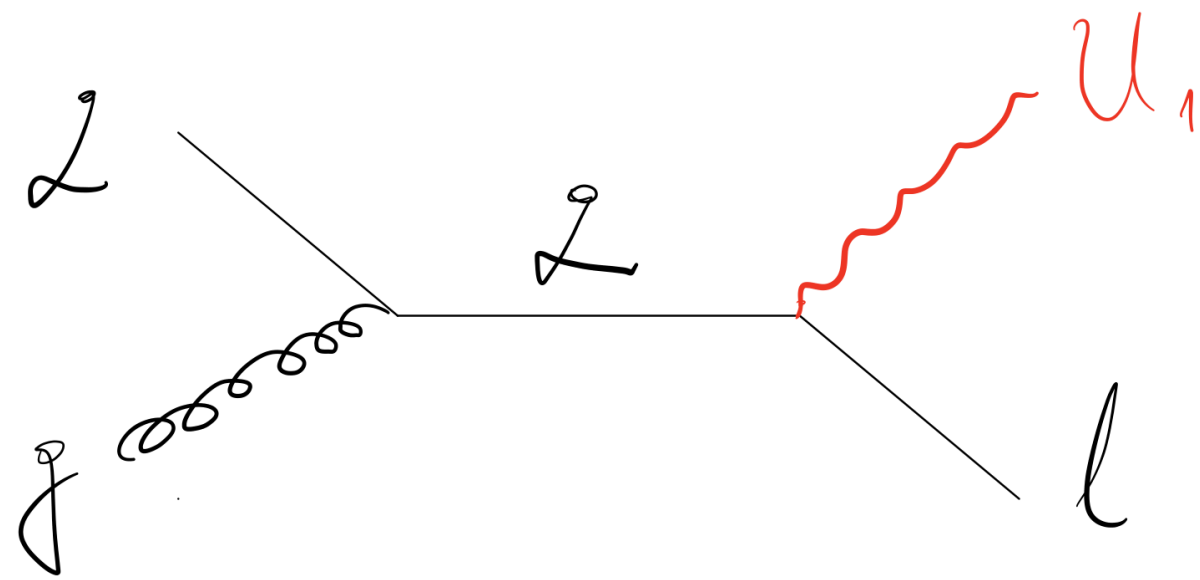
[Buonocore et al, arXiv: 2005.06475]

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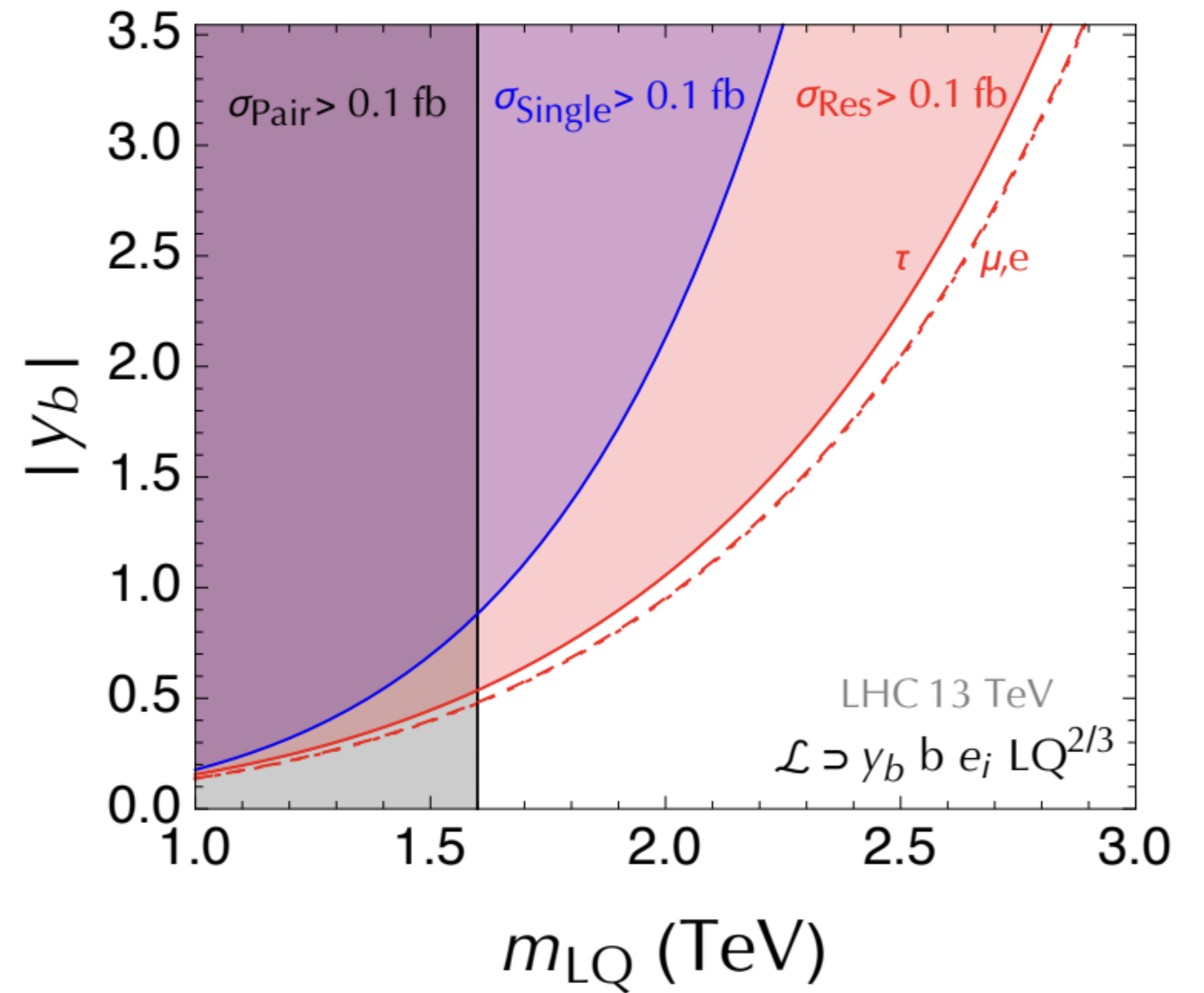
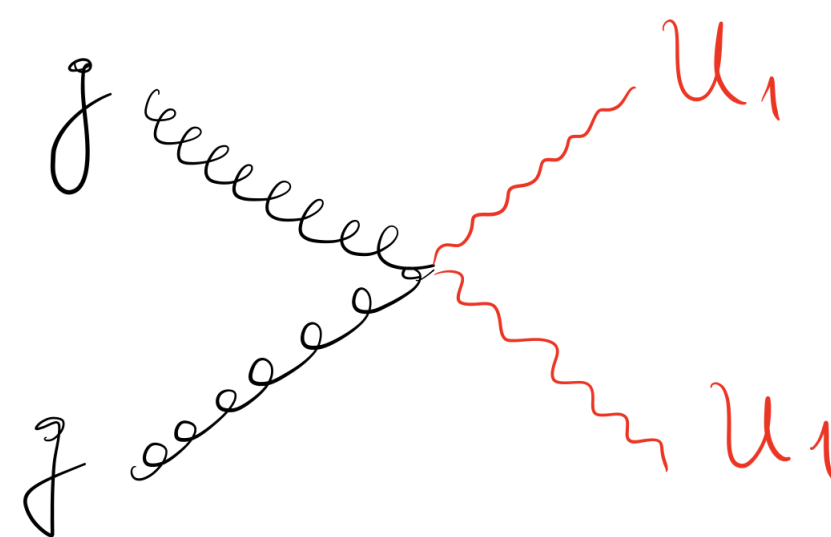
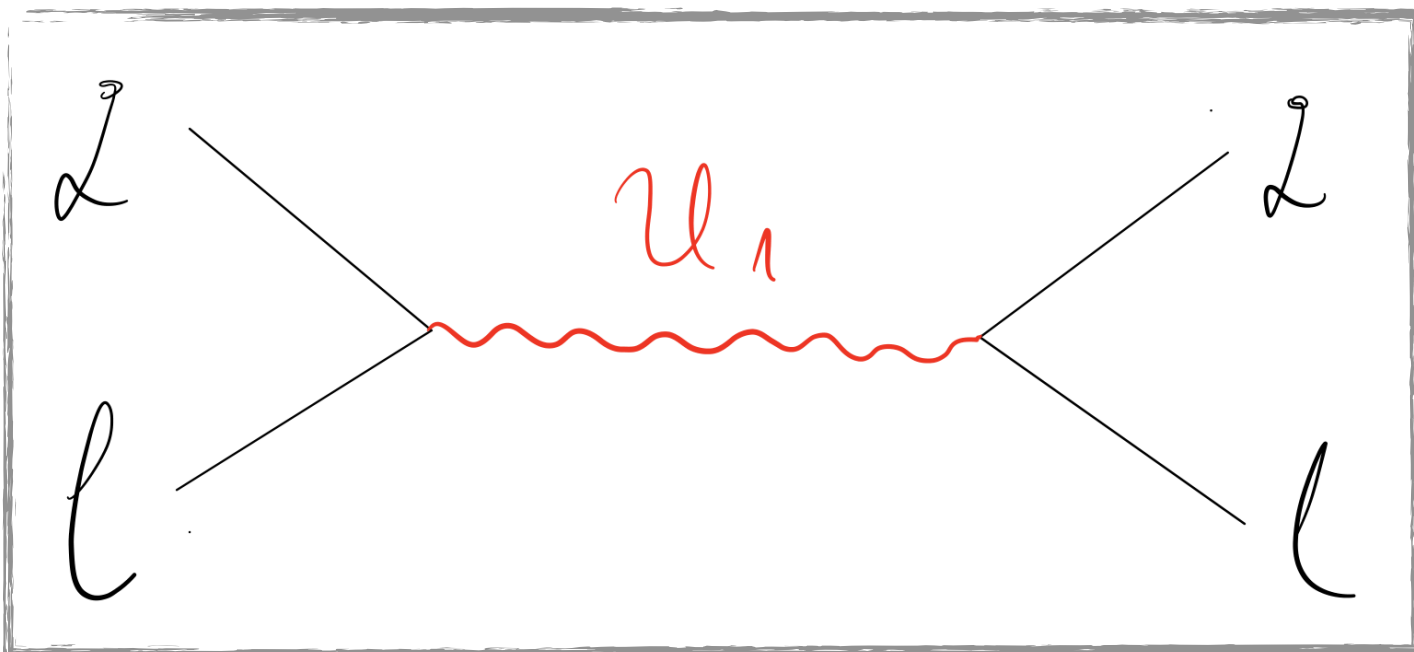
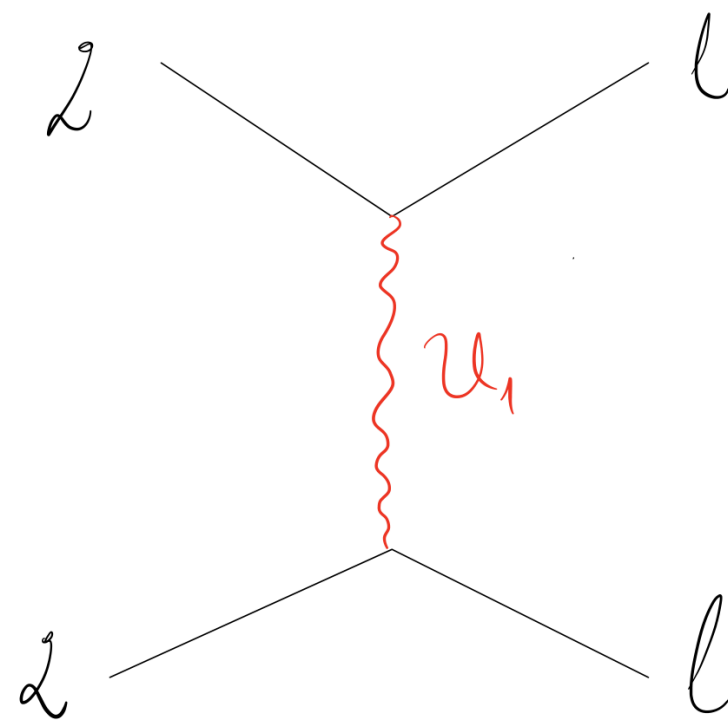
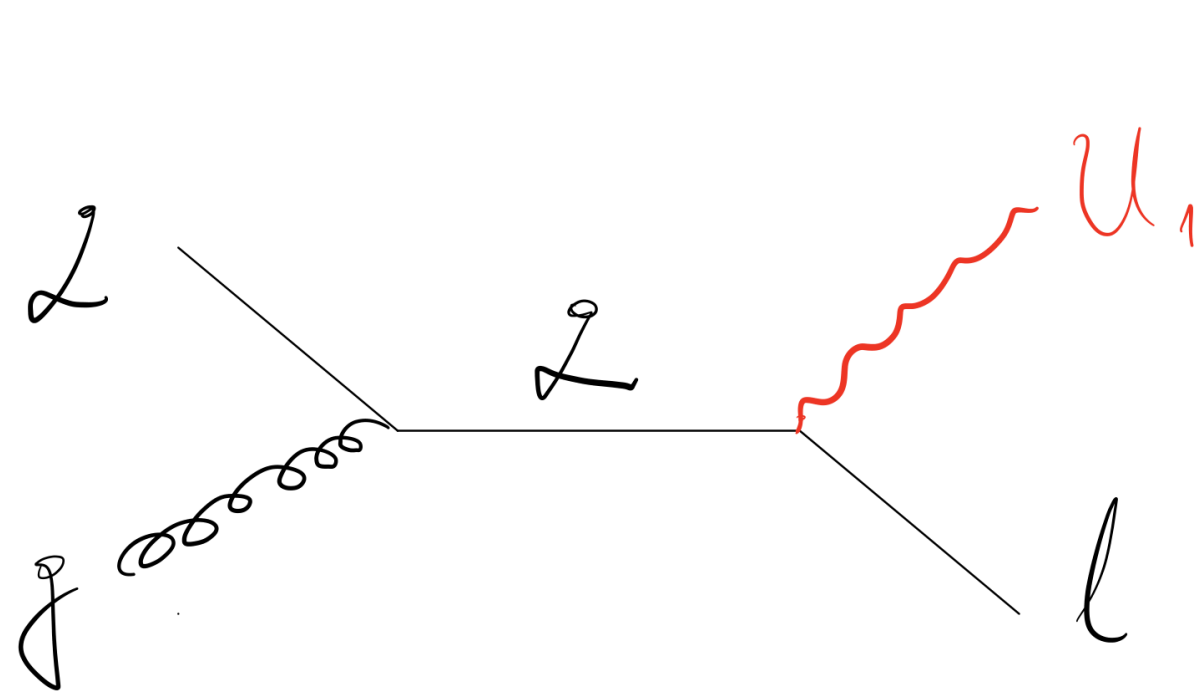
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Can be extended to couplings with light-gen fermions; additional fields required



A. Greljo and N. Selimović, 2012.02092

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# NLO QCD + QED results - Power counting

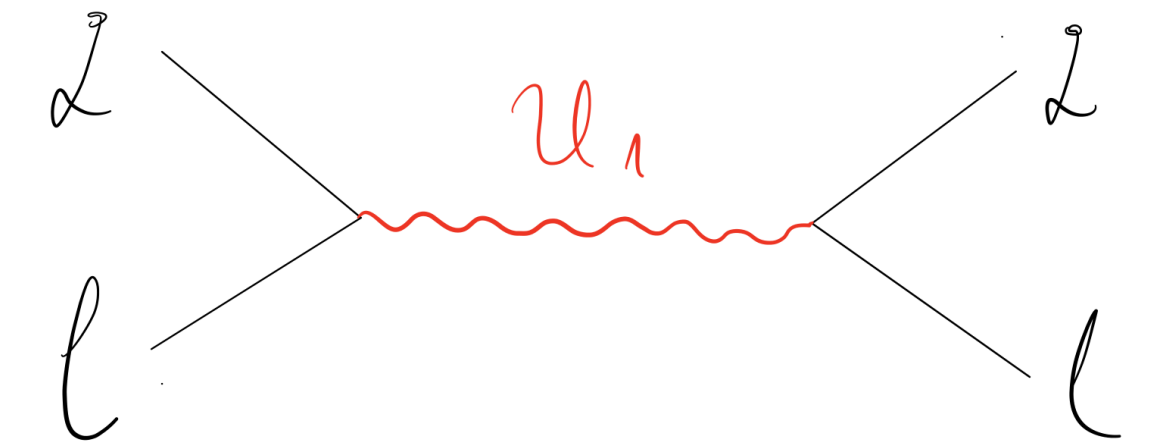
[Manohar, Nason, Salam, Zanderighi, arXiv: 1607.04266]

[A. Greljo and N. Selimović, arXiv: 2012.02092]

- The QCD and (some) QED corrections are of the same order:

$$f_q \sim f_g \sim \mathcal{O}\left(\sum_n (\alpha_s L)^n\right) \sim \mathcal{O}(1) \quad \alpha_s \sim 1/L \quad L = \log(\mu_F^2/\Lambda^2) \quad \Lambda - \text{hadronic scale}$$

- QED coupling size:  $\alpha_{\text{em}} \sim \alpha_s^2$ 
  - $\gamma$  - PDF, 1st order QED effect:  $f_\gamma \sim \mathcal{O}(\alpha_{\text{em}} L) \sim \mathcal{O}(\alpha_s)$
  - $\ell$  - PDF, 2nd order QED effect:  $f_\ell \sim \mathcal{O}((\alpha_{\text{em}} L)^2) \sim \mathcal{O}(\alpha_s^2)$



- Including real corrections:

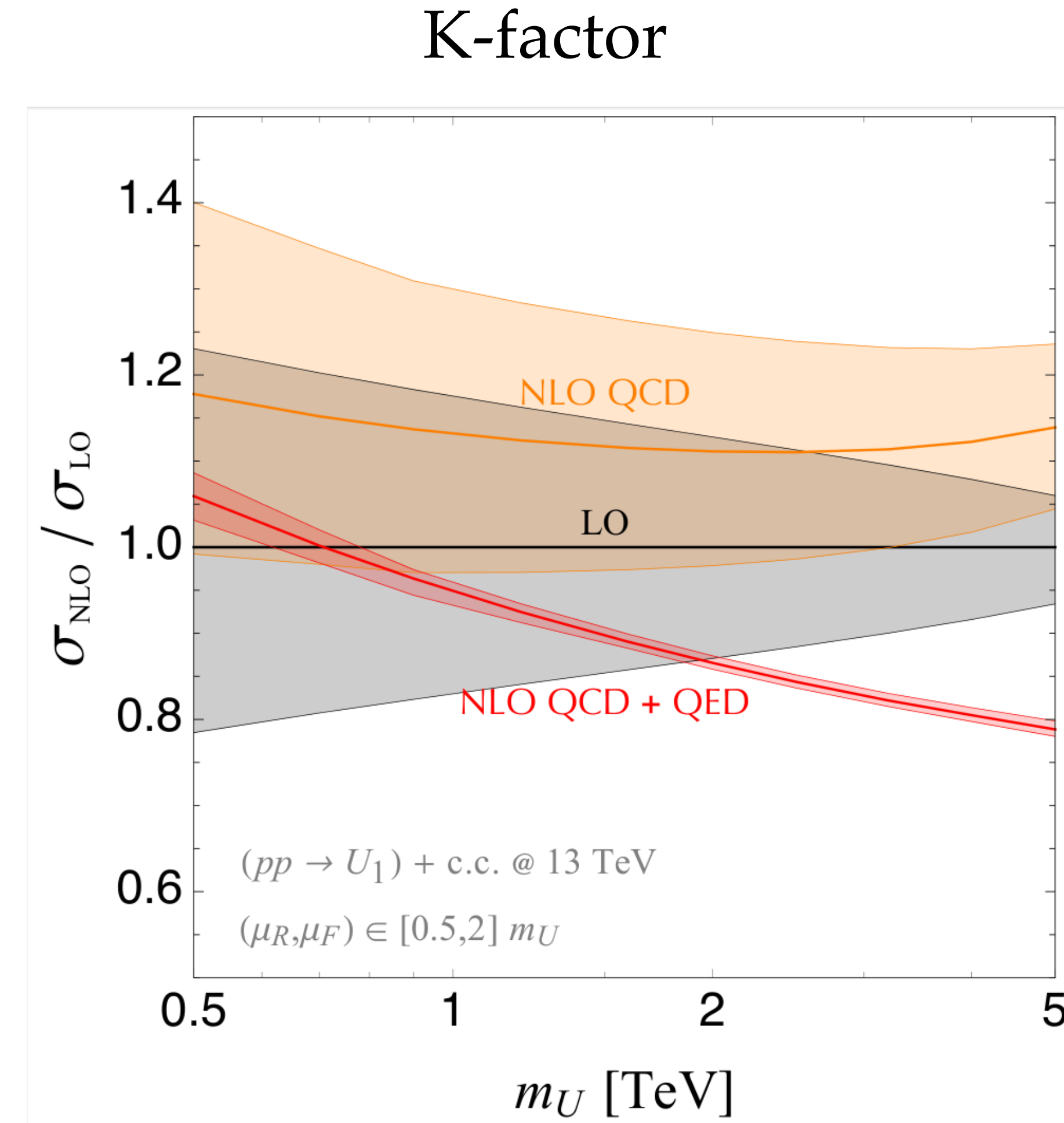
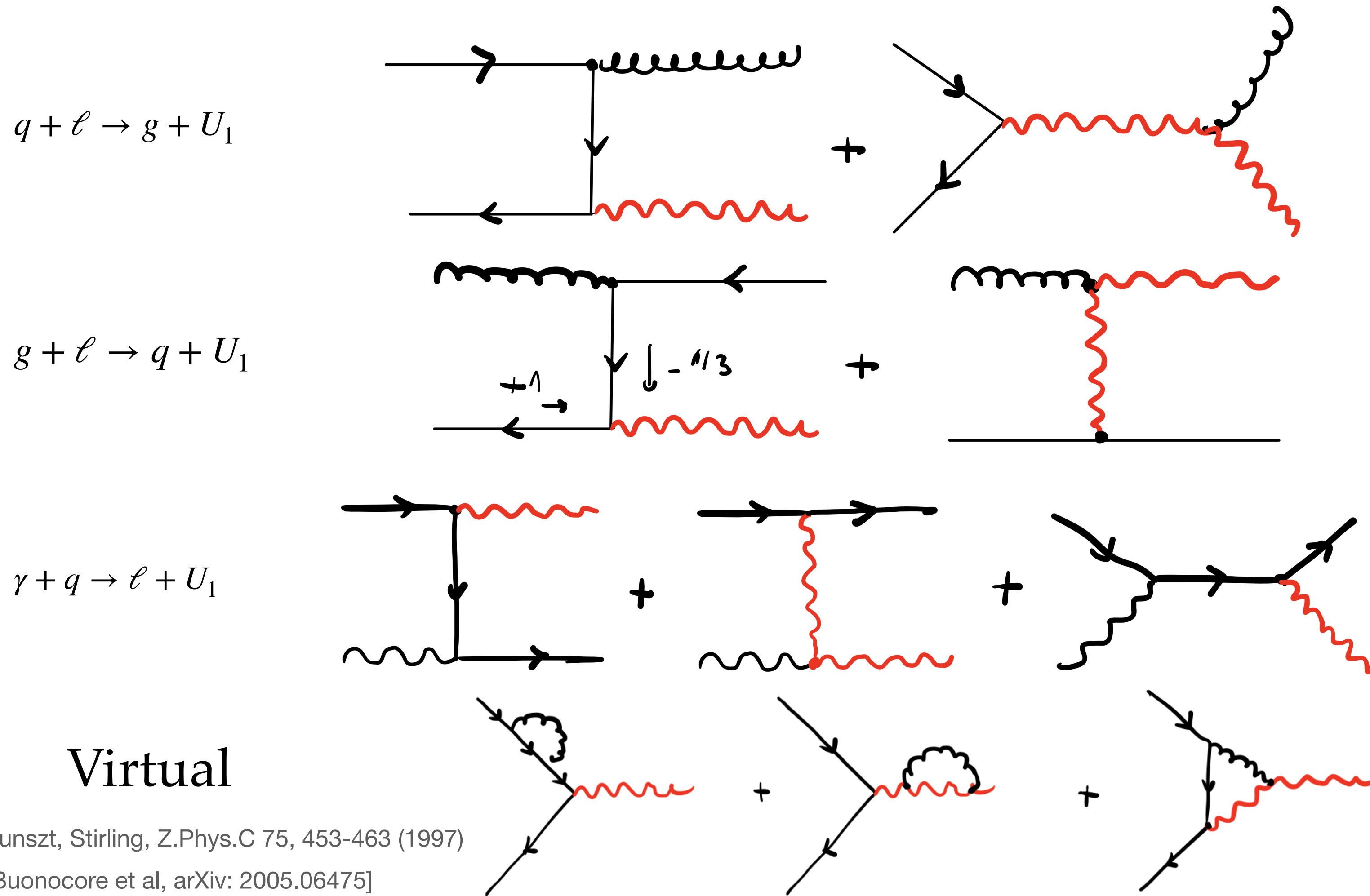
$$q + \ell \rightarrow g + U_1 \quad \sigma_{q\ell}(s) \sim \int (f_q \otimes f_\ell) \hat{\sigma}_{q\ell} \sim 1 \times \alpha_s^2 \times \alpha_s \sim \alpha_s^3$$

$$g + \ell \rightarrow q + U_1 \quad \sigma_{q\gamma}(s) \sim \int (f_q \otimes f_\gamma) \hat{\sigma}_{q\gamma} \sim 1 \times \alpha_s \times \alpha_{\text{em}} \sim \alpha_s^3$$

$$\gamma + q \rightarrow \ell + U_1$$



# NLO QCD + QED results



$g_4$ , LQ propagator - OS ren. scheme



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# Simulation details & exclusion limits

- With the matrix elements squared and the finite renormalization piece, build a POWHEG-BOX event generator

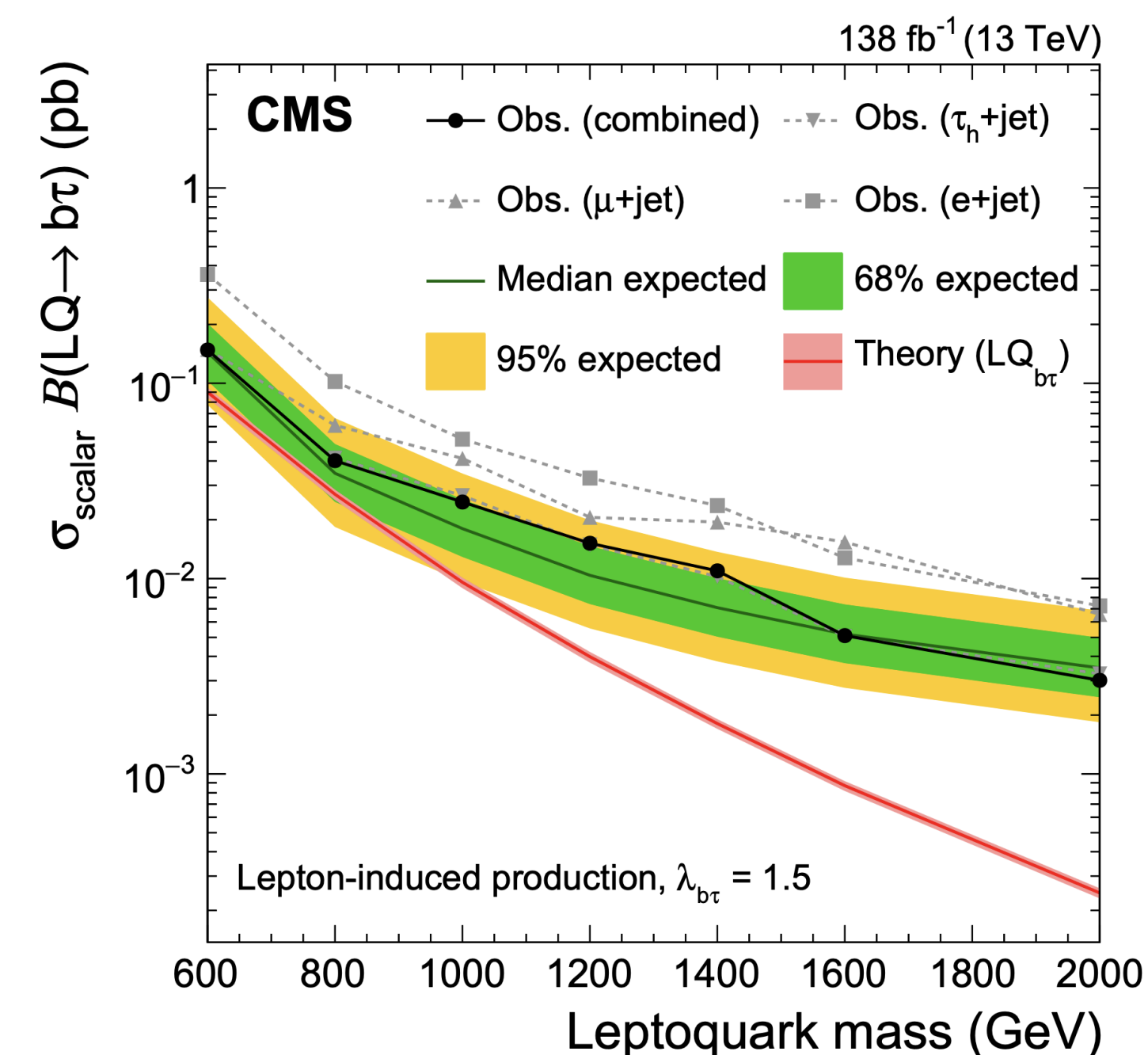
[Frixione et al, arXiv: 0707.3088]

[Buonocore et al, arXiv: 2209.02599]

- HERWIG has implemented lepton showering algorithm, contrary to Pythia - correction  $\sim 30\%$  on the cross-section
- Using the NNPDF3.1 LUXlep PDF set - **with lepton PDFs included**
- Jet clustering - FastJet
- Detector effects, tagging and reconstruction - Delphes or own POWHEG algorithm
- <https://github.com/peterkrack/3rd-Lepton-Quark-Fusion>

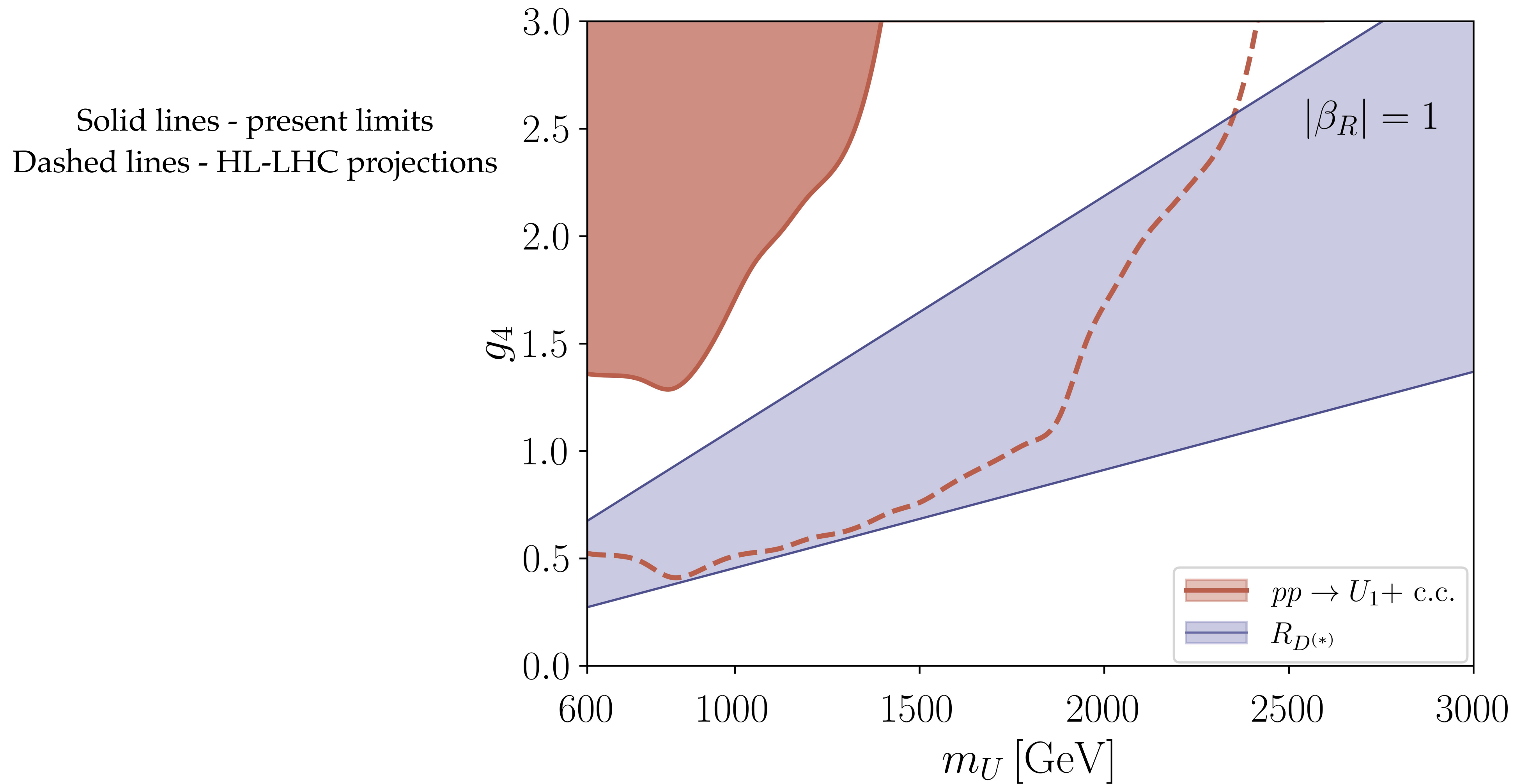
- **First experimental search on resonant scalar LQ production, targeting the  $b\tau$  final state**

[CMS Collaboration, arXiv: 2308.06143]



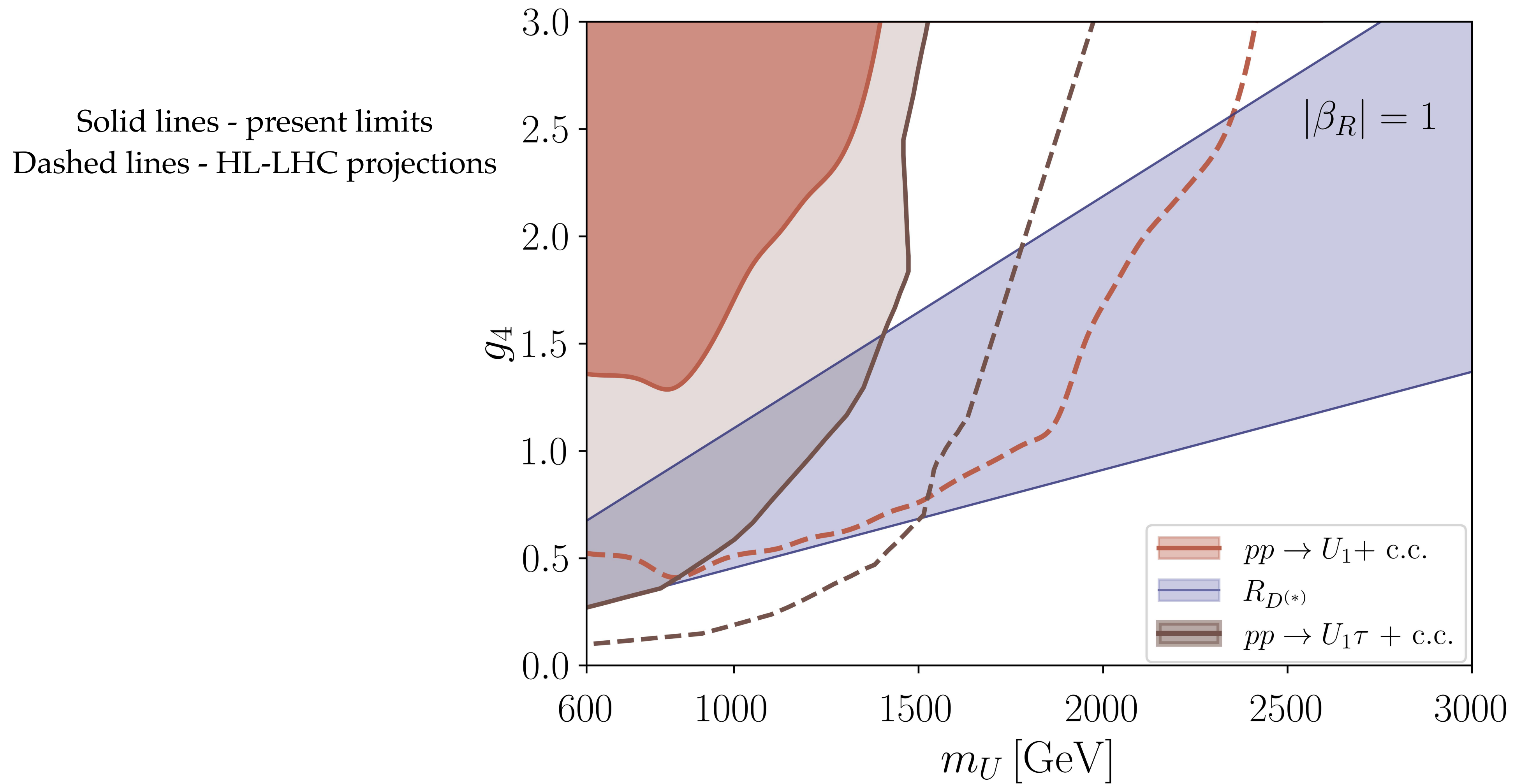
- Data for Brazilian bands (for different couplings) available on HEPData <http://dx.doi.org/10.17182/hepdata.141335>
- Input our signal, check when exclusion bounds are saturated

# Simulation details & exclusion limits



[ATLAS, arXiv: 2305.15962]

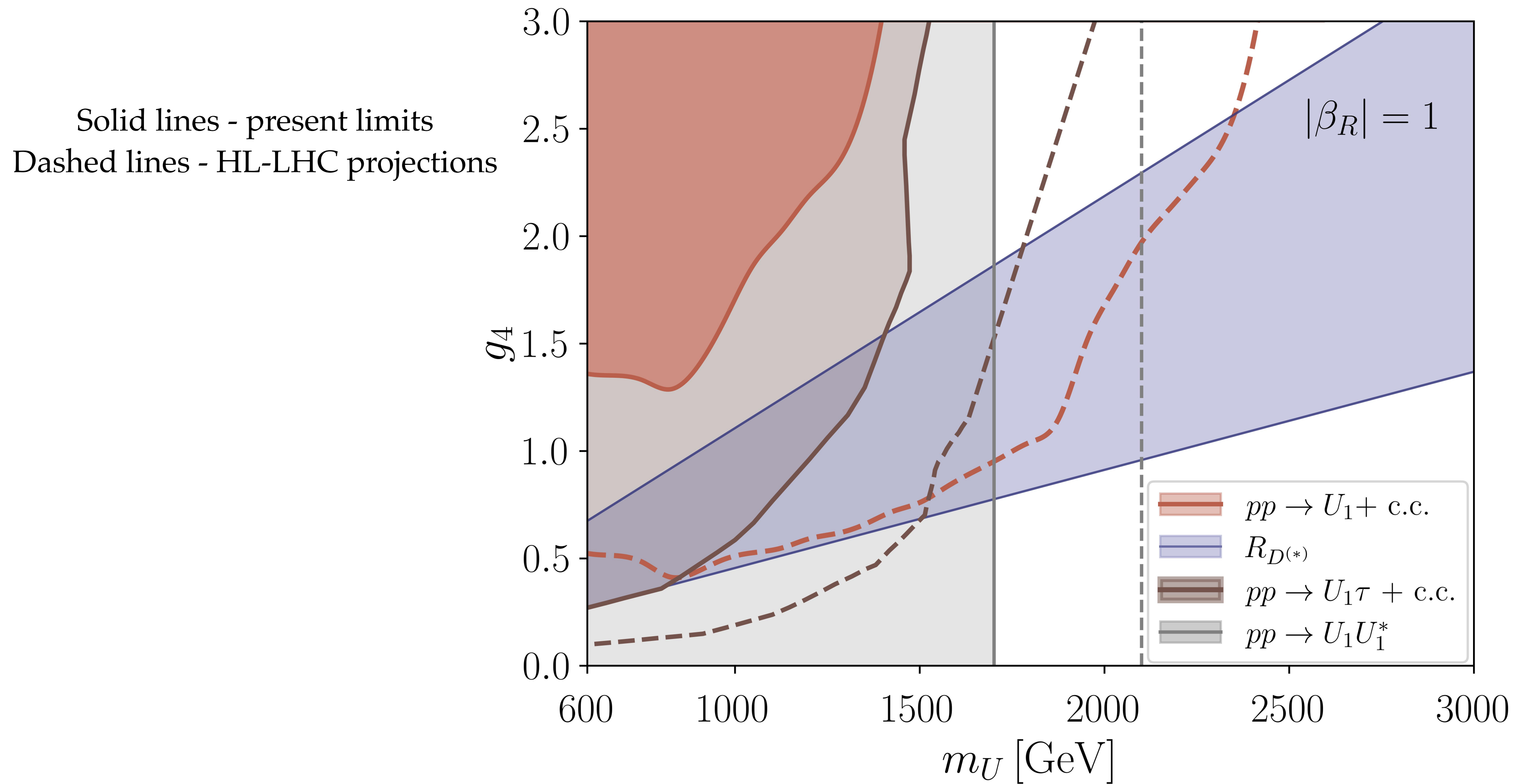
# Simulation details & exclusion limits



[ATLAS, arXiv: 2305.15962]

[CMS, arXiv: 2012.04178]

# Simulation details & exclusion limits



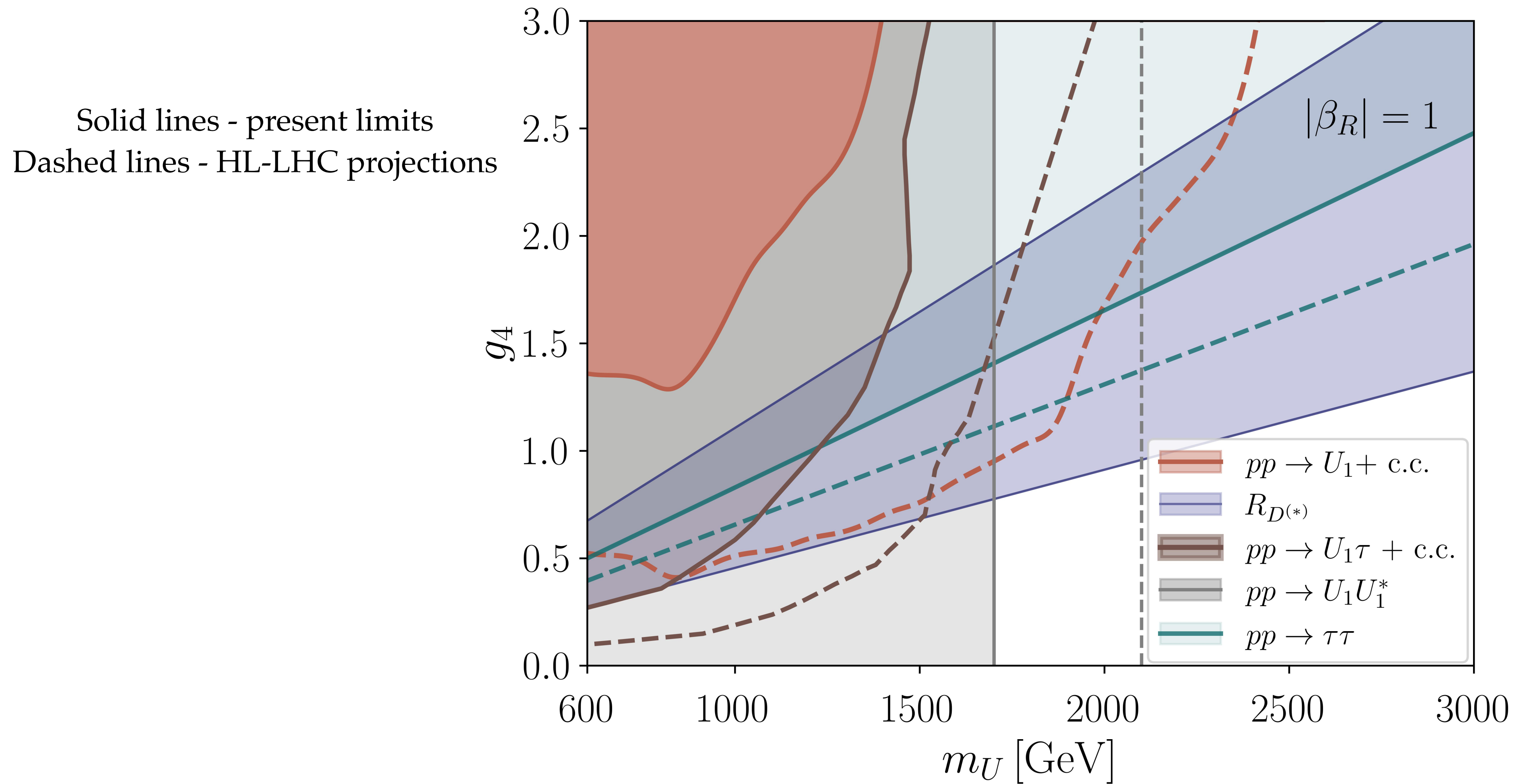
[ATLAS, arXiv: 2305.15962]

[CMS, arXiv: 2012.04178]

[Baker et al, arXiv: 2012.11474]



# Simulation details & exclusion limits



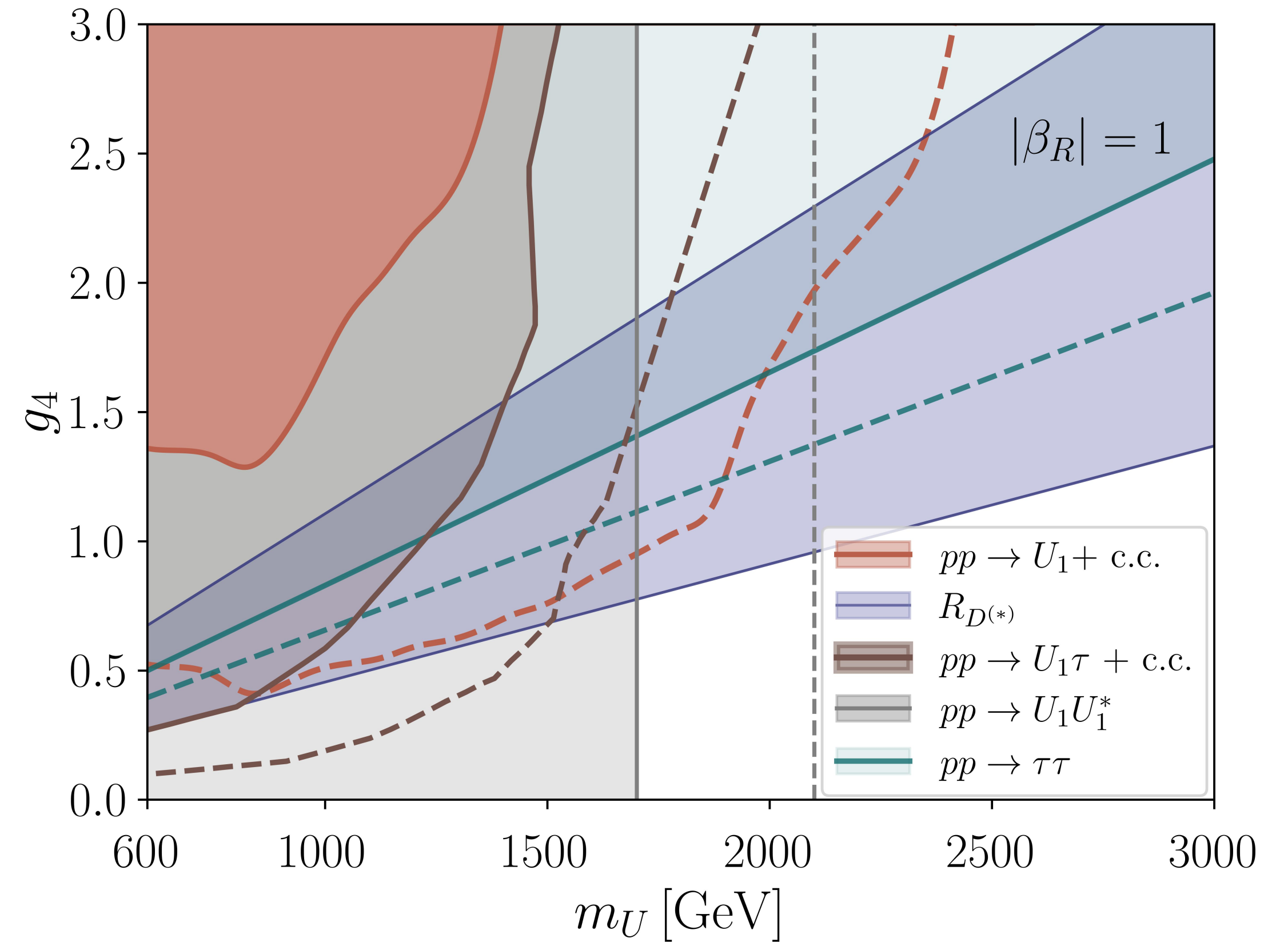
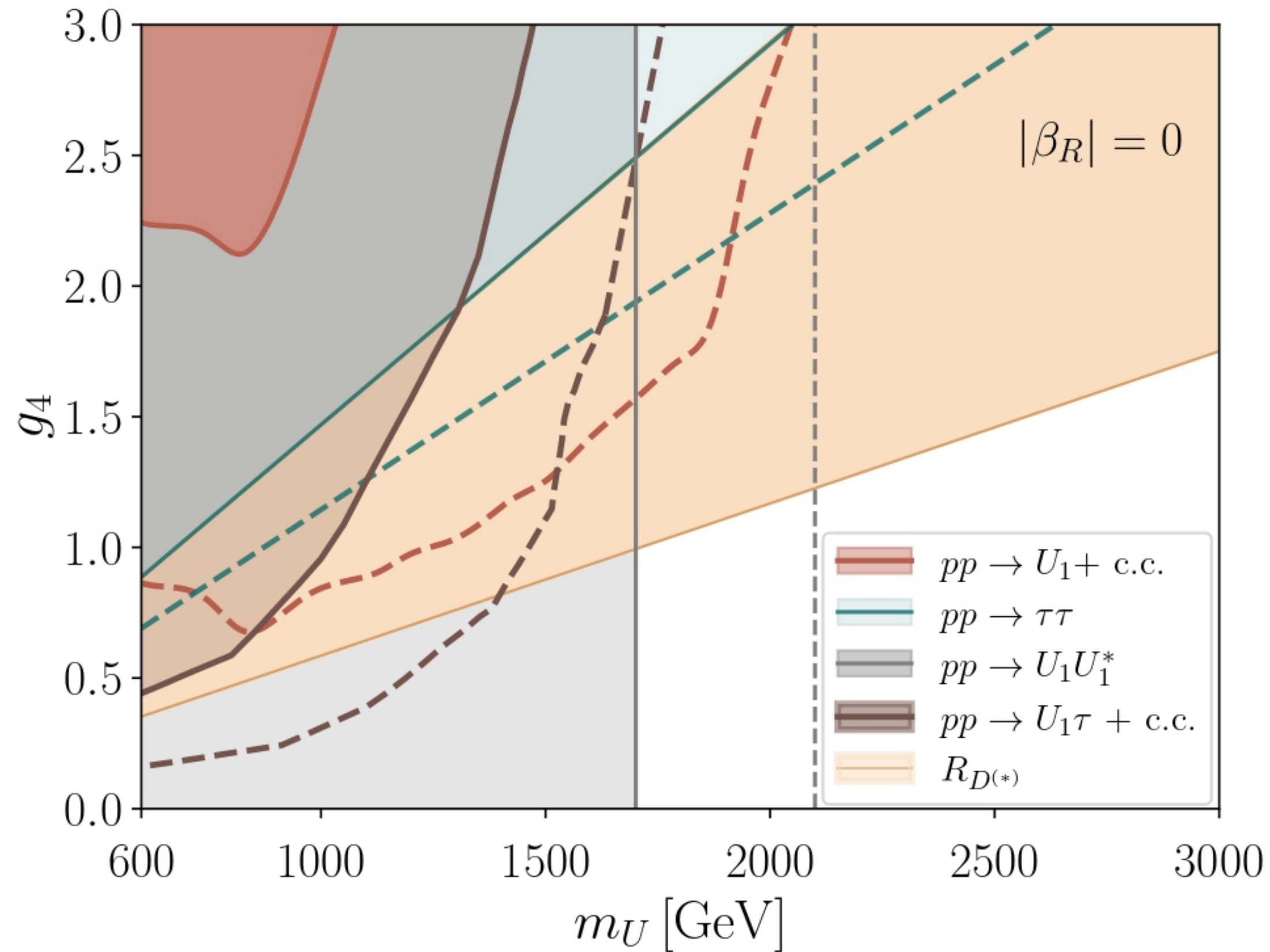
[ATLAS, arXiv: 2305.15962]

[CMS, arXiv: 2012.04178]

[Baker et al, arXiv: 2012.11474]



# Simulation details & exclusion limits



- Current bounds not competitive with other production mechanisms, mainly due to the lepton PDF suppression
- BUT: The enhancement of the COM energy very beneficial for the resonant production
- A discovery channel (opposed to t-channel limits)

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# Conclusions

- Utilized the first resonant LQ search at LHC
- Phenomenology of lepton-initiated production mechanism is well-understood, simulationwise it is controllable
- Tools for simulations steadily improving, there are more products on the market (for hadronic and lepton colliders)

Resonant scalar LQ production @ LHC in POWHEG-BOX - Buonocore et al; arXiv: 2209.02599

LePDF - PDFs for lepton colliders ( $e^+e^-$ ,  $\mu^+\mu^-$ ) - Garosi et al; arXiv: 2303.16964

MadGraph/POWHEG/Herwig compatibility - arXiv: 2108.10261

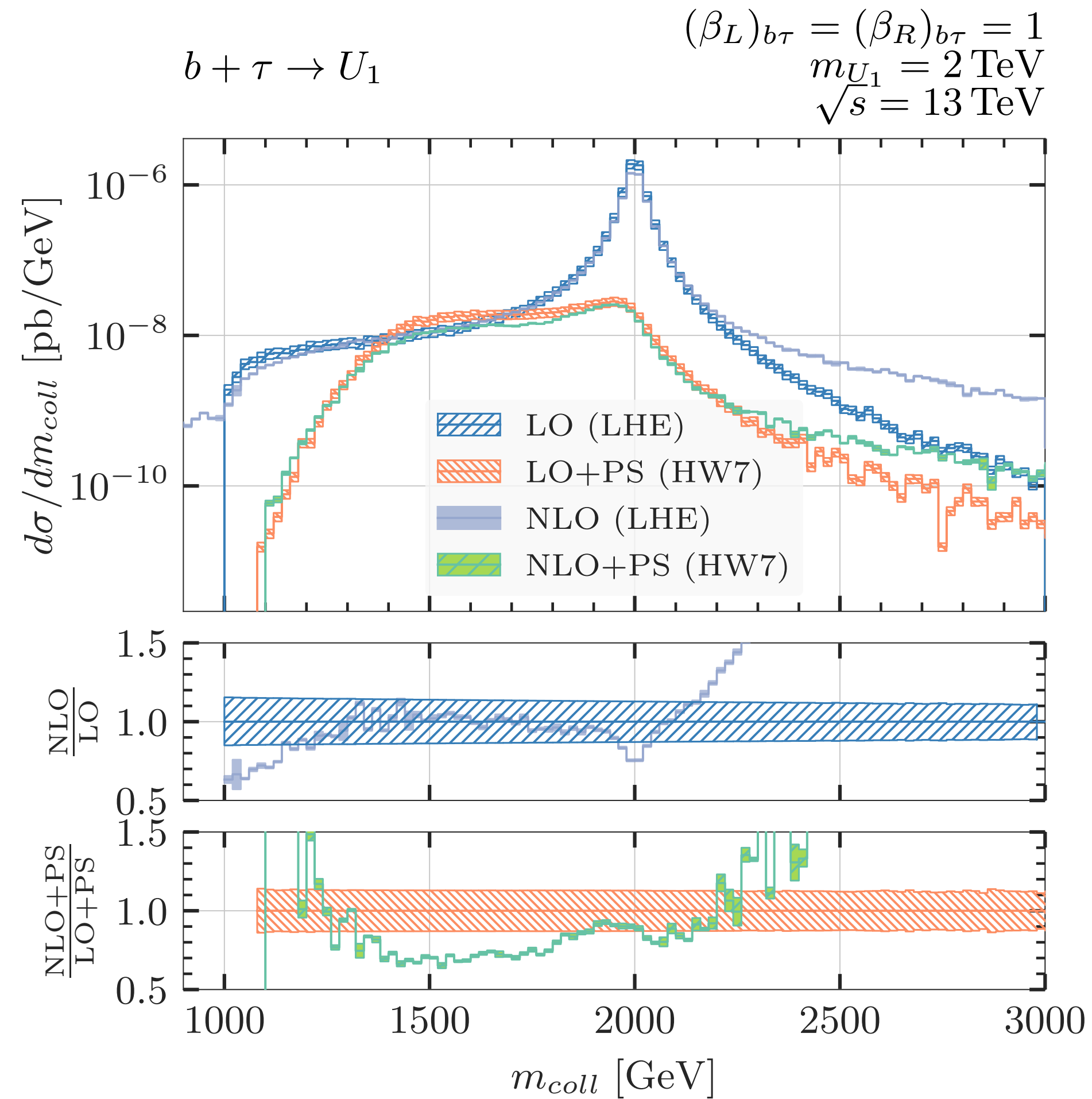
*simulate!*  
Shut up and ~~calculate~~

**Thank you!**

# Backup



# Simulation details - Reconstruction



- Taking into account the decay of  $\tau$  in HERWIG
- Implementation of the CMS cuts in the post-processing:

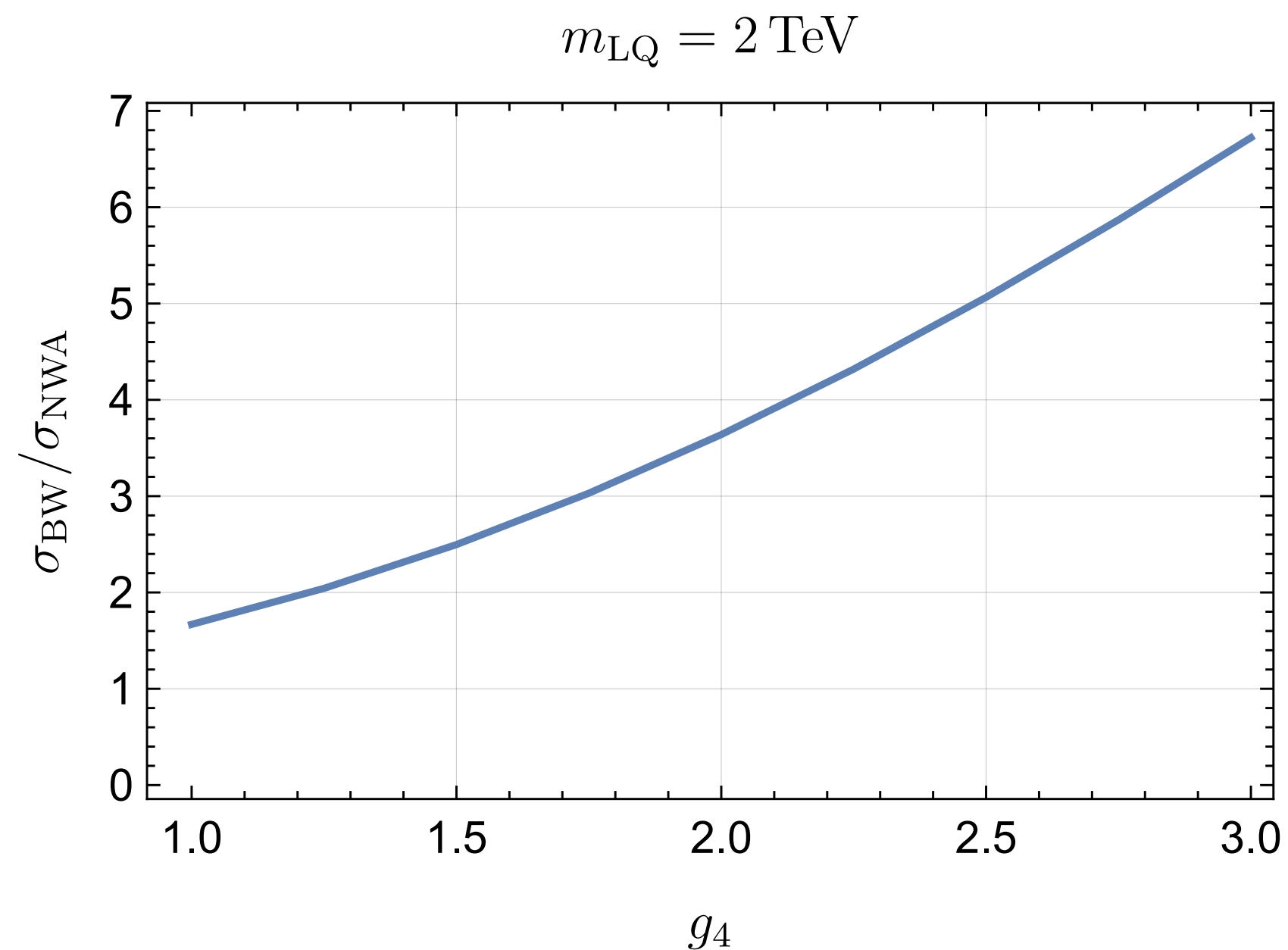
Variable	$\tau_h + \text{jet}$	$e + \text{jet}$	$\mu + \text{jet}$
$p_T^\ell (\text{GeV})$	$>200$	$>100$	$>100$
$ \eta^\ell $	$<2.1$	$<2.1$	$<2.1$
$p_T^{\text{jet}} (\text{GeV})$	$>300$	$>200$	$>200$
$ \eta^{\text{jet}} $	$<2.4$	$<2.4$	$<2.4$
$p_T^{\text{miss}} (\text{GeV})$	$>100$	$>150$	$>150$
$p_T(\vec{\ell} + \vec{p}_T^{\text{miss}}) (\text{GeV})$	$>100$	$>100$	$>100$
$ \Delta\phi(\ell, \vec{p}_T^{\text{miss}})  (\text{radians})$	$<0.3$	$<0.2$	$<0.2$
$\Delta R(\ell, \text{jet})$	$>0.5$	$>0.5$	$>0.5$

[CMS Collaboration, arXiv: 2308.06143]

- Using the quantity  $m_{\text{coll}} = m_{\text{vis}} \sqrt{1 + p_T^{\text{invis}}/p_T^{\text{vis}}}$  to reconstruct the peak

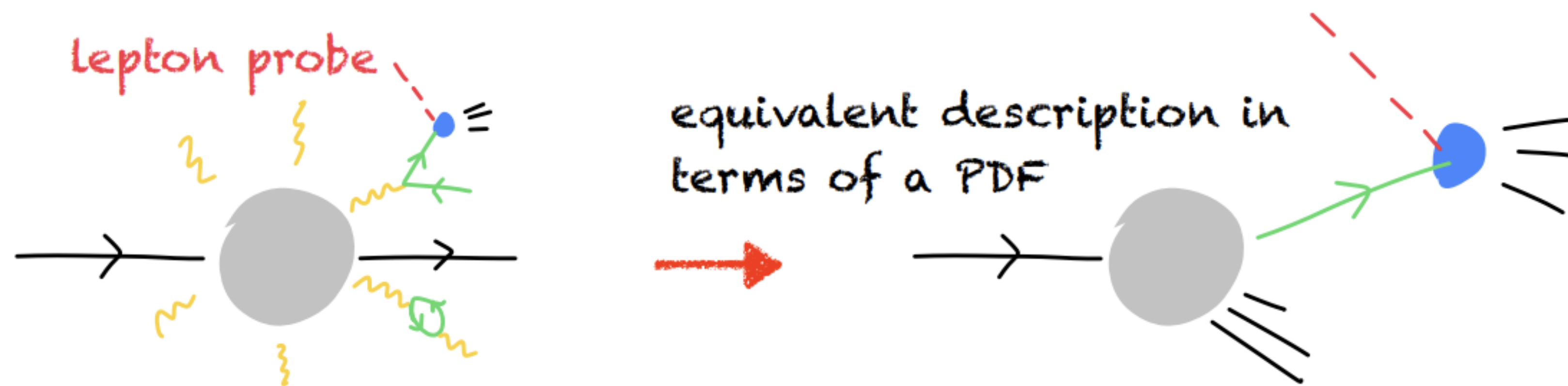
# Finite-width effects

- The narrow-width approximation (NWA) can have an  $\mathcal{O}(1)$  relative error compared to full Breit-Wigner result for  $m_{LQ} \gtrsim 2 \text{ TeV}$  and  $g_4 \gtrsim 1.5$ !
- The Breit-Wigner prescription for the propagator can be turned on by setting the flag `BWgen = 1`.
- For larger values of coupling, the convolutional area between the PDFs and the Breit-Wigner is also larger, hence we are more sensitive to the low- $x$  region of the PDFs (enhanced by orders of magnitude compared to the values of the PDF at the mass pole)



# Lepton PDFs

Two equivalent descriptions of probing the proton through DIS:



Equate the results of these two methods!

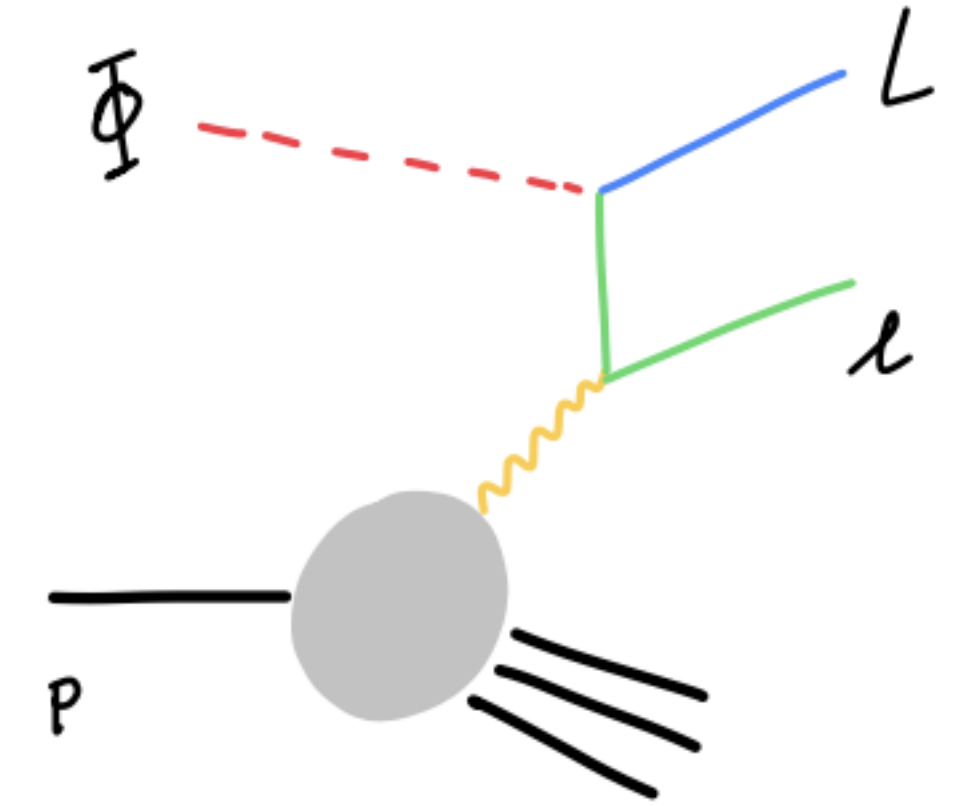
H. Anlauf et al , Comput. Phys. Commun. 70, 97 (1992)

Manohar et al <https://arxiv.org/pdf/1607.04266>

# Lepton PDFs

Method 1 (DIS-style):

$$\sigma = \frac{1}{4p \cdot r} \int \frac{d^4q}{(2\pi)^4} \frac{1}{Q^4} L^{\mu\nu}(r, q) (4\pi) W_{\mu\nu}(p, q)$$



$$W_{\mu\nu}(p, q) = F_1 \left( -g_{\mu\nu} + \frac{q_\mu q_\nu}{q^2} \right) + \frac{F_2}{p \cdot q} \left( p_\mu - \frac{p \cdot q}{q^2} q_\mu \right) \left( p_\nu - \frac{p \cdot q}{q^2} q_\nu \right)$$

$F_1(x_{bj}, Q^2), F_2(x_{bj}, Q^2)$  are the proton structure functions

# Lepton PDFs

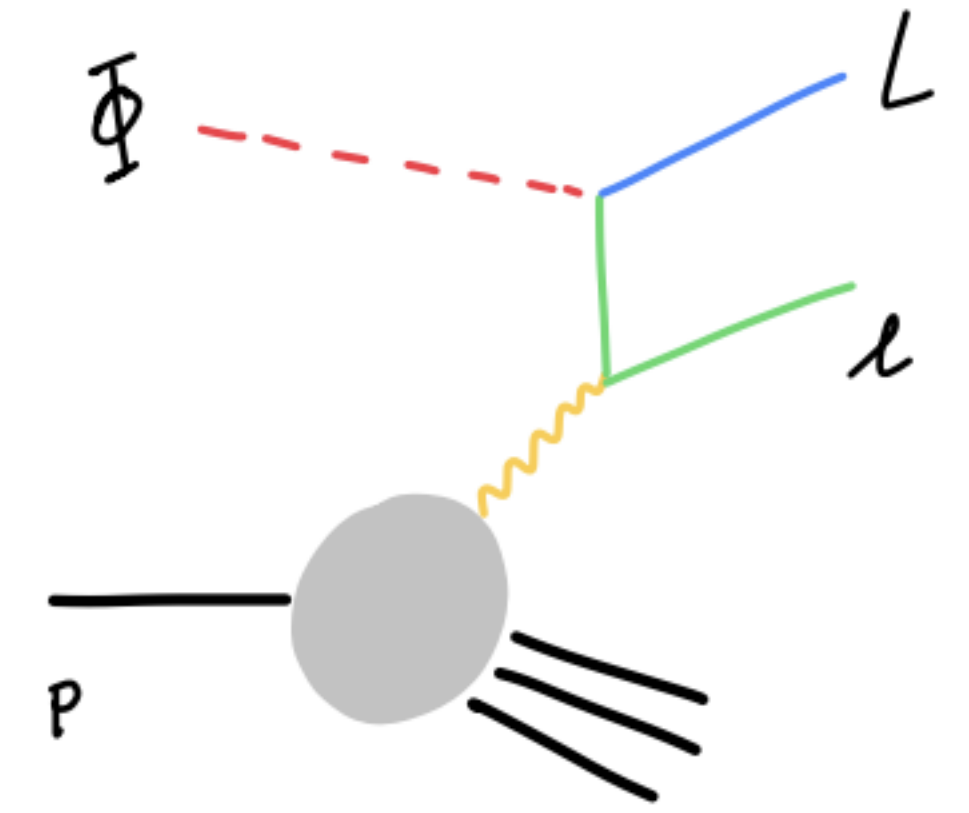
Method 2 (proton PDF):

**LO**

$$\frac{\sigma}{\sigma_B} = \int dx f_\ell(x, \mu_F^2) \delta(Sx - M^2)$$

$$+ \frac{\alpha}{2\pi} \frac{1}{M^2} \int_{\frac{M^2}{S}}^1 dx f_\gamma(x, \mu_F^2) \left\{ z_\ell P_{l\gamma}(z_\ell) \left[ \log \frac{M^2}{\mu_F^2} + \log \frac{(1 - z_\ell)^2}{z_\ell^2} \right] + 4z_\ell^2(1 - z_\ell) \right\}$$

**NLO**



$$z_\ell = \frac{M^2}{xS}$$