



HOW CHARMING CAN ^{THE HIGGS} _{BOSON} BE?

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arXiv:2410.05236 (about to appear in JHEP)

How charming can the Higgs be?

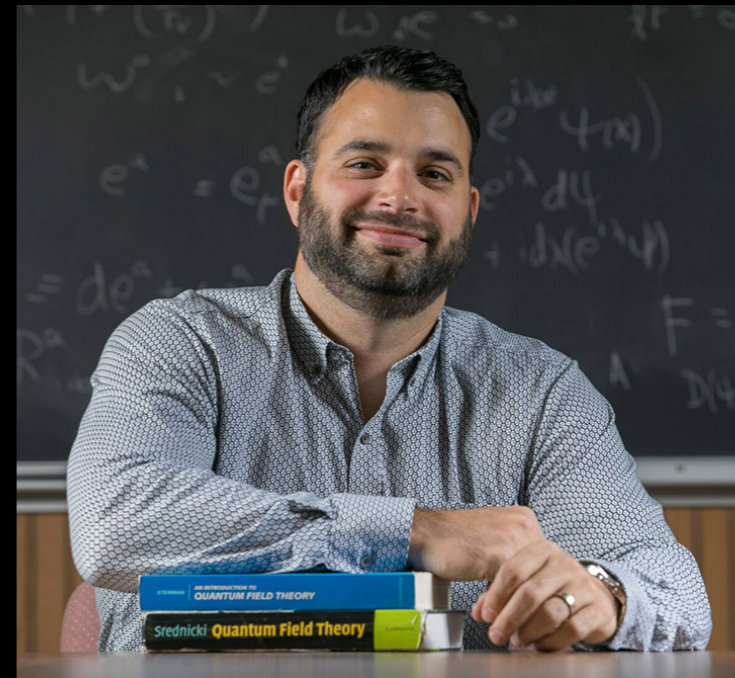
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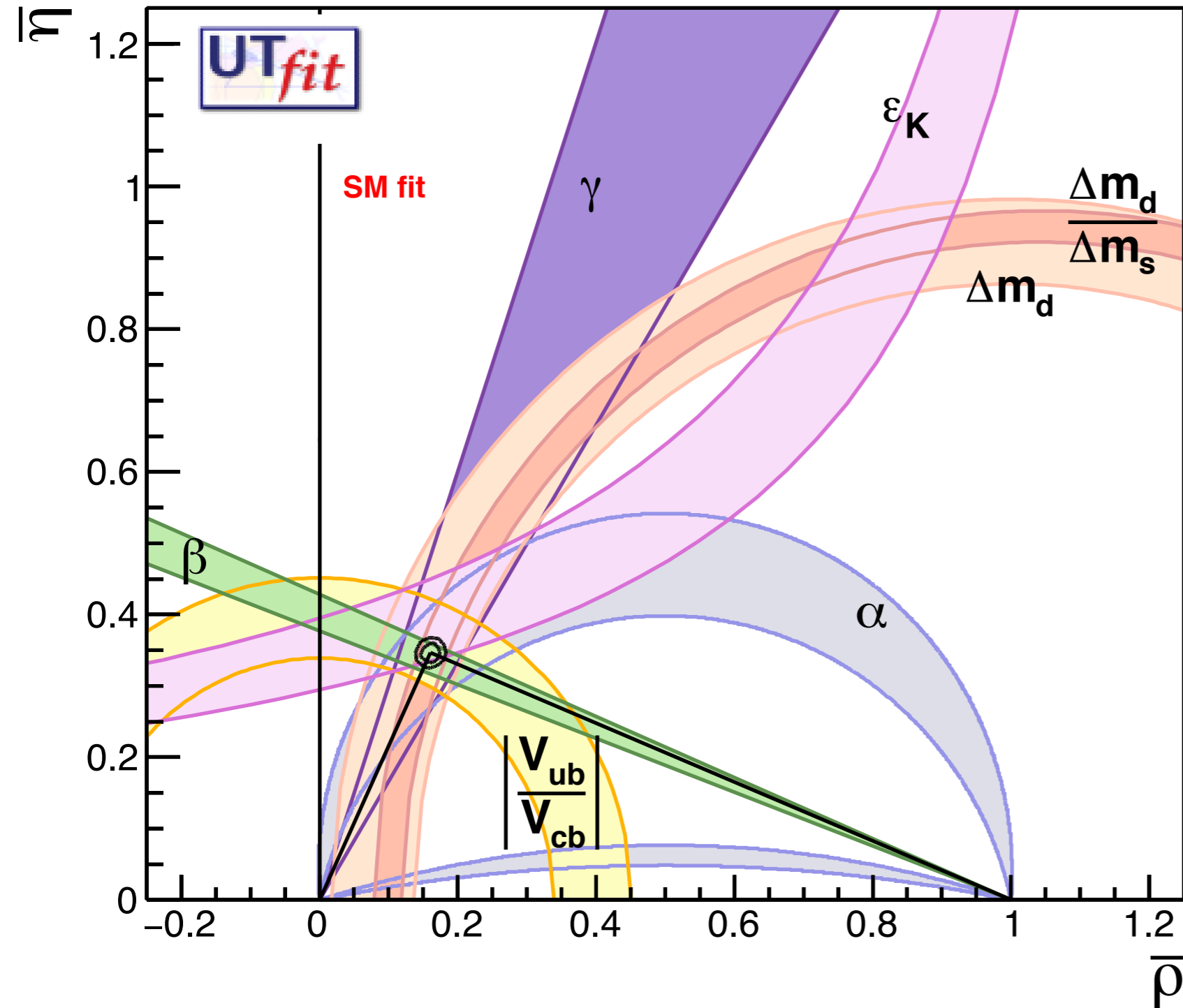
Artemis Giannakopoulou



UTA: Unitarity Triangle Analysis

@ 95% prob

Rend.Lincei Sci.Fis.Nat. 34 (2023)



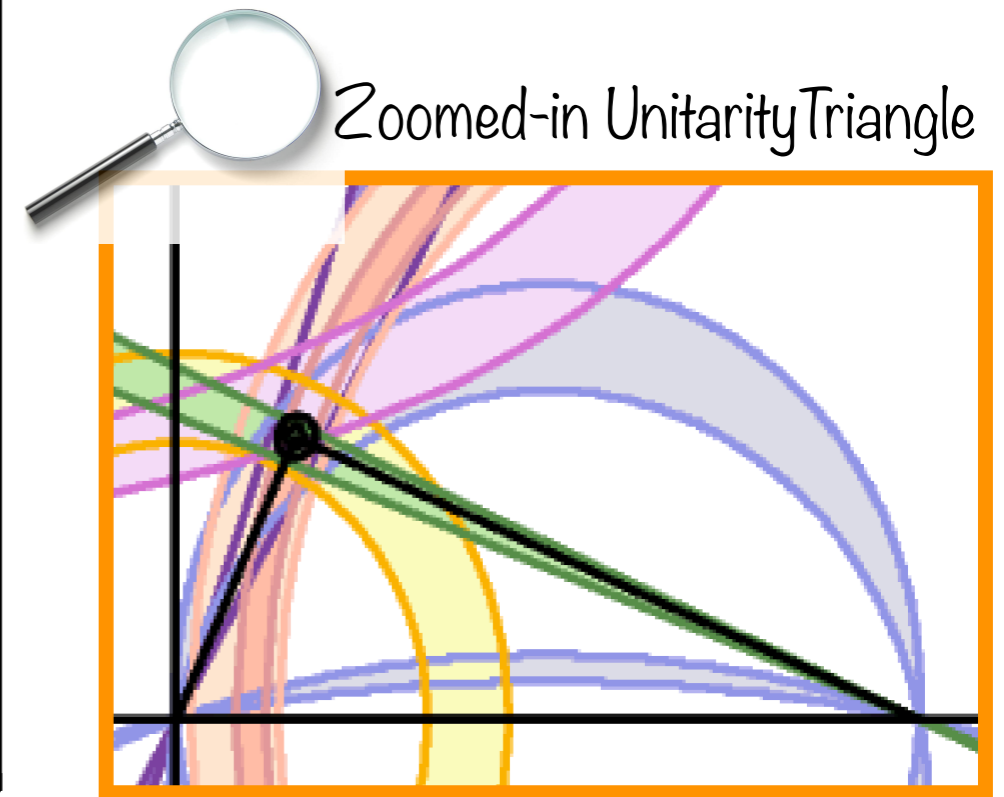
$$\bar{\rho} = 0.160 \pm 0.009 \sim 6\%$$

$$\bar{\eta} = 0.346 \pm 0.009 \sim 3\%$$

UT determination towards % precision

$$\lambda = 0.2251 \pm 0.0008$$

$$A = 0.827 \pm 0.010$$



Flavor & BSM Physics

PoS WIFAI2023 (2024) 007

SM

$$O_1^{q_i q_j} = \bar{q}_{jL}^\alpha \gamma_\mu q_{iL}^\alpha \bar{q}_{jL}^\beta \gamma^\mu q_{iL}^\beta$$

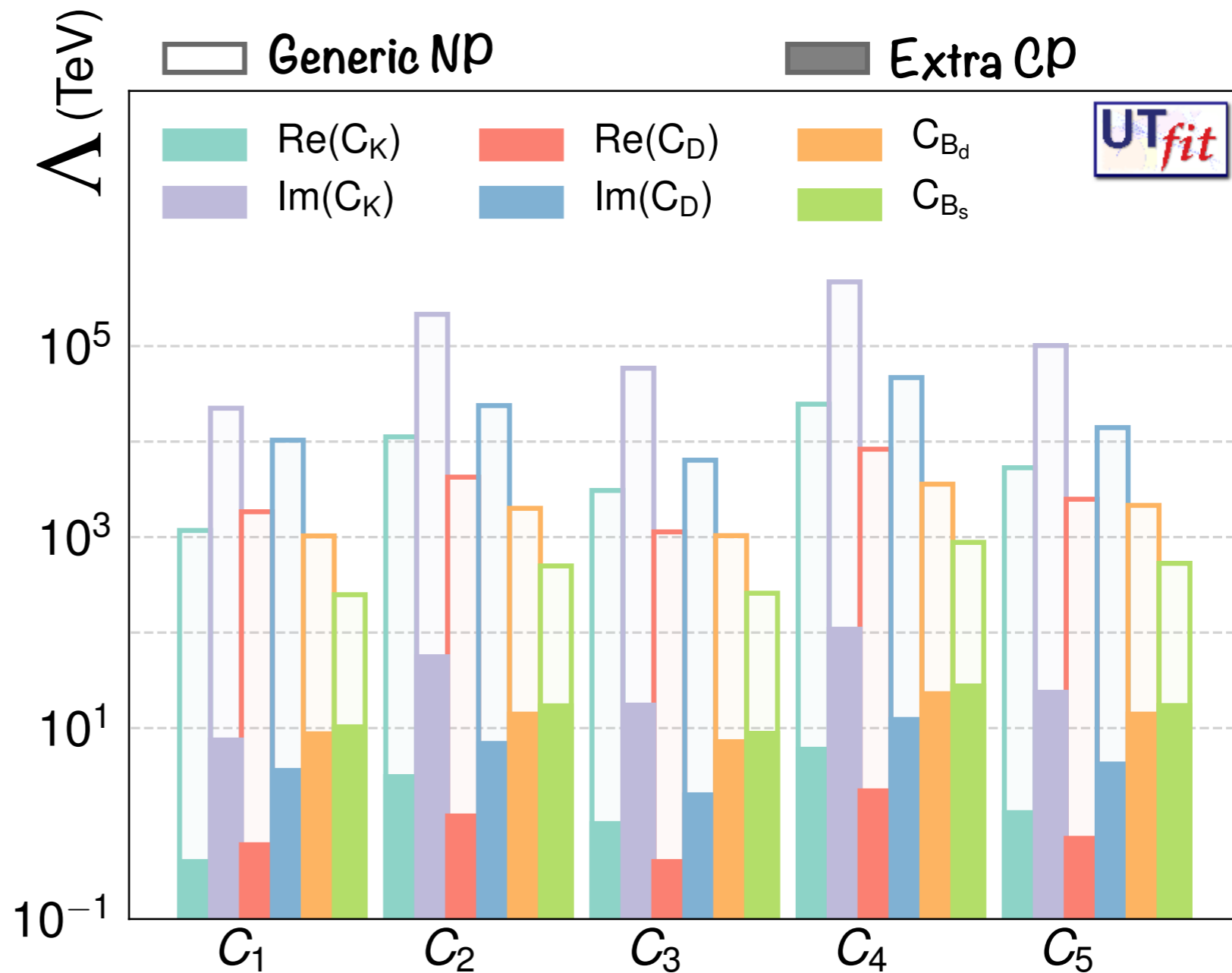
$$O_2^{q_i q_j} = \bar{q}_{jR}^\alpha q_{iL}^\alpha \bar{q}_{jR}^\beta q_{iL}^\beta$$

$$O_3^{q_i q_j} = \bar{q}_{jR}^\alpha q_{iL}^\beta \bar{q}_{jR}^\beta q_{iL}^\alpha$$

$$O_4^{q_i q_j} = \bar{q}_{jR}^\alpha q_{iL}^\alpha \bar{q}_{jL}^\beta q_{iR}^\beta$$

$$O_5^{q_i q_j} = \bar{q}_{jR}^\alpha q_{iL}^\beta \bar{q}_{jL}^\beta q_{iR}^\alpha$$

+ chirally flipped $\tilde{O}_{1,2,3}^{q_i q_j}$



○ **Generic NP** = no SM protection, i.e.: $C(\Lambda) \sim 1/\Lambda^2$ \Rightarrow $\Lambda > 4.7 \times 10^5$ TeV

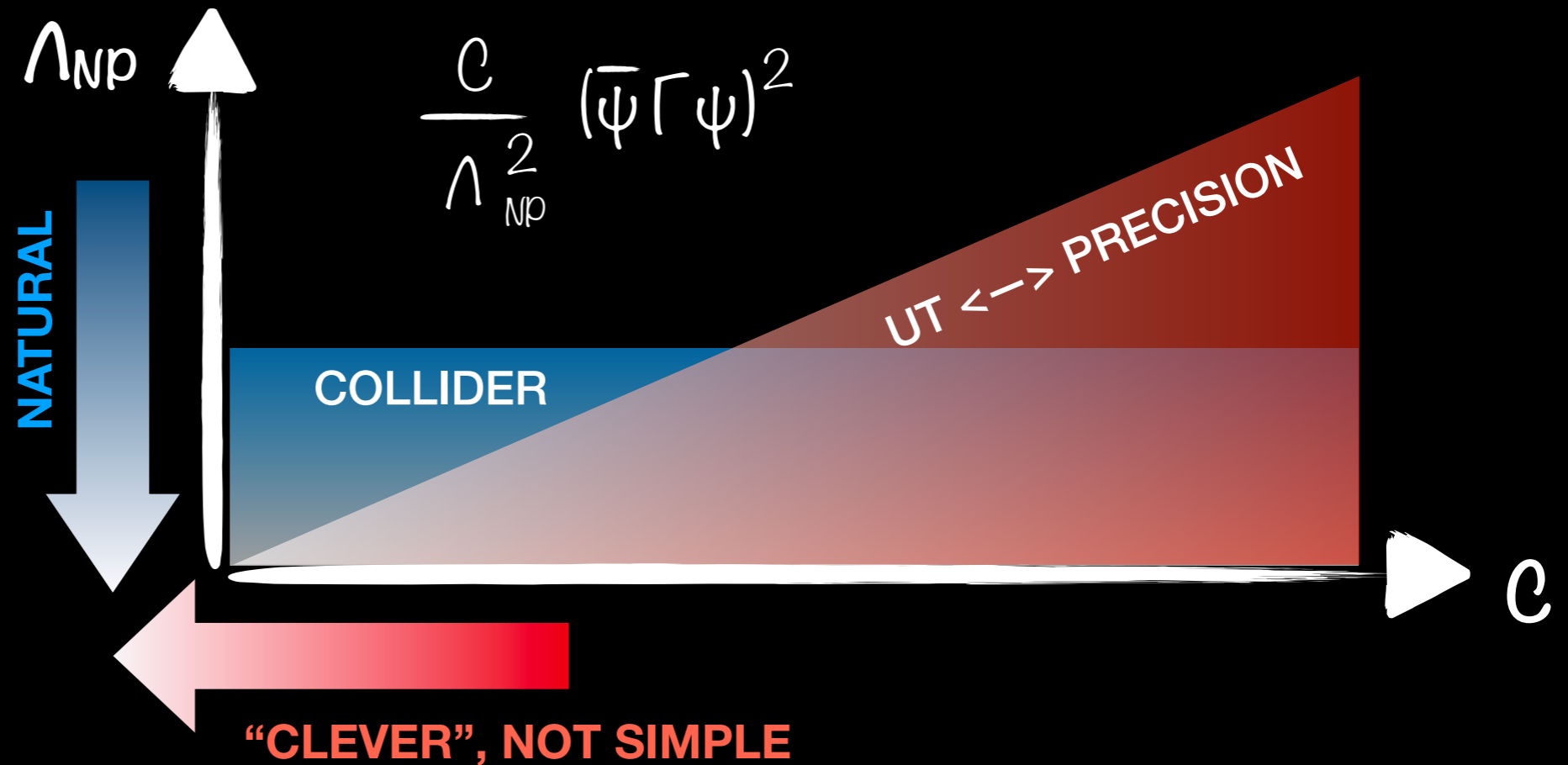
● **Extra CP** = SM-like protection but new $O(1)$ phases \Rightarrow $\Lambda > 108$ TeV



Lessons from Precision

- **SM UT:** Towards % precision ... overall remarkable consistency!
—> *in the HL-LHC era we can hope for a permil test*

- **NP UT:**



**BOTTOM
LINE**

A Theory of Flavor is either **VERY CLEVER** or **"JUST" UNNATURAL**

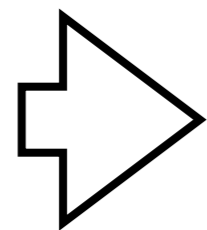
How To Be Clever ...

SEIBERGOLOGY: Yukawa are spurions breaking the **flavor group**

$$U(3)^5 = SU(3)_Q \times SU(3)_d \times SU(3)_u \times SU(3)_L \times SU(3)_e \times U(1)^5$$

$$U(1)^5 = U(1)_B \times U(1)_L \times U(1)_Y \times U(1)_{PQ} \times U(1)_e$$

E.g.: $-\mathcal{L}_{\text{SM}} \supset Y_{ij}^d \bar{Q}_i H d_j + Y_{ij}^u \bar{Q}_i \tilde{H} u_j + h.c.$



$$Y^d \sim (3, \bar{3}, 1) \quad Y^u \sim (3, 1, \bar{3})$$

under $SU(3)_Q \times SU(3)_d \times SU(3)_u$

ASSUMPTION: only spurions that break **flavor** are **SM Yukawas**

—> **MINIMAL FLAVOR VIOLATION** [*Nucl.Phys.B* 645 (2002) 155]

aka **MAXIMAL FLAVOR CONSERVATION** (*in Lausanne*)

MFV leaves room for **New Physics** just above a few TeV.

Model-independent bounds on the standard model effective theory from flavour physics

Physics Letters B

Volume 799, 10 December 2019, 135062

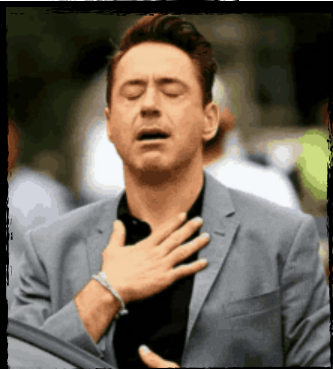
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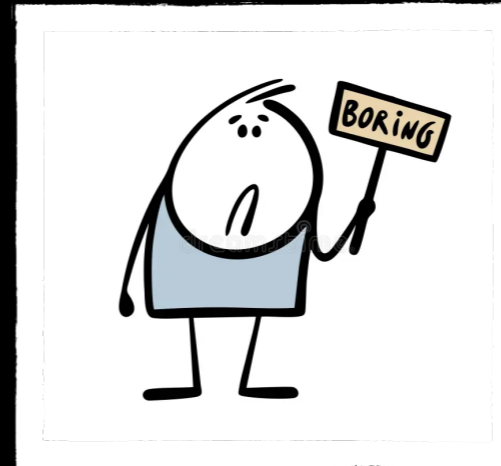
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Little Hierarchy problem possibly not too bad ...

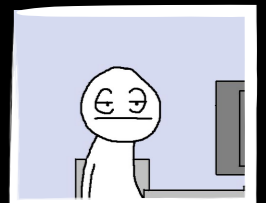


But ... isn't MFV ...



, though?

BSM pheno driven by couplings to 3rd gen

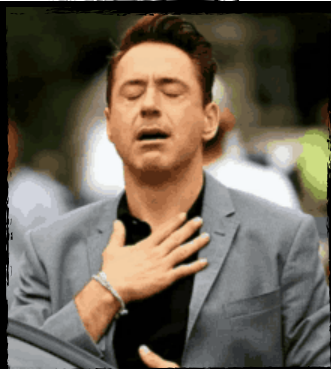


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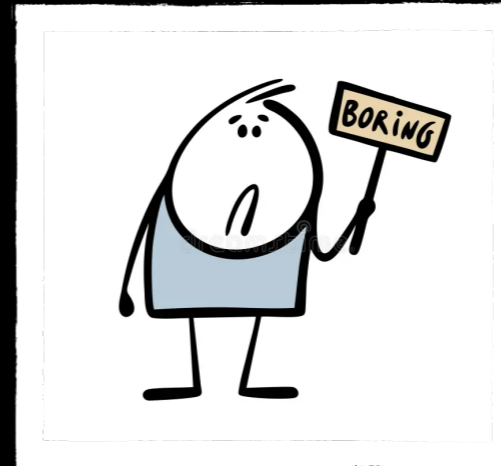
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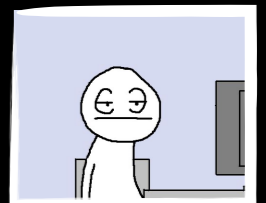


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Can we avoid large FCNCs effects & have BSM
with a radically different flavor than the SM one



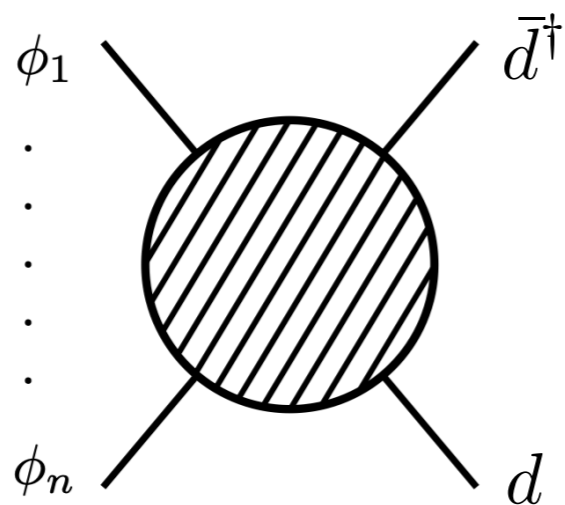
Another Way to Be Clever: Alignment!

— ANSATZ FOR ALIGNMENT IN THE UP-QUARK SECTOR —

Require (at least at the renormalizable level) the UV theory to be such that:

- **NO** breaking of the family number $U(1)_f^3$ and of CP other than the wave-function renormalization Z^d of d
- **NO** fields / spurions non-trivial under $U(3)_d$ but d and Z^d

Mixing with
Spontaneously
Broken Flavor
Vacuum



$$\mathcal{L} \supset i Z_{ij}^d \bar{d}_i^\dagger \bar{\sigma}^\mu d_j + i \bar{u}_i^\dagger \bar{\sigma}^\mu u_i + i \bar{Q}_i^\dagger \bar{\sigma}^\mu Q_i + \text{Yukawa terms all flavor diagonal}$$

Getting the kinetic term canonical one gets: $(Z^d)^{-\frac{1}{2}} \sim V y^d$

The up-quark sector is unaffected and **remains aligned**.

New spurions allowed in this sector: $\lambda^u = \text{diag}(\lambda_u, \lambda_c, \lambda_t)$

EXAMPLE: VECTOR-LIKE QUARK

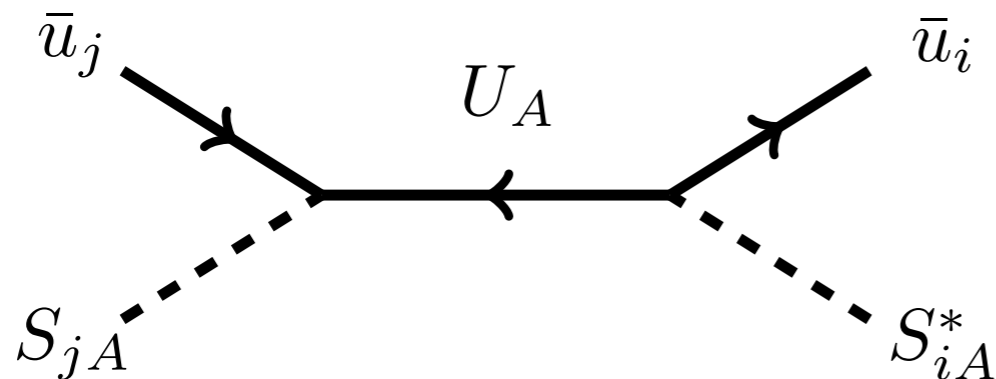
$$\mathcal{L} \supset M_{AB} U_A \bar{U}_B + \xi S_{iA} \bar{u}_i U_A$$

$$- [\eta_{ij}^u Q_i H \bar{u}_j - \eta_{ij}^d Q_i H^c \bar{d}_j + \text{h.c.}] + \mathcal{L}_{\text{BSM}}$$

No additional spurions/fields transforming under $U(3)_{\bar{u}}$

Introduce mixing between up-quark and heavy VLQs in a flavor breaking vacuum

	$U(3)_U$	$U(3)_{\bar{U}}$	$U(3)_{\bar{u}}$	$U(1)_B$	\mathbb{Z}_2
U	3			1/3	-1
\bar{U}		3		-1/3	-1
S	$\bar{3}$		$\bar{3}$		-1



Integrating out heavy quarks leads to wave-function renormalization of the SM up-quarks

$$Z_{ij}^u = \delta_{ij} + \frac{\xi^* \xi}{M_A^* M_A} S_{iA}^* S_{jA}$$

The source of all flavor-breaking!
CKM matrix arises from returning to canonical basis

S.Homiller@PPG2021

2HDM SFV

SFV physics may have generic family non-universal couplings to SM.

A simple framework: **SFV Two Higgs Doublet Models (2HDM).**

$$|D_\mu H_a|^2 - V(H_1, H_2) - \left(\mathcal{Y}_{aij}^u \bar{Q}_{Li} H_a U_{Rj} + \mathcal{Y}_{aij}^d \bar{Q}_{Li} H_a^c D_{Rj} + \mathcal{Y}_{aij}^\ell \bar{L}_{Li} H_a^c \ell_{Rj} + h.c. \right)$$

($a = 1, 2; i, j = 1, 2, 3$)

2HDM features neutral 2 CP-even H:
 $\cos(\beta - \alpha) \neq 0$ **interesting!**

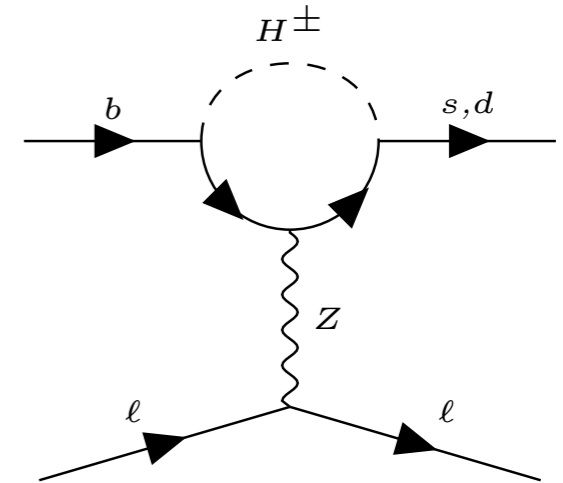
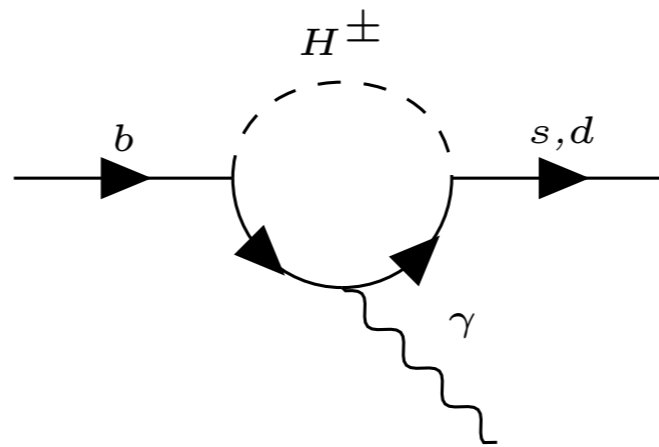
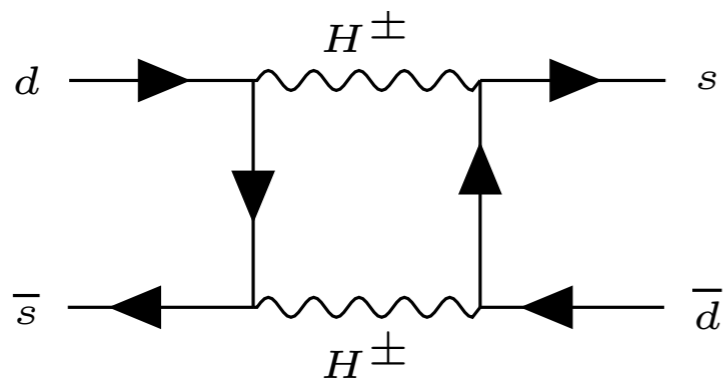
$$\begin{aligned} h &= H_1^0 \sin(\beta - \alpha) + H_2^0 \cos(\beta - \alpha) \\ H &= H_1^0 \cos(\beta - \alpha) - H_2^0 \sin(\beta - \alpha) \end{aligned}$$

+ a CP-odd and charged ones ... for simplicity:

$$m_H = m_{H^\pm} = m_A$$

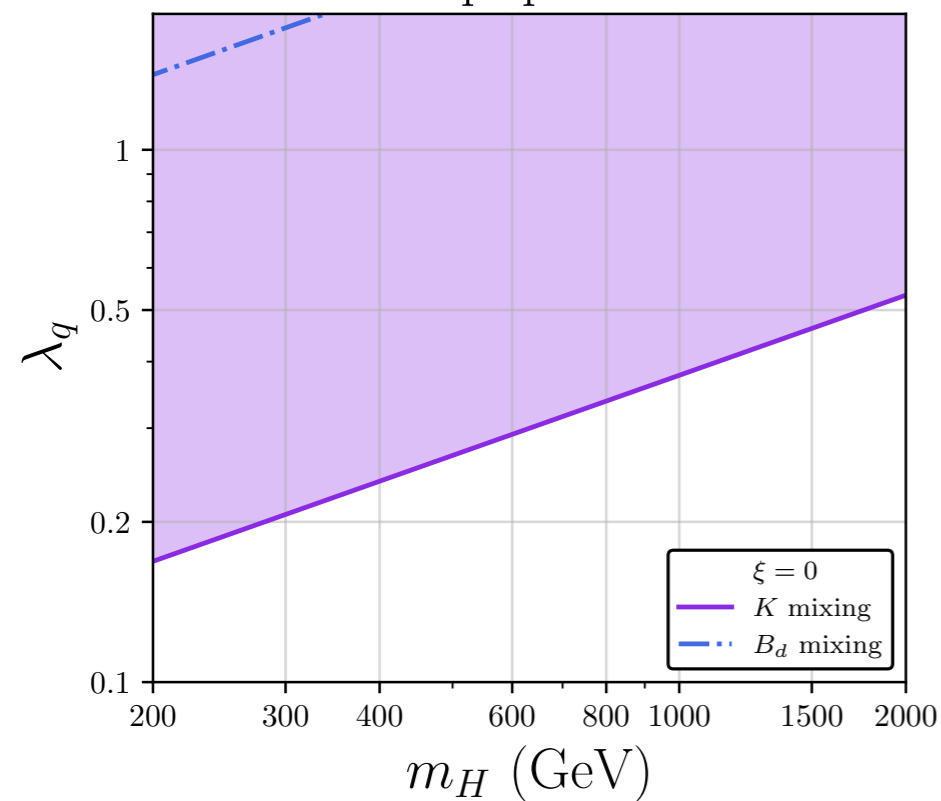
DOWN-TYPE SFV ANSATZ ON 2HDM

$$\begin{aligned} \mathcal{Y}_1^u &= y^u = \text{diag}(y_u, y_c, y_t) & \mathcal{Y}_1^d &= V y^d & \mathcal{Y}_1^\ell &= y^\ell \\ \mathcal{Y}_2^u &= \lambda^u = \text{diag}(\lambda_u, \lambda_c, \lambda_t) & \mathcal{Y}_2^d &= \xi V y^d & \mathcal{Y}_2^\ell &= \xi^\ell y^\ell \end{aligned}$$

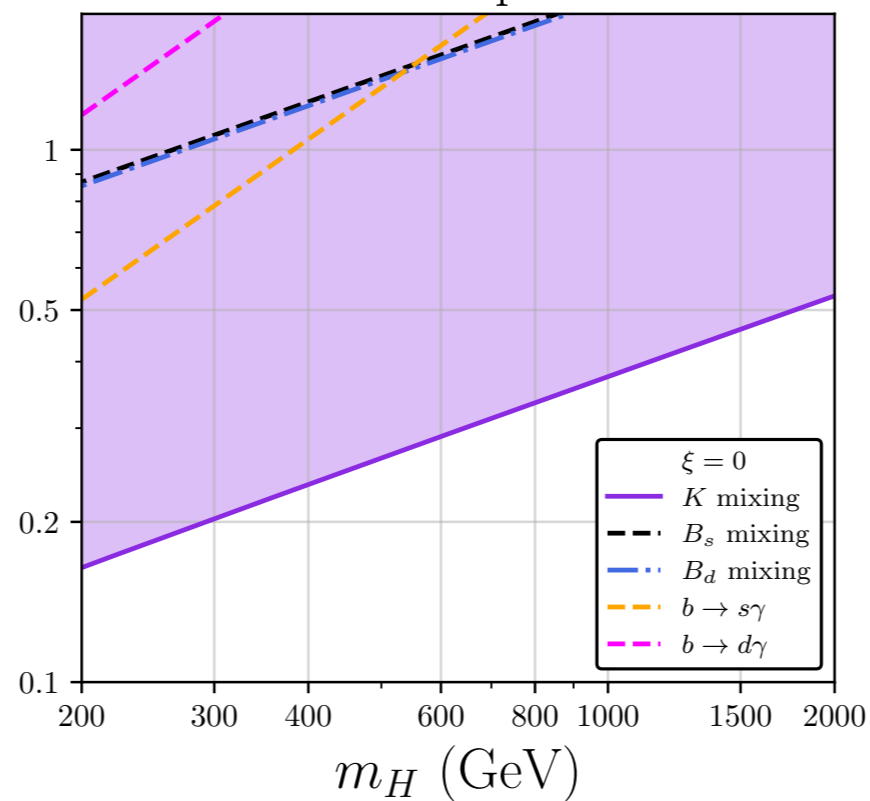


FLAVOR PROTECTION IN DOWN-TYPE SFV ALLOWS NEW HIGGSES
@ THE TEV TO COUPLE TO UP-TYPE QUARKS WITH ~ 0.1 STRENGTH.

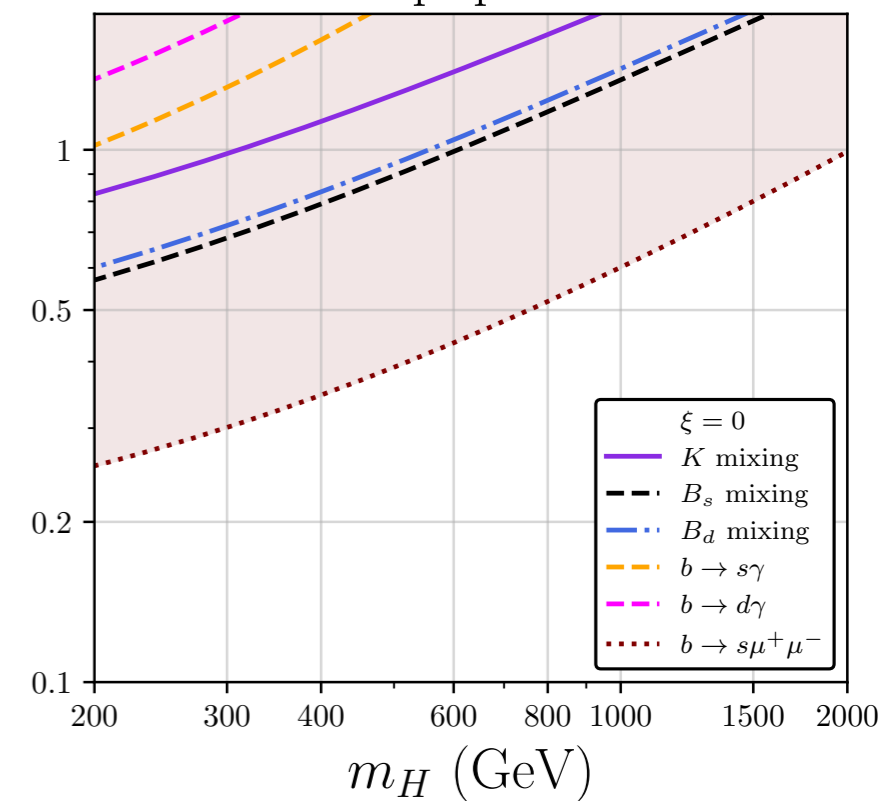
Up quark



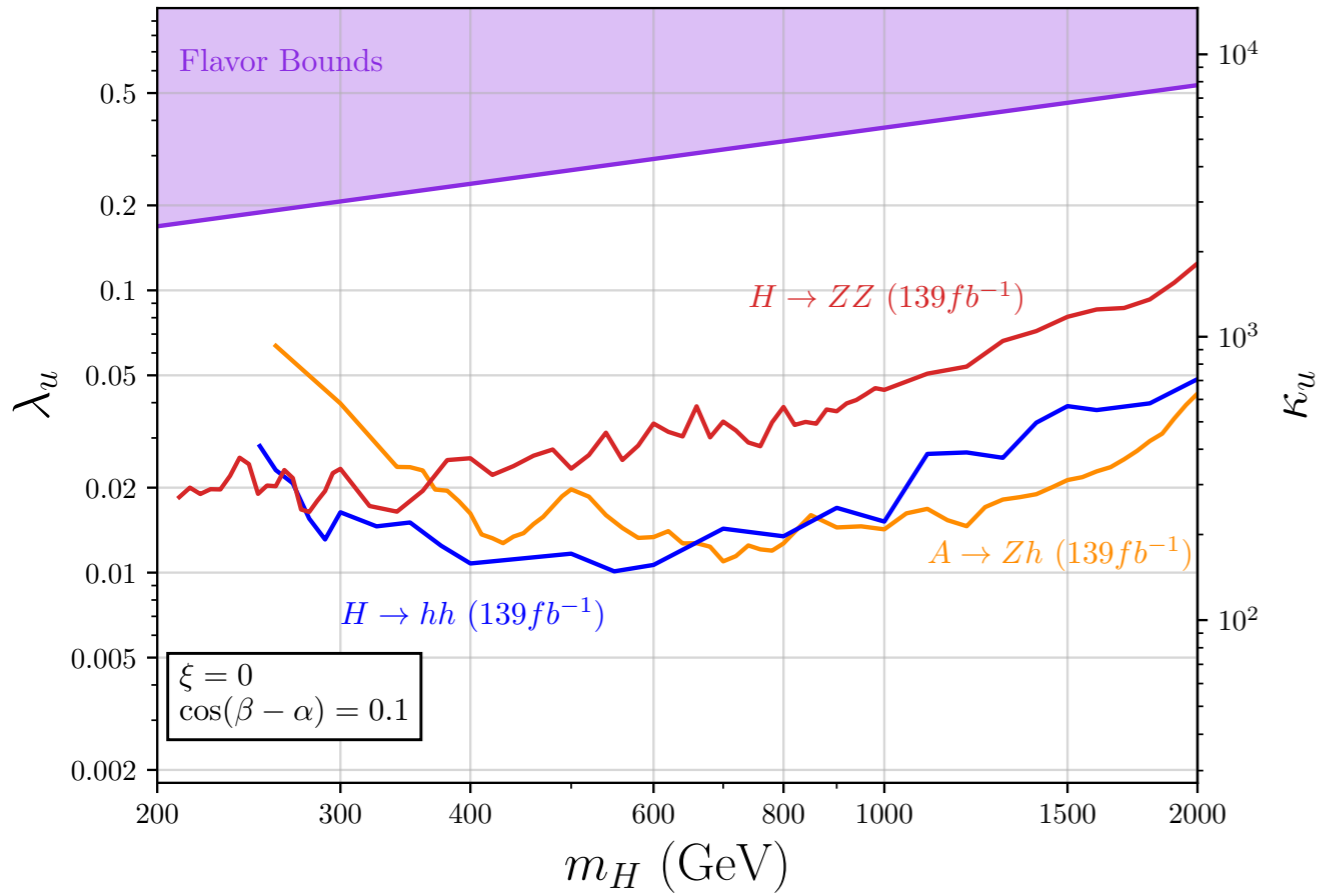
Charm quark



Top quark



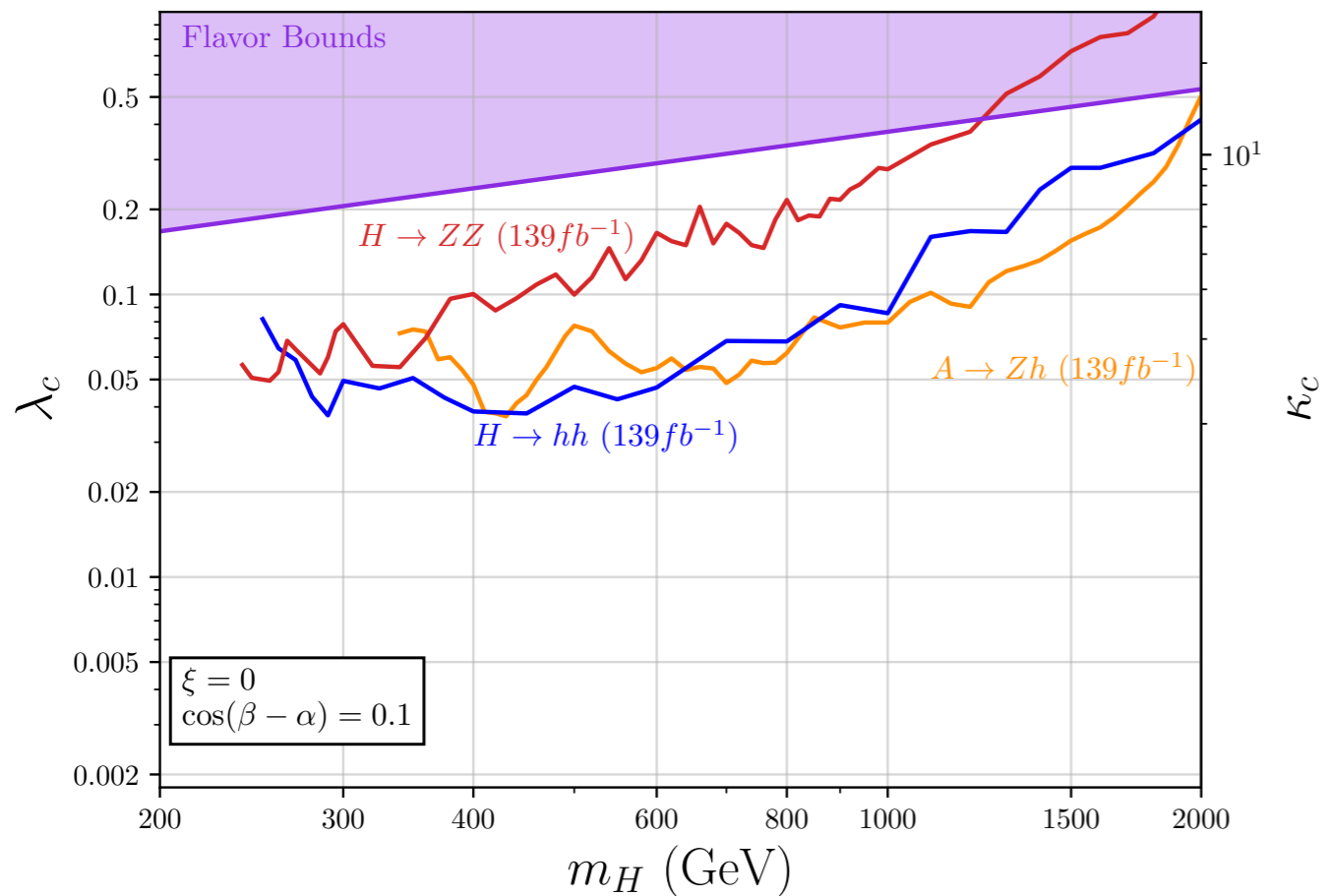
An analogous message holds for the up-type SFV 2HDM as well.



Di-Higgs production is a very sensitive probe of enhanced Higgs couplings to light quarks

arXiv:2101.04119

DOMINANT COLLIDER BOUNDS FROM SEARCH FOR NEW HIGGSES



- Mainly from quark fusion, due to possibly large Yukawa
- Mostly into di-boson & di-Higgs due to non-zero mixing angle

OBS.

$\kappa_{u,c} \equiv y_{u,c}/y_{u,c}^{\text{SM}}$ is the Yukawa modifier

LHC starting to probe Charm Yukawa!

N. Berger @ ICHEP 24

ATLAS and CMS $VH \rightarrow cc$

ATLAS $VH \rightarrow cc$

Simultaneous fit with $VH \rightarrow bb$

$$\mu_{VH \rightarrow cc} < 11.3 @ 95\% CL \text{ (10.4 exp.)}$$

Best limit to date

Factor 2.5 improvement over previous limit !

More in
Francesco Di Bello's talk

$$\Rightarrow |\kappa_c| < 4.2 @ 95\% CL$$

Factor 2 improvement over previous



CMS $VH \rightarrow cc$:

Includes boosted $H \rightarrow cc$ ($p_T^H > 300$ GeV)

$$\mu_{VH \rightarrow cc} < 14 (7.6) @ 95\% CL \text{ best sensitivity}$$

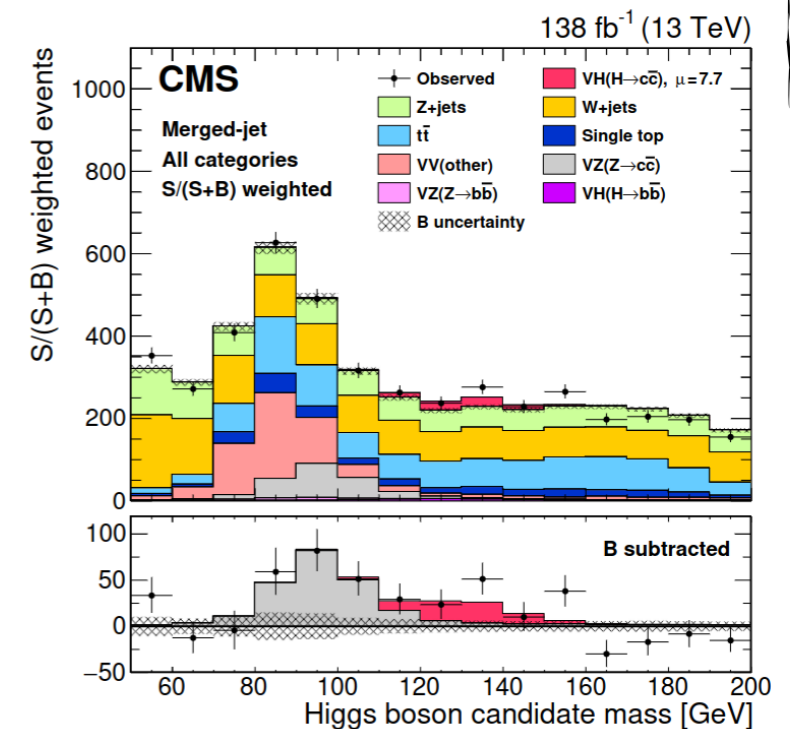
$$\Rightarrow 1.1 < |\kappa_c| < 5.5$$

First observation of $Z \rightarrow cc$ in hadronic collisions.

PRL 131 (2023) 041801,

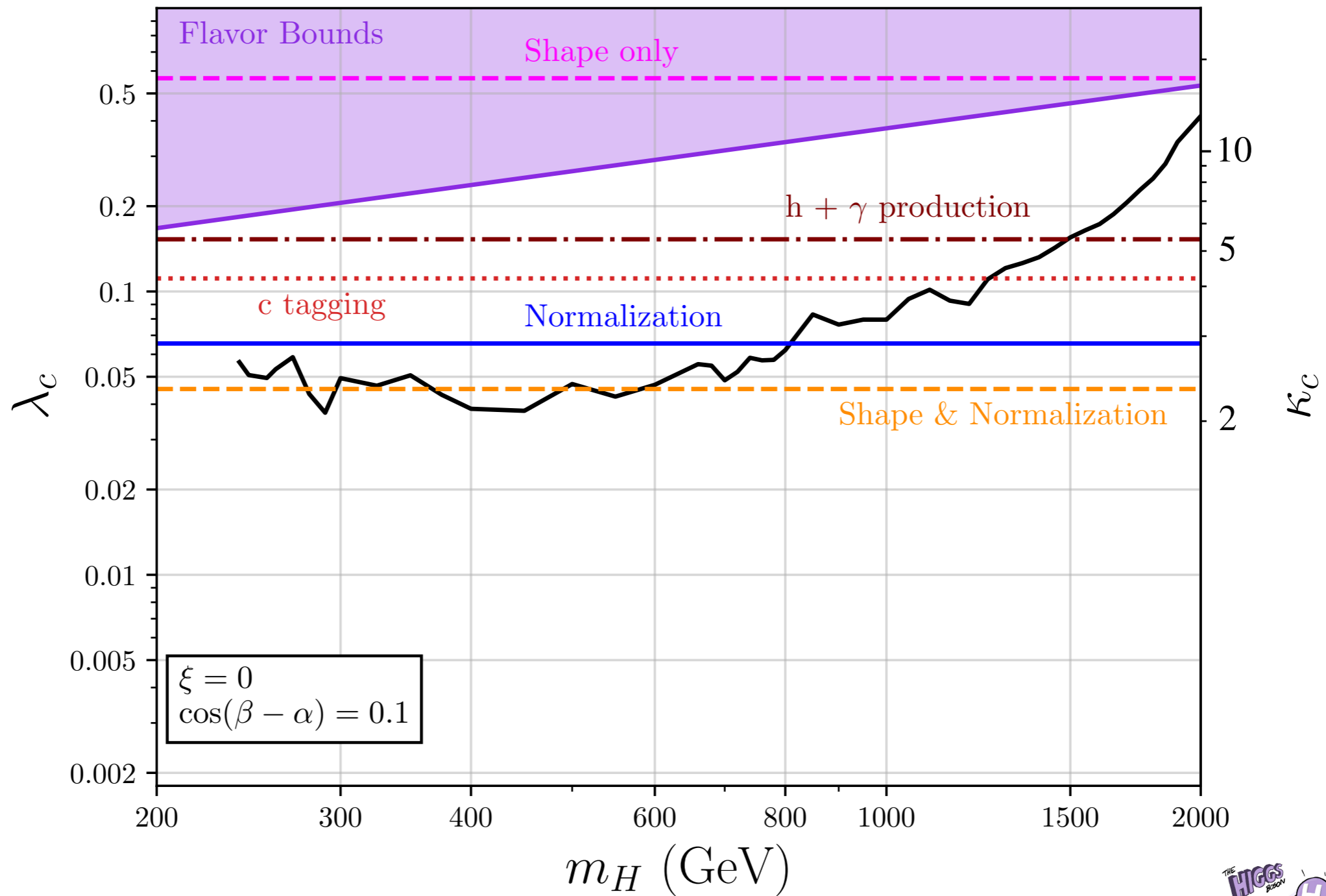
PRL 131 (2023) 061801

More in [Andrea Cardini's talk](#) and [Maarten de Coen's poster](#)



The advent of Deep Learning has been a game changer for ATLAS & CMS sensitivity to 2nd gen Yukawa couplings.

Charming Higgs @ LHC





● **Program of precision with flavor is still very rich, but keeps pointing to very high scales for NP.**

No clear tension emerging from flavor data

● **In absence of signs of NP, the leading paradigms we relied on for BSM should be reconsidered.**

MFV ansatz **vs** SFV models

● **TH + EXP progress can lead to interesting new targets.** *E.g., ~~This Charming Man~~*

[See also 2410.08272] *Higgs!*




TPPC

**LET'S
DO IT
AGAIN**

