



ALMA MATER STUDIORUM
UNIVERSITÀ DI BOLOGNA

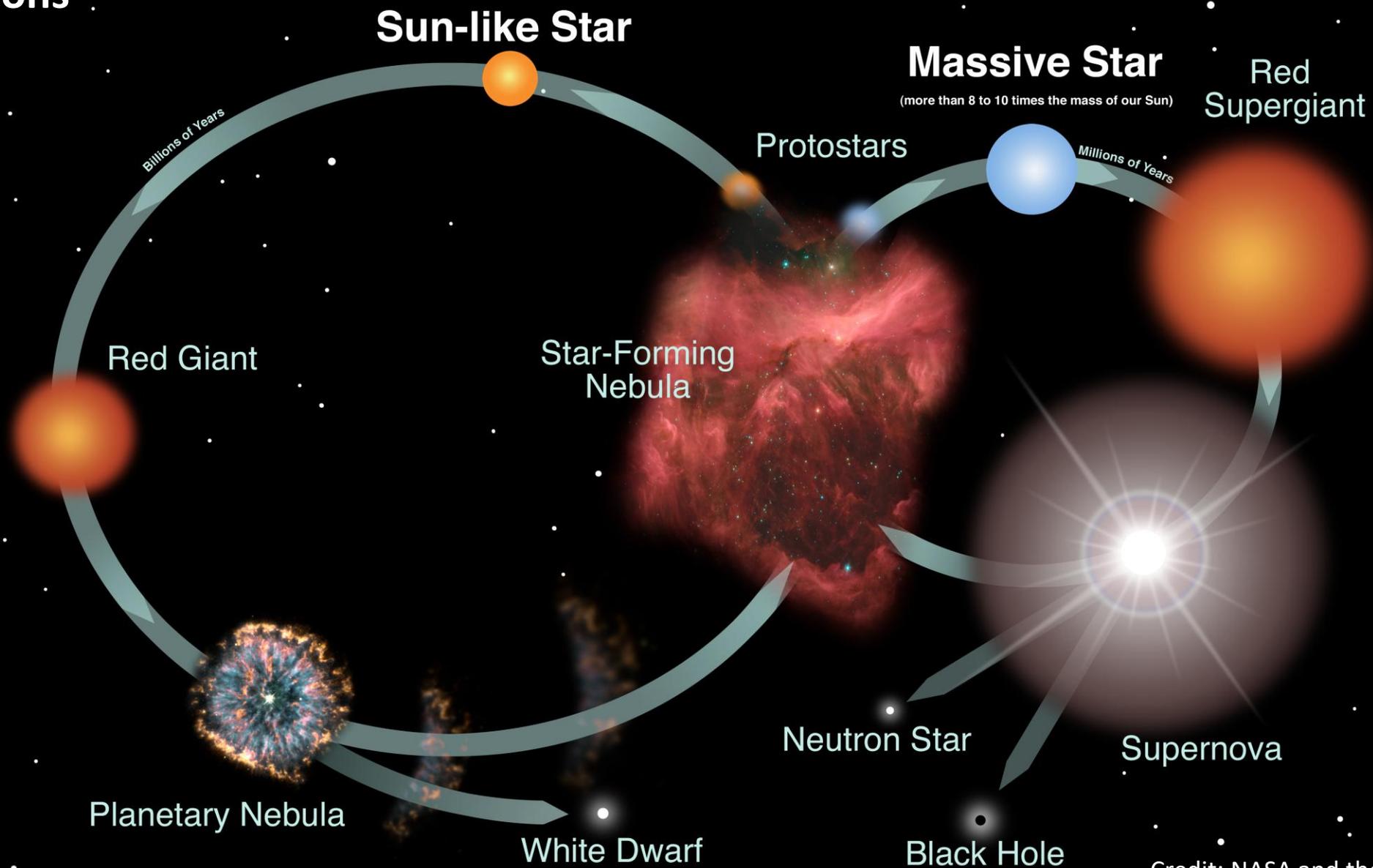


Neutron spectroscopy of ^{26}Mg states: constraining the $^{22}\text{Ne}(\alpha, n)^{25}\text{Mg}$ reaction rate

Nicholas Pieretti on behalf of the n_TOF Collaboration

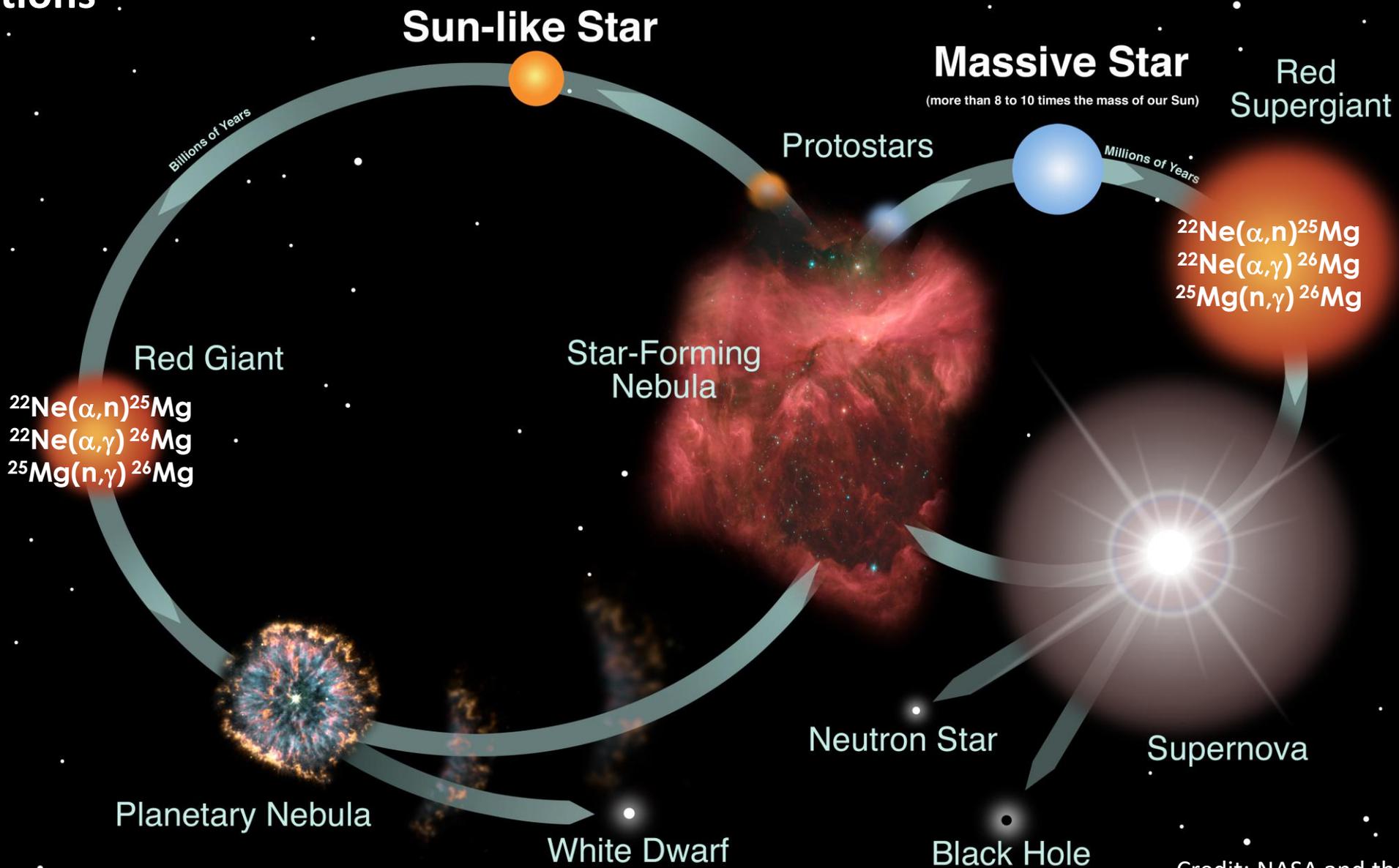
Department of Physics and Astronomy

Motivations



Credit: NASA and the Night Sky Network

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Neutron spectroscopy of ^{26}Mg states: constraining the $^{22}\text{Ne}(\alpha, n)^{25}\text{Mg}$ reaction rate

Motivations

Neutron poison:

- $^{25,26}\text{Mg}$ are key **neutron poisons** during the s-process.
- They compete with ^{56}Fe (basic s-process seed) via $^{25,26}\text{Mg}(n,\gamma)$.

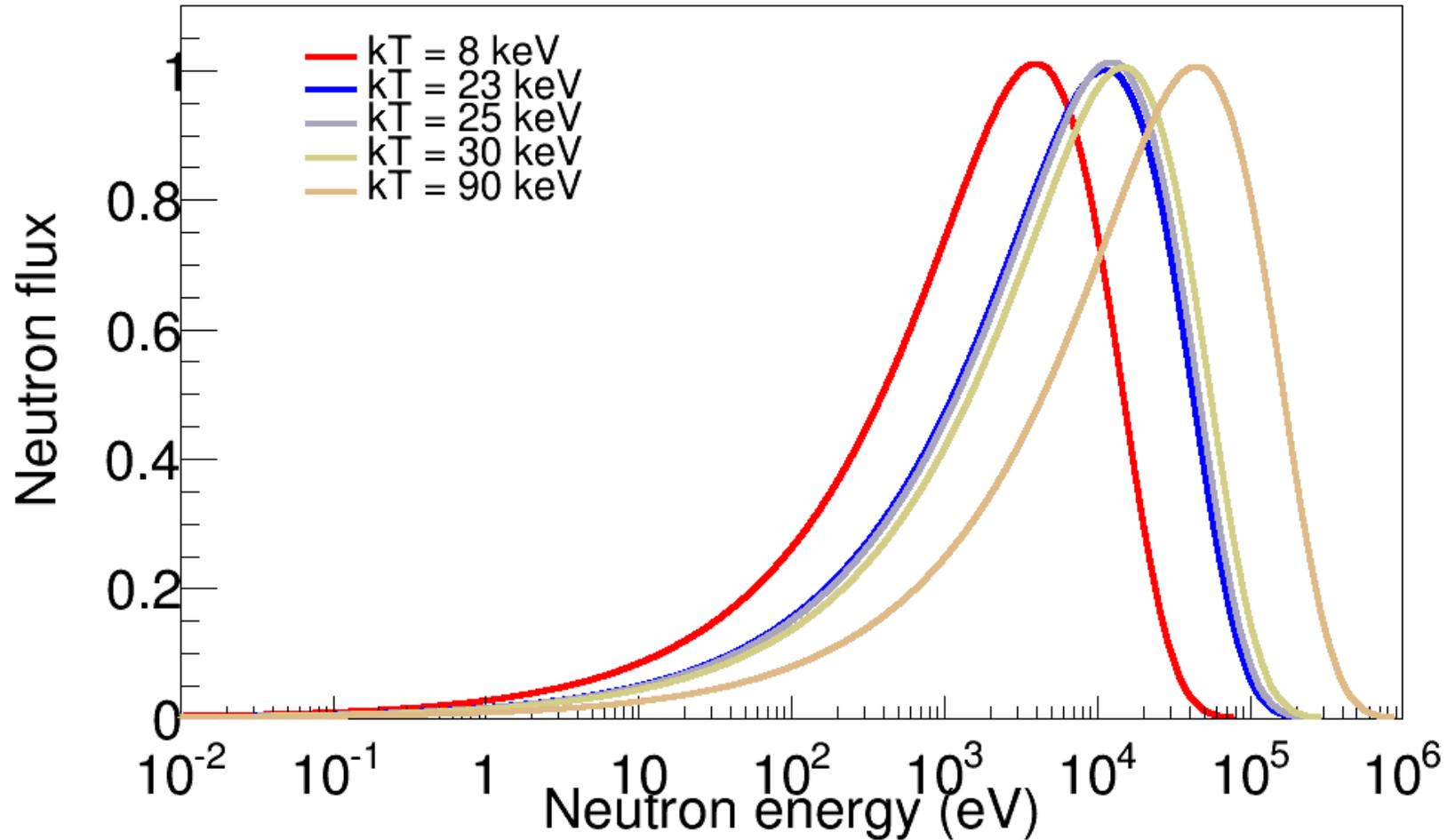
Constraints for $^{22}\text{Ne}(\alpha,n)^{25}\text{Mg}$ and $^{22}\text{Ne}(\alpha,\gamma)^{26}\text{Mg}$:

- $^{22}\text{Ne}(\alpha,n)^{25}\text{Mg}$: primary **neutron source** in Red Giant stars
- Its reaction rate is very uncertain because of poor knowledge of ^{26}Mg **states**.
- From neutron measurements the **energy** and J^π of ^{26}Mg states can be deduced, together with Γ_γ and Γ_n .



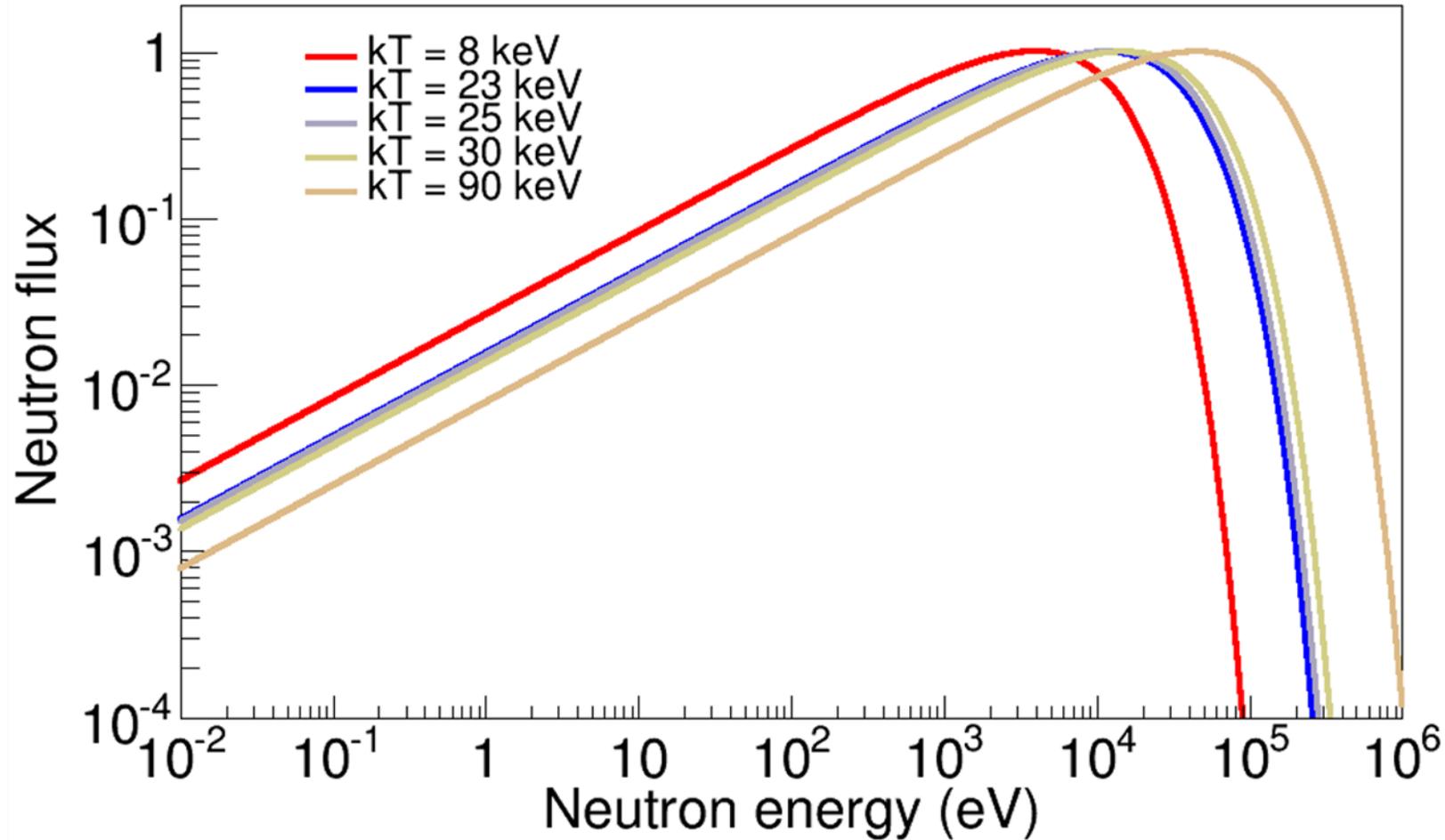
Measurement of $^{25}\text{Mg}(n,\gamma) \leftrightarrow ^{22}\text{Ne}(\alpha,n)^{25}\text{Mg}$

Stellar spectra: AGB (8, 23 keV) and Massive stars (25, 90 keV)



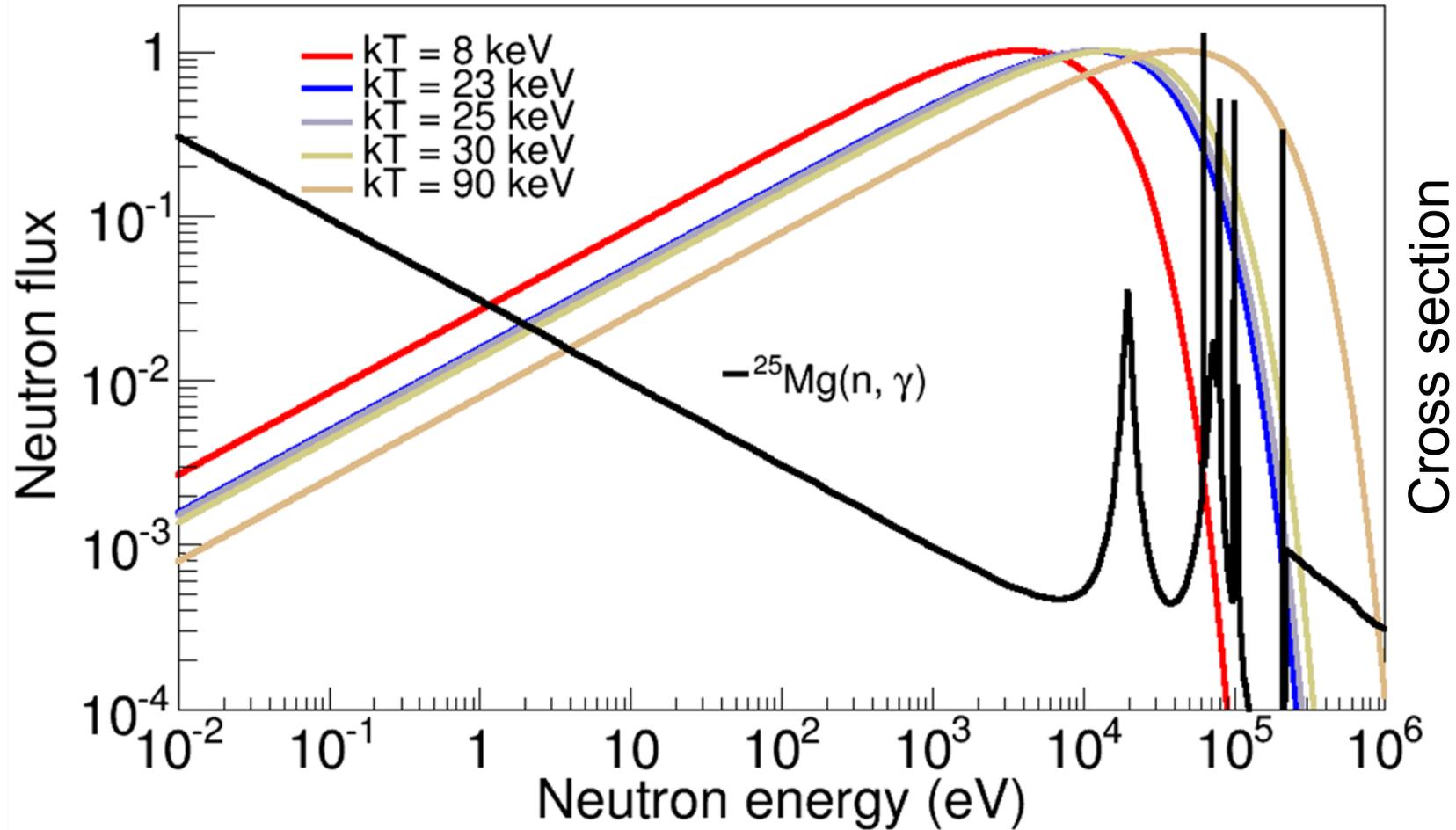
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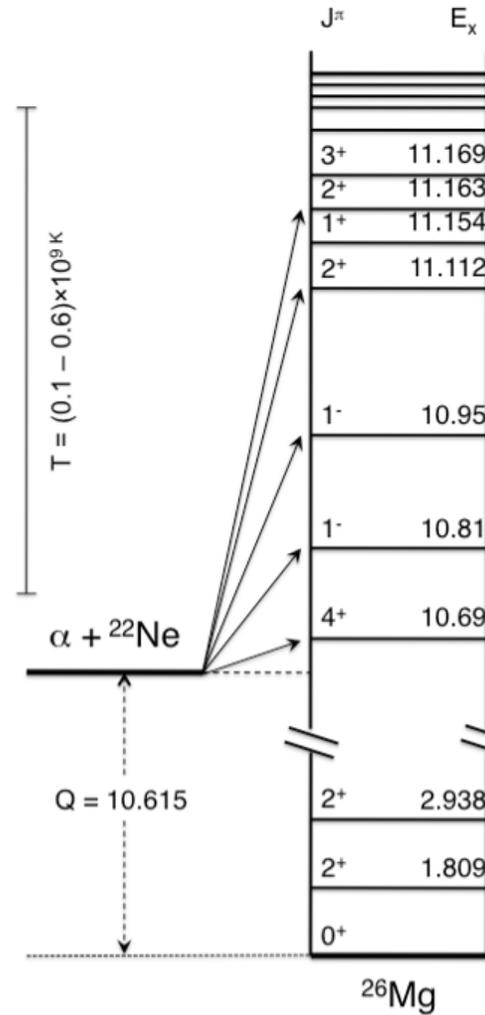
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$\alpha + ^{22}\text{Ne}$

α	$J^\pi = 0^+$
^{22}Ne	$J^\pi = 0^+$

Only natural-parity states in ^{26}Mg can participate in the $^{22}\text{Ne}(\alpha,n)^{25}\text{Mg}$ reaction:

$J^\pi = 0^+, 1^-, 2^+, 3^-, 4^+ \dots$



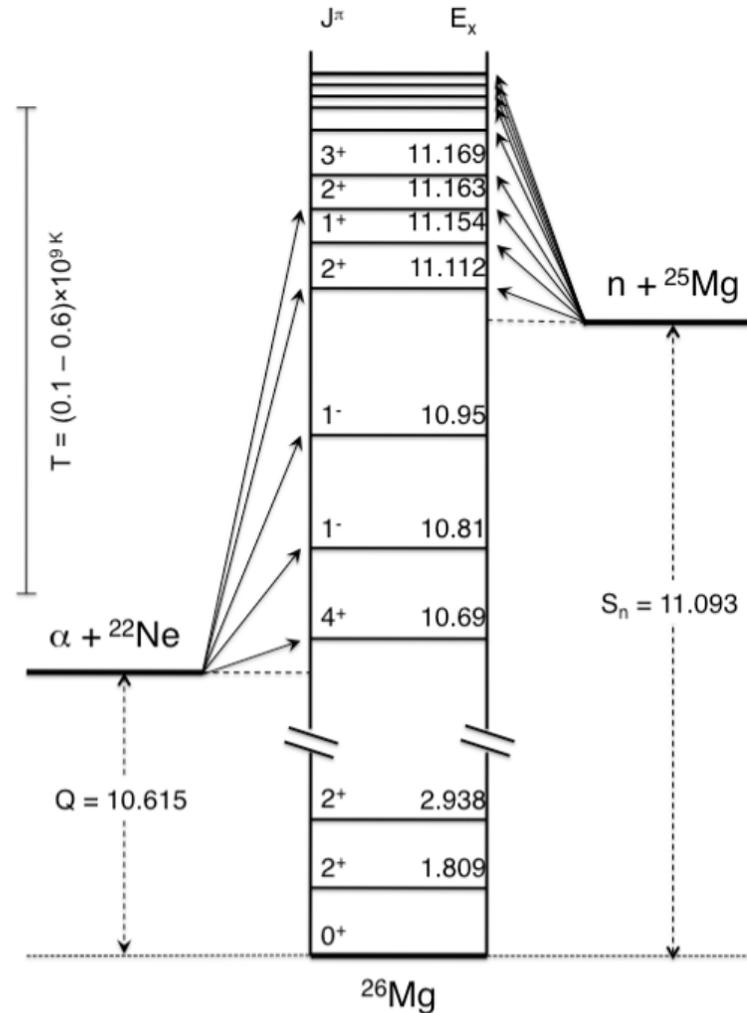
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$n + ^{25}\text{Mg}$

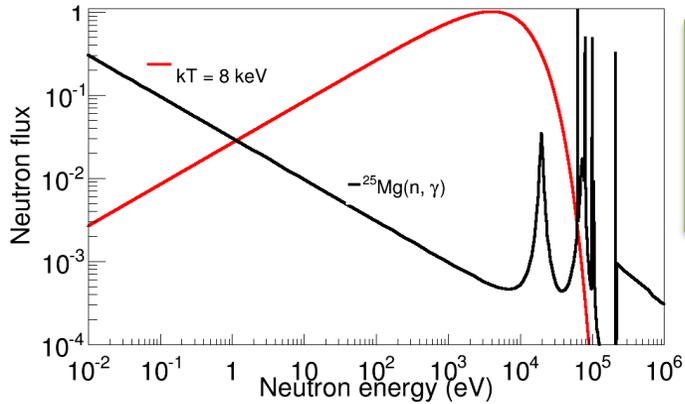
n	$J^\pi = 1/2^+$
^{25}Mg	$J^\pi = 5/2^+$

All states in ^{26}Mg can participate in the $^{25}\text{Mg}(n,\gamma)^{26}\text{Mg}$ reaction:

$J^\pi = 0^+, 0^-, 1^+, 1^-, 2^+, 2^-, \dots$



Neutron-induced reactions



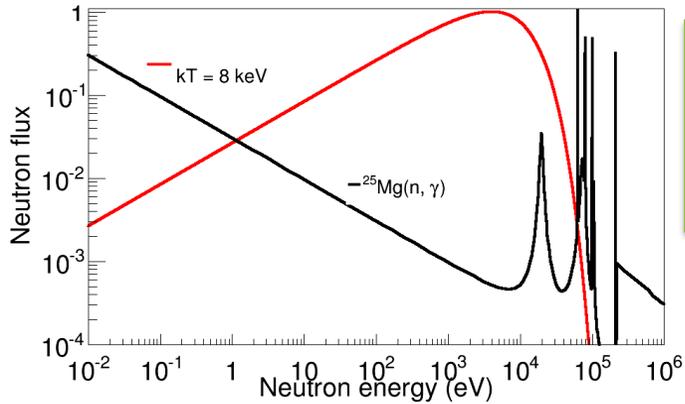
For nuclear astrophysics, what is important is the **Maxwellian Averaged Cross-Sections (MACS)** at various **temperatures** (kT depends on stellar site).

Reaction rate ($\text{cm}^{-3}\text{s}^{-1}$): $r = N_A N_n v \sigma(v) \longrightarrow r = N_A N_n \langle \sigma \cdot v \rangle$

$$MACS \equiv \frac{\langle \sigma \cdot v \rangle}{v_T} = \frac{2}{\sqrt{\pi}(kT)^2} \int_0^{\infty} \sigma(E) E e^{-E/(kT)} dE$$



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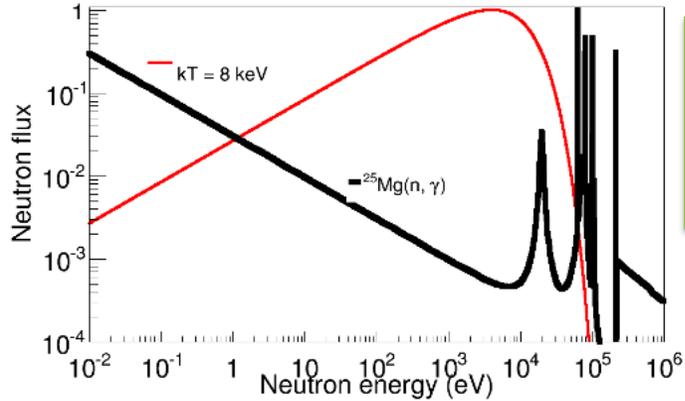
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Two methods to determine MACS:



Neutron-induced reactions



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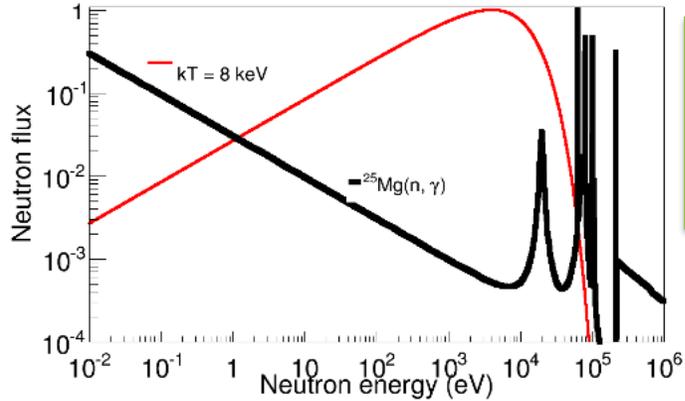
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Two methods to determine MACS:

1. measurement of **energy dependent** neutron capture cross-sections \rightarrow EAR1 & EAR2



Neutron-induced reactions



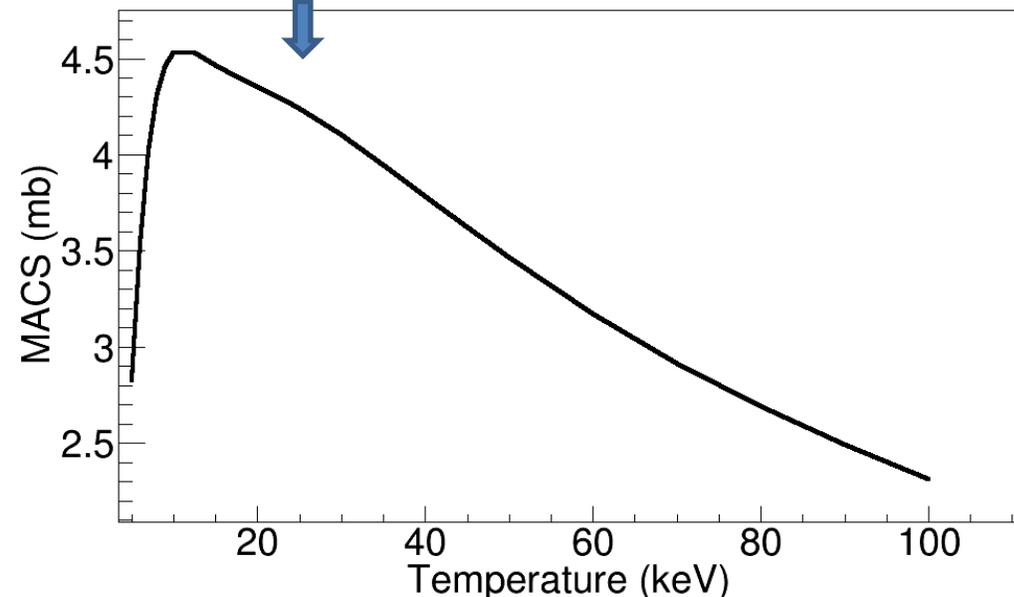
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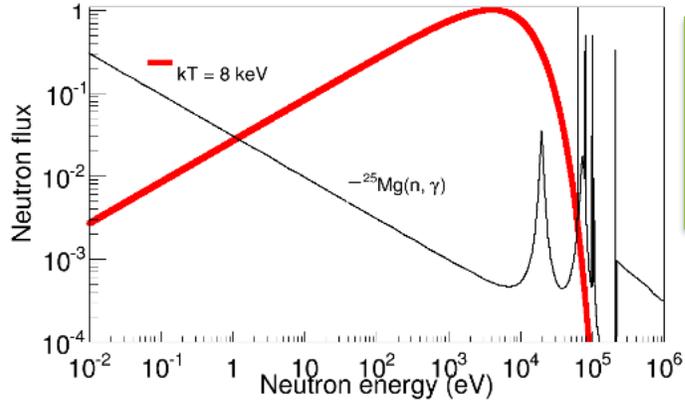
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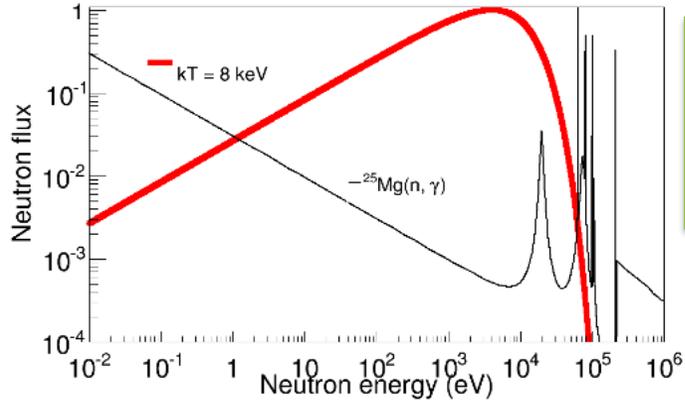
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2. **integral measurement** (energy integrated) using neutron beams with suitable energy \rightarrow **NEAR**



Neutron-induced reactions



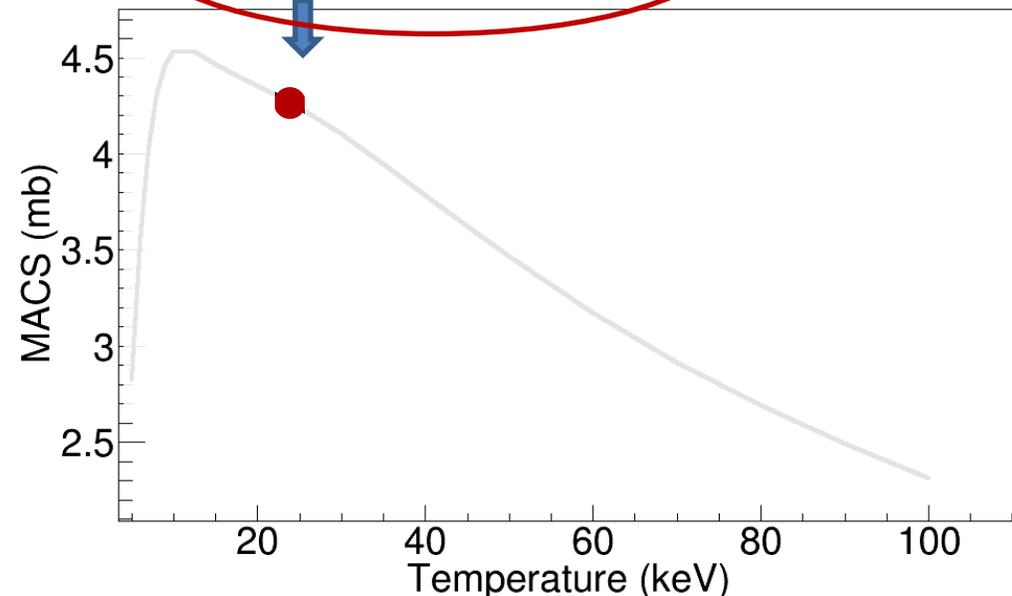
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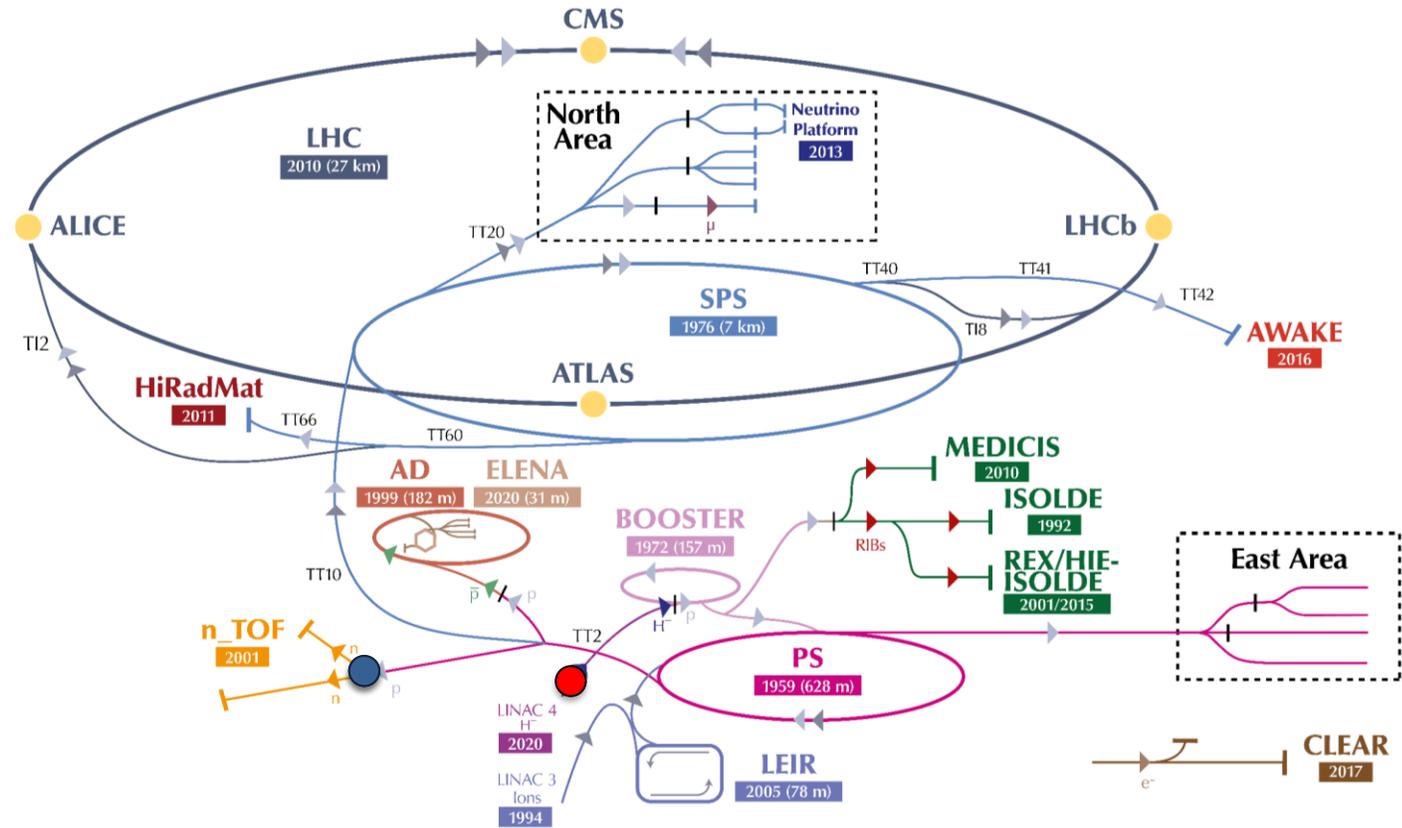
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The CERN accelerator complex

Complexe des accélérateurs du CERN



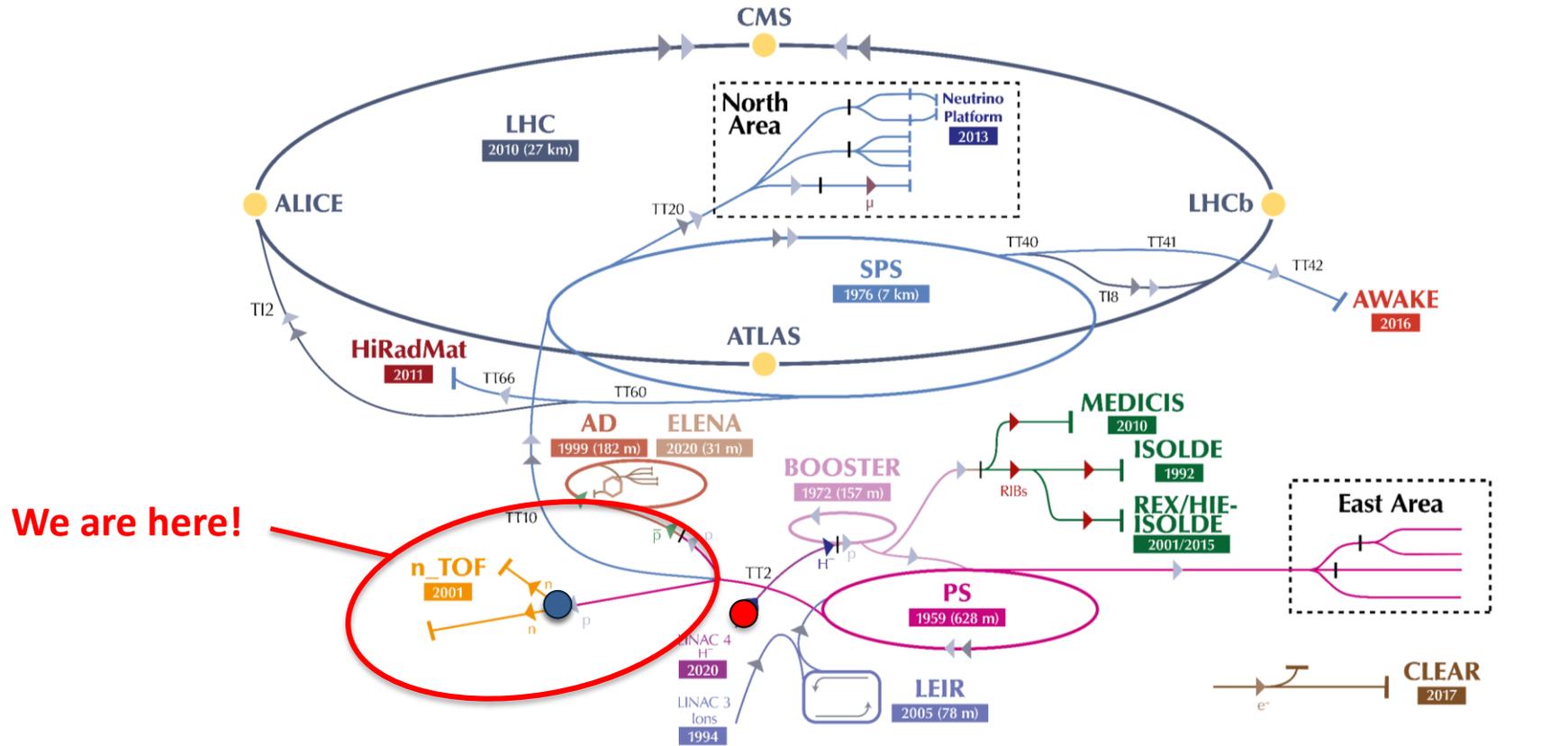
▶ H^- (hydrogen anions) ▶ p (protons) ▶ ions ▶ RIBs (Radioactive Ion Beams) ▶ n (neutrons) ▶ \bar{p} (antiprotons) ▶ e^- (electrons) ▶ μ (muons)

LHC - Large Hadron Collider // SPS - Super Proton Synchrotron // PS - Proton Synchrotron // AD - Antiproton Decelerator // CLEAR - CERN Linear Electron Accelerator for Research // AWAKE - Advanced WAKEfield Experiment // ISOLDE - Isotope Separator OnLine // REX/HIE-ISOLDE - Radioactive Experiment/High Intensity and Energy ISOLDE // MEDICIS // LEIR - Low Energy Ion Ring // LINAC - LINear ACcelerator // n_TOF - Neutrons Time Of Flight // HiRadMat - High-Radiation to Materials // Neutrino Platform



The CERN accelerator complex

Complexe des accélérateurs du CERN



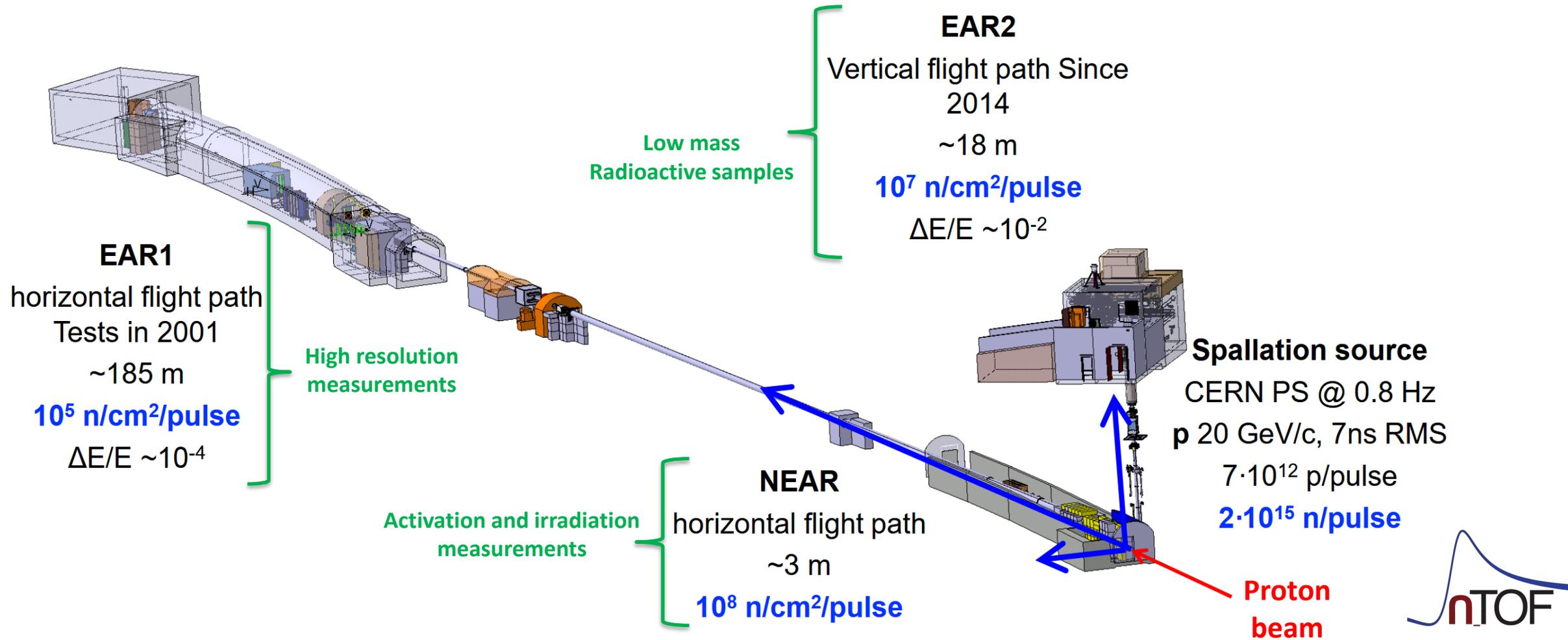
We are here!

▶ H^- (hydrogen anions) ▶ p (protons) ▶ ions ▶ RIBs (Radioactive Ion Beams) ▶ n (neutrons) ▶ \bar{p} (antiprotons) ▶ e^- (electrons) ▶ μ (muons)

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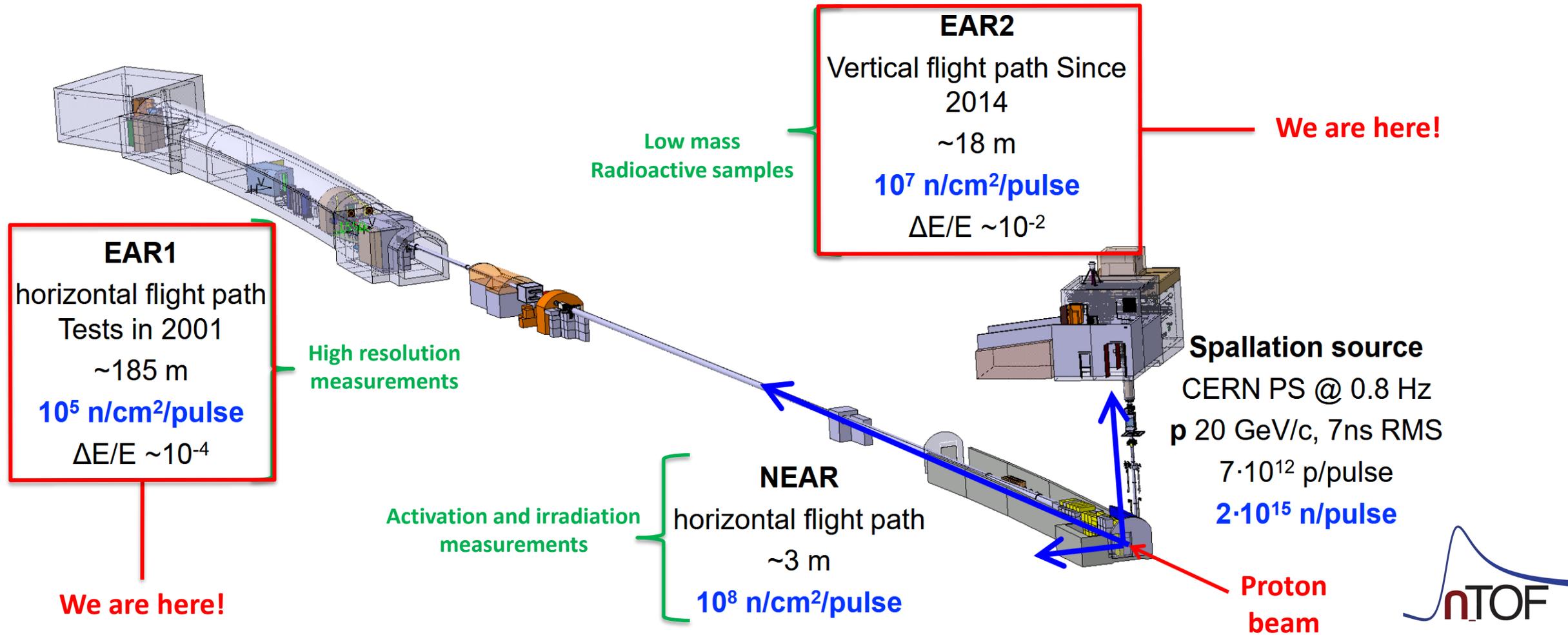
Why n_TOF?



From E. Mendoza, APRENDE WP2-WP4 Workshop



Why n_TOF?

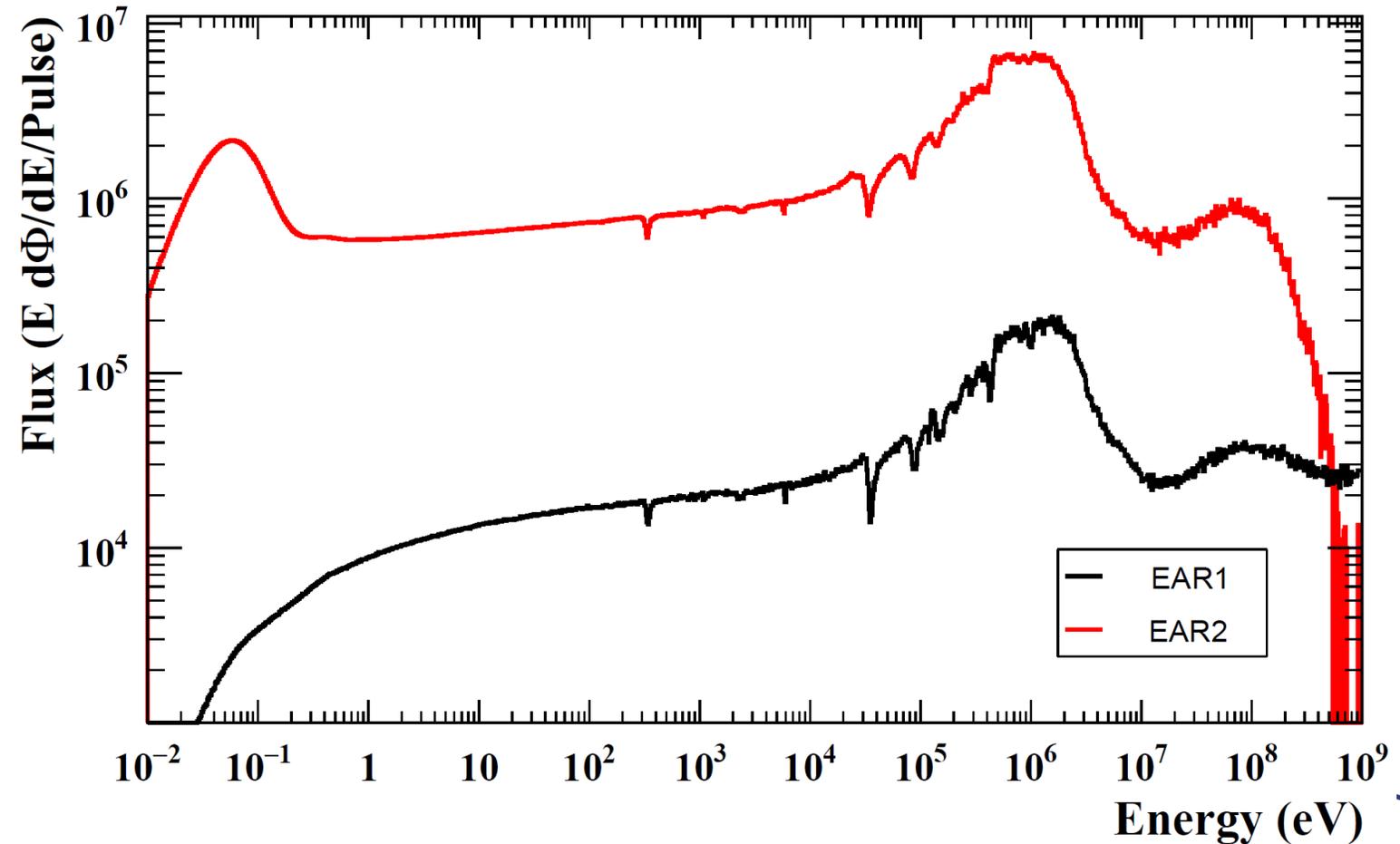


From E. Mendoza, APRENDE WP2-WP4 Workshop



Why n_TOF?

- Wide energy range:
 $10 \text{ meV} \leq E_n \leq 1 \text{ GeV}$
- High current:
 $8.5 \times 10^{12} \text{ p/bunch}$
 $\rightarrow 10^6 \text{ n/pulse}$
- Energy resolution at EAR1:
 $\frac{\Delta E_n}{E_n} = 0.03\% (1 \text{ eV})$
 $\frac{\Delta E_n}{E_n} = 0.5\% (1 \text{ MeV})$



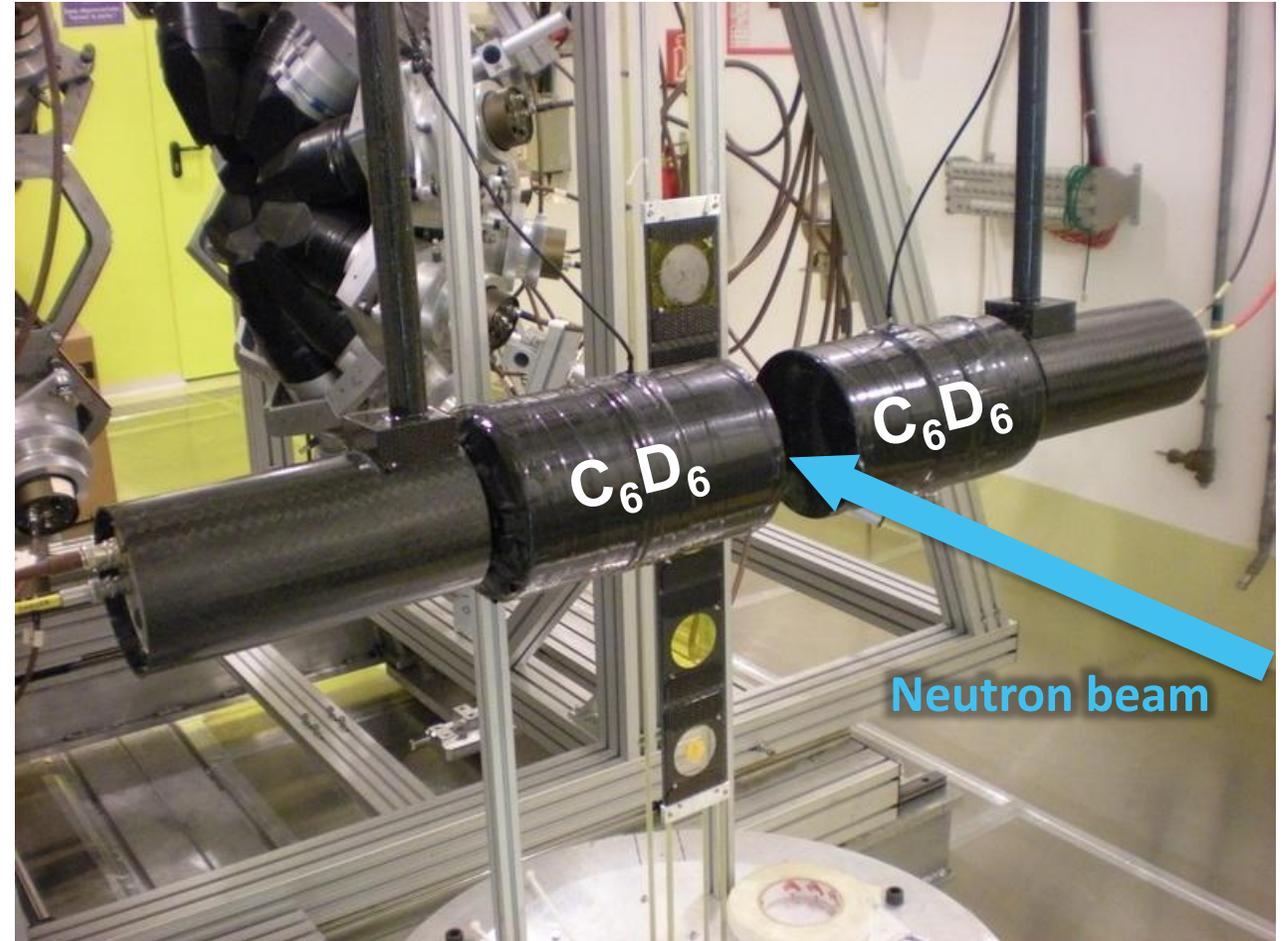
Previous measurements of $^{25}\text{Mg}(n,\gamma)$ at n_TOF

Capture setup:

- 2 C_6D_6 liquid scintillators
- Total Energy Detection System based on PHWT

Mg Sample:

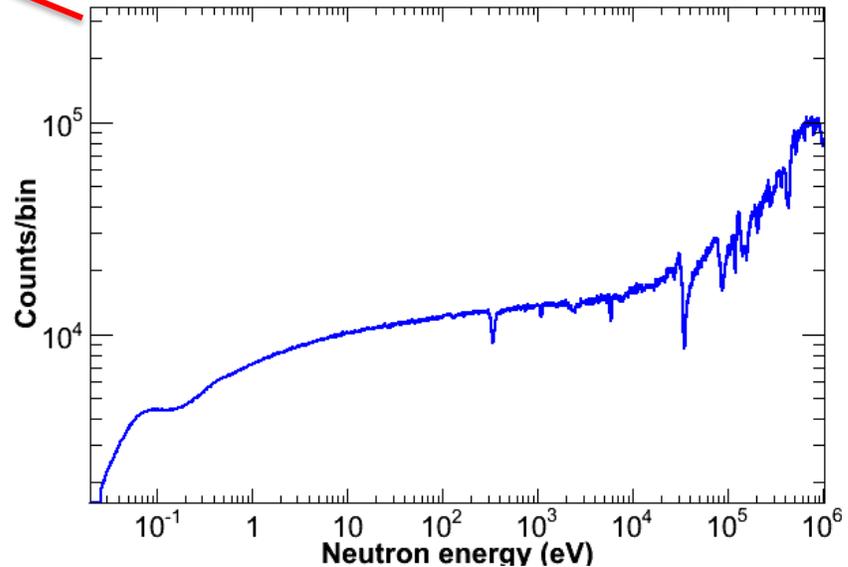
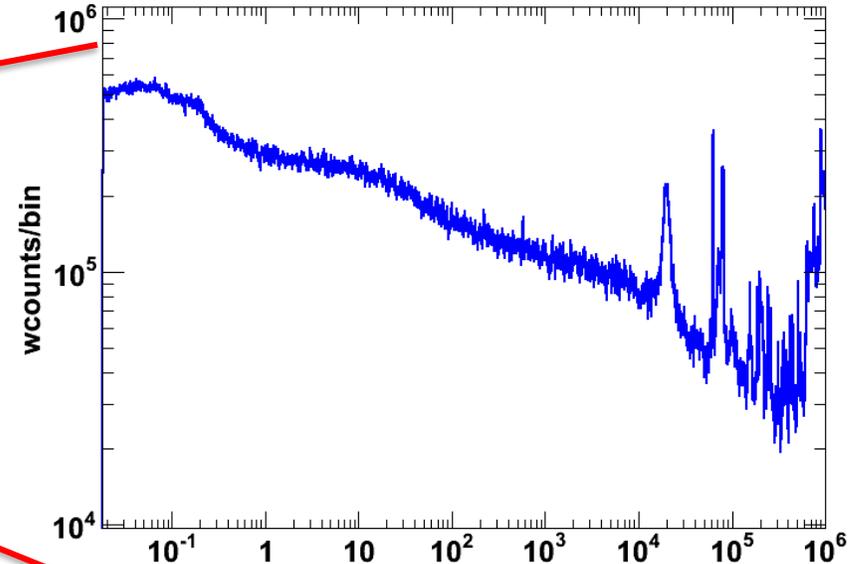
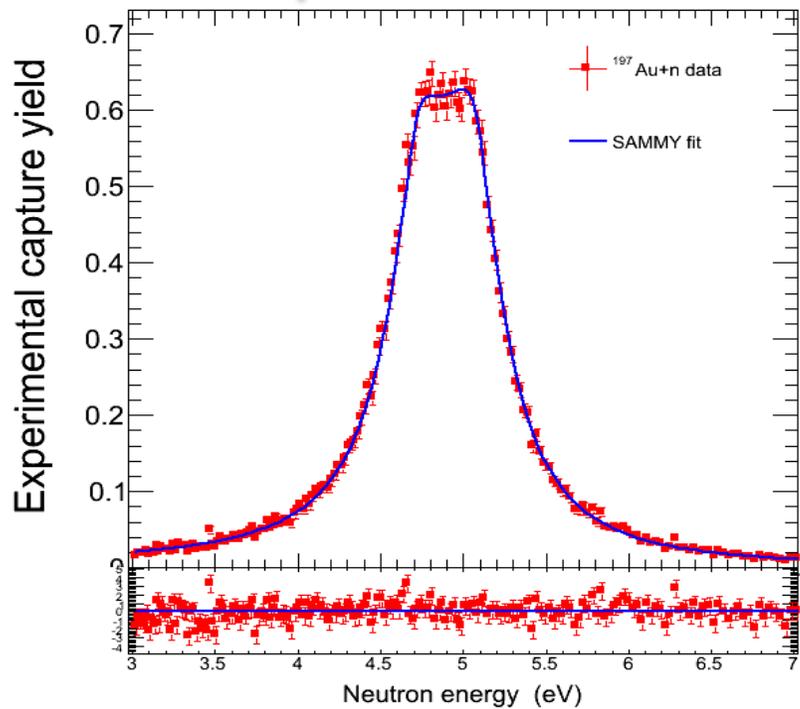
- 3.94 g, 2 cm diameter
- Enrichment 97.86 %
- 3.00×10^{-2} at/b



Previous measurements of $^{25}\text{Mg}(n,\gamma)$ at n_TOF

Experimental capture yield

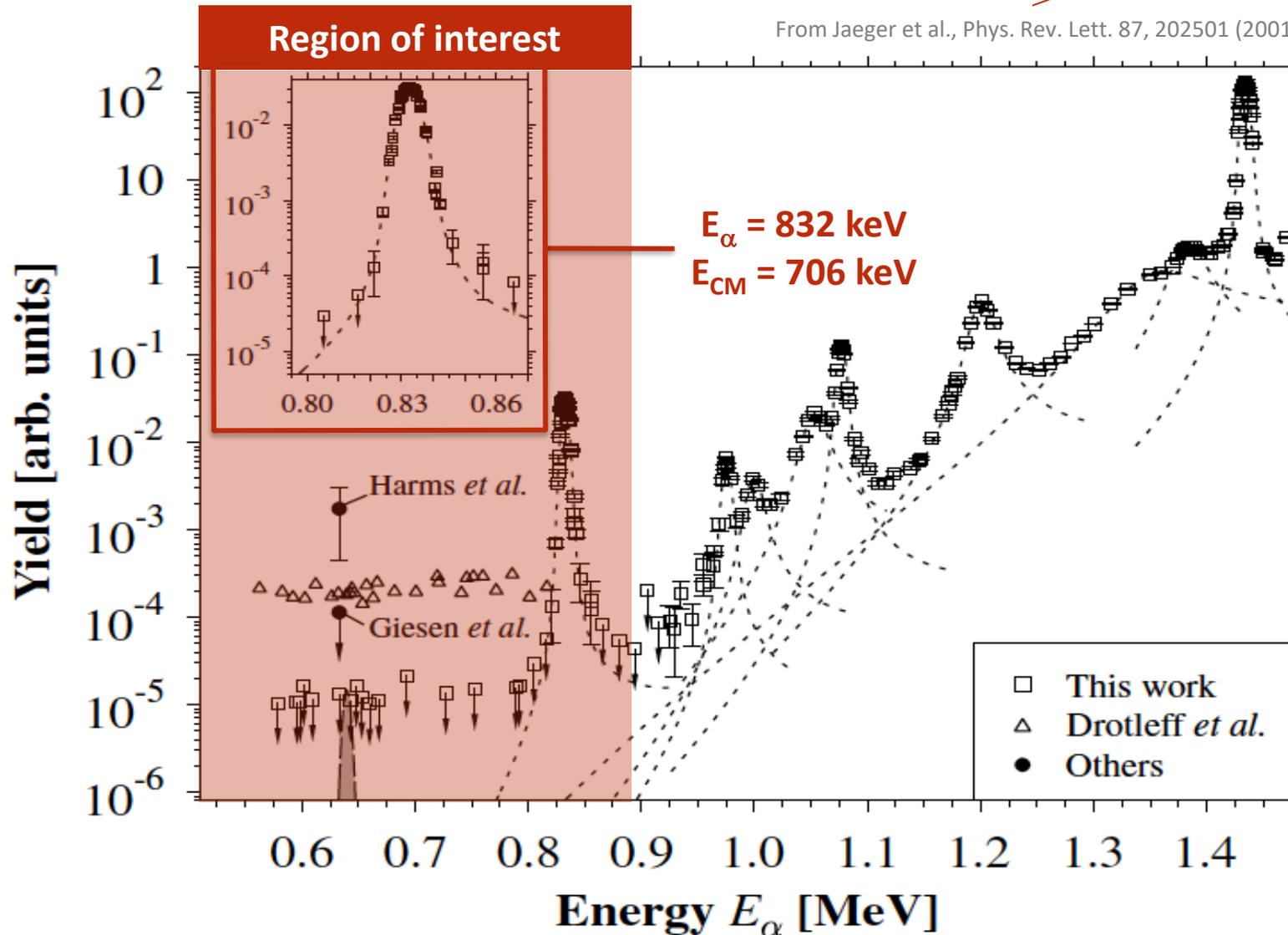
$$Y(E_n) = N \frac{C_w(E_n)}{\varphi_n(E_n)} \propto (1 - e^{-n\sigma_{tot}}) \frac{\sigma_\gamma}{\sigma_{tot}}$$



Measurement of $^{25}\text{Mg}(n,\gamma) \leftrightarrow ^{22}\text{Ne}(\alpha,n)^{25}\text{Mg}$

- Direct measurement of $^{22}\text{Ne}(\alpha,n)^{25}\text{Mg}$ in Gamow Window
- 100-150 μA He^+ beam incident on a ^{22}Ne gas jet target

From Jaeger et al., Phys. Rev. Lett. 87, 202501 (2001)



Data in the literature

$^{22}\text{Ne}(\alpha,n)^{25}\text{Mg}$



Stellar cross sections (MACS) for the s-process



Reevaluation of the $^{22}\text{Ne}(\alpha, \gamma)^{26}\text{Mg}$ and $^{22}\text{Ne}(\alpha, n)^{25}\text{Mg}$ reaction rates

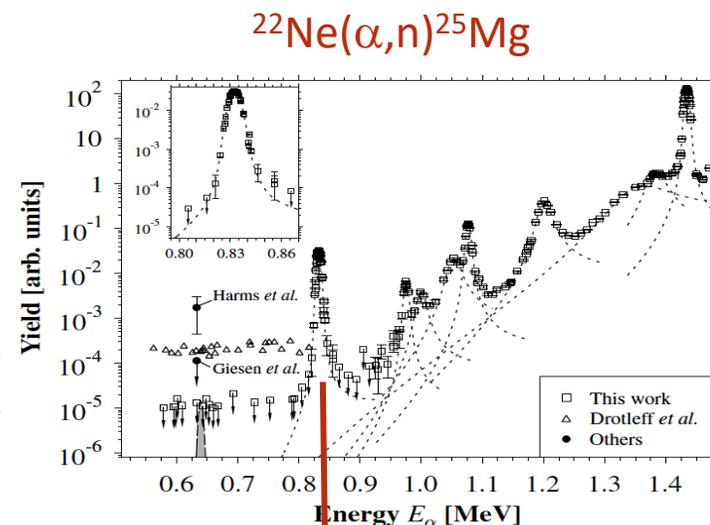
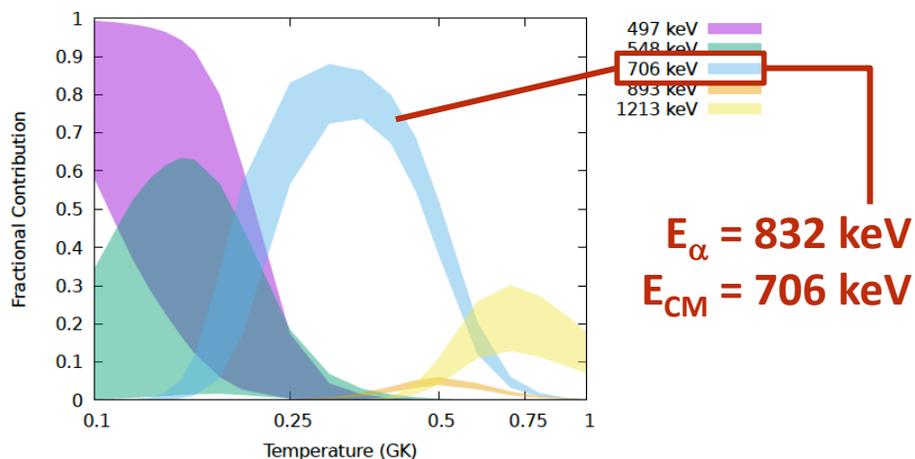
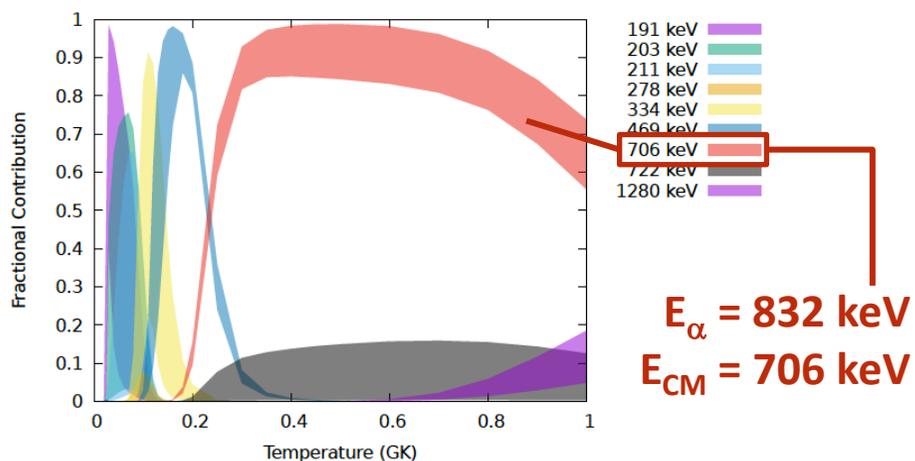
Philip Adsley^{1,2,3,*}, Umberto Battino^{4,†}, Andreas Best^{5,6}, Antonio Cacioli^{7,8}, Alessandra Guglielmetti⁹, Gianluca Imbriani^{5,6}, Heshani Jayatissa¹⁰, Marco La Cognata¹¹, Livio Lamia^{12,11,13} *et al.*

Show more

Phys. Rev. C **103**, 015805 – Published 19 January, 2021

DOI: <https://doi.org/10.1103/PhysRevC.103.015805>

FIG. 1. Fractional contributions of selected resonances to the (top) $^{22}\text{Ne}(\alpha, \gamma)^{26}\text{Mg}$ and (bottom) $^{22}\text{Ne}(\alpha, n)^{25}\text{Mg}$ reaction rates. These fractional contributions are for the recommended reaction rates, which incorporate the Texas A&M results. The shaded region gives the 68% coverage limit for the contribution of each resonance. Note that only the most significant resonances are included in the figure; the sum of the contributions may not reach 100% due to contributions from omitted resonances.



$E_\alpha = 832 \text{ keV}$
 $E_{\text{CM}} = 706 \text{ keV}$



$^{25}\text{Mg}(n,\gamma)$ for neutron source reaction in stars

Resonance strength $^{22}\text{Ne}(\alpha,n)^{25}\text{Mg}$:

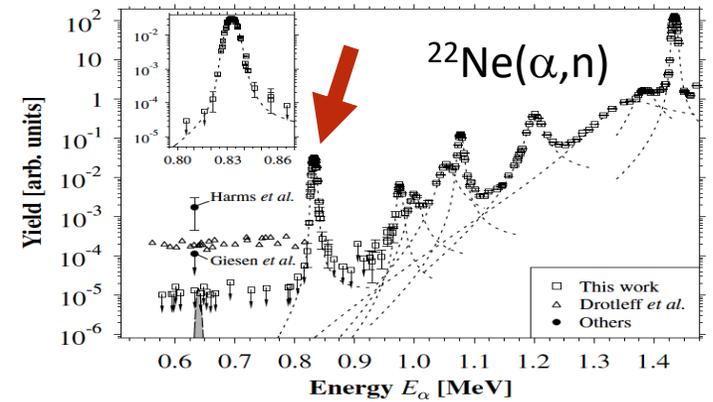
$$\omega_\alpha = g \Gamma_\alpha \Gamma_n / (\Gamma_\alpha + \Gamma_\gamma + \Gamma_n)$$

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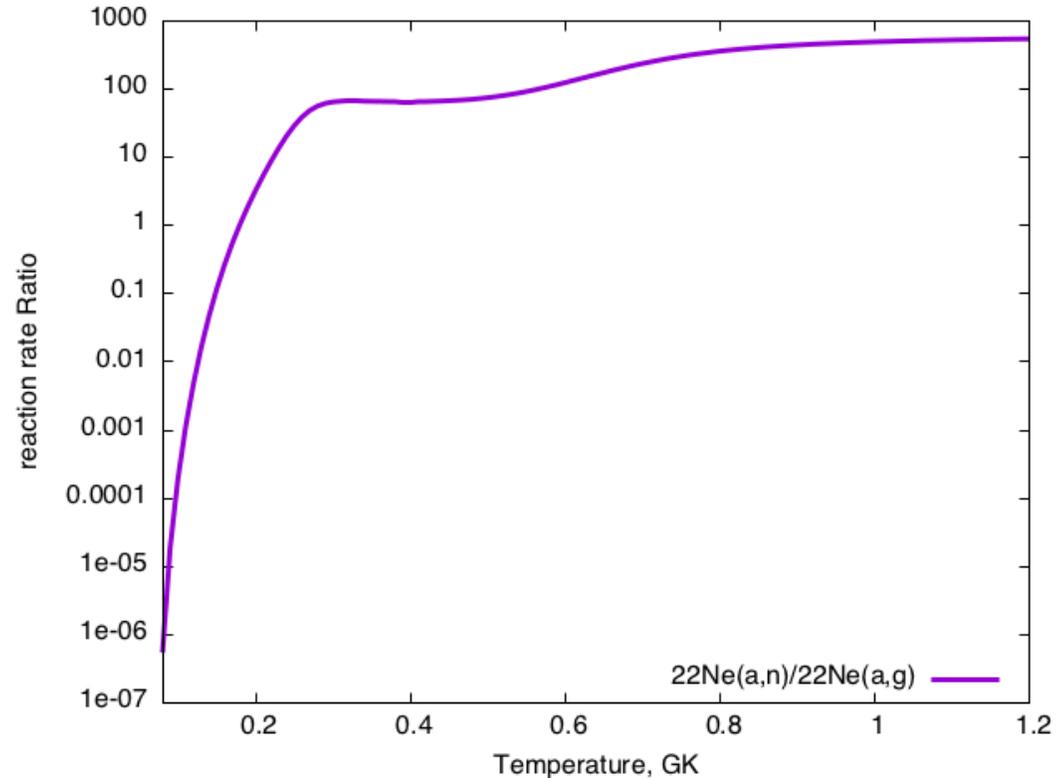
$$\omega_\gamma = g \Gamma_\alpha \Gamma_\gamma / (\Gamma_\alpha + \Gamma_\gamma + \Gamma_n)$$

$$\frac{\omega_\alpha}{\omega_\gamma} = \frac{\Gamma_n}{\Gamma_\gamma}$$

Resonance strength ratio
and Reaction rate ratio
independent of Γ_α



$^{22}\text{Ne}(\alpha,n) / ^{22}\text{Ne}(\alpha,\gamma)$



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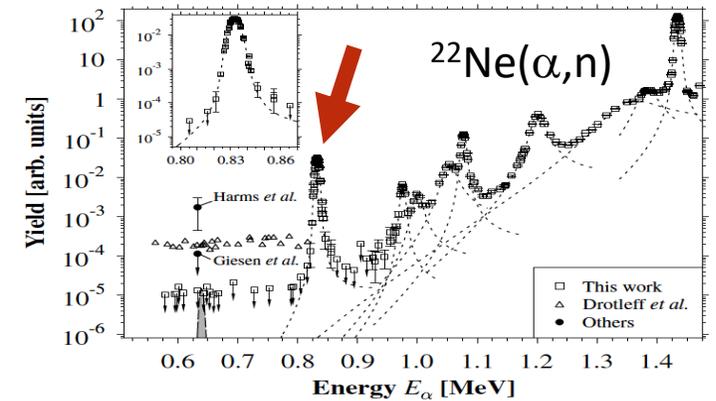
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Resonance strength $^{22}\text{Ne}(\alpha,\gamma)^{26}\text{Mg}$:

$$\omega_{\gamma} = g \Gamma_{\alpha} \Gamma_{\gamma} / (\Gamma_{\alpha} + \Gamma_{\gamma} + \Gamma_n)$$

$$\frac{\omega_{\alpha}}{\omega_{\gamma}} = \frac{\Gamma_n}{\Gamma_{\gamma}}$$

Publication	YEAR	Result	comment
Shahina, PRC	2024	$\Gamma_n / \Gamma_{\gamma} = 2.85(71)$	ω_{α} res. strength
M. Wiescher, EPJA	2023	$\Gamma_n = 0.4 - 1.0$ eV $\Gamma_{\gamma} = 1.33$ eV	Re-evaluation
Y. Chen, PRC	2021	$\Gamma_n = 0.4$ eV $\Gamma_{\gamma} = 1.33$ eV	$^{25}\text{Mg}(d,p)^{26}\text{Mg}$ transfer
S. Ota, PLB	2020	$\Gamma_n / \Gamma_{\gamma} = 1.14(26)$	transfer



$^{22}\text{Ne}(\alpha,n) / ^{22}\text{Ne}(\alpha,\gamma)$



$^{25}\text{Mg}(n,\gamma)$ for neutron source reaction in stars

Resonance strength $^{22}\text{Ne}(\alpha,n)^{25}\text{Mg}$:

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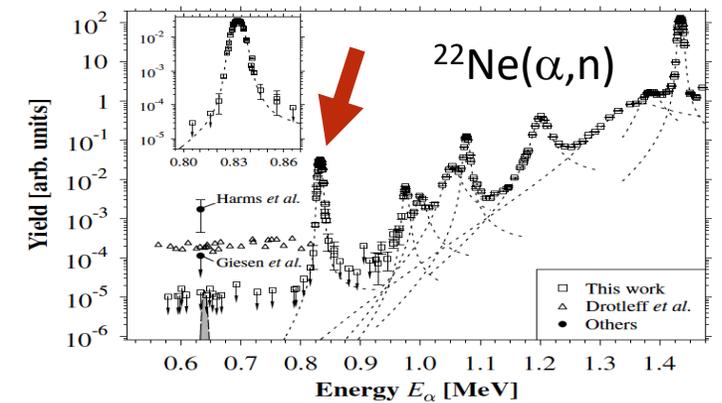
Resonance strength $^{22}\text{Ne}(\alpha,\gamma)^{26}\text{Mg}$:

$$\omega_\gamma = g \Gamma_\alpha \Gamma_\gamma / (\Gamma_\alpha + \Gamma_\gamma + \Gamma_n)$$

$$\frac{\omega_\alpha}{\omega_\gamma} = \frac{\Gamma_n}{\Gamma_\gamma}$$

Neutron width Γ_n and γ -ray width Γ_γ can be deduced from $n + ^{25}\text{Mg}$ experiments

$^{22}\text{Ne}(\alpha,n) / ^{22}\text{Ne}(\alpha,\gamma)$

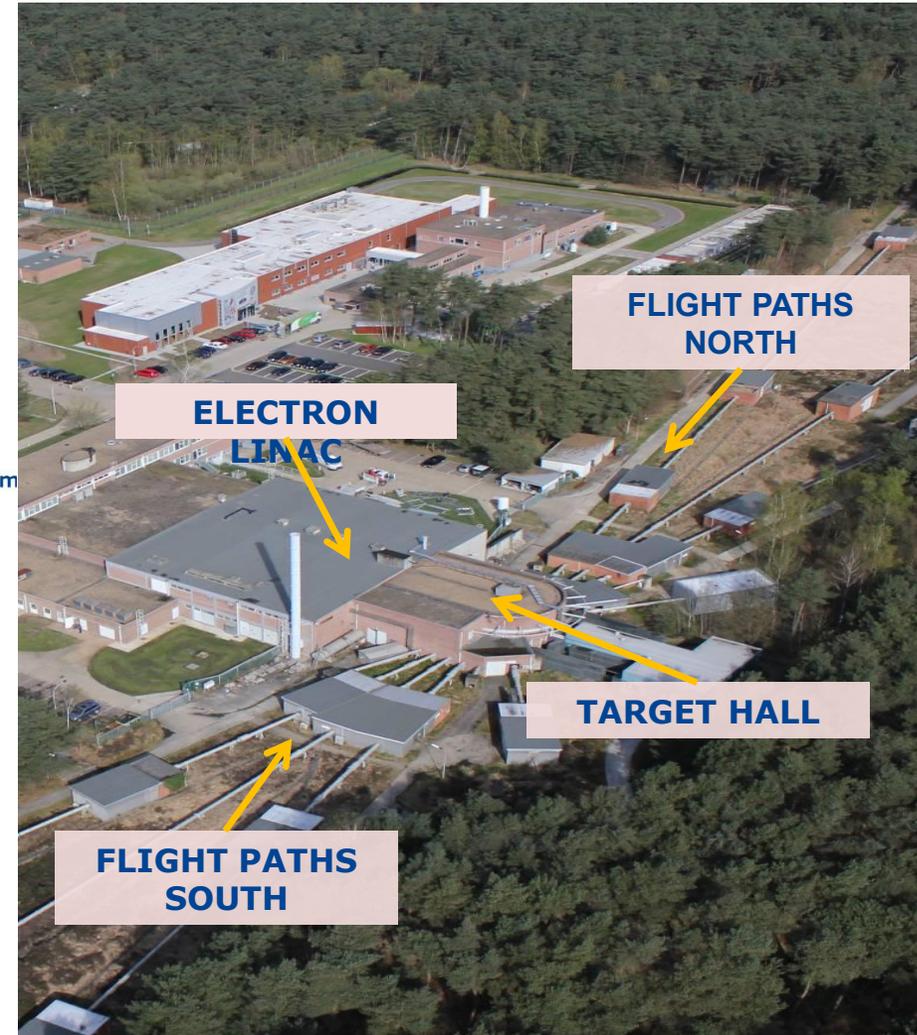
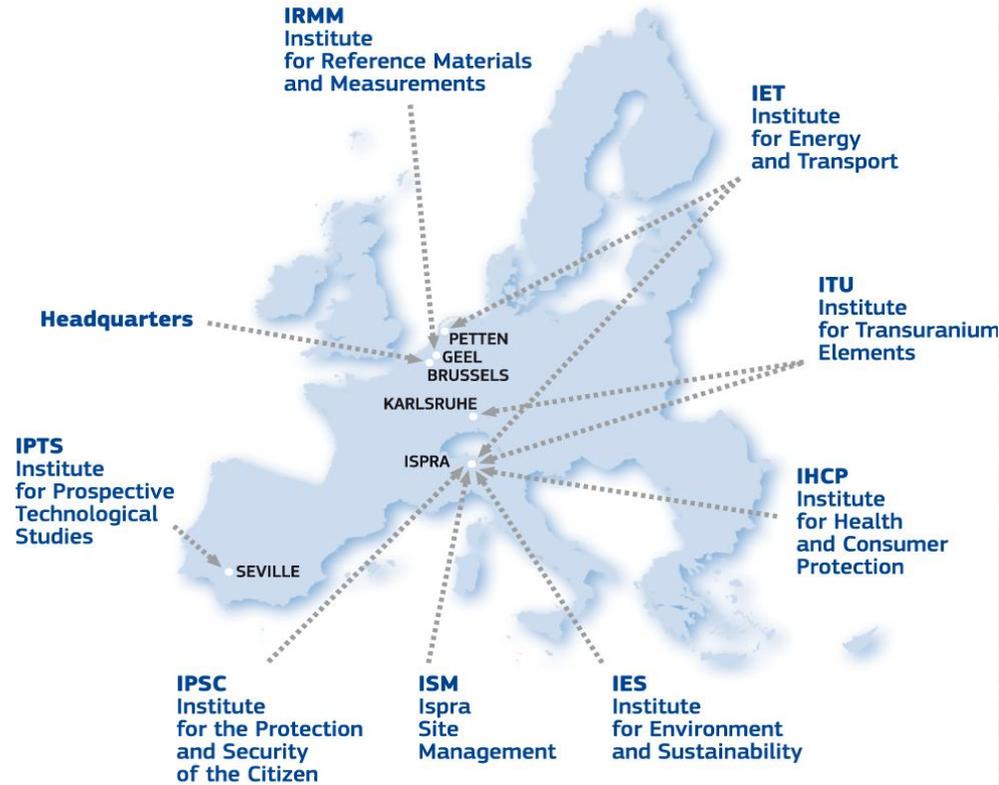


$$\sigma_\gamma(E_n) = g \frac{\pi}{k_n^2} \frac{\Gamma_n \Gamma_\gamma}{(E_n - E_R)^2 + (\Gamma/2)^2}$$

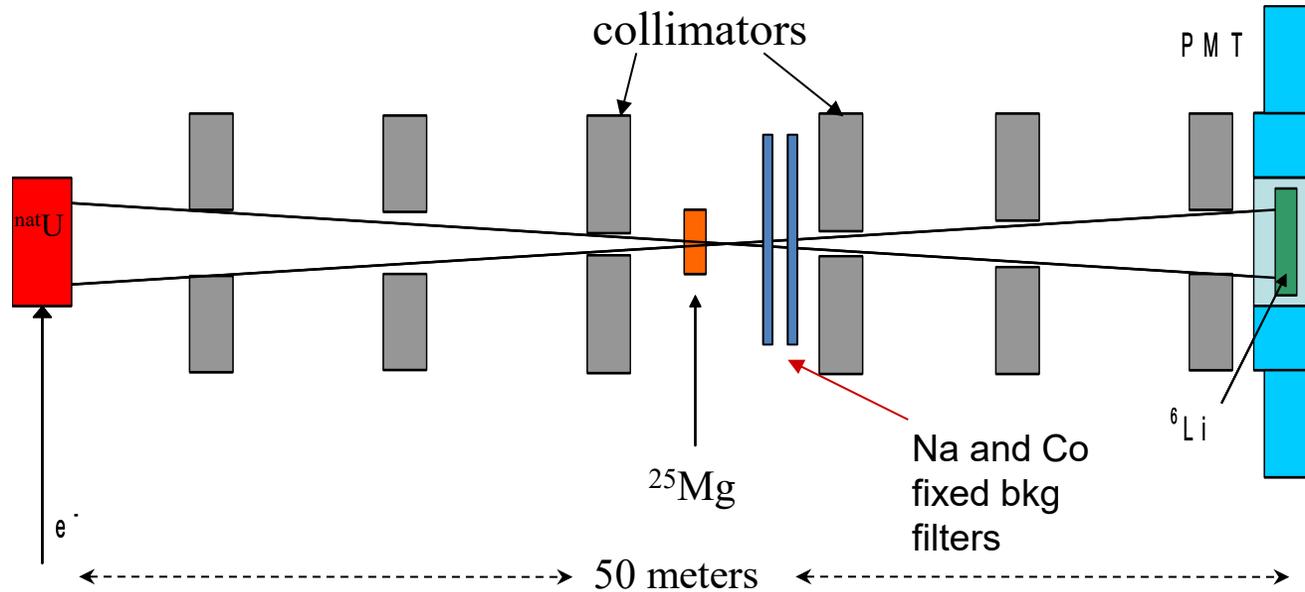
$^{25}\text{Mg}(n,\gamma)$ cross section



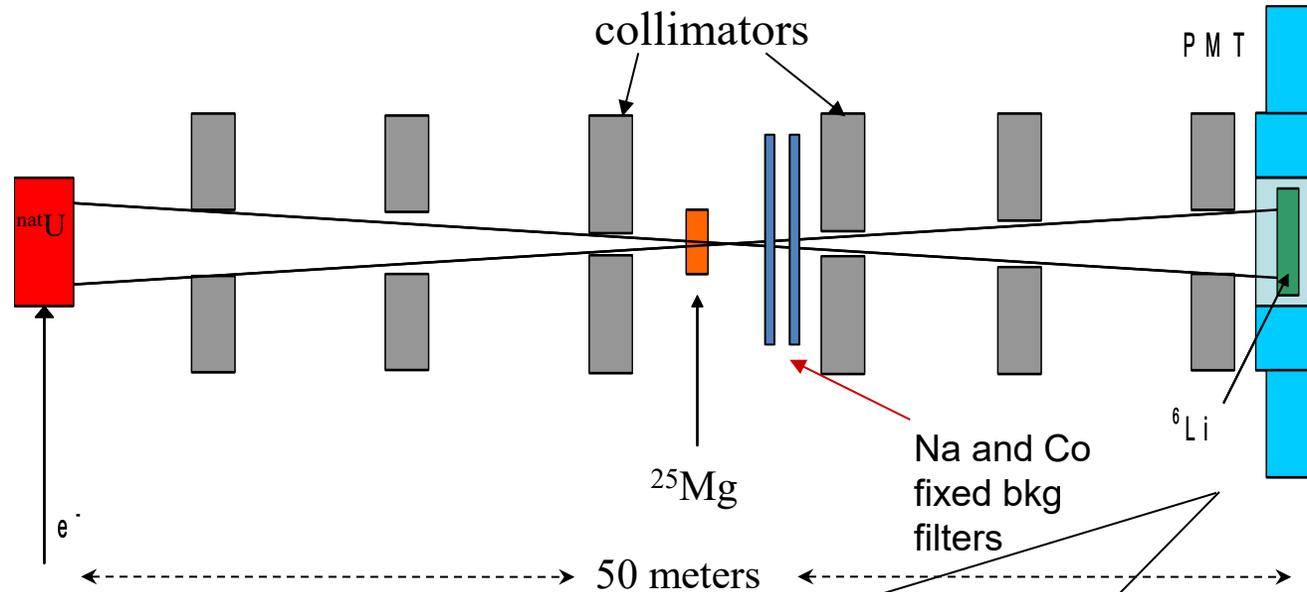
Measurement of $^{25}\text{Mg}(n,\text{tot})$ @ GELINA



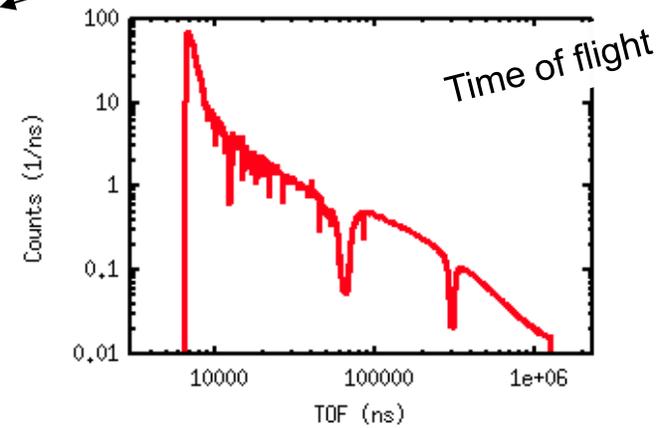
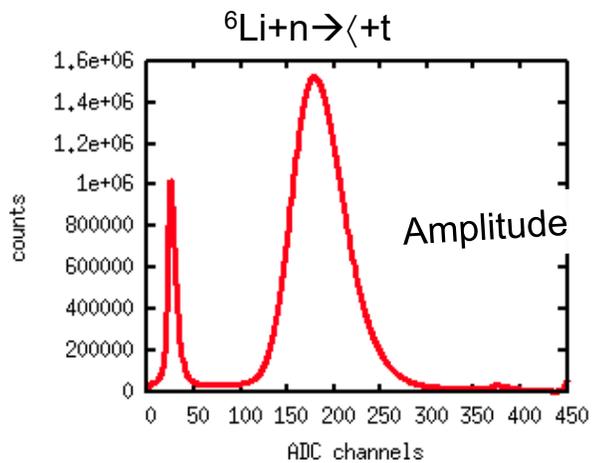
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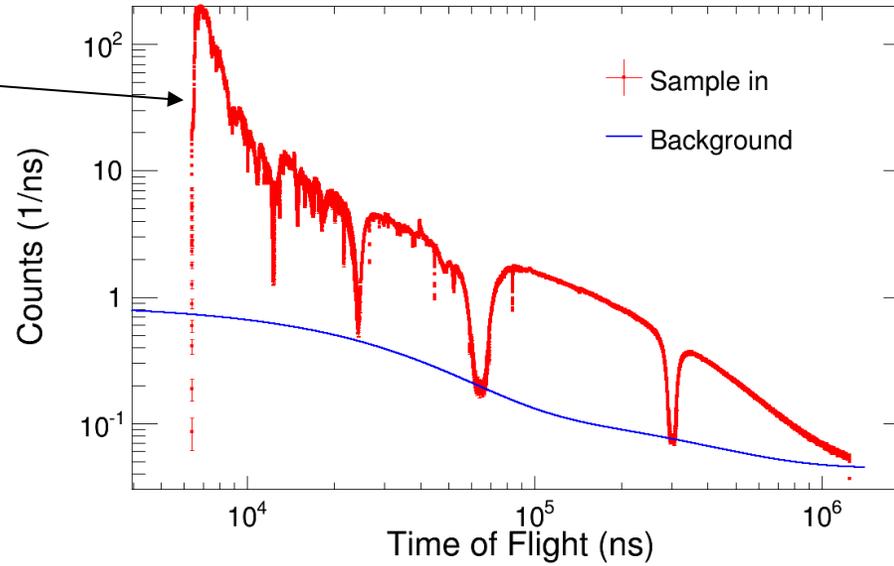
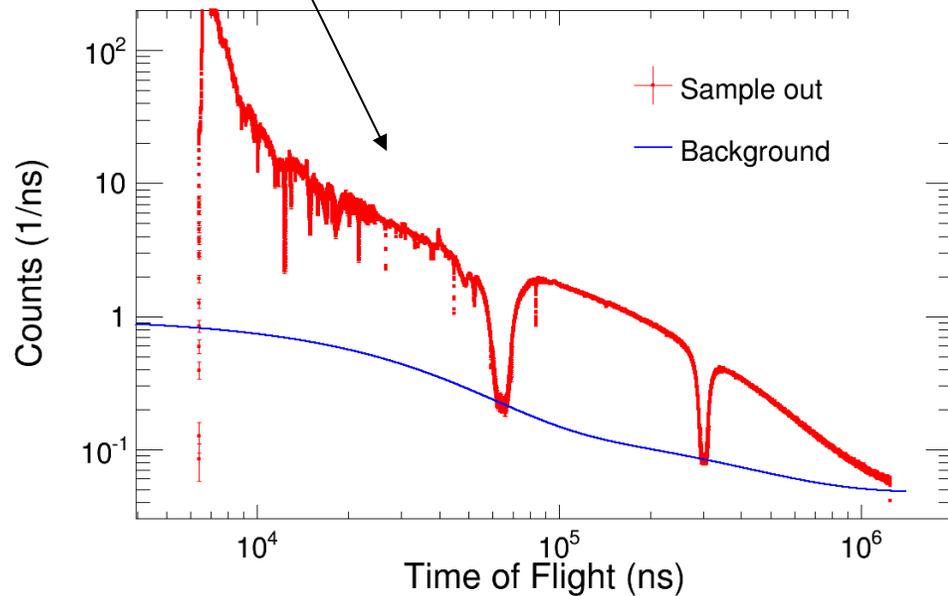


Transmission setup:
 ➤ 2 ^6Li -glass scintillators



Measurement of $^{25}\text{Mg}(n,\text{tot})$ @ GELINA

$$T = \frac{C_{\text{in}}}{C_{\text{out}}} \propto e^{-n\sigma_{\text{tot}}}$$



Background determined by **black resonance** technique:

$$B(t) = b_0 + b_1 e^{-\lambda_1 t} + b_2 e^{-\lambda_2 t} + b_3 e^{-\lambda_3(t+t_0)}$$

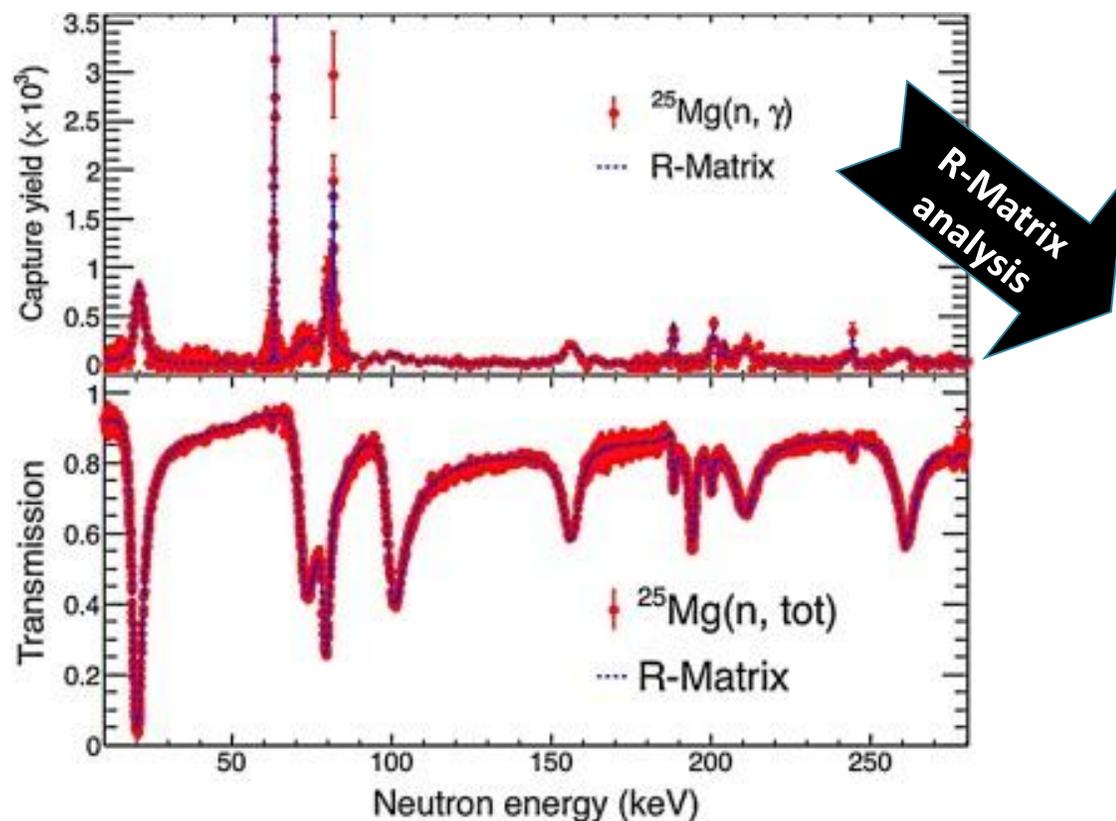
P. Schillebeeckx, *et al.*, Nucl. Data Sheets **113** (2012) 3054



$^{25}\text{Mg}(n,\text{tot})$ and $^{25}\text{Mg}(n,\gamma)$ R-Matrix analysis

C Massimi *et al.*, *Phys. Rev. C* **85**, 044615 (2012)

C Massimi *et al.*, *Phys. Lett. B* **768**, 1 (2017)



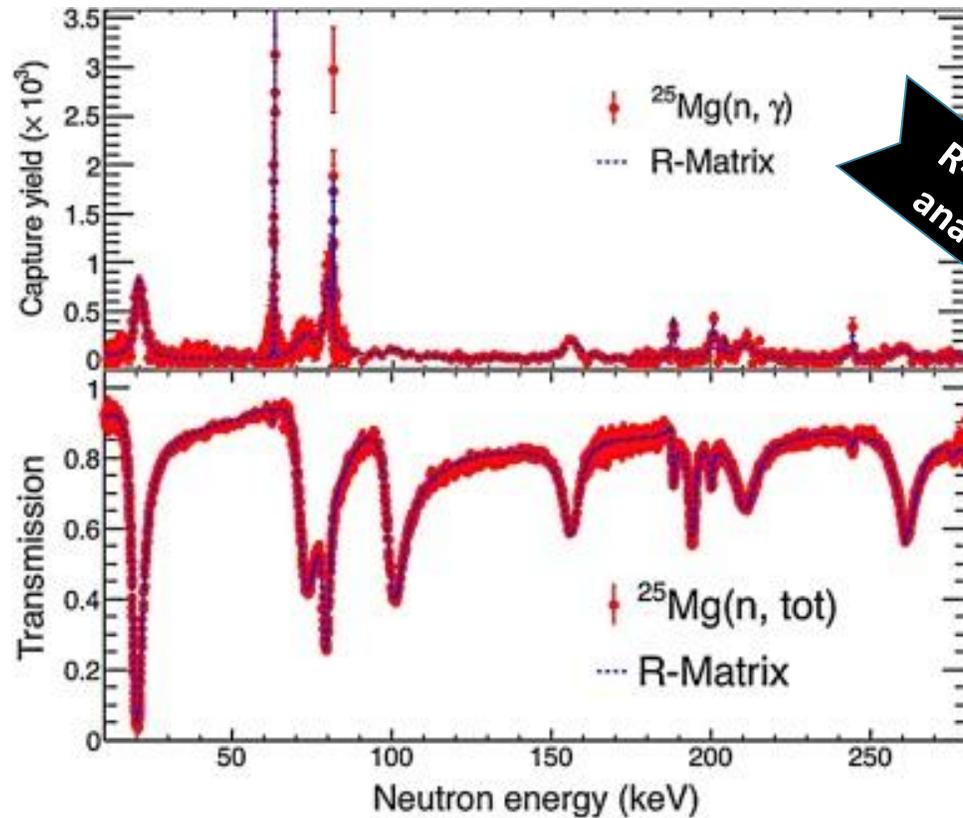
E_n (keV)	E_x (keV)	E_{α}^{Lab} (keV)	J^{π} (\hbar)	Γ_{γ} (eV)	Γ_n (eV)
19.92(1)	11112	589	2^+	1.37(6)	2095(5)
62.73(1)	11154		1^+	4.4(5)	7(2)
72.82(1)	11163	649	2^+	2.8(2)	5310(50)
79.23(1)	11169	656	$3^{- (a)}$	3.3(2)	1940(20)
81.11(1)	11171			5(1)	1 – 30
100.33(2)	11190		3^+	1.3(2)	5230(30)
155.83(2)	11243		2^-	4.7(5)	5950(50)
187.95(2)	11274	779	2^+	2.2(2)	410(10)
194.01(2)	11280	786	$3^{- (a)}$	0.3(1)	1810(20)
199.84(2)	11285		2^-	4.8(4)	1030(30)
203.88(4)	11289			0.9(3)	3 – 20
210.23(3)	11295		2^-	6.6(6)	7370(60)
243.98(2)	11328	843	$2^+ (b)$	2.2(3)	171(6)
260.84(8)	11344			1.0(2)	300 – 3900
261.20(2)	11344		> 3	3.0(3)	6000 – 9000

Energy, J^{π} , Γ_{γ} , Γ_n



$^{25}\text{Mg}(n,\text{tot})$ and $^{25}\text{Mg}(n,\gamma)$ R-Matrix analysis

C Massimi *et al.*, *Phys. Rev. C* **85**, 044615 (2012)
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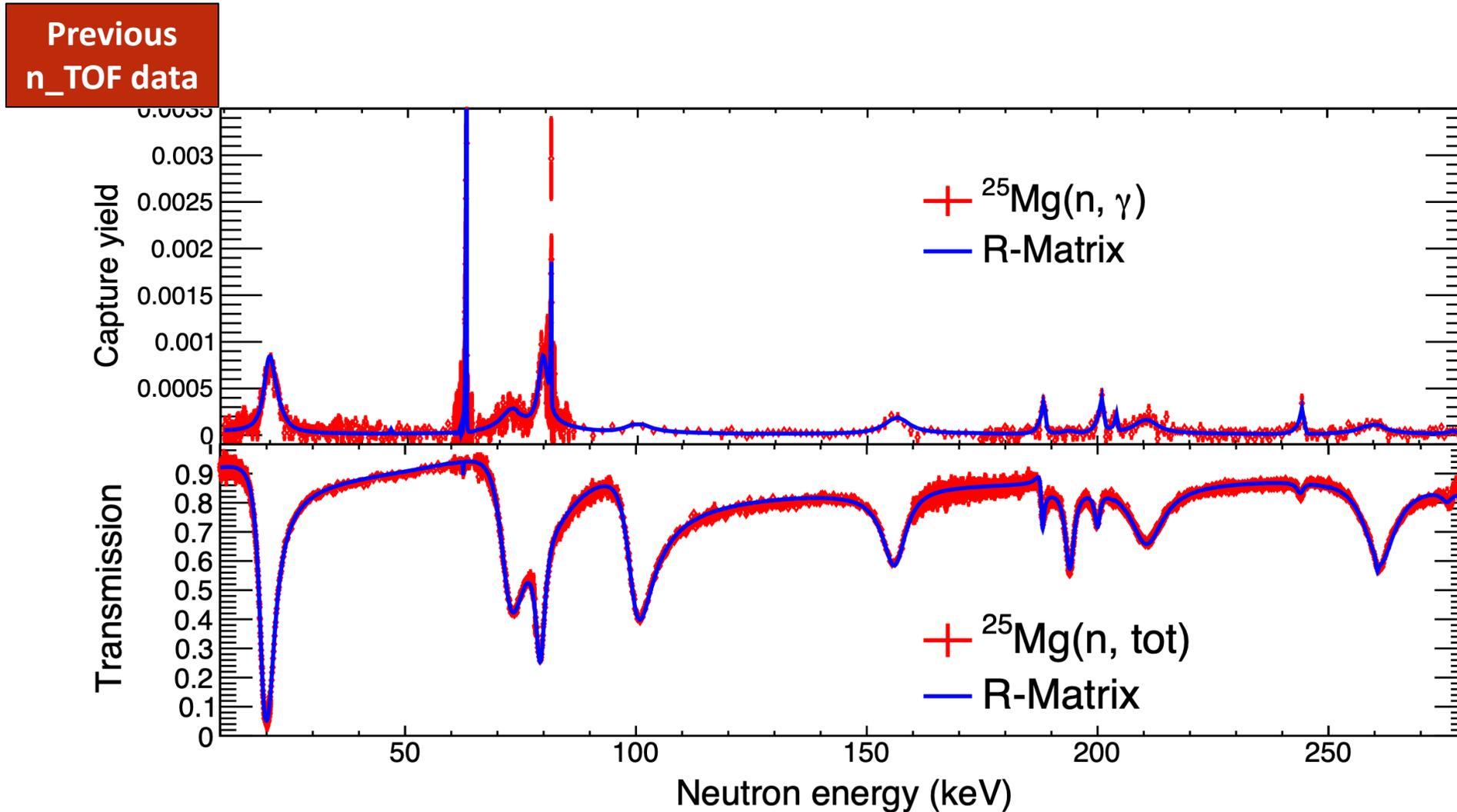


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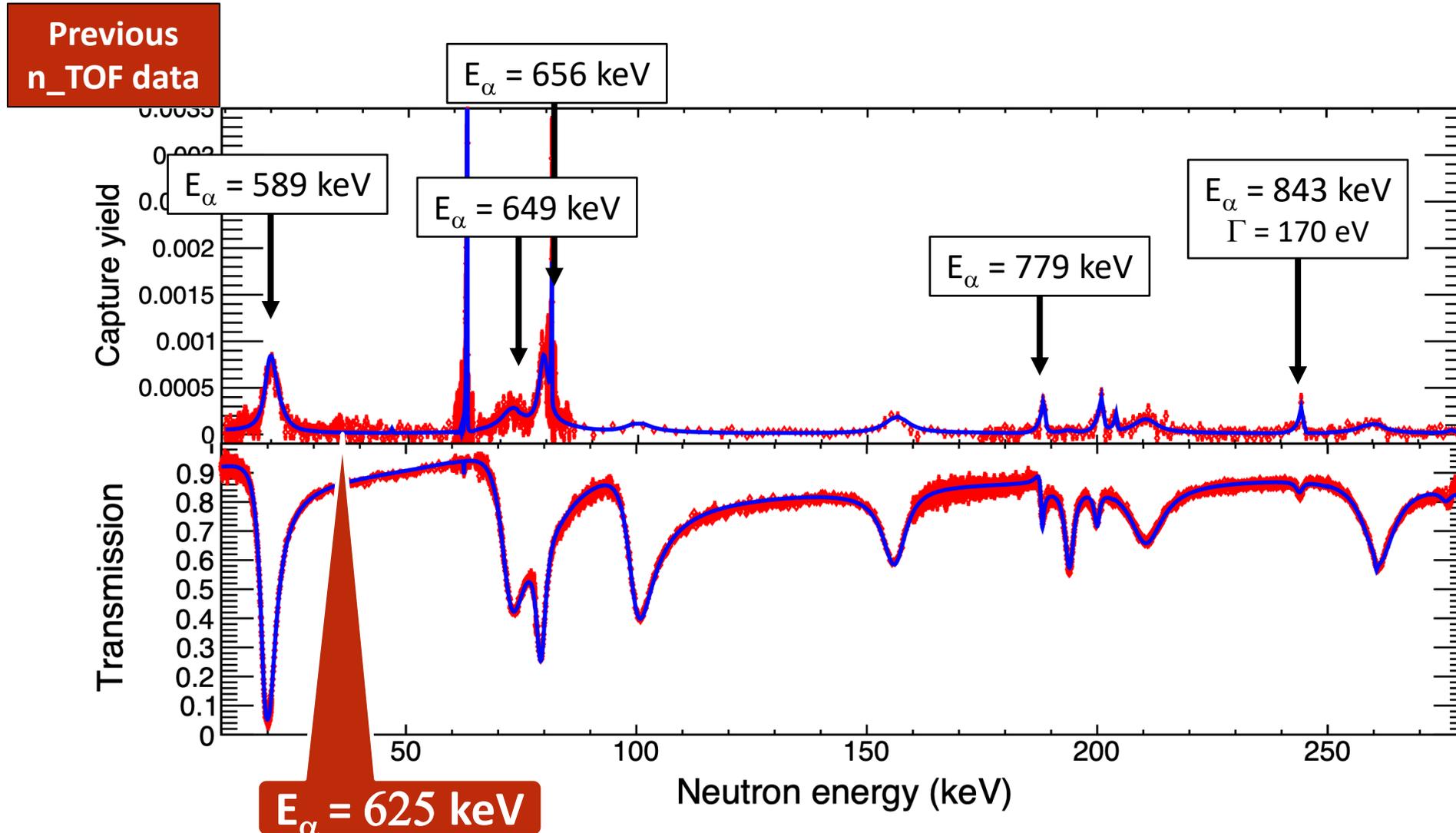
Energy, J^{π} , Γ_{γ} , Γ_n



$^{25}\text{Mg}(n,\gamma)$ for neutron source reaction in stars

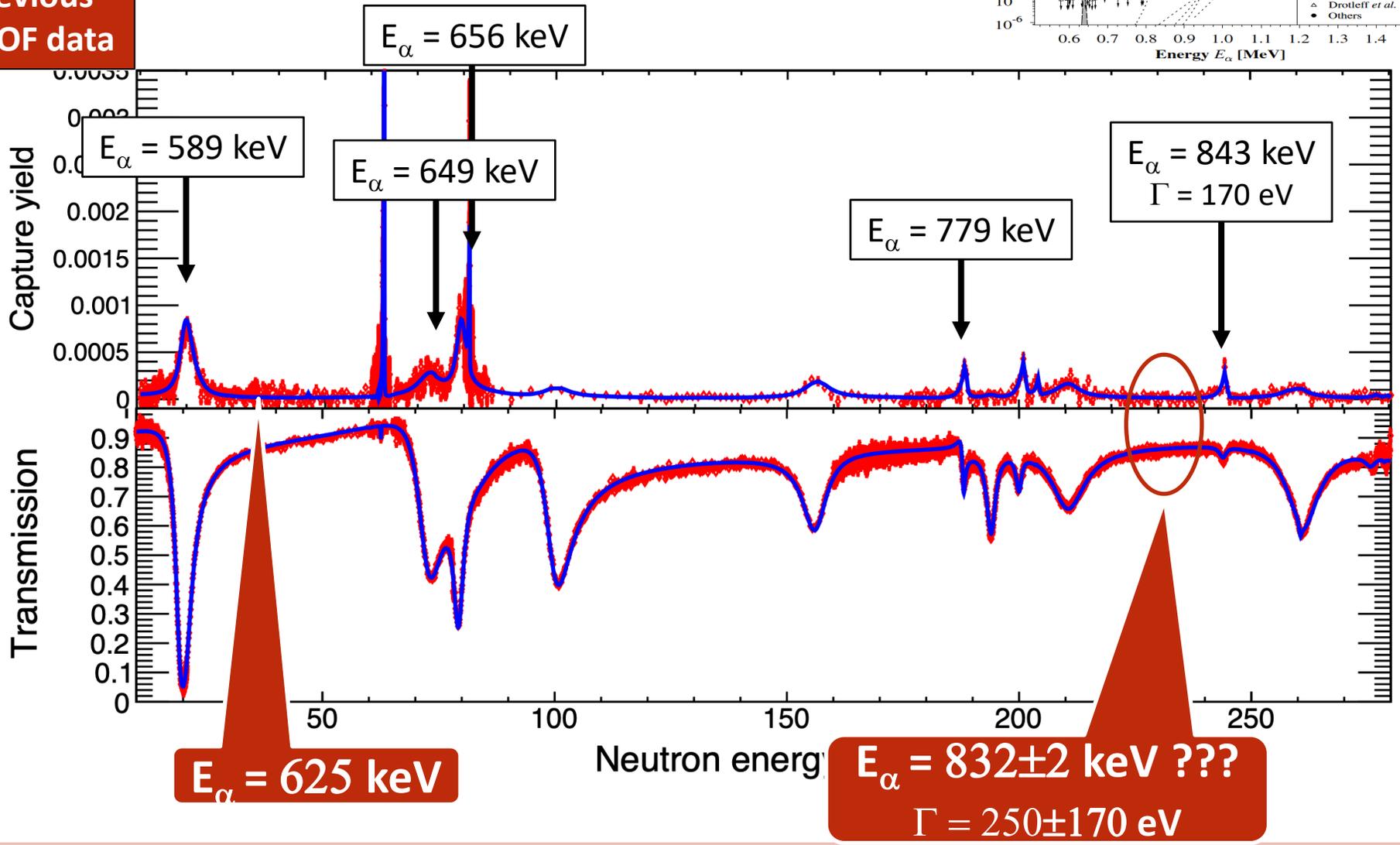


$^{25}\text{Mg}(n,\gamma)$ for neutron source reaction in stars



$^{25}\text{Mg}(n,\gamma)$ for neutron source reaction in stars

Previous
n_TOF data



Conclusions...?

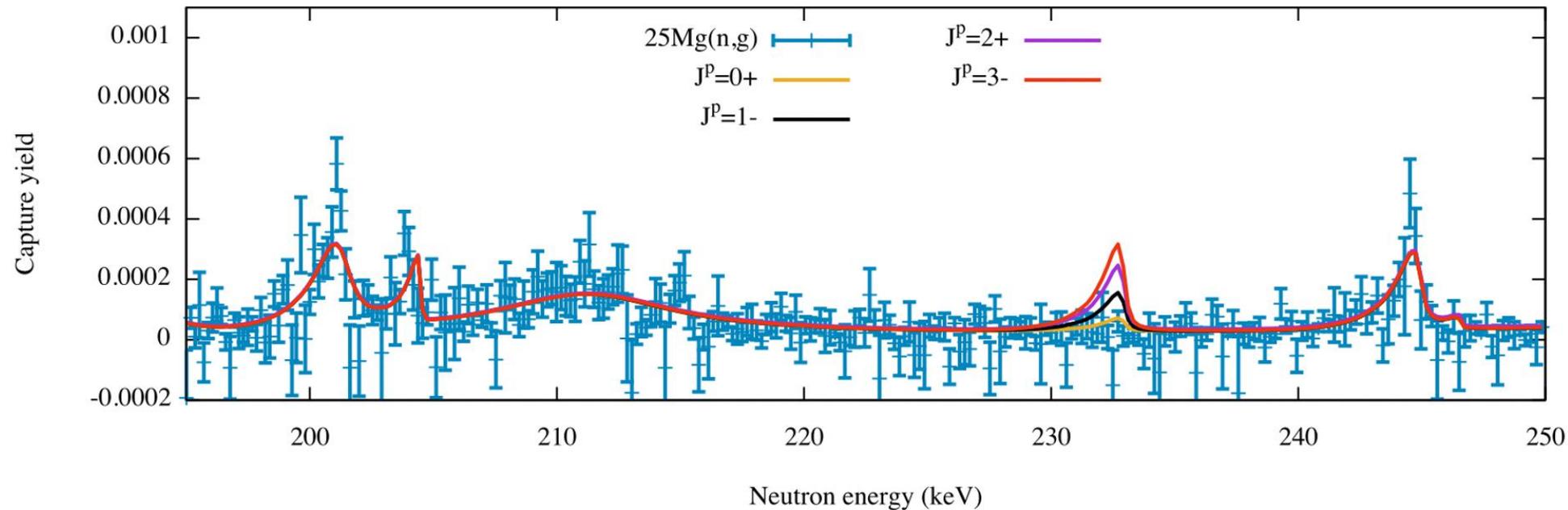
- $^{22}\text{Ne}(\alpha, n)$ and (α, γ) represent a long-standing “problem” in nuclear astrophysics
- Measurements of $^{25}\text{Mg}(n, \text{tot})$ and $^{25}\text{Mg}(n, \gamma)$ were performed at the GELINA facility and the n_TOF facility, respectively, to study excited states in ^{26}Mg
- Simultaneous resonance shape (R-Matrix) analysis of capture and transmission resulted in:
 - accurate $^{25}\text{Mg}(n, \gamma)$ cross section;
 - **energy** and **J^π** determination of ^{26}Mg levels: evidence for natural states;
 - constraints for the competing $^{22}\text{Ne}(\alpha, \gamma)$ reaction;
 - doubts on the **$E_\alpha = 832 \text{ keV}$** resonance.



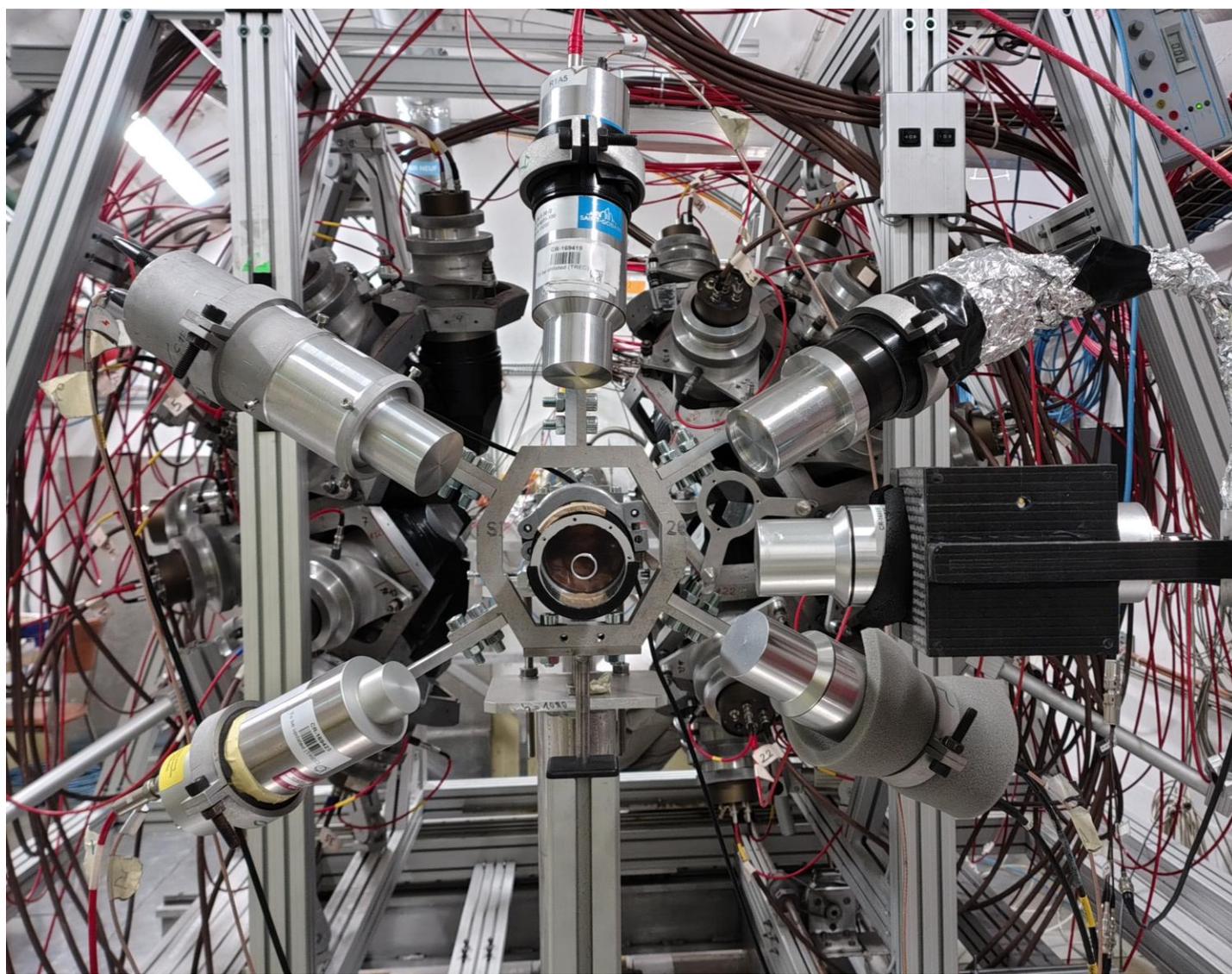
Proposal: $^{25}\text{Mg}(n,\gamma)^{26}\text{Mg}$ @ n_TOF

Our proposal is to **repeat the measurement in EAR1** with a factor 4 higher statistics and with some improvements:

- Combined use of LaBr_3 and C_6D_6 detectors
- Use of a thicker enriched ^{25}Mg sample
- Combine with a capture measurement in EAR2



Proposal: $^{25}\text{Mg}(n,\gamma)^{26}\text{Mg}$ @ n_TOF



- More protons (4×10^{18})
- Thicker Mg sample



Acknowledgments

This project has received funding from the European Union's Horizon Europe Research and Innovation programme under Grant Agreement No 101057511 (EURO-LABS).



Thank you for your attention!



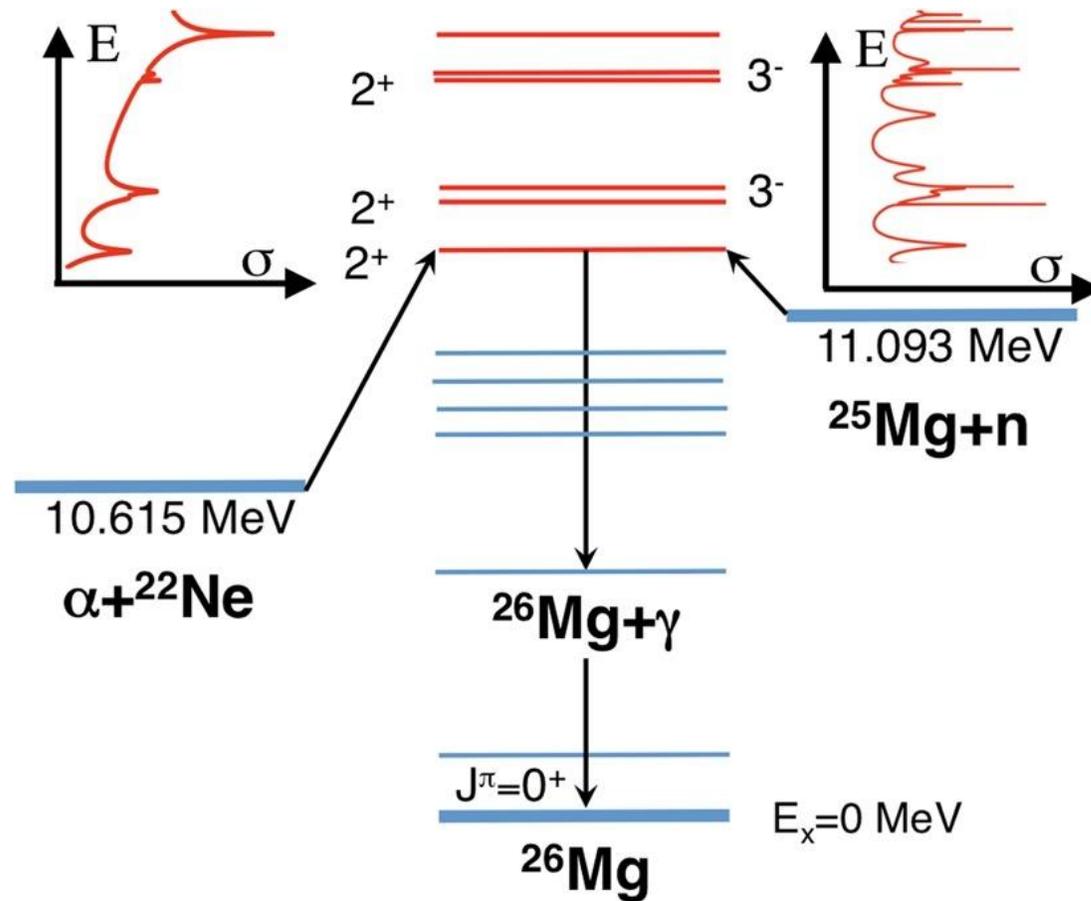
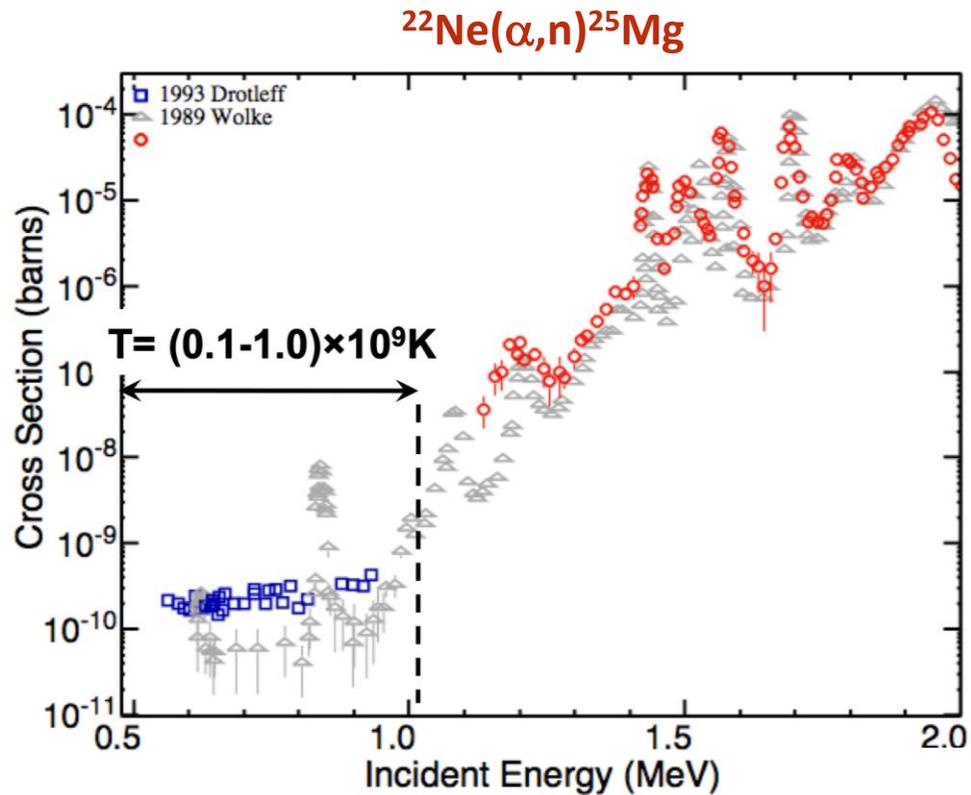
ALMA MATER STUDIORUM
UNIVERSITÀ DI BOLOGNA

Credits:

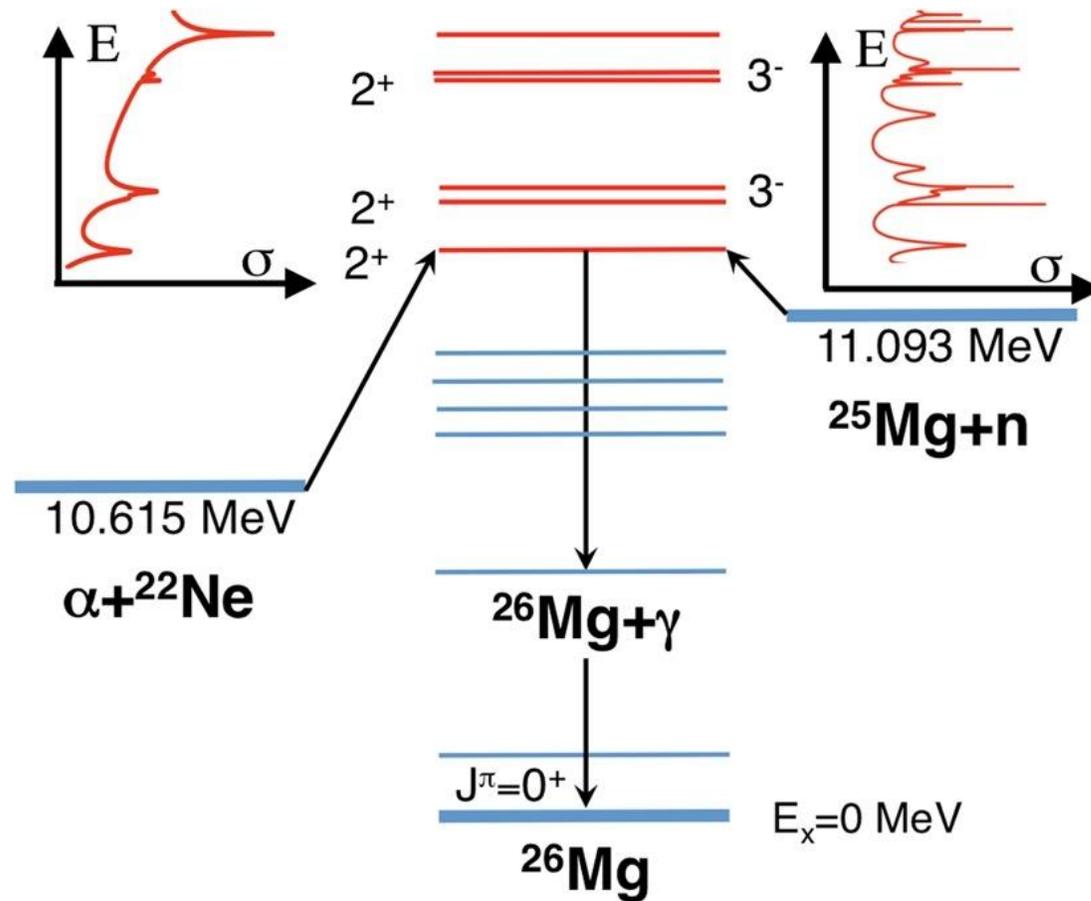
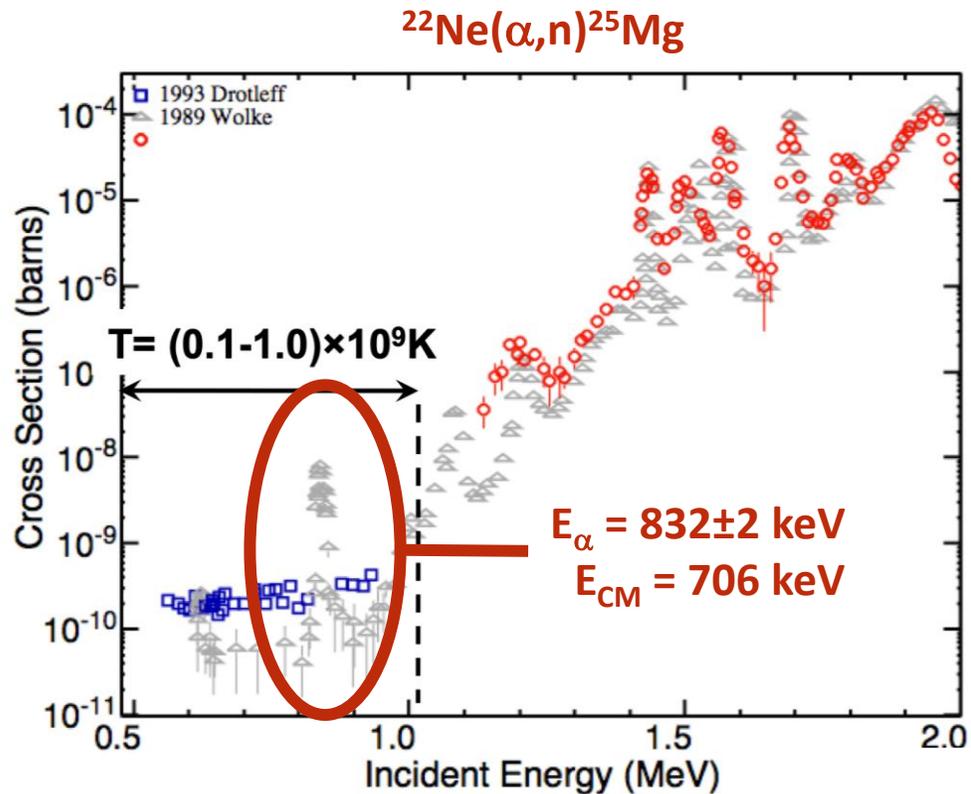
Motivations

- NEUTRON POISON:
 - $^{25,26}\text{Mg}$ are the most important neutron poisons due to neutron capture on Mg stable isotopes, i.e. $^{25,26}\text{Mg}(n,\gamma)$, in competition with neutron capture on ^{56}Fe (the basic s-process seed for the production of heavier isotopes).
- CONSTRAINTS for $^{22}\text{Ne}(\alpha,n)^{25}\text{Mg}$ and $^{22}\text{Ne}(\alpha,\gamma)^{26}\text{Mg}$:
 - $^{22}\text{Ne}(\alpha,n)^{25}\text{Mg}$ is one of the most important neutron source in Red Giant stars. Its reaction rate is very uncertain because of the poorly known property of the states in ^{26}Mg . From neutron measurements the energy, J^π and **energy** of ^{26}Mg states can be deduced, in addition to Γ_γ and Γ_n .

Measurement of $^{25}\text{Mg}(n,\gamma) \leftrightarrow ^{22}\text{Ne}(\alpha,n)^{25}\text{Mg}$



Measurement of $^{25}\text{Mg}(n,\gamma) \leftrightarrow ^{22}\text{Ne}(\alpha,n)^{25}\text{Mg}$



- ^{26}Mg levels via $n + ^{25}\text{Mg}$
- $^{25}\text{Mg}(n,\gamma)$ is not conclusive enough, **need for other reaction channels**



Extra slides

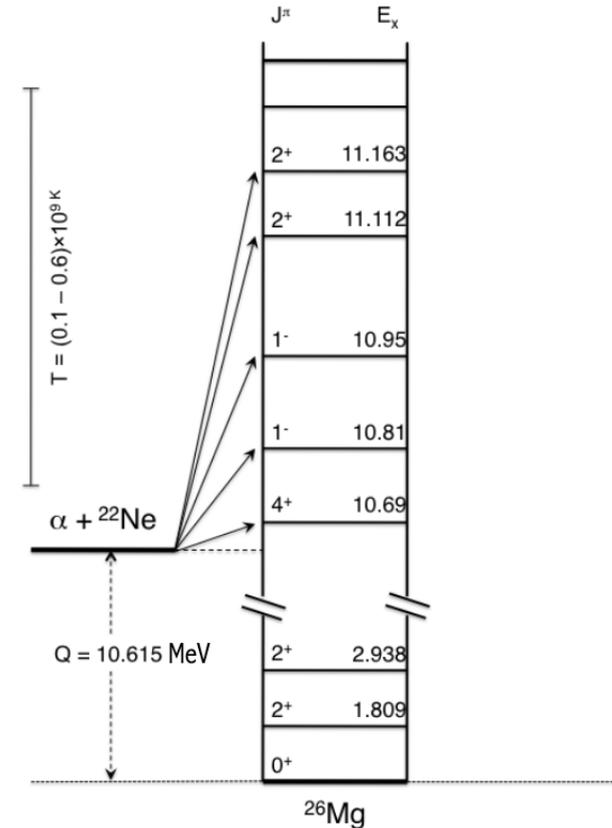
Constraints for the $^{22}\text{Ne}(\alpha, n)^{25}\text{Mg}$ reaction

Element	Spin/ parity
^{22}Ne	0^+
^4He	0^+

$$\vec{J} = \underbrace{\vec{I} + \vec{i}} + \vec{l}$$

$$\vec{J} = \mathbf{0} + \vec{l}$$

Only **natural-parity** (0^+ , 1^- , 2^+ , 3^- , 4^+ , ...) states in ^{26}Mg can participate in the $^{22}\text{Ne}(\alpha, n)^{25}\text{Mg}$ reaction



Extra slides

Constraints for the $^{22}\text{Ne}(\alpha, n)^{25}\text{Mg}$ reaction

Element	Spin/parity
^{25}Mg	$5/2^+$
n	$1/2^+$

$$\vec{J} = \vec{I} + \vec{i} + \vec{l}$$

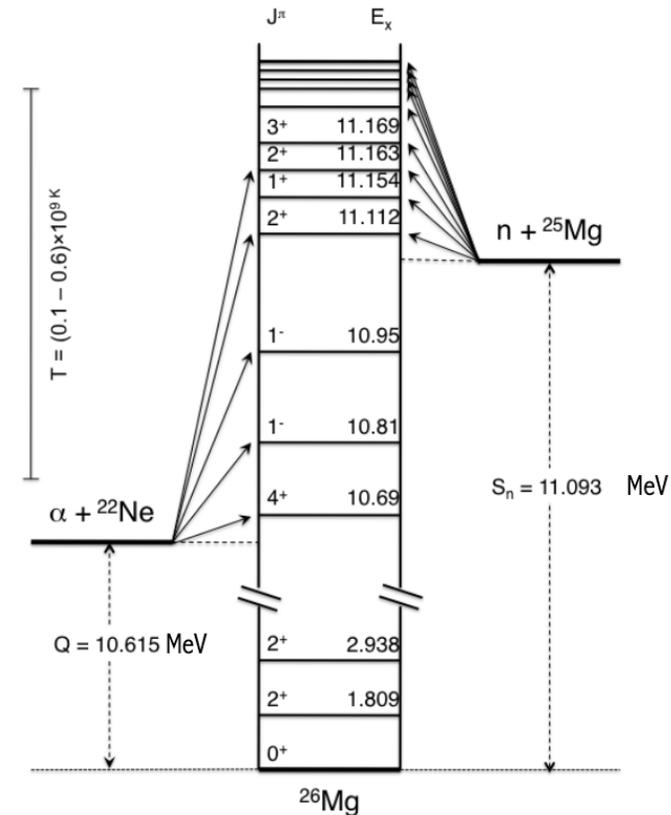
$$\vec{J} = 2 + \vec{l} \quad \vec{J} = 3 + \vec{l}$$

s-wave $\rightarrow J^\pi = \underline{2}^+, 3^+$

p-wave $\rightarrow J^\pi = \underline{1}^-, 2^-, \underline{3}^-, 4^-$

d-wave $\rightarrow J^\pi = \underline{0}^+, 1^+, \underline{2}^+, 3^+, \underline{4}^+, 5^+$

States in ^{26}Mg populated by $^{25}\text{Mg}(n, \gamma)$ reaction



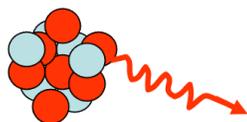
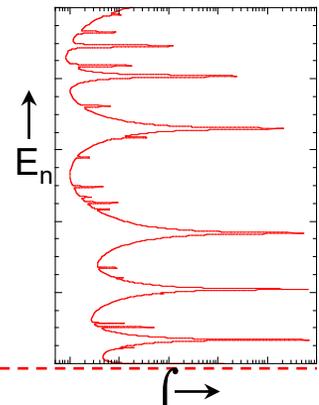
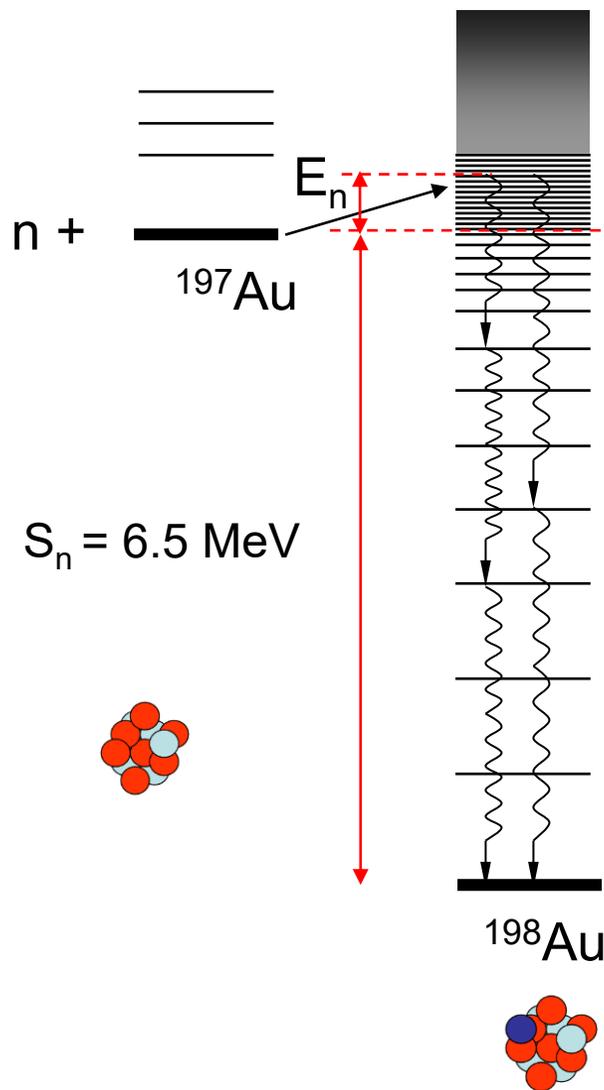
The n_TOF facility

The advantages of n_TOF are a direct consequence of the characteristics of the **PS proton beam**: **high energy, high peak current, low duty cycle.**

proton beam momentum	20 GeV/c
intensity (dedicated mode)	$\sim 10^{13}$ protons/pulse
repetition frequency	1 pulse/1.2s
pulse width	6 ns (rms)
n/p	300
lead target dimensions	80x80x60 cm ³
cooling & moderation material	N ₂ & (H ₂ O + ¹⁰ B)
moderator thickness in the exit face	5 cm
neutron beam dimension in EAR-1 (capture mode)	2 cm (FWHM)



Extra slides: (n,γ)

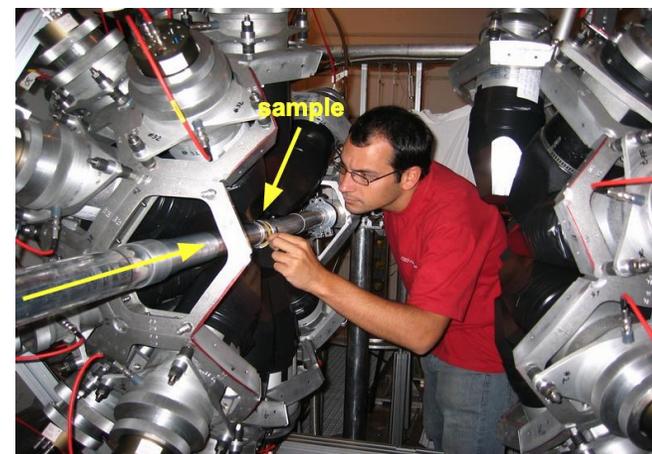
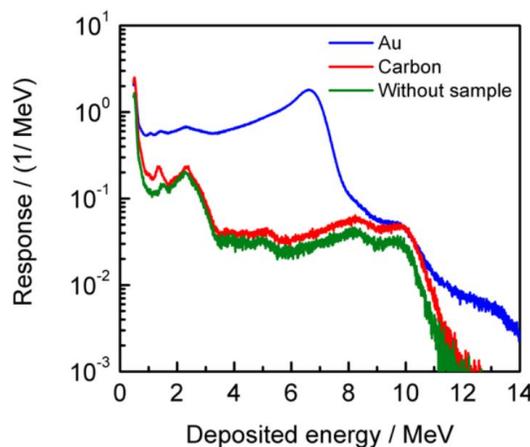


The C_6D_6 Total Energy Detectors (TED)

C_6D_6 scintillators at 135°

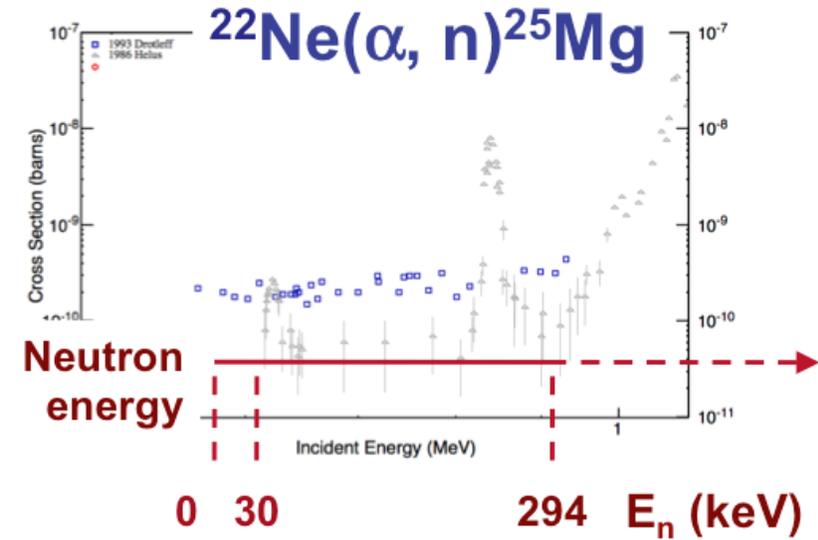
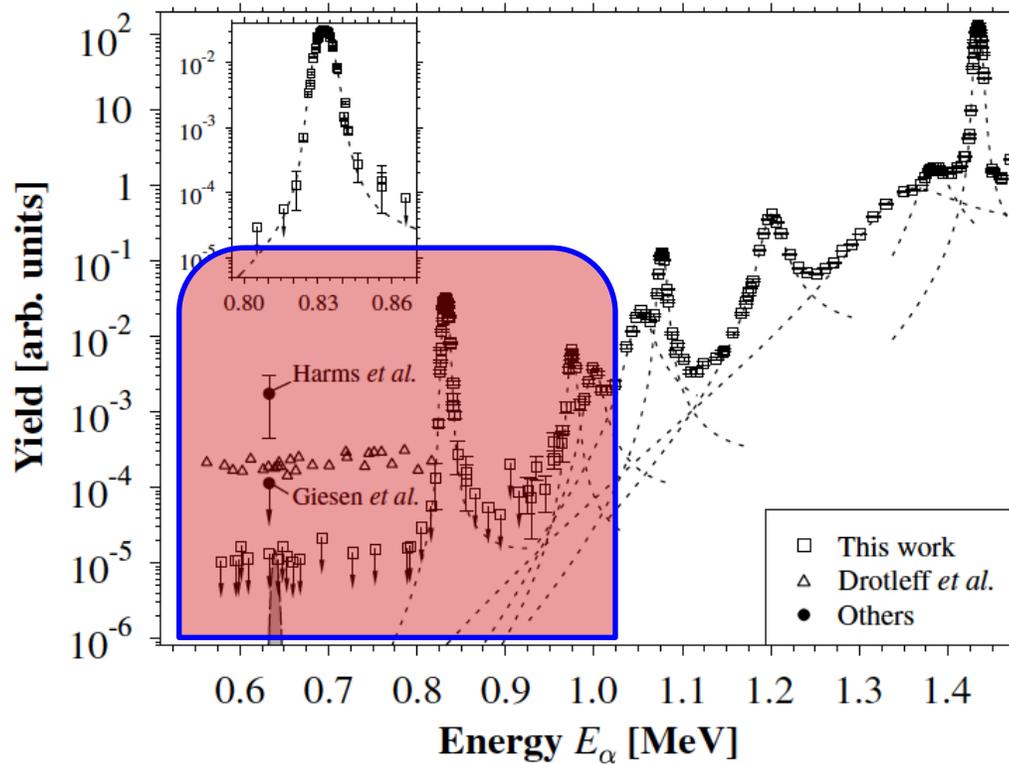


The BaF_2 Total γ -ray Absorption Detector

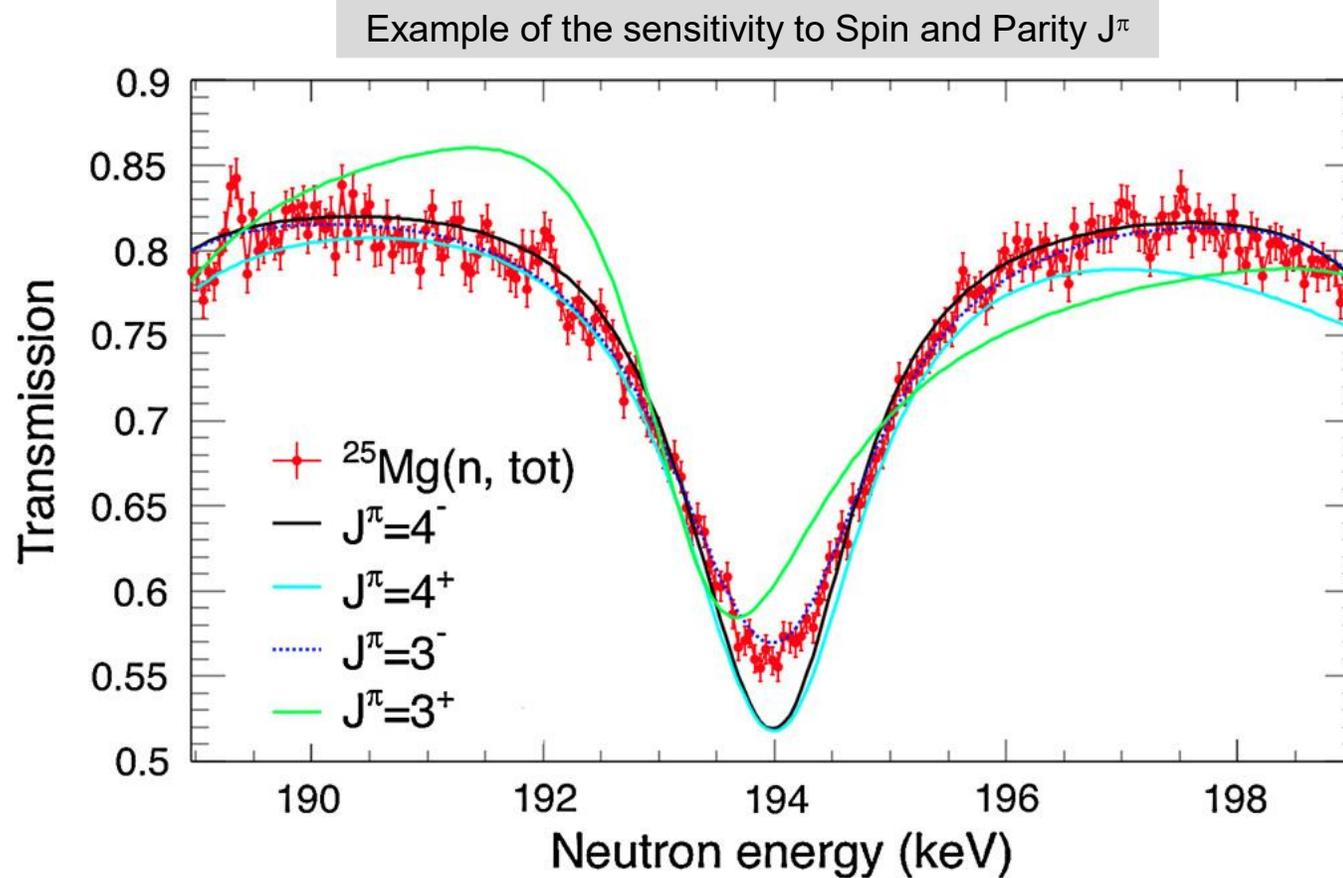


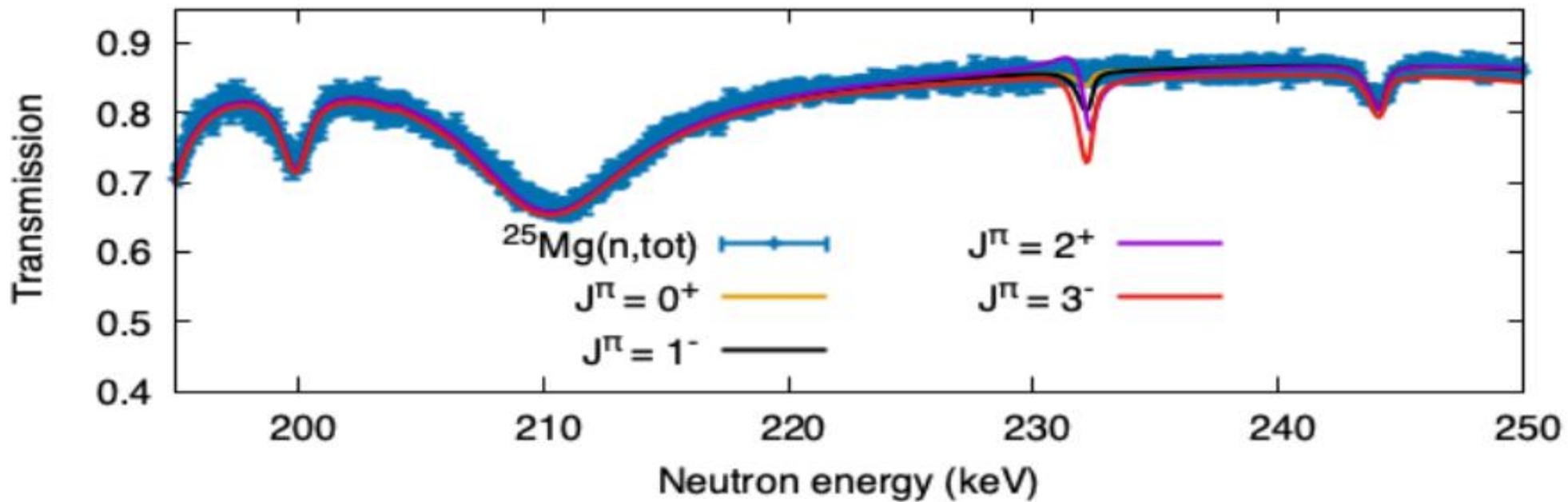
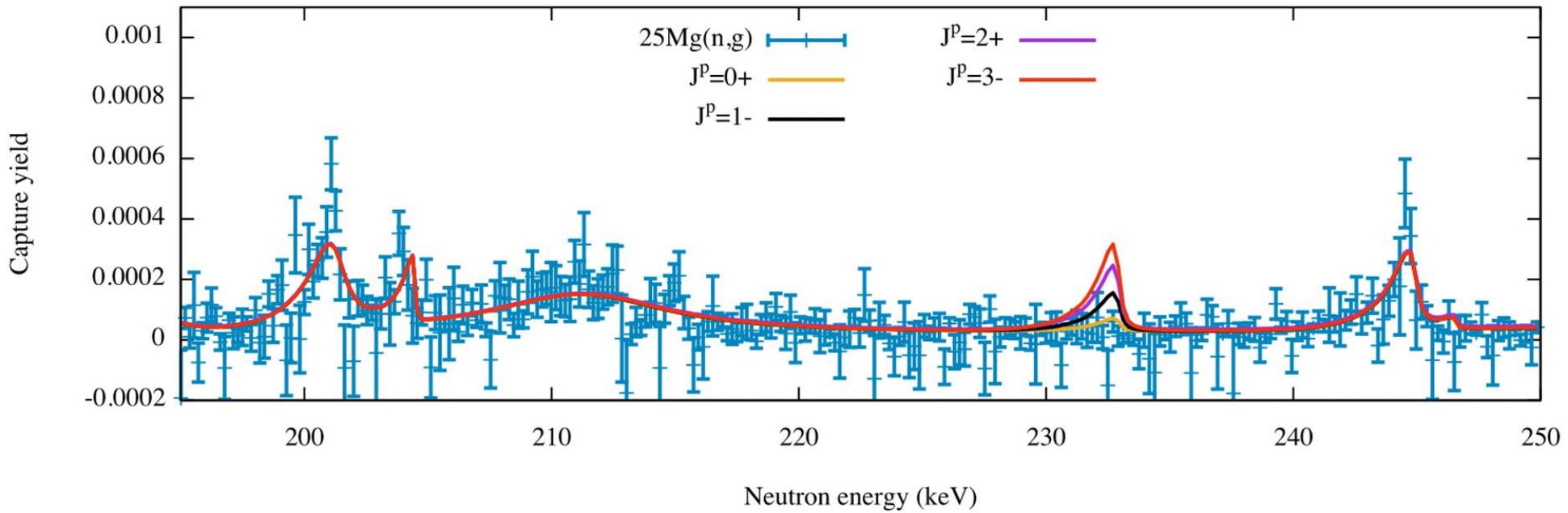
Extra slides

M. Jaeger, *et al.*, Phys Rev. Lett. **87** (2001) 20

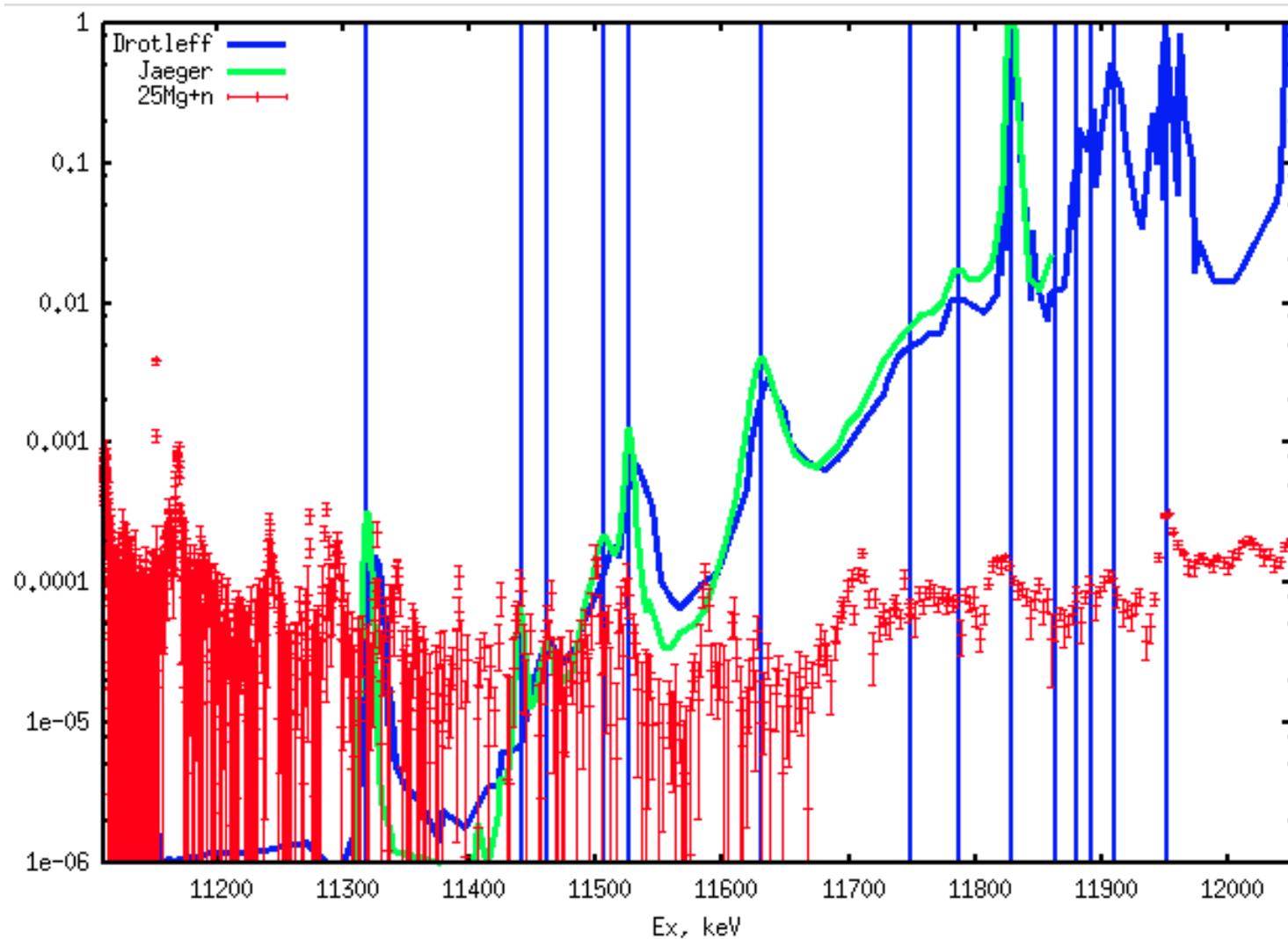


Measurement of $^{25}\text{Mg}(n,\text{tot})$ @ GELINA

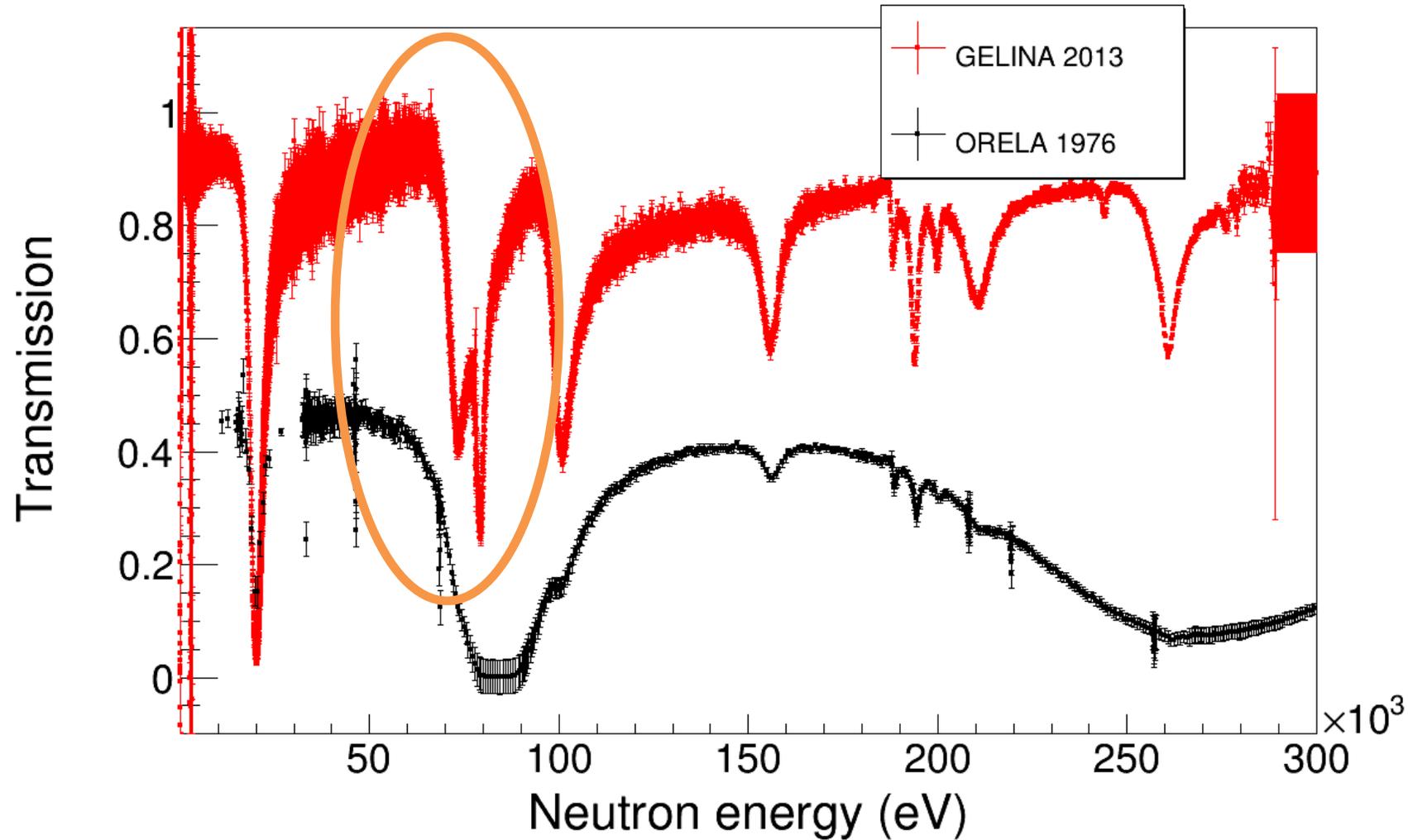




Extra slides



Extra slides



Stellar cross sections (MACS) for the s-process

Some cross sections measured in 2001 - 2024

- ❖ Branching point isotopes:

^{151}Sm , ^{63}Ni , ^{147}Pm , ^{171}Tm , ^{203}Tl , ^{79}Se

- ❖ Abundances in presolar grains:

$^{28,29,30}\text{Si}$, $^{91,92,93,94,96}\text{Zr}$, $^{94,96}\text{Mo}$, ^{146}Nd

- ❖ Magic Nuclei and end-point:

^{139}La , ^{140}Ce , ^{90}Zr , ^{89}Y , ^{88}Sr , $^{204,206,207,208}\text{Pb}$, ^{209}Bi

- ❖ Seeds isotopes:

$^{54,56,57}\text{Fe}$, $^{58,60,62,64}\text{Ni}$, $^{59}\text{Ni}(n,a)$

- ❖ Isotopes of special interest:

$^{186,187,188}\text{Os}$ (cosmochronometer), ^{197}Au (reference cross section), $^{24,25,26}\text{Mg}$, $^{33}\text{S}(n,a)$, $^{14}\text{N}(n,p)$, $^{35}\text{Cl}(n,p)$, $^{26}\text{Al}(n,p)$,
 $^{26}\text{Al}(n,a)$ (neutron poison), ^{154}Gd (s-only isotopes), $^{40}\text{K}(n,p)$, $^{40}\text{K}(n,a)$, $^{63,65}\text{Cu}$, $^{93,94}\text{Nb}$, ^{68}Zn , $^{69,71}\text{Ga}$, $^{70,72,73,74,76}\text{Ge}$,
 $^{77,78,80}\text{Se}$ (weak component), $^{155,157,160}\text{Gd}$, $^7\text{Li}(n,p)$, $^7\text{Li}(n,a)$ BBN

- ❖ Neutron Sources $^{22}\text{Ne}(a,n)^{25}\text{Mg}$ and $^{13}\text{C}(a,n)^{16}\text{O}$:

$n+^{25}\text{Mg}$, $n+^{16}\text{O}$

