



Event 351483885

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Beyond the Flavour Anomalies Workshop  
Roma, 9-11 April 2025

## CP violation in charm: experiment

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*CERN*

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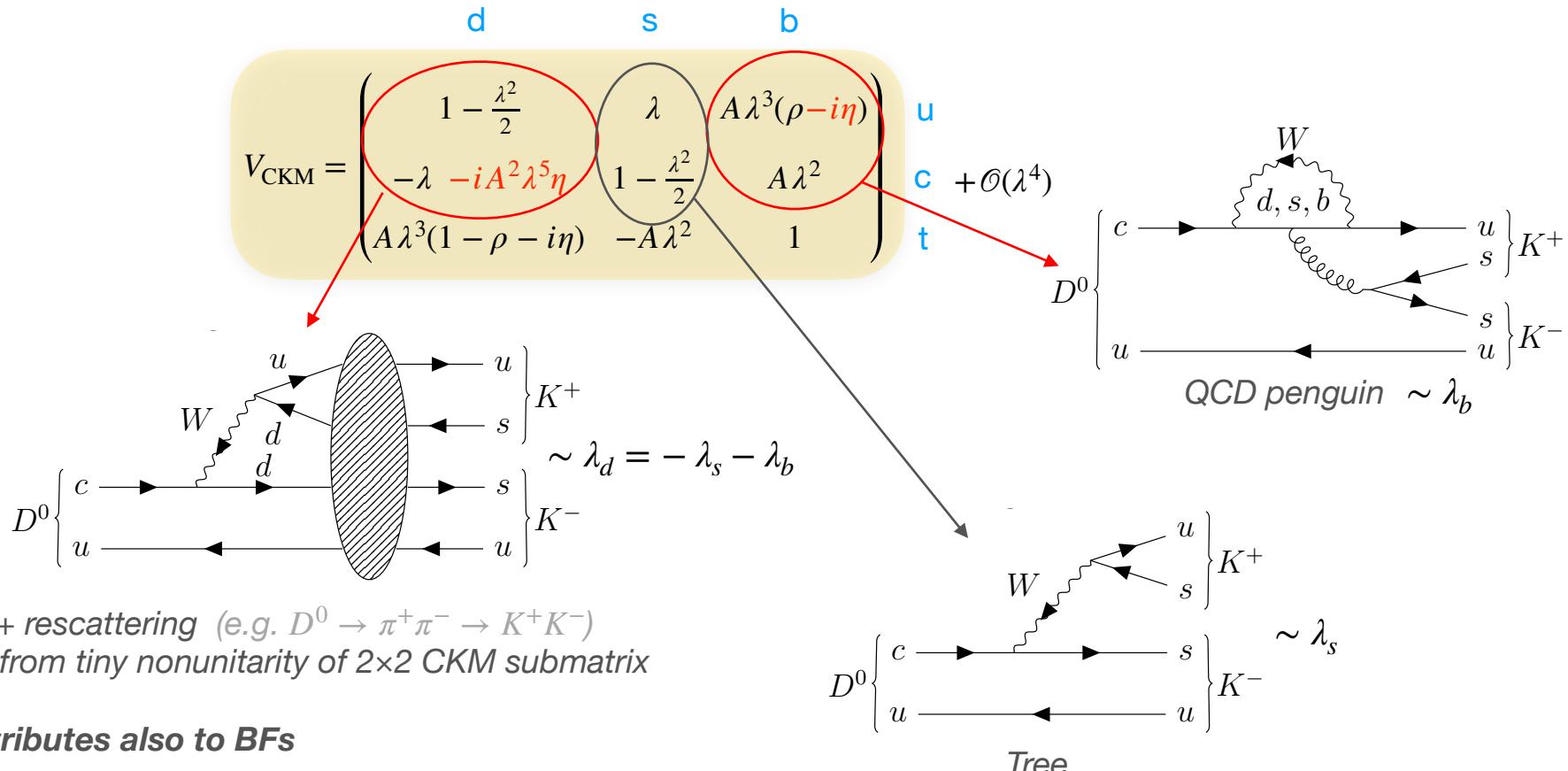
# CPV in charm

$$A(D^0 \rightarrow K^- K^+) \equiv \lambda_s T + \lambda_b P$$

Detectable in Cabibbo-suppressed decays only

$$a_{K^+ K^-}^d \approx -2 \operatorname{Im}\left(\frac{\lambda_b}{\lambda_s}\right) \left| \frac{P}{T} \right| \sin(\delta_P - \delta_T)$$

$1.3 \times 10^{-3}$



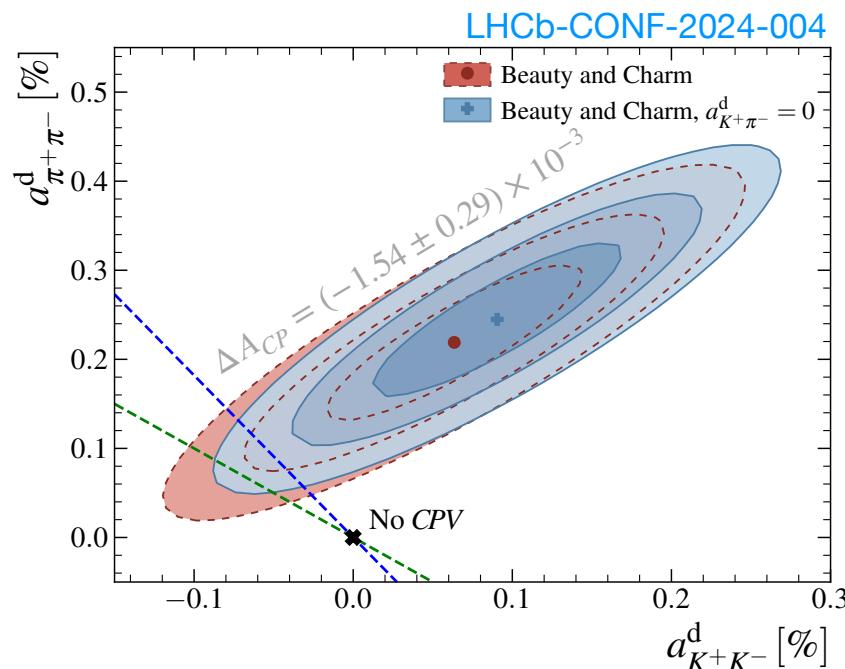
# The $D^0 \rightarrow h^+h^-$ puzzle

## 1. Branching fractions

$$\left| \frac{1}{\tan^2 \theta_C} \frac{A(D^0 \rightarrow K^+ \pi^-)}{A(D^0 \rightarrow K^- \pi^+)} \right| = 1.15 \pm 0.03 \quad O(20\%) \text{ U-spin breaking}$$

$$\left| \frac{A(D^0 \rightarrow K^+ K^-)}{A(D^0 \rightarrow \pi^+ \pi^-)} \right| = 1.81 \pm 0.02 \quad O(1) \text{ U-spin breaking?}$$

## 2. CP asymmetries



$$a_{K^-K^+}^d = (0.9 \pm 0.5) \times 10^{-3}$$
$$a_{\pi^-\pi^+}^d = (2.4 \pm 0.6) \times 10^{-3}$$

$\downarrow$

$$\left| \frac{P}{T} \right| \sin(\delta_P - \delta_T) = O(1)$$

Strict U-spin limit:  $\frac{a_{K^+K^-}^d}{a_{\pi^+\pi^-}^d} = -1$

Improved U-spin limit:  $\frac{a_{K^+K^-}^d}{a_{\pi^+\pi^-}^d} \left| \frac{A(D^0 \rightarrow K^+ K^-)}{A(D^0 \rightarrow \pi^+ \pi^-)} \right| = -1$

# What about theory?

$$\left| \frac{P}{T} \right| \sin(\delta_P - \delta_T) = O(1)$$

- **Light cone sum rules (LCSR):**  $\left| \frac{P}{T} \right| \sim 0.1$

[Khodjamirian Petrov 1706.07780](#)

[Chala Lenz Rusov Scholtz 1903.10490](#)

[Lenz Piscopo Rusov 2312.13245](#)

- **Predictions based on  $\pi\pi/KK$  rescattering data:** under debate

[Franco Mishima Silvestrini 1203.3131](#)

[Bediaga Frederico Magalhaes 2203.04056](#)

[Pich Solomonidi Vale Silva 2305.11951](#)

- **Working hypothesis of large nonperturbative effects** leading to  $\left| \frac{P}{T} \right| = O(1)$

can explain both CP asymmetries and BFs at the same time,  
with O(20%) U-spin breaking

- U-spin average of the penguin contraction is responsible for CPV
- broken penguin is responsible for U-spin breaking in the BFs

[Brod Grossman Kagan Zupan 1203.6659](#)

[Grossman Schacht 1903.10952](#)

Example of concrete model: enhancement through  $f_0(1710)$  and  $f_0(1790)$

[Schacht Soni 2110.07619](#)

# What about theory?

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Experiment

$$\left| \frac{P}{T} \right| \sin(\delta_P - \delta_T) = O(1)$$

Can we measure the phase?  
If  $\ll 1$ , will disfavour SM explanations

- **Predictions based on  $\pi\pi/KK$  rescattering data:** under debate

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How useful would further experimental data be?

- **Working hypothesis of large nonperturbative effects** leading to  $\left| \frac{P}{T} \right| = O(1)$

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Example of concrete model: enhancement through  $f_0(1710)$  and  $f_0(1790)$

Schacht Soni 2110.07619

Coupled channel analysis, e.g. in CEP?

Can we test U-spin symmetry more thoroughly in CP asymmetries and BFs?

Can we say something more about the possible sources of enhancement,  
making predictions in other decay channels?

- E.g.  $f_0$  could enhance also  $\text{BF}(D^0 \rightarrow \gamma\gamma)$ , but not CPV in  $D^+ \rightarrow K^+K^+K^-$  (due to phase space)

# Measuring $\arg(P/T)$ from time-dependent CPV?

Kagan Silvestrini 2001.07207  
LHCb 2105.09889

$$A_{CP}^f(t) \approx a_f^d + \Delta Y_f \frac{t}{\tau_{D^0}}$$

$$\Delta Y_f \approx -x_{12}\phi_2^M + y_{12}a_f^d + \boxed{x_{12}a_f^d} \cot \delta_f$$

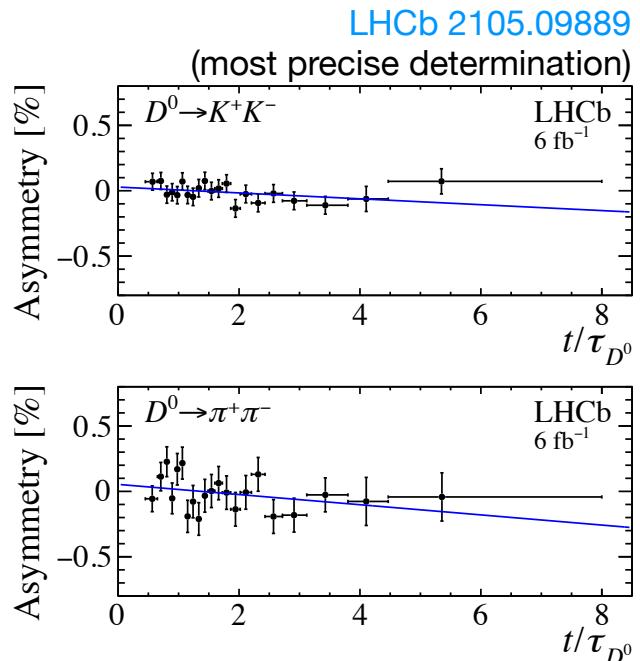
can be determined from

$$D^0 \rightarrow K^+ \pi^-$$

$$D^0 \rightarrow K^+ \pi^- \pi^+ \pi^-$$

$$D^0 \rightarrow K_S^0 \pi^+ \pi^-$$

$0.08 \times 10^{-4}$   
(for  $a_f^d = 2 \times 10^{-3}$ )



WA 2025 (charm-fitter)

$$\Delta Y_{K^+ K^-} = (-0.20 \pm 1.28 \pm 0.32) \times 10^{-4}$$

$$\Delta Y_{\pi^+ \pi^-} = (-3.5 \pm 2.4 \pm 0.4) \times 10^{-4}$$

Projections at LHCb Upgrade II 1808.08865

$$\sigma(\Delta Y_{K^+ K^-}) = 0.14 \times 10^{-4}$$

$$\sigma(\Delta Y_{\pi^+ \pi^-}) = 0.25 \times 10^{-4}$$

LHCb-CONF-2024-004

$$x_{12} = (4.1 \pm 0.5) \times 10^{-3}$$

$$y_{12} = (6.2 \pm 0.2) \times 10^{-3}$$

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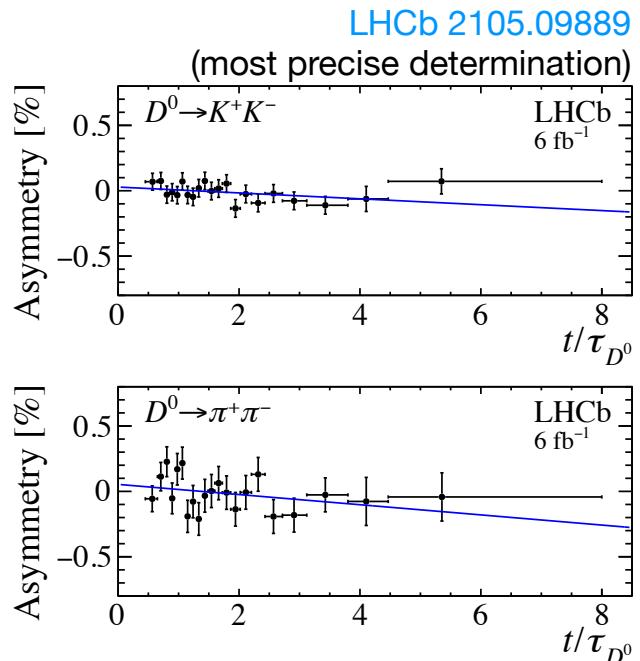
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can be determined from  
 $D^0 \rightarrow K^+\pi^-$   
 $D^0 \rightarrow K^+\pi^-\pi^+\pi^-$   
 $D^0 \rightarrow K_S^0\pi^+\pi^-$

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can get some sensitivity only if  $|\delta_f| \ll 1$   
 (unlikely, but would strengthen the tension  
 with SM expectations)



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$$A_{CP}(f) \approx a_f^d + \Delta Y_f \frac{\langle t \rangle_f}{\tau_{D^0}}$$

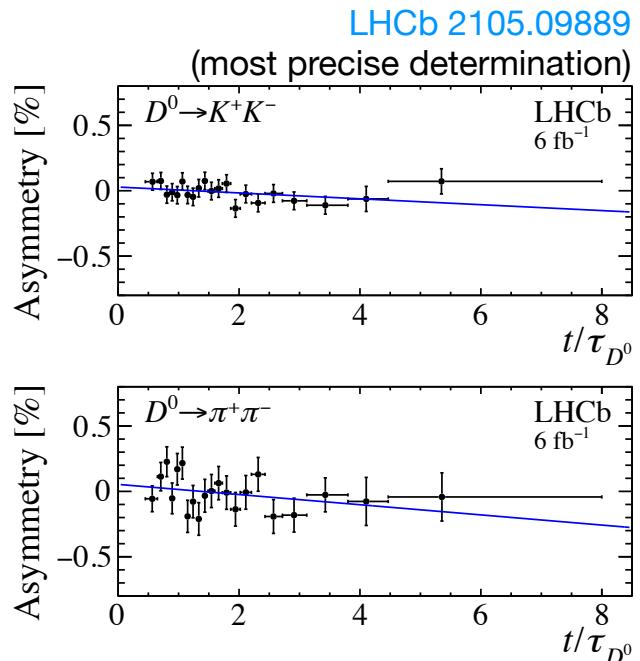
$\sim 1.7$

LHCb 1903.08726

$$\Delta A_{CP} = A_{CP}(K^+ K^-) - A_{CP}(\pi^+ \pi^-) = (15.4 \pm 2.9) \times 10^{-4}$$

All results by LHCb/HFLAV neglect the  $x_{12}a_f^d \cot \delta_f$  contribution to  $\Delta Y_f$  (achieve better uncertainty, but prone to bias in case  $|\delta_f| \ll 1$ )

→ alternative fit where this approximation is removed foreseen for the next update of LHCb-CONF-2024-004



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LHCb-CONF-2024-004  
 $x_{12} = (4.1 \pm 0.5) \times 10^{-3}$   
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# Measuring $\arg(P/T)$ from a Dalitz plot?

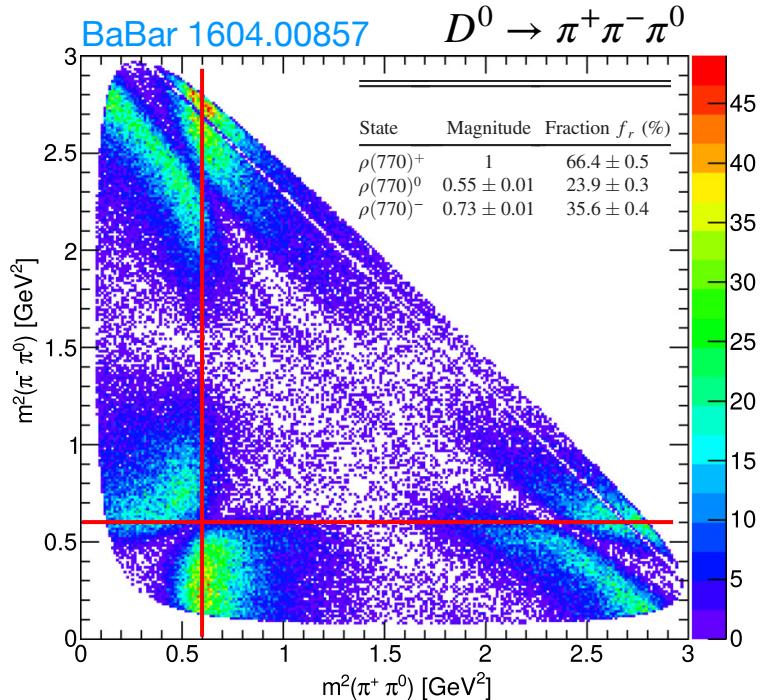
$\delta_{QCD}$  varies across the Dalitz plot. However, it “dresses” T and P in the same way

- not visible in the asymmetry of the fit fractions
- visible in the interference with other resonances

Dery Grossman Schacht Soffer 2101.02560

Ideal scenario:

- significant overlap of two resonances in a region where  $\delta_{QCD}$  change fast
- amplitudes similar in magnitude



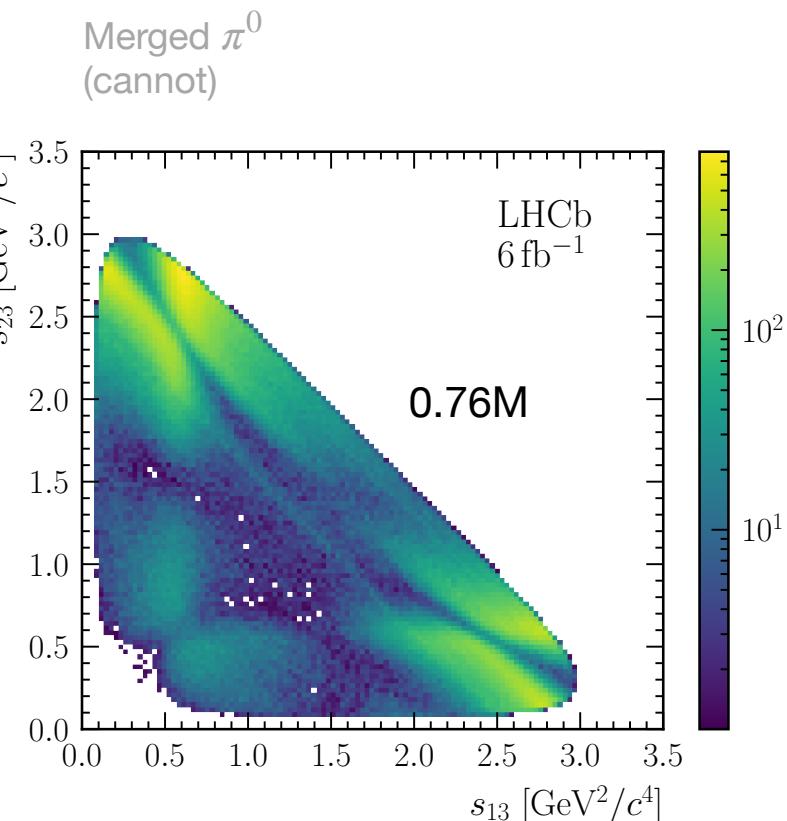
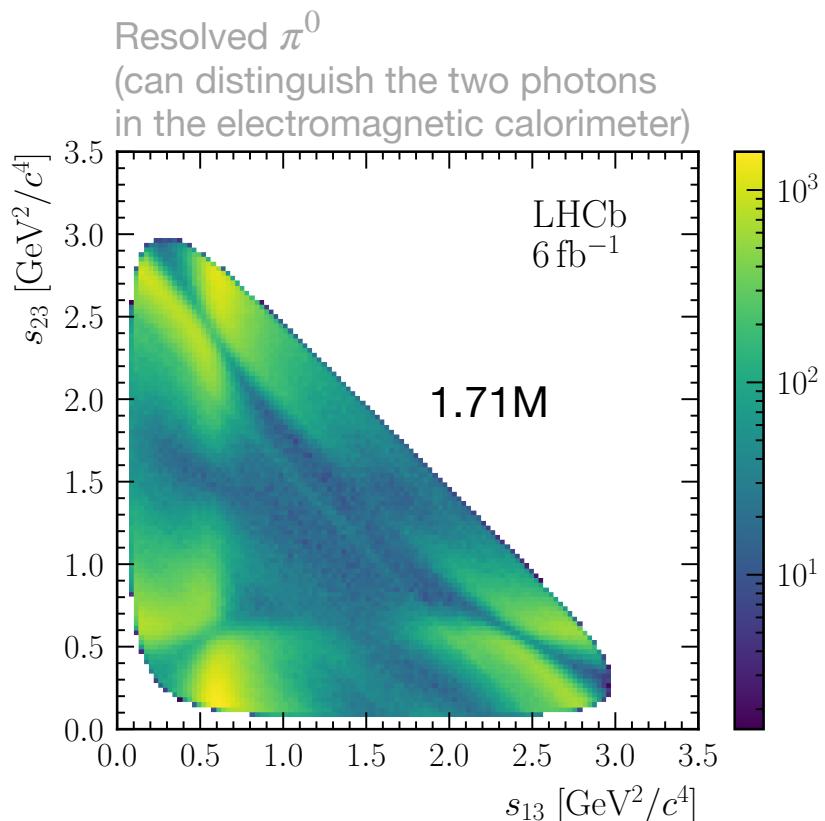
$$A_1(s_{12}, s_{13}) \sim \left[ T_1 + \left( \frac{\lambda_b}{\lambda_d} \right) P_1 e^{i\delta_1} \right] r_1(s_{12}, s_{13}) e^{i\delta_{r_1}(s_{12}, s_{13})}$$

$$A_2(s_{12}, s_{13}) \sim \left[ T_2 + \left( \frac{\lambda_b}{\lambda_d} \right) P_2 e^{i\delta_2} \right] r_2(s_{12}, s_{13}) e^{i\delta_{r_2}(s_{12}, s_{13})}$$

# $D^0 \rightarrow \pi^+ \pi^- \pi^0$ at LHCb

Search for CPV based on a model-independent method (Energy Test) published with Run 2 data

- consistent with CP conservation



Amplitude analysis based on the same data sample ongoing

# Clean discovery channels

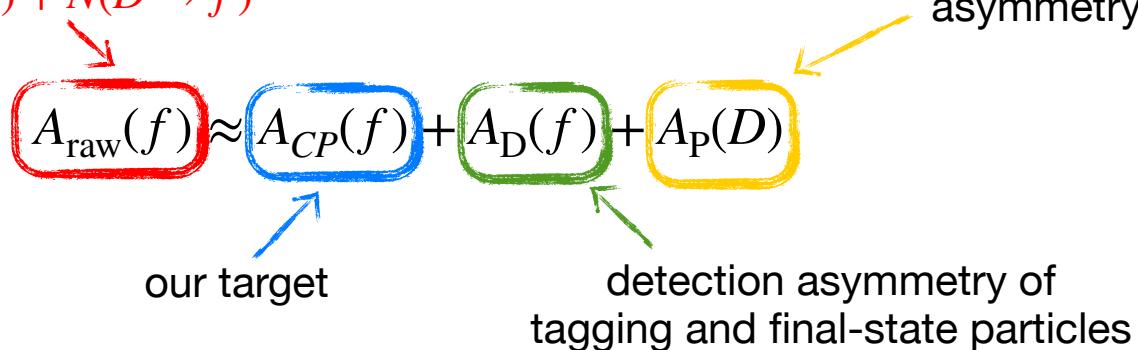
- Isospin selection rules, e.g.  $A_{CP}(D^+ \rightarrow \pi^+\pi^0) = 0$  ( $\Delta I = 3/2$ )
- $SU(3)_F$ -based sum rules  
(if you believe that U-spin symmetry is broken only at  $O(20\%)$  level)
- $BF(D^0 \rightarrow \mu^+\mu^-)$ , angular/CP asymmetries in  $c \rightarrow u\ell^+\ell^-$
- LFU ( $D^0 \rightarrow \mu^+e^-$  etc.)

## However

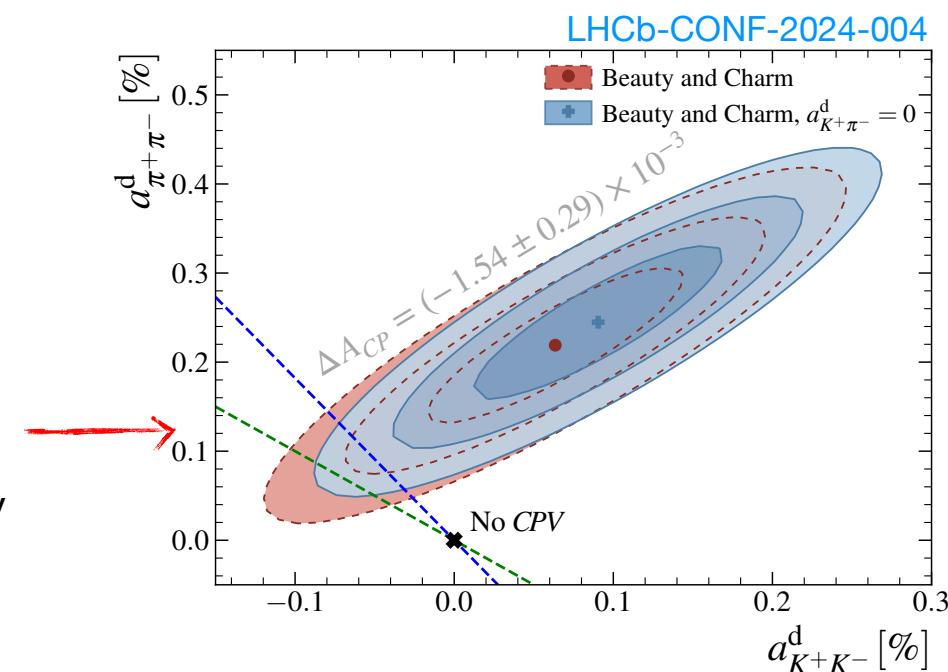
- most of  $SU(3)$  sum rules include at least one  $\pi^0$ 
  - lower precision than final states with charged particles only
  - even worse if coming from  $D_{(s)}^+$  rather than  $D^0$  (no  $m(D^0\pi^+) - m(D^0)$  trick)
- it is possible that we observe CPV in many channels, but won't be able to test the corresponding sum rules with comparable precision
  - improved understanding of long-distance contributions is imperative to eventually assess the compatibility of CPV with the SM

# A side note on control channels

$$A_{\text{raw}}(f) = \frac{N(D \rightarrow f) - N(\bar{D} \rightarrow \bar{f})}{N(D \rightarrow f) + N(\bar{D} \rightarrow \bar{f})}$$



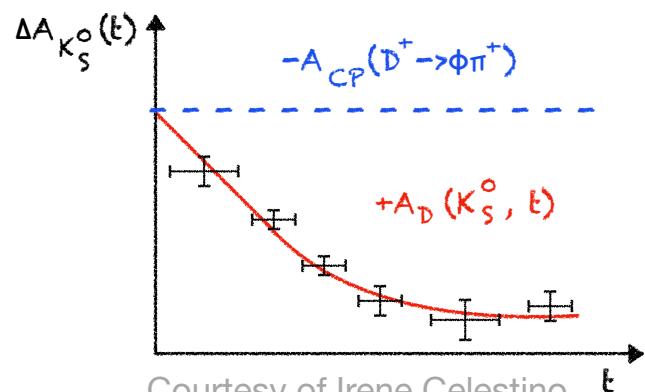
- Few observables are self-calibrating:  
 $\Delta A_{CP}, \Delta Y(D^0 \rightarrow h^+h^-), y_{CP}(D^0 \rightarrow h^+h^-)$   
 $\frac{\Gamma(D^0 \rightarrow K^+\pi^-(\pi^+\pi^-))}{\Gamma(D^0 \rightarrow K^-\pi^+(\pi^+\pi^-))}, D^0 \rightarrow K_S^0\pi^+\pi^-$   
variation of CPV across the phase space of multi-body decays
- Otherwise we rely on calibration channels
  - Cabibbo-favoured decays ( $A_{CP} = 0$ )
  - often include  $K_S$  and limit our sensitivity (e.g.  $D^+ \rightarrow K_S^0\pi^+$  or  $D_s^+ \rightarrow K_S^0K^+$ )



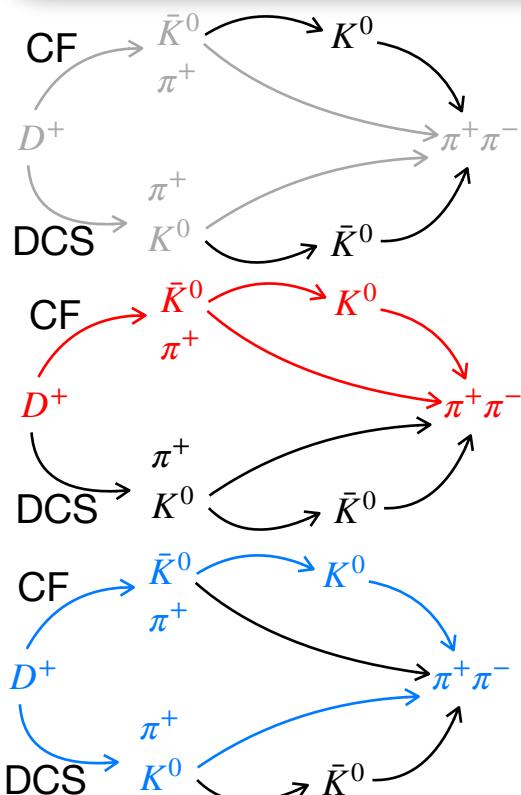
# Example: $A_{CP}(D^+ \rightarrow \phi\pi^+)$

$$\Delta A_{K_S^0}(t) \equiv A_{raw}(D^+ \rightarrow K_S^0(t)\pi^+) - A_{raw}(D^+ \rightarrow \phi\pi^+)$$

$$= \boxed{A_D(K_S^0, t)} - \boxed{A_{CP}(D^+ \rightarrow \phi\pi^+)}$$



$$A_D(K_S^0, t) = A_{CP}^{\text{decay}} + A_{CP}^{\bar{K}^0}(t) + A_{CP}^{\text{int}}(t) + \text{regeneration (and interference)}$$



CPV in charm decay

$$\text{Im}\left(\frac{V_{cd}^* V_{us}}{V_{cs}^* V_{ud}}\right) \approx 10^{-5}$$

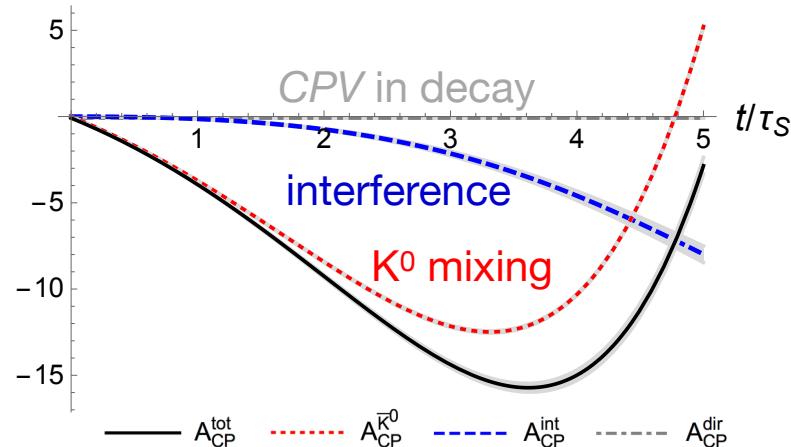
CPV in kaon mixing  
 $\approx \text{Re}(\epsilon) \approx 10^{-3}$

CPV in interference between  
 $D^+$  decay and kaon mixing

$$\text{Im}(\epsilon) \text{Re}\left(\frac{V_{cd}^* V_{us}}{V_{cs}^* V_{ud}}\right) \approx 10^{-3}$$

$$A_{CP}(t)[\times 10^{-3}]$$

Yu Wang Li 1707.09297



Interference crucially depends on:

$$\frac{A(D^+ \rightarrow K^0\pi^+)}{A(D^+ \rightarrow \bar{K}^0\pi^+)} \equiv r_\pi e^{i\delta_\pi}$$

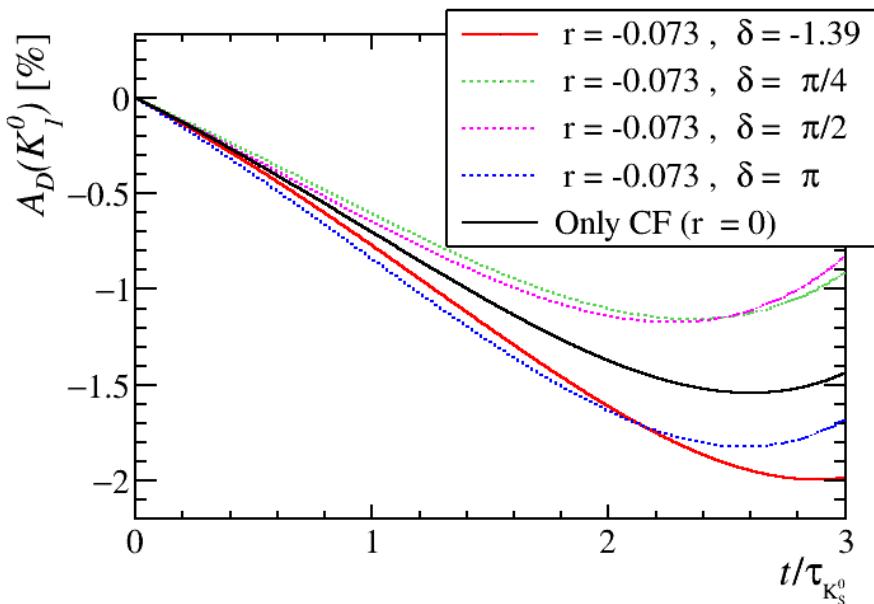
# Kind request to our theory colleagues

$$\frac{A(D^+ \rightarrow K^0\pi^+)}{A(D^+ \rightarrow \bar{K}^0\pi^+)} \equiv r_\pi e^{i\delta_\pi}$$

The values of these parameters have been extracted from data only in the FAT approach:

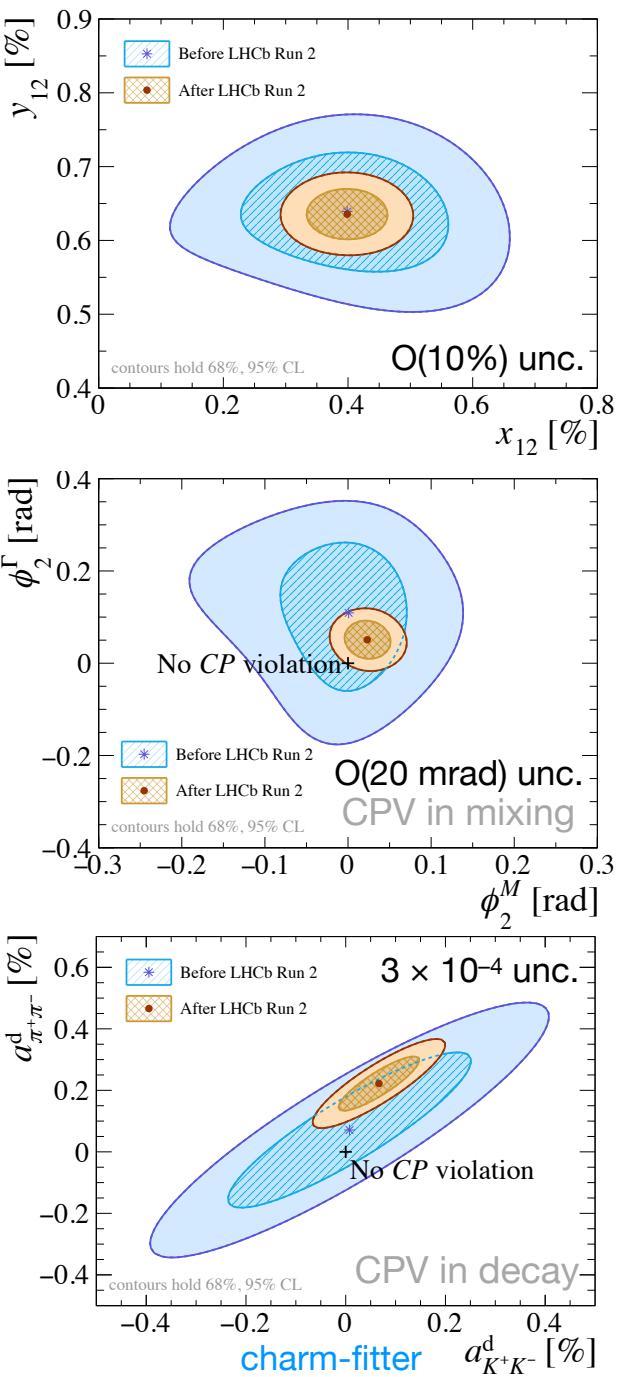
$$r_{\pi^+} = -0.073 \pm 0.004, \quad \delta_{\pi^+} = -1.39 \pm 0.05, \quad \text{Yu Wang Li 1707.09297}$$

Systematic uncertainties are debated. Can alternative estimates be provided? Otherwise treating them as nuisance parameters will limit our precision in Upgrade I



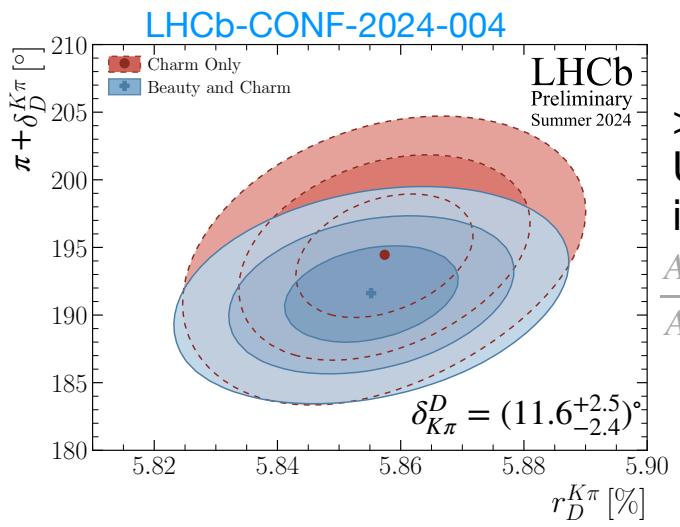
# Conclusions

- After 50 years and huge recent experimental progress, charm is still a “young” field with many open questions and promises many developments in the near future
- Theory input essential for the success of our upgrades
  - if your favourite measurement is not being performed at LHCb, the time to complain is now
    - e.g. we are adding  $D^0 \rightarrow \eta^{(\prime)}\eta$  in Run 3
  - do not forget “boring” CP-conserving measurements (BFs etc.)



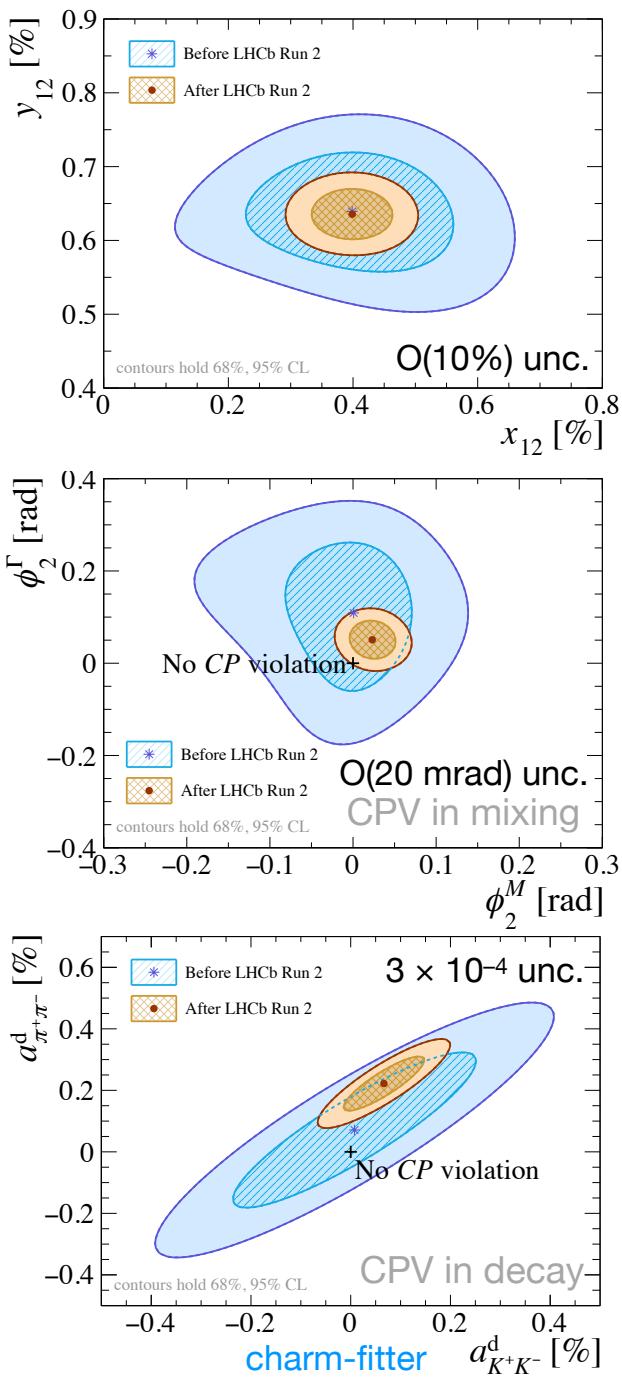
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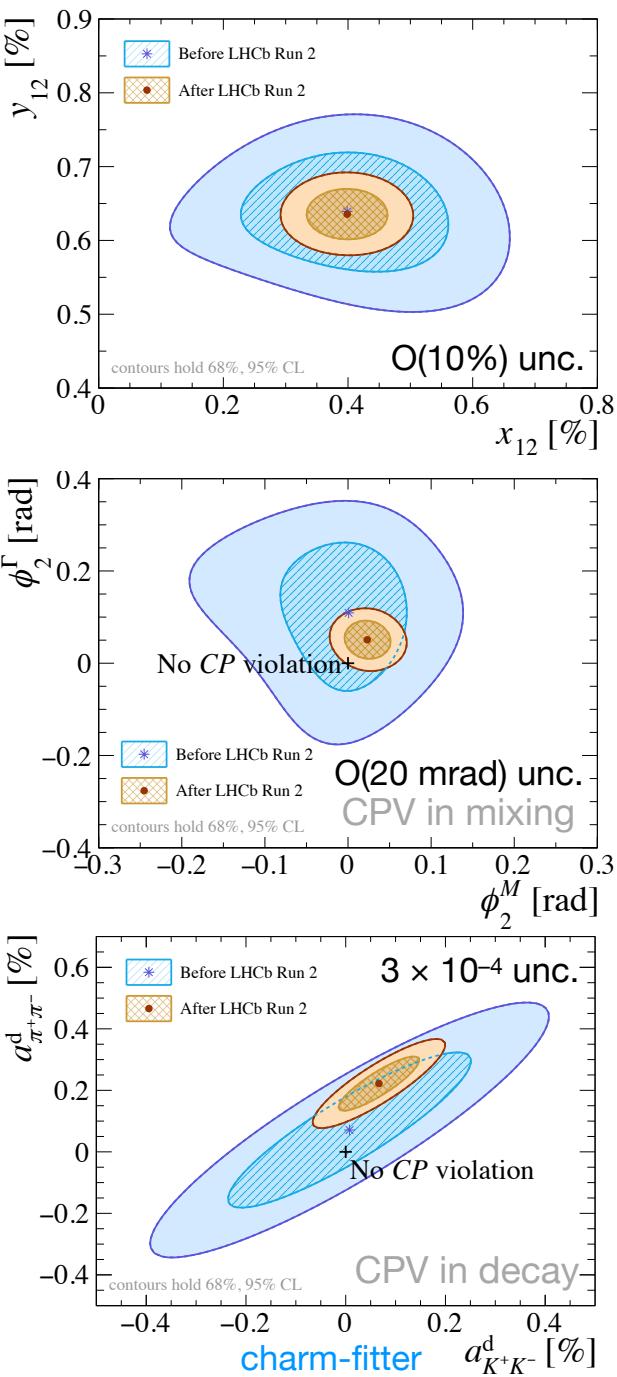
> 4 $\sigma$  evidence for  
U-spin breaking  
in the phase

$$\frac{A(D^0 \rightarrow K^+\pi^-)}{A(\bar{D}^0 \rightarrow K^+\pi^-)} \approx -r_D^{K\pi} e^{-i\delta_D^{K\pi}}$$



# Conclusions

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    - e.g. we are adding  $D^0 \rightarrow \eta^{(\prime)}\eta$  in Run 3
  - do not forget “boring” CP-conserving measurements (BFs etc.)
- Multibody decays are becoming more and more important. We welcome feedback on how to best tackle them (besides amplitude analyses) to make the results most useful for theorists while waiting for new observations
  - binned measurements
  - triple-product asymmetries
  - etc.



# Backup slides

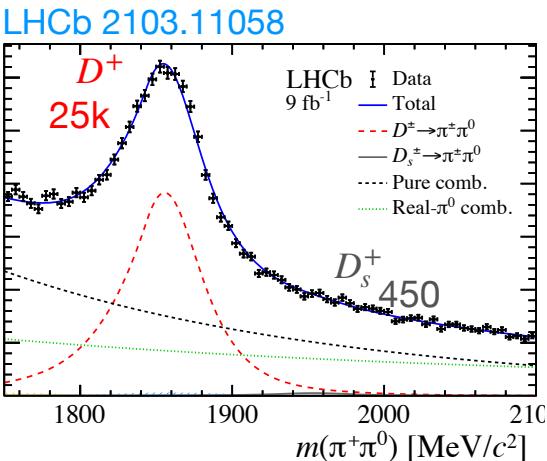


# Neutral efficiencies

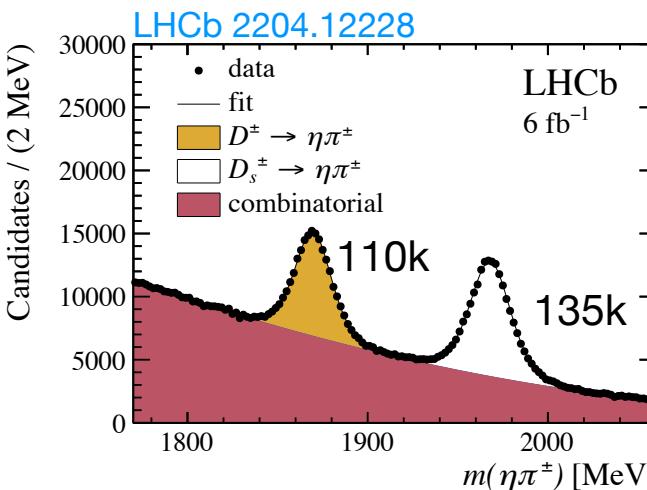
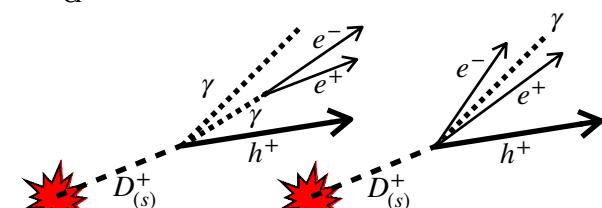
$$\epsilon(\pi^0/\gamma) < \epsilon(\eta^{(\prime)}) < \epsilon(K_S^0)$$

my personal, rough rule of thumb for Run 2:

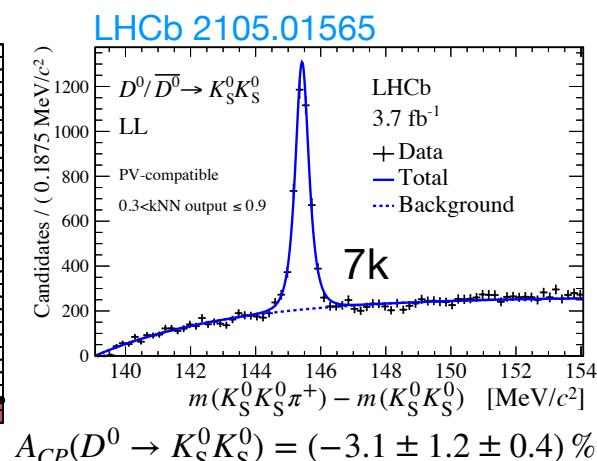
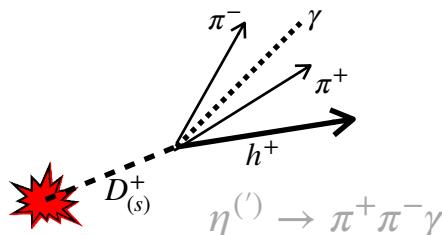
~3% ... ~10%



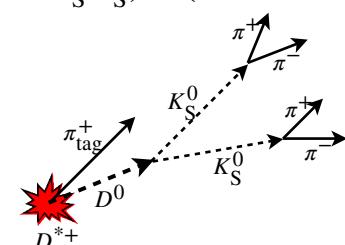
$$A_{CP}(D^+ \rightarrow \pi^+ \pi^0) = (-1.3 \pm 0.9 \pm 0.6)\%$$



$$A_{CP}(D^+ \rightarrow \eta \pi^+) = (0.34 \pm 0.66 \pm 0.14)\%$$



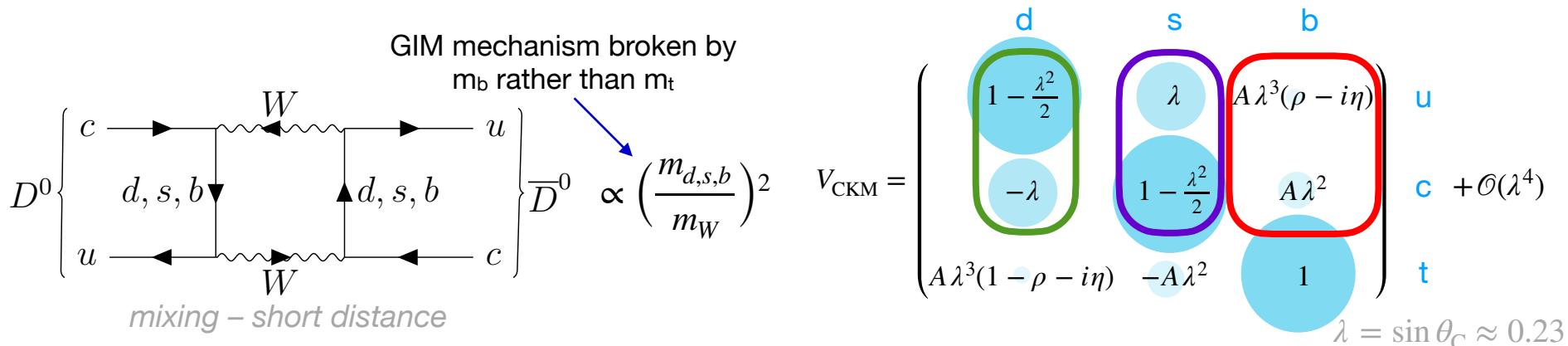
$$A_{CP}(D^0 \rightarrow K_S^0 \bar{K}_S^0) = (-3.1 \pm 1.2 \pm 0.4)\%$$



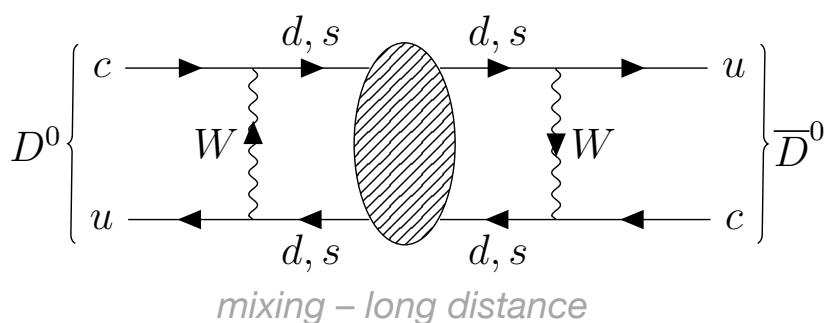
Many KS decay outside of the vertex tracker

# The peculiar phenomenology of charm

FCNC are extremely suppressed



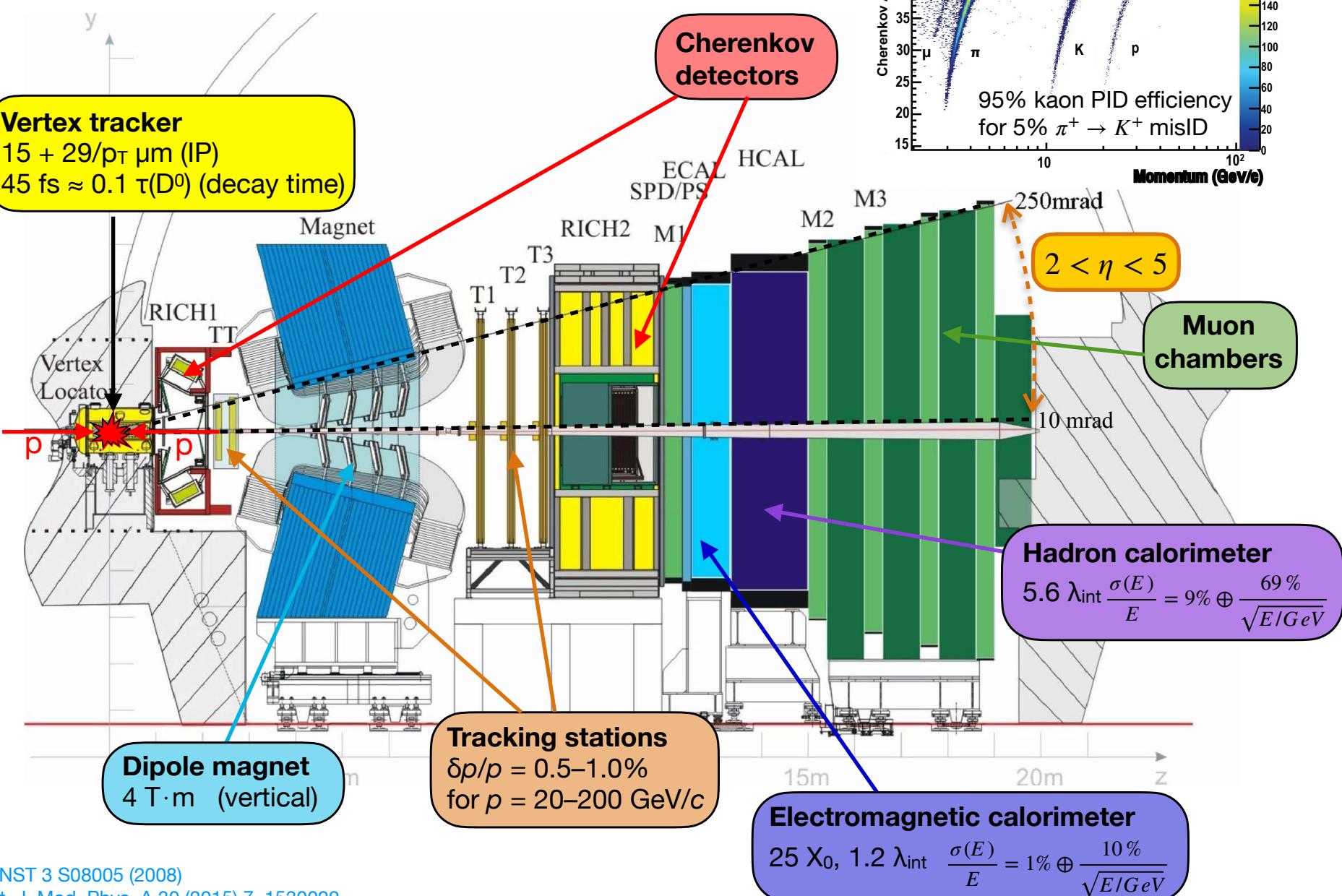
- **CKM suppression** → third generation nearly decouples from the first two
- (**d** + **s** contribution) → 0 in the limit of  $U$ -spin symmetry, i.e.  $m_s = m_d$



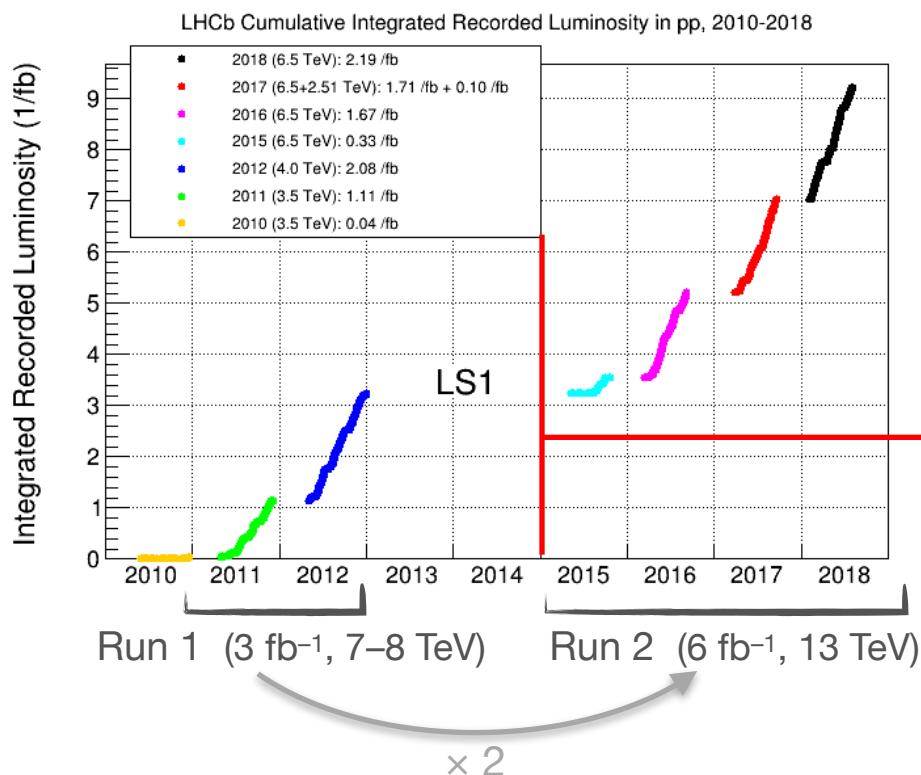
Main contributions from low-energy QCD.  
Charm is:

- too heavy for ChPT or exclusive analysis
- arguably too light for HQET ( $\Lambda_{QCD}/m_c \approx 1$ )

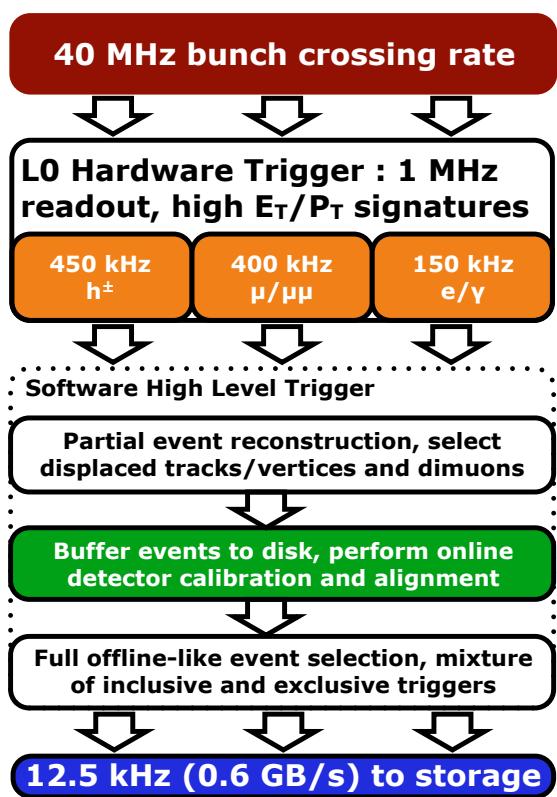
# The LHCb detector



# LHCb trigger and data-taking strategy



Pioneered by charm triggers

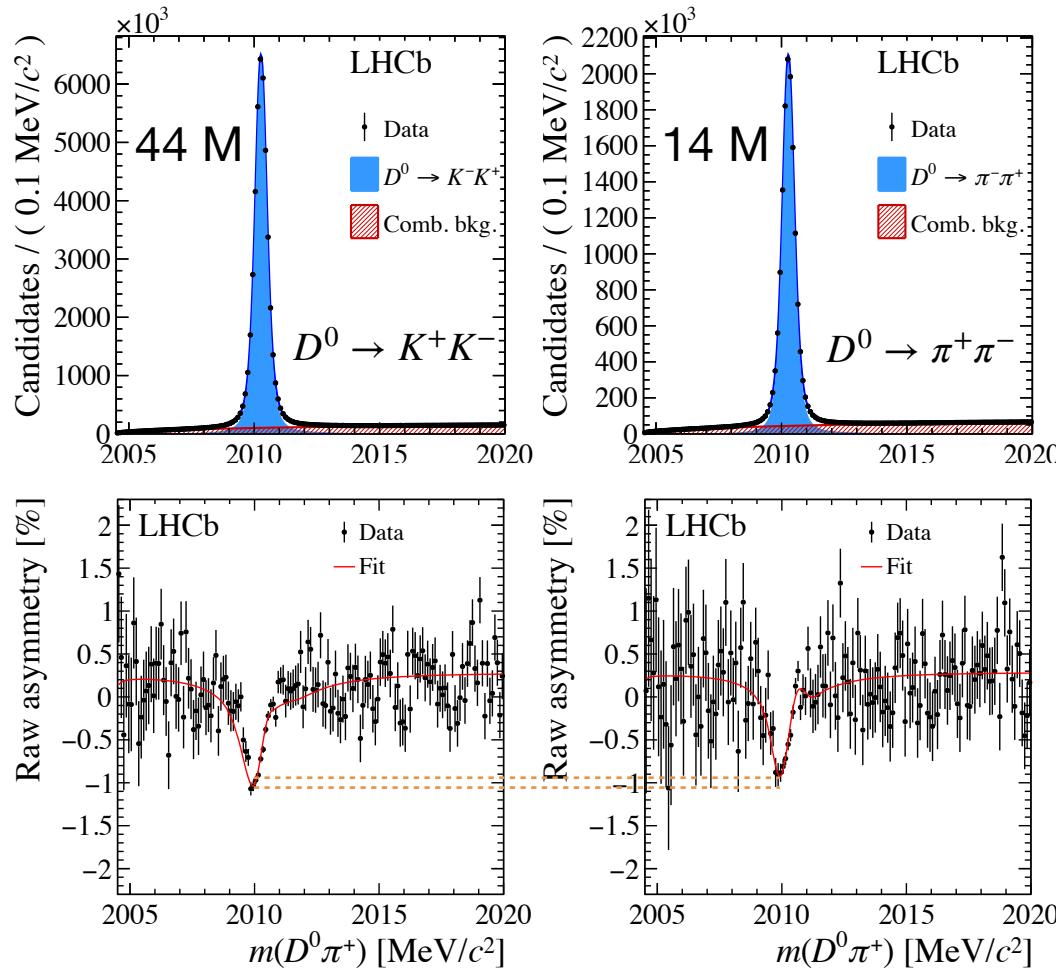


Comput. Phys. Commun. 208 (2016) 35

- Yield has more than doubled between Run 1 and 2:
- $\times 1.7$  increase in production cross-section;
  - new "Turbo" data taking paradigm
    - only the signal candidates are recorded, rest of the event is discarded;
    - improved efficiency, higher rate recorded.

# First observation of CPV

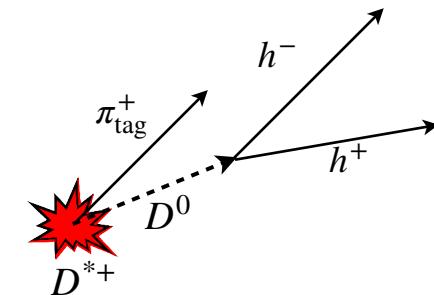
$$D^*(2010)^+ \rightarrow D^0\pi^+$$



$$\Delta A_{CP} = A_{\text{raw}}(K^+K^-) - A_{\text{raw}}(\pi^+\pi^-)$$

$$\approx a_{K^+K^-}^d - a_{\pi^+\pi^-}^d = (-1.54 \pm 0.29) \times 10^{-3} \quad (5.3\sigma)$$

PRL 122 (2019) 211803  
6 fb<sup>-1</sup>, 2015–2018



$$\frac{N(D \rightarrow f) - N(\bar{D} \rightarrow \bar{f})}{N(D \rightarrow f) + N(\bar{D} \rightarrow \bar{f})}$$

$$A_{\text{raw}}(f) \approx A_{CP}(f) + A_{\text{det}}(\pi^+) + A_{\text{prod}}(D^{*+})$$

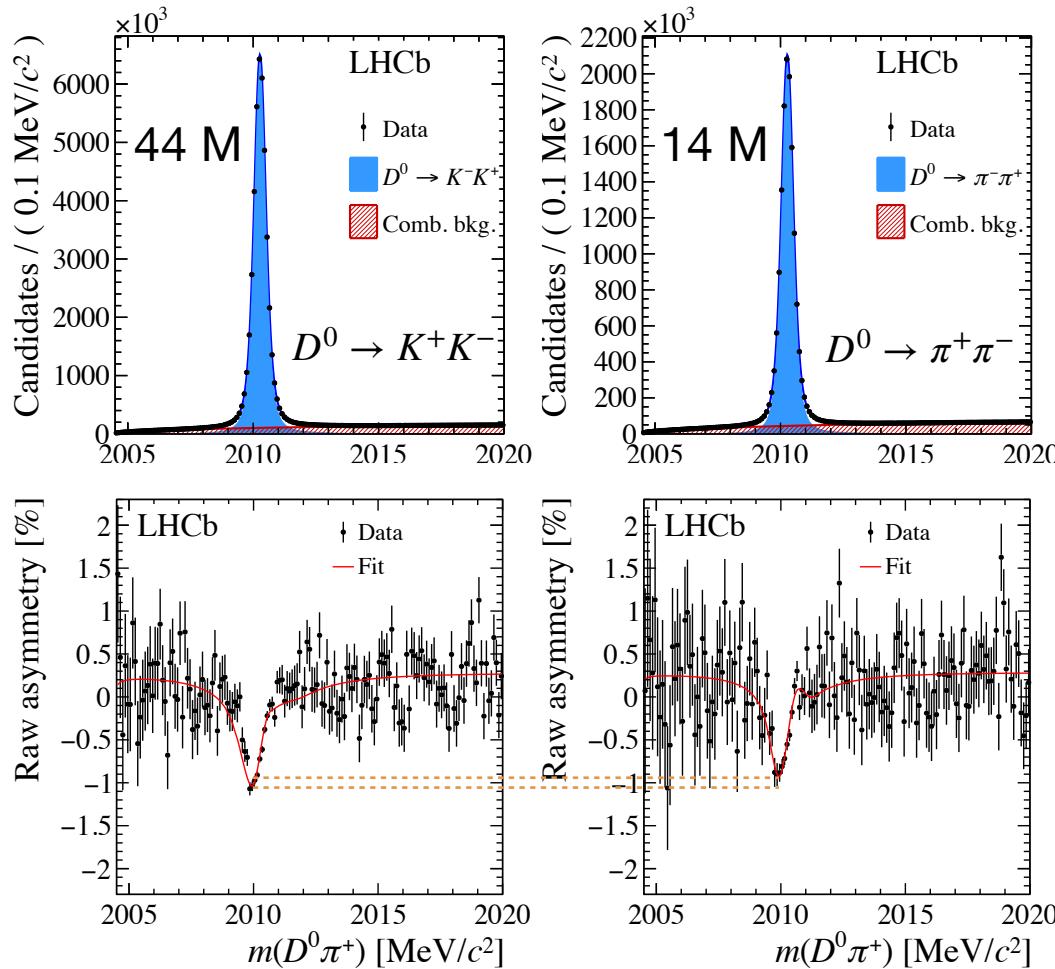
$a_f^d$  (CPV in decay)  
+ small time-dependent contributions

production cross-section

detection of tagging  $\pi^+$

# First observation of CPV

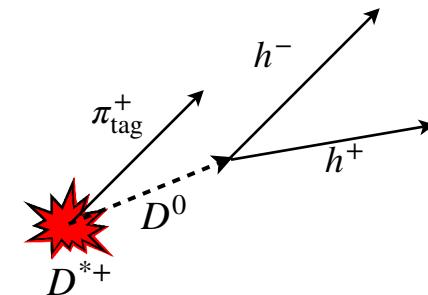
$$D^*(2010)^+ \rightarrow D^0\pi^+$$



$$\Delta A_{CP} = A_{\text{raw}}(K^+K^-) - A_{\text{raw}}(\pi^+\pi^-)$$

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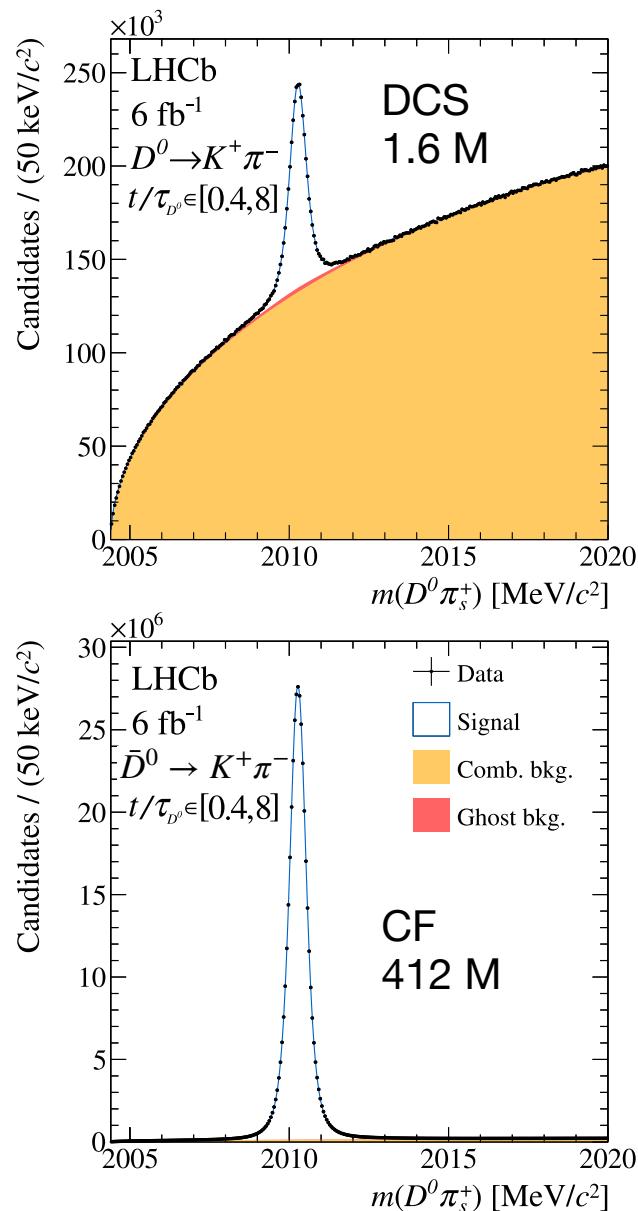
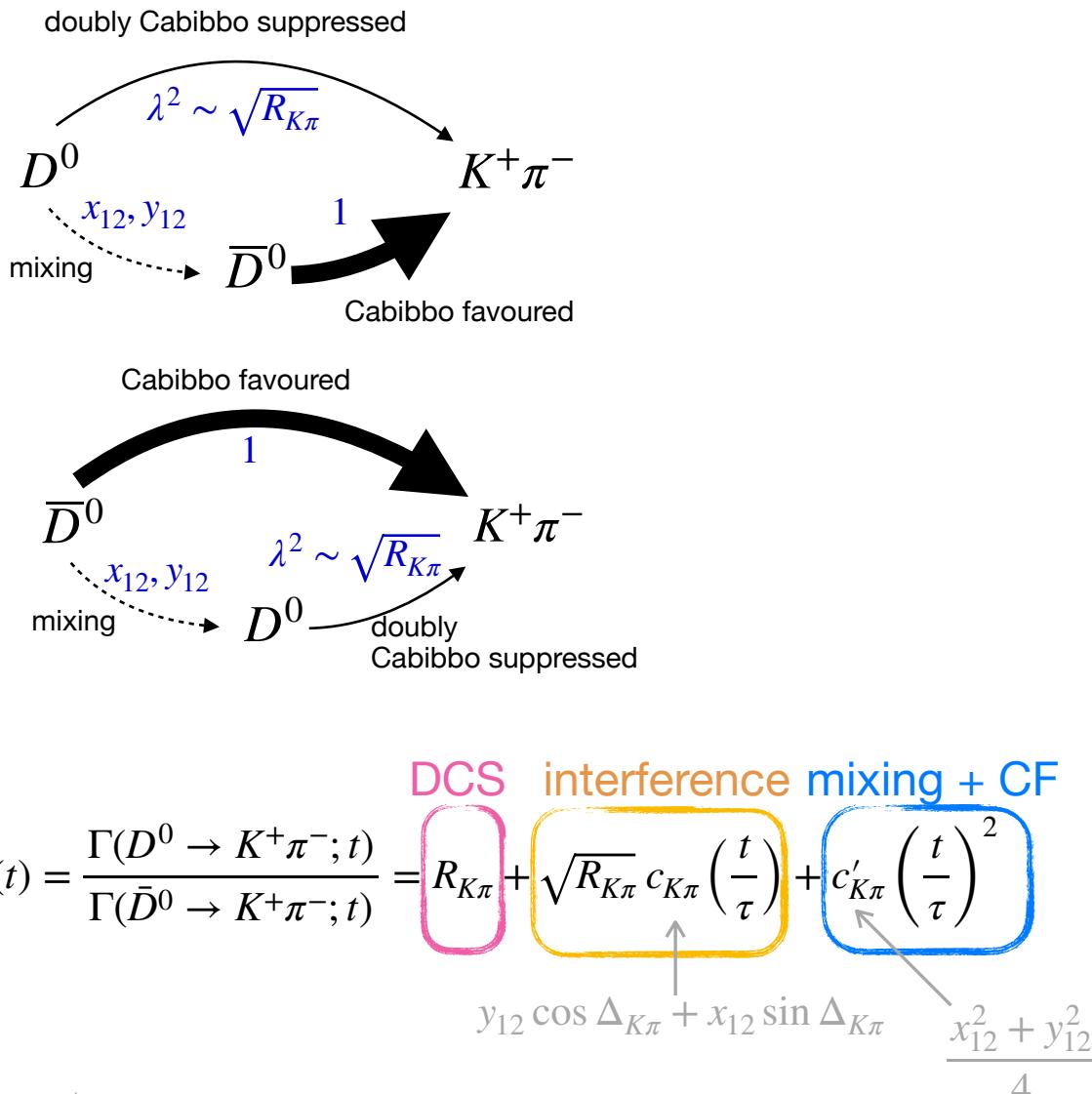
- ▶ Larger than most SM predictions

- Franco, Mishima & Silvestrini, JHEP 05 (2012) 140  
Li et al, Phys. Rev. D 86 (2012) 036012  
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Lenz, Piscopo & Rusov, JHEP 03 (2024) 151

- ▶ O(1–10) enhancement of QCD rescattering or BSM?

- Chala et al, JHEP 07 (2019) 161  
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Schacht & Soni, Phys. Lett. B 825 (2022) 136855  
Dery & Nir, JHEP 12 (2019) 104  
Wang et al, JHEP 09 (2021) 126  
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Dery et al, JHEP 05 (2021) 179  
Cheng & Chiang, Phys. Rev. D 104 (2021) 7, 073003  
Gavrilova, Grossman & Schacht, Phys. Rev. D 109 (2024) 3, 033011

# Mixing and CPV with $D^0 \rightarrow K^+ \pi^-$



# Measuring $A_{CP}(D^0 \rightarrow K^+K^-)$

Nuisance asymmetries subtracted through Cabibbo-favoured ( $c \rightarrow us\bar{d}$ ) decay channels

- no QCD penguin, no chromomagnetic dipole  $\rightarrow$  negligible CPV

$$A_{CP}(D^0 \rightarrow K^+K^-) \approx A_{\text{raw}}(\textcolor{red}{D}^{*+} \rightarrow D^0(\rightarrow K^+K^-)\textcolor{red}{\pi}^+) - A_{\text{raw}}(\textcolor{red}{D}^{*+} \rightarrow D^0(\rightarrow K^-\pi^+)\textcolor{red}{\pi}^+) \\ + A_{\text{raw}}(\textcolor{violet}{D}_s^+ \rightarrow \textcolor{magenta}{K}^+K^-\pi^+) - A_{\text{raw}}(\textcolor{violet}{D}_s^+ \rightarrow \textcolor{magenta}{K}^+K_S^0) + A_{\text{det}}(\textcolor{green}{K}^0)$$

bottleneck to final precision

regeneration and CPV in  
mixing explicitly calculated  
(cf. next slide)  
[CERN-THESIS-2014-274](#)