

News on $D - \bar{D}$

6th edition of “Beyond the flavour anomalies”

Roma, 10/04/2025



Roberto Di Palma

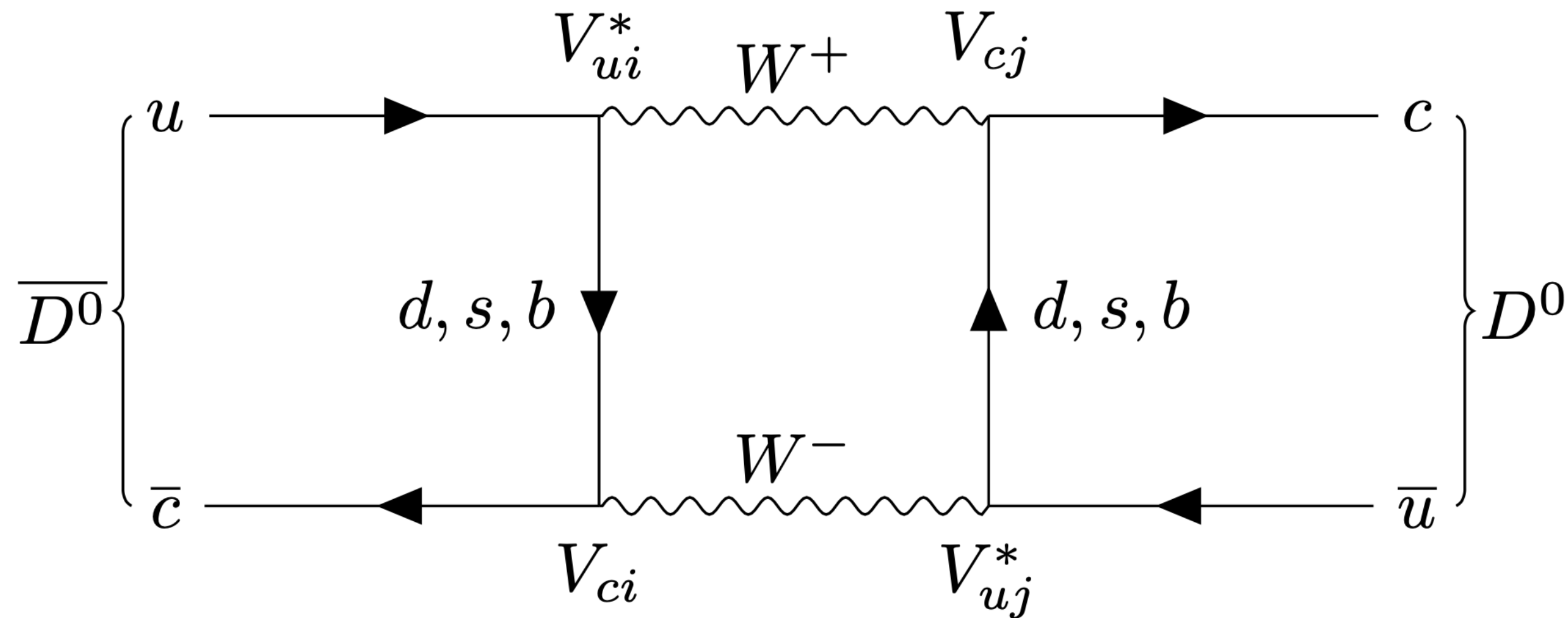
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Why charm mixing?

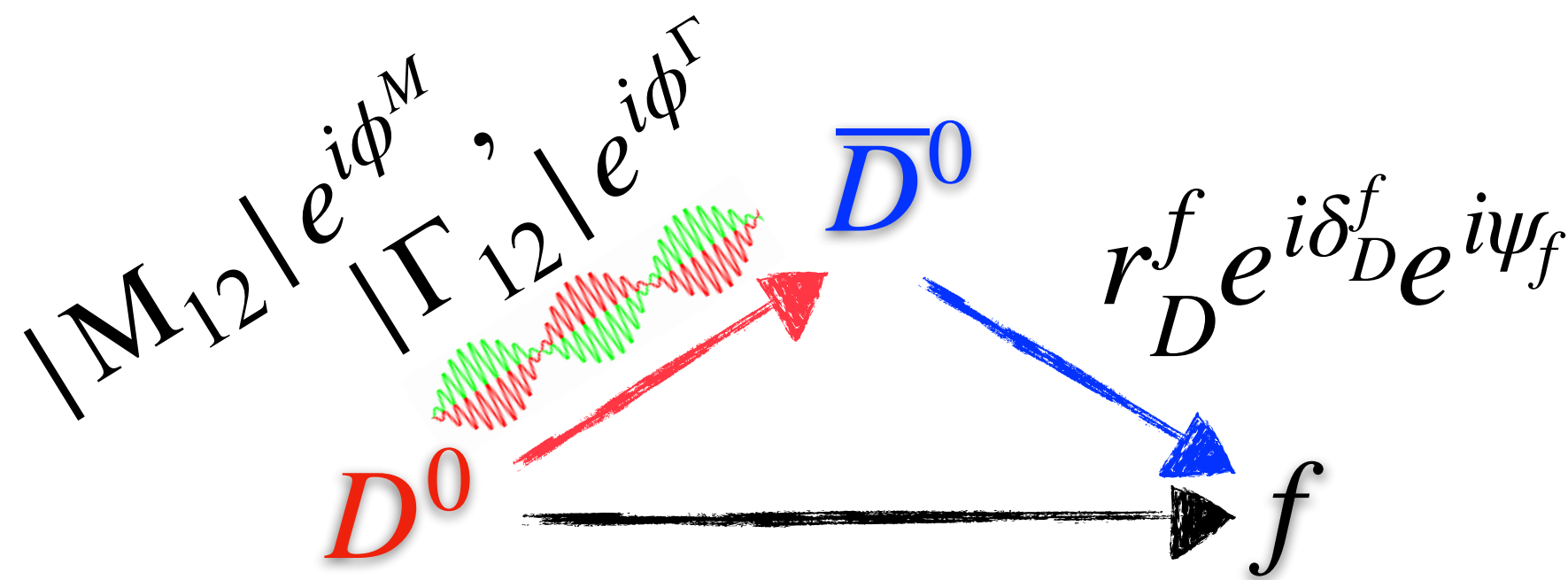
- NO FCNC at the tree-level in the SM
- From GIM + CKM: CPV is very suppressed $O(\theta_c^4)$



$$H = M - i/2\Gamma$$

Description of indirect CPV

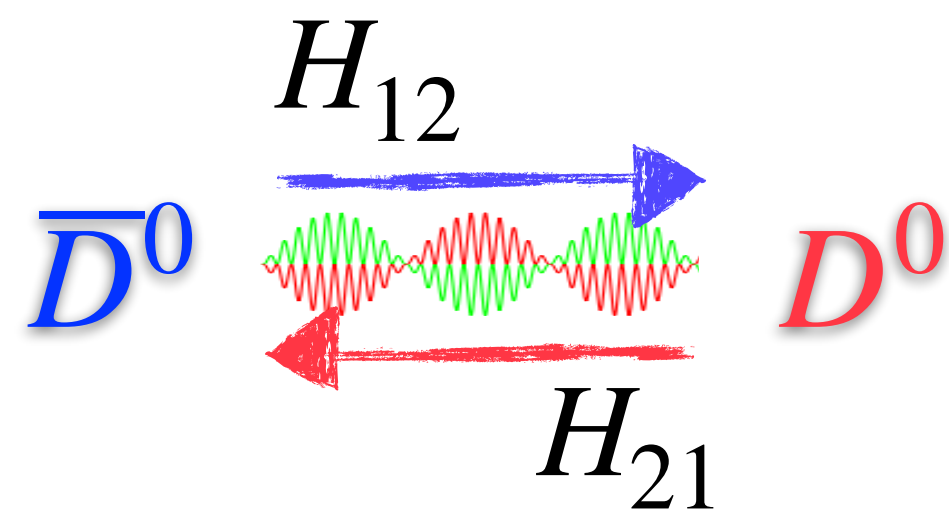
- CPV in the interference between mixing and decay to a state f



$$\phi_f^M = \phi^M + \psi_f$$

$$\phi_f^\Gamma = \phi^\Gamma + \psi_f$$

- CPV in pure mixing



$$\phi_{12} = \phi_f^M - \phi_f^\Gamma$$

Experimental measurements of CPV

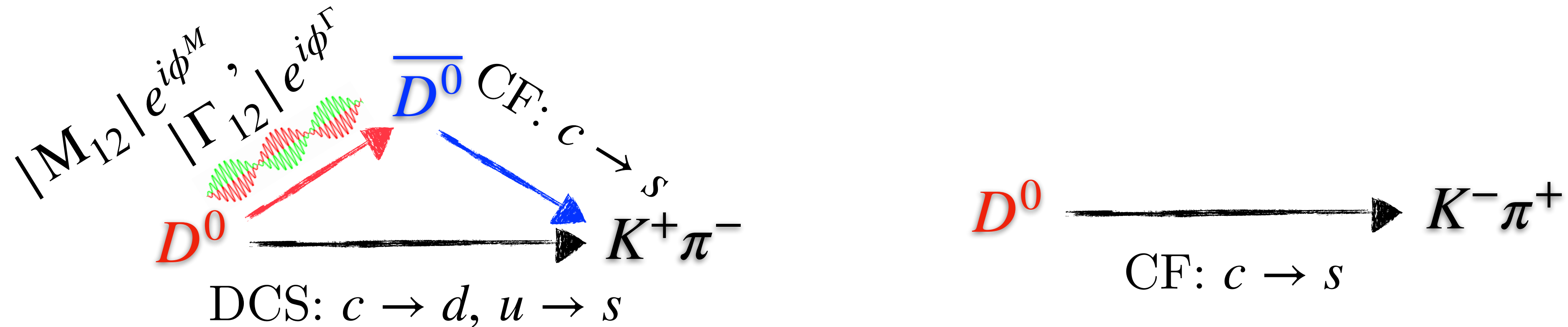
More than 15 years of experiments...



- CF/DCS decays to $K^{\mp}\pi^{\pm}(\pi^0)$
- Phase-space analyses of the three-body mode $K_S^0\pi^+\pi^-$
- SCS decays to $\pi^+\pi^-$, K^+K^- (also direct CPV: $|A_D^f| \neq |A_{\bar{D}}^f|$)

CF/DCS decays to $K^{\mp}\pi^{\pm}$

- Exploiting the interference between the CF and the DCS decays



- They fit the ratio of these two modes in time

(e.g. Belle PRL(2006), Babar PRL(2007), LHCb(2024, 2025))

$$\frac{d\Gamma(D^0 \rightarrow K^+\pi^-)}{d\Gamma(D^0 \rightarrow K^-\pi^+)}(t) = (r_D^f)^2 + r_D^f L^+(\phi_{K\pi}^M, \phi_{K\pi}^\Gamma)t/\tau + Q^+(\phi_{K\pi}^M, \phi_{K\pi}^\Gamma)(t/\tau)^2$$

Three-body decays

- They measure the Dalitz distribution of the events

$$d\Gamma_i(D^0 \rightarrow K_S^0 \pi^+ \pi^-)(t) \quad d\Gamma_i(\bar{D}^0 \rightarrow K_S^0 \pi^+ \pi^-)(t)$$

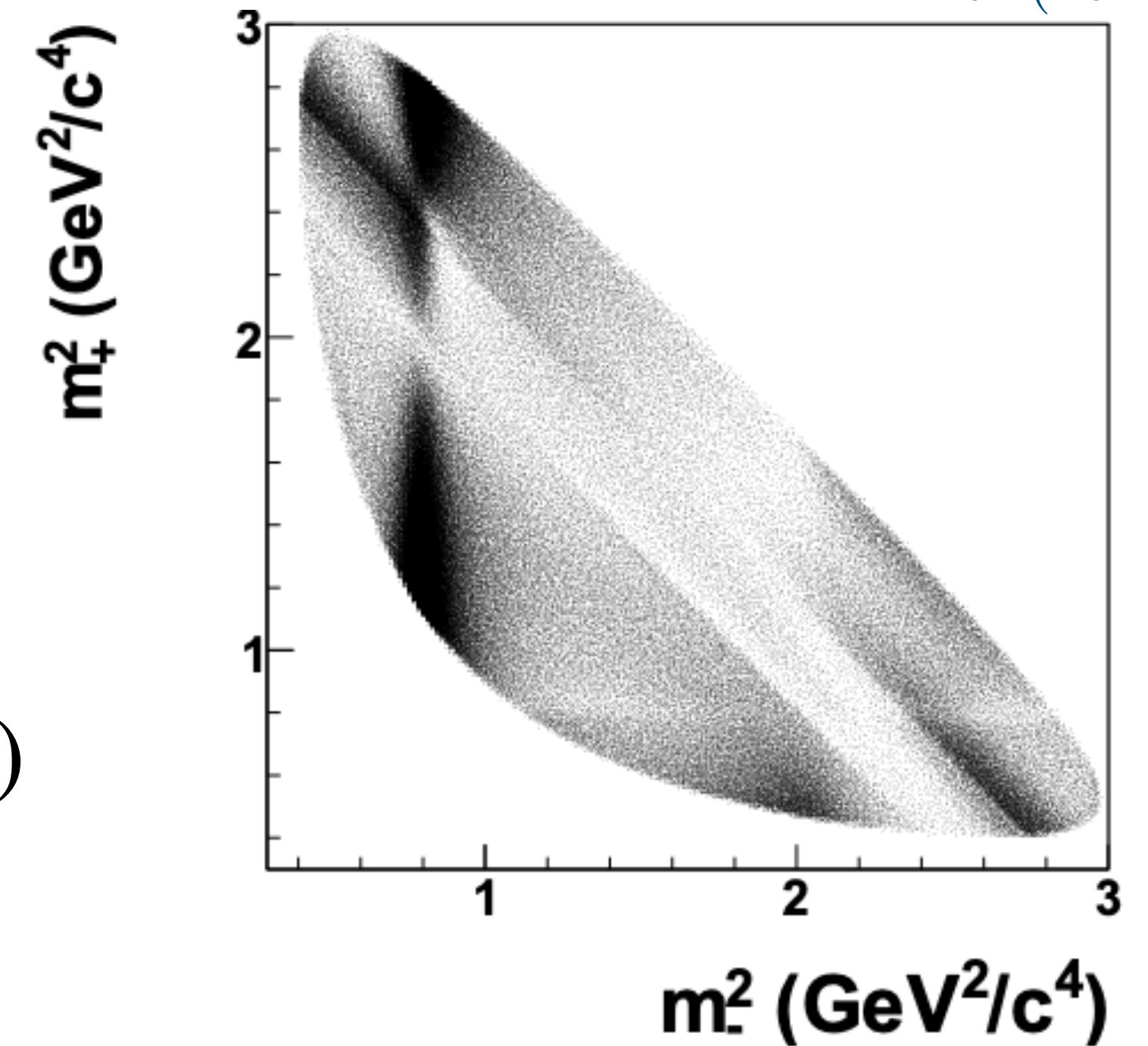
- Fitting CP-conserving obs.

$$x_{CP}^f = 2 |M_{12}| / \Gamma \cos(\phi_{K_S^0 \pi \pi}^M) \quad y_{CP}^f = |\Gamma_{12}| / \Gamma \cos(\phi_{K_S^0 \pi \pi}^\Gamma)$$

- Fitting CPV obs.

$$\Delta x^f = - |\Gamma_{12}| / \Gamma \sin(\phi_{K_S^0 \pi \pi}^\Gamma) \quad \Delta y^f = 2 |M_{12}| / \Gamma \sin(\phi_{K_S^0 \pi \pi}^M)$$

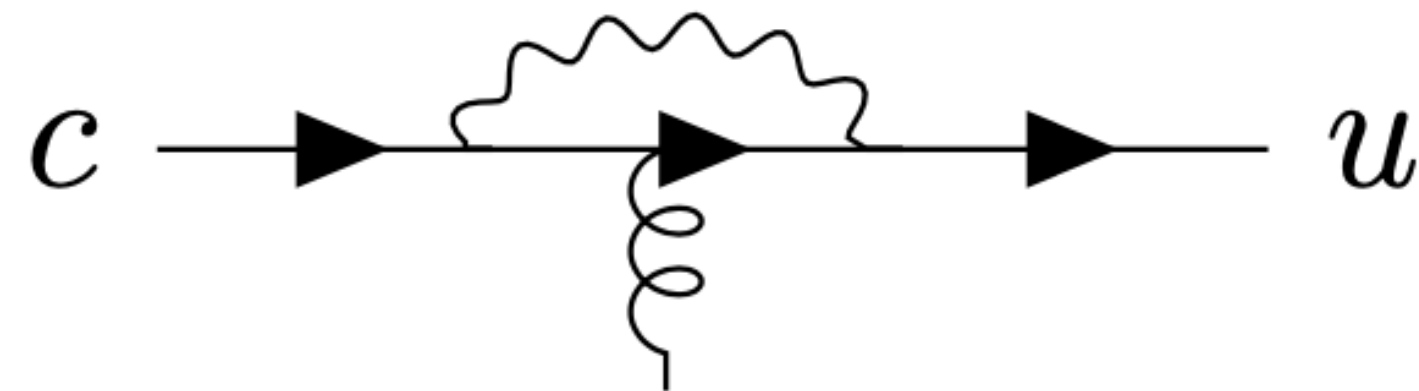
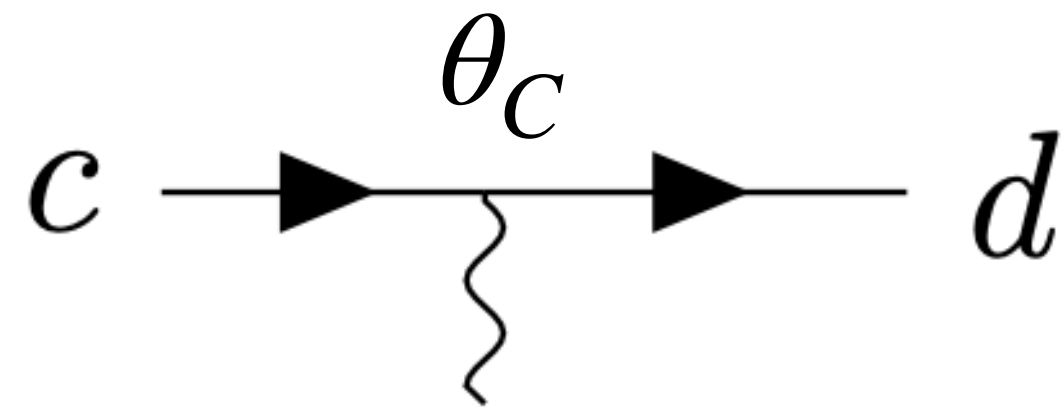
From Belle PRD(2014)



(e.g. Babar PRL(2010), LHCb PRD(2023))

SCS decays to CP eigenstates

- SCS decays to $\pi^+\pi^-(K^+K^-)$ in the SM has tree-level + penguin



- Measurements of the following time-dependent CP asymmetry

(e.g. CDF PRD(2014), LHCb PRD(2021))

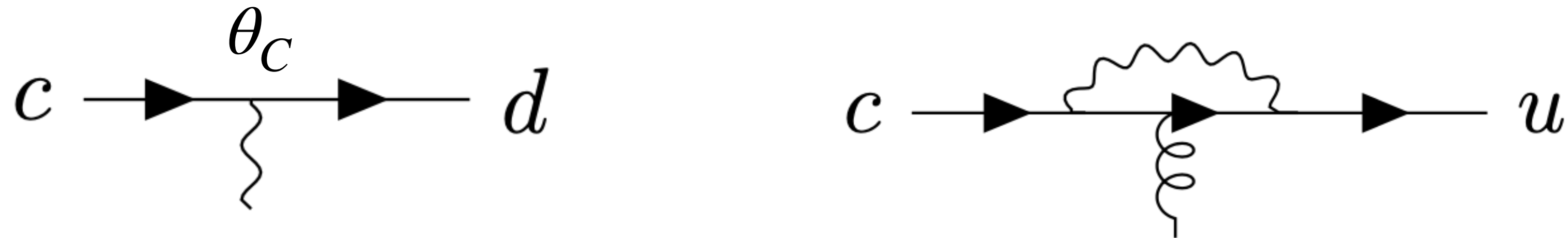
$$A_f(t) = \frac{d\Gamma(D^0 \rightarrow f) - d\Gamma(\bar{D}^0 \rightarrow f)}{d\Gamma(D^0 \rightarrow f) + d\Gamma(\bar{D}^0 \rightarrow f)}(t) = a_f + \Delta Y_f t / \tau$$

- Fitting the following CPV observables

$$a_f \quad \Delta Y_f = (-2 |M_{12}| / \Gamma \sin \phi_f^M + a_f |\Gamma_{12}| / \Gamma)$$

SCS decays to CP eigenstates

- SCS decays to $\pi^+\pi^-(K^+K^-)$ in the SM has tree-level + penguin



- Measurements of the following time-integrated CP asymmetry
(e.g. CDF PRD(2012), Babar PRL(2008), Belle Phys. Lett. B (2008), LHCb PRL(2023))

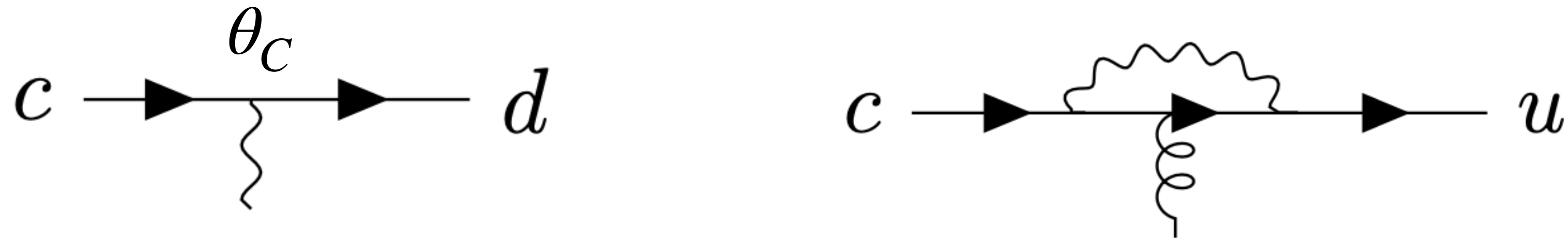
$$A_f = \frac{\Gamma(D^0 \rightarrow f) - \Gamma(\bar{D}^0 \rightarrow f)}{\Gamma(D^0 \rightarrow f) + \Gamma(\bar{D}^0 \rightarrow f)} = a_f + \Delta Y_f \langle t/\tau \rangle$$

- Fitting the following CPV observables

$$A_f \quad \Delta A^{\text{CP}} = A_{KK} - A_{\pi\pi}$$

SCS decays to CP eigenstates

- SCS decays to $\pi^+\pi^-(K^+K^-)$ in the SM has tree-level + penguin



- Measurements of the following time-dependent CP asymmetry

(e.g. E791 PRL(1999), CLEO PRD(2002), Babar PRD(2013), BESIII Phys. Lett. B (2015), Belle PRD(2020), LHCb PRD(2022))

$$R^f(t) = \frac{\Gamma(D^0 \rightarrow f) + \Gamma(\bar{D}^0 \rightarrow f)}{\Gamma(D^0 \rightarrow K^-\pi^+) + \Gamma(\bar{D}^0 \rightarrow K^+\pi^-)} \propto 1 - \tilde{y}_{\text{CP}}^f t/\tau$$

- Fitting the following CP conserving observable

$$\tilde{y}_{\text{CP}}^f = |\Gamma_{12}|/\Gamma \cos \phi_f^\Gamma + r_D^{K\pi}(\dots)$$

Theoretical framework

- Two CPV phases for each of the final states: how can we fit them?



Approximate
universality



Kagan, Silvestrini PRD(2020)

$$\phi_2^M, \phi_2^{\Gamma}$$

Approximate universality

- In the SM, the contributions to H_{12} reads

$$\Gamma_{12}^{\text{SM}} = \sum_{i,j=d,s} \lambda_{uc}^i \lambda_{uc}^j \Gamma_{ij} \qquad \mathbf{M}_{12}^{\text{SM}} = \sum_{i,j=d,s,b} \lambda_{uc}^i \lambda_{uc}^j \mathbf{M}_{ij}$$

- Employing CKM unitarity + U-spin decomposition, we get

$$\Gamma_{12}^{\text{SM}} = \frac{(\lambda_{uc}^s - \lambda_{uc}^d)^2}{4} \Gamma_2 \times \left[1 + \mathcal{O}(10^{-3}) \left(\frac{0.3}{\varepsilon} \right) + \mathcal{O}(10^{-7}) \left(\frac{0.3}{\varepsilon} \right)^2 \right]$$

CPV in Approximate universality

- Two universal CPV phases can be defined w.r.t. the dominant terms

$$\phi_2^{\text{M}} = \arg \left[\frac{\text{M}_{12}}{\text{M}_2(\lambda_{uc}^s - \lambda_{uc}^d)^2/4} \right], \quad \phi_2^{\Gamma} = \arg \left[\frac{\Gamma_{12}}{\Gamma_2(\lambda_{uc}^s - \lambda_{uc}^d)^2/4} \right]$$

- Dispersive mixing M_{12} receives contributions from both
SM + possible short distance New Physics

Is it a good approximation?

- They provide very good approximations for:

- ✱ CF/DCS decays to $K^-\pi^+$: $\phi_{K\pi}^{M,\Gamma} - \phi_2^{M,\Gamma} = O(\lambda_{uc}^b/\lambda_{uc}^s)^2 \longrightarrow O(10^{-6})$

- ✱ Three-body mode $K_S^0\pi\pi$: $\phi_{K_S^0\pi\pi}^{M,\Gamma} - \phi_2^{M,\Gamma} - 2\epsilon_I - 2|\lambda_{uc}^b/\lambda_{uc}^s|\sin(\gamma) = -2\text{Im}[r_0] \longrightarrow O(10^{-4})$

- ✱ SCS decays to $\pi^+\pi^-(K^+K^-)$: In the U-spin symmetric limit $(\phi_{KK}^{M,\Gamma} - \phi_2^{M,\Gamma}) = -(\phi_{\pi\pi}^{M,\Gamma} - \phi_2^{M,\Gamma}) \approx \phi_2^{M,\Gamma} O(\epsilon)$.
Then, we get $(\phi_{KK}^{M,\Gamma} + \phi_{\pi\pi}^{M,\Gamma})/2 = \phi_2^{M,\Gamma}(1 + O(\epsilon^2))$

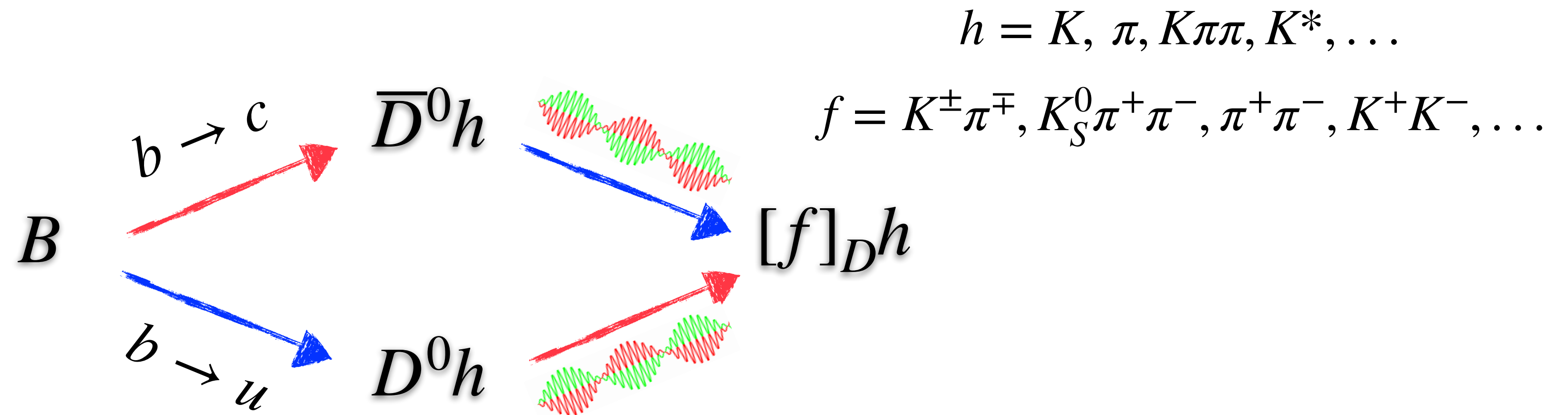
- In this framework, the CPV phases are estimated to be

$$\phi_2^M \sim \phi_2^\Gamma \sim \phi^{\text{U-spin}} = 0.13^\circ \qquad |\phi_2^\Gamma| < 0.3^\circ$$

Beauty observables?

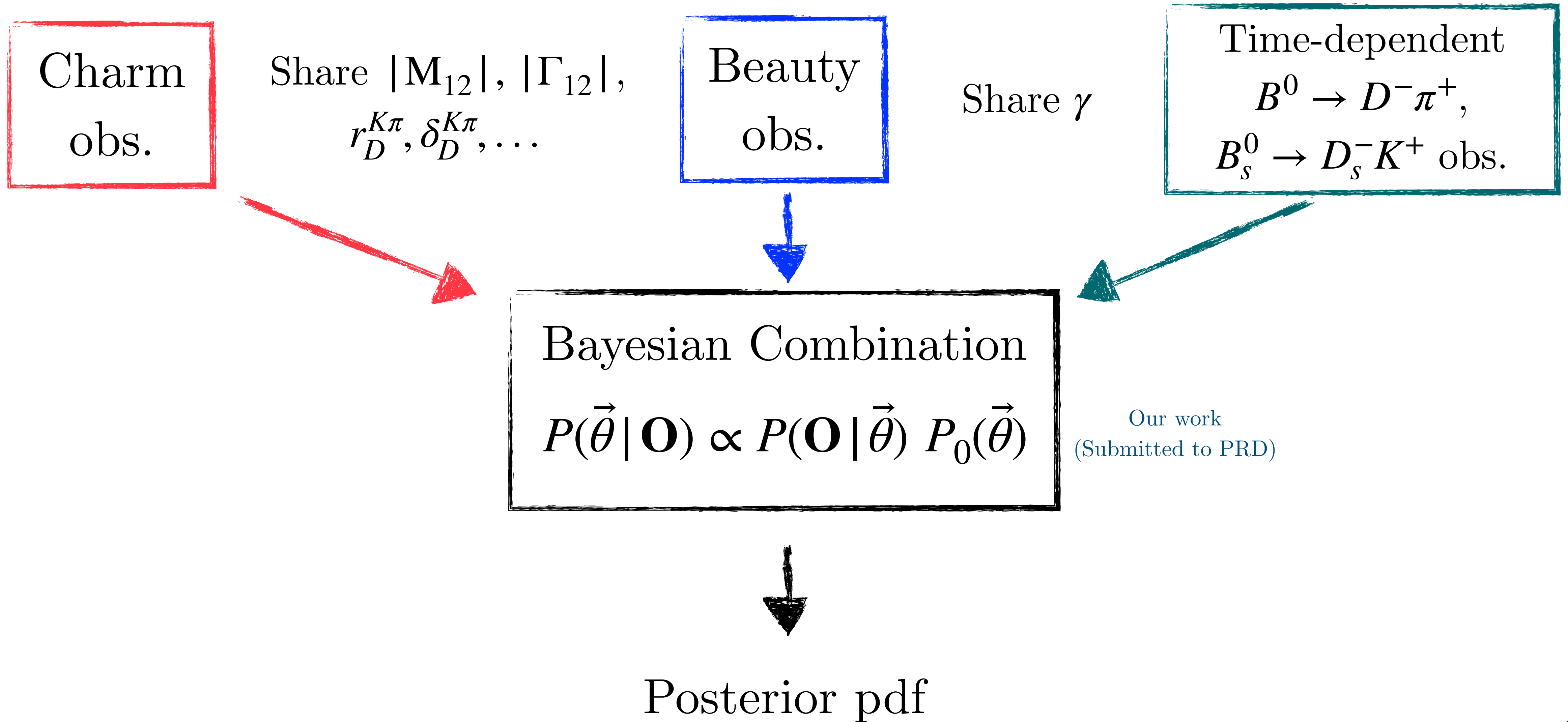
- Determining $\phi_2^{\text{M},\Gamma}$ requires to know $r_D^{K\pi(\pi^0)}$, $\delta_D^{K\pi(\pi^0)}$, ...

- Additional information are provided by [LHCb JHEP\(2021\)](#)

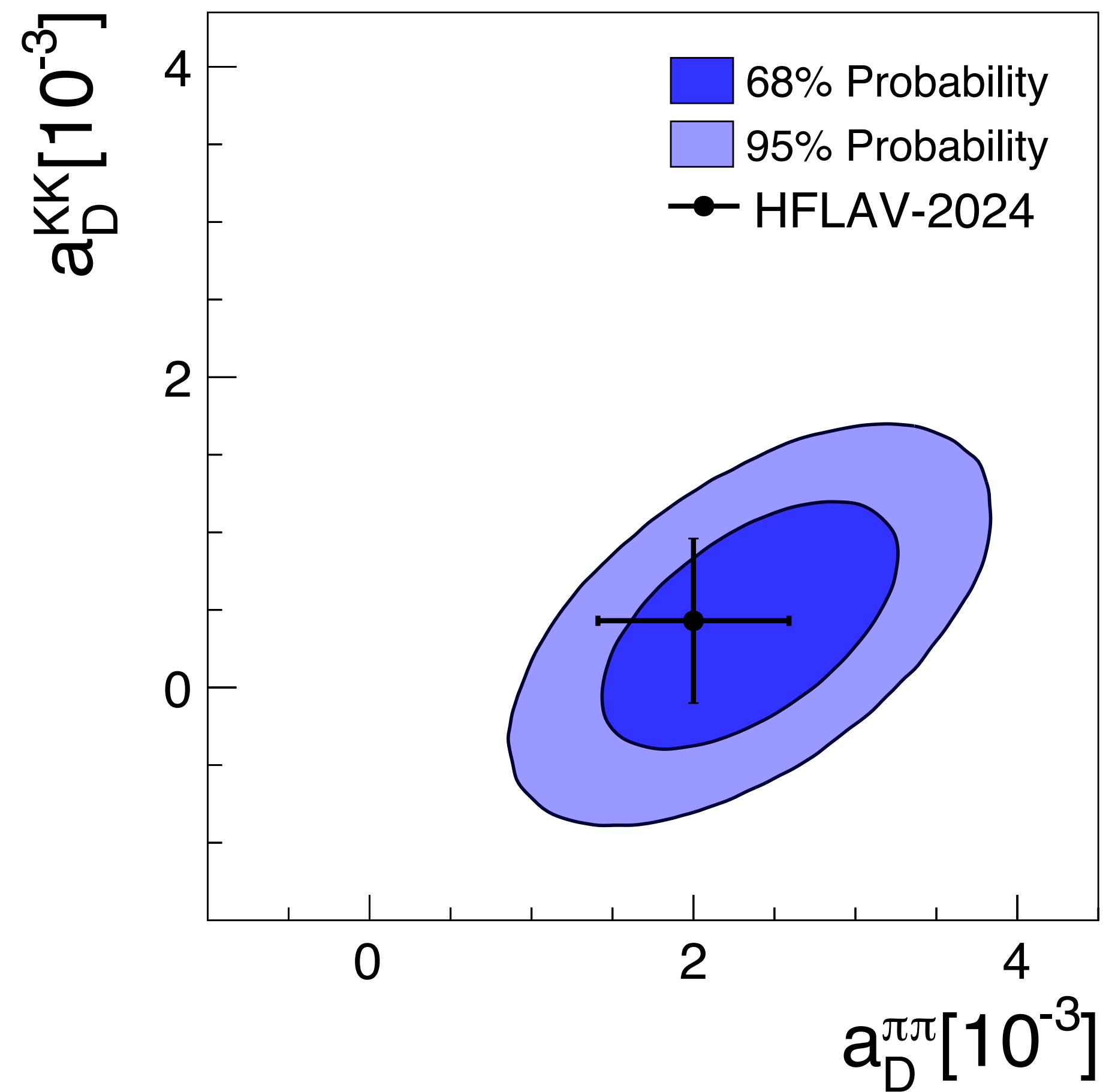
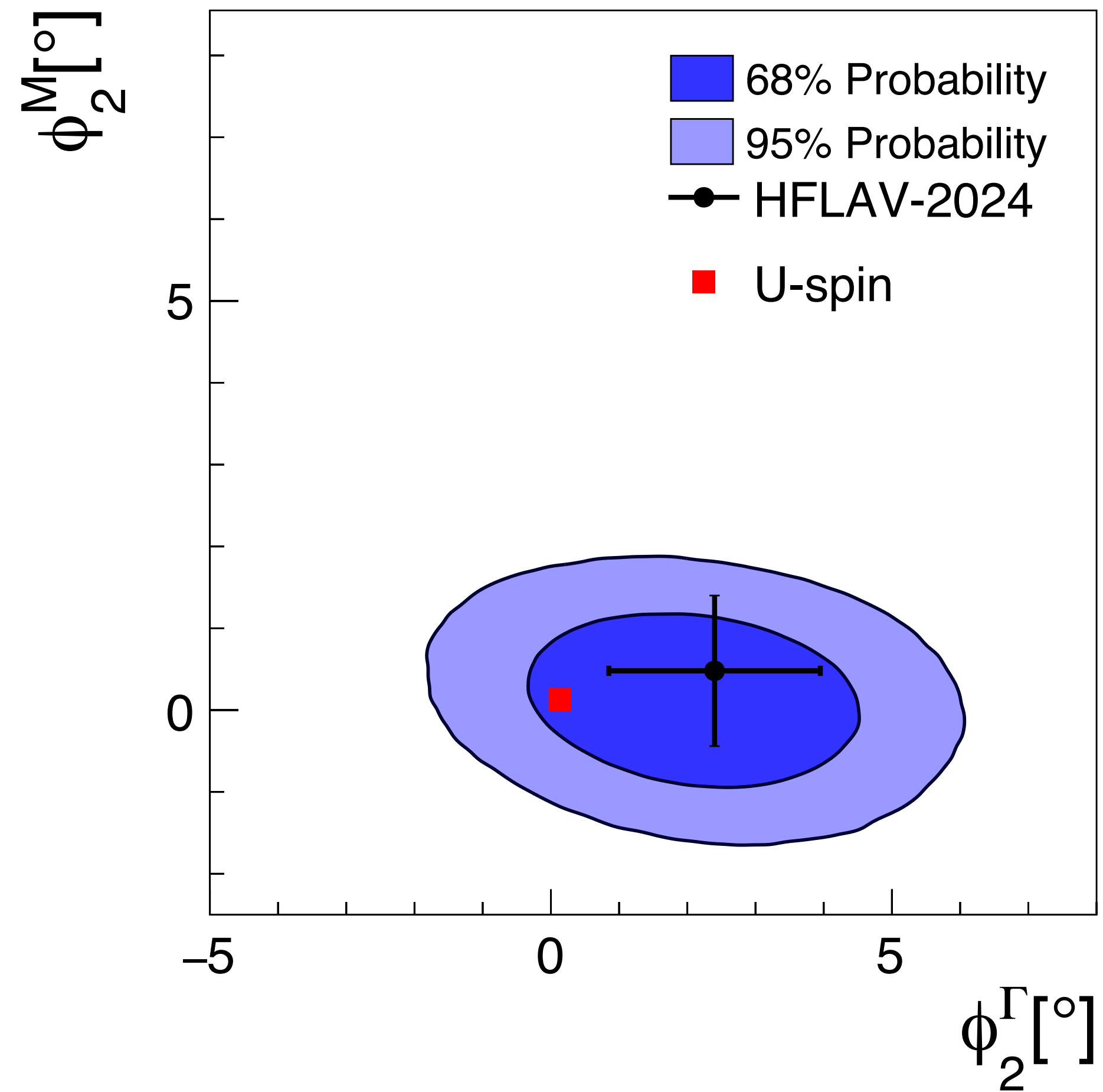


- We can determine the CKM angle $\gamma = \arg[-V_{ud}V_{ub}^*V_{cb}V_{cd}^*]$

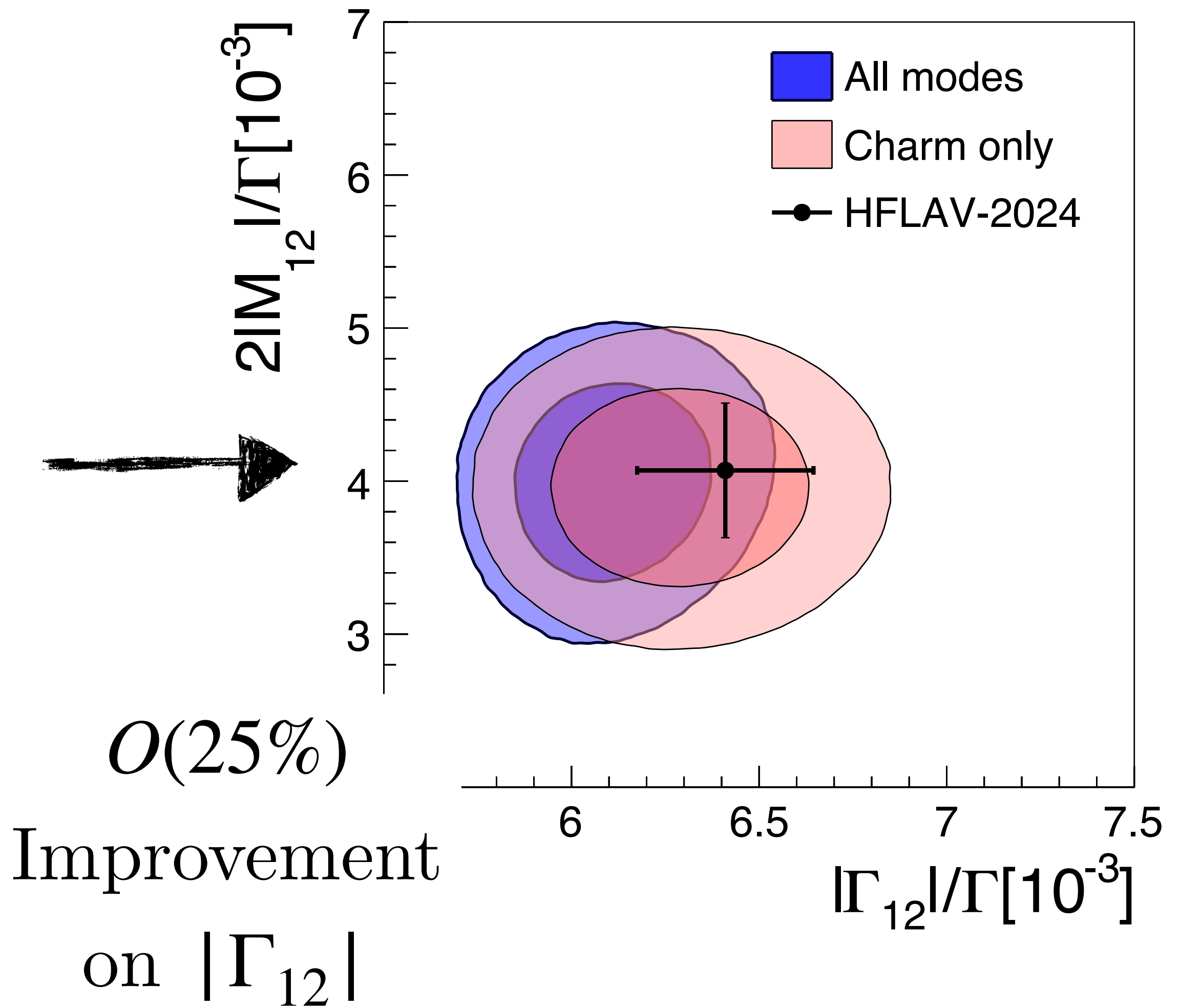
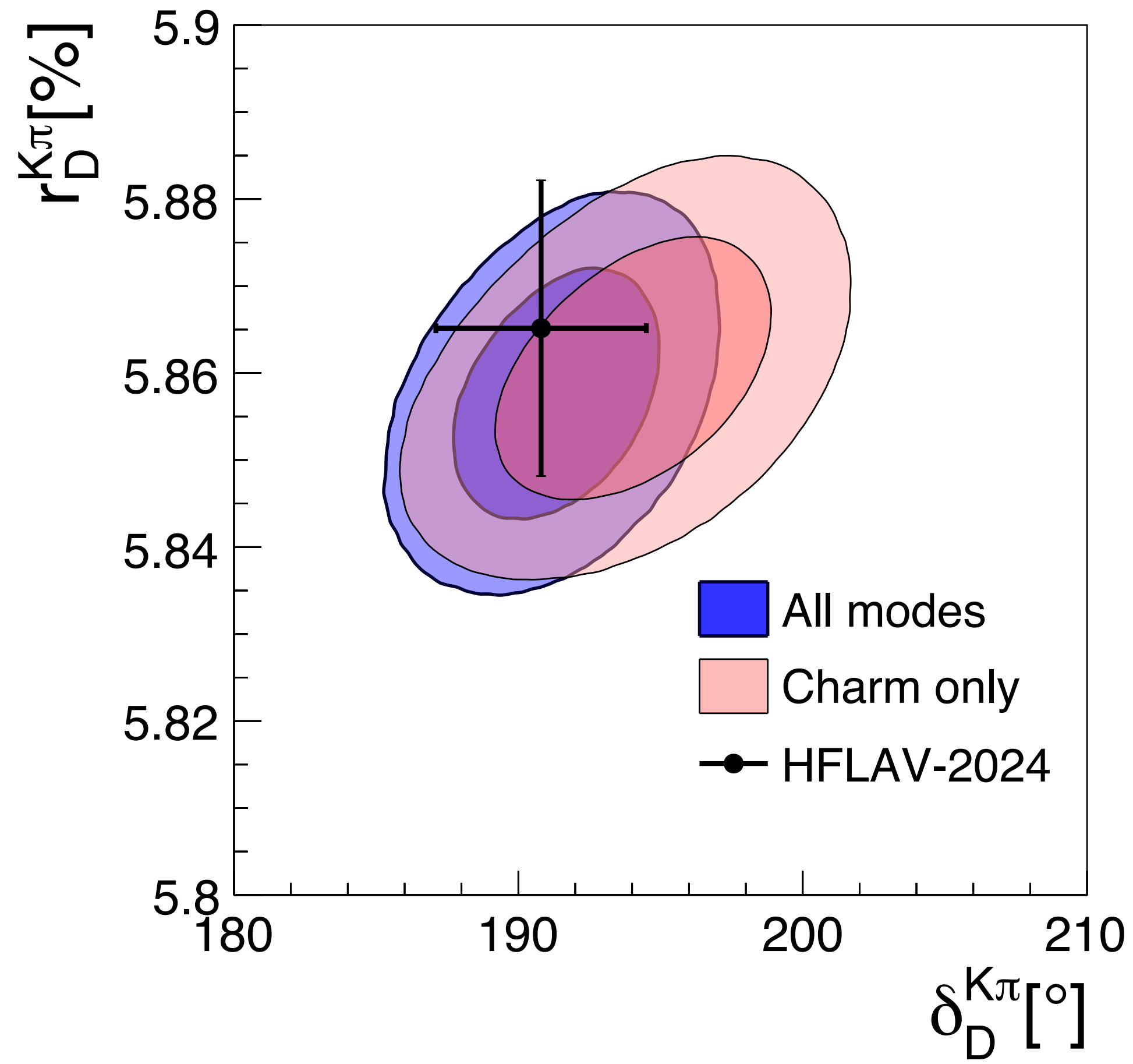
Combination



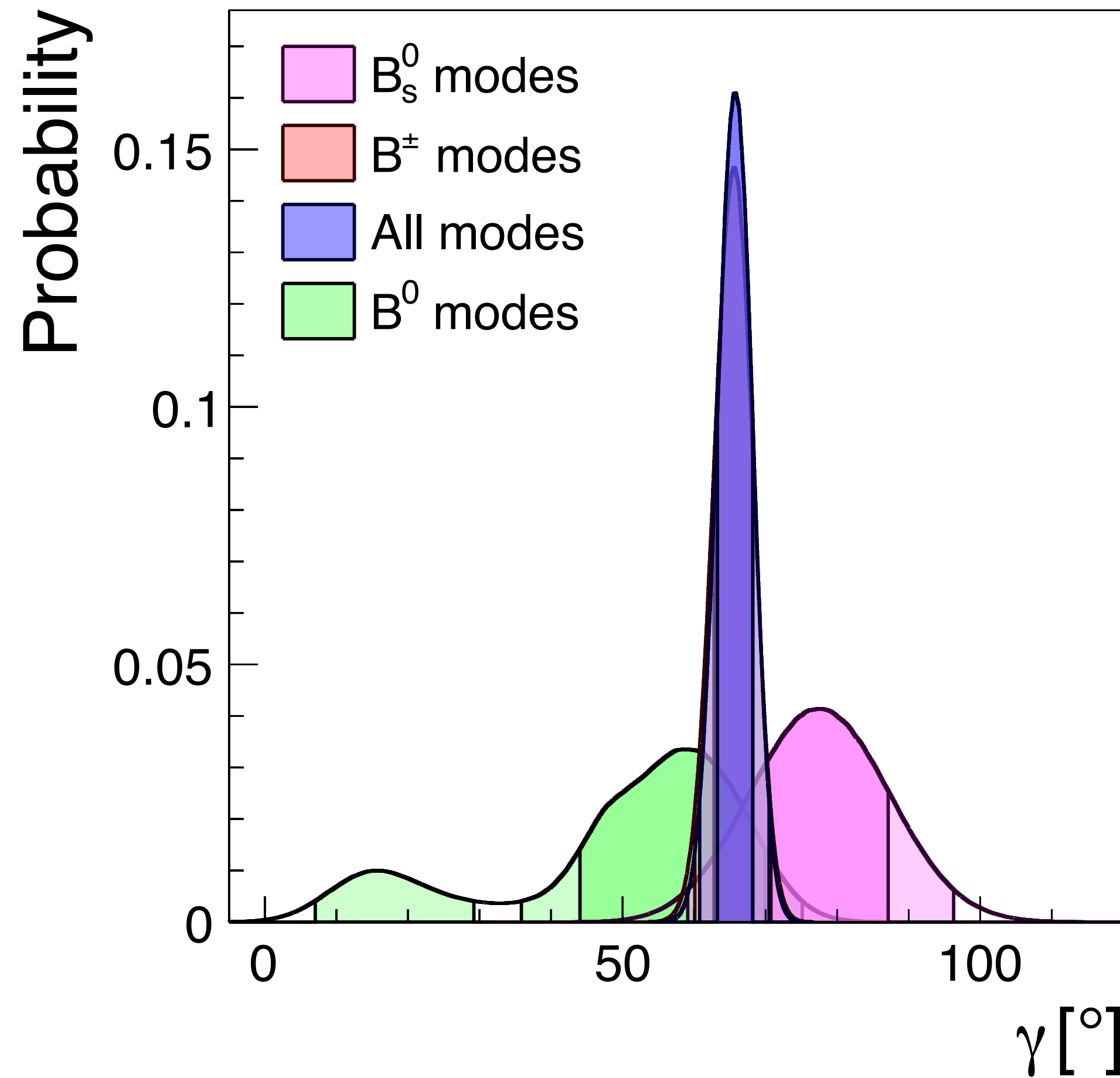
CPV parameters



Impact of the combination



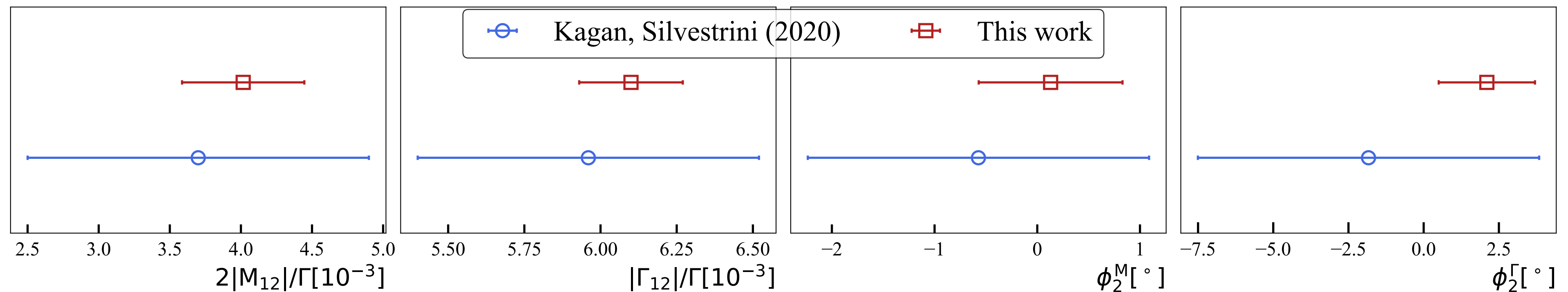
CKM angle γ



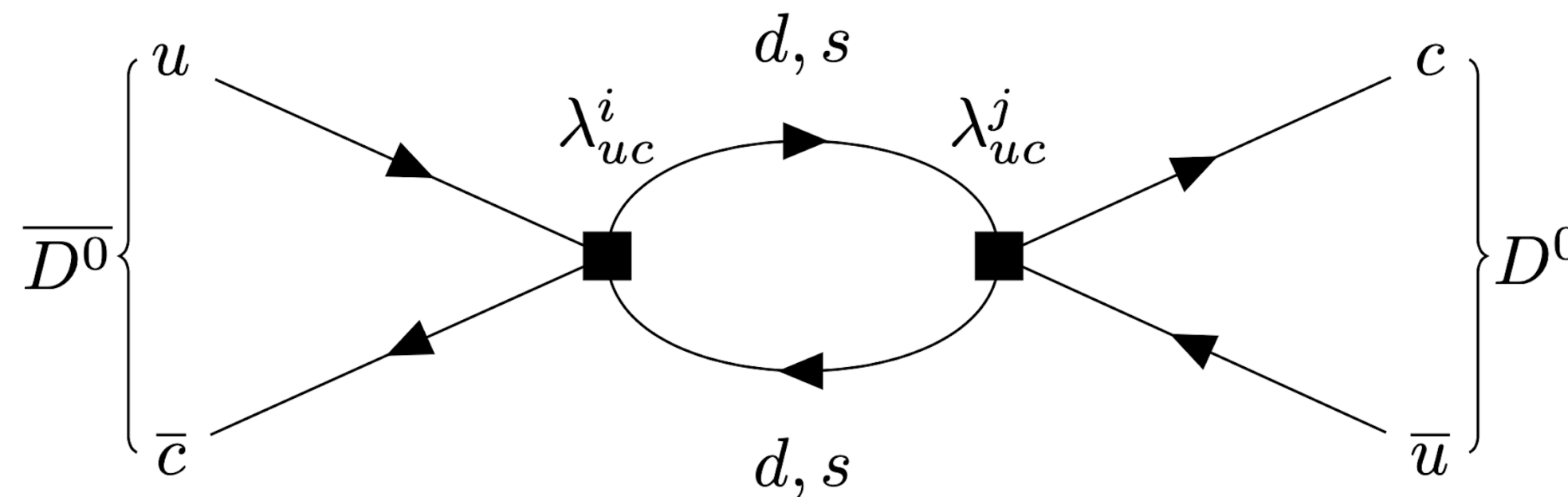
$$\gamma = 65.8(2.5)^\circ$$

Next steps?

- Experimental uncertainties have been reduced by more than a factor of 2(3) on dispersive (absorptive) mixing parameters



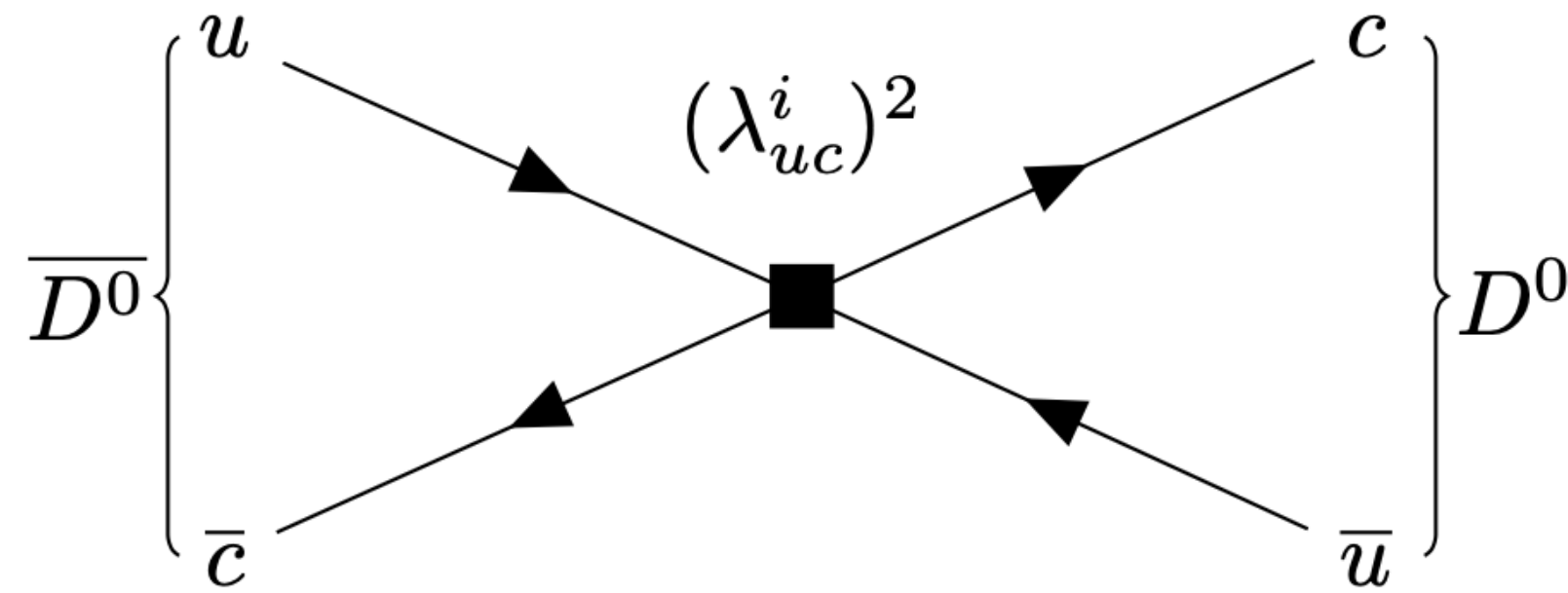
- Can we get SM estimates for CPV in charm mixing?



Backup

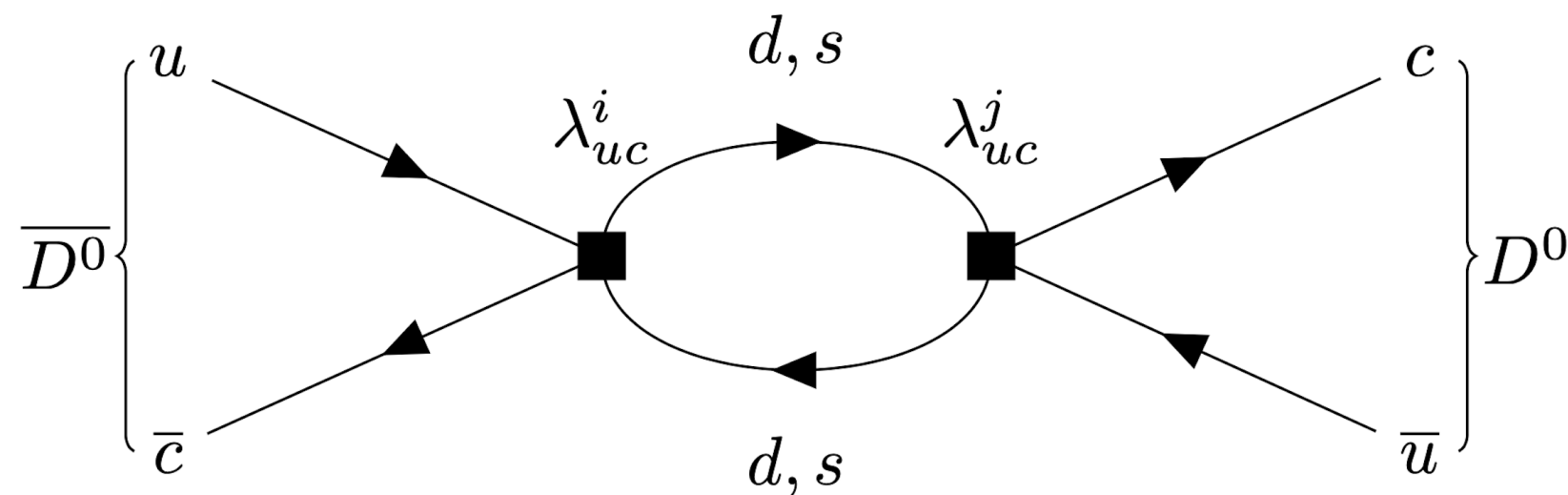
Short vs long distance

- The short distance contribution is given by



$$\text{SM: } \propto (\lambda_{uc}^b m_b)^2 \approx (\theta_C^5 m_b)^2$$

- The long distance contribution is given by



$$\text{SM: } \propto (\lambda_{uc}^s m_s)^2 \approx (\theta_c m_s)^2 \approx 10^2 \times \text{SD}$$

Amplitude decomposition

- The absorptive part of the mixing hamiltonian reads

$$\Gamma_{12}^{\text{SM}} = \frac{(\lambda_{uc}^s - \lambda_{uc}^d)^2}{4} \Gamma_2 + \frac{(\lambda_{uc}^s - \lambda_{uc}^d) \lambda_{uc}^b}{2} \Gamma_1 + \frac{(\lambda_{uc}^b)^2}{4} \Gamma_0$$

- The amplitudes and CKM matrix elements satisfy

$$\Gamma_0 = (\bar{s}s + \bar{d}d)^2 = \mathcal{O}(1) \quad \Gamma_1 = (\bar{s}s - \bar{d}d)(\bar{s}s + \bar{d}d) = \mathcal{O}(\varepsilon) \quad \Gamma_2 = (\bar{s}s - \bar{d}d)^2 = \mathcal{O}(\varepsilon^2)$$

$$\lambda_{uc}^s - \lambda_{uc}^d \approx 0.44 - i1.2 \times 10^{-4} \quad \lambda_{uc}^b \approx (5.7 + i12) \times 10^{-5}$$

- We get the expansion

$$\Gamma_{12}^{\text{SM}} = \frac{(\lambda_{uc}^s - \lambda_{uc}^d)^2}{4} \Gamma_2 \times \left[1 + (0.86 + i1.8) \times 10^{-3} \left(\frac{0.3}{\varepsilon} \right) + (-6.4 + i7.8) \times 10^{-7} \left(\frac{0.3}{\varepsilon} \right)^2 \right]$$

Estimates

- Estimates of $\phi_2^{\text{M},\Gamma}$ can be obtained by using the SM definitions

$$\left. \phi_2^\Gamma \right|_{SM} = \arg \left[1 + \frac{2\lambda_{uc}^b}{\lambda_{uc}^s - \lambda_{uc}^d} \frac{\Gamma_1}{\Gamma_2} \right] = \arg \left[1 - \frac{V_{ub}^* V_{ud}}{V_{cb}^* V_{cd}} \times \left(\frac{2}{1 - \frac{V_{us}^* V_{cs}}{V_{ud}^* V_{cd}}} \right) \varepsilon^{-1} \right]$$
$$\simeq \left| \frac{\lambda_{uc}^b}{\lambda_{uc}^d} \right| \sin(\gamma) \varepsilon^{-1} \approx (2.2 \times 10^{-3}) \times \left[\frac{0.3}{\varepsilon} \right]$$

Upper bound

$$|\phi_2^\Gamma| = \left| \frac{\lambda_{uc}^b}{\lambda_{uc}^d} \right| \sin(\gamma) \frac{\Gamma_1}{\Gamma_2}$$

- Now, we have that $|\Gamma_2| = y_{12}\Gamma/(\lambda_{uc}^d)^2$ and

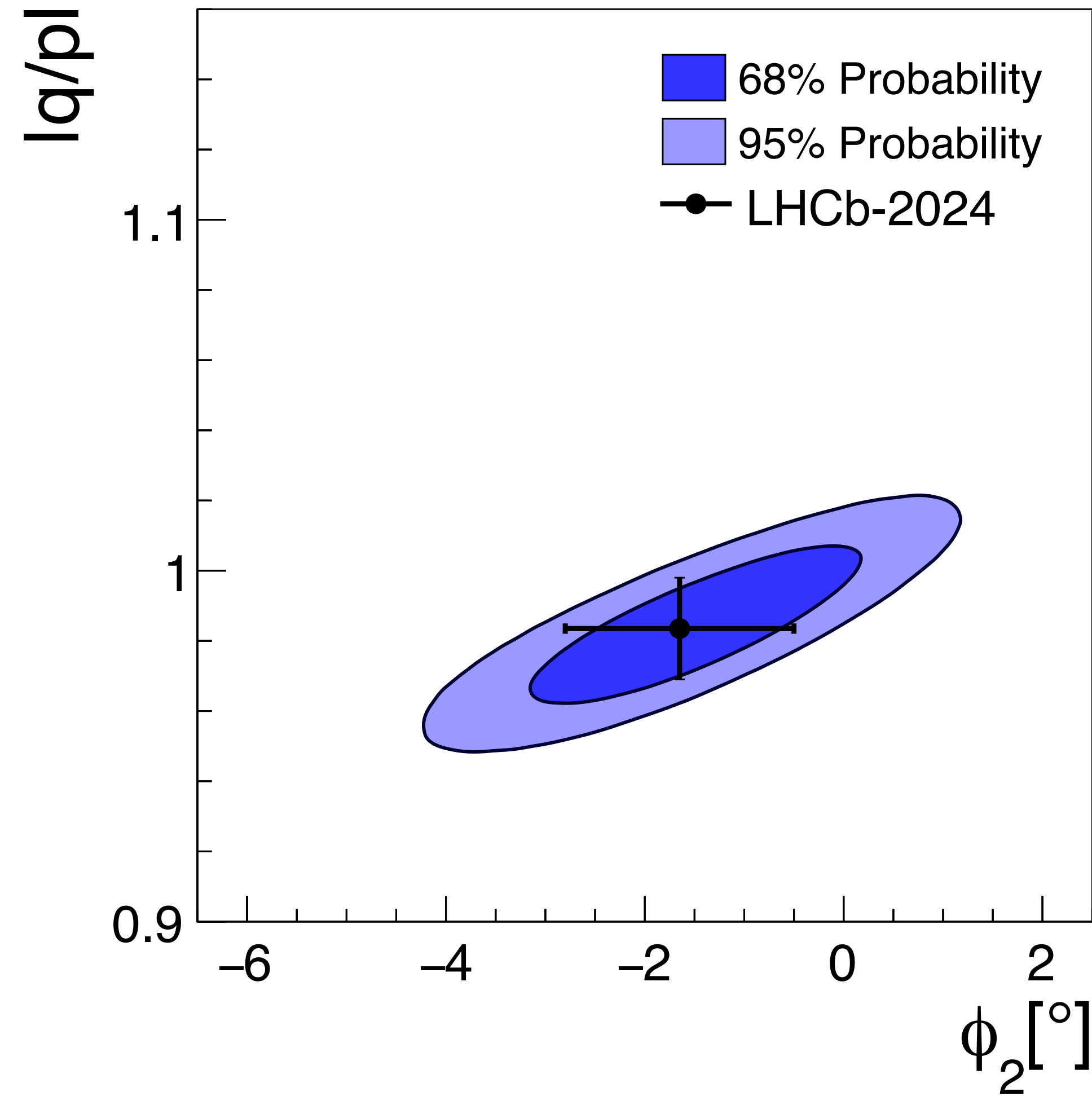
$$|\phi_2^\Gamma| = \left| \frac{\lambda_{uc}^b \lambda_{uc}^d}{y_{12}} \right| \sin(\gamma) \frac{|\Gamma_{sd}|}{\Gamma} \frac{|\Gamma_{ss} - \Gamma_{dd}|}{|\Gamma_{sd}|} \rightarrow O(\varepsilon)$$

$$< 1 + O(\varepsilon) \quad \swarrow_{SU(3)_f}$$

Also see [M. Bobrowski, A. Lenz, J. Riedl, J. Rohrwild\(2010\)](#)

$$|\phi_2^\Gamma| < 5 \times 10^{-3} \varepsilon (1 + O(\varepsilon)) < 0.3^\circ$$

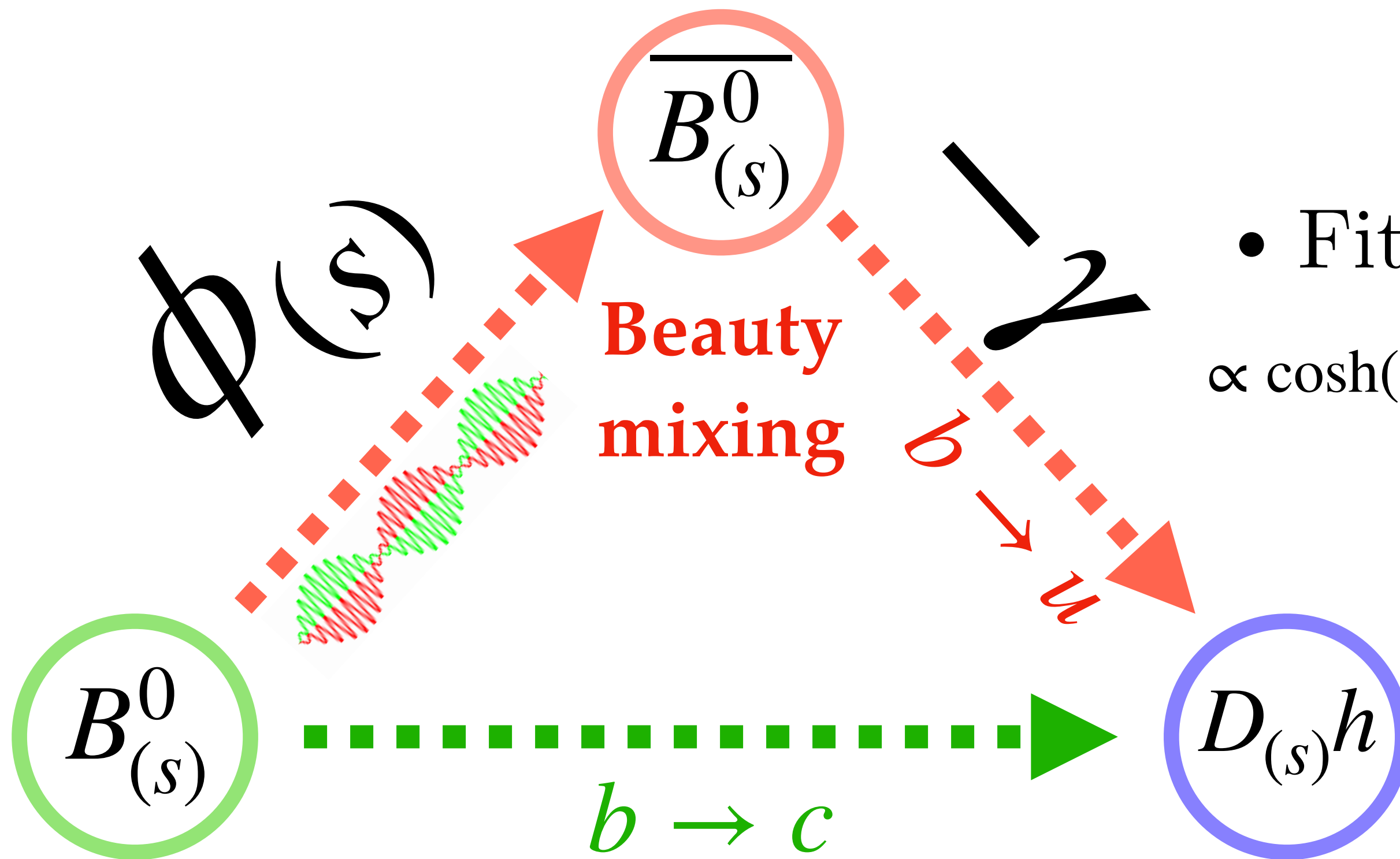
Alternative parametrisation



Neutral B meson obs.

- Exploiting the **CPV** phase of the interference between $B_{(s)}^0$ mixing and decay to charmed mesons $D_{(s)}^\mp h^\pm$

Mixing $\phi = -2\beta$
phases $\phi_s = 2\beta_s$



- Fitting the time-dependent decay rates
 $\propto \cosh(\Delta\Gamma_{(s)}t/2) - G_f \sinh(\Delta\Gamma_{(s)}t/2) + C_f \cos(\Delta m_{(s)}t) - S_f \sin(\Delta m_{(s)}t)$

Observables!!

$$C_f \quad G_f \propto \cos(\delta_{B_{(s)}^0}^f + (\phi_{(s)} - \gamma))$$

$$S_f \propto \sin(\delta_{B_{(s)}^0}^f + (\phi_{(s)} - \gamma))$$