

INFN Bari Theory Group Xmas Workshop Bari, 17 December 2024

High-frequency GWs shining in photons in Galactic magnetic fields

Based on: AL, F. Calore, P. Carenza, A. Mirizzi, 2406.17853 [hep-ph] Phys.Rev.D 110 (2024) 8, 083042

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GWs from known sources

High-frequency GW High-frequency GWs

Exotic sources in the early universe could produce a GW background at high frequencies $f \gtrsim MHz$

Electromagnetism in a curved space-time couples the photon and the gravitational fields

$$
\mathcal{L}_{\rm em} = -\frac{1}{4}g^{\mu\alpha}g^{\nu\beta}F_{\mu\nu}F_{\alpha\rho} + \int d^4x' A_\mu(x)\Pi^{\mu\nu}(x,x')A_\nu(x')
$$

Photon polarization tensor in the Euler-Heisenberg limit ($B \ll B_{cr} \sim 10^{13}$ G)

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Deriving the EoM, in the limit $\omega \simeq k$ *[Raffelt & Stodolsky, Phys.Rev.D 37 (1988) 1237]*

$$
\begin{pmatrix} i\frac{d}{dz} - \omega \end{pmatrix} \begin{pmatrix} h_+ \\ h_\times \\ A_x \\ A_y \end{pmatrix} = H \begin{pmatrix} h_+ \\ h_\times \\ A_y \end{pmatrix} \qquad H = \begin{pmatrix} 0 & \mathcal{H}_{g\gamma} \\ \mathcal{H}_{g\gamma} & \mathcal{H}_{\gamma\gamma} \end{pmatrix}
$$

• Photon dispersion relation:

$$
\mathcal{H}_{\gamma\gamma} = \begin{pmatrix} \Delta_x c_\phi^2 + \Delta_y s_\phi^2 & [\Delta_y - \Delta_x] c_\phi s_\phi \\ [\Delta_y - \Delta_x] c_\phi s_\phi & \Delta_y c_\phi^2 + \Delta_x s_\phi^2 \end{pmatrix}
$$

$$
\Delta_{\lambda}=\boxed{\Delta_{\rm pl}+\Delta_{\rm QED}^{\lambda}+\Delta_{\rm CMB}},
$$

$$
\begin{split} \Delta_{\rm pl} & = -\,\frac{\omega_{\rm pl}^2}{2\,\omega} \\ & \simeq -\,1.1\times 10^{-3}\,\left(\frac{\omega}{1\,\,{\rm MeV}}\right)^{-1} \left(\frac{n_e}{10^{-2}\,\,{\rm cm}^{-3}}\right)\,{\rm kpc}^{-1}\,, \end{split}
$$

$$
\begin{split} \Delta^{\lambda}_{\rm QED}=&\kappa_{\lambda}\,\frac{4\,\alpha^2\,B_T^2\,\omega}{45\,m_e^4} \\ &\simeq 4.5\times 10^{-12}\,\kappa_{\lambda}\,\left(\frac{\omega}{1\ {\rm MeV}}\right)\left(\frac{B_T}{1\mu{\rm G}}\right)^2\,{\rm kpc}^{-1}\,, \end{split}
$$

$$
\Delta_{\rm CMB} = \frac{44\pi^2 \, \alpha^2 \, T_{\rm CMB}^4 \, \omega}{2025 \, m_e^4} \nonumber \\ \simeq 8.7 \times 10^{-11} \left(\frac{\omega}{1 \; \rm MeV} \right) \; \rm kpc^{-1} \,,
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$$

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$$

$$
\Delta_{g\gamma} = \frac{B_{\rm T}}{\sqrt{2}\,M_{\rm P}} \simeq 8.8\times 10^{-10} \left(\frac{B_{\rm T}}{1\,\mu{\rm G}}\right)\,{\rm kpc}^{-1}
$$

• Mixing terms:

$$
\mathcal{H}_{g\gamma} = \begin{pmatrix} \Delta_{g\gamma} s_{\phi} & \Delta_{g\gamma} c_{\phi} \\ \Delta_{g\gamma} c_{\phi} & -\Delta_{g\gamma} s_{\phi} \end{pmatrix}
$$

Large-Scale Galactic Magnetic Field

Jansson-Farrar model for the large-scale Galactic magnetic field [Jansson & Farrar, Astrophys.J. 757 (2012) 14].

- $|B| \sim 1 \mu G$
- Correlated over scales $l_{\text{corr}} \sim 1 \text{ kpc}$
- Disk field + a large halo field
- Reproduces Galactic synchrotron emission

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Model for the electron number density from *[Cordes & Lazio,* astro-ph/0207156] .

- $n_e \sim 10^{-3} 10^{-2} \, cm^{-3}$
- Correlated over scales $l_{corr} \sim 1$ kpc
- Describes large-scale fluctuations in electron density

Conversion Probability

- The oscillation probability can be obtained by solving numerically the EoM
- Assuming a constant and homogenous field along the line of sight

$$
\langle B_T \rangle^2 = \frac{1}{L^2} \left(\left| \int_0^L B_x(z, l, b) dz \right|^2 + \left| \int_0^L B_y(z, l, b) dz \right|^2 \right)
$$

$$
P_{g\gamma} = \frac{4\Delta_{g\gamma}^2}{\Delta_{\rm osc}^2} \sin^2 \left(\frac{\Delta_{\rm osc} z}{2} \right)
$$

 $P_{g\gamma}$ [×10⁻¹⁶]

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• The two approaches match when $l_{\rm osc} \gg l_{\rm corr}$

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$$

- The two approaches match when $l_{osc} \gg l_{corr}$
- Dramatic loss of precision in the numerical approach at $l_{\rm osc} \ll l_{\rm corr}$

Observation of the CPB

Measurements of the Cosmic Photon Background from [R. Hill et al., Appl. Spectrscop. 72 (2018) 663]

- CRB-CMB: blackbody spectrum at $T_{\text{CMB}} \simeq 2.73 \text{ K}$ [D. J. Fixsen, Astrophys. J. 707 (2009) 916]
- CIB-COB-CUB: EBL emission [Dominguez et al., Mon. Not. Roy. Astron. Soc. 410 (2011) 2556]
- **CXB:** AGN accretion disk emission [Comastri et al., Astron. Astrophys. 296 (1995) 1]
- CGB: Blazar emission [M. Fornasa et al., Phys. Rept. 598 (2015)1]

Constraints on the GW background

To avoid an excess in the CPB we must require:

$$
\left.\frac{1}{4\pi}\frac{dF_{g\to\gamma}}{df}\right|_{f=f_i} \lesssim \left|\frac{d^2F^{\rm exp}_{\gamma,i}}{dfd\Omega} - \frac{d^2F^{\rm th}_{\gamma,i}}{dfd\Omega}\right|
$$

- Constraints comparable to axion experiments sensitivities [Ejlli et al., Eur. Phys J. C79 (2019)] [Domcke & Garcia-Cely, PRL 126 (2019)]
- Improvement of 1-2 orders of magnitude with respect to previous Astro bounds

[Ito et al., PTEP 2024 (2024)] [Dandoy et al., arXiv:2402.14092] [Ellis & McDonald, arXiv: 2406.18634]

Summary and conclusions

- High-frequency GW backgrounds are a smoking gun for new physics
- Graviton-photon conversion could probe the ultra-high frequency range
- Graviton-photon mixing in the Galactic magnetic field provides strong constraints on high-frequency GW backgrounds.
- Still far away from cosmological constraints. What's next?

Thank you for your attention

Turbulent magnetic field

The Galactic magnetic field shows evidences for a turbulent component with $l_{corr} \sim 20 - 200$ pc [Iacobelli eta al., Astron. Astrophys. 556 (2013) A72]

• For $l_{osc} \gg l_{corr}$ cell approximation with hard edges holds

 $P_{\text{turb}} \simeq N \Delta P = \langle B_{\text{T}} \rangle^2 L l_{\text{corr}}$

- Suppression of a factor $\sim N$ with respect to the regular component
- For $l_{\rm osc} \ll l_{\rm corr}$ oscillations sensible to fine structure of the field
	- Strong dependency on the field power spectrum