



QUANTUM EFFECTS IN GRAVITY FOM A DELOCALISED QUANTUM SOURCE

- **SNSF Ambizione Fellow, ETH Zürich**







Flaminia Giacomini

The Low-Energy Frontier of Particle Physics Laboratori Nazionali di Frascati 10-12 February 2025

QUANTUM THEORY



Image credits: Perimeter Institute

Flaminia Giacomini - ETH Zurich

GENERAL RELATIVITY







QUANTUM THEORY



Entanglement, superposition...

Image credits: Perimeter Institute

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GENERAL RELATIVITY







QUANTUM THEORY



Entanglement, superposition...

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GENERAL RELATIVITY



Relation between gravity and matter

Matter tells spacetime how to curve; spacetime tells matter how to move.





QUANTUM THEORY



Entanglement, superposition...

Spacetime is the stage

Image credits: Perimeter Institute

Flaminia Giacomini - ETH Zurich

GENERAL RELATIVITY



Relation between gravity and matter

Matter tells spacetime how to curve; spacetime tells matter how to move.

Spacetime is the actor





QUANTUM THEORY



Entanglement, superposition...

Spacetime is the stage

All experiments are compatible with these theories.

Image credits: Perimeter Institute

Flaminia Giacomini - ETH Zurich

GENERAL RELATIVITY



Relation between gravity and matter

Matter tells spacetime how to curve; spacetime tells matter how to move.

Spacetime is the actor





WHERE SHALL WE LOOK FOR QUANTUM EFFECTS IN GRAVITY?

HIGH ENERGIES: **STRONG GRAVITATIONAL AND QUANTUM EFFECTS**



Image credits: Perimeter Institute

Flaminia Giacomini - ETH Zurich

QUANTUM GRAVITY

 $\ell_P \approx 10^{-35} m$



WHERE SHALL WE LOOK FOR QUANTUM EFFECTS IN GRAVITY?

HIGH ENERGIES: **STRONG GRAVITATIONAL AND QUANTUM EFFECTS**

LOW ENERGIES: WEAK-FIELD GRAVITY QUANTUM PARTICLES



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QUANTUM GRAVITY

 $\ell_P \approx 10^{-35} m$



NONCLASSICAL SPACETIME FROM A QUANTUM SOURCE

Article | Published: 10 March 2021 GRAVITY SOURCE: 90 mg **Measurement of gravitational coupling between** millimetre-sized masses

Tobias Westphal 🖂, Hans Hepach, Jeremias Pfaff & Markus Aspelmeyer 🖂

Nature 591, 225–228 (2021) Cite this article



Flaminia Giacomini - ETH Zurich

Schrödinger cat states of a 16-microgram mechanical SUPERPOSED MASS: $10^{-5}g$ oscillator



20 Apr 2023 Vol 380, Issue 6642 pp. 274-278 Dol: 10.1126/science.adf7553

Published: 23 December 201 QUANTUM SUPERPOSITION: 0.5 m Quantum superposition at the half-metre scale T. Kovachy, P. Asenbaum, C. Overstreet, C. A. Donnelly, S. M. Dickerson, A. Sugarbaker, J. M. Hogan & M. A. Kasevich

Nature 528, 530–533 (2015) Cite this article









"... it seems to me that we are in trouble if we believe in quantum mechanics but do not quantize gravitational theory"

R. Feynman, Chapel Hill Conference (1957)

Flaminia Giacomini - ETH Zurich

The Role of Gravitation in Physics

Report from the 1957 Chapel Hill Conference

Cécile M. DeWitt and Dean Rickles (eds.)



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"If you believe in quantum mechanics up to any level then you have to believe in gravitational quantization in order to describe this experiment."

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SIDE REMARK: WHAT PROVES THAT ELECTROMAGNETISM IS QUANTUM?

- 1905: photoelectric effect (Einstein)
- 1923: photoelectric effect does not require quantum (Millikan)
- **1923: Compton effect**
- **1960s: semiclassical theory of radiation (Jaynes)**

$$g^{(2)}(\tau) = \frac{\langle I(t)I(t+\tau)\rangle}{\langle I(t)\rangle^2}$$
$$\tau = 0 \qquad g^{(2)}(0) = \frac{\langle I^2(t)\rangle}{\langle I(t)\rangle^2} \ge 1$$

For a single photon source
$$g^{(2)}(0) = 0 \ngeq 1$$

Signature of nonclassicality





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$\sim (|x_1\rangle_A + |x_2\rangle_A) \otimes (|x_1 + d\rangle_B + |x_2 + d\rangle_B)$ **PRODUCT STATE**

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$$\hat{H}_I = \hat{V}_N = -G \frac{m_A m_B}{|\hat{x}_A - \hat{x}_B|}$$

 $\sim (|x_1\rangle_A + |x_2\rangle_A) \otimes (|x_1 + d\rangle_B + |x_2 + d\rangle_B)$









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DSTATE ~
$$\sum_{i,j=1,2} e^{i\phi_{ij}} |x_i\rangle_A |x_j + d\rangle_B$$

$$\hat{H}_I = \hat{V}_N = -G \frac{m_A m_B}{|\hat{x}_A - \hat{x}_B|}$$

 $\sim (|x_1\rangle_A + |x_2\rangle_A) \otimes (|x_1 + d\rangle_B + |x_2 + d\rangle_B)$









Many people contributed! E.g. Anastopoulos, Aspelmeyer, Barker, Belenchia, Bengyat, Bhatar, Blencowe, Bose, Brukner, Carney, Castro-Ruiz, Chen, Christodoulou, Cooper, Di Biagio, Galley, Geraci, Hackermüller, Howl, Hu, Huggett, Iyer, Kent, Kim, Krisnanda, Lami, Linneman, Liu, Mahesh, Marletto, Marshman, Martín-Martínez, Mazumdar, Milburn, Morley, Müller, Mummery, Naik, Pal, Paterek, Paternostro, Pedernales, Perche, Pitalúa-García, Plenio, Qvarfort, Rovelli, Schneider, Schut, Selby, Serafini, Sillanpää, Tam, Taylor, Toros, Ulbricht, Vedral, Wald, Yant...

ED STATE
$$\sim \sum_{i,j=1,2} e^{i\phi_{ij}} |x_i\rangle_A |x_j + d\rangle_B$$

$$\hat{H}_I = \hat{V}_N = -G \frac{m_A m_B}{|\hat{x}_A - \hat{x}_B|}$$

STATE ~
$$(|x_1\rangle_A + |x_2\rangle_A) \otimes (|x_1 + d\rangle_B + |x_2 + d\rangle_B)$$















ENTANGLEMENT RATE

$$\Gamma_{ent} = \frac{d}{dt}\Delta\phi = \frac{G}{\hbar}\frac{m^2\sigma^2}{d^3}$$

$$m \approx 10^{-5} g$$

 $d \approx 100 \, \mu m$
 $\sigma \approx 1 \, nm$





ENTANGLEMENT RATE

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$$\tau_{ent} = \Gamma_{ent}^{-1} \approx 0.1 \, s$$





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ENTANGLEMENT RATE $G m^2 \sigma^2$ Γ_{ent} ħ

$$m \approx 10^{-5} g$$

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$$\tau_{ent} = \Gamma_{ent}^{-1} \approx 0.1 \, s$$

QUESTION/OBJECTION: You only used the Newton potential to generate entanglement!





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Belenchia, Wald, Giacomini, Castro-Ruiz, Brukner, Aspelmeyer, PRD (2018) 2019 First prize Essay of the Gravity Research Foundation

Along the lines of Baym, Ozawa (2009) and Mari, De Palma, Giovannetti (2016) For QED

> G is a field (GR) No interaction at a distance Linearized quantum gravity

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TAKE-HOME MESSAGE

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TAKE-HOME MESSAGE

Quantum properties of the gravitational field

QUANTIZED RADIATION VACUUM FLUCTUATIONS

are essential to obtain a consistent description of the experiment



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Belenchia, Wald, Giacomini, Castro-Ruiz, Brukner, Aspelmeyer, PRD (2018) 2019 First prize Essay of the Gravity Research Foundation **QUANTIZED RADIATION:** limits observation of interference **VACUUM FLUCTUATIONS:** limits "which-path information" (i.e. measure of entanglement) **NEWTON INTERACTION** B










NEWTON POTENTIAL AS QUANTUM INFORMATION CARRIER

TAKE-HOME MESSAGE

Quantum properties of the gravitational field

QUANTIZED RADIATION VACUUM FLUCTUATIONS

are essential to obtain a consistent description of the experiment

ARGUMENT: Newtonian potential has a quantum information content

Belenchia, Wald, Giacomini, Castro-Ruiz, Brukner, Aspelmeyer, PRD (2018) 2019 First prize Essay of the Gravity Research Foundation **QUANTIZED RADIATION:** limits observation of interference **VACUUM FLUCTUATIONS:** limits "which-path information" (i.e. measure of entanglement) **NEWTON INTERACTION** B











NEWTON POTENTIAL AS QUANTUM INFORMATION CARRIER

TAKE-HOME MESSAGE

Quantum properties of the gravitational field

QUANTIZED RADIATION VACUUM FLUCTUATIONS

are essential to obtain a consistent description of the experiment

ARGUMENT: Newtonian potential has a quantum information content

If instead we want to keep a classical description of gravity, we need to drastically modify our basic principles.

See also Danielson, Satishchandran, Wald PRD (2022)



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Belenchia, Wald, Giacomini, Castro-Ruiz, Brukner, Aspelmeyer, PRD (2018) 2019 First prize Essay of the Gravity Research Foundation **QUANTIZED RADIATION:** limits observation of interference **VACUUM FLUCTUATIONS:** limits "which-path information" (i.e. measure of entanglement) **NEWTON INTERACTION** B















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EMBEDDING INTO GR







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EMBEDDING INTO GR

1. Classical GR + QM does not generate entanglement THEORY-INDEPENDENT NO-GO THEOREM

Galley, F.G., Selby Quantum (2022)







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EMBEDDING INTO GR

1. Classical GR + QM does not generate entanglement THEORY-INDEPENDENT NO-GO THEOREM

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2. Newton interaction + no faster-than-light principle

⇒ vacuum fluctuations and gravitational radiation in a quantum state

Belenchia, Wald, F.G., Castro-Ruiz, Brukner, Aspelmeyer, PRD (2018)







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EXPERIMENT

Newton potential is compatible with the weak-field, non relativistic limit of GR







NEW RESULT: MORE GENERAL EFFECT THAN NEWTON POTENTIAL IN TABLE-TOP EXPERIMENTS

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EMBEDDING INTO GR

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Newton potential is compatible with the weak-field, non relativistic limit of GR

X -

Chen, F.G., 2402.10288 (2024)































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$$\hat{H} = \frac{\hat{p}^2}{2m} + \frac{m\omega^2}{2}\hat{x}^2 - m\gamma\hat{x}$$

Particular solution of classical EoM

$$x_{\gamma} = \frac{\gamma}{\omega^2}$$

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$$x_{\gamma} = \frac{\gamma}{\omega^2}$$

Change of coordinates $\hat{x}' = \hat{x} - x_{\gamma}$

$$\hat{H}' = \frac{\hat{p}'^2}{2m} + \frac{m\omega^2}{2}\hat{x}'^2 - \frac{m\gamma^2}{2\omega^2}$$

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$$\hat{H} = \frac{\hat{p}^2}{2m} + \frac{m\omega^2}{2}\hat{x}^2 - m\gamma\hat{x}$$

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Particular solution of classical EoM

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Chen, F.G., 2402.10288 (2024)

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linearized quantum gravity

 $g_{\mu\nu} = \eta_{\mu\nu} + h_{\mu\nu}$ $h_{\mu\nu} \to \hat{h}_{\mu\nu}$



 $\hat{H}_G(\hat{h}_{ij}, \hat{\pi}_{ij}) |\Psi\rangle_{G+S} = E_0(x) |\Psi\rangle_{G+S}$ (ground state)

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Free gravity Hamiltonian

linearized quantum gravity

 $g_{\mu\nu} = \eta_{\mu\nu} + h_{\mu\nu}$ $h_{\mu\nu} \to \hat{h}_{\mu\nu}$



$$\begin{split} \hat{H}_{G}(\hat{h}_{ij},\hat{\pi}_{ij}) \,|\,\Psi\rangle_{G+S} &= E_{0}(x) \,|\,\Psi\rangle_{G+S} & \text{Free} \\ \text{(ground state)} \\ \\ [\partial_{i}\partial^{i}\hat{h}^{T}(x) + \hat{T}_{00}(x)] \,|\,\Psi\rangle_{G+S} &= 0 & \text{Scalar} \\ \partial_{j}\,\hat{\pi}^{ij}(x) \,|\,\Psi\rangle_{G+S} &= 0 & \text{GAUGE} & \text{Verture} \\ \end{split}$$

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e gravity Hamiltonian

linearized quantum gravity

constraint: Gauss law

 $g_{\mu\nu} = \eta_{\mu\nu} + h_{\mu\nu}$ $h_{\mu\nu} \to \hat{h}_{\mu\nu}$

ector constraint: sversality condition



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$$\hat{H}_{G}(\hat{h}_{ij},\hat{\pi}_{ij}) |\Psi\rangle_{G+S} = E_{0}(x) |\Psi\rangle_{G+S}$$
(ground state)
$$\begin{bmatrix}\partial_{i}\partial^{i}\hat{h}^{T}(x) + \hat{T}_{00}(x)\end{bmatrix} |\Psi\rangle_{G+S} = 0$$
Scalar
$$\partial_{j}\hat{\pi}^{ij}(x) |\Psi\rangle_{G+S} = 0$$
GAUGE
$$\pi_{i}$$

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gravity Hamiltonian

linearized quantum gravity

 $g_{\mu\nu} = \eta_{\mu\nu} + h_{\mu\nu}$ $h_{\mu\nu} \to \hat{h}_{\mu\nu}$

constraint: Gauss law

ector constraint: sversality condition



GROUND STATE: coherent state displaced by the eigenvalue solution

- Localised source: Newton potential h - Delocalised source: general function







THE QUANTUM STATE OF GRAVITY OF A GENERAL QUANTUM SOURCE



Constraint changes the gravitational energy

$$\hat{C} = \partial_i \partial^i \hat{h}^T(x)$$

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teraction

$$\hat{H}_{I} = -\frac{1}{2} \int d^{3}x \hat{h}_{\mu\nu}(\vec{x}) [\hat{T}^{\mu\nu}_{A}(\vec{x}) + \hat{T}^{\mu\nu}_{B}(\vec{x})]$$
miltonian

... zero in the temporal gauge ($h^{0\mu} = 0$)!

 $\hat{T}(x) + \hat{T}^{A}_{00}(x) + \hat{T}^{B}_{00}(x)$





THE QUANTUM STATE OF GRAVITY OF A GENERAL QUANTUM SOURCE



Constraint changes the gravitational energy

$$\hat{C} = \partial_i \partial^i \hat{h}^T(x) + \hat{T}^A_{00}(x) + \hat{T}^B_{00}(x)$$

Entangling phase:

$$\Delta \phi = -\frac{G}{c^4 \hbar} \int d^3 x d^3 y \frac{E_A(\vec{x}) E_B(\vec{y})}{|\vec{x} - \vec{y}|}$$

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THE QUANTUM STATE OF GRAVITY OF A GENERAL QUANTUM SOURCE



$$\hat{C} = \partial_i \partial^i \hat{h}^T(x) + \hat{T}^A_{00}(x) + \hat{T}^B_{00}(x)$$

$$\Delta \phi = -\frac{G}{c^4 \hbar} \int d^3 x d^3 y \frac{E_A(\vec{x}) E_B(\vec{y})}{|\vec{x} - \vec{y}|}$$
 depend on the matter distribution

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teraction

$$\hat{H}_{I} = -\frac{1}{2} \int d^{3}x \hat{h}_{\mu\nu}(\vec{x}) [\hat{T}^{\mu\nu}_{A}(\vec{x}) + \hat{T}^{\mu\nu}_{B}(\vec{x})]$$
miltonian

... zero in the temporal gauge ($h^{0\mu} = 0$)!

Constraint changes the gravitational energy

Entangling phase:





COMPARISON TO CLASSICAL GRAVITY

 $\Delta \phi = -\frac{G}{c^4 \hbar} \int d^3 x d^3 y \frac{E_A(\vec{x}) E_B(\vec{y})}{|\vec{x} - \vec{v}|}$

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COMPARISON TO CLASSICAL GRAVITY

$$\Delta \phi = -\frac{G}{c^4 \hbar} \int d^3 x d^3 y \frac{E_A(\vec{x}) E_B(\vec{y})}{|\vec{x} - \vec{y}|}$$

Quantum state: coherent semiclassical state

$$|\psi\rangle \approx |\alpha(x_i, p_i)\rangle$$

$$\Delta \phi = -\frac{G}{\hbar} \frac{m_A m_B}{|\vec{x}_A - \vec{x}_B|}$$

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LIMIT TO THE NEWTON POTENTIAL

Classical mass density

$$E^{Z}(\vec{x}) = m_{Z}c^{2}\delta(\vec{x} - \vec{x}_{i})$$



COMPARISON TO CLASSICAL GRAVITY

$$\Delta \phi = -\frac{G}{c^4 \hbar} \int d^3 x d^3 y \frac{E_A(\vec{x}) E_B(\vec{y})}{|\vec{x} - \vec{y}|}$$
LIMIT TO THE NEWTON POTENTIAL
Quantum state: coherent semiclassical state
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CANNOT BE REPRODUCED WITH:

- 1. Newton potential
- 2. Schrödinger-Newton equation
- 3. Classical-quantum coupling (Semiclassical gravity)

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Total Hamiltonian

 $\hat{H}_{tot} = \hat{H}_S + \hat{H}_P + \hat{H}_G + \hat{H}_I$







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Total Hamiltonian

$$\hat{H}_{tot} = \hat{H}_S + \hat{H}_P + \hat{H}_G + \hat{H}_I$$

Commutator of gravity operators

$$[\hat{h}_{ij}(\vec{x}), \hat{\pi}^{kl}(\vec{x}')] = i\hbar\alpha\delta^k_{(i}\delta^l_{j)}\delta(\vec{x} - \vec{x}')$$







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$$[\hat{H}_G, \hat{H}_I] |\Psi\rangle_{SGP} \neq 0$$







Correction terms to relative phase $\Delta \phi \sim \frac{Gt}{c^4 \hbar} t^2 [\alpha f_1(T_P^{ij}) + \alpha^2 f_2(T_P^{ij})]$

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QUANTUM COMMUTATOR OF THE GRAVITATIONAL FIELD



Correction terms to relative phase

$$\Delta \phi \sim \frac{Gt}{c^4 \hbar} t^2 [\alpha f_1(T_P^{ij}) + \alpha^2 f_2(T_P^{ij})] \\ \propto t^2 \Delta E_P^2$$

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Total Hamiltonian

$$\hat{H}_{tot} = \hat{H}_S + \hat{H}_P + \hat{H}_G + \hat{H}_I$$

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QUANTUM COMMUTATOR OF THE GRAVITATIONAL FIELD



Correction terms to relative phase

$$\Delta \phi \sim \frac{Gt}{c^4 \hbar} t^2 [\alpha f_1(T_P^{ij}) + \alpha^2 f_2(T_P^{ij})] \\ \propto t^2 \Delta E_P^2$$

Chen, Giacomini, 2402.10288 (2024)

Total Hamiltonian

$$\hat{H}_{tot} = \hat{H}_S + \hat{H}_P + \hat{H}_G + \hat{H}_I$$

Commutator of gravity operators

$$[\hat{h}_{ij}(\vec{x}), \hat{\pi}^{kl}(\vec{x}')] = i\hbar\alpha\delta^k_{(i}\delta^l_{j)}\delta(\vec{x} - \vec{x}')$$

$$[\hat{H}_G, \hat{H}_I] |\Psi\rangle_{SGP} \neq 0$$

Stronger indication that gravity is a quantum field





Flaminia Giacomini - ETH Zurich

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 - We can do better!

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- when we consider a static quantum source of gravity in a delocalized state:
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 - **OPEN QUESTIONS:** concrete implementation and estimates



Flaminia Giacomini - ETH Zurich

BACK-UP SLIDES



	Electromagnetism	Linearized Gra
Temporal gauge	$A_0=0$	$h_{0\mu}=0$
Canonical variables	$\{A_i(\vec{x}), E_j(\vec{x}')\}$	$\{h_{ij}(\vec{x}), \pi^{kl}(\vec{x}')\}$
No. of constraints	1	4
Similar constraints	Gauss law in A basis	Vector constrain
(without matter)	$\partial_j \frac{\delta}{\delta A_j(\vec{x})} \Psi[A] = 0$	$\left \partial_i rac{\delta}{\delta h_{ij}(ec x)} \Psi[h_{ij}] = ight $
Similar constraints	Gauss law in E basis with charge	Scalar constraint
(with matter)	$\nabla \cdot E = \Delta \phi = \rho$	$\Delta h^T = -\rho$
Vacuum state	Gaussian of transverse mode	Gaussian of mode with zero
The d.o.f activated with a static source	Longitudinal mode A_L	Trace of transv h_T

Flaminia Giacomini - ETH Zurich

WHAT IS THE QUANTUM STATE OF GRAVITY ASSOCIATED TO A QUANTUM SOURCE?

Chen, Giacomini, Rovelli Quantum (2023)

Then the quantum state of the Coulomb/Newton field is the ground state $|h_i^0\rangle_G$ of the Hamiltonian with the charge/mass in the quantum state $|\Phi_i\rangle$

ized Gravity

$$|\Psi\rangle_{G+M} = \sum_{i} \alpha_{i} |\Phi_{i}, h_{i}^{0}\rangle_{MG}$$

constraint in h basis

 $\Psi[h_{ij}] = 0$

constraints

of transverse n vith zero trace

of transverse mode





THE QUANTUM STATE OF GRAVITY OF A GENERAL QUANTUM SOURCE



$$|\psi\rangle = \int d\mu(\alpha) \,\psi(\alpha) \,|\,\alpha\rangle$$

 $\langle \Delta x \rangle_{\alpha} \ll \operatorname{Exp.resol}.$ $\langle \Delta p \rangle_{\alpha} \ll \operatorname{Exp.resol}.$

...still the Newton potential!

Flaminia Giacomini - ETH Zurich

Chen, Giacomini, 2402.10288 (2024)

$$\langle \alpha_x | \alpha_{x+\epsilon} \rangle \neq 0$$

$$\langle \alpha_x, h_\alpha | \alpha_{x+\epsilon}, h_\alpha \rangle = 0$$

The shift by the classical solution makes the states perfectly distinguishable

$$|\psi\rangle = \int d\mu(E)\,\psi(E)\,|E\rangle$$

$$00(x)\,|E\rangle = E(\vec{x})\,|E\rangle \qquad E(\vec{x}) \neq mc^2\delta(\vec{x} - \langle\psi, h_{\psi}|\phi, h_{\phi}\rangle = \langle\psi|\phi\rangle$$





