Particles acceleration at astrophysical shocks Exercises

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Particle spectrum from the distribution function

• We want to derive the spectrum of particle accelerated at a plane parallel shock using the distribution function in momentum

Definition
$$f(t, \vec{x}, \vec{p}) \equiv \frac{dN}{dV d^3 p}$$

• Derivation of the transport equation

We can use the Liuville's theorem because we can neglect collisions (otherwise we should use the Boltzman equation)

$$\begin{cases} \frac{Df}{Dt} = \frac{\partial f}{\partial t} + \vec{u} \cdot \nabla_x f + \dot{\vec{p}} \nabla_p f = 0 \quad \rightarrow \text{ Vlasov equation} \\ \dot{\vec{p}} = q \frac{\vec{u}}{c} \times \vec{B} \quad \text{Lorentz force} \end{cases}$$

$$\frac{\partial f}{\partial t} + \vec{u} \dot{\nabla}_x f = \nabla \left[D_{xx} \nabla f \right] + \frac{p}{3} \nabla \cdot \vec{u} \frac{\partial f}{\partial p} + Q(x, p)$$

See e.g. Vietri's book sec. 4.3

 $\mathsf{D}_{\mathsf{x}\mathsf{x}}$ is the spatial diffusion coefficient

Particle spectrum from the distribution function

• Step 1: plane parallel (1D) stationary system: $\partial_{\partial t}$, $\nabla_x \to \partial_x$

$$u\frac{\partial f}{\partial x} = \frac{\partial}{\partial x} \left[D\frac{\partial f}{\partial x} \right] + \frac{p}{3}\frac{du}{dx}\frac{\partial f}{\partial p} + Q(x,p)$$

- Note: assume D spatially constant upstream and downstream
- Velocity profile:

 $u(x) = -u_2 + (u_2 - u_1)\theta(x)$

• Injection occurs only at the shock discontinuity:

$$\begin{cases} Q(x,p) = q_0 \,\delta(p-p_0) \,\delta(x) \\ q_0 = \frac{\eta_{\text{inj}} n_1 u_i}{4\pi p_0^2} \end{cases}$$

- Shock reference frame Downstream Upstream $-u_2$ $-u_1$
- Questions: derive the distribution function in the whole region (upstream, downstream and at the shock)
 - use the boundary condition: $f(x = +\infty) = 0$
 - \bullet assume that a fraction $\xi_{\rm cr} \approx 0.1$ of the shock bulk pressure is converted into CR energy

Particle spectrum from the distribution function

 Hint 1: start solving the equation upstream/downstream using the boundary condition f(x=0,p)=f₀(p)

$$\frac{\partial}{\partial x} \left[D \frac{\partial f}{\partial x} \right] - u \frac{\partial f}{\partial x} = 0$$

• Hint 2: solve equation for $f_0(p)$ by integrating across the shock discontinuity

$$\int_{0^{-}}^{0^{+}} \left\{ \frac{\partial}{\partial x} \left[D \frac{\partial f}{\partial x} \right] - u \frac{\partial f}{\partial x} + \frac{p}{3} \frac{\partial u}{\partial x} \frac{\partial f}{\partial p} + Q(x, p) \right\} = 0$$

• Hint 3: the normalisation is obtained from $P_{\rm CR} = \xi_{\rm CR} \rho_1 u_1^2$

$$P_{\rm CR} = \frac{1}{3} \int 4\pi \, p^2 dp f_0(p) pc \ \Rightarrow f_0(p) = \frac{3\xi_{\rm CR} \rho_1 u_1^2}{4\pi c (m_p c)^4 \Lambda} \left(\frac{p}{m_p c}\right)^{-s}$$

$$\Lambda' = \int_{1}^{p_{\max}/p_0} y^{3-s} dy$$

Particle spectrum from the distribution function

• Solution:

$$f_{0}(p) = \frac{3 \xi_{cr} \rho_{0} u_{sh}^{2}}{4\pi (mc)^{4} c \Lambda} \left(\frac{p}{mc}\right)^{-s}$$

$$s = \frac{3u_{1}}{u_{1} - u_{2}}$$

$$f_{1}(x, p) = f_{0}(p) e^{-u_{1} x/D_{1}}$$

$$f_{2}(x, p) = f_{0}(p)$$

CR spectrum escaping from a SNR

Determine the CR spectrum escaping from a SNR shock during the Sedov-Taylor phase, assuming that, during this phase, the shock converts a constant fraction of energy in escaping CR and the maximum energy decreases as a power low in time.

CR spectrum escaping from a SNR

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Solution

• Energy flux escaping the SNR shock:
$$4\pi p^2 dp \ Q_{esc}(p) \ pc = \xi_{esc}(t) \frac{1}{2} \rho_0 v_{sh}^3 \ 4\pi R_{sh}^2 dt$$

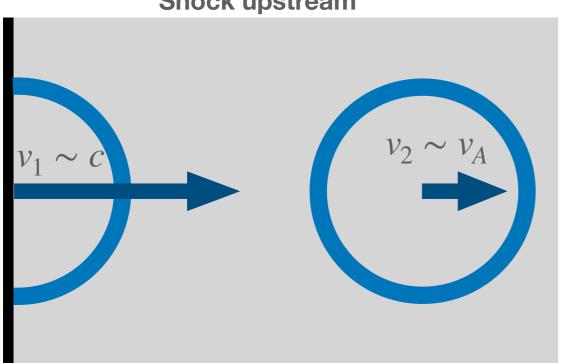
• Maximum energy in time: $p_{max}(t) = p_M \left(t/t_{ST} \right)^{-\delta} \Rightarrow \frac{dp_{max}}{dt} = -\delta \frac{p_{max}}{t}$
Evolution during the Sedov-Taylor phase
$$\begin{cases} R_{sh} = R_{ST} \left(\frac{t}{t_{ST}} \right)^{2/5} \\ v_{sh} = \frac{dR_{sh}}{dt} = \frac{2}{5} v_{ST} \left(\frac{t}{t_{ST}} \right)^{-3/5} \Rightarrow v_{sh}^3 R_{sh}^2 \propto t^{-1} \end{cases}$$

• Total released spectrum:
$$Q_{esc}(p) = \delta \left(\frac{2}{5} \right)^3 \frac{\rho_0}{c} R_{ST}^3 v_{ST}^2 t \ \xi_{esc}(t) p^{-4} \propto \xi_{esc}(t) p^{-4} \end{cases}$$
NOT RELATED TO FERMI ACCELERATION!!!

Additional question: how we can get a spectrum steeper that p-4?

Magnetic field amplification

- Estimate the amplification of magnetic field due to the streaming of CR particles upstream of a plane parallel shock assuming that $\xi_{
 m cr}$ is the fraction of shock kinetic energy converted into CR energy
- Hint: use the idea that particles upstream are fully isotropized after travelling a diffusion length



Shock upstream

Alfvén speed:
$$v_A = \frac{B_0}{\sqrt{4\pi\rho}}$$

Momentum of Alfvén waves $p_m \simeq \frac{1}{v_A} \frac{\delta B^2}{4\pi}$

Exercise n.3 - solution

Magnetic field amplification

• Assume that the momentum lost by particles goes into magnetic field

Initial momentum $p_1 \simeq n_{\rm cr} \gamma m v_1$; final momentum $p_2 \simeq n_{\rm cr} \gamma m v_2$

Momentum of Alfvén waves $p_m \simeq \frac{1}{v_A} \frac{\delta B^2}{4\pi}$

$$\frac{\delta B^2}{4\pi v_A} = p_1 - p_2 = n_{\rm cr} m\gamma (c/3 - v_A) \Rightarrow \delta B^2 \simeq 4\pi n_{\rm cr} m\gamma v_A c/3$$

$$\Rightarrow \frac{\delta B^2}{B_0^2} \simeq \frac{4\pi n_{\rm cr} m\gamma c v_A}{B_0^2} = \frac{4\pi n_{\rm cr} m\gamma c/3 v_A}{v_A^2 4\pi\rho} = \frac{n_{\rm cr}}{n_{\rm ism}} \frac{c/3}{v_A}$$

Using the distribution function from Ex. 1

$$n_{\rm cr} = \int f_0(p) d^3 p \rightarrow \frac{n_{\rm cr}}{n_{\rm ism}} = 3\xi_{\rm cr} \left(\frac{v_{\rm sh}}{c}\right)^2 \frac{\Lambda'}{\Lambda}$$

$$\Rightarrow \frac{\delta B}{B_0} \simeq \left(\xi_{\rm cr} \frac{\Lambda'}{\Lambda} \frac{v_{\rm sh}}{c} \frac{v_{\rm sh}}{v_A}\right)^{1/2} \sim 1 \qquad \qquad \Lambda' = \int_1^{p_{\rm max}/p_0} y^{2-s} dy$$