A REVIEW ON THE THEORY OF EXOTIC HADRONS

Angelo Esposito





WIFAI, Bari, November 2025

OUTLINE

I will give an overview of the current understanding of the nature of exotic hadrons

- What's up with exotic hadrons?
- EFT formulation of the problem
- Possible hints from "crowded" QCD
- Models and frameworks for the internal structure

<u>Disclaimer</u>: this will be a critical review, not free of personal opinions and likely incomplete



QCD is the theory of strong interactions

$$\mathcal{L} = \sum_{q} \bar{q}_i \left(i \gamma \cdot D_{ij} - m_q \delta_{ij} \right) q_j - \frac{1}{4} G^a_{\mu\nu} G^{\mu\nu,a}$$

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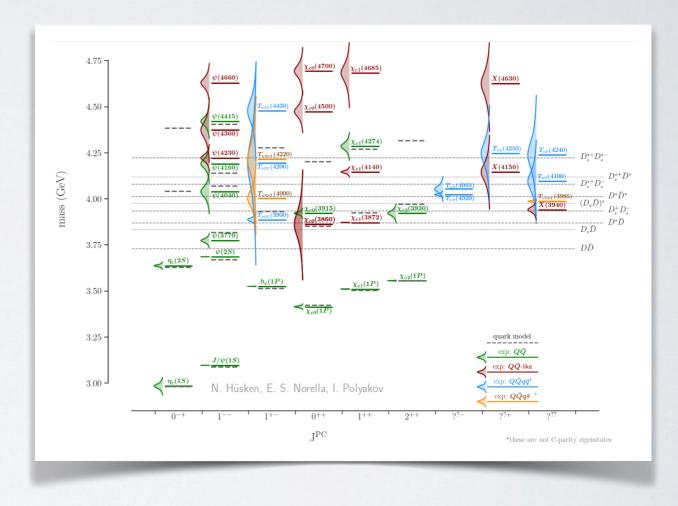
A priori, not only mesons and baryons

$$3 \otimes \bar{3} = 1 \oplus \dots$$
 meson $3 \otimes 3 \otimes 3 \otimes 3 = 1 \oplus \dots$ baryon $3 \otimes 3 \otimes \bar{3} \otimes \bar{3} = 1 \oplus \dots$ tetraquark $8 \otimes 8 \otimes \dots \otimes 8 = 1 \oplus \dots$ glueball

And many more...



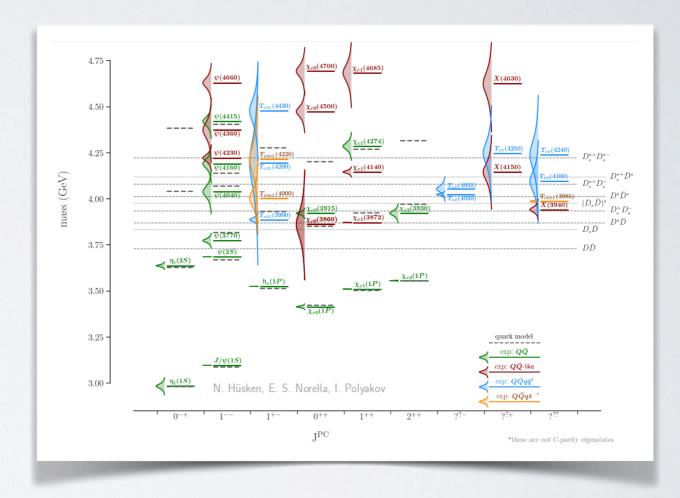
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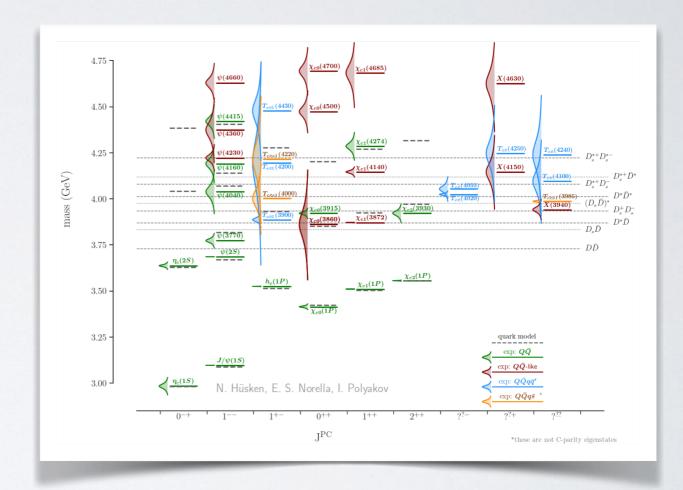
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- Isospin violation
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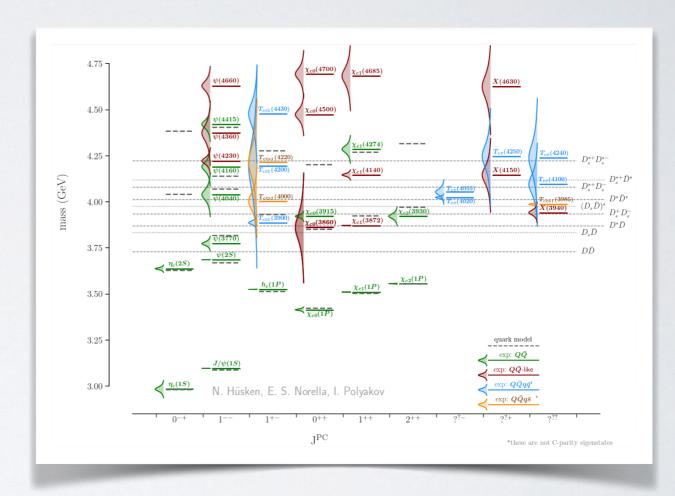
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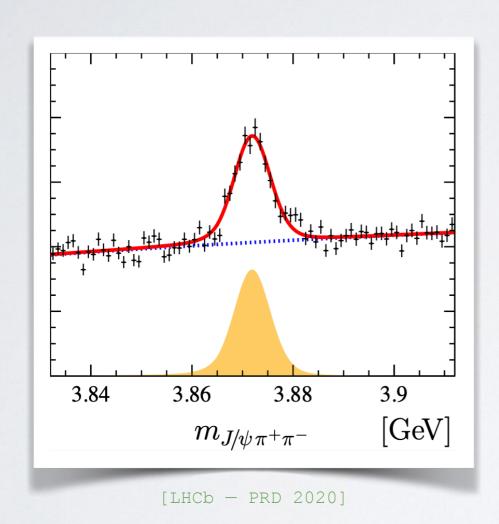
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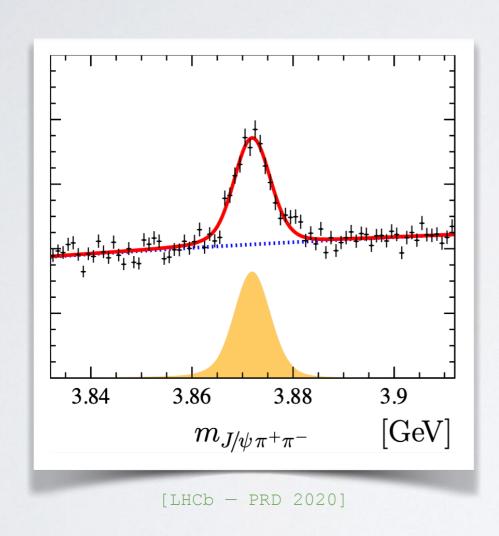


$$\left| m_X - m_{D^0} - m_{\bar{D}^{*0}} \right| \lesssim 0.1 \ \mathrm{MeV}$$
 very close to threshold

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very narrow

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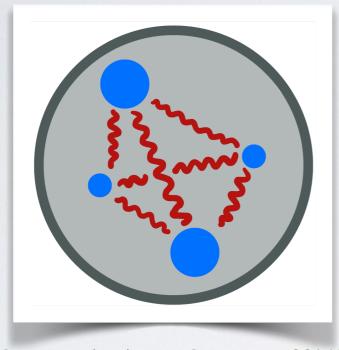
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Can be explained by a $Q \bar Q q \bar q$ structure with remarkable fine tunings

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[e.g. Maiani et al. - PRD 2014]

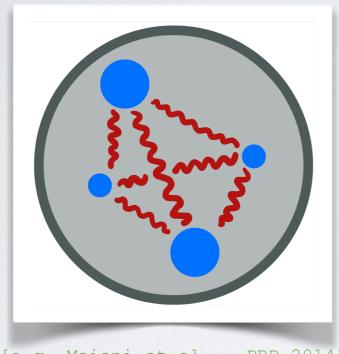
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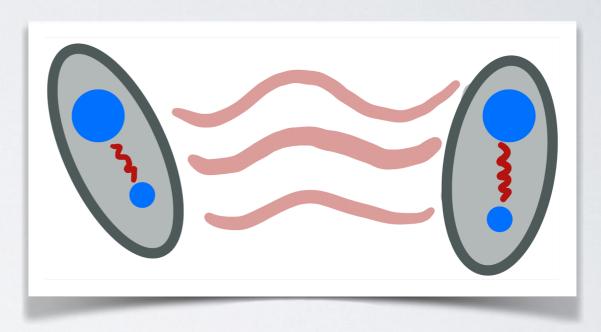
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Short distance QCD

Hadronic size $(r \leq 1 \text{ fm})$

Analogue to mesons/baryons



[e.g. Guo et al. - Rev.Mod.Phys. 2018]

Hadron molecule

Long distance QCD

Very large $(r \gg 1 \text{ fm})$

Analogue to deuteron

SHARP EFT FORMULATION

[Braaten, AE, Glioti, Guo, Hanhart, Hyodo, Kaplan, Kinugawa, Mikhasenko, Pilloni, Polosa, Rattazzi, and many more...]

Consider an EFT valid at distances $r \gg \Lambda_{\rm QCD}^{-1} \simeq 1$ fm. How does the exotic hadron arise?

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$$+ \cdots = \sum_{\text{pole}} = -\frac{2\pi^2}{\lambda_{\text{phys}}^2 \mu^3}$$

Option 2: tetraquark

Input degrees of freedom: two standard hadrons and one tetraquark

$$\mathcal{L}_{\text{tetra}} \sim \mathcal{L}_0(h_1, h_2) + \phi^{\dagger} \left(i \partial_t + \frac{\nabla^2}{2M} \right) \phi - \varepsilon \phi^{\dagger} \phi + g \phi h_1^{\dagger} h_2^{\dagger} + \text{h.c.}$$

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$$+ \cdots = \sum_{\text{pole at}} \text{pole at}$$

$$E_{\text{pole}} = -\varepsilon_{\text{phys}}$$

Both options need fine tuning to have the state near threshold

$$\frac{1}{\lambda_{\rm phys}} = \frac{1}{\lambda} + \mu \Lambda \ll \mu \Lambda$$
 or $\varepsilon_{\rm phys} = \varepsilon - g^2 \mu \Lambda \ll \frac{\Lambda^2}{\mu}$

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$$\mathcal{A}_{\text{mol}}(E) \sim \frac{\lambda_1 \sqrt{-2\mu E - i\epsilon} + \lambda_2}{1 - i\lambda_3 \left[\sqrt{2\mu E} + \sqrt{2\mu (E - \Delta)}\right] - \lambda_4 \sqrt{2\mu E} \sqrt{2\mu (E - \Delta)}}$$

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$$\mathcal{A}_{\text{tetra}}(E) \sim \frac{g^2}{E - \varepsilon + ig^2 \left[\sqrt{2\mu E} + \sqrt{2\mu(E - \Delta)} \right]}$$

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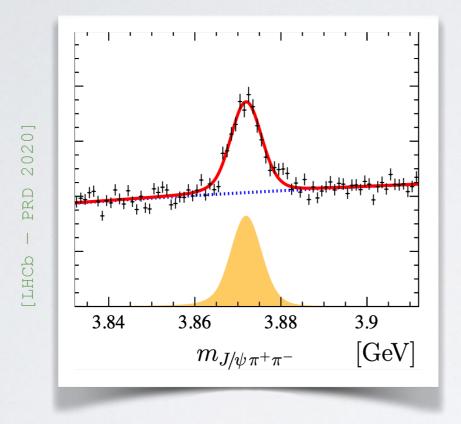
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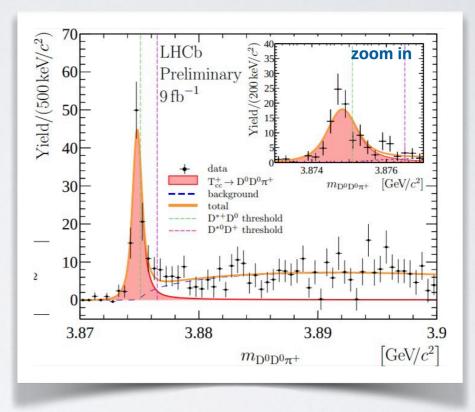
· ...and so on...



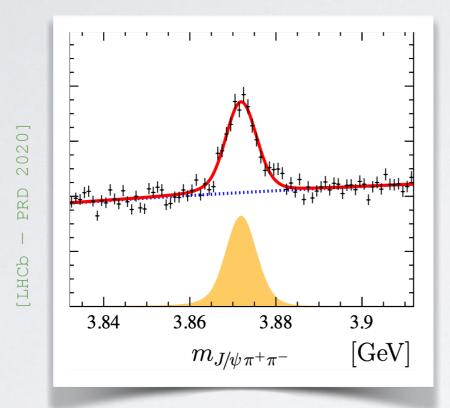
DATA

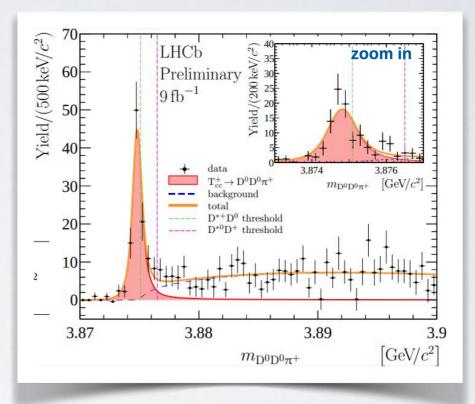
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[LHCb - Nature Phys. 2022]



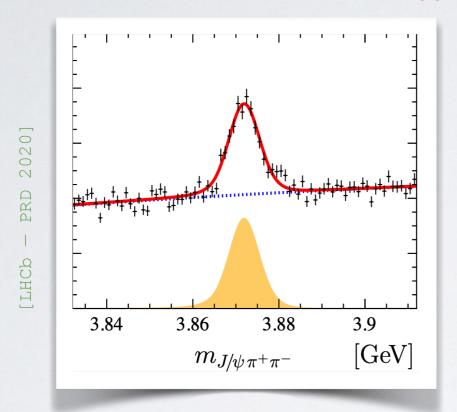


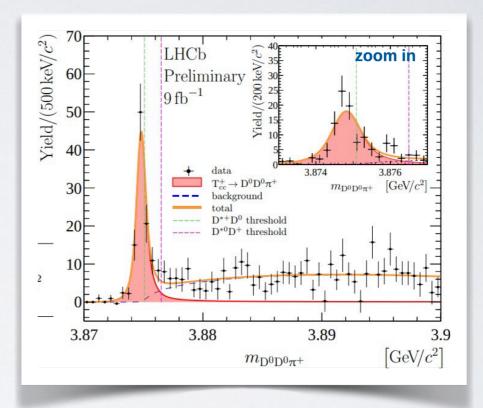
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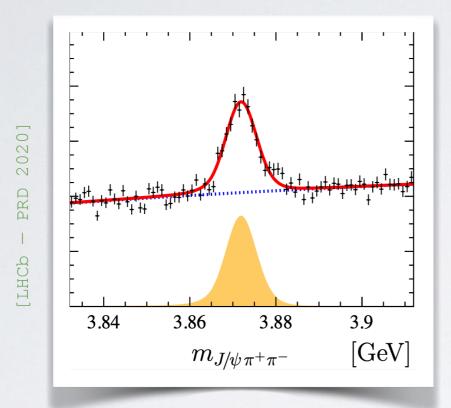
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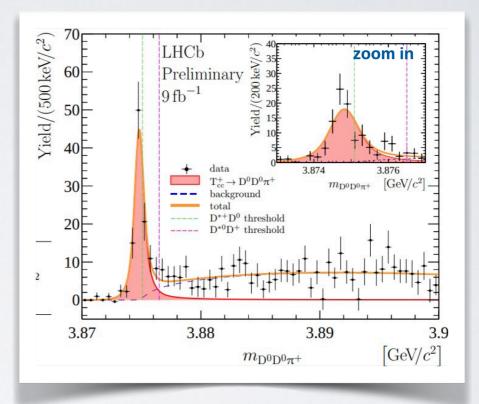
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...but the analysis is hard and the topic is still very controversial

[e.g., Baru et al. - PLB 2022]

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However, are current data sufficiently precise to even discriminate between different options?

[Abreu, Albaladejo, Braaten, AE, Ferreiro, Ingles, Liao, Navarra, Nieves, Pilloni, Polosa, Zang, and many more...]

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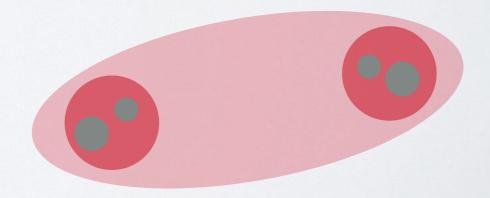
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Molecule

Binding of two $Q\bar{q}$ in a volume $V \sim 1 \text{ fm}^3 + \text{their coalescence}$



Generally believed that loosely bound molecules should be hard to

form in prompt high-energy collisions [e.g., Bignamini et al. - PRL 2009; Artoisenet, Braaten - PRD 2010; Meng, Han, Chao - PRD 2017]

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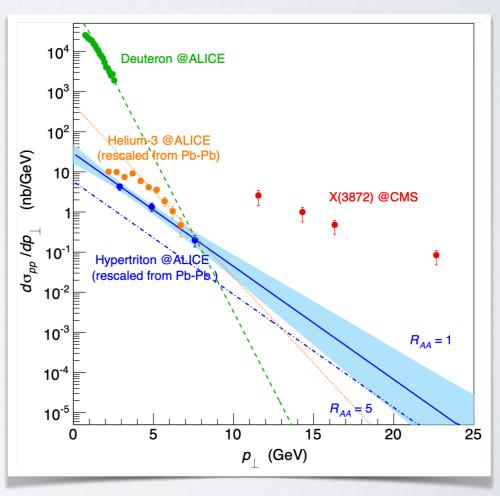
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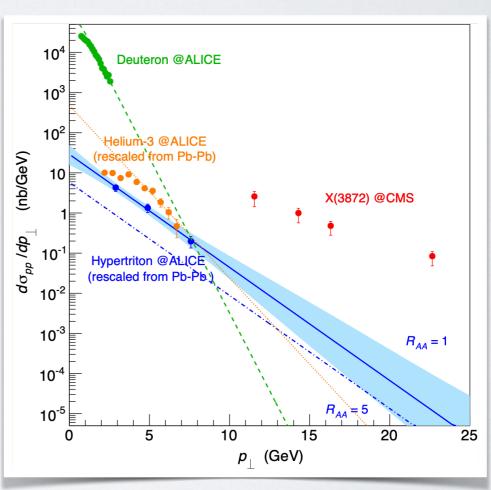
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... although this intuition has been [e.g., Artoisenet, Braaten - PRD 2010] challenged



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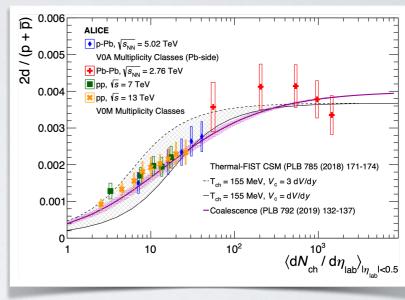
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On the contrary, tetraquarks should dissociate more due to higher number of collisions

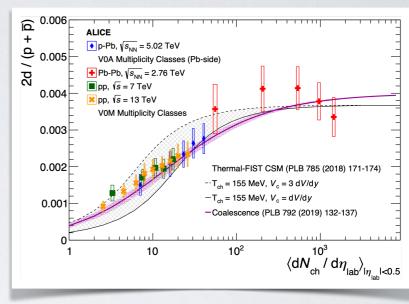
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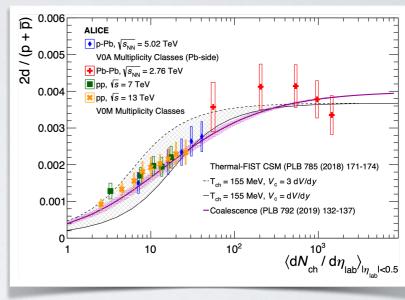
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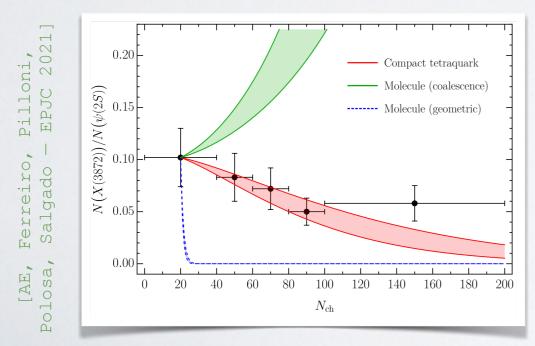
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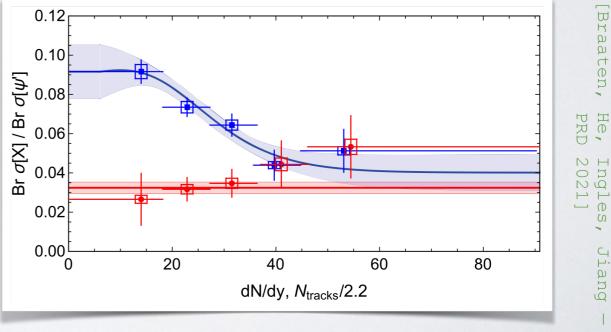
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[ALICE - PLB 2019, EPJC 2020]

This points to a tetraquark behavior... but adding one more fitting parameter molecules can also be accommodated



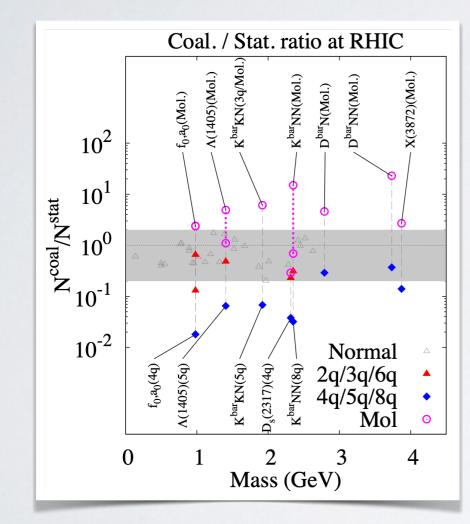


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HEAVY ION

A similar trend (more molecules, less tetraquarks) is also expected in heavy ion collisions

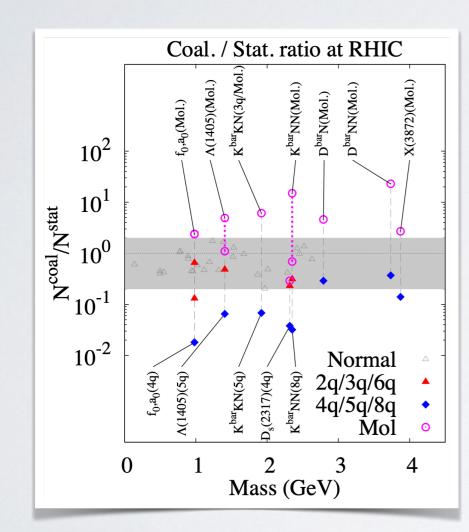
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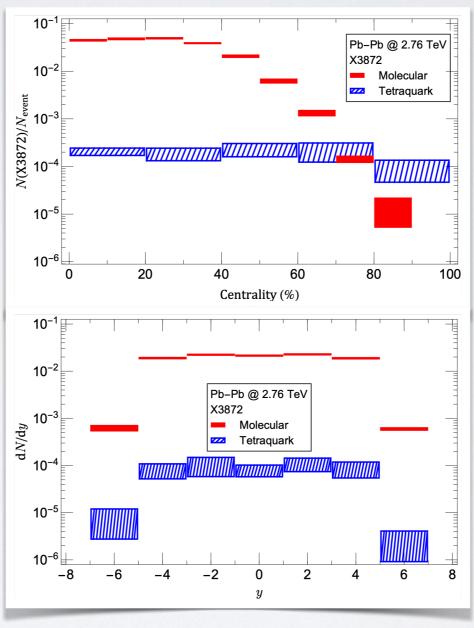
[ExHIC coll. - PRL 2011]

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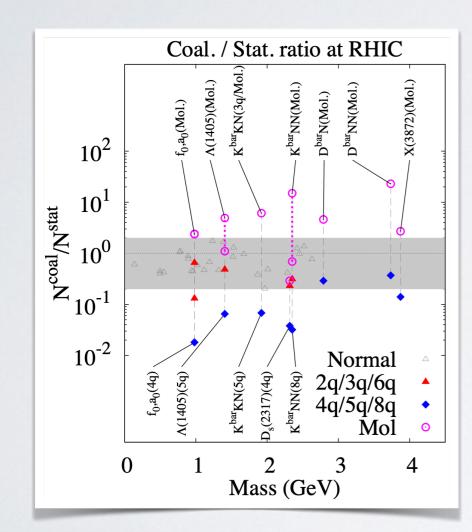
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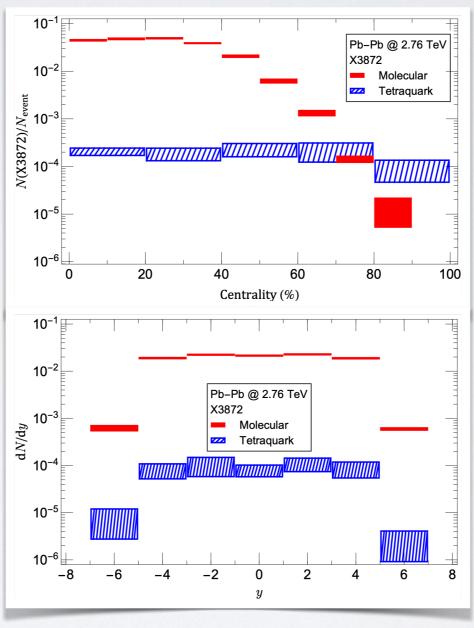
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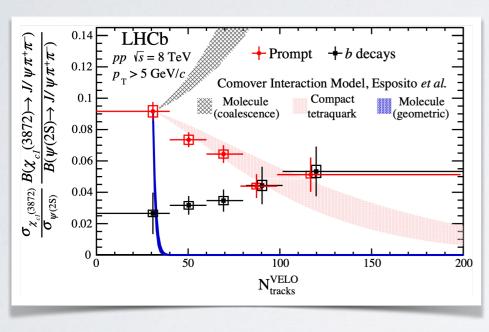
[Chen et al. - PRC 2022]

Although, not

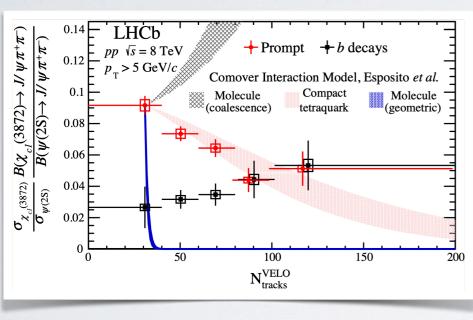
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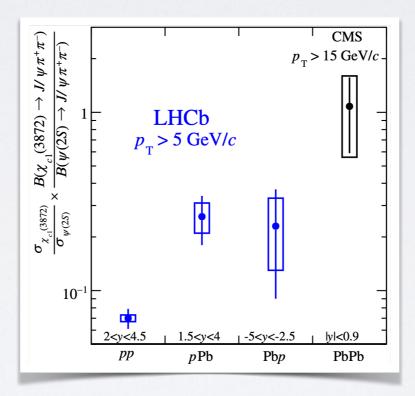




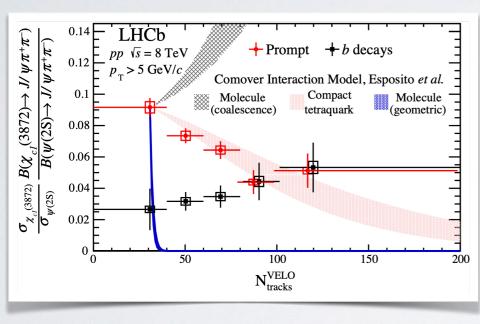
[LHCb - PRL 2021]



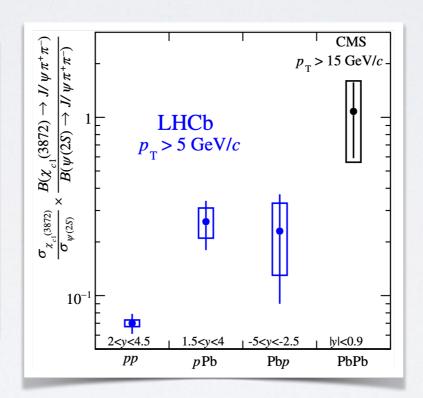
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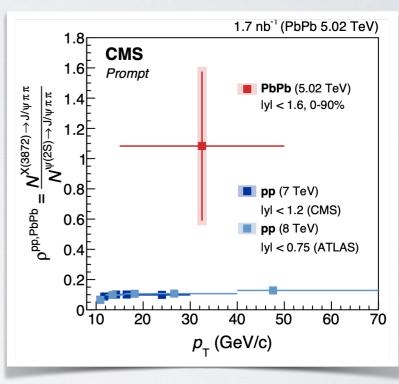
[LHCb - PRL 2024]



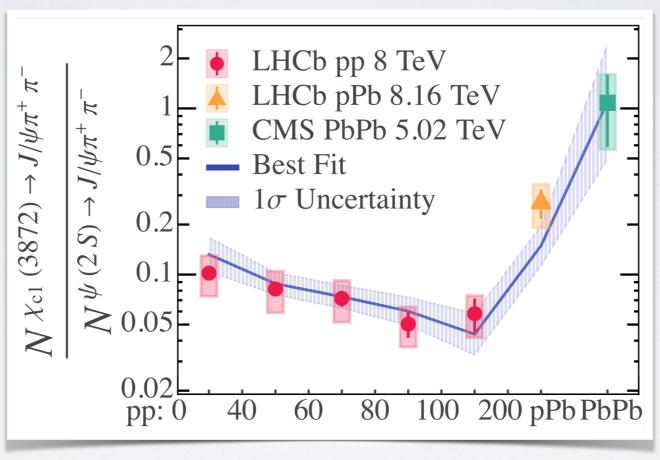
[LHCb - PRL 2021]



[LHCb - PRL 2024]

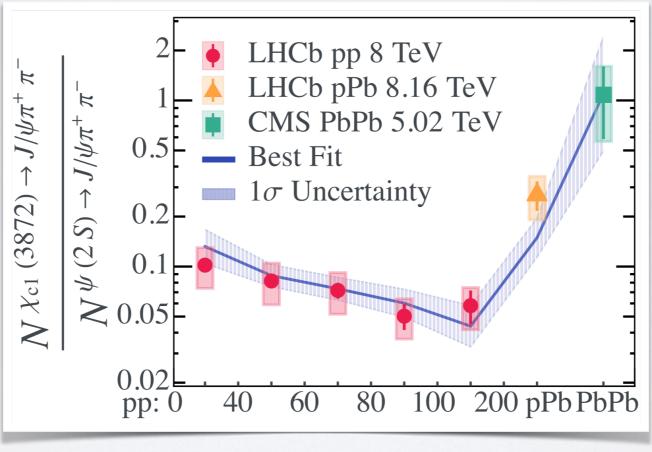


[CMS - PRL 2022]



[Guo, Guo, Liao, Wang, Xing - PRC 2024]

However, experimental data show a curious trend when going from small to large systems $(pp \rightarrow pPb \rightarrow PbPb)$



[Guo, Guo, Liao, Wang, Xing - PRC 2024]

Could be due to new creation mechanisms kicking in for larger systems

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However: are we sure we can learn anything about puzzling particles from systems we do not fully understand yet?

FRAMEWORKS FOR THE INTERNAL STRUCTURE

[Brambilla, Brodsky, Germani, Grinstein, Guo, Hanart, Lebed, Maiani, Piccini, Polosa, Vairo, and many more...]

The most natural thing to do (ideally) would be to study the actual internal structure of these states

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Tetraquark

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Very first attempt:

$$H = \sum_{i} m_{i} + \sum_{i,j} \kappa_{i,j} \mathbf{S}_{i} \cdot \mathbf{S}_{j}$$

[Maiani, Piccini, Polosa, Riquer - PRD 2005; Maiani, Piccini, Polosa, Riquer - PRD 2014]

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Molecule

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$$V_{\pi}(r) \sim \frac{e^{-\mu r}}{m_{\pi}r}$$

[Swanson - PLB 2004; Tornqvist - PLB 2004]

In both instances, models, parameters and mechanism have quickly proliferated

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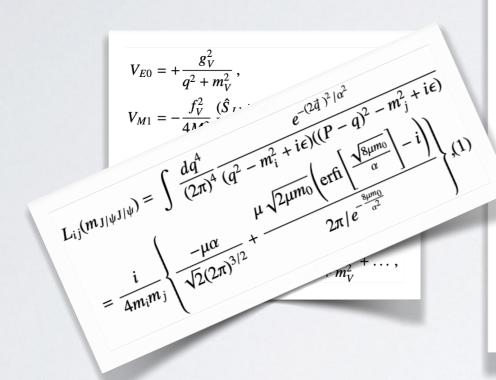
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$$\begin{split} V_{E0} &= + \frac{g_V^2}{q^2 + m_V^2} \,, \\ V_{M1} &= - \frac{f_V^2}{4M^2} \, \frac{(\hat{S}_{L1} \times \vec{q}) \cdot (\hat{S}_{L2} \times \vec{q})}{q^2 + m_V^2} \\ &= - \frac{2}{3} \, \frac{f_V^2}{4M^2} \, \hat{S}_{L1} \cdot \hat{S}_{L2} \, \frac{q^2}{q^2 + m_V^2} + \dots \,, \\ V_{E2} &= + \frac{h_V^2}{4M^4} \, \frac{(\hat{Q}_{L1,ij} \, q_i q_j) \, (\hat{Q}_{L2,lm} \, q_l q_m)}{q^2 + m_V^2} \\ &= + \frac{h_V^2}{36M^4} \, (\hat{Q}_{L1,ij} \, \hat{Q}_{L2,ij}) \, \frac{q^4}{q^2 + m_V^2} + \dots \,, \end{split}$$

Molecule	J^P	V	Attractive'
$Dar{D}$	0++	C_0	Yes
$D^*ar{D}$	1++	$C_0 + C_1$	Most
$D^*ar{D}$	1+-	$C_0 - C_1$	Likely
$D^*\bar{D}^*$	0++	$C_0 - 2C_1$	Likely
$D^*\bar{D}^*$	1+-	$C_0 - C_1$	Likely
$D^*\bar{D}^*$	2++	$C_0 + C_1$	Most
Molecule	J^P	V	Attractive'
$ar{D}\Sigma_c$	1 -	D_0	Yes
$\bar{D}\Sigma_c^*$	$\frac{1}{2}$ $\frac{3}{2}$ $\frac{1}{2}$ $\frac{1}{2}$ $\frac{1}{2}$ $\frac{1}{2}$ $\frac{1}{2}$ $\frac{1}{2}$ $\frac{1}{2}$ $\frac{1}{2}$ $\frac{1}{2}$	D_0	Yes
$\bar{D}^*\Sigma_c$	1 -	$D_0 - \frac{4}{3} D_1$	Likely
$\bar{D}^*\Sigma_c$	$\frac{3}{2}$	$D_0 + \frac{2}{3} D_1$	Most
$\bar{D}^*\Sigma_c^*$	1 -	$D_0 - \frac{5}{3} D_1$	Likely
$\bar{D}^*\Sigma_c^*$	$\frac{3}{2}$	$D_0 - \frac{2}{3} D_1$	Likely
$ar{D}^*\Sigma_c^*$	<u>5</u> -	$D_0 + D_1$	Most
Molecule		V	Attractive'
$\Sigma_c \bar{\Sigma}_c$	0-+	$E_0 - \frac{4}{3}E_1$	Likely
$\Sigma_car{\Sigma}_c$	1		Yes
$\Sigma_c^* \bar{\Sigma}_c$	1-+	$E_0 - E_1 - \frac{15}{2}E_2$	Likely
$\Sigma_c^* \bar{\Sigma}_c$	1	$E_0 - \frac{11}{9}E_1 + \frac{15}{2}E_2$	Likely
$\Sigma_c^* \bar{\Sigma}_c$	2-+	$E_0 + \frac{1}{3}E_1 - \frac{3}{2}E_2$	Likely
$\Sigma_c^* \bar{\Sigma}_c$	2	$E_0 + E_1 + \frac{3}{2}E_2$	Most
$\Sigma_c^* \bar{\Sigma}_c^*$	0^{-+}	$E_0 - \frac{15}{9}E_1 + \frac{15}{2}E_2$	Likely
$\Sigma_c^* \bar{\Sigma}_c^*$	1	$E_0 - \frac{11}{9}E_1 + \frac{3}{2}E_2$	Likely
$\Sigma_c^* \bar{\Sigma}_c^*$	2-+	$E_0 - \frac{1}{3}E_1 - \frac{9}{2}E_2$	Likely
$\Sigma_c^* \bar{\Sigma}_c^*$	3	$E_0 + E_1 + \frac{3}{2}E_2$	Most

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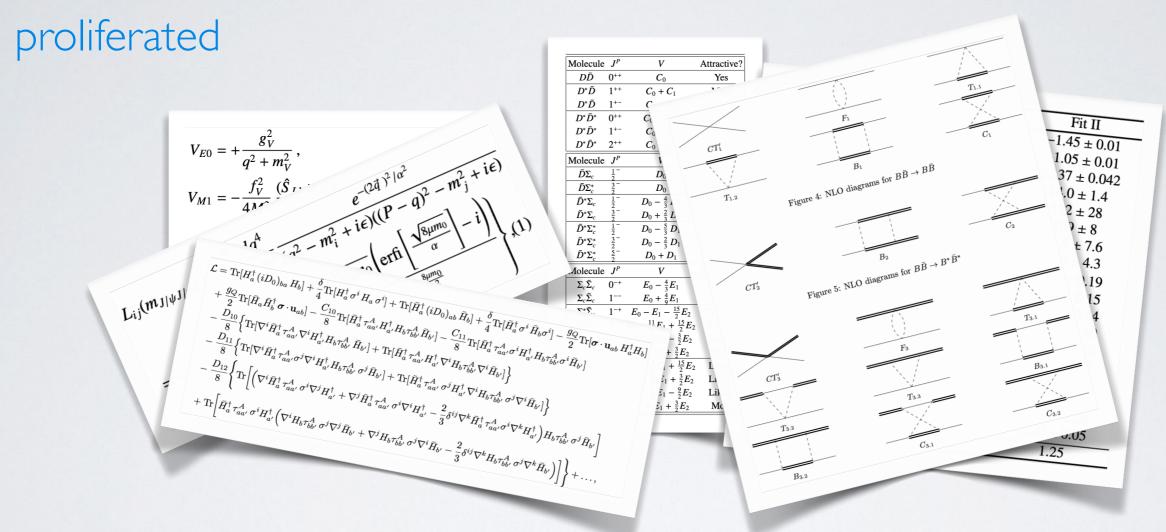
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	$ar{D}\Sigma_c$	$\frac{\frac{1}{2}^{-}}{\frac{3}{2}^{-}}$ $\frac{1}{2}^{-}$ $\frac{1}{2}^{-}$ $\frac{3}{2}^{-}$	D_0	Yes
	$\bar{D}\Sigma_c^*$	$\frac{3}{2}^{-}$	D_0	Yes
	$\bar{D}^*\Sigma_c$	1-	$D_0 - \frac{4}{3} D_1$	Likely
	$\bar{D}^*\Sigma_c$	$\frac{3}{2}^{-}$	$D_0 + \frac{2}{3} D_1$	Most
	$\bar{D}^*\Sigma_c^*$	$\frac{1}{2}$ $\frac{1}{3}$ $ \frac{1}{2}$ $\frac{1}{2}$ $ \frac{1}{2}$ $ \frac{1}{2}$ $ -$	$D_0 - \frac{5}{3} D_1$	Likely
V	$\bar{D}^*\Sigma_c^*$	$\frac{3}{2}^{-}$	$D_0 - \frac{2}{3} D_1$	Likely
1	$\bar{D}^*\Sigma_c^*$		$D_0 + D_1$	Most
	Molecule	J^P	V	Attractive?
	$\Sigma_c \bar{\Sigma}_c$	0-+	$E_0 - \frac{4}{3}E_1$	Likely
	$\Sigma_c \bar{\Sigma}_c$	1	$E_0 + \frac{4}{9}E_1$	Yes
	$\Sigma_c^* \bar{\Sigma}_c$	1-+	$E_0 - E_1 - \frac{15}{2}E_2$	Likely
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	$\Sigma_c^* \bar{\Sigma}_c$	2^{-+}	$E_0 + \frac{1}{3}E_1 - \frac{3}{2}E_2$	Likely
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		_		

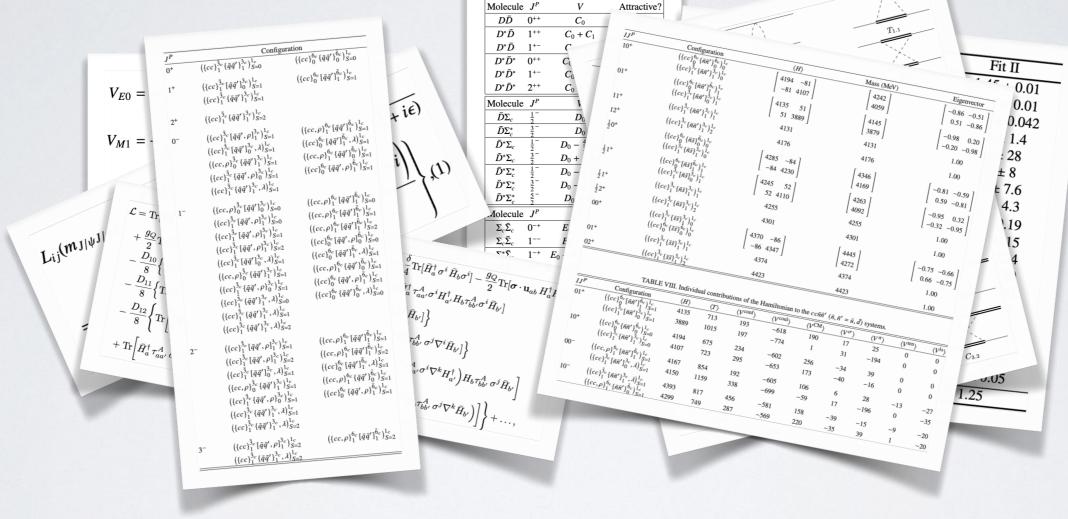
Parameters		
c_0 (GeV ⁻¹)	Fit I	
C' (GeV-1)	-1.52 ± 0.02	Fit II
c_0' (GeV ⁻¹)	-0.946 ± 0.058	-1.45 ± 0.01
$ g'_{direct}/g_{direct} $	0.0767 ± 0.0204	-1.05 ± 0.01
$g_{J/\psi J/\psi}/g_{direct}$	8.53 ± 3.64	0.137 ± 0.042
$ g_{\eta_{cX_{c1}}}/g' $	91.6 ± 3.64	14.0 ± 1.4
$ \mathcal{S}J/\psi_{h_c}/\mathcal{G}_J $	91.6 ± 75.4	112 ± 28
$\delta_{\chi_{c0}\chi_{c1}}/g_{direct}$	69.7 ± 16.1	109 ± 8
$ g_{\chi_{c0\chi'_{c1}}}/g_{direct} $	33.3 ± 8.2	38.5 ± 7.6
$\phi_{J/\psi J/\psi}$ (rad)	25.8 ± 10.6	10.0 ± 7.6
$\phi_{\eta_{c\chi_{c1}}}$ (rad)	1.53 ± 0.51	19.0 ± 4.3
$\phi_{J/\psi h_c}$ (rad)	2.69 ± 0.20	3.16 ± 0.19
$\phi_{\chi_{c0}\chi_{c1}}$ (rad)	4.40 ± 0.33	2.80 ± 0.15
$\phi_{\chi_{c0}\chi'_{c1}}$ (rad)	2.14 ± 0.18	2.95 ± 0.24
$\alpha_{v_1v_2}$ (C.13)	2.00 ± 0.33	2.89 ± 0.20
$\alpha_{J/\psi J/\psi}$ (GeV)	1.71 ± 0.01	3.23 ± 0.20
$\alpha_{\eta_c\chi_{c1}}$ (GeV)	1.71 ± 0.01	2.30 ± 0.21
$\alpha_{J/\psi h_c}$ (GeV)	1.71 ± 0.01	1.20 ± 0.21
$\alpha_{\chi_{c0}\chi_{c1}}$ (GeV)	1.71 ± 0.01 1.71 ± 0.01	1.20 ± 0.03
$\alpha_{\chi_{c0\chi'_{c1}}}$ (GeV)	1.71 ± 0.01	1.73 ± 0.26
$\chi^2/d.o.f$	1.71 ± 0.01	5.20 ± 0.05
	1.41	
		1.25

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proliferated



Beside a few exceptions, properties and spectrum are extremely model

dependent -- hardly the right way to shed light on the issue

A more systematic framework leverages the existing hierarchies

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Since $m_Q\gg m_q$, $\Lambda_{\rm QCD}$ one develops a Born-Oppenheimer approximation for the 4-quark constituents

[e.g., Braaten, Langmack, Smith - PRD 2014; Berwein, Brambilla, Castellà, Vairo - PRD 2015; Maiani, Polosa, Riquer - PRD 2019; Allaman, Ekterachian, Nardi, Rattazzi, Stelz - JHEP 2024; Braaten, Bruschini - PLB 2025]

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Light quark energy as effective potential for the heavy quarks

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Light quark energy as effective potential for the heavy quarks

The separation of scales makes the problem clean and with corrections that can be estimated systematically

The systematic EFT for the Born-Oppenheimer approximation has

been developed -> see Antonio Vairo's talk!

[Berwein, Brambilla, Castellà, Vairo - PRD 2015; Brambilla, Krein, Castellà, Vairo - PRD 2018; Soto, Castellà - PRD 2020; Berwein, Brambilla, Mohapatra, Vairo - PRD 2024; Brambilla, Mohapatra, Vairo - PRL 2025]

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• Large-N and large m_Q expansion

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- Production in ultra-periferal heavy ion collisions

[e.g., Moreira, Bertulani, Gonçalves, Navarra — PRD 2016; AE, Manzari, Pilloni, Polosa — PRD 2021; Gonçalves, Moreira — PLB 2021; Niu, Ding, Wang, Yang — CPL 2025]

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Thank you for the attention!