





Analysis of $D^0 \to K_S^0 \pi^+ \pi^-$ decays in Run 3

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Overview

Introduction

- Mixing and CPV
- Status of the art
- $D^0 \to K_S^0 \pi^+ \pi^- \text{channel}$
- Bin-flip method
- LHCb experiment

Run 3 analysis

- HLT1 Selection
- Datasets and Selection
- Signal extraction
- Efficiency Maps

Conclusions

Neutral meson mixing

A quantum mechanical phenomenon in which neutral mesons can oscillate between their particle and anti-particle state

The D^0 system has the lowest frequency of oscillations

This is due to the D^0 system having a very small mass difference between the mass eigenstates D_1 and D_2

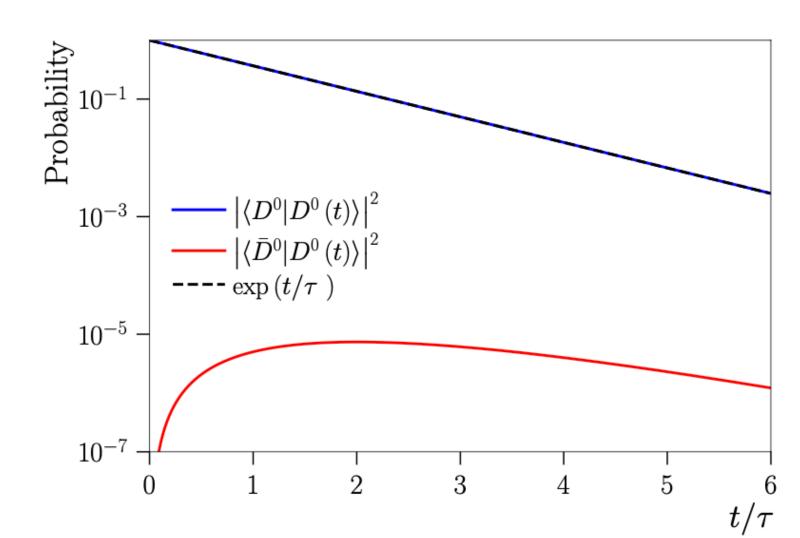
Their mass eigenstates, related to the flavor eigenstates, are:

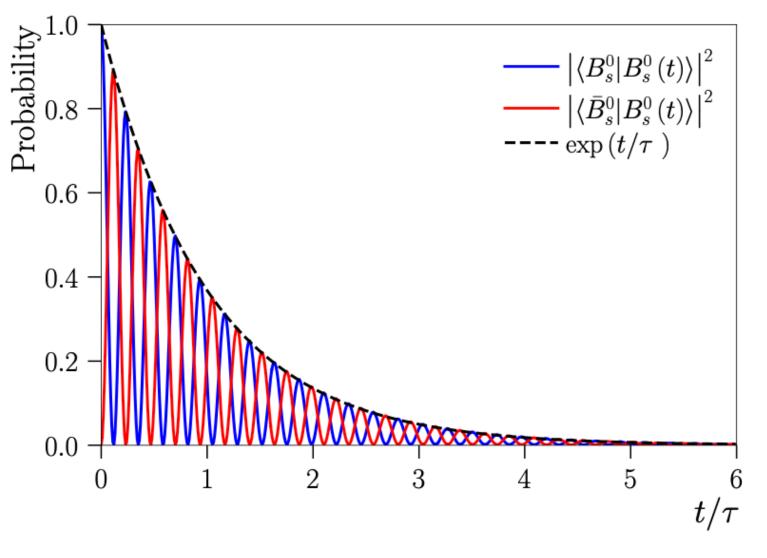
$$|D_{1,2}\rangle = p |D^0\rangle \pm q |\bar{D}^0\rangle$$
 with $|p|^2 + |q|^2 = 1$

In the limit of CP symmetry, q = p and the oscillations characterized by two dimensionless parameters:

$$x \equiv \frac{\Delta m}{\Gamma} = \frac{2(m_1 - m_2)}{\Gamma_1 + \Gamma_2} , \qquad \qquad y \equiv \frac{\Delta \Gamma}{2\Gamma} = \frac{\Gamma_1 - \Gamma_2}{\Gamma_1 + \Gamma_2}$$

where $m_{1,2}$ and $\Gamma_{1,2}$ are the mass and decay width of the CP-even/odd eigenstate $D_{1,2}$ and Γ is the average decay width





Types of CP

CPV in mixing:

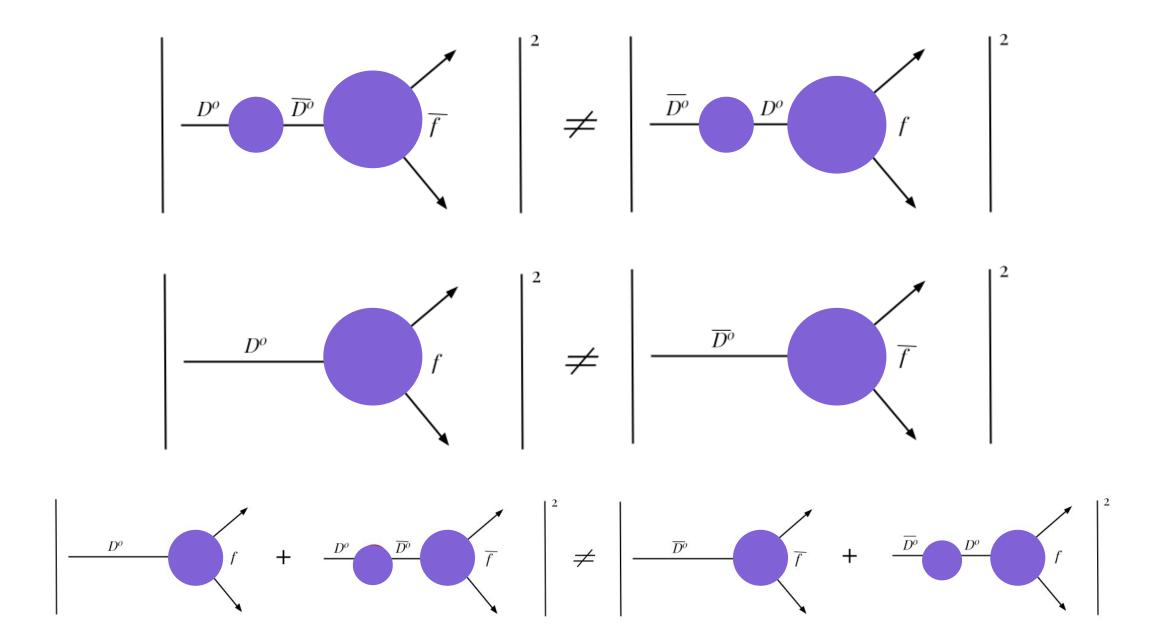
 $^{\square}$ Different probabilities for $D^0 \leftrightarrow \overline{D}{}^0$ oscillations

CPV in decay:

□ Difference between decay rates $D^0 \to f$ and $\overline{D}{}^0 \to \overline{f}$

CPV in interference (mixing-decay):

 ${}^{\Box}$ When D^0 and \overline{D}^0 decay to the same final state, amplitudes interfere



Note that decay channels have different sensitivities to the variables $x = \frac{\Delta m}{\Gamma}$ and $y = \frac{\Delta \Gamma}{2\Gamma}$

Multibody States (e.g. $D^0 \to K_S^0 \pi^+ \pi^-$):

- \Box Dalitz plot \Longrightarrow access to local variation of decay amplitudes and strong phases vs phase space
- \Box Sensitive to strong-phase differences across bins \Longrightarrow requires an amplitude model or external inputs

$$\Rightarrow x, \Delta x, y, \Delta y$$
$$\Rightarrow x, \Delta x$$

Status of the Art

LHCb:

$$D^0 \to K_S^0 \pi^+ \pi^- \text{in } 2021 \text{ (PRL127,111801)}$$

$$x = (3.98^{+0.56}_{-0.54}) \times 10^{-3}$$

$$y = (4.6^{+1.5}_{-1.4}) \times 10^{-3}$$

□ first observation of $\Delta m \neq 0$ in D^0 system (significance exceeding 7σ)

Belle & BelleII:

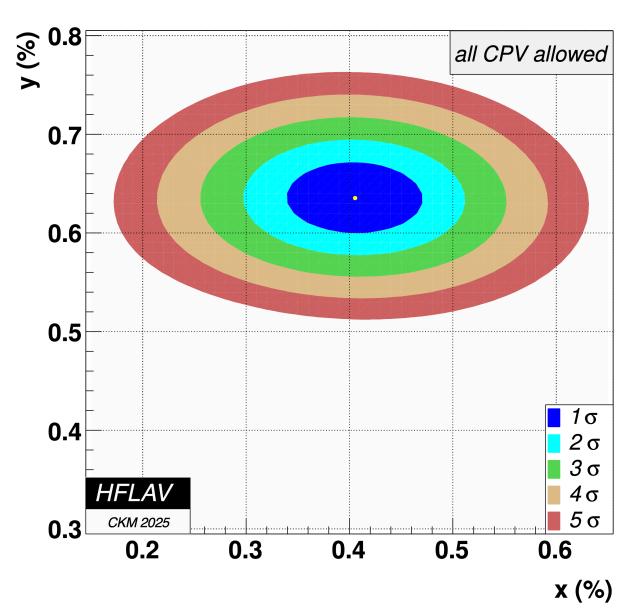
$$D^0 \to K_S^0 \pi^+ \pi^- \text{in } 2025 \text{ (PRD111,112011)}$$

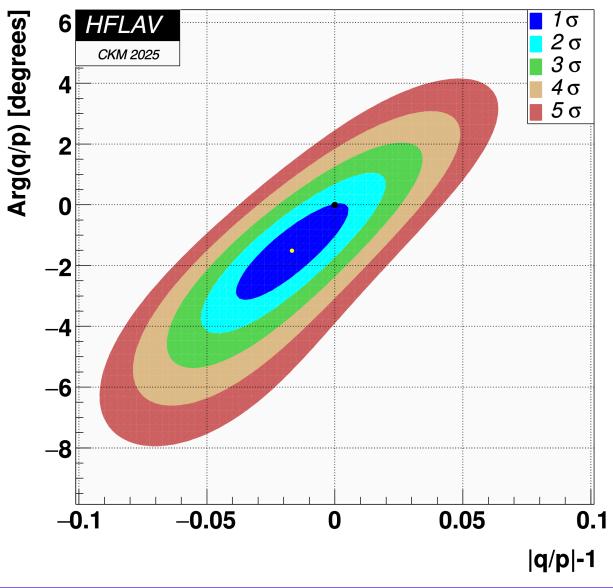
a Assuming no CPV:

$$x = (4.0 \pm 1.7 \pm 0.4) \times 10^{-3}$$

$$y = (2.9 \pm 1.4 \pm 0.3) \times 10^{-3}$$

- Results are consistent with previous determinations (PRD 89, 091103(R))
 - 20% and 14% improvement compared to previous model-dependent determinations (PRD 89, 091103(R))





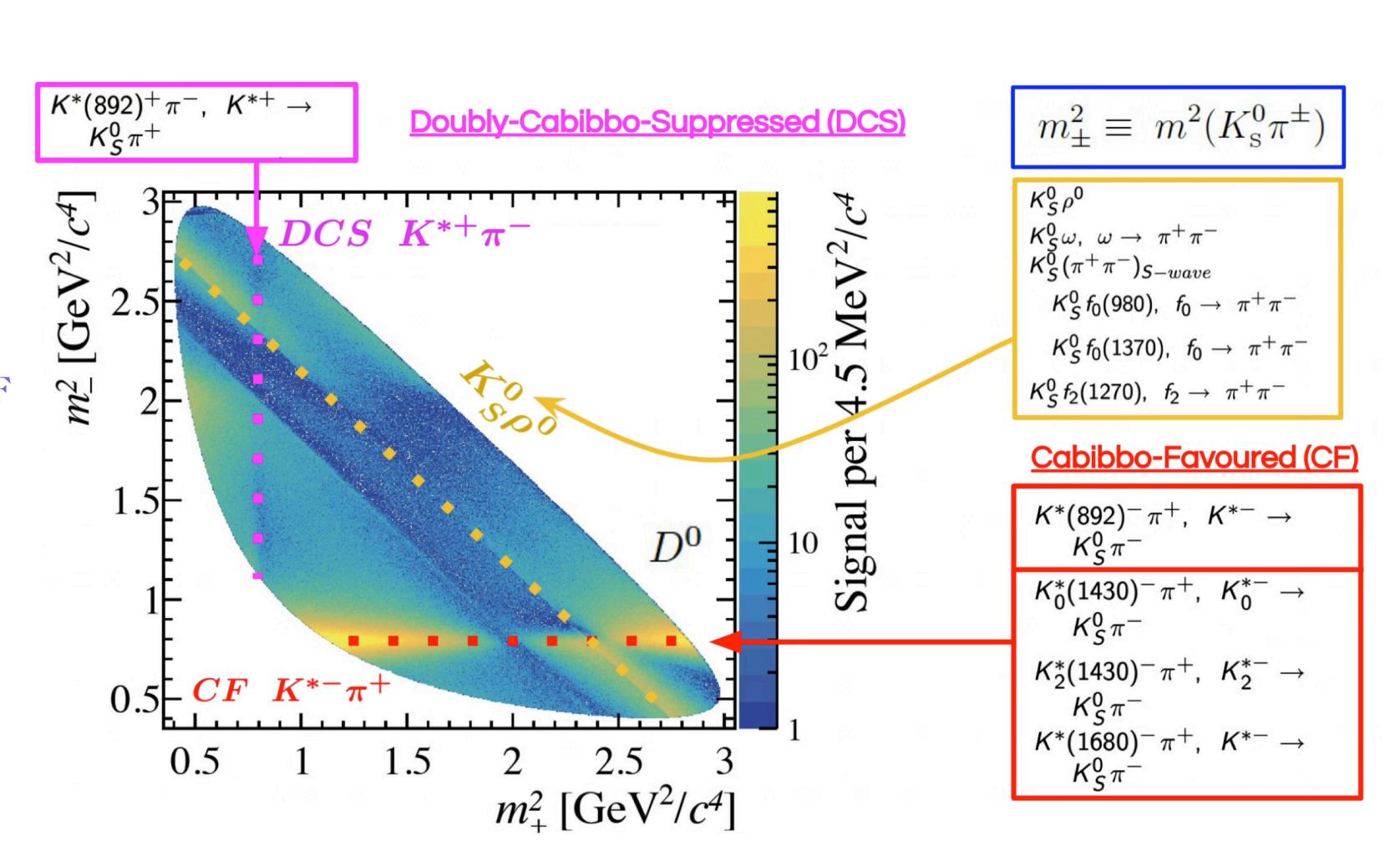
$$D^0 \to K_S^0 \pi^+ \pi^-$$
 channel

- [□] Since Run 3, with the removal of LO hardware trigger, the selection efficiency of $D^0 \to K_S^0 \pi^+ \pi^-$ decays improved
- Run 2 trigger-induced decay-time/Dalitz correlations
 - main source of systematics
- Removal of LO allows us to consider a single selection stage
 - which is easier to simulate compared to the hardware one
 - reduced trigger-related systematics and improved control of trigger effects

Bin-flip method (I)

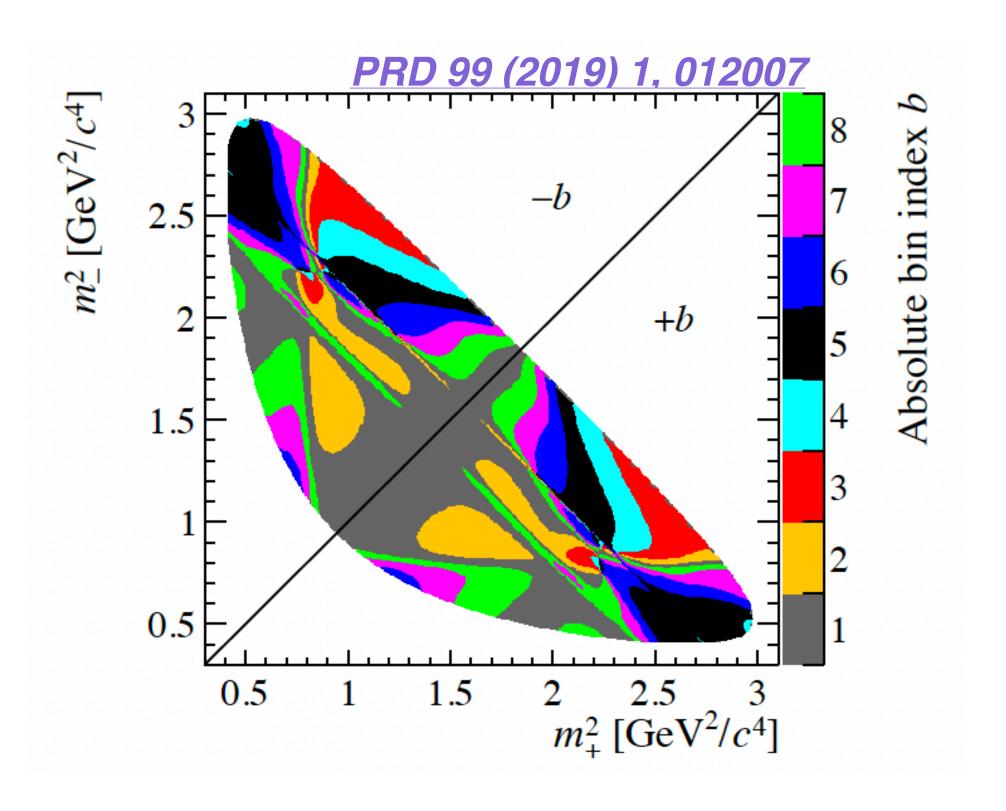
- □ $D^0 \to K_S^0 \pi^+ \pi^-$ allows to calculate mixing and CP violating parameters
- This decay receives contributions from the CF and DCS decays:

$$m_{\pm}^{2} \equiv \begin{cases} m^{2}(K_{S}^{0}\pi^{\pm}) & \text{for } D^{0} \to K_{S}^{0}\pi^{\pm}\pi^{\mp} \\ m^{2}(K_{S}^{0}\pi^{\mp}) & \text{for } \bar{D}^{0} \to K_{S}^{0}\pi^{\mp}\pi^{\pm} \end{cases}$$



Bin-flip method (II)

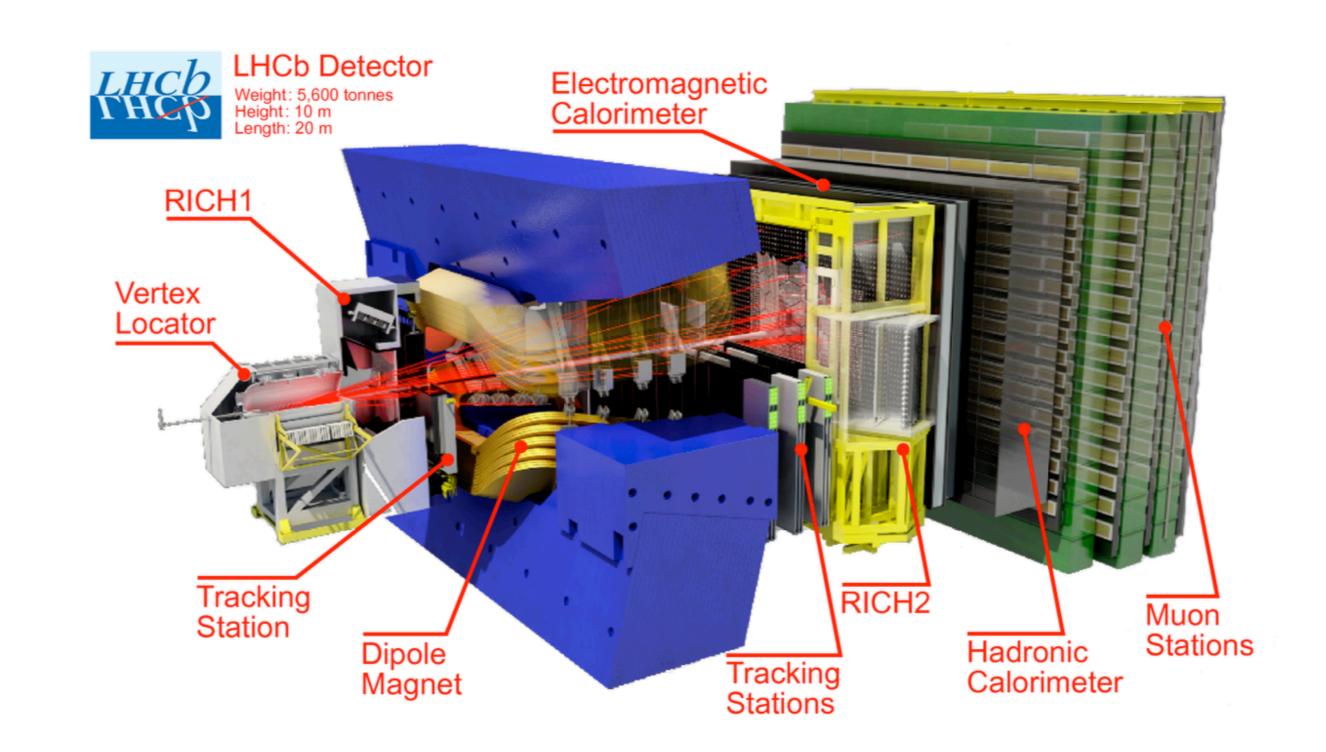
- Model-independent method
- Data are partitioned into bins of the Dalitz plot
 - defined to preserve nearly constant strong phase differences $\delta(m_-^2, m_+^2)$ between the D^0 and \bar{D}^0 amplitudes within each bin: subscript b
- Dalitz plane divided into two regions:
 - $m_{+}^{2} > m_{-}^{2}$: contribution from CF: +b
 - □ $m_{+}^2 < m_{-}^2$: contribution from DCS: -b
- At first order, it is not affected by uncertainties due to time acceptance effects



LHCb detector

LHCb (Ref 1, Ref 2)

- LHC:pp collision
- A single-arm forward spectrometer designed for the study of particles containing b or c quarks
- □ Pseudorapidity range: $2 < \eta < 5$
- Highly efficient PID
- Very good momentum and IP resolution
- Excellent primary and secondary vertex reconstruction



HLT1 Selection

□ LHCb Run3 ⇒ removal of hardware trigger for the HLT1 software one

One Track:

- Uses one pion for selection
- MVA cuts based on PT, IPChi2

Two Tracks:

- Uses the two pions for selection
- MVA cuts based on PT(hh) and IPCHI2(hh)

D2Kshh:

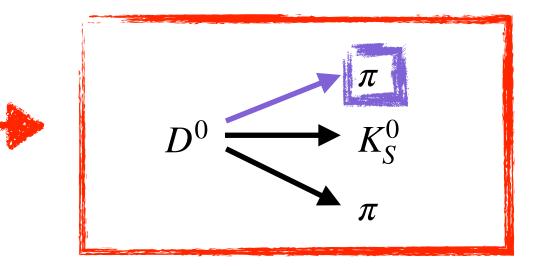
- Uses the 4 tracks
- ^{\square} Selection on the masses of K_S^0 and D^0 , PT and IP

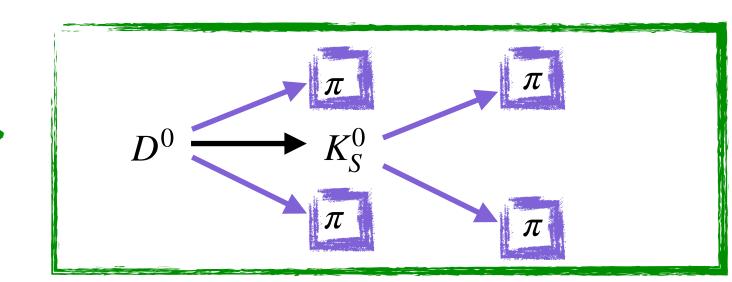
KS

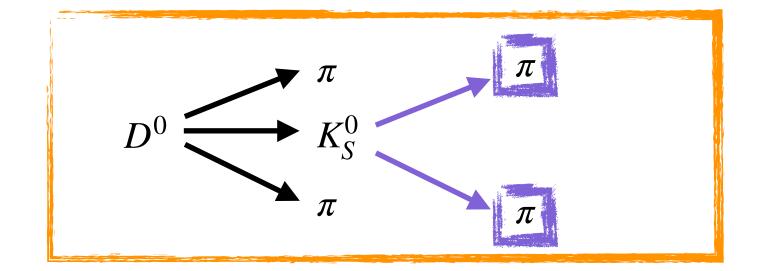
- Uses the K_S^0 daughters
- □ Selection on the mass of K_S^0 , PT and IP

TIS:

No selected track participates in the HLT1 selections





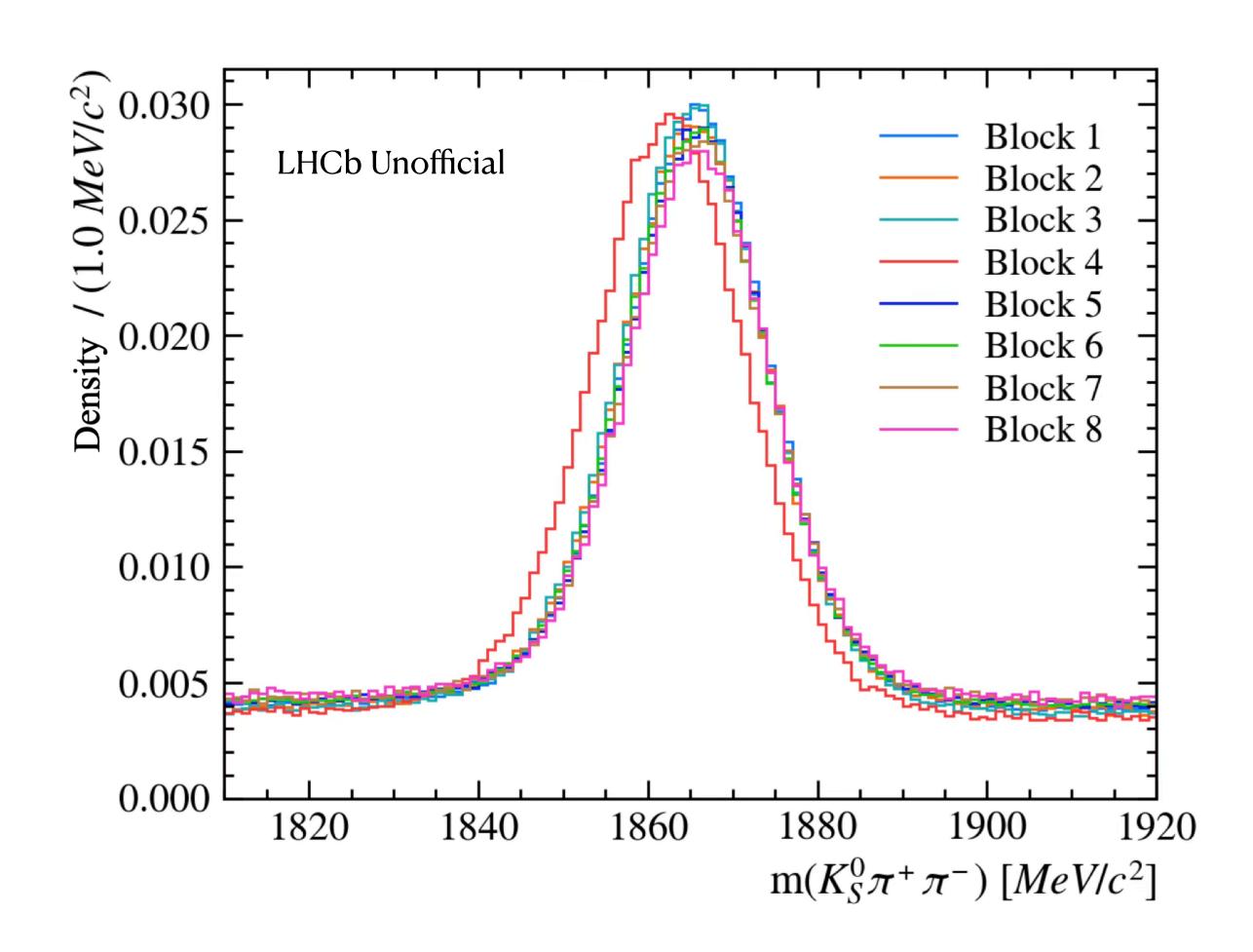


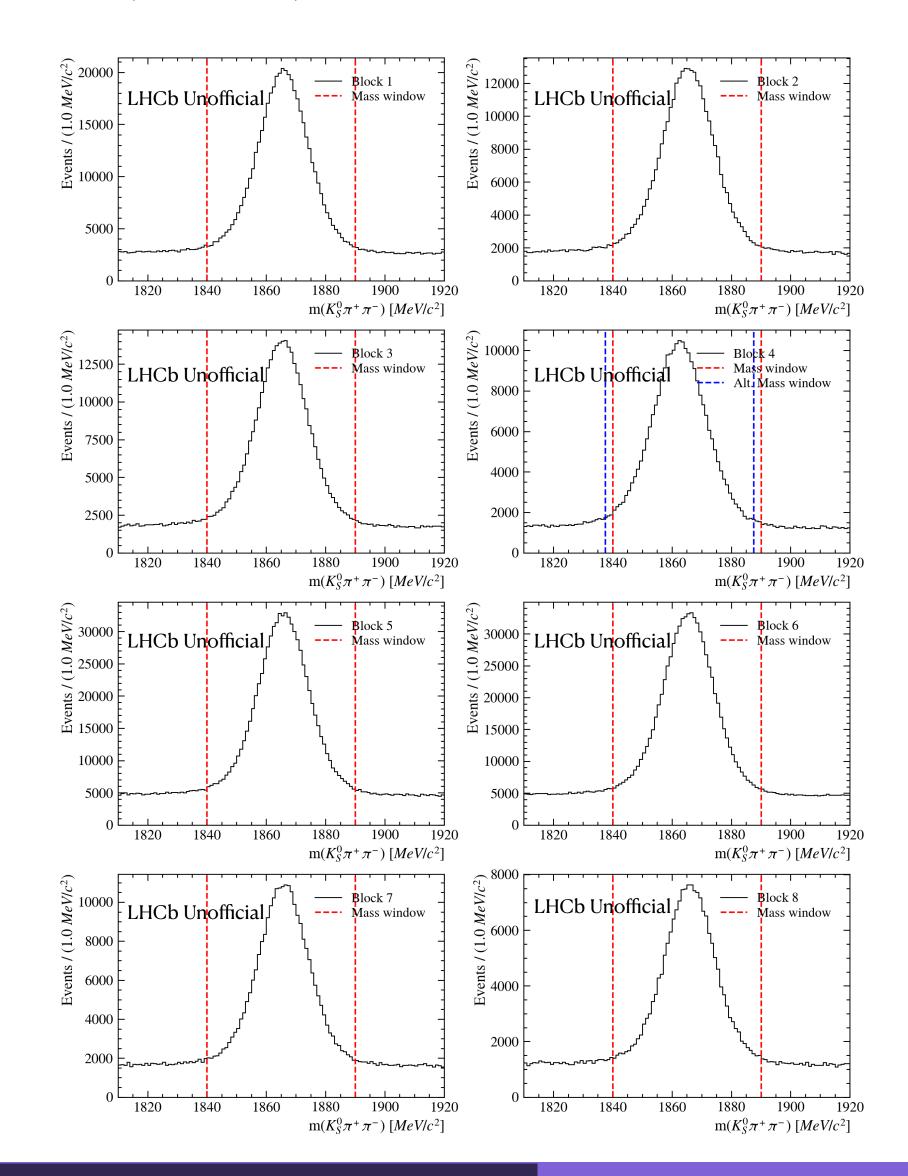
Dataset & Selection

- Data
 - ^o 2024 data mag up and down (9/fb)
 - Both long (LL) and downstream (DD) tracks
 - Data divided by Blocks
- Selection:
 - $p_T(D^0)$ different for each trigger line
 - $\log(\chi_{IP}^2(D^0)) < 2.2$ to remove secondaries
 - $\Box DOCA(\pi\pi) < 0.14 \,\mathrm{mm}$
 - $z_{V_{tx}}(K_S^0) z_{V_{tx}}(D^0) > 10 \,\mathrm{mm}$
 - m_{D^0} in [1840,1890] MeV/c²
 - □ Introduced a shift of 2.5 MeV in block 4 due to data reconstructed with preliminary alignment \Rightarrow will be updated
 - Different modes to select those candidates at trigger level
 - offer different acceptances

Hlt	P _T (D ⁰) >
OneTrack	4000
TwoTrack	2500
TIS	3750
D2Kshh	2500
KS	4500

Selection: m(D⁰)



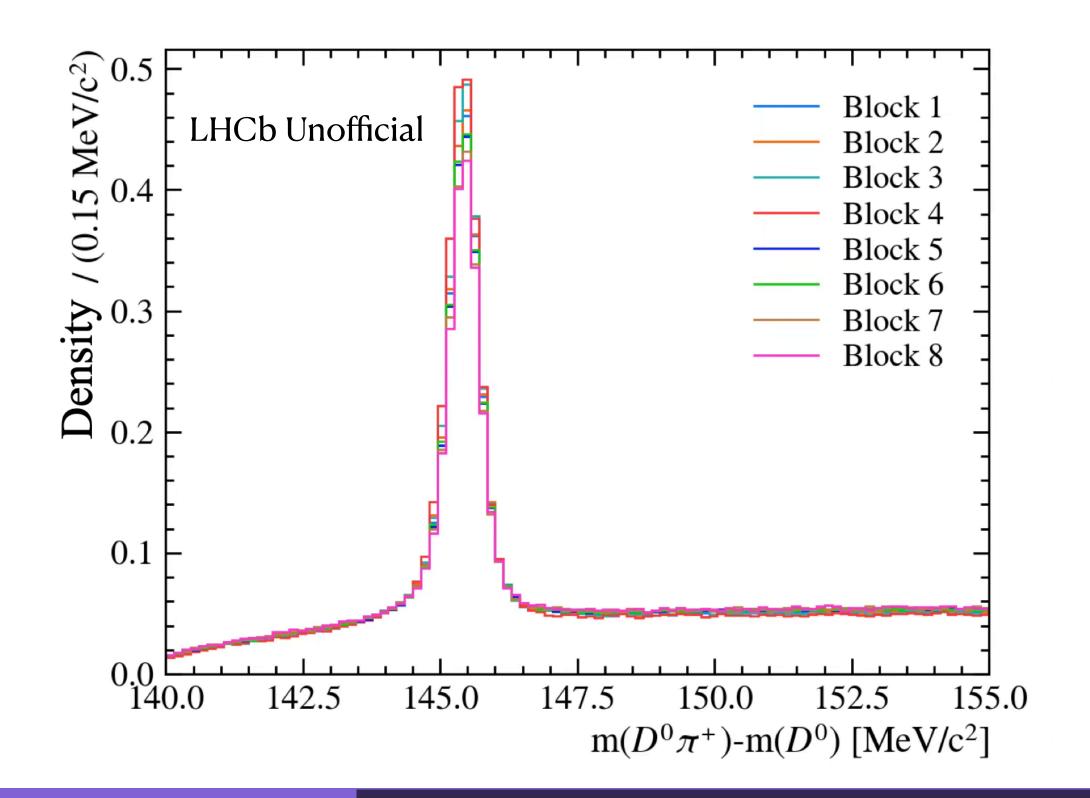


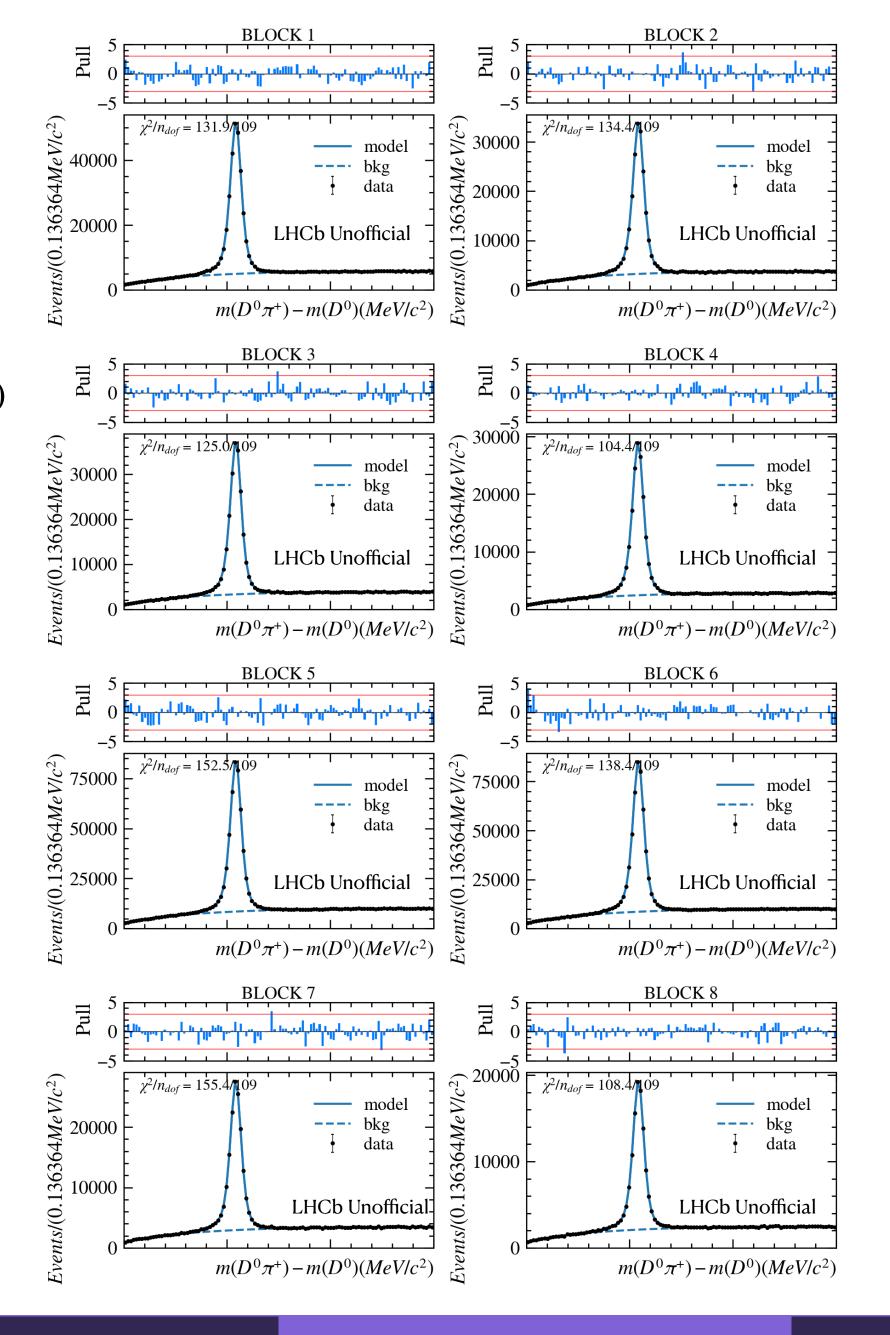
Signal extraction: Am

- To fit the Δm distribution:
 - for signal: a Johson SU + Gaussian
 - for background: Threshold function

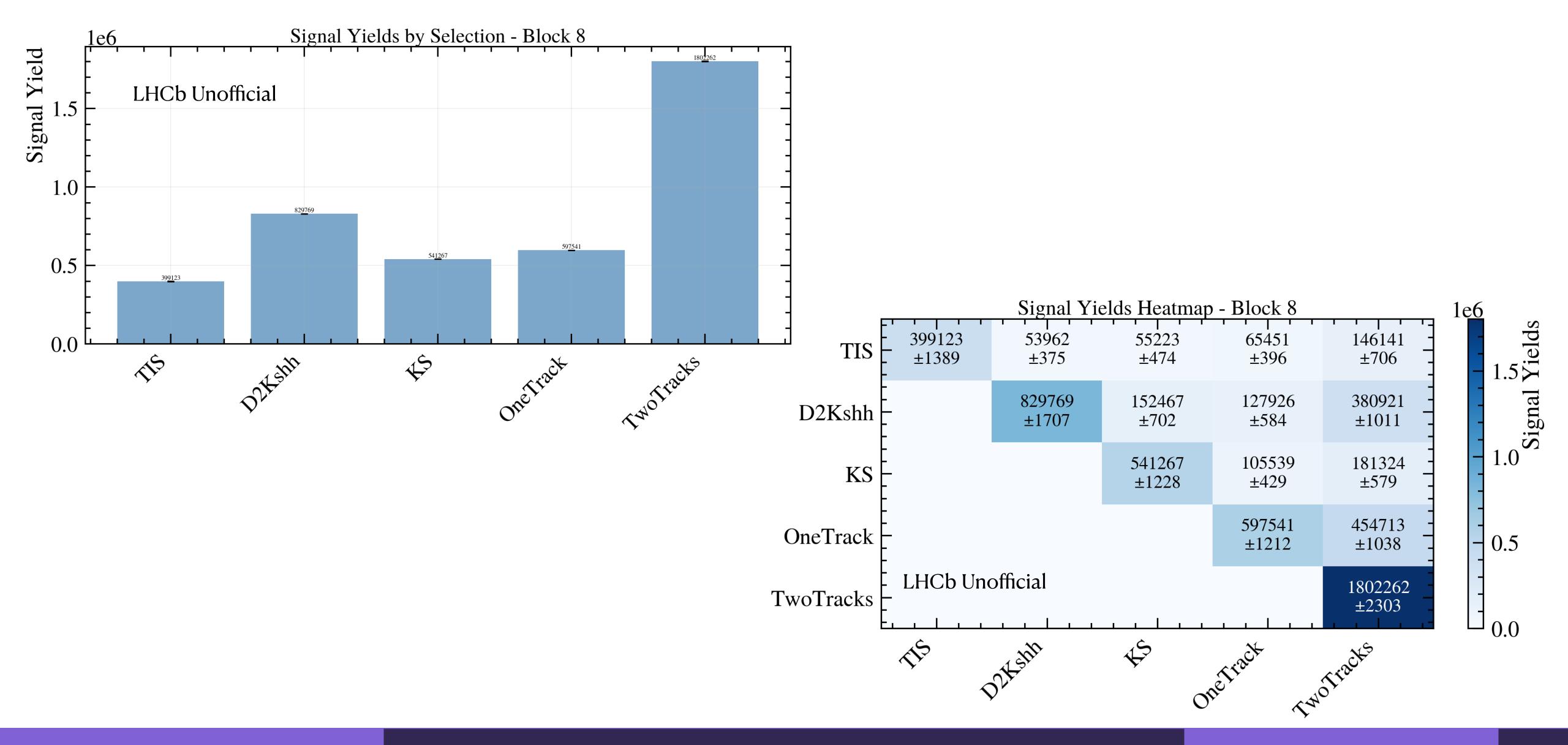
$$DmModel(x) = f_{sig} \left(f_{JSU} \cdot JSU(x) + (1 - f_{JSU}) \cdot \left(w_1 \cdot G_1(x) \right) \right) + (1 - f_{sig}) \cdot RooThresholdDAcp(x; \theta, a, b)$$

$$with RooThresholdDAcp(x; \theta, a, b) = \Theta(x - x_{th})(x - x_0)^a e^{-b(x - x_0)}$$

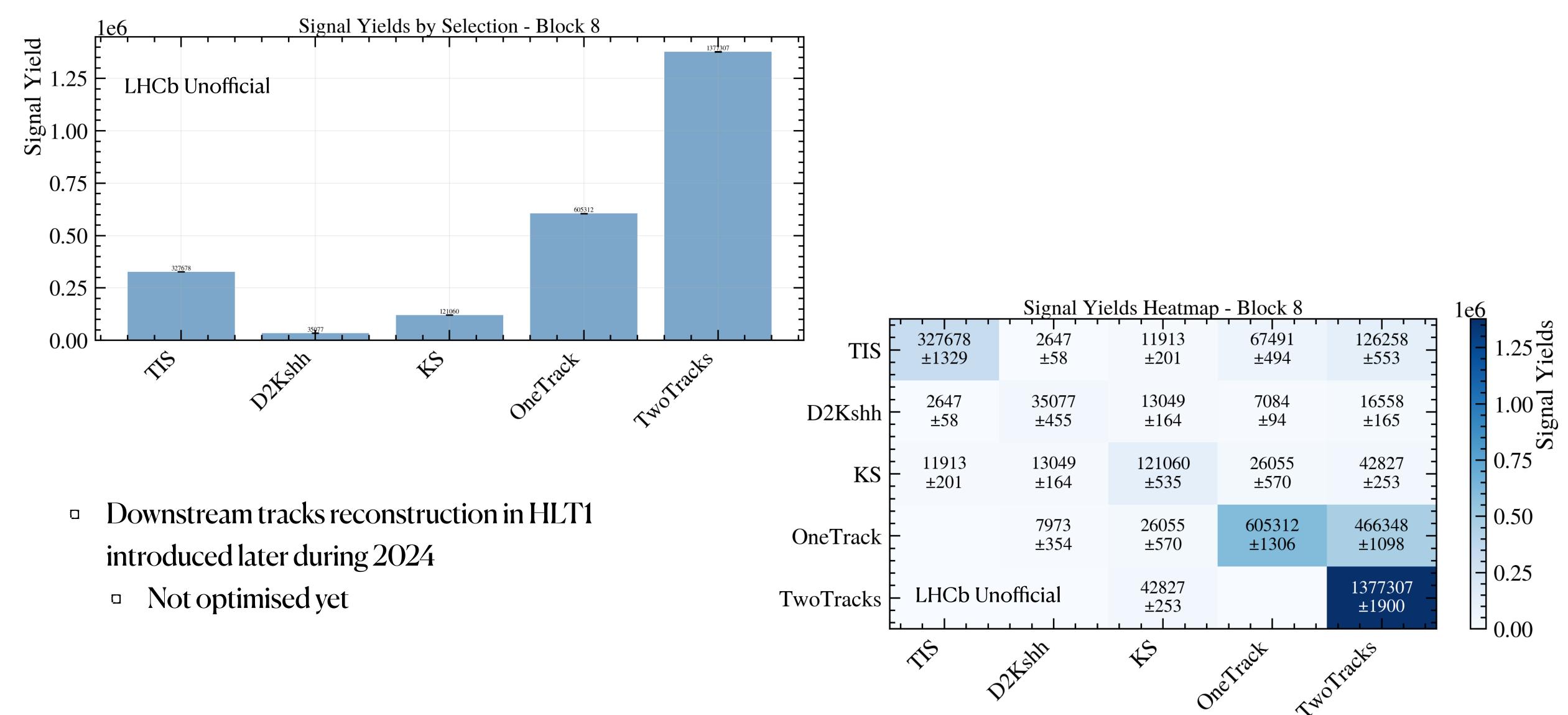




Signal Yields by Selection - LL



Signal Yields by Selection - DD



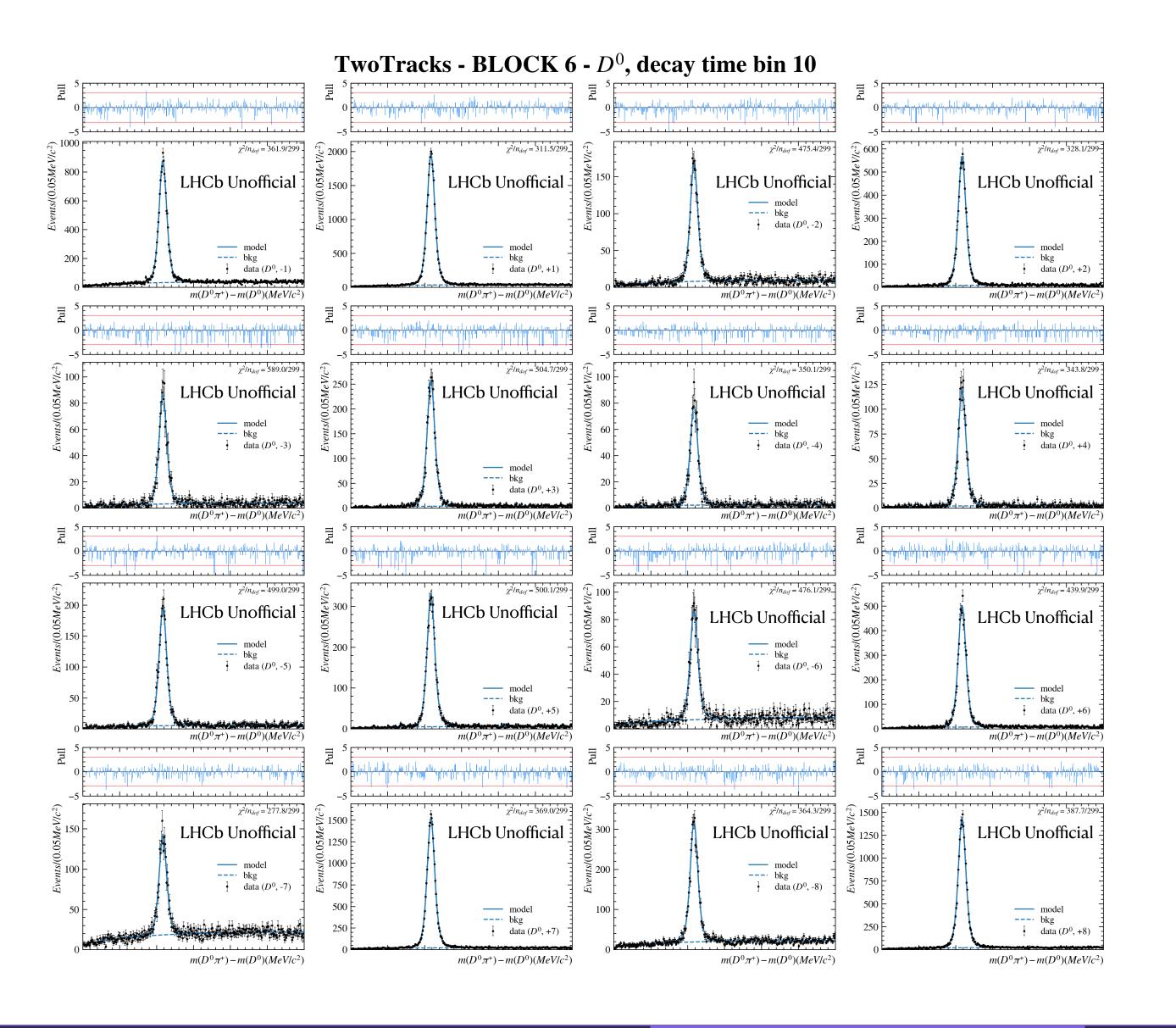
Run 3 2024 data divided per dataset

"Yields obtained through a fit of the Δm variable ($\Delta m = m(D^{*+}) - m(D^{0})$)

	Lumi (MU+MD)	Yield/lumi		
Data	[pb ⁻¹]	LL	DD	
Run2	5400	1826.4 ± 1.2	3837.31 ± 1.66	
Block 4 [9808, 9910]	1131.8	4912.04 ± 3.04	2854.37 ± 2.24	
Block 3 [9911, 9943]	670.4	5509.18 ± 4.32	3377.92 ± 3.27	
Block 2 [9945, 9978]	601	6510.27 ± 4.94	3919.64 ± 3.74	
Block 1 [9982, 10056]	1161.1	6407.83 ± 3.60	3823.57 ± 2.61	
Block 5 [10059, 10102]	1125.8	6651.64 ± 3.72	4797.59 ± 3.02	
Block 6 [10104, 101381]	949.1	6537.33 ± 4.06	4674.80 ± 3.18	
Block 7 [10197, 10213]	717.6	5810.06 ± 4.28	4019.62 ± 3.42	
Block 8 [10214, 10232]	417.6	5787.06 ± 5.65	3979.94 ± 4.45	
Run3	6774.4	6020.6 ± 1.4	3913.7 ± 1.1	

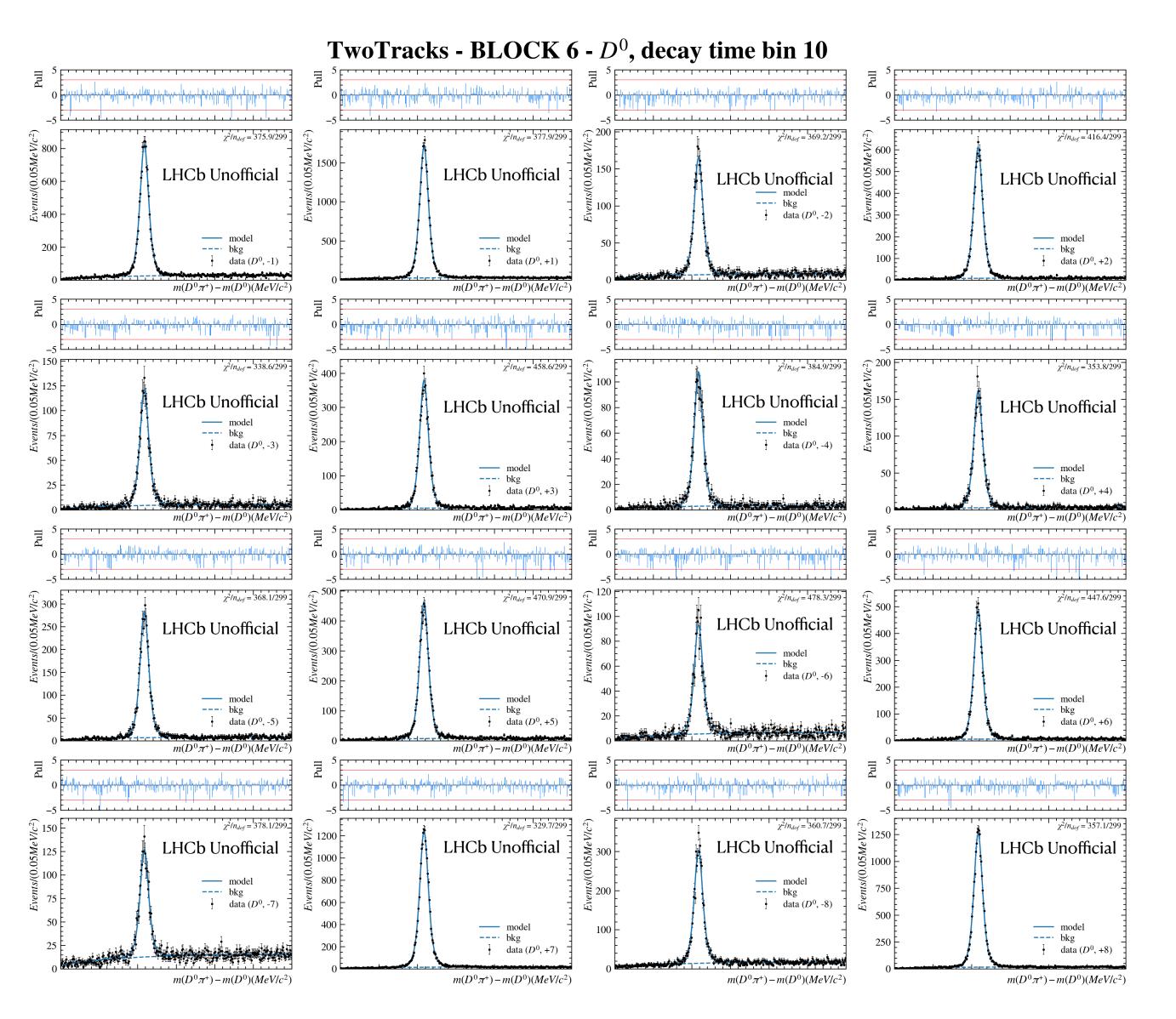
Simultaneous fits-LL

- Simultaneous fits in phase-space bins per lifetime bins
 - ^{\square} To directly extract from the fit the R_b ratios in each bin



Simultaneous fits-DD

- Simultaneous fits in phase-space bins per lifetime bins
 - ^{\square} To directly extract from the fit the R_b ratios in each bin



Efficiency Maps - Two Tracks

Initial Problem

- Correct data by removing experimentally induced correlations between decay time and phase-space coordinates
- Full correction would require knowledge of the 3D (phase-space + decay time)
 efficiency variation

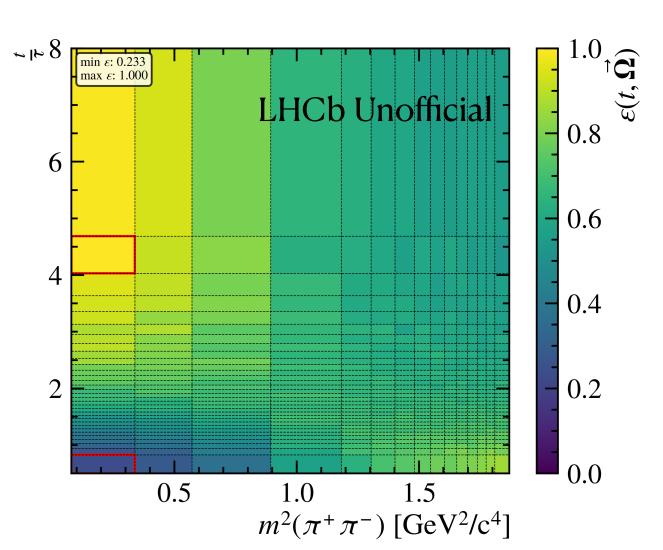
Adopted Approach

- No need to correct for overall efficiency variation
 - bin-flip method handles independent variations of decay time and phase-space
- Sufficient to make decay-time acceptance uniform across the phase space

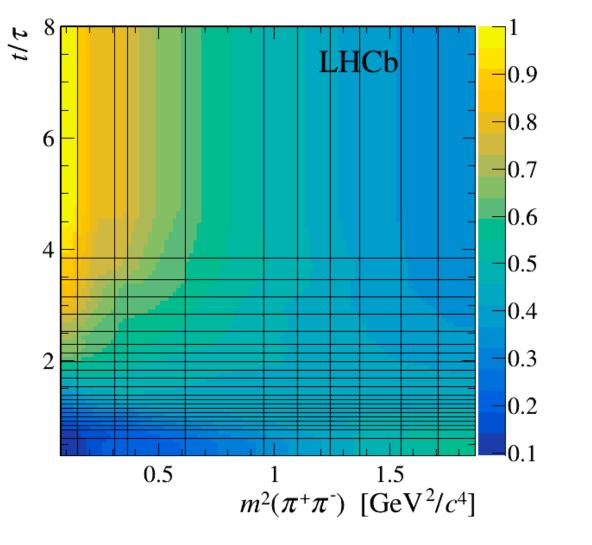
Estimation of Relative Efficiency

- Efficiency estimated directly from data using background subtraction via sWeights
- Decay-time differences across phase space attributed only to efficiency variations
- $^{\square}$ D^0 and \overline{D}^0 without charge conjugation, assuming negligible direct CPV in bin-flip ratios
- \square Run 3 maps have higher minimum values \Longrightarrow smaller efficiency variation

Run3 Block 8 - LL



Run2



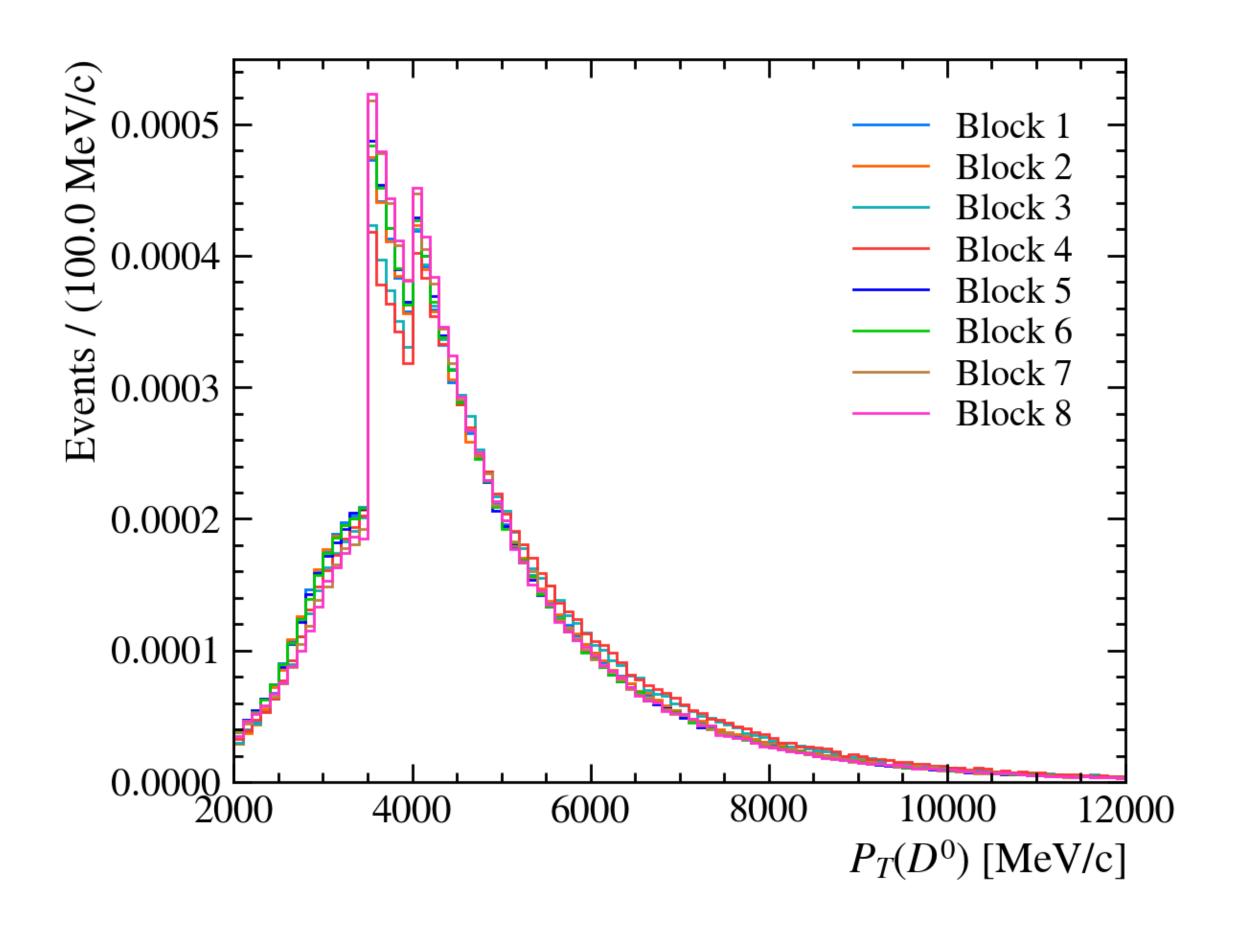
Conclusions

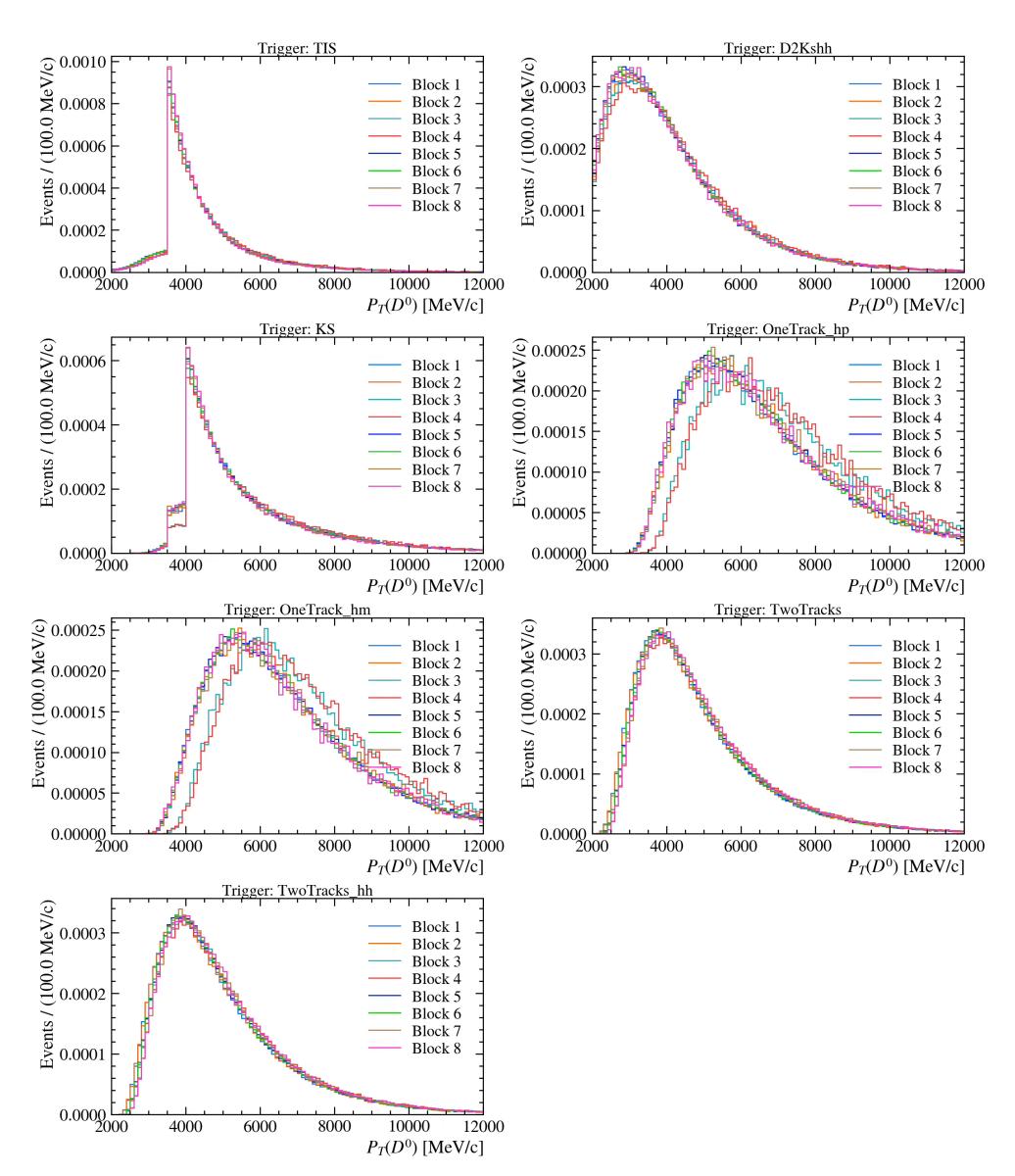
- □ With OneTrack and TwoTracks datasets \Rightarrow threefold increase in signal yields in Run 3 2024 with respect to Run 2
- Aim for publication by 2026 with the 2024 dataset

Thank you for your attention

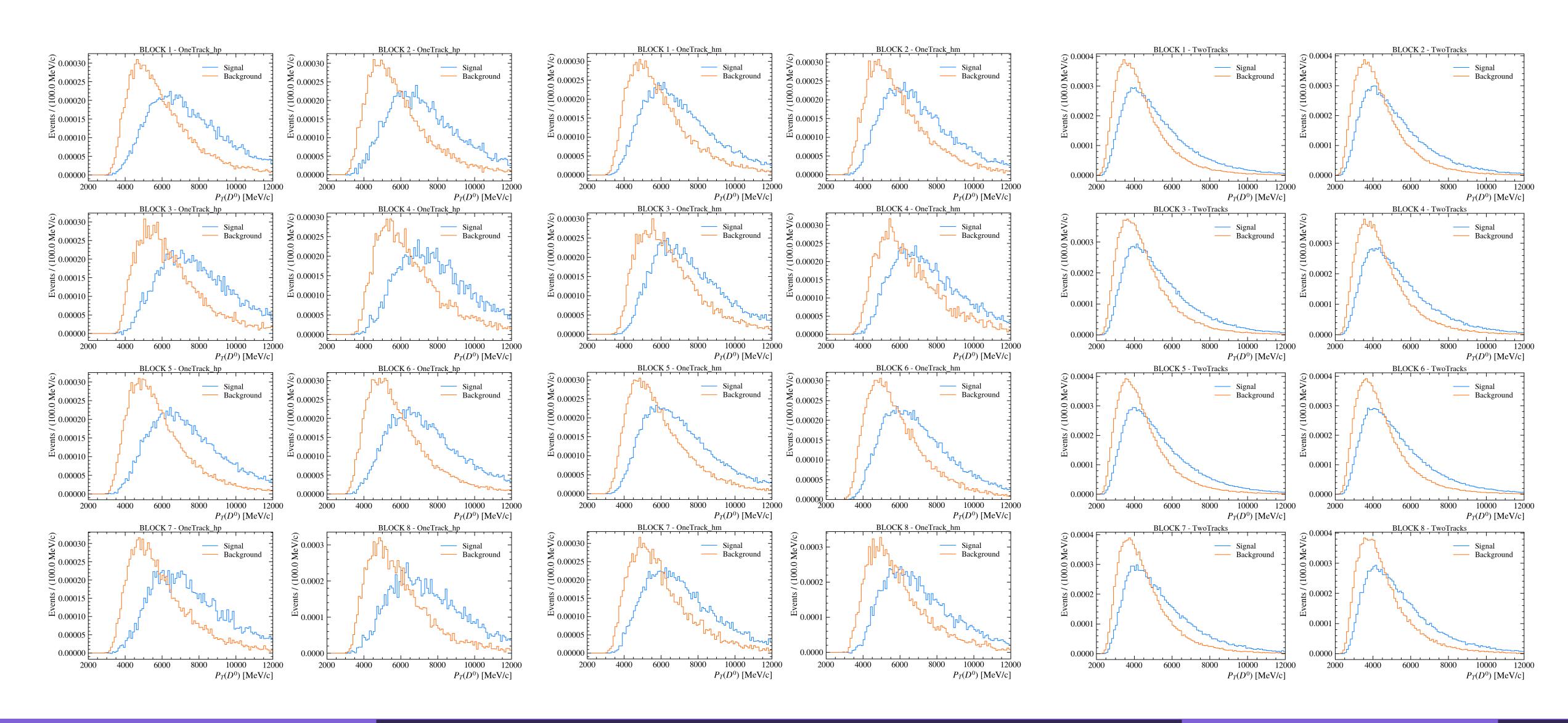
Backup



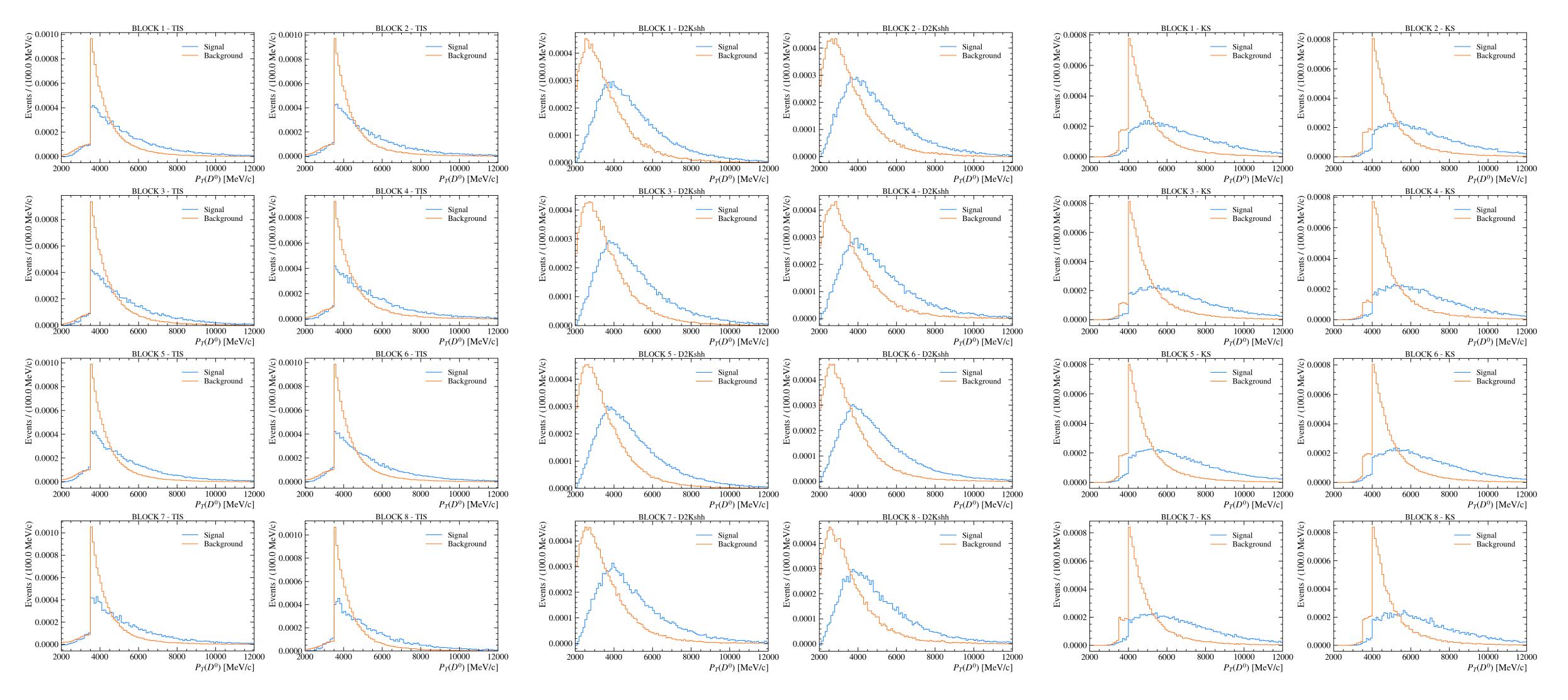




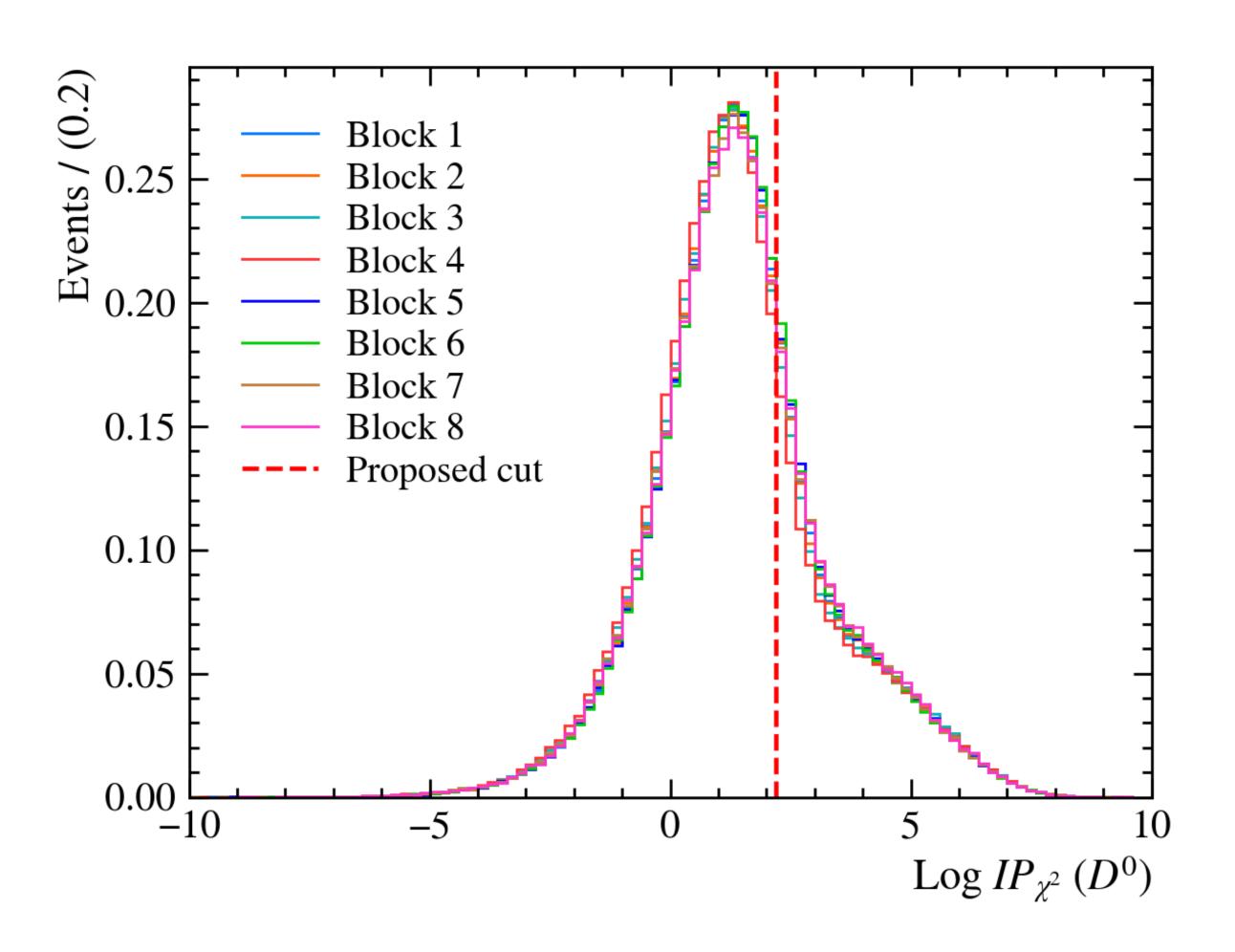


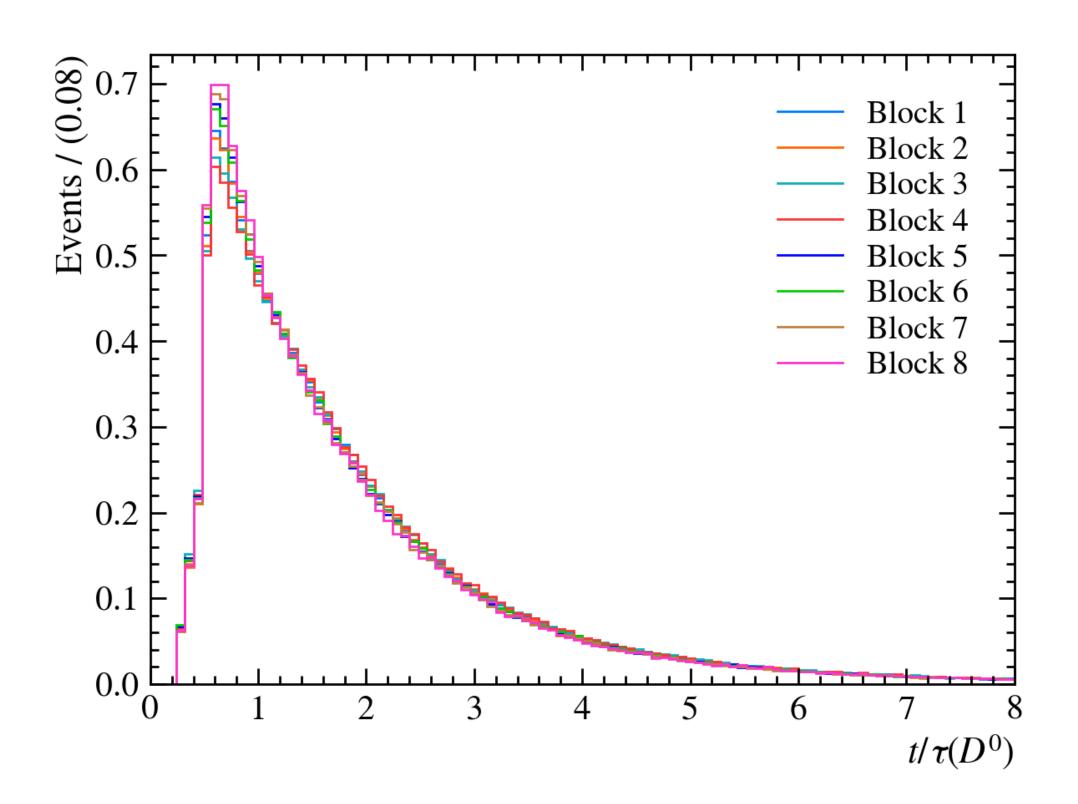


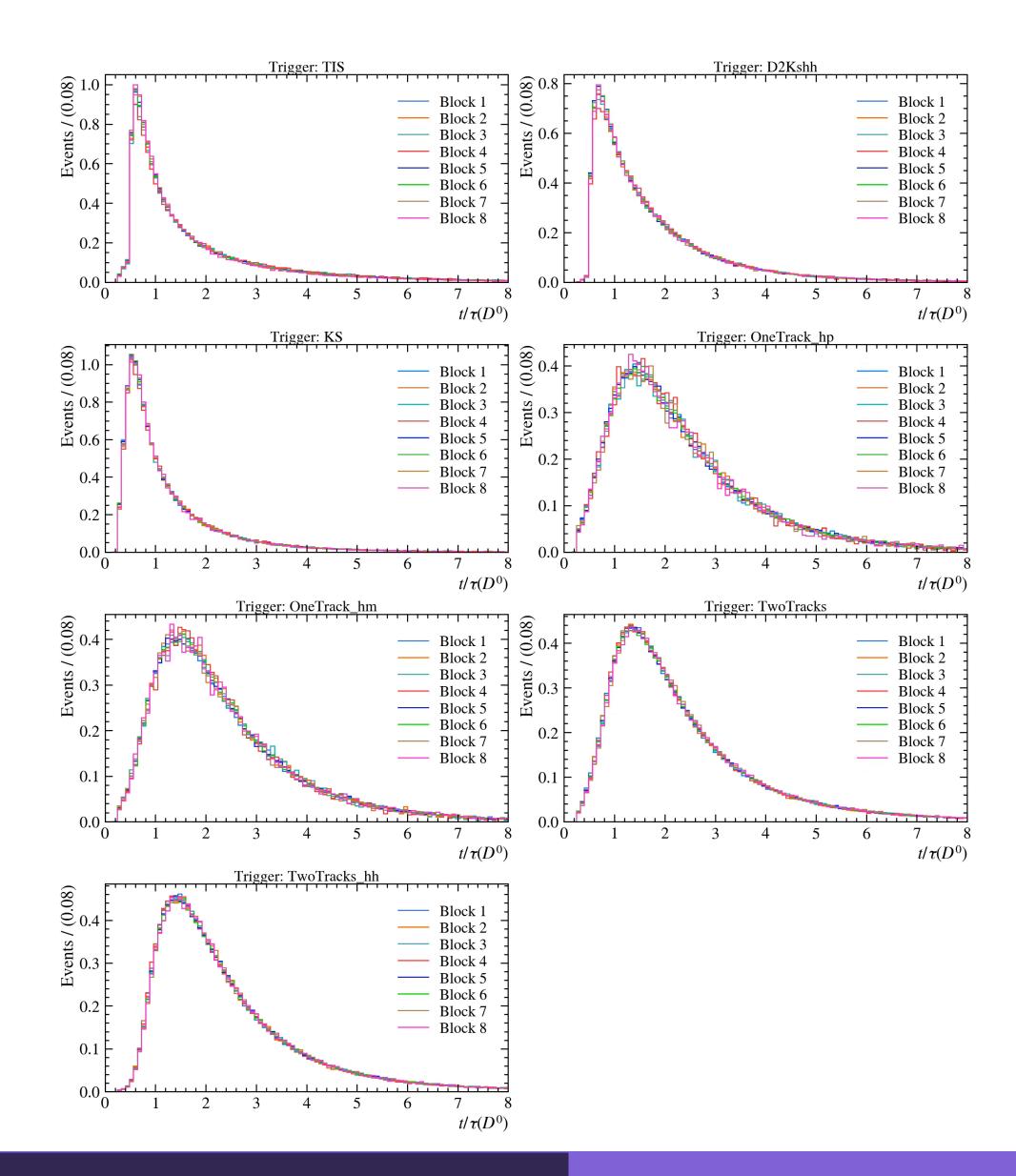




$log(IP\chi^2(D^0))$







Efficiency Maps (II)

Practical Implementation

- Data binned in $(\overrightarrow{\Omega}, t)$ with variable bin widths
- Finer bins at low decay times to capture rapid efficiency changes

$$\epsilon_{ijk} = \frac{\frac{n_{ijk}}{\sum_{l} n_{ijl}}}{\frac{\sum_{l,m} n_{lmk}}{\sum_{l,m,n} n_{lmn}}}$$

- \bar{a} i, j, k bins of $m_{\pi^+\pi^-}^2$, $|\cos\theta_{\pi^+\pi^-}|$, decay-time
- Then corrected to that accounts for mixing, assuming specific values of y and strong phases $\delta: \varepsilon_{y,\delta}(t \mid \overrightarrow{\Omega})$

Efficiency Maps (III)

- To properly account for mixing, it is necessary to remove the effect of mixing when building the weights
- This is achieved by correcting the data using the ratio: $\frac{\varepsilon_{\text{data}}(t|\overrightarrow{\Omega})}{f_{y,\delta}(t|\overrightarrow{\Omega})}$
- Repeating the original calculation $\epsilon_{ijk} = \frac{\frac{n_{ijk}}{\sum_{l,m} n_{lmk}}}{\frac{\sum_{l,m} n_{lmk}}{\sum_{l,m} n_{lmn}}}$ with entries weighted by: $\frac{1}{f_{y,\delta}(t \mid \overrightarrow{\Omega})}$
- The result is a correction function that accounts for mixing, assuming specific values of y and strong phases δ : $\varepsilon_{y,\delta}(t\mid \overrightarrow{\Omega})$
- These $\varepsilon_{y,\delta}(t\mid \overrightarrow{\Omega})$ are used to construct correction factors for the ratios of fitted yields in each Dalitz and decay-time bin

2025 vs 2024

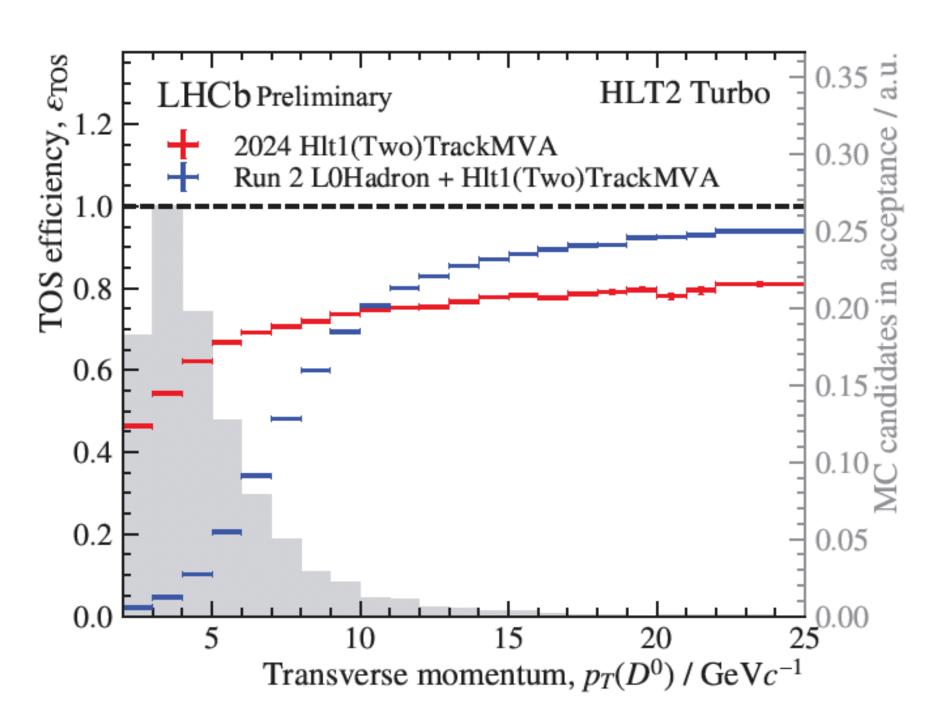
^o Comparing 2025 data from s25c1 with 2024 data from s24c4, since they have the same $\mu \sim 5.3$

Data	Lumi (pb ⁻¹)		OneTrackMVA	TwoTracksMVA	D2Kshh	KS	TIS
	s25c1	~ ~	827254 ± 1405	1865682 ± 2072	855762 ± 1474	662107 ± 1304	461950±1365
	359	9 LL	1523058±1931	3660205 ± 2810	1510928±1834	1236102 ± 1765	897573 ± 1925
MD	s24c4	DD	1056735±1400	2182209 ± 2135	22569 ± 198	17682±196	540685 ± 976
	717.7	DD	1171289±1500	2606830 ± 2337	34060 ± 211	114126 ± 538	644636±1533
	s25c1 658	~ ~	1532064±1990	3446050 ± 2860	1541683±1910	1249487 ± 1781	891057 ± 2008
		658 LL	878312 ± 1403	2121200 ± 2173	868819 ± 1362	713746±1335	526476 ± 1478
MU	s24c4 417.6	DD	1939167 ± 1967	3985785 ± 2863	38150 ± 253	31162 ± 261	1054181 ± 2073
		DD	673808±1168	1500036±1752	20090 ± 215	64553 ± 481	373641 ± 1268

Run 3 LHCb Prospects

LHCB-FIGURE-2024-030

- Combined HLT1 (Two) TrackMVA TOS efficiencies in D decays to light hadrons, computed in bins of B transverse momentum in 2024 data
- For comparison, the combined TOS efficiencies of L0Hadron and HLT1TrackMVA/
 HLT1TwoTrackMVA calculated in Run 2 (from 2019 JINST 14 PO4013)
- ^{\square} Generator-level distributions of $p_T(D)$ for MC events within the LHCb acceptance in grey
- At lower $p_T(D)$ efficiencies in 2024 data are significantly higher than those of Run 2, whilst at higher $p_T(D)$ the reverse is true
- The latter is driven by differences in tracking algorithms between HLT1 and HLT2 in 2024, wherein higher proportions of tracks found in HLT2 are not found in HLT1 at higher $p_T(D)$
- In Run1+Run2, efficiency at low momentum was gained by triggering on tracks not belonging to the signal candidate



TOS efficiencies in $D^0 \to K^-\pi^+$