Production at SPES of a ⁹⁴Nb radioactive source for studying the quenching of the weak interaction axial-vector coupling constant

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The new cyclotron at LNL

International Workshop on future research program with the high power cyclotron of SPES-LNL

Call for Expression of Interest



Challenges in beta decays spectral shapes

- Recent theoretical studies have highlighted the sensitivity of beta decay parameters in complex decays
- ✓ A pertinent example is the study of the axial-vector coupling constant g_A in nuclear decays, which is expected to be different from that extracted from neutron decay $\frac{g_A}{g_V}$ =

-1.27641(45)stat(33)sys

- ✓ For allowed beta decays, namely Fermi ($\Delta l = 0$; $\Delta s = 0$) and Gamow-Teller ($\Delta l = 0$; $\Delta s = 1$) g_A controls the half-life of the decay, being the spectral shape unaffected. Typically, the weak interaction operators need to be renormalized, by quenching the g_A , to account for missing physics in many body calculations
- ✓ For forbidden ($\Delta l > 0$) beta decays also the spectral shape depends on g_A , especially for highly forbidden non-unique decays



M. Haaranen et al., Phys. Rev. C 93, 034308 (2016)

Challenges in beta decays spectral shapes

- ✓ Recent realistic shell model calculations have emphasized a challenging situation for a simultaneous description of the decay half-life and spectral shape for highly forbidden non unique decays
- ✓ Renormalizing the g_A to get an accurate description of the half-life spoils the theoretical description of the measured spectra
- ✓ This topic is quite hot since g_A enters with its 4th power in the expression of the decay rate for $0\nu\beta\beta$ decay. A factor of 2 of "quenching" would reduce the rate of factor 16!!!
 - ✓ Experimentally the challenge is to provide energy spectra, free from distortions due to background, down to very low energy ($T_e < 100 \ keV$), where the sensitivity to g_A is the highest



G. De Gregorio et al., PRC 110, 014324 (2024)

The ASPECT-BET project

This proposal is part of the PRIN-2022 ASPECT-BET (An sdd-SPECTrometer for BETa decay studies) project [A. Nava et al., Sensors 24 (2024) 8202]

We have developed an innovative and versatile spectrometer to enable state of the art measurements of beta energy spectra, almost background free

The spectrometer uses Silicon Drift Detectors (SDD) technology to detect electrons in the 10 keV - 1 MeV range, a crucial energy region for beta decay studies









The ASPECT-BET project

The use of forward and backward **VETO detectors allows** to suppress the background

The careful characterization of the radioactive source allows to correct for tricky distortion effects due to selfabsorption of slowly moving electrons





⁹⁴Nb beta decay

⁹⁴Nb is unstable against β^- decay to ⁹⁴Mo



92Mo	93Mo	94Mo	95Mo	96Mo	97Mo	98Mo
STABLE	4000 y	STABLE	STABLE	STABLE	STABLE	STABLE
14.649%	ε=100%	9.187%	15.873%	16.673%	9.582%	24.292%
91Nb	92Nb	93Nb	94Nb	95Nb	96Nb	97Nb
680 y	3.47e+7 y	Stable	2.04e+4 y	34.991 d	23.35 h	72.1 min
ε+β+=100%	ε+β+=100%	100%	β~=100%	β~=100%	β~=100%	β=100%
90Zr STABLE 51.45%	912 r STABLE 11.22%	92Zr STABLE 17.15%	93Zr 1.6le+6 y β=100%	94Zr STABLE 17.38%	952r 64.032 d β=100%	96Zr 2.29e+19 y 2.8% 2β=100%
89Y	90Y	91Υ	92Y	93Y	94Y	95Y
STABLE	64.046 h	58.56 d	3.54 h	10.17 h	18.7 min	10.4 min
100%	β=100%	β=100%	β=100%	β=100%	β~=100%	β~=100%
885r	89Sr	90Sr	91Sr	92Sr	93Sr	94Sr
STABLE	50.56 d	28.91 y	9.68 h	2.61 h	7.43 min	75.3 s
82.58%	β=100%	β=100%	β=100%	β=100%	β=100%	β ⁻ =100%

The decay proceeds almost fully from the 6⁺ ground state of ⁹⁴Nb to the excited 4⁺ state of ⁹⁴Mo, thus resulting in a non unique angular momentum transfer dominated by $\Delta J = 2^+$ (second forbidden decay)

The measured half life is $T_{1/2} = 2.04 * 10^4 y$. The end-point energy is 471.5 keV.

Procurement of the ⁹⁴Nb isotope

⁹⁴Nb is not easy to find in nature as it is radioactive and it is not at all commercially provided so it should be produced in laboratory

- ✓ The possibility to produce ⁹⁴Nb from neutron irradiation of the natural isotope ³⁹Nb has been halted due to the very low cross section for neutron capture reaction, with a concern about the self absorption of the resulting thick source
- ✓ ⁹⁴Nb is not present in the REX-ISOLDE or TRIUMF databases, so it cannot be just implanted on a target backing (<u>https://isoyields2.web.cern.ch/</u><u>https://yield.targets.triumf.ca/search/yield/data</u>)
- ✓ A new perspective is explored from the use of (p,n) charge exchange reaction to produce ⁹⁴Nb from the stable ⁹⁴Zr

94Mo	95Mo	96Mo
STABLE	STABLE	STABLE
9.187%	15.873%	16.673%
93Nb	94Nb	95Nb
STABLE	2.04e+4 y	34.991 d
100%	β~=100%	β ⁻ =100%
92Zr	932r	94Zr
STABLE	1.61e+6 y	STABLE
17.15%	β~=100%	17.38%

The ⁹⁴Zr(p,n)⁹⁴Nb reaction



Note that the (p,n) cross section is maximum ($\sigma(p,n) \sim 630 \text{ mb}$) at $E_p \sim 8 \text{ MeV}$ reducing to ($\sigma(p,n) \sim 30 \text{ mb}$) at $E_p \sim 35 \text{ MeV}$

Experiment strategy

✓ Due to the long half-life of the ⁹⁴Nb decay, $T_{1/2} = 6.4 * 10^{11} s$, the number of ⁹⁴Nb to be generated by (p,n) reaction should be at least of the order of 10^{12} to get source activity of the order of Bq

Assuming:

- Proton beam energy: 35 MeV
- Beam current: **300 uA**
- Target thickness: 0.2 mg/cm² (300 nm)
- Energy loss in the ZrO₂ target: ~ 2 keV
- (p,n) cross section: ~ 30 mb



One gets:

- Source activity after one day: $dN/dt \sim 6 Bq$
- Power dissipated in the ZrO_2 target: $\dot{Q} \sim 0.6 W$

The heat dissipation challenge

- ✓ Although ⁹⁴ZrO₂ features a high melting temperature $T_{fus} \sim 2680$ °C, it is a bad thermal conductor $k \sim 1.7 \frac{W}{m*K}$, making out target melting with $\dot{Q} \sim 0.6 W$ dissipated power
- An efficient heat dissipation can be guaranteed adopting the multilayer graphene (MLG) technology, developed by the NUMEN collaboration [F. Cappuzzello et al., Int. Jour. of Mod. Phys. A 36, 30 (2021) 2130018].
- ✓ The MLG are highly ordered thin stacks of Graphene layers used as the backing support where the ${}^{94}\text{ZrO}_2$ target is evaporated on. Thanks to the very high thermal conductivity of MLG ($k \sim 2200 \frac{W}{m*K}$), the heating of the ${}^{94}\text{ZrO}_2$ is maintained in the order of 300 K, thus safely below the melting point, if a MLG foil of 2000 nm is used.





Source purity for evaporation approach

- Since we are not using a pure ⁹⁴Zr target and other processes can compete with (p,n) charge exchange reaction we can in principle build a manyfold radioactive source whose decay pattern needs to be known with high accuracy.
- $\underline{\mathbf{N}}$
- First, we need to consider the role of Oxigen atoms from the ⁹⁴ZrO₂ oxide and the Carbon atoms from the MLG backing. An insight on this problem shows that Carbon and Oxygen impurities in the ⁹⁴Zr target are not an issue since the only radioactive nuclide with a long half-life (5700 y) is ¹⁴C. However, this isotope cannot be generated from C and can be hardly generated by O as it requires very suppressed processes such ¹⁶O(p, 3p)¹⁴C or ¹⁷O(p,3pn)¹⁴C or ¹⁸O(p,3p2n)¹⁴C with a negligible cross section.
- Then we need to consider the limited purity, in terms of 94 Zr, of commercial 94 ZrO₂, which is about 85-90%, the main contaminants being other Zr isotopes, namely 90 Zr, 91 Zr, 92 Zr and 96 Zr, with about 1-2% presence each one (depending on the producer). The main challenge comes from 91 Nb, which decays by electron capture with a half-life $T1/2=6.80*10^2 y$, thus potentially generating a 91 Nb activity comparable to the 94 Nb one. Nevertheless, the decay products from 91 Nb are Auger electrons with 2.02 keV and 13.4 keV, the latter of which may be visible just above the detector threshold (~10 keV).





Beam time request

- ✓ In conclusion, a promising technique to produce for the first time a ⁹⁴Nb radioactive source is at reach, using the (p,n) reaction with the very intense proton beam from LNL SPES facility
- \checkmark If successful, the same technique could be adopted to build other unique radioactive sources
- ✓ For the purpose of the ASPECT BET project, which is committed to measure the beta decay full spectrum with high energy resolution, low energy threshold and very low systematic uncertainties, a source activity of more than 20 Bq would be ideal, which in turn means three days irradiation time