

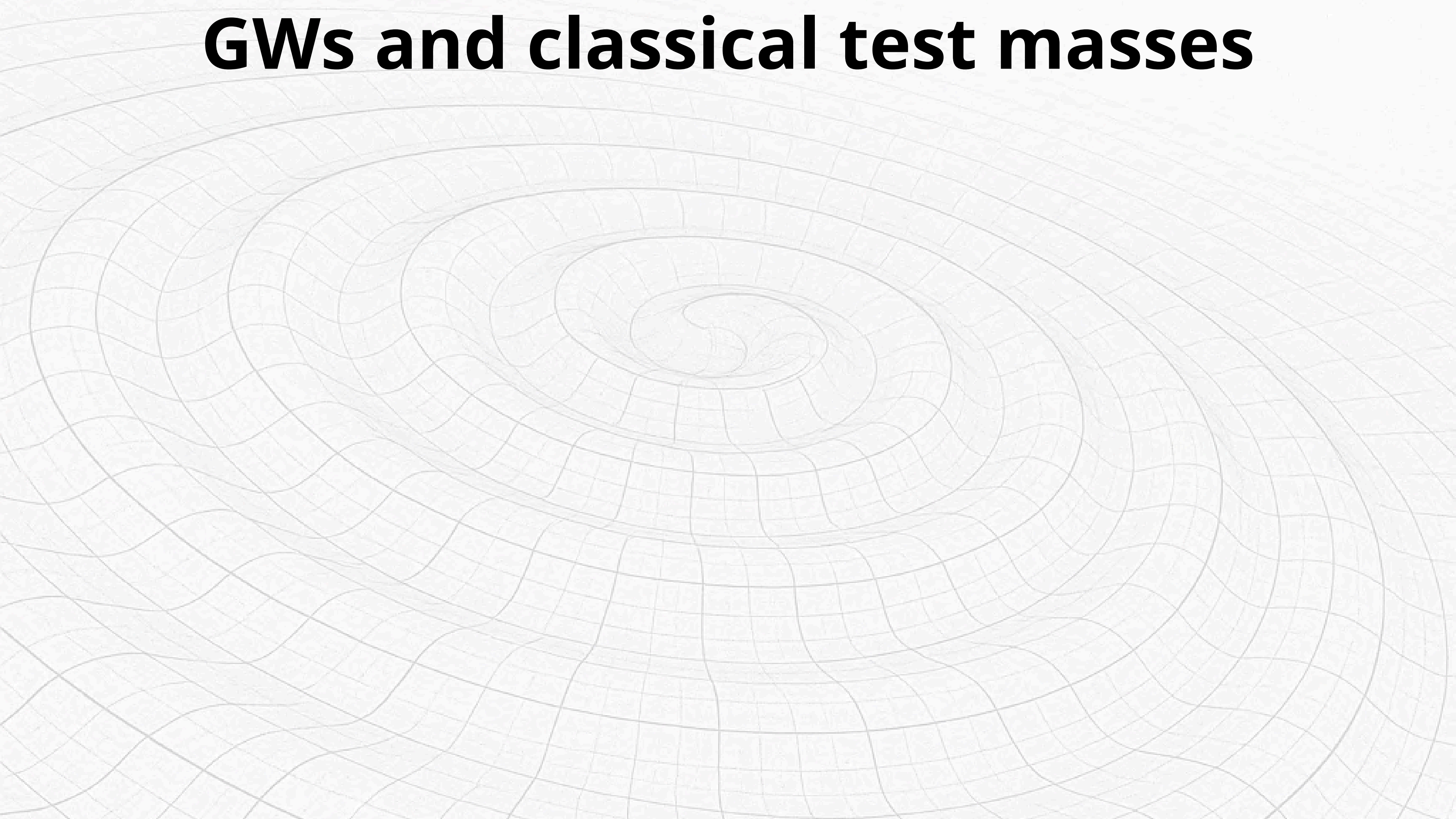
Gravitational wave imprints on spontaneous emission

Jerzy Paczos

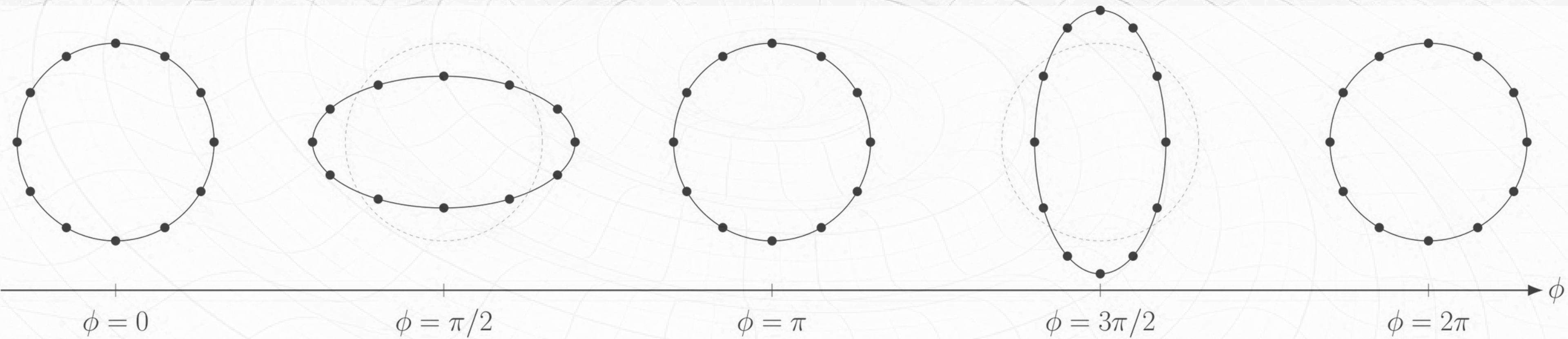
Navdeep Arya, Sofia Qvarfort, Daniel Braun, and Magdalena Zych

RQI North, Naples 2025

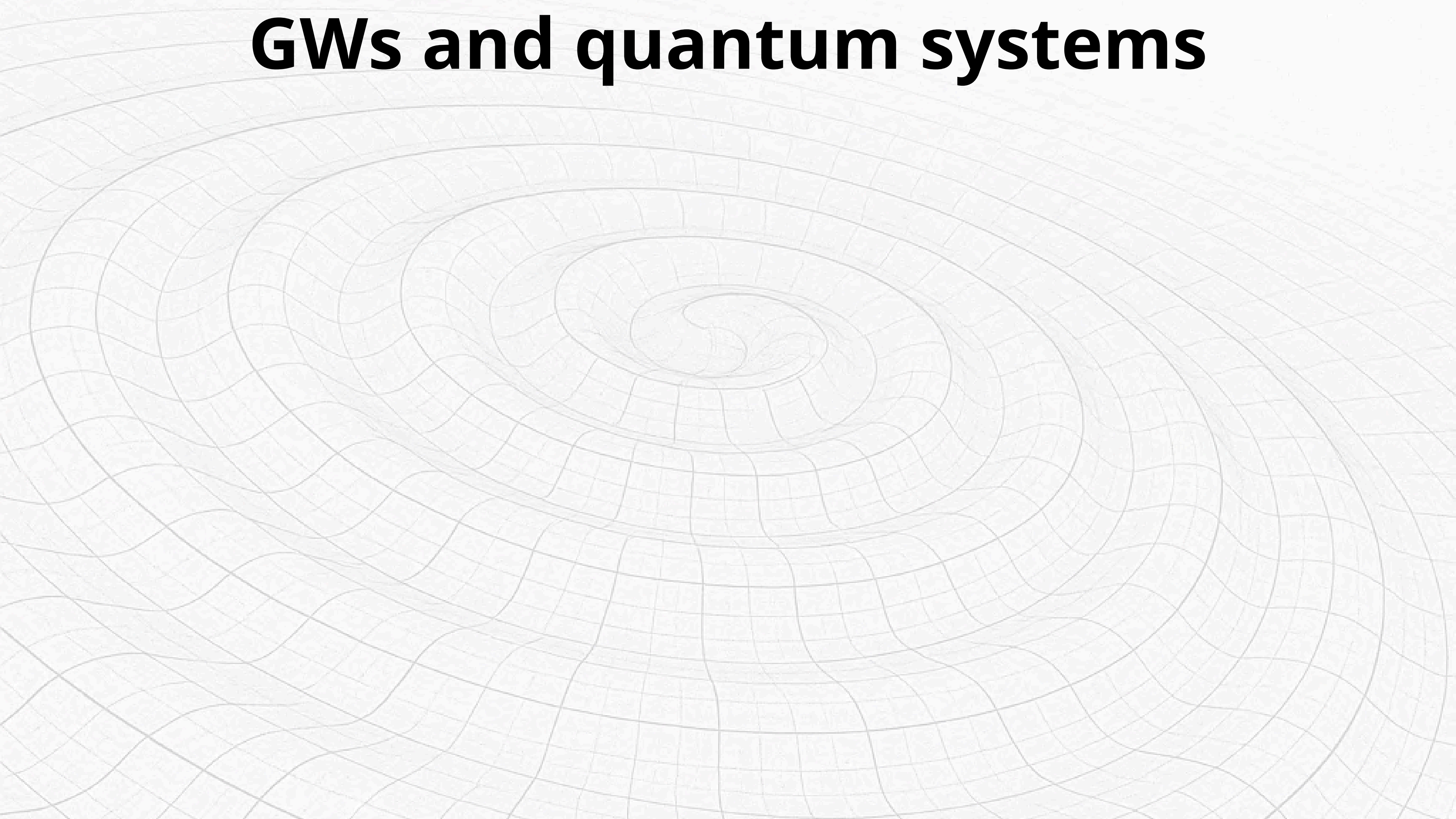
GWs and classical test masses



GWs and classical test masses

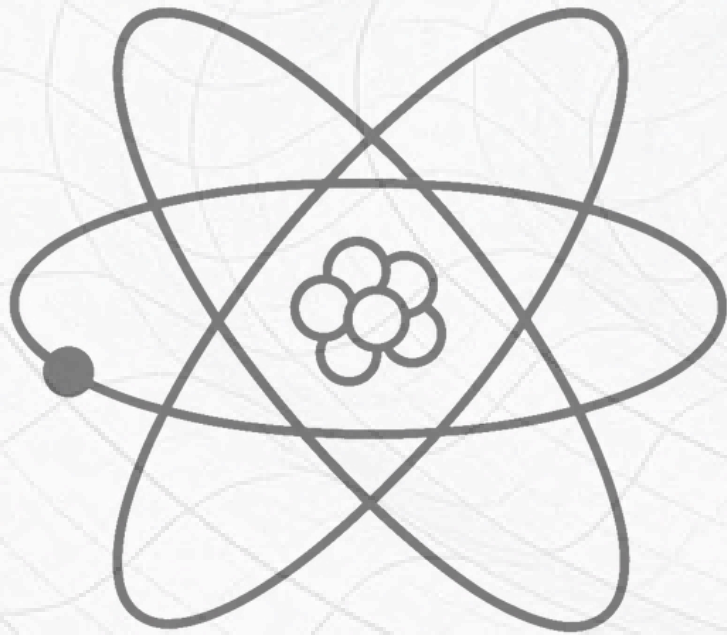


GWs and quantum systems



GWs and quantum systems

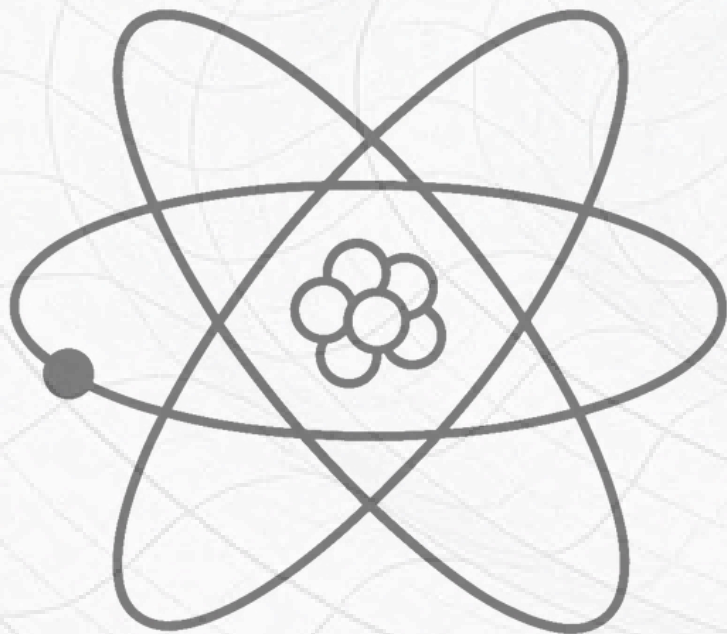
Atoms



- energy shifts
- resonant transitions

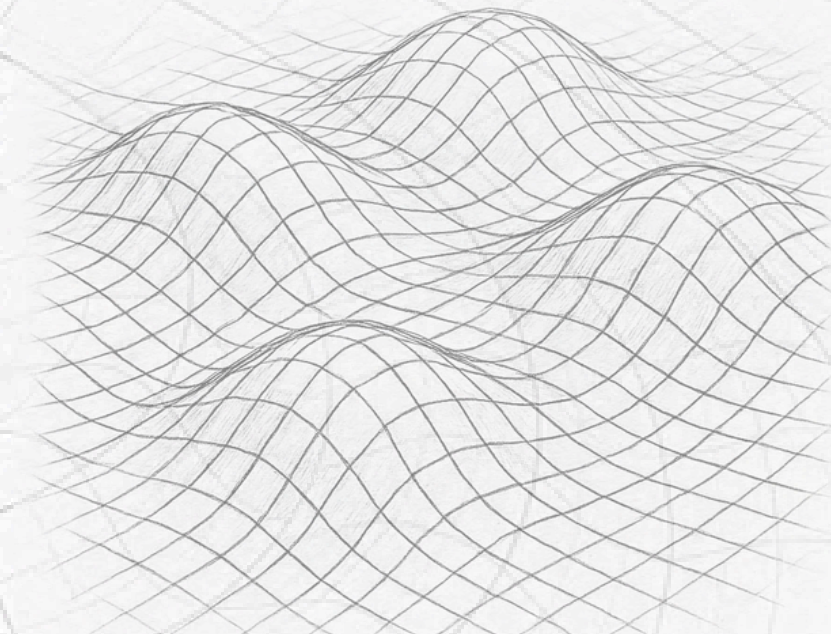
GWs and quantum systems

Atoms



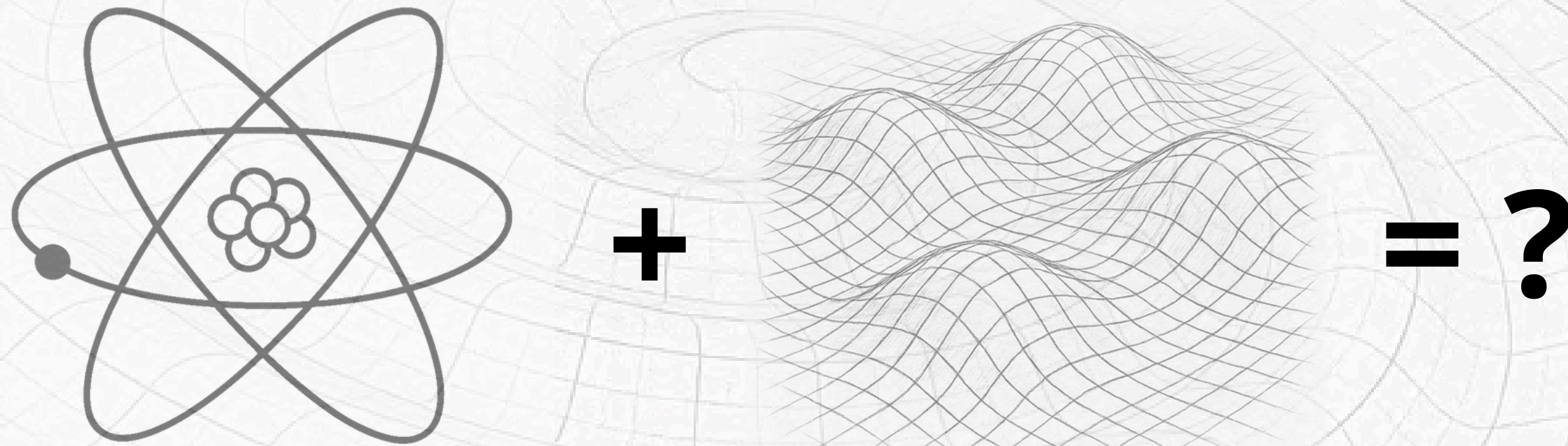
- energy shifts
- resonant transitions

Quantum fields

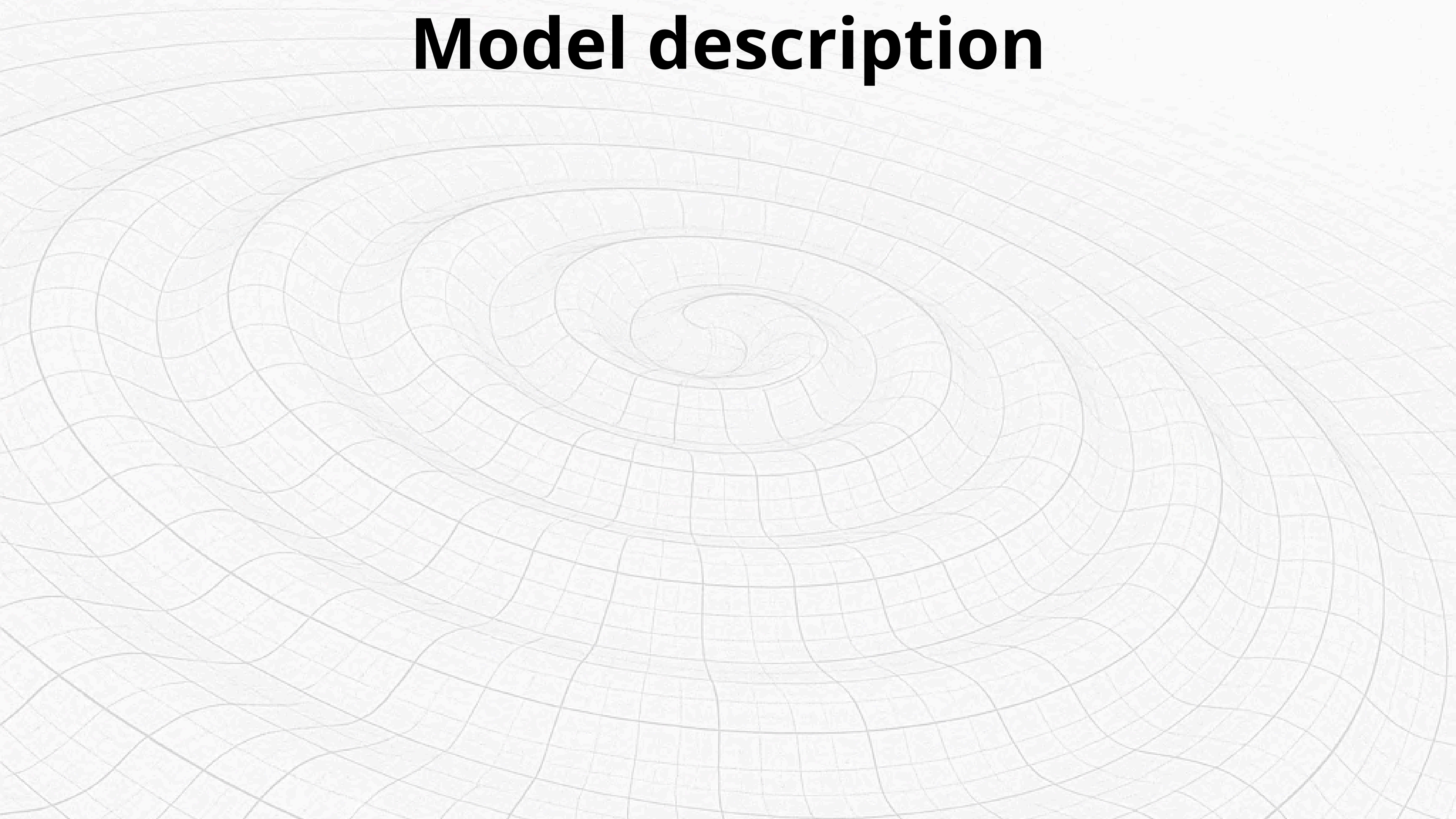


- particle production
- particle scattering

GWs and spontaneous emission



Model description



Model description

Gravitational wave:

- plane wave; traveling along z ; amplitude $\mathcal{A} \ll 1$; frequency ω
- $ds^2 = -dt^2 + dz^2 + (1 + \mathcal{A} \cos[\omega(t - z)])dx^2 + (1 - \mathcal{A} \cos[\omega(t - z)])dy^2$

Model description

Gravitational wave:

- plane wave; traveling along z ; amplitude $\mathcal{A} \ll 1$; frequency ω
- $ds^2 = -dt^2 + dz^2 + (1 + \mathcal{A} \cos[\omega(t - z)])dx^2 + (1 - \mathcal{A} \cos[\omega(t - z)])dy^2$

Atom:

- point-like two-level system: $|g\rangle$ and $|e\rangle$; energy gap ω_0
- monopole operator: $\hat{m}(t) = e^{i\omega_0 t}|e\rangle\langle g| + e^{-i\omega_0 t}|g\rangle\langle e|$

Model description

Gravitational wave:

- plane wave; traveling along z ; amplitude $\mathcal{A} \ll 1$; frequency ω
- $ds^2 = -dt^2 + dz^2 + (1 + \mathcal{A} \cos[\omega(t - z)])dx^2 + (1 - \mathcal{A} \cos[\omega(t - z)])dy^2$

Atom:

- point-like two-level system: $|g\rangle$ and $|e\rangle$; energy gap ω_0
- monopole operator: $\hat{m}(t) = e^{i\omega_0 t} |e\rangle\langle g| + e^{-i\omega_0 t} |g\rangle\langle e|$

Quantum field:

- massless, real, scalar
- quantized on a GW background
- field operator: $\hat{\phi}(x) = \int d^3\mathbf{k} \left[u_{\mathbf{k}}(x) \hat{a}_{\mathbf{k}} + u_{\mathbf{k}}^*(x) \hat{a}_{\mathbf{k}}^\dagger \right]$

Model description

Gravitational wave:

- plane wave; traveling along z ; amplitude $\mathcal{A} \ll 1$; frequency ω
- $ds^2 = -dt^2 + dz^2 + (1 + \mathcal{A} \cos[\omega(t - z)])dx^2 + (1 - \mathcal{A} \cos[\omega(t - z)])dy^2$

Atom:

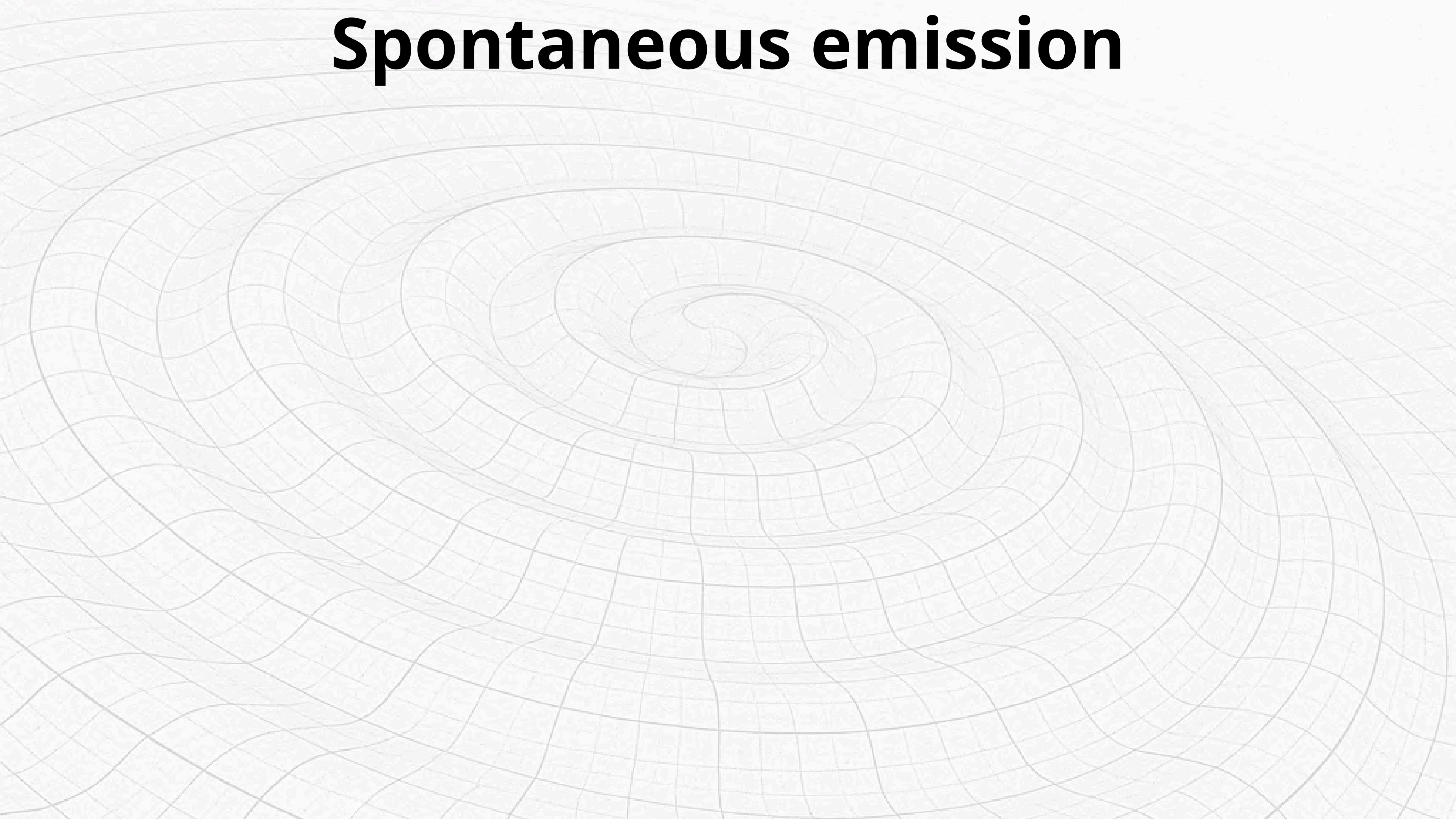
- point-like two-level system: $|g\rangle$ and $|e\rangle$; energy gap ω_0
- monopole operator: $\hat{m}(t) = e^{i\omega_0 t} |e\rangle\langle g| + e^{-i\omega_0 t} |g\rangle\langle e|$

Quantum field:

- massless, real, scalar
- quantized on a GW background
- field operator: $\hat{\phi}(x) = \int d^3\mathbf{k} \left[u_{\mathbf{k}}(x) \hat{a}_{\mathbf{k}} + u_{\mathbf{k}}^*(x) \hat{a}_{\mathbf{k}}^\dagger \right]$

Interaction: $\hat{H}_I(t) = \varepsilon \hat{m}(t) \hat{\phi}(x(t))$

Spontaneous emission



Spontaneous emission

Initial state: $|\psi_0\rangle = |e\rangle \otimes |0\rangle$

- atom in the excited state
- field in the vacuum state

Spontaneous emission

Initial state: $|\psi_0\rangle = |e\rangle \otimes |0\rangle$

- atom in the excited state
- field in the vacuum state

Atom's trajectory: $x(t) = 0$

Spontaneous emission

Initial state: $|\psi_0\rangle = |e\rangle \otimes |0\rangle$

- atom in the excited state
- field in the vacuum state

Atom's trajectory: $x(t) = 0$

State at time t : $|\psi(t)\rangle = |\psi_0\rangle - i \int_0^t dt' \hat{H}_I(t') |\psi_0\rangle - \int_0^t dt' \int_0^{t'} dt'' \hat{H}_I(t') \hat{H}_I(t'') |\psi_0\rangle$

Spontaneous emission

Initial state: $|\psi_0\rangle = |e\rangle \otimes |0\rangle$

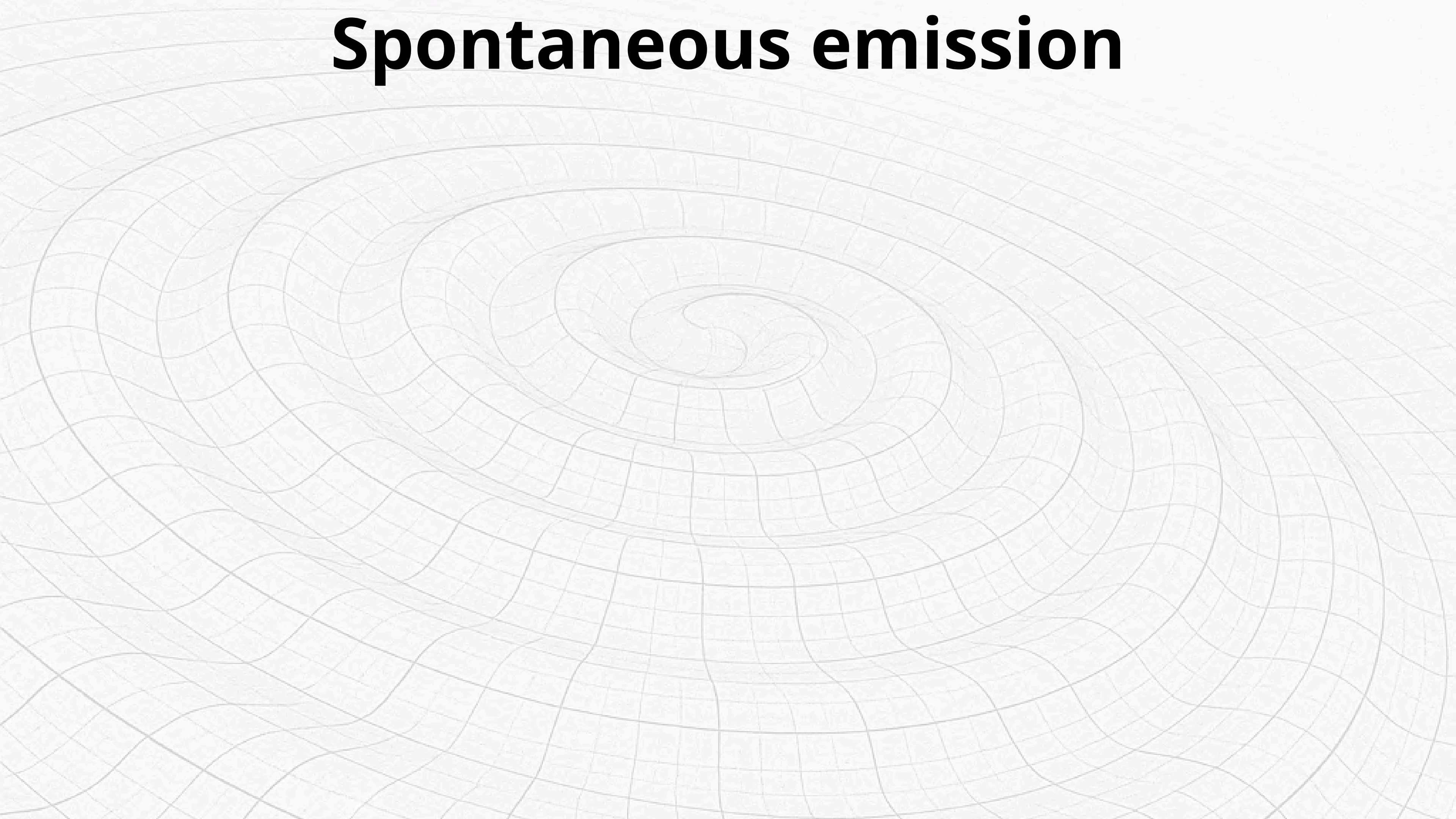
- atom in the excited state
- field in the vacuum state

Atom's trajectory: $x(t) = 0$

State at time t : $|\psi(t)\rangle = |\psi_0\rangle - i \int_0^t dt' \hat{H}_I(t') |\psi_0\rangle - \int_0^t dt' \int_0^{t'} dt'' \hat{H}_I(t') \hat{H}_I(t'') |\psi_0\rangle$

Expected number of emitted photons: $\langle n_{\mathbf{k}}(t) \rangle \equiv \langle \psi(t) | \hat{a}_{\mathbf{k}}^\dagger \hat{a}_{\mathbf{k}} | \psi(t) \rangle$

Spontaneous emission



Spontaneous emission

Expected number of emitted photons:

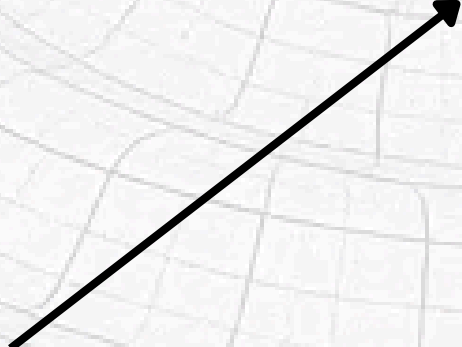
$$\langle n_{\mathbf{k}}(t) \rangle = \langle \tilde{n}_{\mathbf{k}}(t) \rangle + \langle \delta n_{\mathbf{k}}(t) \rangle$$

Spontaneous emission

Expected number of emitted photons:

$$\langle n_{\mathbf{k}}(t) \rangle = \langle \tilde{n}_{\mathbf{k}}(t) \rangle + \langle \delta n_{\mathbf{k}}(t) \rangle$$

Flat spacetime contribution



Spontaneous emission

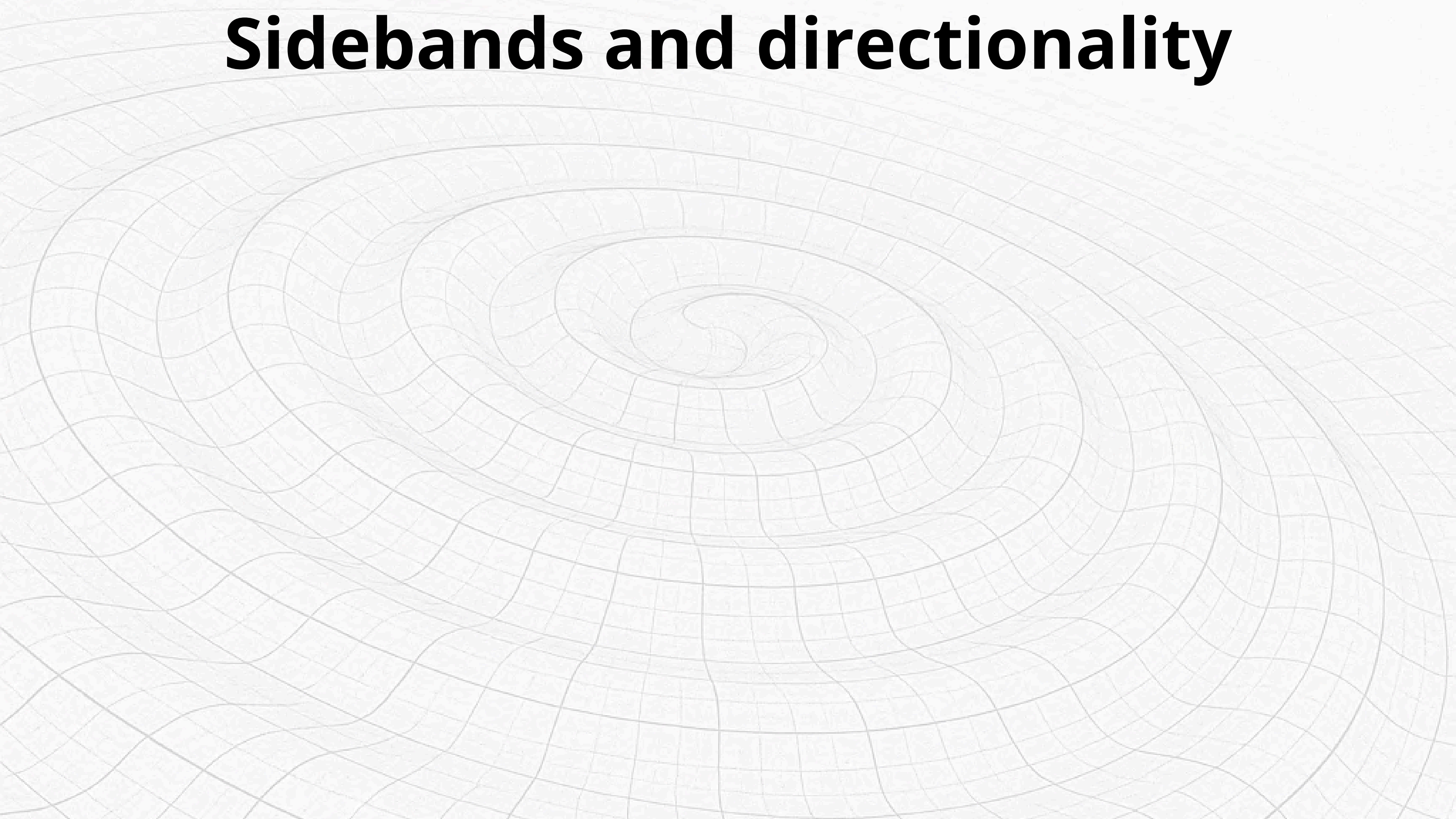
Expected number of emitted photons:

$$\langle n_{\mathbf{k}}(t) \rangle = \langle \tilde{n}_{\mathbf{k}}(t) \rangle + \langle \delta n_{\mathbf{k}}(t) \rangle$$

Flat spacetime contribution

GW correction

Sidebands and directionality



Sidebands and directionality

Frequency dependence:

$$\langle \tilde{n}_{\mathbf{k}}(t) \rangle \propto \text{sinc}^2(\delta_{\mathbf{k}} t/2) \quad \langle \delta n_{\mathbf{k}}(t) \rangle \propto \text{sinc}(\delta_{\mathbf{k}} t/2) (\text{sinc}[(\delta_{\mathbf{k}} - \omega)t/2] - \text{sinc}[(\delta_{\mathbf{k}} + \omega)t/2])$$

Sidebands in the spectrum!

$$\delta_{\mathbf{k}} \equiv |\mathbf{k}| - \omega_0$$

Sidebands and directionality

Frequency dependence:

$$\langle \tilde{n}_{\mathbf{k}}(t) \rangle \propto \text{sinc}^2(\delta_{\mathbf{k}} t/2) \quad \langle \delta n_{\mathbf{k}}(t) \rangle \propto \text{sinc}(\delta_{\mathbf{k}} t/2) (\text{sinc}[(\delta_{\mathbf{k}} - \omega)t/2] - \text{sinc}[(\delta_{\mathbf{k}} + \omega)t/2])$$

Sidebands in the spectrum!

$$\delta_{\mathbf{k}} \equiv |\mathbf{k}| - \omega_0$$

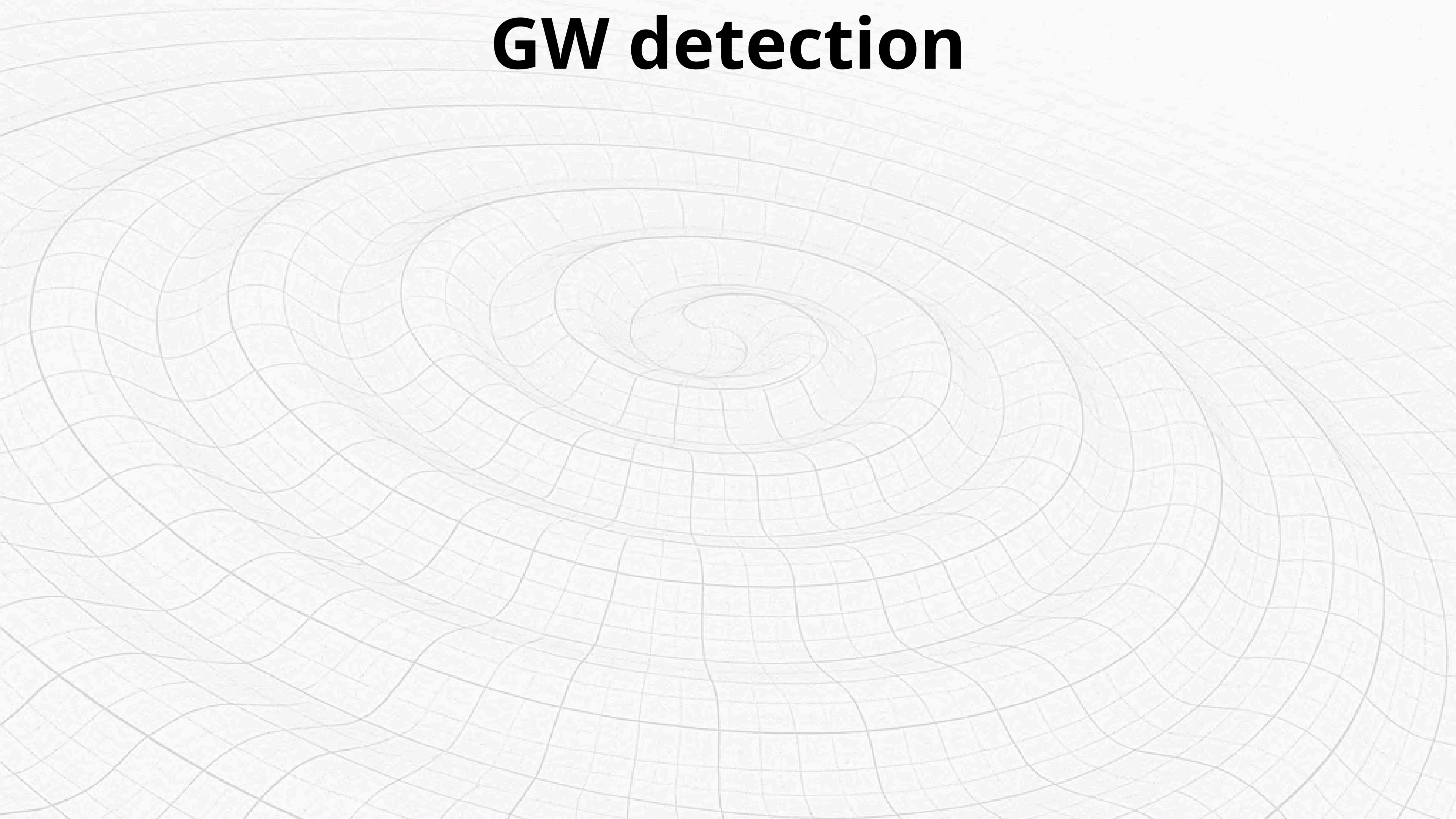
Angular dependence:

$$\langle \tilde{n}_{\mathbf{k}}(t) \rangle \propto 1 \quad \langle \delta n_{\mathbf{k}}(t) \rangle \propto \cos^2(\theta/2) \cos(2\varphi)$$

Directionality of the emission!

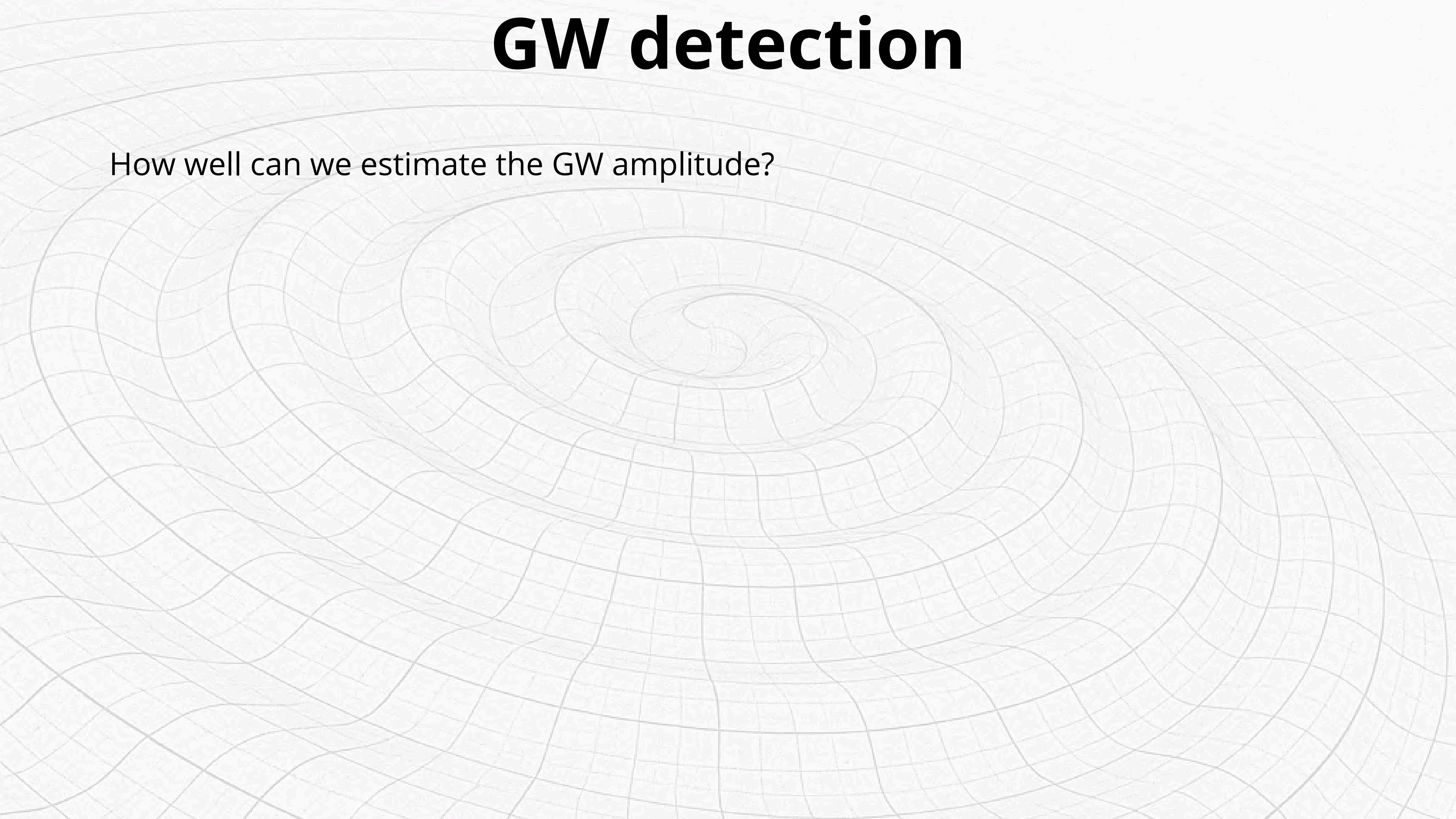
$$\mathbf{k} = (|\mathbf{k}| \sin \theta \cos \varphi, |\mathbf{k}| \sin \theta \sin \varphi, |\mathbf{k}| \cos \theta)$$

GW detection



GW detection

How well can we estimate the GW amplitude?



GW detection

How well can we estimate the GW amplitude?

Cramer-Rao bound: $\delta\mathcal{A} \geq \frac{1}{\sqrt{M\mathcal{I}(t)}}$

$\delta\mathcal{A}$ - estimation uncertainty

M - independent repetitions

$\mathcal{I}(t)$ - Fisher information

GW detection

How well can we estimate the GW amplitude?

Cramer-Rao bound: $\delta\mathcal{A} \geq \frac{1}{\sqrt{M\mathcal{I}(t)}}$

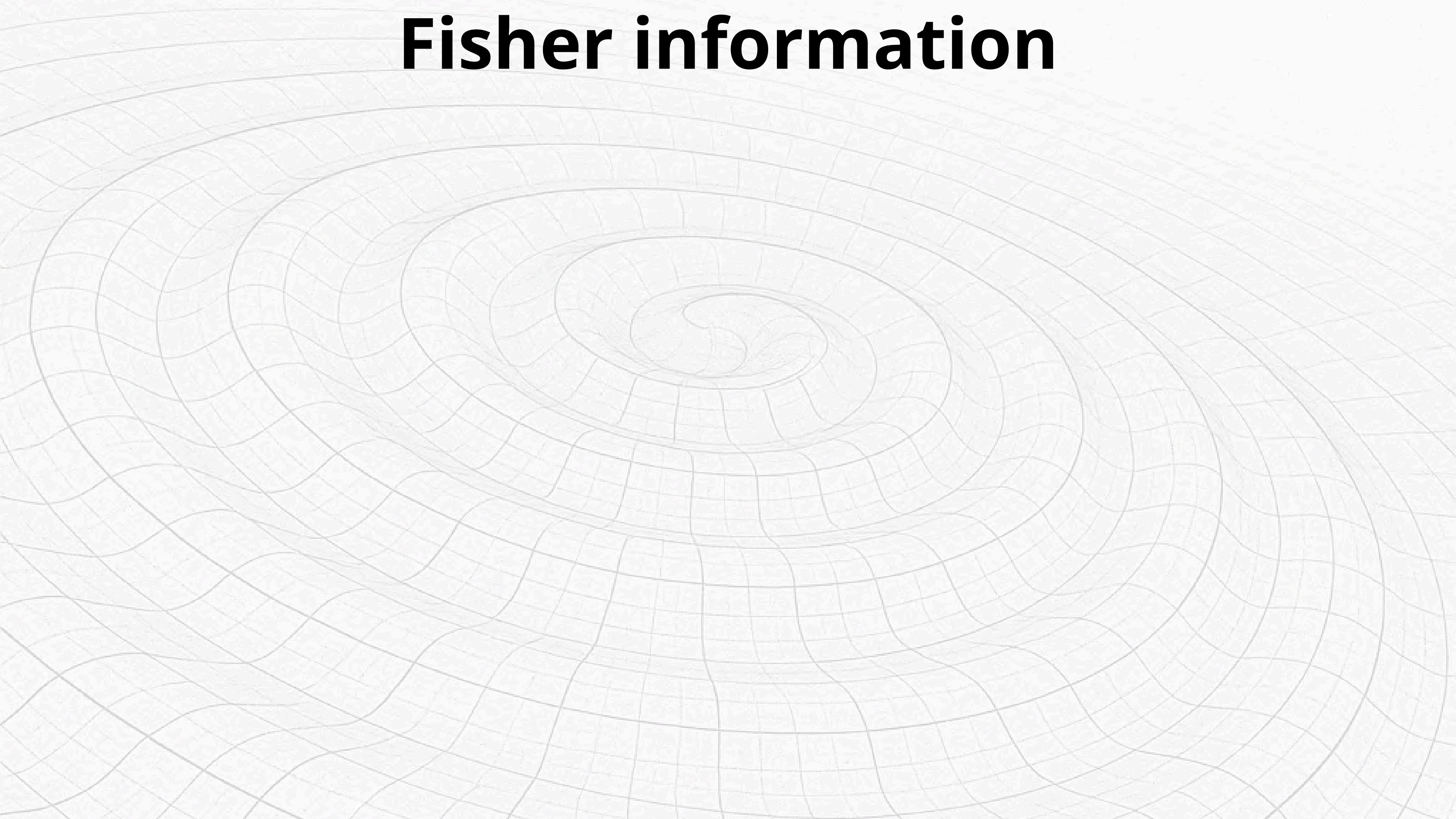
$\delta\mathcal{A}$ - estimation uncertainty

M - independent repetitions

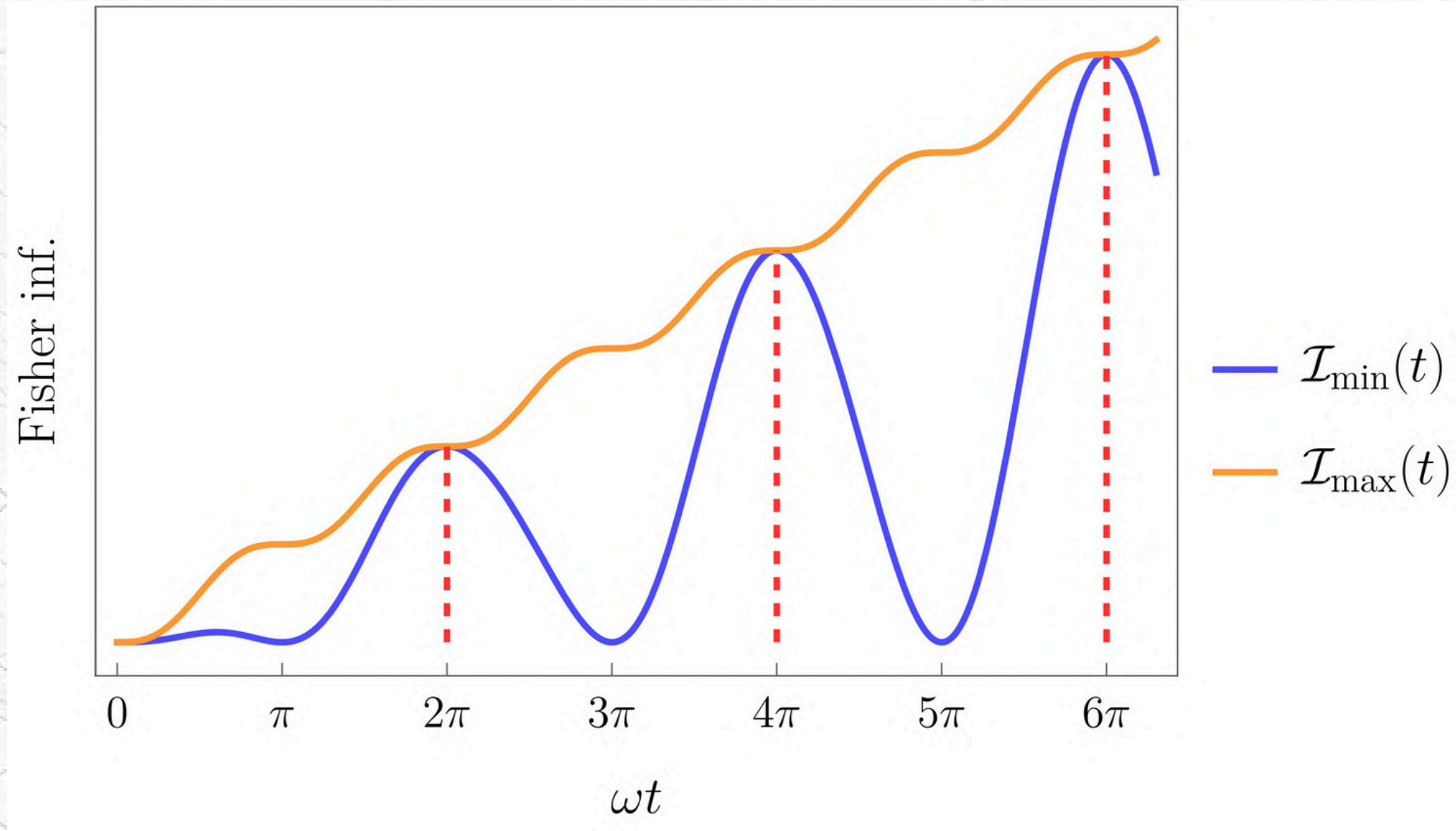
$\mathcal{I}(t)$ - Fisher information

Measure $\langle n_{\mathbf{k}}(t) \rangle$ for all $\mathbf{k}...$ What is the corresponding Fisher information?

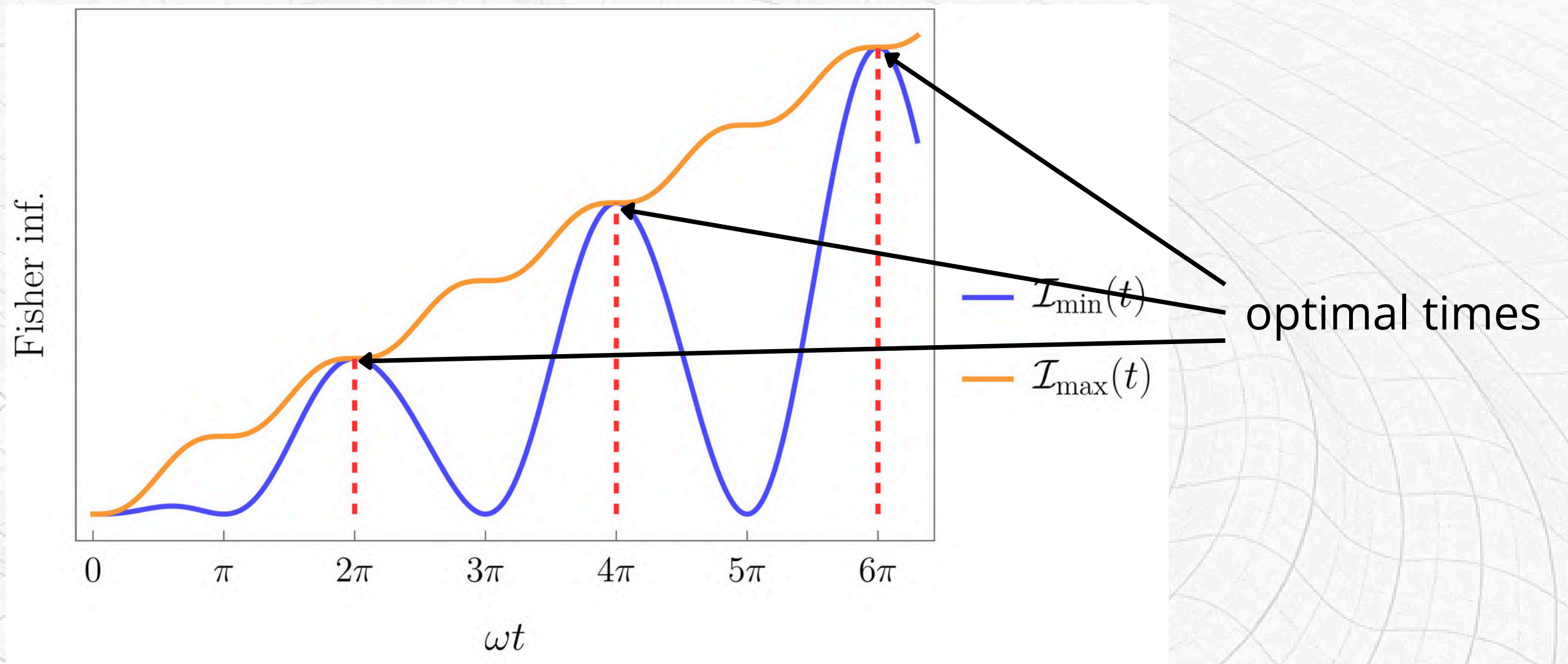
Fisher information



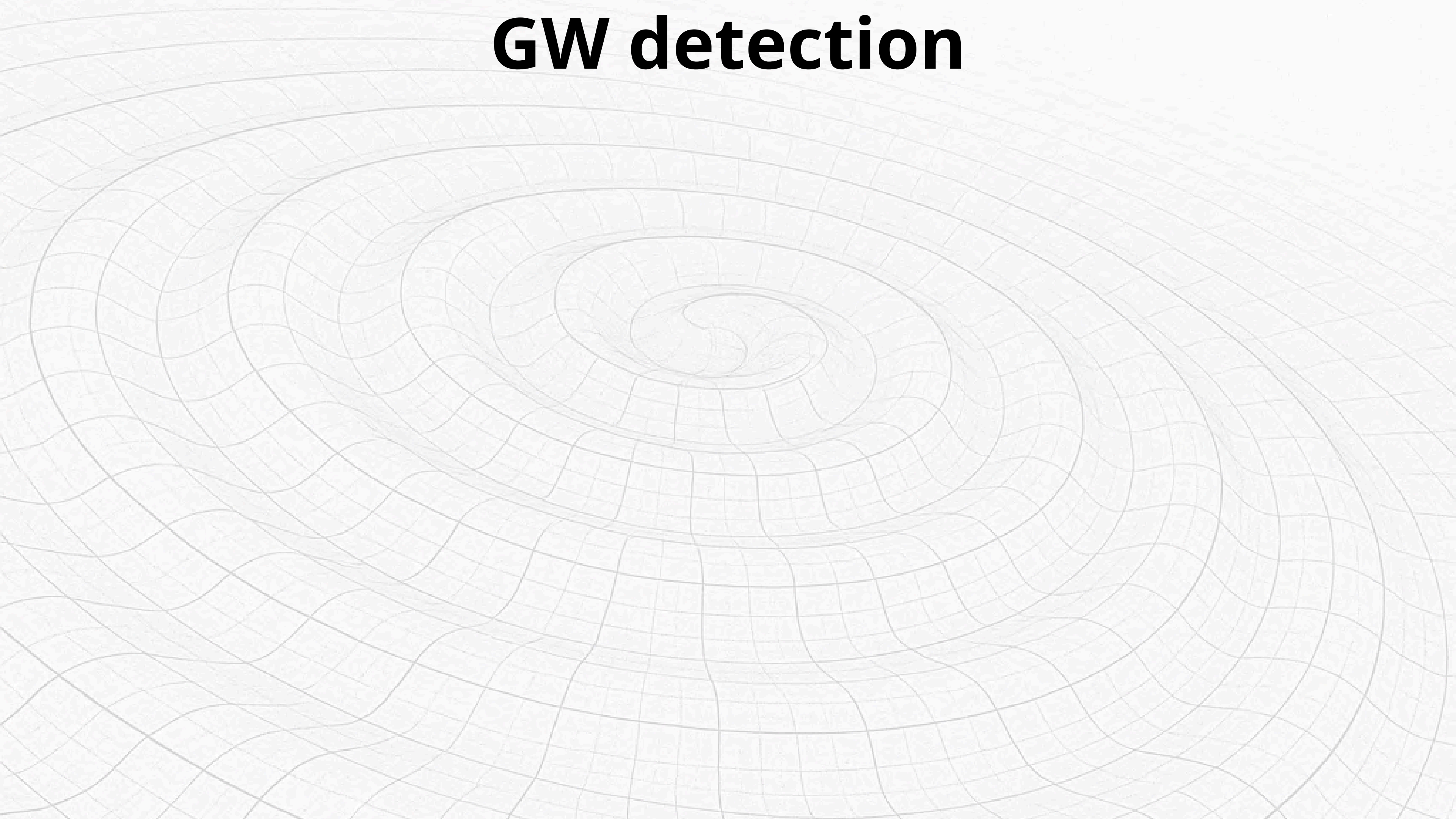
Fisher information



Fisher information



GW detection



GW detection

Cramer-Rao bound: $\delta\mathcal{A} \geq \frac{1}{\sqrt{M\mathcal{I}(t)}}$ M - number of atoms

GW detection

Cramer-Rao bound: $\delta\mathcal{A} \geq \frac{1}{\sqrt{M\mathcal{I}(t)}}$ M - number of atoms

Detection possible (in principle) if $\delta\mathcal{A} \leq \mathcal{A} \implies 1 \leq \mathcal{A}\sqrt{M\mathcal{I}(t)}$

GW detection

Cramer-Rao bound: $\delta\mathcal{A} \geq \frac{1}{\sqrt{M\mathcal{I}(t)}}$ M - number of atoms

Detection possible (in principle) if $\delta\mathcal{A} \leq \mathcal{A} \implies 1 \leq \mathcal{A}\sqrt{M\mathcal{I}(t)}$

Minimal required number of atoms: $M \geq (\mathcal{A}\omega_0/\omega)^{-2}$

GW detection

Cramer-Rao bound: $\delta\mathcal{A} \geq \frac{1}{\sqrt{M\mathcal{I}(t)}}$ M - number of atoms

Detection possible (in principle) if $\delta\mathcal{A} \leq \mathcal{A} \implies 1 \leq \mathcal{A}\sqrt{M\mathcal{I}(t)}$

Minimal required number of atoms: $M \geq (\mathcal{A}\omega_0/\omega)^{-2}$

$$\mathcal{A} \sim 10^{-21}, \quad \omega_0 \sim 10^{14} \text{ Hz}$$

GW detection

Cramer-Rao bound: $\delta\mathcal{A} \geq \frac{1}{\sqrt{M\mathcal{I}(t)}}$ M - number of atoms

Detection possible (in principle) if $\delta\mathcal{A} \leq \mathcal{A} \implies 1 \leq \mathcal{A}\sqrt{M\mathcal{I}(t)}$

Minimal required number of atoms: $M \geq (\mathcal{A}\omega_0/\omega)^{-2}$

$$\mathcal{A} \sim 10^{-21}, \quad \omega_0 \sim 10^{14} \text{ Hz}$$

Lower LIGO limit: $\omega \sim 10 \text{ Hz} \implies M \geq 10^{16}$

GW detection

Cramer-Rao bound: $\delta\mathcal{A} \geq \frac{1}{\sqrt{M\mathcal{I}(t)}}$ M - number of atoms

Detection possible (in principle) if $\delta\mathcal{A} \leq \mathcal{A} \implies 1 \leq \mathcal{A}\sqrt{M\mathcal{I}(t)}$

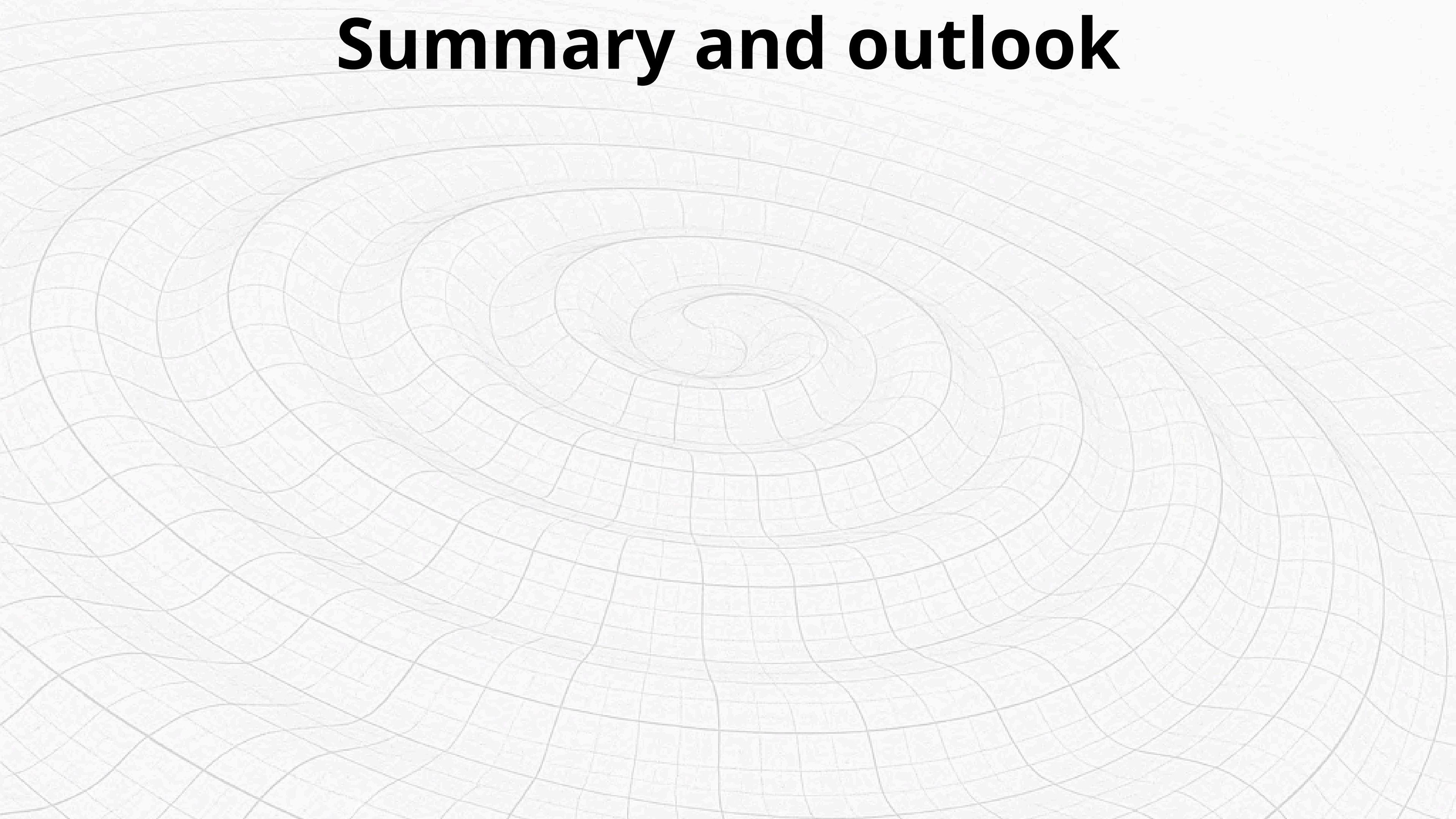
Minimal required number of atoms: $M \geq (\mathcal{A}\omega_0/\omega)^{-2}$

$$\mathcal{A} \sim 10^{-21}, \quad \omega_0 \sim 10^{14} \text{ Hz}$$

Lower LIGO limit: $\omega \sim 10 \text{ Hz} \implies M \geq 10^{16}$

Strontium linewidth: $\omega \sim 10^{-3} \text{ Hz} \implies M \geq 10^8$

Summary and outlook



Summary and outlook

Gravitational waves induce sidebands and directionality of the emission.

Summary and outlook

Gravitational waves induce sidebands and directionality of the emission.

It might be detectable - the requirements are not daunting.

Summary and outlook

Gravitational waves induce sidebands and directionality of the emission.

It might be detectable - the requirements are not daunting.

The method is general. Can we use spontaneous emission to probe the geometry of spacetime?

More details:

arXiv:2506.13872

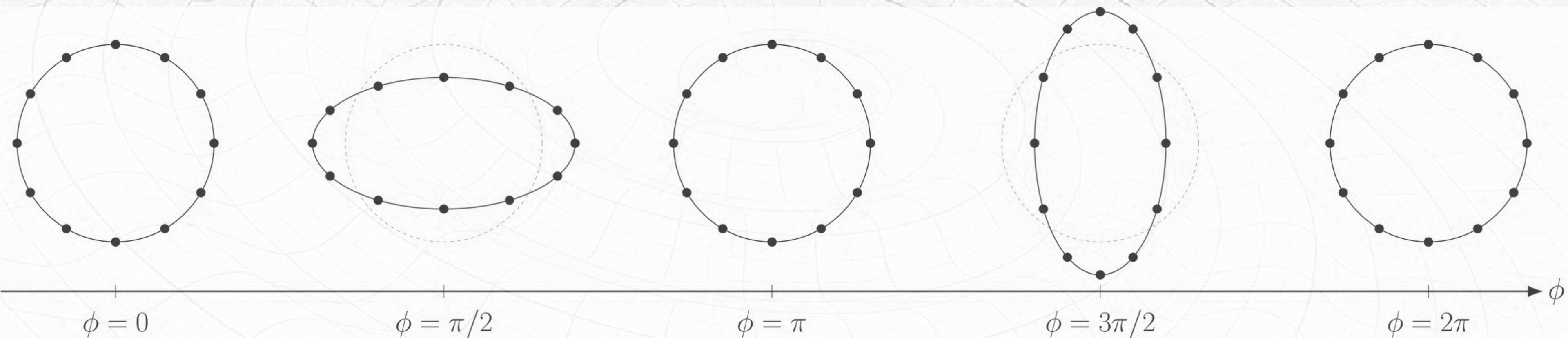
Gravitational wave imprints on spontaneous emission

Jerzy Paczos

Navdeep Arya, Sofia Qvarfort, Daniel Braun, and Magdalena Zych

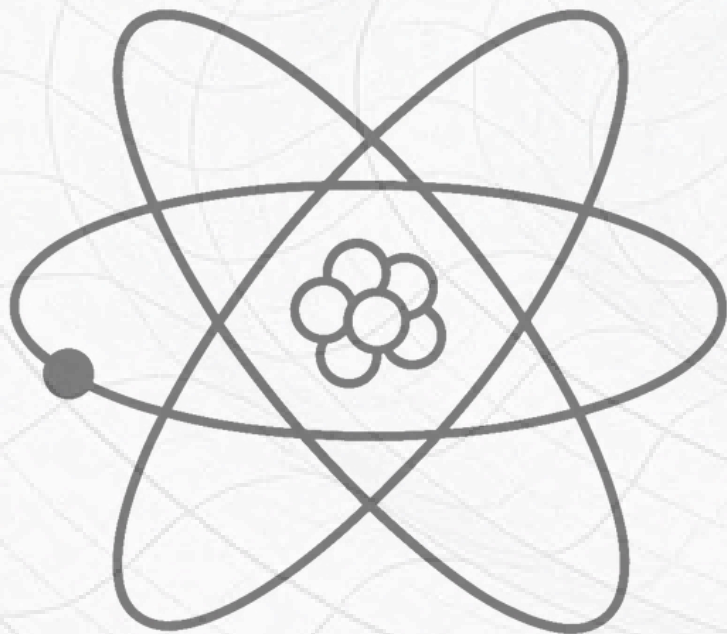
RQI North, Naples 2025

GWs and classical test masses



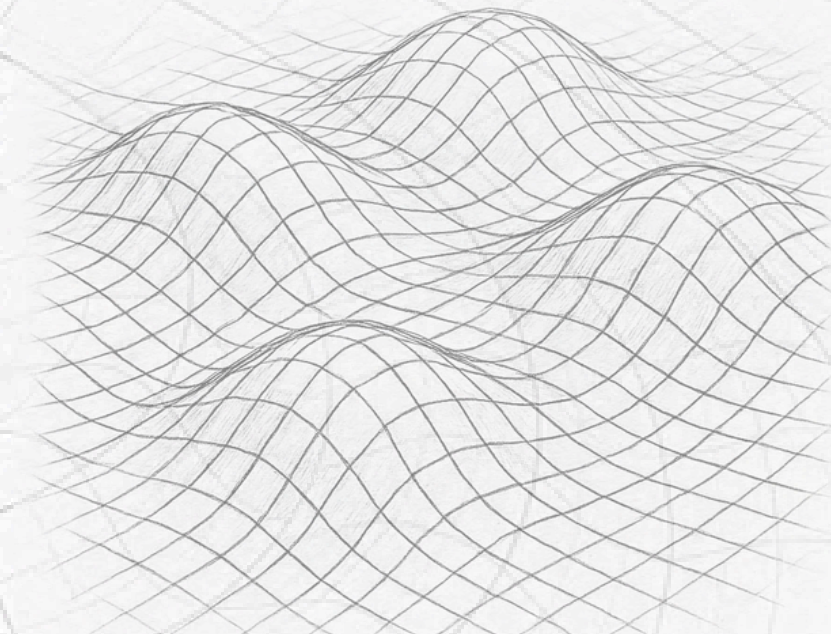
GWs and quantum systems

Atoms



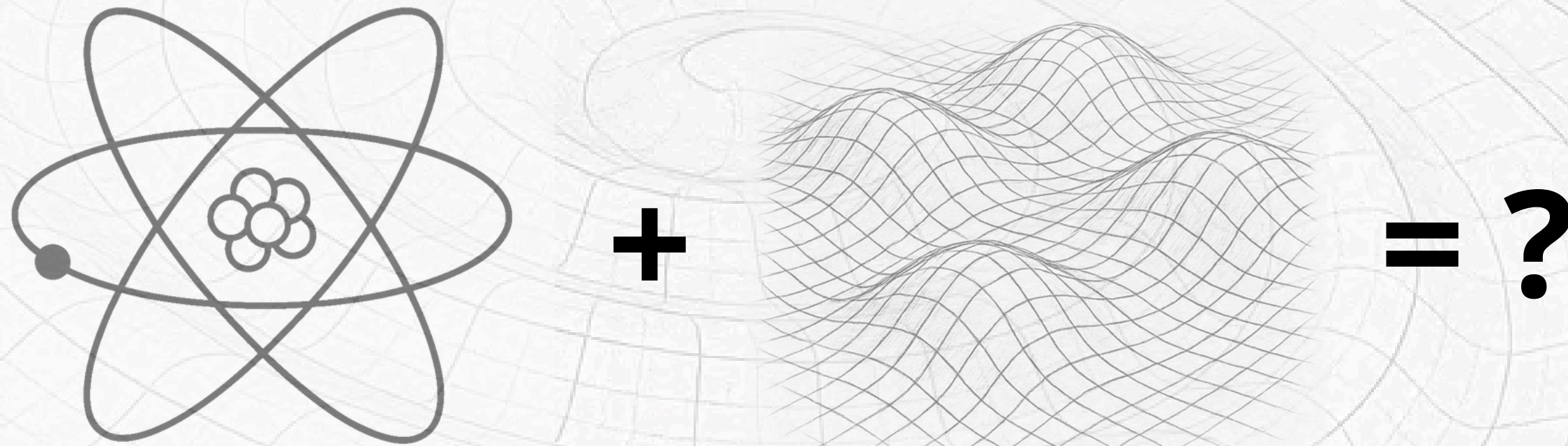
- energy shifts
- resonant transitions

Quantum fields



- particle production
- particle scattering

GWs and spontaneous emission



Model description

Gravitational wave:

- plane wave; traveling along z ; amplitude $\mathcal{A} \ll 1$; frequency ω
- $ds^2 = -dt^2 + dz^2 + (1 + \mathcal{A} \cos[\omega(t - z)])dx^2 + (1 - \mathcal{A} \cos[\omega(t - z)])dy^2$

Atom:

- two-level system: $|g\rangle$ and $|e\rangle$; energy gap ω_0
- monopole operator: $\hat{m}(t) = e^{i\omega_0 t} |e\rangle\langle g| + e^{-i\omega_0 t} |g\rangle\langle e|$

Quantum field:

- massless, real, scalar
- quantized on a GW background
- field operator: $\hat{\phi}(x) = \int d^3\mathbf{k} \left[u_{\mathbf{k}}(x) \hat{a}_{\mathbf{k}} + u_{\mathbf{k}}^*(x) \hat{a}_{\mathbf{k}}^\dagger \right]$

Interaction: $\hat{H}_I(t) = \varepsilon \hat{m}(t) \hat{\phi}(x(t))$

Spontaneous emission

Initial state: $|\psi_0\rangle = |e\rangle \otimes |0\rangle$

- atom in the excited state
- field in the vacuum state

Atom's trajectory: $x(t) = 0$

State at time t : $|\psi(t)\rangle = |\psi_0\rangle - i \int_0^t dt' \hat{H}_I(t') |\psi_0\rangle - \int_0^t dt' \int_0^{t'} dt'' \hat{H}_I(t') \hat{H}_I(t'') |\psi_0\rangle$

Expected number of emitted photons: $\langle n_{\mathbf{k}}(t) \rangle \equiv \langle \psi(t) | \hat{a}_{\mathbf{k}}^\dagger \hat{a}_{\mathbf{k}} | \psi(t) \rangle$

Spontaneous emission

Expected number of emitted photons:

$$\langle n_{\mathbf{k}}(t) \rangle = \langle \tilde{n}_{\mathbf{k}}(t) \rangle + \langle \delta n_{\mathbf{k}}(t) \rangle$$

Flat spacetime contribution

GW correction

Spontaneous emission

Flat spacetime contribution:

$$\langle \tilde{n}_k(t) \rangle = \frac{\varepsilon^2 t^2}{(2\pi)^3 8|k|} \text{sinc}^2(\delta_k t/2) \quad \delta_k \equiv |k| - \omega_0$$

GW correction:

$$\langle \delta n_k(t) \rangle = \frac{\varepsilon^2 t^2}{(2\pi)^3 8|k|} \mathcal{A} \frac{|k|}{\omega} f(\delta_k, t) g(\theta, \varphi) \quad k = (|k| \sin \theta \cos \varphi, |k| \sin \theta \sin \varphi, |k| \cos \theta)$$

$$f(\delta_k, t) = \text{sinc}(\delta_k t/2) \cos(\omega t/2) (\text{sinc}[(\delta_k - \omega)t/2] - \text{sinc}[(\delta_k + \omega)t/2])$$

$$g(\theta, \varphi) \equiv \cos^2(\theta/2) \cos(2\varphi)$$

Sidebands and directionality

Frequency dependence:

$$\langle \tilde{n}_{\mathbf{k}}(t) \rangle \propto \text{sinc}^2(\delta_{\mathbf{k}} t/2) \quad \langle \delta n_{\mathbf{k}}(t) \rangle \propto \text{sinc}(\delta_{\mathbf{k}} t/2) (\text{sinc}[(\delta_{\mathbf{k}} - \omega)t/2] - \text{sinc}[(\delta_{\mathbf{k}} + \omega)t/2])$$

Sidebands in the spectrum!

$$\delta_{\mathbf{k}} \equiv |\mathbf{k}| - \omega_0$$

Angular dependence:

$$\langle \tilde{n}_{\mathbf{k}}(t) \rangle \propto 1 \quad \langle \delta n_{\mathbf{k}}(t) \rangle \propto \cos^2(\theta/2) \cos(2\varphi)$$

Directionality of the emission!

Equivalence principle

The atom is point-like... Why is there any effect at all?

The atom-field system is extended.

No problems with EP!

The total emission rate remains unchanged.

No information about GW in the atomic state!

GW detection

How well can we estimate the GW amplitude?

Cramer-Rao bound: $\delta\mathcal{A} \geq \frac{1}{\sqrt{M\mathcal{I}(t)}}$

$\delta\mathcal{A}$ - estimation uncertainty

M - independent repetitions

$\mathcal{I}(t)$ - Fisher information; measurement-dependent

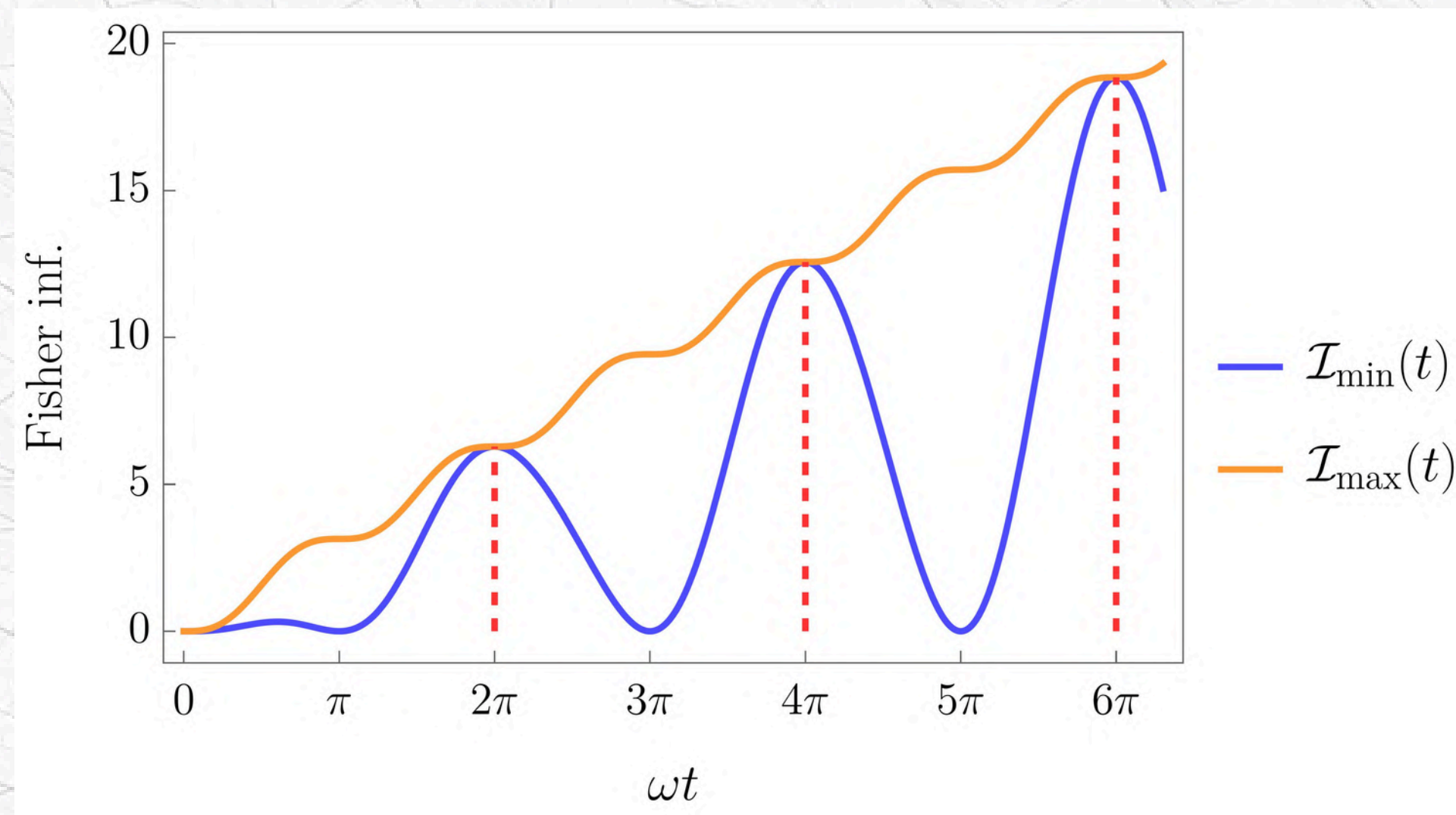
Measure $\langle n_{\mathbf{k}}(t) \rangle$ for all $\mathbf{k}$... What is the corresponding Fisher information?

Fisher information

Bounds: $\mathcal{I}_{\min}(t) = \frac{\bar{n}(t)}{3} \left(\frac{\omega_0}{\omega} \right)^2 \cos^2(\omega t/2) [1 - \text{sinc}(\omega t)]$

$$\mathcal{I}_{\max}(t) = \frac{\bar{n}(t)}{3} \left(\frac{\omega_0}{\omega} \right)^2 [1 - \cos(\omega t) \text{sinc}(\omega t)]$$

$\bar{n}(t)$ - total number of emitted photons

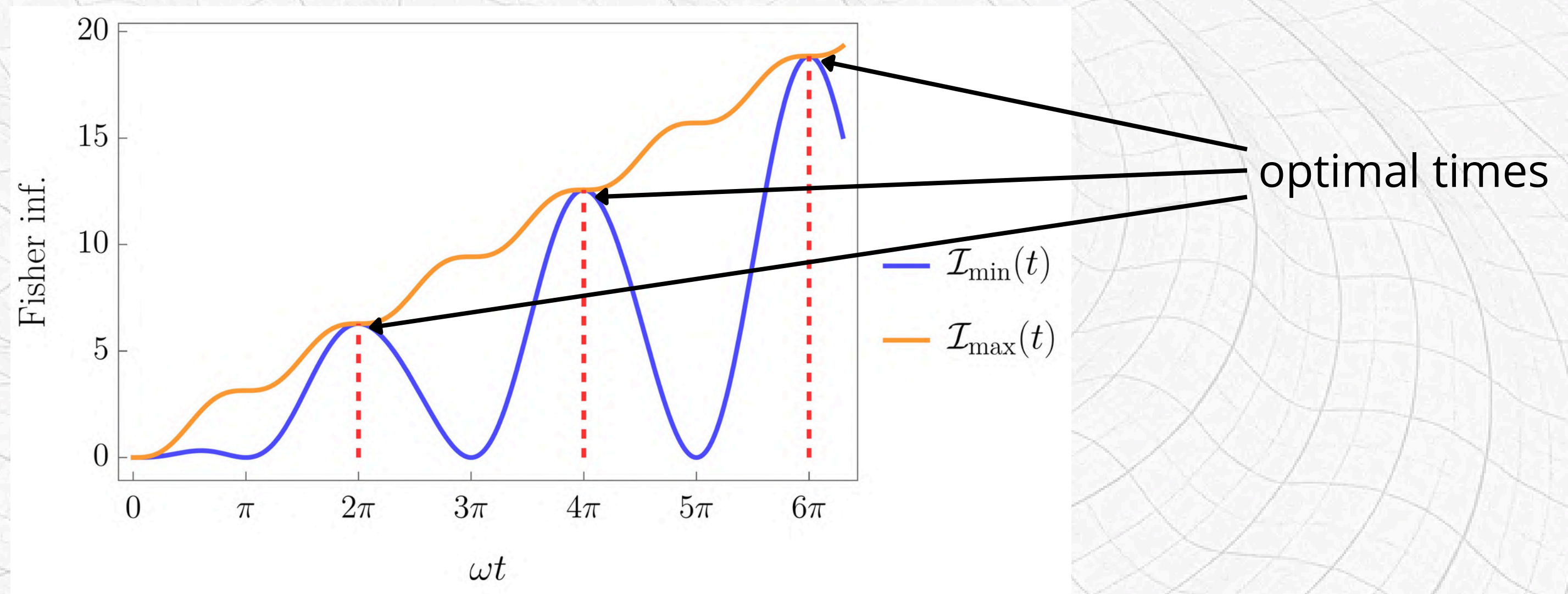


Fisher information

Bounds: $\mathcal{I}_{\min}(t) = \frac{\bar{n}(t)}{3} \left(\frac{\omega_0}{\omega} \right)^2 \cos^2(\omega t/2) [1 - \text{sinc}(\omega t)]$

$$\mathcal{I}_{\max}(t) = \frac{\bar{n}(t)}{3} \left(\frac{\omega_0}{\omega} \right)^2 [1 - \cos(\omega t) \text{sinc}(\omega t)]$$

$\bar{n}(t)$ - total number of emitted photons



GW detection

Cramer-Rao bound: $\delta\mathcal{A} \geq \frac{1}{\sqrt{M\mathcal{I}(t)}}$ M - independent repetitions
(number of atoms)

Detection possible (in principle) if $\delta\mathcal{A} \leq \mathcal{A} \implies 1 \leq \mathcal{A}\sqrt{M\mathcal{I}(t)}$

$$t = \frac{2\pi m}{\omega}, m \in \mathbb{N}: \quad \mathcal{I}(t) = \frac{\bar{n}(t)}{3} \left(\frac{\omega_0}{\omega} \right)^2$$

$$\bar{n}(t) \lesssim 1 \implies (\mathcal{A}\omega_0/\omega)^{-2} \leq M \quad \mathcal{A} \sim 10^{-21}, \quad \omega_0 \sim 10^{14} \text{ Hz}$$

$$\text{Lower LIGO limit: } \omega \sim 10 \text{ Hz} \implies M \geq 10^{16}$$

$$\text{Strontium linewidth: } \omega \sim 10^{-3} \text{ Hz} \implies M \geq 10^8$$

Summary

Gravitational waves induce sidebands and directionality of the emission.

The effect is consistent with the equivalence principle, even though the atom is point-like.

It might be detectable - the requirements are not daunting.

Outlook

A more realistic model of atom-light interaction is needed for accurate experimental predictions.

An analysis of noise is required.

The method is general. Can we use spontaneous emission to probe the geometry of spacetime?