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[with Joseph Balsells, arXiv:2506.14705]





Geodesic motion of quantum objects

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Quantum object on classical Riemannian space-time, metric g_{ab} .

Extended by quantum fluctuations: Experiences tidal forces.

New formalism:

Geometric optics approximation of QFT on curved space-time.

Implications:

- \rightarrow Entropy and purity acquire weight.
- → Complete description requires non-Riemannian geometry.





Evaluated with quantum-information methods, focusing on time dilation.

[Smith, Ahmadi: Nat. Comm. 11 (2020) 5360; Zych, Costa, Pikovski, Brukner: Nat. Comm. 2 (2022) 505]

New strategy: Geometrical formulation of quantum mechanics combines quantum and space-time physics.

Geometric formulation



Equivalent to Hamiltonian dynamics:

→ Poisson bracket

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$$\{\langle \hat{O}_1 \rangle_{\rho}, \langle \hat{O}_2 \rangle_{\rho}\} = \frac{\langle [\hat{O}_1, \hat{O}_2] \rangle_{\rho}}{i\hbar}$$

extended by Leibniz rule. Phase-space geometry.

- → Effective Hamiltonian $H_{\text{eff}}(\rho) = \langle \hat{H} \rangle_{\rho}$ on state space.
- \rightarrow Hamilton's equations

$$\frac{\mathrm{d}\langle\hat{O}\rangle_{\rho}}{\mathrm{d}t} = \{\langle\hat{O}\rangle_{\rho}, H_{\mathrm{eff}}(\rho)\}$$





Basic expectation values $q = \langle \hat{q} \rangle$ and $p = \langle \hat{p} \rangle$, central moments

$$\Delta(q^a p^b) = \langle (\hat{q} - q)^a (\hat{p} - p)^b \rangle_{\text{symm}} ,$$

in completely symmetric ordering: phase-space coordinates.

Not canonical, for instance $\{\Delta(q^2), \Delta(p^2)\} = 4\Delta(qp),$ $\{\Delta(q^2), \Delta(qp)\} = 2\Delta(q^2),$ $\{\Delta(qp), \Delta(p^2)\} = 2\Delta(p^2).$

Casimir–Darboux theorem:

Every Poisson manifold has local coordinates q_i , p_i , C_I such that $\{q_i, p_j\} = \delta_{ij}$ and $\{q_i, C_I\} = 0 = \{p_i, C_I\}$.

Casimir functions C_I conserved for any Hamiltonian.

Second-order moments

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$$\{\Delta(q^2), \Delta(p^2)\} = 4\Delta(qp)$$

$$\{\Delta(q^2), \Delta(qp)\} = 2\Delta(q^2)$$

$$\{\Delta(qp), \Delta(p^2)\} = 2\Delta(p^2)$$

Casimir–Darboux variables:

$$\Delta(q^2) = s^2 \quad , \quad \Delta(qp) = sp_s \quad , \quad \Delta(p^2) = p_s^2 + \frac{C}{s^2}$$

with canonical $\{s, p_s\} = 1$ and $\{s, C\} = 0 = \{p_s, C\}$.

 $C = \Delta(q^2)\Delta(p^2) - \Delta(qp)^2 \ge \hbar^2/4$

[Jackiw, Kerman: PLA 71 (1979) 158; Jalabert, Pastawski: PRL 86 (2001) 2490; Prezhdo: Theo Chem Acc 116 (2006) 206; Vachaspati, Zahariade: PRD 98 (2018) 065002]

Two classical degrees of freedom

[Baytaş, MB, Crowe, arXiv:1810.12127, arXiv:1811.00505]

 $\Delta(x_1^2) = s_1^2$, $\Delta(x_2^2) = s_2^2$, $\Delta(x_1x_2) = s_1s_2\cos\beta$

with new canonical variable β : Correlation as an angle.

Momentum p_{β} appears in $\Delta(x_1p_2)$ and in

$$\Delta(p_1^2) = p_{s_1}^2 + \frac{U_1}{s_1^2}$$

where

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 $U_{1} = (p_{\alpha} - p_{\beta})^{2} + \frac{1}{2\sin^{2}\beta} \left((C_{1}^{2} - 4p_{\alpha}^{2}) - \sqrt{C_{2}^{4} - C_{1}^{4} + (C_{1}^{2} - 4p_{\alpha}^{2})^{2}} \sin(\alpha + \beta) \right)$ with additional canonical pair (α, p_{α}) and Casimirs C_{1}, C_{2} . Purity $\operatorname{Tr}(\rho^{2})$ conserved for any Hamiltonian: Function of C_{1}, C_{2} .

Hamiltonian proper time

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[Balsells, MB, arXiv:2410.08156]

Canonical coordinates $\{x^a, \pi_b\} = \delta_b^a$. Hamiltonian

$$H = -\frac{1}{2m} \left(g^{ab}(x)\pi_a \pi_b + m^2 c^2 \right)$$

implies geodesic motion, proper time $\Delta \tau = \int \sqrt{1 + 2H/(mc^2)} \, d\tau$.

$$H_{\text{eff}} = -\frac{1}{2m} \left(\langle g^{ab}(\hat{x}) \hat{\pi}_a \hat{\pi}_b \rangle + m^2 c^2 \right)$$

$$= -\frac{1}{2m} \left(g^{ab} \pi_a \pi_b + m^2 c^2 \right)$$

$$+ g^{ij} \Delta(p_i p_j) + \frac{\partial g^{aj}}{\partial x^k} \pi_a \Delta(x^k p_j) + \frac{1}{2} \frac{\partial^2 g^{ab}}{\partial x^k \partial x^l} \pi_a \pi_b \Delta(x^k x^l) \right)$$

Quantize position \hat{x} in g^{ab} , but no quantization of gravity.





[Balsells, MB, arXiv:2503.06667]

Inserting canonical variables, crucial expression

$$\sqrt{C_2^4 - C_1^4 + (C_1^2 - 4p_{\alpha}^2)^2}$$

from U_1 appears in effective Hamiltonian/proper time.

- → In general, non-quadratic in momenta: Non-Riemannian (Finsler) geometry in combined space-time/state space.
- \rightarrow Quadratic if

 $C_1 = C_2$ forbidden $p_{\alpha} = 0$ Gaussian

The weight of purity



 $C_1 \text{ and } C_2 \text{ conserved: } \dot{C}_{1,2} = \{C_{1,2}, H\} = 0 \text{ for any } H.$

Related to symplectic eigenvalues ν_{\pm} of covariance matrix $\sigma_{ij} = \Delta(x_i x_j)$:

$$C_1^2 = \nu_+^2 + \nu_-^2$$
, $C_2^2 = \nu_+^2 - \nu_-^2$.

Uncertainty relation implies $\nu_+, \nu_- \geq \hbar/2$, thus $C_1 \neq C_2$.

Purity

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$$\mu(\sigma) = \frac{\hbar^2}{4\sqrt{\det(\sigma)}} = \frac{\hbar^2}{4\nu_+\nu_-} = \frac{\hbar^2}{2\sqrt{C_1^4 - C_2^4}}$$

appears in $\sqrt{C_2^4 - C_1^4 + (C_1^2 - 4p_{\alpha}^2)^2}$.

Logarithmic negativity

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[Balsells, MB, arXiv:2506.14705]

Complicated expression simplifies for $\beta = \pi/2$:

$$e^{-2E_{\mathcal{N}}}|_{\beta=\pi/2} = C_1^2 + 4p_{\beta}^2 - 4p_{\alpha}^2 - \sqrt{(C_1^2 + 4p_{\beta}^2 - 4p_{\alpha}^2)^2 - C_1^4 + C_2^4}$$

 $p_{-} = \frac{1}{2}(\nu_{+} - \nu_{-})$: limits regions forbidden by reality conditions

Entangled states above and below red hyperbolae







- → Detailed relationship between time dilation effects and quantum-information parameters.
- \rightarrow Also applicable to tidal forces and geodesic deflection.
- → Specific indications for modified space-time structure, related to entropy and purity.
- → Remaining canonical parameter

$$\cos \alpha = -\sin \beta \frac{(\Delta_{x_2 p_2}^{x_1 x_2})^2 - (\Delta_{x_1 p_2}^{x_1 x_2})^2 + \Delta_{x_1 x_2}^{x_1 x_2} \Delta_{x_1 p_1}^{x_1 p_1} - \Delta_{x_1 x_2}^{x_1 x_2} \Delta_{x_2 p_2}^{x_2 p_2}}{\Delta_{x_1 x_2}^{x_1 x_2} \sqrt{(C_1^2 - 4p_\alpha^2)^2 - C_1^4 + C_2^4}}$$

where

$$\Delta_{A_1A_2}^{B_1B_2} = \Delta(A_1B_1)\Delta(A_2B_2) - \Delta(A_1B_2)\Delta(A_2B_1)$$

Distinguished by canonical nature of variables. Specific quantum-information meaning?