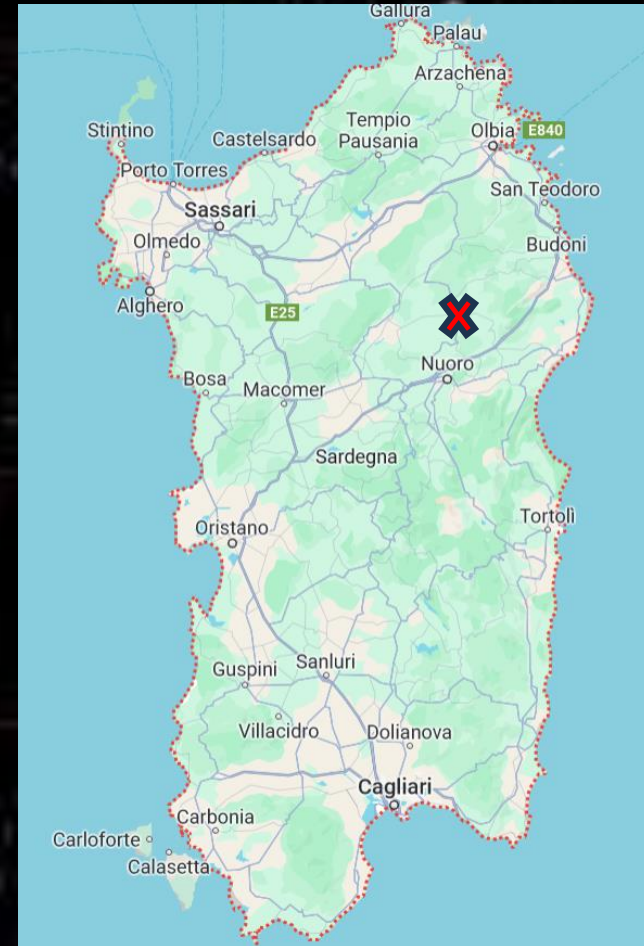


The Weight of Vacuum



Archimedes: an interdisciplinary experiment

ARCHIMEDES is a truly interdisciplinary experiment:

- Scientific motivations involve GR and QFT;
- The expected signal is produced through Superconductivity and Thermodynamics;
- The signal detection is based in high precision opto-mechanics.

And finally, there are some spin-off implication related to DM search and Geophysics.

- Gravity
- Quantum Field Theory
- Superconductivity
- Cryogenics
- Precision Mechanics
- Optics
- Electronics
- Dark Matter search
- Geophysics

Scientific Motivations



Since the birth of Quantum Mechanics the question rised if the zero-point energy gravitates (Nerst, Pauli...) –The first attempt of by Pauli



$$\sum \frac{1}{2} \hbar \omega \quad \longrightarrow \quad \infty$$

Pauli inserted a cut-off on the minimal length (electron classical radius) and inserted the value of the energy density in the static Einstein solution

The expected radius of the Universe was: 31 Km!

Cosmological constant problem: “why the universe exhibits a vacuum energy density much smaller than the one resulting from application of quantum mechanics and equivalence principle?” (Weinberg **Rev.Mod.Phys. 61 (1989) 1-23**)

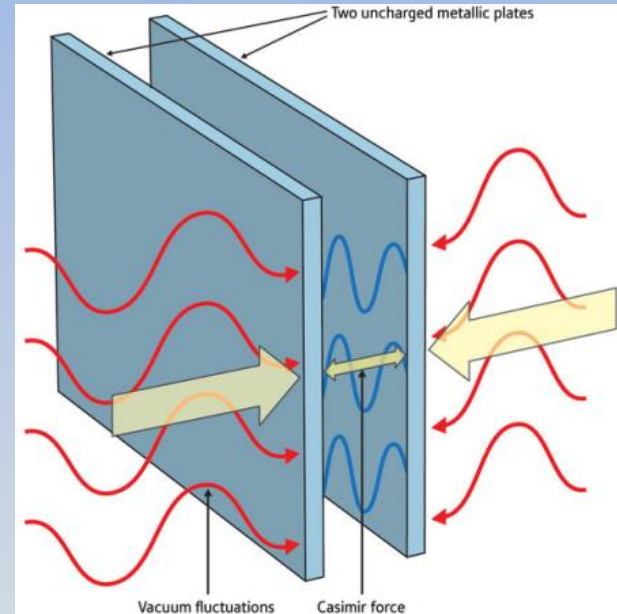
Main question still open with no experimental answer

Does vacuum fluctuations gravitate or not?

Does vacuum pressure exhibits the red-shift ?

The reality of macroscopic vacuum fluctuation.

The Casimir effect It is derived considering the zero point e.m. energy contained in a Casimir cavity, i.e. in the volume defined by two perfectly reflecting parallel plates



$$E = \sum \frac{1}{2} \hbar \omega$$

If the plates are perfectly reflecting the modes that can oscillate must have discrete wavenumbers on vertical axes $k_z = n\pi/a$ while all values are allowed for k_x e k_y

$$E = \frac{hcL^2}{2} \sum_{n=-\infty}^{n=\infty} \int \frac{d^2k}{(2\pi)^2} \sqrt{k^2 + \left(\frac{n\pi}{a}\right)^2} \longrightarrow \infty$$

The regularization is made by determining the Casimir Energy as the change in energy when the plates are at distance “a” with respect to the plates having $a \rightarrow \infty$

$$E_{\text{reg}} = E(a) - E(\infty)$$

- **Casimir Energy**
$$E_{\text{reg}} = -\frac{\pi^2 L^2 \hbar c}{720 a^3}$$

- **Casimir Pressure**
$$P_C = \frac{1}{L^2} \frac{\partial U}{\partial a} = -\frac{\pi^2 \hbar c}{240 a^4} = 1.3 \times 10^{-3} \text{ N/m}^2 (1 \mu\text{m}/a^4)$$

First prediction: Casimir 1948

First measure (force): Sparnay 1956

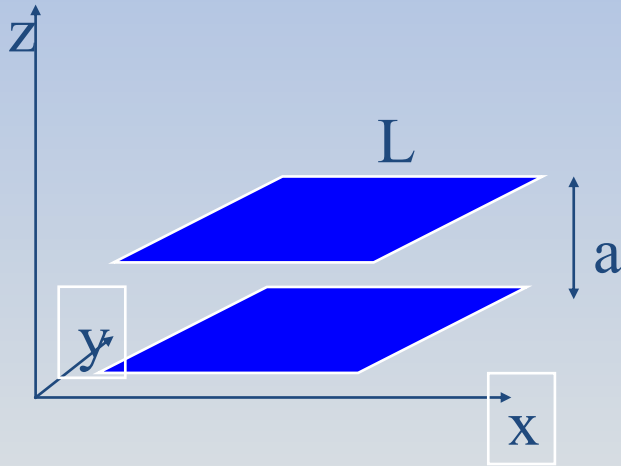
First measure (force) in the original flat-flat configuration: Carugno: 2002

Presently tested (force) with an accuracy of 0.5% (Mohideen: 2005)

(No problems in QFT in flat space-time)

Weighing the vacuum

The idea is to weigh a **rigid** Casimir cavity when the vacuum energy is modulated by changing the reflectivity of the plates. The forces along z are

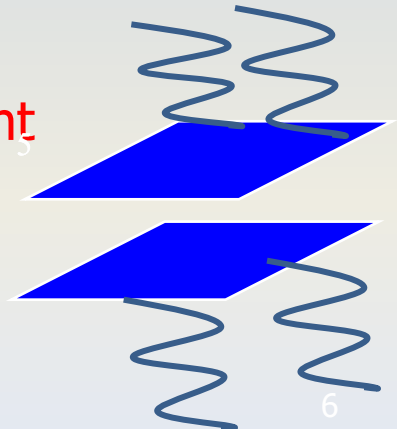


$$f_v = -\frac{1}{\sqrt{-\det g}} \frac{\partial}{\partial x^\mu} (\sqrt{-\det g} T_v^\mu) + \frac{1}{2} \frac{\partial g_{\rho\sigma}}{\partial x^\nu} T^{\rho\sigma},$$

$$\vec{F}_{tot} = \frac{|E_C|}{c^2} g \hat{z}$$

The total force is directed upward and it is equal to the weight of the vacuum modes that are removed from the cavity

IN ANALOGY WITH ARCHIMEDES FORCE



Vacuum fluctuation force on a rigid Casimir cavity in a gravitational field

E. Calloni^{a,b}, L. Di Fiore^b, G. Esposito^{b,*}, L. Milano^{a,b}, L. Rosa^{a,b}

^a *Università Federico II di Napoli, Complesso Universitario di Monte Sant' Angelo, Via Cintia, 80126 Naples, Italy*

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Received 24 September 2001; received in revised form 4 April 2002; accepted 5 April 2002

PHYSICAL REVIEW D **74**, 085011 (2006)

Energy-momentum tensor for a Casimir apparatus in a weak gravitational field

Giuseppe Bimonte,^{*} Enrico Calloni,[†] Giampiero Esposito,[‡] and Luigi Rosa[§]

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INFN, Sezione di Napoli, Complesso Universitario di Monte S. Angelo, Via Cintia, Edificio N', 80126 Napoli, Italy

(Received 6 June 2006; published 10 October 2006)

PHYSICAL REVIEW D **76**, 025008 (2007)

Relativistic mechanics of Casimir apparatuses in a weak gravitational field

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and INFN, Sezione di Napoli, Napoli, Italy

(Received 7 March 2007; published 6 July 2007)

Experimental problem: modulate Casimir energy without exchanging too much energy with the system (to not destroy the possibility of measurement and control) and measure it.

LASER



The energy E sent to the film is about $5 \times 10^5 \text{ J}$

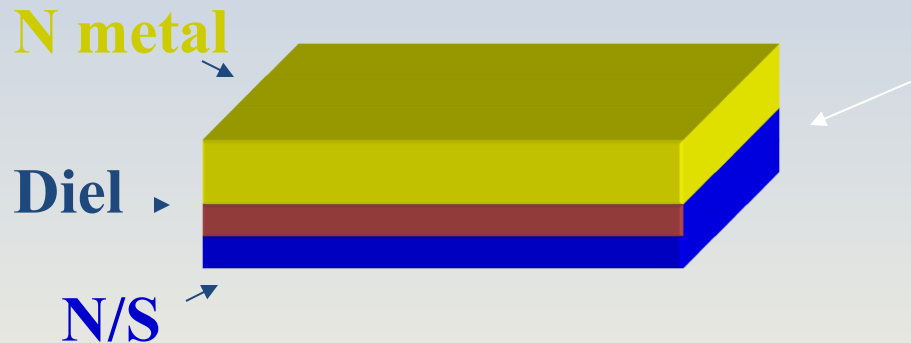
The variation of casimir energy ΔE_{cas} is about $2.5 \times 10^{-19} \text{ J}$

The efficiency is $\varepsilon = \Delta E_{\text{cas}}/E = 10^{-14}$

Modulation of Casimir energy with superconductors (type I)

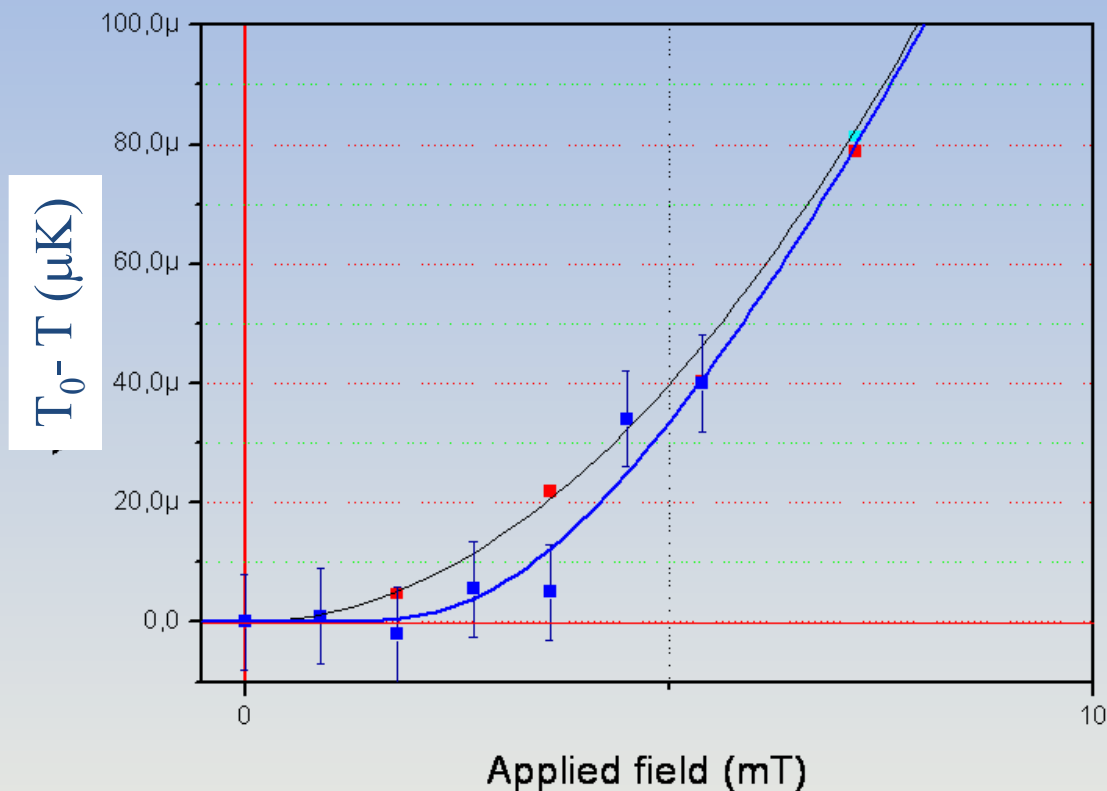
Use of superconductors

- The condensation energy is very small so it can be expected that the variation of Casimir energy at the transition for a superconductor inside a cavity can be of the same order, or even dominates, the total transition energy



$$\frac{V}{2\mu_0} \left(\frac{B_{c\parallel}^{\text{cav}}(T)}{\rho} \right)^2 = \mathcal{E}_{\text{cond}}(T) + \Delta E_{\text{cas}}(T)$$

Results and references on energy modulation – Aladin Experiment



The data are compatible with the theory and the region of energy of different behaviour is the expected one

G. Bimonte et Al. - J. Phys. A: Math. Theor. **41** 164023 (2008)

A. Allocca et Al. Jour. Of. Supercond. And Novel Magnetism. **25**, 2557-2565 (2012)

Use of Stratified high T_c superconductors

IOP PUBLISHING

JOURNAL OF PHYSICS A: MATHEMATICAL AND THEORETICAL

J. Phys. A: Math. Theor. **41** (2008) 164038 (5pp)

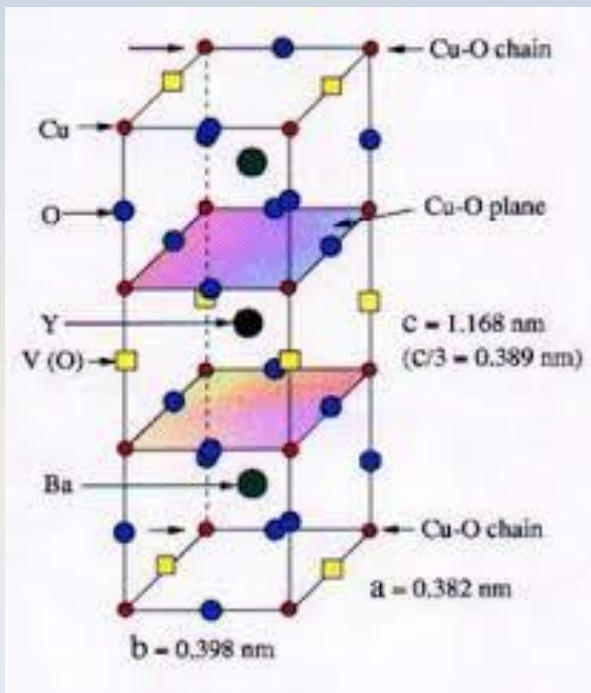
doi:10.1088/1751-8113/41/16/164038

On the Casimir effect in the high- T_c cuprates

Achim Kempf

Departments of Applied Mathematics and Physics, University of Waterloo, ON N2L 3G1, Canada

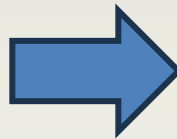
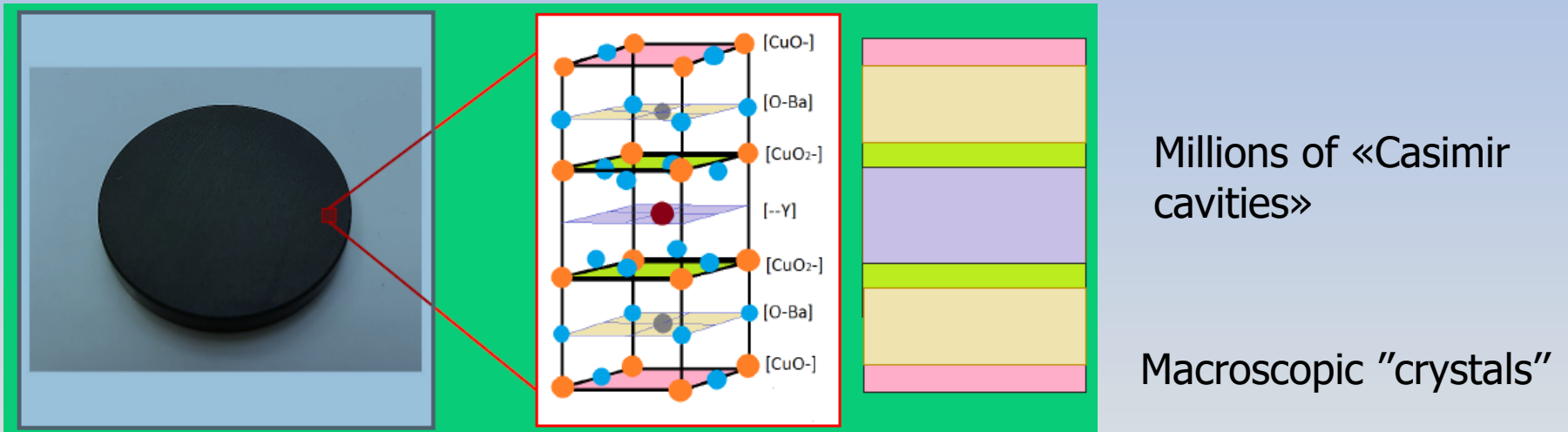
In normal state the plane (that will become superconducting) is a very poor conductor \rightarrow good variation of Casimir energy at the transition



$$\Delta\eta_E = \frac{\Delta E_C}{E_C} \approx 10^{-4}$$

The variation of casimir energy can be a significant fraction of the whole energy condensation

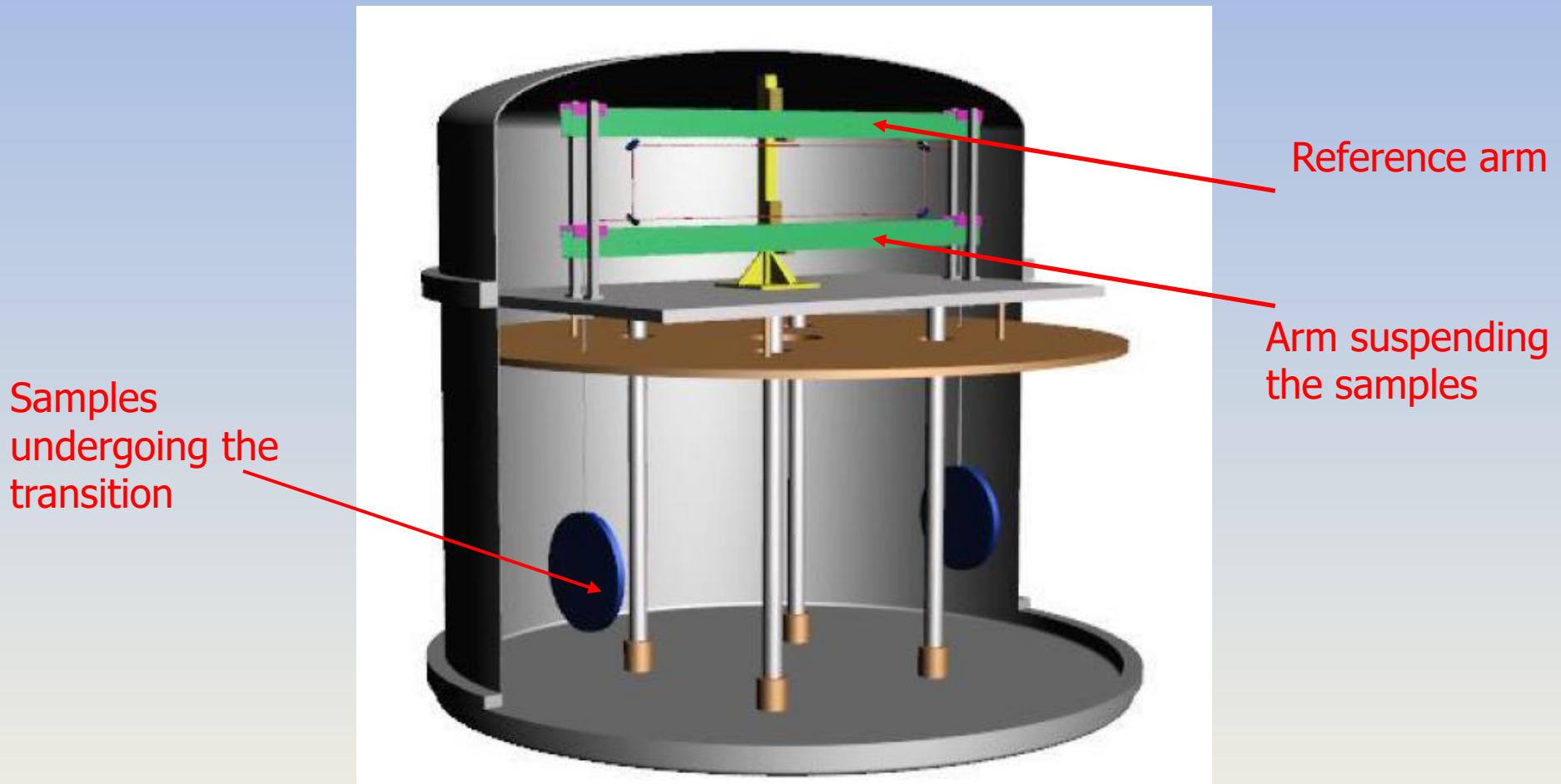
Experiment method : weighing the condensation energy of Type II superconductors and modulate the transition to modulate the weight



Expected weight modulation

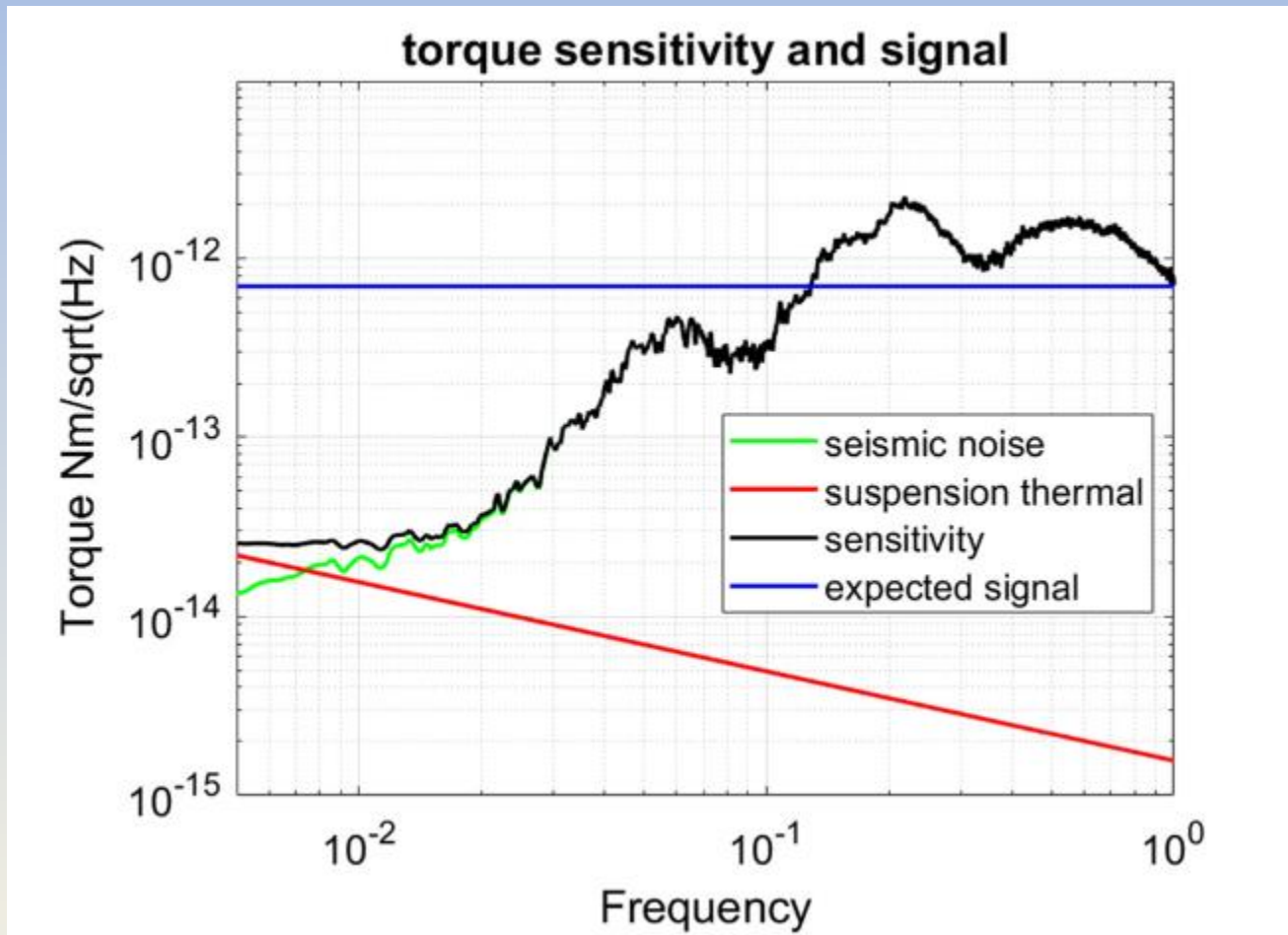
$$F = 5 \cdot 10^{-16} \text{ N}$$

Measurement scheme



Use a beam-balance → modulate the force by modulating the temperature of the superconductor so that it makes transitions between Normal and superconducting state - Expected modulation of force $F = 5 \cdot 10^{-16} \text{ N}$

Expected sensitivity



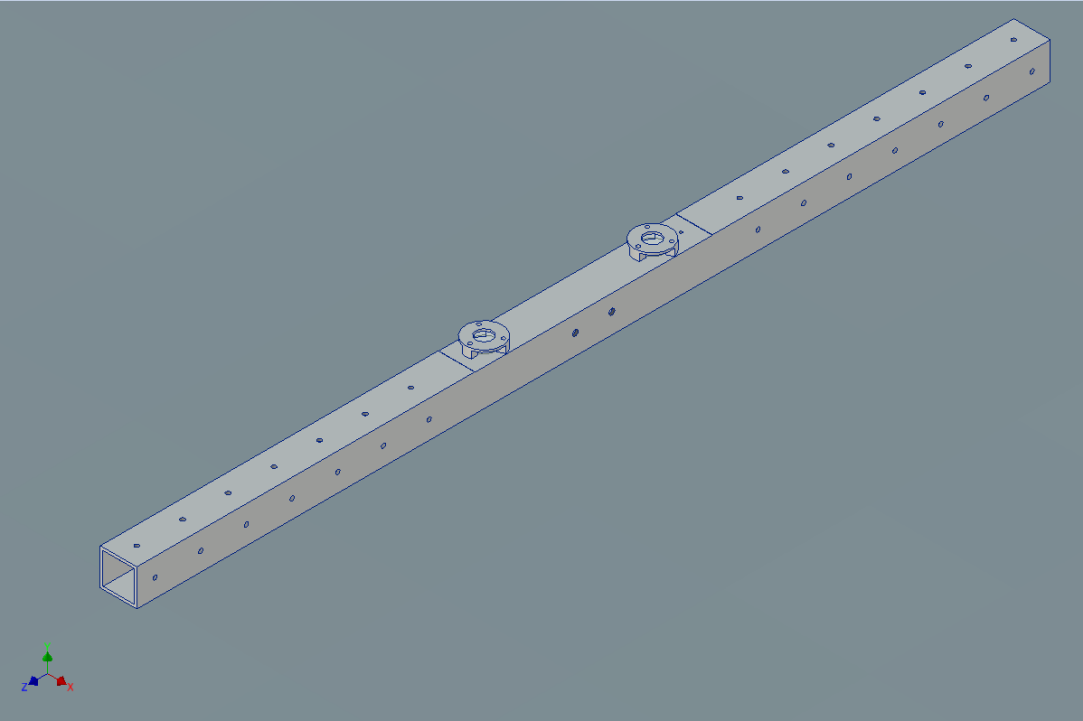
Sensitivity and signal in one-month integration time

The experiment

What is needed

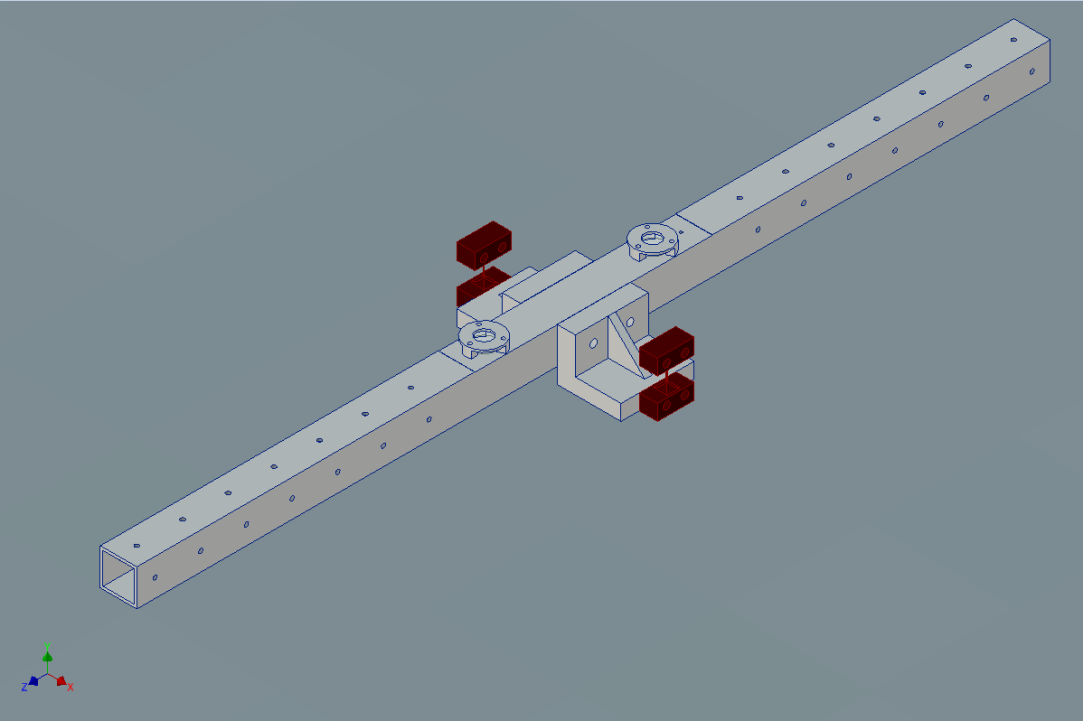
- ☐ An extremely sensitive cryogenic balance
- ☐ A suitable modulation technique & sample choice
- ☐ A very low seismic noise site

Balance prototype



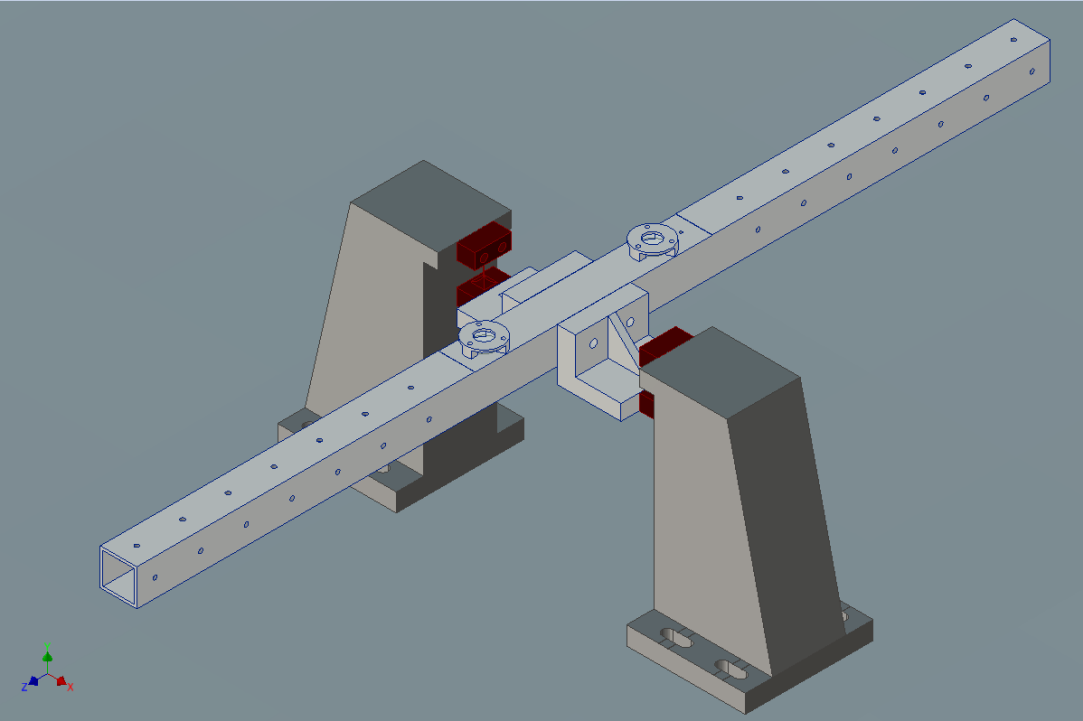
- 50 cm long arm with low momentum of inertia

Balance prototype



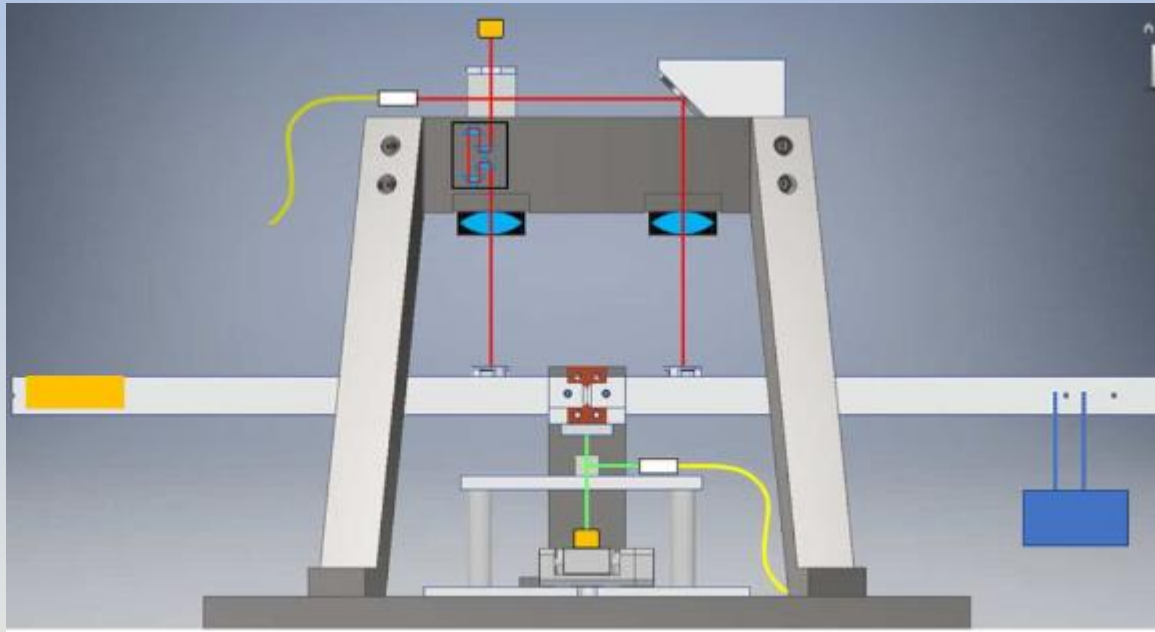
- 50 cm long arm with low momentum of inertia
- Suspended through **thin flexible joints** (Cu-Be, $100\ \mu\text{m} \times 500\ \mu\text{m}$), very similar in design to LIGO tiltmeters (Venkateswara et al., 2014)

Balance prototype



- 50 cm long arm with low momentum of inertia
- Suspended through **thin flexible joints** (Cu-Be, $100\ \mu\text{m} \times 500\ \mu\text{m}$), very similar in design to LIGO tiltmeters (Venkateswara et al., 2014)
- The balance **center of mass** is positioned **as close as possible to the bending point** ($\approx 10\ \mu\text{m}$)

Optical Read-Out

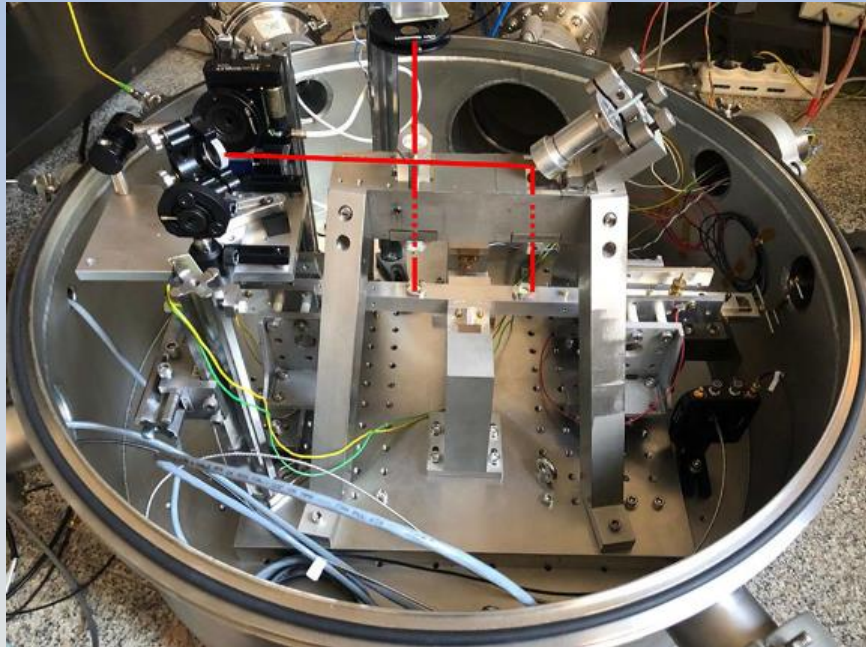


The interferometer is kept in the working point with a feed-back system

Tilts w.r.t. the ground
are read finely with a
**Michelson
interferometer**

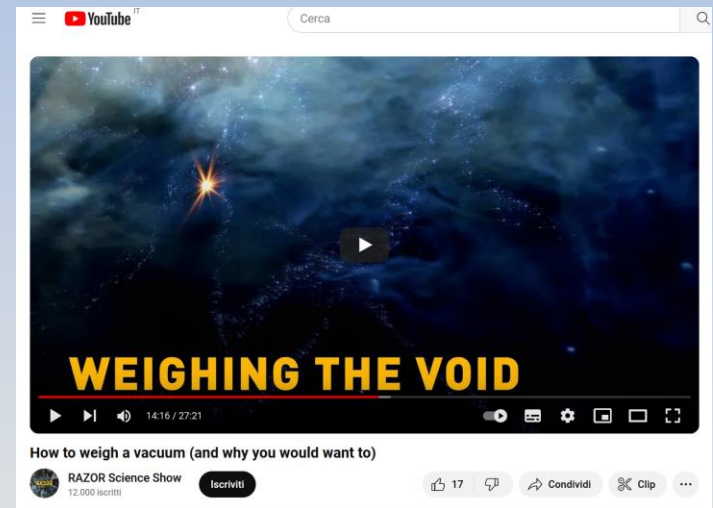
Optical lever only as
coarse positioning
sensor

The Balance prototype



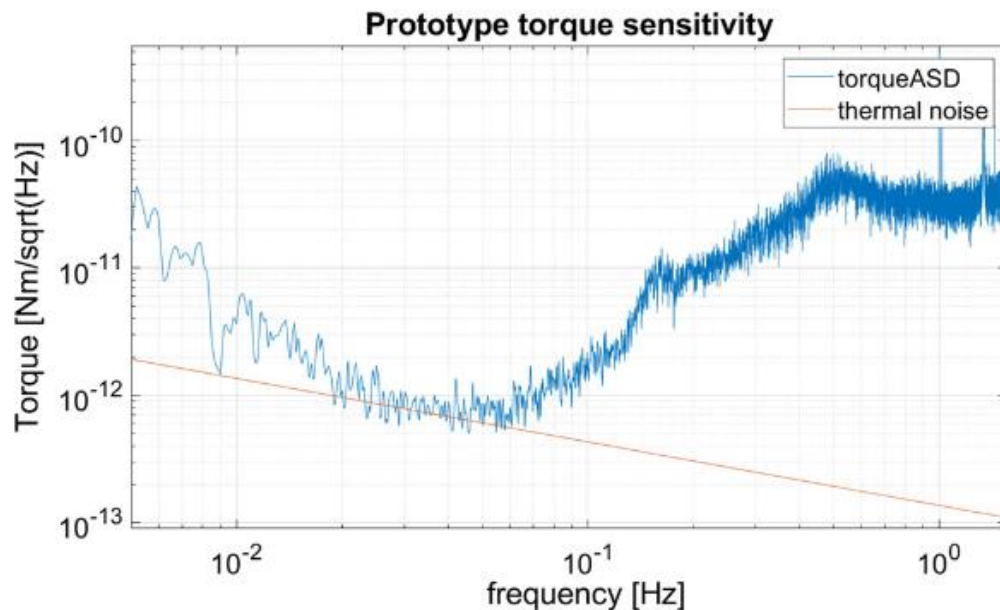
Laser light path in red

[How to weigh a vacuum \(and why you would want to\) \(youtube.com\)](https://www.youtube.com/watch?v=oWDZBUeBQzk)

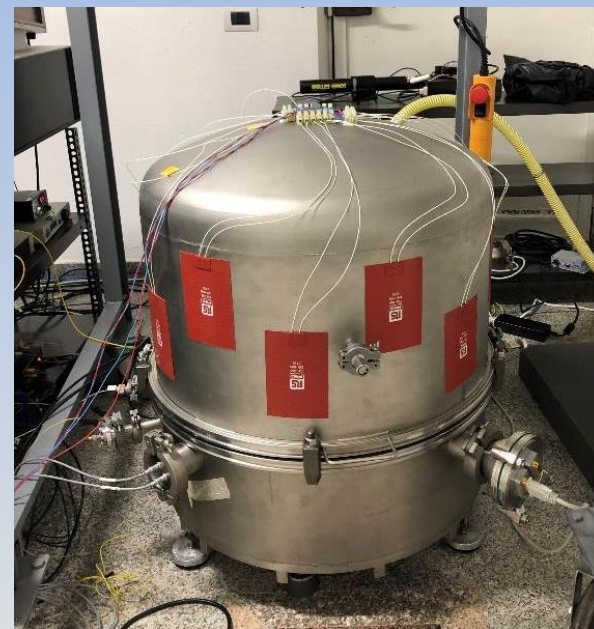


<https://youtu.be/oWDZBUeBQzk>
14.56

Prototype results



Best sensitivity of the balance Prototype – compatible with thermal noise limit



The prototype balance with thermal patches

Lessons Learned

- 1) Verification of efficiency of optical read-out and control systems
- 2) Importance of low environmental noises (pumps, air conditioning, seismic noise)
- 3) Importance of environmental temperature stability

The final apparatus: cryostat

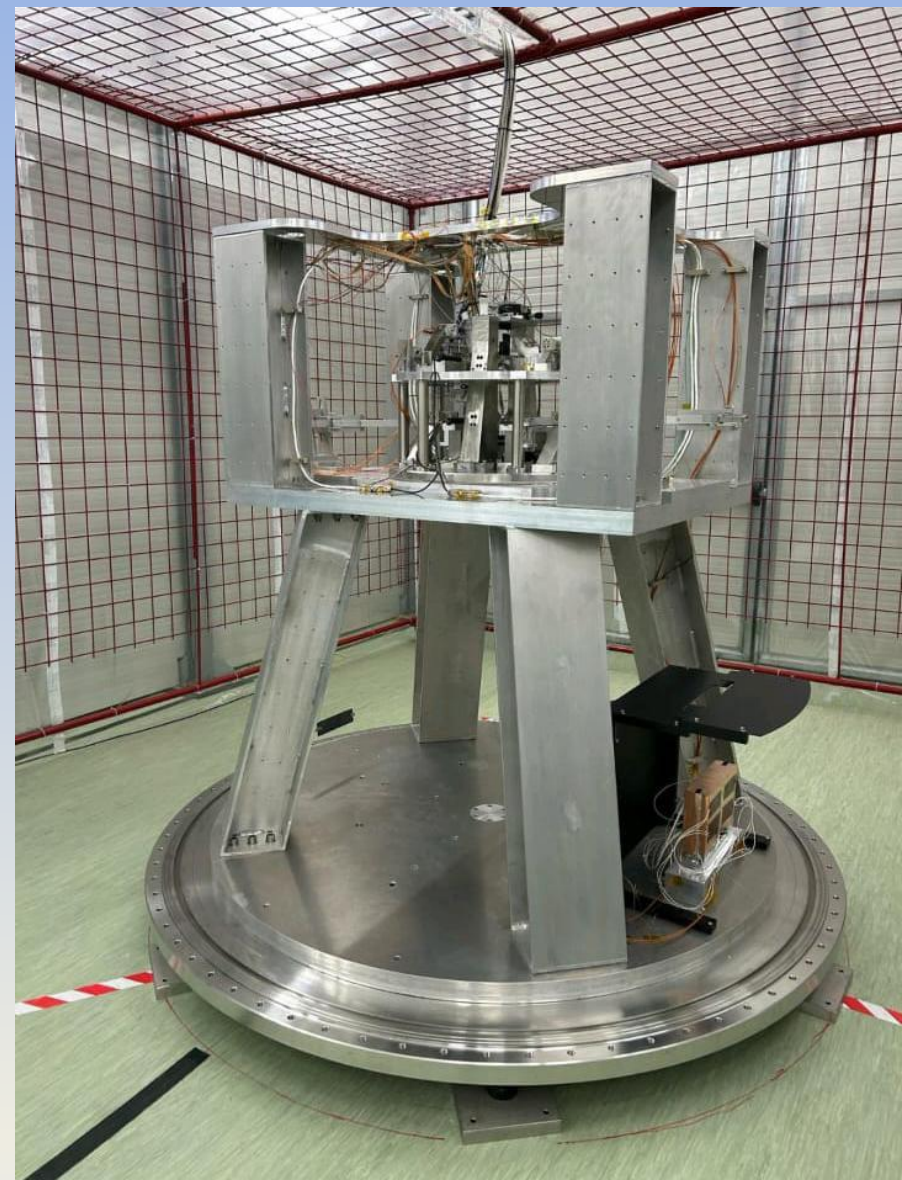


The cryostat has been completed and tested by the end of July. It is now ready to be sent to the site. Presently we are considering where to have it sent. Indeed timescale requires that we start the cryogenic activities in the temporary building, so for the moment the cryostat is hosted in the factory. 23

The final balance and the start of the commissioning



Search for spurious electromagnetic signals



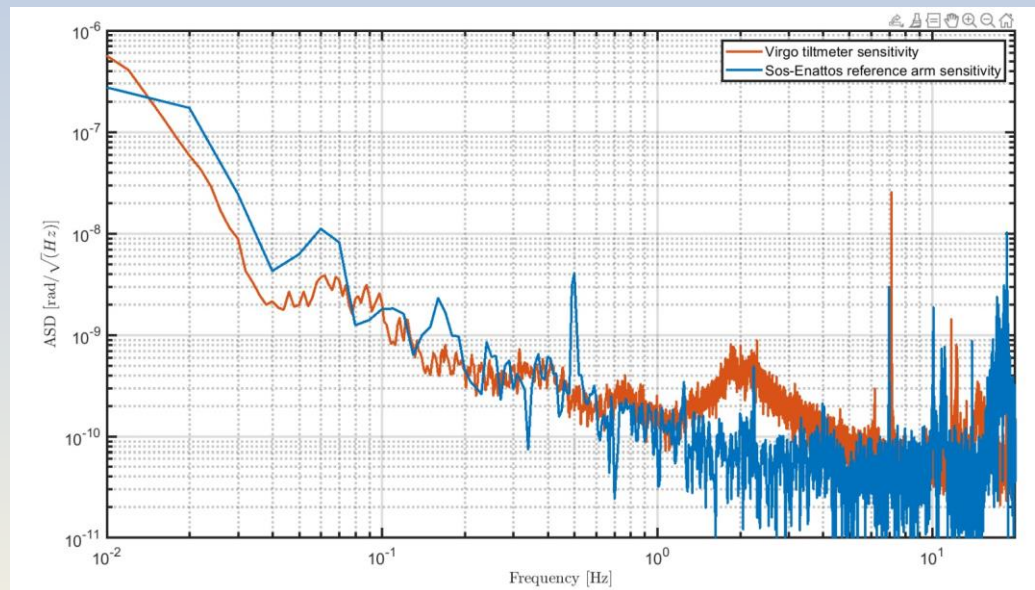
The balance suspending two different test samples – one sample is inside the oven

The very first data of the balance under vacuum

- The lock is acquired and is stable with the reference arm suspended (not obvious because the reference arm is optically more complex with respect to measurement arm)
- All connections tested – no major noise problems
- Accident....re-opened last week – new closure expected for mid July



Experimental «accident» : the T of the measurement arm «jumped» over the fork



Good noise of the reference – equal or better than in Virgo – but just a hint, because it is affected by the noise of the measurement arm with the T touching the fork

Casimir Energy Modulation

The method is to have the sample in a thermal bath, modulate «slowly» the temperature of the bath so that the sample makes a reversible transformation from normal to superconductor and viceversa

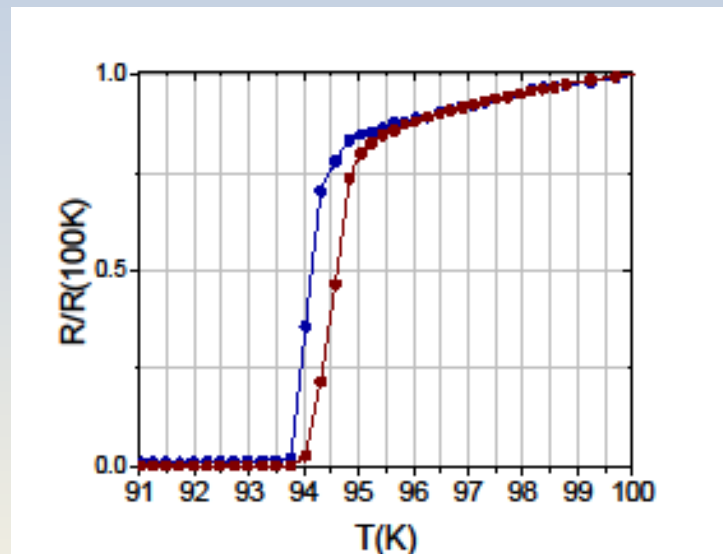


Superconductor choice: GdBCO

- GdBCO: **single crystal** already realized in discs of 10 cm diameter and tests on 15 cm diameter - acquired 1) a 10 cm disc and 2) several smaller discs.
- BSCCO: pressed powder (polycrystalline) realized in discs/rings 10 cm diameter – acquired several small samples tested with negative results – pressed powder not suitable for the experiment



GdBCO sample



GdBCO: transitions of two pieces of the same sample $\Delta T_c = 0.5 K$

Need of a quiet site!

Low seismic – No antropic noise

Sindaco Mario Calia (lista civica)

dall'11-6-2012

Territorio

Coordinate 40°28'N 9°29'E

Altitudine 521 m s.l.m.

Superficie 148,72 km²

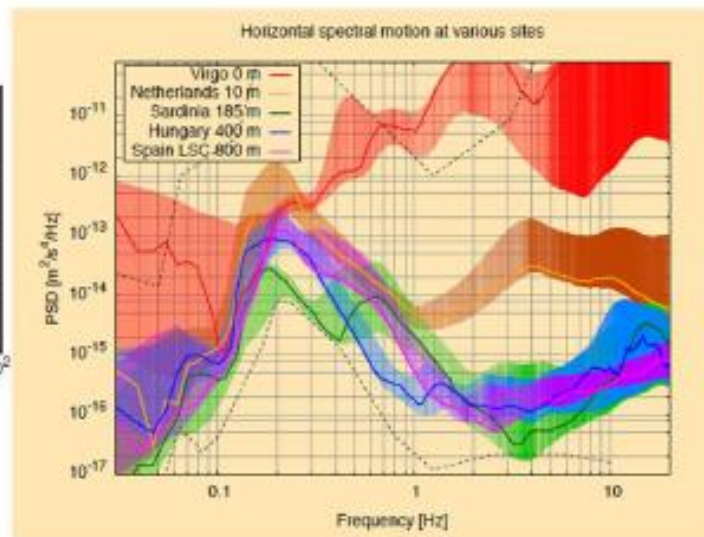
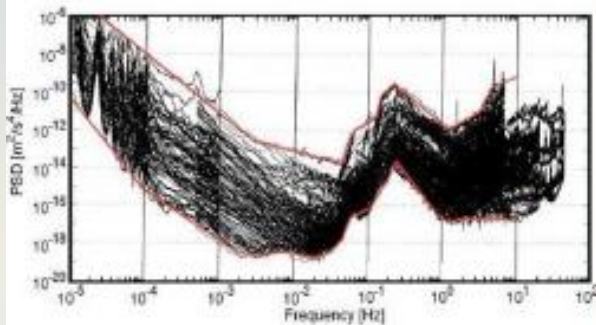
Abitanti 1 407^[1] (31-7-2016)

Densità 9,46 ab./km²

Comuni confinanti Bitti, Dorgali, Galtelli, Irgoli,
Loculi, Lodè, Onani, Orune,
Siniscola



SOS-Enattos Mine



Seismic
Measurements
By Virgo and
ET collaborations

The present lab in Sos Enattos



The need of finding an alternative site

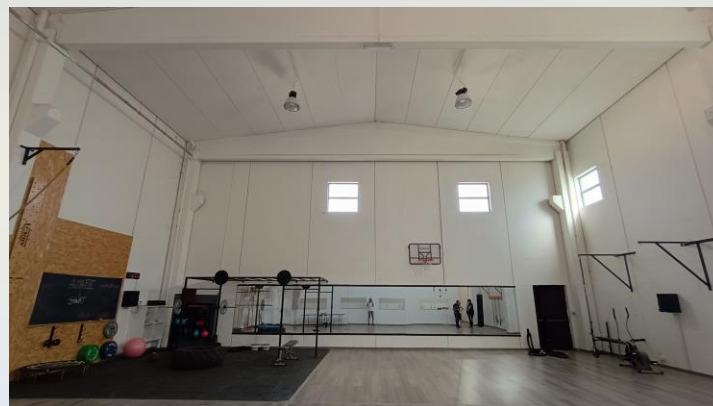


The present
building

The future
ET_SUnLab

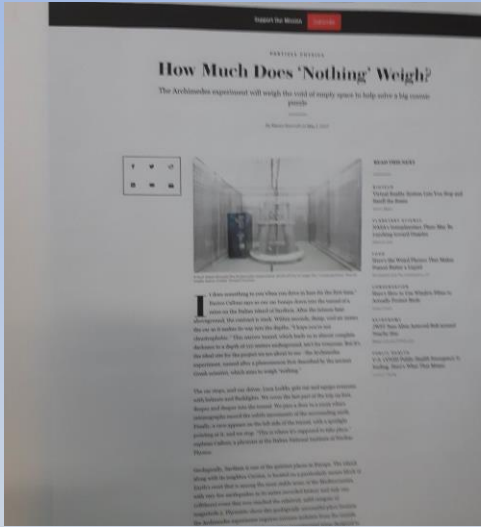


Next year in a temporary building in Orosei



Archimedes Falls out

Make Sos Enattos a known site for Einstein Telescope



Scientific American

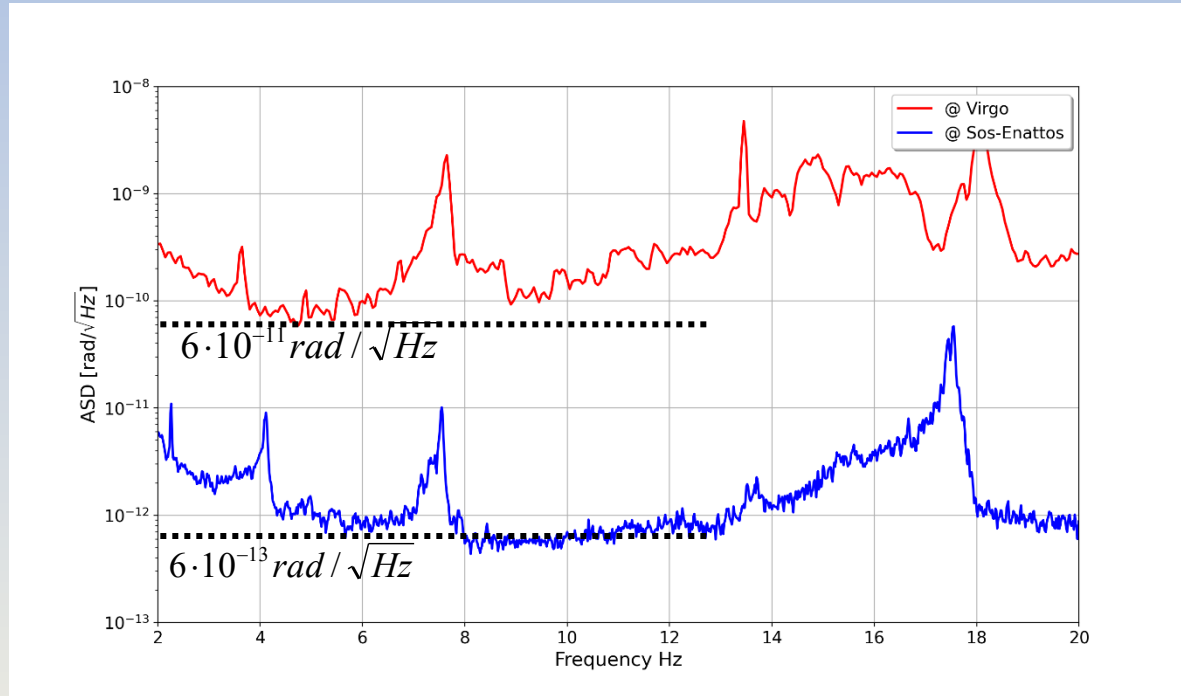


Le figaro



Focus

Tiltmeter – the best tiltmeter in few Hz regime



❑ Best sensitivity in the world for a tiltmeter in the region 2 Hz – 20 Hz

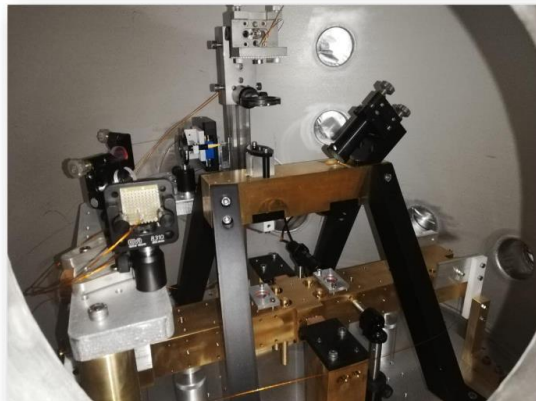
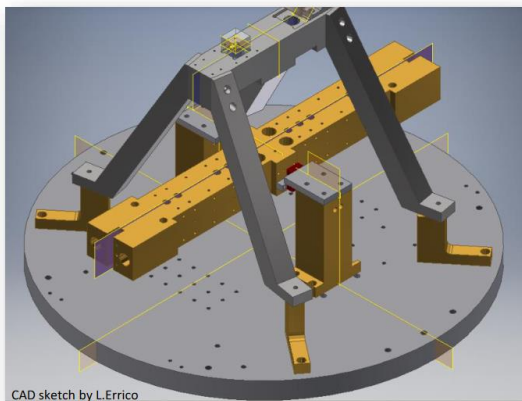
Tiltmeter installed in the Virgo GW observatory

A new tiltmeter has been realized with Virgo funds and installed in Virgo at the NE tower

ῬΑκίνητος: the new tiltmeter

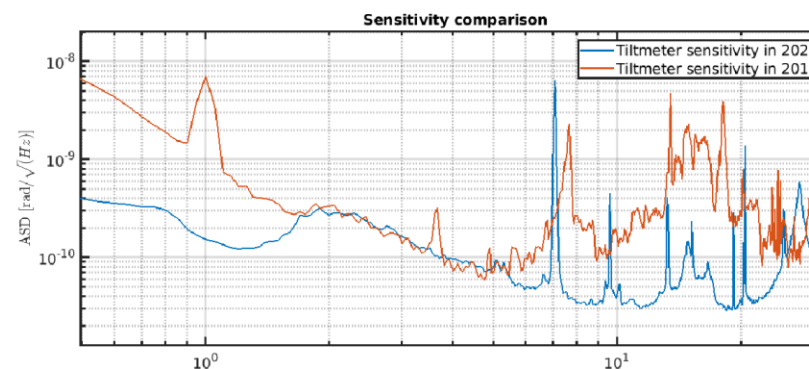
The new tiltmeter exploits the same working principles as the prototype.

Main improvement: arm with much higher momentum of inertia: 13 kg of brass, $I = 0.33 \text{ kg} \cdot \text{m}^2$, more than one order of magnitude bigger than the previous version, joint size: $0.1 \times 3 \text{ mm}$



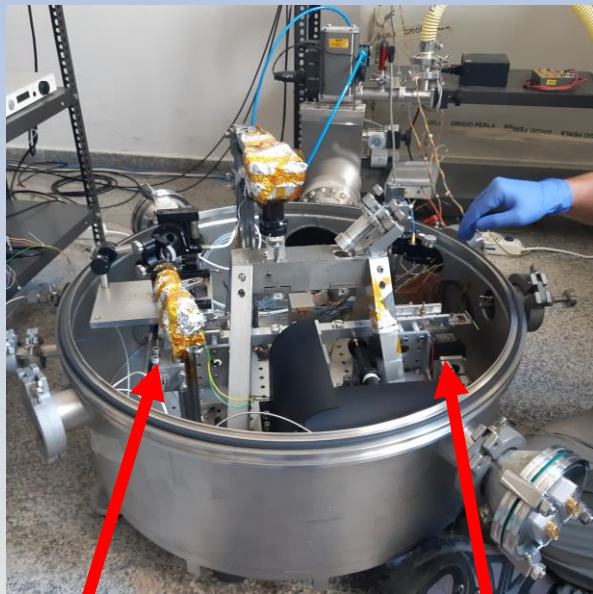
Comparison between tiltmeters sensitivity in 2019 vs 2023

About $8 \text{ e-}11 \text{ @} 1 \text{ Hz}$,
 $3 \text{ e-}11 \text{ @} 10 \text{ Hz}$

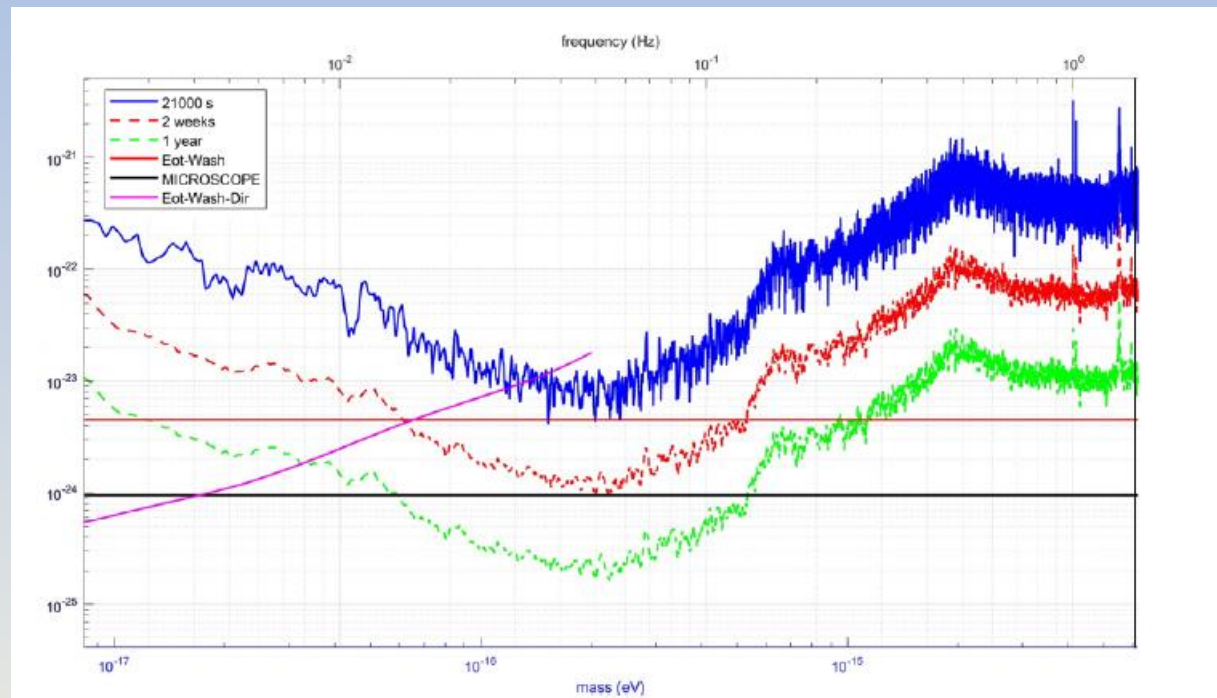


Fall-out activities: direct search for ultralight B-L dark photon

The expected signal on the balance is a monochromatic noise at the frequency $f = mc^2/h$



Brass (now Pb) aluminum



1 night of measurement !

*Notice that the EOT-WASH and MICROSCOPE are testing the eventual existence of the dark photon as a particle but they are not testing the fact that it is a dark matter component
On the other hand the direct search can not disprove the existence of such a particle but only tests if it is the constituent of the dark matter*

Conclusion

- Archimedes will run at room temperature in the next months
- Hopefully it will perform the final measurement in the next 3-4 years

Thanks

SPARES

Vacuum fluctuation force on a rigid Casimir cavity in a gravitational field

E. Calloni ^{a b}, L. Di Fiore ^b, G. Esposito ^b ✉, L. Milano ^{a b}, L. Rosa ^{a b}



Torre del Greco - NA

sive gravitational mass). In particular, we agree with the statement that a system with a given rest energy momentum tensor $T^{\bar{0}\bar{0}}$ has the inertial mass tensor $m^{ij} = T^{\bar{0}\bar{0}}\delta^{ij} + T^{\bar{i}\bar{j}}$ [12].

Unfortunately, at present, there is not yet a direct calculation of inertial mass of a Casimir cavity in 4 dimensions; nevertheless direct calculations of inertial mass on a two-dimensional space–time [13] are in agreement with our calculation if “extended”

to four dimensions. Furthermore, it is very important to note that, when considering the total force acting on the real cavity, which is an isolated system, the contribution to the force resulting from the spatial part of the energy-momentum tensor is balanced by the contribution of the mechanical energy-momentum tensor, and hence should not be considered for experimental evaluation. The resulting force

is then the Newtonian force on the sum of the rest Casimir energy and rest mechanical mass: the contribution of vacuum fluctuations leads to a gravitational push on the Casimir apparatus expressed by the formula

$$\vec{F} \cong \frac{\pi^2 L^2 \hbar c}{720 a^3} \frac{g}{c^2} \vec{e}_r, \quad (9)$$

Simulations

If the sample has the geometry of a ring and the radiating surface is increased by a thin foil then the modulation times are compatible with expectations --- **Simulations on going with Comsol and Ansys**

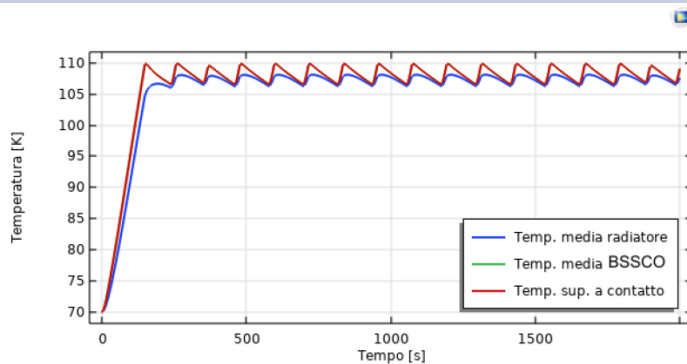


Figura 1.13: Risultati simulazione per la configurazione in Figura 1.1 e con parametri descritti nella sezione 1.2.3. La linea blu è la temperatura media del radiatore, in verde la temperatura media del campione mentre in rosso la temperatura media superficiale della corona a contatto termico.

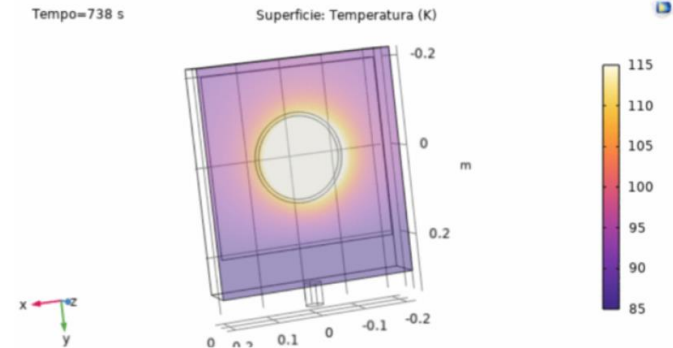
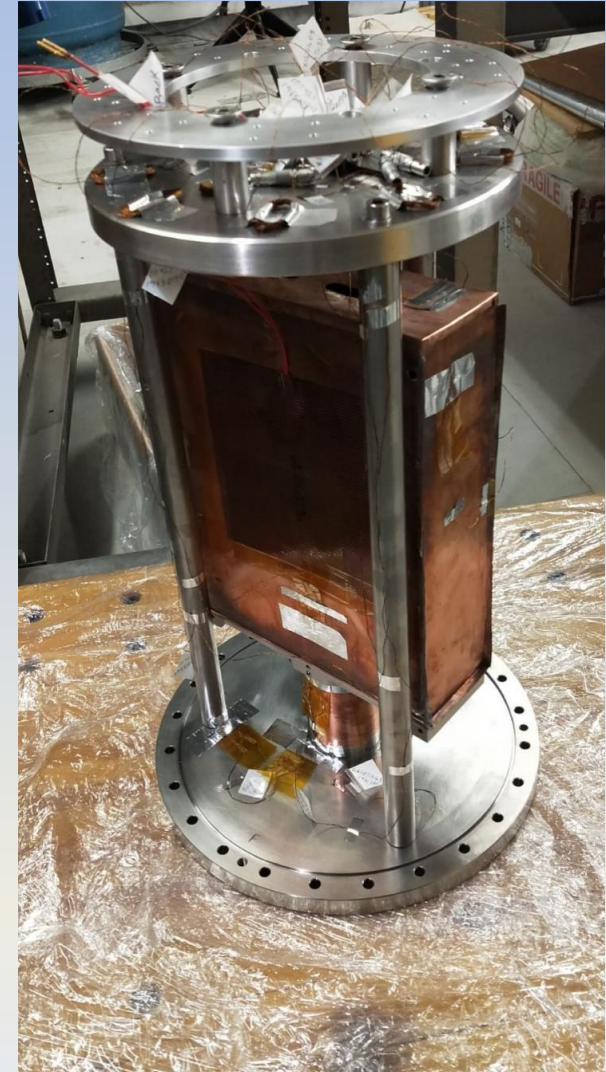
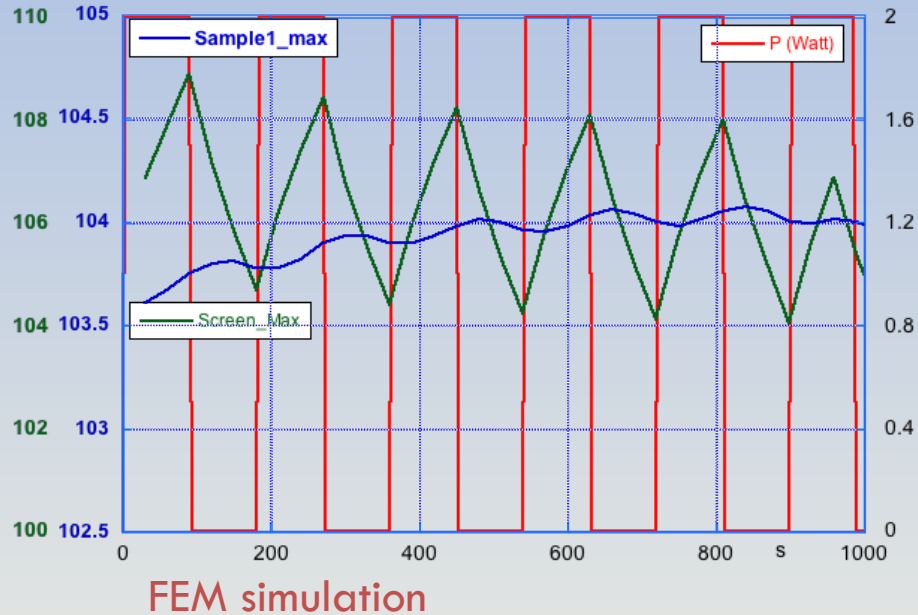


Figura 1.14: Profilo in temperatura della superficie del forno ad un certo istante di tempo $t = 738$ s corrispondente alla fase di riscaldamento. La figura mostra il profilo del fascio incidente sul forno.

BSSCO $\rightarrow T_c = 110$ K ; Modulation Period $T = 120$ s; Laser Power $P = 10$ W;
Sample volume 20 cm^3 -- Similar for GdBCO but more critical due to $T_c = 94$ K₁

Temperature modulation: test in Rome cryogenic temperature



Temperature modulation: test at **EGO** (European gravitational Observatory) at room temperature



*Test of the Archimedes
modulation system at room
temperature*



EGO site hosting the Virgo GW detector

Pressure red-shift

A simple summation of the lower force and upper force on the plates would bring to a somewhat unsuspected result:

$$F_{\text{inf}} + F_{\text{sup}} = F_{\text{cas}} (1 + \delta\phi) + \frac{|E_C|}{c^2} g - F_{\text{cas}} = 4 \frac{|E_C|}{c^2} g$$

$$F_{\text{cas}} = -L^2 \frac{\pi \hbar c}{240 a^4}$$

$$E_{\text{cas}} = -L^2 \frac{\pi \hbar c}{720 a^3}$$

The lower vacuum «photons» must exert a bigger force because the force will be red-shifted when reaching the same level of upper plate → in the experiment the sum must be done taking into account the red-shift because the cavity is rigid and hanged in a unique point - **(for this effect our measurement is a null measurement on pressure red-shift)**

$$\left\{ \begin{array}{l} F_{\text{sup}} = F_C \\ F_{\text{inf}} = -F_C (1 + \delta\phi) + \frac{|E_C|}{c^2} g \end{array} \right. \longleftrightarrow \vec{F}_{\text{tot}} = \frac{|E_C|}{c^2} g \hat{z}$$

E. Calloni et.al. Phys. Letters A, 297, 328-333, (2002)

G. Bimonte, E. Calloni, G. Esposito, L. Rosa - Phys. Rev D 74, 085011 (2006)

G. Bimonte, E. Calloni et. al. Phys.Rev.D76:025008, (2007)

G. Bimonte, E. Calloni, L. Rosa, Phys.Rev.D77:044026, (2008)

5

On interpretation of Tolman-Ehrenfest effect:

C. Rovelli, M. Smerlak Class. Quant. Grav. 28 (2011) 075007

Casimir Energy variation in real material

$$\Delta E_{\text{cas}}^0(a, d) = A \frac{\hbar}{2} \int \frac{dk_1 dk_2}{(2\pi)^2} \left\{ \sum_p (\omega_{\mathbf{k}_\perp, p}^{(n, TM)} + \omega_{\mathbf{k}_\perp, p}^{(n, TE)}) - \sum_p (\omega_{\mathbf{k}_\perp, p}^{(s, TM)} + \omega_{\mathbf{k}_\perp, p}^{(s, TE)}) \right\} \quad (1)$$

where $A \gg a^2$ is the area of the cavity, $\mathbf{k}_\perp = (k_1, k_2)$ denotes the two-dimensional wave vector in the xy plane, while $\omega_{\mathbf{k}_\perp, p}^{(n/s, TM)}$ ($\omega_{\mathbf{k}_\perp, p}^{(n/s, TE)}$) denote the proper frequencies of the TM (TE) modes, in the n/s states of the film, respectively.

Cauchy integral formula and renormalization



$$\left(\sum_p \omega_{\mathbf{k}_\perp, p}^{(n, TM)} - \sum_p \omega_{\mathbf{k}_\perp, p}^{(s, TM)} \right)_{\text{ren}} = \frac{1}{2\pi} \int_{-\infty}^{\infty} d\zeta \left(\log \frac{\Delta_n^{(1)}(i\zeta)}{\tilde{\Delta}_{n\infty}^{(1)}(i\zeta)} - \log \frac{\Delta_s^{(1)}(i\zeta)}{\tilde{\Delta}_{s\infty}^{(1)}(i\zeta)} \right)$$

where

$$\Delta_{jl}^{TE} = \frac{K_j - K_l}{K_j + K_l}, \quad \Delta_{jl}^{TM} = \frac{K_j \epsilon_l(i\zeta) - K_l \epsilon_j(i\zeta)}{K_j \epsilon_l(i\zeta) + K_l \epsilon_j(i\zeta)},$$

$$K_j = \sqrt{\epsilon_j(i\zeta) - 1 + p^2}, \quad I = n, s; \quad j, l = 1, 2, n, s.$$

And inserting it in (1)

$$\Delta E_{\text{cas}} = A \frac{\hbar}{2} \int \frac{d\mathbf{k}_{\perp}}{(2\pi)^2} \int_{-\infty}^{\infty} \frac{d\zeta}{2\pi} \left(\log \frac{Q_n^{TE}}{Q_s^{TE}} + \log \frac{Q_n^{TM}}{Q_s^{TM}} \right)$$

where

$$Q_I^{TE/TM}(\zeta, p)$$

$$= \frac{(1 - \Delta_{1I}^{TE/TM} \Delta_{12}^{TE/TM} e^{-2\zeta K_1 L/c})^2 - (\Delta_{1I}^{TE/TM} - \Delta_{12}^{TE/TM} e^{-2\zeta K_1 L/c})^2 e^{-2\zeta K_I D/c}}{1 - (\Delta_{1I}^{TE/TM})^2 e^{-2\zeta K_I D/c}}$$

Ragguagli teorici

L. Rosa et al.. Casimir energy for two and three superconducting coupled cavities
Submitted to JPLP (2017)

E. Calloni et al Towards weighing the condensation energy to ascertain the Archimedes force of vacuum
Phys.Rev. D90 (2014) no.2, 022002

G. Bimonte, E. Calloni et. al. Relativistic mechanics of Casimir apparatuses in a weak gravitational field
Phys.Rev.D76:025008, (2007)

G. Bimonte, E. Calloni , G. Esposito, L. Rosa - Energy-momentum tensor for a Casimir apparatus in a weak gravitational field Phys. Rev D 74, 085011 (2006)

G. Bimonte, E. Calloni et al Towards measuring variations of Casimir energy by a superconducting cavity
Phys.Rev.Lett. 94 (2005) 180402

E. Calloni et.al. Vacuum fluctuation force on a rigid Casimir cavity in a gravitational field Phys. Letters A, 297, 328-333, (2002)

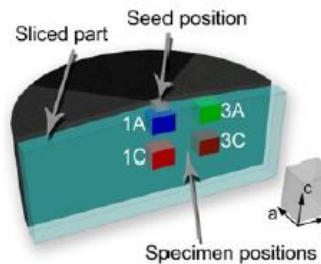
Da questo punto di vista, ritengo che la misura proposta da ARCHIMEDES sia di eccezionale interesse. Si tratta dell'unico esperimento di laboratorio di cui io sia a conoscenza che possa oggi dare informazioni sperimentali dirette sull'interazione fra energia di vuoto e gravità, e in particolare confermare l'esistenza di effetti fisici per i quali abbiamo oggi solo previsioni teoriche, in un contesto generale di notevole confusione teorica.

Superconduttività

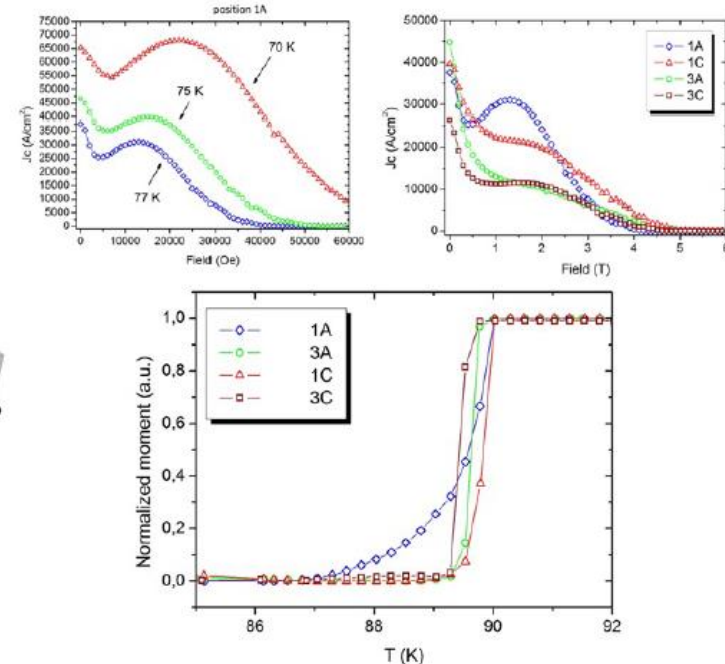
- ❑ E' stato aggiunto un WP all'esperimento ed aggiunti FTE efficaci.
- ❑ Varie ditte ri-contattate per YBCO



High T_c
for Pract



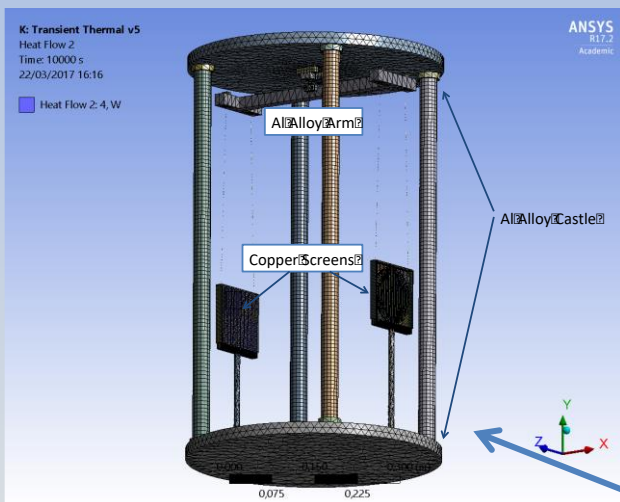
characterization for YBCO levitation bulk 14mm dia, 6mm



Inserita tra le milestone la realizzazione di un apparato per la misura di energia di transizione in dischi di grande dimensione. Questo è necessario per mantenere il controllo nel laboratorio delle caratteristiche dei materiali e perché ad una prima indagine non risulta che esista un tale apparato. Le indagini continueranno.

Attuazione termica (Roma)

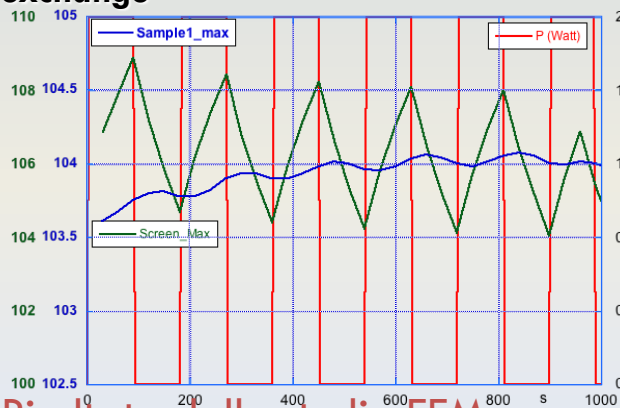
- Modulazione della temperatura del campione superconduttore attraverso il solo scambio radiativo per non aggiungere altra energia diversa da quella legata alle fluttuazioni del vuoto



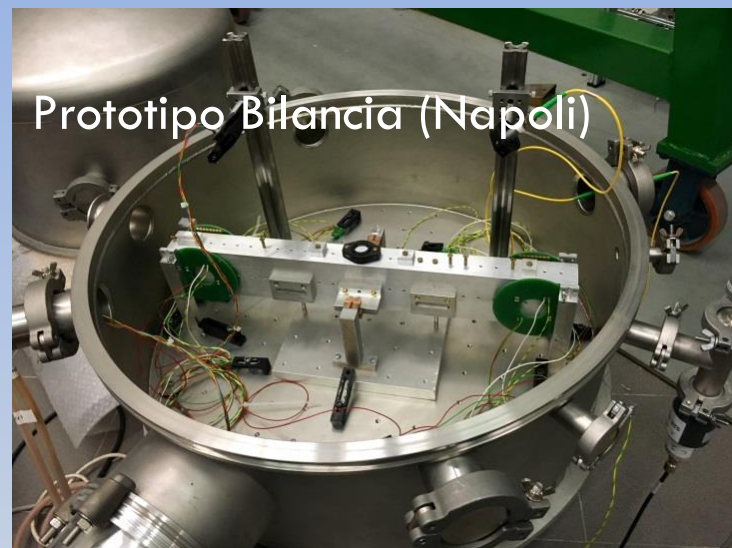
Criostato Virgo lab



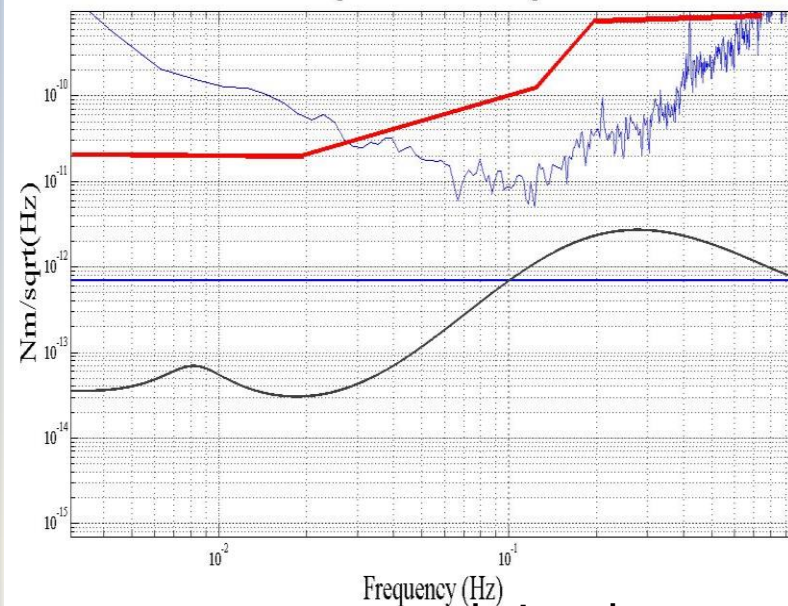
Case with only radiative heat exchange



Risultato dello studio FEM



Torque sensitivities and signal



- expected signal
- prototype
- Wash
- Archimedes noise

A side-measurement (more difficult): the weight of the entropy*T

$$\Delta U = G_n(T) + TS_n(T) - G(T, 0) - TS_s(T, 0)$$

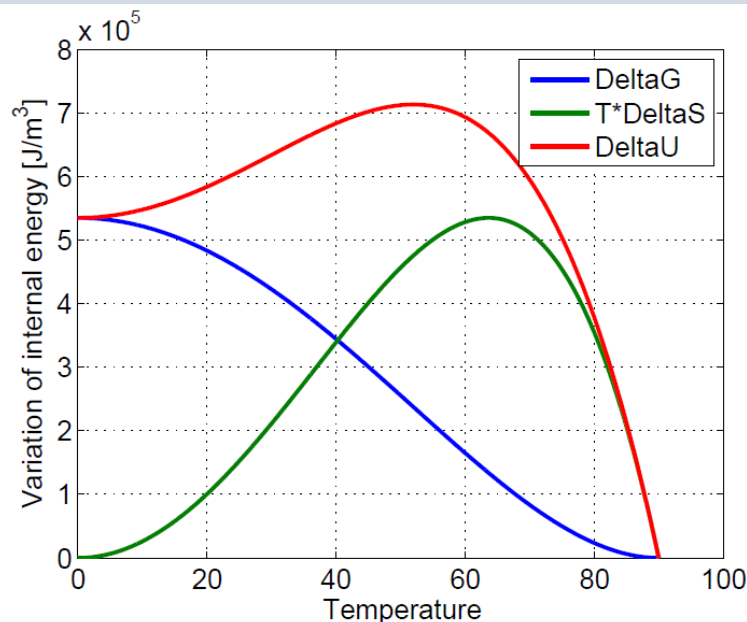
Difference in internal energy for a transition at fixed temperature

$$G_n(T) - G_s(T, 0) = \frac{1}{2\mu_0} B_c(T)^2$$

$$B_c(T) = B_0 \left[1 - \left(\frac{T}{T_c} \right)^2 \right]$$

$$S_n(T) - S_s(T) = 2 \frac{B_0^2}{\mu_0} \left(\frac{T}{T_c^2} \right) \left[1 - (T/T_c)^2 \right]$$

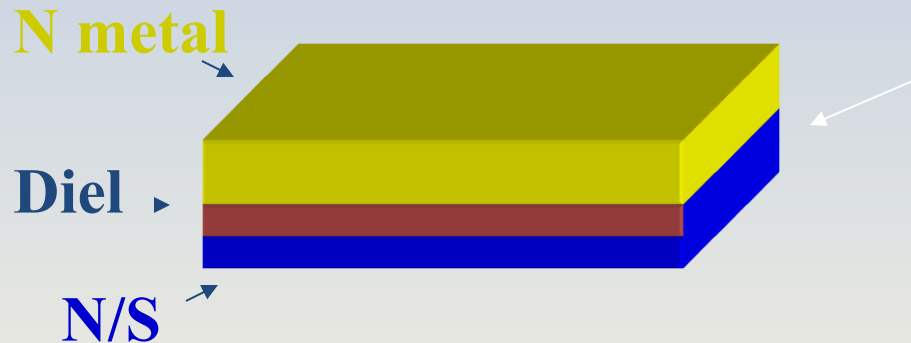
Difference in entropy for a transition at fixed temperature valid for BCS – approximatively for layered type II



*Disregarding in this particular discussion the contribution of Casimir effect the weight of the entropy (times Temperature) can be considered as an interesting **side-measurement** of the final experiment. This weight is classical (no question on it) but never measured*

Use of superconductors

- The condensation energy is very small so it can be expected that the variation of Casimir energy at the transition for a superconductor inside a cavity can be of the same order, or even dominates, the total transition energy

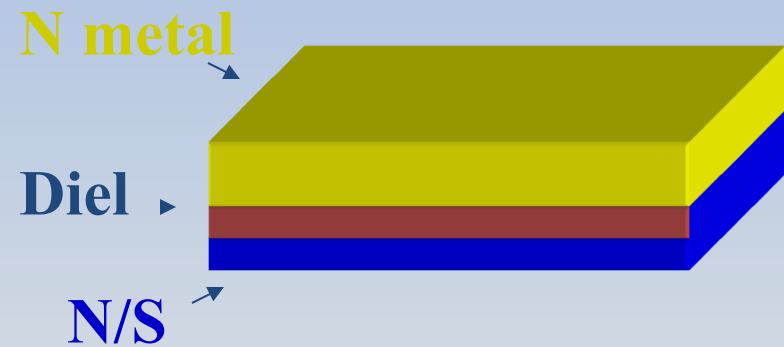


$$\frac{V}{2\mu_0} \left(\frac{B_{c\parallel}^{\text{cav}}(T)}{\rho} \right)^2 = \mathcal{E}_{\text{cond}}(T) + \Delta E_{\text{cas}}(T)$$

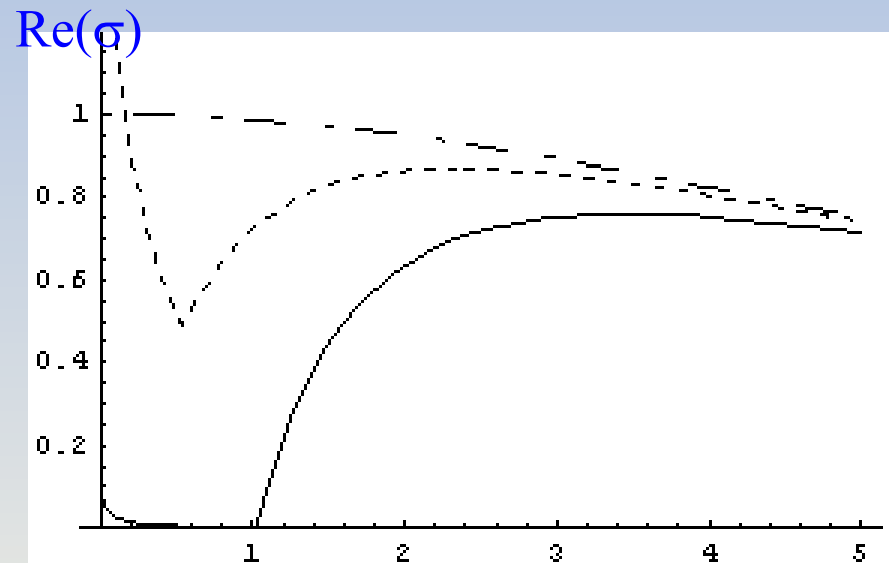
The change in energy can be calculated following the Casimir energy calculation in case of real plates with complex conductivity σ

$$\Delta E_C = -\Delta\eta_E \frac{\hbar c \pi^2}{720} \frac{A}{L^3}$$

$\Delta\eta_E$: modulation factor with respect perfect reflectivity



$$\Delta\eta_E = \frac{\Delta E_C}{E_C} \approx \frac{kT_C}{h\nu_C} = \frac{kT_C}{hc/L} \approx 10^{-6}$$



$$x = \hbar\omega/(2kT_c)$$

Plot of real part of conductivity σ normalized to zero frequency Drude conductivity σ_0 for different temperatures:

— · — $T = T_c$ (Drude) $T/T_c = 0.9$ — $T/T_c = 0.3$

The conductivity changes only in the very low frequency region (microwave) so the modulation depth (if T_c is of the order of 1 K) is expected to be small for small T_c ??.

..but also the energy exchanged with the system, besides the vacuum energy, is expected to be small being linked to the condensation energy which is (roughly) proportional to T_c^2 . Better to use low T_c superconductors. If the two energy variations are comparable then it is expected that vacuum fluctuations modifies the transition

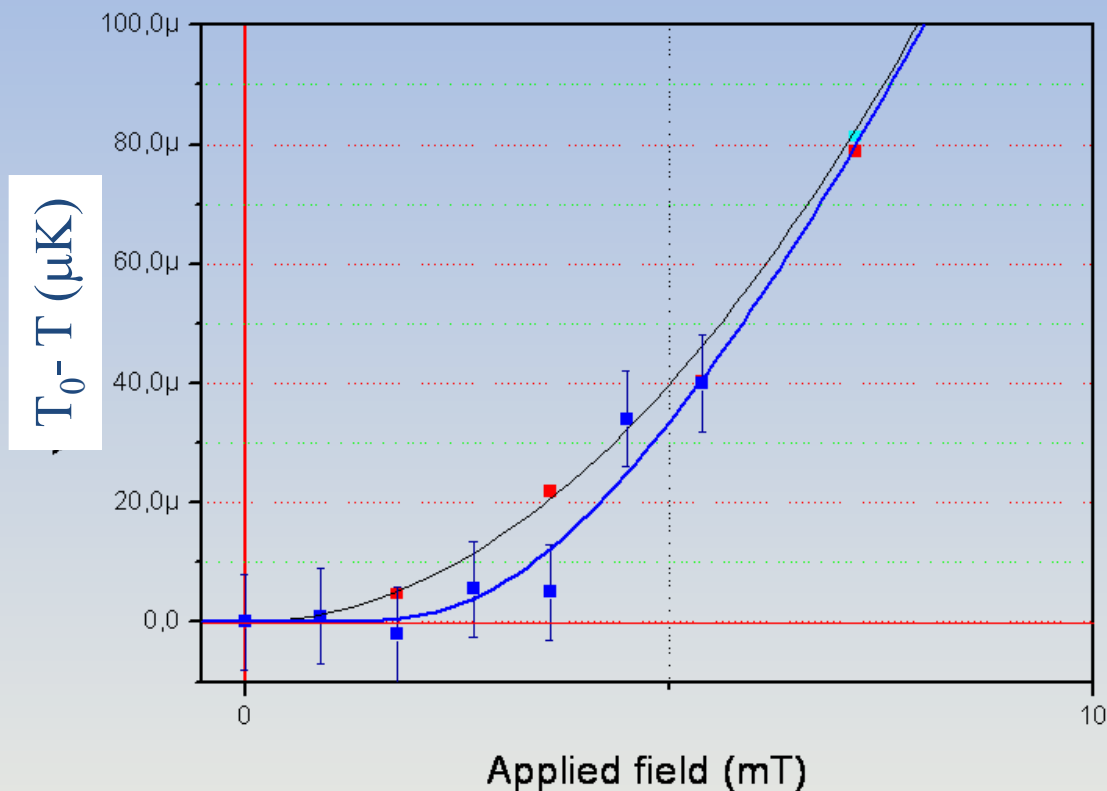
Is there a way to measure ΔF_c ?

Superconductivity is destroyed by a critical magnetic field .



The proposed way to measure ΔF_c consists in placing the cavity in a parallel magnetic field and measuring the critical field that destroys the superconductivity of the film.

Results and references on energy modulation – Aladin Experiment



The data are not in contrast with the theory and the region of energy of different behaviour is the expected one

2008

G. Bimonte, et Al. - Nucl.Phys. B726 (2005) 441-463

G. Bimonte et Al. - Phys.Rev.Lett. 94 (2005) 180402

G. Bimonte et Al. - J. Phys. A: Math. Theor. **41** 164023 (2008)

A. Allocca et Al. Jour. Of. Supercond. And Novel Magnetism. **25**, 2557-2565 (2012)

In the case of a thin film, of thickness $D \ll \lambda, \xi$ (with λ the penetration depth and ξ the correlation length).

$$\rho \approx \sqrt{24} \frac{\lambda}{D} \left(1 + \frac{9D^2}{\pi^6 \xi^2} \right)$$

Pressure red-shift

A simple summation of the lower force and upper force on the plates would bring to a somewhat unsuspected result:

$$F_{\text{inf}} + F_{\text{sup}} = F_C (1 + \delta\phi) + \frac{|E_C|}{c^2} g - F_C = 4 \frac{|E_C|}{c^2} g$$

$$F_C = L^2 \frac{\pi \hbar c}{240 a^4}$$

$$E_C = -L^2 \frac{\pi \hbar c}{720 a^3}$$

The lower vacuum «photons» must exert a bigger force because the force will be red-shifted when reaching the same level of upper plate → in the experiment the sum must be done taking into account the red-shift because the cavity is rigid and hanged in a unique point - (for this effect our measurement is a null measurement)

$$\left\{ \begin{array}{l} F_{\text{sup}} = F_C \\ F_{\text{inf}} = -F_C (1 + \delta\phi) + \frac{|E_C|}{c^2} g \end{array} \right. \longleftrightarrow \vec{F}_{\text{tot}} = \frac{|E_C|}{c^2} g \hat{z}$$

E. Calloni et.al. Phys. Letters A, 297, 328-333, (2002)

•G. Bimonte, E. Calloni, G. Esposito, L. Rosa - Phys. Rev D 74, 085011 (2006) *Phys.Rev.D* 74 (2006) 085011, *Phys.Rev.D* 75 (2007) 049904 (erratum),

•*Phys.Rev.D* 75 (2007) 089901 (erratum), *Phys.Rev.D* 77 (2008) 109903 (erratum)

On interpretation of Tolman-Ehrenfest effect:

C. Rovelli, M. Smerlak *Class. Quant. Grav.* 28 (2011) 075007

Hal M. Haggard and Carlo Rovelli, arXiv:1302.0724

S. A. Fulling et al. *Phys. Rev. D* 76:025004 (2007)

K.A. Milton et al. *J. Phys. A* 41:164052 (2008)

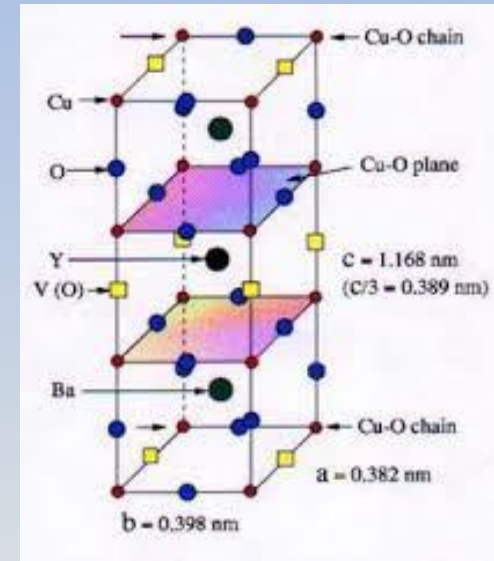
Progress on Casimir contribution to transition energy

Starting point: Kempf hypothesis: plasma sheet no dissipation
– zero temperature - order of magnitude estimation

Coupled cavities and low thickness limit regime with descriptions of each layer as a very thin superconductor which undergoes conductivity variation at the transition

Coupled cavities and plasma sheet description of each layer with dissipation and actual temperature

$$\Delta\eta_E \approx 10^{-4}$$



R. Bimonte, E. Calloni et al -Towards measuring variations of Casimir energy by a superconducting cavity
Phys.Rev.Lett. 94 (2005) 180402

L. Rosa et al - Casimir energy for two and three superconducting coupled cavities: Numerical calculations
Eur.Phys.J.Plus 132 (2017) no.11, 478

L. Rosa et al – Casimir energy for layered superconductors (in preparation -2021)

Theoretical issues

- ❑ Casimir effects in Layered type II superconductors
- ❑ Casimir effects in multi-layer coupled cavities
- ❑ Is plasma sheet approximation correct type II layered superconductors ?

Eur. Phys. J. Plus (2017) **132**: 478
DOI 10.1140/epjp/i2017-11750-y

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Regular Article

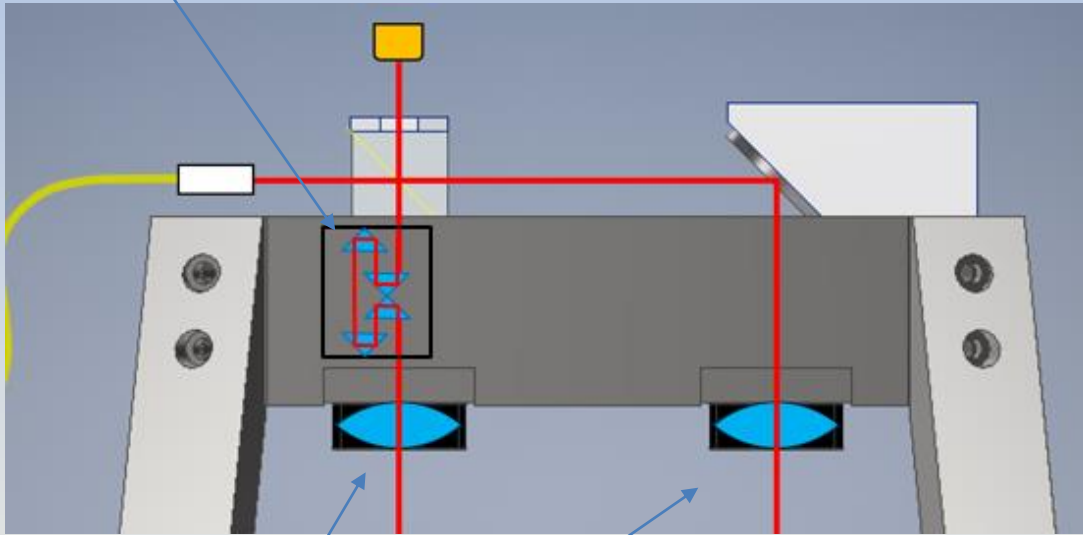
Casimir energy for two and three superconducting coupled cavities: Numerical calculations

L. Rosa^{1,2,a}, S. Avino^{2,3}, E. Calloni^{1,2}, S. Caprara^{4,5}, M. De Laurentis¹, R. De Rosa^{1,2}, Giampiero Esposito², M. Grilli^{4,5}, E. Majorana⁶, G.P. Pepe⁷, S. Petrarca^{4,6}, P. Puppo⁶, P. Rapagnani^{4,6}, F. Ricci^{4,6}, C. Rovelli^{8,9}, P. Ruggi¹⁰, N.L. Saini⁴, C. Stornaio², and F. Tafuri¹

From experimental point of view: even if the Casimir contribution would be of few percent the measurement can assess if vacuum fluctuations interacts with gravity or not

Optical Read-Out: prisms and lenses

prisms



Lenses with focal length equal
to arm distance

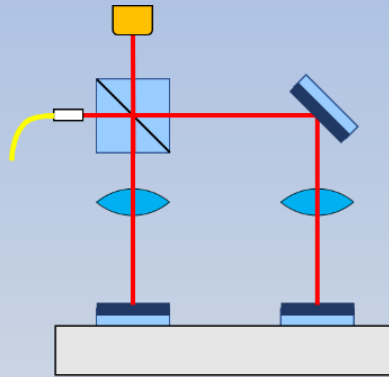
Prisms are used to equalize the length of the optical paths ($\Delta L \approx 2mm$). In this way **frequency noise is highly suppressed.**

Also amplitude noise is reduced by measuring a pick-off on the input beam and performing a signal normalization.

Optical Read-Out: prisms and lenses

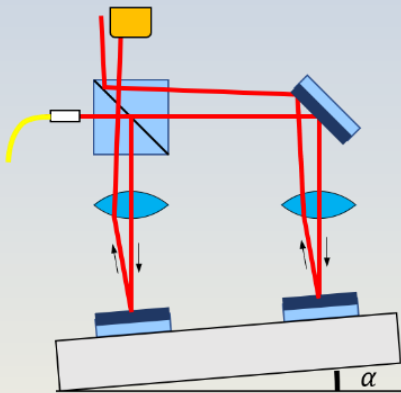
Lenses are used to maintain good alignment and contrast against static tilt.

a)

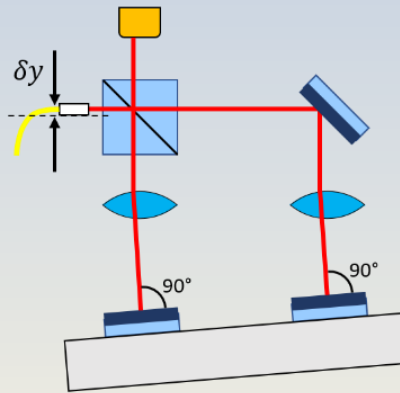


(a) The interferometer is aligned while the balance arm is horizontal

b)



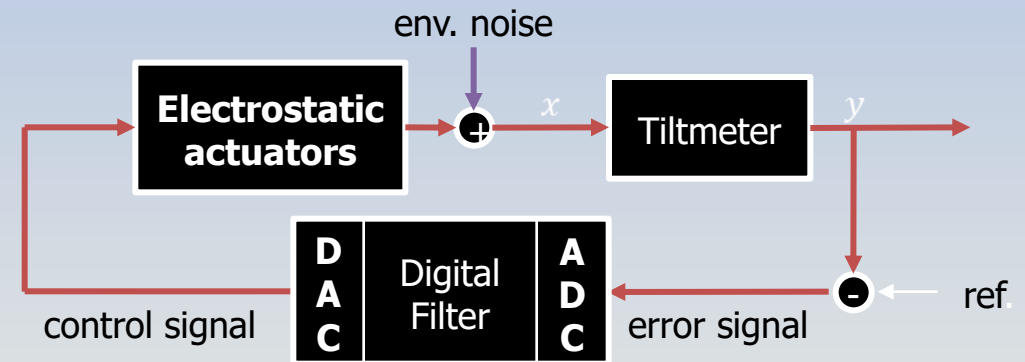
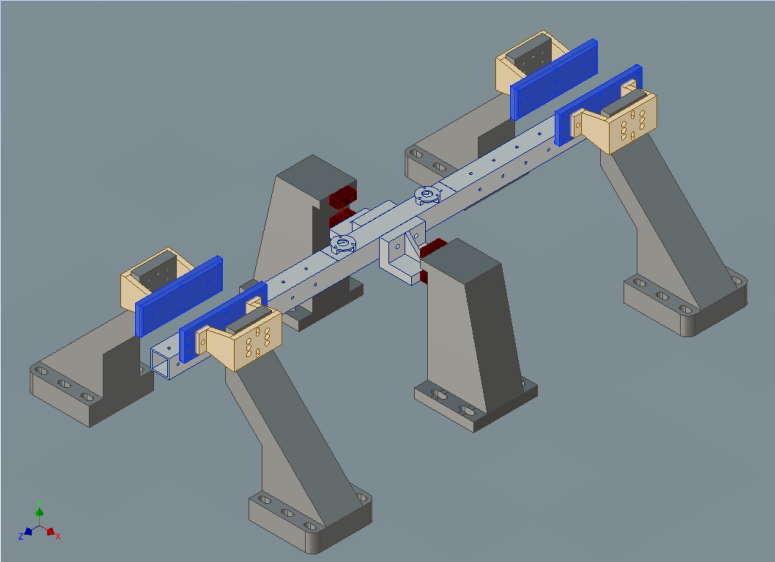
c)



(b) An arm tilt α would misalign the interferometer, but **contrast changes at second order thanks to the lenses**

(c) Alignment can be restored by moving vertically the input laser beam by $\delta y = L_f \alpha$ (L_f lens focal length), **without moving any optics on the suspended arm**

Acquisition & Control



NI-cRIO

