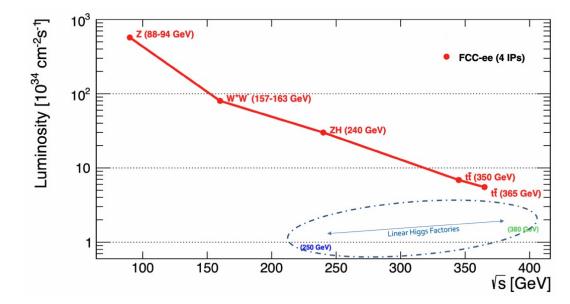
New Physics at FCC-ee

G. Polesello (INFN Pavia)

#### FCC-ee

- Energy range: Z-pole to ttbar production
- Very high luminosity yielding huge statistics of SM high mass particles, in particular
  - 6x10<sup>12</sup> Z bosons
  - ~2.5x10<sup>6</sup> Higgs bosons
- Very clean (e<sup>+</sup>e<sup>-</sup>) experimental environment as compared to hadron machines

 $\rightarrow$  see opportunities for BSM discovery in this environment



Working point	Z pole	WW thresh.	ZH	$t\overline{t}$	
$\sqrt{s}$ (GeV)	88, 91, 94	157, 163	240	340-350	365
Lumi/IP $(10^{34} \mathrm{cm}^{-2} \mathrm{s}^{-1})$	140	20	7.5	1.8	1.4
Lumi/year (ab <sup>-1</sup> )	68	9.6	3.6	0.83	0.67
Run time (year)	4	2	3	1	4
Integrated lumi. $(ab^{-1})$	205	19.2	10.8	0.42	2.70
			$2.2 \times 10^6  \mathrm{ZH}$	$2 \times 10$	$^{6}{ m t}{ar { m t}}$
Number of events	$6\times 10^{12}{\rm Z}$	$2.4\times 10^8{\rm WW}$	+	$+370 \mathrm{k} \mathrm{ZH}$	
			$65k~{\rm WW} \to {\rm H}$	$+92k \; WW \rightarrow H$	

#### BSM

After (HL-)LHC, generic new physics excluded up to scale of few TeV How to go beyond:

- Loopholes in LHC searches:
  - BSM particles with masses <few GeV</li>
  - Compressed BSM spectra
- High statistics and clean environment to explore higher mass scales:
  - Deviations from SM prediction of precision measurements
  - Direct detection of rare decays of H,Z,top

	Model	$\ell, \gamma$	Jets†	$E_{T}^{miss}$	∫£ dt[fb	-1]	Lim	nit	$\int \mathcal{L} dt = (3)$		Reference
	ADD $G_{KK} + g/q$ ADD non-resonant $\gamma\gamma$ ADD QBH ADD BH multiget RS1 $G_{KK} \rightarrow \gamma\gamma$ Bulk RS $G_{KK} \rightarrow tt$ Bulk RS $g_{KK} \rightarrow tt$	$\begin{array}{c} 0 \ e, \mu, \tau, \gamma \\ 2 \ \gamma \\ - \\ 2 \ \gamma \\ multi-channe \\ 1 \ e, \mu \\ 1 \ e, \mu \end{array}$	1 - 4j 2j $\ge 3j$ $\ge 1 b, \ge 1J/2$ $\ge 2 b, \ge 3j$		139 36.7 139 3.6 139 36.1 36.1 36.1 36.1	MD Ms Mth GKK mass GKK mass KKK mass KK mass		2.3 1.8 TeV	8.6 TeV 9.4 TeV 9.55 TeV 4.5 TeV TeV 3.8 TeV		2102.10874 1707.04147 1910.08447 1512.02586 2102.13405 1808.02380 1804.10823 1803.09678
	$\begin{array}{l} \mathrm{SSM} \ Z' \to \ell\ell \\ \mathrm{SSM} \ Z' \to \tau\tau \\ \mathrm{Leptophobic} \ Z' \to tt \\ \mathrm{Leptophobic} \ Z' \to tt \\ \mathrm{SSM} \ W' \to \tau\nu \\ \mathrm{SSM} \ W' \to \tau\nu \\ \mathrm{SSM} \ W' \to \tau\nu \\ \mathrm{HVT} \ W' \to WZ \ \mathrm{model} \ B \\ \mathrm{HVT} \ W' \to WZ \to \ell\nu \ \ell' \ \mathrm{model} \ B \\ \mathrm{HVT} \ W' \to WZ \to \ell\nu \ \ell' \ \mathrm{model} \ B \\ \mathrm{HVT} \ W' \to WZ \to \mu \ NR \\ \mathrm{HSM} \ W_R \to \mu \ NR \end{array}$	$\begin{array}{c} 2 \ e, \mu \\ 2 \ \tau \\ 0 \ e, \mu \\ 1 \ e, \mu \\ 1 \ \tau \\ 0.2 \ e, \mu \\ e \ l \ C \ 3 \ e, \mu \\ 1 \ e, \mu \\ 2 \ \mu \end{array}$	- 2 b ≥1 b, ≥2 J - 2 j/1 J 2 j/1 J 2 j/1 J 2 j/1 J 1 J	Yes	139 36.1 36.1 139 139 139 139 139 139 139 80	Z' mass Z' mass Z' mass W' mass W' mass W' mass W' mass W' mass Z' mass Z' mass	340 GeV	2.42 2.1 Te	5.1 TeV TeV 50 5.0 TeV 5.0 TeV 4.4 TeV 4.3 TeV 3.9 TeV 5.0 TeV	$\Gamma/m = 1.2\%$ $g_V = 3$ $g_V c_H = 1, g_F = 0$ $g_V = 3$ $m(N_R) = 0.5$ TeV, $g_L = g_R$	1903.06248 1709.07242 1805.09299 2005.05138 1906.05609 ATLAS-CONF-2021 2004.14636 2207.03925 2004.14636 1904.12879
5	Cl qqqq Cl llqq Cl eebs Cl µµbs Cl tttt	2 e,μ 2 e 2 μ ≥1 e,μ	2 j - 1 b ≥1 b, ≥1 j	- - - Yes	37.0 139 139 139 36.1	Λ Λ Λ Λ		1.8 TeV 2.0 Te 2.5		21.8 TeV $\eta_{LL}^-$ 35.8 TeV $\eta_{LL}^-$ $g_* = 1$ $g_* = 1$ $ C_{4t}  = 4\pi$	1703.09127 2006.12946 2105.13847 2105.13847 1811.02305
5	Axial-vector med. (Dirac DM) Pseudo-scalar med. (Dirac DM) Vector med. Z'-2HDM (Dirac D Pseudo-scalar med. 2HDM+a		2 j 1 - 4 j 2 b	- Yes Yes	139 139 139 139	m <sub>med</sub> m <sub>med</sub> m <sub>Z'</sub> m <sub>a</sub>	376 GeV	800 GeV	3.8 TeV 3.0 TeV	$\begin{array}{c} g_q = 0.25, \ g_{\chi} = 1, \ m(\chi) = 10 \ {\rm TeV} \\ g_q = 1, \ g_{\chi} = 1, \ m(\chi) = 1 \ {\rm GeV} \\ {\rm tan} \beta = 1, \ g_{\chi} = 0.8, \ m(\chi) = 100 \ {\rm GeV} \\ {\rm tan} \beta = 1, \ g_{\chi} = 1, \ m(\chi) = 10 \ {\rm GeV} \end{array}$	ATL-PHYS-PUB-20 2102.10874 2108.13391 ATLAS-CONF-202
3	Scalar LQ 1 <sup>st</sup> gen Scalar LQ 2 <sup>rd</sup> gen Scalar LQ 3 <sup>rd</sup> gen Scalar LQ 3 <sup>rd</sup> gen Scalar LQ 3 <sup>rd</sup> gen Scalar LQ 3 <sup>rd</sup> gen Vector LQ mix gen Vector LQ mix gen	$\begin{array}{c} 2 \ e \\ 2 \ \mu \\ 1 \ \tau \\ 0 \ e, \mu \\ \geq 2 \ e, \mu, \geq 1 \ \tau \\ 0 \ e, \mu, \geq 1 \ \tau \\ \text{multi-channe} \\ 2 \ e, \mu, \tau \end{array}$	0-2j,2b	Yes	139 139 139 139 139 139 139 139	LO mass LO mass LO <sup>2</sup> mass LO <sup>3</sup> mass LO <sup>3</sup> mass LO <sup>4</sup> mass LO <sup>4</sup> mass LO <sup>4</sup> mass		1.8 TeV 1.7 TeV 1.49 TeV 1.24 TeV 1.43 TeV 1.26 TeV 2.0 Te' 1.96 TeV		$\begin{array}{l} \beta = 1 \\ \beta = 1 \\ \mathcal{B}(LQ_2^o \rightarrow b\tau) = 1 \\ \mathcal{B}(LQ_2^o \rightarrow t\tau) = 1 \\ \mathcal{B}(LQ_2^o \rightarrow t\tau) = 1 \\ \mathcal{B}(LQ_2^o \rightarrow b\tau) = 1 \\ \mathcal{B}(LQ_2^o \rightarrow b\tau) = 1 \\ \mathcal{B}(LQ_1^o \rightarrow t\mu) = 1, \text{Y-M coupl.} \\ \mathcal{B}(LQ_1^o \rightarrow b\tau) = 1, \text{Y-M coupl.} \end{array}$	2006.05872 2006.05872 2303.01294 2004.14060 2101.11582 2101.12527 ATLAS-CONF-202 2303.01294
fermions	$ \begin{array}{l} VLQ \ TT \to Zt + X \\ VLQ \ BB \to Wt/Zb + X \\ VLQ \ T_{5/7} T_{5/3} T_{5/3} \to Wt + X \\ VLQ \ T \to Ht/Zt \\ VLQ \ Y \to Wb \\ VLQ \ B \to Hb \\ VLL \ t' \to Z\tau/H\tau \end{array} $	1 e, μ 1 e, μ	el	Yes Yes Yes	139 36.1 36.1 139 36.1 139 139	T mass B mass T <sub>5/3</sub> mass T mass Y mass B mass r' mass		1.46 TeV 1.34 TeV 1.64 TeV 1.8 TeV 1.85 TeV 2.0 Te 898 GeV		$\begin{array}{l} {\rm SU(2)\ doublet} \\ {\rm SU(2)\ doublet} \\ {\mathcal B}(T_{5/3} \to Wt) = 1, \ c(T_{5/3}Wt) = 1 \\ {\rm SU(2)\ singlet,} \ \kappa_T = 0.5 \\ {\mathcal B}(Y \to Wb) = 1, \ c_R(Wb) = 1 \\ {\rm SU(2)\ doublet,} \ \kappa_{B^0} = 0.3 \\ {\rm SU(2)\ doublet} \end{array}$	2210.15413 1808.02343 1807.11883 ATLAS-CONF-202 1812.07343 ATLAS-CONF-202 2303.05441
ferm	Excited quark $q^* \rightarrow qg$ Excited quark $q^* \rightarrow q\gamma$ Excited quark $b^* \rightarrow bg$ Excited lepton $\tau^*$	- 1γ - 2τ	2j 1j 1b,1j ≥2j		139 36.7 139 139	q* mass q* mass b* mass τ* mass			6.7 TeV 5.3 TeV 3.2 TeV 4.6 TeV	only $u^*$ and $d^*, \Lambda = m(q^*)$ only $u^*$ and $d^*, \Lambda = m(q^*)$ $\Lambda = 4.6 \text{ TeV}$	1910.08447 1709.10440 1910.08447 2303.09444
2000	Type III Seesaw LRSM Majorana v Higgs triplet $H^{\pm\pm} \rightarrow W^{\pm}W^{\pm}$ Higgs triplet $H^{\pm\pm} \rightarrow \ell \ell$ Multi-charged particles Magnetic monopoles	2,3,4 e, µ 2 µ 2,3,4 e, µ (SS 2,3,4 e, µ (SS 	≥2j 2j 5) various 5) - - - -	Yes - Yes - -	139 36.1 139 139 139 34.4	N <sup>0</sup> mass N <sub>R</sub> mass H <sup>±±</sup> mass H <sup>±±</sup> mass multi-charged particle monopole mass	350 GeV	910 GeV 1.08 TeV 1.59 TeV 2.37	3.2 TeV	$\begin{split} m(W_R) &= 4.1 \text{ TeV}, g_L = g_R \\ \text{DY production} \\ \text{DY production} \\ \text{DY production},  q  = 5e \\ \text{DY production},  g  = 1g_D, \text{ spin } 1/2 \end{split}$	2202.02039 1809.11105 2101.11961 2211.07505 ATLAS-CONF-202 1905.10130

+Small-radius (large-radius) iets are denoted by the letter i (J.

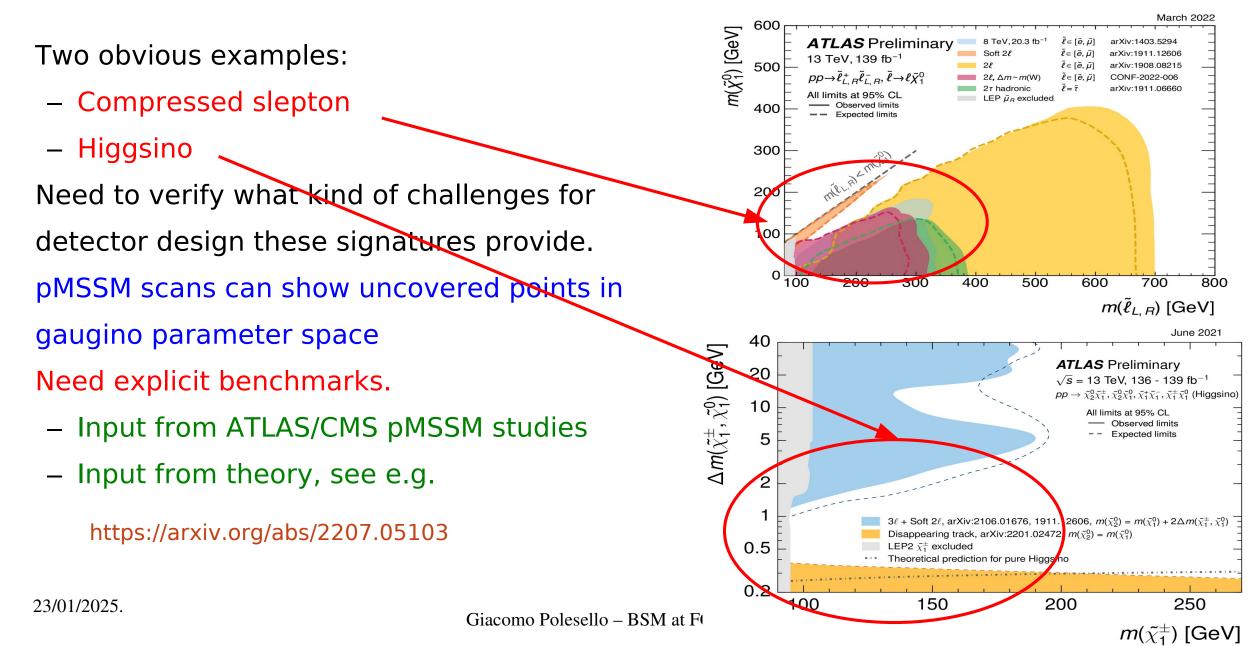
#### Talk by A. Valenti today



Main thrust of this talk, based on results shown at recent: ECFA workshop FCC workshop

#### 23/01/2025.

### LHC loopholes: SUSY



#### The role of EFTs

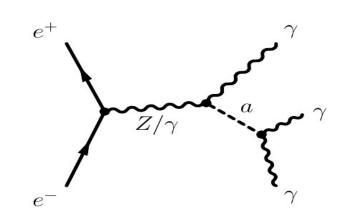
Before LHC clear theory font-runner: SUSY

For FCC no strong theoretical guidance: rely on EFT approach:

- Postulate a new BSM particle a
- Add to SM Lagrangian terms for coupling of a to relevant SM particles suppressed by scale of new physics

Example: axion-like particle coupling to vector bosons

$$\mathcal{L}_{\text{eff}} \ni e^2 C_{\gamma\gamma} \frac{a}{\Lambda} F_{\mu\nu} \tilde{F}^{\mu\nu} + \frac{2e^2}{s_w c_w} C_{\gamma Z} \frac{a}{\Lambda} F_{\mu\nu} \tilde{Z}^{\mu\nu}$$



Achievable scale with 10<sup>12</sup> Z:

BR(Z→aγ)=(C<sub>γZ</sub>/Λ)<sup>2</sup>x8.6e<sup>-4</sup> (Λ in TeV) → For BR(Z→aγ)=1e<sup>-12</sup>: C<sub>γZ</sub>/Λ ~ 3e<sup>-5</sup> TeV<sup>-1</sup>

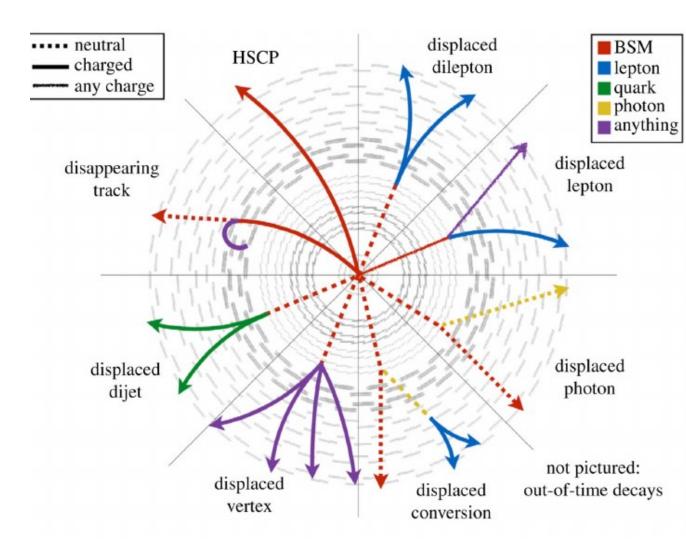
With  $10^{12}$  Z, a new physics scale of  $\sim 10^4$  TeV can be explored by looking for rare decays

#### Lifetime of new particles

Width of particle a of mass  $m_a$  decaying only through vertex (C/ $\Lambda$ ):

#### $\Gamma_a \sim m_a^3 (C/\Lambda)^2$

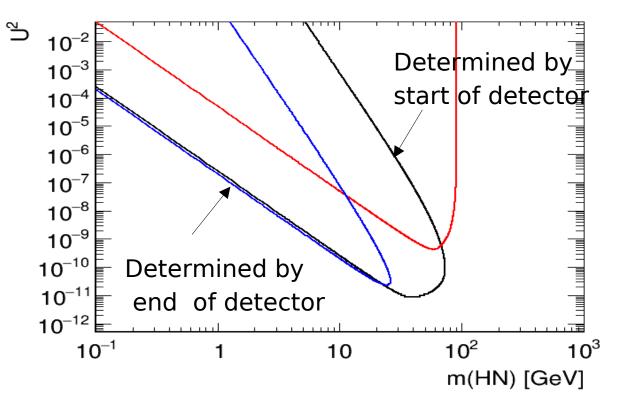
- For low masses and low couplings small BSM particle width  $\rightarrow$  long lifetimes, LLPs
- Wealth of signatures with little/no SM background
- LHC detectors designed without thinking of LLPs (although they are doing pretty well on it!)
- Establish requirements for FCC detectors enabling them to fully exploit physics opportunity of LLPs



23/01/2025.

### Prompt vs LLP

Generically reach is defined in m(new physics)-coupling plane True e.g for ALP, HNL



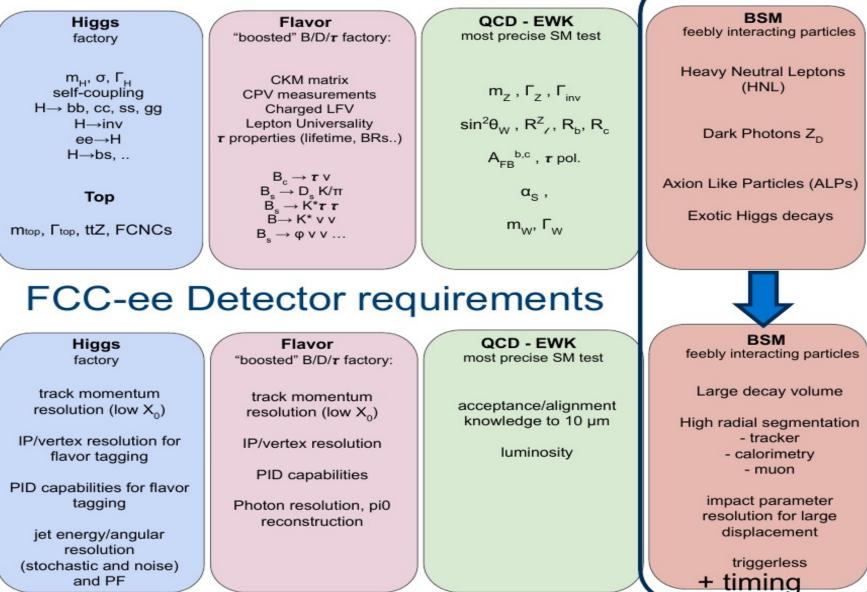
Complementary reach of three different signatures:

- Prompt
- Decay in inner detector
- Decay in calo/muon detector

Study of coverage for a given model should address all three signatures. Very different experimental requirements

#### FCC-ee: physics vs detector requirements

#### **FCC-ee Physics landscape**



FCC has very large menu of physics topics

Each of these poses a specific experimental challenge and pushes detector optimisation

Unique challenges from BSM: long-lived particles

#### Benchmark studies in FCC-ee PED BSM group

- Z-pole (extendable to all Fcc-ee runs)
  - Axion-like particle searches
  - Heavy Neutral Lepton (HNL) searches
- Higgs sector
  - Higgs decay to long-lived scalars
  - Searches for additional higgses

Reach studies based on parametrised simulation: define requirements on detectors for full exploitation of FCC-ee direct BSM potential.

Next step will be moving to detailed GEANT4 simulation 23/01/2025.

### Recent workshops

Today personal choice out of large material, more info in three recent workshops



2ND "FEE ITHLY & FRANCE WORKSHOP"

VENICE, PALAZZO FRANCHETTI - NOVEMBER 4 - 6, 2024

#### ECFA workshop link

See talk of R. Franceschini there for topics I will not touch Link

#### Venice workshop link



CERN workshop link

#### Workflow of experimental analyses



•Background files produced centrally based on FCC software.

- •Signal files produced either centrally or by analysis group.
- •DELPHES output stored in EDM4HEP format
- •Use FCCsoftware to produce ntuples for analysis based on FCCanalysis package

•Two large production campaigns, spring2021 and winter2023

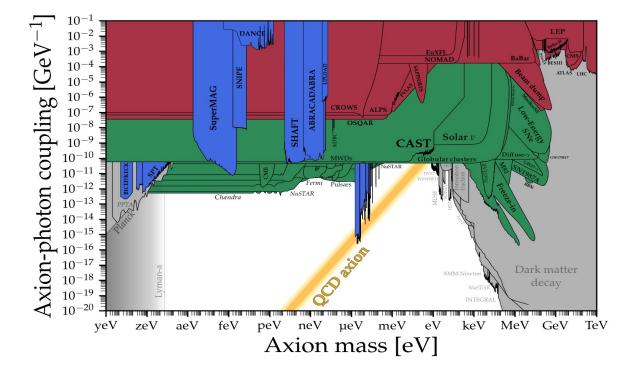
•Main limitation: statistics at peak 23/01/2025.

## Z-pole studies

#### ALPs

Axion Like Particles (ALP): hypothetical pseudoscalar with similar interactions as the QCD axion, appearing naturally in many extensions of the SM

Couples to Z/photon, can be abundantly produced at FCC-ee

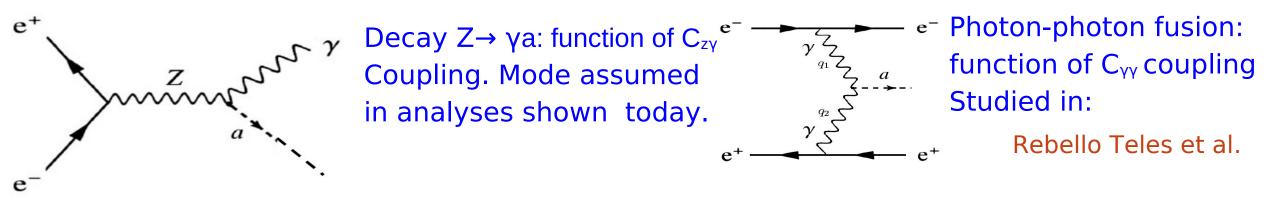


High statistics of FCC-ee Z-pole run allows exploration of much lower couplings to photons than tested to date in mass range 0.1-90 GeV In BSM group ongoing studies for different ALP decay modes:  $a \rightarrow \gamma \gamma$ ,  $a \rightarrow gluon gluon$ ,  $a \rightarrow \mu \mu$ 

#### Typical vector boson part of ALP Lagrangian Bauer et al:arXv:1808.10323

$$\mathcal{L}_{eff} \ni e^2 C_{\gamma\gamma} \frac{a}{\Lambda} F_{\mu\nu} \tilde{F}^{\mu\nu} + \frac{2e^2}{s_w c_w} C_{\gamma Z} \frac{a}{\Lambda} F_{\mu\nu} \tilde{Z}^{\mu\nu} + \frac{e^2}{s_w^2 c_w^2} C_{ZZ} \frac{a}{\Lambda} Z_{\mu\nu} \tilde{Z}^{\mu\nu}$$
with  $C_{\gamma\gamma} = C_{WW} + C_{BB}, \qquad C_{\gamma Z} = c_w^2 C_{WW} - s_w^2 C_{BB}, \qquad C_{ZZ} = c_w^4 C_{WW} + s_w^4 C_{BB}$ 
Benchmark:Assume a couples to hypercharge and not to SU2 (C<sub>ww</sub>=0)  $C_{\gamma Z} = -s_w^2 C_{\gamma\gamma}$ 

2-d parameter space:  $(m_a, C_{\gamma\gamma})$ Production in FCC-ee Z-pole run



#### $a \rightarrow \gamma \gamma$

Regions of interest:

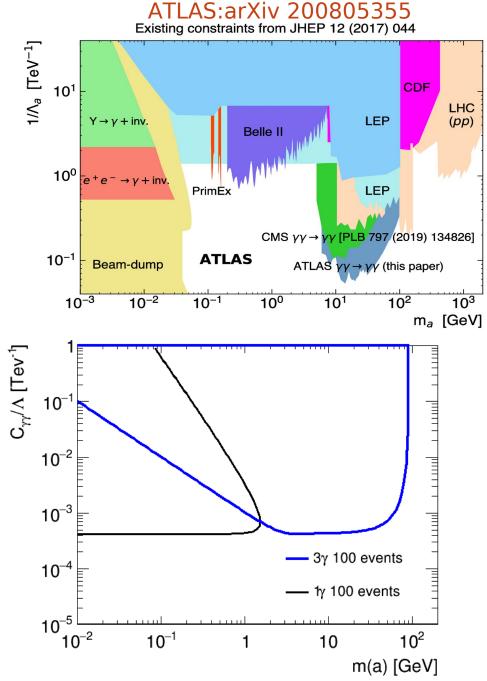
•0.1< ma < 10 GeV:

loose limits from previous e+e- searches, out of reach of beam dump

•10 <  $m_a$  < 90 GeV:

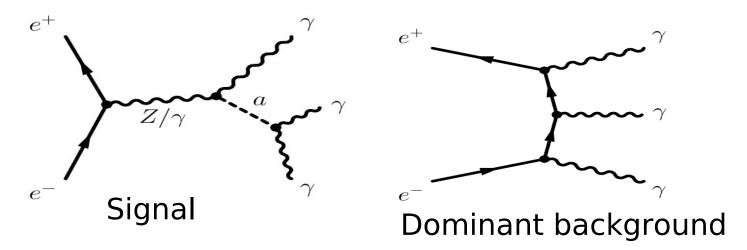
dominated by LHC photon-photon fusion potential for FCC-ee Z pole run

Depending on a lifetime consider two cases
•Three photons are observed in detector
•The ALP decays outside the detector: only a monocromatic photon in the event
Two different regions in parameter space
covered.



23/01/2025.

#### 3γ ALP analysis

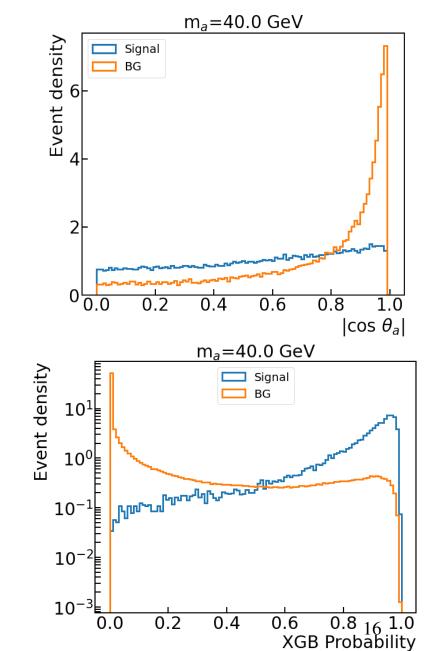


•3 photons within  $|\eta|{<}2.6$  and energy>1 GeV

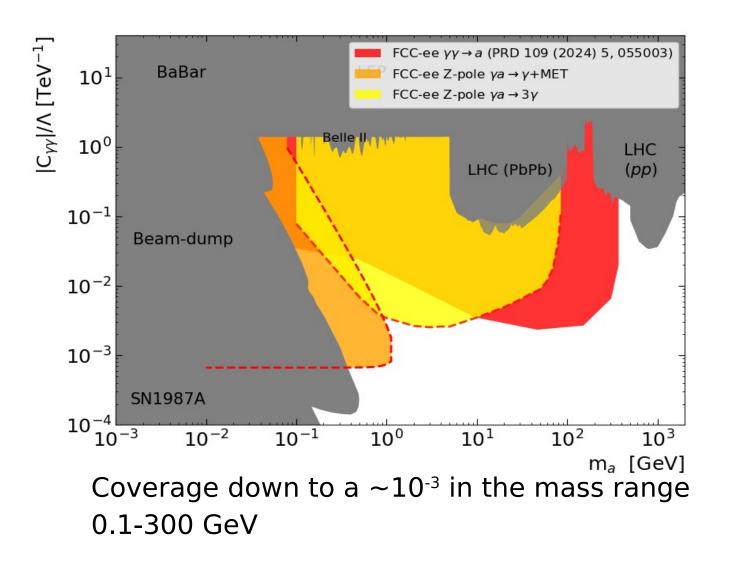
•Scan test masses *M* between 0.1 and 85 GeV

- •Assign 2 photons to ALP decay based on kinematic compatibility ( $\gamma_1, \gamma_2$ ), third photon to Z decay ( $\gamma_3$ )
- •Build BDT probability based on 3 angular variables+  $m(\gamma_1\gamma_2)$ ,  $E_{\gamma_3}$ , and  $E_{\gamma_2}/E_{\gamma_1}$

#### G.P. Talk at CERN



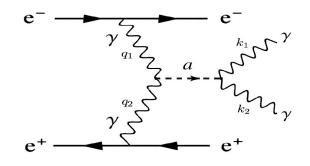
# Results



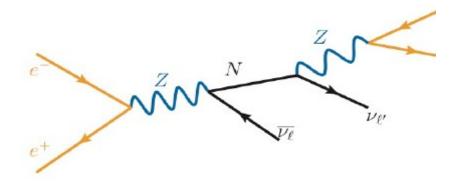
Grey areas : existing exclusions

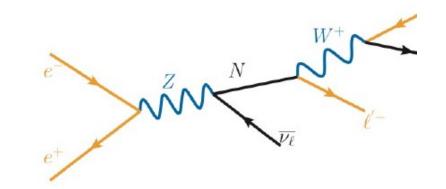
Yellow and orange: areas with >2 $\sigma$  significance respectively For the 3-photon and 1-photon analysis

Red area is analysis of Rebello Teles et al. addressing ALP production in photon-photon fusion



#### **HNL Models**





Production in Z decay via mixing with light neutrinos

$$BR(Z \to \nu N) = \frac{2}{3} |U_N|^2 BR(Z \to \text{invisible}) \left(1 + \frac{m_N^2}{2m_Z^2}\right) \left(1 - \frac{m_N^2}{m_Z^2}\right) \left|U_N|^2 \equiv \sum_{\ell=e,\mu,\tau} |U_{\ell N}|^2$$

Decay: three-body decay into 3 fermions via virtual W/Z

Decay length:

$$L_{N_i} = \simeq \frac{1.6}{U_i^2} \left(\frac{M_i}{GeV}\right)^{-6} \left(1 - (M_i/M_Z)^2\right) \ cm$$

23/01/2025.

For each HNL, phenomenology determined by 4 parameters: mass, mixing with three lepton flavours

#### Benchmark models

#### Two models:

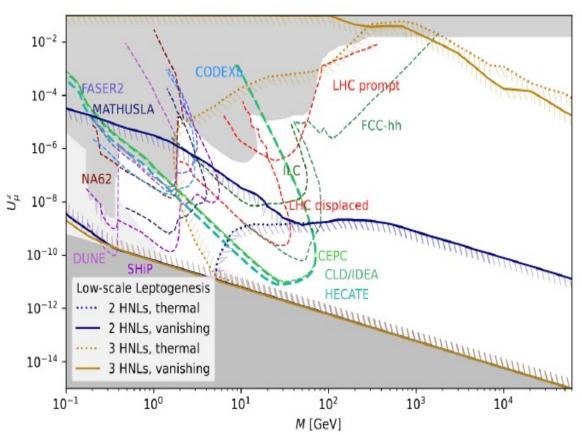
#### •Minimal realistic seesaw scenario:

Pseudo Dirac pair of semidegenerate Majorana HNLs
Coupling to all leptons
Parameter choices compatible with leptogenesis and oscillation data

#### •Single low-mass HNL

mixing with one lepton flavour I: only 2 parameters  $m_N$  and  $U_I$ , useful for comparing experiments or accelerators

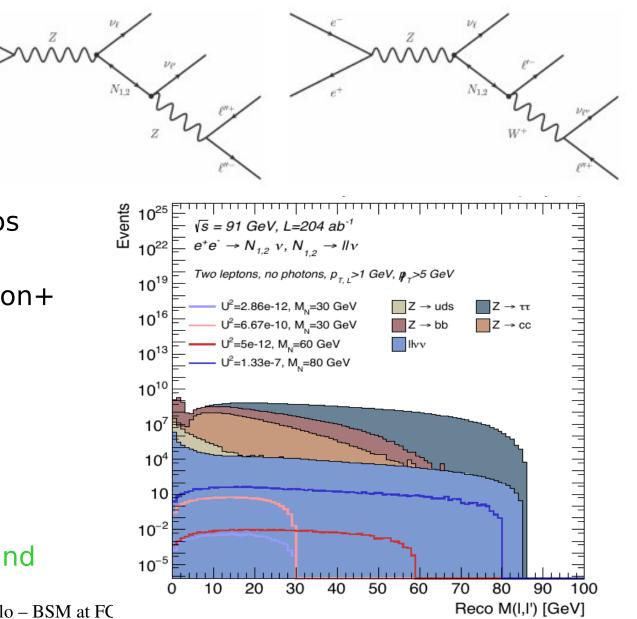
#### ArXiv:2203.05502



For single HNL several analysis with coupling to both e and  $\mu$  and considering fully leptonic and semileptonic N decay: show only semileptonic  $\mu$  case

### **Two HNLs**

#### S. Giappichini et al. arXiv:2410-03615



Consider only decay of N into 3 leptons :

 $e^+e^- \rightarrow N_{1,2}\nu$ ,  $N_{1,2} \rightarrow II\nu$ 

For  $\tau$  consider only leptonic decays Final state with two leptons and two neutrinos

Backgrounds: Z decays from official production+ 4-fermion irreducible:  $e^+e^- \rightarrow |^+|^-\nu\nu$ 

#### **Preselection:**

2 reco leptons pT>1 GeV  $p_T^{miss}$ >5 GeV veto photons and additional tracks

+ kinematic selections to suppress background

23/01/2025.

Giacomo Polesello – BSM at FC

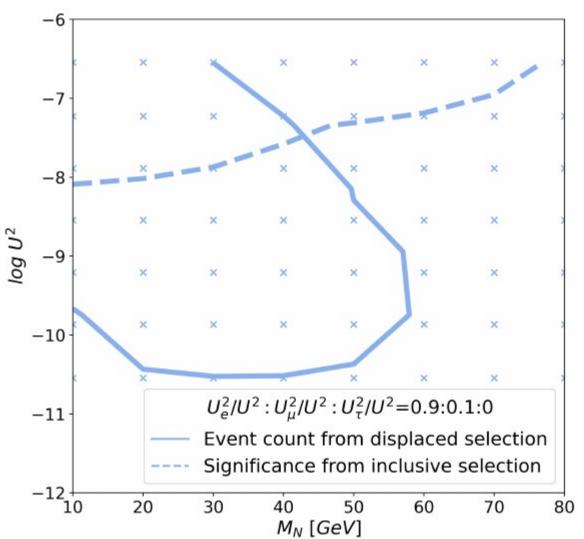
# Results

Two different kinematic selections: inclusive and displaced, separated by requirement on impact parameter of leptons Displaced:

|d<sub>0</sub><sup>|</sup>|>0.64 mm

For displaced selection achieve zero background, sensitive up to  $\sim$ 60 GeV

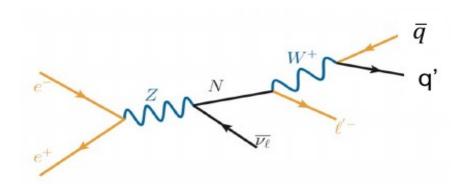
 $N_{1,2} \rightarrow \ell \ell v \text{ at FCC} - ee, \sqrt{s} = 91 \text{ GeV}, \mathcal{L}_{int} = 204 \text{ ab}^{-1}$ 



Giacomo Polesello – BSM at FCC-ee

#### GP, Nicolò Valle

# Single HNL→ µjj

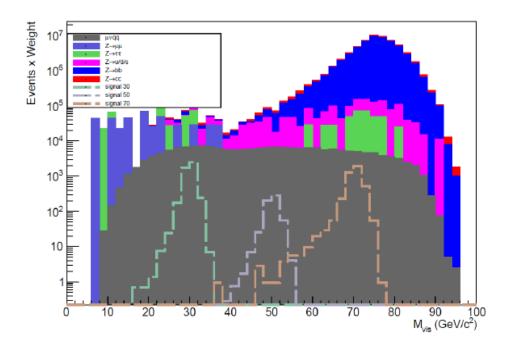


Most favourable decay: 50% cross-section
Full reconstruction of HNL possible.
Momentum of neutrino recoiling against HNL fixed by recoil formula:

$$p_{\nu}(M_{N_1}) = \frac{M_Z^2 - M_{N_1}^2}{2 M_Z}$$

Strong kinematic constraints allow efficient Background suppression

23/01/2025.



Prompt analysis at Z peak: Reducible backgrounds: Z decays, dominated by Zbb Irreducible background: 4-body μνqq

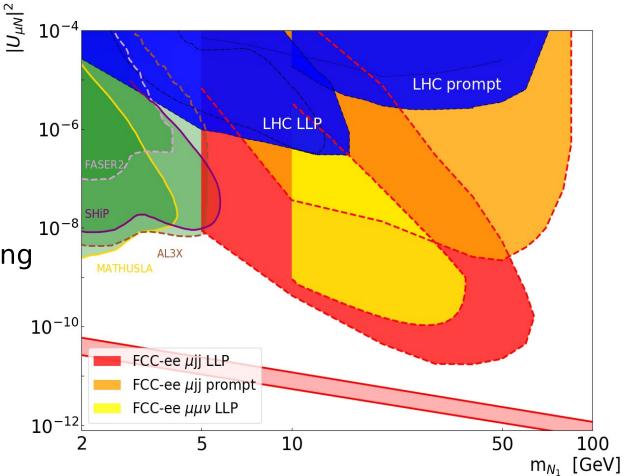
# Results



Reconstruct good vertex from all tracks in event, require most tracks connected to vertex

- Two different kinematic selections depending on radial position of vertex  $r_{vx:}$
- •Prompt (orange)  $r_{vx} < 0.5$ mm •LLP (red):  $r_{vx} > 0.5$ mm  $\rightarrow$  zero background

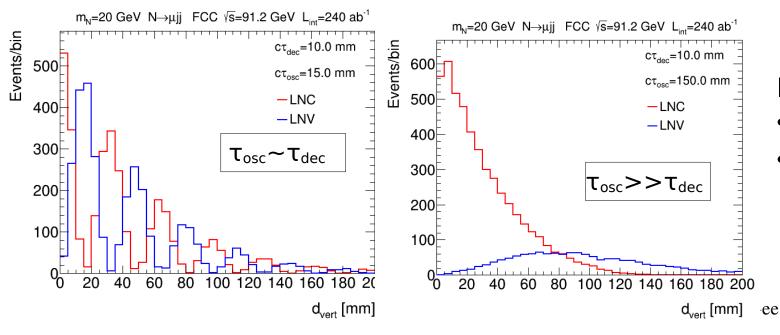
Also shown LLP μμνν final state (yellow)

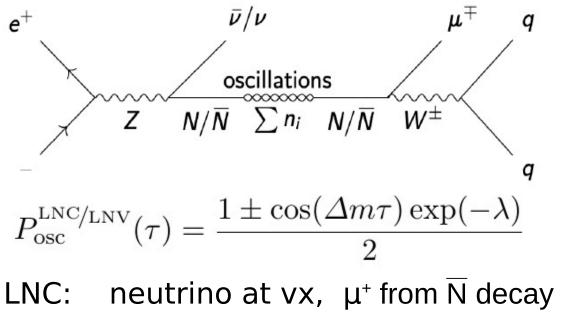


## HNL oscillations

Pseudo-Dirac pair, decay  $N \rightarrow \mu j j$ Implemented in pSPSS(phenomenologically symmetry protected seesaw) model S. Antusch et al. JHEP 10 (2023) 129

Two HNL oscillate into each other as they propagate before decaying





LNV: antineutrino at vx,  $\mu^2$  from N decay

Interplay of two times:
•τ<sub>osc</sub>: oscillation period ~Δm
•τ<sub>dec</sub>: determined by mass and mixing angle
GP, Nicolò Valle

ECFA Talk

24

# Oscillation variables

#### **Production:**

Asymmetry in HNL production angle from EWK Z polarization

$$A_{\ell}^{\mathrm{FB}}(\tau) \coloneqq \frac{A_{\ell^{-}}^{\mathrm{FB}}(\tau) - A_{\ell^{+}}^{\mathrm{FB}}(\tau)}{2} = A_{N}^{\mathrm{FB}} \Delta P_{\mathrm{osc}}(\tau)^{-1}$$

W

Ζ

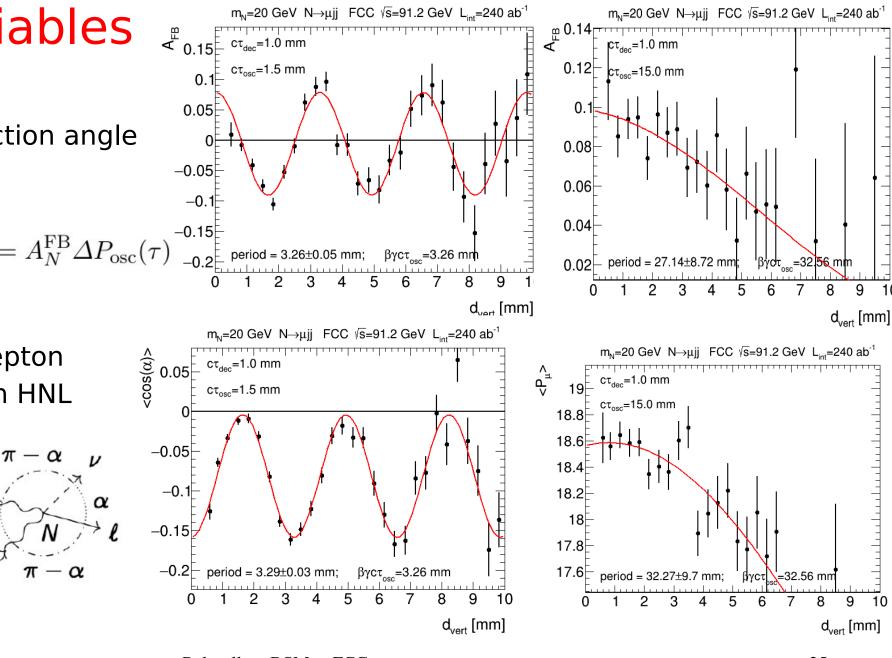
 $\alpha$ 

#### Decay:

Opening angle neutrino-lepton in HNL rest frame ( $\alpha$ ) from HNL polarisation

 $e^{-}$ 

JHEP 11 (2024) 102



# Dirac-Majorana

Investigate Dirac or Majorana nature of HNL

HNL production:

Dirac HNL:

HNL production angle  $\vartheta_{vis}$  different <sup>50</sup> for positive and negative  $\mu$ Two-state Majorana:

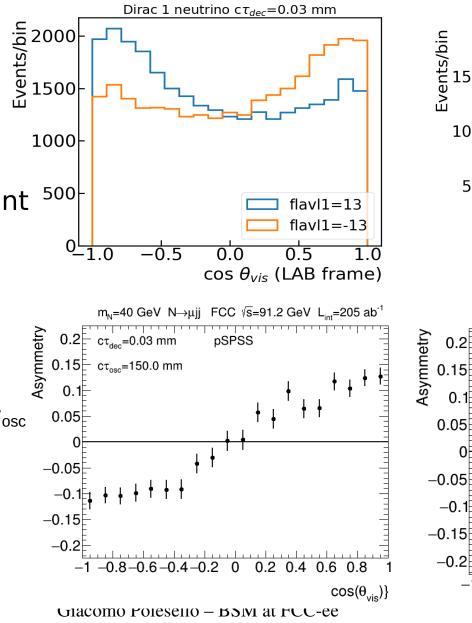
No difference

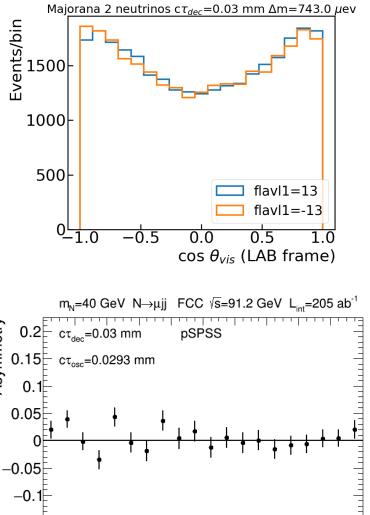
```
In pSPPS model:
```

Majorana behaviour for  $c\tau_{dec}=c\tau_{osc}$ Dirac behaviour for  $c\tau_{osc} >> c\tau_{dec}$ 

Plot  $\mu^+\mu^-$  asymmetry in bins of  $\cos\theta_{vis}$  $\frac{\#(\mu^+)-\#(\mu^-)}{\#(\mu^+)+\#(\mu^-)}$ 

# Seminal work:Blondel et al arXiv:2105.06576





-0.8 - 0.6 - 0.4 - 0.2

0.8

0.2 0.4 0.6

0

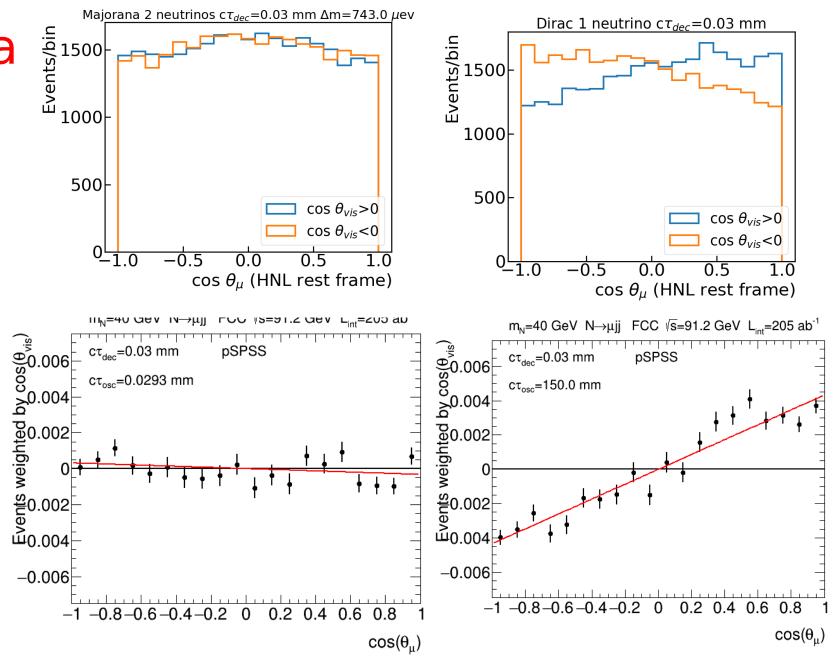
# Dirac-Majorana

#### HNL decay:

 $\vartheta_{\mu}$ , angle of  $\mu$  with respect to HNL direction of flight in HNL rest frame sensitive to Dirac-Majorana nature

Different distributions for forward and backward produced HNL for Dirac case

Plot number of events in bins of  $\cos \theta_{\mu}$  weighted by  $\cos \theta_{vis}$ 



# Higgs run

# Higgs decay into long-lived scalars

- Hidden sector model with scalar portal
  - New dark scalar mixes with SM Higgs via angle sin $\vartheta$
- Exotic decay of SM Higgs, , into new scalar then decays into SM states
- Small mixing angles yields long-lived scalars:
  - LLPs  $\rightarrow$  Displaced Vertex (DV) search
- Targets Zh stage; 240 GeV & 10.8 ab-1
  Signature generated with HAHM model:

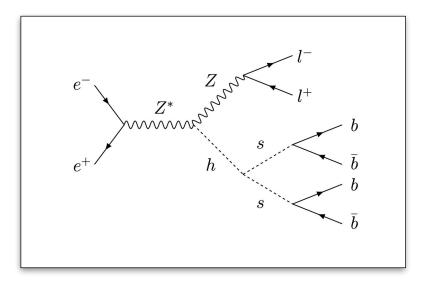
 $e^+e^- \rightarrow Z \rightarrow Zh, Z \rightarrow l^+l^-, h \rightarrow ss, s \rightarrow b^+b^-$ 

• Main backgrounds:WW, ZZ, WZ

G. Ripellino, M. Vande Voorde, A. Gallén, R. Gonzalez Suarez arXiv:2412.10141

23/01/2025.

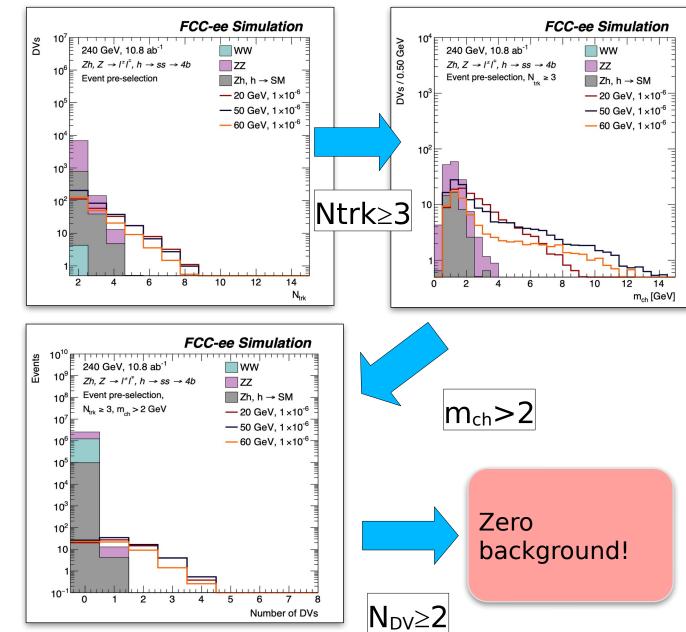
Giacomo Polesello – BSM at FCC-ee



$c\tau \ [{\rm mm}]$	$BR(h \rightarrow ss)$
3.4	$8.1  imes 10^{-4}$
38	$8.1  imes 10^{-4}$
340	$8.1  imes 10^{-4}$
34000	$8.1  imes 10^{-4}$
1.4	$10.2  imes 10^{-4}$
140	$10.2  imes 10^{-4}$
14000	$10.2  imes 10^{-4}$
12	$10.9  imes 10^{-4}$
110	$10.9  imes 10^{-4}$
1200	$10.9  imes 10^{-4}$
0.9	$7.4  imes 10^{-4}$
88	$7.4  imes 10^{-4}$
8800	$7.4  imes 10^{-4}$
	$\begin{array}{c c} 3.4\\ 3.4\\ 38\\ 340\\ 34000\\ \hline 1.4\\ 140\\ 14000\\ \hline 12\\ 110\\ 1200\\ \hline 0.9\\ 88\end{array}$

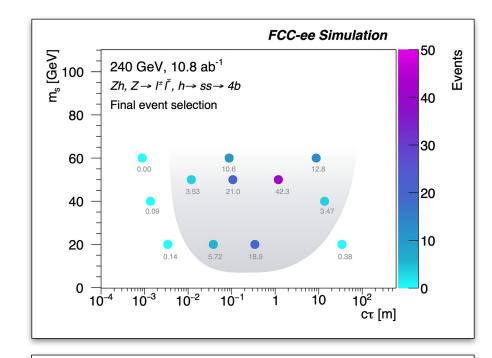
# Analysis

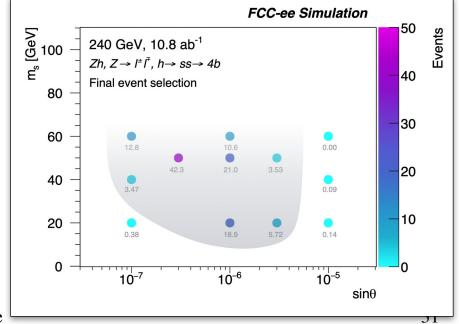
- Secondary Vertex finder of <u>LCFIPlus algorithm</u> used
  - Custom track selection:  $p_T>1$  GeV &  $|d_0|>2$  mm
- Selections:
  - Event selection:
    - 2 iso. leptons ( $\mu$  or e), oppositesign, same flavour
    - 70 GeV <  $m_{\parallel}$  < 110 GeV
    - At least 2 DVs passing the full DV selection
  - DV selection:
    - $N_{trk} \ge 3$
    - m<sub>ch</sub>>2 GeV



### Results

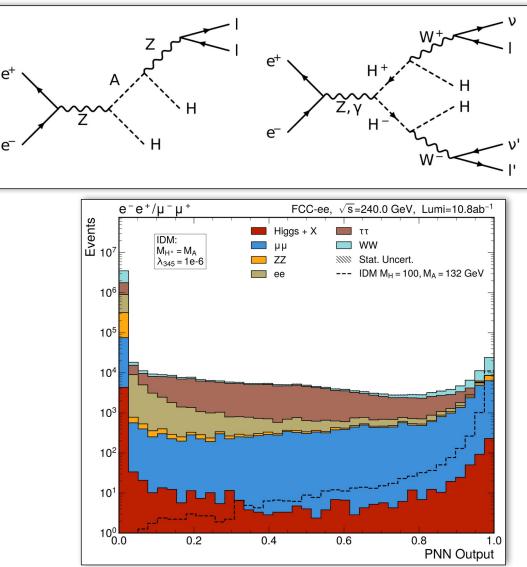
- SM background free search
- Based on generated grid draw contour of signal points with 3 events in two planes
  - ms vs cτ
  - ms vs sinθ
- Successfully performed sensitivity analysis
  - BR(h $\rightarrow$ ss) probed to 1e-4 for ct~1m





# Additional Higgs bosons

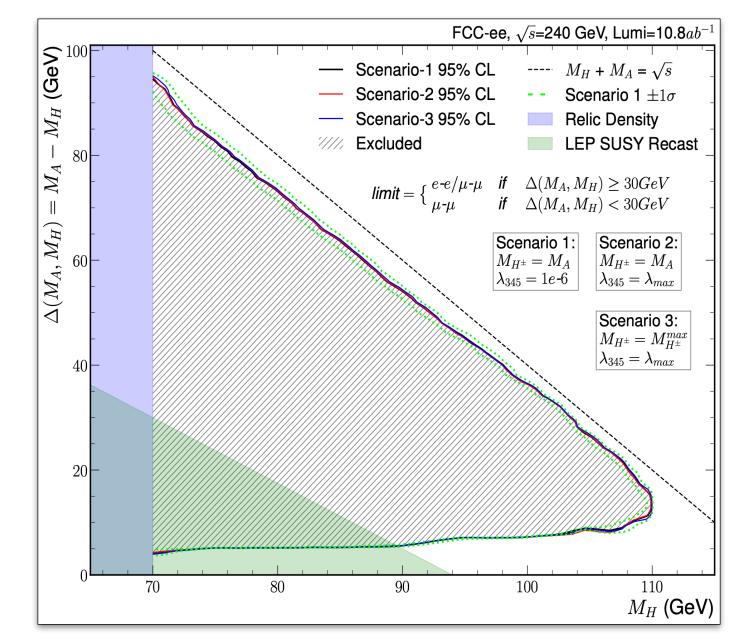
- Inert Two-Higgs-Doublet model (IDM)
  - Five Higgs bosons, h is SM Higgs
  - BSM Higgs do not couple to fermions and are pair-produced
  - Five new free parameters:  $m_{H}, m_{A}, m_{H+}, quartic couplings \lambda_2, \lambda_{345}$
- One new Dark matter candidate: H (invisible)
- Final state: I<sup>+</sup>I<sup>-</sup>HH ( signature 2 leptons + missing energy)
- Backgrounds: inclusive e<sup>+</sup>e<sup>-</sup>→l<sup>+</sup>l<sup>-,</sup> WW,ZZ,ZH, top when open
- Require 2e (2µ) with p>5 GeV + some  $E^{miss}$  and nothing else in event
- Insert kinematic variables for event in parametric neural network (PNN)
   23/01/2025.



#### E. Curtis, A.-M. Magnan & Tania Robens <u>ECFA 2024 presentation</u>, <u>CDS note</u>

### Results

- Results plotted in plane defined by DM mass M<sub>H</sub>, and mass difference M<sub>A</sub>-M<sub>H</sub>
- Three different scenarios with different values of M<sub>H+</sub> and quartic couplings
- Little sensitivity to additional parameters
- Sensitivity dominated by ZA channel
- Most of kinematically accessible parameter space covered



### Conclusions

The FCC-ee has a large potential for new physics through direct searches

- High statistics at Z pole gives access to very rare decays: new physics suppressed by high scales
- Exotics Higgs decays open a portal to dark physics

Exploration of low masses and couplings addresses long-lived signatures which have low SM backgrounds

Vigorous effort in FCC PED group to assess reach of FCC-ee data for most relevant benchmark models, and to extract detector requirements as an input to detector design

# Backup

## List of topics

Cross-check a recently circulated list of ECFA-WRG1-SRCH (Rebeca G-S. is convener both for us and for that group)

- Heavy Neutral Leptons \*
- Exotic Higgs boson decays \*
- Light SUSY scenarios and scenarios with light scalars
- Axion-like particles (ALP) \*
- Z', dark photons and other light mediator scenarios

For items with \* organised activity in our community, addressed in talks by

G. Ripellino and S. Kulkarni

23For SUSY I'll discuss some benchmark possibilities

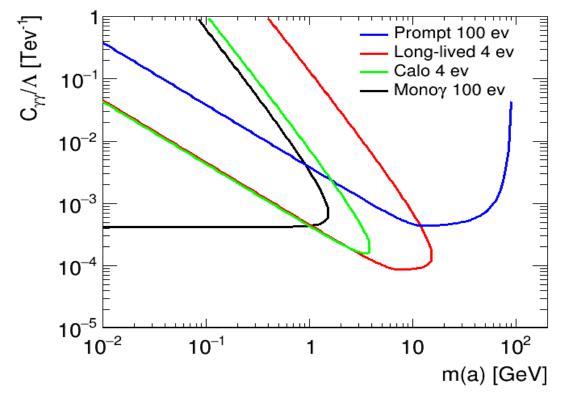
### Long Lived searches: History

Rend. Fis. Acc. Lincei s. 9, v. 12:5-18 (2001)

Fisica. — SUSY Long-Lived Massive Particles: Detection and Physics at the LHC. Nota di Sandro Ambrosanio, Barbara Mele, Aleandro Nisati, Silvano Petrarca, Giacomo Polesello, Adele Rimoldi e Giorgio Salvini, presentata (\*) dal Socio G. Salvini.

	Model	Signature	∫£ dt [fb	<sup>-1</sup> ]	Life	etime limit						Reference
	RPV $\tilde{t} \rightarrow \mu q$	displaced vtx + muon	136	t lifetime	- · · · ·				0.003-6.0 m		$m(\tilde{t})=1.4$ TeV	2003.11956
	$RPV\tilde{\chi}^0_1 \to eev/e\mu v/\mu\mu v$	displaced lepton pair	32.8	${\widetilde \chi}_1^0$ lifetime				0.003-1.0 m			$m(\tilde{q})=$ 1.6 TeV, $m(\tilde{\chi}_1^0)=$ 1.3 TeV	1907.10037
	$\operatorname{RPV} \widetilde{\chi}^0_1 \to q q q$	displaced vtx + jets	139	${\widetilde \chi}_1^0$ lifetime			-		0.00135-9.0 m		$m(\widetilde{\chi}_1^0) = 1.0 \text{ TeV}$	2301.13866
	$\operatorname{GGM} \tilde{\chi}^0_1 \to Z  \tilde{G}$	displaced dimuon	32.9	${\widetilde \chi}_1^0$ lifetime			-		0.029-18.	0 m	$m( ilde{g}){=}$ 1.1 TeV, $m( ilde{\chi}_1^0){=}$ 1.0 TeV	1808.03057
	GMSB	non-pointing or delayed y	139	${\widetilde \chi}_1^0$ lifetime				0.3	24-2.4 m		$m(\tilde{\chi}^0_1, \tilde{G})$ = 60, 20 GeV, $\mathcal{B}_H$ = 2%	2209.01029
	GMSB $\tilde{\ell} \to \ell \tilde{G}$	displaced lepton	139	$\tilde{\ell}$ lifetime			-	6-750 mm			$m(\tilde{\ell}) = 600 \text{ GeV}$	2011.07812
1000	GMSB $\tilde{\tau} \rightarrow \tau \tilde{G}$	displaced lepton	139	<sup>∓</sup> lifetime			9-270	mm			<i>m</i> (ℓ)= 200 GeV	2011.07812
)	AMSB $pp \rightarrow \tilde{\chi}_1^\pm \tilde{\chi}_1^0, \tilde{\chi}_1^+ \tilde{\chi}_1^-$	disappearing track	136	$\tilde{\chi}_1^{\pm}$ lifetime				0	.06-3.06 m		$m(\tilde{\chi}_1^{\pm}) = 650 \text{ GeV}$	2201.02472
	AMSB $pp \rightarrow \tilde{\chi}_1^{\pm} \tilde{\chi}_1^0, \tilde{\chi}_1^{+} \tilde{\chi}_1^{-}$	large pixel dE/dx	139	$\tilde{\chi}_1^{\pm}$ lifetime				0.3-30.	0 m		$m(\tilde{\chi}_1^{\pm}) = 600 \text{ GeV}$	2205.06013
	Stealth SUSY	2 MS vertices	36.1	S lifetime			0	l.1-519 m			$\mathcal{B}(\tilde{g} \to \tilde{S}g) = 0.1, \ m(\tilde{g}) = 500 \text{ GeV}$	1811.07370
	Split SUSY	large pixel dE/dx	139	ĝ lifetime				> 0	.45 m		$m({ ilde g}){=}$ 1.8 TeV, $m({ ilde \chi}_1^0){=}$ 100 GeV	2205.06013
	Split SUSY	displaced vtx + $E_{\rm T}^{\rm miss}$	32.8	ĝ lifetime					0.03-13.2 n	1	$m( ilde{g}) =$ 1.8 TeV, $m( ilde{\chi}_1^0) =$ 100 GeV	1710.04901
	Split SUSY	0 $\ell,$ 2 – 6 jets $+ E_T^{miss}$	36.1	g lifetime	-	-	-	0.0	)-2.1 m		$m( ilde{g}) =$ 1.8 TeV, $m( ilde{\chi}_1^0) =$ 100 GeV	ATLAS-CONF-201
	$H \rightarrow s s$	2 MS vertices	139	s lifetime				0.31-72	2.4 m		<i>m</i> ( <i>s</i> )= 35 GeV	2203.00587
0/01	$H \rightarrow s s$	2 low-EMF trackless jets	139	s lifetime					0.19-6.94 m		m(s)=35 GeV	2203.01009
	$VH$ with $H \rightarrow ss \rightarrow bbbb$	2l + 2 displ. vertices	139	s lifetime		4-85	i mm				<i>m</i> ( <i>s</i> )= 35 GeV	2107.06092
	FRVZ $H  ightarrow 2\gamma_d + X$	2 µ-jets	139	$\gamma_d$ lifetime				0.654-939 mm			$m(\gamma_d) = 400 \text{ MeV}$	2206.12181
066	FRVZ $H  ightarrow 4 \gamma_d + X$	2 µ-jets	139	$\gamma_d$ lifetime			2.	.7-534 mm			$m(\gamma_d) = 400 \text{ MeV}$	2206.12181
	$H \rightarrow Z_d Z_d$	displaced dimuon	32.9	Z <sub>d</sub> lifetime		0.009-24.0 r	n				$m(Z_d) = 40 \text{ GeV}$	1808.03057
	$H \rightarrow ZZ_d$ 2	e, µ + low-EMF trackless	jet 36.1	Z <sub>d</sub> lifetime					0.21-5.2 m		$m(Z_d) = 10 \text{ GeV}$	1811.02542
	$\Phi(200 \text{ GeV}) \rightarrow ss$	ow-EMF trk-less jets, MS v	tx 36.1	s lifetime				0.41	-51.5 m		$\sigma \times B = 1 \text{ pb}, m(s) = 50 \text{ GeV}$	1902.03094
oraiai	$\Phi(600 \text{ GeV}) \rightarrow ss$	ow-EMF trk-less jets, MS v	tx 36.1	s lifetime			0.04-21.5	m			$\sigma \times \mathcal{B} = 1 \text{ pb}, m(s) = 50 \text{ GeV}$	1902.03094
ด้	$\Phi(1 \text{ TeV}) \rightarrow ss$	ow-EMF trk-less jets, MS v	tx 36.1	s lifetime			0.06-5	52.4 m	_		$\sigma \times \mathcal{B} = 1 \text{ pb, } m(s) = 150 \text{ GeV}$	1902.03094
_	$W \to N\ell, N \to \ell\ell\nu$	displaced vtx (µµ,µe, ee) +	μ 139	N lifetime		0.74-42 mm					m(N)= 6 GeV, Dirac	2204.11988
	$W \to N\ell, N \to \ell\ell\nu$	displaced vtx (μμ,μe, ee) +	μ 139	N lifetime		3.1-33 mm					m(N)= 6 GeV, Majorana	2204.11988
	$W \to N\ell, N \to \ell\ell\nu$	displaced vtx (µµ,µe, ee) +	e 139	N lifetime	-	0.49-81	mm				m(N)= 6 GeV, Dirac	2204.11988
	$W \to N\ell, N \to \ell\ell\nu$	displaced vtx (µµ,µe, ee) +	e 139	N life <mark>time</mark>	1	0.39-51 mm					m(N)= 6 GeV, Majorana	2204.11988
		√s = 13 TeV √s = 13			0.001	0.01	0.1		1 10		<sup>100</sup> cτ [m]	

### Parameter space coverage for $e^+e^- \rightarrow \gamma a \rightarrow \gamma \gamma \gamma$



4 experimental regions depending on decay length L of ALP

- •100 events for L<10 mm (prompt)
- •4 events for 10<L<2000 mm (Long lived) Decay in ID
- •4 events for 2000<L<4500 mm (Calo) Decay in calorimeter
- •100 events for L>4500 mm: ALP decays outside the detector, only accompanying photon detected (monophoton)

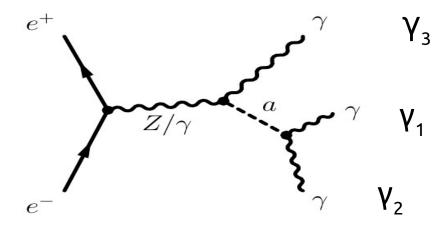
Experimental distinction of  $3\gamma$  prompt analysis and LLP analyses depends on how well one can detect a ALP decay away from vertex  $\rightarrow$  today show  $3\gamma$  analysis making no assumptions on vertex detection. In addition study very long-lived ALP resulting in a single photon recoiling against MET from undetected ALP

23/01/2025.

# 3γ analysis

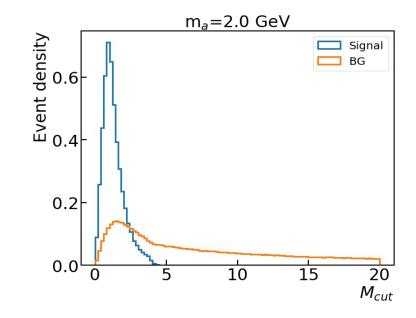
•3 photons within detector acceptance ( $|\eta| < 2.6$ ) and energy>1 GeV •Scan test masses M between 0.1 and 85 GeV For each M and  $E_{CM}$  photon produced alongside ALP has energy Assign three photons to ALP or to Z decay: For given test mass and assignment: Measure compatibility with expected kinematics

$$M_{cut} = \sqrt{(M_a - M)^2 / \sigma_{M_a}^2 + (E_{\gamma_3} - E_{\gamma})^2 / \sigma_{E_{\gamma_3}}^2)}$$



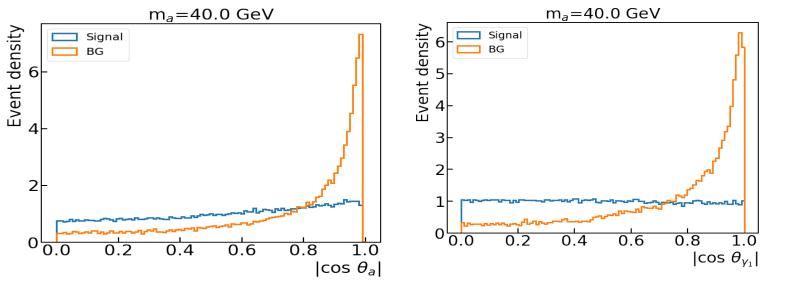
Choose assignment minimising M<sub>cut</sub>

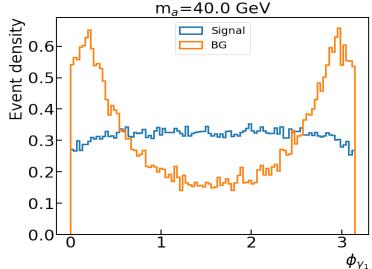
m(γ1, γ2)≡M<sub>a</sub>



 $E_{\gamma} = \frac{E_{CM}^2 - M^2}{2E_{CM}}$ 

## **Discriminant variables**





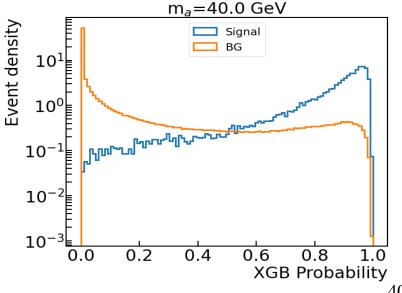
Require that event only contains three photons. For a fixed mass, signal fully defined by three variables, after rotation such that  $\varphi_{\gamma3}=0$ :

•Polar angle of ALP in lab system  $|\cos \theta_{\alpha}|$ 

•Polar angle of  $\gamma_1$  in ALP rest system  $|\cos \theta_{\gamma_1}|$ 

•Azimuthal angle of  $\gamma_1$  in ALP rest system  $\phi_{\gamma 1}$ 

Train a boosted decision tree (XGB) on 5 variables, the three above+  $E_{\gamma 2}/E_{\gamma 1}$  and  $M_{cut}$ 

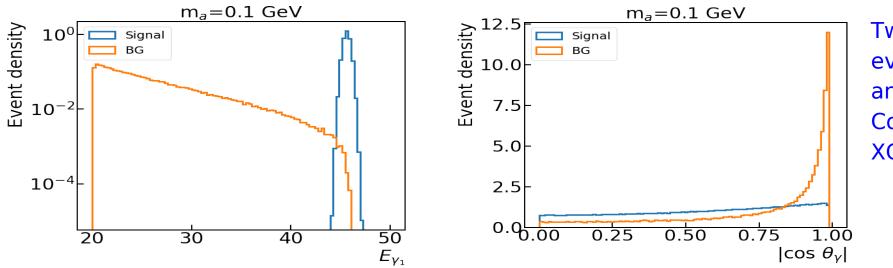


# $\gamma$ +MET analysis

Relevant mass range below ~2~GeV  $\rightarrow$  signature is a monochromatic photon of energy ~45.5 GeV and nothing else in the detector Consider two backgrounds: irreducible: e<sup>+</sup>e<sup>-</sup> $\rightarrow\gamma\nu\nu$ 

reducible: e+e-→γe+e- where the electron and positron are outside detector acceptance (|η|>3).

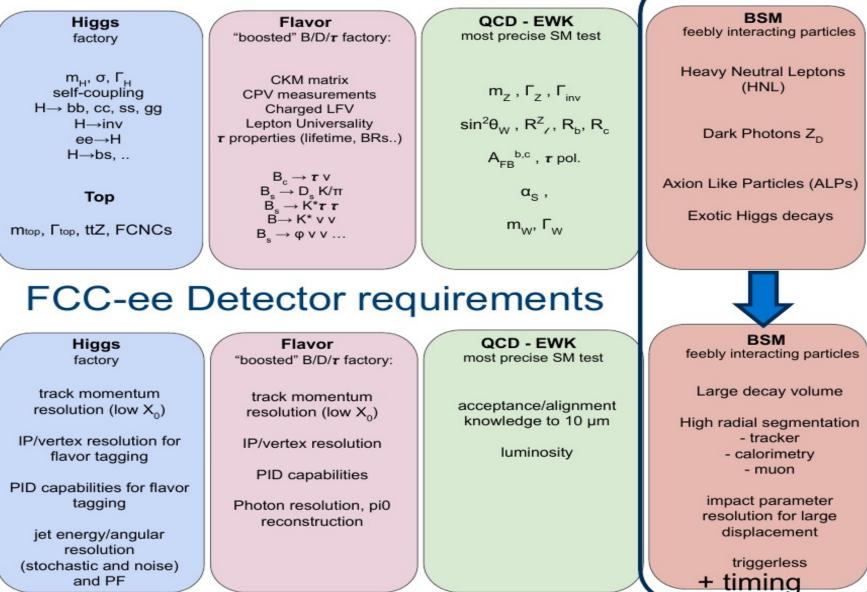
Signal and backgrounds produced with MG5MC@NLO and passed through the usual PYTHIA-DELPHES chain



Two variables characterise event, energy and polar angle of photon. Combine them through XGB as for prompt analysis

### FCC-ee: physics vs detector requirements

#### **FCC-ee Physics landscape**



FCC has very large menu of physics topics

Each of these poses a specific experimental challenge and pushes detector optimisation

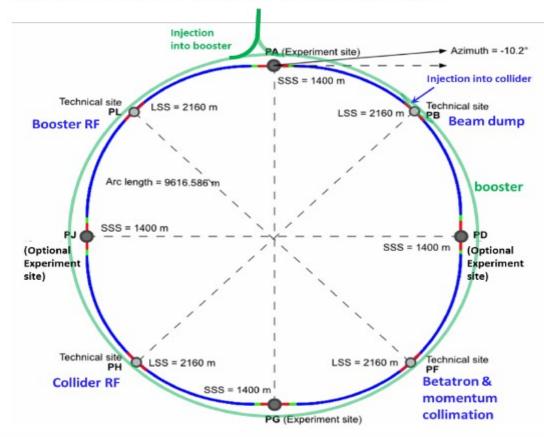
Unique challenges from BSM: long-lived particles

# Luminosity scenario

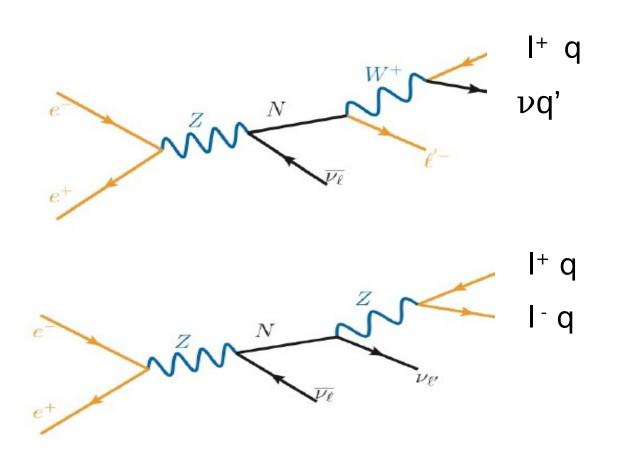
#### Lowest-risk baseline: 90.7 km ring, 8 surface points, 4-fold superperiodicity, possibility of 2 or 4 IPs

Whole project now adapted to this placement

4 Interaction points (IP): For Z pole run assume an Integrated luminosity of 205 ab-<sup>1</sup>, corresponding to 6 10<sup>12</sup> Z



# Decay signatures



Analysis matrix: for HNL •Decay final state  $(I=e,\mu)$ :

- j j l ~50% \*
- j j' nu ~20%
- I I nu ~5% \*
- I l' nu ~9%
- •lτnu ~9%
  - (BRs for  $m_{HN}$  < 80 GeV)
- Decay lengths
  - Prompt
  - LL decay In ID
  - LL decay in Calo

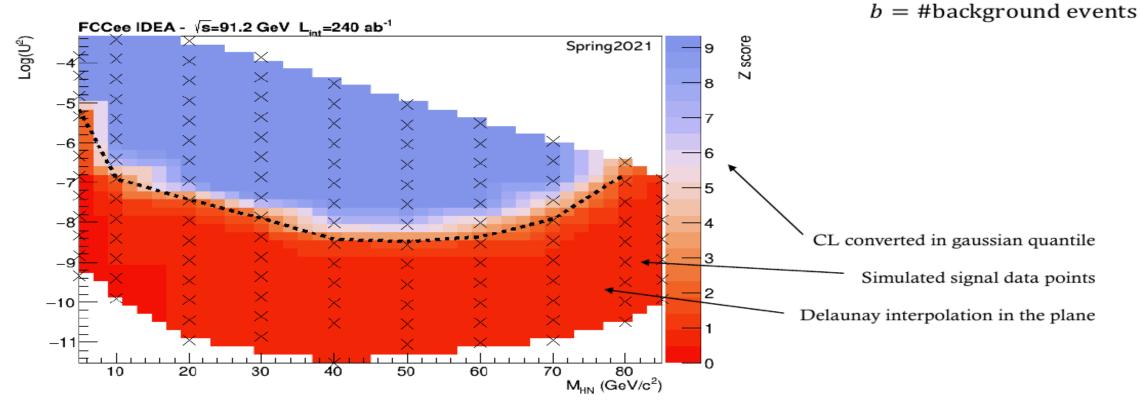
Signatures with \* studied in group

## Prompt results

#### GP, Nicolò Valle

- Baseline: Integrated Lumi =  $240 \text{ ab}^{-1} \leftrightarrow 8 \times 10^{12}$  Z boson events
- Looking for  $U^2$  producing 95% CL excess of events

For each HNL mass *M*:  $P[n < b | HNL(M, U^2)] = 1 - CL$ 

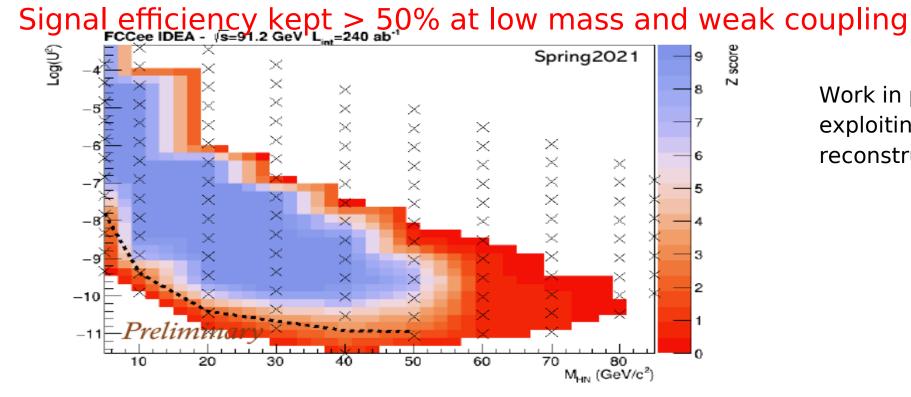


# LL results

Low mass (  $\leq$  40 GeV/c 2 ) HNL long-lived for couplings of interest, loss of efficency when requiring muon prompt

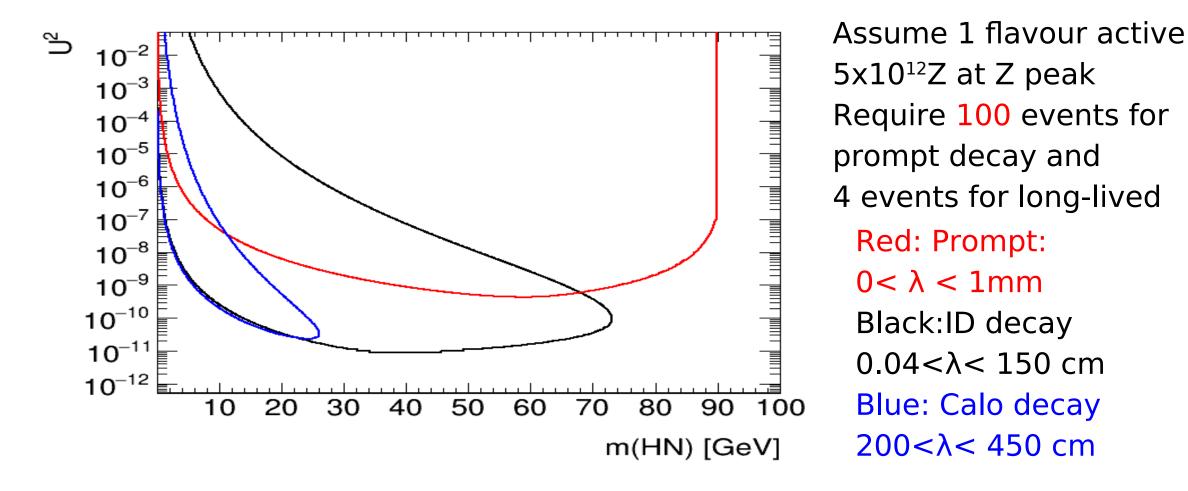
Background highly suppressed

Use detailed parameterization of IDEA tracking performance in DELPHES-FCC Kinematic selection not modified, prompt background suppressed by  $D_{\parallel} > 1 \text{ mm}$ 



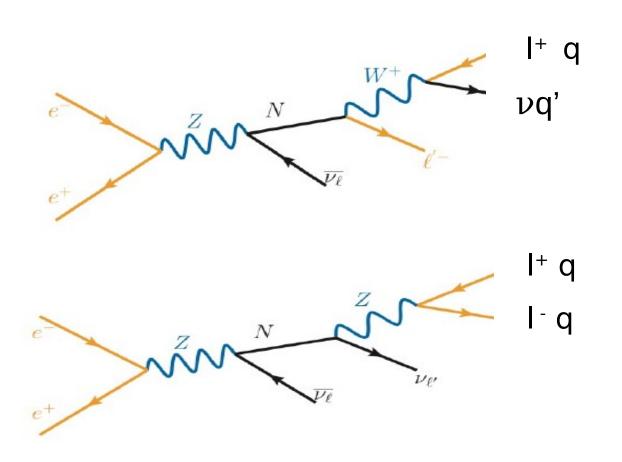
Work in progress on approach exploiting detailed HNL vertex reconstruction

# Linear scale



Prompt decays dominate for  $m_{HNL}$  >70 GeV

# Decay signatures



Analysis matrix: for HNL •Decay final state  $(I=e,\mu)$ :

- j j l ~50%
- •jjnu ~20%
- •llnu ~5%
- I l' nu ~9%
- •lτnu ~9%
  - (BRs for  $m_{HN}$  < 80 GeV)
- •Decay lengths
  - Prompt
  - LL decay In ID
  - LL decay in Calo

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- high precision measurement at the end of tracker
- σ<sub>rΦ</sub>
  - finely segmented vertex detector
- Challenging requirements for detector materials

Physics process	Measurands	Detector subsystem	Performance requirement
$ZH, Z \to e^+e^-, \mu^+\mu^-$ $H \to \mu^+\mu^-$	$m_H, \sigma(ZH)$ BR( $H \to \mu^+ \mu^-$ )	Tracker	$\Delta(1/p_T) = 2 \times 10^{-5} \oplus \frac{0.001}{p(\text{GeV}) \sin^{3/2} \theta}$
$H  ightarrow b ar{b} / c ar{c} / g g$	$BR(H \rightarrow b\bar{b}/c\bar{c}/gg)$	Vertex	$\sigma_{r\phi} = 5 \oplus rac{10}{p({ m GeV})  imes \sin^{3/2} heta}(\mu{ m m})$
$H \rightarrow q\bar{q}, WW^*, ZZ^*$	$\begin{array}{c} BR(H \to q\bar{q}, \\ WW^*,  ZZ^*) \end{array}$	ECAL HCAL	$\sigma_E^{ m jet}/E=3\sim4\%$ at 100 GeV
$H \to \gamma \gamma$	$\mathrm{BR}(H\to\gamma\gamma)$	ECAL	$\frac{\Delta E/E}{\frac{0.20}{\sqrt{E(\text{GeV})}} \oplus 0.01}$

Slide by R.Ferrari

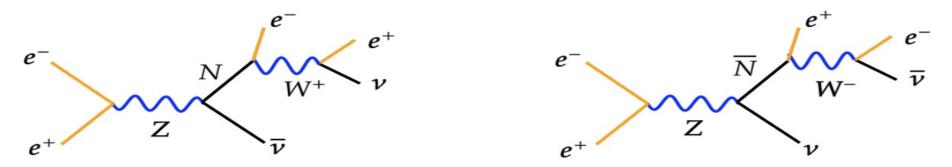
#### DELPHES setup for Spring 2021:

- •Detailed parametrisation of IDEA tracker, including covariance matrices
- •Calo resolution: EM 11%/sqrt(E), HAD: 30%/sqrt(E), 1% constant term
- •Particle flow approach to jet reconstruction

# Dirac versus Majorana

Blondel et al arXiv:2105.06576

No same-sign lepton signature as for  $W \rightarrow I HNL$ , rely on final state kinematics



• Dirac neutrinos  $(e^+e^- \rightarrow Z \rightarrow \nu \bar{N}; e^+e^- \rightarrow Z \rightarrow \bar{\nu}N)$ 

$$\frac{1}{\sigma_{N,\bar{N}}} \frac{d\sigma_{N,\bar{N}}}{d\cos\theta} \propto \left( g_R^2 (1 \mp \cos\theta)^2 + g_L^2 (1 \pm \cos\theta)^2 + \frac{M_N^2}{m_Z^2} (g_L^2 + g_R^2) \sin^2\theta \right)$$

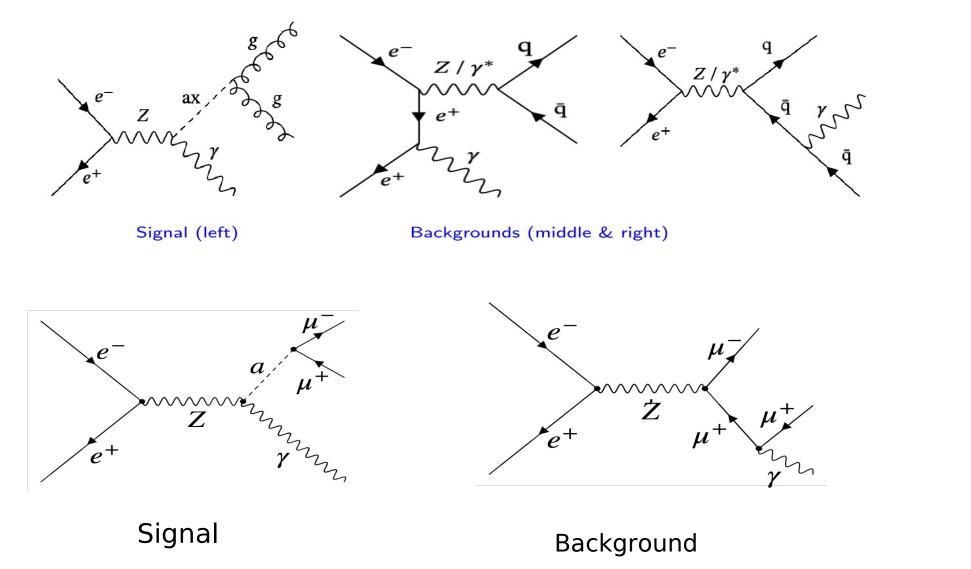
• Majorana neutrinos ( $e^+e^- \rightarrow Z \rightarrow \nu N$ )

$$\frac{1}{\sigma_N} \frac{d\sigma_N}{d\cos\theta} \propto \left(1 + \cos^2\theta + \frac{M_N^2}{m_Z^2}\sin^2\theta\right)$$

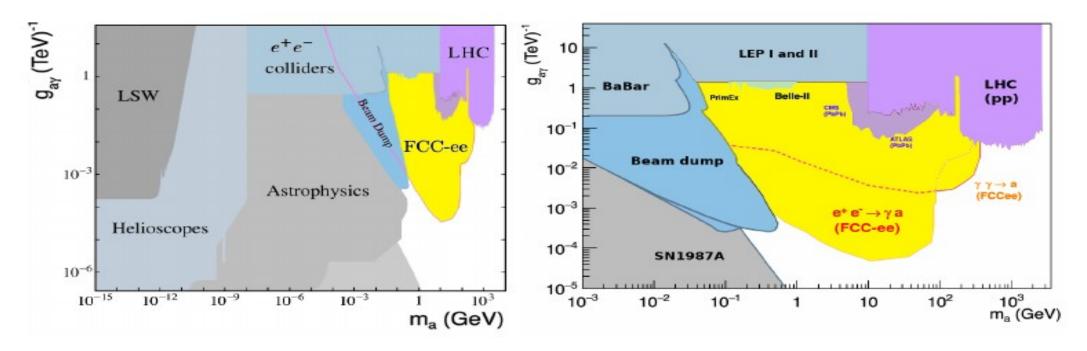
Relevant both for prompt and LLP, LLP has additional handle in lifetime 23/01/2025.

Giacomo Polesello – BSM at FCC-ee

### Additional ALP decay modes considered



#### Parameter space coverage

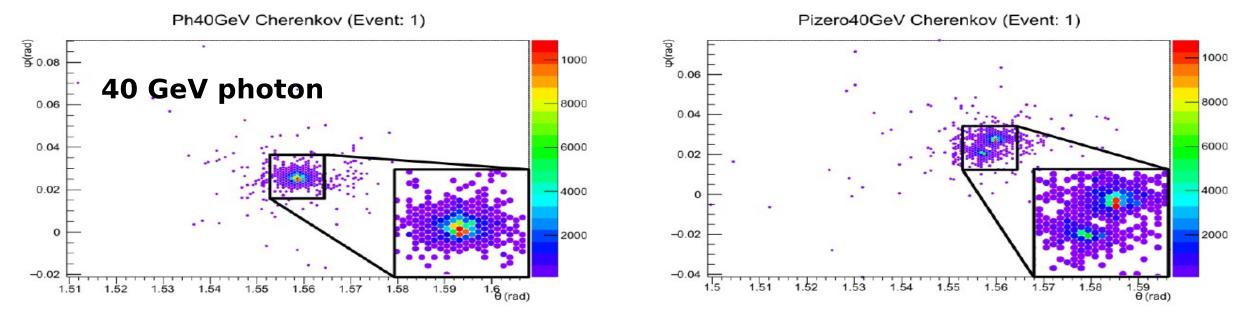


Plot in the MT report:  $e^+e^- \rightarrow \gamma a$  line is theory calculation requiring 4 ALP decays inside detector. 4 events might work for long-lived but prompt analysis has a huge irreducible background  $e^+e^+ \rightarrow \gamma \gamma \gamma$ , requiring detailed background analysis

Plots originally from Rebello Teles et al.

### Example: exploiting the full granularity of IDEA DR Calo

With Silicon PMs it is possible to read one by one all of the fibers in the calorimeter  $\rightarrow$  possibility to separate very close photons and to precisely measure invariant mass



Ideal field of application for ML image recognition, work ongoing in Pavia (master thesis A. Villa) 23/01/2025.

Giacomo Polesello - BSM at FCC-ee

### **Calorimeter parametrisation**

Take truth stable photons from PYTHIA tree in edm4hep, and smear them according to:

For DR fiber: performance figures from full simulation of testbeam prototype. Shown e.g in talk at ICHEP

$$\frac{\sigma(E)}{E} = \frac{0.139}{\sqrt{E}} + 0.006$$

$$\sigma(x) = \frac{4.05}{\sqrt{E}} + 0.0$$
  $\sigma(y) = \frac{3.23}{\sqrt{E}} + 0.0055$ 

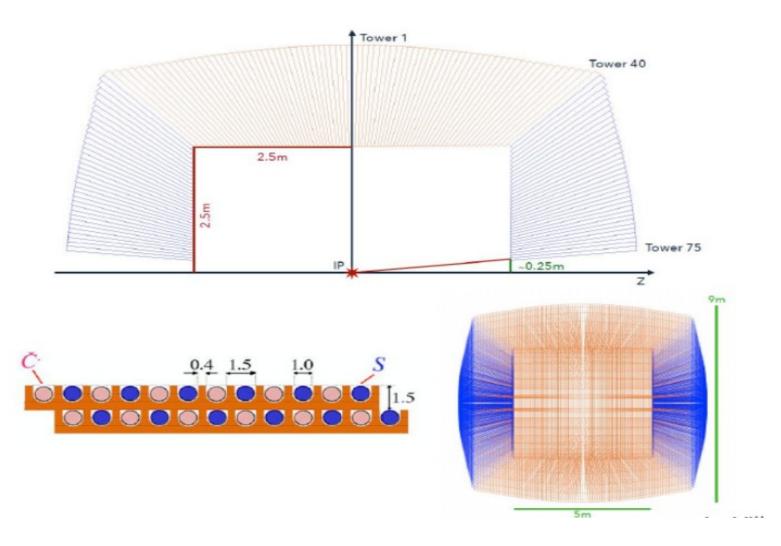
For crystal: energy resolution as in DELPHES card, Position resolution from Lucchini et al. paper

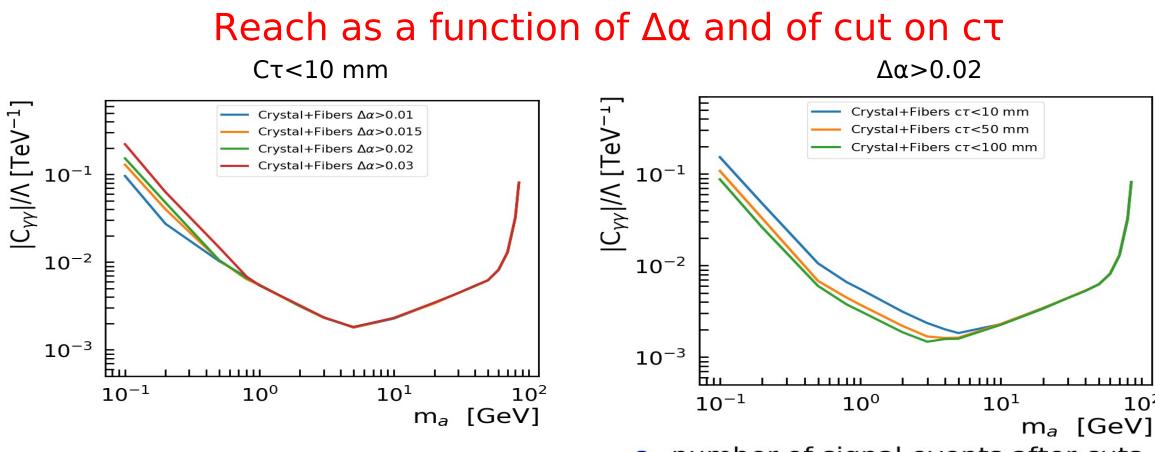
$$\frac{\sigma(E)}{E} = \frac{0.03}{\sqrt{E}} \oplus 0.005 \oplus \frac{0.002}{E}$$
$$\sigma(\theta) = \frac{1.5}{\sqrt{E}} \oplus 0.33$$

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## IDEA DR Calorimeter, old version





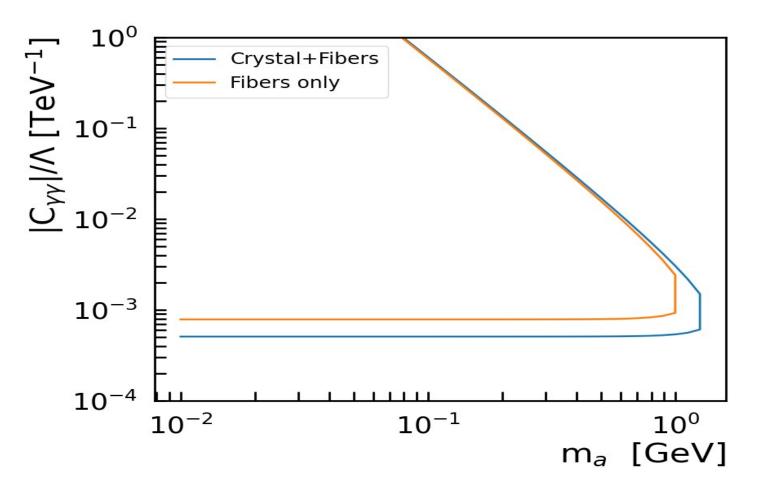
Plot  $2\sigma$  reach as function of mass and coupling, assuming 0.1% systematics Define significance as:

s=number of signal events after cuts b=background events after cuts n=s+b,  $\sigma =$  systematic uncertainty on b

$$Z = \sqrt{2\left(n \ln[\frac{n(b+\sigma^2)}{b^2 + n\sigma^2}] - \frac{b^2}{\sigma^2} \ln[1 + \frac{\sigma^2(n-b)}{b(b+\sigma^2]}]\right)}$$

10<sup>2</sup>

## Results



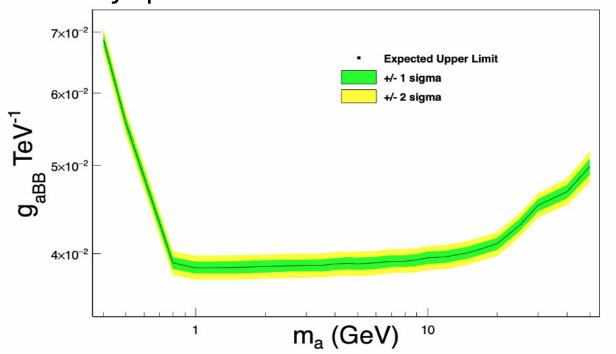
Irreducible background small at 45.6 GeV, but it increases very fast as energy goes down.

Smaller energy window determined by better resolution significantly increases reach

## A similar exercise

Recent paper: Steinberg, Wells, arXiv:2101.00520

Addressing the same model in the framework of ILC GigaZ ILC detector: R(ECal)~1.85 m. GARLIC photon reco: require photons with  $\Delta R > 0.035$  and with less than 10% of energy in reconstructed cone from nearby photon



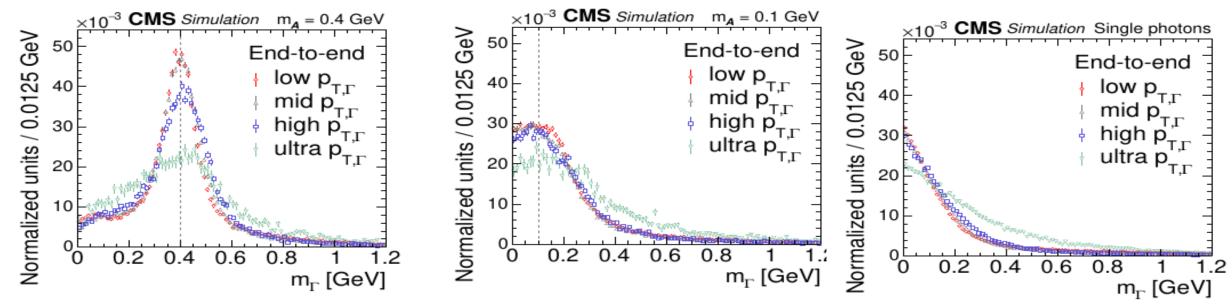
Simple analysis, require: •3 non-overlapping photons E>2 GeV •Εγ-Εγ<sup>recoil</sup><5 GeV

$$E_{\text{recoil}}^{\gamma}(m_a) = (M_Z^2 - m_a^2)/2M_Z$$

Significant loss in sensitivity, but in this setup search extended down to ALP masses if few hundred MeV

### An encouraging example from CMS

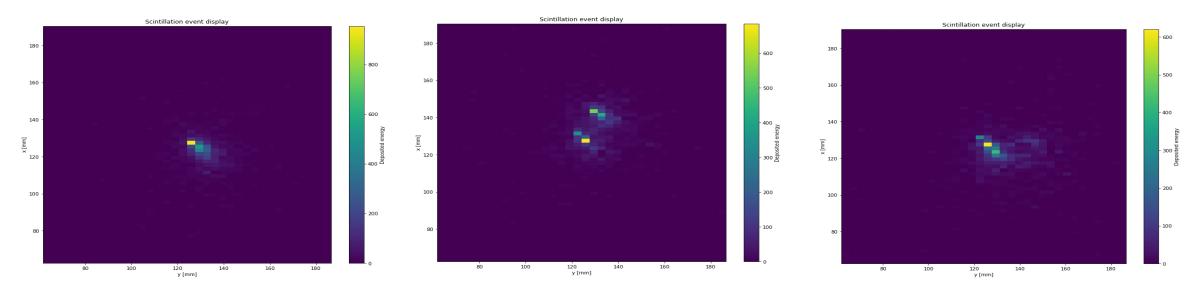
#### PRD 108 (2023) 052002



Using a CNN-based algorithm, reconstruct peak of 100 MeV particle. CMS granularity: 2.3 cm, IDEA Crystal: 1 cm IDEA Fiber: 2 mm Can probably improve on CMS result

### Two photons in fiber calorimeter

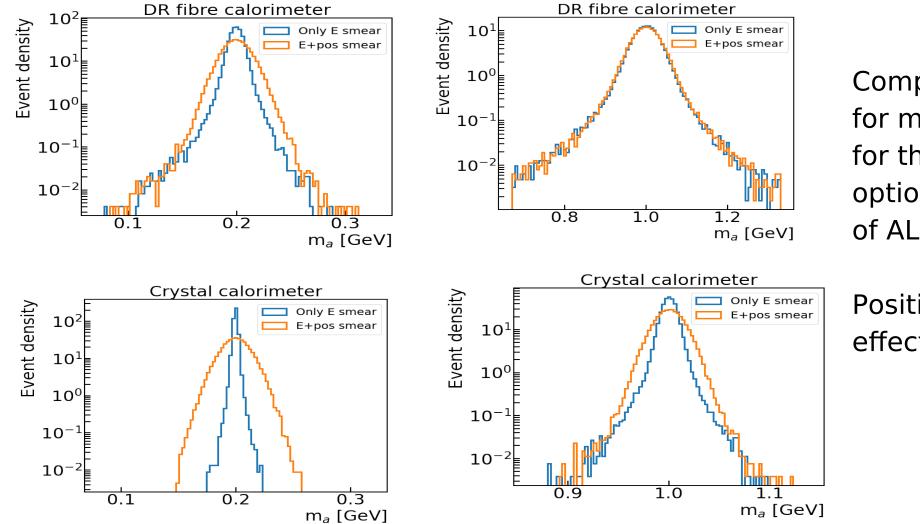
#### One fiber every 2 mm read with SiPMs



G4 simulation of energy deposition of a 40 GeV photon (left), and of two examples 40  $\pi^{\rm o}\,$  produced at 2m from a fiber calorimeter prototype

(Master thesis G.Salsi) Very high granularity can be exploited to measure the two clusters using image reconstruction techniques  $\rightarrow$  start work soon on that Waiting for results becoming available, reject events where  $\Delta \alpha$  between two PMONTOF is smaller than 0.01, 0.015, 0.02 0.03 and study reach as a function of cut

### Mass resolution



Compare mass resolution for  $m_a=0.2$ , 1 GeV for the two calorimeter options, for prompt decays of ALP

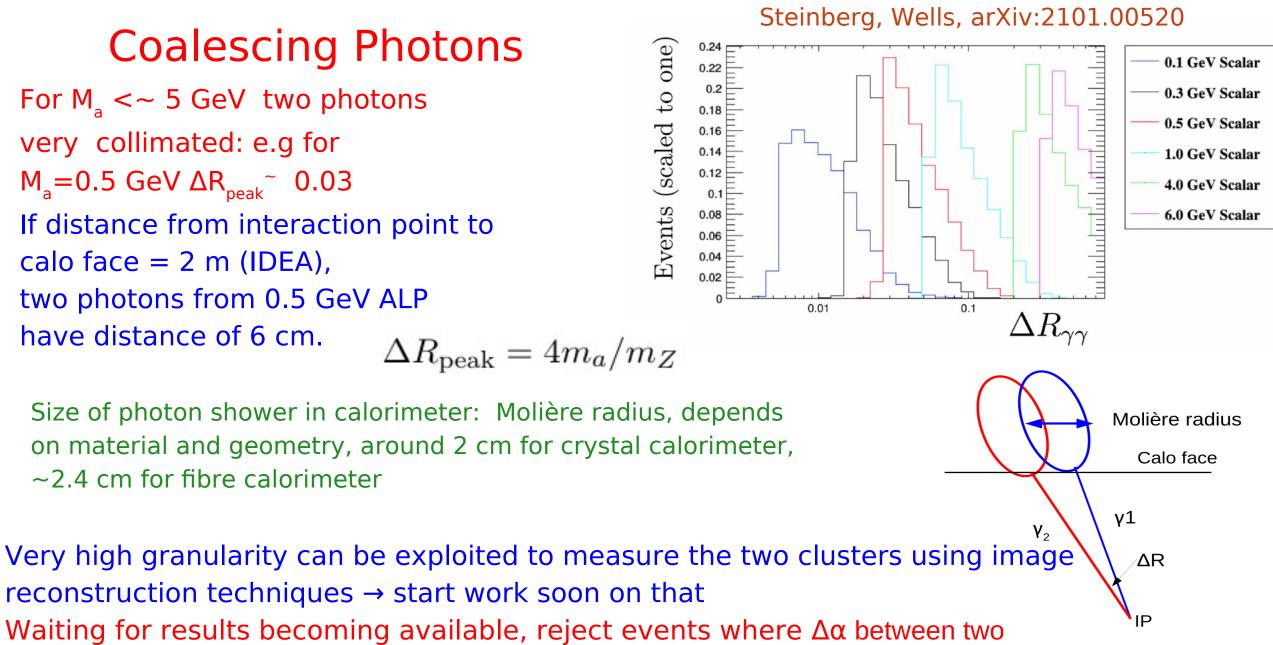
Position resolution dominant effect up to  $\sim 1 \text{ GeV}$ 

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## **Coalescing Photons**

For  $M_a < \sim 5$  GeV two photons very collimated: e.g for  $M_a = 0.5 \text{ GeV } \Delta R_{\text{peak}} \sim 0.03$ If distance from interaction point to calo face = 2 m (IDEA), two photons from 0.5 GeV ALP have distance of 6 cm.

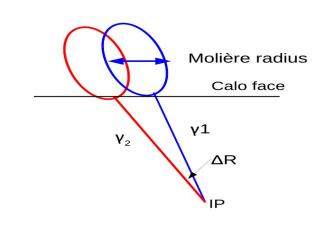
 $\sim$ 2.4 cm for fibre calorimeter



photons, smaller than 0.01, 0.015, 0.02, 0.03 and study reach as a function of cut Giacomo Polesello – BSM at FCC-ee

Experimental issues at low masses (~<5 GeV)

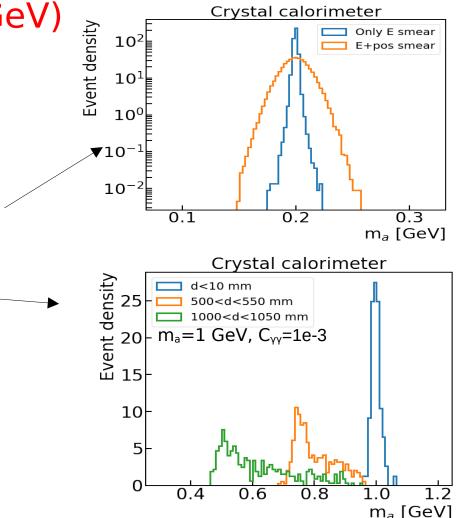
- Signal acceptance strongly affected by width of measured ALP mass
- At low masses three geometrical effects:
  - •γγ Mass resolution of dominated by uncertainty on measured photon impact point
    •ALP decaying far from interaction point: mass reconstruction assumes photons produced in centre of detector. If long decay path, γγ angle Δα and mass underestimated
    •γγ from ALP decay coalesce in calorimeter:



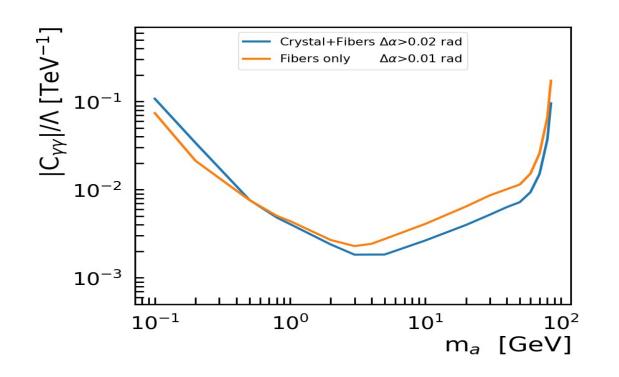
$$\Delta R_{\rm peak} = 4m_a/m_Z$$

Need full simulation for separation of nearby photons For this study assume two photons reconstructed as one If  $\Delta \alpha > 02$  (0.1) for crystal (fibre) EM calo

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## Results



For each signal and background sample events after cuts normalised to FCC-ee lumi s=number of signal events after cuts b=background events after cuts n=s+b,  $\sigma =$  systematic uncertainty on b Find cut on XGB output maximising significance calculated as:

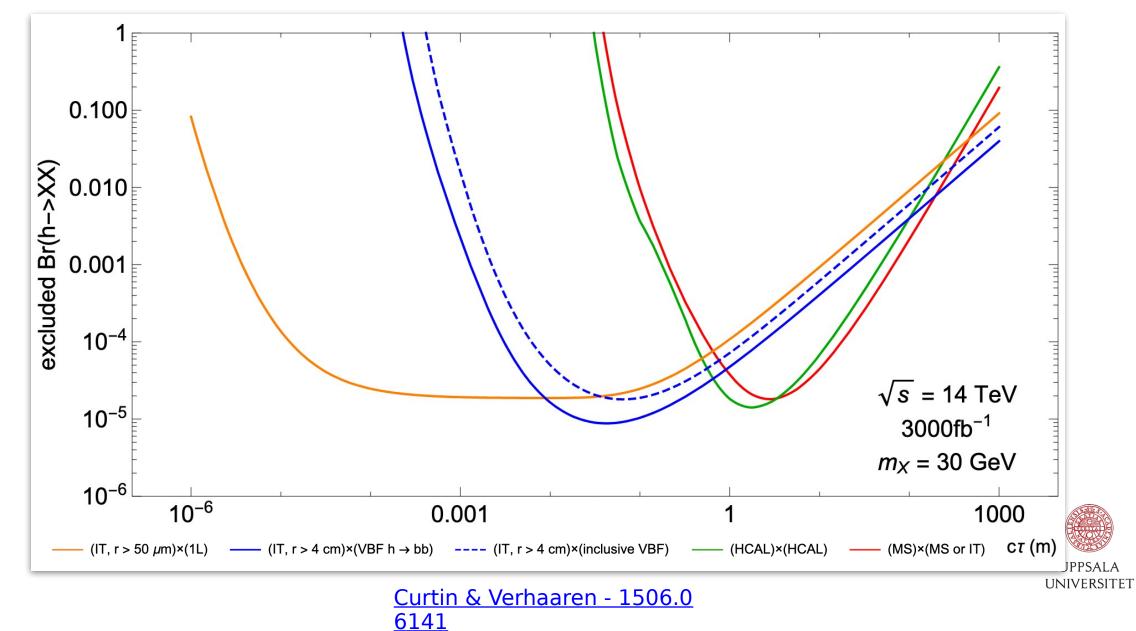
$$Z = \sqrt{2\left(n \ln[\frac{n(b+\sigma^2)}{b^2 + n\sigma^2}] - \frac{b^2}{\sigma^2} \ln[1 + \frac{\sigma^2(n-b)}{b(b+\sigma^2]}]\right)}$$

Significant advantage of better energy resolution at high masses At low masses better granularity should allow better separation of close-by photons

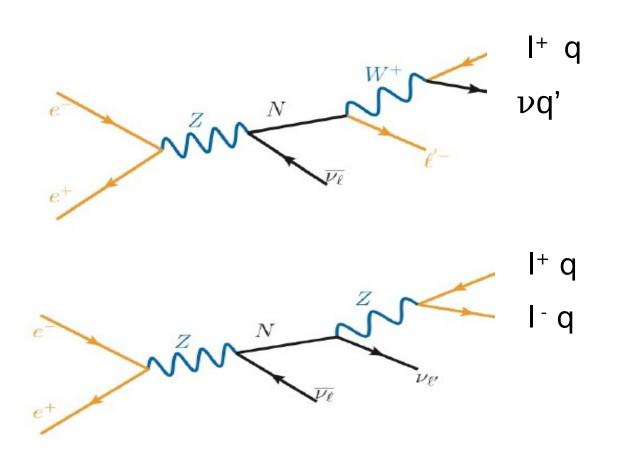
Cross-section proportional to  $C_{\gamma\gamma}^{2}$ For each test mass plot  $C_{\gamma\gamma}$  such that Z=2

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### Projected HL-LHC limits for exotic Higgs decays



# Decay signatures



Analysis matrix: for HNL •Decay final state  $(I=e,\mu)$ :

- j j l ~50% \*
- j j' nu ~20%
- I I nu ~5% \*
- I l' nu ~9%
- •lτnu ~9%
  - (BRs for  $m_{HN}$  < 80 GeV)
- Decay lengths
  - Prompt
  - LL decay In ID
  - LL decay in Calo

Signatures with \* studied in group

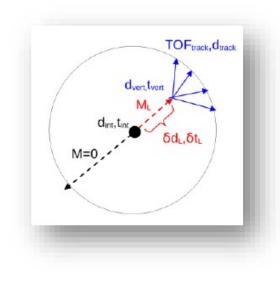
## Mass measurement through timing

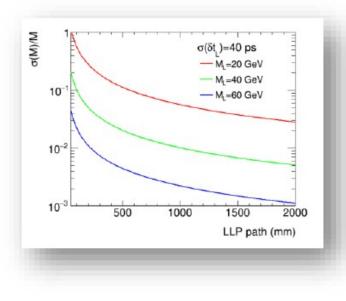
$$m_N = E_{cm} \sqrt{\frac{1 - \beta_N}{1 + \beta_N}} = E_{cm} F(\beta_N) \qquad \qquad \sigma(m_N) \sim E_{cm} F'(\beta_N) \sigma(\beta_N) \qquad \qquad \beta_N = \frac{\delta d_N}{\delta t_N}$$

The HNL mass can be constrained by measuring its decay timing and path

Resolution controlled by the uncertainty on HNL decay time and on the undetected interaction point \*

\*  $\sigma_x$  = 5.96 µm,  $\sigma_y$  = 23.8 nm,  $\sigma_z$  = 0.397 mm,  $\sigma_z$  = 36.3 ps

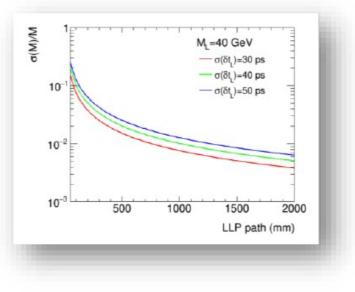




Measurement below the percent level is possible with plausible detector performance,

for sufficiently high masses

and long lifetimes

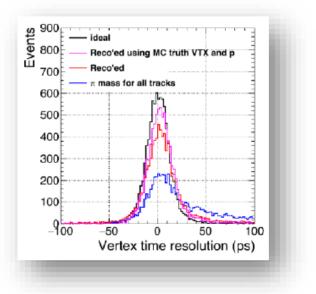


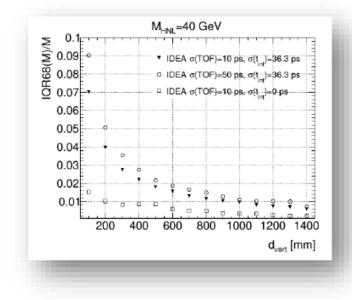
# Mass measurement through timing

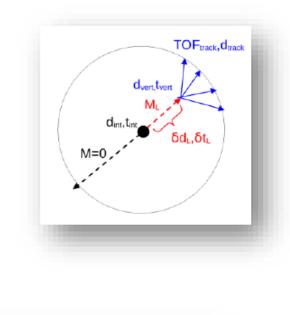
#### arXiv:2406.05102

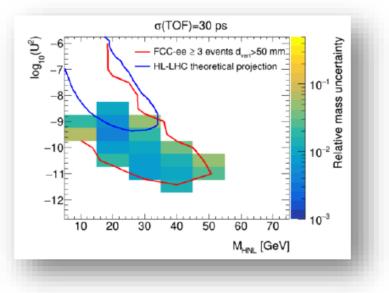
#### **Realistic conditions** simulated in IDEA, using the $N \rightarrow \mu j j$ channel

- arphi  $\sigma$  (TOF) determined only by detector technology
- $\checkmark$  The HNL vertex is known and its flight distance is computed
- Iterative procedure set up to optimize the mass hypotheses, possibly spoiled by the long HNL flight distance
- $\checkmark$  Timing resolution roughly scaling with sqrt of number of tracks
- $\sphericalangle~200 \mu m \simeq \sigma(d_{vert})$  dominated by the uncertainty on the interaction point
- $\checkmark$  Dependence on HNL yield vs ( $m_N$ , U): evaluated with MC for the expected Z-pole run luminosity

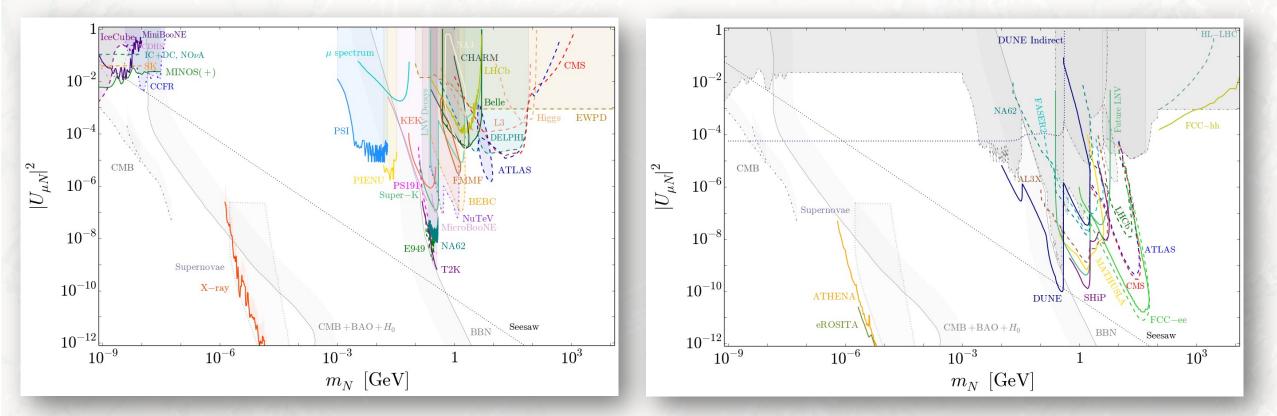








## Existing limits and projections



arXiv:1912.03058

#### Dependence on hadronic resolution

- Window for baseline study from DELPHES
- Assume signal efficiency unchanged after enlarging mass window according to resolution 2.
- 3. Calculate number of background events for enlarged window and calculate significance

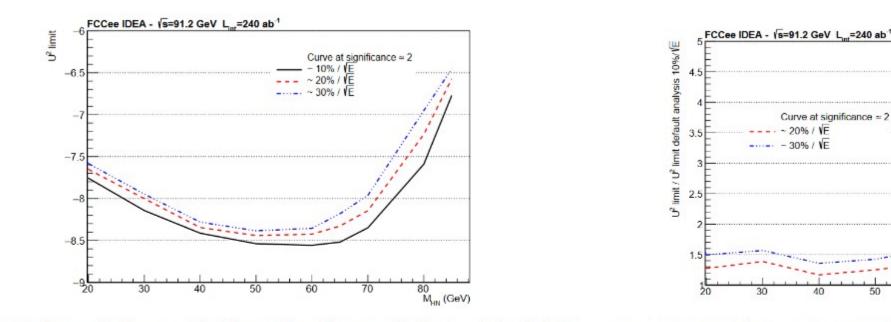


Fig. 24 Curves at Significance = 2 for different values of the assumed hadronic resolution. Each line is a linear interpolation of Z vs.  $\log(U)$  at the value Z = 2.

Fig. 25 Ratio of the  $U^2$  limit obtained with 20% and 30% resolutions with respect to the nominal resolution as a function of  $M_{N_1}$ .

50

60

70

80

M<sub>HN</sub> (GeV)