

Montecarlo and tools for lepton colliders

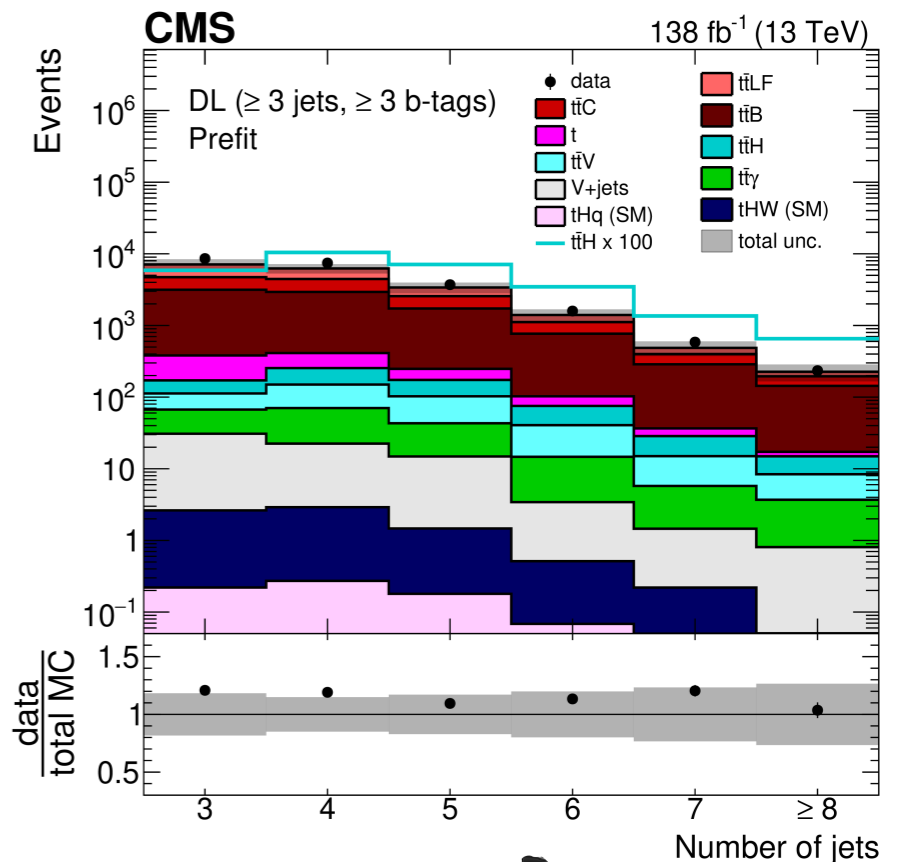
Marco Zaro, Milan U. & INFN



Istituto Nazionale di Fisica Nucleare
Sezione di Milano

From the LHC to a future collider

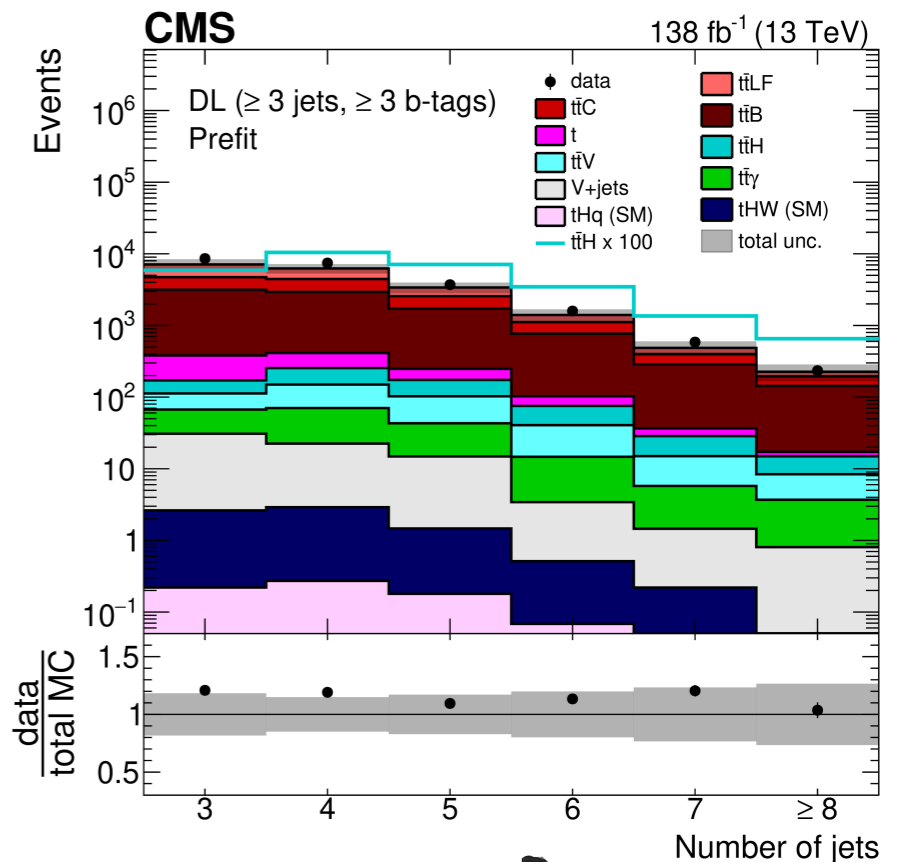
- MonteCarlo generators are fundamental workhorses for collider physics
 - Assess the reach of new machines
 - Extract results and compare theory vs data
- Lepton colliders demand $<0.1\%$ precision
 - Prominent role of EW corrections
 - NLO mandatory, often not enough
- In general, two kind of tools:
 - Process specific, with the highest achievable accuracy
 - General-purpose (mostly relying on automatic techniques), with a satisfactory accuracy



From the LHC to a future collider

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What I'll present focuses on the latter, but may be relevant also for the former



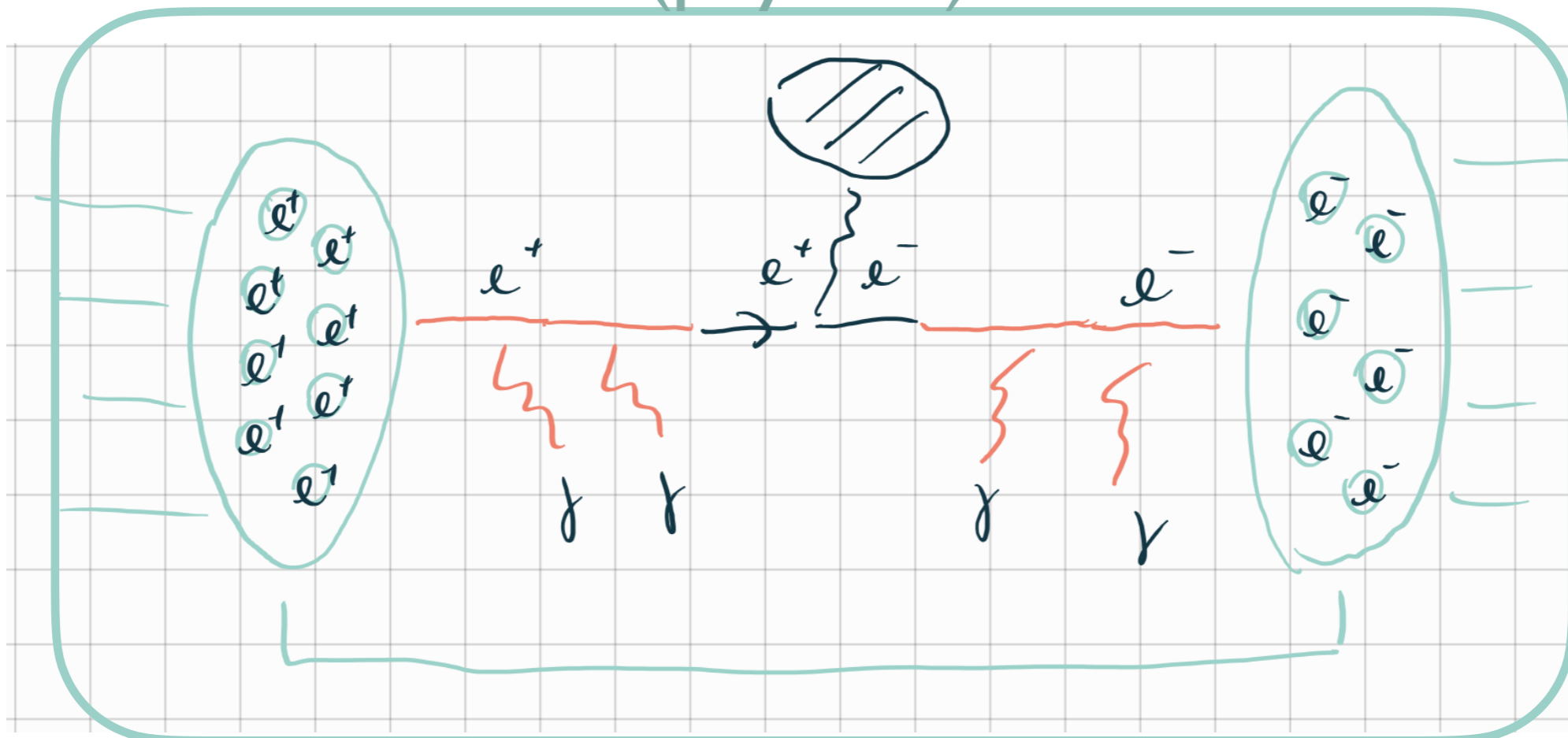
Lepton collisions

(In a collinear-factorisation inspired picture)

$$d\Sigma_{e^+e^-}(P_{e^+}, P_{e^-}) = \sum_{kl} \int dy_+ dy_- \mathcal{B}_{kl}(y_+, y_-) d\sigma_{kl}(y_+ P_{e^+}, y_- P_{e^-})$$

$$d\sigma_{kl}(p_k, p_l) = \sum_{ij} \int dz_+ dz_- \Gamma_{i/k}(z_+, \mu, m) \Gamma_{j/l}(z_-, \mu, m) d\hat{\sigma}_{ij}(z_+ p_k, z_- p_l, \mu)$$

Beam-level (physical) cross section



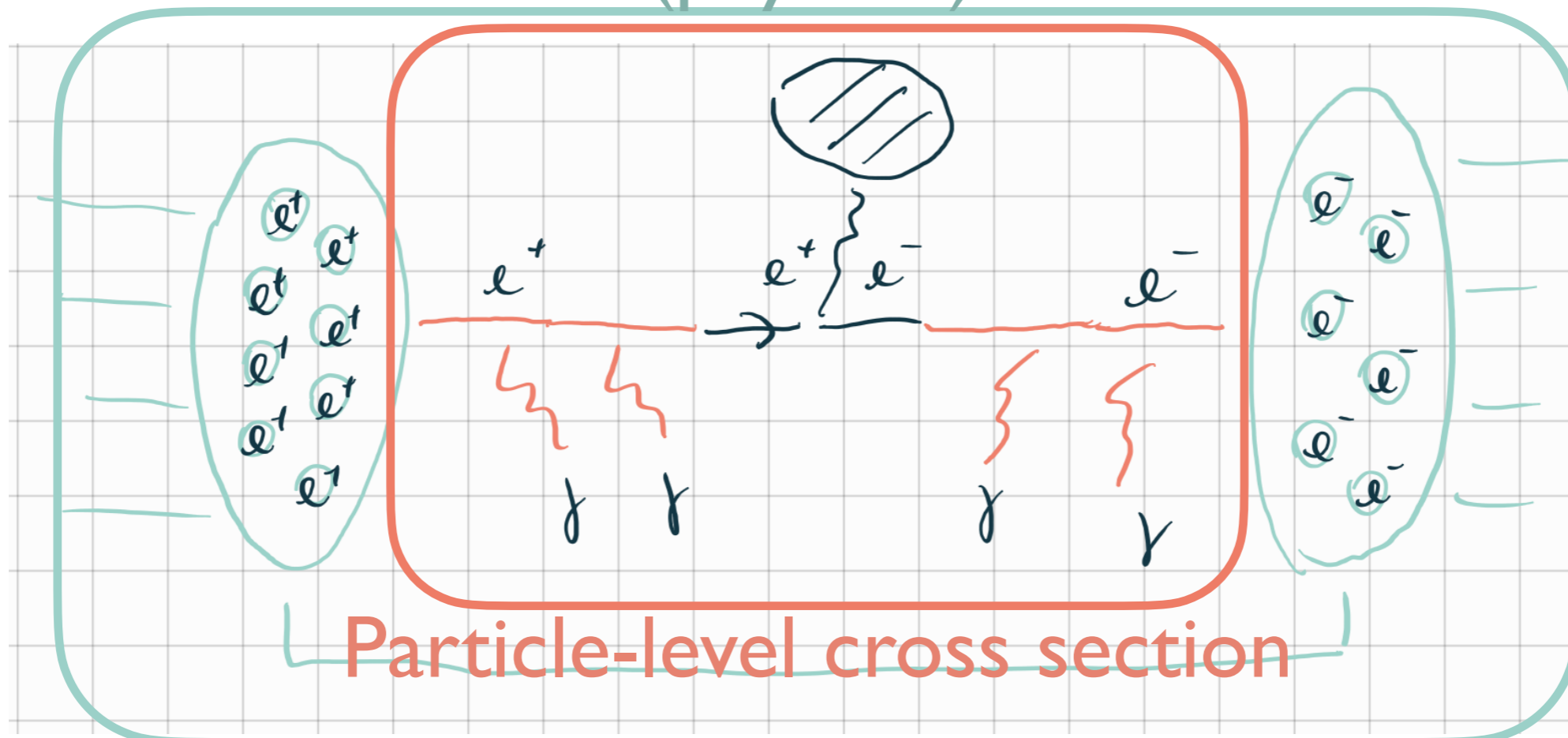
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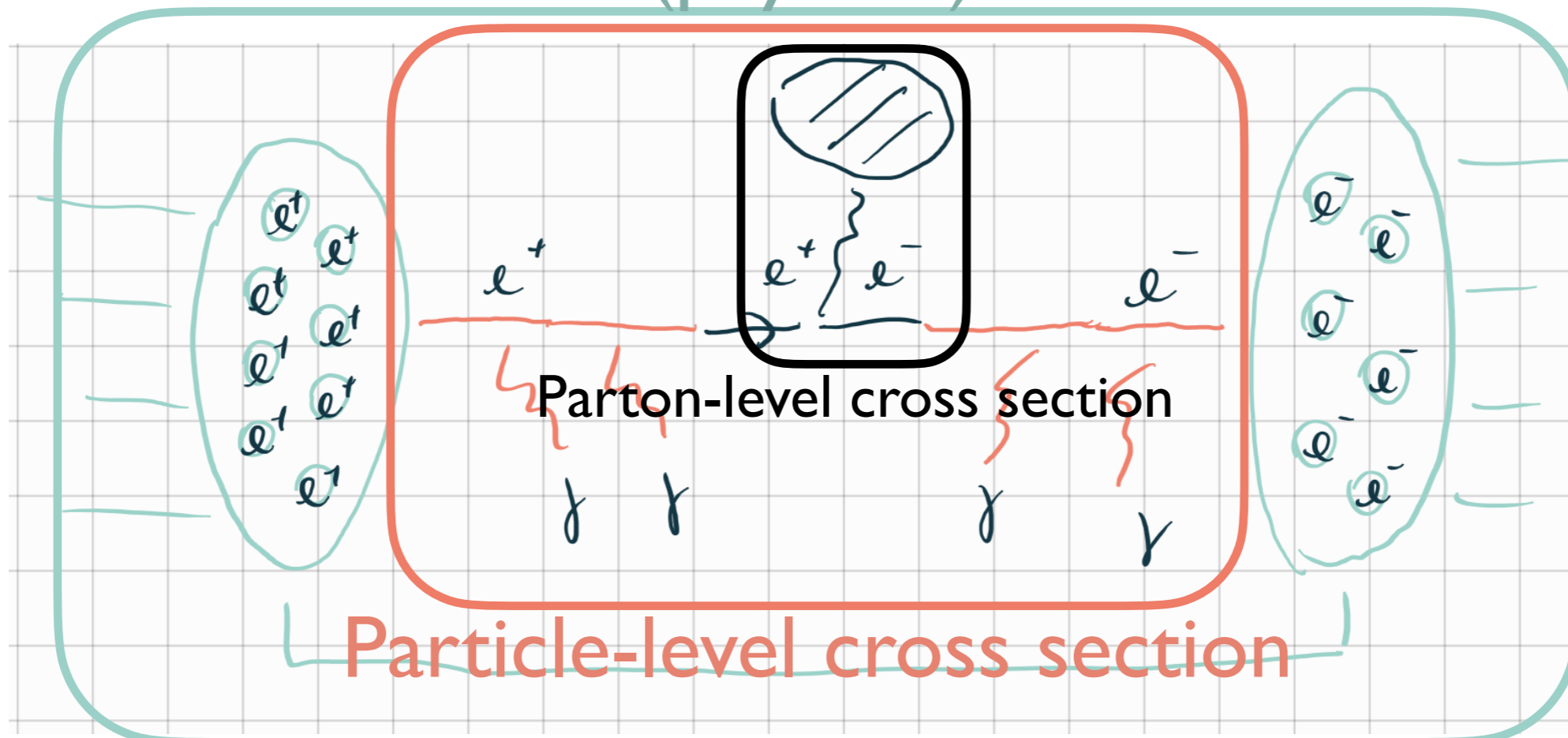
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Beam-level (physical) cross section





Lepton collisions

- **Beam-beam interactions** (aka beamstrahlung) are machine-dependent collective effects. Can be fitted with ad-hoc tools (e.g. GuineaPig, Circe, ...). Less important for circular colliders than for linear ones
- **ISR** is universal (like hadronic PDFs), and can be computed perturbatively (*unlike* hadronic PDFs)
- The **partonic cross section** is the usual one. Because of the form of PDFs, needs new phase-space and momentum mappings



Beamstrahlung

- Description of beam-beam interactions outside the scope of MC generators. Specific tools exist
- E.g.: simulated beamstrahlung with GuineaPig, and fit the spectra (energy spread functions) with Circe
Circe: Ohi, hep-ph/9607454; GuineaPig: Schulte, 1998
- Specific to collider parameters
- Effects are usually tabulated as grids, pre-convoluted with PDFs
- Need special care when matching with parton-shower:
 - The PS enters between ISR and beamstrahlung

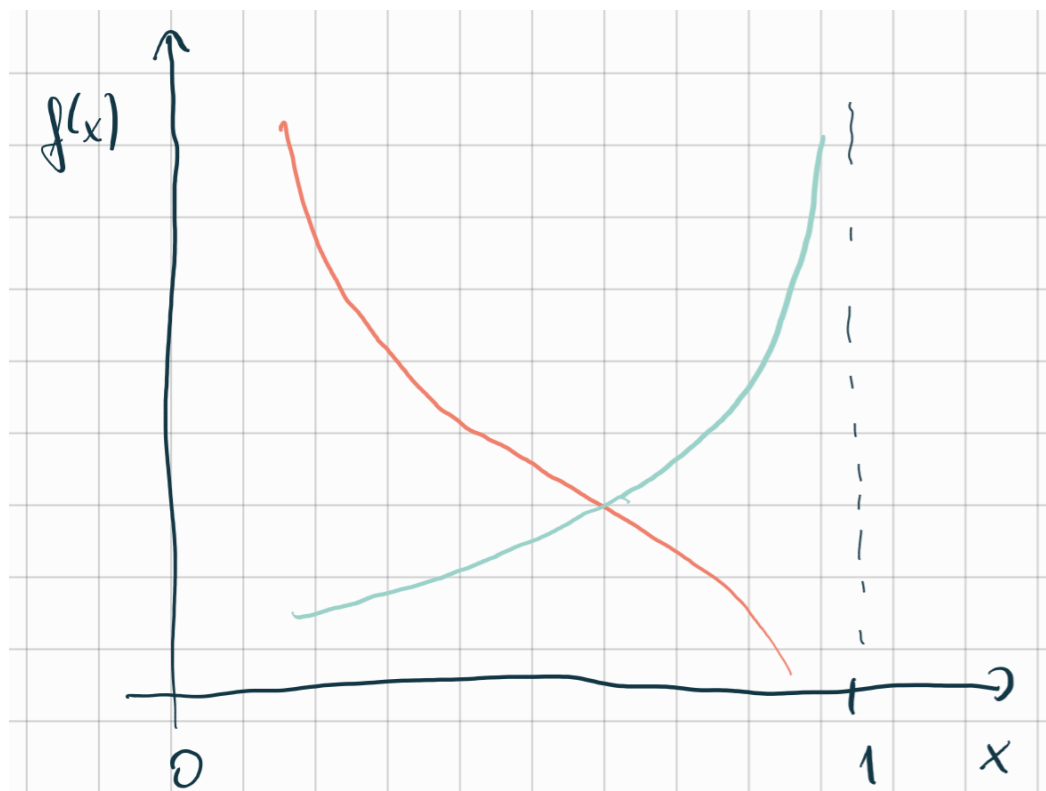


ISR

at LL: GRIBOV V. N. and LIPATOV L. N.
Sov. J. Nucl. Phys., 15 (1972) 438

$$D_{GL}(x, Q^2) = \frac{\exp[(1/2)\eta(3/4 - \gamma_E)]}{\Gamma(1 + (1/2)\eta)} \frac{1}{2} \eta (1-x)^{(1/2)\eta-1} \simeq \frac{1}{(1-x)^{1-\eta/2}} = \delta(1-x) + \mathcal{O}(\eta)$$

$$\left(\eta = \frac{2\alpha}{\pi} \log \frac{Q^2}{m^2} \right) \sim 0.05 \text{ for } Q=100 \text{ GeV}$$



- **Hadronic PDFs** vanish at large x (divergence at small-x avoided by cuts)
- **Leptonic PDFs** diverge (but are integrable) at large x
- While leptonic PDFs have been substantially improved since 1972, the asymptotic behaviour is unchanged
- A different phase space mapping is required wrt pp collisions



ISR for e^+e^- :

State of the art

- Recently, the ISR structure function was obtained at NLL accuracy¹. This required the NLO initial conditions²
- A new factorisation scheme (alternative to $\overline{\text{MS}}$) has been proposed (Δ -scheme)³, which improves the behaviour of the evolved PDF at large x
- PDFs available with α in three ren. scheme: G_μ , $\alpha(m_Z)$, $\overline{\text{MS}}$ (with proper treatment of all thresholds)⁴. All available within eMELA⁵
- Photon and e^+ -in- e^- densities available as well

¹ Bertone, Cacciari, Frixione, Stagnitto, [1911.12040](#)

² Frixione [1909.03886](#)

³ Frixione [2105.06688](#)

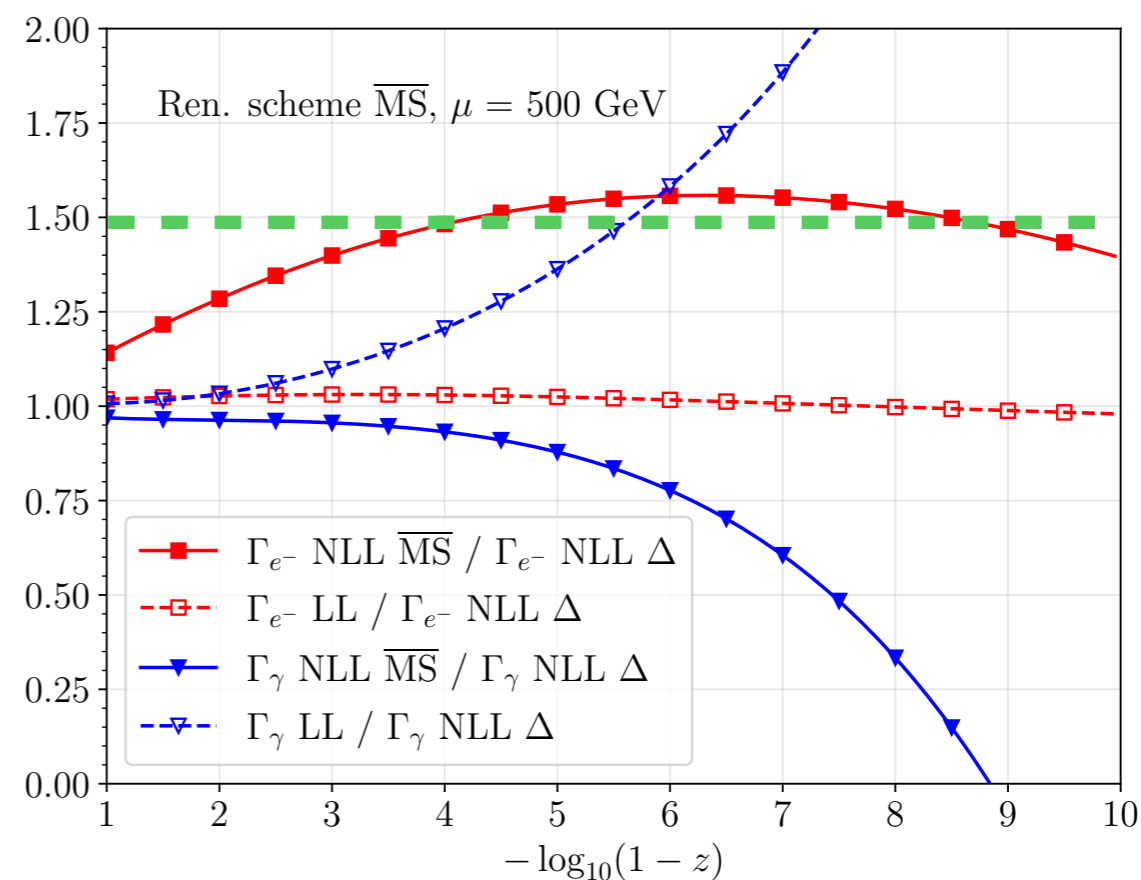
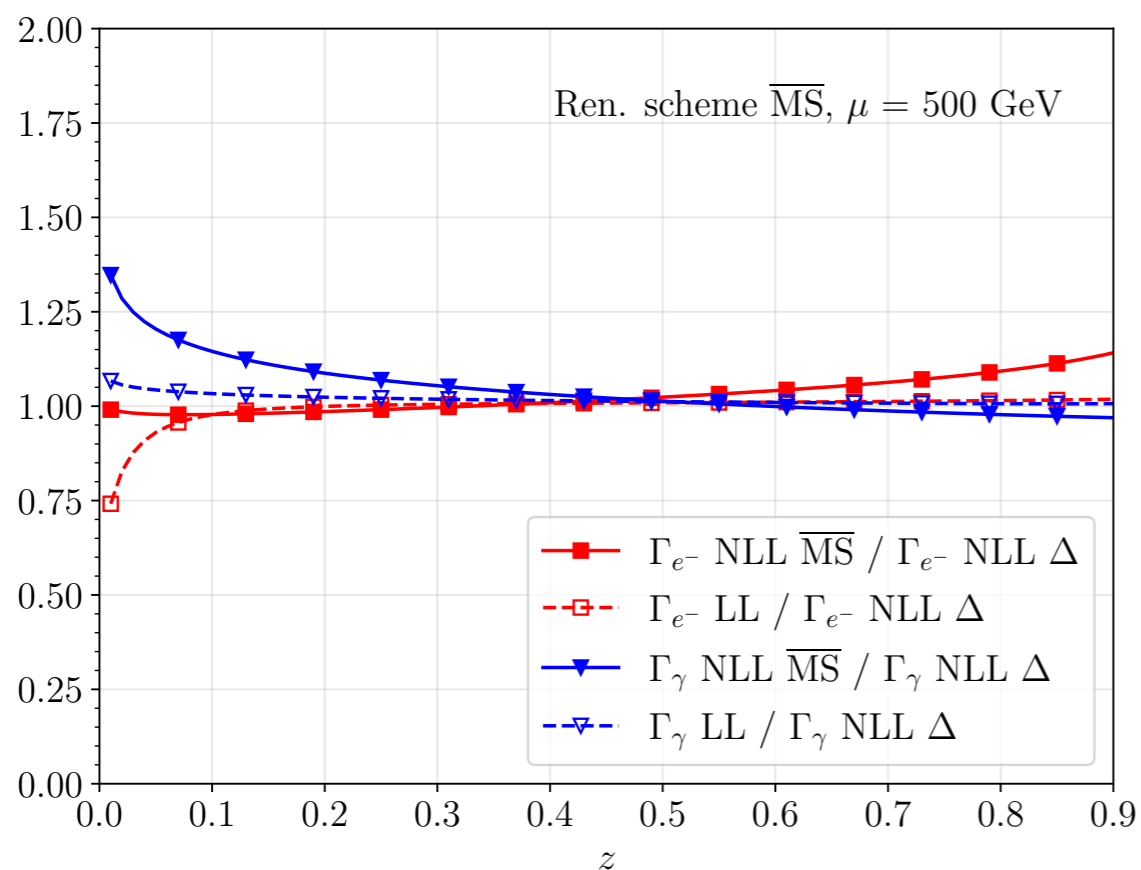
⁴ Bertone, Cacciari, Frixione, Stagnitto, MZ, Zhao [2207.03265](#)

⁵ <https://github.com/gstagnit/eMELA>



ISR at NLL

Bertone, Cacciari, Frixione, Stagnitto, MZ, Zhao [2207.03265](#)

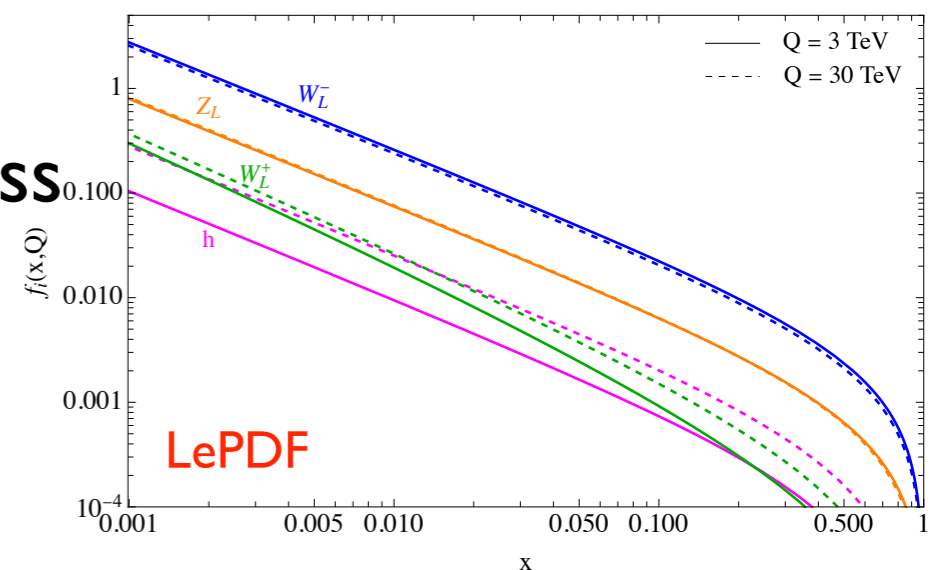
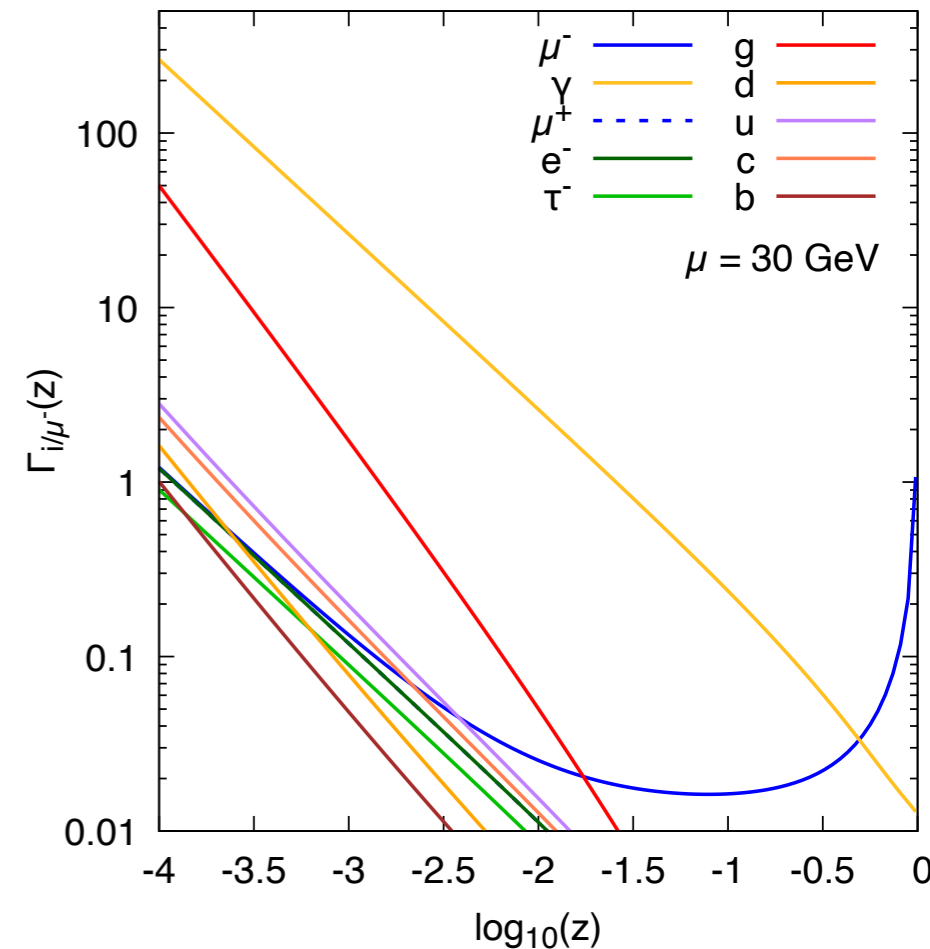


- NLL- $\overline{\text{MS}}$ seems very different from LL, while NLL- Δ is closer
- Differences between NLL- $\overline{\text{MS}}$ and NLL- Δ are 10-50% at large x
- Physical cross sections (NLO-accurate) will display much smaller discrepancies



ISR for $\mu^+\mu^-$

- In principle, a muon is just a heavier electron, so most of what has been said so far should apply
- However, muon colliders can reach far higher energy, lower values of x are probed
 - The photon, and even quarks and gluons are important in this region
- Muon density including all massless/light particles available in eMELA at LL (QED+QCD), WIP at NLL (+small x)
[Frixione, Stagnitto, 2309.07516, +Bonvini WIP](#)
- In principle, one can also treat W/Z as massless particles in those regions where $m_W \ll \hat{s}$
[Han, Ma, Xe, 2007.14300,](#)
[Garosi, Marzocca, Trifinopoulos, 2303.16964 \(LePDF\)](#)
 See Davide Pagani's talk





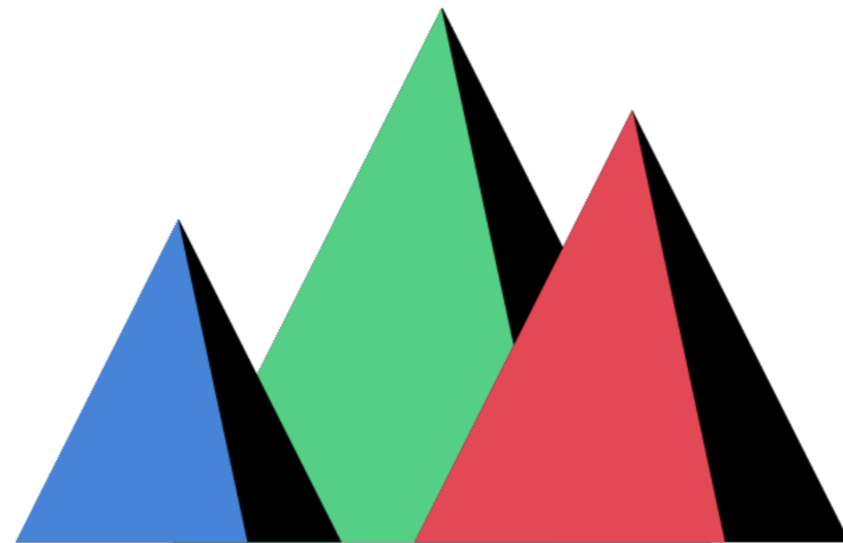
General-purpose MCs for lepton colliders



MG5_aMC



Whizard



SHERPA



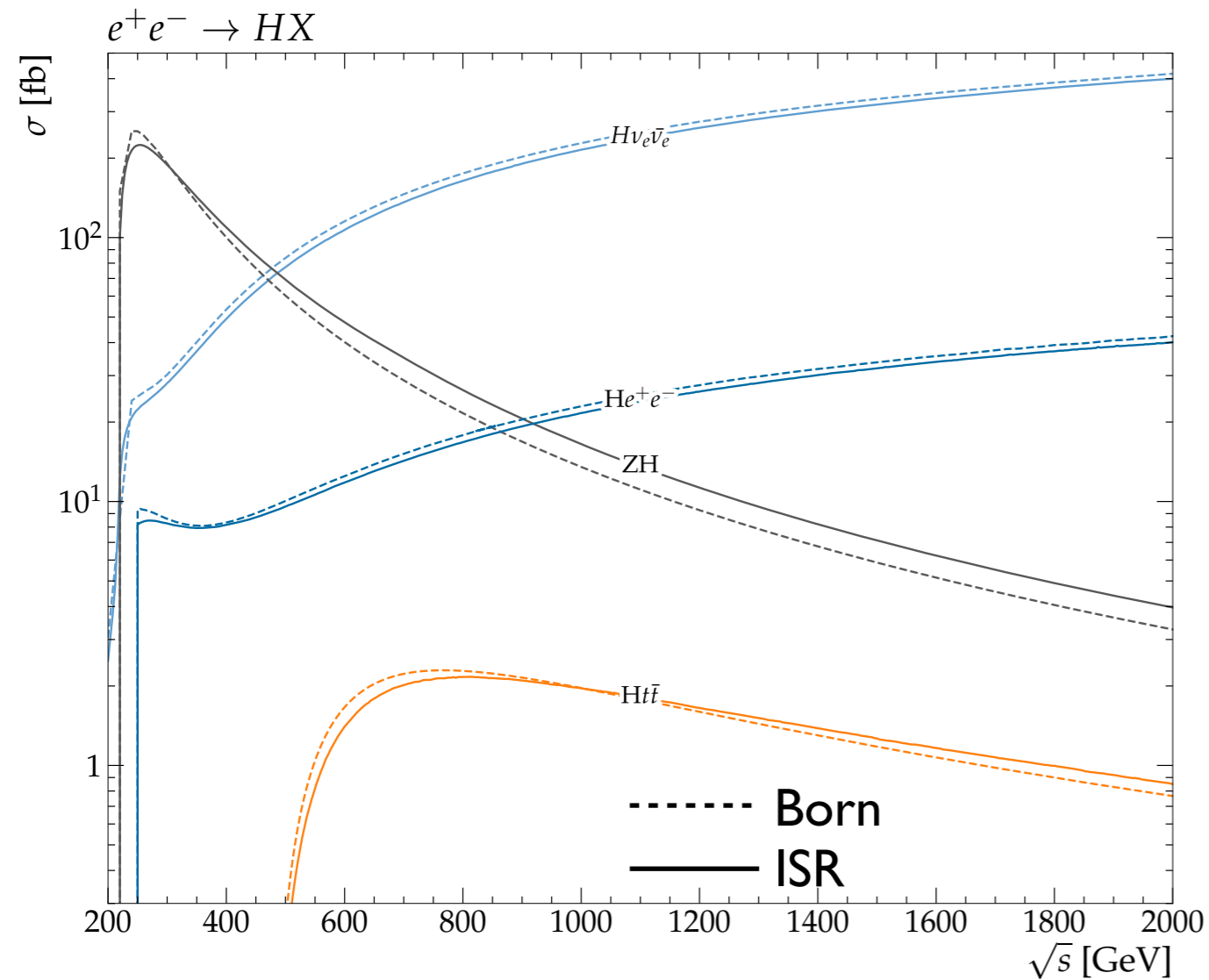
Sherpa3 at lepton colliders

Bothmann et al, 2410.22148; Krauss et al, 2203.10948

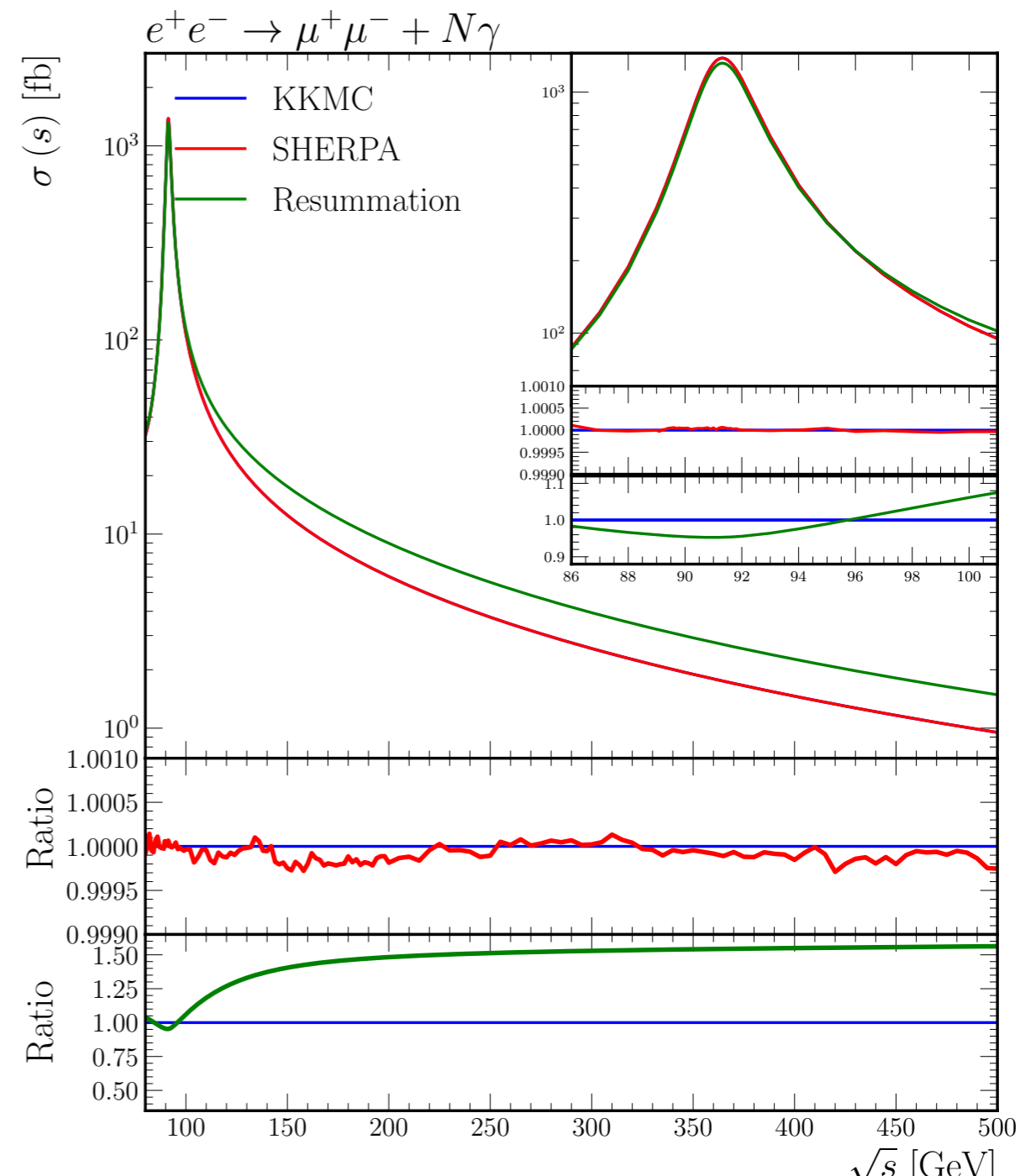
- Treatment of polarised beams
- Photon resummation available both with PDFs (LL) and with YSF
- EW corrections for lepton collider in the future
- Photoproduction with the EPA [Hoche et al, 2310.18674](#)
- QED effects in parton shower, including $\gamma \rightarrow f\bar{f}$ [Flower et al, 2210.07007](#)

Results with Sherpa

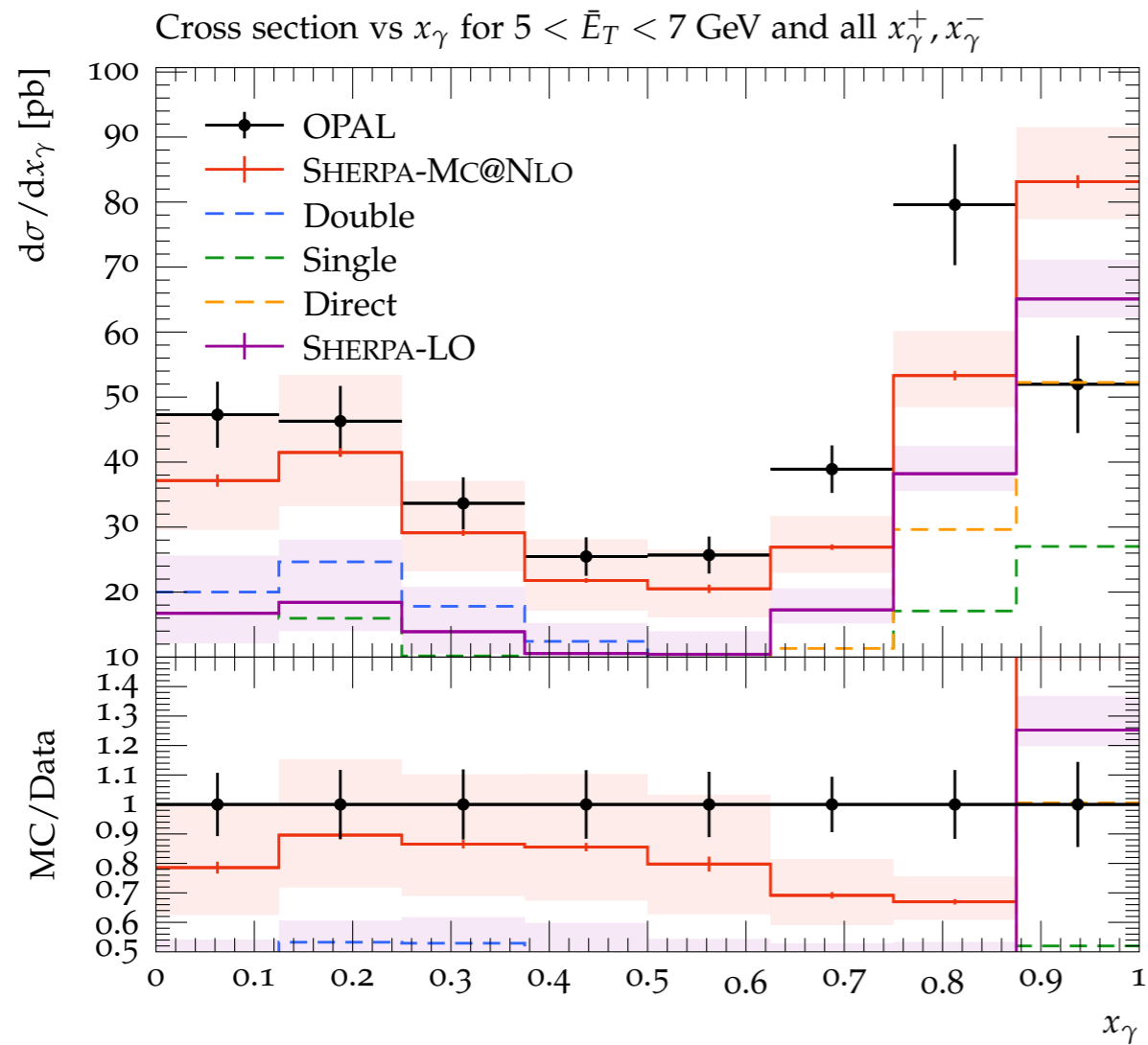
Krauss et al, 2203.10948



$$d\sigma = \sum_{n_\gamma=0}^{\infty} \frac{e^{Y(\Omega)}}{n_\gamma!} d\Phi_Q \left[\prod_{i=1}^{n_\gamma} d\Phi_i^\gamma \tilde{S}(k_i) \Theta(k_i, \Omega) \right] \left(\tilde{\beta}_0 + \sum_{j=1}^{n_\gamma} \frac{\tilde{\beta}_1(k_j)}{\tilde{S}(k_j)} + \sum_{\substack{j,k=1 \\ j < k}}^{n_\gamma} \frac{\tilde{\beta}_2(k_j, k_k)}{\tilde{S}(k_j) \tilde{S}(k_k)} + \dots \right)$$

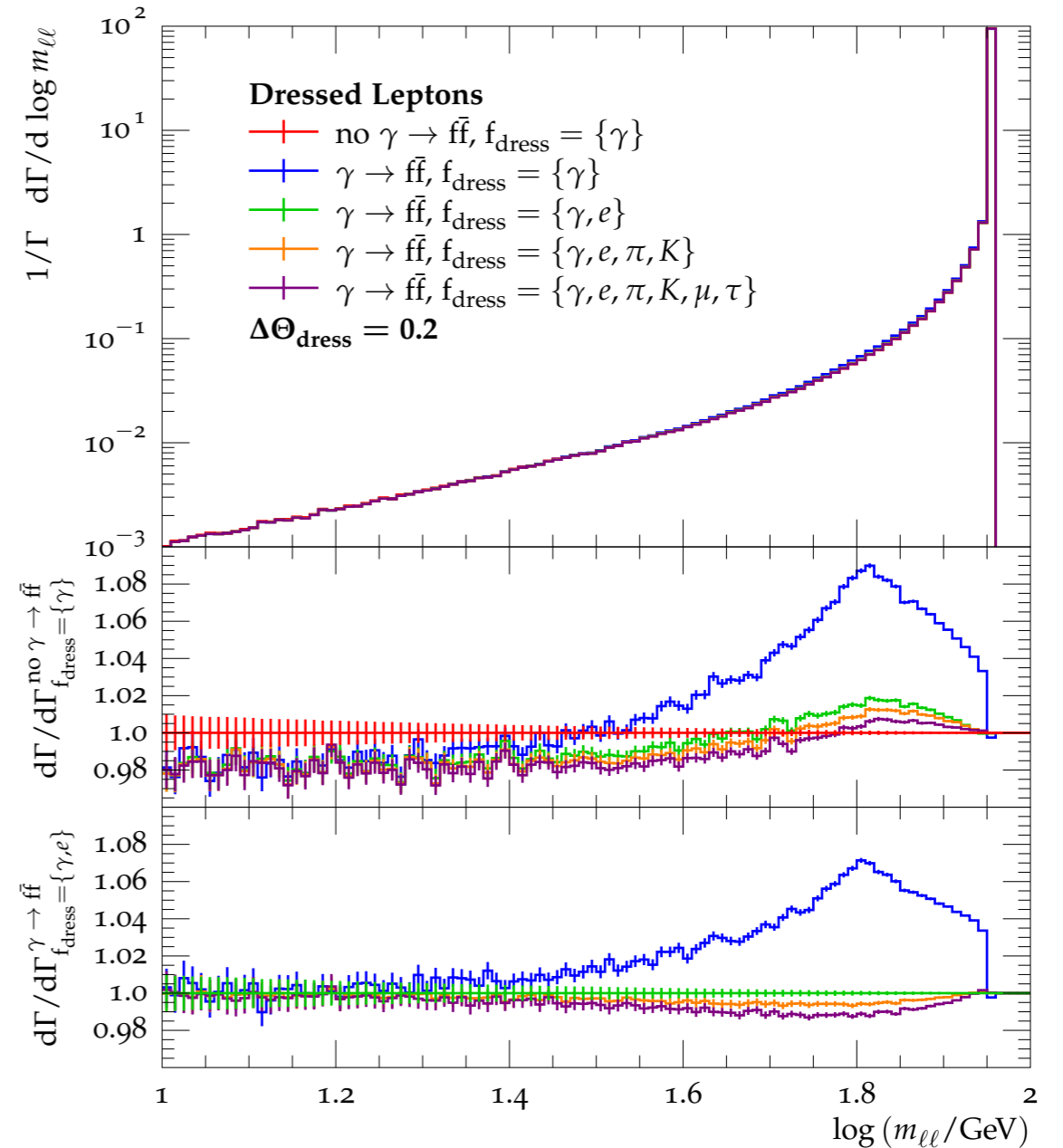


Results with Sherpa



Photoproduction at LEP:

Hoche et al, 2310.18674



$\gamma \rightarrow f\bar{f}$ splittings:

Flower et al, 2210.07007



Whizard at lepton colliders

- NLO QCD/EW accuracy for arbitrary processes
- NLO QCD matched à la Powheg
- Beamstrahlung included via Circe
- ISR at LL, wip for NLL (independent from eMELA)
- Full support of polarised beams
- WIP for EW PDFs



Results with Whizard

ee @ 1 TeV, NLO QCD

Process	WHIZARD+OpenLoops	
	σ_{LO} [fb]	σ_{NLO} [fb]
$e^+e^- \rightarrow jj$	622.737(8)	639.39(5)
$e^+e^- \rightarrow jjj$	340.6(5)	317.8(5)
$e^+e^- \rightarrow jjjj$	105.0(3)	104.2(4)
$e^+e^- \rightarrow jjjjj$	22.33(5)	24.57(7)
$e^+e^- \rightarrow jjjjjj$	3.583(17)	4.46(4)
<hr/>		
$e^+e^- \rightarrow t\bar{t}$	166.37(12)	174.55(20)
$e^+e^- \rightarrow t\bar{t}j$	48.12(5)	53.41(7)
$e^+e^- \rightarrow t\bar{t}jj$	8.592(19)	10.526(21)
$e^+e^- \rightarrow t\bar{t}jjj$	1.035(4)	1.405(5)
$e^+e^- \rightarrow t\bar{t}t\bar{t}$	$0.6388(8) \cdot 10^{-3}$	$1.1922(11) \cdot 10^{-3}$
$e^+e^- \rightarrow t\bar{t}t\bar{t}j$	$2.673(7) \cdot 10^{-5}$	$5.251(11) \cdot 10^{-5}$
<hr/>		
$e^+e^- \rightarrow t\bar{t}H$	2.020(3)	1.912(3)
$e^+e^- \rightarrow t\bar{t}Hj$	$2.536(4) \cdot 10^{-1}$	$2.657(4) \cdot 10^{-1}$
$e^+e^- \rightarrow t\bar{t}Hjj$	$2.646(8) \cdot 10^{-2}$	$3.123(9) \cdot 10^{-2}$
$e^+e^- \rightarrow t\bar{t}Z$	4.638(3)	4.937(3)
$e^+e^- \rightarrow t\bar{t}Zj$	$6.027(9) \cdot 10^{-1}$	$6.921(11) \cdot 10^{-1}$
$e^+e^- \rightarrow t\bar{t}Zjj$	$6.436(21) \cdot 10^{-2}$	$8.241(29) \cdot 10^{-2}$
$e^+e^- \rightarrow t\bar{t}W^\pm jj$	$2.387(8) \cdot 10^{-4}$	$3.716(10) \cdot 10^{-4}$
$e^+e^- \rightarrow t\bar{t}HZ$	$3.623(19) \cdot 10^{-2}$	$3.584(19) \cdot 10^{-2}$
$e^+e^- \rightarrow t\bar{t}ZZ$	$3.788(6) \cdot 10^{-2}$	$4.032(7) \cdot 10^{-2}$
$e^+e^- \rightarrow t\bar{t}HH$	$1.3650(15) \cdot 10^{-2}$	$1.2168(16) \cdot 10^{-2}$
$e^+e^- \rightarrow t\bar{t}W^+W^-$	$1.3672(21) \cdot 10^{-1}$	$1.5385(22) \cdot 10^{-1}$

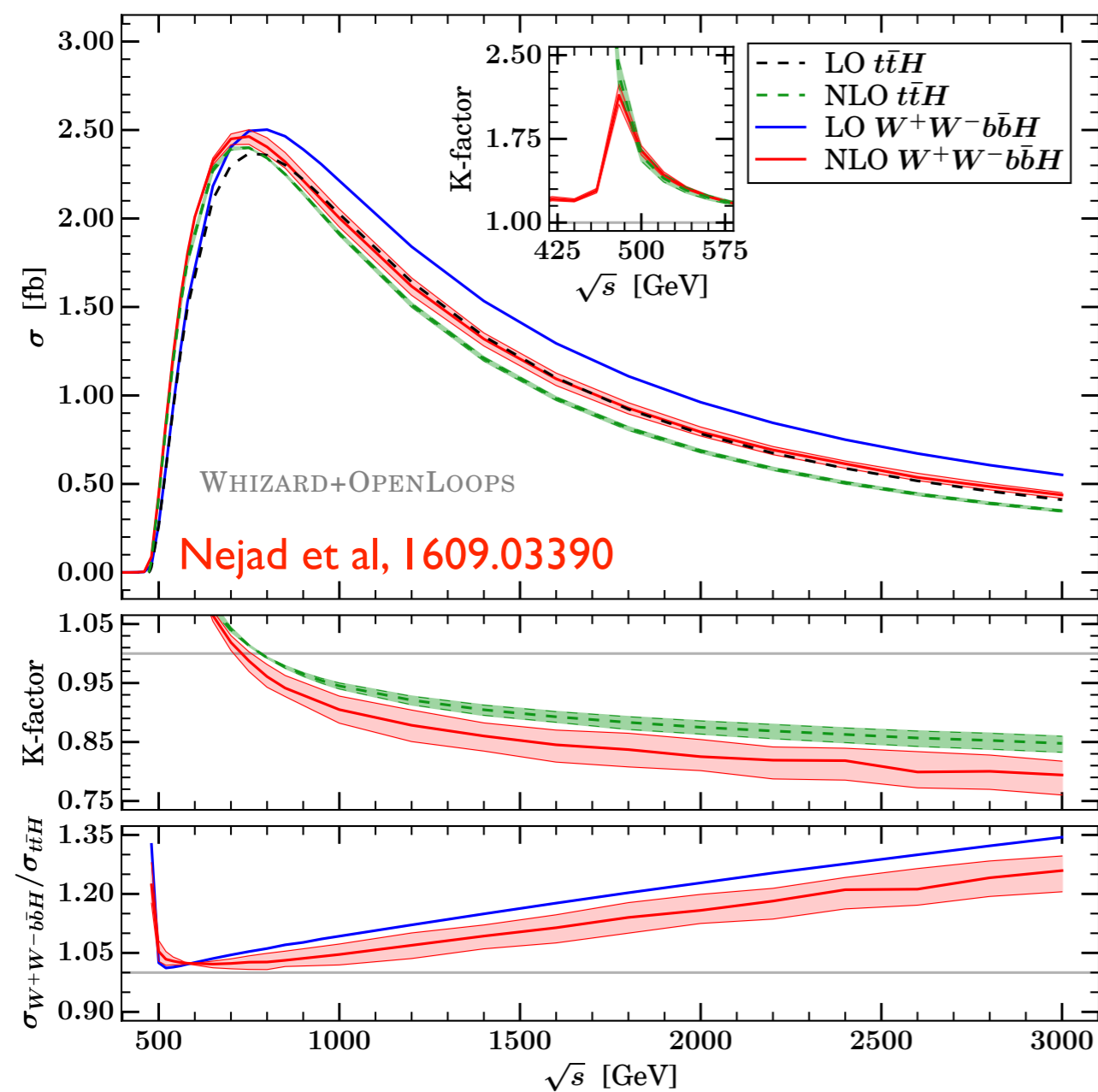
$\mu\mu$ @ 3 TeV, NLO EW

$\mu^+\mu^- \rightarrow X, \sqrt{s} =$	$\sigma_{\text{LO}}^{\text{incl}}$ [fb]	$\sigma_{\text{NLO}}^{\text{incl}}$ [fb]	δ_{EW} [%]
W^+W^-	$4.6591(2) \cdot 10^2$	$4.847(7) \cdot 10^2$	+4.0(2)
ZZ	$2.5988(1) \cdot 10^1$	$2.656(2) \cdot 10^1$	+2.19(6)
HZ	$1.3719(1) \cdot 10^0$	$1.3512(5) \cdot 10^0$	-1.51(4)
HH	$1.60216(7) \cdot 10^{-7}$	$5.66(1) \cdot 10^{-7} *$	
<hr/>			
W^+W^-Z	$3.330(2) \cdot 10^1$	$2.568(8) \cdot 10^1$	-22.9(2)
W^+W^-H	$1.1253(5) \cdot 10^0$	$0.895(2) \cdot 10^0$	-20.5(2)
ZZZ	$3.598(2) \cdot 10^{-1}$	$2.68(1) \cdot 10^{-1}$	-25.5(3)
HZZ	$8.199(4) \cdot 10^{-2}$	$6.60(3) \cdot 10^{-2}$	-19.6(3)
HHZ	$3.277(1) \cdot 10^{-2}$	$2.451(5) \cdot 10^{-2}$	-25.2(1)
HHH	$2.9699(6) \cdot 10^{-8}$	$0.86(7) \cdot 10^{-8} *$	
<hr/>			
$W^+W^-W^+W^-$	$1.484(1) \cdot 10^0$	$0.993(6) \cdot 10^0$	-33.1(4)
W^+W^-ZZ	$1.209(1) \cdot 10^0$	$0.699(7) \cdot 10^0$	-42.2(6)
W^+W^-HZ	$8.754(8) \cdot 10^{-2}$	$6.05(4) \cdot 10^{-2}$	-30.9(5)
W^+W^-HH	$1.058(1) \cdot 10^{-2}$	$0.655(5) \cdot 10^{-2}$	-38.1(4)
$ZZZZ$	$3.114(2) \cdot 10^{-3}$	$1.799(7) \cdot 10^{-3}$	-42.2(2)
$HZZZ$	$2.693(2) \cdot 10^{-3}$	$1.766(6) \cdot 10^{-3}$	-34.4(2)
$HHZZ$	$9.828(7) \cdot 10^{-4}$	$6.24(2) \cdot 10^{-4}$	-36.5(2)
$HHHZ$	$1.568(1) \cdot 10^{-4}$	$1.165(4) \cdot 10^{-4}$	-25.7(2)

Results with Whizard

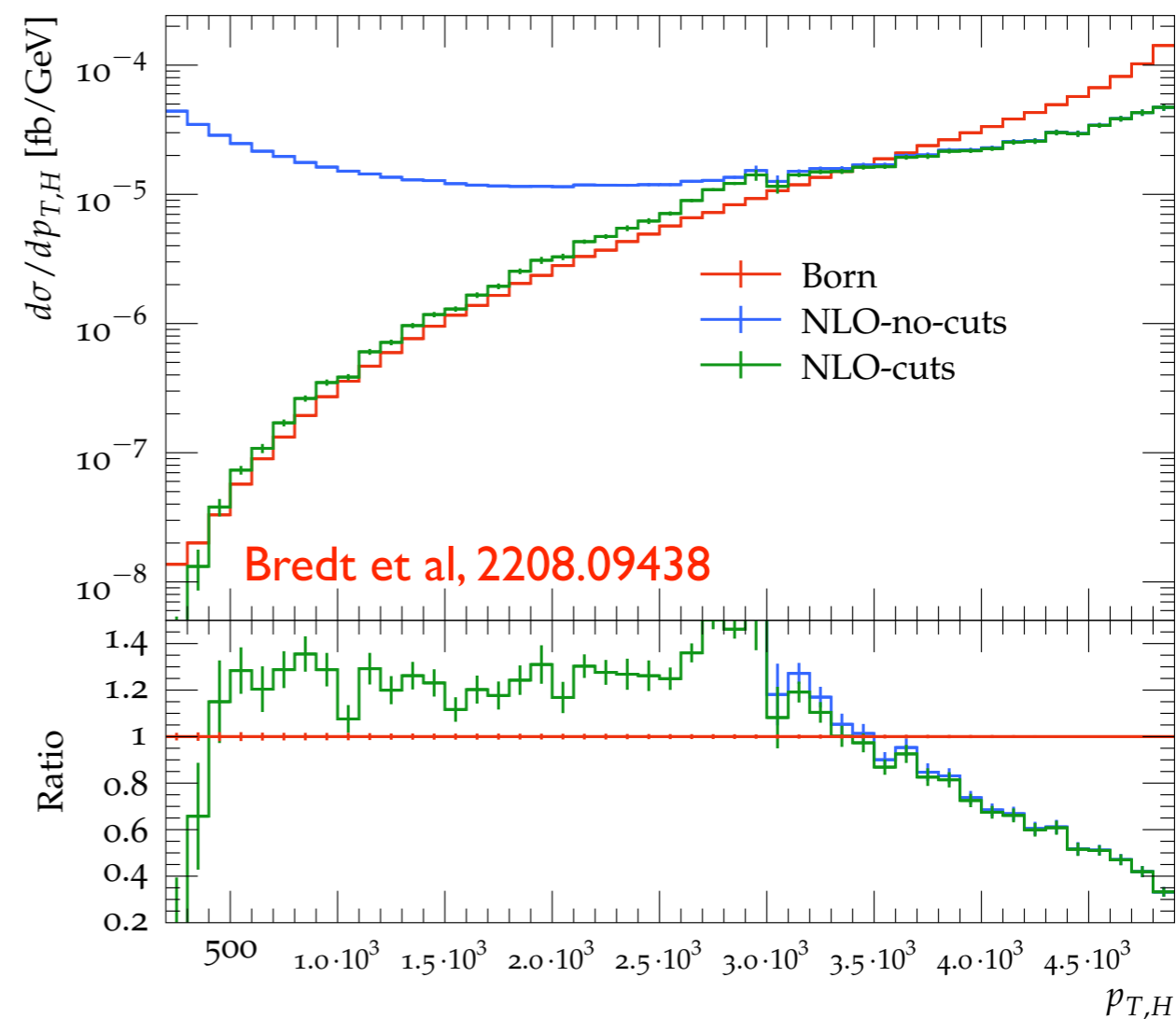
$ee \rightarrow ttH$ offshell at NLO QCD

$$e^+e^- \rightarrow t\bar{t}H \text{ and } e^+e^- \rightarrow W^+W^-b\bar{b}H$$



$\mu\mu \rightarrow ZH$ at NLO EW

$$\sqrt{s} = 10 \text{ TeV}$$





MG5_aMC at lepton colliders

- NLO EW and QCD corrections for (almost) all processes
- Beam polarisation for LO simulations
- Via eMELA, ISR (+beamstrahlung) in different ren/fact schemes
- The code automatically takes care to add to the short-distance xsection those terms necessary for consistency
 - Factorisation-scheme kernels included in the cross-section for Δ scheme and LL PDFs
 - Virtuals are corrected in order to account for different ren. scheme in model and PDFs ($\alpha(m_Z) \rightarrow \overline{MS}$)
- For details and how-to, see <https://answers.launchpad.net/mg5amcnlo/+faq/3324>

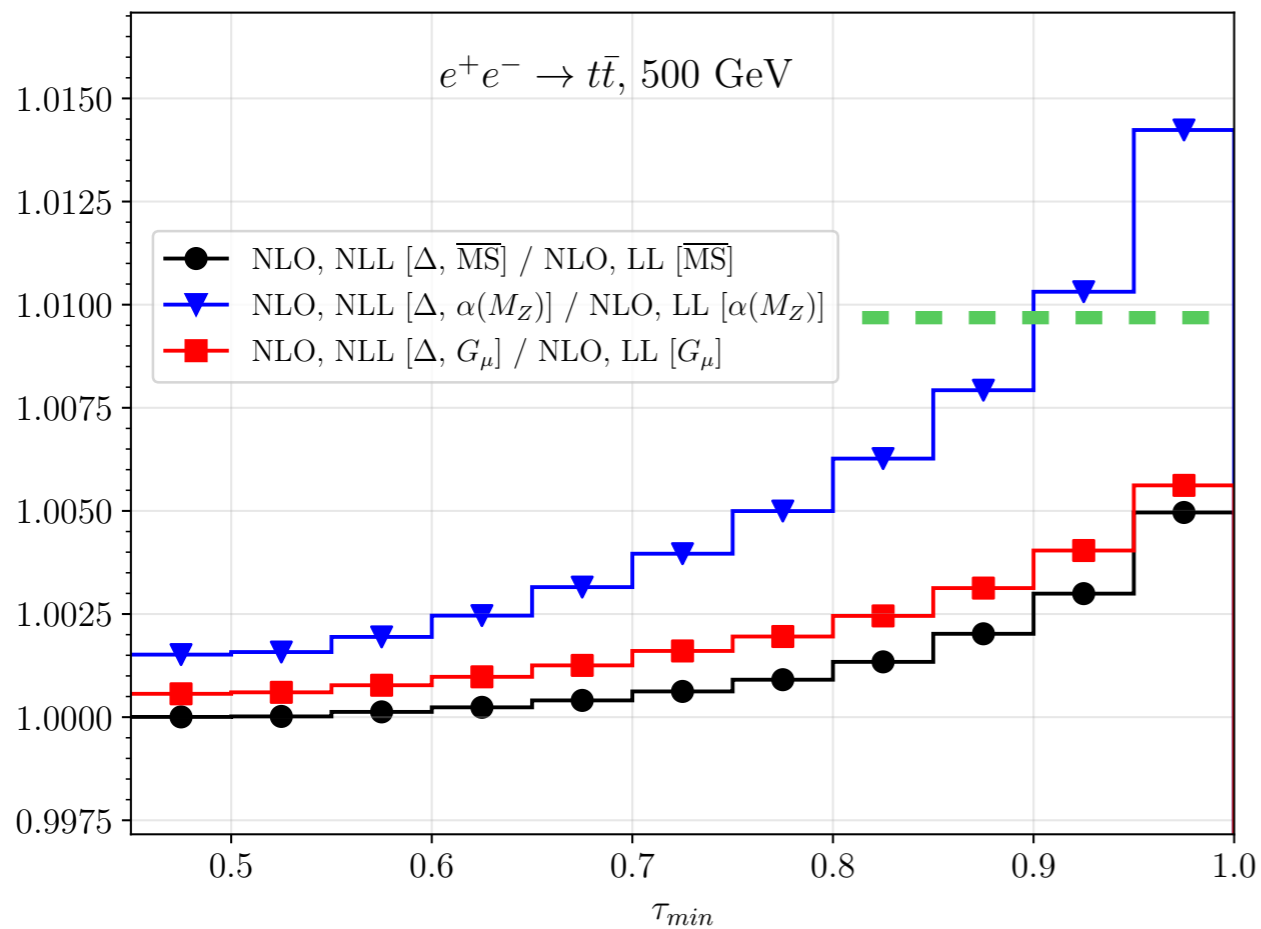


Some results:

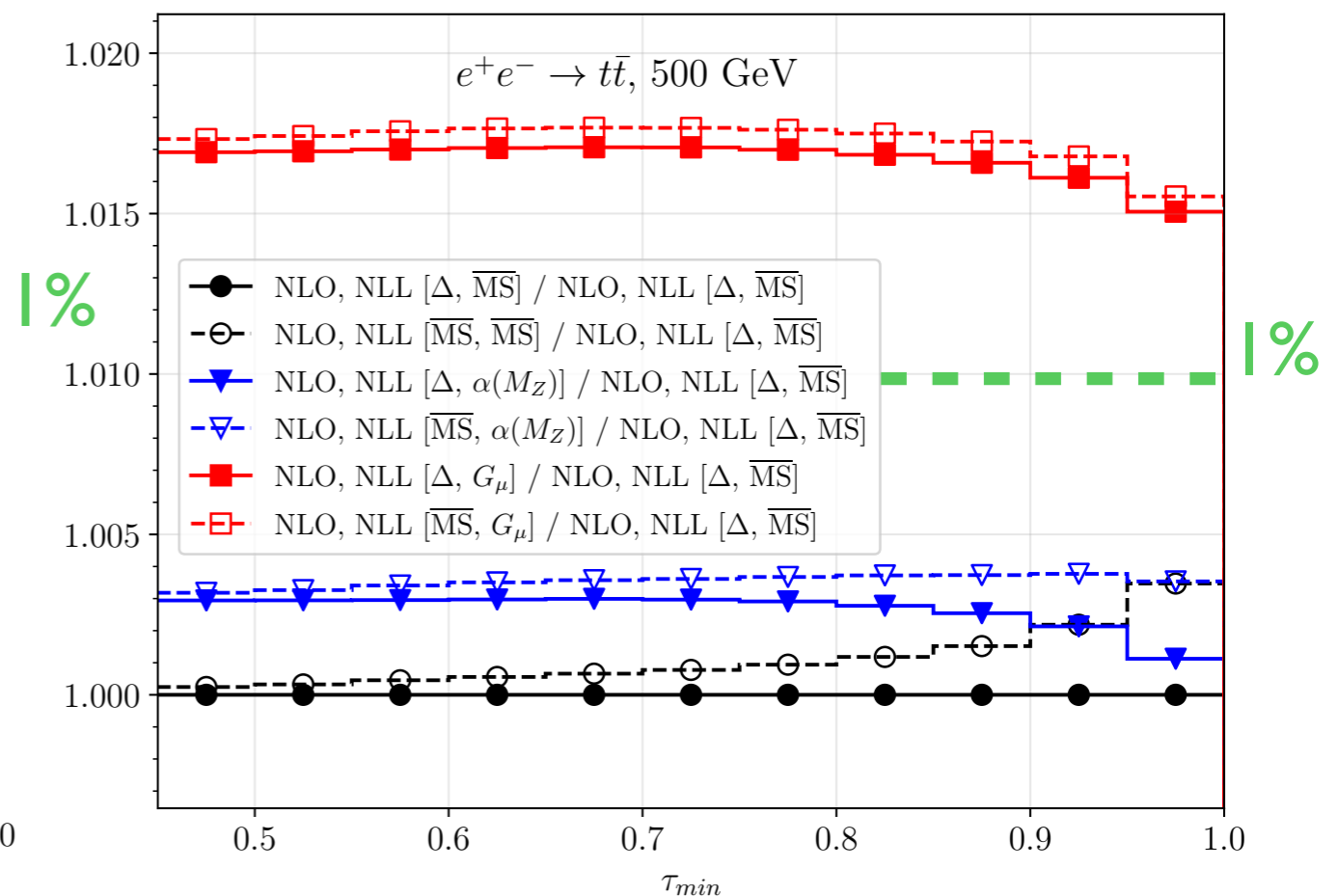
Bertone, Cacciari, Frixione, Stagnitto, MZ, Zhao [2207.03265](#)

- $e^+e^- \rightarrow t\bar{t}$ @500 GeV; Observable: $\sigma(\tau_{min}) = \int d\sigma \Theta\left(\tau_{min} \leq \frac{M_{p\bar{p}}^2}{s}\right)$

LL vs NLL



NLL different fact. schemes



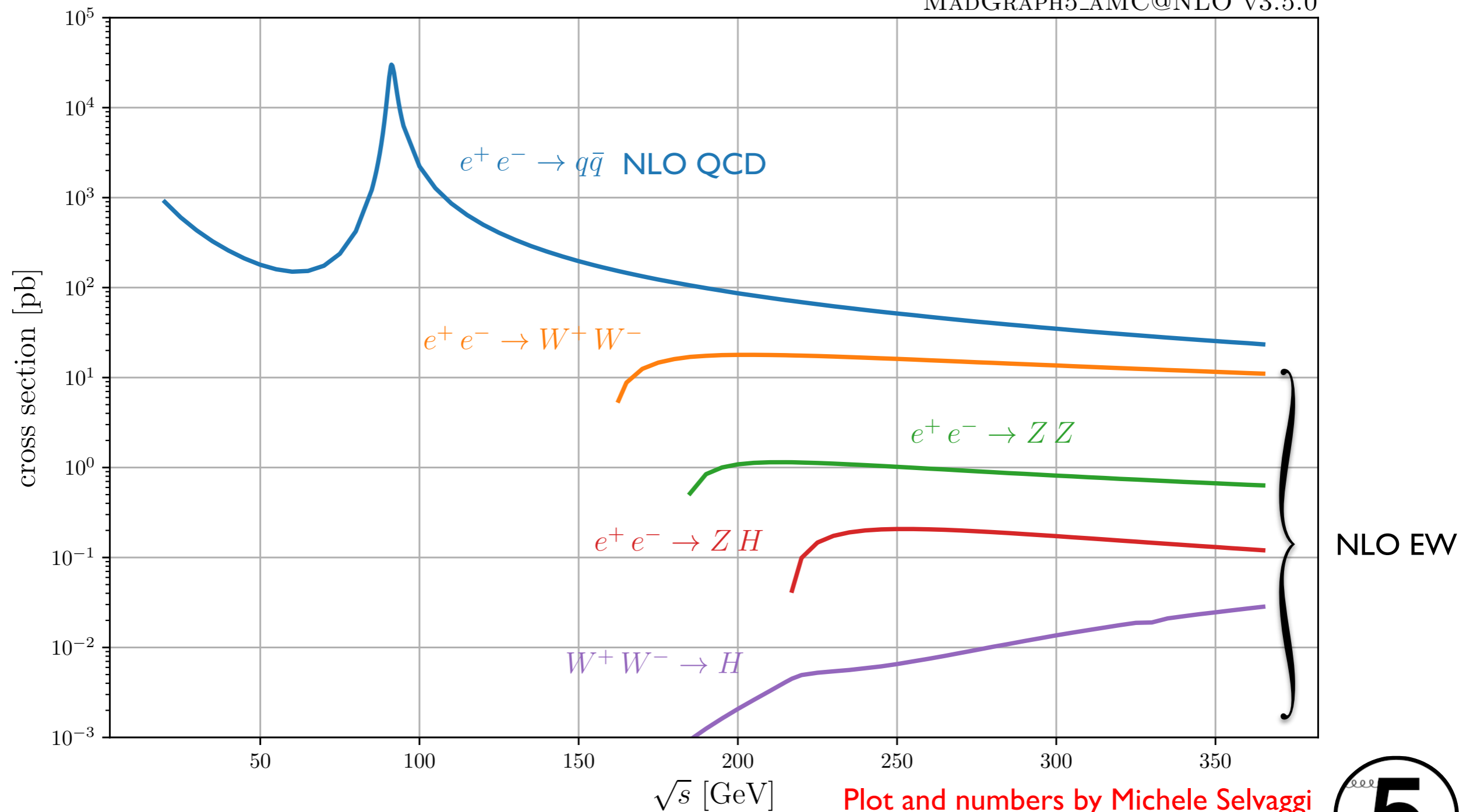
- NLL-MSbar vs NLL- Δ is (at most) at the few-per-mil level





More results:

MADGRAPH5_AMC@NLO v3.5.0

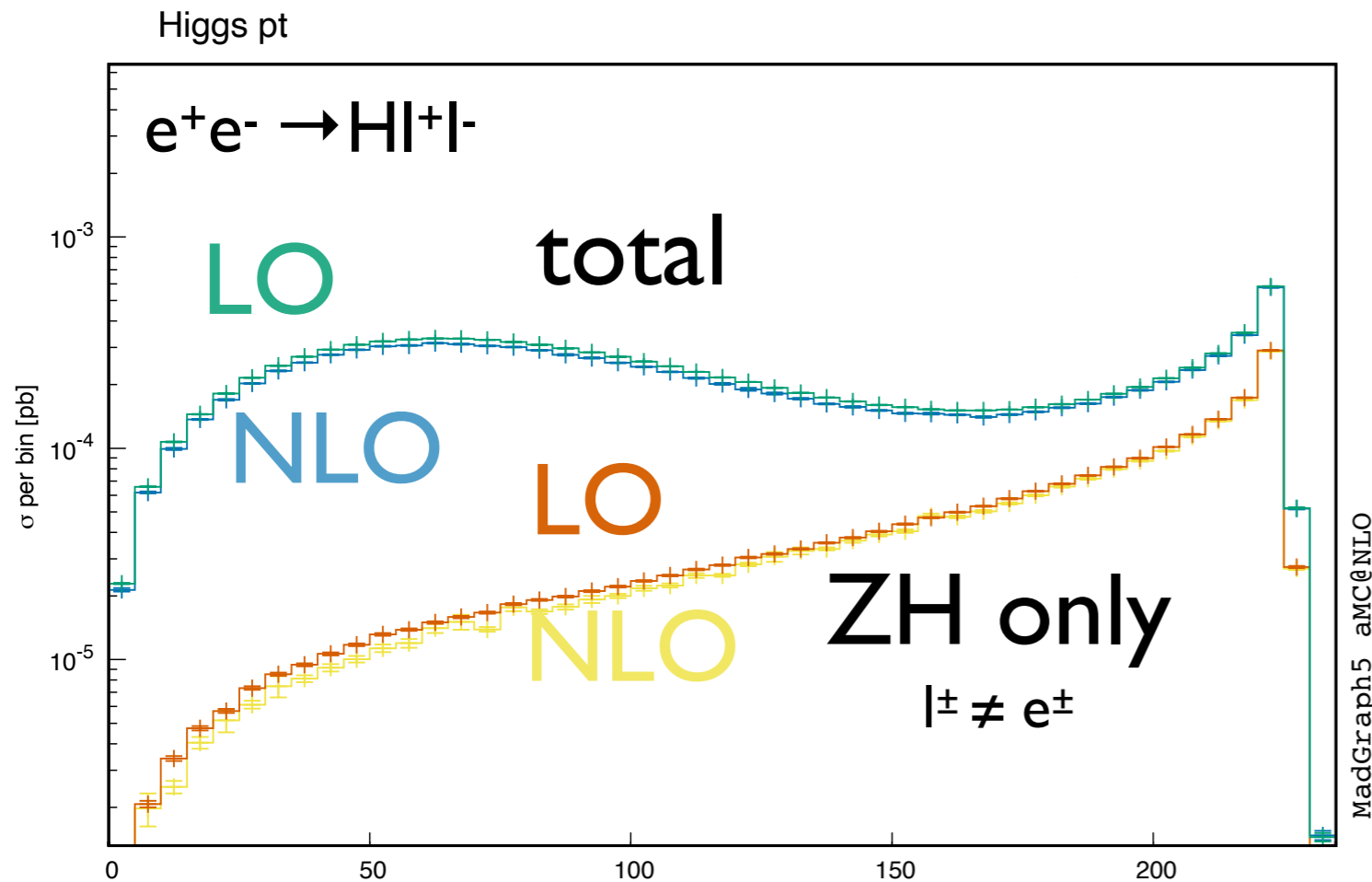




Complex-mass scheme:



$$e^+e^- \rightarrow H|^{+}|^{-}$$



- Qualitatively similar results to Denner, Dittmaier, Roth, Weber, [hep-ph/0302198](https://arxiv.org/abs/hep-ph/0302198)
- Results obtained in 15mins (on a cluster) @ 0.1%

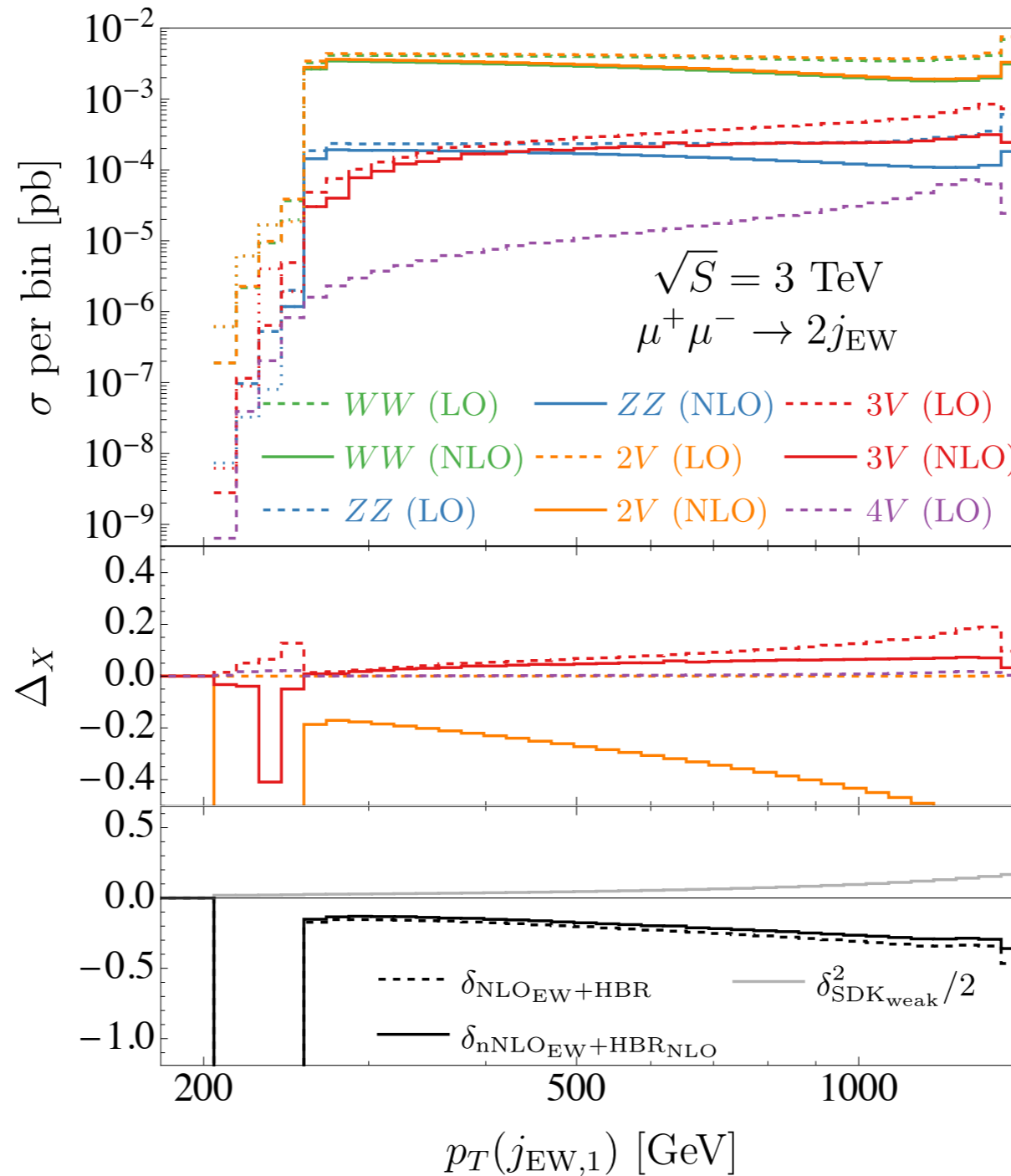
Inclusive timing profile :

Overall slowest channel	0:06:15
Average channel running time	0:03:42
Aggregated total running time	8:05:57



EW corrections at muon colliders

Pagani, Ma, MZ, 2409.09129, see also Davide's talk

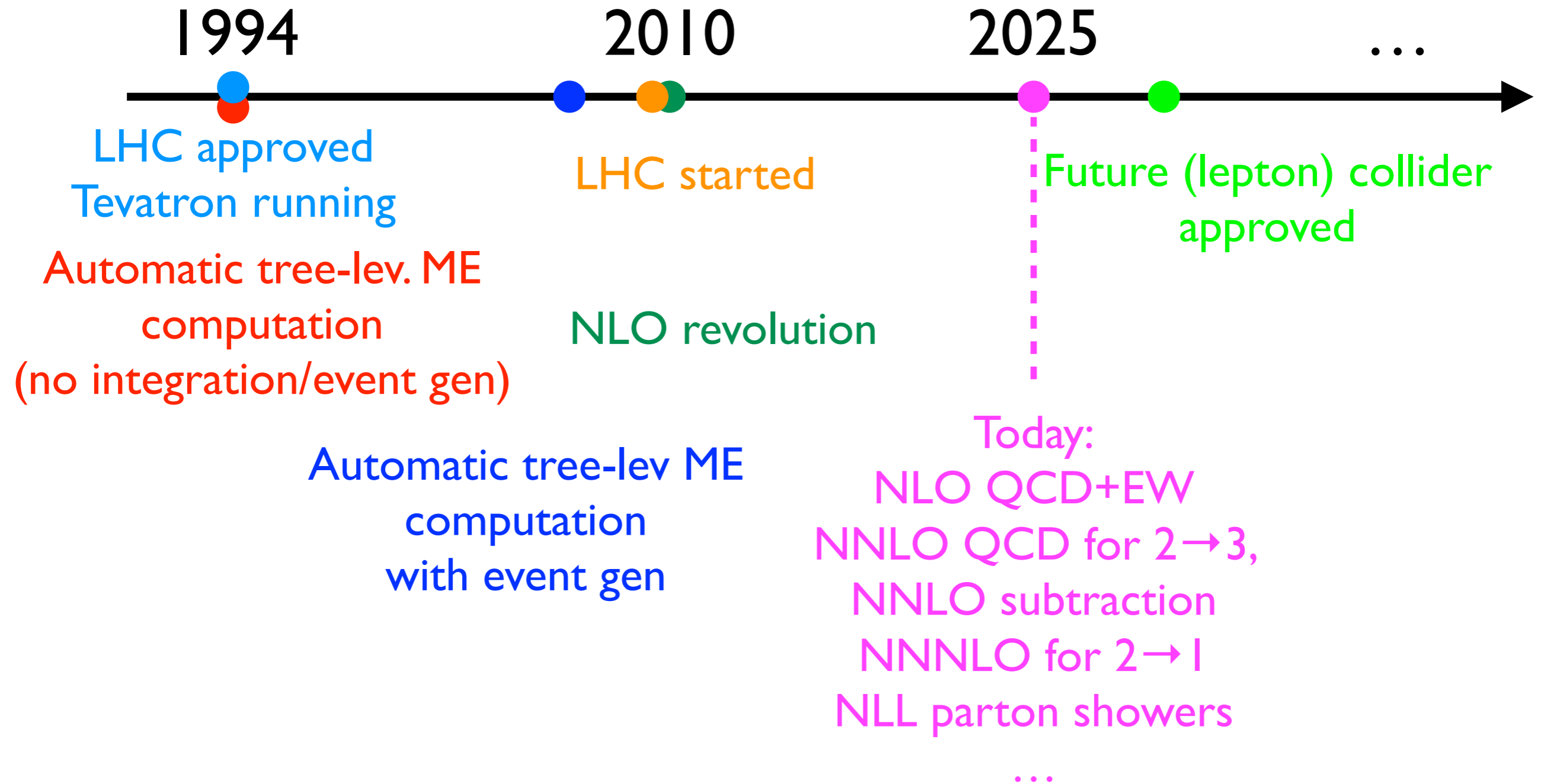


Investigate several aspect of EW corrections, real radiation, etc...





The food drives your appetite





Outlook

- Not covered in this talk, but still relevant: KKMC-ee for $f\bar{f}$ production, process specific tools, e.g. BabaYaga
- The LHC has taught us many lessons:
 - General-purpose MC tools are fundamental workhorses for all physics analyses
 - The possibility of using different tools is key for validation and proper estimate of uncertainties
- Progress in MCs is driven by the destination: we are in a much better shape than 30 years ago, when LHC started



Thank you!



Matching with PS: ISR at NLL?

QED Parton Shower

see for instance review in 0912.0749

It allows for exclusive photon emission in the context of collinear factorisation.

Photon energies dictated by distribution in z , whereas angles are generated independently according to the YFS formula, valid in the soft limit:

$$\cos \theta_l \propto - \sum_{i,j=1}^N \eta_i \eta_j \frac{1 - \beta_i \beta_j \cos \theta_{ij}}{(1 - \beta_i \cos \theta_{il})(1 - \beta_j \cos \theta_{jl})}$$

with η_i a charge factor and β_i the speed of the emitting particle.

Algorithm adopted in BabaYaga [$e^+e^- \rightarrow e^+e^-$, $e^+e^- \rightarrow \mu^+\mu^-$, $e^+e^- \rightarrow \gamma\gamma$]

hep-ph/0003268, hep-ph/0103117, hep-ph/0312014, hep-ph/0801.3360, hep-ph/0607181

Balossini, Bignamini, Carloni Calame, Lunardini, Montagna, Nicrosini, Piccinini

BabaYaga also includes a matching to NLO QED in the short distance cross section

slide by G. Stagnitto

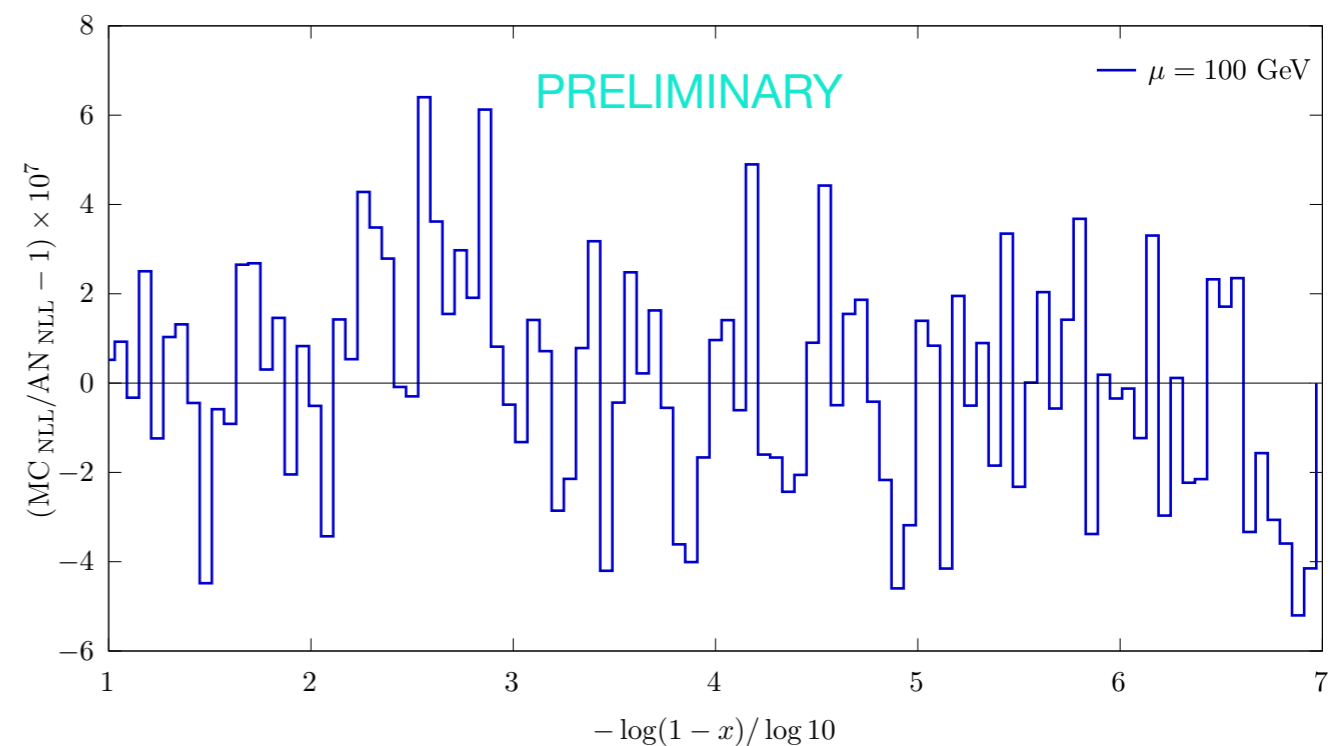
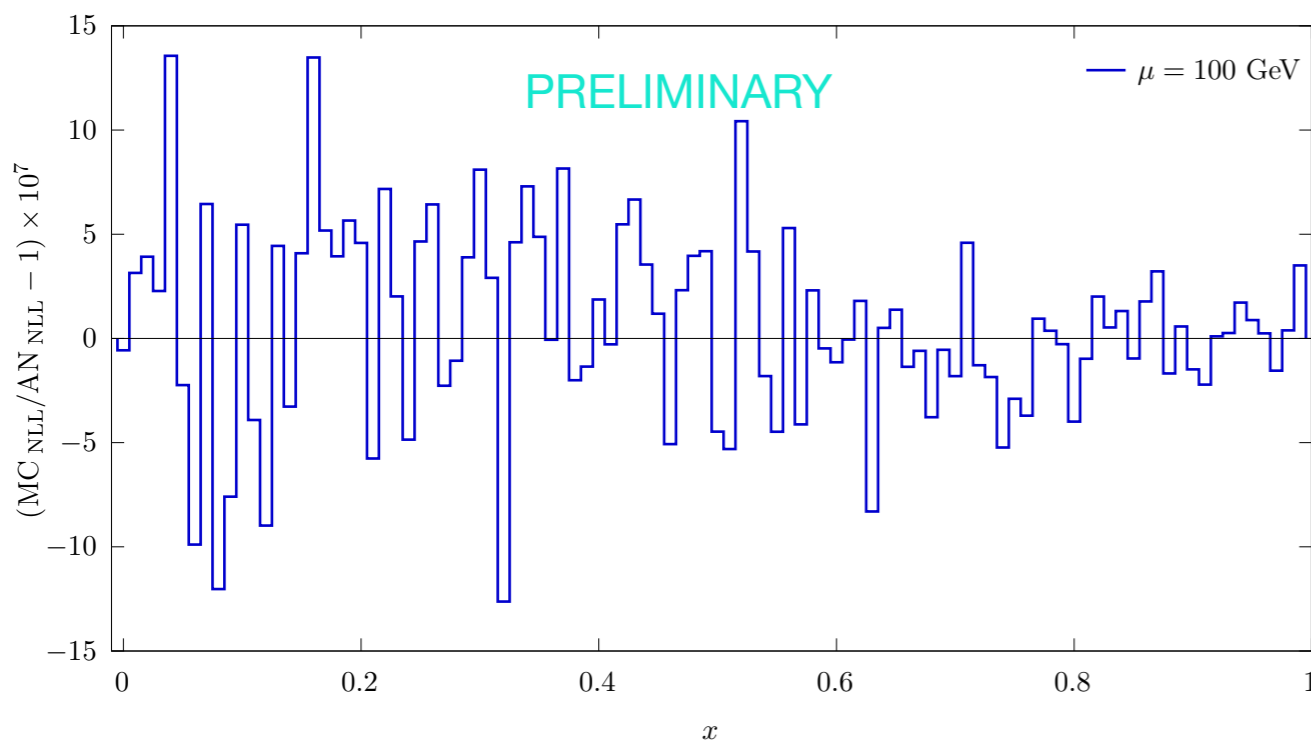
Matching with PS: ISR at NLL?

Towards a “NLL” QED Parton Shower

C. M. Carloni Calame, M. Chiesa, S. Frixione, G. Montagna, F. Piccinini, GS

$$D(x, s) = \sum_{n=0}^{\infty} \prod_{i=1}^n \left\{ \int_{m_e^2}^{s_{i-1}} \frac{ds_i}{s_i} \Pi(s_{i-1}, s_i) \frac{\alpha}{2\pi} \int_{x/(z_1 \cdots z_{i-1})}^{x_+} \frac{dz_i}{z_i} P(z_i) \right\} \Pi(s_n, m_e^2) D\left(\frac{x}{z_1 \cdots z_n}, m_e^2\right)$$

With a NLL iterative solution, we recover the known (non-singlet) NLL PDFs



WIP towards exclusive kinematics of final-state photons and singlet components
slide by G. Stagnitto