

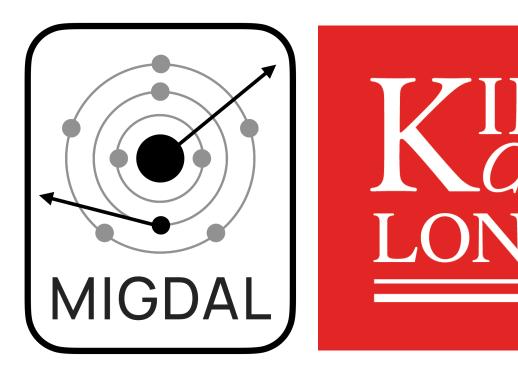
The Migdal effect (with neutral projectiles)

Christopher M^c**Cabe**

With Peter Cox, Matthew Dolan and Harry Quiney (Univ. of Melbourne) and the MIGDAL Collaboration









LDMA 2025, 9 April 2025

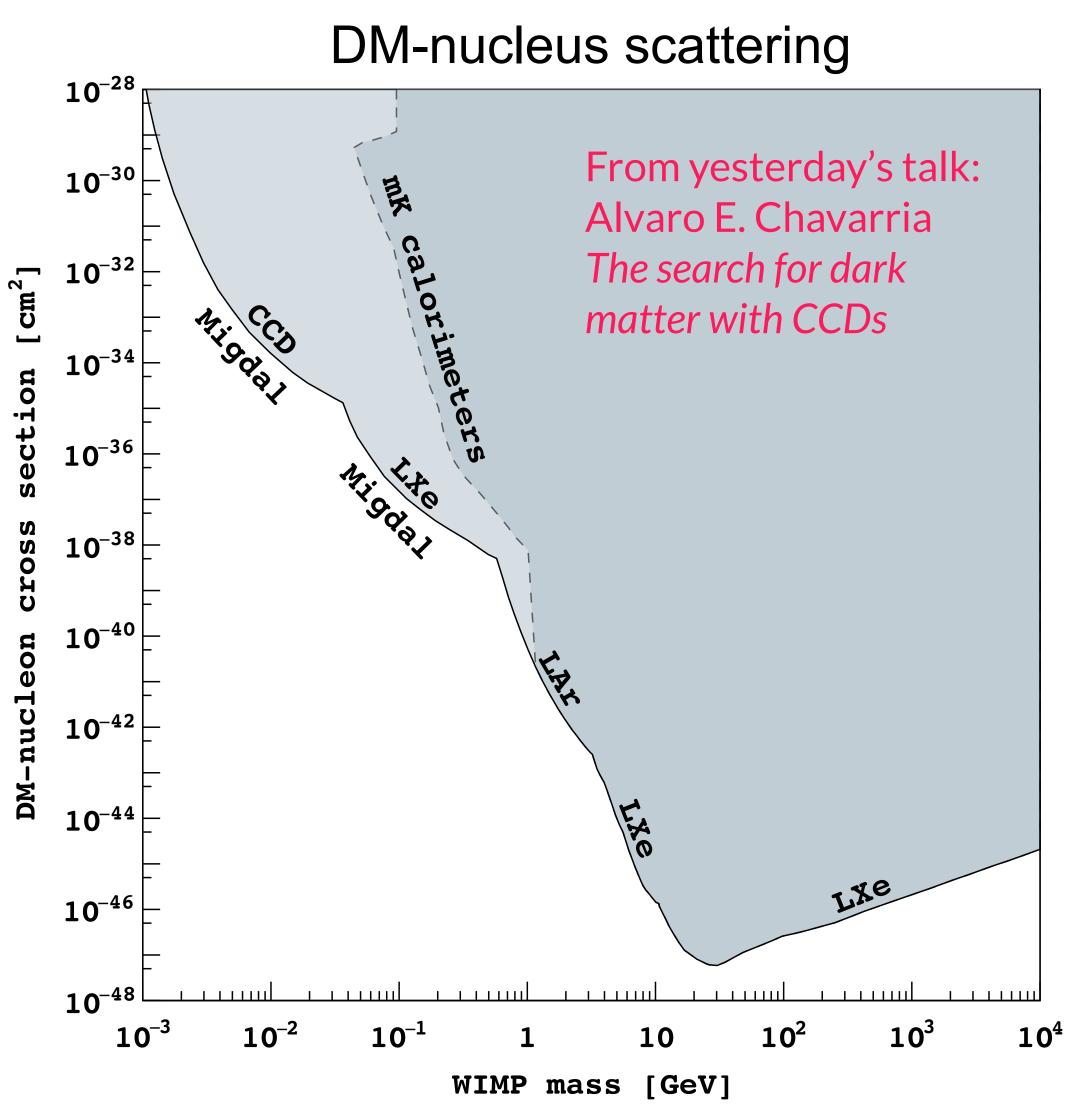






Motivation

Migdal searches now dominate searches for DM lighter than ~ 1 GeV





What is the Migdal effect?

Intuition: Neutral projectile scattering on helium

Neutral projectile (Dark matter or neutron) $v_e \sim \alpha c$

Helium atom

Fine-structure constant: $\alpha = 1/137$



Intuition: Neutral projectile scattering on helium

Neutral projectile (Dark matter or neutron) $v_e \sim \alpha c$

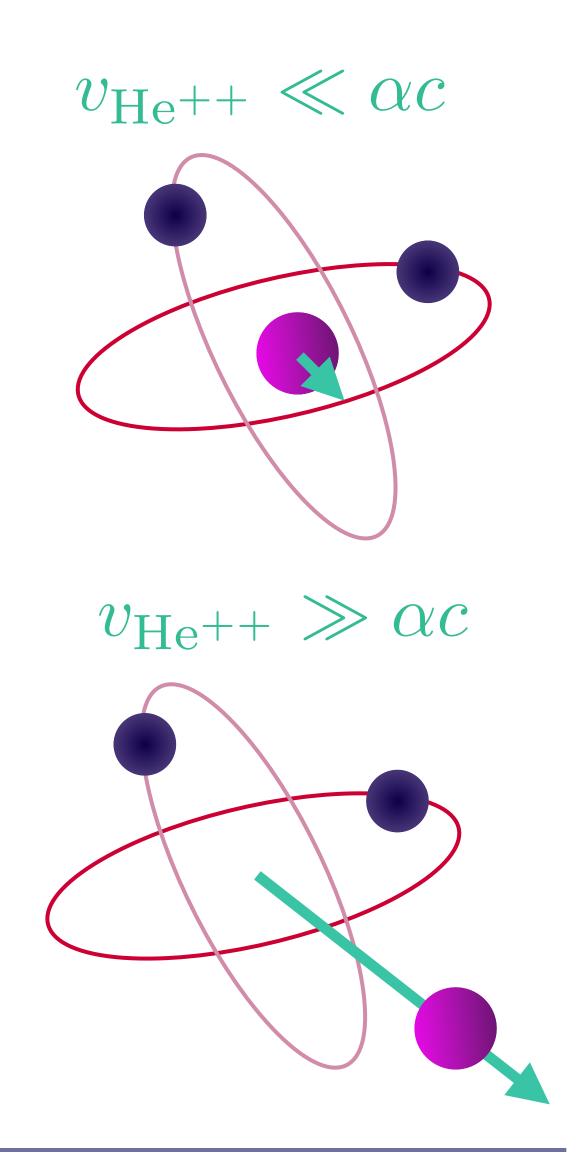
Helium atom

Fine-structure constant: $\alpha = 1/137$

Christopher McCabe

Low speed recoil:
 remain in ground state

2. High speed recoil:- double ionisation(electrons 'left behind')



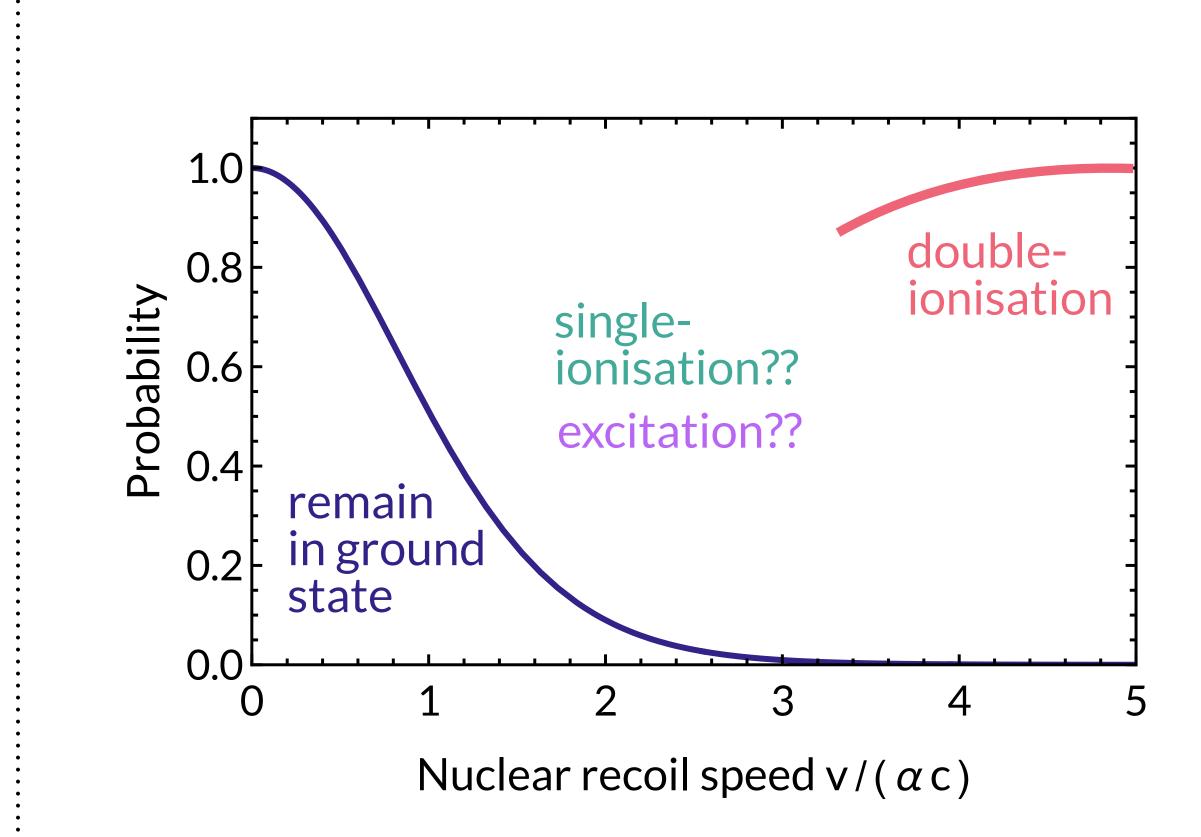
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[*In the rest of this talk c=1]

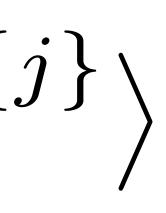


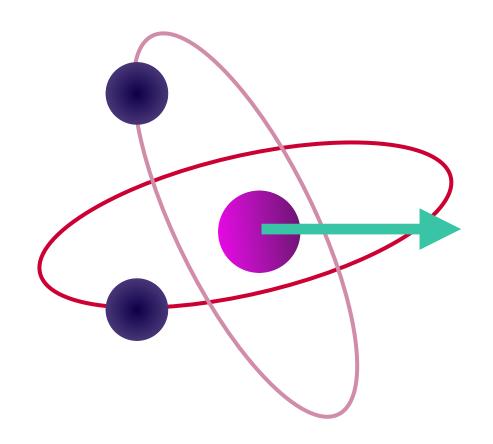
Migdal transition element

$\left\langle \Psi_{f}^{\{k\}} \left| e^{im_{e}\mathbf{v}\cdot\sum_{a}\mathbf{r}_{a}} \left| \Psi_{i}^{\{j\}} \right\rangle \right.$

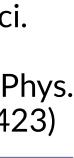
 $|\Psi_i^{\{j\}}\rangle$ describes the bound atomic-electrons wavefunction $e^{im_e \mathbf{v} \cdot \sum_a \mathbf{r}_a}$ accounts for Galilean boost $\mathbf{v} =$ Nuclear recoil velocity $|\Psi_{f}^{\{k\}}
angle$

A. Migdal, J. Phys. Acad. Sci. USSR 4 (1941) 449-453 (See also E. L. Feinberg, J. Phys. Acad. Sci. USSR 4 (1941) 423)





describes the final state wavefunction (excitation, ionisation, etc)



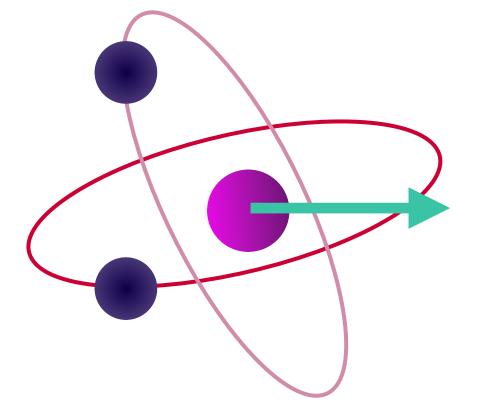
Migdal transition element

$$\langle \Psi_f^{\{k\}} | e^{im_e \mathbf{v} \cdot \sum_a \mathbf{r}_a} | \Psi_i^{\{m_e \mathbf{v} \cdot \sum_a \mathbf{r}_a | \Psi_i^{\{m_e \mathbf{v} \cdot \sum_a \mathbf{r}_a | \Psi_i^{\{m_e \mathbf{v} \cdot \sum_a \mathbf{v}_a \mathbf{v}_a | \Psi_i^{\{m_e \mathbf{v} \cdot \sum_a \|\Psi_i^{\{m_e \mathbf{v} \in \mathbb{v} | \Psi_i^{\{m_e \mathbf{v} \cdot \sum_a \|\Psi_i^{\{m_e \mathbf{v} \cdot \sum_a \|\Psi_i^{\{m_e \mathbf{v} \in \mathbb{v} | \Psi_i^{\{m_e \mathbf{v} \in \mathbb{v} \mid \Psi_i^{\{m_e \mathbf{v} \in \mathbb{v} | \Psi_i^$$

Previous calculations utilise the 'dipole approximation':

$$\exp\left(im_e \mathbf{v} \cdot \sum_{a=1}^{N} \mathbf{r}_a\right) \approx 1 + im_e \mathbf{v} \cdot \sum_{a=1}^{N} \mathbf{r}_a$$

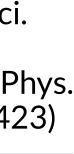
A. Migdal, J. Phys. Acad. Sci. USSR 4 (1941) 449-453 (See also E. L. Feinberg, J. Phys. Acad. Sci. USSR 4 (1941) 423)



Dipole approximation good for: (i) small v scattering processes and (ii) single ionisation processes

In our work, we keep the full exponential factor (sounds easy but lots of extra work!)

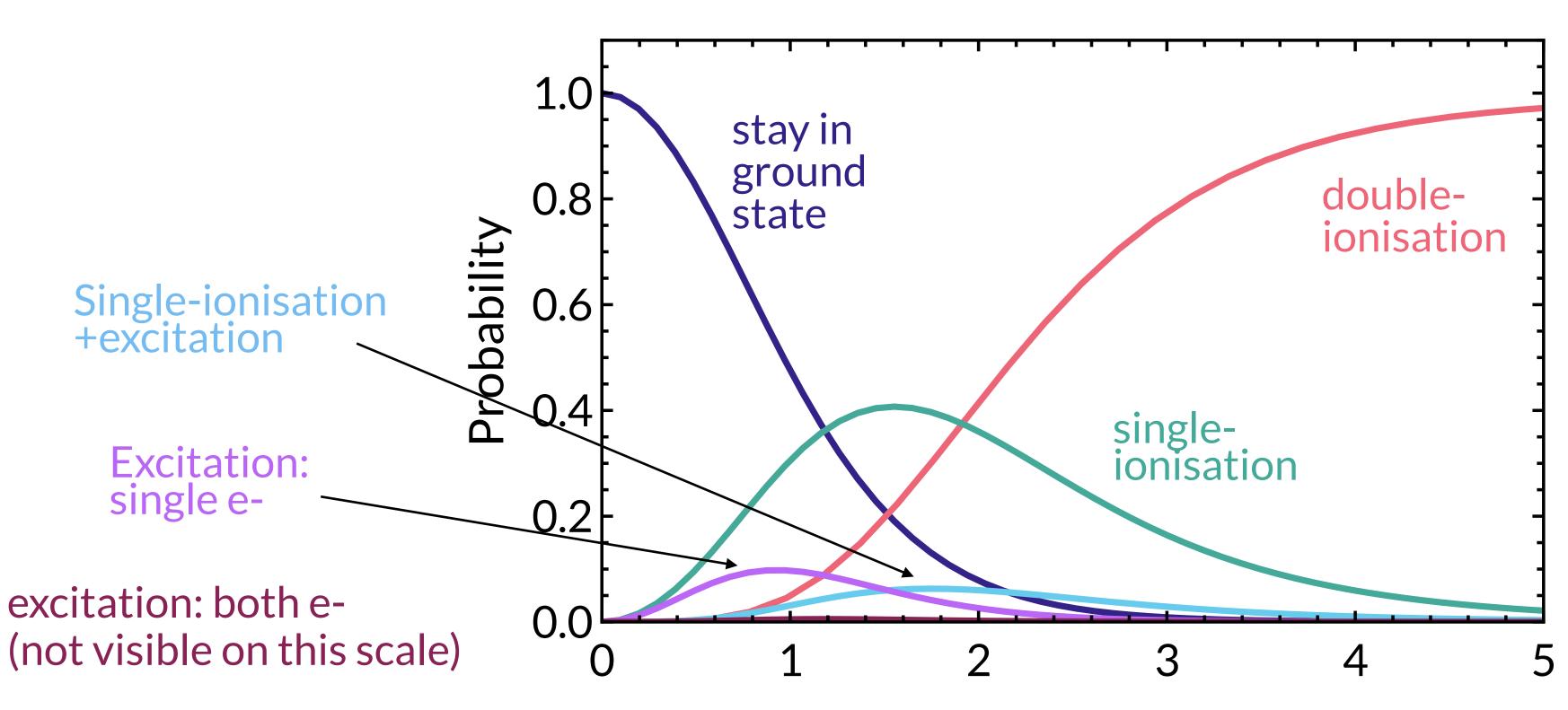
> Cox, Dolan, CM, Quiney, arXiv:2208.12222, PRD











Our results are valid for any v/α for He, C, F, Ne, Si, Ar, Ge, Kr, Xe

Intuition: Helium results

GRASP+RATIP Cox, Dolan, CM, Quiney, arXiv:2208.12222, PRD

Nuclear recoil speed v/(α c)

Previous calculations could only give the single-ionisation curve for $\,v/lpha\ll 1$



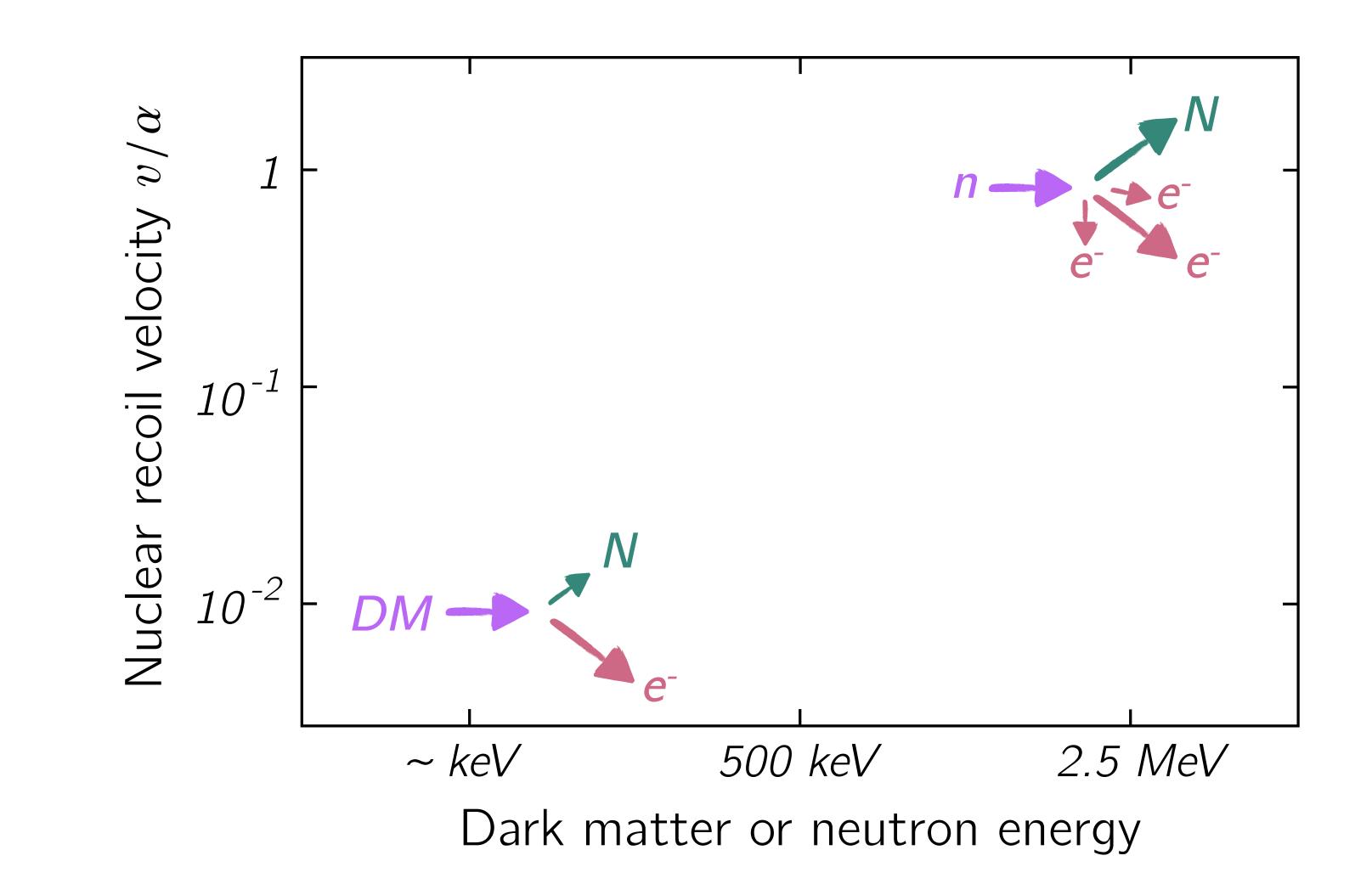


'Migdal effect' electrons and the nucleus are coupled in atomic systems: perturbation of the nucleus can induce electronic transitions

Transition probability depends on the speed of the recoiling nucleus

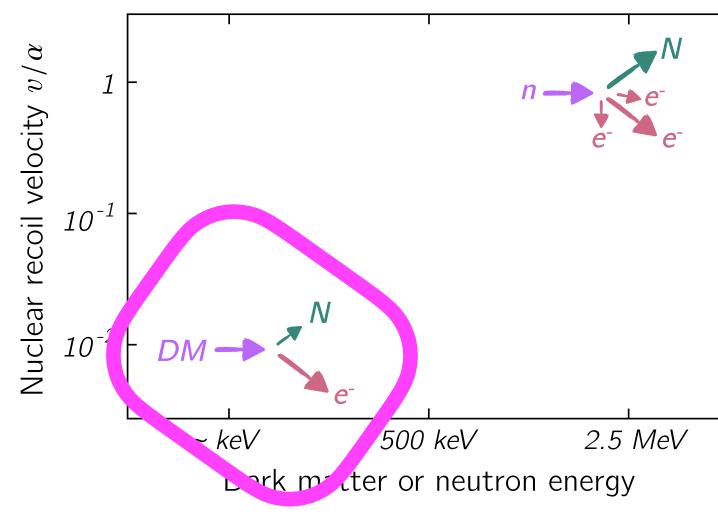


Migdal effect: regimes characterised by v/α



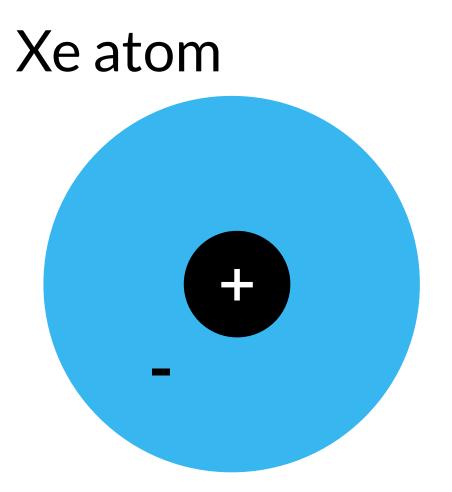


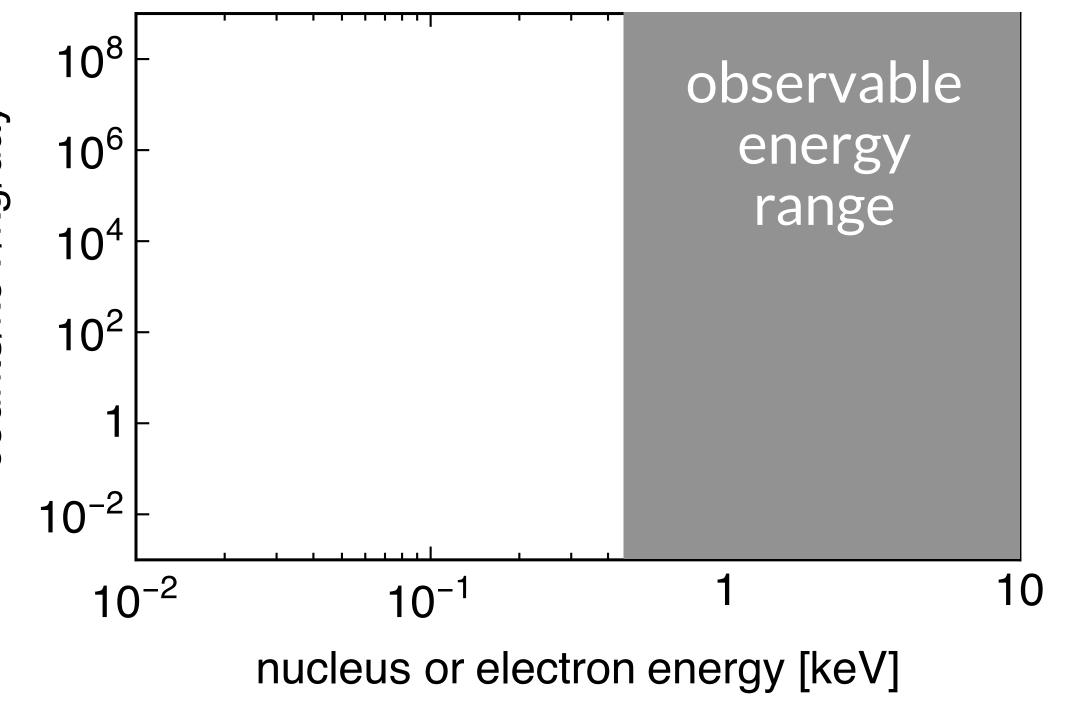
Small v regime: dark matter searches





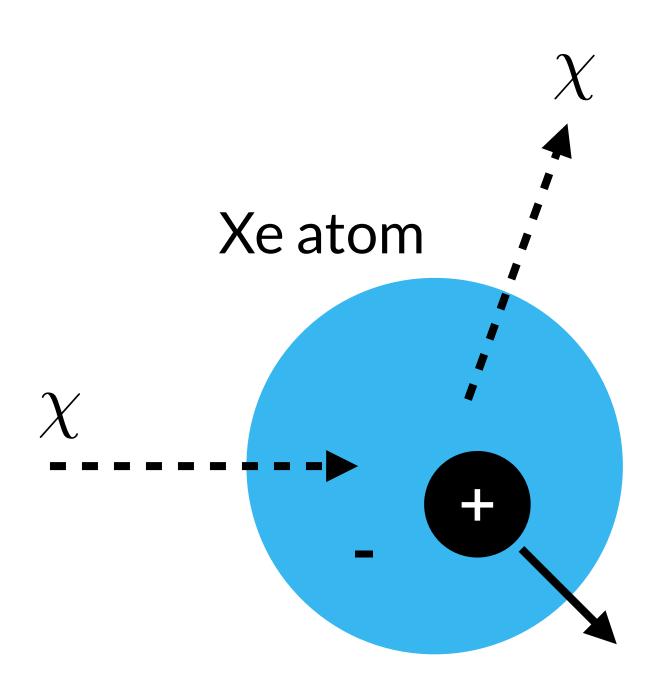
Benefits: consider DM scattering with xenon



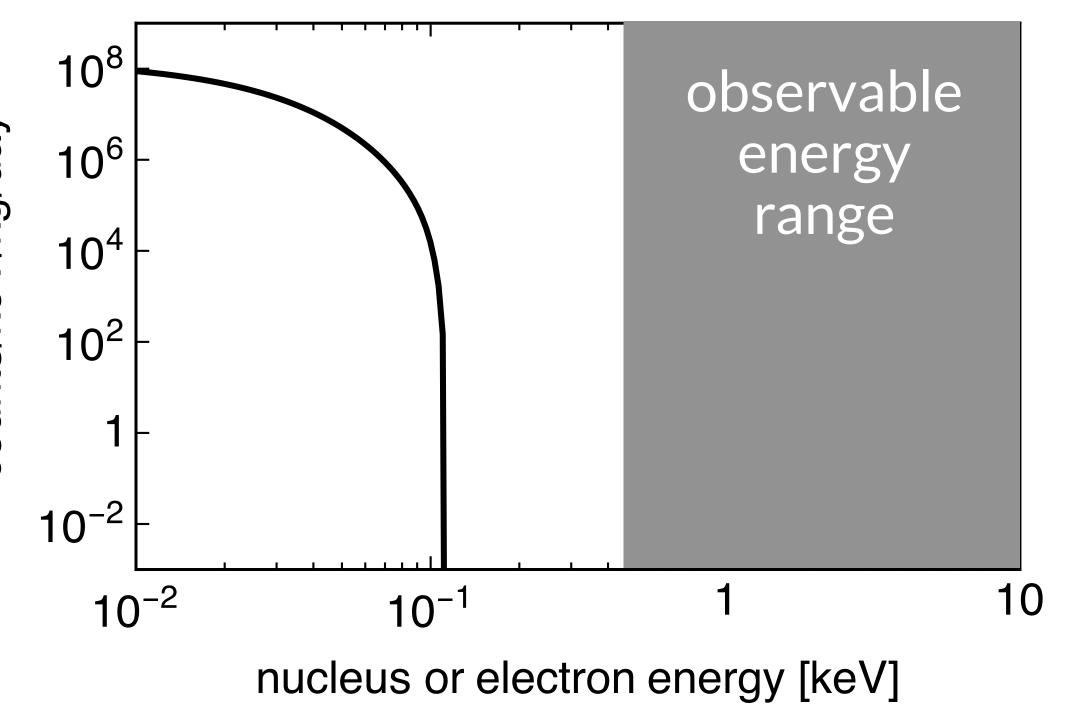




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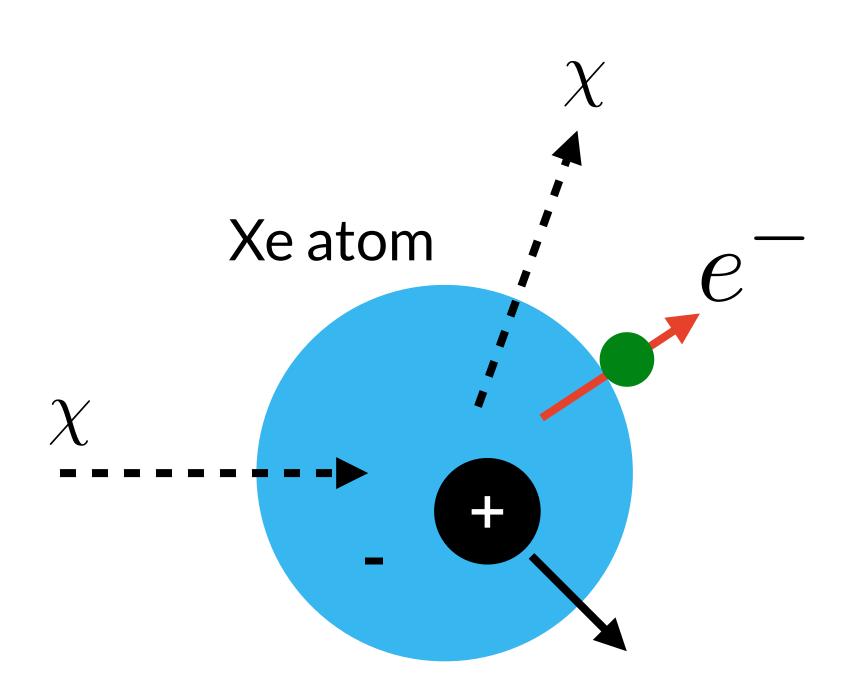


m_{DM} = 1 GeV 'Normal' nuclear scattering

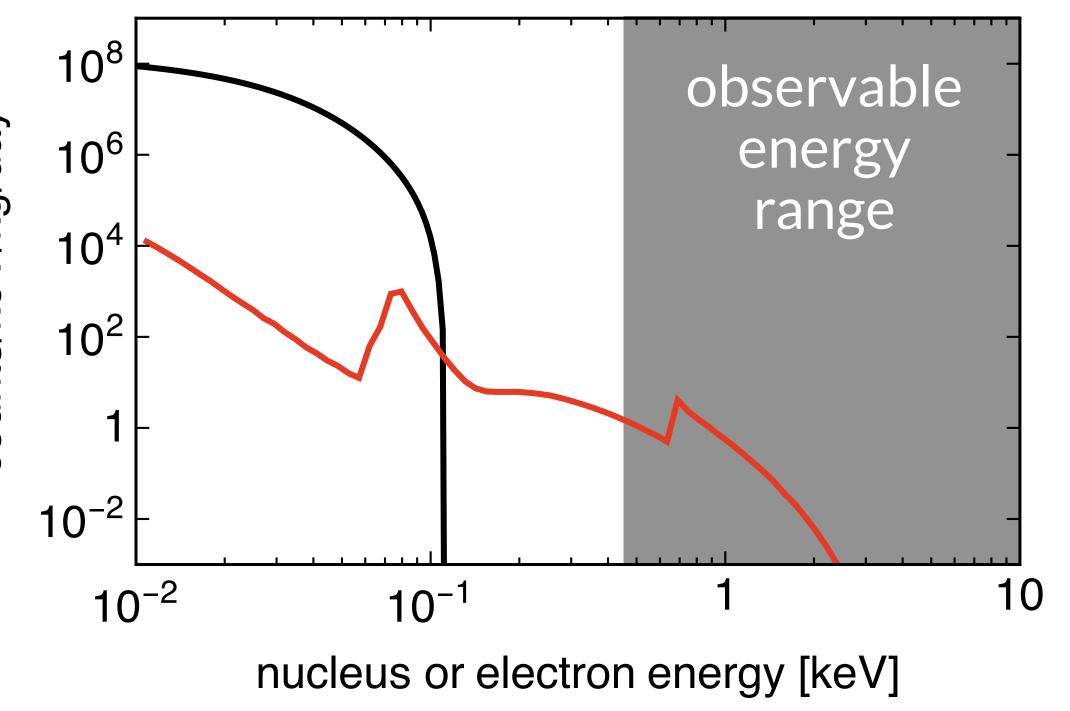




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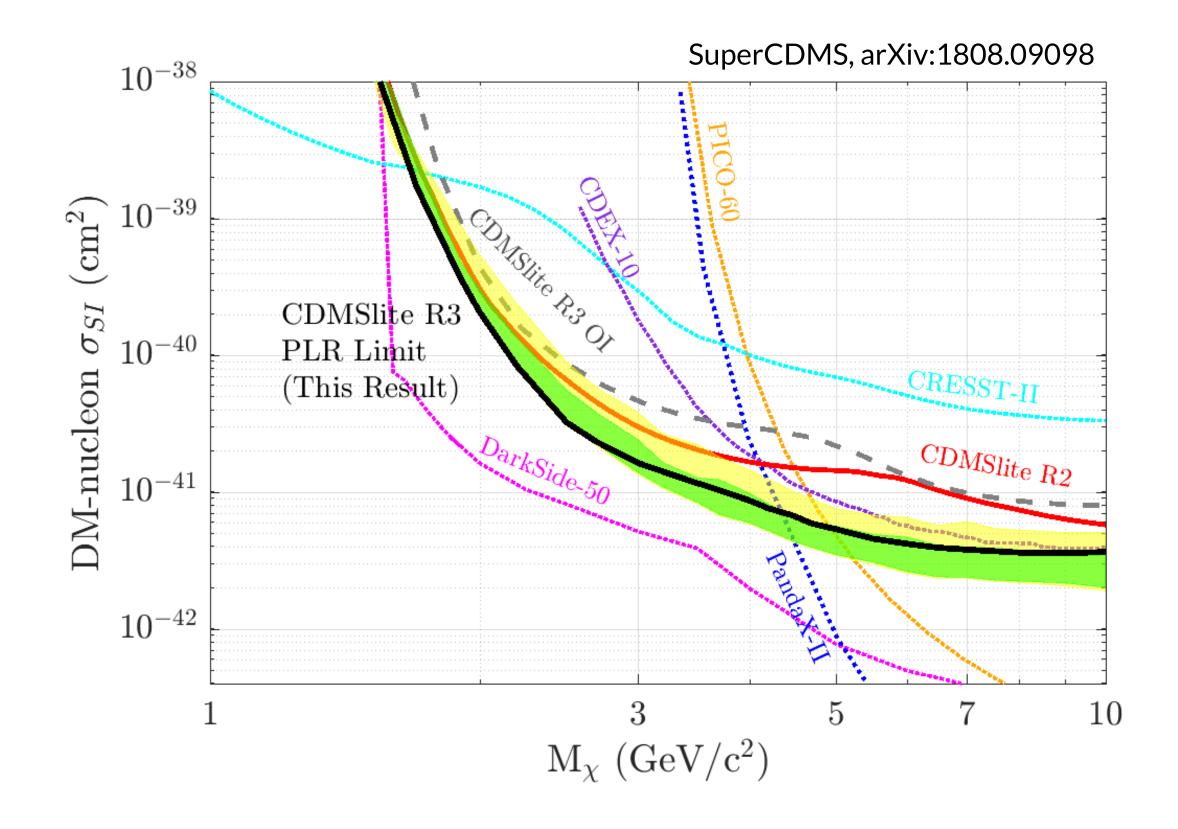


m_{DM} = 1 GeV 'Normal' nuclear scattering + Migdal effect (ionisation of 1 electron)

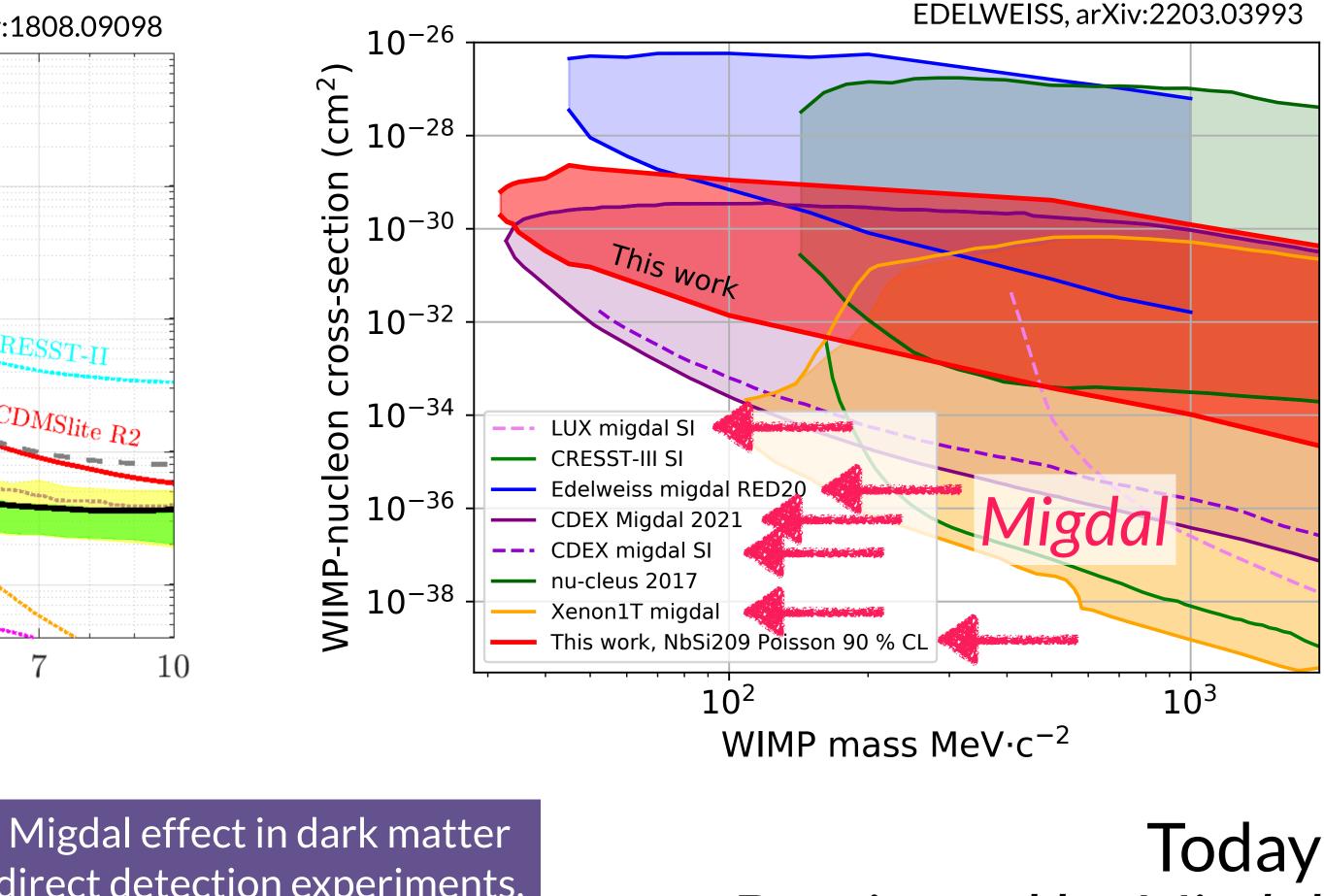




Sub-GeV searches increasingly dominated by Migdal



Pre-2018 No Migdal limits direct detection experiments, Ibe et al arXiv:1707.07258



Dominated by Migdal



Is there evidence for the Migdal effect?

Evidence? Yes, but...

Т. 9 Журнал экспериментальной и теоретической физики Вып. 10

A.B. Migdal's papers date bacl

Predicted effect in:

- 1. α , β decay
- 2. Neutral scattering

В работе дачей большой

T.9 M

При яд передачей бо Пон малых и нонизация вылетает из энергиях отд оболочках. При сте низация атомов при ядерных реакциях

A. Munza

она с электроном, крание порадка 10⁻²³ см², во втором — п много меньш

etne 1 + P. (r., r. ... Fr) , I' представляет собой Ф-функцик

движется со скоо

дается функцией Ф1 (r1, г2...г.). Так кал

 $W = \left| \int \overline{\Psi}_1 e^{\eta \cdot r_i} \Psi_0 dr_1 \dots dr_f \right|^2,$

ственным, пр зация, обусловленная магнитным и специфическим ядерным взаимодействием нейтрона с электроном, крайне мала — соответствующее сечение в первом случае порядка 10⁻²⁸ см², во втором — порядка 10⁻³⁶ см²).

Effect has been observed скорость такой нонизация может быть очень просто рассчитана. Так как интересси случай боющих энергий отдачи и, следовательно, больших скорость пасесць Средского отдачи и, следовательно, больших электронных периодов. Следовательно, измессние скорости ядра происходит

M.S. Rapaport, F. Asaro and I. Pearlman K-shell electron shake-off accompanying alpha decay, PRC 11, 1740-1745 (1975) M.S. Rapaport, F. Asaro and I. Pearlman L- and M-shell electron shake-off accompanying alpha decay, PRC 11, 1746-1754 (1975) C. Couratin et al., First Measurement of Pure Electron Shakeoff in the β Decay of Trapped 6He+Ions, PRL 108, 243201 (2012)

передаче энергии Р много меньше размеров электронных осолочек, то каро можно считать не сместившимся за время удара. Для получения вероятности возбуждения или нонизации нужно исходную Ф-функцию атома разложить по собственным функциям движущегося ядра. Можно поступить несколько иначе, а именно перейти к системе координат, в которой ядро покоится; тогда собственными функциями задачи будут обыч-

Effect has not been observere superious apa. Havanbuas pyukuus Ψ_0 при втом преобра-

Действительно, миожитель е на 1 ч представляет собой Ф-функцию центра инерции оболочки, который в старой системе координат покоился, а в новой движется со скоростью v, равной по величине и противоположной по направлению скорости ядра.

Пусть конечное состояние атома в рассматриваемой системе координат дается функцией Ф1 (r1, г2...г). Так как ядро за время удара не сместилось, то координаты влектронов в Ф, отсчитаны от той же точки, что и в Ф. Вероятность перехода в консчное состояние дается выражением:

 $W = \int \overline{\Psi}_1 e^{\eta \tau_1} \Psi_0 d\tau_1 \dots d\tau_f,$



JOURNAL of PHYSICS

IONIZATION OF ATOMS ACCOMPANYING a- and S-DECAY

By A. MIGDAL (Received November 15, 1940)

The probability of ionization of the inner electron shells accompanying a- and 8-decay culated. Also an estimation of the order of magnitude of ionization of the outer shells

1. Ionization accomanying β-decay

1. The probability of ionization of an ntom as a result of the β -decay can be without difficulty calculated if one makes use of the fact that the velocity of a β -electron is usually great as compared with velocities of atomic electrons. It is easily seen that in this case one can neglect the direct interaction of the β -decay electron with the atomic ones. The ionization is due to the fact that the neu-lear charge is changed within a time in-terval which is short comparing to atomic periods. $W \sim \frac{1}{h^2} \left(\frac{a}{a} + \frac{a}{y_c}\right) = \left(\frac{e}{h_c}\right)^2$ (the quantity $\gamma = E/mc^2$ disappears because the Lorentz contraction of the field is compensated by an increase of the latter. On the other hand, the probability of ionization by a suddens change of nuc-lear charge, as will be shown, is of the direct interaction to be small $\left(\frac{Z_{eff}e^2}{2}\right)^2 \ll 1$ ntom as a result of the β -decay can be

The following estimation shows that the direct interaction can be actually neglect-ed. The probability of an electron tran-sition due to the direct interaction is according to perturbation theory:

 $\int V_{01} e^{i\omega_{01}t} dt |^2$ W = ______ "on is here the matrix element of the peray electron traverses electron shells is lity is equal to the much smaller than the atomic periods.

Journal of Physics, Vol. IV. No. 5

 $W \sim \frac{V^{a} \tau^{2}}{\hbar^{2}} \sim \frac{1}{\hbar^{2}} \left(\frac{\gamma e^{2}}{a} \cdot \frac{a}{\gamma c} \right)^{2} = \left(\frac{e^{2}}{\hbar c} \right)$

 $\left(\frac{Z_{\text{eff}}e^{s}}{\hbar c}\right)^{2} \ll 1.$

The condition (2) has a simple meaning

in the case of a K-electron, becaus $(Ze^{z}/\hbar c)^{z} = (V_{h}/c)^{z}$. Therefore, the direct tivistic correction. The condition (2) is pproximately valid even for K-elec 2. One can calculate the probability of

ionization by means of a sudden change is here the matrix element of the per-phation energy; $\omega_{s1} = (E_1 - E_4)/\hbar$ —the fre-ency corresponding to the electron transi-til its of the order of atomic frequencies. The W-function of atomic electrons does not it its of atomic frequencies. tion; it is of the order of atomic frequencies. not change when the decay electrons doe The time interval τ within which the de-emitted. Therefore, the transition probabi

резко неадиабатически, так что Ф - функция электронов-не может измениться

(1)

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The following estimation shows that the direct interaction can be actually neglected. The probability of an electron tran-sition due to the direct interaction is according to perturbation theory:

$$W = \frac{\left| \int_{0}^{\infty} V_{02} e^{i\omega_{01}t} dt \right|^{2}}{\hbar^{2}} .$$
 (1)

5 Journal of Physics, Vol. IV. No. 5

Hence the transition probability is of the

$$W \sim \frac{\gamma^{a} \tau^{2}}{\hbar^{2}} \sim \frac{1}{\hbar^{2}} \left(\frac{\gamma e^{2}}{a} \cdot \frac{a}{\gamma c} \right)^{2} = \left(\frac{e^{2}}{\hbar c} \right)^{2}$$

On the other hand, the probability of

(the quantity $\gamma = E/mc^2$ disappears because the Lorentz contraction of the field is compensated by an increase of the latter. ioniziation by a «sudden» change of nuclear charge, as will be shown, is of the order of $1/Z_{eff}^{2}$. Hence the condition for the direct interaction to be small

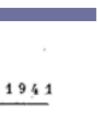
$$\left(\frac{Z_{\rm eff}e^{a}}{\hbar c}\right)^{2} \ll 1.$$

The condition (2) has a simple meaning in the case of a K-electron, because $(Ze^2/\hbar c)^2 = (V_{\hbar}/c)^2$. Therefore, the direct interaction is to be considered as a relativistic correction. The condition (2) is approximately valid even for K-electrons of uranium.

2. One can calculate the probability of ionization by means of a sudden change V_{o1} is here the matrix element of the per- of the nuclear charge in the following turbation energy; $\omega_{o1} = (E_1 - E_0)/\hbar$ the fre- manner. The above estimation shows that quency corresponding to the electron transi- the W-function of atomic electrons does tion; it is of the order of atomic frequencies. not change when the decay electron is The time interval τ within which the de- emitted. Therefore, the transition probabicay electron traverses electron shells is lity is equal to the square of the coeffimuch smaller than the atomic periods. cient of expansion of the Y-function cor-

order

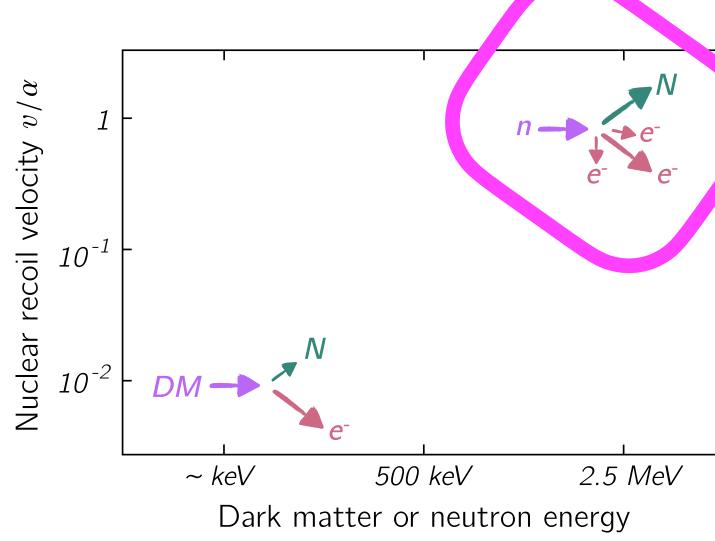
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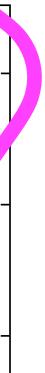


(2)

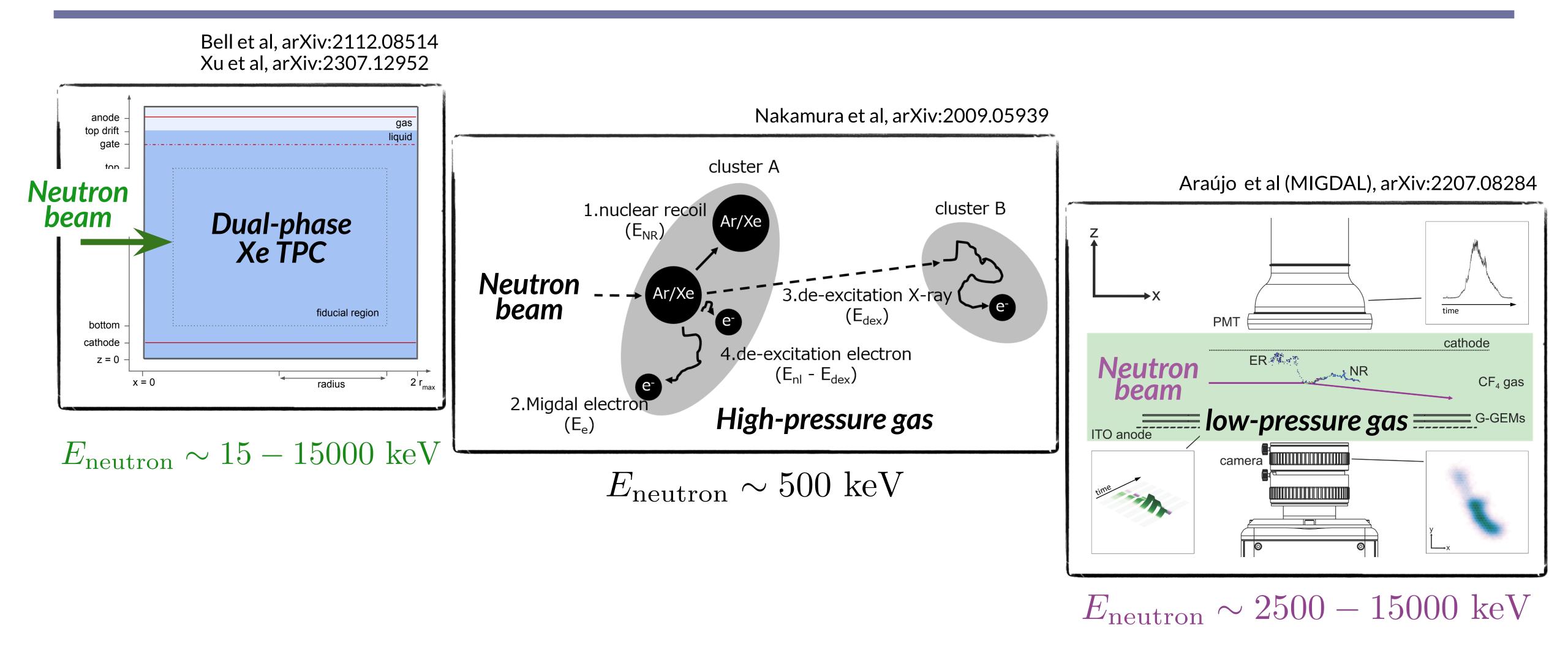


Large v regime: searches for the effect with neutrons





Finding evidence: Proposals with neutrons

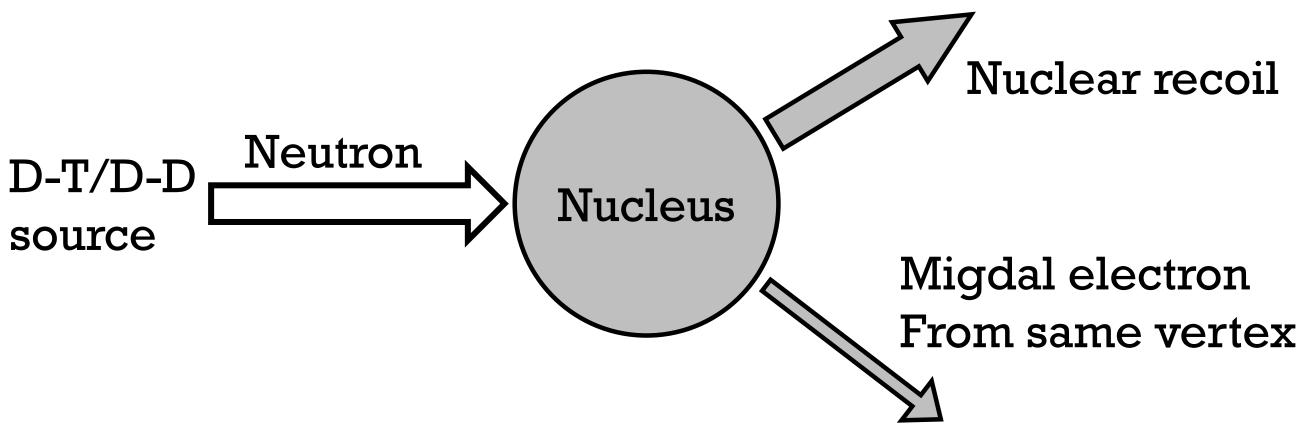








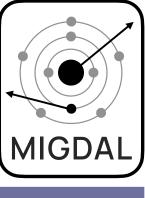
effect in nuclear scattering:



- Phase 1: Observe the effect in CF4 in high energy recoils
- Phase 2: Observe the Migdal effect in CF4 + noble gases



MIGDAL experiment: aims IMPER]



Create a dedicated experiment for the *unambiguous* observation of the Migdal

We are the only experiment aiming to observe the nuclear and electron recoils emerging from a common vertex

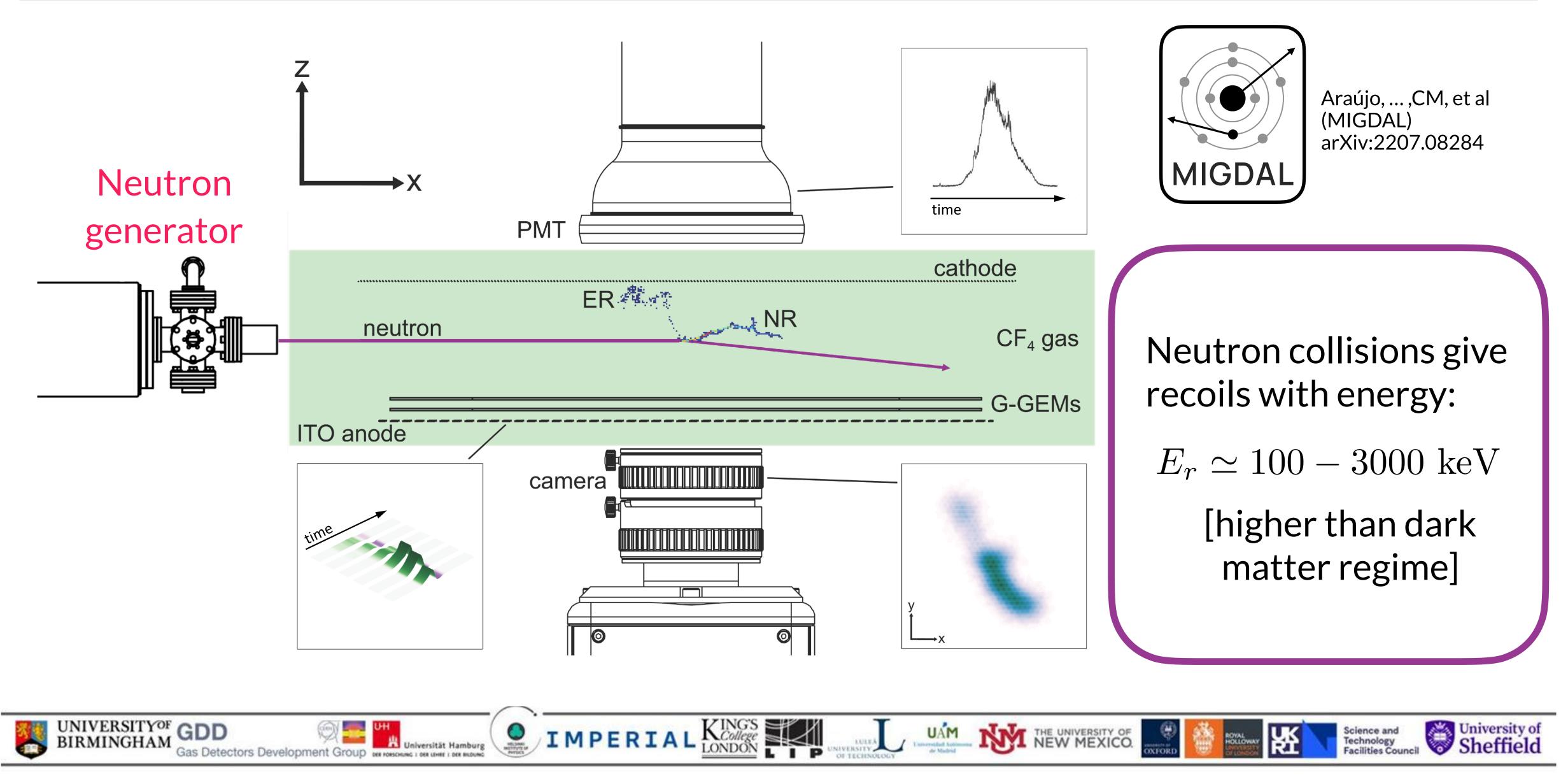








MIGDAL experiment: schematic MPERIAL





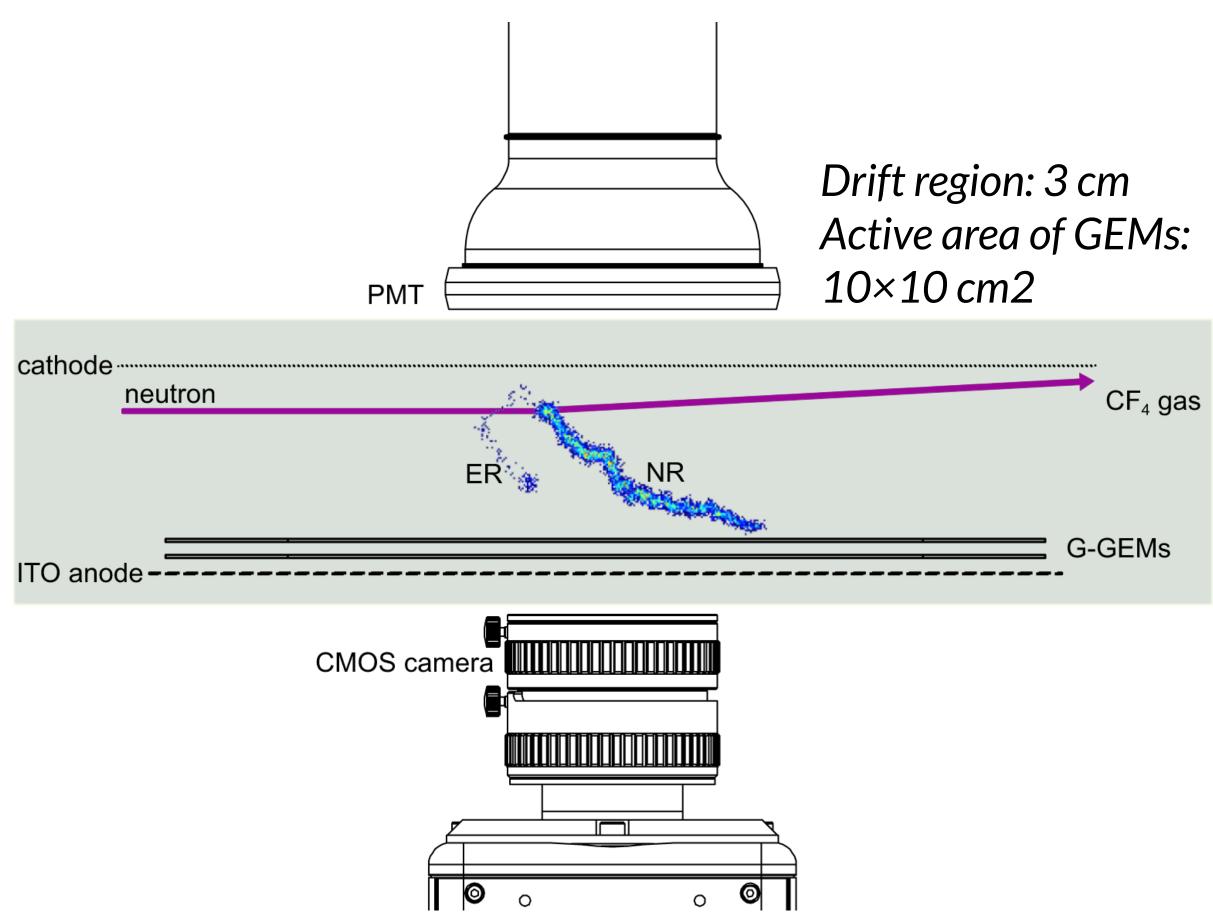
Optical Time Projection Chamber

Camera: images GEM scintillation through viewport behind ITO anode. Readout of (x,y) plane O(10) ms timing resolution

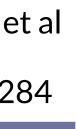
ITO anode: collects charge. Readout of (x,z) plane O(1) ns timing resolution

PMT: Detects primary and secondary (GEM) scintillation Readout of depth (z) coordinate

Setup allows for 3D track reconstruction



Simulated Migdal event with a 10 keV electron & 250 keV fluorine recoil. Scaled-up by a factor of 3.





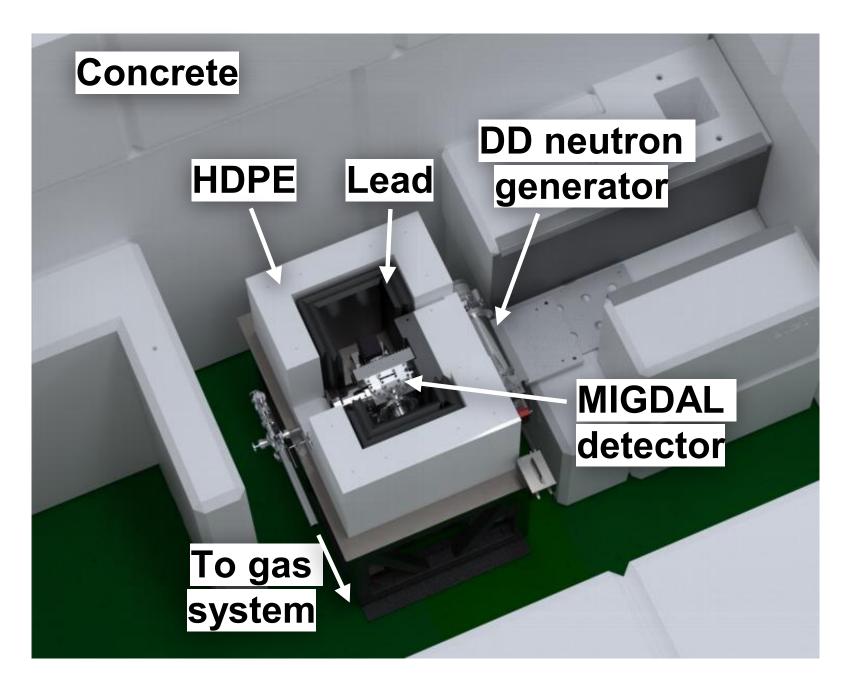
Installed at NILE Facility, Rutherford Appleton Laboratory, UK

High-yield DD - 10⁹ n/s @ 2.

Bespoke DD n within Target Source



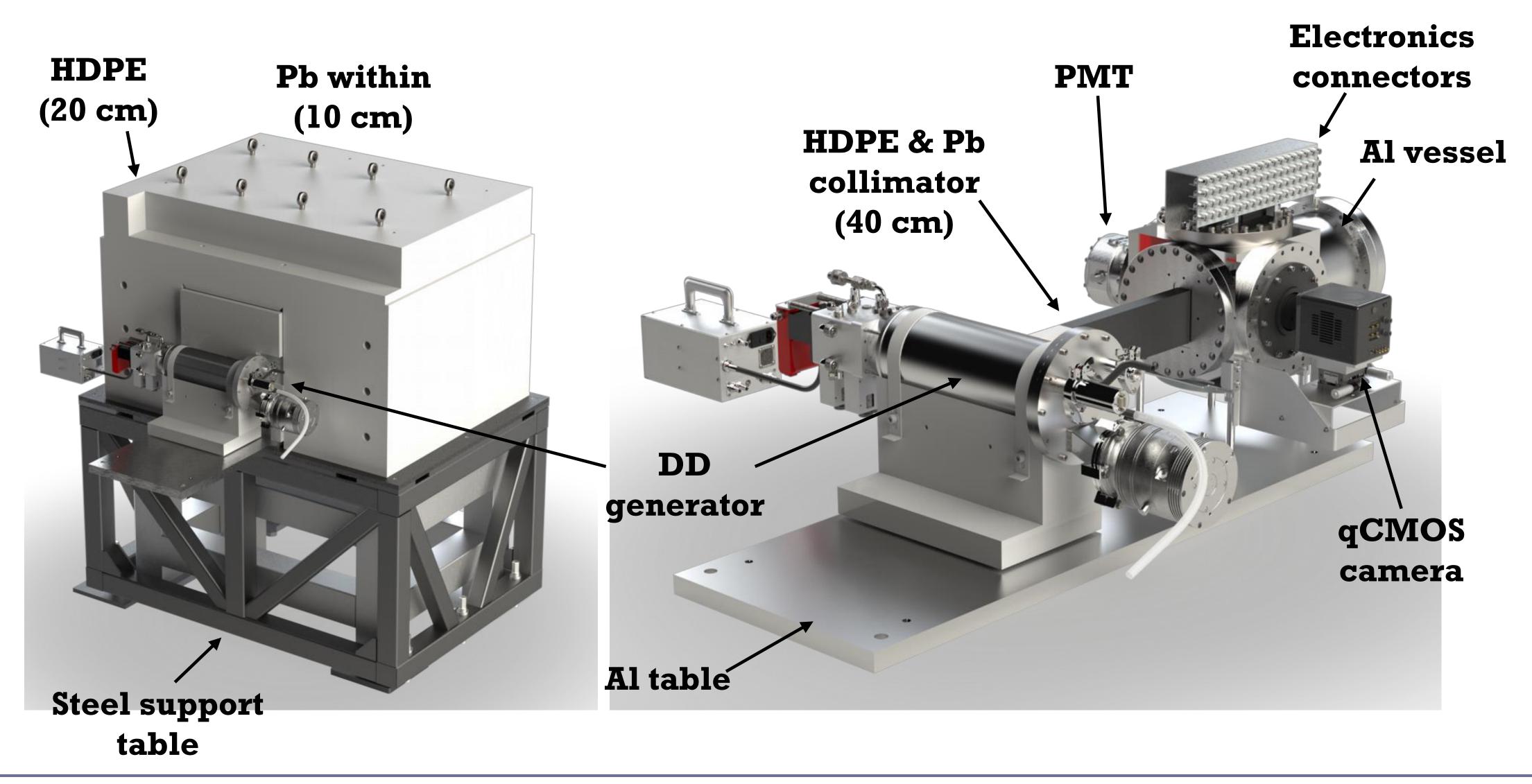
Concrete bunker with interlocked access D-D generator installed in "shielding bunker" MIGDAL detector in the centre of the bunker





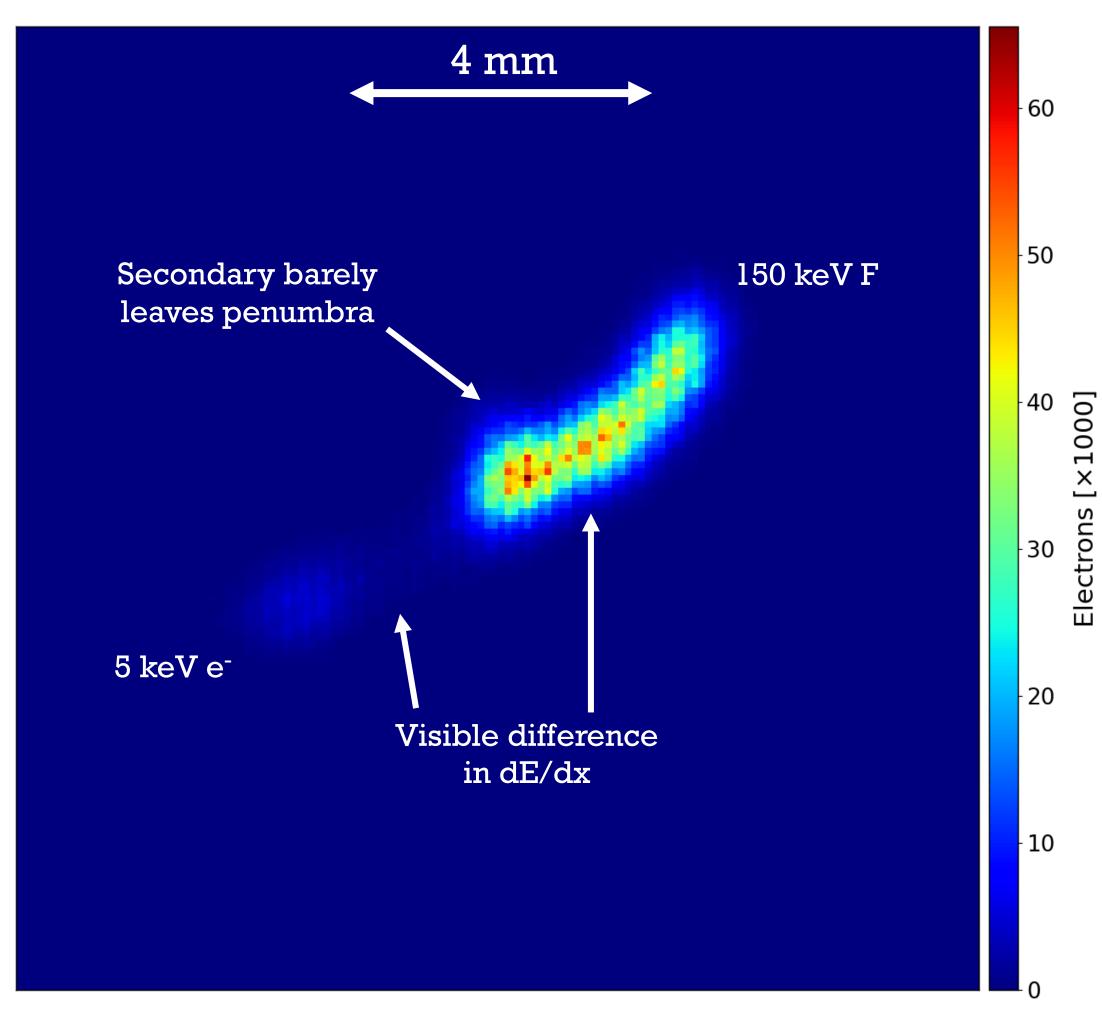


Shielded and Unshielded renders of the experiment



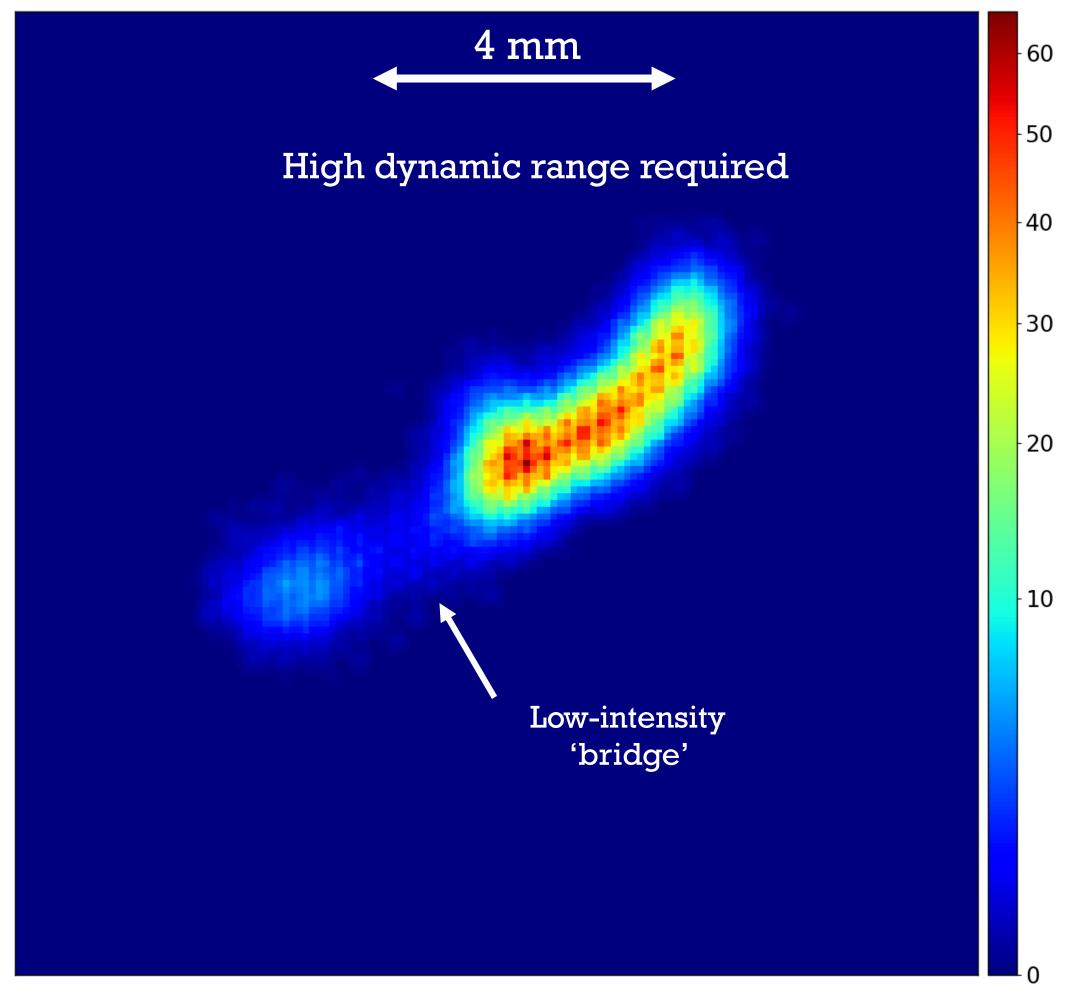


Simulated camera images of Migdal event

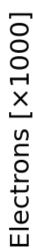


Linear-scale colour map

Christopher McCabe

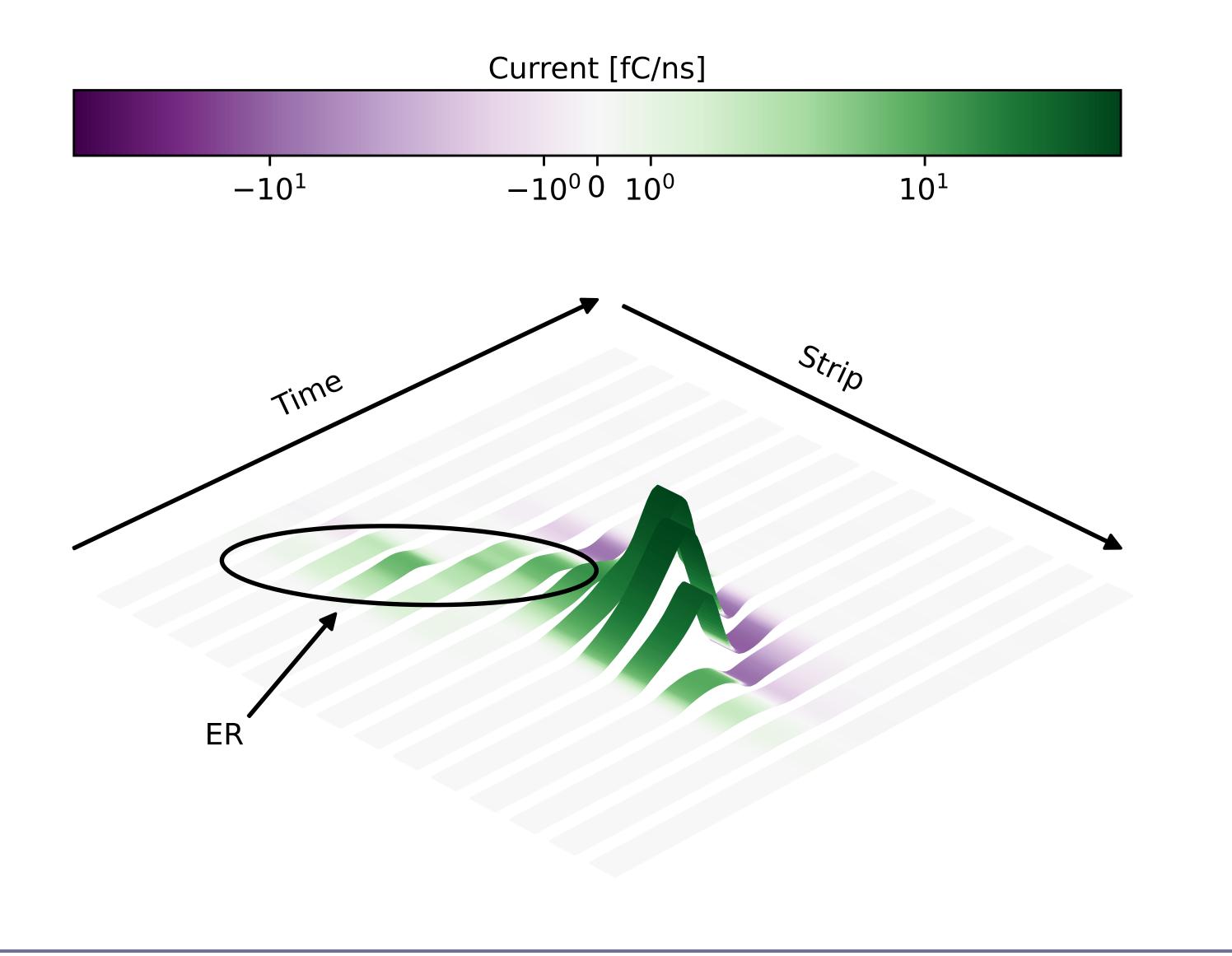


Log-scale colour map

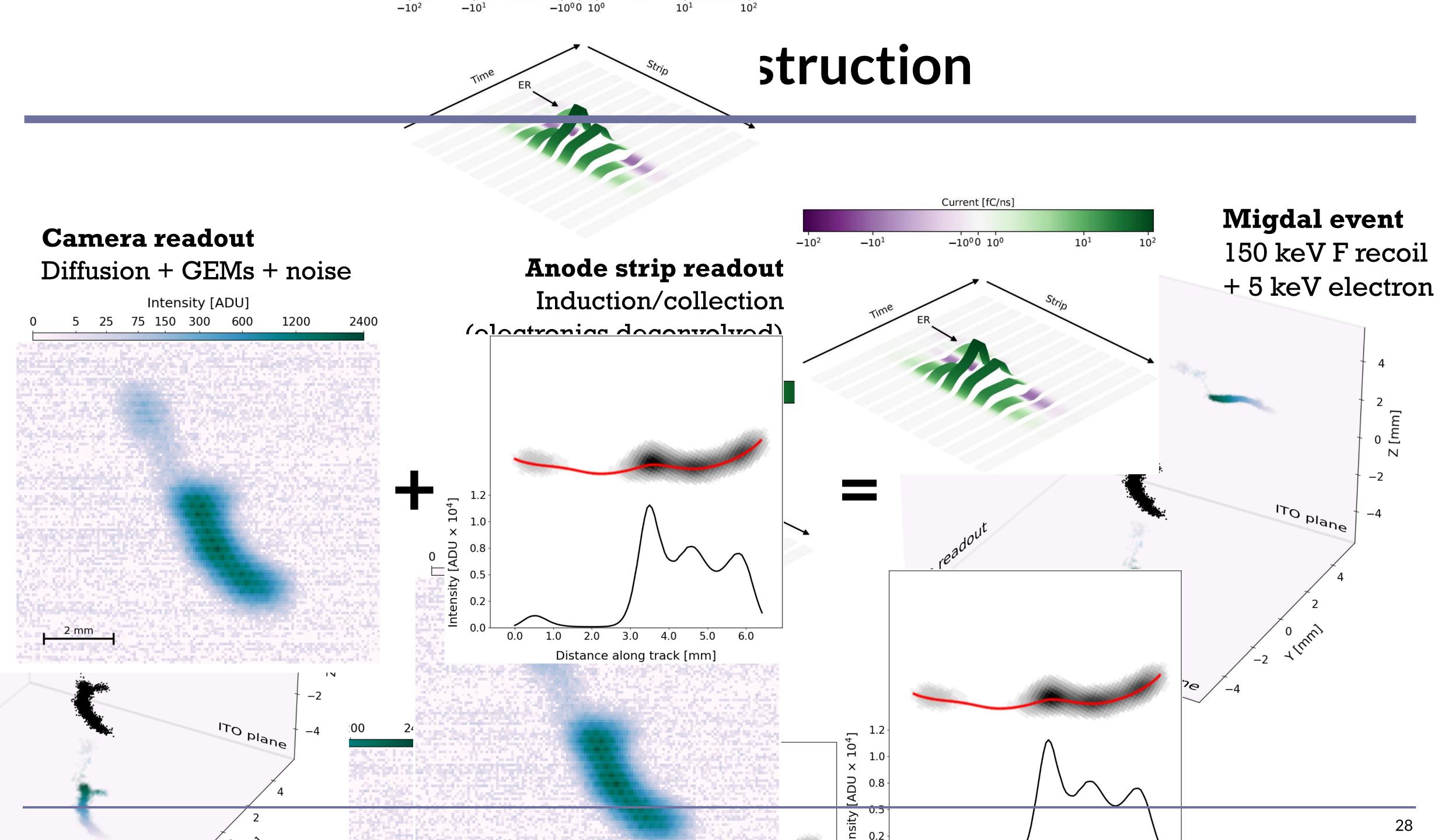




Simulated ITO signals of Migdal event



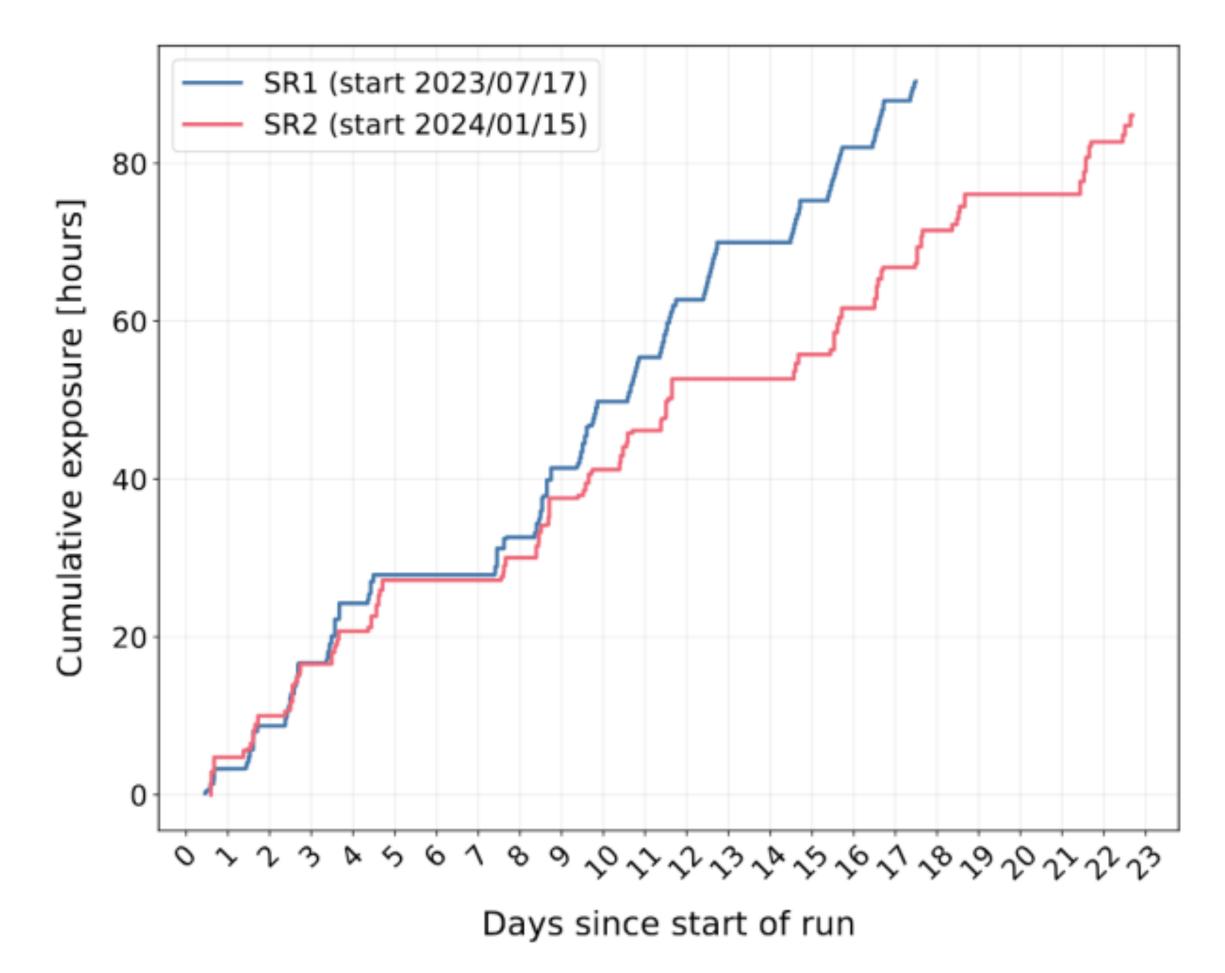




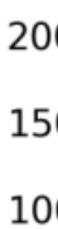


- First Science Run: 17/07/23 03/08/23
- Second Science Run: 15/01/24 06/02/24
- Data taken using D-D neutron generator recorded continuously during 10-hour shifts
- 50% data remained blinded
- Approximately 500,000 NRs in total

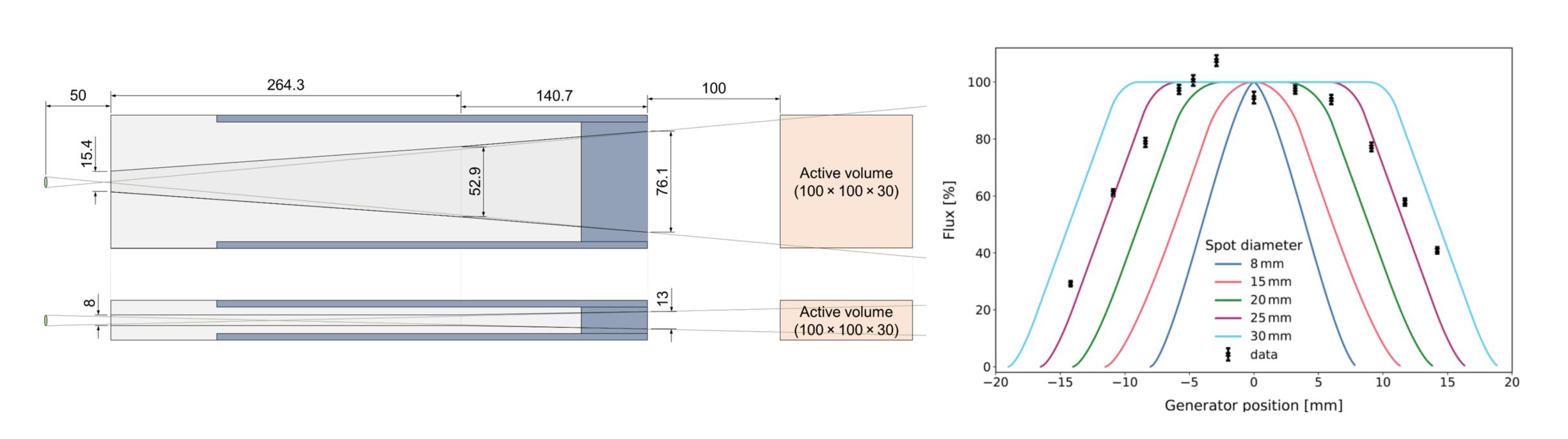
Science operations







Science operations: neutron beam

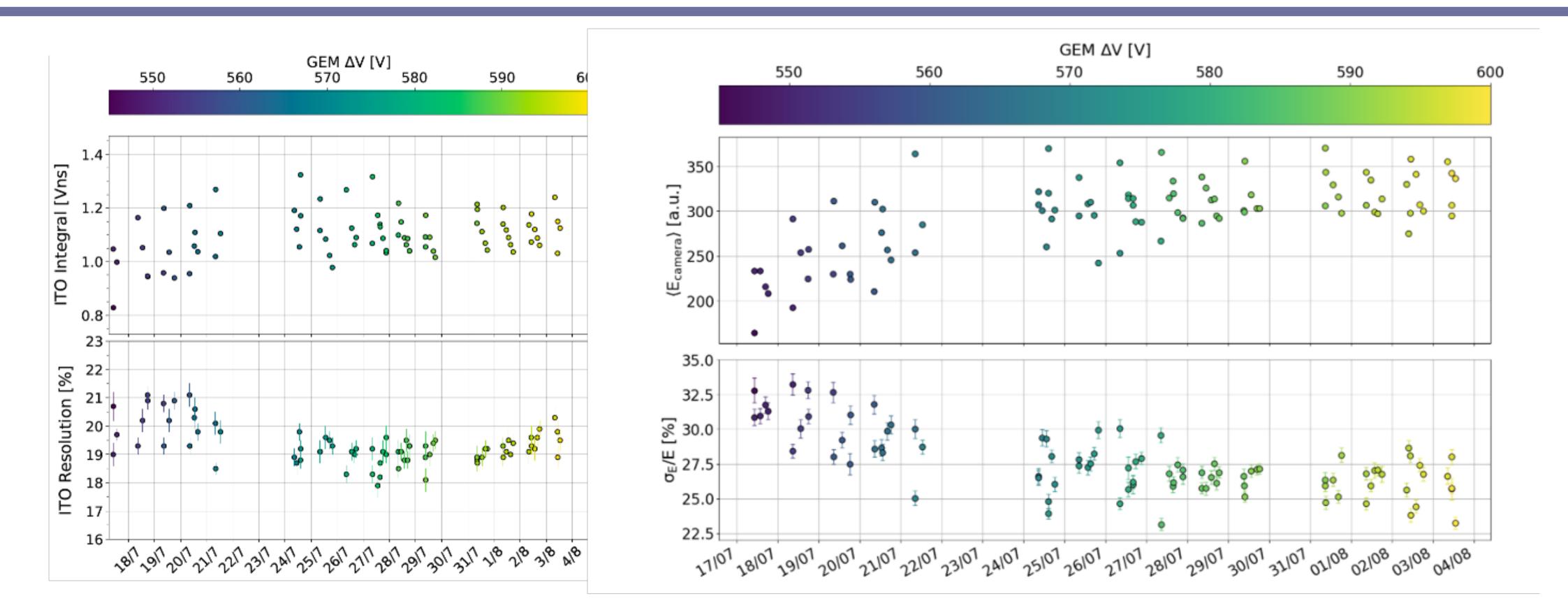


- Expected 2.6×10^5 n/s entering the active volume, but measured 6×10^4 n/s.
- Collimator designed ~8 mm neutron production spot diameter within the DD generator: but measured diameter was ~25 mm.
- This reduced the NR event rate in the active volume from ~15 Hz to ~5 Hz. - The camera was pulled closer to the active volume to capture more light.
 - This further reduced the contained NR rate in the ROI to ~2 Hz, which we observe in the data.





Science operations: detector calibrations



- ⁵⁵Fe calibration performed several times per day
- Energy scale is consistent over the course of the science run with ~20% variation
- Resolution in ITO ~20% and in camera ~ 25 32 % camera readout depending on the gain

day the science run with ~20% variation 32 % camera readout depending on the gain



Backgrounds

| Component | Torologra | D-D neutrons | | D-T neutrons | |
|--|--|--------------|--------------------|--------------|--------------------|
| | Topology | >0.5 | $515~\mathrm{keV}$ | > 0.5 | $515~\mathrm{keV}$ |
| Recoil-induced δ -rays | Delta electron from NR track origin | ≈ 0 | 0 | 541,000 | 0 |
| Particle-Induced X-ray Emission (PIXE) | | | | | |
| X-ray emission | Photoelectron near NR track origin | 1.8 | 0 | 365 | 0 |
| Auger electrons | Auger electron from NR track origin | 19.6 | 0 | $42,\!000$ | 0 |
| ${\rm Bremsstrahlung} \ {\rm processes}^{\dagger}$ | | | | | |
| Quasi-Free Electron Br. (QFEB) | Photoelectron near NR track origin | 112 | ≈ 0 | 288 | ≈ 0 |
| Secondary Electron Br. (SEB) | Photoelectron near NR track origin | 115 | ≈ 0 | 279 | ≈ 0 |
| Atomic Br. (AB) | Photoelectron near NR track origin | 70 | ≈ 0 | 171 | ≈ 0 |
| Nuclear Br. (NB) | Photoelectron near NR track origin | ≈ 0 | ≈ 0 | 0.013 | ≈ 0 |
| Photon interactions | | | | | |
| Neutron inelastic γ -rays (gas) | Compton electron near NR track origin | 1.6 | 0.47 | 0.86 | 0.25 |
| Random track coincidences | Photo-/Compton electron near NR track | ≈ 0 | ≈ 0 | ≈ 0 | ≈ 0 |
| Gas radioactivity | | | | | |
| Trace contaminants | Electron from decay near NR track origin | 0.2 | 0.01 | 0.03 | ≈ 0 |
| Neutron activation | Electron from decay near NR track origin | 0 | 0 | ≈ 0 | ≈ 0 |
| Secondary nuclear recoil fork | NR track fork near track origin | | ~ 1 | | ~ 1 |
| Total background | Sum of the above components | | 1.5 | | 1.3 |
| Migdal signal | Migdal electron from NR track origin | | 32.6 | | 84.2 |

[†] These processes were (conservatively) evaluated at the endpoint of the nuclear recoil spectra.

Araújo, ... ,CM, et al (MIGDAL) arXiv:2207.08284

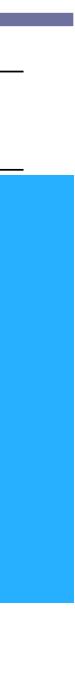


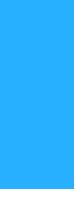
Backgrounds

| Component | | D-D neutrons | | D-T neutrons | |
|---|--|--------------|--------------------|--------------|--------------------|
| Component | Topology | | $515~\mathrm{keV}$ | > 0.5 | $515~\mathrm{keV}$ |
| Recoil-induced δ -rays | Delta electron from NR track origin | ≈ 0 | 0 | $541,\!000$ | 0 |
| Particle-Induced X-ray Emission (PIXE) | | | | | |
| X-ray emission | Photoelectron near NR track origin | 1.8 | 0 | 365 | 0 |
| Auger electrons | Auger electron from NR track origin | 19.6 | 0 | $42,\!000$ | 0 |
| Bremsstrahlung processes [†] Elimi | eshol | d | | | |
| Quasi-Free Electron Br. (QFEB) | nated by applying energy three Photoelectron near NR track origin | 112 | ≈ 0 | 288 | ≈ 0 |
| Secondary Electron Br. (SEB) | Photoelectron near NR track origin | 115 | ≈ 0 | 279 | ≈ 0 |
| Atomic Br. (AB) | Photoelectron near NR track origin | 70 | ≈ 0 | 171 | ≈ 0 |
| Nuclear Br. (NB) | Photoelectron near NR track origin | ≈ 0 | ≈ 0 | 0.013 | ≈ 0 |
| Photon interactions | | | | | |
| Neutron inelastic γ -rays (gas) | Compton electron near NR track origin | 1.6 | 0.47 | 0.86 | 0.25 |
| Random track coincidences | Photo-/Compton electron near NR track | ≈ 0 | ≈ 0 | ≈ 0 | ≈ 0 |
| Gas radioactivity | minated by ITO timing resolu | tion | | | |
| Trace contaminants | Electron from decay near NR Back origin | | 0.01 | 0.03 | ≈ 0 |
| Neutron activation | Electron from decay near NR track origin | 0 | 0 | ≈ 0 | ≈ 0 |
| Secondary nuclear recoil fork | NR track fork near track origin | | ~ 1 | | ~ 1 |
| Total background | Sum of the above components | | 1.5 | | 1.3 |
| Migdal signal | Migdal electron from NR track origin | | 32.6 | | 84.2 |
| | | | | | |

[†] These processes were (conservatively) evaluated at the endpoint of the nuclear recoil spectra.

Araújo, ... ,CM, et al (MIGDAL) arXiv:2207.08284

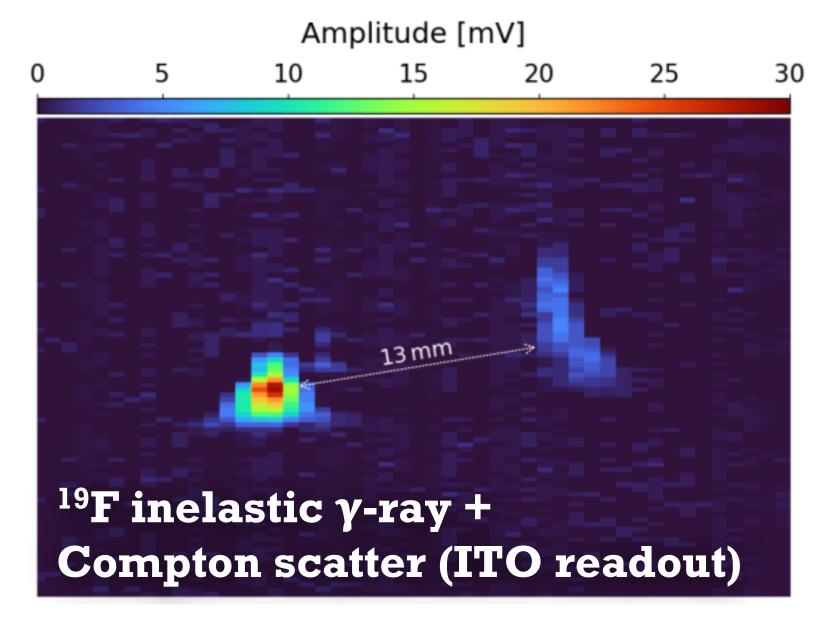


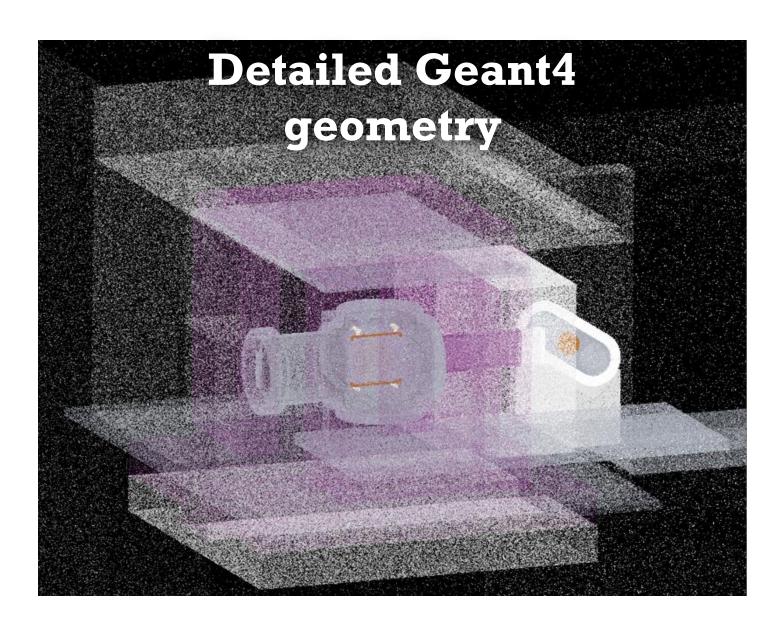






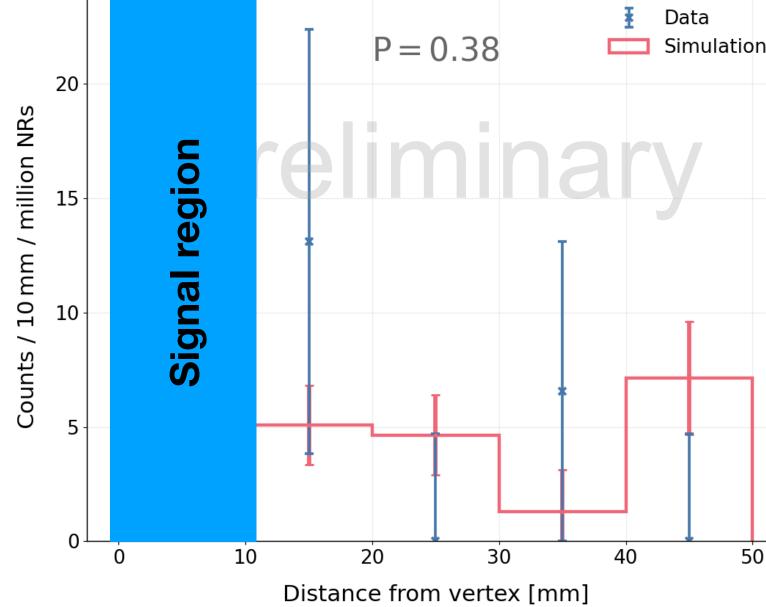
| (PIXE) | Donta chectron monn rene track origin | | 0 | |
|--------|--|-------------|-------------|-------------|
| | Photoelectron near NR track origin | 1.8 | 0 | |
| | Auger electron from NR track origin | 19.6 | 0 | :kgrou |
| EB) | Photoelectron near NR track origin | 112 | ≈ 0 | |
| 3) | Photoelectron near NR track origin | 115 | ≈ 0 | |
| | Photoelectron near NR track origin | 70 | ≈ 0 | |
| | Photoelectron near NR track origin | ≈ 0 | ≈ 0 | |
| | Compton electron near NR track origin | 1.6 | 0.47 | |
| | | | | eometry to |
| | Photo-/Compton electron near NR track | ≈ 0 | ≈ 0 | - |
| | Electron from decay near NR track origin | 0.2 | 0.01 | d NR + ER d |
| | Electron from decay near NR track origin | 0 | 0 | |
| | Delta electron near NR track origin | ≈ 0 | ≈ 0 | f ERs produ |
| | NR track fork near track origin | _ | ≈ 1 | I LINS PIOU |
| | Sum of the above components | | 1.5 | |
| | Migdal electron from NR track origin | | 32.6 | |
| | | - | | |





und simulations

- calculate the expected number of γ -rays
- coincidences is consistent
- luced within 3 mm of an NR vertex is very small (good news)



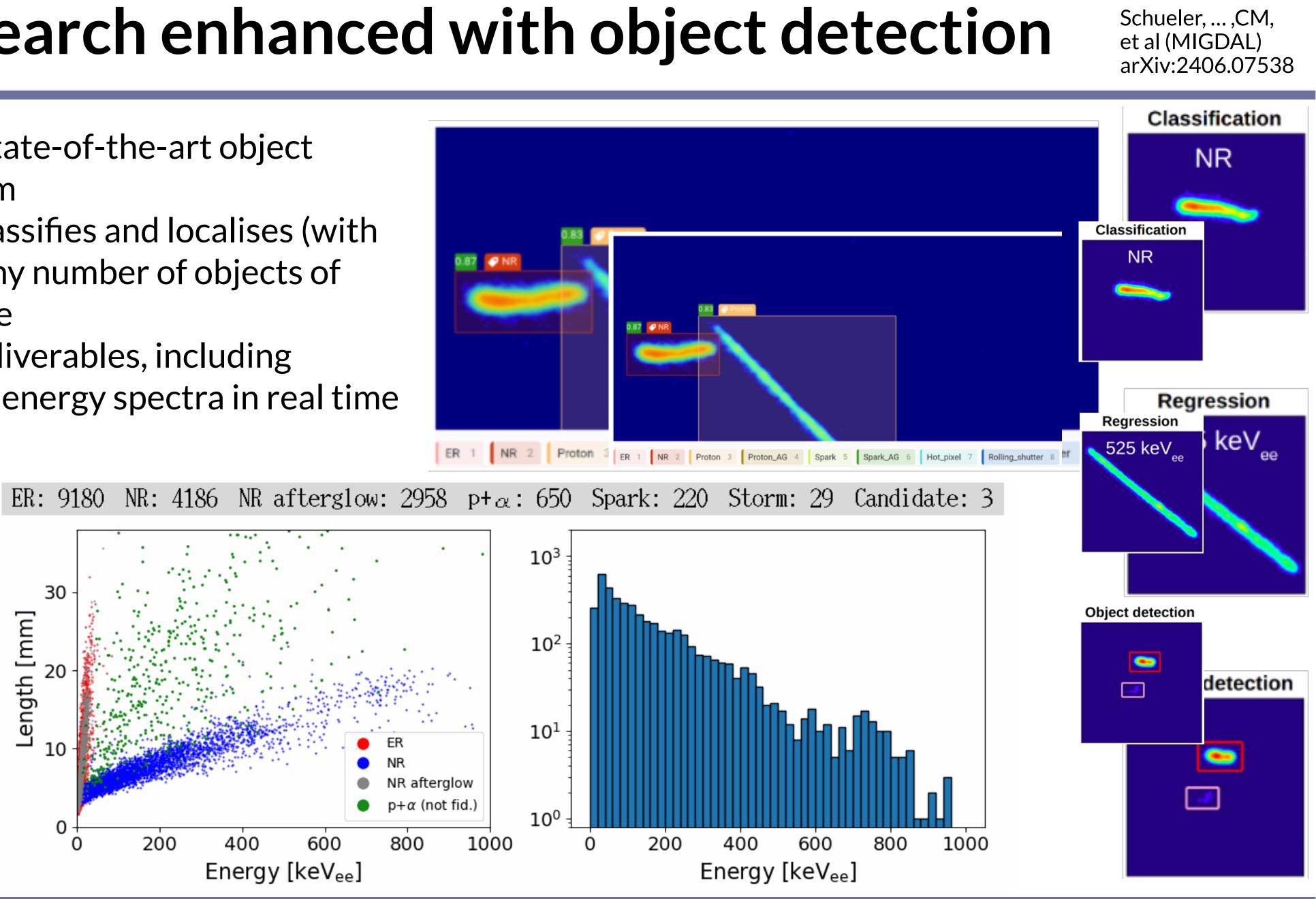


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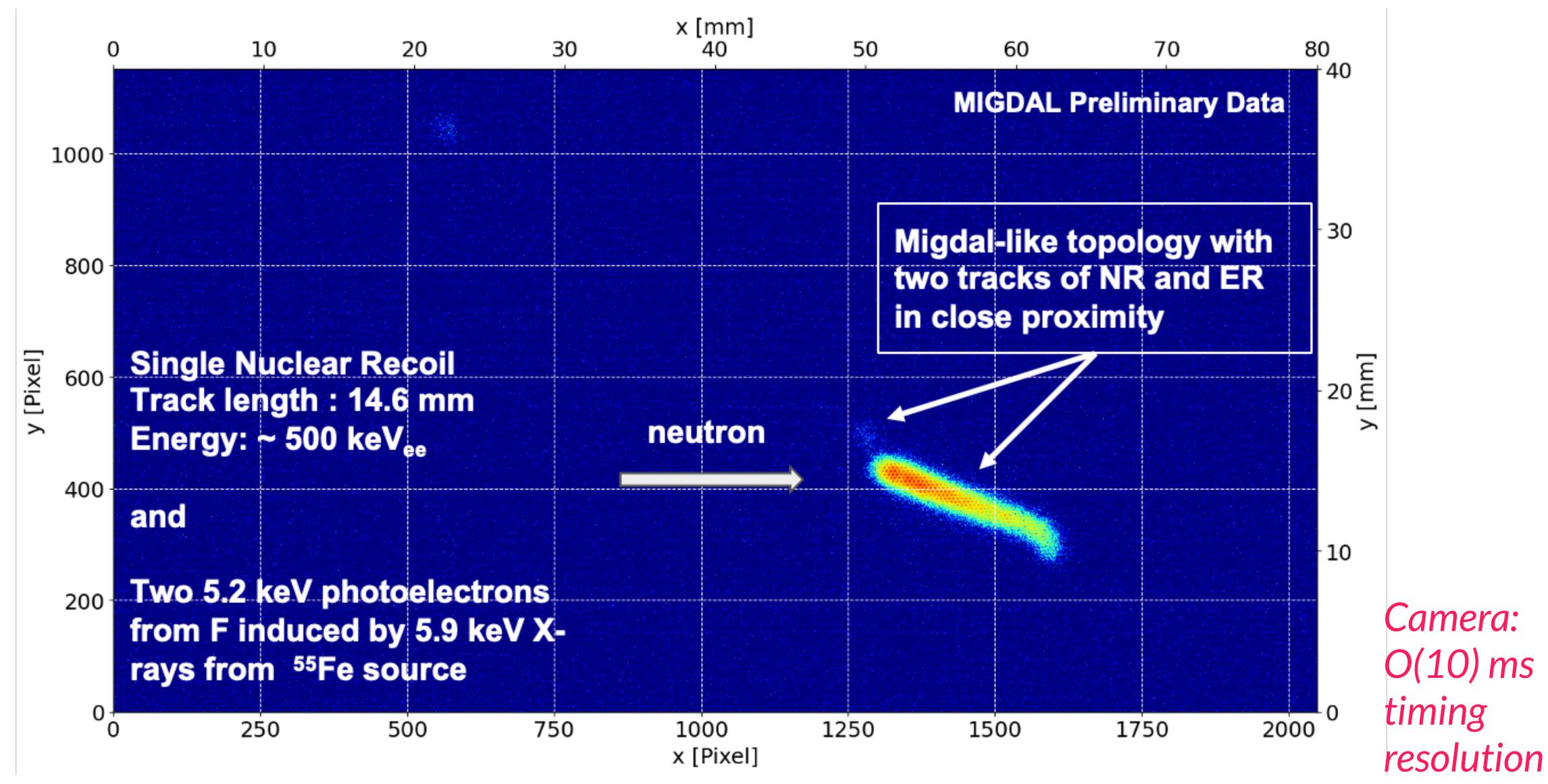
Search enhanced with object detection

- Utilise YOLOv8: state-of-the-art object detection algorithm
- Simultaneously classifies and localises (with bounding boxes) any number of objects of interest in an image
- Provides online deliverables, including particle ID and NR energy spectra in real time





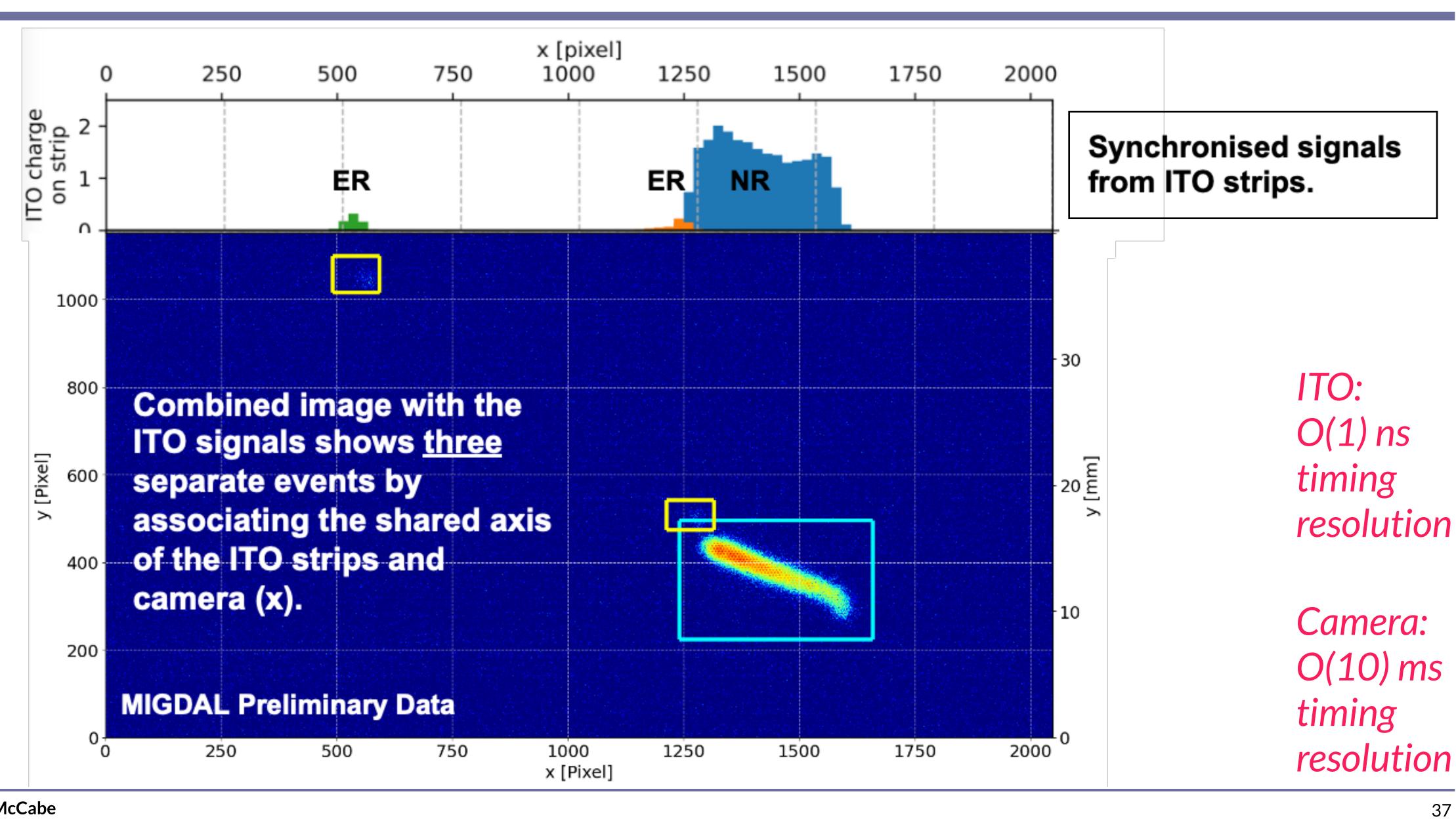
Example Migdal-topology event



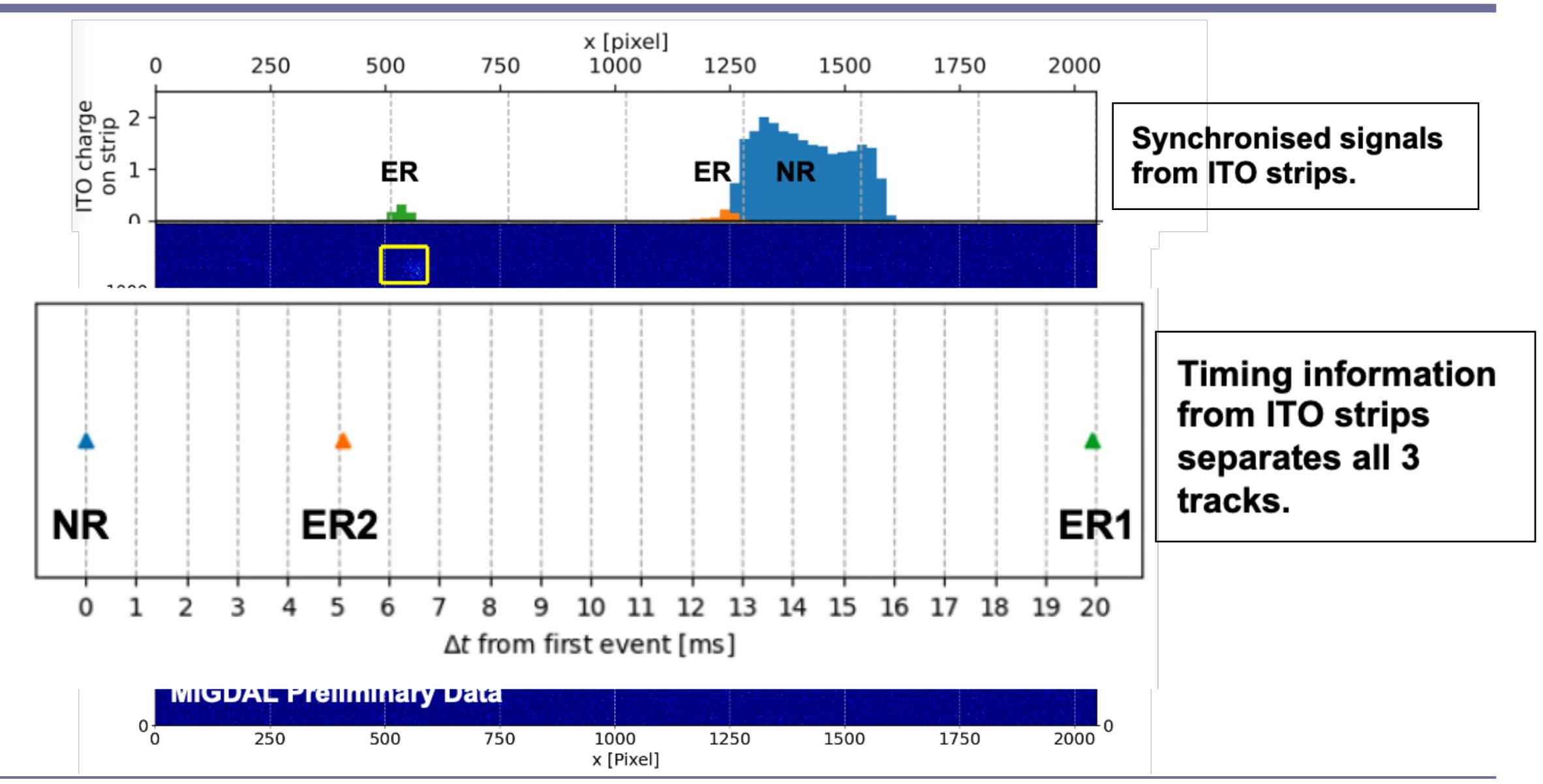


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Example Migdal-topology event



Not a Migdal event: no temporal overlap!





Summary

The Migdal effect is an old effect (from 1940s) that is used for dark matter sub-GeV searches

On the theory side, we have extended previous calculations to the high nuclear-recoil speed regime & confirmed the accuracy of existing calculations (lbe et al) for DM searches

In the UK...

- we are building a detection platform to characterise the effect in multiple elements



regime & confirmed the accuracy of existing calculations (lbe et al) for DM searches

In the UK...

- we are building a detection platform to characterise the effect in multiple elements
- Detector performed as designed during Science Runs 1 and 2
- We have several weeks of stable DD data: approximately 500,000 NRs in total
- Backgrounds appear to be as expected
- Data analysis of the two science runs is ongoing (stay tuned)

- The Migdal effect is an old effect (from 1940s) that is used for dark matter sub-GeV searches
- On the theory side, we have extended previous calculations to the high nuclear-recoil speed

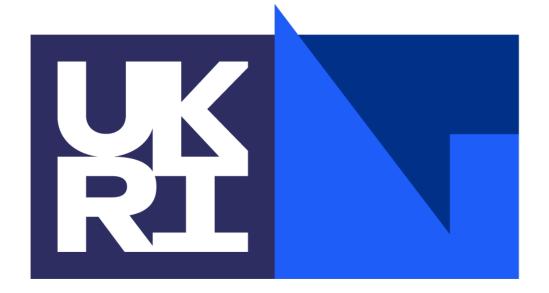
- Science Run 3 planned for this year with several improvements:
 - Higher resolution digitiser (CAEN V1730)
 - Increased spatial resolution in the ITO subsystem
 - Testing addition of a third GEM to provide additional amplification stage
 - Recommissioned DD generator with smaller spot size & redesigned collimator
 - Plans to run with CF4 + Ar mixture





Thank you

"The MIGDAL experiment: Measuring a rare atomic process to aid the search for "Precise Predictions and New Insights for Atomic Ionisation from the Migdal Effect" dark matter" H.M. Araújo et al (MIGDAL) Peter Cox, Matthew Dolan Christopher McCabe and Harry Quiney arXiv:2207.08284, Astroparticle Phys (2023) arXiv:2208.12222, PRD (2023) "Transforming a rare event search into a not-so-rare event search in real-time with Data files of probabilities available now: <u>https://petercox.github.io/Migdal/</u> deep learning-based object detection" J. Schueler et al (MIGDAL) arXiv:2406.07538, PRD (to appear)



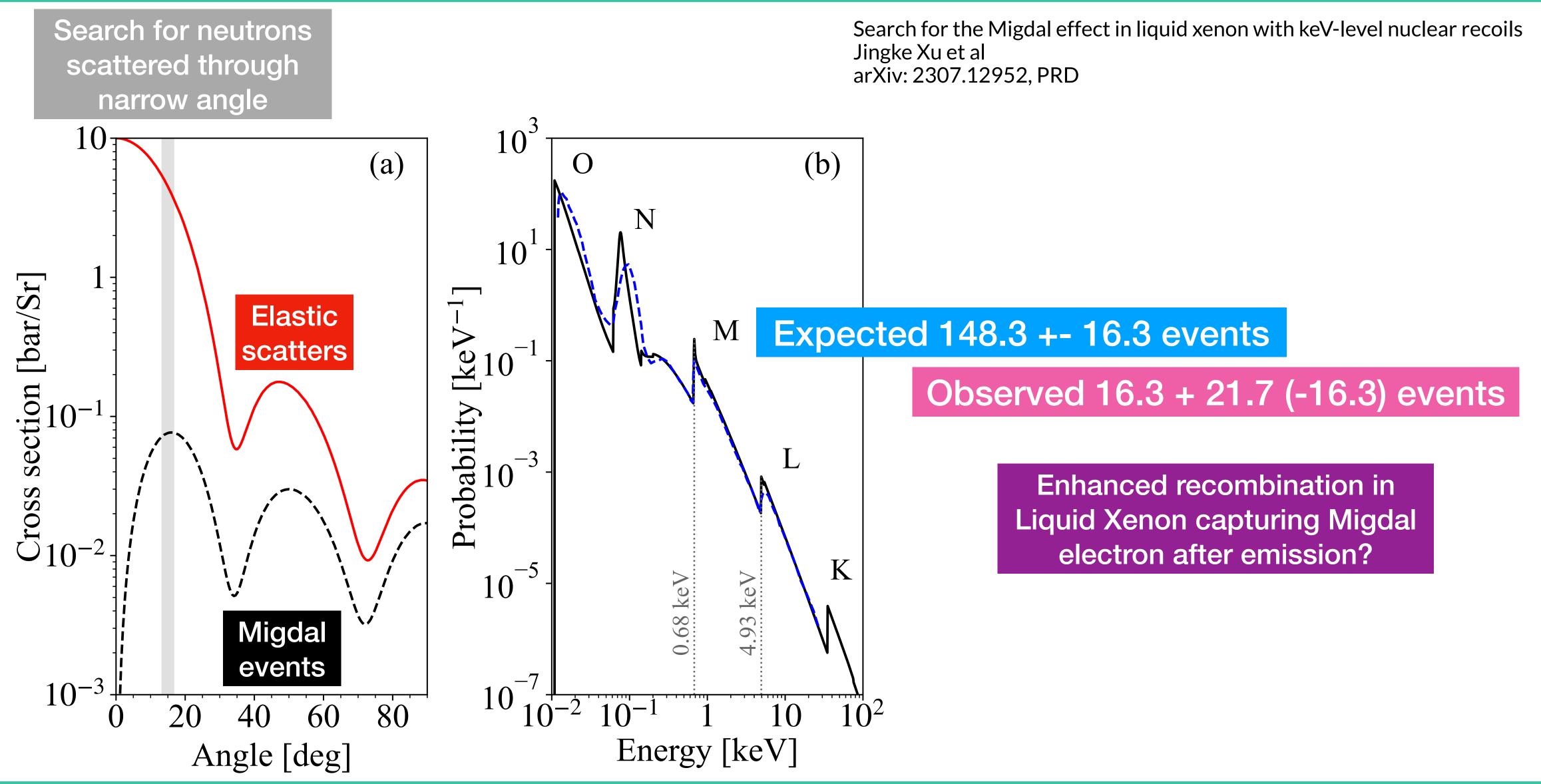
Science and Technology **Facilities Council**



Backup



Search in LXe with DT neutron generator





Search in LXe with DT neutron generator

