

LIGHT DARK MATTER DIRECT DETECTION WITH ANTI- FERROMAGNETS

Angelo Esposito



SAPIENZA
UNIVERSITÀ DI ROMA



Istituto Nazionale di Fisica Nucleare

Light Dark Matter 2025, Genova

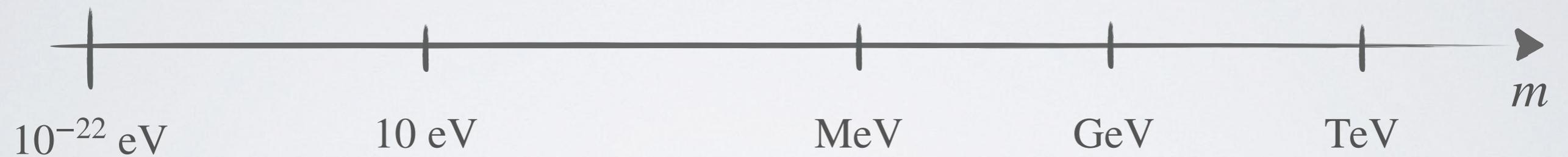
SUB-MEV DARK MATTER

SUB-MEV DARK MATTER

- Possible dark matter mass range is huge

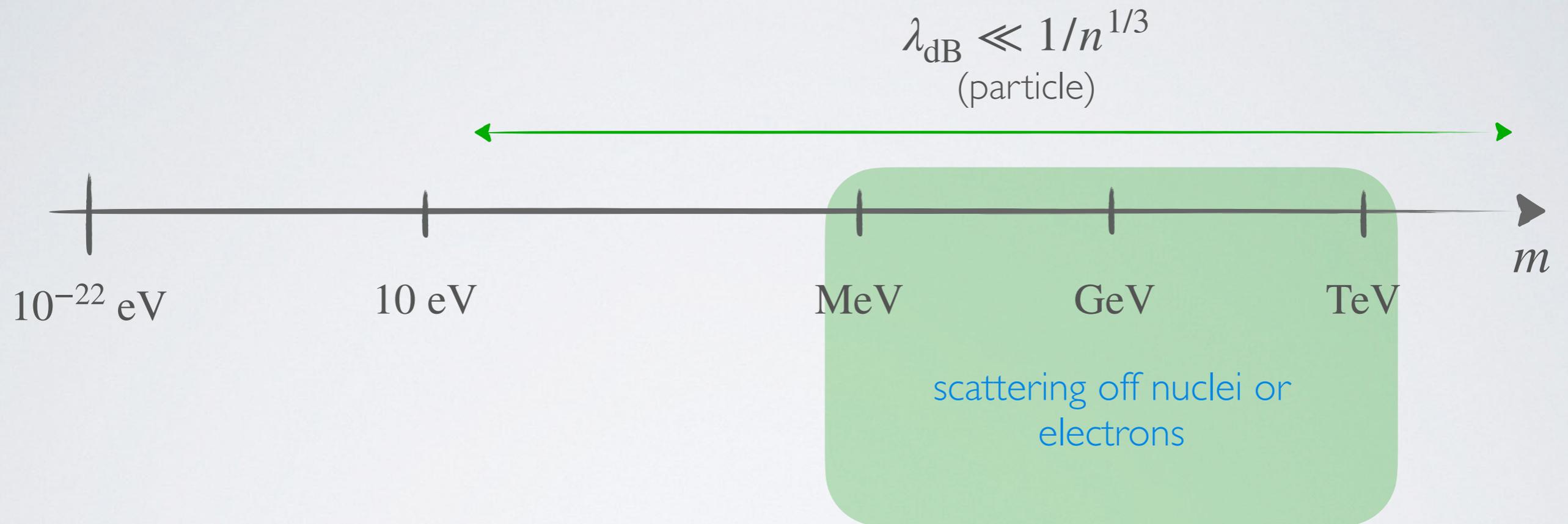
SUB-MEV DARK MATTER

- Possible dark matter mass range is huge



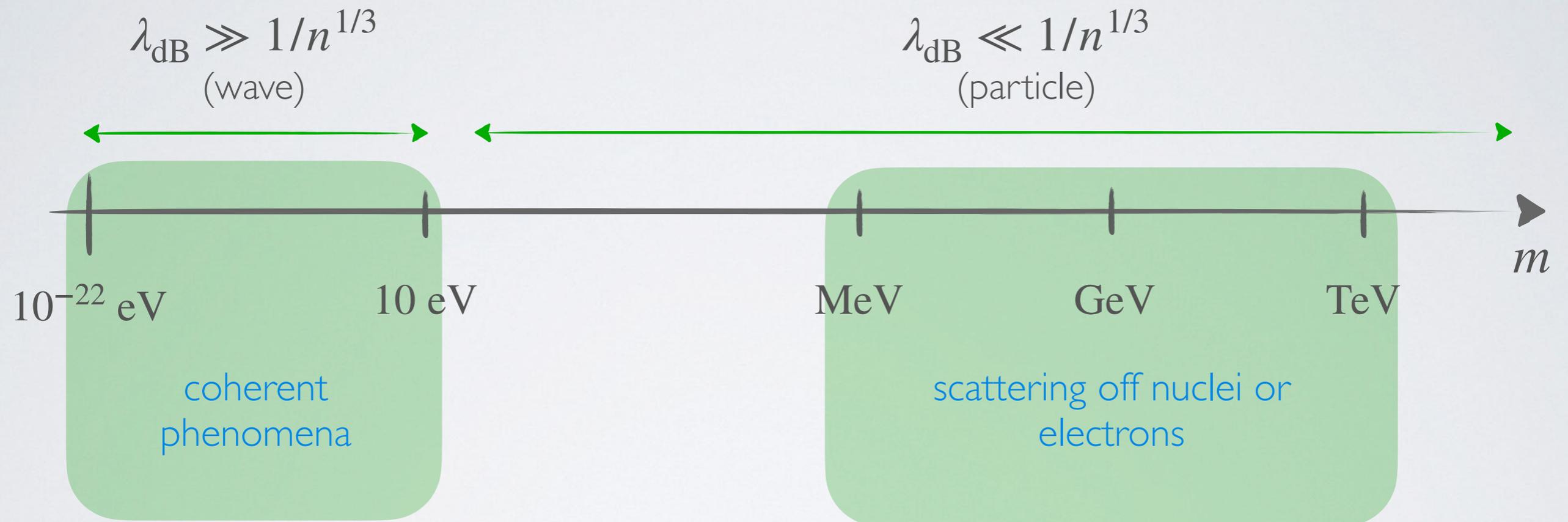
SUB-MEV DARK MATTER

- Possible dark matter mass range is huge



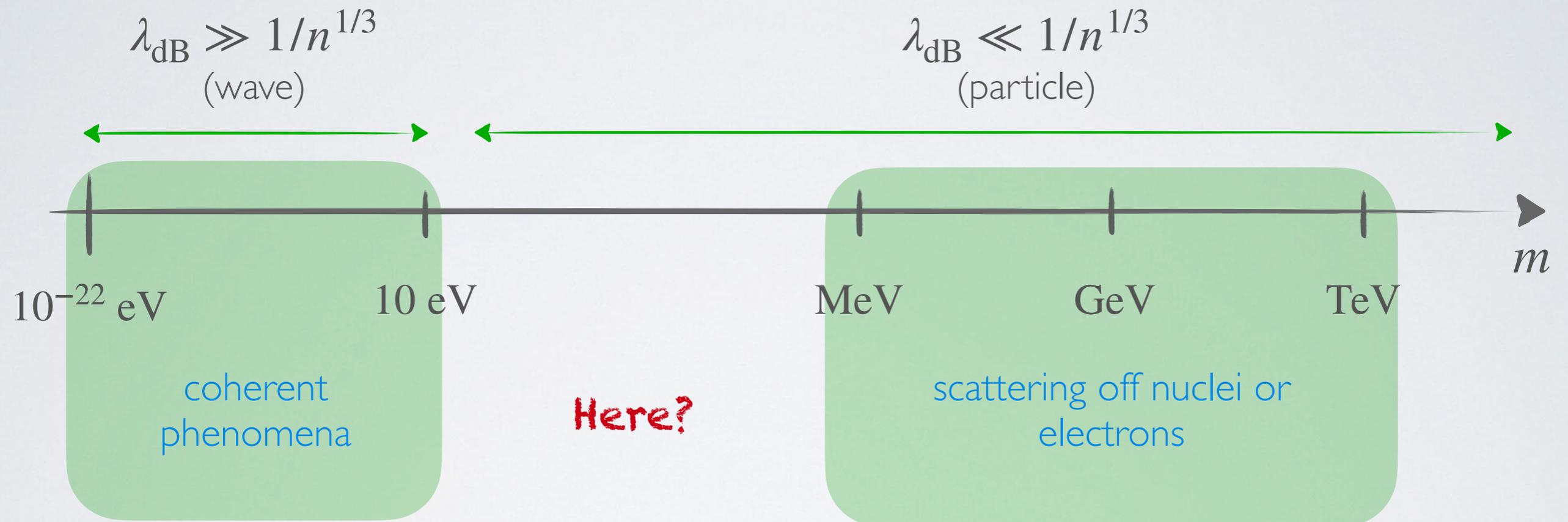
SUB-MEV DARK MATTER

- Possible dark matter mass range is huge



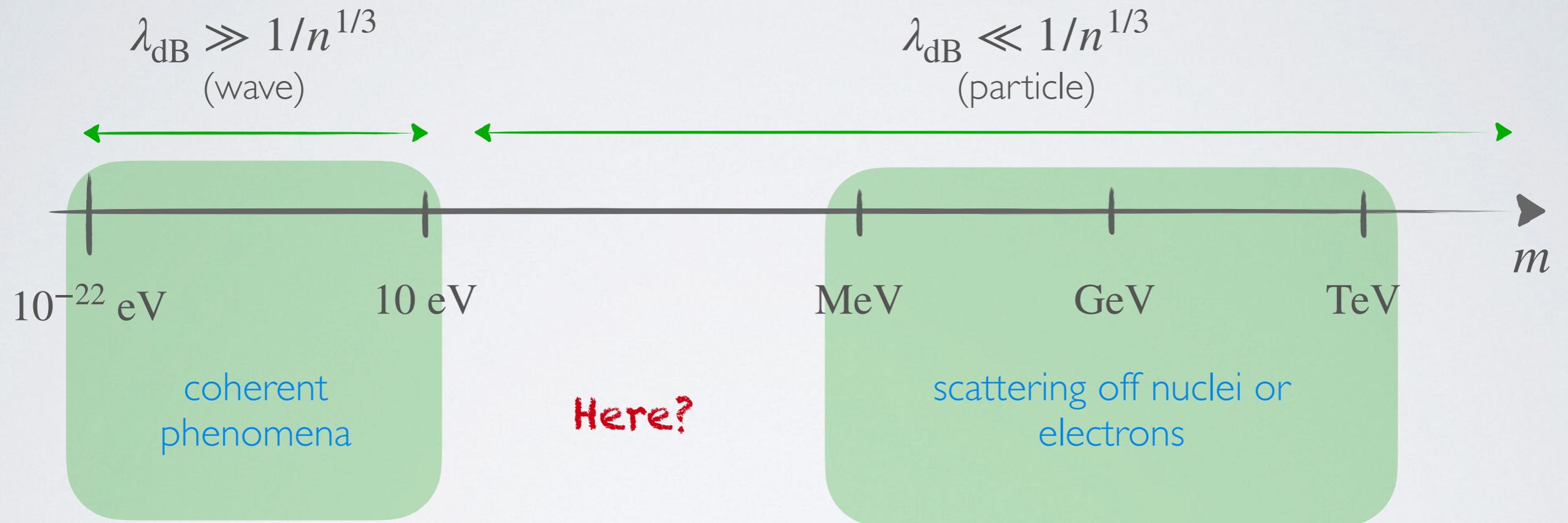
SUB-MEV DARK MATTER

- Possible dark matter mass range is huge



SUB-MEV DARK MATTER

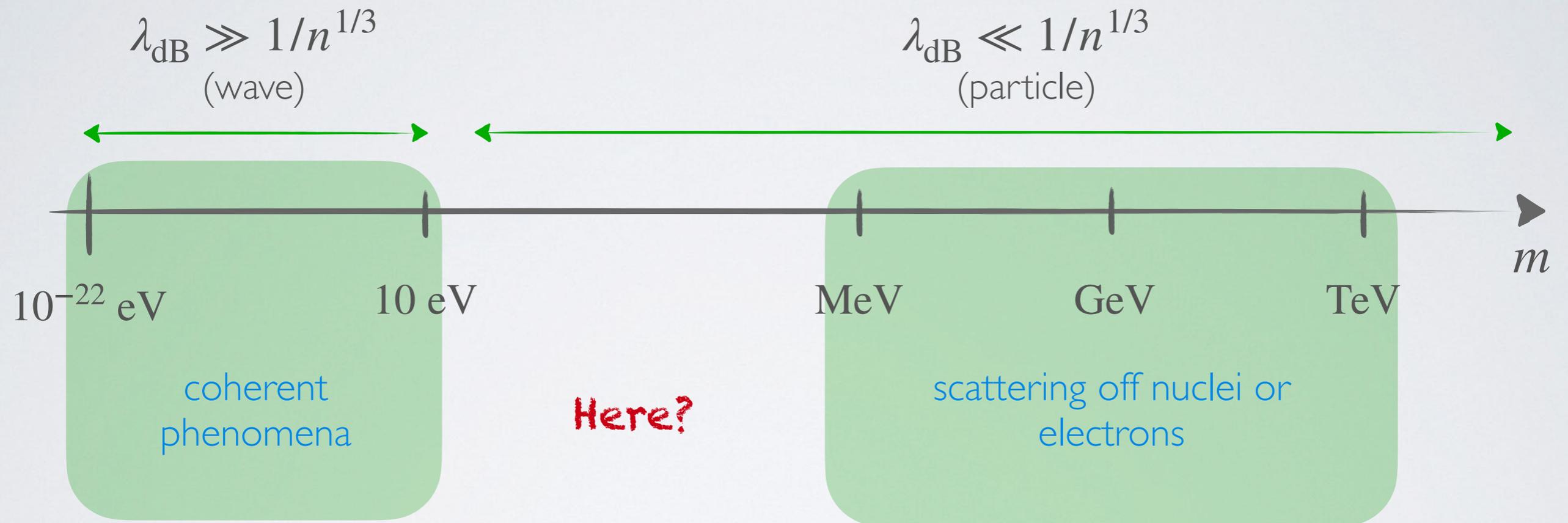
- Possible dark matter mass range is huge



- Dark matter is a particle but too light for recoil off nuclei / electrons

SUB-MEV DARK MATTER

- Possible dark matter mass range is huge



- Dark matter is a particle but too light for recoil off nuclei / electrons
- Need new materials and/or observables

SUB-MEV DARK MATTER

SUB-MEV DARK MATTER

- For an elastic scattering

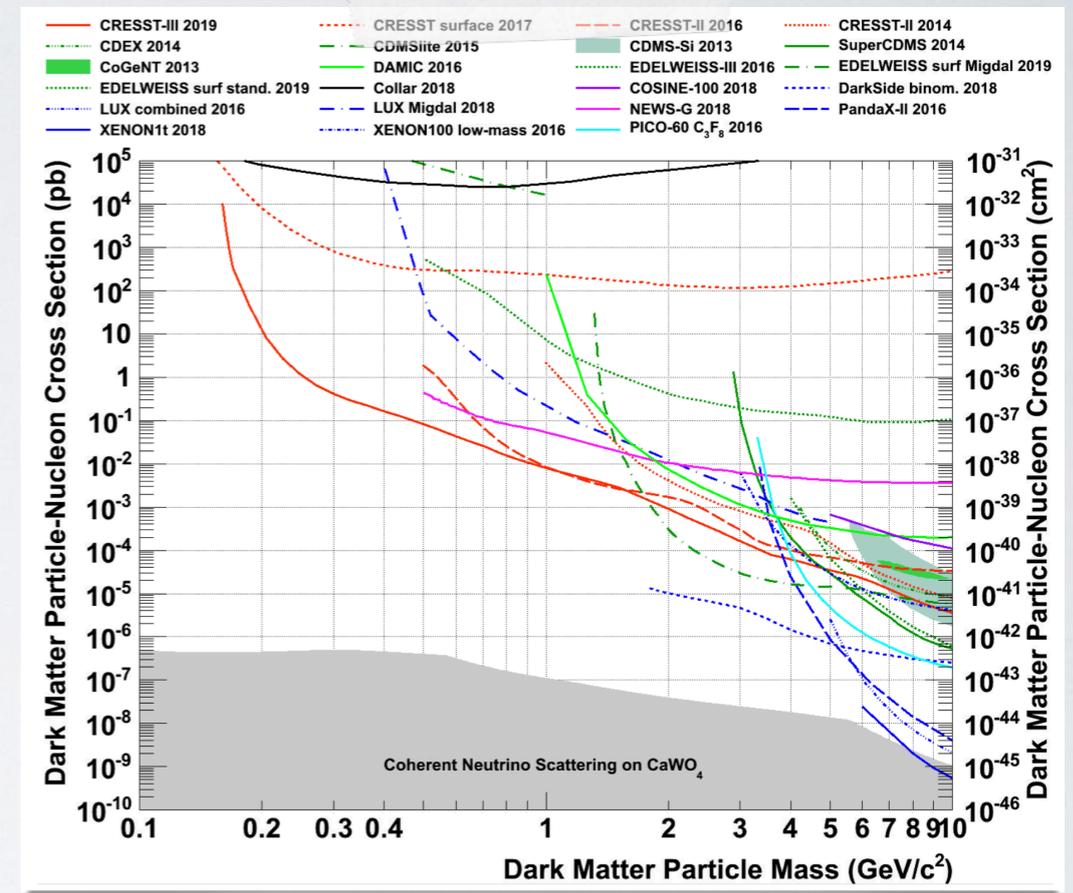
$$E_T = 4 \frac{m_\chi/m_T}{(1 + m_\chi/m_T)^2} E_\chi \gtrsim E_{\text{exp}}$$

SUB-MEV DARK MATTER

- For an elastic scattering

$$E_T = 4 \frac{m_\chi / m_T}{(1 + m_\chi / m_T)^2} E_\chi \gtrsim E_{\text{exp}}$$

- For $m_\chi \lesssim 1 \text{ MeV}$ elastic scattering off nuclei or electrons is very inefficient



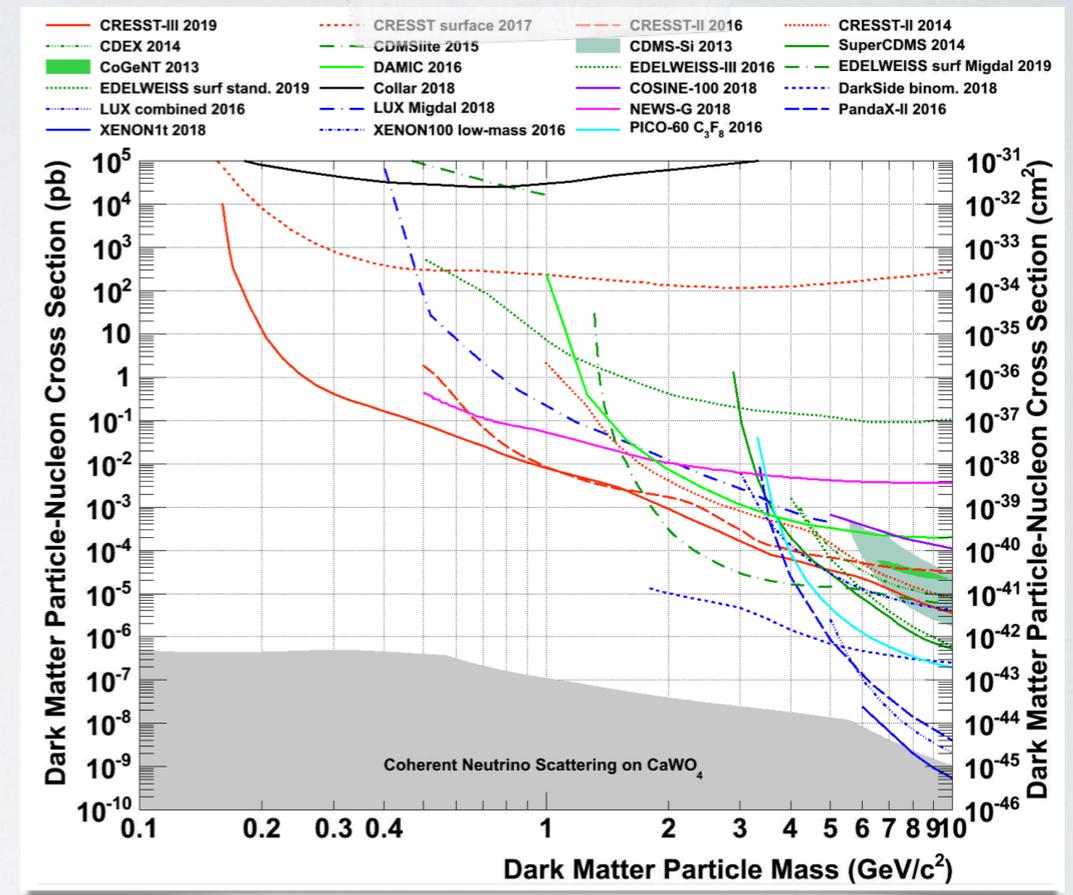
[CRESST PRD 2019]

SUB-MEV DARK MATTER

- For an elastic scattering

$$E_T = 4 \frac{m_\chi / m_T}{(1 + m_\chi / m_T)^2} E_\chi \gtrsim E_{\text{exp}}$$

- For $m_\chi \lesssim 1 \text{ MeV}$ elastic scattering off nuclei or electrons is very inefficient



[CRESST PRD 2019]

- To evade this we must look into inelastic processes

COLLECTIVE EXCITATIONS

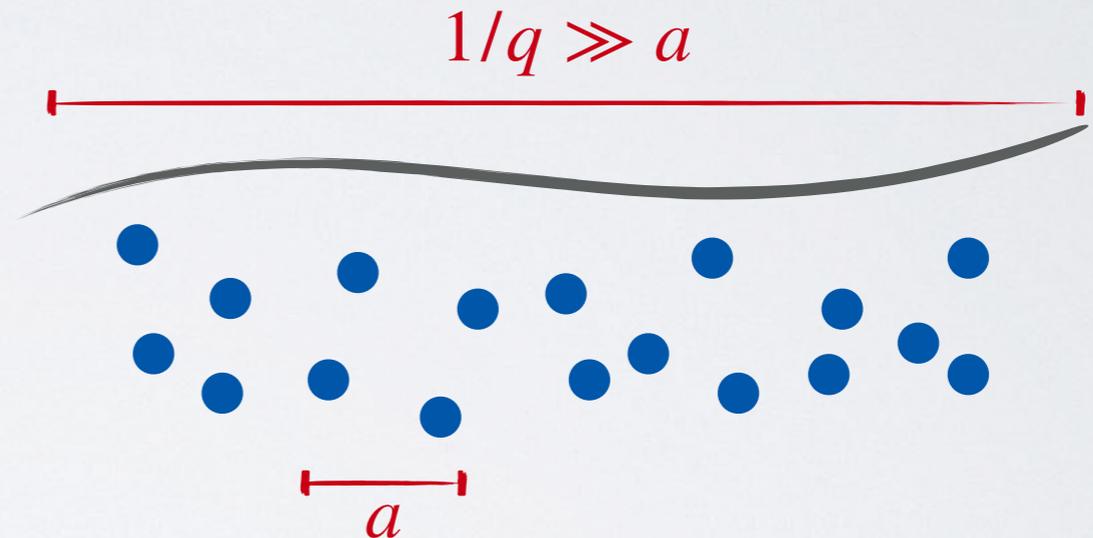
COLLECTIVE EXCITATIONS

- For $m_\chi \lesssim 1 \text{ MeV}$, we cannot talk about isolated particles in the detector

COLLECTIVE EXCITATIONS

- For $m_\chi \lesssim 1 \text{ MeV}$, we cannot talk about isolated particles in the detector
- Dark matter will rather interact with collective excitations

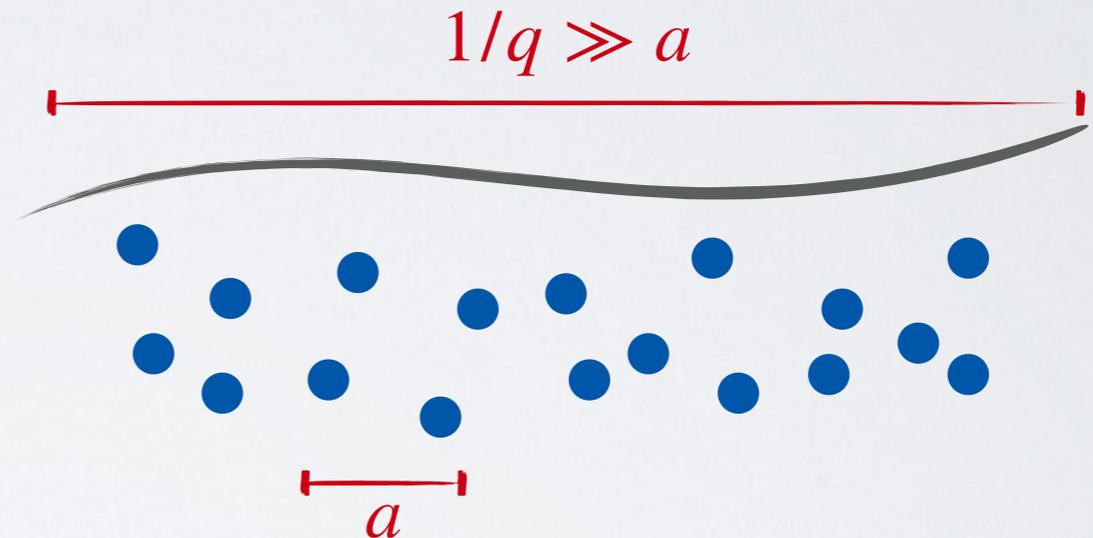
$$\frac{1}{q} \sim \frac{1}{m_\chi v_\chi} \gtrsim 1 \text{ \AA}$$



COLLECTIVE EXCITATIONS

- For $m_\chi \lesssim 1 \text{ MeV}$, we cannot talk about isolated particles in the detector
- Dark matter will rather interact with collective excitations

$$\frac{1}{q} \sim \frac{1}{m_\chi v_\chi} \gtrsim 1 \text{ \AA}$$



- Interesting and active field with plenty of ideas

* phonons on solid crystals

[e.g., Knapen et al. PLB 2018; Griffin et al. PRD 2021]

* phonons in superfluid ^4He

[e.g., Guo, McKinsey PRD 2013; Schutz, Zurek PRL 2016; Caputo, **AE**, Polosa PRD 2019]

* many others...

[for a review, Kahn, Lin Rept.Prog.Phys. 2022]

(ANTI-)FERROMAGNETS

(ANTI-)FERROMAGNETS

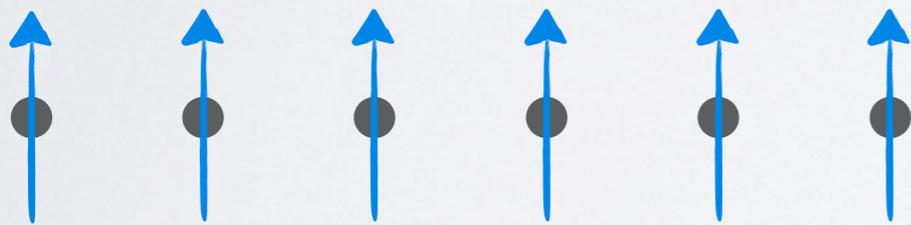
- A unexplored area is dark matter with *spin-dependent interactions*

(ANTI-)FERROMAGNETS

- A unexplored area is dark matter with *spin-dependent interactions*
- A possibility is to look for the interaction between dark matter and *spin-ordered systems* [Trickle, Zhang, Zurek PRL 2020; Mitridate et al. PRD 2020; Chigus, Moroi, Nakayama PRD 2020; Trickle, Zhang, Zurek PRD 2022]

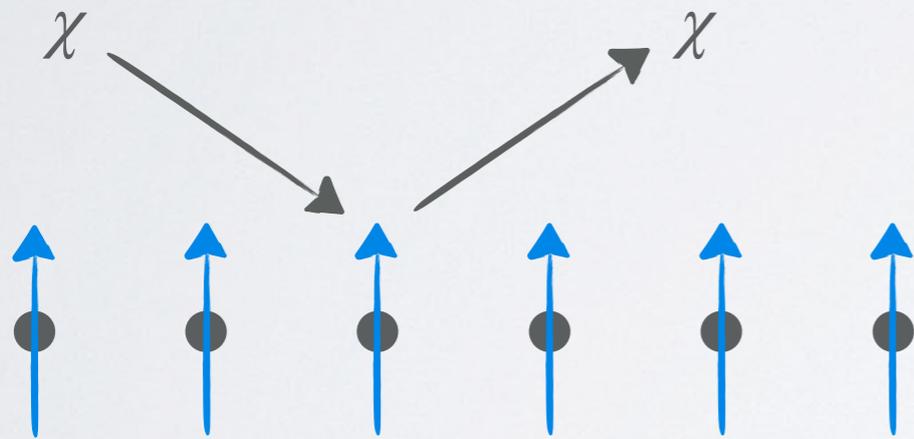
(ANTI-)FERROMAGNETS

- A unexplored area is dark matter with *spin-dependent interactions*
- A possibility is to look for the interaction between dark matter and *spin-ordered systems* [Trickle, Zhang, Zurek PRL 2020; Mitridate et al. PRD 2020; Chigus, Moroi, Nakayama PRD 2020; Trickle, Zhang, Zurek PRD 2022]



(ANTI-)FERROMAGNETS

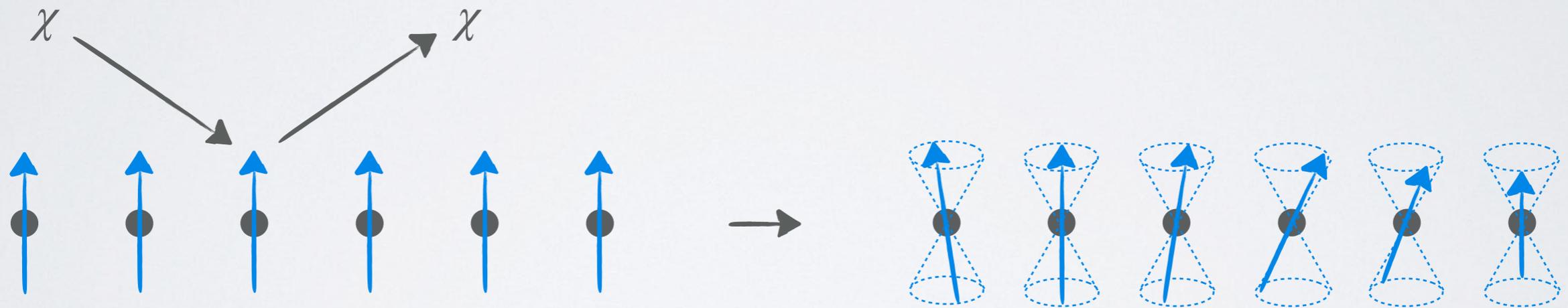
- A unexplored area is dark matter with *spin-dependent interactions*
- A possibility is to look for the interaction between dark matter and *spin-ordered systems* [Trickle, Zhang, Zurek PRL 2020; Mitridate et al. PRD 2020; Chigus, Moroi, Nakayama PRD 2020; Trickle, Zhang, Zurek PRD 2022]



(ANTI-)FERROMAGNETS

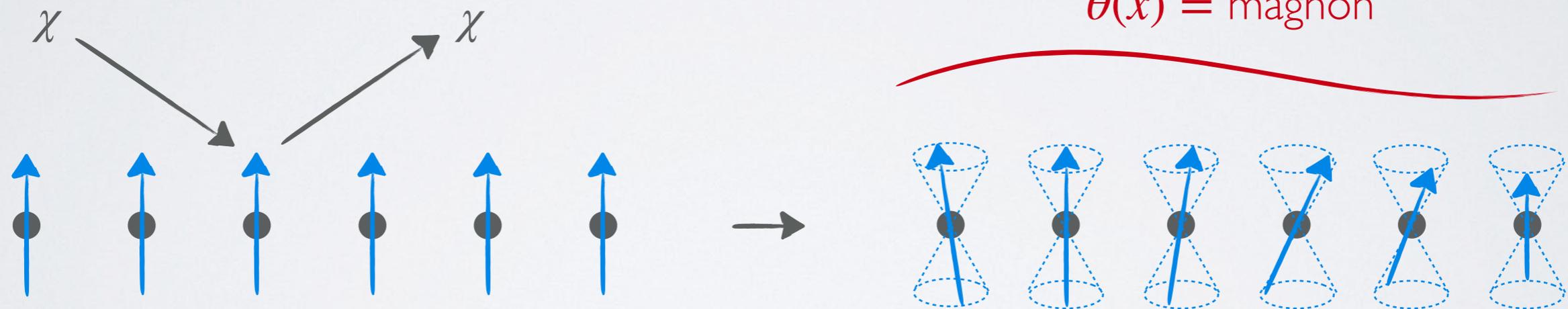
- A unexplored area is dark matter with *spin-dependent interactions*
- A possibility is to look for the interaction between dark matter and *spin-ordered systems*

[Trickle, Zhang, Zurek PRL 2020; Mitridate et al. PRD 2020; Chigus, Moroi, Nakayama PRD 2020; Trickle, Zhang, Zurek PRD 2022]



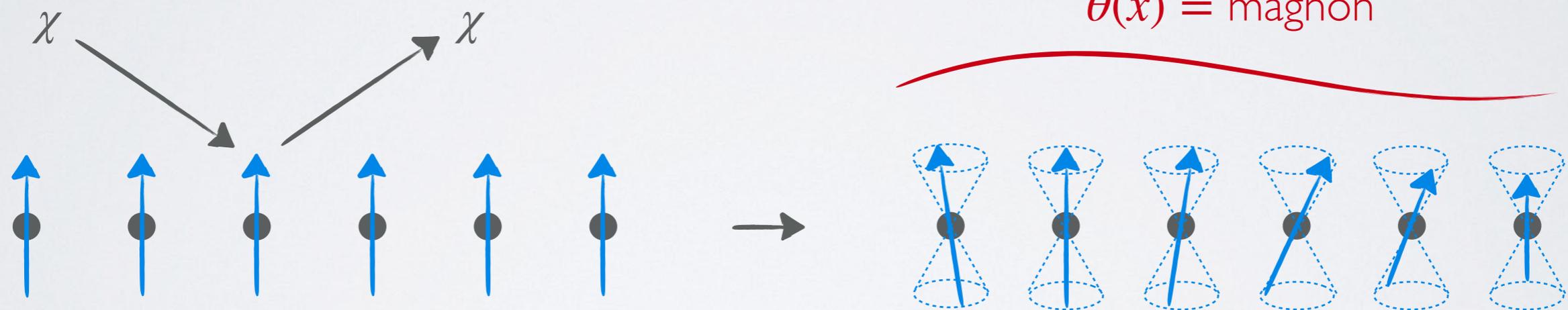
(ANTI-)FERROMAGNETS

- A unexplored area is dark matter with *spin-dependent interactions*
- A possibility is to look for the interaction between dark matter and *spin-ordered systems* [Trickle, Zhang, Zurek PRL 2020; Mitridate et al. PRD 2020; Chigus, Moroi, Nakayama PRD 2020; Trickle, Zhang, Zurek PRD 2022]



(ANTI-)FERROMAGNETS

- A unexplored area is dark matter with *spin-dependent interactions*
- A possibility is to look for the interaction between dark matter and *spin-ordered systems* [Trickle, Zhang, Zurek PRL 2020; Mitridate et al. PRD 2020; Chigus, Moroi, Nakayama PRD 2020; Trickle, Zhang, Zurek PRD 2022]



- Ways to detect few magnons have been proposed (TES? SQUIDs? quantum sensors? cavities?) [Trickle, Zhang, Zurek PRL 2020; Lachance-Quirion et al. Science Advances 2017; Lachance-Quirion et al. Science 2020]

DM-SPIN INTERACTION

DM-SPIN INTERACTION

- At low energies dark matter couples to *spin density field*

DM-SPIN INTERACTION

- At low energies **dark matter couples to *spin density field***
- Two benchmark models:

$$\mathcal{L}_{\text{m.d.}} \sim V_{\mu\nu} \bar{\chi} \sigma^{\mu\nu} \chi + V_{\mu} \bar{e} \gamma^{\mu} e$$

[e.g., Sigurdson et al. PRD 2004; Chang, Weiner, Yavin PRD 2010]

$$\mathcal{L}_{\text{p.m.}} \sim \phi \bar{\chi} \chi + \phi \bar{e} i \gamma^5 e$$

[e.g., Banks, Fortin, Thomas 1007.5515; Bagnasco, Dine, Thomas PLB 1994]

DM-SPIN INTERACTION

- At low energies **dark matter couples to spin density field**
- Two benchmark models:

$$\mathcal{L}_{\text{m.d.}} \sim V_{\mu\nu} \bar{\chi} \sigma^{\mu\nu} \chi + V_{\mu} \bar{e} \gamma^{\mu} e \quad [\text{e.g., Sigurdson et al. PRD 2004; Chang, Weiner, Yavin PRD 2010}]$$

$$\mathcal{L}_{\text{p.m.}} \sim \phi \bar{\chi} \chi + \phi \bar{e} i \gamma^5 e \quad [\text{e.g., Banks, Fortin, Thomas 1007.5515; Bagnasco, Dine, Thomas PLB 1994}]$$

- For a non-relativistic system, at low energies:

$$\mathcal{L}_{\text{m.d.}} \xrightarrow{\text{NR}} \chi^{\dagger} \sigma^i \chi (\delta^{ij} - \nabla^{-2} \nabla^i \nabla^j) e^{\dagger} \sigma^j e \xrightarrow{\text{IR}} \chi^{\dagger} \sigma^i \chi (\delta^{ij} - \nabla^{-2} \nabla^i \nabla^j) s_j$$

$$\mathcal{L}_{\text{p.m.}} \xrightarrow{\text{NR}} \chi^{\dagger} \chi \nabla^{-2} \nabla_i e^{\dagger} \sigma^i e \xrightarrow{\text{IR}} \chi^{\dagger} \chi \nabla^{-2} \nabla_i s^i$$

DM-SPIN INTERACTION

- At low energies **dark matter couples to spin density field**
- Two benchmark models:

$$\mathcal{L}_{\text{m.d.}} \sim V_{\mu\nu} \bar{\chi} \sigma^{\mu\nu} \chi + V_{\mu} \bar{e} \gamma^{\mu} e \quad [\text{e.g., Sigurdson et al. PRD 2004; Chang, Weiner, Yavin PRD 2010}]$$

$$\mathcal{L}_{\text{p.m.}} \sim \phi \bar{\chi} \chi + \phi \bar{e} i \gamma^5 e \quad [\text{e.g., Banks, Fortin, Thomas 1007.5515; Bagnasco, Dine, Thomas PLB 1994}]$$

- For a non-relativistic system, at low energies:

$$\mathcal{L}_{\text{m.d.}} \xrightarrow{\text{NR}} \chi^{\dagger} \sigma^i \chi (\delta^{ij} - \nabla^{-2} \nabla^i \nabla^j) e^{\dagger} \sigma^j e \xrightarrow{\text{IR}} \chi^{\dagger} \sigma^i \chi (\delta^{ij} - \nabla^{-2} \nabla^i \nabla^j) s_i$$

$$\mathcal{L}_{\text{p.m.}} \xrightarrow{\text{NR}} \chi^{\dagger} \chi \nabla^{-2} \nabla_i e^{\dagger} \sigma^i e \xrightarrow{\text{IR}} \chi^{\dagger} \chi \nabla^{-2} \nabla_i s^i$$

spin density

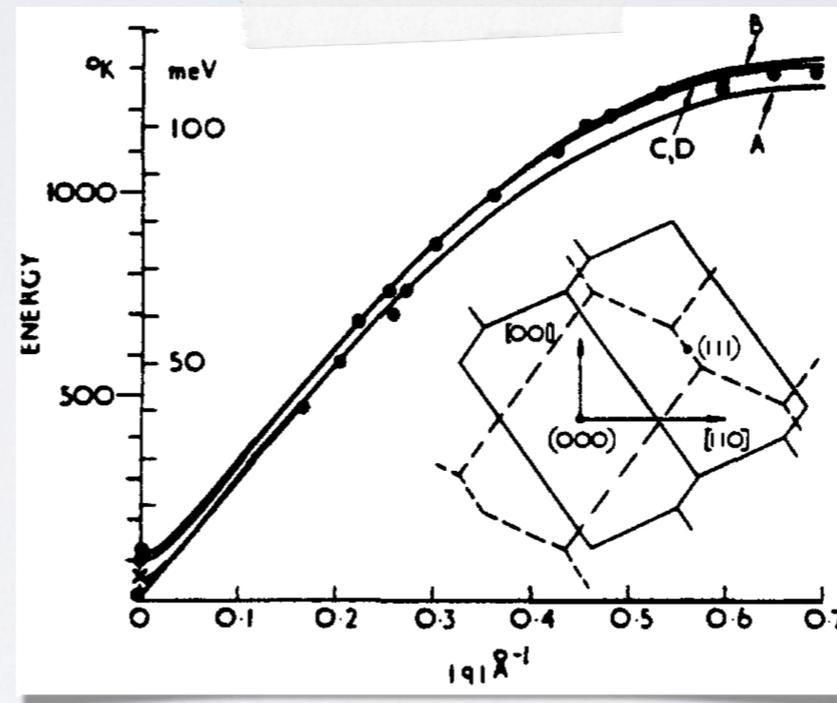
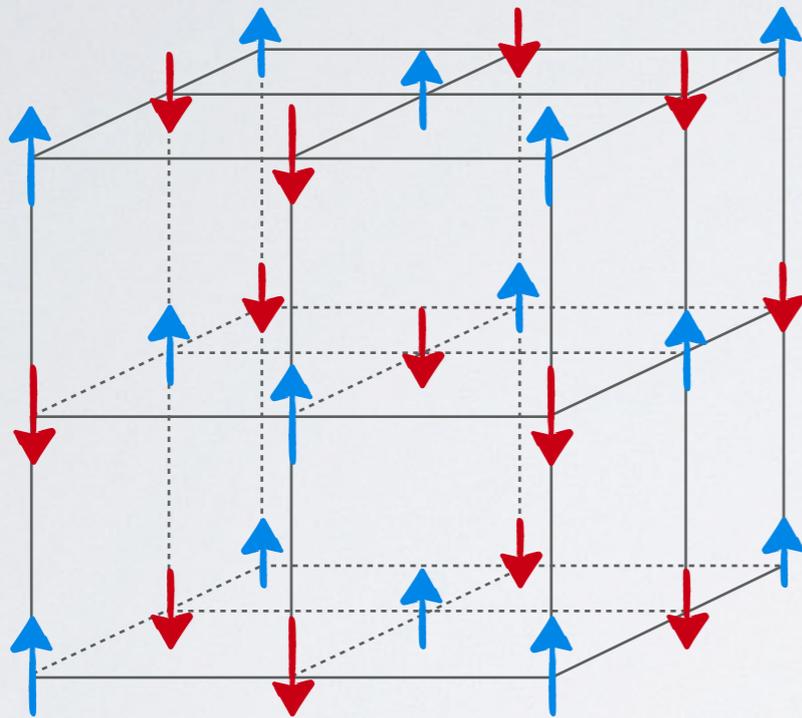
ANTI-FERROMAGNETS

ANTI-FERROMAGNETS

- An optimal class of materials turns out to be *anti-ferromagnets*

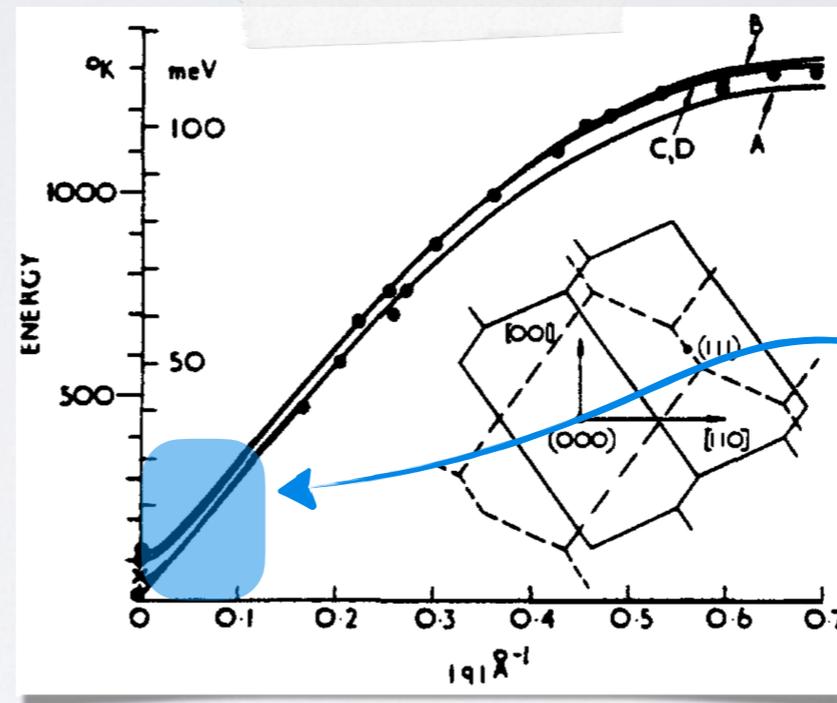
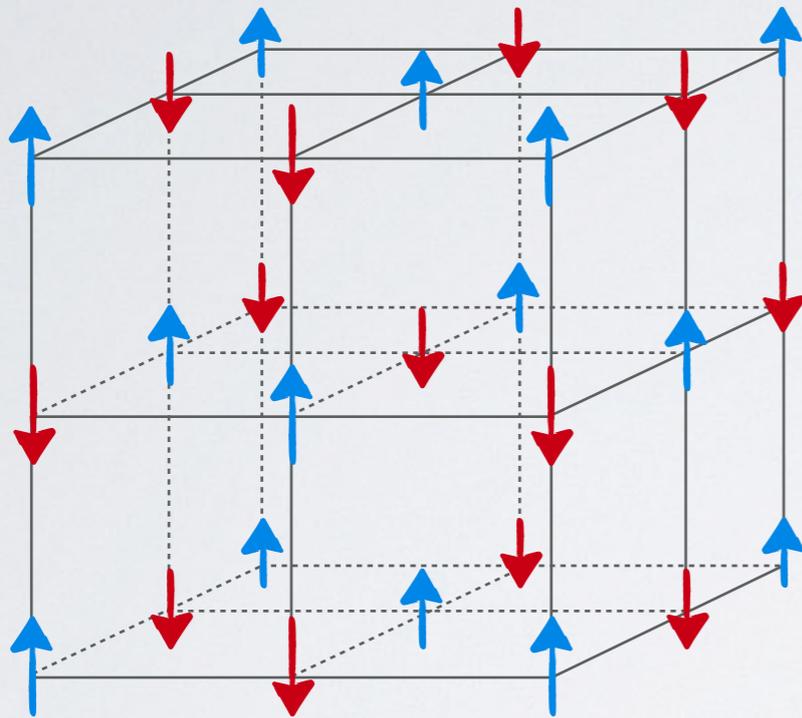
ANTI-FERROMAGNETS

- An optimal class of materials turns out to be *anti-ferromagnets*



ANTI-FERROMAGNETS

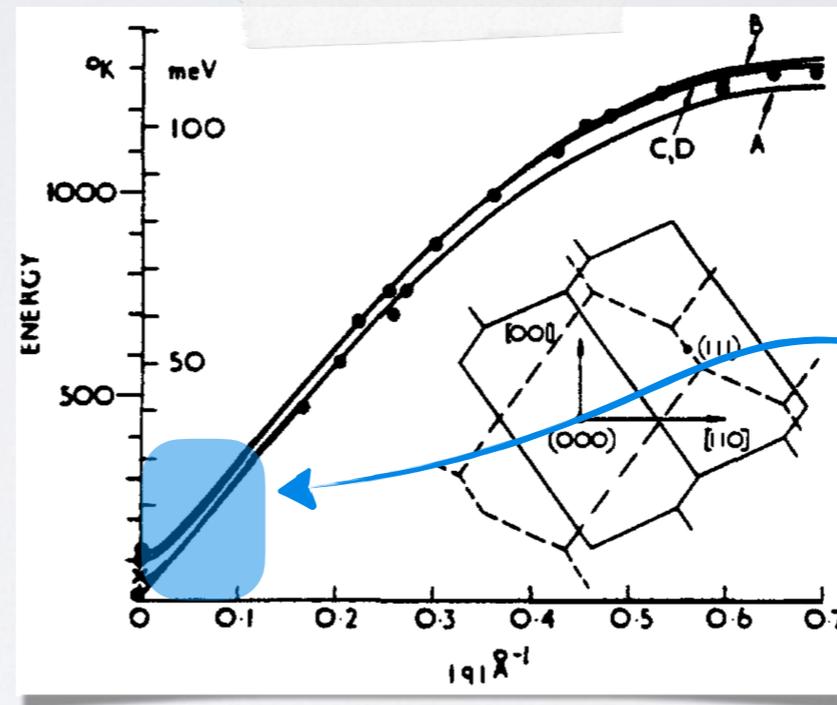
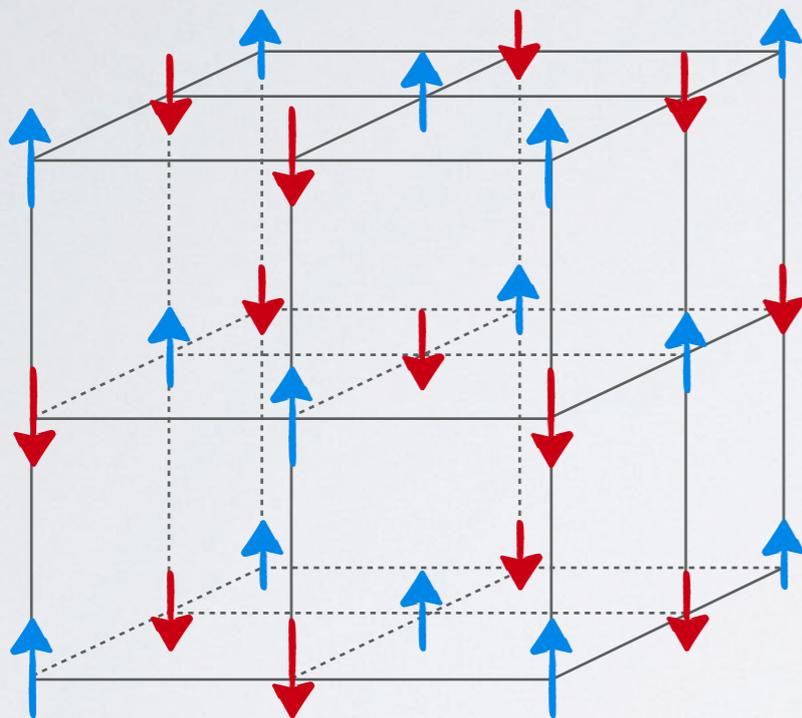
- An optimal class of materials turns out to be *anti-ferromagnets*



gapless magnons have
 $\omega(q) = v_{\theta} q$

ANTI-FERROMAGNETS

- An optimal class of materials turns out to be *anti-ferromagnets*

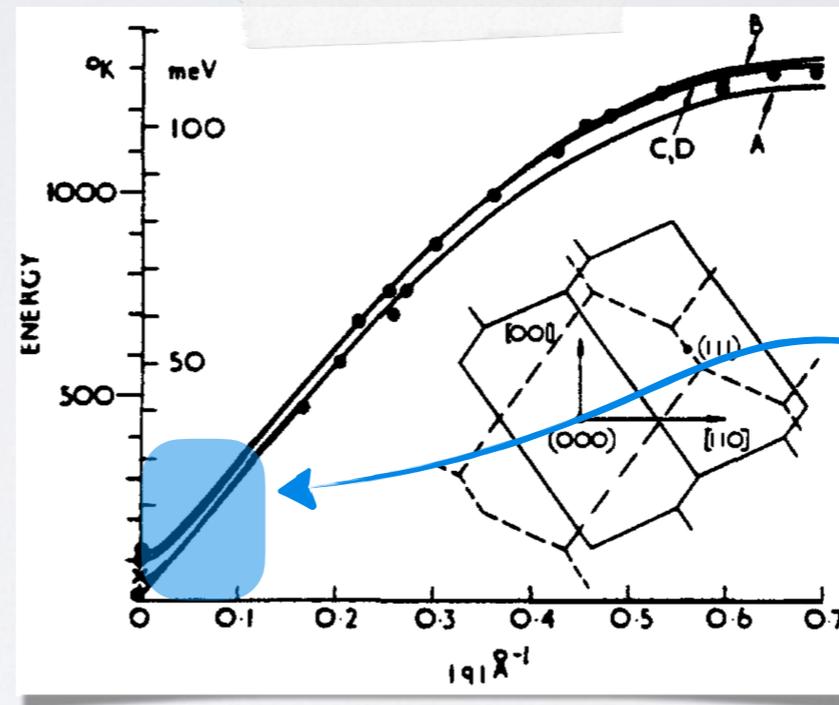
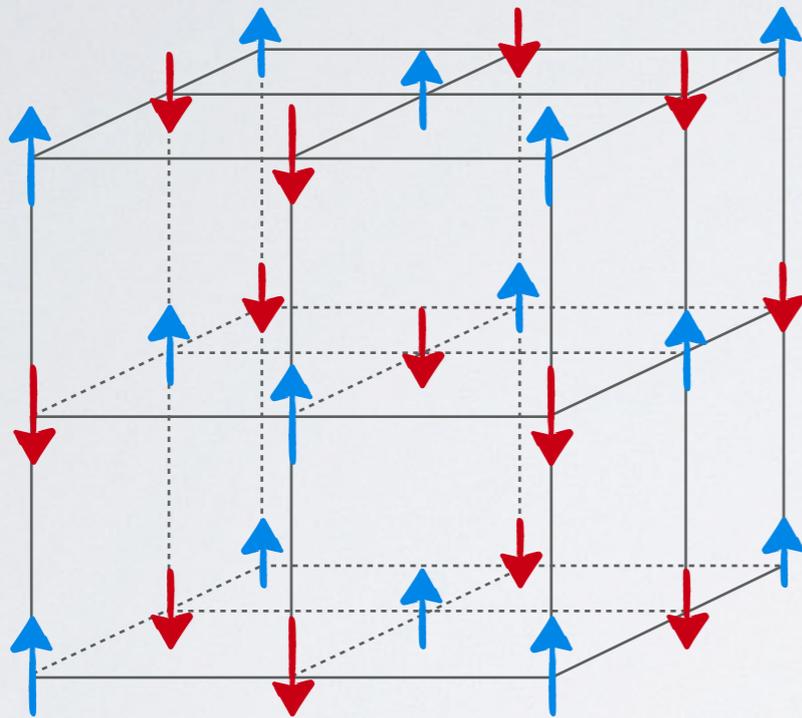


gapless magnons have
 $\omega(q) = v_{\theta} q$

- For single-magnon emission: $\omega_{\max} \simeq 4 \left(v_{\theta} / v_{\chi} \right) E_{\chi}$

ANTI-FERROMAGNETS

- An optimal class of materials turns out to be *anti-ferromagnets*



gapless magnons have
 $\omega(q) = v_{\theta} q$

- For single-magnon emission: $\omega_{\max} \simeq 4 \left(v_{\theta} / v_{\chi} \right) E_{\chi}$
- Nickel-oxide has $v_{\theta} \sim 0.1 v_{\chi}$ \rightarrow very efficient at absorbing dark matter energy [AE, Pavaskar PRD 2023]

ANTI-FERROMAGNETS

ANTI-FERROMAGNETS

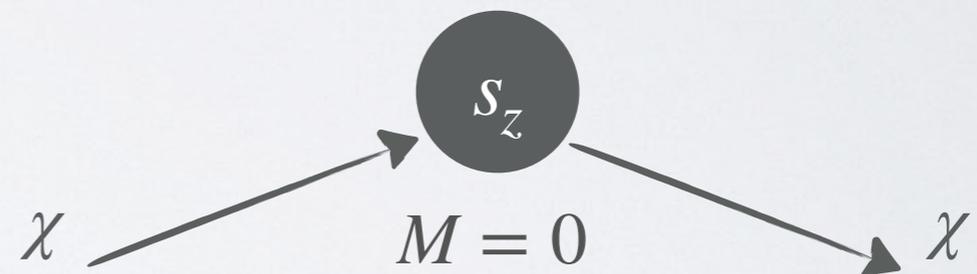
- Magnons in anti-ferromagnets have **two polarizations**, analogous to particle and anti-particle in a relativistic theory

ANTI-FERROMAGNETS

- Magnons in anti-ferromagnets have **two polarizations**, analogous to particle and anti-particle in a relativistic theory
- This allows to emit **magnon and anti-magnon pairs** while preserving magnetization

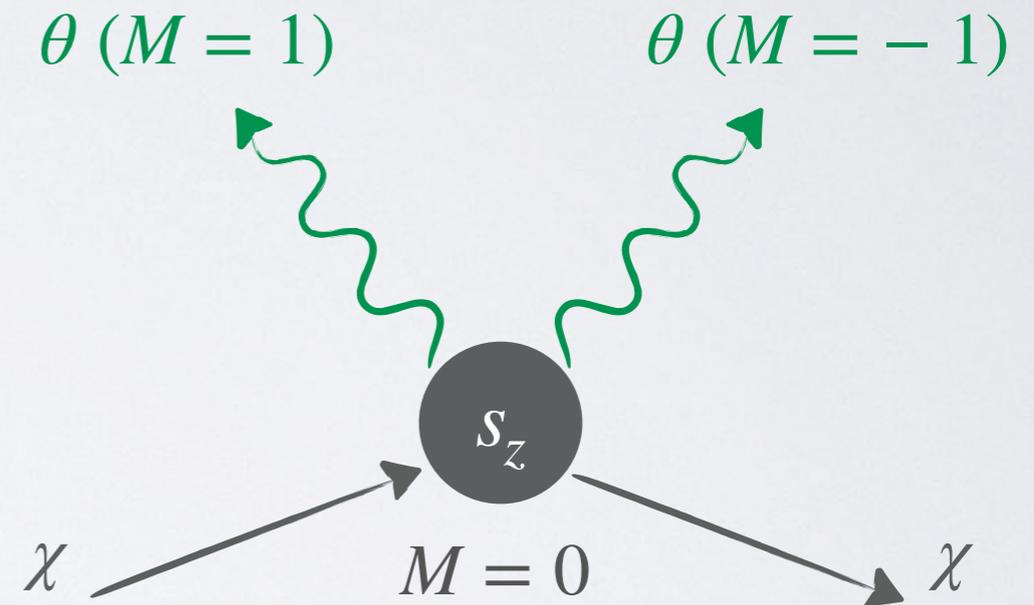
ANTI-FERROMAGNETS

- Magnons in anti-ferromagnets have **two polarizations**, analogous to particle and anti-particle in a relativistic theory
- This allows to emit **magnon and anti-magnon pairs** while preserving magnetization



ANTI-FERROMAGNETS

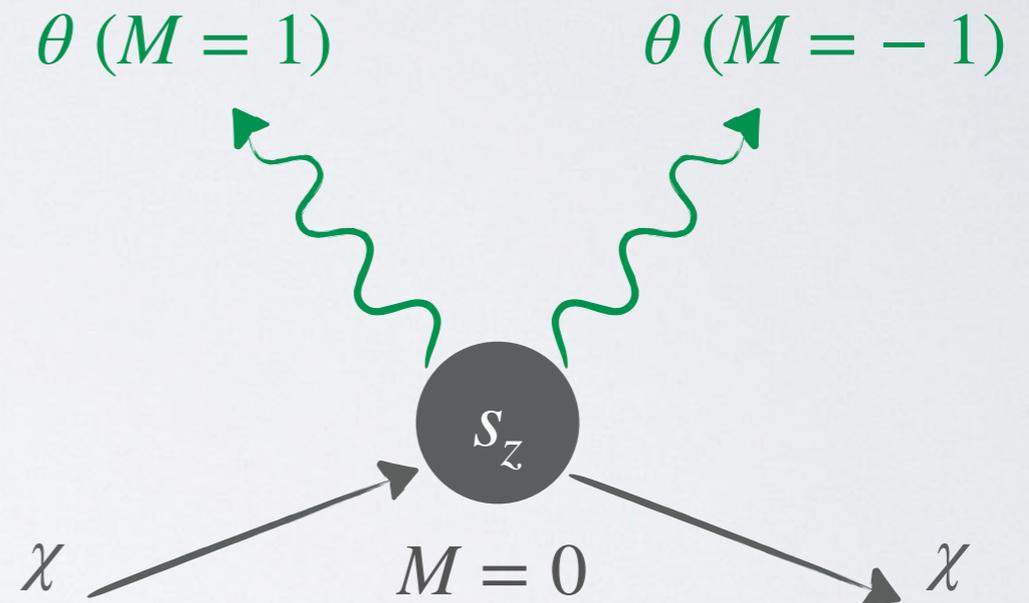
- Magnons in anti-ferromagnets have **two polarizations**, analogous to particle and anti-particle in a relativistic theory
- This allows to emit **magnon and anti-magnon pairs** while preserving magnetization



ANTI-FERROMAGNETS

- Magnons in anti-ferromagnets have **two polarizations**, analogous to particle and anti-particle in a relativistic theory

- This allows to emit **magnon and anti-magnon pairs** while preserving magnetization



- **Multi-magnon emission process** evades the kinematical constraints and **get down to $m_\chi \sim \mathcal{O}(\text{keV})$**

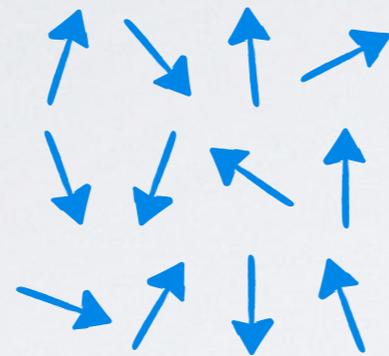
MAGNONS

MAGNONS

- Anti-ferromagnets spontaneously break internal spin symmetry

MAGNONS

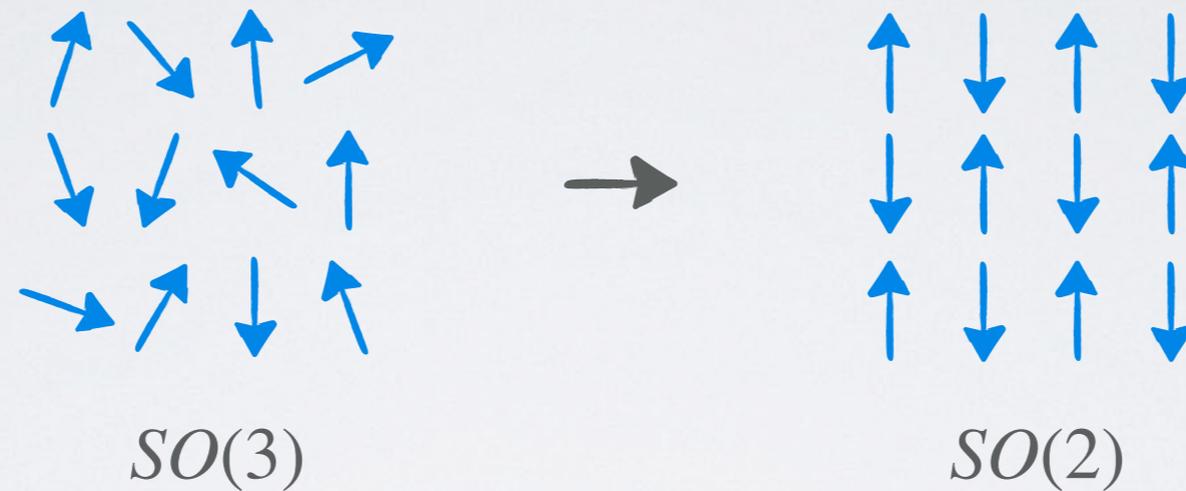
- Anti-ferromagnets spontaneously break internal spin symmetry



$SO(3)$

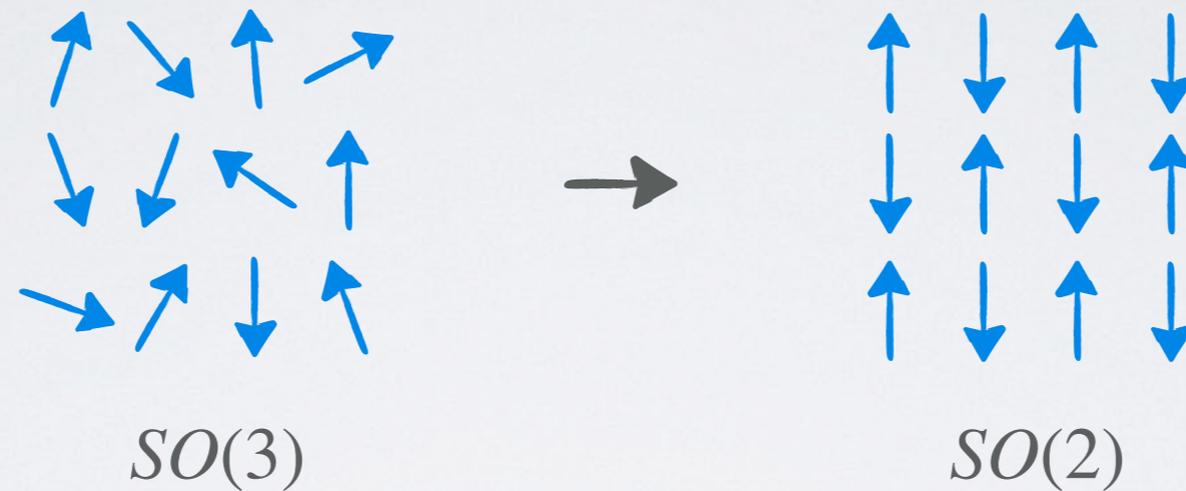
MAGNONS

- Anti-ferromagnets spontaneously break internal spin symmetry



MAGNONS

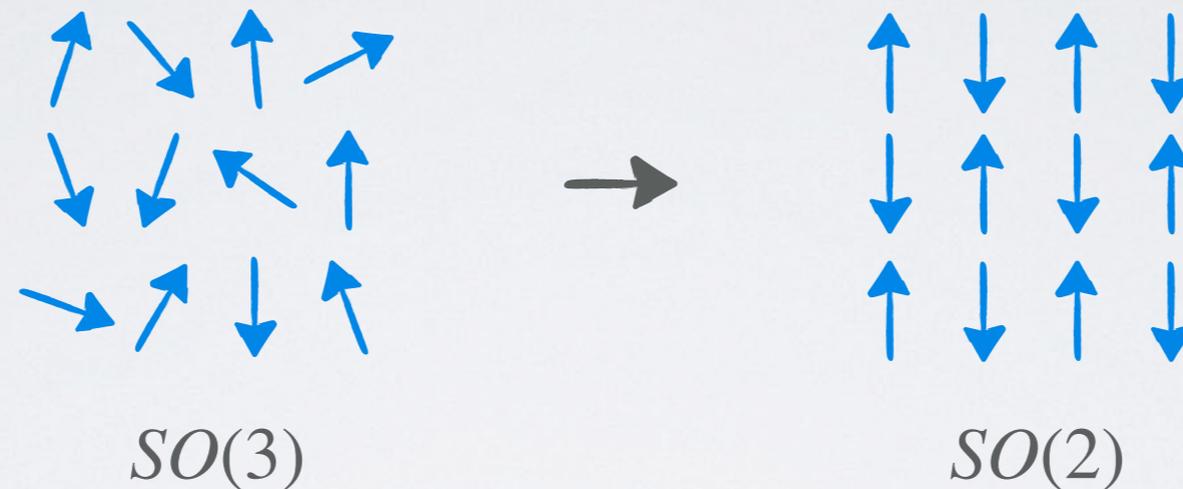
- Anti-ferromagnets spontaneously break internal spin symmetry



Gapless magnon = Goldstone

MAGNONS

- Anti-ferromagnets spontaneously break internal spin symmetry



Gapless magnon = Goldstone

- At low energies/momenta magnons can be described by an EFT:
 - * invariant under the full symmetry group
 - * organized in a derivative expansion

MAGNONS

MAGNONS

- Very similar to the non-linear σ -model:

MAGNONS

- Very similar to the non-linear σ -model:

$$\mathbf{n}(x) = e^{i[\theta^1(x)J_1 + \theta^2(x)J_2]} \cdot \hat{\mathbf{z}} \xrightarrow{SO(3)} R \cdot \mathbf{n}(x)$$

MAGNONS

- Very similar to the non-linear σ -model:

$$\mathbf{n}(x) = e^{i[\theta^1(x)J_1 + \theta^2(x)J_2]} \cdot \hat{\mathbf{z}} \xrightarrow{SO(3)} R \cdot \mathbf{n}(x)$$

- At lowest order in the derivative expansion, the **most general invariant Lagrangian (density)** is

MAGNONS

- Very similar to the non-linear σ -model:

$$\mathbf{n}(x) = e^{i[\theta^1(x)J_1 + \theta^2(x)J_2]} \cdot \hat{\mathbf{z}} \xrightarrow{SO(3)} R \cdot \mathbf{n}(x)$$

- At lowest order in the derivative expansion, the **most general invariant Lagrangian (density)** is

$$\mathcal{L} = c_1 \dot{\mathbf{n}}^2 - c_2 (\nabla_i \mathbf{n})^2$$

MAGNONS

- Very similar to the non-linear σ -model:

$$\mathbf{n}(x) = e^{i[\theta^1(x)J_1 + \theta^2(x)J_2]} \cdot \hat{\mathbf{z}} \xrightarrow{SO(3)} R \cdot \mathbf{n}(x)$$

- At lowest order in the derivative expansion, the **most general invariant Lagrangian (density)** is

$$\begin{aligned}\mathcal{L} &= c_1 \dot{\mathbf{n}}^2 - c_2 (\nabla_i \mathbf{n})^2 \\ &= c_1 (\dot{\theta}^a)^2 - c_2 (\vec{\nabla} \theta^a)^2 + \dots\end{aligned}$$

MAGNONS

- Very similar to the non-linear σ -model:

$$\mathbf{n}(x) = e^{i[\theta^1(x)J_1 + \theta^2(x)J_2]} \cdot \hat{\mathbf{z}} \xrightarrow{SO(3)} R \cdot \mathbf{n}(x)$$

- At lowest order in the derivative expansion, the **most general invariant Lagrangian (density)** is

$$\begin{aligned} \mathcal{L} &= c_1 \dot{\mathbf{n}}^2 - c_2 (\nabla_i \mathbf{n})^2 \\ &= c_1 (\dot{\theta}^a)^2 - c_2 (\vec{\nabla} \theta^a)^2 + \dots \end{aligned}$$

can be extracted from
dispersion relation +
neutron scattering data

$$v_\theta = c_2/c_1$$

$$\sigma_n \propto c_1$$

[Pavaskar, Penco, Rothstein SciPost Phys. 2022; **AE**, Pavaskar PRD 2023]

MAGNONS

MAGNONS

- Recall that the dark matter interacts via **spin density, $\mathbf{s}(x)$**

MAGNONS

- Recall that the dark matter interacts via **spin density, $\mathbf{s}(x)$**
- Easily computed as **$SO(3)$ Noether current in the EFT:**

MAGNONS

- Recall that the dark matter interacts via **spin density, $\mathbf{s}(x)$**
- Easily computed as **$SO(3)$ Noether current in the EFT:**

$$s_i = c_1 (\mathbf{n} \times \dot{\mathbf{n}})_i = c_1 \left[\delta_{ia} \dot{\theta}^a + \delta_{i3} \epsilon_{ab} \theta^a \dot{\theta}^b + \dots \right]$$

[AE, Pavaskar PRD 2023]

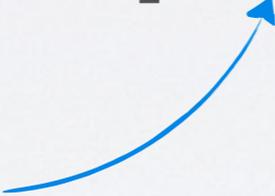
MAGNONS

- Recall that the dark matter interacts via **spin density, $\mathbf{s}(x)$**
- Easily computed as **$SO(3)$ Noether current in the EFT:**

$$s_i = c_1 (\mathbf{n} \times \dot{\mathbf{n}})_i = c_1 \left[\delta_{ia} \dot{\theta}^a + \delta_{i3} \epsilon_{ab} \theta^a \dot{\theta}^b + \dots \right]$$

[AE, Pavaskar PRD 2023]

one-magnon
emission



MAGNONS

- Recall that the dark matter interacts via **spin density, $\mathbf{s}(x)$**
- Easily computed as **$SO(3)$ Noether current in the EFT:**

$$s_i = c_1 (\mathbf{n} \times \dot{\mathbf{n}})_i = c_1 \left[\delta_{ia} \dot{\theta}^a + \delta_{i3} \epsilon_{ab} \theta^a \dot{\theta}^b + \dots \right]$$

[AE, Pavaskar PRD 2023]

one-magnon emission

two-magnons emission

MAGNONS

- Recall that the dark matter interacts via **spin density, $\mathbf{s}(x)$**
- Easily computed as **$SO(3)$ Noether current in the EFT:**

$$s_i = c_1 (\mathbf{n} \times \dot{\mathbf{n}})_i = c_1 \left[\delta_{ia} \dot{\theta}^a + \delta_{i3} \epsilon_{ab} \theta^a \dot{\theta}^b + \dots \right] \quad [\text{AE, Pavaskar PRD 2023}]$$

one-magnon emission \rightarrow $\delta_{ia} \dot{\theta}^a$

$\theta^a \dot{\theta}^b$ \rightarrow two-magnons emission

- Structure **completely dictated by symmetry** \rightarrow just need c_1

MAGNONS

- Recall that the dark matter interacts via **spin density, $\mathbf{s}(x)$**
- Easily computed as **$SO(3)$ Noether current in the EFT:**

$$s_i = c_1 (\mathbf{n} \times \dot{\mathbf{n}})_i = c_1 \left[\delta_{ia} \dot{\theta}^a + \delta_{i3} \epsilon_{ab} \theta^a \dot{\theta}^b + \dots \right] \quad [\text{AE, Pavaskar PRD 2023}]$$

one-magnon emission \rightarrow $\delta_{ia} \dot{\theta}^a$

$\theta^a \dot{\theta}^b$ \rightarrow two-magnons emission

- Structure **completely dictated by symmetry** \rightarrow just need c_1
- Much more “HEP friendly” than standard language

IDEAL REACH

IDEAL REACH

- Local QFT Lagrangian → use [standard QFT methods](#) to compute event rates [AE, Pavaskar PRD 2023]

IDEAL REACH

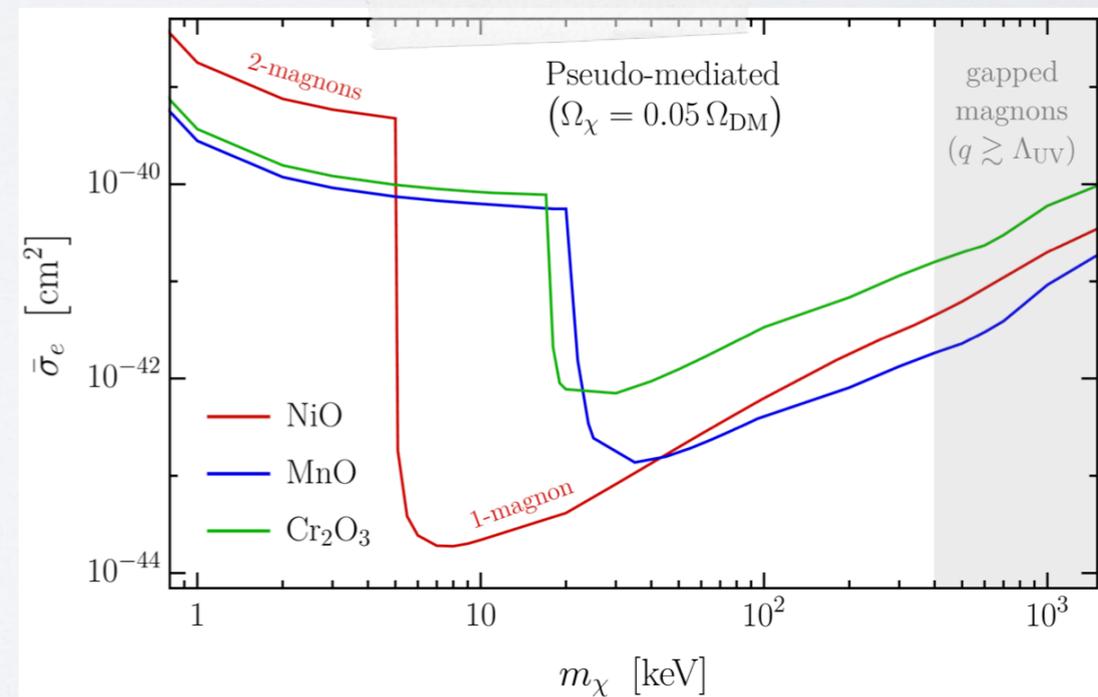
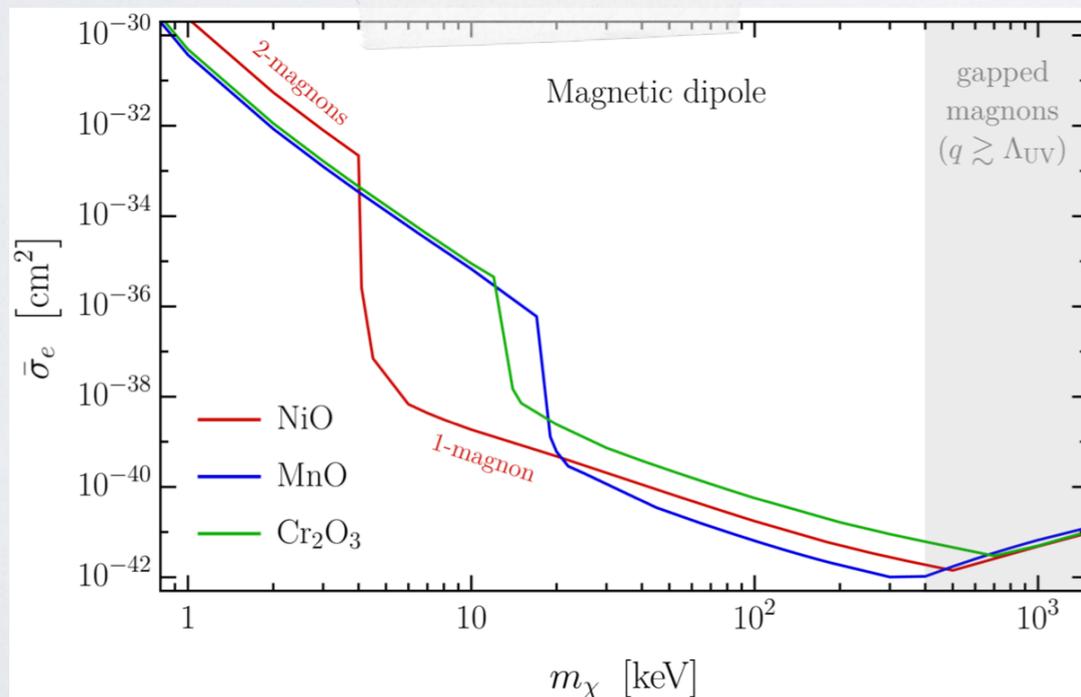
- Local QFT Lagrangian \rightarrow use [standard QFT methods](#) to compute event rates [AE, Pavaskar PRD 2023]

$$\begin{aligned}
 \begin{array}{c} a, \lambda_1 \\ \vdots \\ s \rightarrow \bullet \rightarrow s' \end{array} &= -\frac{g_\chi g_e \sqrt{c_1}}{m_e} \omega \times \begin{cases} \frac{4}{\Lambda_x} P_{ia}(\mathbf{q}) \sigma^i & \text{m.d.} \\ q^a / q^2 & \text{p.m.} \end{cases}, \\
 \begin{array}{c} a, \lambda_1 \quad b, \lambda_2 \\ \vdots \quad \vdots \\ s \rightarrow \bullet \rightarrow s' \end{array} &= \frac{g_\chi g_e}{m_e} (\omega_1 - \omega_2) \epsilon_{ab} \times \begin{cases} \frac{4}{\Lambda_x} P_{iz}(\mathbf{q}) \sigma^i & \text{m.d.} \\ q^z / q^2 & \text{p.m.} \end{cases}
 \end{aligned}$$

IDEAL REACH

- Local QFT Lagrangian \rightarrow use **standard QFT methods** to compute event rates [AE, Pavaskar PRD 2023]

$$\begin{aligned}
 \begin{array}{c} a, \lambda_1 \\ \uparrow \\ s \rightarrow \bullet \rightarrow s' \end{array} &= -\frac{g_\chi g_e \sqrt{c_1}}{m_e} \omega \times \begin{cases} \frac{4}{\Lambda_\chi} P_{ia}(\mathbf{q}) \sigma^i & \text{m.d.} \\ q^a / q^2 & \text{p.m.} \end{cases}, \\
 \begin{array}{c} a, \lambda_1 \quad b, \lambda_2 \\ \uparrow \quad \uparrow \\ s \rightarrow \bullet \rightarrow s' \end{array} &= \frac{g_\chi g_e}{m_e} (\omega_1 - \omega_2) \epsilon_{ab} \times \begin{cases} \frac{4}{\Lambda_\chi} P_{iz}(\mathbf{q}) \sigma^i & \text{m.d.} \\ q^z / q^2 & \text{p.m.} \end{cases}
 \end{aligned}$$



AXIONS

AXIONS

- The **very same material** can be used to look for **axion dark matter**,
coupling to electrons [Catinari, **AE**, Pavaskar 2411.11971]

AXIONS

- The **very same material** can be used to look for **axion dark matter, coupling to electrons** [Catinari, **AE**, Pavaskar 2411.11971]
- The axion can be absorbed by the anti-ferromagnet:

AXIONS

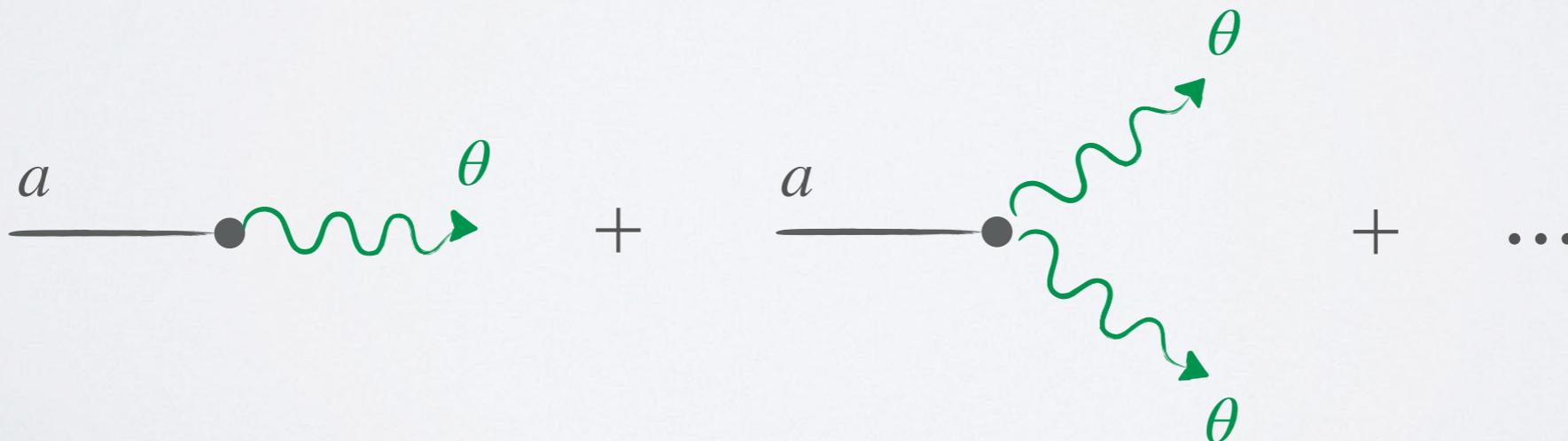
- The **very same material** can be used to look for **axion dark matter, coupling to electrons** [Catinari, **AE**, Pavaskar 2411.11971]
- The axion can be absorbed by the anti-ferromagnet:

$$\mathcal{L} = \frac{g_{aee}}{2m_e} \partial_\mu a (\bar{e} \gamma^\mu \gamma^5 e) \longrightarrow \frac{g_{aee}}{m_e} \vec{\nabla} a \cdot \vec{s}$$

AXIONS

- The **very same material** can be used to look for **axion dark matter, coupling to electrons** [Catinari, **AE**, Pavaskar 2411.11971]
- The axion can be absorbed by the anti-ferromagnet:

$$\mathcal{L} = \frac{g_{aee}}{2m_e} \partial_\mu a (\bar{e} \gamma^\mu \gamma^5 e) \longrightarrow \frac{g_{aee}}{m_e} \vec{\nabla} a \cdot \vec{s}$$



AXIONS

AXIONS

- For axion conversion to a single magnon, energy and momentum conservation implies

$$m_a = \omega(m_a \mathbf{v}_a) \simeq \omega(0)$$

AXIONS

- For **axion conversion to a single magnon**, energy and momentum conservation implies

$$m_a = \omega(m_a \mathbf{v}_a) \simeq \omega(0)$$

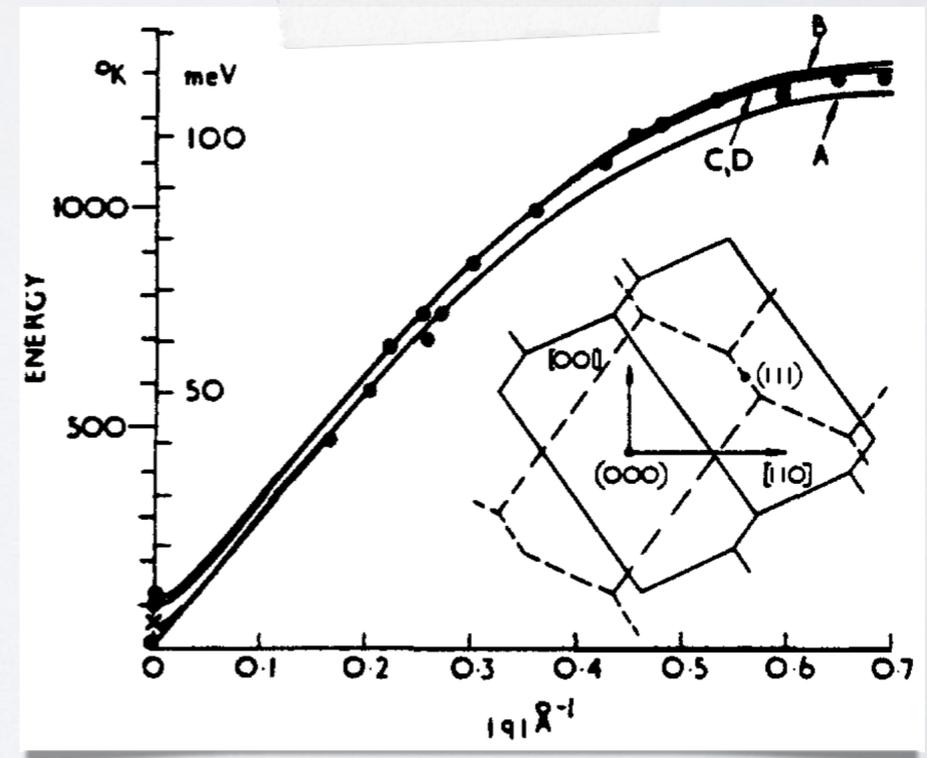
- It is necessary for magnons to have a **finite gap**, $\omega(0) > 0$

AXIONS

- For axion conversion to a single magnon, energy and momentum conservation implies

$$m_a = \omega(m_a \mathbf{v}_a) \simeq \omega(0)$$

- It is necessary for magnons to have a finite gap, $\omega(0) > 0$
- Fear not! This is indeed what happens in reality



AXIONS

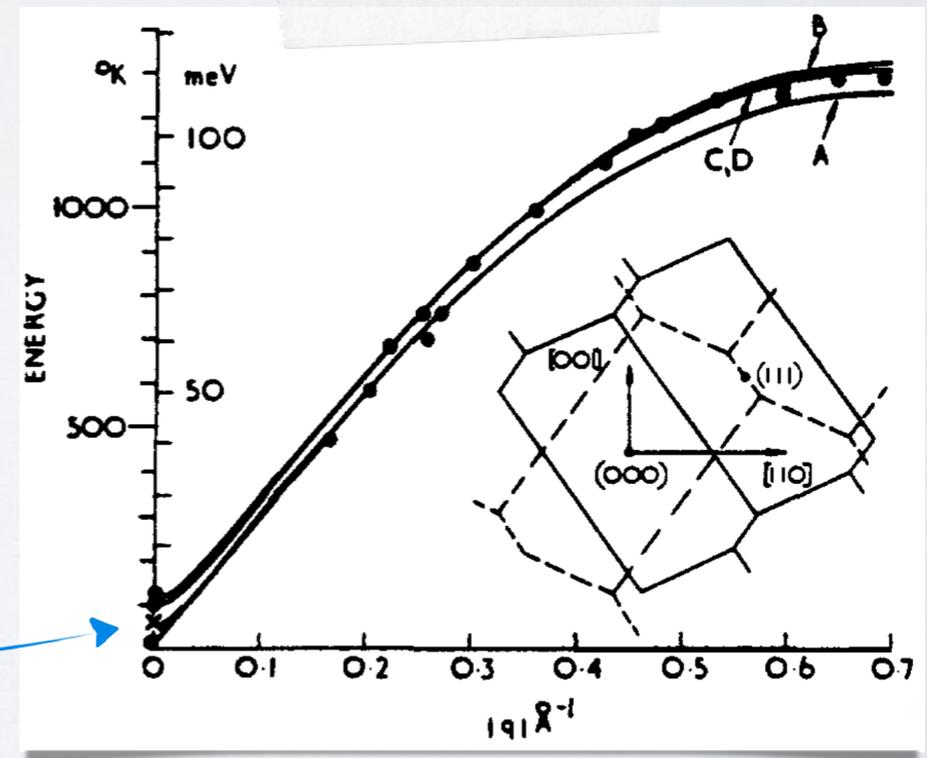
- For axion conversion to a single magnon, energy and momentum conservation implies

$$m_a = \omega(m_a \mathbf{v}_a) \simeq \omega(0)$$

- It is necessary for magnons to have a finite gap, $\omega(0) > 0$

- Fear not! This is indeed what happens in reality

real magnons have non-zero gap



AXIONS

AXIONS

- To study this case we extended the EFT to include more details of the structure of NiO:

AXIONS

- To study this case we extended the EFT to include more details of the structure of NiO:
 - A. intrinsic magnetic anisotropy (induces a gap)

AXIONS

- To study this case we extended the EFT to include more details of the structure of NiO:
 - A. intrinsic magnetic anisotropy (induces a gap)
 - B. external magnetic field (makes the gap *tunable*)

AXIONS

- To study this case we extended the EFT to include **more details of the structure of NiO**:
 - A. intrinsic magnetic anisotropy (induces a gap)
 - B. external magnetic field (makes the gap *tunable*)
- The EFT is slightly more complicated, but much richer:

AXIONS

- To study this case we extended the EFT to include **more details of the structure of NiO**:
 - A. intrinsic magnetic anisotropy (induces a gap)
 - B. external magnetic field (makes the gap *tunable*)
- The EFT is slightly more complicated, but much richer:

$$\mathcal{L} = c_1 \left[(\dot{\mathbf{n}} + \mu \mathbf{B} \times \mathbf{n})^2 - v_\theta^2 (\nabla_i \mathbf{n})^2 + \lambda_z n_z^2 - \lambda_x n_x^2 \right]$$

[Catinari, **AE**, Pavaskar 2411.09761]

AXIONS

- To study this case we extended the EFT to include **more details of the structure of NiO**:
 - A. intrinsic magnetic anisotropy (induces a gap)
 - B. external magnetic field (makes the gap *tunable*)
- The EFT is slightly more complicated, but much richer:

$$\mathcal{L} = c_1 \left[(\dot{\mathbf{n}} + \mu \mathbf{B} \times \mathbf{n})^2 - v_\theta^2 (\nabla_i \mathbf{n})^2 + \lambda_z n_z^2 - \lambda_x n_x^2 \right]$$

[Catinari, **AE**, Pavaskar 2411.09761]

- Rather **non-trivial field theory!**

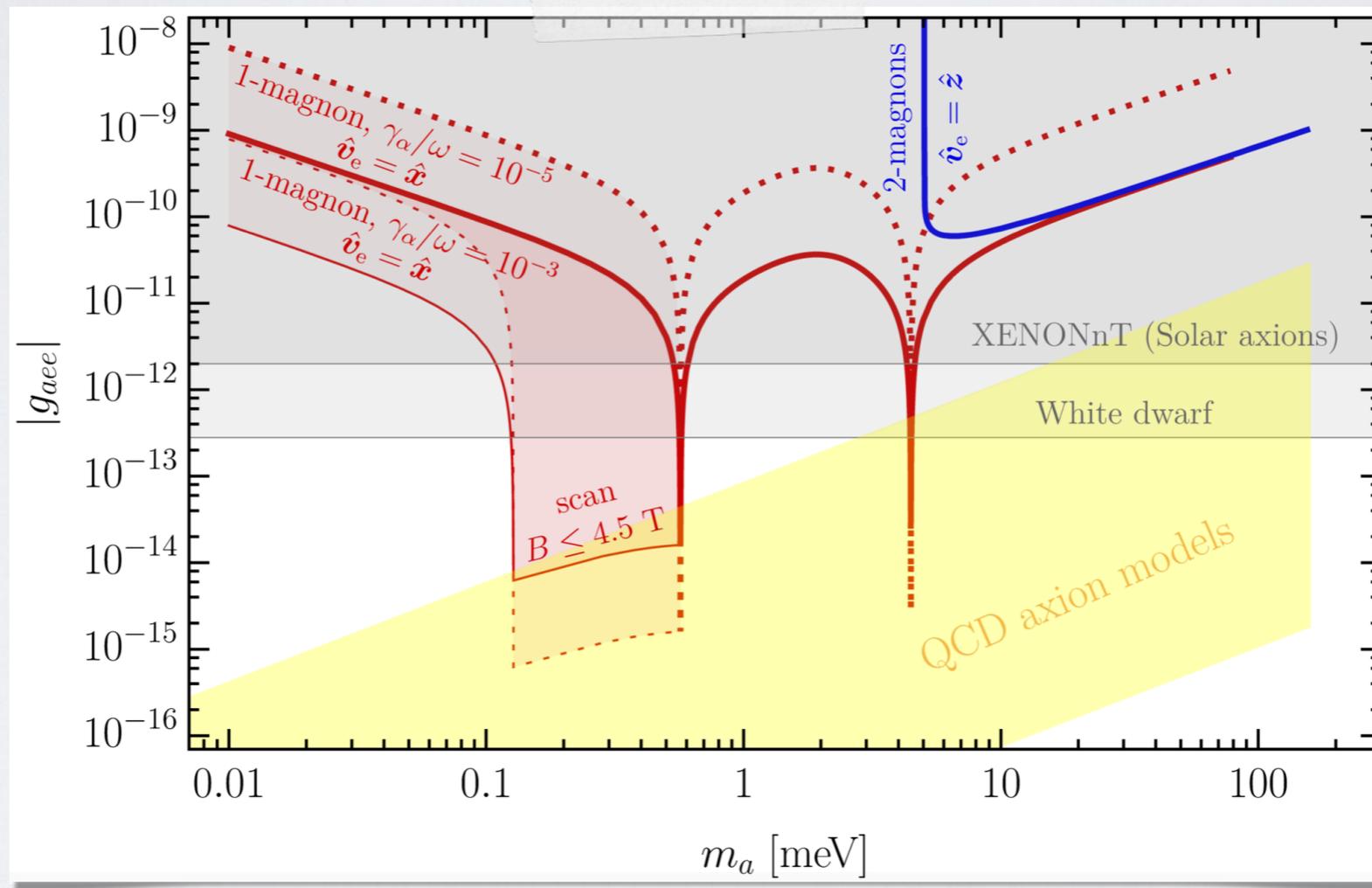
AXIONS

AXIONS

- Nickel-oxide is *very promising also for axion searches:*

AXIONS

- Nickel-oxide is **very promising** also for axion searches:



[Catinari, **AE**, Pavaskar 2411.09761]

WHAT NEXT?

- A plethora of open questions:
 - ▶ is any other good material out there? [Marocco, Wheeler 2501.18120]
 - ▶ what is the **actual observable**? How do we see magnons?

$$H = \mu\mathbf{B} \cdot \mathbf{S} \quad \Rightarrow \quad |\text{phys}\rangle = |\theta\rangle + \alpha|\gamma\rangle$$

produce



detect



- ▶ What the actual **magnon lifetime**? [work in progress w/ Carugno, Catinari, Pavaskar]

$$\frac{\Gamma(\theta \rightarrow \gamma)}{\Gamma(\theta \rightarrow \text{phonons})} = ?$$

WHAT NEXT?

WHAT NEXT?

- The search for sub-MeV dark matter requires new ideas

WHAT NEXT?

- The search for sub-MeV dark matter requires new ideas
- Condensed matter effects (or signatures!) become unavoidable and one must find a way to incorporate them in a theoretical framework

WHAT NEXT?

- The search for sub-MeV dark matter requires new ideas
- Condensed matter effects (or signatures!) become unavoidable and one must find a way to incorporate them in a theoretical framework
- Many particle physics programs find themselves at the edge with condensed matter. This can be a challenge... but also a great asset!

WHAT NEXT?

- The search for sub-MeV dark matter requires new ideas
- Condensed matter effects (or signatures!) become unavoidable and one must find a way to incorporate them in a theoretical framework
- Many particle physics programs find themselves at the edge with condensed matter. This can be a challenge... but also a great asset!

Thank you for the attention!